



From Raw data to Physics Results (3/3)

Paul Laycock

July 5th 2022



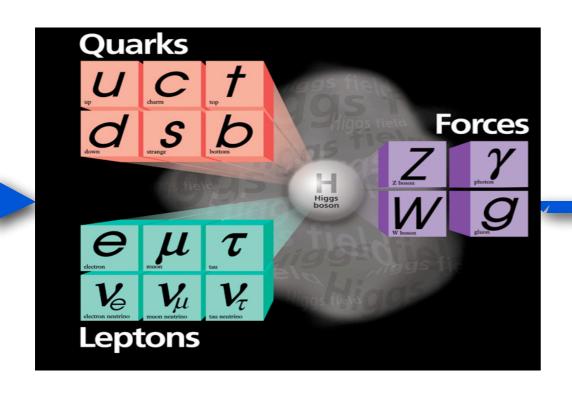


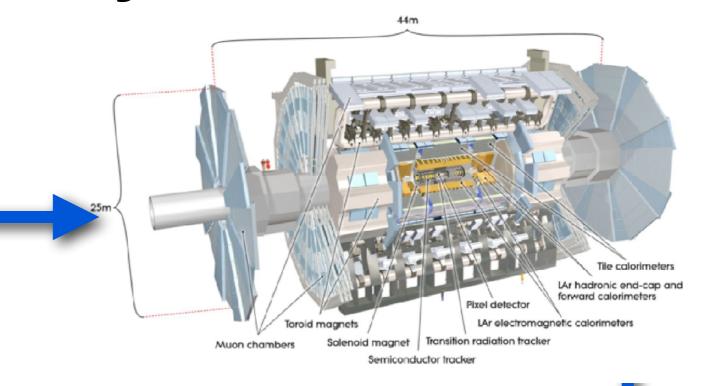


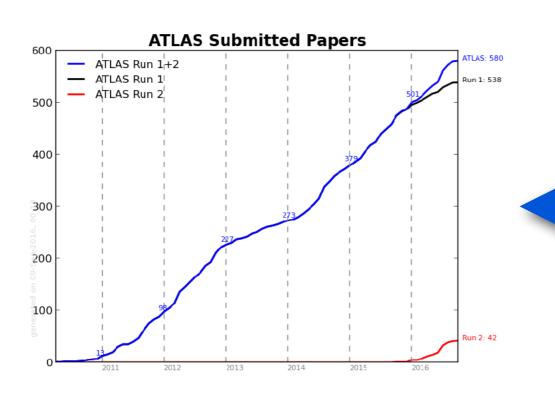


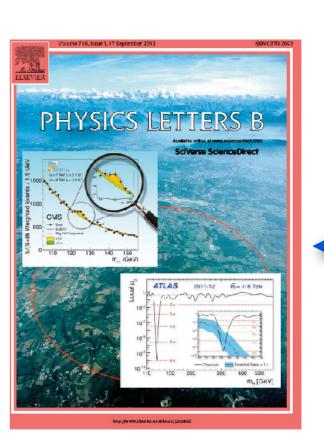
@BrookhavenLab

The particle physics cycle









Course outline

Lecture 1

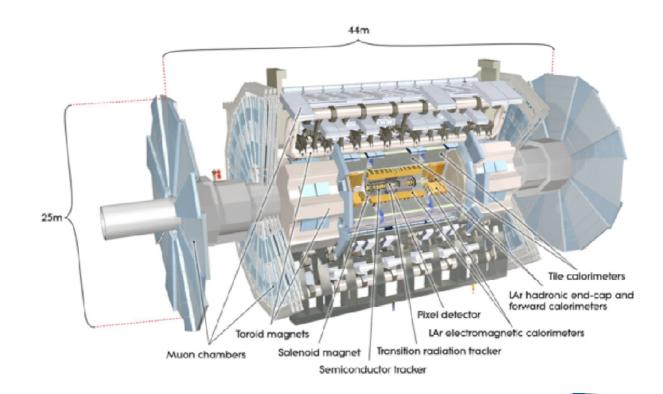
The journey of raw data from the detector to a publication

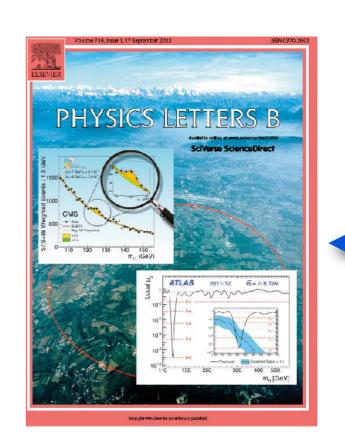
Lecture 2

 How we reconstruct fundamental physics processes from raw detector data

Lecture 3

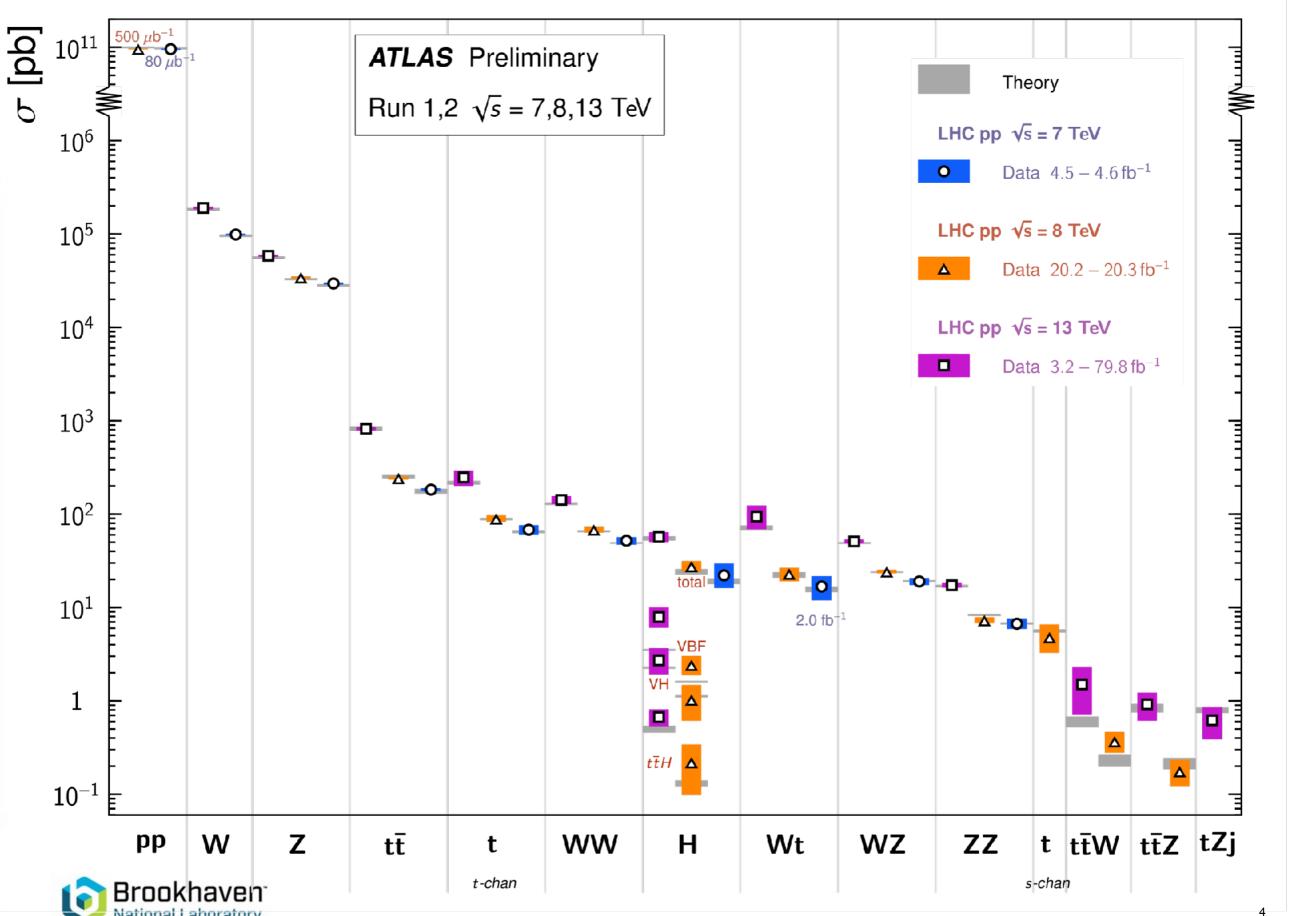
 How we extract our signals from the mountain of data, finding needles in the haystack







Standard Model Total Production Cross Section Measurements Status: July 2018



Measuring cross sections

$$\sigma = \frac{N}{L}$$

The cross section for a process is defined as the number of events divided by luminosity



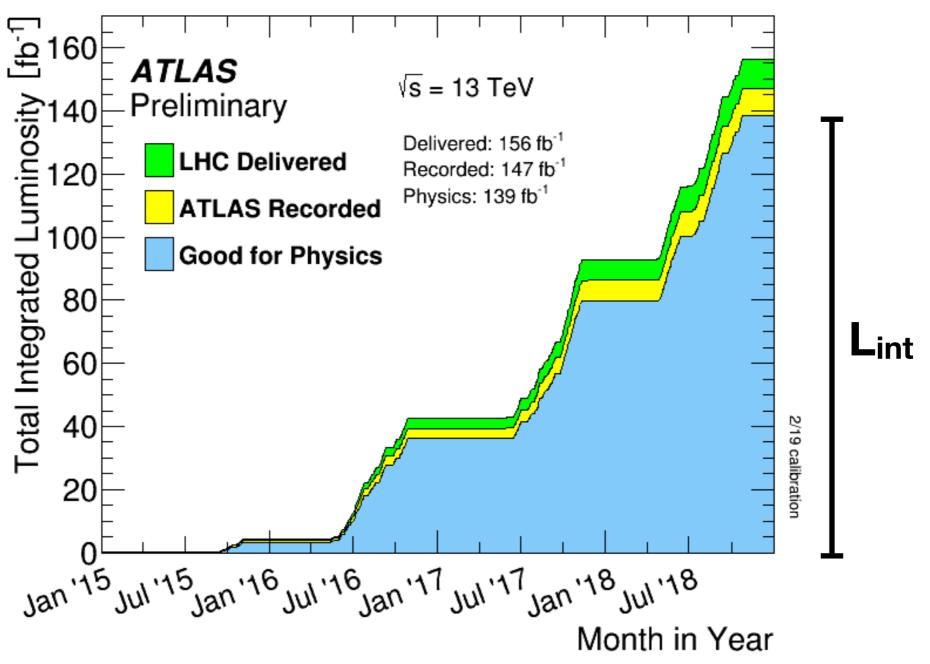
Measuring cross sections

$$\sigma = \frac{N}{L_{int}}$$

The cross section for a process is defined as the number of events divided by the integrated luminosity, *L*_{int}, which measures how much data we have collected



ATLAS Luminosity

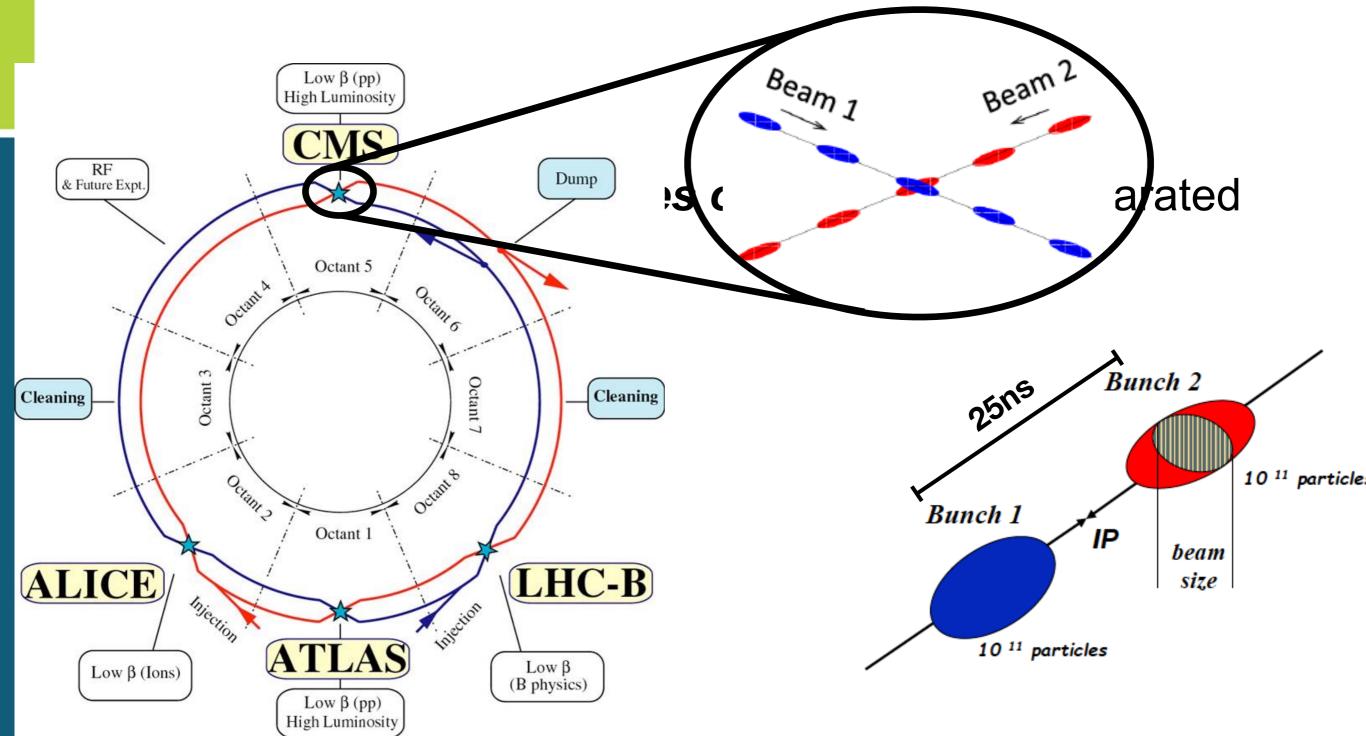


Question: Why does ATLAS record less data than the LHC delivers? How do we know the integrated luminosity delivered?



