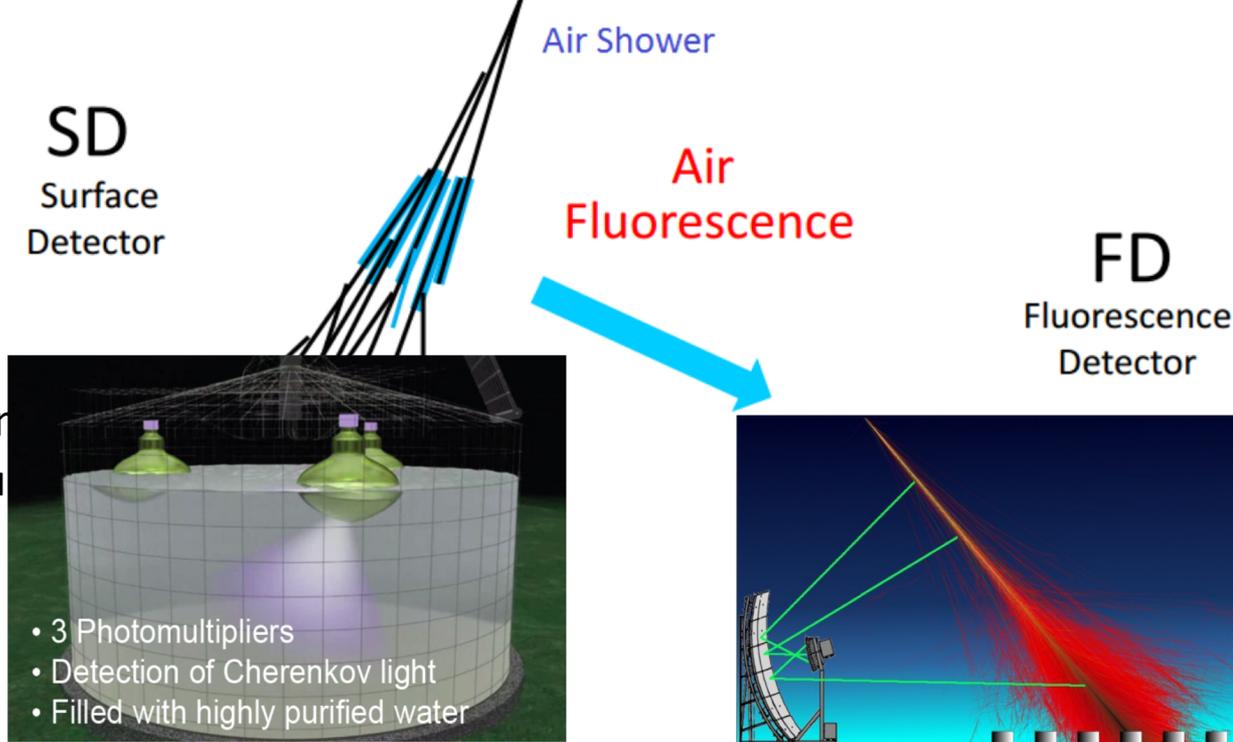
The background of the slide is a night sky featuring a vibrant green aurora borealis. In the foreground, a dark landscape is visible with a prominent red light tower and a white cylindrical structure. The Milky Way galaxy is also visible in the upper right portion of the sky.

# Search for anisotropies in the arrival directions of Ultrahigh Energy Cosmic Rays

Rogério M. de Almeida, for the Pierre Auger Collaboration  
Universidade Federal Fluminense

# Pierre Auger Observatory

- ☐ Particles with  $E > 10^{18}$  eV
- ☐ **Pierre Auger Observatory**
  - ★ Malargue - Mendoza, Argentina
  - ★ Largest UHECR experiment built to date
- ☐ **Main techniques:**
  - **Fluorescence detector (FD)**
  - **Surface detector (SD)**



- 1660 water Cherenkov
- 100% duty cycle
- Lateral profile

- 27 telescopes
- 15% duty cycle
- Longitudinal profile

# The study of anisotropies is important to understand the origin of UHECR

- **Large Scale Anisotropies**
  - Science 357 (2017) 1266
  - Astrophys. J. 868 (2018) 4
  - Last update: R. M. de Almeida [Auger Collab.] ICRC 2021
- **Intermediate-scale anisotropies**
  - ApJ.Lett. 853:L29 (2018)
  - Accepted for publication in ApJS

# Large scale: weighted harmonic analysis

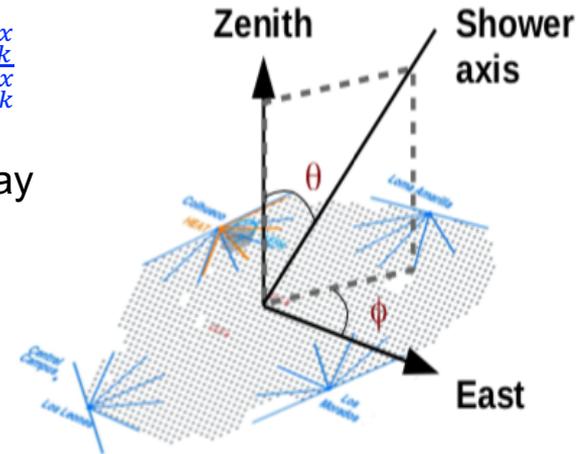
- Search for harmonic modulation in right ascension and azimuth:  $x = \alpha$  or  $\phi$
- Fourier coefficients of order  $k$  (1 or 2):  $a_x^k = \frac{2}{N} \sum_{i=1}^N w_i \cos(kx_i)$ ,  $b_x^k = \frac{2}{N} \sum_{i=1}^N w_i \sin(kx_i)$
- amplitude  $r_k^x = \sqrt{(a_k^x)^2 + (b_k^x)^2}$ , phase  $\varphi_k^x = \frac{1}{k} \tan^{-1} \frac{b_k^x}{a_k^x}$
- weights: small variations in coverage and tilt of the array

$$w_i = [\Delta N_{cell}(\alpha_i^0)(1 + 0.003 \tan \theta_i \cos(\phi_i - \phi_0))]^{-1}$$

number of active detector cells

right ascension of the zenith of the observatory

average tilt of the array  $\phi_0 = -30^\circ$



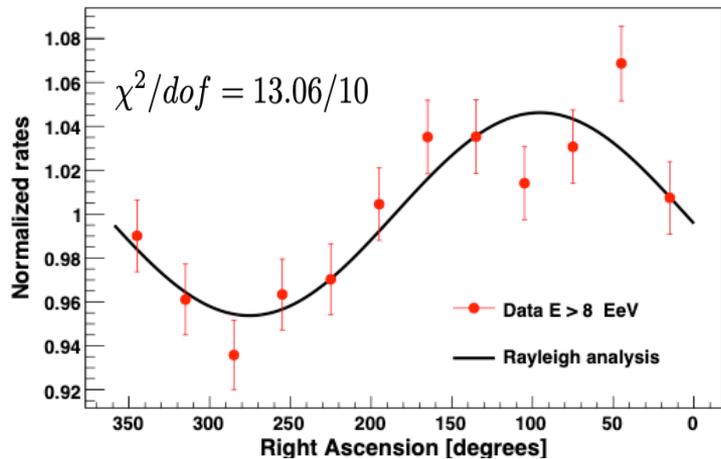
For dipolar modulation,  $d_{\perp} \simeq \frac{r_1^{\alpha}}{\langle \cos \delta \rangle}$  and  $d_z \simeq \frac{b_1^{\phi}}{\langle \sin \theta \rangle \cos l_{obs}}$

