

Preliminary results of $K_{\mu 4}^{00}$ decay first observation (conference slides)

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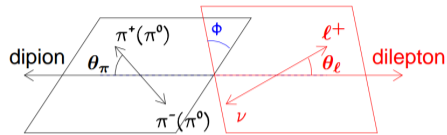
NA48/2 collaboration meeting, 25.02.2022

Outline

- 1 Theoretical framework
- 2 NA48/2 setup
- 3 Selection
- 4 Acceptance
- 5 Residual background
- 6 Signal extraction
- 7 Decay model-related systematic uncertainty
- 8 Result
- 9 Conclusion
- 10 Spare slides

$K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$ ($K_{\mu 4}^{00}$) theoretical framework

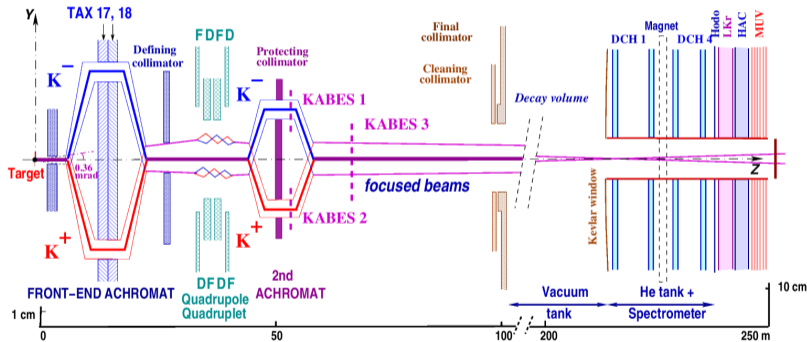
In general, K_{l4} decay width depends on four form-factors: F, G, R, H and is a function of five Cabibbo-Maksymowicz variables: S_π (dipion mass squared), S_l (dilepton mass squared) and angles θ_π (in the dipion frame), θ_l (in the dilepton frame), ϕ .



- For $K_{\mu 4}^{00}$, in s-wave approximation for $\pi^0 \pi^0$, there are no dependence on $\cos \theta_\pi, \phi$, and only $F \neq 0, R \neq 0$.
- Unlike $K_{e 4}^{00}$ case, R plays some role due to μ mass;
- Motivation for $K_{\mu 4}^{00}$ measurement: first observation, ChPT test, check of R presence, potential study of $\pi\pi$ rescattering effects in the $F(S_\pi)$.
- According to lepton universality, experimental $F(S_\pi, S_l)$ parameterization from $K_{e 4}^{00}$ [J.R. Batley et al. JHEP 08 (2014) 159] may be used for $K_{\mu 4}^{00}$.
- The only available source of $R(S_\pi, S_l)$ is ChPT calculation [J.Bijnens, G.Colangelo, J.Gasser, Nucl.Phys.B 427, 427 (1994)].
- In our MC we use experimental $F(K_{e 4}^{00})$ scaled down by about 18% (close to the 1-loop F), theoretical $R(1\text{-loop})$ and zero phase between F and R .
- Other versions were tested for systematics.

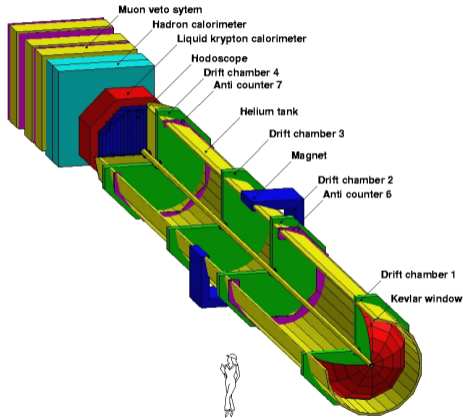
NA48/2 beamline (CERN SPS, 2003-2004)

NA48/2 main goal were $K_{3\pi}$ charge asymmetry studies; additional rare decays program.



- Two charged beams, 6% of K^\pm , from 400 GeV/c protons on target. Beams central momentum 60 GeV/c, momentum band $\pm 3.8\%$ (*rms*), width ≈ 1 cm.
- KABES (Kaon Beam Spectrometer), spatial resolution $800 \mu\text{m}$, momentum resolution 1% and time resolution 600 ps.

NA48/2 setup (CERN SPS, 2003-2004)



- Magnetic spectrometer (drift chambers DCH1–DCH4): Spatial resolution $\sim 90 \mu\text{m}$ per chamber, momentum resolution $\sigma_p/p = (1.02 \oplus 0.044 \cdot p)\%$ (momentum p in GeV/c).
- Scintillator hodoscope (HOD), time resolution ~ 150 ps.
- Liquid Krypton EM calorimeter (LKr). High granularity, quasi-homogeneous
 $\sigma_x = \sigma_y = (0.42/\sqrt{E} \oplus 0.06)$ cm.
 $\sigma_E/E = (3.2/\sqrt{E} \oplus 9.0/E \oplus 0.42)\%$ (energy E in GeV).
- Hadronic calorimeter, muon system MUV.

Selection

NA48/2 data collected in 2003-2004.