LHC collisions

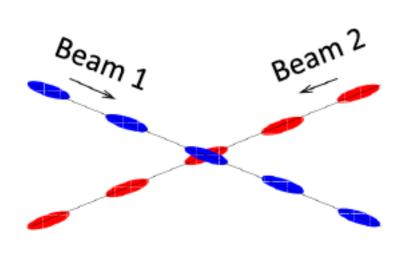
Figures adapted from Michaela Schaumann's <u>third lecture</u> (11/07/19) on "Particle Accelerators and Beam Dynamics"

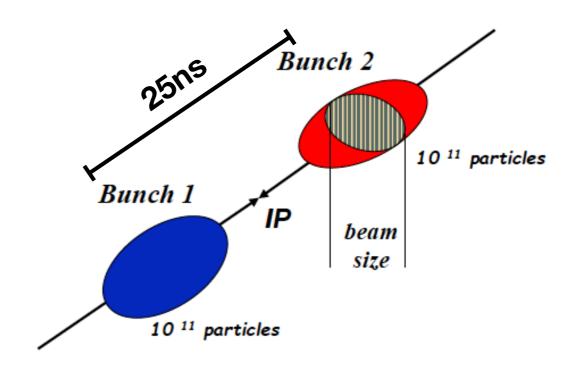


• The LHC accelerates *bunches of 10¹¹ protons* separated by 25ns gaps



Measuring Luminosity at the LHC



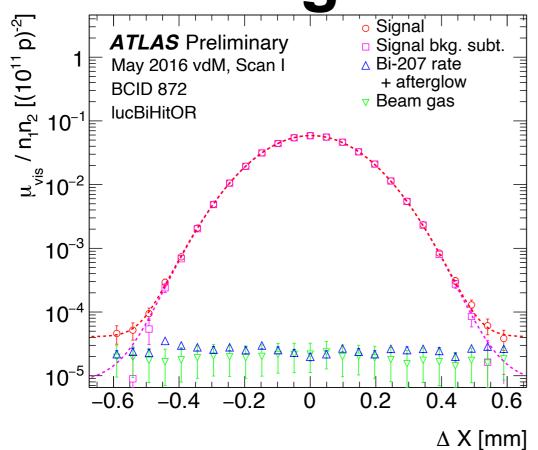


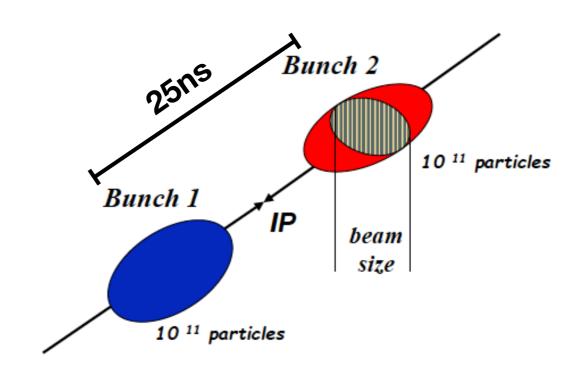
Ingredients for a measurement of the luminosity

- Measuring the size of the beams (for a certain LHC configuration)
 - This requires a dedicated measurement where we scan the beams across each other in the horizontal and vertical directions - a van der Meer scan
- Measuring the beam currents in each bunch
 - This is done during collisions, integrating all of the bunch currents and knowing their size, we can calculate the luminosity
- Make many cross checks because this is such a crucial measurement



Measuring Luminosity at the LHC





Ingredients for a measurement of the luminosity

- Measuring the size of the beams (for a certain LHC configuration)
 - This requires a dedicated measurement where we scan the beams across each other in the horizontal and vertical directions - a van der Meer scan
- Measuring the beam currents in each bunch
 - This is done during collisions, integrating all of the bunch currents and knowing their size, we can calculate the luminosity
- Make many cross checks because this is such a crucial measurement



Measuring cross sections

$$\sigma = \frac{N}{L_{int}}$$

The cross section for a process is defined as the number of events divided by the integrated luminosity, *L*_{int}, which measures how much data we have collected

$$\sigma = \frac{N_{obs}}{A \cdot \epsilon \cdot L_{int}}$$

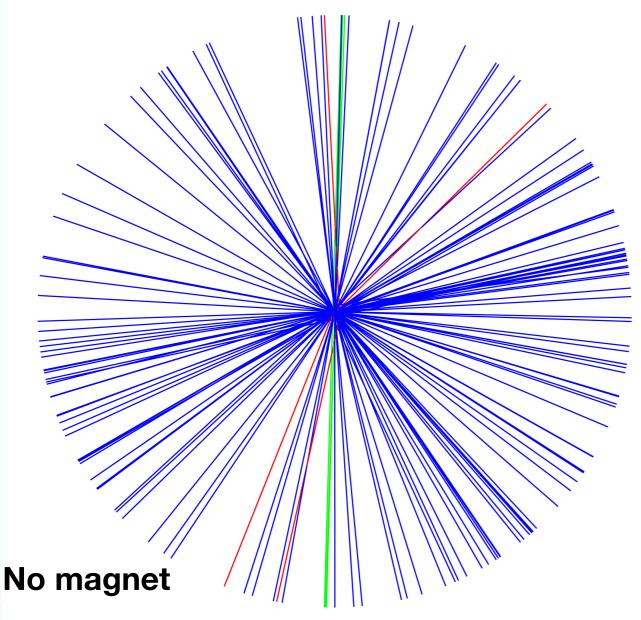
 N_{obs} in data needs to be corrected for the detector acceptance, A, for selecting those events. The reconstruction efficiency, ε , is a product of all of the efficiencies that we need to measure and ensure that they are the same in our data and simulation - how?



Before the detector, came the simulation

When designing detectors, we *simulate detector response* to physics of interest Adding a *solenoid magnet* makes it possible to measure momentum (and charge) in our tracker by measuring curvature in the transverse plane

Interesting physics is often at *high momentum*, e.g. four high momentum muon tracks here

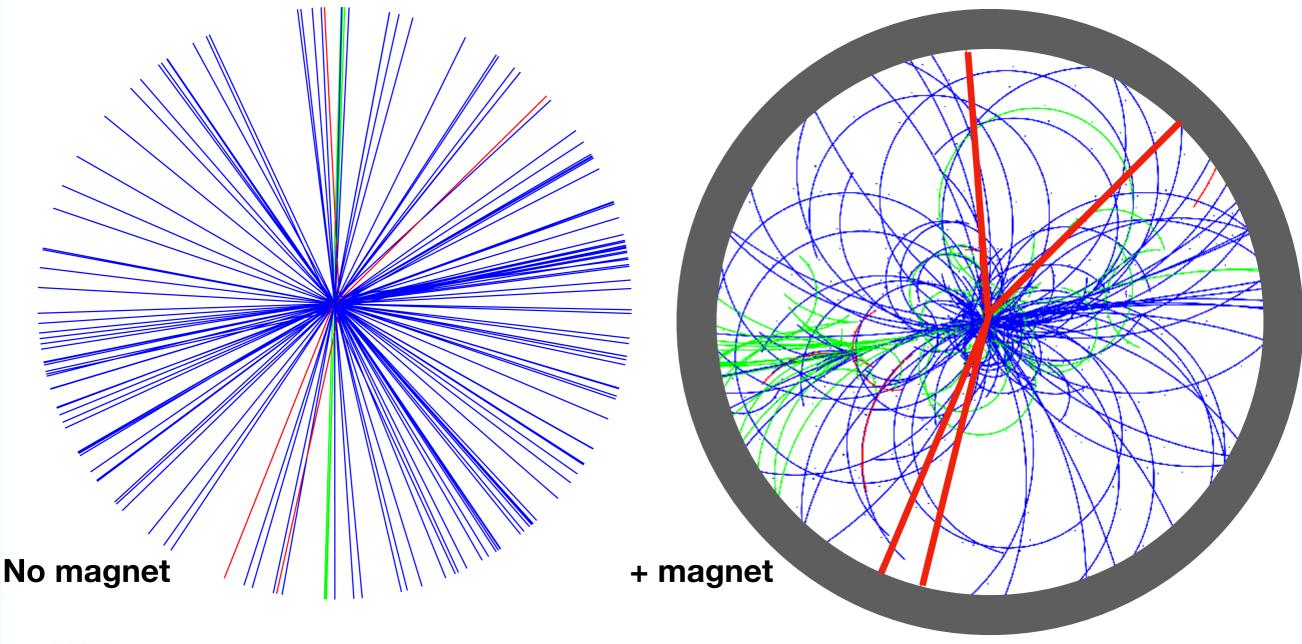




Before the detector, came the simulation

When designing detectors, we *simulate detector response* to physics of interest Adding a *solenoid magnet* makes it possible to measure momentum (and charge) in our tracker by measuring curvature in the transverse plane

Interesting physics is often at *high momentum*, e.g. four high momentum muon tracks here



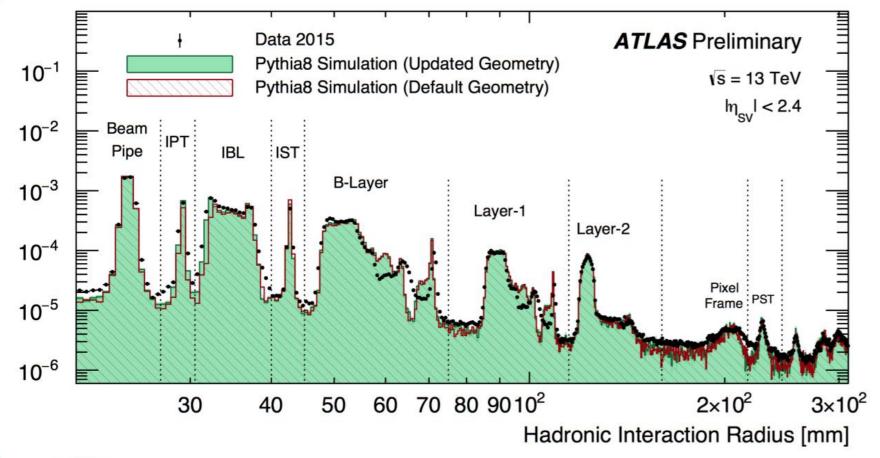


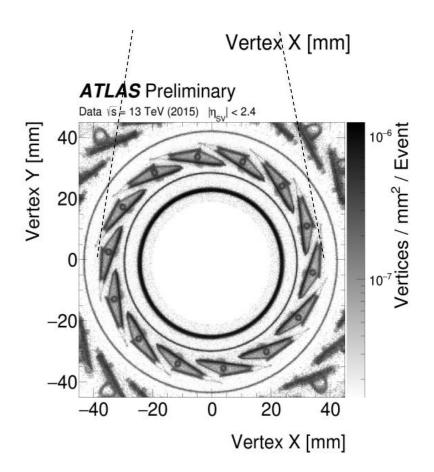
Simulation and understanding detectors

We use **simulations** to model the detector as **accurately** and **precisely** as possible

We then *test* that our simulations are accurate *using real data*We correct our simulations if necessary

Once our simulation is an *accurate model* of our detector, we can use it to *correct the data for detector response*







Real vs perfect tracking detectors

The perfect tracking detector

- is constructed from zero mass material
- has no noise
- is 100% efficiency
- and has infinite resolution

A real tracking detector

- is constructed from real material
 - particles interact with the detector and scatter, altering the particle trajectory
- suffers from noise
 - noise can be confused with particle tracks
- has less than 100% efficiency
 - particles are not always detected and there can even be dead regions
- has finite resolution
 - it may not always be possible to resolve two particle trajectories



Simulation chain

Event Generation

simulate the physics process.

