



Effective Models of Nonsingular Black Holes

Speaker: Andrea Pierfrancesco Sanna

Based on: Phys. Rev. D **106** (2022), arXiv:2204.09444

In collaboration with: Mariano Cadoni, Mauro Oi

INFN, Sezione di Cagliari, Università degli Studi di Cagliari



The classical unavailability of spacetime singularities

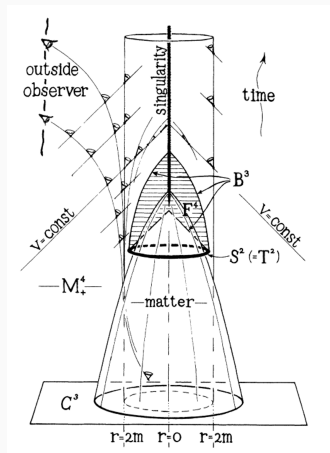
Gravitational collapse in general relativity leads to **unavoidable** spacetime singularities^[Penrose (1964)]

The classical unavoidability of spacetime singularities

Gravitational collapse in general relativity leads to **unavoidable** spacetime singularities [Penrose (1964)]

Very few ingredients

- Validity of Einstein's equations
- Energy conditions
- Global Hyperbolicity



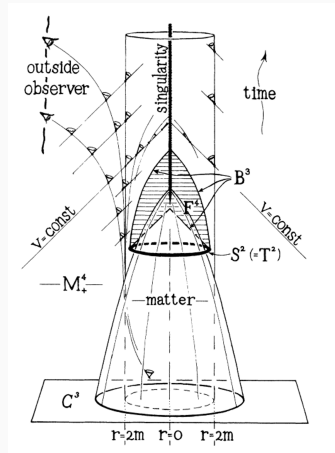
From Penrose (1964)

The classical unavailability of spacetime singularities

Gravitational collapse in general relativity leads to **unavoidable** spacetime singularities [\[Penrose \(1964\)\]](#)

Very few ingredients

- Validity of Einstein's equations
- Energy conditions
- Global Hyperbolicity



From Penrose (1964)

Relaxing one (or more) assumptions allows to circumvent the theorem and obtain regular spacetimes [\[Carballo-Rubio et al. \(2019\)\]](#)

Common assumption: corrections to GR at the [ultraviolet](#) scale

$$l_P = \sqrt{\frac{\hbar G}{c^3}} \approx 10^{-35} \text{ m} \quad \text{Planck length}$$

Common assumption: corrections to GR at the **ultraviolet** scale

$$l_P = \sqrt{\frac{\hbar G}{c^3}} \approx 10^{-35} \text{ m} \quad \text{Planck length}$$

Strong indications of the relevance of quantum gravity effects at **infrared** scales

Common assumption: corrections to GR at the **ultraviolet** scale

$$l_P = \sqrt{\frac{\hbar G}{c^3}} \approx 1.6 \times 10^{-35} \text{ m} \quad \text{Planck length}$$

Strong indications of the relevance of quantum gravity effects at **infrared** scales

- Nontrivial quantum structure of the horizon: Fuzzball^[Mathur (2005)], Firewall^[Almheiri et al. (2013)]

Common assumption: corrections to GR at the **ultraviolet** scale

$$\ell_P = \sqrt{\frac{\hbar G}{c^3}} \approx 10^{-35} \text{ m} \quad \text{Planck length}$$

Strong indications of the relevance of quantum gravity effects at **infrared** scales

- Nontrivial quantum structure of the horizon: Fuzzball^[Mathur (2005)], Firewall^[Almheiri et al. (2013)]
- Nonlocal modifications of QFT at horizon scale^[Giddings (2013, 2021)]

Common assumption: corrections to GR at the **ultraviolet** scale

$$l_P = \sqrt{\frac{\hbar G}{c^3}} \approx 10^{-35} \text{ m} \quad \text{Planck length}$$

Strong indications of the relevance of quantum gravity effects at **infrared** scales

- Nontrivial quantum structure of the horizon: Fuzzball^[Mathur (2005)], Firewall^[Almheiri et al. (2013)]
- Nonlocal modifications of QFT at horizon scale^[Giddings (2013, 2021)]
- QNMs spectrum
Black hole = quantum harmonic oscillator with R_S ^[Maggiore (2008)]

Common assumption: corrections to GR at the **ultraviolet** scale

$$l_P = \sqrt{\frac{\hbar G}{c^3}} \approx 10^{-35} \text{ m} \quad \text{Planck length}$$

Strong indications of the relevance of quantum gravity effects at **infrared** scales