# Dipole reconstruction

$E$ (EeV)	$N$	$d_{\perp}$	$d_z$	$d$	$\alpha_d [^{\circ}]$	$\delta_d [^{\circ}]$	$P(\geq r_1^{\alpha})$
4-8	106, 290	$0.01^{+0.006}_{-0.004}$	$-0.012 \pm 0.008$	$0.016^{+0.008}_{-0.005}$	$97 \pm 29$	$-48^{+23}_{-22}$	$1.4 \times 10^{-1}$
8-16	32, 794	$0.055^{+0.011}_{-0.009}$	$-0.03 \pm 0.01$	$0.063^{+0.013}_{-0.009}$	$95 \pm 10$	$-28^{+12}_{-13}$	$3.1 \times 10^{-7}$
16-32	9, 156	$0.072^{+0.021}_{-0.016}$	$-0.07 \pm 0.03$	$0.10^{+0.03}_{-0.02}$	$81 \pm 15$	$-43^{+14}_{-14}$	$7.5 \times 10^{-4}$
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$\geq 32$	2, 448	$0.11^{+0.04}_{-0.03}$	$-0.12 \pm 0.05$	$0.16^{+0.05}_{-0.04}$	$139 \pm 19$	$-47^{+16}_{-15}$	$1.0 \times 10^{-2}$

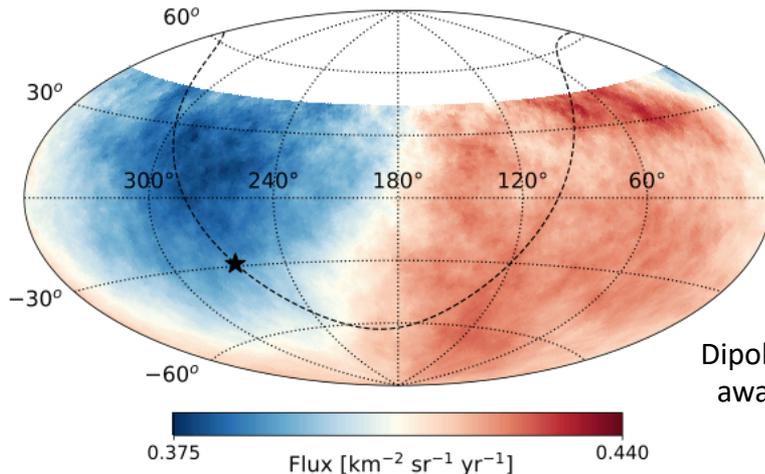
was  $1.4 \times 10^{-9}$  (ApJ 2020) and  $2.6 \times 10^{-8}$  (Science 2017)

Corresponds to  $6.6\sigma$



$E \geq 8$  EeV

Smoothed by a top-hat window with  $45^{\circ}$  of radius



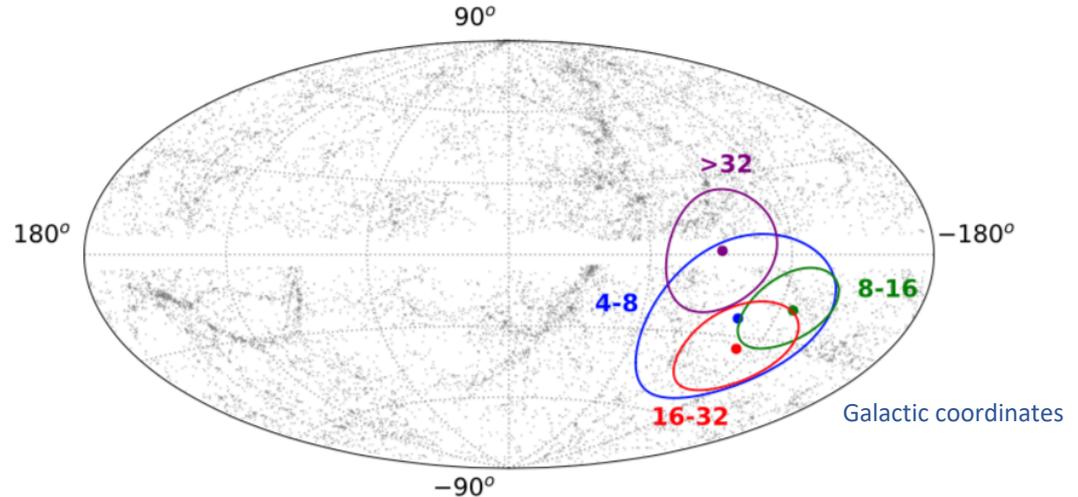
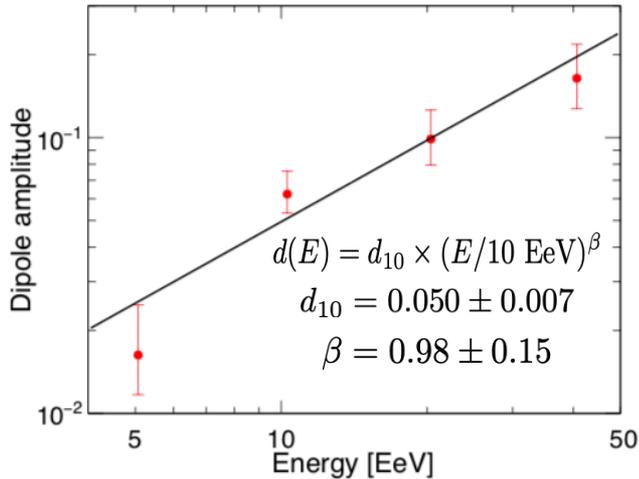
Dipole points  $\sim 115^{\circ}$  away from the GC