Detector Simulation

simulate the interaction of the particles with the detector material.

→ Course on Detectors

Digitization

Translate interactions with detector into realistic signals.

→ Course on Electronics /Trigger/DAQ

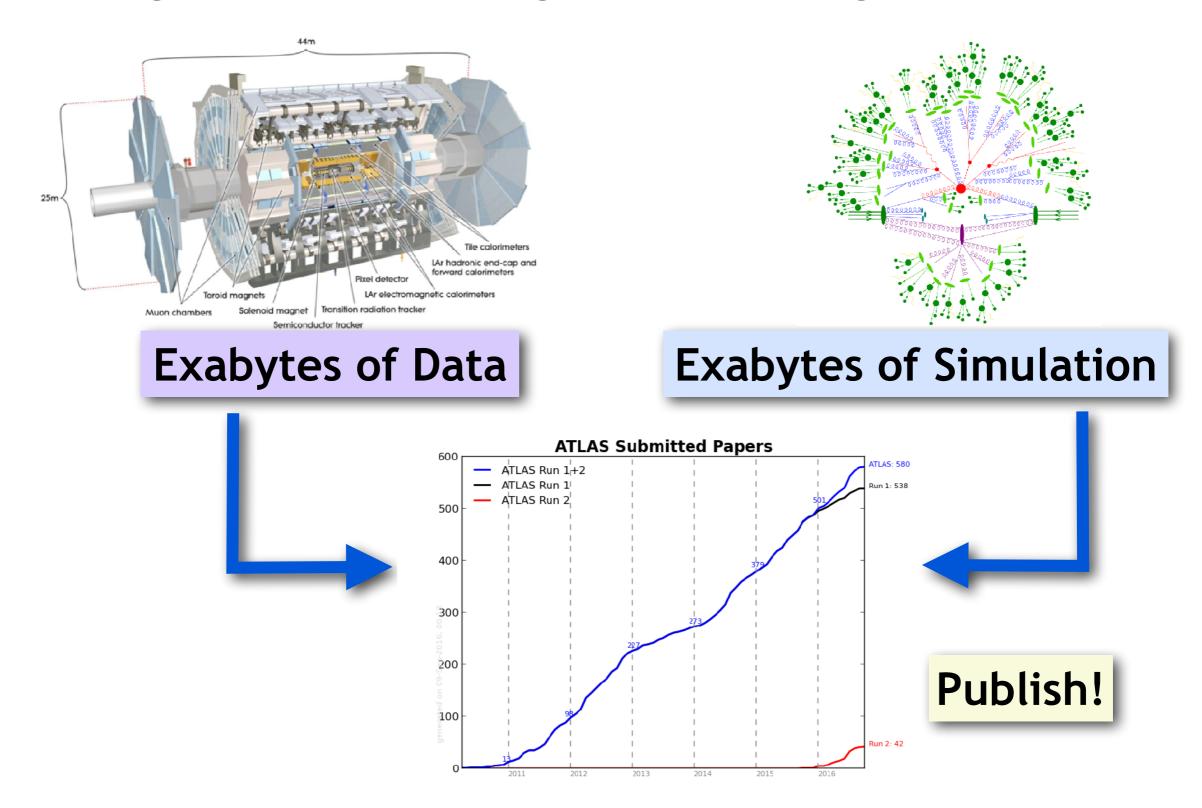
Reconstruction

Go from signals back to particles, as for real data.

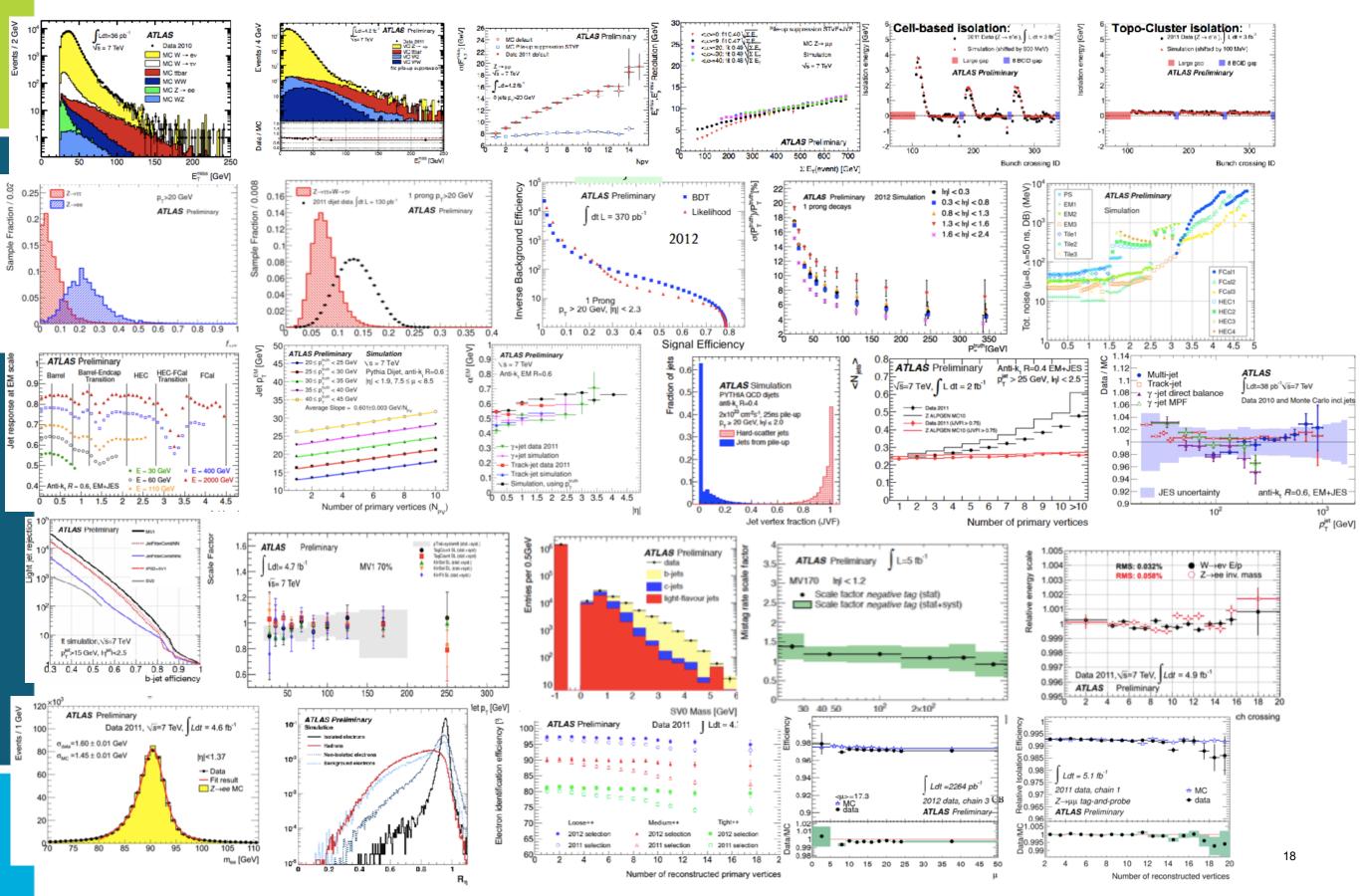
→ This Course



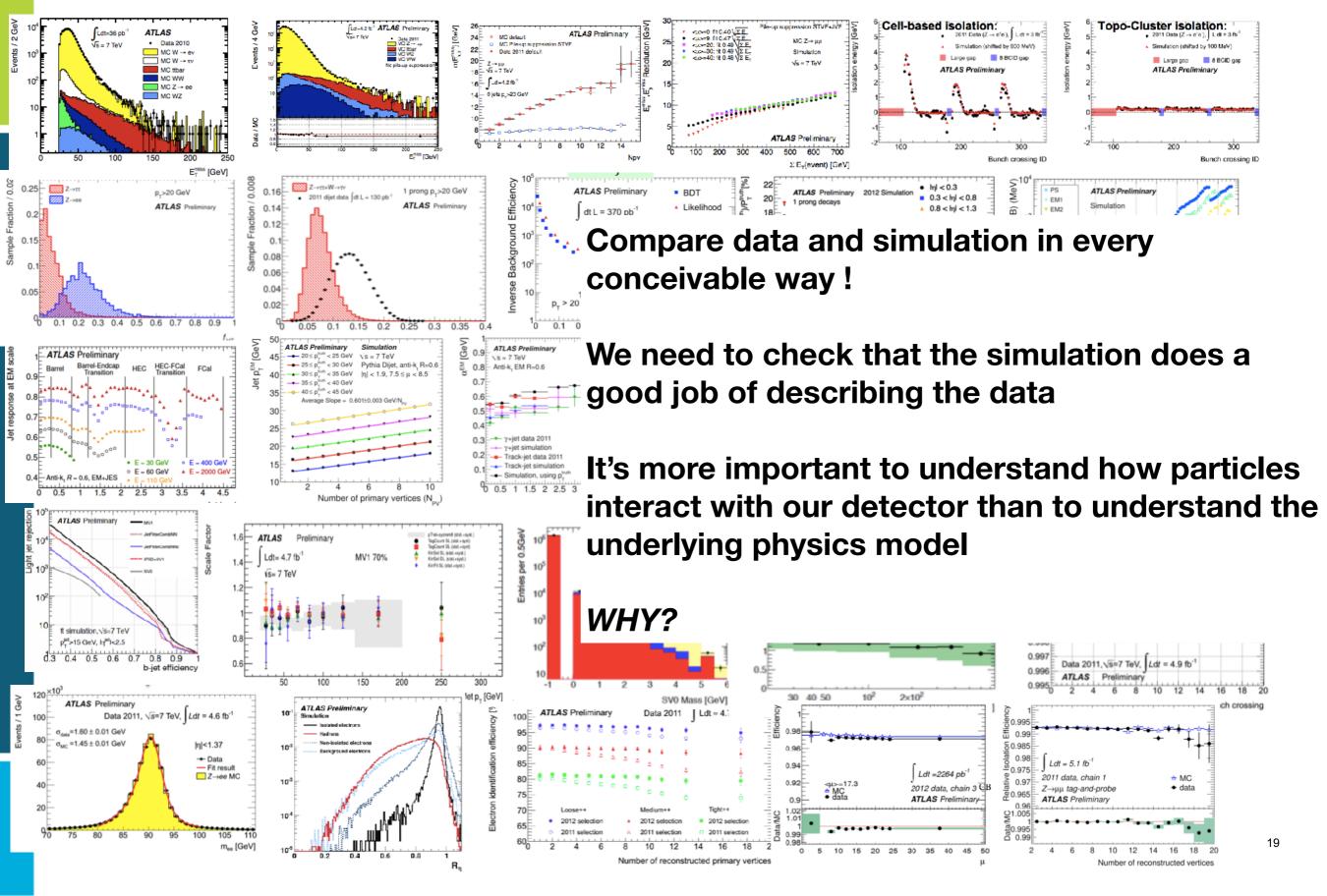
Exabyte-scale physics analysis



Ingredients to the ATLAS physics program



Ingredients to the ATLAS physics program

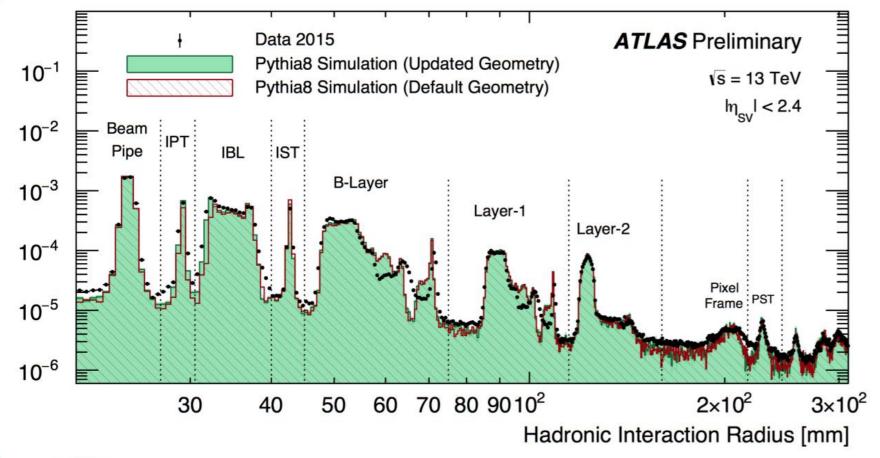


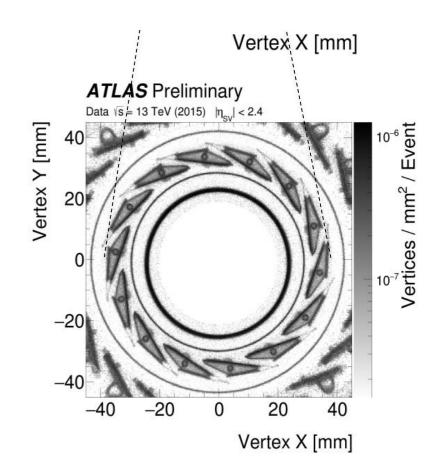
Simulation and understanding detectors

We use **simulations** to model the detector as **accurately** and **precisely** as possible