Signal (S) $K_{\mu 4}$ is $K^{\pm} \rightarrow \pi^0 \pi^0 \mu^{\pm} \nu$. Normalization (N) $K_{3\pi}$ is $K^{\pm} \rightarrow \pi^{\pm} \pi^0 \pi^0$.

A common trigger chain: L1 trigger using HOD and LKr, followed by L2 trigger using DCH for online momentum calculation.

Combination: 4 isolated photons in time with a KABES track and with a DCH track.

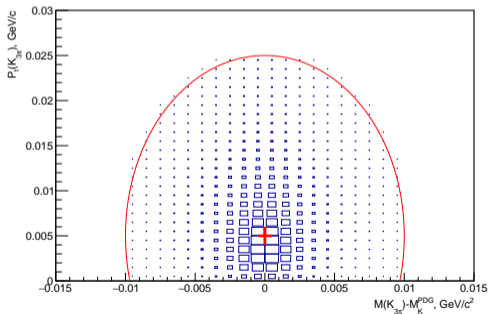


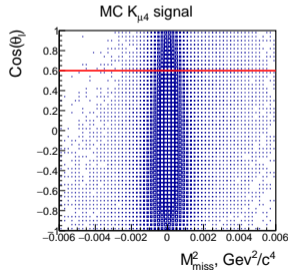
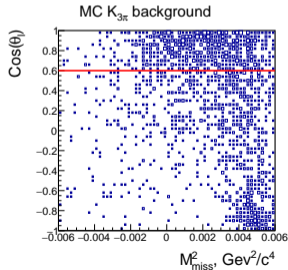
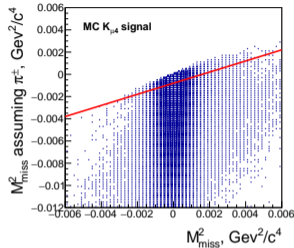
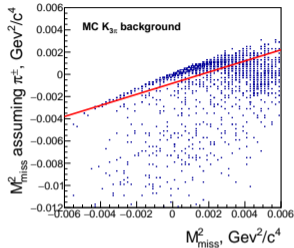
Figure: $K_{3\pi}$ kinematic selection area

Normalization $K^{\pm} \rightarrow \pi^{\pm} \pi^0 \pi^0$ channel selection
No missing momentum is expected.

$K_{3\pi}$ kinematic selection ellipse:

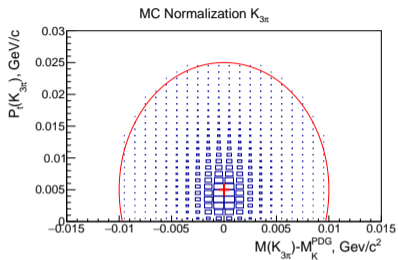
- center:
 - $M(K_{3\pi}) = M_K^{PDG}$;
 - $P_t = 5 \text{ MeV}/c$;
- semi-axes:
 - $\Delta M(K_{3\pi}) = 10 \text{ MeV}/c^2$;
 - $\Delta P_t = 20 \text{ MeV}/c$.
- 72.99×10^6 $K_{3\pi}$ selected data events.

$K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$ signal events selection

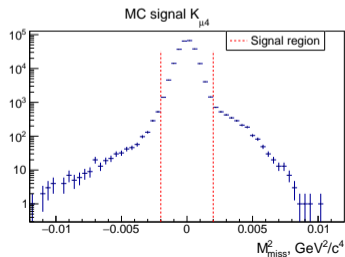


- Off the $K_{3\pi}$ kinematic ellipse;
- DCH track has associated MUV response;
- M_{miss}^2 assuming $\pi^\pm < 0.5M_{miss}^2 - 0.0008 \text{ GeV}^2/c^4$;
- $\cos(\Theta_l) < 0.6$;
- $S_l > 0.03 \text{ GeV}^2/c^4$ (to reject $\pi^\pm \rightarrow \mu^\pm \nu$).
- 3718 $K_{\mu 4}$ data candidates selected (any M_{miss}^2).

Acceptances



- Acceptance of the $K^\pm \rightarrow \pi^0 \pi^0 \pi^\pm$ normalization channel is $(4.477 \pm 0.002)\%$,



- M_{miss}^2 signal region $[-0.002, 0.002] \text{ GeV}^2/c^4$ contains 98.2% of selected MC events.
- Acceptance of the $K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$ signal is $(0.651 \pm 0.001)\%$,
- Acceptance of the $K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$ signal for the restricted phase space $S_i^{\text{true}} > 0.03 \text{ GeV}^2/c^4$ is $(3.453 \pm 0.007)\%$.

Residual background

Main background: $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ followed by $\pi^\pm \rightarrow \mu\nu$ before MUV with a probability $\approx 10\%$ for $P(\pi^\pm) \approx 10 \text{ GeV}/c$.

