

PBH assisted search for QCD axion Dark Matter

Work in collaboration with E. Schiappacasse
arXiv : 2205.02255

Gongjun Choi (CERN)

CERN-CKC workshop, 7/Jun/2022

Can $f_{\text{PBH}} = \mathcal{O}(0.1) - \mathcal{O}(1)\%$ be of a help in
search for
(large decay constant) QCD axion DM?

Outline

- Large decay constant axion & EMD era
- Higher chance of PBH formation
- Ultra Compact Minihalo (UCMH) formation
- PBH-assisted search for QCD axion DM

Large decay constant axion

- Correct axion DM relic abundance with $\theta_i \sim \mathcal{O}(1)$

$$\Omega_a h^2 \simeq (2 \times 10^4) \times \theta_i^2 \times \left(\frac{F_a}{10^{16} \text{GeV}} \right)^{\frac{7}{6}}$$

$$\rightarrow F_a < 10^{12} \text{ GeV}$$

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G. Dvali (9505253)

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EMD era & Entropy production

- Not really a contrived set-up
if we can naturally accommodate a heavy particle
- In SUSY models or string theory
→ saxion or moduli (σ)
- If σ decays before BBN & produces enough radiation
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→ Ω_a can be sufficiently diluted
- Necessarily accompanied by
early matter dominated (EMD) era!

EMD era & Entropy production

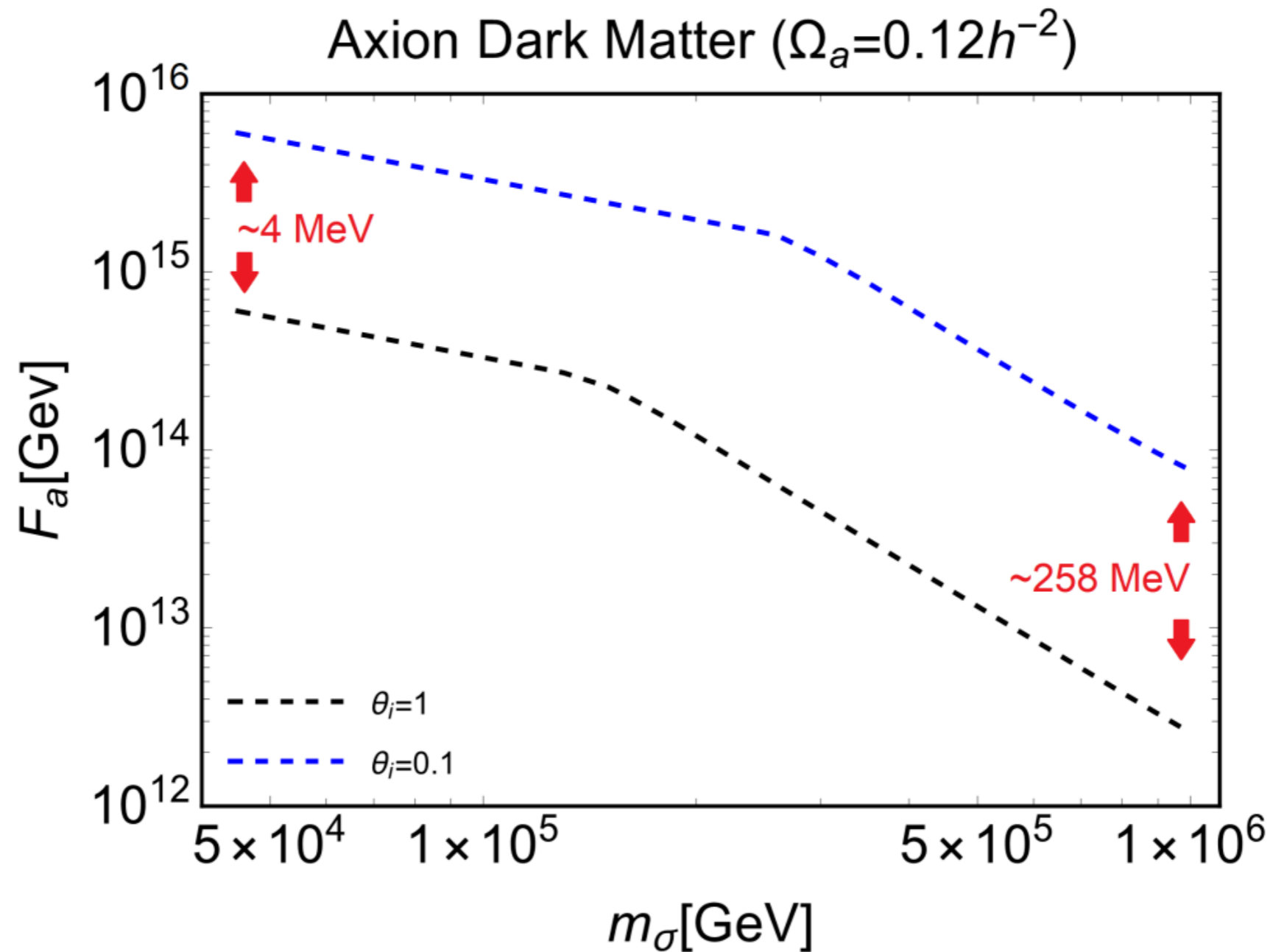
- Quantifying the newly produced entropy

$$\Delta \equiv \frac{S_{\text{new}}}{S_{\text{old}}} = \frac{s_{\text{new}}(R_{\text{rh}2}) R_{\text{rh}2}^3}{s_{\text{old}}(R_{\text{osc}}) R_{\text{osc}}^3} = \frac{s_{\text{new}}(R_{\text{rh}2})}{s_{\text{old}}(R_{\text{rh}2})}$$

$$\begin{aligned} \Omega_a h^2 &= \frac{m_a s_0 h^2}{\rho_{\text{cr},0}} \left(\frac{n_a(R_{\text{rh}2})}{s_{\text{new}}(R_{\text{rh}2})} \right) \\ &= \frac{m_a s_0 h^2}{\rho_{\text{cr},0}} \left(\frac{n_a(R_{\text{rh}2})}{s_{\text{old}}(R_{\text{rh}2})} \right) \left(\frac{s_{\text{old}}(R_{\text{rh}2})}{s_{\text{new}}(R_{\text{rh}2})} \right) & \frac{4\rho_\sigma(R_{\text{rh}2})}{3T_{\text{rh}2}} = s_{\text{new}}(R_{\text{rh}2}) \\ &= \frac{m_a s_0 h^2}{\rho_{\text{cr},0}} \left(\frac{n_a(R_{\text{osc}})}{s_{\text{old}}(R_{\text{osc}})} \right) \Delta^{-1} \\ &= \Omega_{a,\text{old}} h^2 \Delta^{-1} \end{aligned}$$

EMD era & Entropy production

- Diluting axion abundance by σ -decay



PBH formation probability

- PBH formation probability ($\beta \equiv \rho_{\text{PBH}}/\rho_{\text{tot}}$)
→ much higher in MD era than RD era

$$\beta_{\text{MD}} \simeq 0.056 \sigma^5 \quad \gg \quad \beta_{\text{RD}}(M) \simeq \text{erfc}[\delta_c/\sqrt{2}\sigma(M)]$$

