



Tokyo Tech

16<sup>th</sup> Varenna Conference on Nuclear Reaction Mechanisms

# Prompt-fission observable and fission yield calculations for actinides by TALYS

Tokyo Tech.<sup>1</sup>, Uppsala University<sup>2</sup>, CEA<sup>3</sup>, IAEA<sup>4</sup>

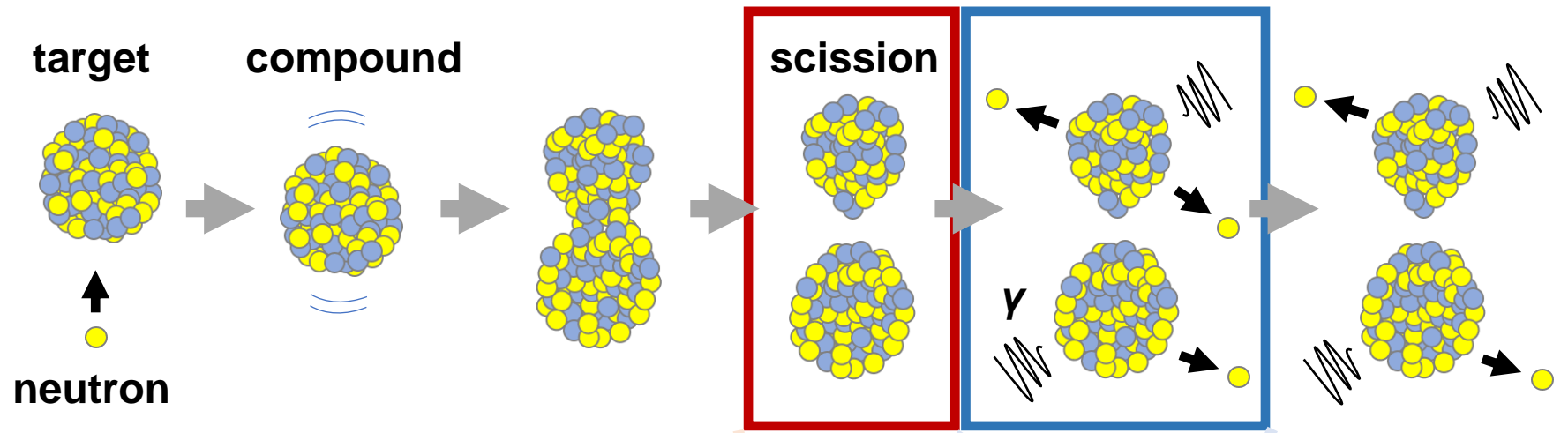
**Kazuki Fujio**<sup>1</sup>, Ali Al-Adili<sup>2</sup>, Fredrik Nordström<sup>2</sup>, Jean-François Lemaître<sup>3</sup>, Shin Okumura<sup>4</sup>, Satoshi Chiba<sup>1</sup>, Arjan Koning<sup>4</sup>

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# Outline

- **Hauser-Feshbach statistical decay calculation implemented in TALYS**
- **Sensitivity study of fission observables based on  $^{235}\text{U}$  + thermal neutron reaction**
- **Application to neutron-induced fission of  $^{235}\text{U}$  from thermal up to 5 MeV**
- **Global study on actinides at 1 MeV incident neutron reaction**
- **Conclusions**

# New approach implemented in TALYS



## Fission fragment information

- Fission fragment yield  $Y_{ff}(Z, A)$
- Total Kinetic Energy TKE
- Mean excitation energy  $\bar{E}_x$
- Width of the excitation energy distribution  $\sigma E_x$

## Hauser-Feshbach statistical decay

- Independent fission product yield
- Isomeric yield ratio
- Neutron /  $\gamma$ -ray multiplicity
- Prompt Fission Neutron Spectrum / Prompt Fission  $\gamma$ -ray Spectrum

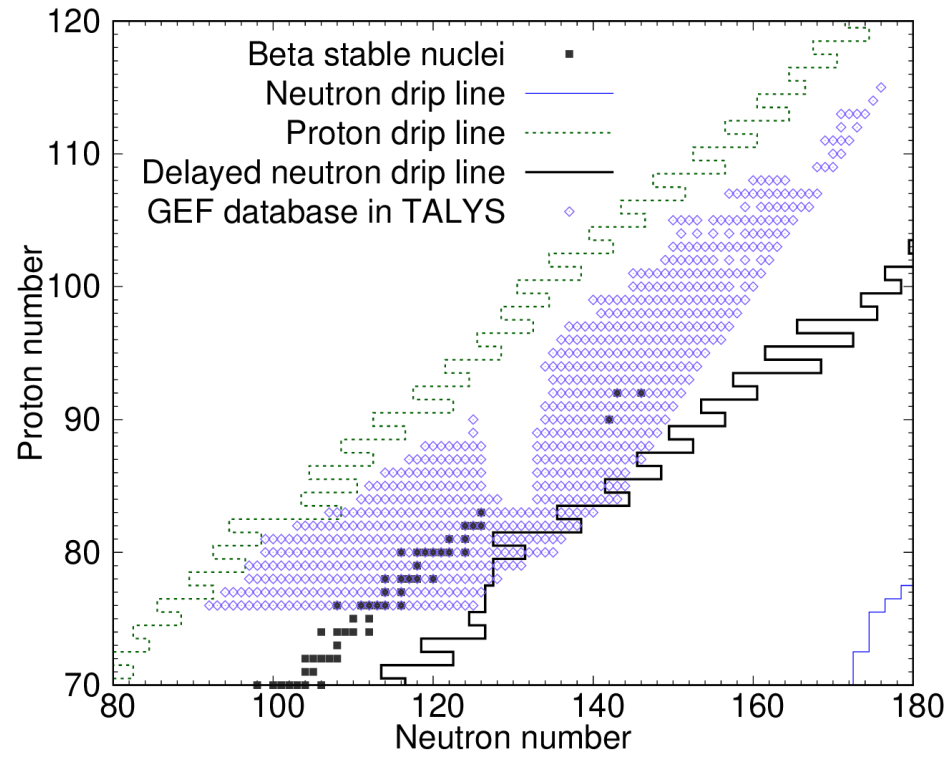
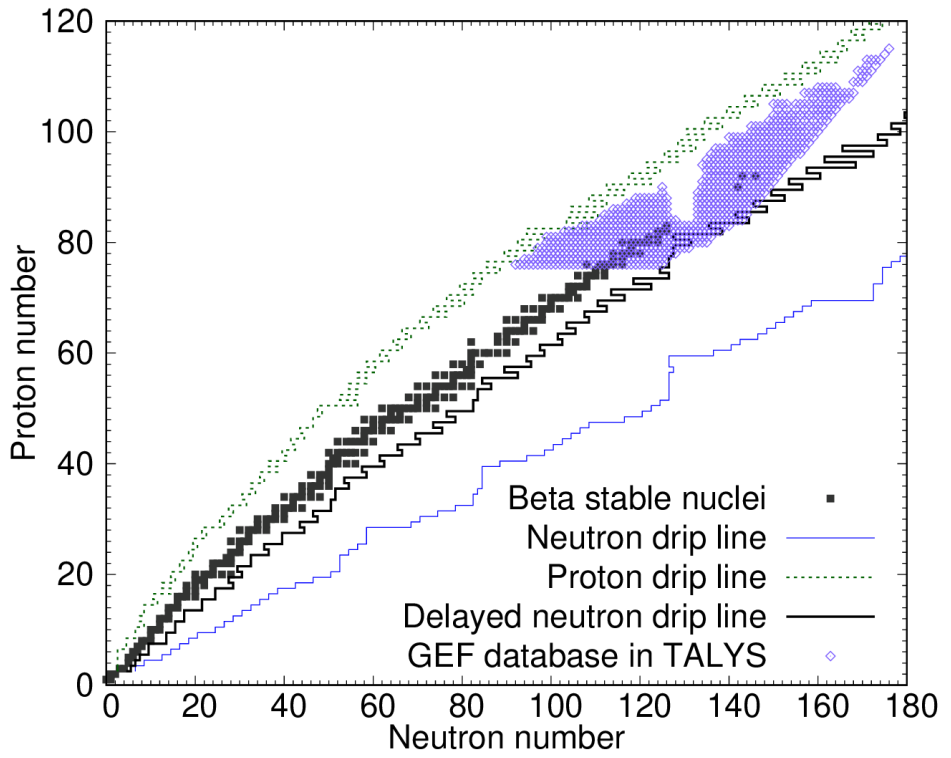
TALYS contains fission fragment information from several theoretical codes.

# Fission fragment database in TALYS

**GEF** Monte Carlo-based phenomenological fission model gives fission observables data not only after decays but also pre-neutron data.

[K. -H. Schmidt, B. Jurado, C. Amouroux, C. Schmitt, Nuclear Data Sheets, 131, 107-221 \(2016\).](#)

- TALYS contains 737 fissioning nuclei ranging from  $_{76}\text{Os}$  to  $_{115}\text{Mc}$   
[F. Nordström, Technical Report UPTeC ES21016, Uppsala university \(2021\).](#)



# Fission fragment database in TALYS

**HF<sup>3</sup>D** Designed with a fully deterministic technique with fitting functions  
[S. Okumura, T. Kawano, P. Jaffke, P. Talou, and S. Chiba, JNST, 55\(9\), 1009-1023 \(2018\).](#)

- For neutron-induced fission of <sup>235</sup>U, <sup>238</sup>U, and <sup>239</sup>Pu

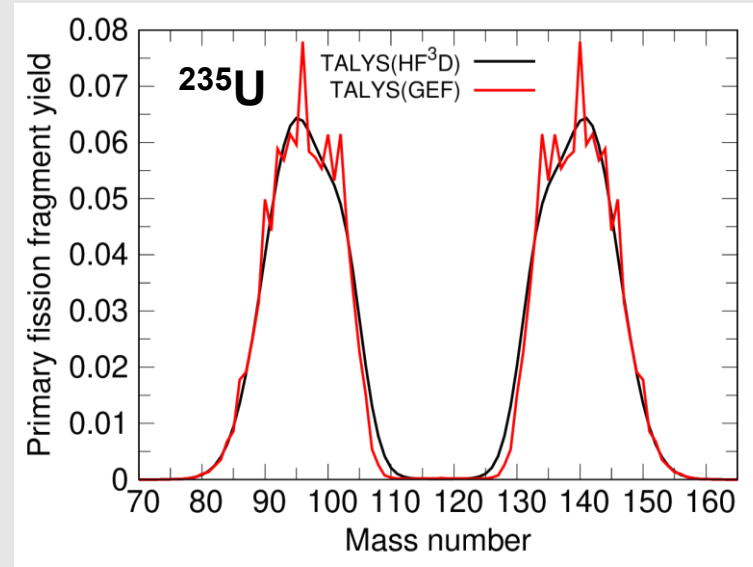
## Mass distribution of primary fragments

$$Y(A) = \sum_{i=1}^5 \frac{Y_i}{\sqrt{2\pi}\sigma_i} \exp \left\{ -\frac{(A - A_m + \Delta_i)^2}{2\sigma_i^2} \right\}$$

$$A_m = A_c/2 \quad A_c : \text{mass number of compound nucleus}$$

$\sigma_i, \Delta_i$  : Gaussian parameters

$Y_{1,5}$  : yield ( $i = 1 \sim 5$ )



## TKE distribution of primary fragments

$$\text{TKE}(A_h) = (p_1 - p_2 A_h) \left\{ 1 - p_3 \exp \left( -\frac{(A_h - A_m)^2}{p_4} \right) \right\} + \varepsilon_{\text{TKE}}$$

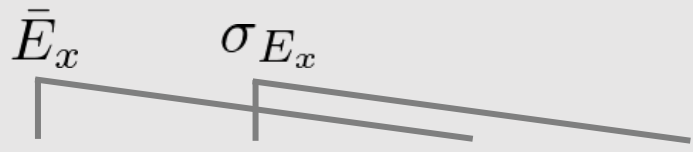
$p_i$  : fitting parameters  $\varepsilon_{\text{TKE}}$  : correction term to ensure the average TKE

- Charge distribution is obtained from Wahl's  $Z_p$  model
- $R_T$  model is used for the excitation energy partition  
 $R_T = 1.29$  at thermal energy and  $R_T = 1.00$  as  $R_T$  approaches 5 MeV.