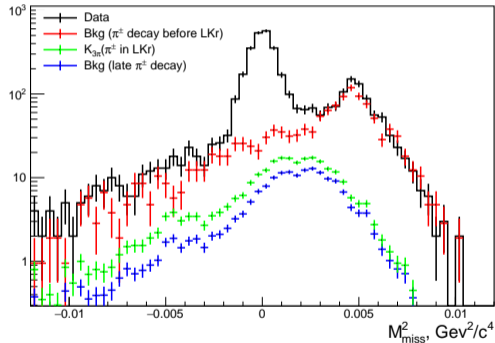
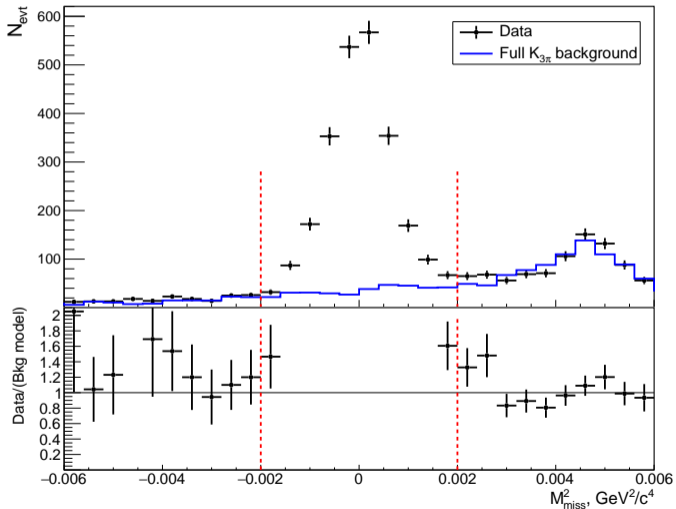


Figure: Simulated background with π^\pm decay before LKr (normalized from the kaon flux). Background models for $K_{3\pi}$ with late π^\pm decay are with arbitrary normalization.

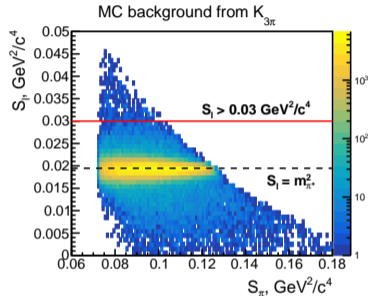
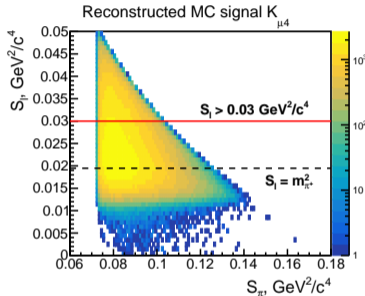
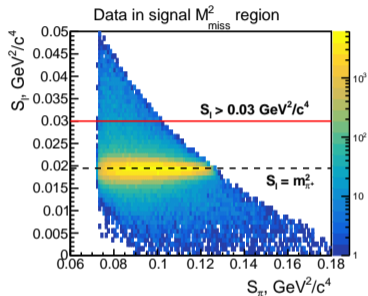
- $K_{3\pi}$ background with π^\pm decay before LKr, from MC simulation.
- $K_{3\pi}$ background with late π^\pm decay or muon emission in a late hadron shower:
 - A data sample is selected as $K_{\mu 4}^{00}$ without MUV hits requirement and with $0.3 < E_{LKr}/P_{DCH} < 0.8$ (π^\pm shower in LKr).
 - It corresponds to random MUV response (used for systematic uncertainty tests).
 - $K_{3\pi}$ background with late π^\pm decay: a weight $WP \propto \frac{1}{P_{DCH}}$ for each event to take into account the pion decay probability.

$K_{\mu 4}^{00}$ signal extraction fit



- 2437 candidates are found in the signal region.
- Fit in the M_{miss}^2 interval $[-0.003, 0.006]$ GeV^2/c^4 , ignoring the signal region to decrease sensitivity to the imperfect MC resolution.
- Data fit by a linear combination of the two background components and the MC signal tails.
- $354 \pm 33_{stat} \pm 62_{syst}$ background events.
- The background-related systematics are determined by varying the way the background is estimated.

Signal versus S_{π^+}, S_l



- The branching ratio is measured for the restricted phase space $S_l^{true} > 0.03 \text{ GeV}^2/c^4$.
- Extrapolation to the full phase space depends on the theory.

Signal versus S_{π}, S_l

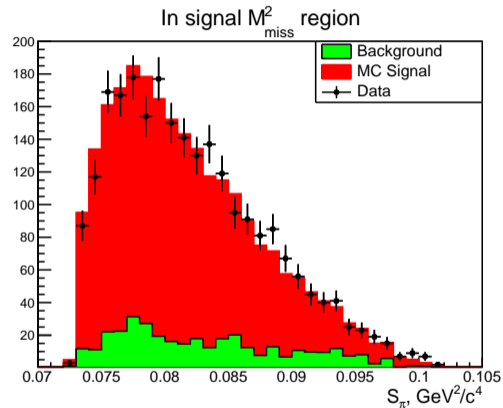
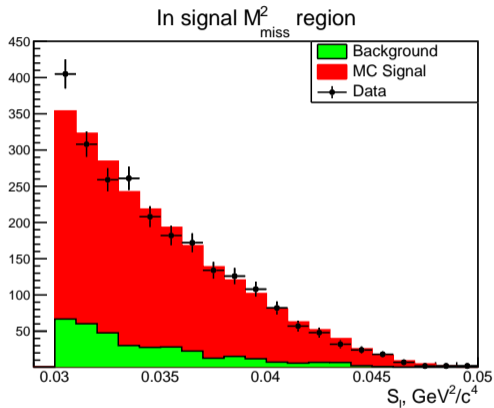