T. Harada et al. (Astrophys.J, 2016)

B. Carr (Astrophys. J. 201, 1 (1975)), T. Harada et al. (PRD, 2013)

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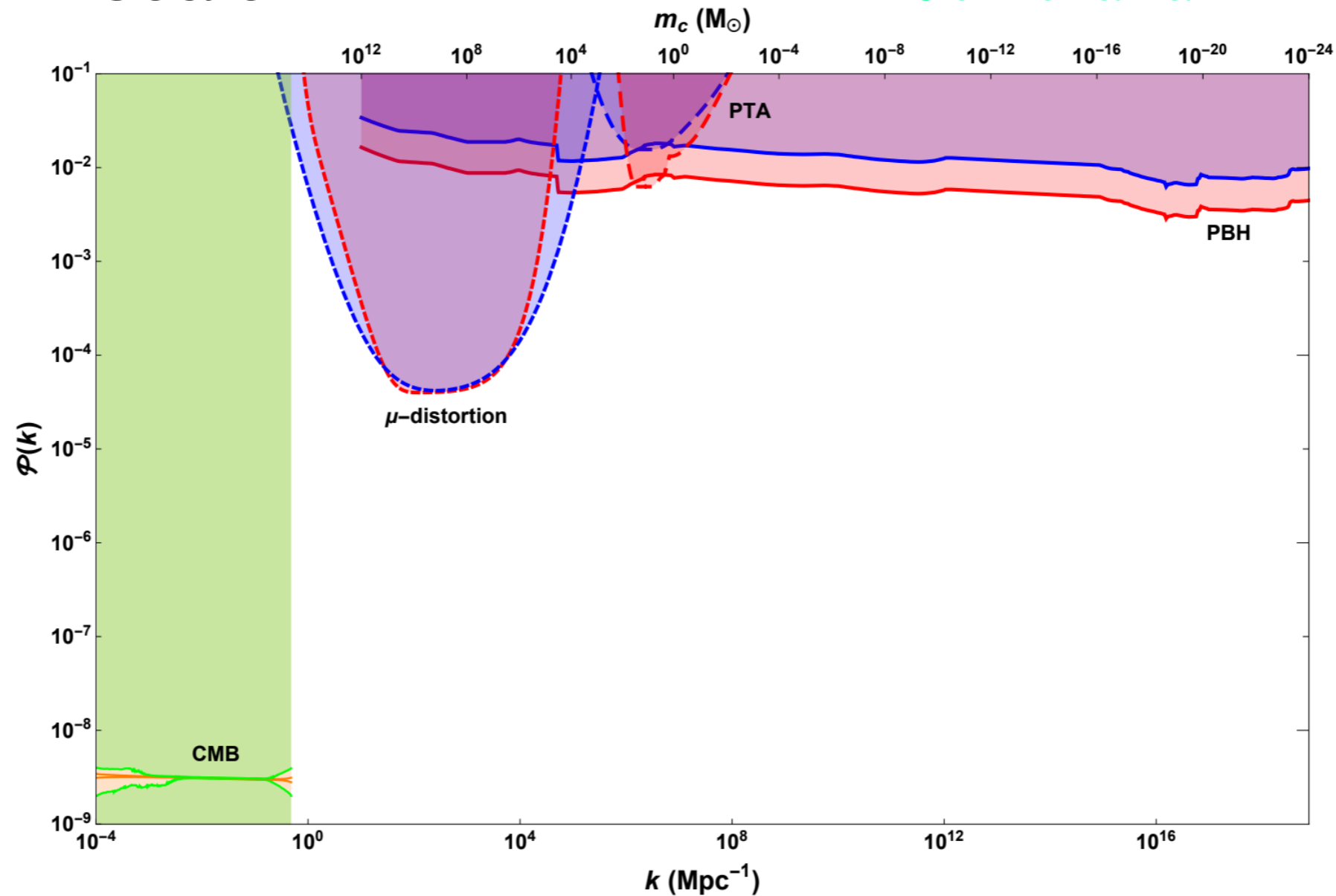
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EMD era motivated by large $F_a > 10^{12}\text{GeV}$ can serve as the environment facilitating formation of PBHs

Can we have a large enough $P_{\zeta}(k)$?

- In fact, $P_{\zeta}(k)$ in small scales are much less constrained than CMB scale