# Dipole reconstruction

$E$ (EeV)	$N$	$d_{\perp}$	$d_z$	$d$	$\alpha_d [^{\circ}]$	$\delta_d [^{\circ}]$	$P(\geq r_1^{\alpha})$
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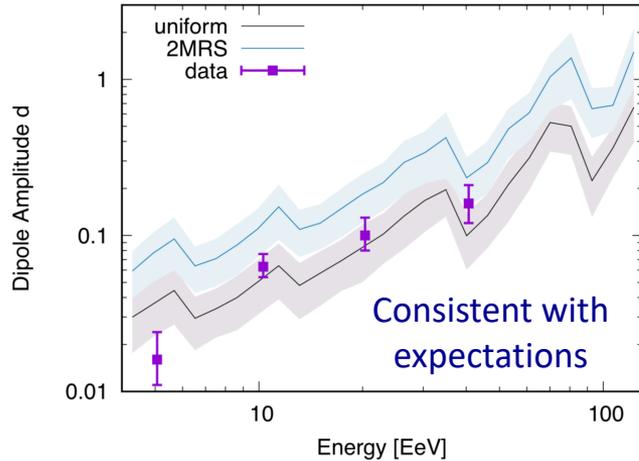
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Corresponds to  $6.6\sigma$

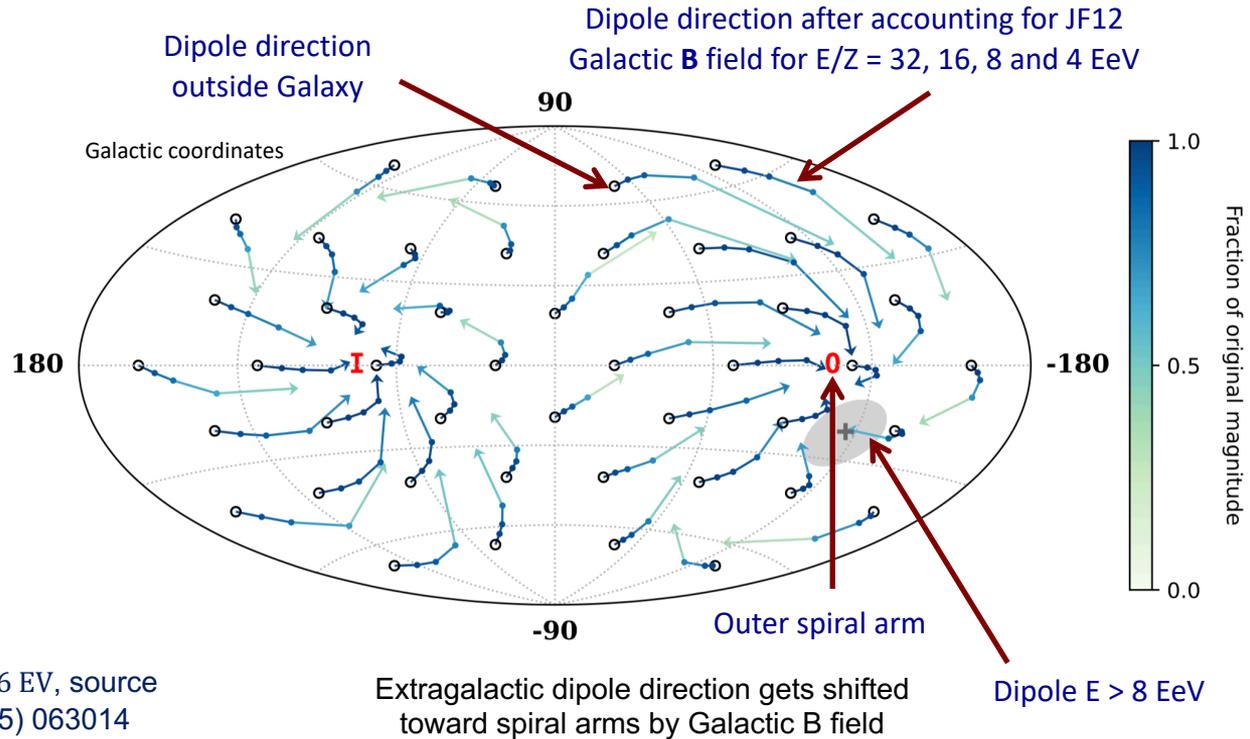


No clear trend in the evolution of dipole direction with energy

# Dipole interpretation



Models with mixed composition,  $R_{max} = 6$  EV, source density  $10^{-4} \text{ Mpc}^{-3}$  from PRD 92 (2015) 063014



**Possibly due to the larger relative contribution from nearby sources to the flux at higher energies: extensive theoretical work (prediction and interpretation)**

Giler et al. (JPhG 1980), Berezhinsky et al. (A&A 1990), Harari et al. (PRD 2014), Harari et al. (PRD 2015), Globus & Piran (ApJL 2017), Wittkowski & Kampert (ApJL 2018), di Matteo et al. (MNRAS 2018), Hackstein et al. (MNRAS 2016), Hackstein et al. (MNRAS 2018), Ding et al. (ApJL 2021)

# Search for Intermediate scale anisotropies

## Analysis strategy

- Sky model probability maps:

- Null hypothesis  $H_0$ : isotropy  $n^{H_0}(\mathbf{u}) = \frac{\omega(\mathbf{u})}{\sum_i \omega(\mathbf{u}_i)}$

- Single population signal model  $H_1$ :

$$s_j(\mathbf{u}; \Theta) = \omega(\mathbf{u}) \times \phi_j a(d_j) \times \exp\left(\frac{\mathbf{u} \cdot \mathbf{u}_j}{2(1 - \cos \Theta)}\right)$$

$\Phi_i$  = flux model x attenuat. model x ang. smearing ( $\Theta$ ) x exposure

$$n^{H_1}(\mathbf{u}) = (1 - \alpha) \times n^{H_0}(\mathbf{u}) + \alpha \times \frac{\sum_j s_j(\mathbf{u}; \Theta)}{\sum_i \sum_j s_j(\mathbf{u}_i; \Theta)}$$

(free parameters:  $\alpha$  and  $\Theta$ )

- Test statistics:  $TS = 2 \log(H_1/H_0)$   $\longrightarrow$   $TS = 2 \sum_i k_i \times \ln \frac{n^{H_1}(\mathbf{u}_i)}{n^{H_0}(\mathbf{u}_i)}$

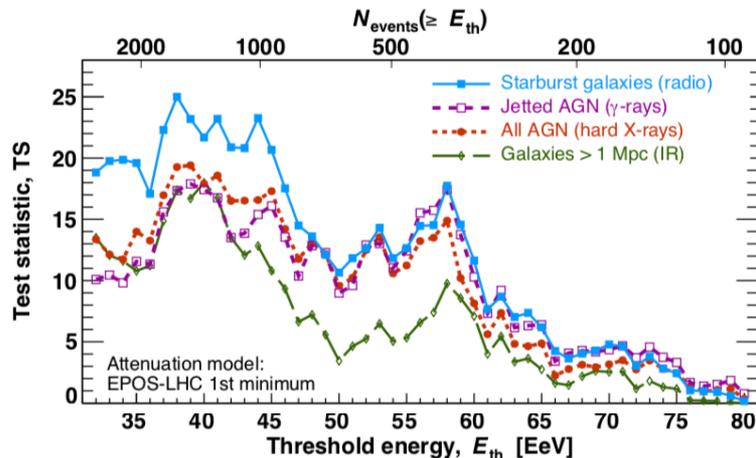
- p-value from Wilk's theorem:  $p(TS) = p_{\chi^2}(TS, ndf)$

