

# Dark Matter Axions in the Early Universe with a Period of Increasing Temperature

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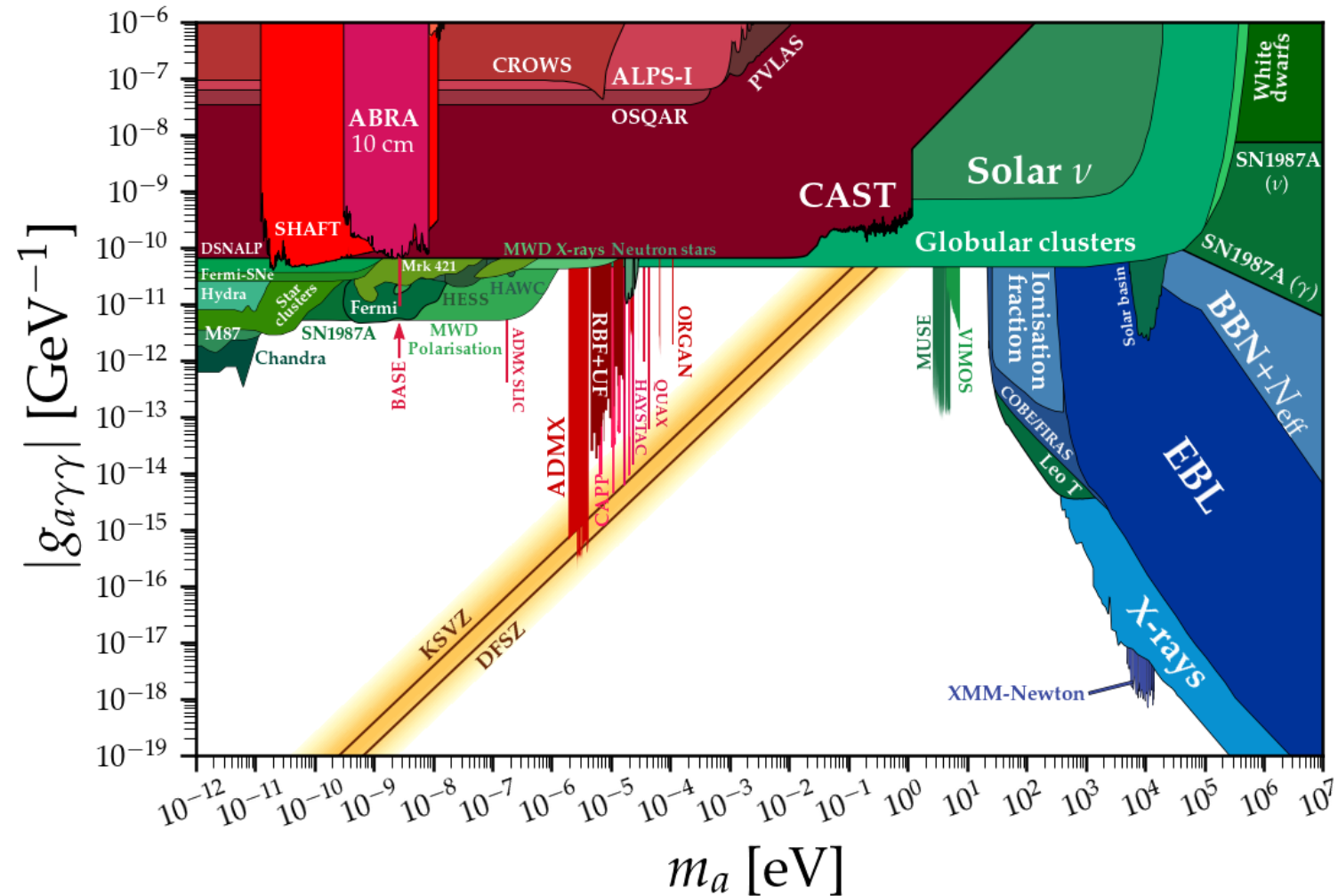
## ASTROCENT



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# Introduction

- QCD axions are a promising dark matter (DM) candidate
- Many ongoing observational/experimental searches targeting a variety of couplings
- Part of their success as DM candidates is the misalignment mechanism



# Axion misalignment mechanism

- Initial value of angle  $\theta$  fixed after Peccei-Quinn (PQ) breaking at a high scale  $f_a$
- Axion field ( $a$ ) frozen as long as Hubble rate  $>$  axion mass