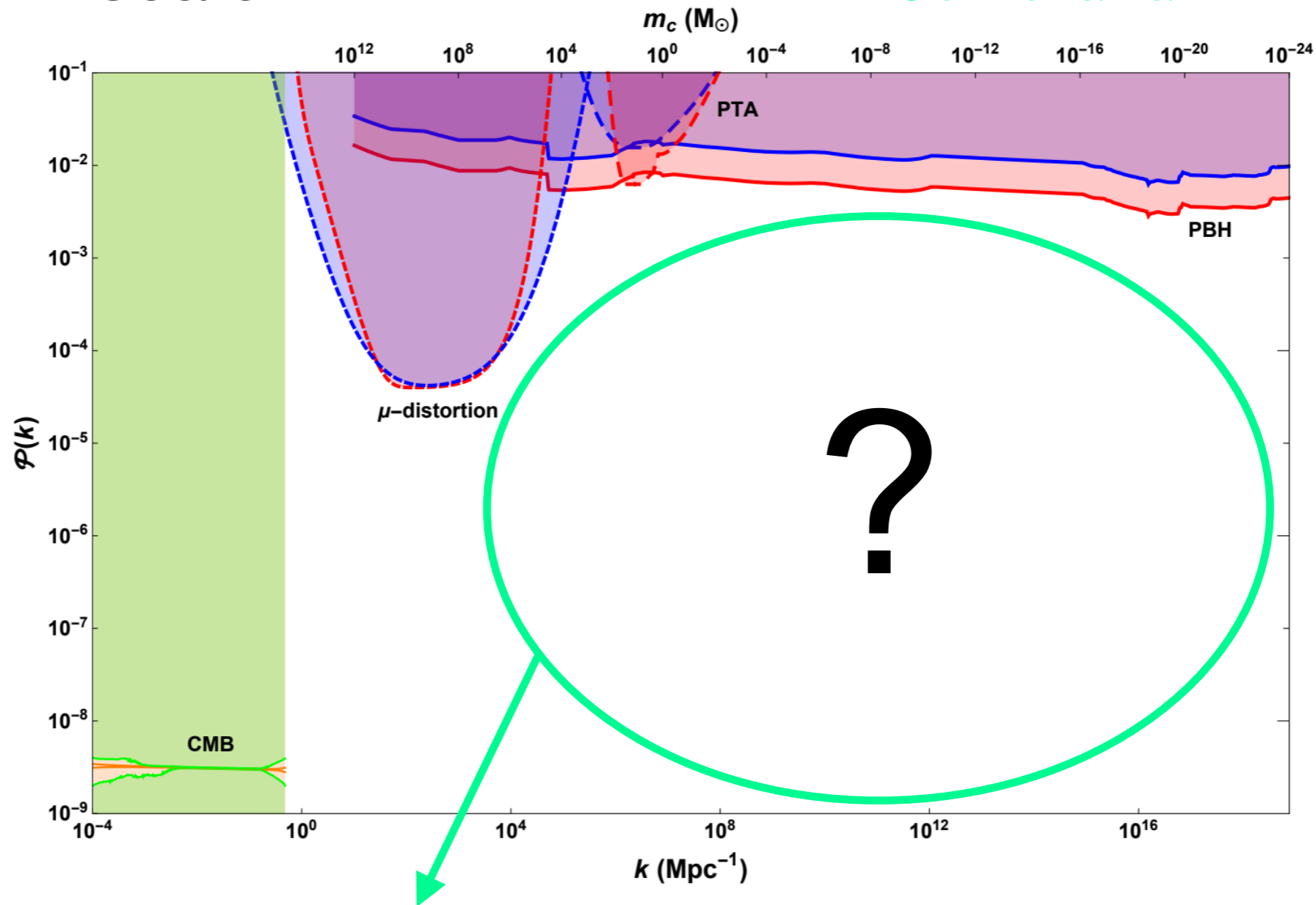
A. Gow et al. arxiv: 2008.03289



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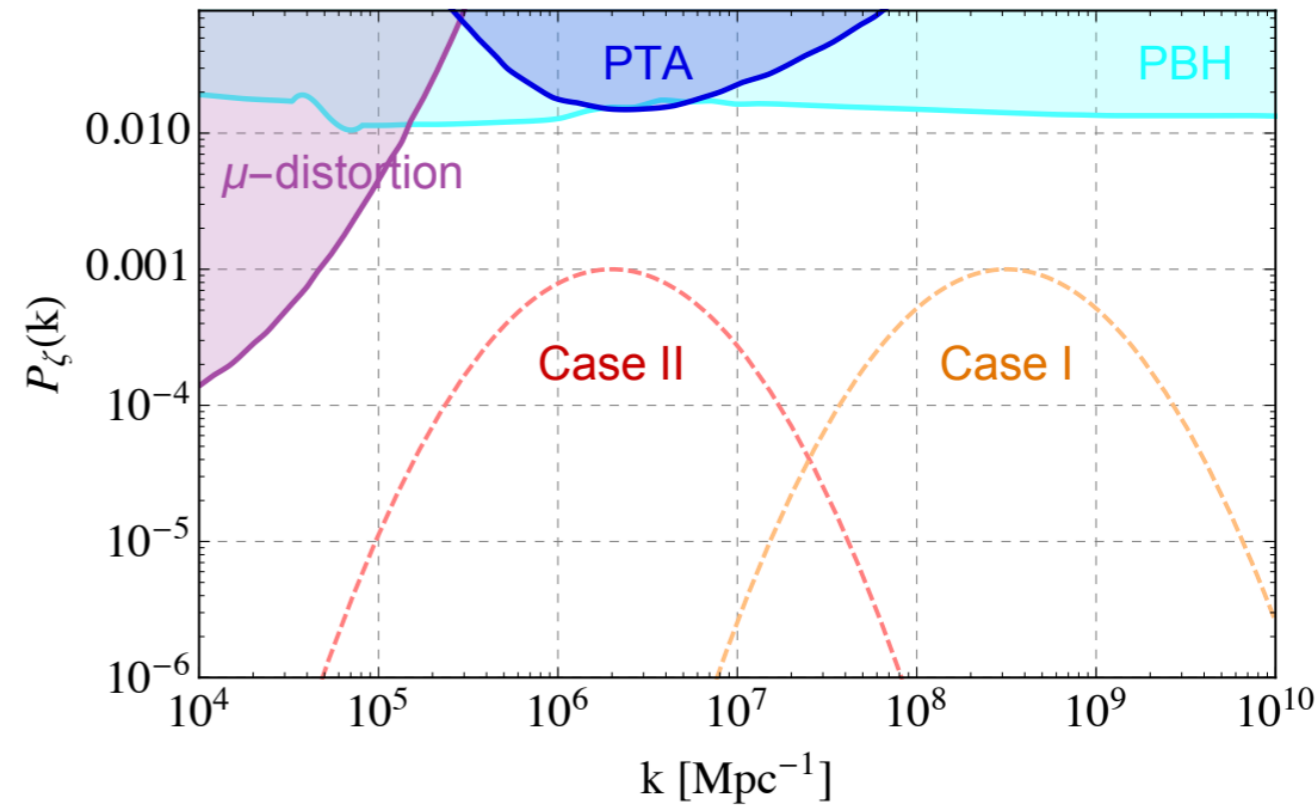
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No reason why we can't assume enhanced $P_{\zeta}(k)$ on small scales

$P_\zeta(k)$ on small scales & f_{PBH}

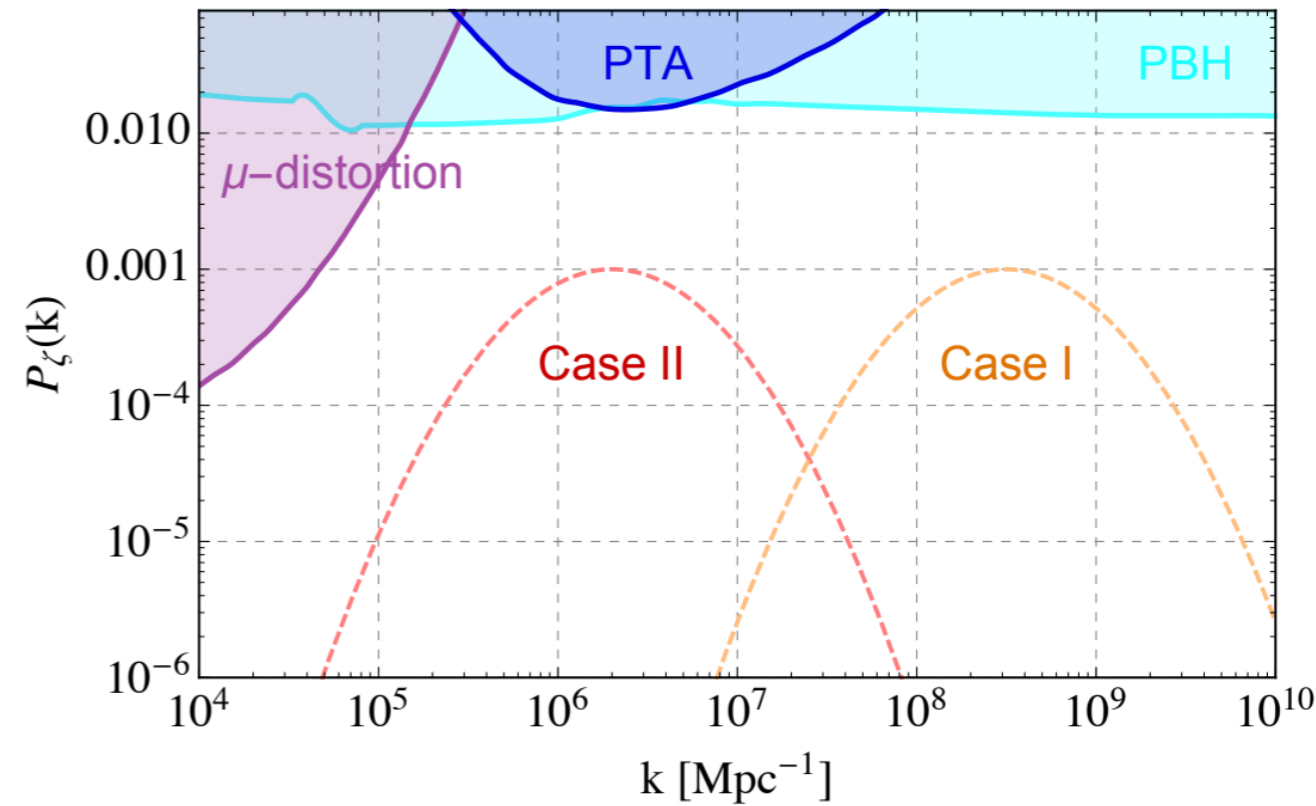
$$P_\zeta(k) = A_s \left(\frac{k}{k_{\text{CMB}}} \right)^{n_s - 1} + A_p e^{-\frac{(N_k - N_p)^2}{2\sigma_p^2}}$$
$$\simeq A_s \left(\frac{k}{k_{\text{CMB}}} \right)^{n_s - 1} + A_p e^{-\frac{(\log(k/k_p))^2}{2\sigma_p^2}}$$



$P_\zeta(k)$ on small scales & f_{PBH}

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$$f_{\text{PBH}}(M) \simeq (2.35 \times 10^5)$$