# From catalogs to UHECR sky models

We study four cosmic rays sky models obtained from:

- **Near-infrared emission of galaxies (2MASS)**
  - Assuming UHECR luminosity proportional to stellar mass
- **Radio emission from starburst galaxies**
  - Assuming UHECR luminosity proportional to starforming activity
- **X-rays from AGNs (Swift-BAT)**
  - Assuming UHECR luminosity is driven by accretion onto super-massive black holes
- **$\gamma$ -rays from jetted AGNs (Fermi-LAT)**
  - Assuming UHECR luminosity is proportional to the inner jet activity

# Results

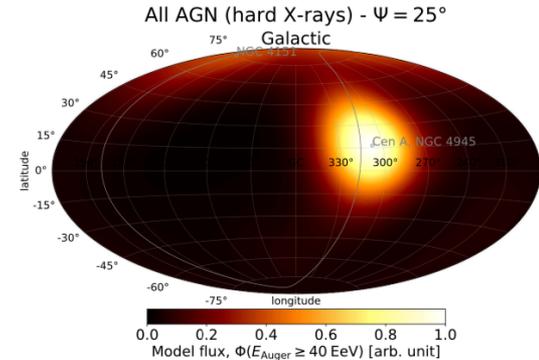
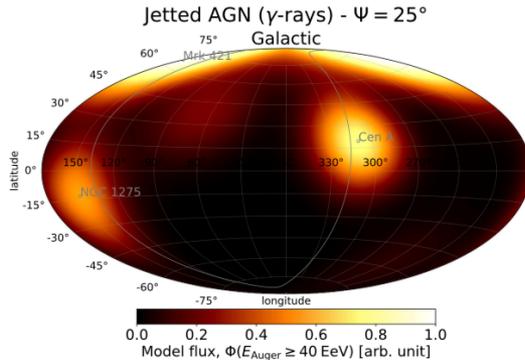
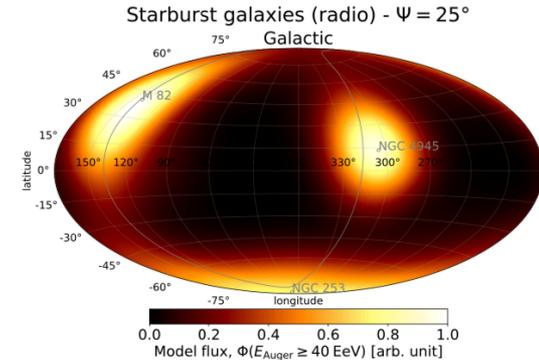
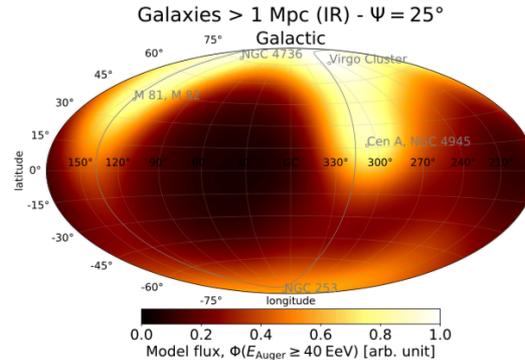
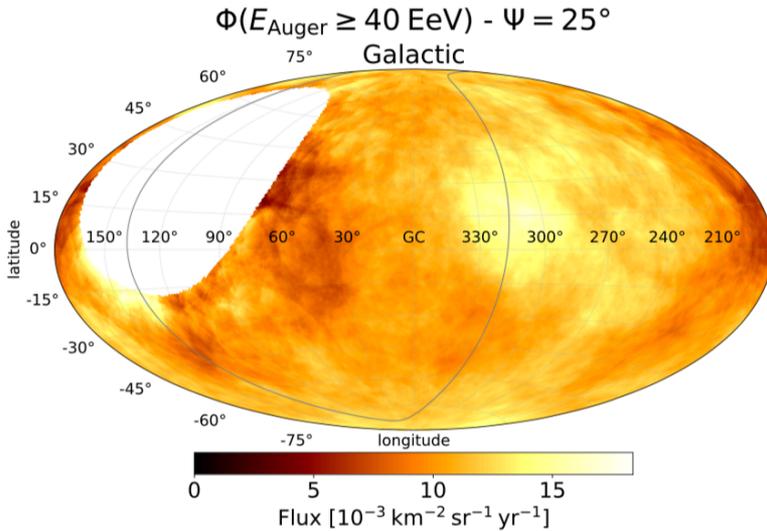


Catalog	$E_{\text{th}}$ [EeV]	Fisher search radius, $\Theta$ [deg]	Signal fraction, $\alpha$ [%]	$\text{TS}_{\text{max}}$	Post-trial $p$ -value
All galaxies (IR)	40	$16^{+11}_{-6}$	$16^{+10}_{-6}$	18.0	$7.9 \times 10^{-4}$
Starbursts (radio)	38	$15^{+8}_{-4}$	$9^{+6}_{-4}$	25.0	$3.2 \times 10^{-5}$
All AGNs (X-rays)	39	$16^{+8}_{-5}$	$7^{+5}_{-3}$	19.4	$4.2 \times 10^{-4}$
Jetted AGNs ( $\gamma$ -rays)	39	$14^{+6}_{-4}$	$6^{+4}_{-3}$	17.9	$8.3 \times 10^{-4}$
All galaxies (IR)	58	$14^{+9}_{-5}$	$18^{+13}_{-10}$	9.8	$2.9 \times 10^{-2}$
Starbursts (radio)	58	$18^{+11}_{-6}$	$19^{+20}_{-9}$	17.7	$9.0 \times 10^{-4}$
All AGNs (X-rays)	58	$16^{+8}_{-6}$	$11^{+7}_{-6}$	14.9	$3.2 \times 10^{-3}$
Jetted AGNs ( $\gamma$ -rays)	58	$17^{+8}_{-5}$	$12^{+8}_{-6}$	17.4	$1.0 \times 10^{-3}$

