

# First measurement of $K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$ ( $K_{\mu 4}^{00}$ ) decay

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on behalf of the NA48/2 Collaboration

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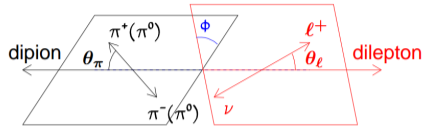
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- 8 Result
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# $K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$ ( $K_{\mu 4}^{00}$ ) state of the art

$K \rightarrow \pi\pi\mu\nu$  ( $K_{I4}$ ) depends on  $F, G, R, H$  form-factors.

Cabibbo-Maksymowicz variables:  $S_\pi$  (dipion mass squared),  $S_l$  (dilepton mass squared) and angles  $\theta_\pi$  (in the dipion frame),  $\theta_l$  (in the dilepton frame),  $\phi$ .



- For  $K_{\mu 4}^{00}$ , s-wave for  $\pi^0 \pi^0$ , there are no dependence on  $\cos \theta_\pi, \phi$ , and only  $F$  and  $R$  contribute.
- Unlike  $K_{e 4}^{00}$  case,  $R$  plays some role due to  $\mu$  mass.

$K_{I4}$ mode	BR [ $10^{-5}$ ]	$N_{cand}$	
$K_{e 4}^\pm$	$4.26 \pm 0.04$	1108941	NA48/2 (2012)
$K_{e 4}^{00}$	$2.55 \pm 0.04$	65210	NA48/2 (2014)
$K_{\mu 4}^\pm$	$1.4 \pm 0.9$	7	Bisi et al. (1967)
$K_{\mu 4}^{00}$	?	0	

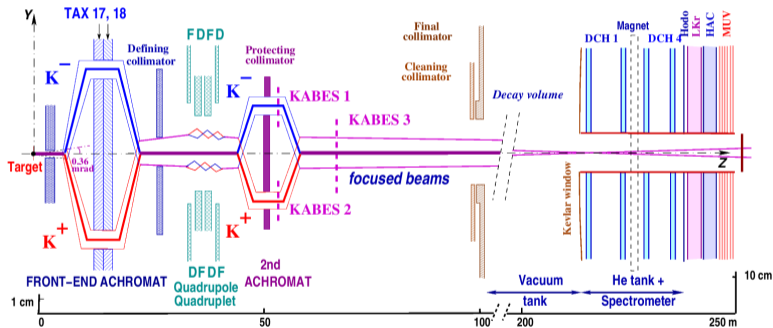
$K_{\mu 4}^{00}$ : first observation, ChPT test, check of  $R$  presence, potential study of  $\pi\pi$  rescattering effects in the  $F(S_\pi)$ .

$K_{\mu 4}$ : huge bkg  $K^\pm \rightarrow \pi\pi(\pi^\pm \rightarrow \mu^\pm \nu)$ .

- According to lepton universality, experimental  $F(S_\pi, S_l)$  parameterization from  $K_{e 4}^{00}$  [NA48/2 JHEP 08 (2014) 159] may be used for  $K_{\mu 4}^{00}$ .
- The only available source of  $R(S_\pi, S_l)$  is ChPT calculation [J.Bijnens, G.Colangelo, J.Gasser, Nucl.Phys.B 427 (1994) 427].

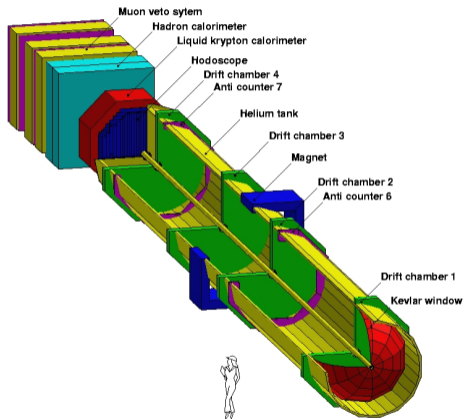
# NA48/2 beamline (CERN SPS, 2003-2004)

NA48/2 main goal were  $K_{3\pi}$  charge asymmetry studies; additional rare decays program.



- Two charged beams:
  - 6% of  $K^\pm$
  - $\langle P_K \rangle \approx 60 \text{ GeV}/c$
  - $\Delta P_K / \langle P_K \rangle \approx \pm 3.8\%$
- KABES (Kaon Beam Spectrometer) resolutions:
  - $\sigma(X, Y) \sim 800 \mu\text{m}$
  - $\sigma(P_K) / P_K \sim 1\%$
  - $\sigma(T) \sim 600 \text{ ps}$

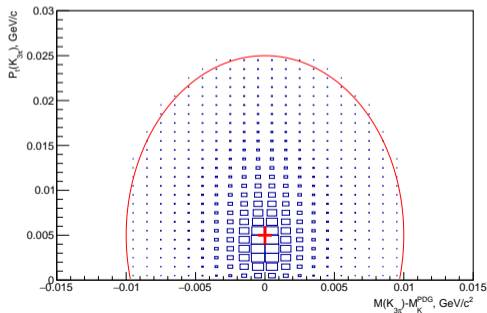
# NA48/2 setup (CERN SPS, 2003-2004)



- Magnetic spectrometer (drift chambers DCH1–DCH4):
  - $\sigma(X, Y) \sim 90 \mu\text{m}$  per chamber
  - $\sigma(P_{DCH})/P_{DCH} = (1.02 \oplus 0.044 \cdot P_{DCH})\%$   
( $P_{DCH}$  in GeV/c)
- Scintillator hodoscope (HOD):
  - $\sigma(T) \sim 150 \text{ ps}$
- Liquid Krypton EM calorimeter (LKr):
  - $\sigma_x = \sigma_y = (0.42/\sqrt{E_\gamma} \oplus 0.06) \text{ cm}$
  - $\sigma(E_\gamma)/E_\gamma = (3.2/\sqrt{E_\gamma} \oplus 9.0/E_\gamma \oplus 0.42)\%$   
( $E_\gamma$  in GeV)
- Hadronic calorimeter, muon system MUV.