- Nontrivial quantum structure of the horizon: Fuzzball^[Mathur (2005)], Firewall^[Almheiri et al. (2013)]
- Nonlocal modifications of QFT at horizon scale^[Giddings (2013, 2021)]
- QNMs spectrum
Black hole = quantum harmonic oscillator with R_S ^[Maggiore (2008)]
- Corpuscular gravity^[Dvali & Gomez (2013), Dvali et al. (2021, 2022)]

A lesson from cosmology: the dark matter problem

Standard Cosmological Model: GR + energy content made of baryonic matter (5%), dark matter (30%) and dark energy (65%)

A lesson from cosmology: the dark matter problem

Standard Cosmological Model: GR + energy content made of baryonic matter (5%), dark matter (30%) and dark energy (65%)

[Verlinde (2016)]: dark matter effects are a manifestation of the competition between baryonic matter and dark energy degrees of freedom

A lesson from cosmology: the dark matter problem

Standard Cosmological Model: GR + energy content made of baryonic matter (5%), dark matter (30%) and dark energy (65%)

[Verlinde (2016)]: dark matter effects are a manifestation of the competition between baryonic matter and dark energy degrees of freedom

Schwarzschild-de Sitter solution

$$f(r) = 1 - \frac{R_S}{r} - \frac{r^2}{L^2}$$

R_S : baryonic-matter scale (**inner horizon**)

L : cosmological scale (**outer horizon**)

$R_0 = \sqrt{\frac{\rho}{R_S L}}$: dark-matter scale

A lesson from cosmology: the dark matter problem

Standard Cosmological Model: GR + energy content made of baryonic matter (5%), dark matter (30%) and dark energy (65%)

[Verlinde (2016)]: dark matter effects are a manifestation of the competition between baryonic matter and dark energy degrees of freedom

Schwarzschild-de Sitter solution

$$f(r) = 1 - \frac{R_S}{r} - \frac{r^2}{L^2}$$

R_S : baryonic-matter scale (**inner horizon**)

L : cosmological scale (**outer horizon**)

$R_0 = \sqrt{R_S L}$: dark-matter scale

Effective model: GR sourced by an *anisotropic* fluid encoding *long-range quantum gravity* effects [Cadoni et al. (2018-2021)]

Building effective nonsingular black-hole models

Main idea: use the same (*but reversed!*) principle adopted for the cosmological solution to address the singularity problem

Building effective nonsingular black-hole models

Main idea: use the same (*but reversed!*) principle adopted for the cosmological solution to address the singularity problem

Framework: GR + Anisotropic Fluid

$$G = 8\pi G T ; \quad T = (\rho + p_r) u u + p_r g + p_k p_r w w$$

Building effective nonsingular black-hole models

Main idea: use the same (*but reversed!*) principle adopted for the cosmological solution to address the singularity problem

Framework: GR + Anisotropic Fluid

$$G = 8\pi G T ; \quad T = (\rho + p_r) u u + p_r g + p_k w w$$

Spherical symmetry

$$ds^2 = e^{2\alpha(r)} dt^2 + e^{2\beta(r)} dr^2 + r^2 d\Omega^2$$

Building effective nonsingular black-hole models

Main idea: use the same (*but reversed!*) principle adopted for the cosmological solution to address the singularity problem

Framework: GR + Anisotropic Fluid

$$G = 8\pi G T ; \quad T = (\rho + p_r) u u + p_r g + p_k w w$$

Spherical symmetry $ds^2 = e^{2\alpha(r)} dt^2 + e^{2\beta(r)} dr^2 + r^2 d\Omega^2$

Equation of State $p_k =$

Building effective nonsingular black-hole models

Main idea: use the same (*but reversed!*) principle adopted for the cosmological solution to address the singularity problem

Framework: GR + Anisotropic Fluid

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} ; \quad T_{\mu\nu} = (\rho + p_r) u_\mu u_\nu + p_r g_{\mu\nu} + p_k w_\mu w_\nu$$

Spherical symmetry $ds^2 = e^{2\alpha(r)} dt^2 + e^{2\beta(r)} dr^2 + r^2 d\Omega^2$

Equation of State $p_k =$

- A very simple choice

Building effective nonsingular black-hole models

Main idea: use the same (*but reversed!*) principle adopted for the cosmological solution to address the singularity problem

Framework: GR + Anisotropic Fluid

$$G = 8\pi G T ; \quad T = (\rho + p_r) u u + p_r g + p_k p_r w w$$

Spherical symmetry $ds^2 = e^{2\alpha(r)} dt^2 + e^{2\beta(r)} dr^2 + r^2 d\Omega^2$