# Results

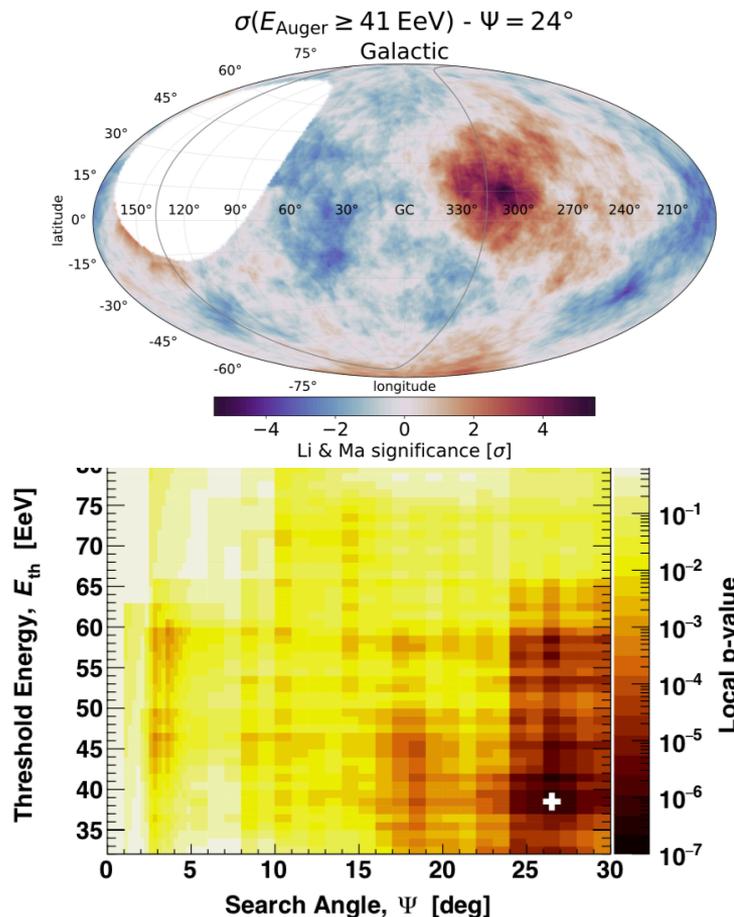
## Best fit sky models above 40 EeV

### Flux Map above 40 EeV

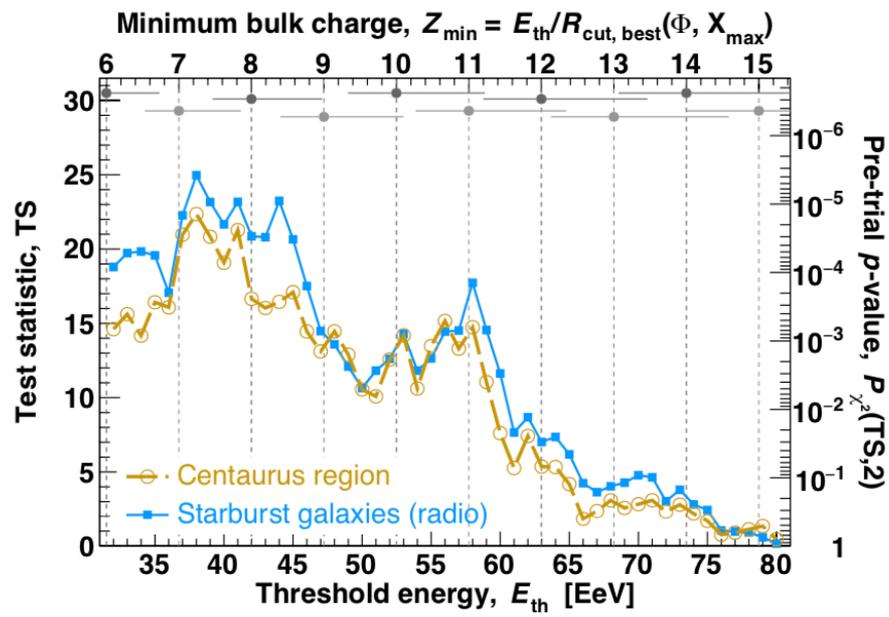


Note: sky maps do not include any isotropic component

# The Centaurus Region



$\log_{10}(R_{\text{cut}}/V) \approx 18.7$  from Auger best fit (2017)



# Conclusions

## Large-scale anisotropies:

- The statistical significance of the large-scale dipolar modulation observed above 8 EeV has increased to  $6.6\sigma$ , (p-value =  $5.1 \times 10^{-11}$ ).
- The dipolar amplitude increases with energies although there is no clear trend in the change of the dipole direction as a function of energy.
  - Possibly due to the larger relative contribution from nearby sources to the flux at higher energies; Interpretation of the rec. dipole directions requires taking into account the magnetic deflections of the particles during their trajectories.

## Intermediate-scale anisotropies:

- Evidence for a deviation in excess of isotropy at intermediate angular scale with  $\sim 15^\circ$  Gaussian spread or  $25^\circ$  top-hat radius obtained at the  $4\sigma$  significance level for cosmic-ray energies above  $\sim 40$  EeV

# Backup

## Dipolar + quadrupolar reconstruction

Energy [EeV]	$d_i$	$Q_{ij}$
4-8	$d_x = -0.008 \pm 0.007$	$Q_{zz} = 0.008 \pm 0.036$
	$d_y = 0.008 \pm 0.007$	$Q_{xx} - Q_{yy} = 0.004 \pm 0.026$
	$d_z = -0.008 \pm 0.021$	$Q_{xy} = -0.01 \pm 0.01$
		$Q_{xz} = -0.02 \pm 0.02$
		$Q_{yz} = -0.008 \pm 0.017$
8-16	$d_x = -0.005 \pm 0.013$	$Q_{zz} = 0.074 \pm 0.064$
	$d_y = 0.045 \pm 0.013$	$Q_{xx} - Q_{yy} = 0.02 \pm 0.05$
	$d_z = 0.01 \pm 0.04$	$Q_{xy} = 0.039 \pm 0.024$
		$Q_{xz} = -0.002 \pm 0.031$
		$Q_{yz} = -0.03 \pm 0.03$

Quadrupolar amplitudes not significant and dipole components consistent with those obtained by assuming a pure dipole

Energy [EeV]	$d_i$	$Q_{ij}$
16-32	$d_x = 0.05 \pm 0.02$	$Q_{zz} = -0.14 \pm 0.14$
	$d_y = 0.09 \pm 0.02$	$Q_{xx} - Q_{yy} = 0.17 \pm 0.09$
	$d_z = -0.15 \pm 0.07$	$Q_{xy} = -0.05 \pm 0.04$
		$Q_{xz} = 0.12 \pm 0.06$
		$Q_{yz} = 0.06 \pm 0.06$
$\geq 32$	$d_x = -0.12 \pm 0.05$	$Q_{zz} = -0.17 \pm 0.26$
	$d_y = 0.11 \pm 0.05$	$Q_{xx} - Q_{yy} = 0.43 \pm 0.17$
	$d_z = -0.22 \pm 0.13$	$Q_{xy} = 0.10 \pm 0.09$
		$Q_{xz} = -0.12 \pm 0.11$
		$Q_{yz} = 0.13 \pm 0.11$
$\geq 8$	$d_x = -0.001 \pm 0.011$	$Q_{zz} = 0.02 \pm 0.06$
	$d_y = 0.06 \pm 0.01$	$Q_{xx} - Q_{yy} = 0.08 \pm 0.04$
	$d_z = -0.03 \pm 0.03$	$Q_{xy} = 0.02 \pm 0.02$
		$Q_{xz} = 0.02 \pm 0.03$
		$Q_{yz} = -0.003 \pm 0.026$