(zero-temp.  
axion mass)

$$\theta(t) \equiv \frac{a(t)}{f_a}$$

$$m_a \approx 5.7 \text{ meV} \left( \frac{10^9 \text{ GeV}}{f_a} \right)$$

$$T_{\text{QCD}} \approx 150 \text{ MeV}$$

Hubble rate:

$$H(T) \propto \frac{T^2}{M_{\text{P}}}$$

(radiation domination)

Axion mass:

$$m(T) \approx m_a \begin{cases} \left( \frac{T_{\text{QCD}}}{T} \right)^4 & T > T_{\text{QCD}} \\ 1 & T < T_{\text{QCD}} \end{cases}$$

# Axion misalignment mechanism

- As temperature of Universe cools, axion mass increases while Hubble rate drops

- Axion oscillation begins:

$$\text{“crossing” condition:} \\ 3 H(T_{\text{osc}}) \approx m(T_{\text{osc}})$$

- Energy density averages to matter  $\rightarrow$  “standard mass window” for correct DM relic abundance assuming standard RD history:

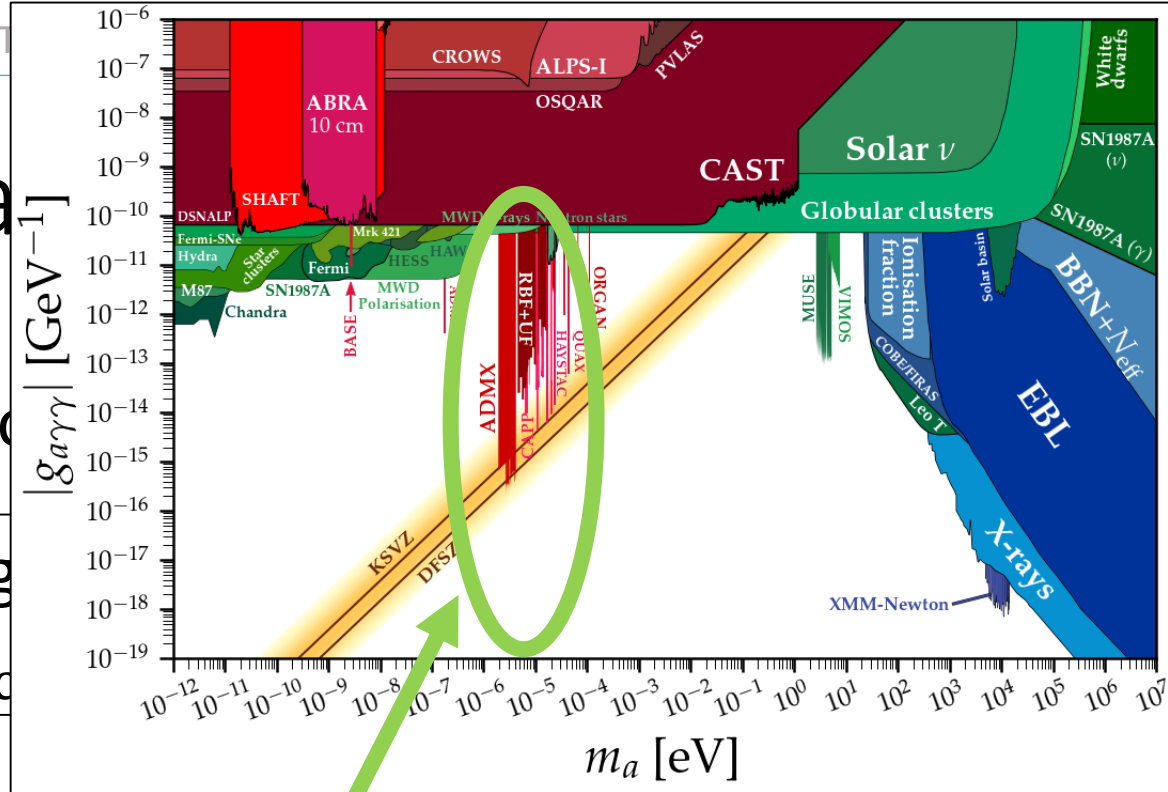
$$10^{-6} \text{ eV} \lesssim m_a \lesssim 10^{-5} \text{ eV} \quad \text{for} \quad 0.5 \lesssim \theta_i \lesssim \pi/\sqrt{3}$$

- Notice that this mechanism depends on thermal history  
 $\rightarrow$  nonstandard cosmologies (NSCs) can alter axion production