# Fission fragment table

# Z = 92  
 # A = 236  
 # Ex (MeV) = 6.55e+00  
 # Ntotal = 207

#	Zl	Al	Zh	Ah	Yield	TKE[MeV]	TXE[MeV]	$\bar{E}_x$	$\sigma_{E_x}$	E1[MeV]	W1[MeV]	Eh[MeV]	Wh[MeV]
	28	72	64	164	1.7222e-06	1.4140e+02	2.1593e+01	9.0083e+00	3.3421e+00	1.2584e+01	4.6688e+00		
	29	73	63	163	3.3249e-06	1.4256e+02	2.2054e+01	9.1763e+00	3.3591e+00	1.2877e+01	4.7139e+00		
	30	73	62	163	1.5739e-06	1.4514e+02	2.0166e+01	8.7278e+00	3.5748e+00	1.1438e+01	4.6847e+00		
	29	74	63	162	1.6362e-06	1.4375e+02	1.9948e+01	7.9037e+00	3.2017e+00	1.2044e+01	4.8788e+00		
	30	74	62	162	1.0661e-05	1.4634e+02	2.2937e+01	1.0050e+01	3.6535e+00	1.2887e+01	4.6847e+00		



## Excitation energy distribution

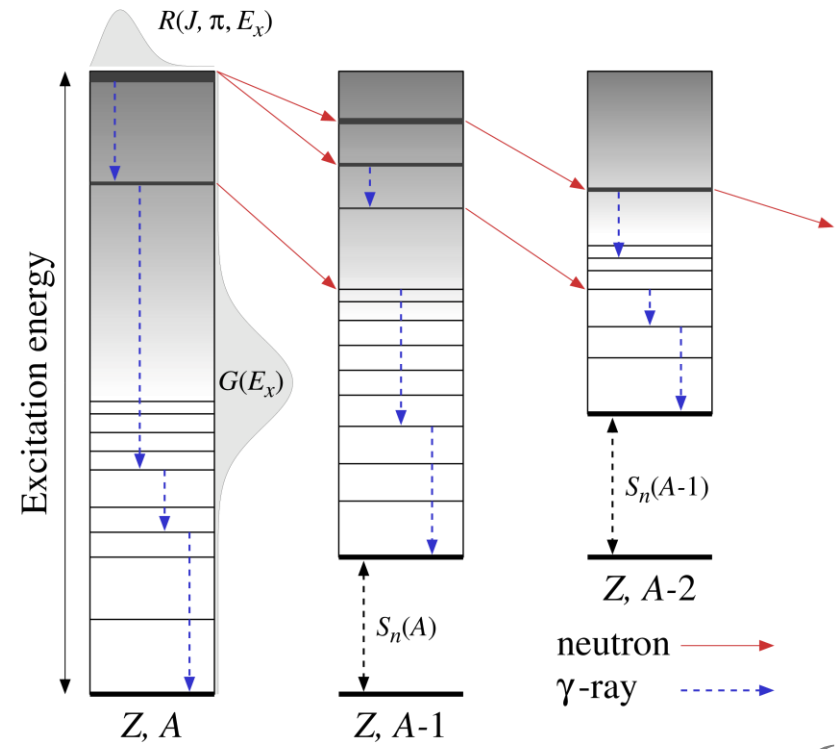
$$G(E_x) = \frac{1}{\sqrt{2\pi}\sigma_{E_x}} \exp \left\{ -\frac{(E_x - \bar{E}_x)^2}{2\sigma_{E_x}^2} \right\}$$

$\bar{E}_x$  : mean excitation energy  
 $\sigma_{E_x}$  : width of the excitation energy distribution

## Spin-parity distribution

$$R(J, \pi, E_x) = \frac{1}{2} \cdot \frac{2J + 1}{2f^2\sigma^2(E_x)} \exp \left\{ -\frac{(J + 1/2)^2}{2f^2\sigma^2(E_x)} \right\}$$

$\sigma^2(E_x)$  : spin-cut off parameter  
 $f^2$  : scaling factor



# Sensitivity study of fission observables

## Spin-parity distribution

$$R(J, \pi, E_x) = \frac{1}{2} \cdot \frac{2J + 1}{2X\sigma^2(E_x)} \exp \left\{ -\frac{(J + 1/2)^2}{2X\sigma^2(E_x)} \right\}$$

$X$  : scaling factor

For primary fission fragments  $\rightarrow X = f^2$

For fission products  $\rightarrow X = f_s$

to assure a reasonable agreement with experimental data.

## The number of continuum states $N$

$$\Delta_{\text{bins}}(Z, A) = (E_x^{\text{max}} - E_x^{\text{level}})/N$$

$\Delta_{\text{bins}}(Z, A)$  : energy width of discretized continuum state

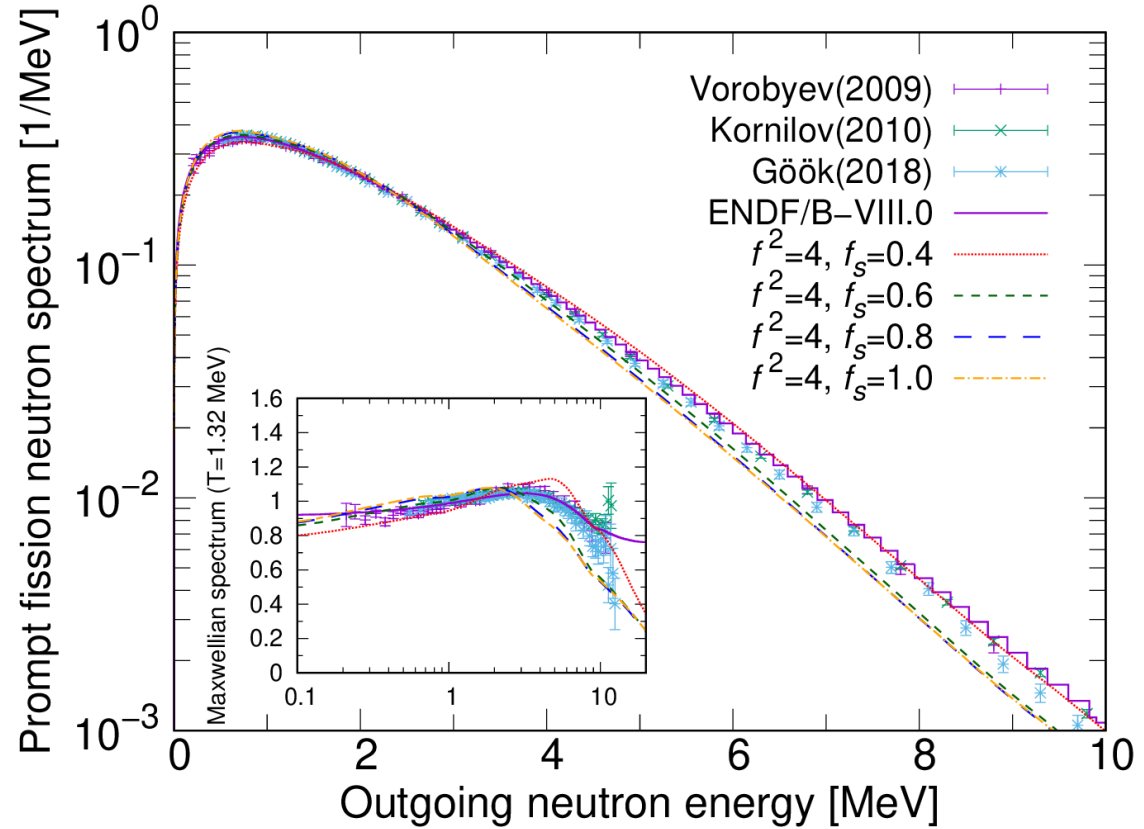
$E_x^{\text{max}}$  : maximum excitation energy

$E_x^{\text{level}}$  : excitation energy at the last discrete level

# Sensitivity on scaling parameters

- Sensitivity study using  $(Y_{ff}(Z, A), TKE, \bar{E}_x, \sigma_{E_x})$  obtained from HF<sup>3</sup>D
- Prioritize to reproduce the neutron multiplicity  $\bar{\nu}_n$  and the shape of PFNS
- Run parameter sensitivity analysis  $3 < f^2 < 6$  and  $0.4 < f_s < 1.0$