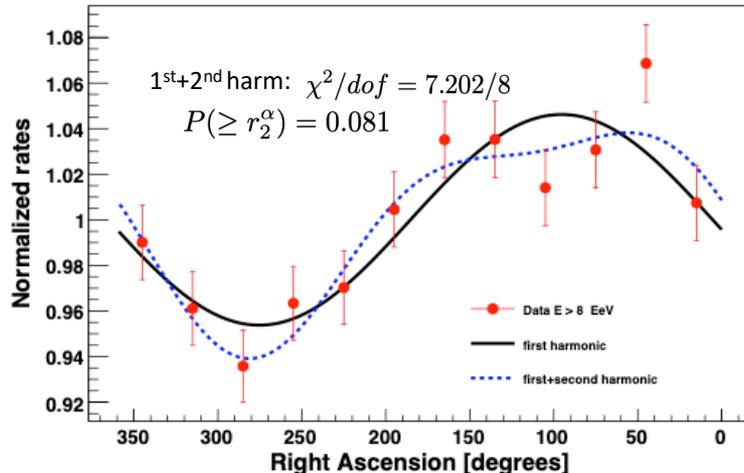
Quadrupolar amplitudes not significant and dipole components consistent with those obtained by assuming a pure dipole

# Dipole reconstruction

$E$ (EeV)	$N$	$d_{\perp}$	$d_z$	$d$	$\alpha_d [^{\circ}]$	$\delta_d [^{\circ}]$	$P(\geq r_1^{\alpha})$
4-8	106, 290	$0.01^{+0.006}_{-0.004}$	$-0.012 \pm 0.008$	$0.016^{+0.008}_{-0.005}$	$97 \pm 29$	$-48^{+23}_{-22}$	$1.4 \times 10^{-1}$
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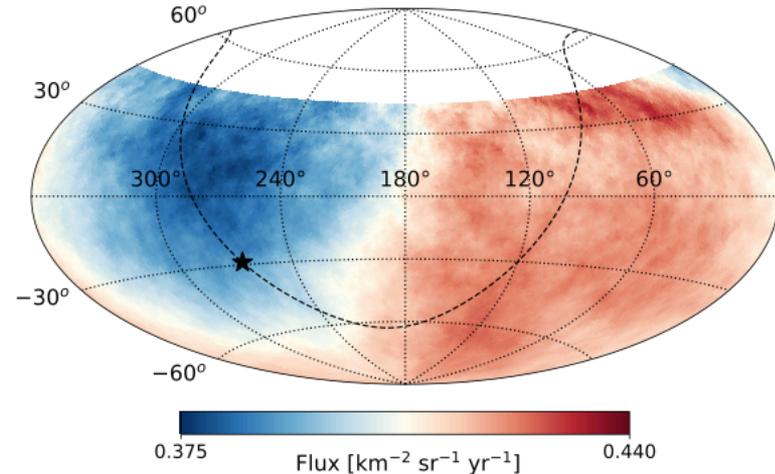
was  $1.4 \times 10^{-9}$  (ApJ 2020) and  $2.6 \times 10^{-8}$  (Science 2017)