Figure: 1D projections comparison for $S_l > 0.03 \text{ GeV}^2/c^4$

Acceptance variations due to $K_{\mu 4}^{00}$ generator modifications

J.R. Batley et al. JHEP 08 (2014) 159:

$$F(K_{e4}) = (1 + aq^2 + bq^4 + c \cdot S_I/4m_{\pi^+}^2), q^2 \geq 0$$

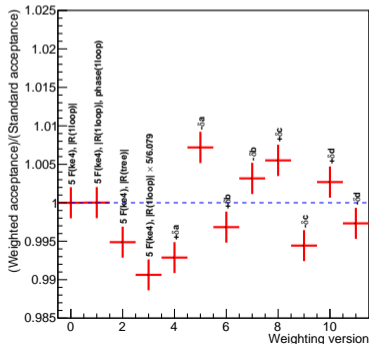
$$F(K_{e4}) = (1 + d\sqrt{|q^2/(1 + q^2)|} + c \cdot S_I/4m_{\pi^+}^2), q^2 \leq 0,$$

$$\text{where } q^2 = S_{\pi}/4m_{\pi^+}^2 - 1.$$

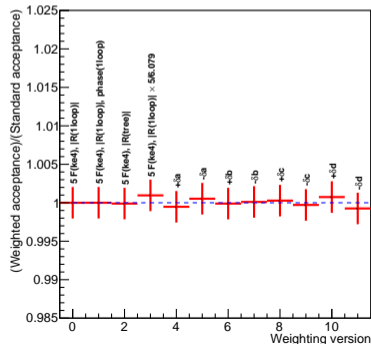
The absolute normalization is also measured $F = f \times F(K_{e4})$,
 $f = 6.079 \pm 0.055$.

- Decay generator was modified by MC events weighting.
- The acceptance spread is taken as systematics.

Full acceptance relative variations



($S_I > 0.03$) acceptance relative variations



Preliminary result: Ingredients

$$BR(K_{\mu 4}^{00}) = \frac{N_S}{N_N} \cdot \frac{A_N}{A_S} \cdot K_{trig} \cdot BR(K_{3\pi}^{00}). \quad (1)$$

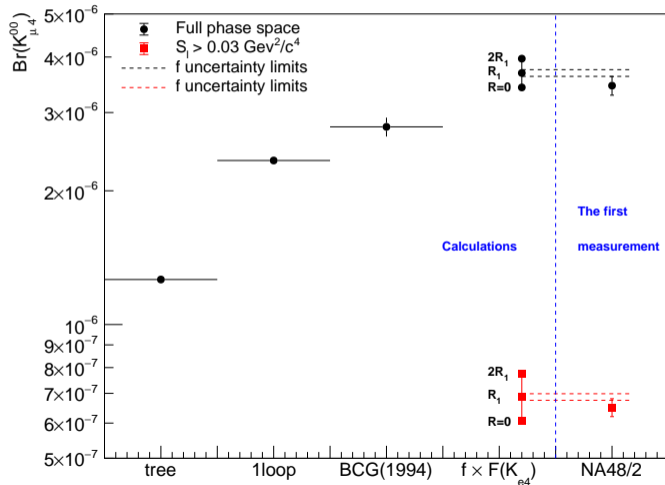
- Extracted signal $N_S = N_{Sign. \text{ cand.}} - N_{Bkg} = 2437 - (354 \pm 33_{stat}) = 2083 \pm 59_{stat}$ events;
- Signal/Background is $5.89 \pm 0.66_{stat}$;
- Number of normalization events $N_N = 72.99 \times 10^6$;
- Normalization acceptance $A_N = (4.477 \pm 0.002)\%$;
- Signal acceptance for the restricted phase space $A_S^r = (3.453 \pm 0.007)\%$;
- Signal acceptance for the full phase space $A_S = (0.651 \pm 0.001)\%$;
- Trigger correction (extracted with control triggers)
 $K_{trig} = K_{CHT} \cdot K_{NUT} = (0.998 \pm 0.002) \cdot (1.0007 \pm 0.0007) = 0.999 \pm 0.002$;
- PDG $BR(K_{3\pi}^{00}) = (1.760 \pm 0.023)\%$;