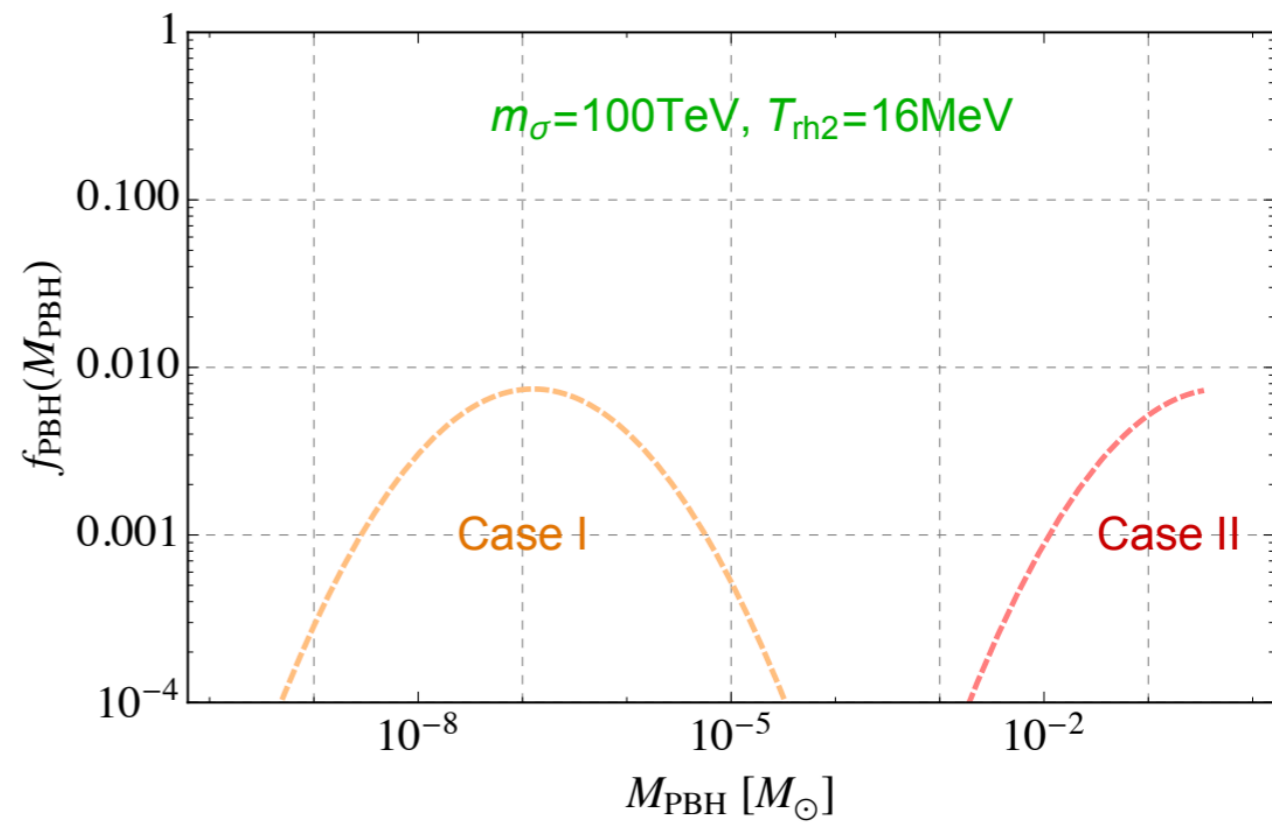
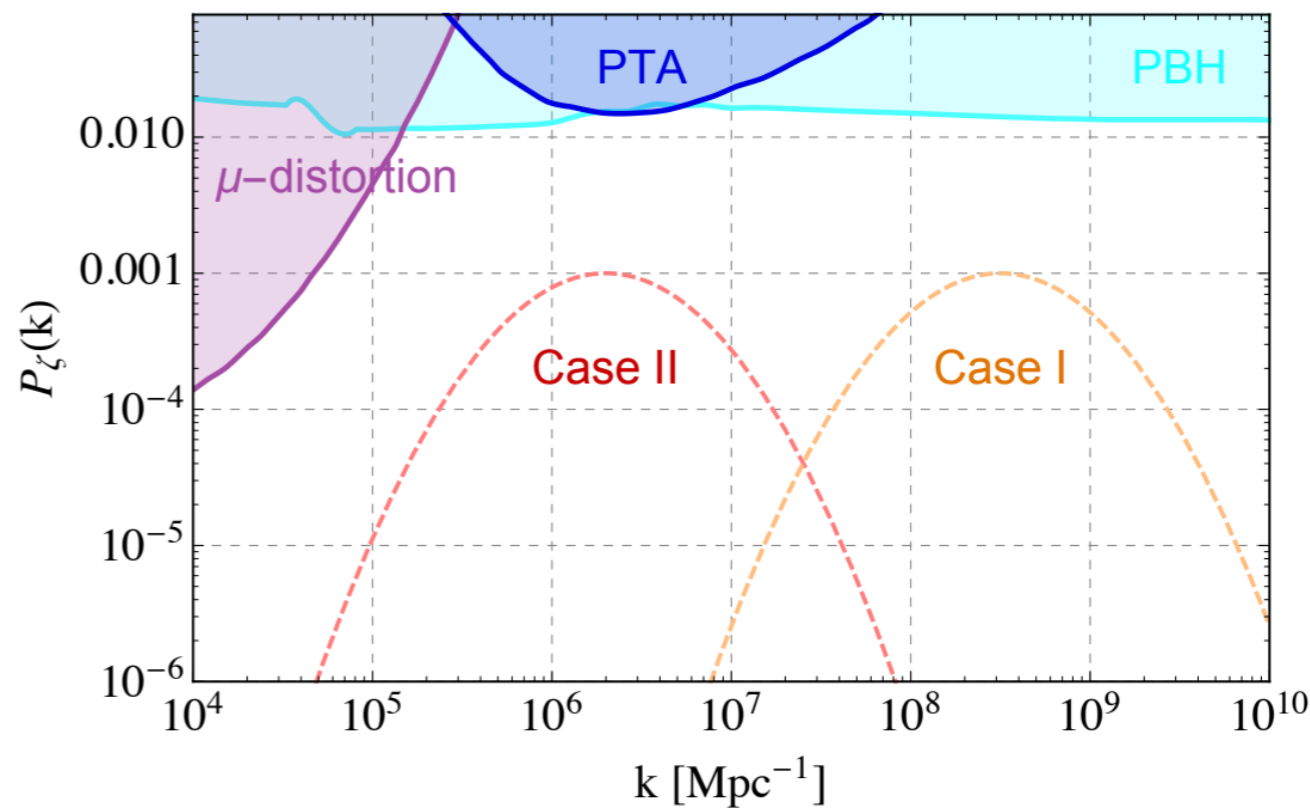
$$\times P_\zeta(k)^{5/2} \times \left(\frac{m_\sigma}{100 \text{ TeV}} \right)^{3/2}$$

$P_\zeta(k)$ on small scales & f_{PBH}

$$P_\zeta(k) = A_s \left(\frac{k}{k_{\text{CMB}}} \right)^{n_s - 1} + A_p e^{-\frac{(N_k - N_p)^2}{2\sigma_p^2}}$$

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$$f_{\text{PBH}}(M) \simeq (2.35 \times 10^5) \times P_\zeta(k)^{5/2} \times \left(\frac{m_\sigma}{100 \text{ TeV}} \right)^{3/2}$$



UCMH (PBH+axion halo) formation

- Axion accretion around (stationary and isolated) PBH
 - UCMH formation
 - efficient during matter-dominated era