# Events selection

- Signal  $K_{\mu 4}$  is  $K^{\pm} \rightarrow \pi^0 \pi^0 \mu^{\pm} \nu$
- Normalization  $K_{3\pi}$  is  $K^{\pm} \rightarrow \pi^{\pm} \pi^0 \pi^0$
- Trigger chain: L1 trigger using HOD and LKr, followed by L2 trigger using DCH for online momentum calculation.
- Event selection: 4 isolated photons consistent with  $2\pi^0$  in time-spatial matching with a KABES beam track and a DCH track.



Normalization  $K_{3\pi}$  kinematic selection ellipse:

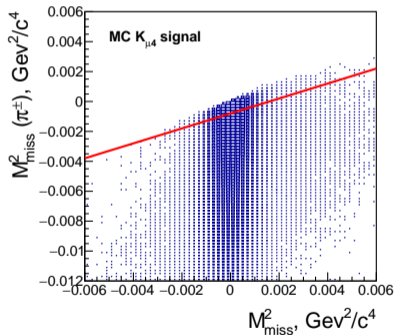
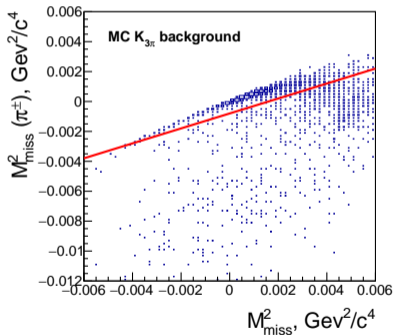
- center:
  - $M(K_{3\pi}) = M_K^{PDG}$
  - $P_t = 5 \text{ MeV}/c$
- semi-axes:
  - $\Delta M(K_{3\pi}) = 10 \text{ MeV}/c^2$
  - $\Delta P_t = 20 \text{ MeV}/c$
- $72.99 \times 10^6$   $K_{3\pi}$  selected data events.

# $K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$ signal events selection

- Off the  $K_{3\pi}$  kinematic ellipse
- DCH track has associated MUV response

$$M_{miss}^2 = (\mathbf{P}_K - \mathbf{P}(\pi_1^0) - \mathbf{P}(\pi_2^0) - \mathbf{P}(\mu^\pm))^2$$

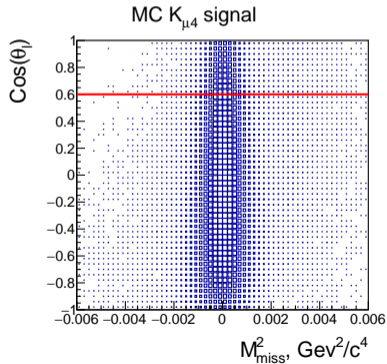
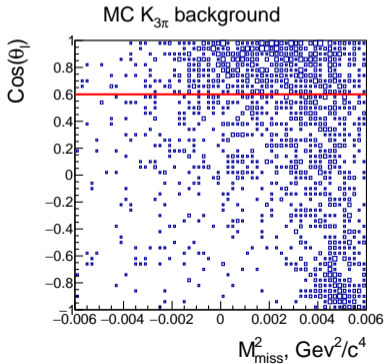
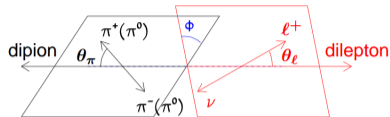
$$M_{miss}^2(\pi^\pm) = (\mathbf{P}_K - \mathbf{P}(\pi_1^0) - \mathbf{P}(\pi_2^0) - \mathbf{P}(\pi^\pm))^2$$



- $M_{miss}^2(\pi^\pm) < 0.5 M_{miss}^2 - 0.0008 \text{ GeV}^2/c^4$

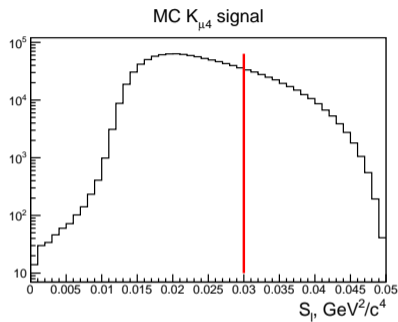
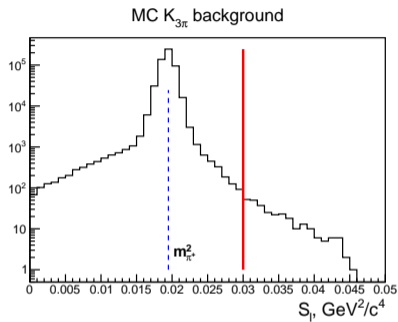
# $K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$ signal events selection

- $\cos(\Theta_l) < 0.6$



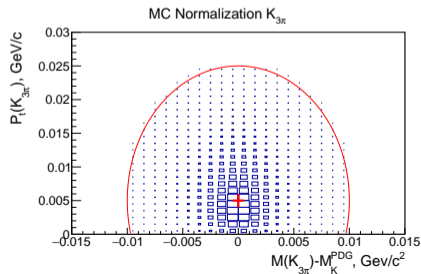
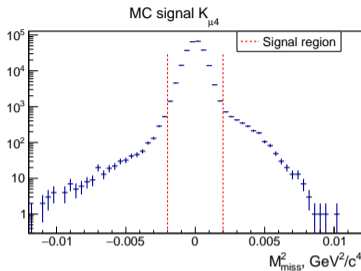