We then *test* that our simulations are accurate *using real data*We correct our simulations if necessary

Once our simulation is an accurate model of our detector, we can use it to correct the data for detector response







Measuring cross sections

$$\sigma = rac{N}{L_{int}}$$

The cross section for a process is defined as the number of events divided by the integrated luminosity, *L*_{int}, which measures how much data we have collected

$$\sigma = \frac{N_{obs}}{A \cdot \epsilon \cdot L_{int}}$$

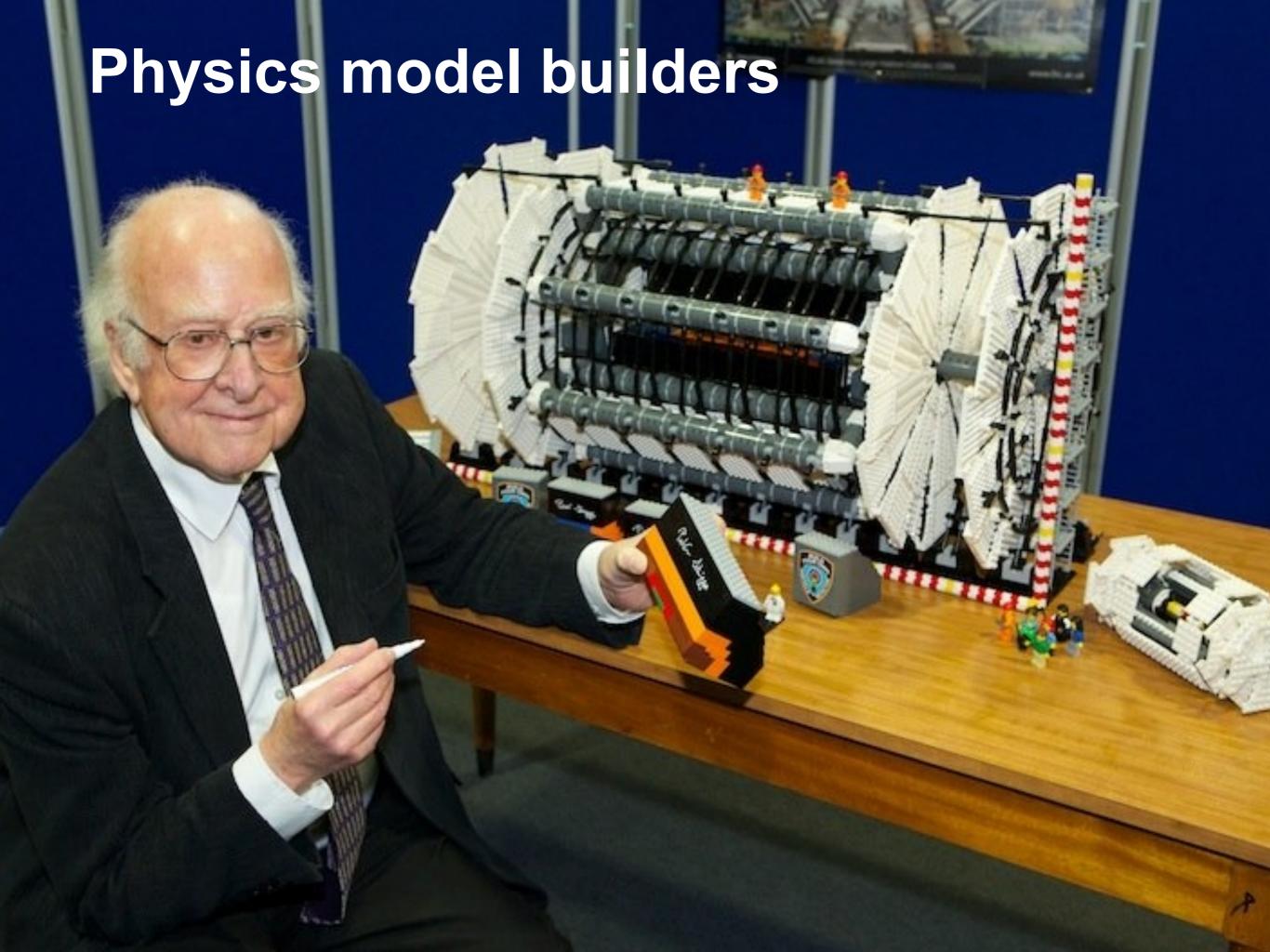
 N_{obs} in data needs to be corrected for the detector acceptance, A, for selecting those events. The reconstruction efficiency, ε , is a product of all of the efficiencies that we need to measure and ensure that they are the same in our data and simulation

$$\sigma = \frac{N_{obs} - N_{bkg}}{A \cdot \epsilon \cdot L_{int}}$$

Finally, in reality we will have some background to our signal and we need to subtract those events that are not part of our signal process



Now we can compare this to the theoretical cross section!



Physics event generators

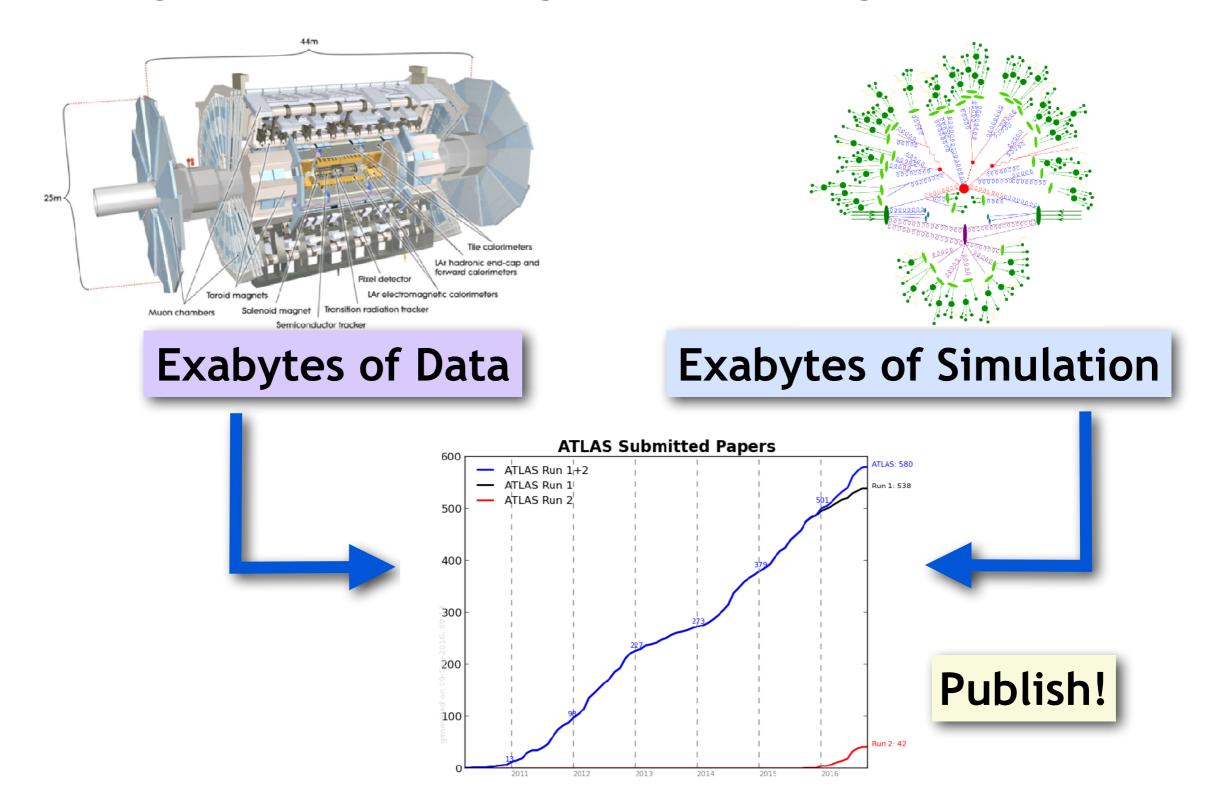


There are lots of different physics models implemented in physics event generators, depending on the type of physics that you're interested in

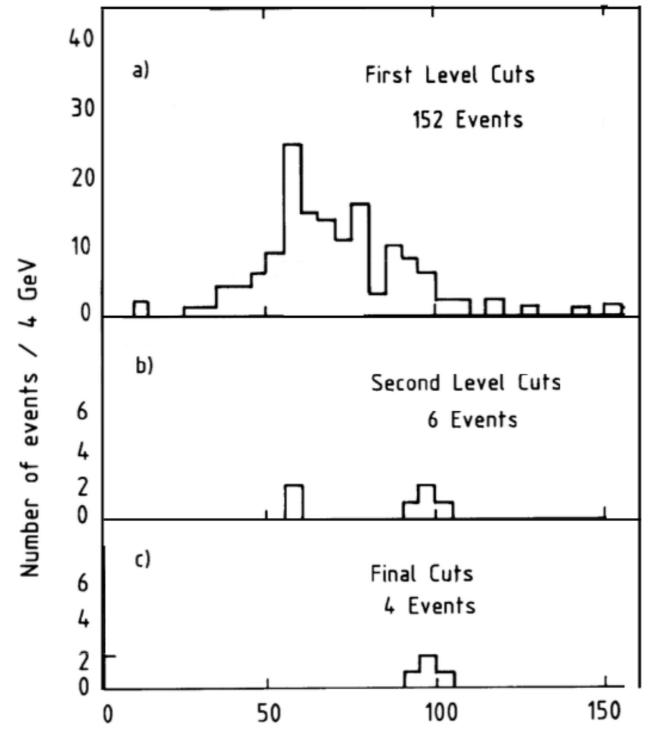
We want to see if reality looks like theory (and which one!)



Exabyte-scale physics analysis



First - measuring the Z boson



Z->ee in UA1

Two EM clusters with E_T >25GeV.

As above plus a track with p_T>7GeV pointing to the cluster. Hadronic and track isolation requirements applied.

A second custer has also an isolated track.

Uncorrected invariant mass cluster pair (GeV/c²)

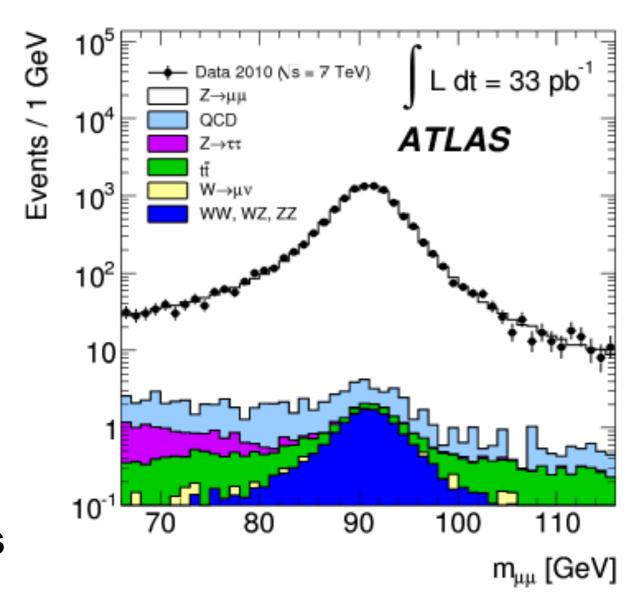


Measuring the Z boson at ATLAS

$$\sigma = \frac{N_{obs} - N_{bkg}}{A \cdot \epsilon \cdot L_{int}}$$

Select events with (here) two muons

Question: what other selections can we apply to the muons?



Here I have only considered events with two muons

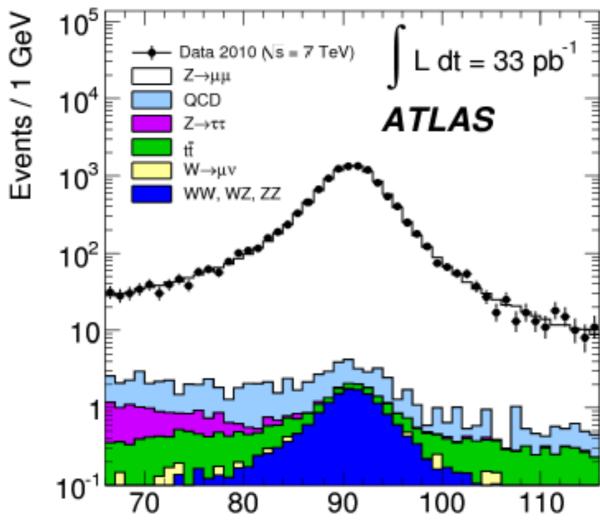
Question: is this the cross section for Z boson production?