E. Bertschinger (APJS 58, 39 (1985))

K. Mack, J. Ostriker, M. Ricotti (Astrophys. J. 665, 1277 (2007))

M. Ricotti, A. Gould (Astrophys. J. 707, 979 (2009))

$$M_{\text{UCMH}}(z) = 3 \left(\frac{1000}{1+z} \right) M_{\text{PBH}},$$

$$R_{\text{UCMH}}(z) = 0.019 \text{ pc} \left(\frac{1000}{z+1} \right) \left(\frac{M_{\text{UCMH}}(z)}{M_{\odot}} \right)^{1/3}$$

$$\rho_{\text{UCMH}}(r) \approx 0.23 M_{\odot} \text{pc}^{-3} \left(\frac{R_{\text{UCMH}}}{r} \right)^{9/4} \left(\frac{10^2 M_{\text{PBH}}}{M_{\text{UCMH}}} \right)^3$$

Axion Tidal Stream

Axion tidal stream (ATS)

- Inside UCMH, $\rho_a > \rho_{\text{DM}}$
 - If disruption happens, this gives axion tidal streams
 - If ρ_a in the tidal stream satisfies $\rho_a > \rho_{\text{DM}}$, this can be of help to direct axion DM searches

P. Tinyakov, I. Tkachev, K. Zioutas (JCAP, 2016)

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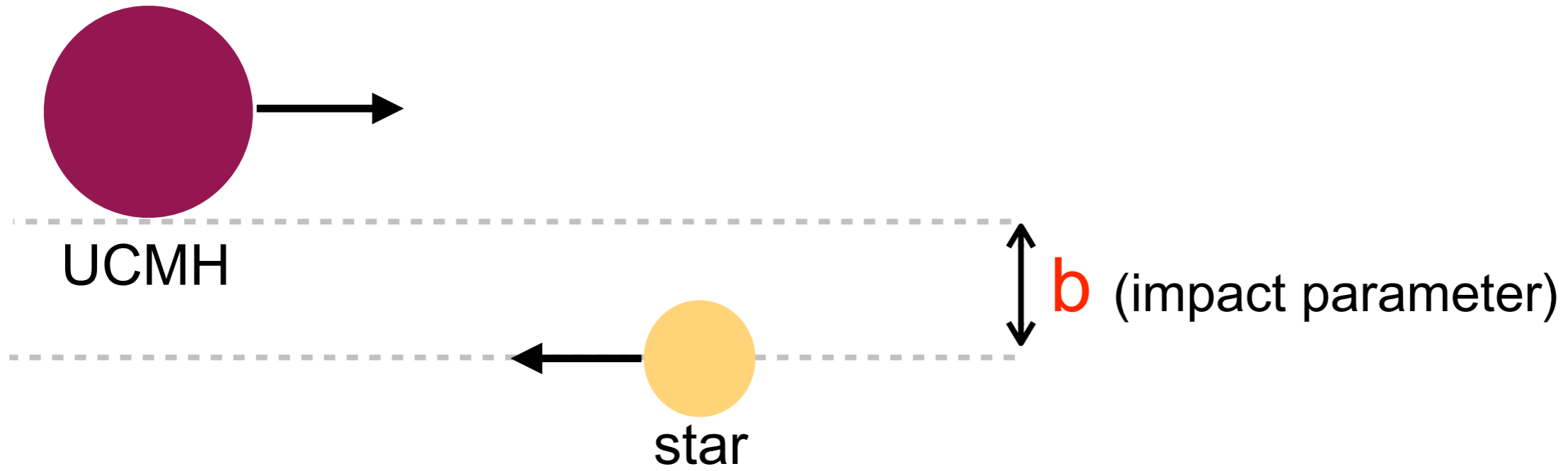
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Q. **For what M_{PBH}** , can ATS be helpful to direct axion DM searches?