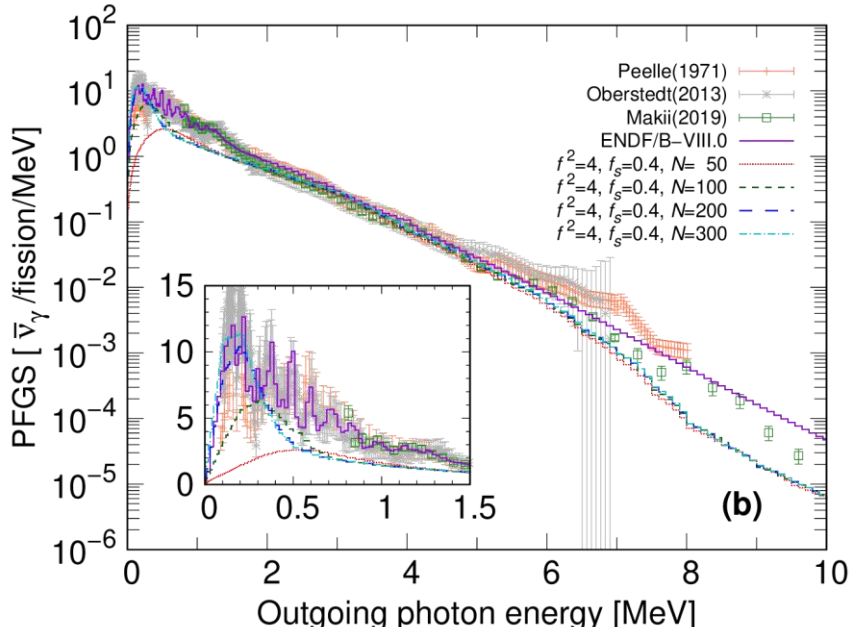
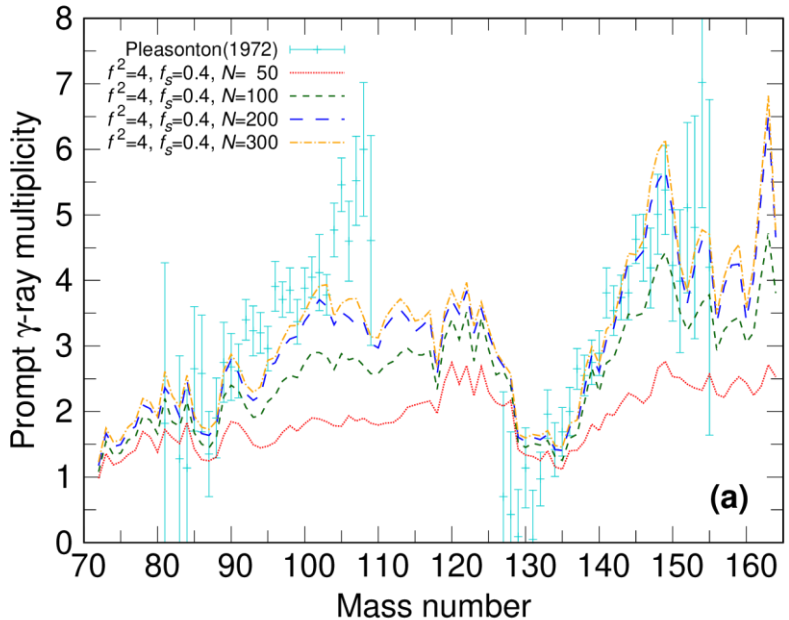
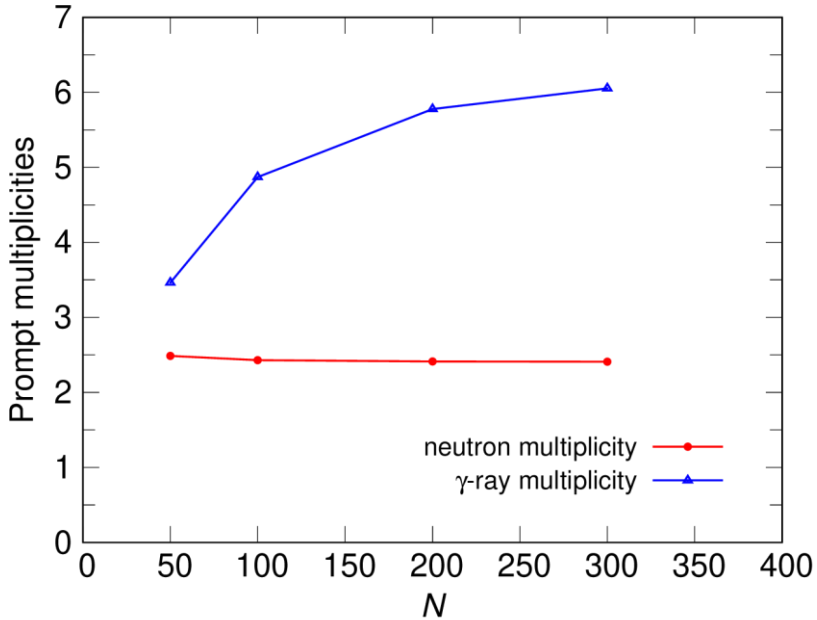
$f^2$	$f_s$	$\bar{\nu}_n$
3	0.4	2.49
3	0.6	2.47
3	0.8	2.45
3	1.0	2.44
4	0.4	2.41
4	0.6	2.40
4	0.8	2.39
4	1.0	2.38
5	0.4	2.35
5	0.6	2.34
5	0.8	2.33
5	1.0	2.33
6	0.4	2.30
6	0.6	2.29
6	0.8	2.29
6	1.0	2.29
ENDF-B/VIII.0		2.41
JEFF-3.3		2.41



- **By comparing with evaluated data, we conclude  $f^2 = 4$  and  $f_s = 0.4$  are the optimal values.**
- Our result is consistent with another recent study ( $f_s = 0.4 - 0.5$ )  
[A. Rodrigo, et al., arXiv:2303.09595 \(2023\)](https://arxiv.org/abs/2303.09595)



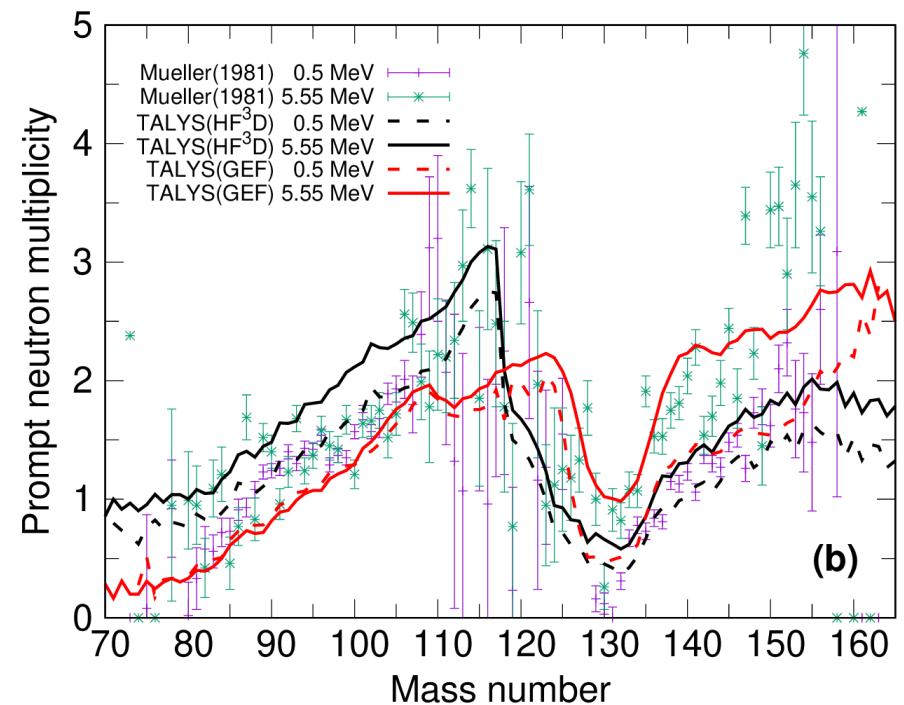
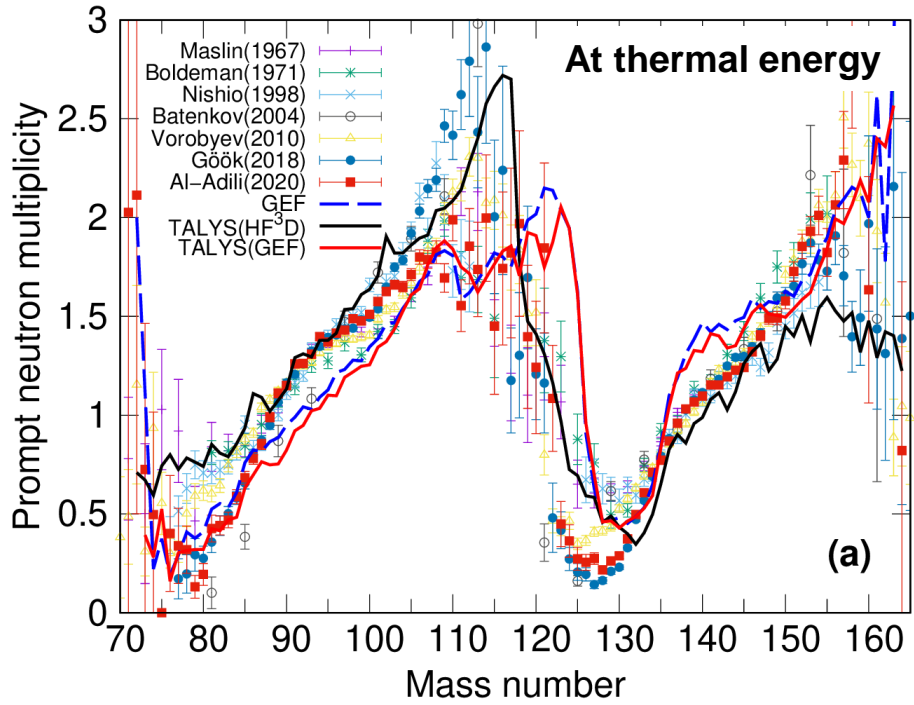
# Sensitivity on the number of continuum states



- The  $\gamma$ -ray multiplicity  $\bar{\nu}_\gamma$  increases with increasing  $N$ .
- A prominent peak appears in PFGS.
- $\therefore$  Transitions between small energy levels increase as  $N$  increases.
- **We chose  $N = 300$  as the optimal value for improved consistency with experimental and evaluated data.**

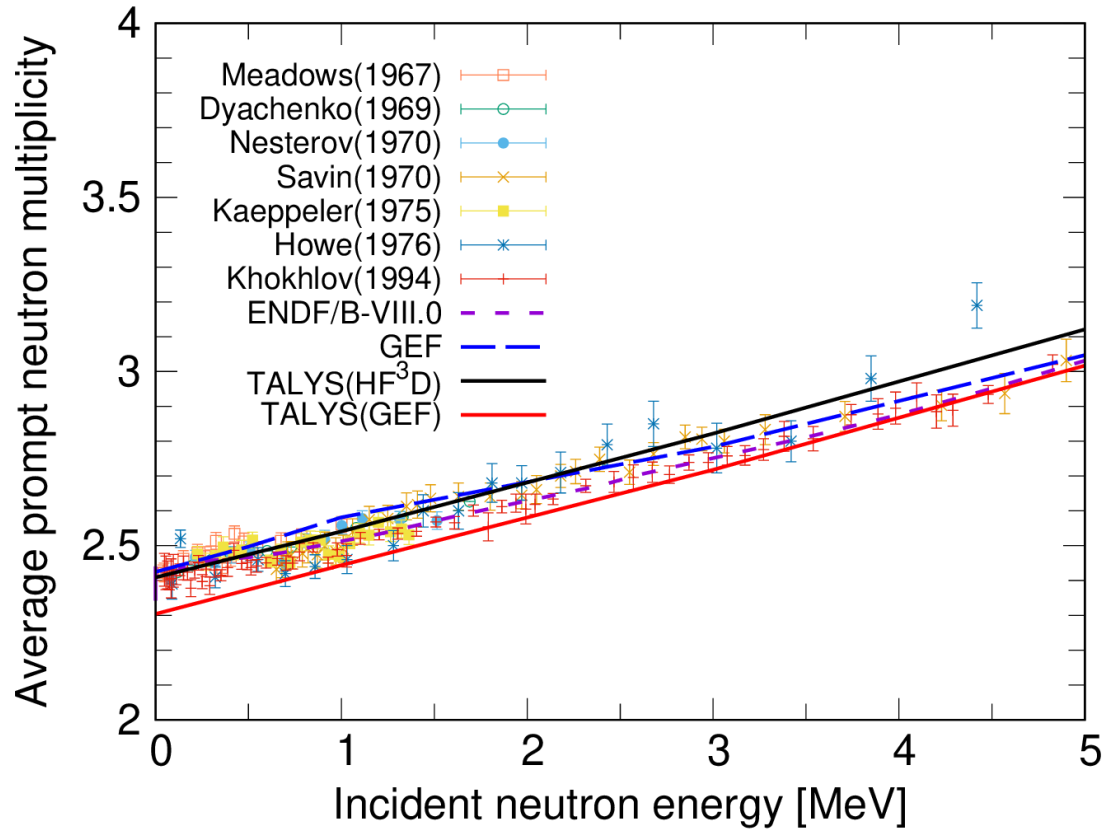
# Application to $^{235}\text{U}(n,f)$

- The optimal values:  $f^2 = 4$ ,  $f_s = 0.4$ , and  $N = 300$
- Comparison between TALYS(GEF), TALYS(HF<sup>3</sup>D), experimental, and evaluated data