Measuring the Z boson at ATLAS

$$\sigma = \frac{N_{obs} - N_{bkg}}{A \cdot \epsilon \cdot L_{int}}$$

Backgrounds are small but still need to be measured and subtracted



We will quote a fiducial cross 100 70 80 section corresponding to good detector acceptance

After making the event selection, applying the same selection to all of the simulations of background processes, and measuring my acceptance and efficiencies (and knowing the luminosity) - am I done?



m_{μμ} [GeV]

Measuring the Z boson at ATLAS

$$\sigma = \frac{N_{obs} - N_{bkg}}{A \cdot \epsilon \cdot L_{int}}$$

Table 5: Measured fiducial $Z \to \ell^+\ell^-$ differential and integrated cross sections for electron and muon channels

		$Z \rightarrow e^+e^-$				$Z \rightarrow \mu^+\mu^-$			
yee min	yee mux	do /d yee	$\delta\sigma_{\rm stat}$	$\delta\sigma_{\rm syst}$	$\delta\sigma_{\text{lumi}}$	do /d yee	$\delta\sigma_{\rm gai}$	$\delta\sigma_{\rm syst}$	$\delta\sigma_{lumi}$
		[pb]	[pb]	[pb]	[pb]	[pb]	[pb]	[pb]	[pb]
0.0	0.5	99.9	2.5	1.6	1.9	105.2	2.4	1.1	2.0
0.5	1.0	100.3	2.7	1.6	1.9	101.9	2.3	1.0	1.9
1.0	1.5	89.2	2.7	1.4	1.7	89.8	2.1	0.8	1.7
1.5	2.0	59.6	2.4	1.2	1.1	61.0	1.8	0.6	1.1
2.0	2.5	19.6	1.3	0.7	0.4	20.3	1.2	0.2	0.4
0.0	2.5	369.0	5.3	4.7	6.9	377.9	4.4	3.4	7.1

No! You would like to publish with the smallest *uncertainties* possible Every ingredient to the analysis comes with an uncertainty

Nobs has a statistical uncertainty

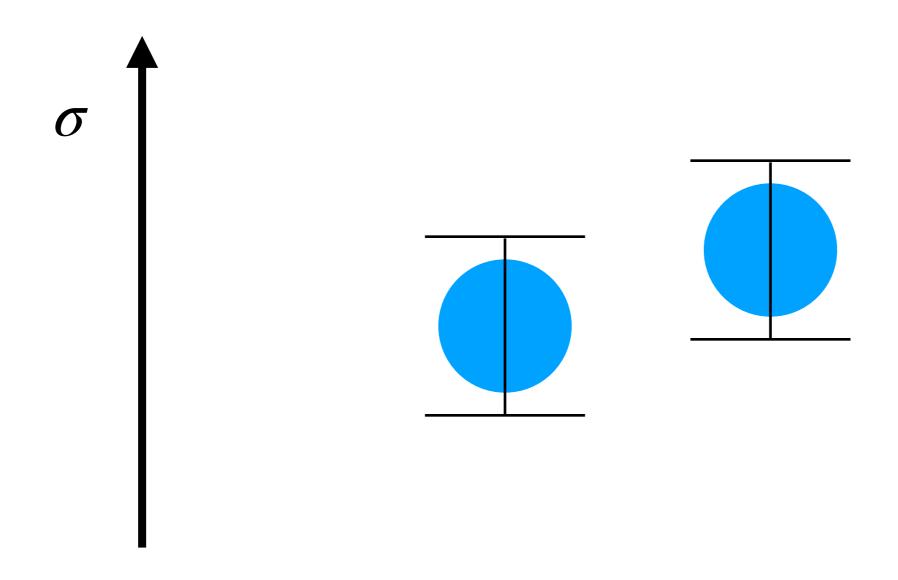
N_{bkg} is typically composed of several sources (different physics processes) with corresponding statistical and systematic contributions to the final uncertainty

A and particularly ε have many systematic components stemming from each reconstruction algorithm that we employed

Finally, *L_{int}* also has an uncertainty that dictates how well we know the absolute scale of the measurement - a *normalisation* uncertainty