# $K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$ signal events selection



- $S_l = m(\mu\nu)^2 > 0.03 \text{ GeV}^2/c^4$  (to reject  $\pi^\pm \rightarrow \mu^\pm \nu$ ).
- 3718  $K_{\mu 4}$  data candidates selected
- 2437 data candidates in  $M_{miss}^2$  signal region  $[-0.002, 0.002] \text{ GeV}^2/c^4$
- The MC  $M_{miss}^2$  signal region contains 98.2% of all selected MC events

# Acceptances



- $K_{\mu 4}^{00}$  signal acceptance is

$$A_S = \frac{N_{\text{Selected in signal region}}^{MC}}{N_{\text{Generated (all } S_i^{\text{true}})}^{MC}} = (0.651 \pm 0.001)\%$$

- However, for the restricted phase space region  $S_i^{\text{true}} > 0.03 \text{ GeV}^2/c^4$ , the signal acceptance is

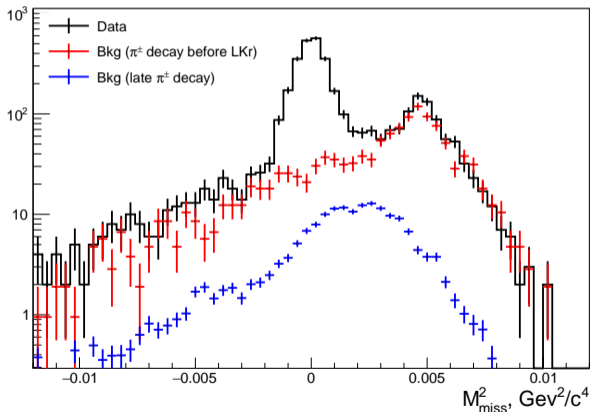
$$A_S^r = \frac{N_{\text{Selected in signal region}}^{MC}}{N_{\text{Generated } (S_i^{\text{true}} > 0.03)}^{MC}} = (3.453 \pm 0.007)\%$$

- $K_{3\pi}$  normalization channel acceptance is

$$A_N = \frac{N_{\text{Selected}}^{MC}}{N_{\text{Generated}}^{MC}} = (4.477 \pm 0.002)\%$$

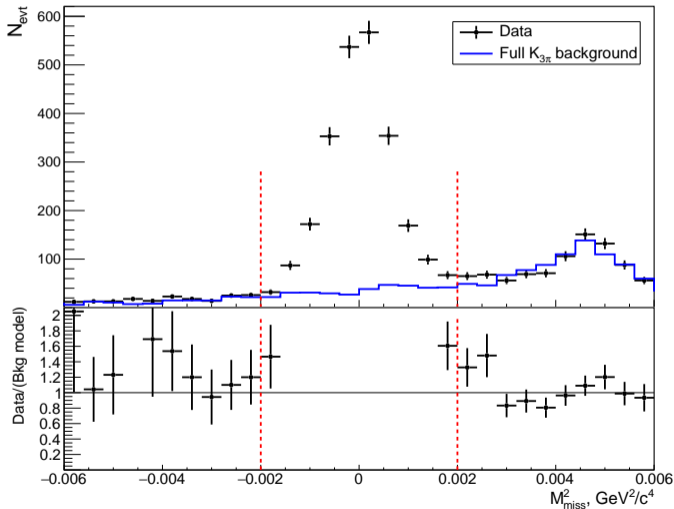
# Residual background

- $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ ,
- followed by  $\pi^\pm \rightarrow \mu\nu$  before MUV with a probability  $\approx 10\%$  for  $P(\pi^\pm) \approx 10$  GeV/c.



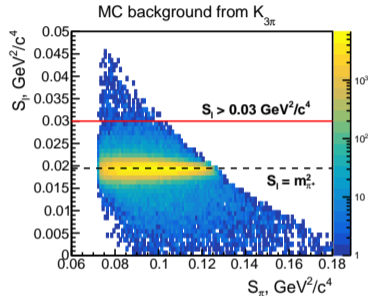
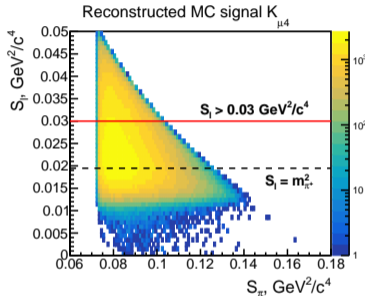
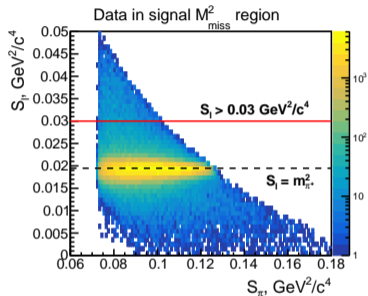
- $K_{3\pi}$  background with  $\pi^\pm$  decay before LKr: from MC.
- $K_{3\pi}$  background with late  $\pi^\pm$  decay or muon emission in a late hadron shower:
  - Can not be easily simulated
  - Data-driven method of estimation
  - Background-enhanced control sample, selected using  $E_{LKr}$  and  $P_{DCH}$

# $K_{\mu 4}^{00}$ signal extraction fit



- 2437 candidates in the signal region.
- Fit in the  $M_{\text{miss}}^2$  interval  $[-0.003, 0.006] \text{ GeV}^2/c^4$ , ignoring the signal region to decrease sensitivity to the imperfect MC resolution.
- Data fit by a linear combination of background and MC signal tails.
- $354 \pm 33_{\text{stat}} \pm 62_{\text{syst}}$  background events.
- The background-related systematics are determined by varying the way the background is estimated.