Preliminary result: Central values and errors budget

	Full phase space		$S_I > 0.03 \text{ GeV}^2/c^4$	
$Br(K_{\mu 4})$ central value	3.45×10^{-6}		0.651×10^{-6}	
	δBR	$\delta BR/BR$	δBR	$\delta BR/BR$
Data stat. error	0.10×10^{-6}	2.85%	0.019×10^{-6}	2.85%
MC stat. error	0.01×10^{-6}	0.21%	0.001×10^{-6}	0.21%
Trigger	0.01×10^{-6}	0.18%	0.001×10^{-6}	0.18%
Background	0.10×10^{-6}	2.96%	0.019×10^{-6}	2.96%
Accidentals	0.01×10^{-6}	0.32%	0.002×10^{-6}	0.32%
MUV inefficiency	0.06×10^{-6}	1.65%	0.011×10^{-6}	1.65%
Form Factor modelling	0.05×10^{-6}	1.37%	0.001×10^{-6}	0.14%
$BR(K_{3\pi})$ error (external)	0.05×10^{-6}	1.31%	0.009×10^{-6}	1.31%
Total error	0.17×10^{-6}	4.83%	0.030×10^{-6}	4.64%

- Accidentals obtained from side bands of time distributions;
- MUV inefficiency uncertainty taken as full inefficiency effect.

Preliminary result: Comparison to theory



Theory:

- J. Bijnens, G. Colangelo, J. Gasser, Nucl. Phys. B, 427 (1994) 427:
 - Tree approximation;
 - 1-loop;
 - BCG(1994): 'beyond 1-loop' with measured F from [Rosselet etc. Phys. Rev. D 15 (1977) 574].
- Re-calculated now:
 - $F(K_{e4})$ from NA48/2 (2015);
 - $R_1 = R(1loop)$;
 - 1-loop (F,R) phase;
 - 2020 PDG constants.

- A first observation and branching fraction measurement of $K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$ decay is performed by NA48/2 experiment at SPS in CERN;
- We observe 2437 signal candidates with an estimated background of $354 \pm 33_{stat} \pm 62_{syst}$ events, Signal/Background ratio is $5.9 \pm 1.4_{tot}$
- Preliminary result for restricted phase space ($S_I > 0.03$) is
$$BR(K_{\mu 4}^{00}, S_I > 0.03) = (0.65 \pm 0.019_{stat} \pm 0.024_{syst}) \times 10^{-6} = (0.65 \pm 0.03) \times 10^{-6};$$
- Preliminary full phase space result is
$$BR(K_{\mu 4}^{00}) = (3.4 \pm 0.10_{stat} \pm 0.13_{syst}) \times 10^{-6} = (3.4 \pm 0.2) \times 10^{-6}.$$
- The results are consistent with a contribution of the R form factor, as computed at 1-loop ChPT.

SPARE SLIDES

Theoretical framework: decay width

$K_{\mu 4}^{00}$ matrix element is [J.Bijnens, G.Colangelo, J.Gasser, Nucl.Phys.B 427, 427 (1994)]:

$$T = \frac{G_f}{\sqrt{2}} \cdot V_{us}^* \cdot \underbrace{\bar{u}(p_\nu) \cdot \gamma_\mu \cdot (1 - \gamma_5) \cdot v(p_l)}_{\text{Lepton part}} \cdot \underbrace{(V_\mu - A_\mu)}_{\text{Hadron part}} \quad (2)$$

where

$$V_\mu = \frac{-H}{M_K^3} \varepsilon_{\mu,\nu,\rho,\sigma} L^\nu P^\rho Q^\sigma, \quad A_\mu = -i \frac{1}{M_K} [P_\mu F + Q_\mu G + L_\mu R], \quad (3)$$

$\varepsilon_{0,1,2,3}=1$ and the four-momenta are defined as $P = p_1 + p_2$, $Q = p_1 - p_2$ (p_1 and p_2 are the two pion momenta) and $L = p_l + p_\nu$. The form factors F, G, R, H are analytic functions of the decay kinematic variables.

In general, decay width is a function of five Cabibbo-Maksymowicz variables. But for $K_{\mu 4}^{00}$, in s-wave approximation for $\pi^0 \pi^0$, there are no dependence on $\cos \theta_\pi, \phi$, and so $G = 0, H = 0$:

$$d\Gamma_3 = \frac{G_f^2 |V_{us}|^2 (1 - z_l)^2 \sigma_\pi X}{2^{11} \pi^5 M_K^5} (I_1 + I_2 (2(\cos \theta_l)^2 - 1) + I_6 \cos \theta_l) dS_\pi dS_l d \cos \theta_l. \quad (4)$$

where $I_1 = \frac{1}{4} \{ (1 + z_l) |F_1|^2 + 2z_l |F_4|^2 \}$, $I_2 = -\frac{1}{4} (1 - z_l) |F_1|^2$, $I_6 = z_l \text{Re}(F_1^* F_4)$, $F_1 = X \cdot F$, $F_4 = -(PL)F - S_l R$, $z_l = \frac{m_l^2}{S_l}$,

$\sigma_\pi = \sqrt{1 - \frac{4M_\pi^2}{s_\pi}}$, $X = \frac{1}{2} \sqrt{\lambda(M_K^2, S_\pi, S_l)}$, $\lambda(x, y, z) = x^2 + y^2 + z^2 - 2(xy + xz + yz)$.

Only three variables ($S_\pi, S_l, \cos \theta_l$) and two form factors (F, R) matter.