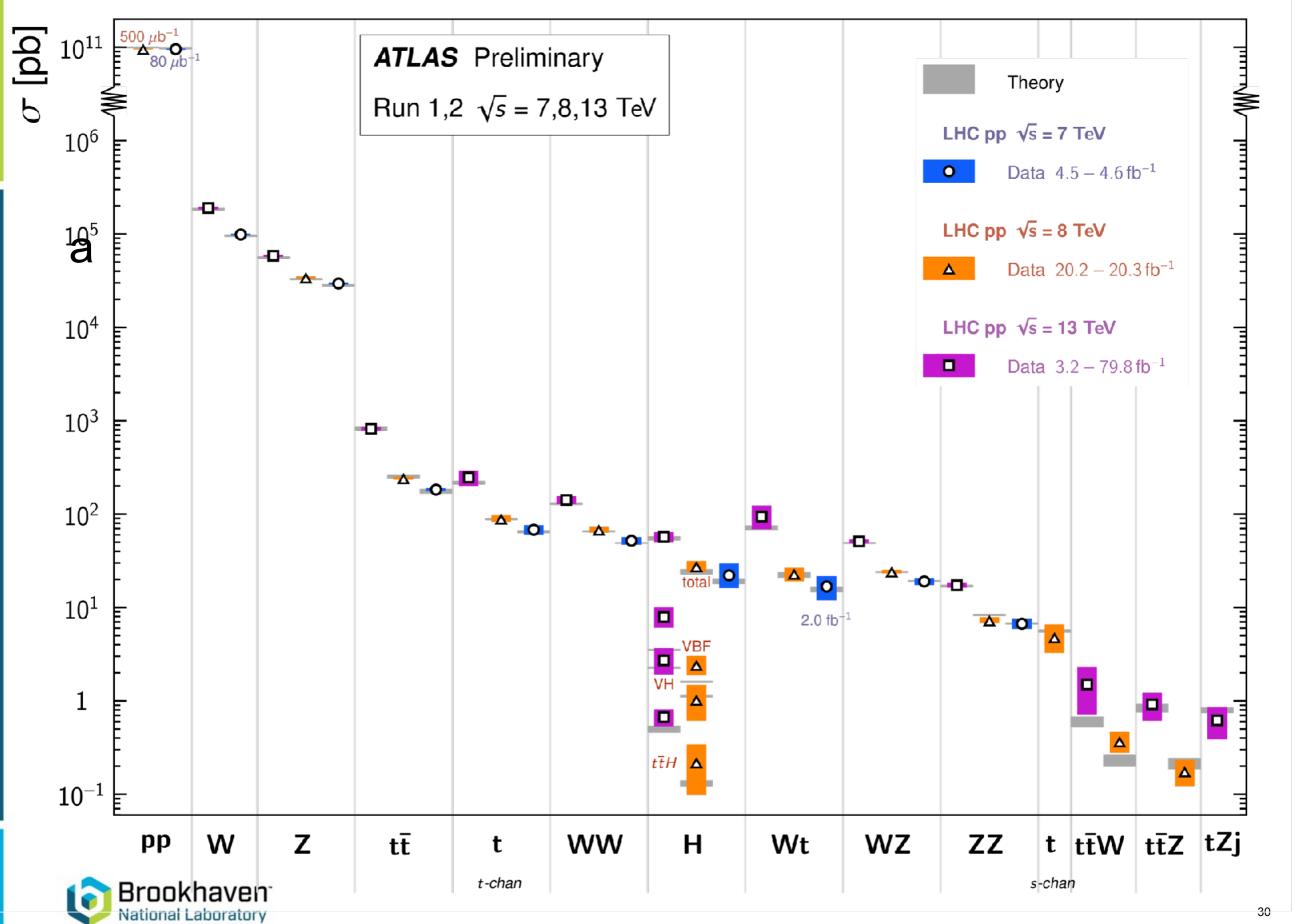
Uncertain about uncertainties

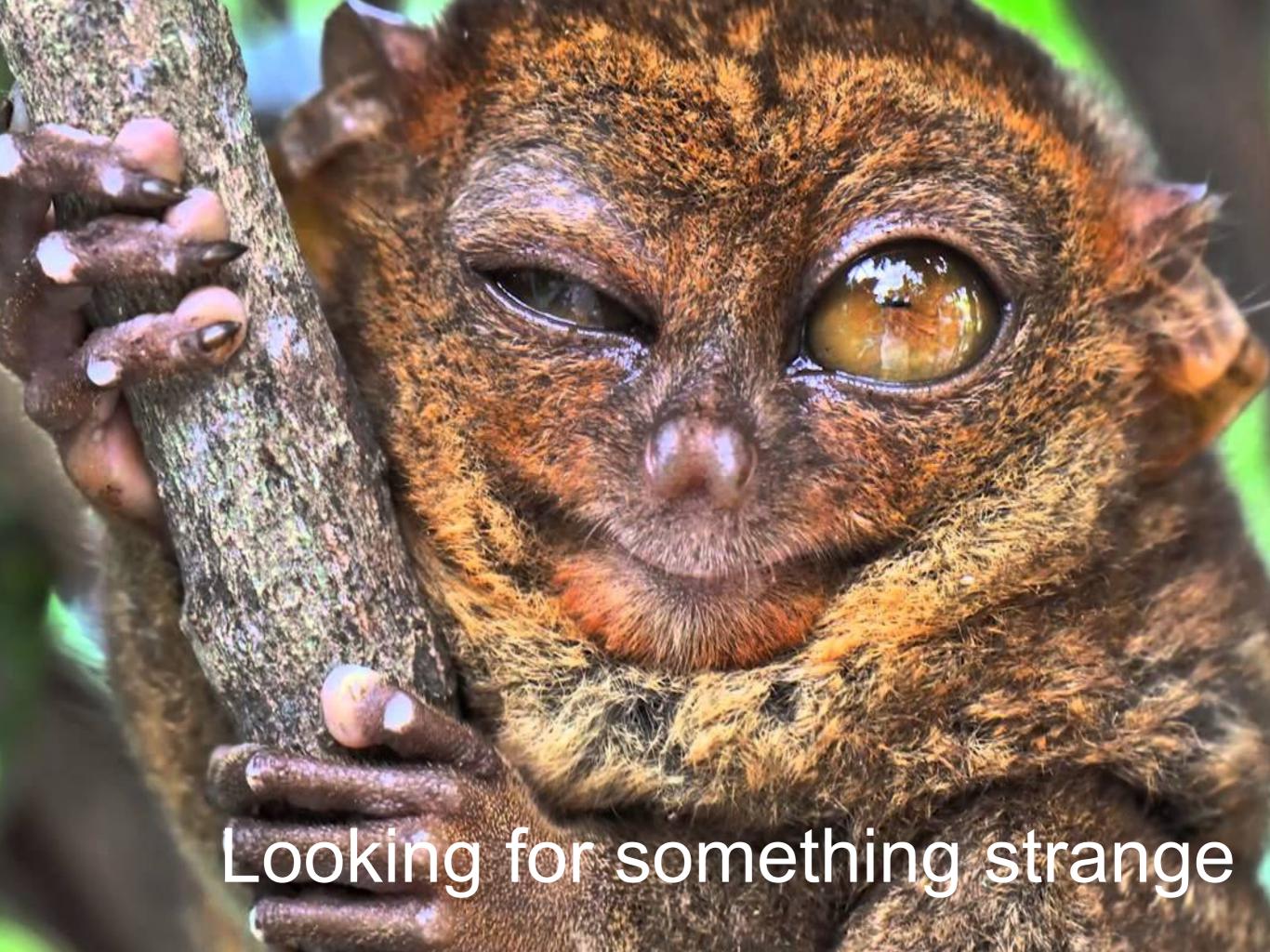


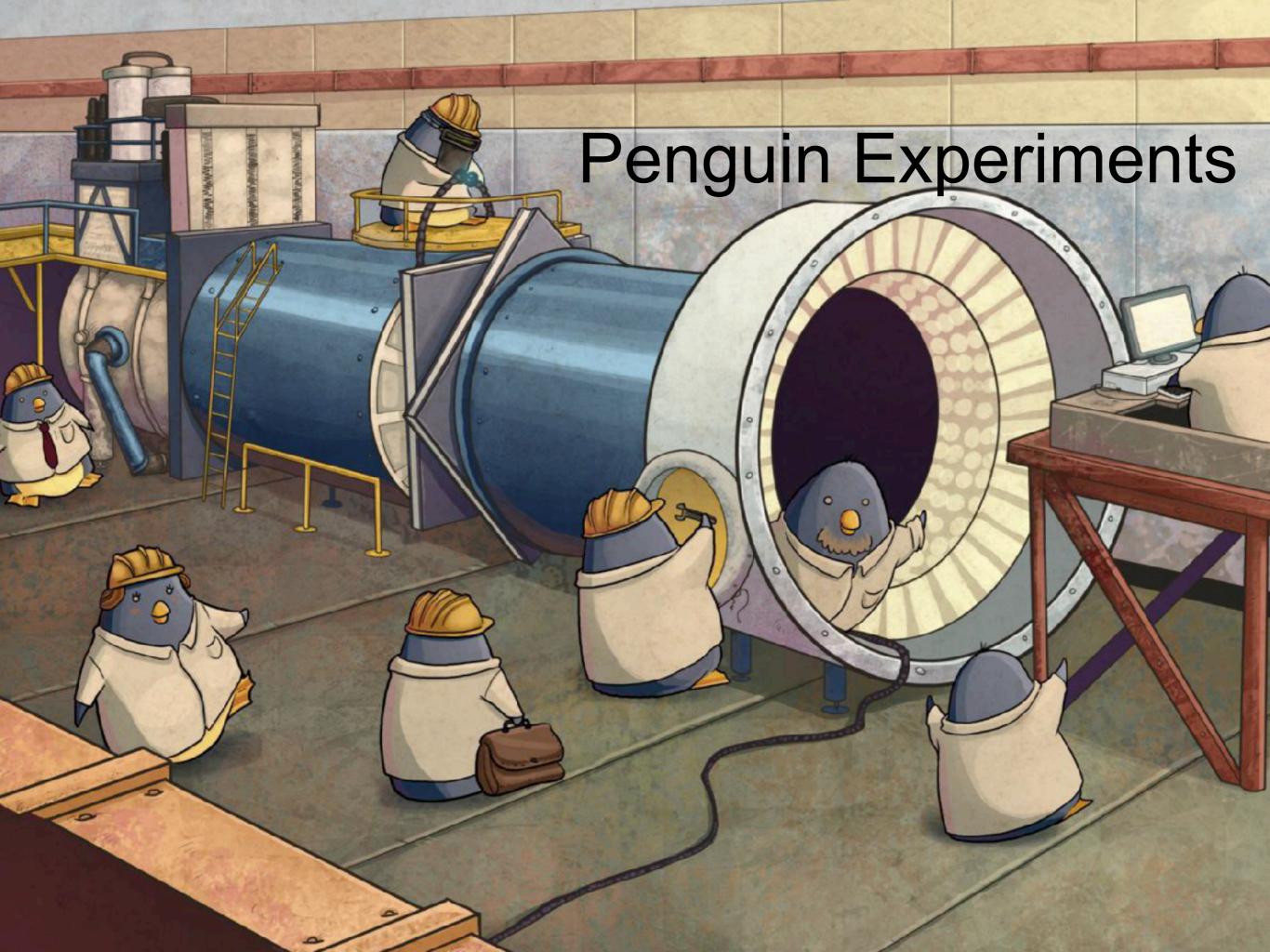
How are *statistical*, *systematic* and *normalisation* uncertainties correlated across the individual measurements?



Standard Model Total Production Cross Section Measurements Status: July 2018

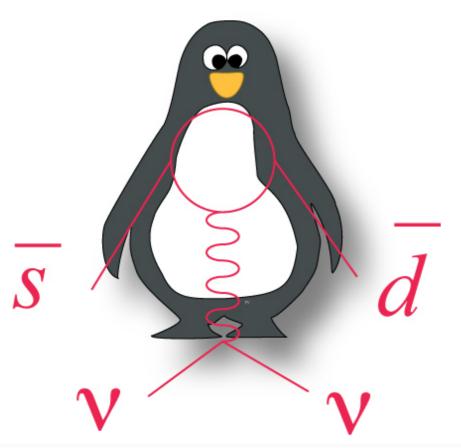


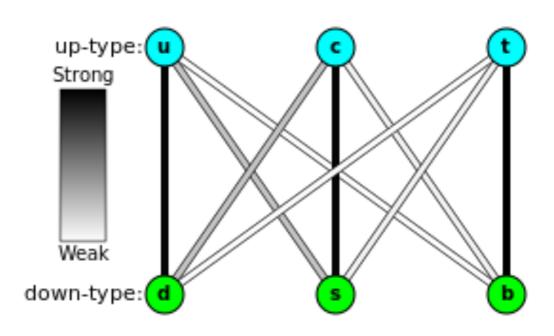






How strange quarks turn into other quarks





BR(K⁺
$$\rightarrow$$
 π⁺υ $\bar{\nu}$) = (8.4 ± 1.0) \times 10⁻¹¹ [Buras et al. JHEP 1511 (2015) 33]

Theory predicts this happens ~100 times in a million million kaon decays:

- 1) Record a million million kaon decays
- 2) Analysis do we count 100 signal events?

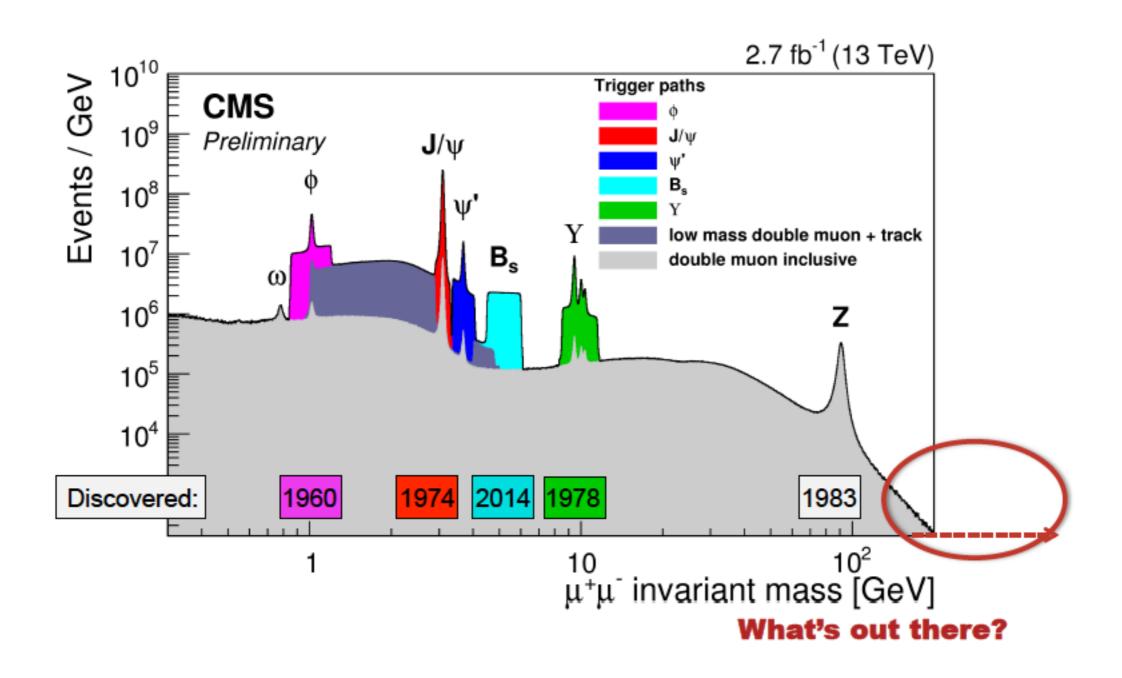
 Yes Congratulations, that's very impressive!

 No Congratulations, you've discovered new physics!

1 in 10 billion is the same probability as finding you in the Earth's population



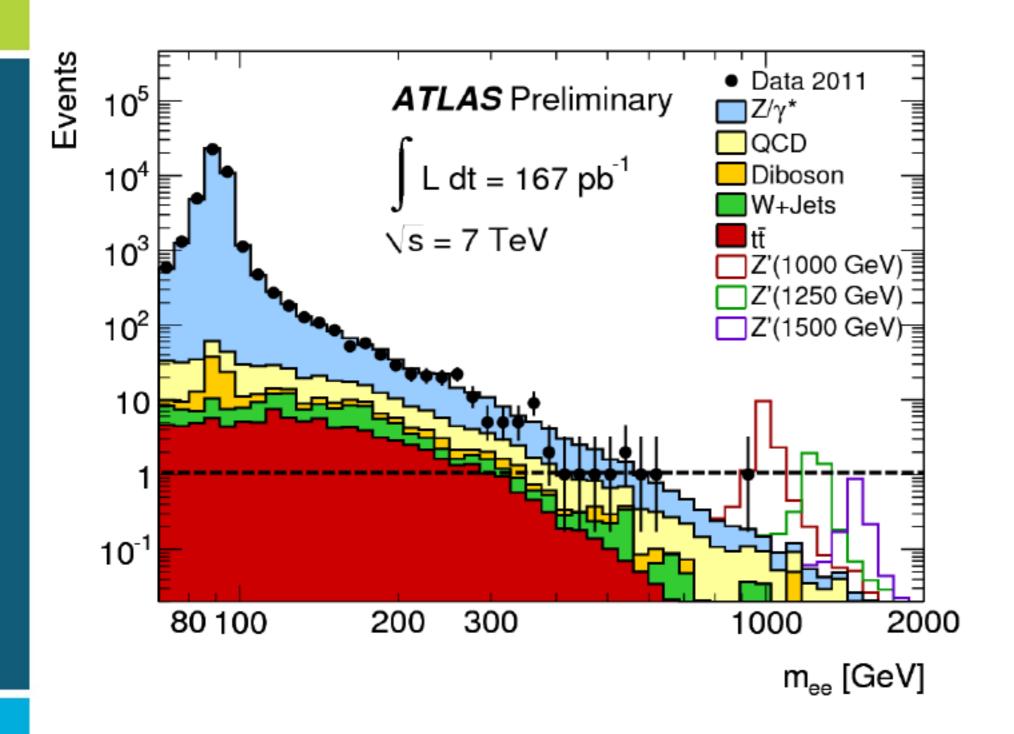
Elements of a search





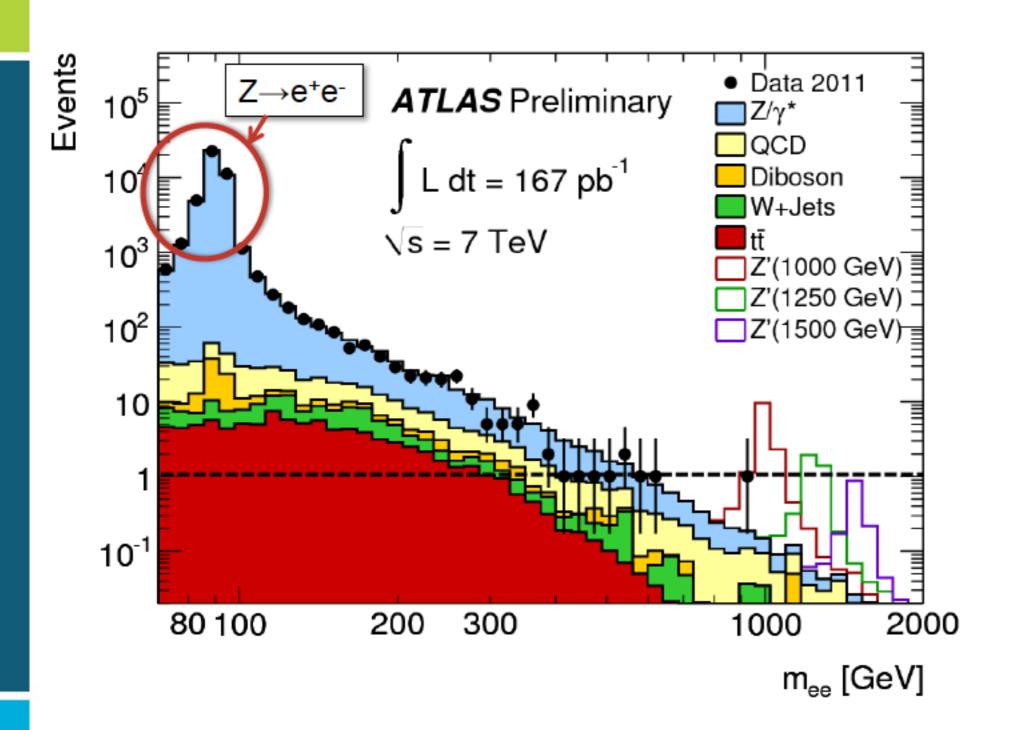
Elements of a search

© Like Z->ee but at higher mass.