# Signal versus $S_{\pi^+}, S_I$



- The branching ratio is measured for the restricted phase space  $S_I^{true} > 0.03 \text{ GeV}^2/c^4$ .
- Extrapolation to the full phase space depends on the theory.

# Signal versus $S_{\pi}, S_l$

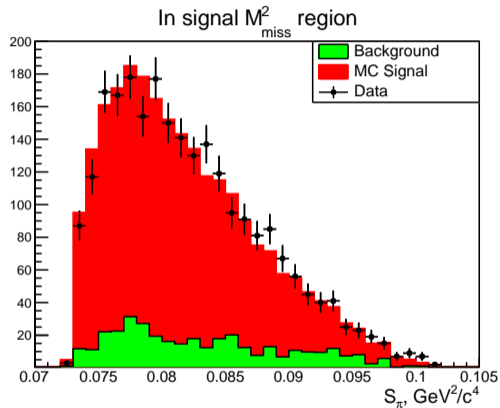
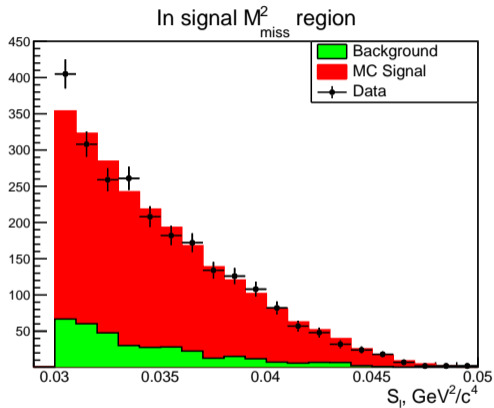


Figure: 1D projections comparison for  $S_l > 0.03 \text{ GeV}^2/c^4$

# Acceptance variations due to $K_{\mu 4}^{00}$ generator modifications

NA48/2 JHEP 08 (2014) 159:

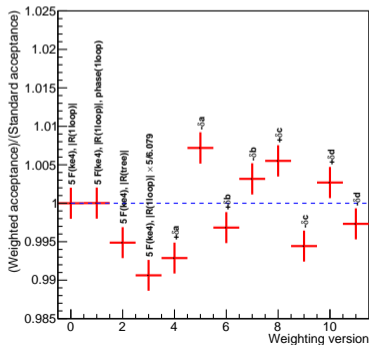
The absolute normalization is also measured  $F = f \times F(K_{e4})$ ,  
 $f = 6.079 \pm 0.055$ .

$$F(K_{e4}) = \begin{cases} (1 + aq^2 + bq^4 + c \cdot S_I/4m_{\pi^+}^2) & \text{for } q^2 \geq 0 \\ (1 + d\sqrt{|q^2/(1+q^2)|} + c \cdot S_I/4m_{\pi^+}^2) & \text{for } q^2 < 0 \end{cases}$$

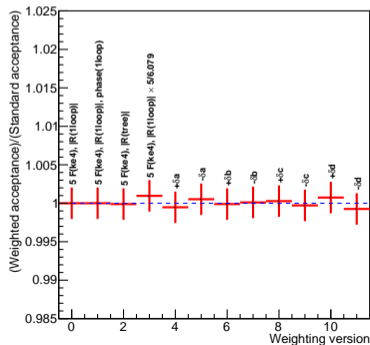
where  $q^2 = S_{\pi}/4m_{\pi^+}^2 - 1$ .

- Decay generator was modified by MC events weighting.
- The acceptance spread is taken as systematics.

Full acceptance relative variations



( $S_1 > 0.03$ ) acceptance relative variations



# Preliminary result: Ingredients

$$BR(K_{\mu 4}^{00}) = \frac{N_S}{N_N} \cdot \frac{A_N}{A_S} \cdot K_{trig} \cdot BR(K_{3\pi}^{00}).$$

- Extracted signal  $N_S = N_{Sign. cand.} - N_{Bkg} = 2437 - (354 \pm 33_{stat}) = 2083 \pm 59_{stat}$  events;
  - Signal/Background is  $5.89 \pm 0.66_{stat}$ ;
- Number of normalization events  $N_N = 72.99 \times 10^6$ ;
- Normalization acceptance  $A_N = (4.477 \pm 0.002)\%$ ;
- Signal acceptance for the restricted phase space  $A_S^r = (3.453 \pm 0.007)\%$ ;
- Signal acceptance for the full phase space  $A_S = (0.651 \pm 0.001)\%$ ;
- Trigger correction (extracted with control triggers)  
 $K_{trig} = K_{CHT} \cdot K_{NUT} = (0.998 \pm 0.002) \cdot (1.0007 \pm 0.0007) = 0.999 \pm 0.002$ ;
- PDG  $BR(K_{3\pi}^{00}) = (1.760 \pm 0.023)\%$ ;