Equation of State $p_k =$

- A very simple choice
- Natural connection with dark energy/cosmological constant

Building effective nonsingular black-hole models

Main idea: use the same (*but reversed!*) principle adopted for the cosmological solution to address the singularity problem

Framework: GR + Anisotropic Fluid

$$G = 8\pi G T ; \quad T = (\rho + p_r) u u + p_r g + p_k p_r w w$$

Spherical symmetry $ds^2 = e^{2\alpha(r)} dt^2 + e^{2\beta(r)} dr^2 + r^2 d\Omega^2$

Equation of State $p_k =$

- A very simple choice
- Natural connection with dark energy/cosmological constant
- Allows to have a stress-energy tensor invariant under radial Lorentz boosts and define a spherically symmetric vacuum [Dymnikova (1992)]

Building effective nonsingular black-hole models

Main idea: use the same (*but reversed!*) principle adopted for the cosmological solution to address the singularity problem

Framework: GR + Anisotropic Fluid

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} ; \quad T_{\mu\nu} = (\rho + p_r) u_\mu u_\nu + p_r g_{\mu\nu} + p_k \delta_{\mu\nu} + p_w w_\mu w_\nu$$

Spherical symmetry $ds^2 = e^{2\alpha(r)} dt^2 + e^{2\beta(r)} dr^2 + r^2 d\Omega^2$

Equation of State $p_k =$

- A very simple choice
- Natural connection with dark energy/cosmological constant
- Allows to have a stress-energy tensor invariant under radial Lorentz boosts and define a spherically symmetric vacuum [\[Dymnikova \(1992\)\]](#)

Solution:

$$e^{2\alpha(r)} = e^{2\beta(r)} \quad A(r) = 1 - \frac{2GM(r)}{r} ; \quad M(r) = 4\pi \int_0^r \rho(r') r'^2 dr'$$

Imposing a de Sitter core

To eliminate the singularity: glue together a Schwarzschild patch at great distances from the center with a de Sitter patch near $r = 0$

$$A(r) \underset{r \rightarrow 0}{\sim} 1 - \frac{r^2}{\ell^2} \qquad A(r) \underset{r \rightarrow \infty}{\sim} 1 - \frac{R_s}{r}$$

Imposing a de Sitter core

To eliminate the singularity: glue together a Schwarzschild patch at great distances from the center with a de Sitter patch near $r = 0$

$$A(r) = 1 - \frac{r_s^2}{r^2} \qquad A(r) = 1 - \frac{R_s}{r}$$

The position of the horizons is reversed with respect to the SdS case

To eliminate the singularity: glue together a Schwarzschild patch at great distances from the center with a de Sitter patch near $r = 0$

$$A(r) \underset{r \rightarrow 0}{\sim} 1 - \frac{r^2}{\hat{L}^2} \qquad A(r) \underset{r \rightarrow \infty}{\sim} 1 - \frac{R_S}{r}$$

The position of the horizons is reversed with respect to the SdS case

R_S : baryonic-matter scale (outer horizon)

\hat{L} : "de Sitter" length (inner horizon)

Imposing a de Sitter core

To eliminate the singularity: glue together a Schwarzschild patch at great distances from the center with a de Sitter patch near $r=0$

$$A(r) \underset{r \rightarrow 0}{\sim} 1 - \frac{r^2}{\hat{\Lambda}^2} \qquad A(r) \underset{r \rightarrow \infty}{\sim} 1 - \frac{R_S}{r}$$

The position of the horizons is reversed with respect to the SdS case

R_S : baryonic-matter scale (outer horizon)

$\hat{\Lambda}$: "de Sitter" length (inner horizon)

$$\hat{\Lambda} = R_S^{1/3} \Lambda^{-2/3}$$

To eliminate the singularity: glue together a Schwarzschild patch at great distances from the center with a de Sitter patch near $r=0$

$$A(r) \underset{r \ll 0}{\approx} 1 - \frac{r^2}{\hat{\Lambda}^2} \qquad A(r) \underset{r \gg 1}{\approx} 1 - \frac{R_S}{r}$$

The position of the horizons is reversed with respect to the SdS case

R_S : baryonic-matter scale (outer horizon)