- TALYS results show the saw-tooth shape that is consistent with experimental data.
- TALYS(GEF) reproduces that the  $\bar{\nu}_n$  increases from heavy fragments.
- TALYS(HF<sup>3</sup>D) does not exhibit the trend.  
 → TALYS reflects the difference in the energy-sorting mechanism from GEF and HF<sup>3</sup>D

# $^{235}\text{U}(n,f)$ : neutron multiplicity



	$\bar{\nu}_n$
TALYS(GEF)	2.30
TALYS(HF <sup>3</sup> D)	2.41
GEF	2.42
ENDF-B/VIII.0	2.41
JEFF-3.3	2.41

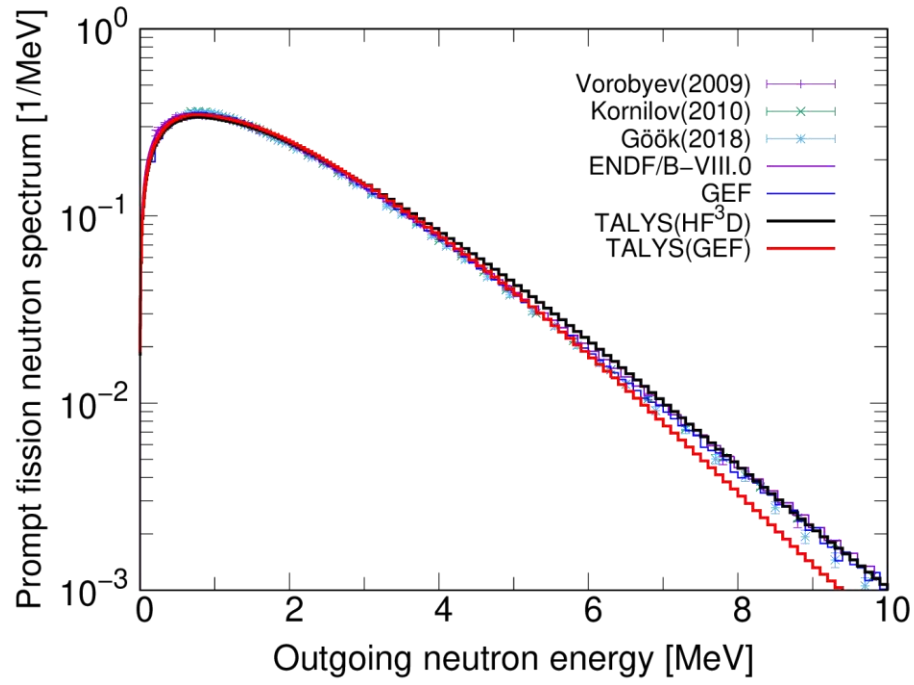
- TALYS(HF<sup>3</sup>D) successfully reproduces the evaluated value at thermal energy as original HF<sup>3</sup>D model.

[S. Okumura, T. Kawano, P. Jaffke, P. Talou, and S. Chiba, JNST, 55\(9\), 1009-1023 \(2018\).](#)

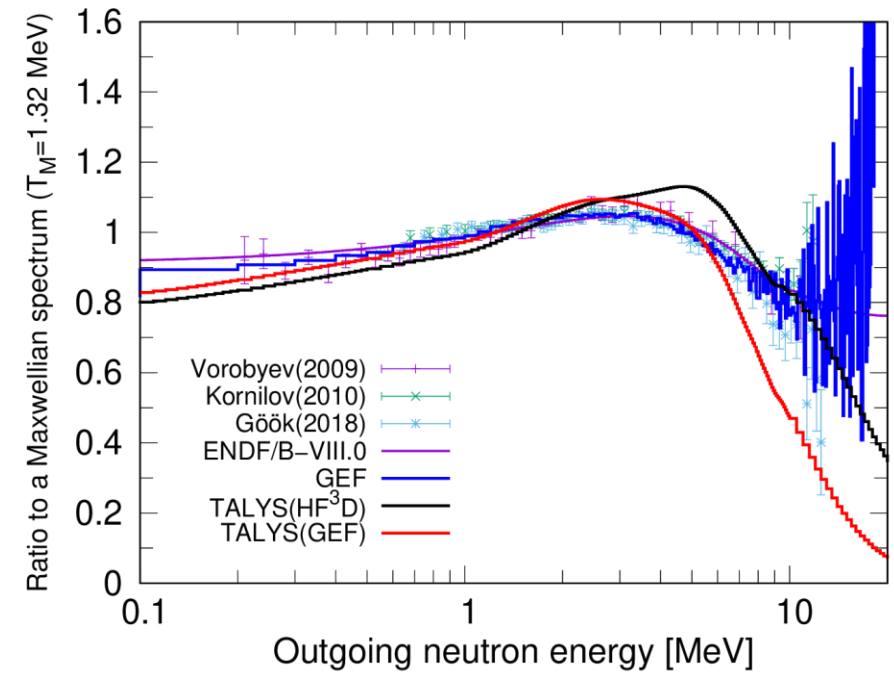
- TALYS(GEF) underestimates the evaluated data at thermal energy by about 0.1, but it agrees with the data as the incident energy increases.

# $^{235}\text{U}(n,f)$ : PFNS

## PFNS in the laboratory frame

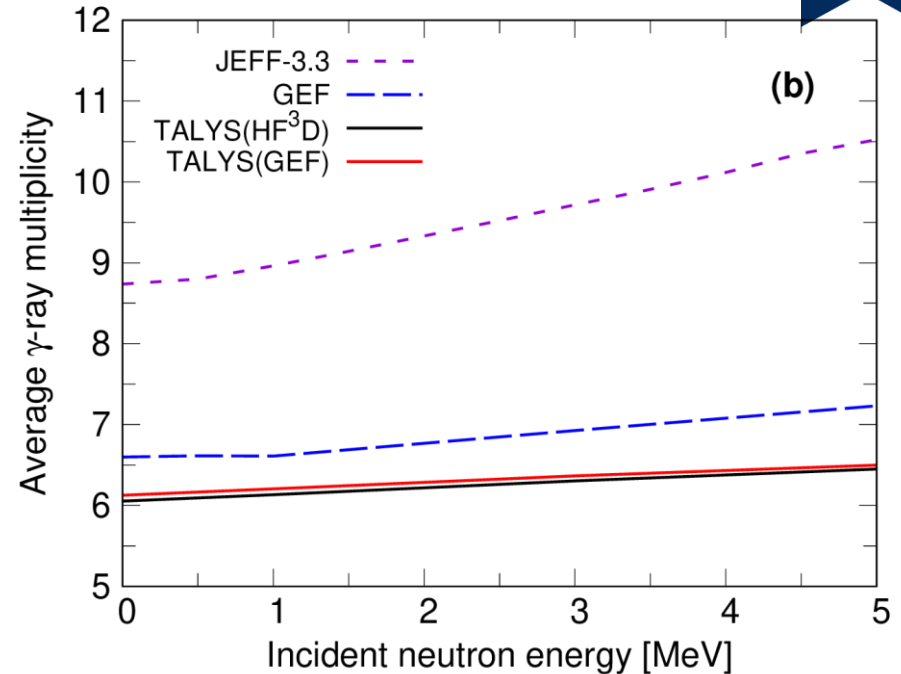
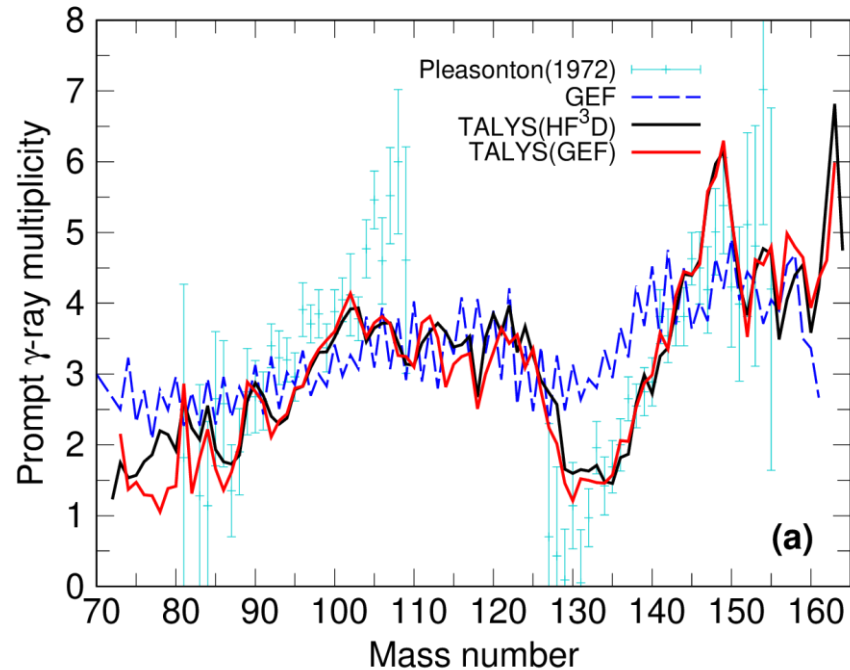


## PFNS as ratio to a Maxwellian spectrum



- TALYS(GEF) is underestimated at higher energies.
- The pronounced peak in the TALYS(HF<sup>3</sup>D) around 6 – 7 MeV is well above the experimental data.