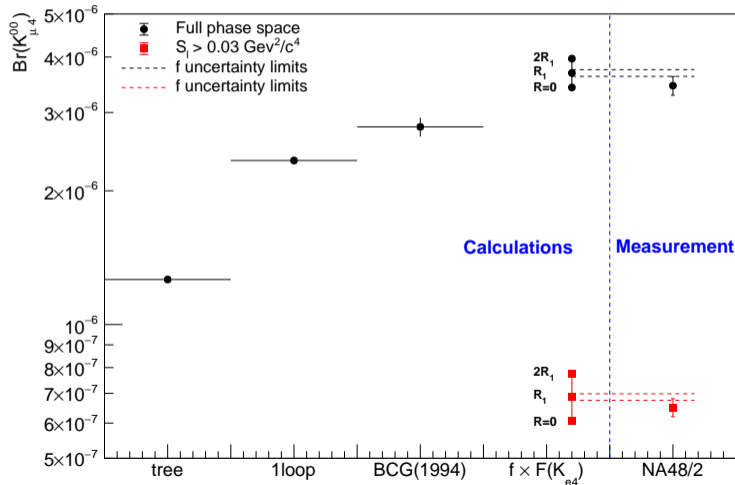


# Preliminary result: Central values and errors budget

	Full phase space		$S_l > 0.03 \text{ GeV}^2/c^4$	
$BR(K_{\mu 4})$ central value [ $10^{-6}$ ]	3.45		0.651	
	$\delta BR [10^{-6}]$	$\delta BR / BR$	$\delta BR [10^{-6}]$	$\delta BR / BR$
Data stat. error	0.10	2.85%	0.019	2.85%
MC stat. error	0.01	0.21%	0.001	0.21%
Trigger	0.01	0.18%	0.001	0.18%
Background	0.10	2.96%	0.019	2.96%
Accidentals	0.01	0.32%	0.002	0.32%
MUV inefficiency	0.06	1.65%	0.011	1.65%
Form Factor modelling	0.05	1.37%	0.001	0.14%
$BR(K_{3\pi})$ error (external)	0.05	1.31%	0.009	1.31%
Total error	0.17	4.83%	0.030	4.64%

- Accidentals obtained from side bands of time distributions;
- MUV inefficiency uncertainty taken as full inefficiency effect.

# Preliminary result: Comparison to theory



## Theory:

- J. Bijnens, G. Colangelo, J. Gasser, Nucl. Phys. B, 427 (1994) 427:
  - Tree approximation;
  - 1-loop;
  - BCG(1994): 'beyond 1-loop' with measured  $F$  from [Rosselet etc. Phys. Rev. D 15 (1977) 574].
- Re-calculated now:
  - $F(K_{e4})$  from NA48/2 (2015);
  - $R_1 = R(1loop)$ ;
  - 1-loop (F,R) phase;
  - 2020 PDG constants.

A first observation and branching fraction measurement of  $K^\pm \rightarrow \pi^0 \pi^0 \mu^\pm \nu$  decay is performed by NA48/2 experiment at SPS in CERN

- We observe 2437 signal candidates with an estimated background of  $354 \pm 33_{stat} \pm 62_{syst}$  events, Signal/Background ratio is  $5.9 \pm 1.4_{tot}$
- Preliminary result for restricted phase space ( $S_I > 0.03$ ) is

$$BR(K_{\mu 4}^{00}, S_I > 0.03) = (0.65 \pm 0.019_{stat} \pm 0.024_{syst}) \times 10^{-6} = (0.65 \pm 0.03) \times 10^{-6};$$

- Preliminary full phase space result is

$$BR(K_{\mu 4}^{00}) = (3.4 \pm 0.10_{stat} \pm 0.13_{syst}) \times 10^{-6} = (3.4 \pm 0.2) \times 10^{-6}.$$

- The results are consistent with a contribution of the R form factor, as computed at 1-loop ChPT.

SPARE SLIDES

# Theoretical framework: decay width

$K_{\mu 4}^{00}$  matrix element is [J.Bijnens, G.Colangelo, J.Gasser, Nucl.Phys.B 427, 427 (1994)]:

$$T = \frac{G_f}{\sqrt{2}} \cdot V_{us}^* \cdot \underbrace{\bar{u}(p_\nu) \cdot \gamma_\mu \cdot (1 - \gamma_5) \cdot v(p_l)}_{\text{Lepton part}} \cdot \underbrace{(V_\mu - A_\mu)}_{\text{Hadron part}} \quad (1)$$

where

$$V_\mu = \frac{-H}{M_K^3} \varepsilon_{\mu,\nu,\rho,\sigma} L^\nu P^\rho Q^\sigma, \quad A_\mu = -i \frac{1}{M_K} [P_\mu F + Q_\mu G + L_\mu R], \quad (2)$$

$\varepsilon_{0,1,2,3}=1$  and the four-momenta are defined as  $P = p_1 + p_2$ ,  $Q = p_1 - p_2$  ( $p_1$  and  $p_2$  are the two pion momenta) and  $L = p_l + p_\nu$ . The form factors  $F, G, R, H$  are analytic functions of the decay kinematic variables.

In general, decay width is a function of five Cabibbo-Maksymowicz variables. But for  $K_{\mu 4}^{00}$ , in s-wave approximation for  $\pi^0 \pi^0$ , there are no dependence on  $\cos \theta_\pi, \phi$ , and so  $G = 0, H = 0$ :

$$d\Gamma_3 = \frac{G_f^2 |V_{us}|^2 (1 - z_l)^2 \sigma_\pi X}{2^{11} \pi^5 M_K^5} (I_1 + I_2 (2(\cos \theta_l)^2 - 1) + I_6 \cos \theta_l) dS_\pi dS_l d \cos \theta_l. \quad (3)$$

where  $I_1 = \frac{1}{4} \{ (1 + z_l) |F_1|^2 + 2z_l |F_4|^2 \}$ ,  $I_2 = -\frac{1}{4} (1 - z_l) |F_1|^2$ ,  $I_6 = z_l \text{Re}(F_1^* F_4)$ ,  $F_1 = X \cdot F$ ,  $F_4 = -(PL)F - S_l R$ ,  $z_l = \frac{m_l^2}{S_l}$ ,  $\sigma_\pi = \sqrt{1 - \frac{4M_\pi^2}{s_\pi}}$ ,  $X = \frac{1}{2} \sqrt{\lambda(M_K^2, S_\pi, S_l)}$ ,  $\lambda(x, y, z) = x^2 + y^2 + z^2 - 2(xy + xz + yz)$ .