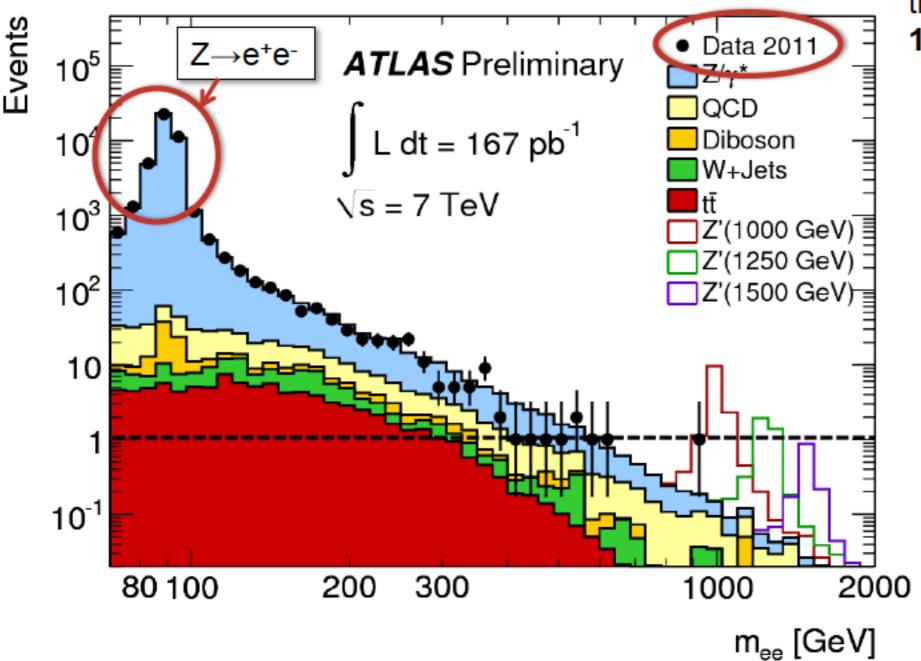


© Like Z->ee but at higher mass.





© Like Z->ee but at higher mass.

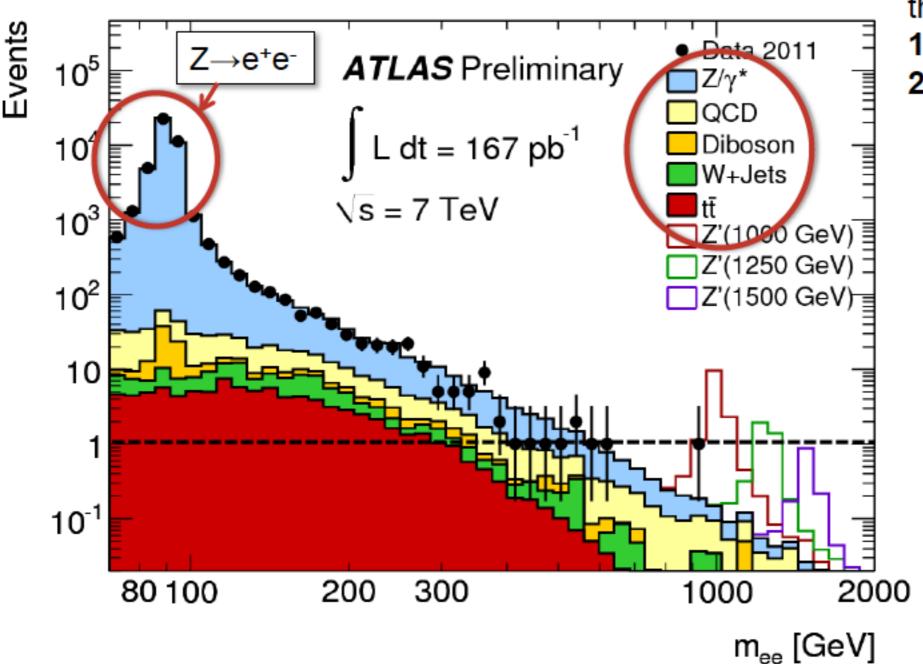


Select 2 electron candidates and plot their invariant mass for:

1. Data



© Like Z->ee but at higher mass.

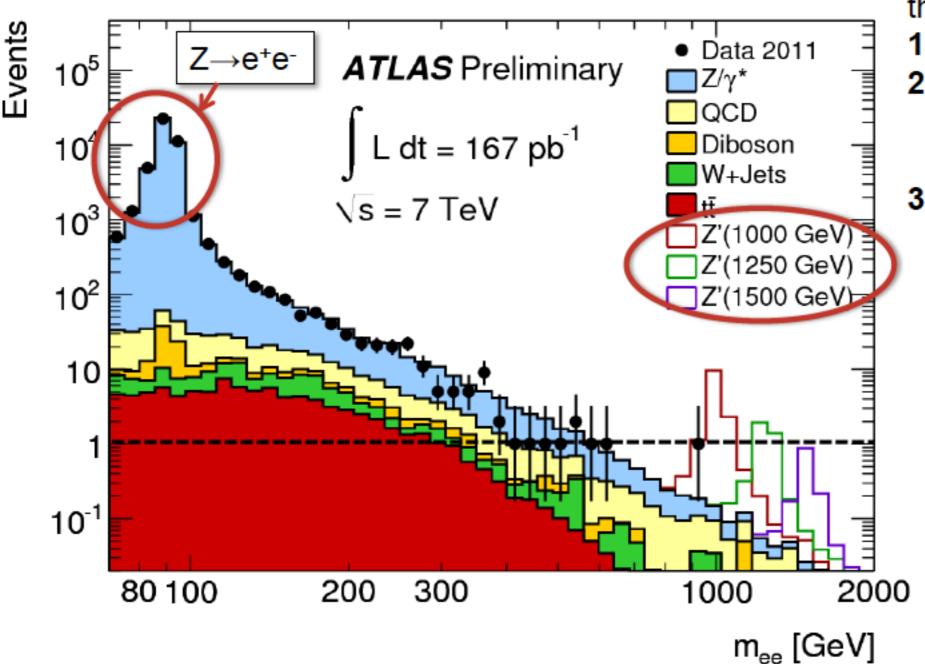


Select 2 electron candidates and plot their invariant mass for:

- 1. Data
- 2. Simulated background events



© Like Z->ee but at higher mass.

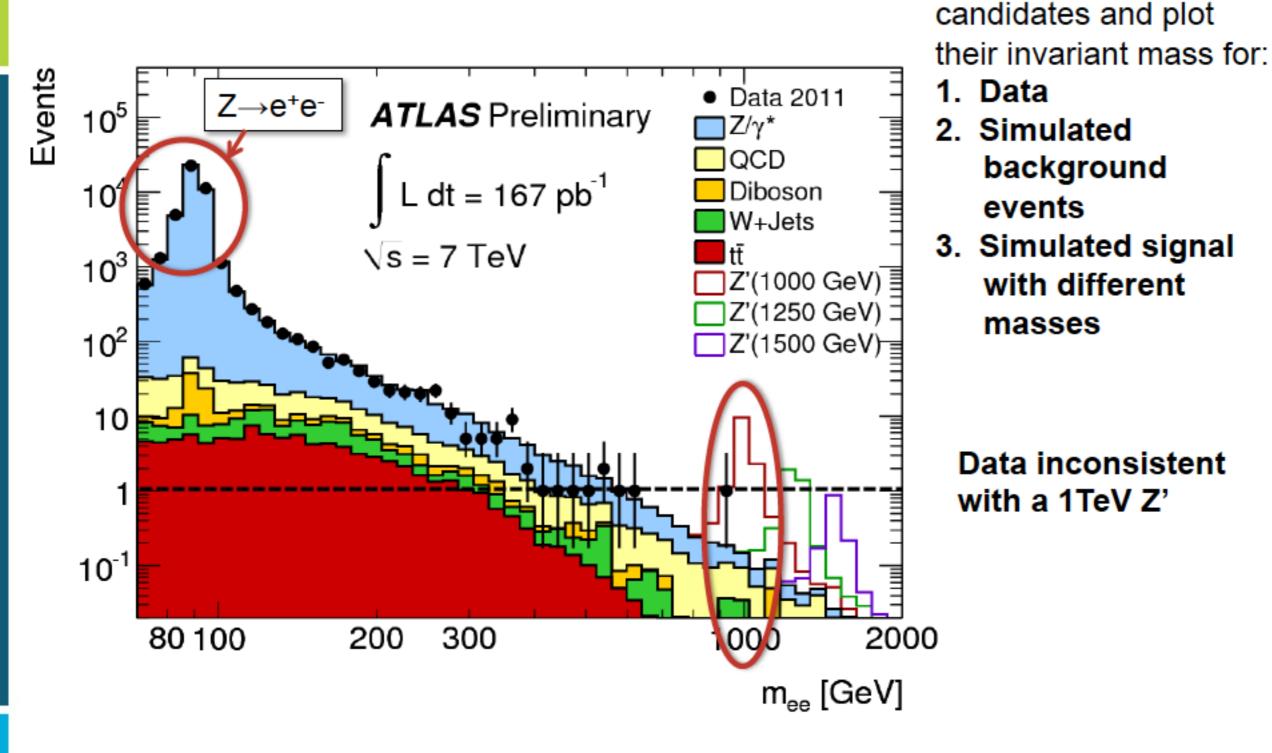


Select 2 electron candidates and plot their invariant mass for:

- 1. Data
- 2. Simulated background events
- 3. Simulated signal with different masses



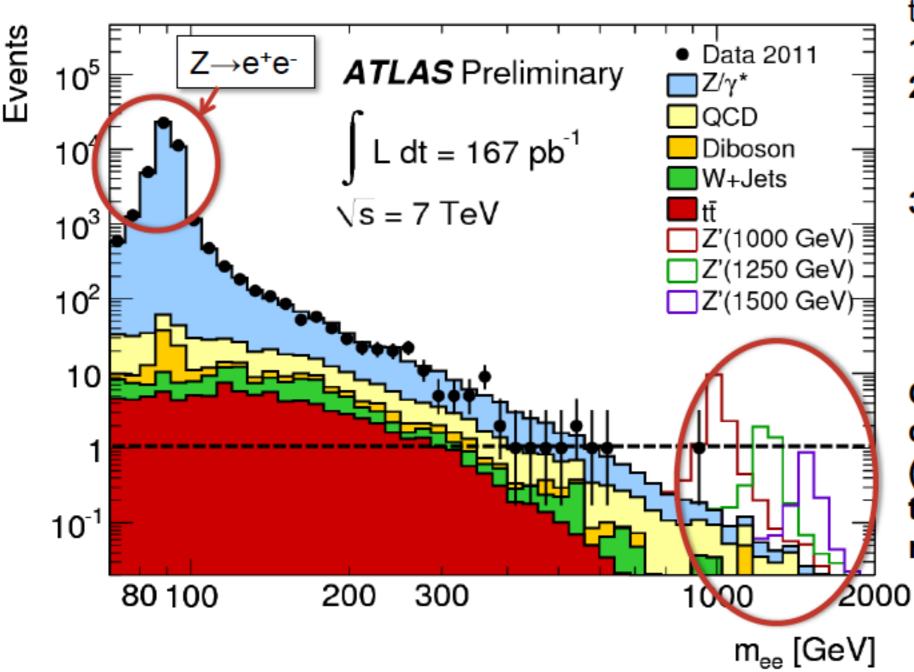
© Like Z->ee but at higher mass.





Select 2 electron

© Like Z->ee but at higher mass.