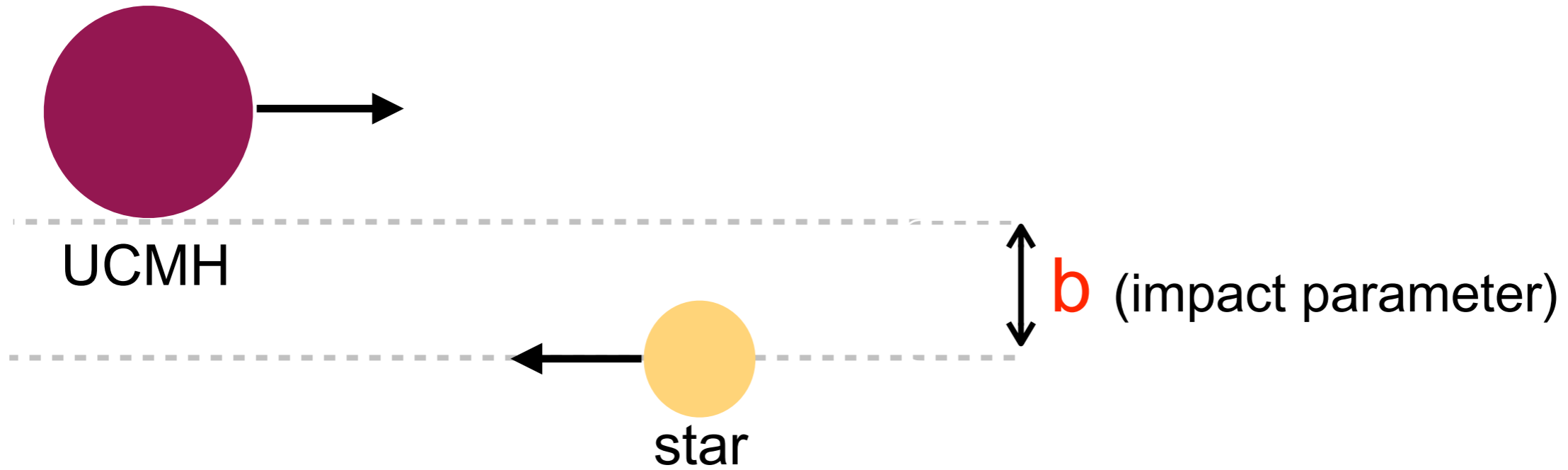
→ need to know

(1) disruption probability & (2) the encounter rate

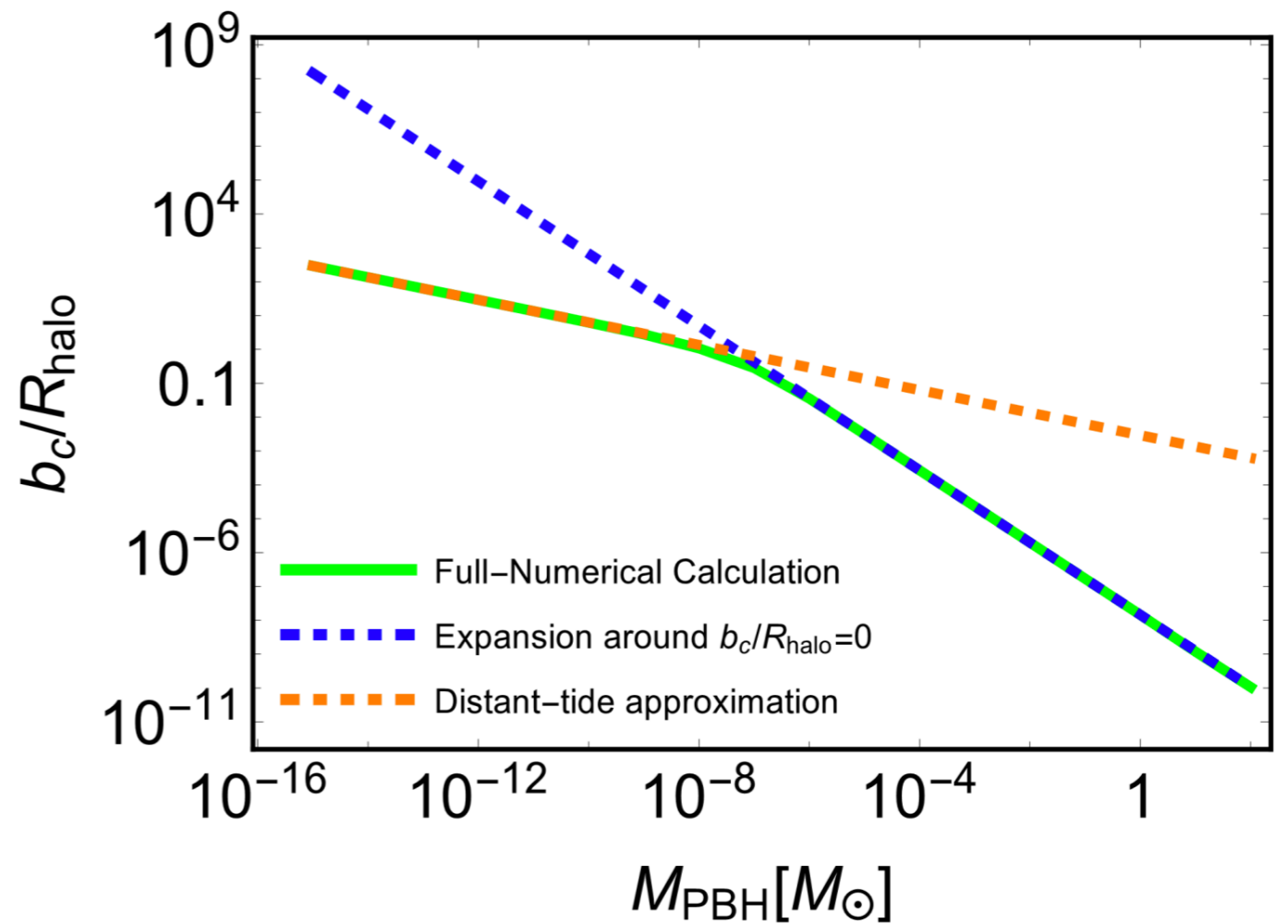
Disruption probability



Disruption probability



- $\Delta E_{\text{int}}(b_c) = E_{\text{binding}}$



Disruption probability

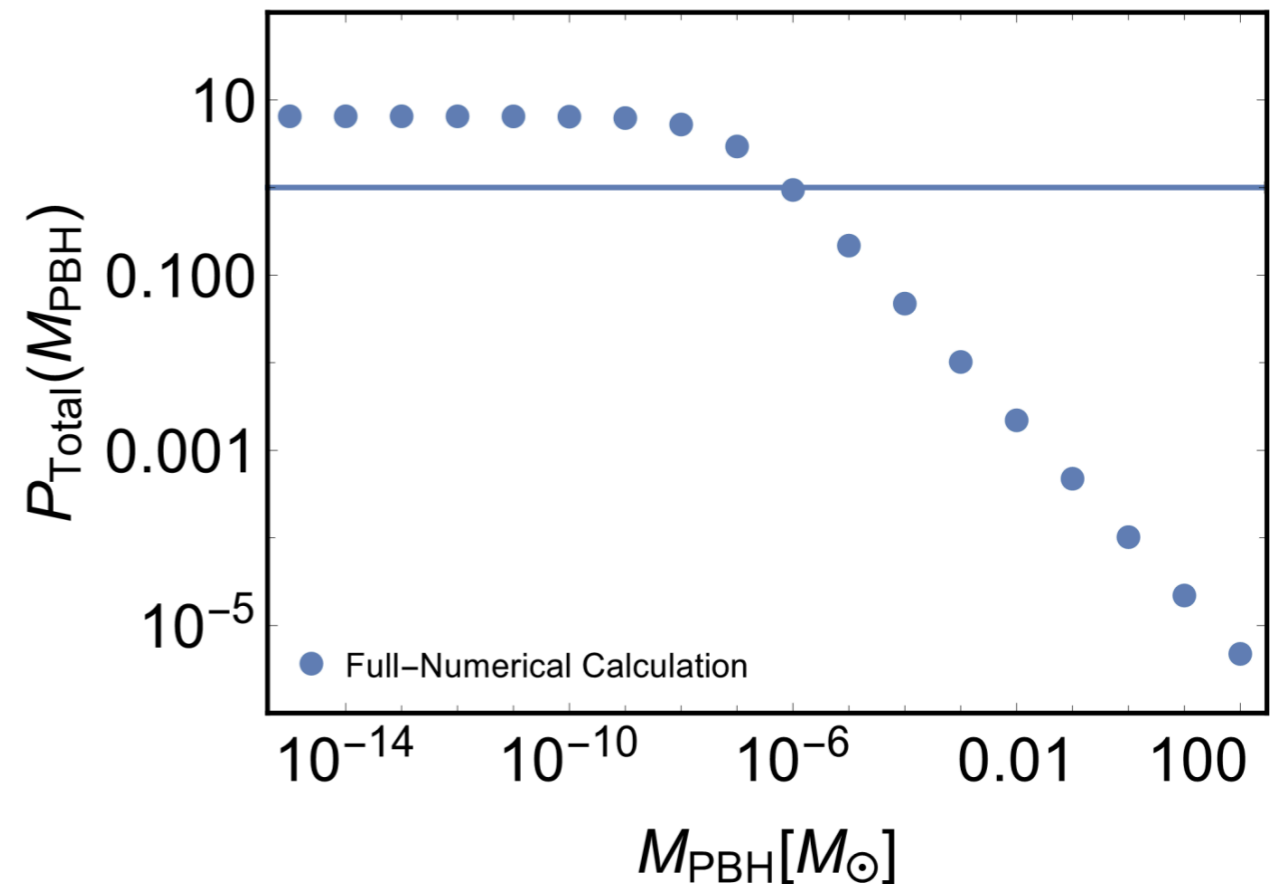
- disruption probability for a single encounter
→ $\Delta E_{\text{int}}/E_{\text{binding}}$
- Total = one-off disruptive events ($b < b_c$)
+ contribution from multiple encounters ($b > b_c$)

$$P_s = 2\pi n_{\star} v_{\text{rel}} t_{\text{cross}} \times \left(\int_0^{b_c} b db + \frac{1}{E_b} \int_{b_c}^{\infty} \Delta E(b) b db \right)$$