Only three variables ( $S_\pi, S_l, \cos \theta_l$ ) and two form factors ( $F, R$ ) matter.

# Theoretical framework: form factors

According to lepton universality,  $F_s(S_\pi, S_l)$  s-wave form factor measured in  $K_{e4}^{00}$  may be used for  $K_{\mu 4}^{00}$  decay MC simulation, while the only available source of  $R_s(S_\pi, S_l)$  is ChPT calculation.

- We have ChPT 1-loop generator code [BCG(1994)] with  $F(1loop), R(1loop)$ .
- For the predictions beyond 1-loop (central value), BCG(1994) used  $F = 5.59(1 + 0.08q^2)$ , where  $q^2 = S_\pi/4m_{\pi^+}^2 - 1$  [L.Rosselet et al., Phys. Rev. D15 (1977) 574].
- Additionally, from the NA48/2  $K_{e4}$  analysis [J.R. Batley et al. JHEP 08 (2014) 159] we have the best measurement of F shape:

$$\begin{aligned} F(K_{e4}) &= (1 + aq^2 + bq^4 + c \cdot S_l/4m_{\pi^+}^2), q^2 \geq 0 \\ F(K_{e4}) &= (1 + d\sqrt{|q^2/(1 + q^2)|} + c \cdot S_l/4m_{\pi^+}^2), q^2 \leq 0, \end{aligned} \quad (4)$$

where  $q^2 = S_\pi/4m_{\pi^+}^2 - 1$ ,  $a = 0.149 \pm 0.033 \pm 0.014$ ,  $b = -0.070 \pm 0.039 \pm 0.013$ ,  $c = 0.113 \pm 0.022 \pm 0.007$ ,  $d = -0.256 \pm 0.049 \pm 0.016$ .

For systematic studies we take into account the combined (both statistical and systematic) uncertainties of these parameters. Moreover, the absolute value was also measured  $F = 6.079F(K_{e4})$ .

Nevertheless, for MC production we have used  $5F(K_{e4})$  (to be close to the 1-loop F),  $R(1loop)$ ,  $H = G = 0$  and zero phase between  $F$  and  $R$ .

# F form factor

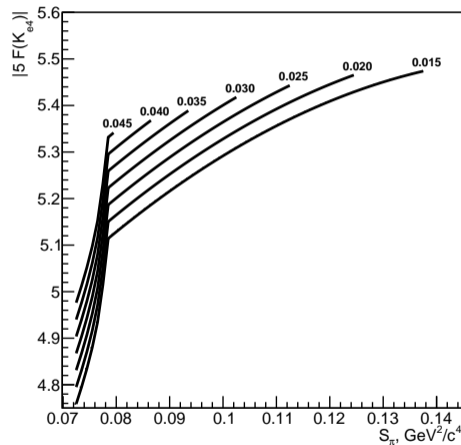
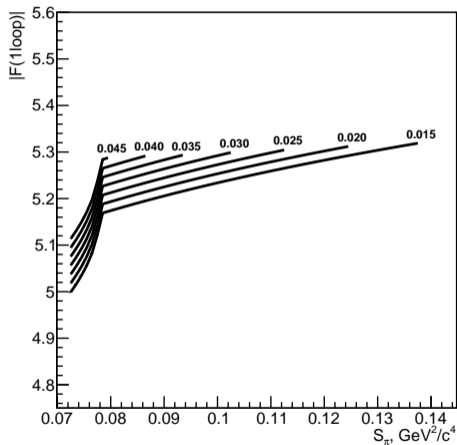


Figure: Values near the lines:  $S_l$  [GeV $^2/c^4$ ]

# R form factor

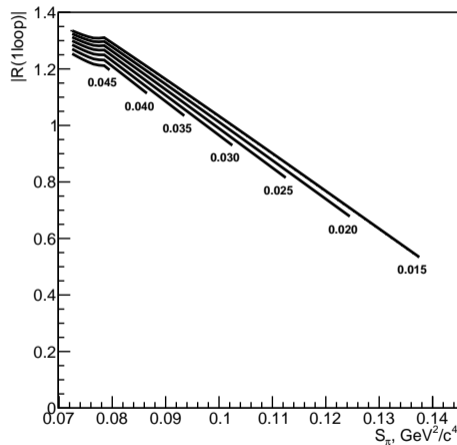
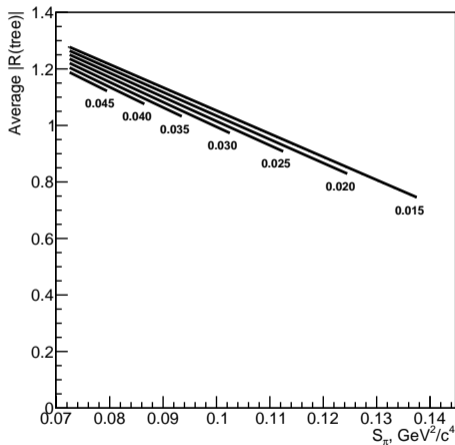