Theoretical framework: form factors

According to lepton universality, $F_s(S_\pi, S_l)$ s-wave form factor measured in K_{e4}^{00} may be used for $K_{\mu 4}^{00}$ decay MC simulation, while the only available source of $R_s(S_\pi, S_l)$ is ChPT calculation.

- We have ChPT 1-loop generator code [BCG(1994)] with $F(1loop), R(1loop)$.
- For the predictions beyond 1-loop (central value), BCG(1994) used $F = 5.59(1 + 0.08q^2)$, where $q^2 = S_\pi/4m_{\pi^+}^2 - 1$ [L.Rosselet et al., Phys. Rev. D15 (1977) 574].
- Additionally, from the NA48/2 K_{e4} analysis [J.R. Batley et al. JHEP 08 (2014) 159] we have the best measurement of F shape:

$$\begin{aligned} F(K_{e4}) &= (1 + aq^2 + bq^4 + c \cdot S_l/4m_{\pi^+}^2), q^2 \geq 0 \\ F(K_{e4}) &= (1 + d\sqrt{|q^2/(1 + q^2)|} + c \cdot S_l/4m_{\pi^+}^2), q^2 \leq 0, \end{aligned} \quad (5)$$

where $q^2 = S_\pi/4m_{\pi^+}^2 - 1$, $a = 0.149 \pm 0.033 \pm 0.014$, $b = -0.070 \pm 0.039 \pm 0.013$, $c = 0.113 \pm 0.022 \pm 0.007$, $d = -0.256 \pm 0.049 \pm 0.016$.

For systematic studies we take into account the combined (both statistical and systematic) uncertainties of these parameters. Moreover, the absolute value was also measured $F = 6.079F(K_{e4})$.

Nevertheless, for MC production we have used $5F(K_{e4})$ (to be close to the 1-loop F), $R(1loop)$, $H = G = 0$ and zero phase between F and R .

Common selection criteria for signal and normalization channels

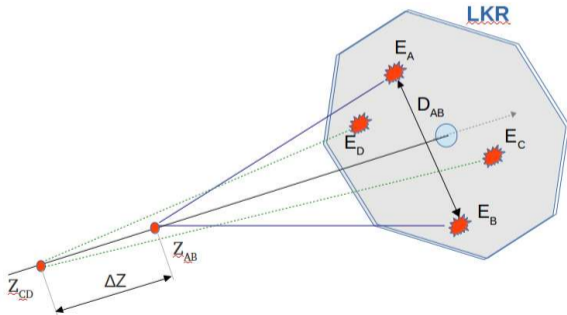
NA48/2 data collected in 2003-2004.

Signal (S) $K_{\mu 4}$ is $K^{\pm} \rightarrow \mu^{\pm} \nu \pi^0 \pi^0$. Normalization (N) $K_{3\pi}$ is $K^{\pm} \rightarrow \pi^{\pm} \pi^0 \pi^0$.

A common trigger chain: L1 trigger using HOD and LKr, followed by L2 trigger using DCH (MBX).

Combination: 4 isolated photons in time with a KABES track and with a DCH track.

Vertex 'charged position' Z_c is taken from the space matching of the KABES track and the DCH track.



- P_{KABES} in [54,67] GeV/c; P_{DCH} in [5,35] GeV/c;
 $E_{\gamma} > 3$ GeV;
- For π_i^0 ($i = 1, 2$):
 - $Z_1 = Z_{AB} = D_{AB} \sqrt{E_A E_B} / m_{\pi^0}$;
 - $Z_2 = Z_{CD} = D_{CD} \sqrt{E_C E_D} / m_{\pi^0}$.
- $-1600 \text{ cm} < Z_n = Z_{LKR} - (Z_1 + Z_2)/2 < 9000 \text{ cm}$
(Z_n is the vertex 'neutral position');
- Flunge cut: $R_{\gamma} @ DCH1 > 11 \text{ cm}$, assuming
 $Z_n + 400 \text{ cm}$ decay position;
- $|Z_1 - Z_2| < 500 \text{ cm}$; $|Z_n - Z_c| < 600 \text{ cm}$.

The best combination of each mode in the event has the most compatible Z_1, Z_2, Z_c .

Normalization $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ channel selection

No missing momentum is expected.