$\hat{\Lambda}$: "de Sitter" length (inner horizon)

$$\hat{\Lambda} = R_S^{1/3} \hat{\Lambda}^{2/3}$$

$\hat{\Lambda}$ is an additional quantum hair

Several models in the literature

- ^ Bardeen model [Bardeen (1968)]
- ^ Dymnikova model [Dymnikova (1992)]
- ^ Hayward model [Hayward (2005)] [First proposed by Poisson & Israel (1988)]
- ^ Gaussian-core model [Nicolini et al. (2006) ; Modesto et al. (2011)]
- ^ ...
- ^ Model with the strongest deviations from Schwarzschild at
in nity [Lan & Wang (2022) ; Cadoni, De Laurentis, De Martino, Della Monica, Oi & APS (2022)]

Several models in the literature

- ^ Bardeen model [Bardeen (1968)]
- ^ Dymnikova model [Dymnikova (1992)]
- ^ Hayward model [Hayward (2005)] [First proposed by Poisson & Israel (1988)]
- ^ Gaussian-core model [Nicolini et al. (2006) ; Modesto et al. (2011)]
- ^ ...
- ^ Model with the strongest deviations from Schwarzschild at
in nity [Lan & Wang (2022) ; Cadoni, De Laurentis, De Martino, Della Monica, Oi & APS (2022)]

In some models, the hair is an UV QG parameter (noncommutative geometry [Nicolini et al. (2006)] , LQG effects [Modesto et al. (2011)] , others [Rovelli et al. (2014) ; Frolov (2014)])

Thus considered to be of the order of ℓ_P

Several models in the literature

- ^ Bardeen model [Bardeen (1968)]
- ^ Dymnikova model [Dymnikova (1992)]
- ^ Hayward model [Hayward (2005)] [First proposed by Poisson & Israel (1988)]
- ^ Gaussian-core model [Nicolini et al. (2006) ; Modesto et al. (2011)]
- ^ ...
- ^ Model with the strongest deviations from Schwarzschild at
in nity [Lan & Wang (2022) ; Cadoni, De Laurentis, De Martino, Della Monica, Oi & APS (2022)]

In some models, the hair is an UV QG parameter (noncommutative geometry [Nicolini et al. (2006)] , LQG effects [Modesto et al. (2011)] , others [Rovelli et al. (2014) ; Frolov (2014)])

Thus considered to be of the order of ℓ_P

What's new here? Wider effective perspective where ℓ can take any value

Several models in the literature

- ^ Bardeen model [\[Bardeen \(1968\)\]](#)
- ^ Dymnikova model [\[Dymnikova \(1992\)\]](#)
- ^ Hayward model [\[Hayward \(2005\)\]](#) [First proposed by Poisson & Israel (1988)]
- ^ Gaussian-core model [\[Nicolini et al. \(2006\) ; Modesto et al. \(2011\)\]](#)
- ^ ...
- ^ Model with the strongest deviations from Schwarzschild at
in nity [\[Lan & Wang \(2022\) ; Cadoni, De Laurentis, De Martino, Della Monica, Oi & APS \(2022\)\]](#)

In some models, the hair is an UV QG parameter (noncommutative geometry [\[Nicolini et al. \(2006\)\]](#) , LQG effects [\[Modesto et al. \(2011\)\]](#) , others [\[Rovelli et al. \(2014\) ; Frolov \(2014\)\]](#))

Thus considered to be of the order of ρ

What's new here? Wider effective perspective where ρ can take any value

$$A(r) = 1 - \frac{R_S}{r} F\left(\frac{r}{\rho}\right)$$

$$\begin{cases} < 8 & F(y) = y^2 & \text{for } y < 0 \\ : & F(y) = \frac{1}{y} & \text{for } y \geq 1 \end{cases}$$

$\hat{\rho} = R_S(\hat{\rho}, \hat{\rho}_P)$ Small short-scale QG corrections

$\hat{\rho}(R_S(\ell_P))$ Small short-scale QG corrections

$\hat{\rho}(R_S)$ Horizonless compact object

- ^ ` $R_S (\ell \ell_P)$) Small short-scale QG corrections
- ^ ` R_S) Horizonless compact object
- ^ ` R_S) QG effects at the horizon scale

- ^ ` R_S (` ` ρ)) Small short-scale QG corrections
- ^ ` R_S) Horizonless compact object
- ^ ` R_S) QG effects at the horizon scale

Extremal model ` c R_S

Second-order phase transition separates two thermodynamic branches

Second-order phase transition separates two thermodynamic branches

- Branch II (\dot{R}_S): **unstable** evaporating models (Hawking branch)

Second-order phase transition separates two thermodynamic branches

- Branch II (\dot{R}_S): **unstable** evaporating models (Hawking branch)
- Branch I (\dot{R}_S): **stable remnants**