Figure: Values near the lines:  $S_l$  [ $\text{GeV}^2/c^4$ ]



# Common selection criteria for signal and normalization channels

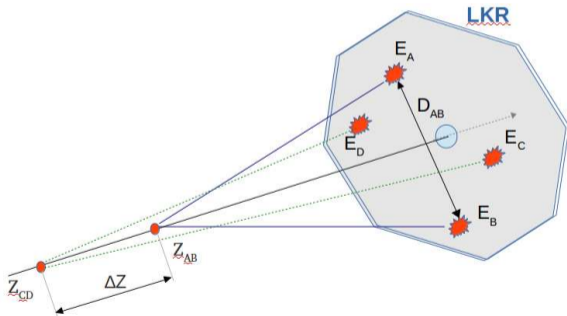
NA48/2 data collected in 2003-2004.

Signal (S)  $K_{\mu 4}$  is  $K^{\pm} \rightarrow \mu^{\pm} \nu \pi^0 \pi^0$ . Normalization (N)  $K_{3\pi}$  is  $K^{\pm} \rightarrow \pi^{\pm} \pi^0 \pi^0$ .

A common trigger chain: L1 trigger using HOD and LKr, followed by L2 trigger using DCH (MBX).

Combination: 4 isolated photons in time with a KABES track and with a DCH track.

Vertex 'charged position'  $Z_c$  is taken from the space matching of the KABES track and the DCH track.



- $P_{KABES}$  in [54,67] GeV/c;  $P_{DCH}$  in [5,35] GeV/c;  
 $E_{\gamma} > 3$  GeV;
- For  $\pi_i^0$  ( $i = 1, 2$ ):
  - $Z_1 = Z_{AB} = D_{AB} \sqrt{E_A E_B} / m_{\pi^0}$ ;
  - $Z_2 = Z_{CD} = D_{CD} \sqrt{E_C E_D} / m_{\pi^0}$ .
- $-1600 \text{ cm} < Z_n = Z_{LKR} - (Z_1 + Z_2)/2 < 9000 \text{ cm}$   
( $Z_n$  is the vertex 'neutral position');
- Flunge cut:  $R_{\gamma} @ DCH1 > 11 \text{ cm}$ , assuming  
 $Z_n + 400 \text{ cm}$  decay position;
- $|Z_1 - Z_2| < 500 \text{ cm}$ ;  $|Z_n - Z_c| < 600 \text{ cm}$ .

The best combination of each mode in the event has the most compatible  $Z_1, Z_2, Z_c$ .

# Normalization $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ channel selection

No missing momentum is expected.

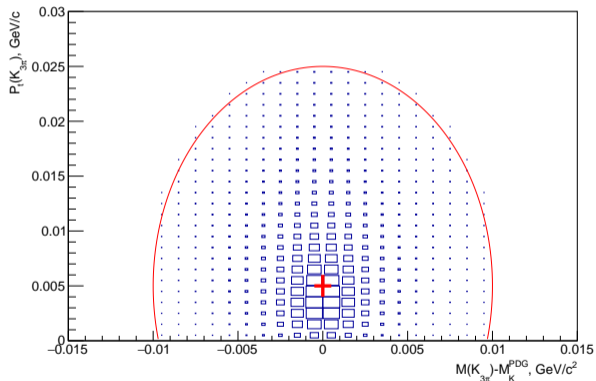


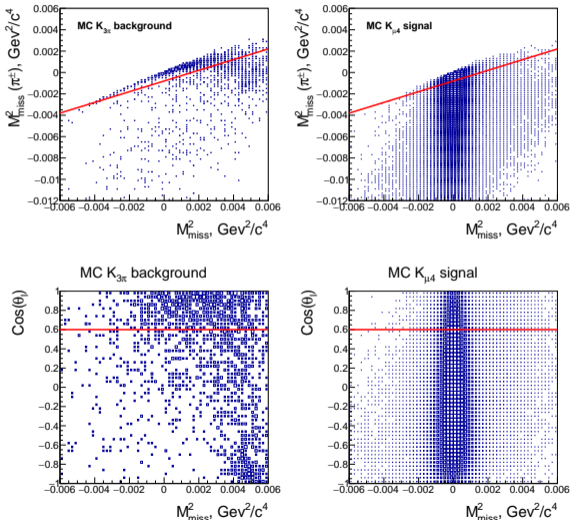
Figure:  $K_{3\pi}$  kinematic selection area

$K_{3\pi}$  kinematic selection ellipse:

- center:
  - $M(K_{3\pi}) = M_K^{PDG}$ ;
  - $P_t = 5 \text{ MeV}/c$ ;
- semi-axes:
  - $\Delta M(K_{3\pi}) = 10 \text{ MeV}/c^2$ ;
  - $\Delta P_t = 20 \text{ MeV}/c$ .

Normalization channel selection output is  $N_N = 72.99 \times 10^6$   $K_{3\pi}$  reconstructed events.

# Signal $K^\pm \rightarrow \mu^\pm \nu \pi^0 \pi^0$ channel selection



- Off the  $K_{3\pi}$  kinematic ellipse;
- Good  $K - \mu$  tracks matching;
- $P_{DCH} > 10 \text{ GeV}/c$  for MUV efficiency;
- DCH track is in restricted MUV acceptance with high MUV efficiency;
- DCH track has associated MUV hits in the first two planes;
- $M_{miss}^2$  (assuming  $\pi^\pm$ )  $< 0.5M_{miss}^2 - 0.0008 \text{ GeV}^2/c^4$ ;
- $\cos(\Theta_I) < 0.6$ ;
- $S_I > 0.03 \text{ GeV}^2/c^4$  to reject  $\pi^\pm \rightarrow \mu^\pm \nu$ .