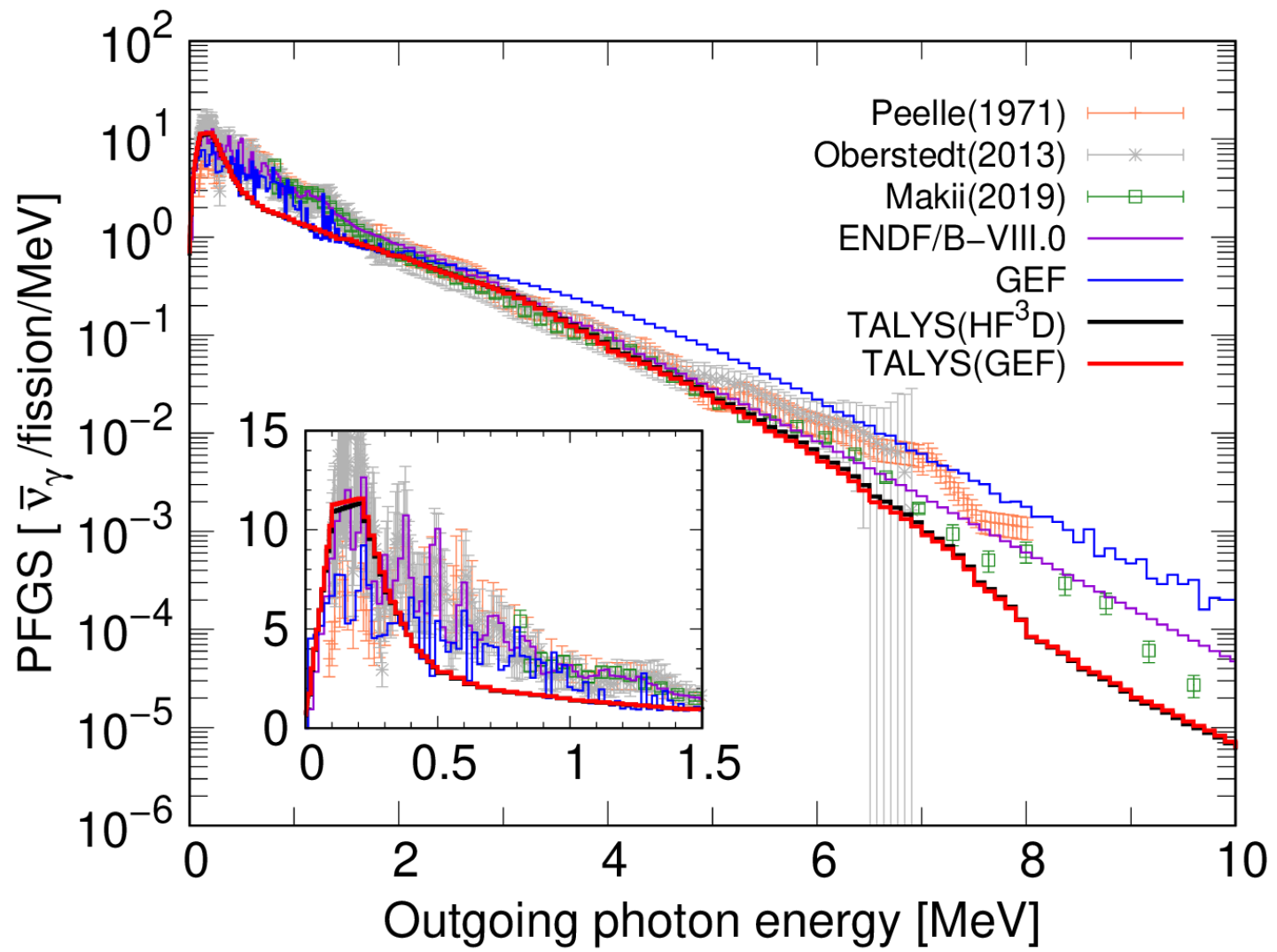
# $^{235}\text{U}(n,f)$ : $\gamma$ -ray multiplicity



	$\bar{\nu}_\gamma$
TALYS(GEF)	6.13
TALYS(HF <sup>3</sup> D)	6.05
GEF	6.61
<hr/>	
Oberstedt (2013)	$8.19 \pm 0.11$
Verbinski (1973)	$6.70 \pm 0.30$
Pleasanton (1972)	$6.51 \pm 0.30$
Peelle (1971)	$7.45 \pm 0.35$
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ENDF-B/VIII.0	8.58
JEFF-3.3	8.74

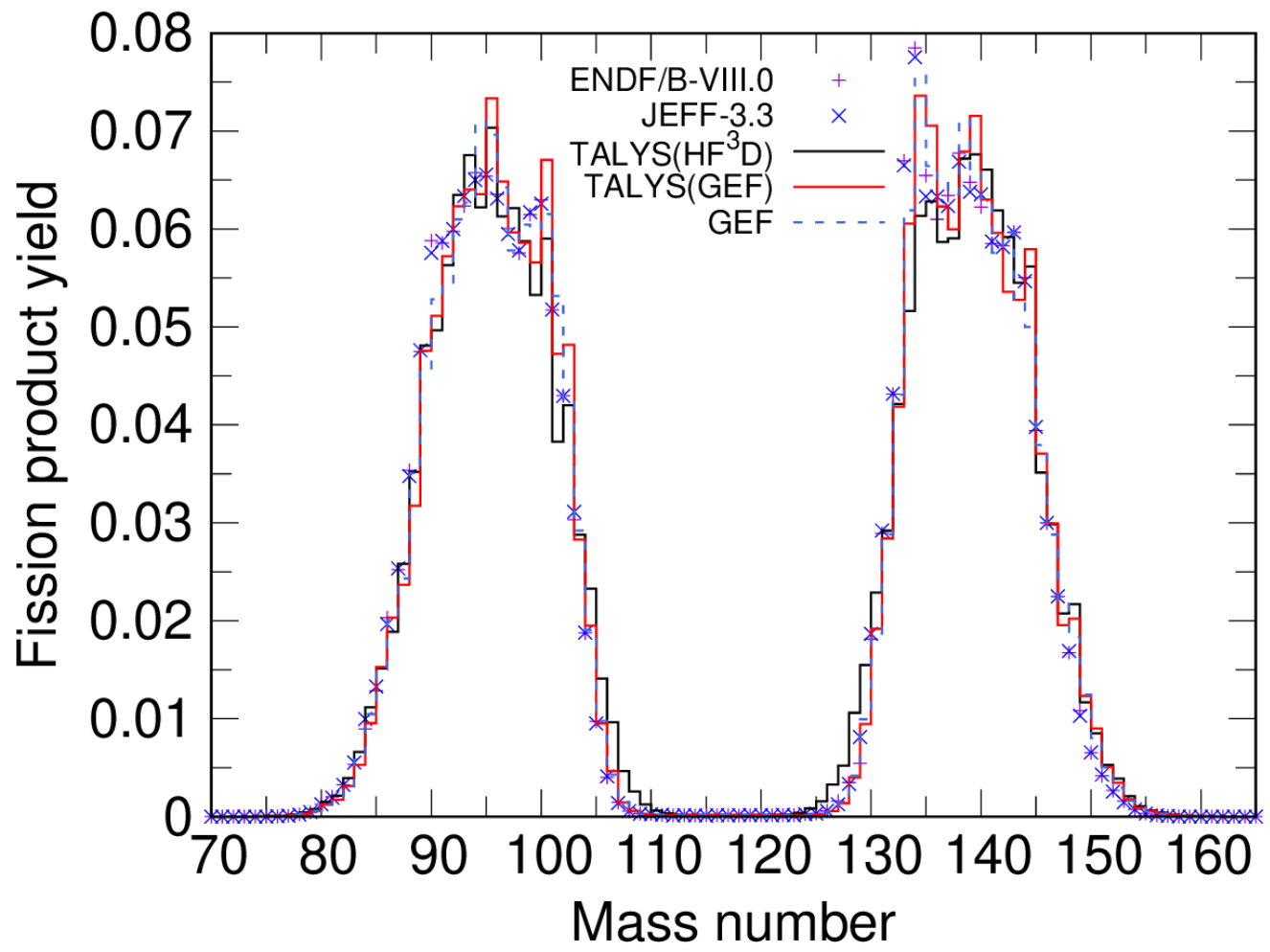
- While stand-alone GEF has a flatter  $\bar{\nu}_\gamma$ , TALYS results exhibit the saw-tooth shape.
- TALYS'  $\bar{\nu}_\gamma$  are smaller around fragment mass number  $A=100$  to  $110$ .
- TALYS underestimates  $\bar{\nu}_\gamma$  compared to experimental and evaluated data.

# $^{235}\text{U}(n,f)$ : PFGS



- A pronounced peak is observed around 0.2 MeV in both TALYS(GEF) and TALYS(HF<sup>3</sup>D).

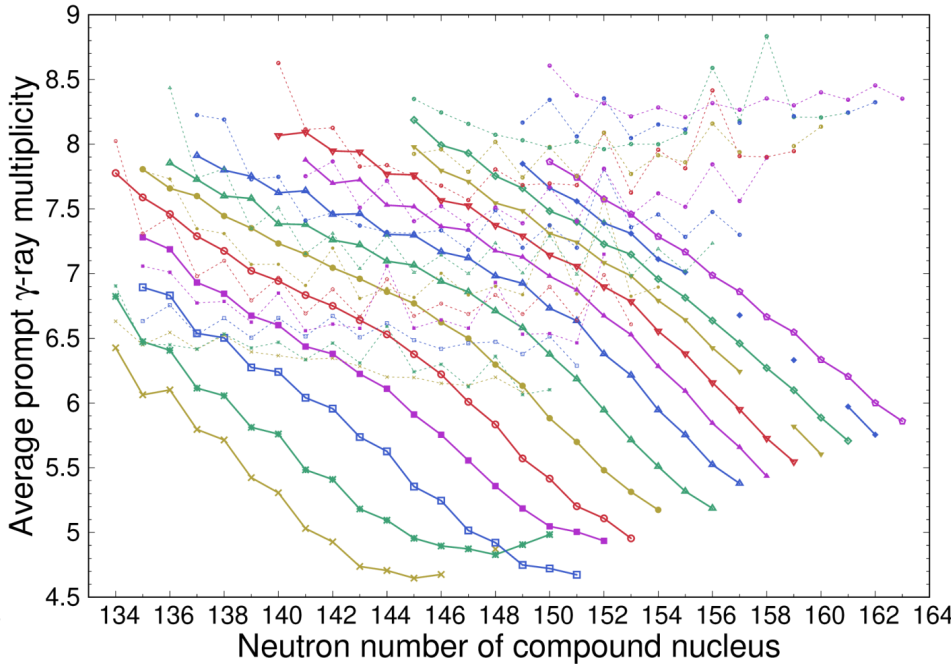
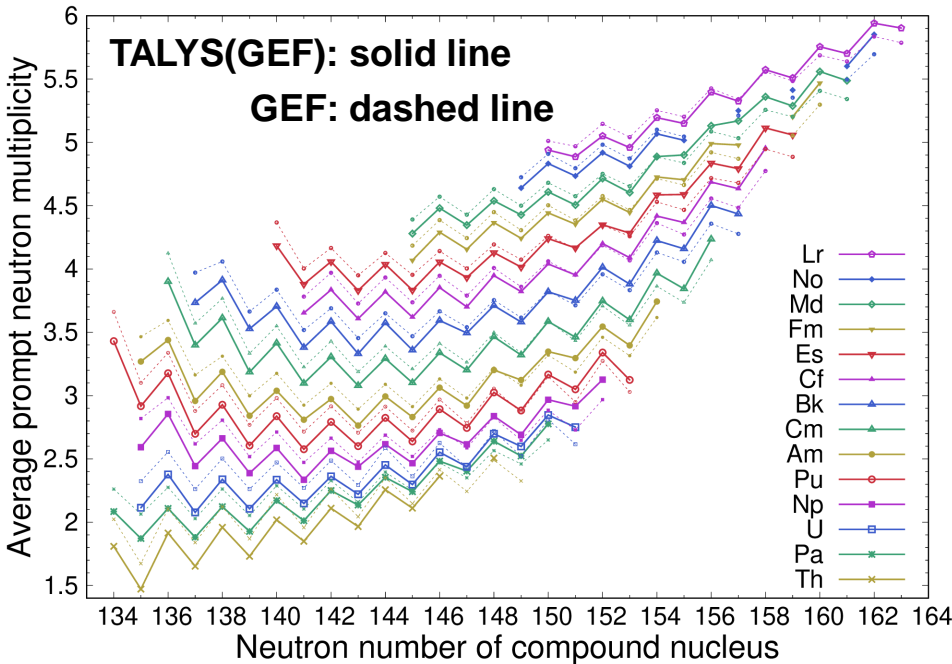
# $^{235}\text{U}(n,f)$ : fission product yield



- Both TALYS results reproduce the prominent peak at  $A = 138$  for heavy fragments and at  $A = 94, 100$  for light ones.
- The prominent peak at  $A = 134$  appears only in TALYS(GEF).