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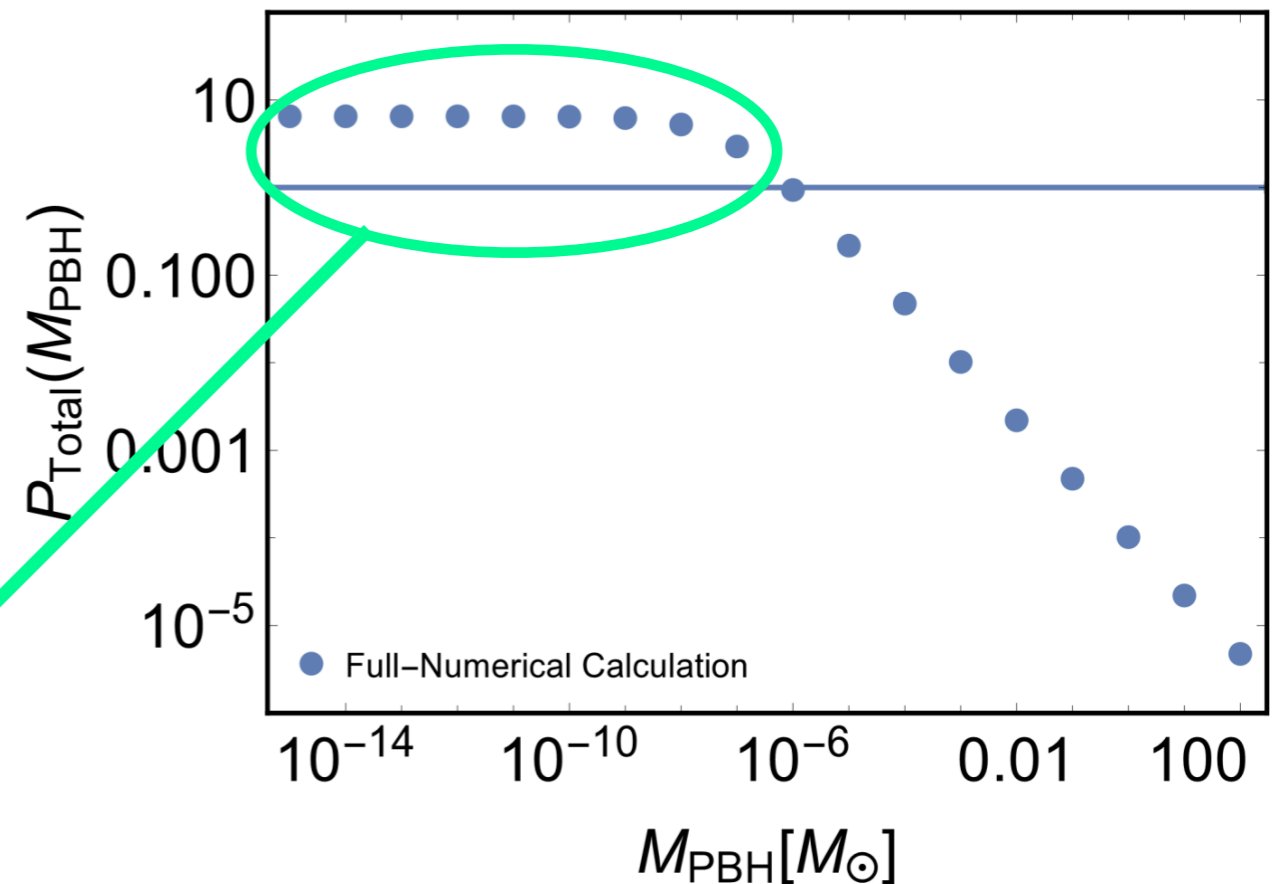
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Total disruption for $M_{\text{PBH}} < 10^{-6} M_{\text{sun}}!$

Encounter rate

- The rate for earth to encounter the axion stream?

$$\Gamma_{\text{st}-\oplus}(r) = \frac{FF(r)}{\tau(r)} = FF(r) \times \frac{v_{\text{rel}}}{2r}$$

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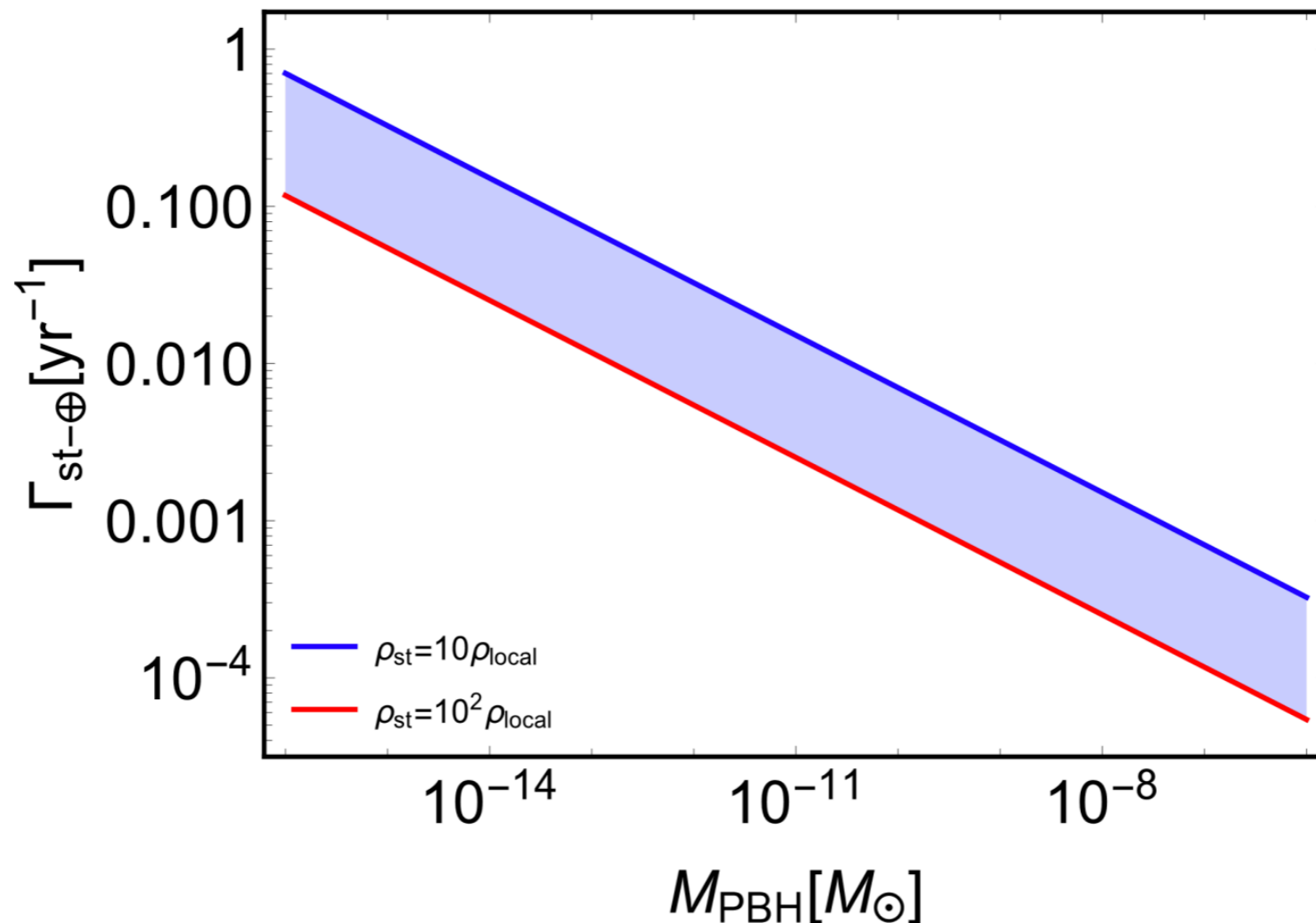
$$\begin{aligned} \text{FF}(r) &= \left(\frac{f_{\text{PBH}} \rho_{\text{local}}}{M_{\text{PBH}}} \right) \times \left(\frac{4\pi r^3 \sigma_a(r) \Delta t}{3 r} \right), \\ &\approx 0.001 \left(\frac{f_{\text{PBH}}}{10^{-2}} \right) \left(\frac{r/R_{\text{UCMH}}}{10^{-2}} \right)^{15/8} \left(\frac{M_{\text{UCMH}}}{10^2 M_{\text{PBH}}} \right)^{5/2} \end{aligned}$$

$$\tau(r) \approx 0.6 \text{ yr} \left(\frac{r/R_{\text{UCMH}}}{10^{-2}} \right) \left(\frac{M_{\text{UCMH}}}{10^2 M_{\text{PBH}}} \right)^{4/3} \left(\frac{M_{\text{PBH}}}{10^{-8} M_{\odot}} \right)^{1/3}$$

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Encounter rate

- Enhancement in ρ_a
→ improve the sensitivity of experimental search

- For instance, in ABRACADABRA

Y. Kahn, B. R. Safdi, J. Thaler (PRL, 2016)

oscillating magnetic field flux through pick-up loop

$$\Phi_a(t) = g_{a\gamma\gamma} |\vec{B}_0| \sqrt{2\bar{\rho}_a} \cos(m_a t)$$

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$100\rho_a \rightarrow$ one order improvement in $g_{a\gamma\gamma}$