3718  $K_{\mu 4}$  candidates selected (any  $M_{miss}^2$ ).

# $K_{\mu 4}^{00}$ signal extraction procedure

We fit  $K_{\mu 4}$  data  $M_{miss}^2$  spectrum with with a linear combination of the simulated signal, simulated background and unsimulated background spectra by minimizing the  $\chi^2$ :

$$\chi^2 = \sum_{i=\min}^{\max} \frac{(Data_i - p_0 \cdot S_i - p_1((1 - p_2) \cdot Bg_i + p_2 \cdot UBg_i))^2}{\delta Data_i^2 + p_0^2 \cdot \delta S_i^2 + p_1^2(1 - p_2)^2 \cdot \delta Bg_i^2 + p_1^2 p_2^2 \cdot \delta UBg_i^2}, \quad (5)$$

where  $i$  is the  $M_{miss}^2$  bin number;  $Data_i$  is the bin content of the data histogram;  $S_i$ ,  $Bg_i$ ,  $UBg_i$  are the bin contents of the simulated signal, simulated background and unsimulated background model, correspondingly.  $\delta Data_i$ ,  $\delta S_i$ ,  $\delta Bg_i$  and  $\delta UBg_i$  are their statistical uncertainties, and  $p_0, p_1, p_2$  are free parameters of the fit.

Prior to fit, the  $S_i$ ,  $Bg_i$ ,  $UBg_i$  are scaled to make their integrals in the signal region equal to one. An interval of  $[0,1]$  is allowed for  $p_2$  values during the fit. In such a way,  $p_0$  represents the best fit MC signal in terms of events amount,  $p_1$  is the total background in the signal region and  $p_2$  is the share of unsimulated background.

But we don't consider  $p_0$  parameter (fit of the peak) as a measured signal size, as it may depend on the peak resolution simulation quality. The signal is extracted as a difference between the data histogram content in the signal region and  $p_1$  representing the total background:

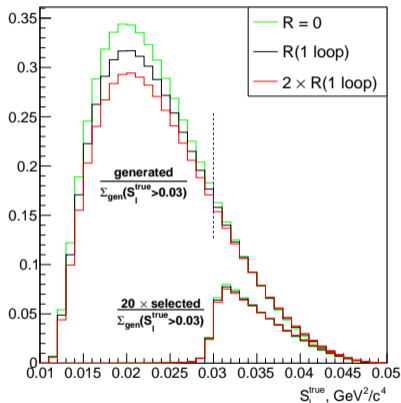
$$N_S = \sum_j Data_j - p_1, \quad (6)$$

where  $j$  bins interval corresponds to the signal region.

For the present result, we avoid the resolution simulation problem by ignoring the signal region in  $\chi^2$  calculation.

The tails of MC simulated signal are taken into account outside the signal region, that requires an estimation of MC signal using the relation (6) during the fit.

# $K_{\mu 4}^{00}$ acceptance vs R form factor contribution



- We change R(1 loop) form factor normalization between 0% and 200% to illustrate different R-sensitivity of
  - the full phase space acceptance
$$A_S = \frac{N_{\text{Selected}}^{\text{MC}}}{N_{\text{Generated}}^{\text{MC}}(\text{all } S_1^{\text{true}})}$$
 and
  - the restricted phase space acceptance
$$A_S^r = \frac{N_{\text{Selected}}^{\text{MC}}}{N_{\text{Generated}}^{\text{MC}}(S_1^{\text{true}} > 0.03)}.$$
- For our systematic uncertainty, we consider only 20% R variation.

Figure: MC acceptance ingredients for different R contributions

# Trigger efficiency

- A common trigger chain for the signal (S) and normalization (N) modes: a first level (L1) trigger using signals from HOD (Q1+Q2) and LKr (NUT), followed by a second level (L2) trigger using DCH (MBX).
- It is possible to measure separately NUT efficiency using the  $Q1 \cdot [E_{LKr} > 10 \text{ GeV}] \cdot MBX$  control trigger and the 'charged trigger chain'  $CHT = (Q1 + Q2) \cdot MBX$  efficiency using a special control trigger  $n_x > 2 || n_y > 2$  for 2003 and  $n_x > 3 || n_y > 3$  for 2004 data.
- The ratio  $K_{trig} = \frac{\mathcal{E}_N}{\mathcal{E}_S}$  is a multiplicative trigger correction to the measured branching ratio. We decompose it to the neutral trigger and charged trigger corrections, that may be measured separately:  $K_{trig} = (K_{NUT} = \frac{\mathcal{E}_N^{NUT}}{\mathcal{E}_S^{NUT}}) \times (K_{CHT} = \frac{\mathcal{E}_N^{CHT}}{\mathcal{E}_S^{CHT}})$ .
- Direct  $\mathcal{E}_S$  measurement has only 2% precision due to small statistics of the control samples for the rare decay.
- $\mathcal{E}_S^{CHT}$  is estimated using MC trigger simulation and taking into account comparison between MC and measured  $\mathcal{E}_N^{CHT}$ ;
- $\mathcal{E}_S^{NUT}$  is recalculated from  $\mathcal{E}_N^{NUT}$  using MC simulated  $(S_\pi, Min(E_\gamma), Z_c)$  distributions.

As a result, we estimate the charged trigger correction factor as  $K_{CHT} = 0.998 \pm 0.002$  and the neutral trigger correction factor as  $K_{NUT} = 1.0007 \pm 0.0007$ .