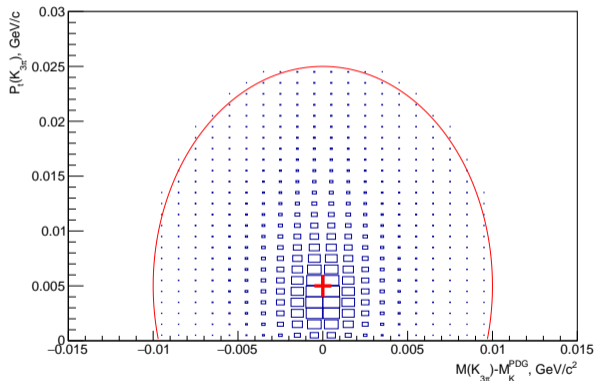


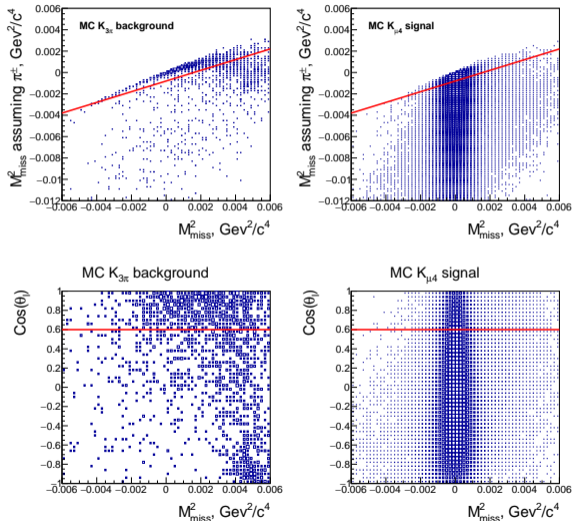
Figure: $K_{3\pi}$ kinematic selection area

$K_{3\pi}$ kinematic selection ellipse:

- center:
 - $M(K_{3\pi}) = M_K^{PDG}$;
 - $P_t = 5 \text{ MeV}/c$;
- semi-axes:
 - $\Delta M(K_{3\pi}) = 10 \text{ MeV}/c^2$;
 - $\Delta P_t = 20 \text{ MeV}/c$.

Normalization channel selection output is $N_N = 72.99 \times 10^6$ $K_{3\pi}$ reconstructed events.

Signal $K^\pm \rightarrow \mu^\pm \nu \pi^0 \pi^0$ channel selection



- Off the $K_{3\pi}$ kinematic ellipse;
- Good $K - \mu$ tracks matching;
- $P_{DCH} > 10 \text{ GeV}/c$ for MUV efficiency;
- DCH track is in restricted MUV acceptance with high MUV efficiency;
- DCH track has associated MUV hits in the first two planes;
- $M_{\text{miss}}^2(\text{assuming } \pi^\pm) < 0.5M_{\text{miss}}^2 - 0.0008 \text{ GeV}^2/c^4$;
- $\cos(\Theta_I) < 0.6$;
- $S_I > 0.03 \text{ GeV}^2/c^4$ to reject $\pi^\pm \rightarrow \mu^\pm \nu$.

3718 $K_{\mu 4}$ candidates selected (any M_{miss}^2).

$K_{\mu 4}^{00}$ signal extraction procedure

We fit $K_{\mu 4}$ data M_{miss}^2 spectrum with with a linear combination of the simulated signal, simulated background and unsimulated background spectra by minimizing the χ^2 :

$$\chi^2 = \sum_{i=\min}^{\max} \frac{(Data_i - p_0 \cdot S_i - p_1((1 - p_2) \cdot Bg_i + p_2 \cdot UBg_i))^2}{\delta Data_i^2 + p_0^2 \cdot \delta S_i^2 + p_1^2(1 - p_2)^2 \cdot \delta Bg_i^2 + p_1^2 p_2^2 \cdot \delta UBg_i^2}, \quad (6)$$

where i is the M_{miss}^2 bin number; $Data_i$ is the bin content of the data histogram; S_i , Bg_i , UBg_i are the bin contents of the simulated signal, simulated background and unsimulated background model, correspondingly. $\delta Data_i$, δS_i , δBg_i and δUBg_i are their statistical uncertainties, and p_0, p_1, p_2 are free parameters of the fit.

Prior to fit, the S_i , Bg_i , UBg_i are scaled to make their integrals in the signal region equal to one. An interval of $[0,1]$ is allowed for p_2 values during the fit. In such a way, p_0 represents the best fit MC signal in terms of events amount, p_1 is the total background in the signal region and p_2 is the share of unsimulated background.

But we don't consider p_0 parameter (fit of the peak) as a measured signal size, as it may depend on the peak resolution simulation quality. The signal is extracted as a difference between the data histogram content in the signal region and p_1 representing the total background:

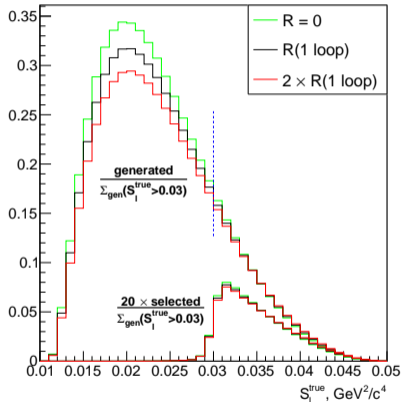
$$N_S = \sum_j Data_j - p_1, \quad (7)$$

where j bins interval corresponds to the signal region.

For the present result, we avoid the resolution simulation problem by ignoring the signal region in χ^2 calculation.

The tails of MC simulated signal are taken into account outside the signal region, that requires an estimation of MC signal using the relation (7) during the fit.

$K_{\mu 4}^{00}$ acceptance vs R form factor contribution



- We change R(1 loop) form factor normalization between 0% and 200% to illustrate different R-sensitivity of
 - the full phase space acceptance
$$A_S = \frac{N_{\text{Selected}}^{\text{MC}}}{N_{\text{Generated}}^{\text{MC}}(\text{all } S_1^{\text{true}})}$$
 and
 - the restricted phase space acceptance
$$A_S^r = \frac{N_{\text{Selected}}^{\text{MC}}}{N_{\text{Generated}}^{\text{MC}}(S_1^{\text{true}} > 0.03)}.$$
- For our systematic uncertainty, we consider only 20% R variation.