Corresponds to  $6.6\sigma$

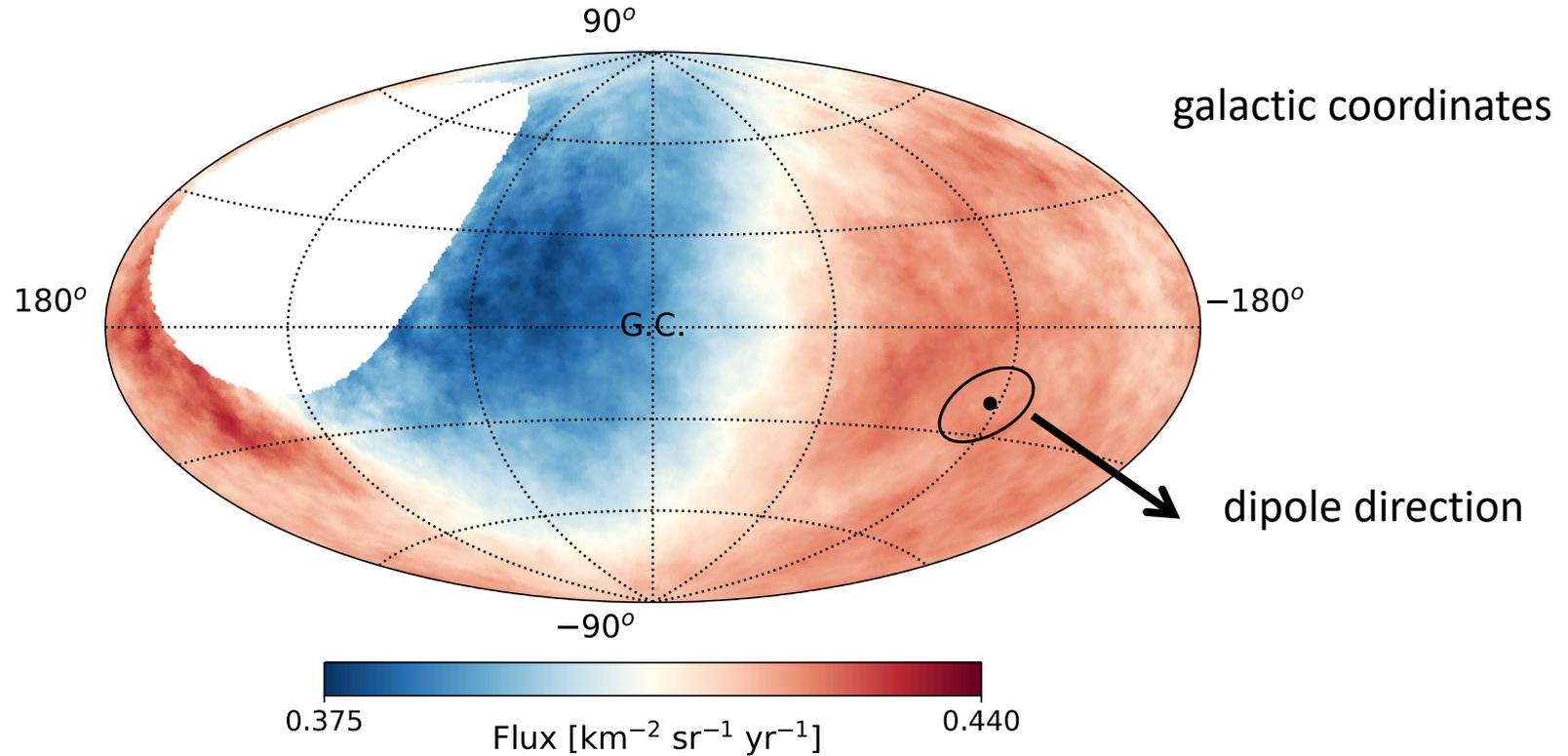


$E \geq 8$  EeV

Smoothed by a top-hat window with  $45^{\circ}$  of radius



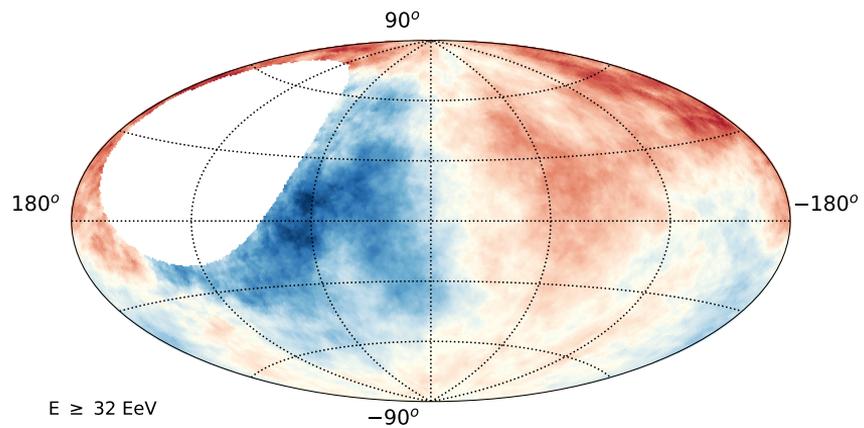
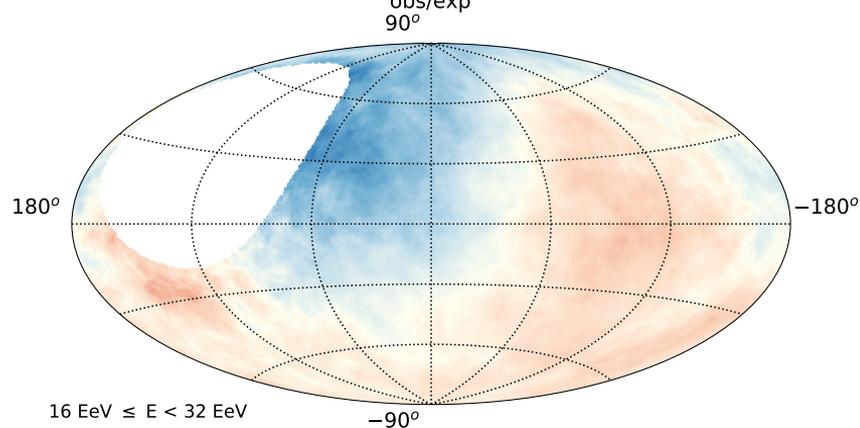
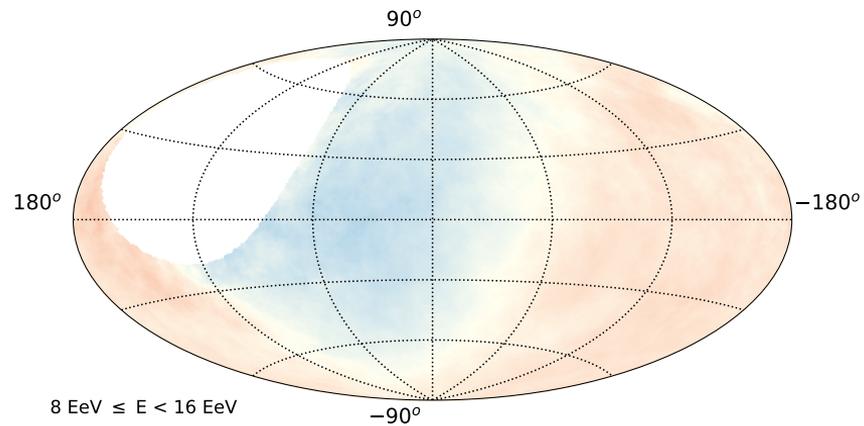
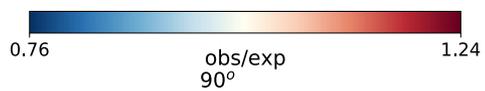
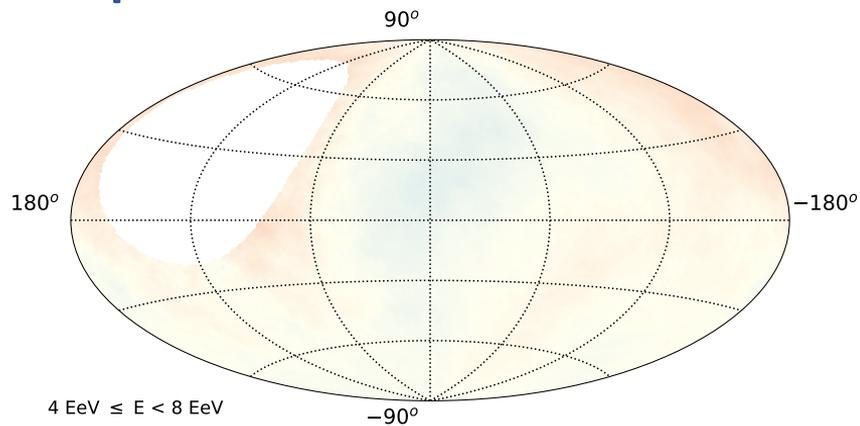
# Dipole reconstruction $E > 8$ EeV