# Global study on actinides

We examined 243 selected actinide isotopes with 1 MeV incident neutron energy using:  $f^2 = 4$ ,  $f_s = 0.4$ , and  $N = 150$ .



**TALYS is now able to perform this kind of global calculation of statistical decay of primary fission fragments, ranging from very neutron deficient to neutron-rich nuclei.**

- $\bar{\nu}_n$  calculated by stand-alone GEF and TALYS(GEF) coincides with each other quite well including the zigzag pattern caused by the pairing effects.
- $\bar{\nu}_\gamma$  for the neutron-rich nuclei differs much between TALYS(GEF) and GEF, which reflect difference of the statistical decay calculation in these codes.



# Conclusions

- TALYS recently has been extended to perform Hauser-Feshbach statistical decay calculation with fission fragment distribution database generated by GEF, HF<sup>3</sup>D, SPY, and user's own data.
- TALYS has a limitation within the energy range up to first-chance fission. In the future, TALYS will be responsible for multi-chance fission.
- The optimal parameters are decided to prioritize to reproduce  $\bar{\nu}_n$  and the shape of PFNS.
- TALYS shows a decent agreement with the experimental and evaluated data of prompt neutron observables and independent fission product yield, especially.
- TALYS is now able to perform this kind of global calculation of statistical decay of primary fission fragments, ranging from very neutron deficient to neutron-rich nuclei.

# Acknowledgements

- Authors thank T. Kawano (Los Alamos National Laboratory) and K.-H. Schmidt for valuable discussions.
- The IAEA-NDS acknowledges the internship program “The nuclear Regulation Human Resource Development Program (ANSET: Advanced Nuclear 3S Education and Training)” entrusted to Tokyo Institute of Technology, Tokyo, Japan by the Nuclear Regulation Agency of Japan, for supporting this work.
- A. Al-Adili would like to acknowledge Liljewalch travel scholarships and Ingegerd Berghs stiftelse for their research grants.
- K. Fujio thanks C. Ishizuka and acknowledges her Grant-in-Aid for Scientific Research (B), MEXT, Japan, and by Japan Society for the Promotion of Science (JSPS) KAKENHI Grant Number 21H01856.

**Thank you for your kind attention!**

# Other important parameters

- **Optical model potentials**

Koning-Delaroche global optical model

- **Level density parameters**

Constant temperature model using the level density parameters and systematics from:

[A. J. Koning, S. Hilaire, S. Goriely, Nucl. Phys. A810, 13-76 \(2008\).](#)

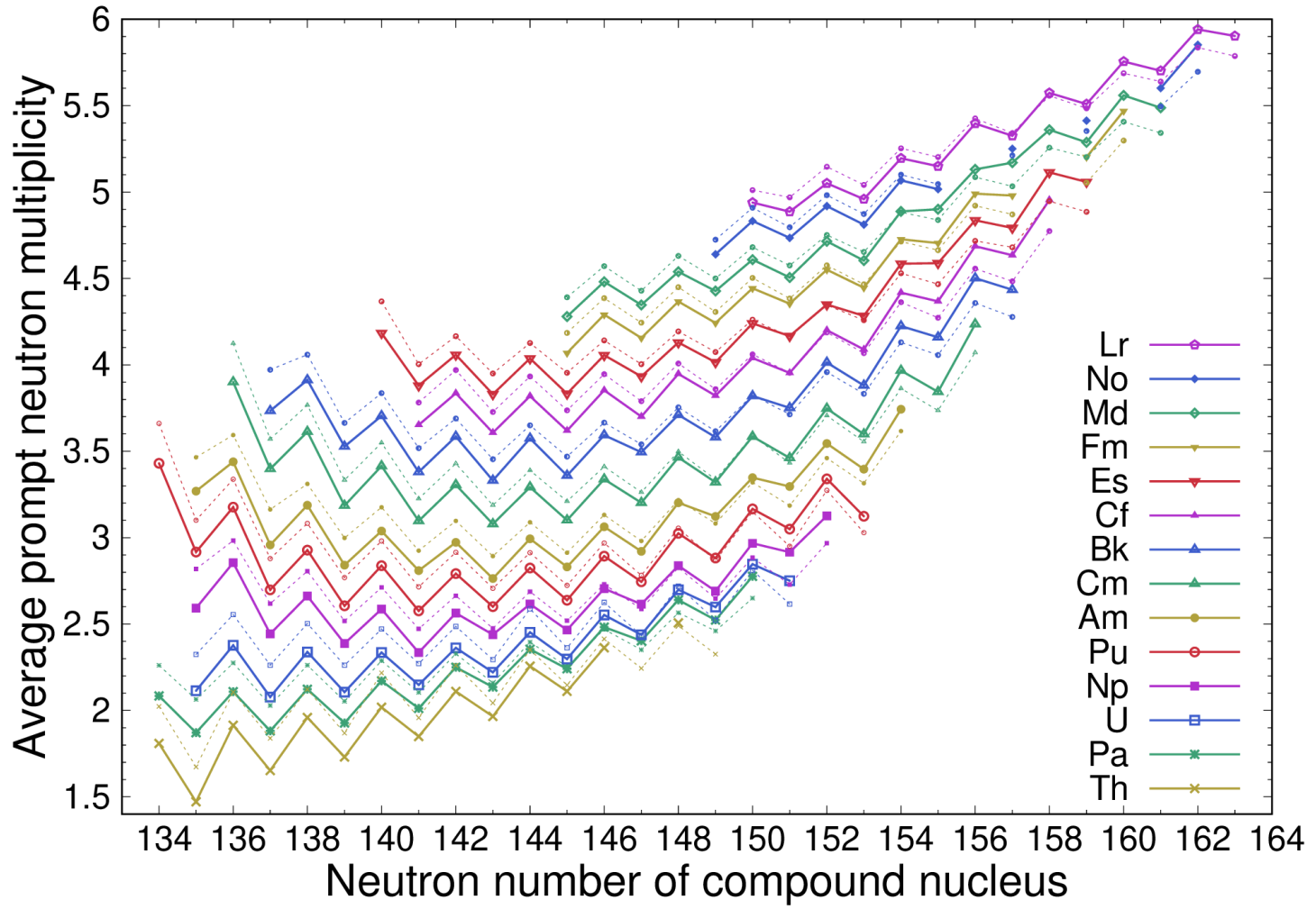
- **E1 and M1  $\gamma$ -ray strength function**

IAEA-CRP SMLO 2019 tables and IAEA GSF CRP 2018

- **Discrete level properties**

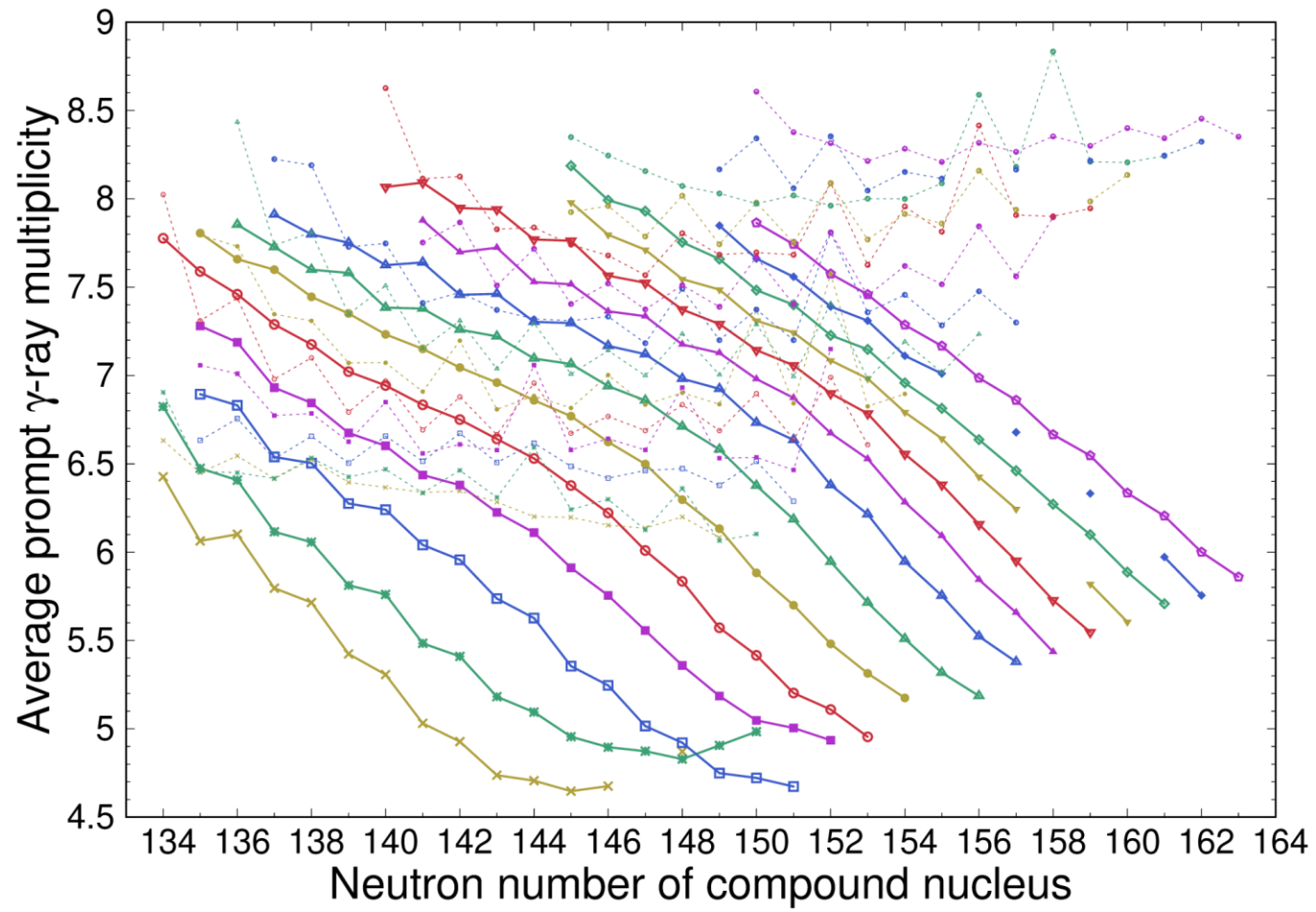
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# Global study on actinides



- The zigzag pattern is attributed to the difference in neutron separation energy of the compound nucleus.
- The even  $N$  compound nucleus gains more excitation energy due to the pairing of the captured incident neutron compared to an odd  $N$  compound nucleus.

# Global study on actinides



- $\bar{\nu}_\gamma$  is known to be sensitive to the angular momentum population in the fragments.
- Different treatments of the fragment angular momenta in GEF and TALYS could be one reason for such discrepancies.