Second-order phase transition separates two thermodynamic branches

- Branch II (\dot{R}_S): **unstable** evaporating models (Hawking branch)
- Branch I (\dot{R}_S): **stable remnants**

For other interesting thermodynamic features, see
Cadoni, Oi, **APS**, arXiv:2204.09444

Observable phenomenology

- Light ring
- QNMs spectrum
- Motion of the S2-star around SgrA

- Light ring
- QNMs spectrum
- Motion of the S2-star around SgrA

Example: model with the strongest deviations from Schwarzschild

[Cadoni, De Laurentis, De Martino, Della Monica, Oi, **APS**, arXiv:2211.11585]

$$A(r) = 1 - \frac{2GM}{r} + \frac{2GMr^2}{(r + r_g)^3} - \frac{6GM}{r^2}$$

- Light ring
- QNMs spectrum
- Motion of the S2-star around SgrA

Example: model with the strongest deviations from Schwarzschild

[Cadoni, De Laurentis, De Martino, Della Monica, Oi, **APS**, arXiv:2211.11585]

$$A(r) = 1 - \frac{2GM}{r} + \frac{2GMr^2}{(r + r_0)^3} + 6GM \frac{r_0}{r^2}$$

With $M = M_{\text{SgrA}}$ at $r = r_{\text{S2}}$

Observable phenomenology

- Light ring
- QNMs spectrum
- Motion of the S2-star around SgrA

Example: model with the strongest deviations from Schwarzschild

[Cadoni, De Laurentis, De Martino, Della Monica, Oi, **APS**, arXiv:2211.11585]

$$A(r) = 1 - \frac{2GM}{r} + \frac{2GM}{r} \left(\frac{r}{r_S} \right)^{-3} + 6GM \frac{r}{r^2}$$

With $M = M_{\text{SgrA}}$ at $r = r_{\text{S2}}$

$$\left(\frac{r}{r_S} \right)^{-3} \approx \frac{6GM}{r^2} \approx 10^{-53}$$

$$\left(\frac{r}{r_S} \right)^{-3} \approx \frac{6GM}{r^2} \approx 10^{-9}$$

(e.g. $\frac{r}{r_S} \approx 0.296 GM$)

Observable phenomenology

- Light ring
- QNMs spectrum
- Motion of the S2-star around SgrA

Example: model with the strongest deviations from Schwarzschild

[Cadoni, De Laurentis, De Martino, Della Monica, Oi, **APS**, arXiv:2211.11585]

$$A(r) = 1 - \frac{2GM}{r} + \frac{2GM}{r} \left(\frac{r}{r_S} \right)^{-3} + 6GM \frac{r}{r^2}$$

With $M = M_{\text{SgrA}}$ at $r = r_{\text{S2}}$

$$\left(\frac{r}{r_S} \right)^{-3} \approx \frac{6GM}{r^2} \approx 10^{53}$$

$$0.2GM \approx \frac{6GM}{r^2} \approx 10^9$$

(e.g. $r_c \approx 0.296 GM$)

Observable phenomenology

- Light ring
- QNMs spectrum
- Motion of the S2-star around SgrA

Example: model with the strongest deviations from Schwarzschild

[Cadoni, De Laurentis, De Martino, Della Monica, Oi, **APS**, arXiv:2211.11585]

$$A(r) = 1 - \frac{2GM}{r} + \frac{2GM}{r} \left(\frac{r}{r_c} \right)^{-3} + 6GM \frac{r}{r_c^2}$$

With $M = M_{\text{SgrA}}$ at $r = r_{\text{S2}}$

$$\left(\frac{r_{\text{P}}}{r_c} \right)^{-3} \frac{6GM}{r^2} \approx 10^{-53}$$

$$\left(\frac{0.2GM}{r_c} \right)^{-3} \frac{6GM}{r^2} \approx 10^{-9}$$

(e.g. $r_c = 0.296 GM$)

S2 Observational data

$\approx 0.47 GM$ with 95 % C. L.

[arXiv:2211.11585]

- The possibility of having IR QG corrections is an alternative interesting way to resolve the singularity problem
- Deviations from standard phenomenology are hopeless to be detected in the near future if $\Lambda \ll R_S$

They are not if $\Lambda \gg R_S$

- The possibility of having IR QG corrections is an alternative interesting way to resolve the singularity problem
- Deviations from standard phenomenology are hopeless to be detected in the near future if $\Lambda \ll R_S$

They are not if $\Lambda \sim R_S$

THANKS FOR YOUR ATTENTION!