Transient Radio Signal

Transient radio signals

- UCMH encounter with a neutron star
→ at r satisfying $m_a = \omega_p$ (conversion radius),
resonant axion conversion to photon

- Power per solid angle

$$\frac{d\mathcal{P}(\theta)}{d\Omega} \approx \frac{g_{a\gamma\gamma}^2 B_0^2 \rho_a(r_c) \pi v_c}{6m_a} \left(\frac{R_{\text{NS}}^2}{r_c} \right)^3 [3\cos^2(\theta) + 1]$$

A. Hook, Y. Kahn, B. R. Safdi, Z. Sun (PRL, 2018)

Transient radio signals

- What matters in terms of detection
→ “spectral flux density”, i.e. $S = (d\mathcal{P}/d\Omega)/(d^2 \Delta\nu)$

$$S(\theta) = \tilde{S} \times \frac{3\cos^2(\theta) + 1}{|3\cos^2(\theta) - 1|^{7/6}}$$

$$\tilde{S} \sim \mu\text{Jy} \left(\frac{\rho_a(r_c)}{0.23 M_\odot \text{pc}^3} \right) \left(\frac{P}{1 \text{ s}} \right)^{7/6} \left(\frac{B_0}{10^{14} \text{ G}} \right)^{5/6} \times$$
$$\left(\frac{g_{a\gamma\gamma}}{10^{-12} \text{ GeV}^{-1}} \right)^2 \left(\frac{R_{\text{NS}}}{10 \text{ km}} \right)^{5/2} \left(\frac{M_{\text{NS}}}{M_\odot} \right)^{1/2} \left(\frac{m_a}{\text{GHz}} \right)^{4/3} \times \left(\frac{100 \text{ pc}}{d} \right)^2 \left(\frac{1 \text{ kHz}}{\Delta\nu} \right)$$

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For example, at $r \sim O(10^{-6}) - O(10^{-5}) R_{\text{UCMH}}, \sim 10^{10} - 10^{11} \rho_{\text{DM}}$

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$$\sim 4 \text{ MHz} \left(\frac{m_a}{\text{GHz}} \right)^{1/3} \left(\frac{\text{s}}{P} \right)^{4/3} \left(\frac{B_0}{10^{14} \text{ G}} \right)^{1/3} \left(\frac{R_{\text{NS}}}{10 \text{ km}} \right) \sin \theta_m$$

Transient radio signals

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We look at $m_a \sim 50\text{MHz} - 1\text{GHz}$ ($F_a = 10^{12} \sim 5 \times 10^{13} \text{ GeV}$)

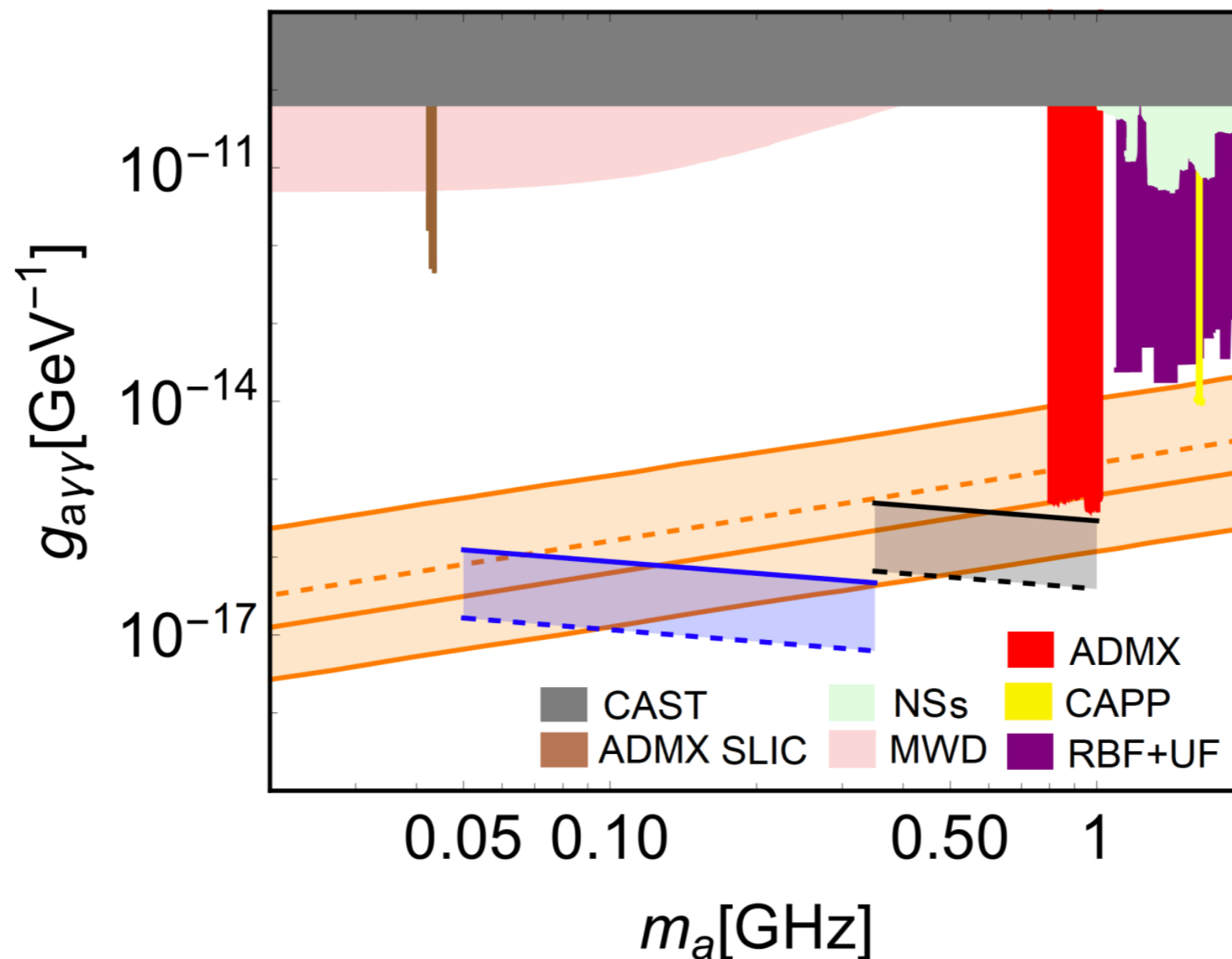
Transient radio signals

- Compare to minimum detectable flux of the chosen radio telescope (SKA1-low, SKA-mid)

$$S_{\min} = \text{SNR}_{\min} \frac{\text{SEFD}}{\eta_s \sqrt{2\Delta B \Delta t_{\text{obs}}}} \quad ,$$
$$\sim 220 \mu\text{Jy} \left(\frac{\text{SNR}_{\min}}{5} \right) \left(\frac{\text{SEFD}}{10 \text{ Jy}} \right) \left(\frac{0.9}{\eta_s} \right) \left(\frac{1 \text{ kHz}}{\Delta B} \right)^{1/2} \left(\frac{1 \text{ yr}}{\Delta t_{\text{obs}}} \right)^{1/2}$$

Transient radio signals

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Summary

- Large decay constant axion DM may imply early matter domination (EMD) era
- EMD era can be the environment facilitating PBH formation
- 0.1-1% PBH + axion DM evolve to UCMH
- $M_{\text{PBH}} \sim 10^{-15} - 10^{-14} M_{\text{sun}}$ can give $O(10)$ enhancement in axion DM energy density via tidal stream
- SKA can probe $m_a \sim 50\text{MHz} - 1\text{GHz}$ ($F_a = 10^{12} \sim 5 \times 10^{13} \text{ GeV}$) via detection of transient radio signal