# Sensitivity on scaling parameters

TALYS(HF <sup>3</sup> D)						
# of $N$	$f^2$	$f_s$	$\bar{\nu}_\gamma$	$\bar{\nu}_n$	$\langle \epsilon_\gamma \rangle [\text{MeV}]$	$\langle \epsilon_n \rangle [\text{MeV}]$
300	3	0.4	5.06	2.49	0.869	2.049
	3	0.6	5.76	2.47	0.818	1.945
	3	0.8	6.21	2.45	0.786	1.915
	3	1.0	6.50	2.44	0.764	1.907
	4	0.4	6.05	2.41	0.772	2.079
	4	0.6	6.92	2.40	0.728	1.941
	4	0.8	7.48	2.39	0.699	1.899
	4	1.0	7.85	2.38	0.677	1.887
	5	0.4	6.85	2.35	0.714	2.107
	5	0.6	7.90	2.34	0.675	1.938
	5	0.8	8.55	2.33	0.646	1.886
	5	1.0	8.96	2.33	0.625	1.869
	6	0.4	7.45	2.30	0.681	2.132
	6	0.6	8.66	2.29	0.646	1.935
	6	0.8	9.36	2.29	0.617	1.876
	6	1.0	9.85	2.29	0.593	1.855
ENDF-B/VIII.0			8.58	2.41	0.85	2.00
JEFF-3.3			8.74	2.41	0.81	

# Sensitivity on $N$

TALYS(HF<sup>3</sup>D)

# of $N$	$f^2$	$f_s$	$\bar{\nu}_\gamma$	$\bar{\nu}_n$	$\langle \epsilon_\gamma \rangle$ [MeV]	$\langle \epsilon_n \rangle$ [MeV]
50	4	0.4	3.46	2.49	1.192	1.990
100	4	0.4	4.87	2.43	0.923	2.045
200	4	0.4	5.78	2.41	0.804	2.072
300	4	0.4	6.05	2.41	0.772	2.079
ENDF-B/VIII.0			8.58	2.41	0.85	2.00
JEFF-3.3			8.74	2.41	0.81	



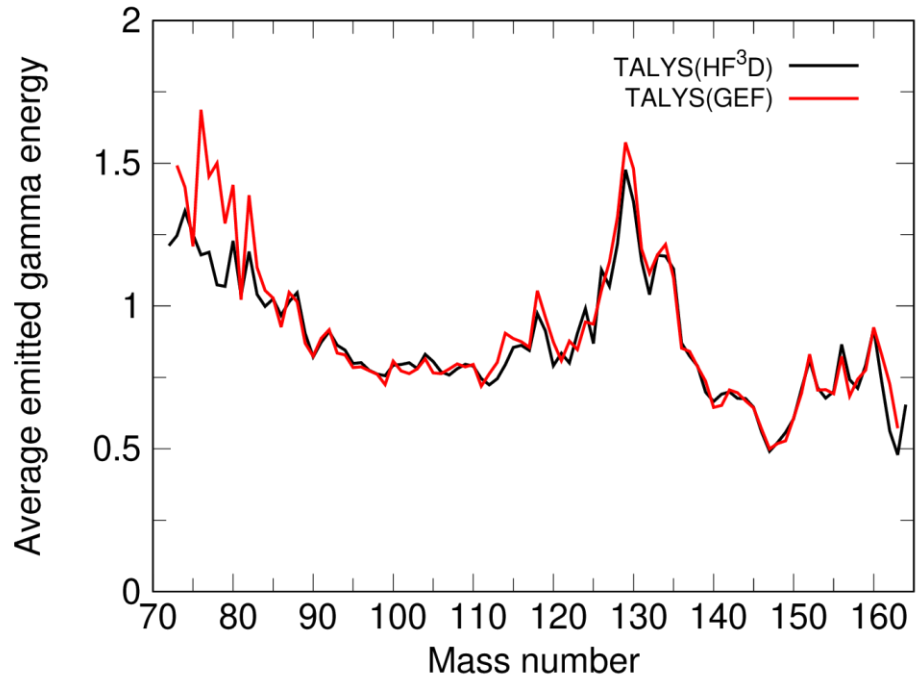
# $^{235}\text{U}(n,f)$ : multiplicities and average energy

• The optimal values:  $f^2 = 4$ ,  $f_s = 0.4$ , and  $N = 300$

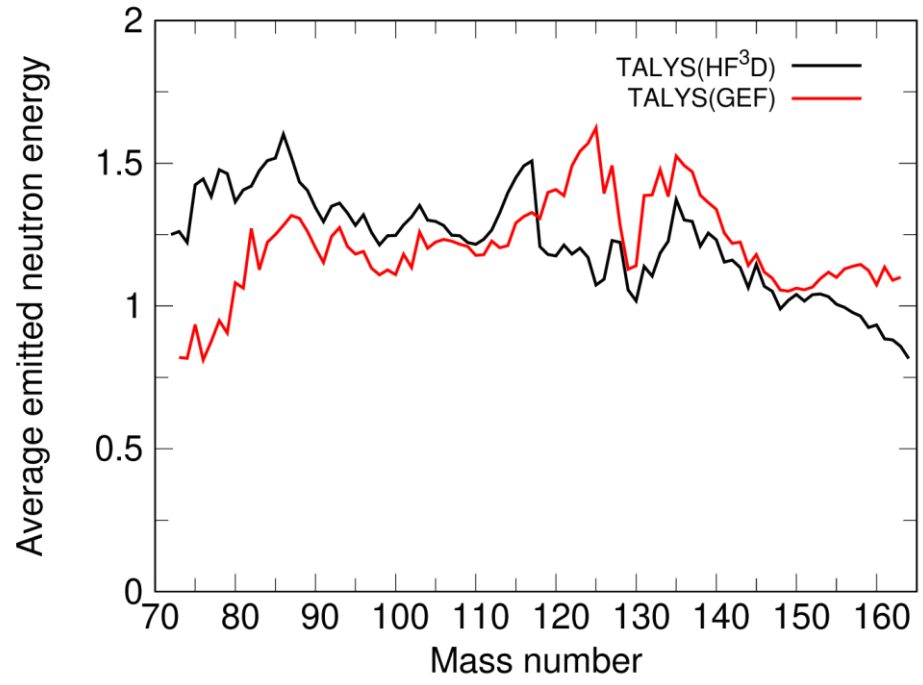
	$\bar{\nu}_\gamma$	$\bar{\nu}_n$	$\langle \epsilon_\gamma \rangle [\text{MeV}]$	$\langle \epsilon_n \rangle [\text{MeV}]$
TALYS(GEF)	6.13	2.30	0.761	1.991
TALYS(HF <sup>3</sup> D)	6.05	2.41	0.772	2.079
GEF	6.61	2.42	0.962	1.997
Oberstedt et al.(2013)	8.19±0.11		0.85±0.02	
Verbinski et al.(1973)	6.70±0.30		0.97±0.05	
Pleasanton et al.(1972)	6.51±0.30		0.99±0.07	
Peelle et al.(1971)	7.45±0.35		0.96±0.05	
ENDF-B/VIII.0	8.58	2.41	0.85	2.00
JEFF-3.3	8.74	2.41	0.81	

# $^{235}\text{U}(n,f)$ : average energies

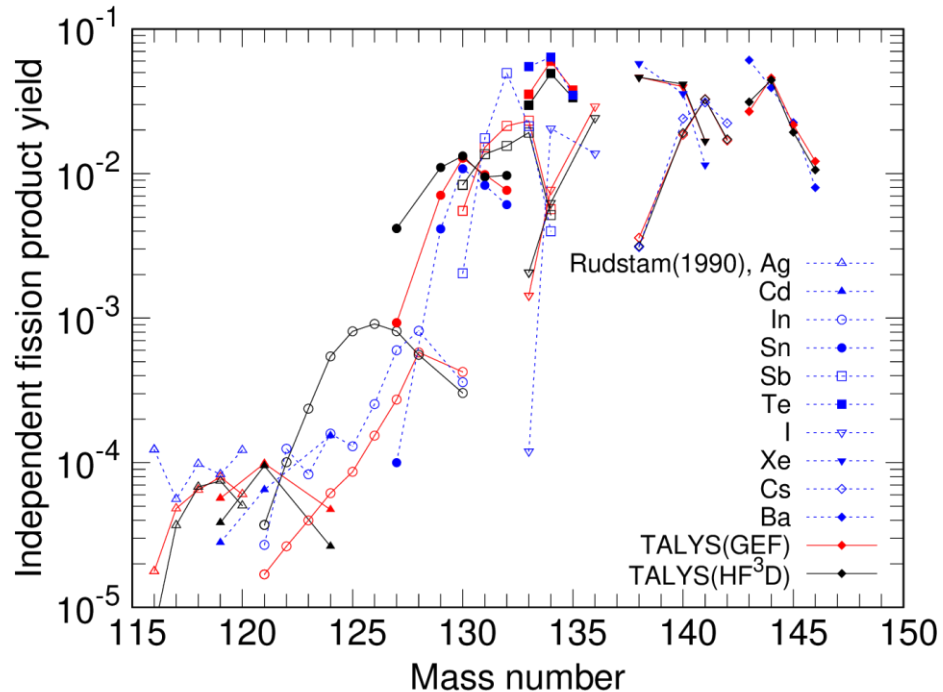
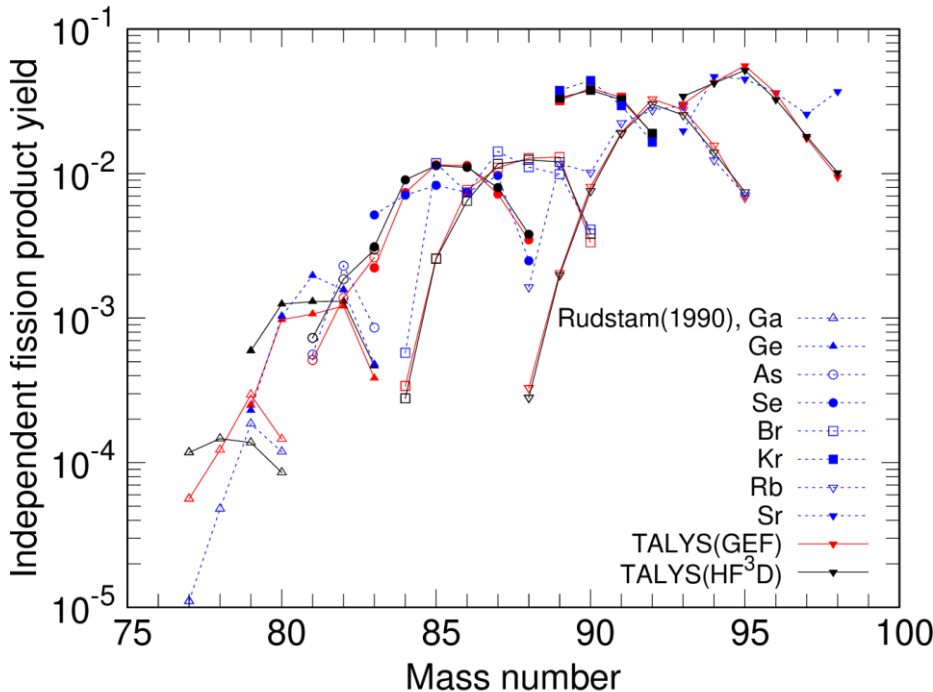
## Average emitted $\gamma$ -ray energy