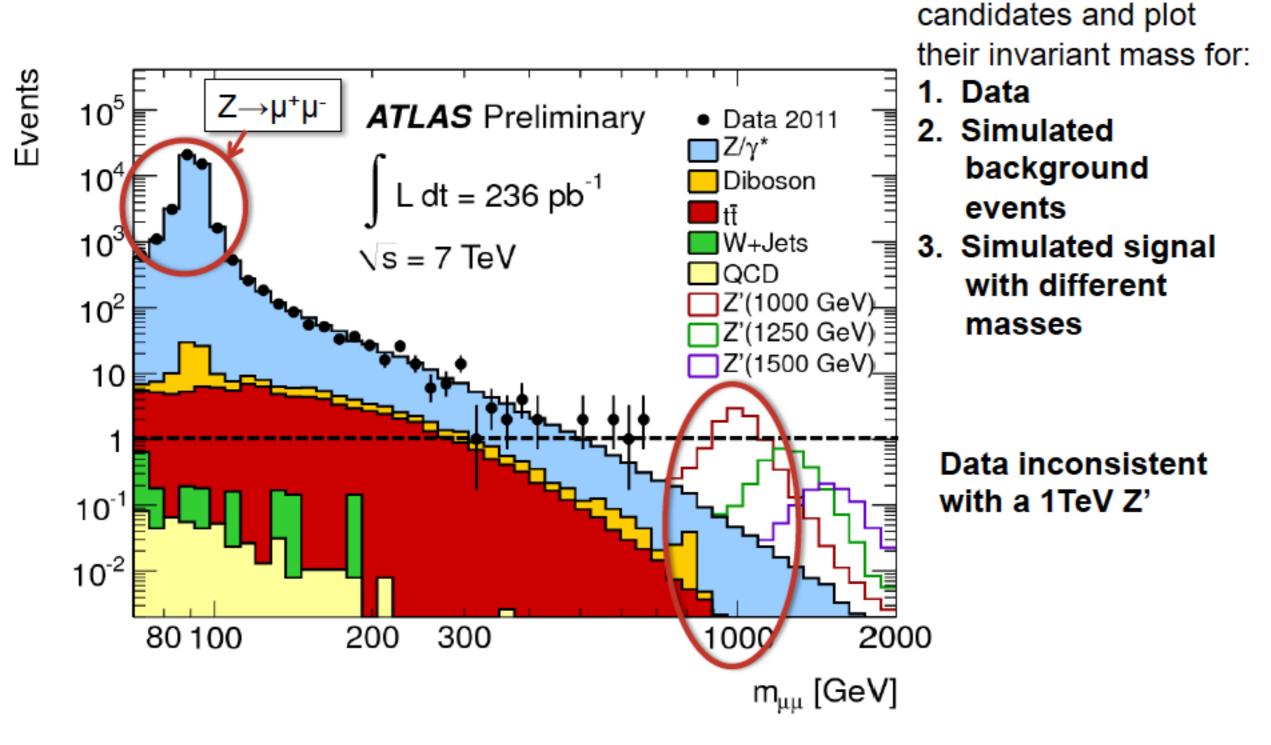
Select 2 electron candidates and plot their invariant mass for:

- 1. Data
- 2. Simulated background events
- 3. Simulated signal with different masses

Cross-section decreases with mass (higher the mass of the Z', the more data needed to discover it)

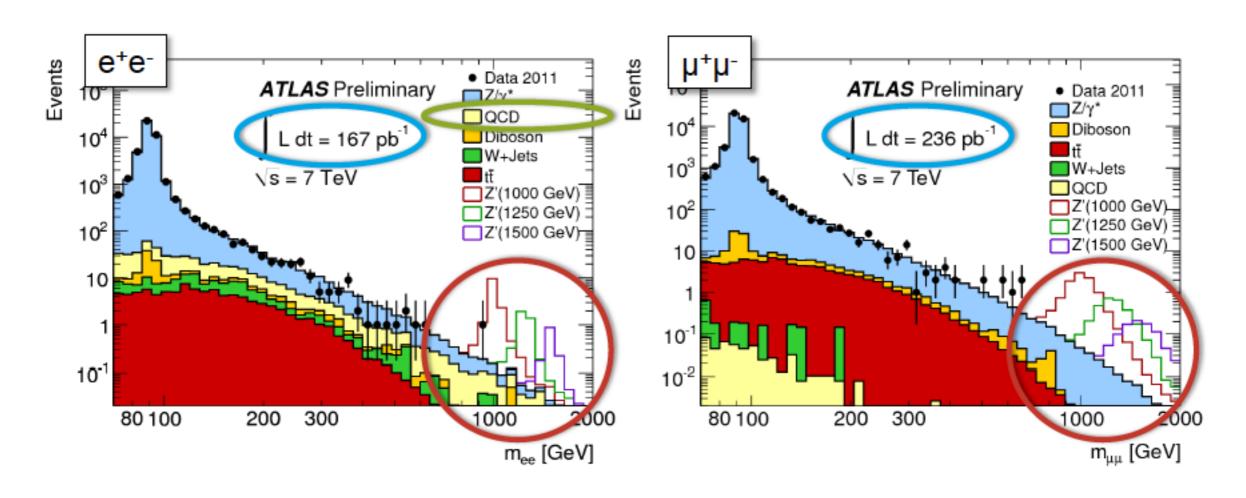


And similar for muons





Select 2 muon



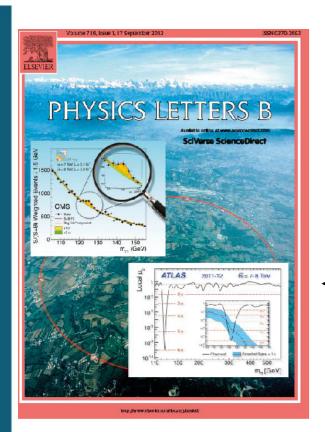
Differences in:

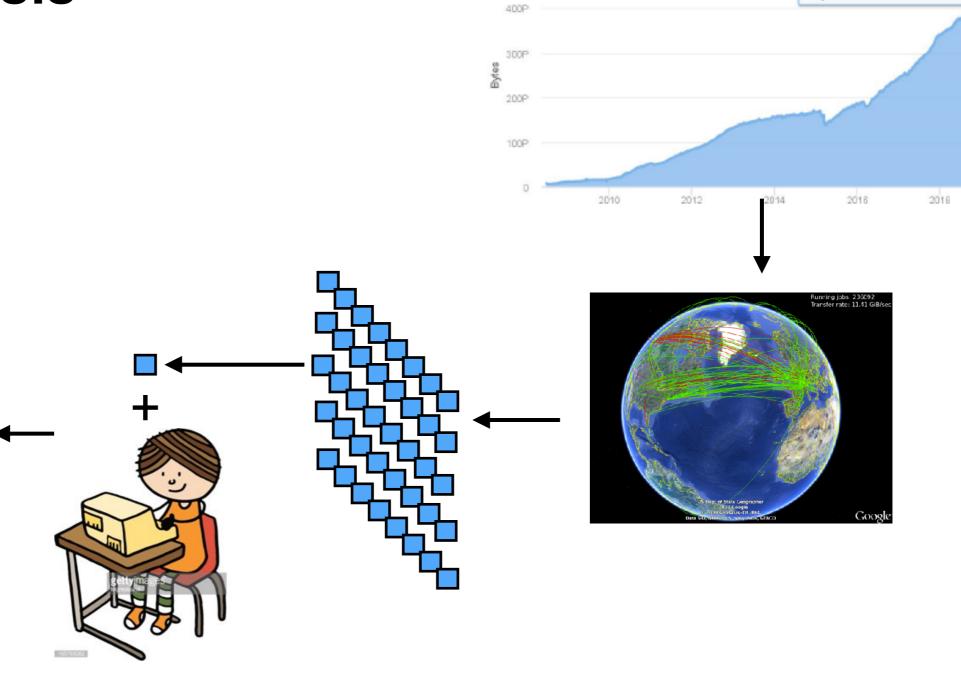
Why is the resolution worse in the muon channel?

- Resolution
- Background composition
- O Dataset



Data analysis





500P



Sunday, Mar 3, 2019 • Bytes: 452 760 696 701 954 000

Needles in haystacks

We record billions of events

• The data are structured but each event is different - unique data

science challenge

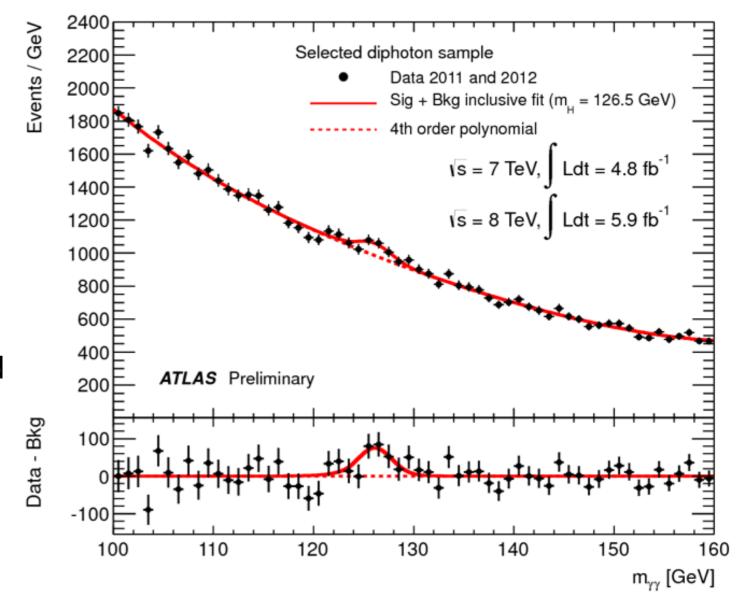
Data reduction proceeds via a twopronged approach:

Select only the events that you are interested in

• e.g. events with two photons

Keep only the information you need Throw away the rest!

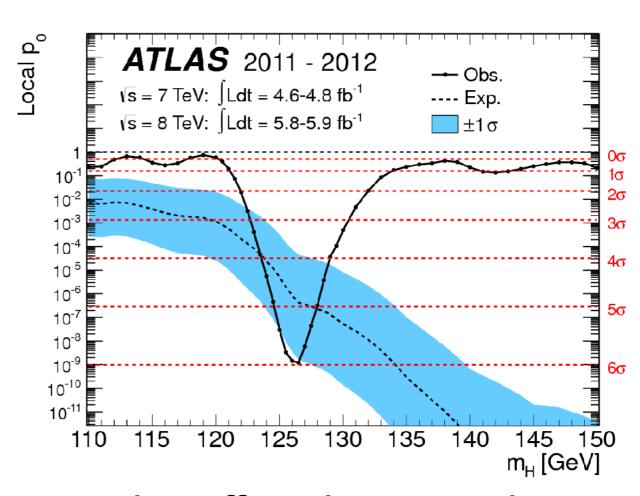
Final statistical inference is only performed on the reduced data





Higgs discovery in 2012

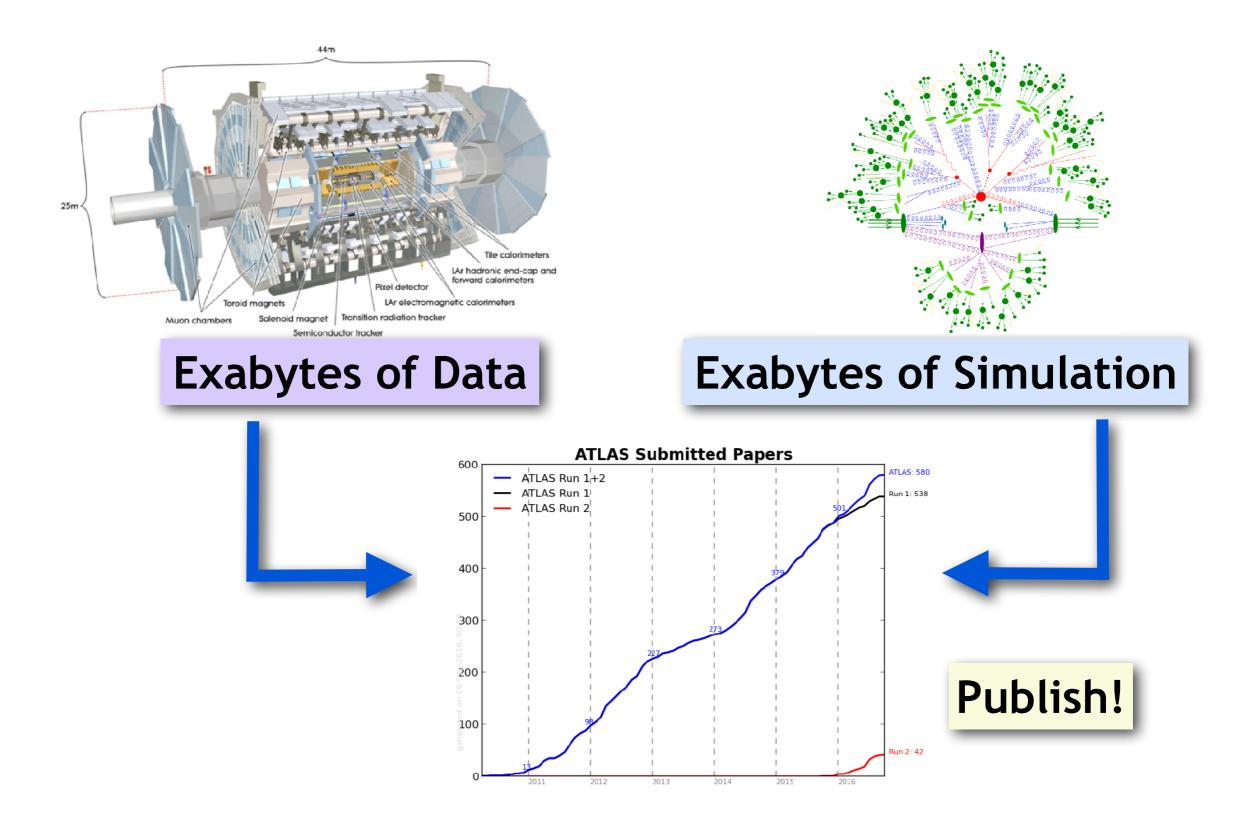




Now you can appreciate a little better the effort that went into this plot - and why we were very happy!



Now you know how to do exabyte-scale physics analysis!



Contact details

Dr Paul Laycock, based at CERN

Office: B40 4-C-16

email: laycock@bnl.gov or paul.james.laycock@cern.ch

In-person Q&A again!:

Tuesday 13:30-15:00 in Salle Anderson, B40 S2-A01