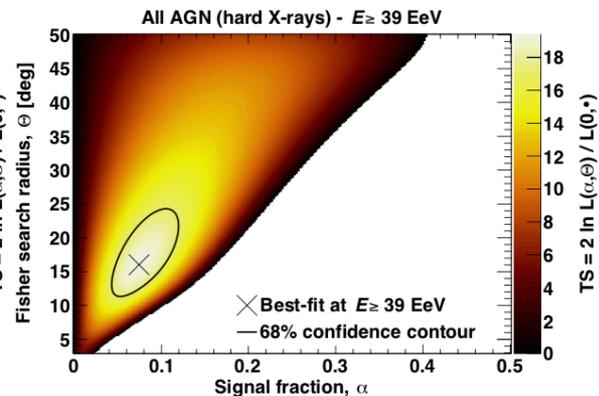
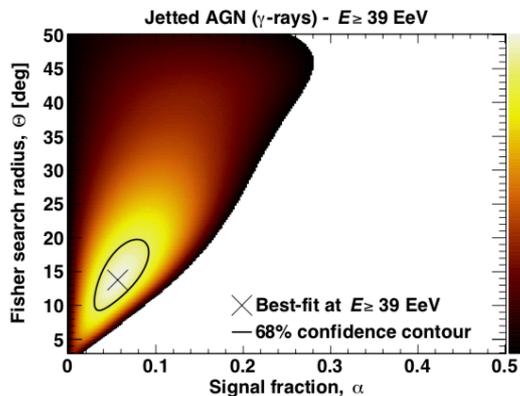
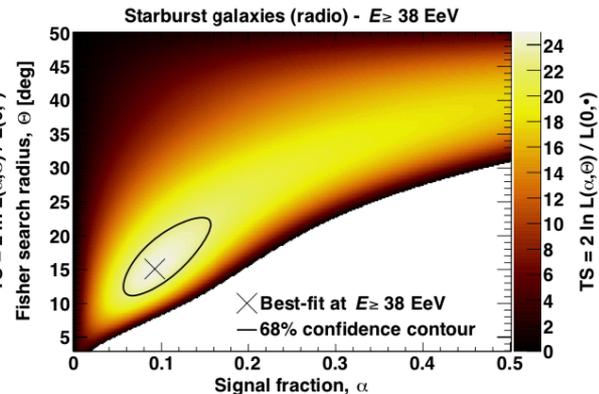
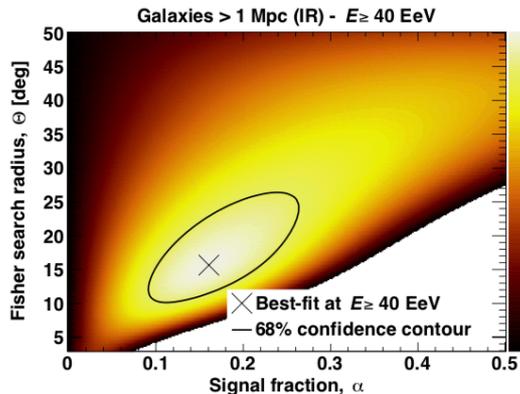
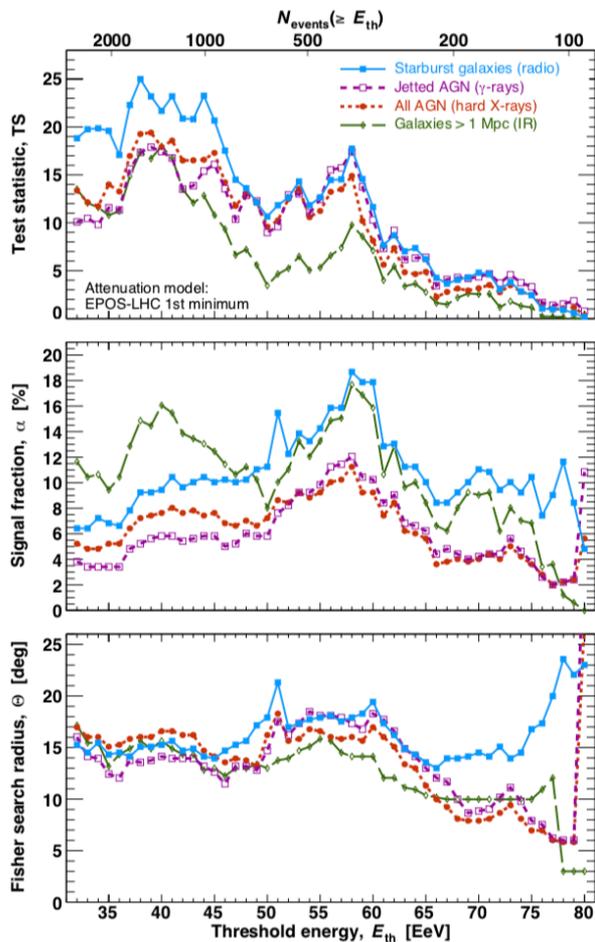
Dipole points to  $(l, b) = (-117^\circ, -21^\circ) \sim 115^\circ$  away from the direction of the Galactic center

# Dipole reconstruction

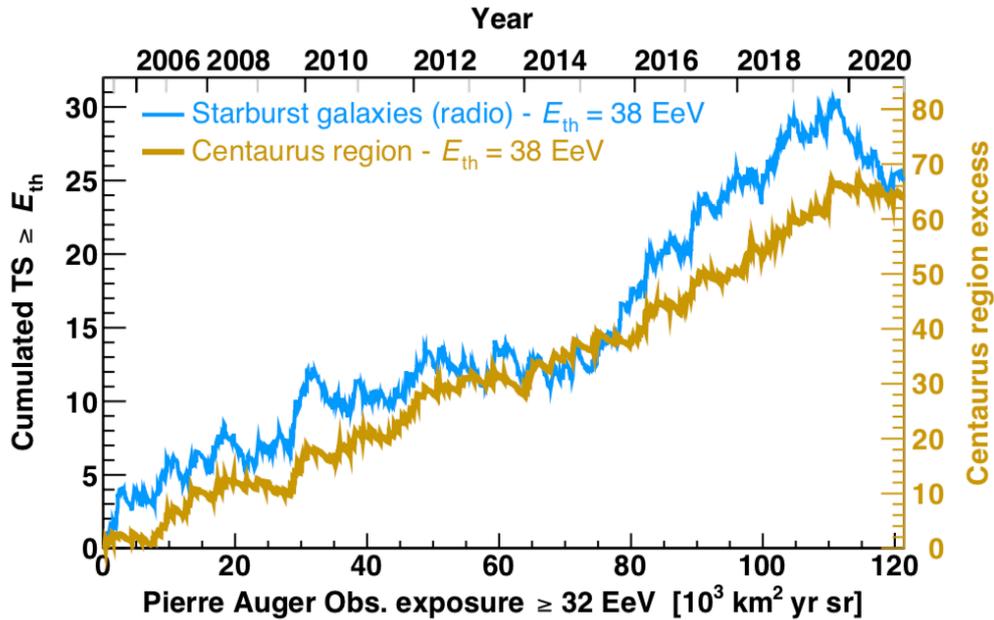
Ratio between observed and expected events in windows of  $45^\circ$



# Results



# Evolution of test statistic



Assuming a fixed top-hat angular scale  $\Psi = 27^\circ$  and a continued growth of the excess, the  $5\sigma$  discovery threshold would be within reach by the end of 2025 ( $\pm 2$  calendar years)