## Average emitted neutron energy

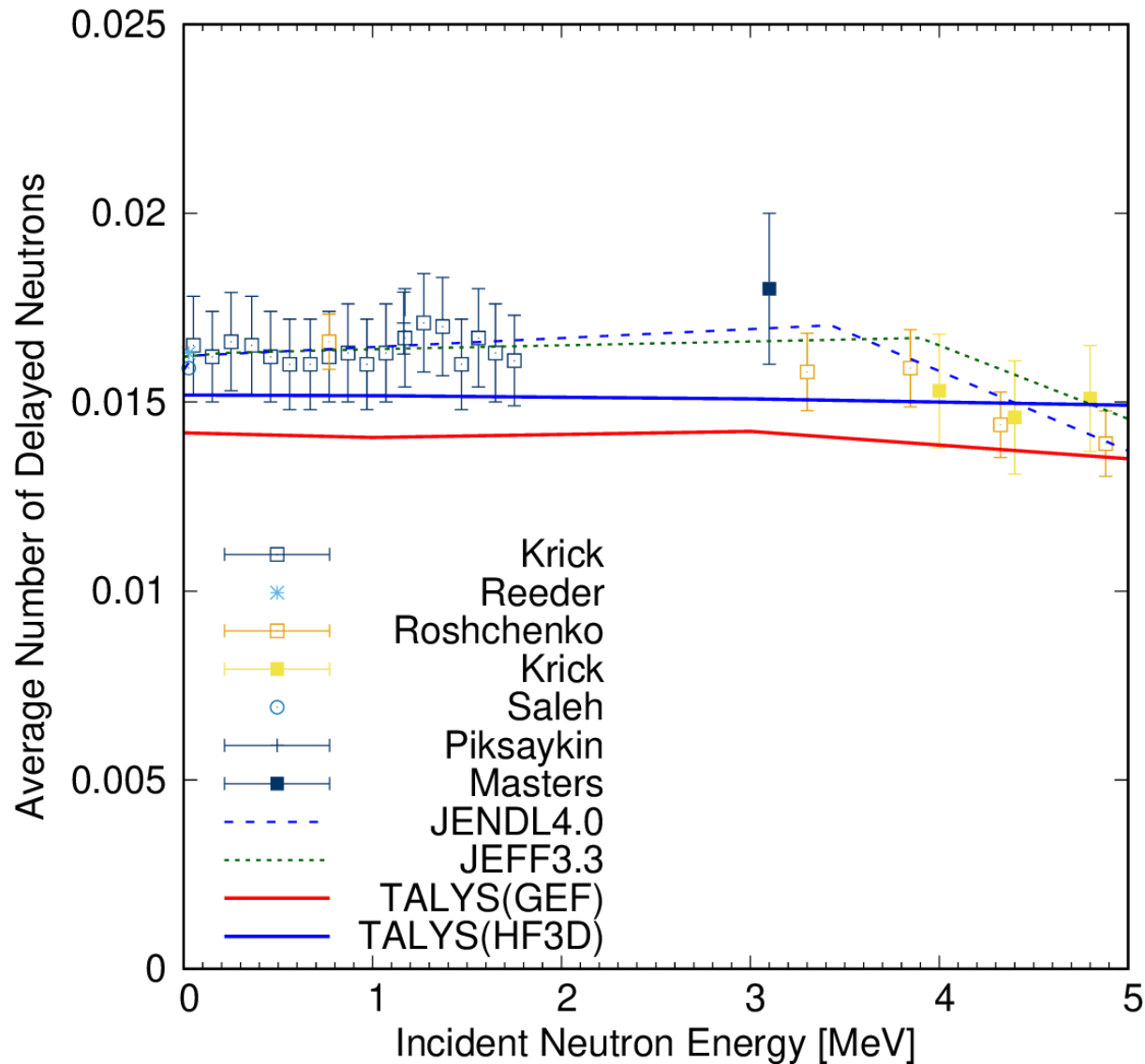


# $^{235}\text{U}(n,f)$ : fission product yield

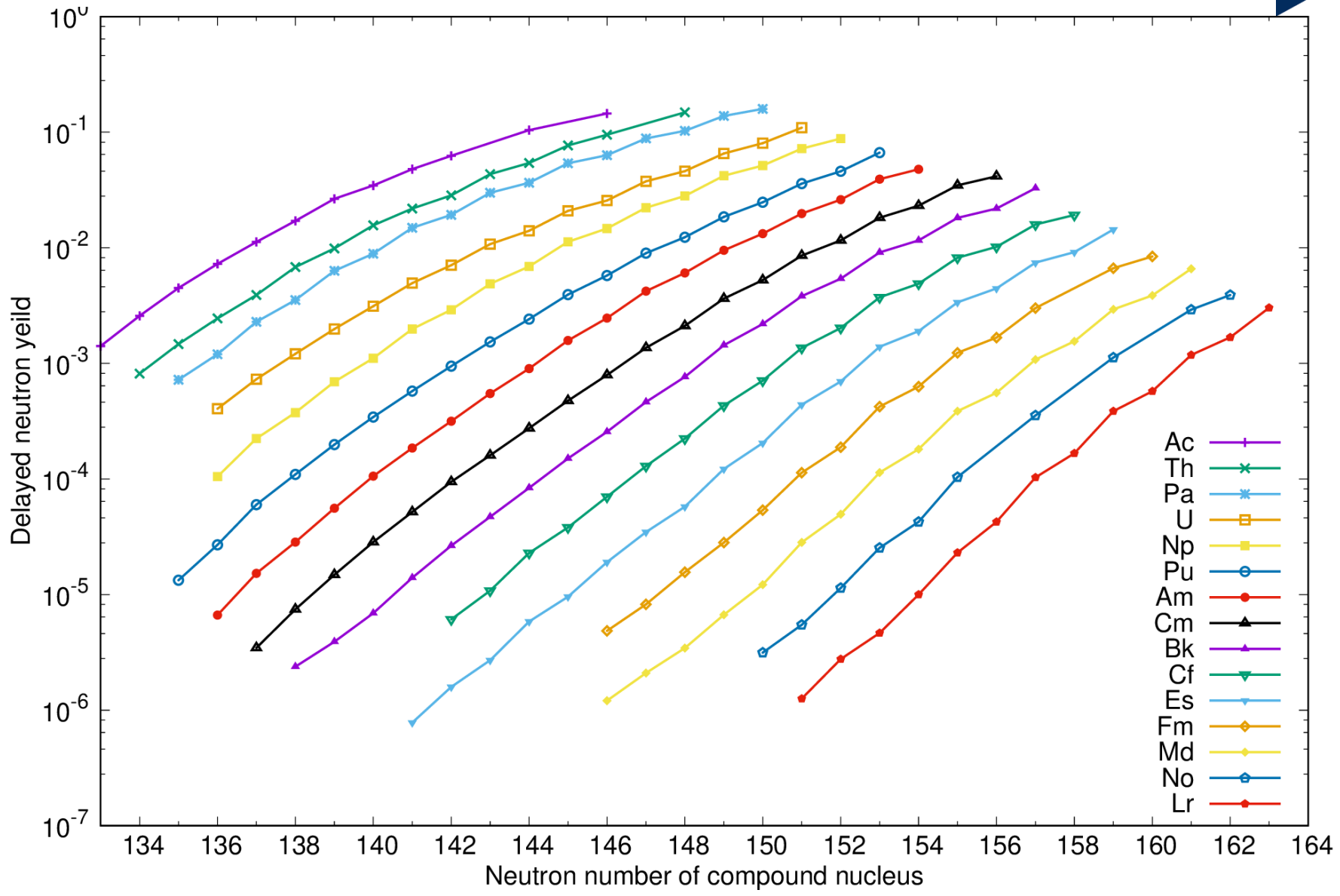


- TALYS results are roughly consistent with the tendency shown in the experimental data.

# $^{235}\text{U}(n,f)$ : average number of delayed neutron



# Global study of delayed neutron yield

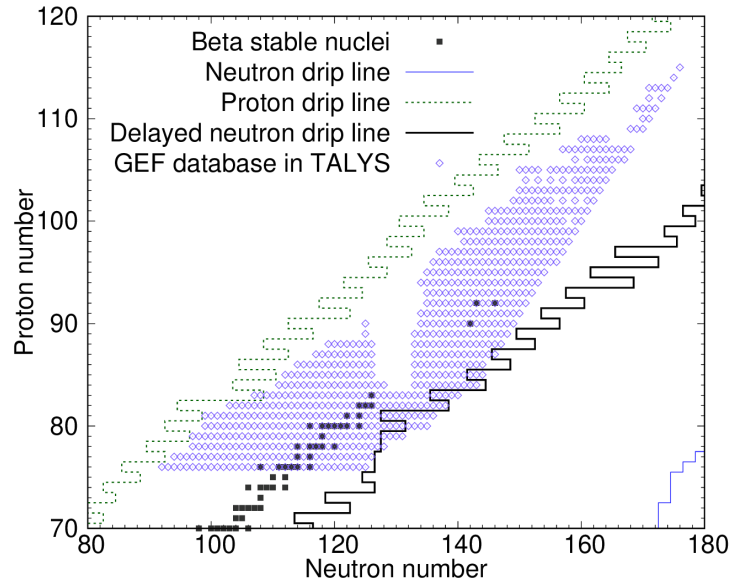
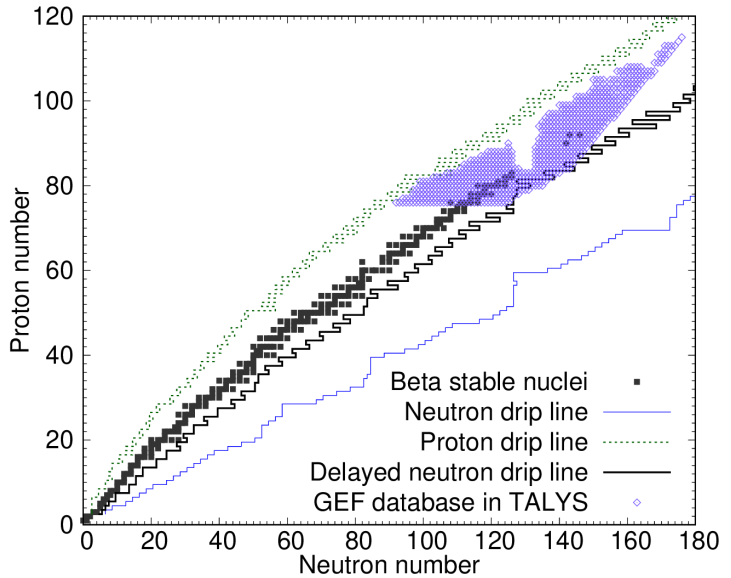


# Fission fragment database in TALYS

**GEF** Monte Carlo-based phenomenological fission model

[K. -H. Schmidt, B. Jurado, C. Amouroux, C. Schmitt, Nuclear Data Sheets, 131, 107-221 \(2016\).](#)

- For 737 fissioning nuclei ranging from  $_{76}\text{Os}$  to  $_{115}\text{Mc}$



**HF<sup>3</sup>D** Designed with a fully deterministic technique with fitting functions

[S. Okumura, T. Kawano, P. Jaffke, P. Talou, and S. Chiba, JNST, 55\(9\), 1009-1023 \(2018\).](#)

- For neutron-induced fission of  $^{235}\text{U}$ ,  $^{238}\text{U}$ , and  $^{239}\text{Pu}$

**SPY** Obtained from a statistical scission point model using microscopic calculation

[J. -F. Lemaître, S. Goriely, S. Hilaire, and J.-L. Sida, Phys. Rev. C99, 034612 \(2019\).](#)

Arbitrary fission fragment data provided by users