Figure: MC acceptance ingredients for different R contributions

Trigger efficiency

- A common trigger chain for the signal (S) and normalization (N) modes: a first level (L1) trigger using signals from HOD (Q1+Q2) and LKr (NUT), followed by a second level (L2) trigger using DCH (MBX).
- It is possible to measure separately NUT efficiency using the $Q1 \cdot [E_{LKr} > 10 \text{ GeV}] \cdot MBX$ control trigger and the 'charged trigger chain' $CHT = (Q1 + Q2) \cdot MBX$ efficiency using a special control trigger $n_x > 2 || n_y > 2$ for 2003 and $n_x > 3 || n_y > 3$ for 2004 data.
- The ratio $K_{trig} = \frac{\mathcal{E}_N}{\mathcal{E}_S}$ is a multiplicative trigger correction to the measured branching ratio. We decompose it to the neutral trigger and charged trigger corrections, that may be measured separately: $K_{trig} = (K_{NUT} = \frac{\mathcal{E}_N^{NUT}}{\mathcal{E}_S^{NUT}}) \times (K_{CHT} = \frac{\mathcal{E}_N^{CHT}}{\mathcal{E}_S^{CHT}})$.
- Direct \mathcal{E}_S measurement has only 2% precision due to small statistics of the control samples for the rare decay.
- \mathcal{E}_S^{CHT} is estimated using MC trigger simulation and taking into account comparison between MC and measured \mathcal{E}_N^{CHT} ;
- \mathcal{E}_S^{NUT} is recalculated from \mathcal{E}_N^{NUT} using MC simulated $(S_\pi, Min(E_\gamma), Z_c)$ distributions.

As a result, we estimate the charged trigger correction factor as $K_{CHT} = 0.998 \pm 0.002$ and the neutral trigger correction factor as $K_{NUT} = 1.0007 \pm 0.0007$.

Signal versus S_{π}, S_l

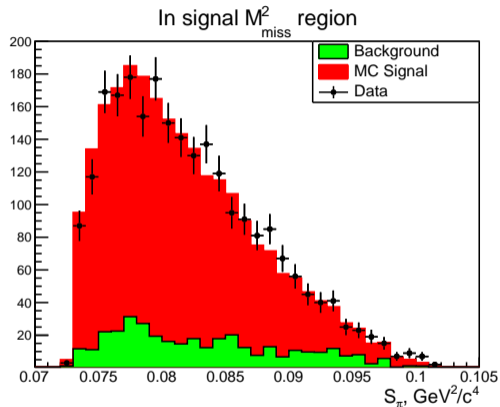
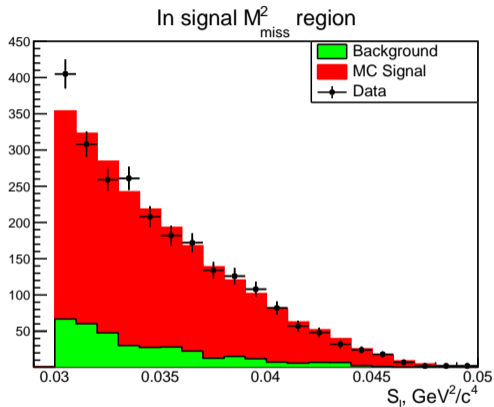


Figure: 1D projections comparison for $S_l > 0.03 \text{ GeV}^2/c^4$

Signal versus S_π, S_l

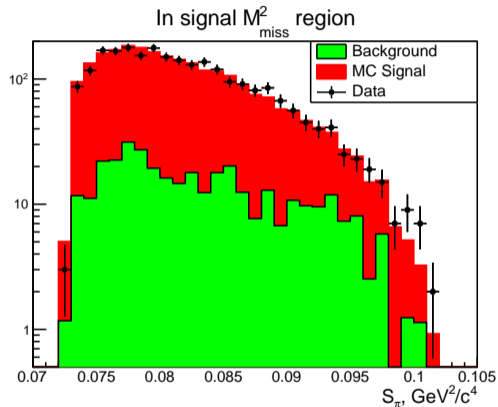
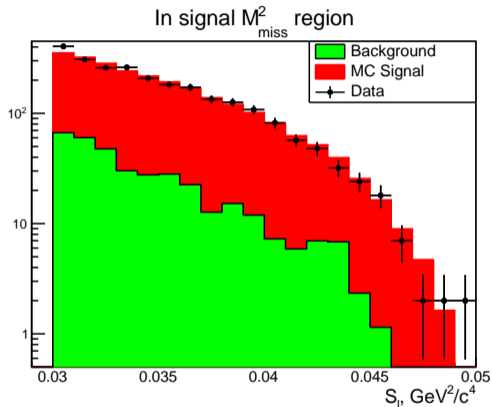


Figure: 1D projections comparison for $S_l > 0.03 \text{ GeV}^2/c^4$