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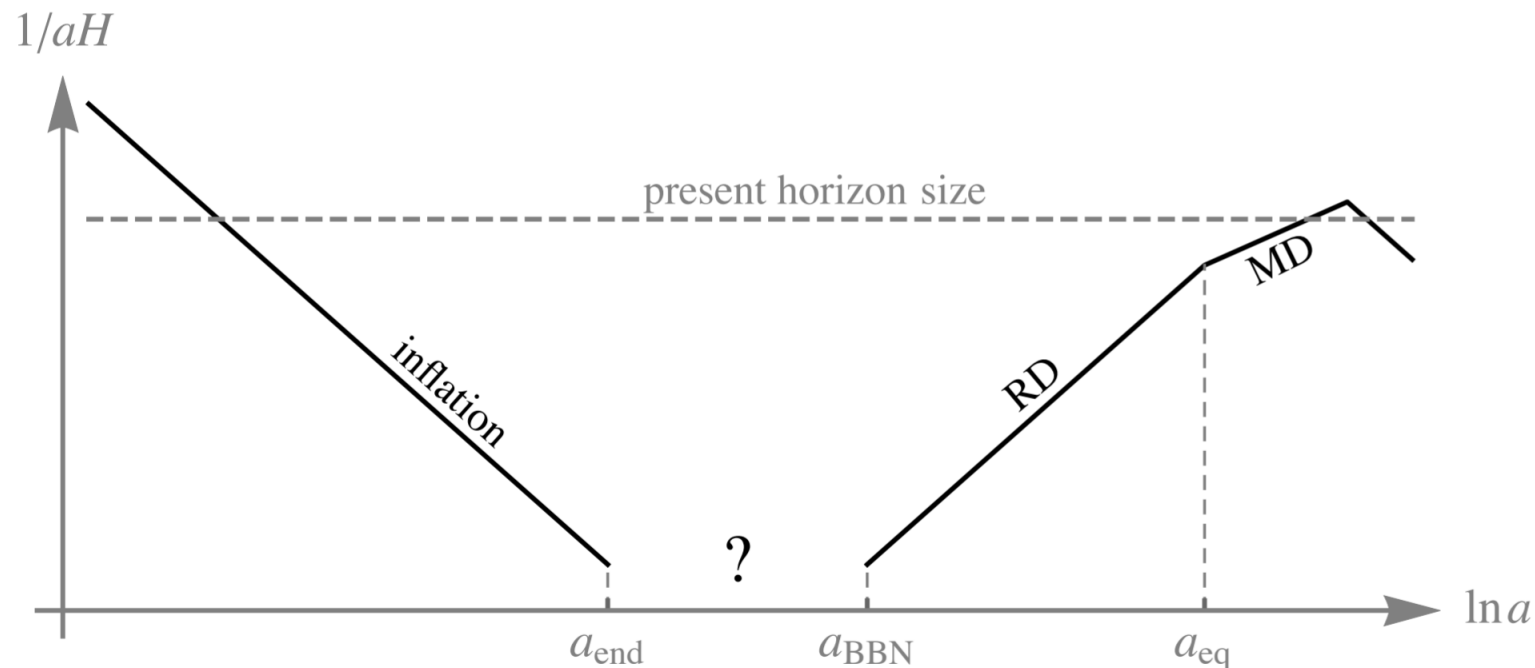
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- Notice that this mechanism depends on thermal history  $\rightarrow$  nonstandard cosmologies (NSCs) can alter axion production

# Nonstandard cosmologies (NSCs)

- Usual assumption is radiation domination (RD) between end of inflation and start of BBN
- Many well-motivated deviations: inflationary reheating, early matter domination (MD) from moduli, kination, etc. (for a review, see Allahverdi et al. 2006.16182)
- Must end before  $T_{\text{BBN}} \sim \text{MeV}$
- Can significantly affect early processes



# NSCs – increasing temperature

- Typically temperature decreases during NSC (though can be at a different rate)
- Particularly interesting class of alternate histories involves a period of increasing temperature
- Can be achieved in a variety of models where the decay rate of dominating component has explicit scale-factor ( $R$ ) or temperature ( $T$ ) dependence:  $\Gamma(R, T)$
- Examples: higher-order operators, shapes of potentials, field-dependent mass, interaction with thermal background

Co et al. 2007.04328; Ahmed et al. 2111.06065;

Garcia et al. 2012.10756; Mukaida+Nakayama 1208.3399

# Our scenario

- NSC history with intervening matter domination (MD) by scalar field  $\phi$ :  
Inflation – RD – MD – RD (BBN)
- Period of increasing temperature from time-dependent decay rate  $\Gamma_\phi$
- Axion production from misalignment in this background

- Define parameter  $x$  during increasing-temperature period:

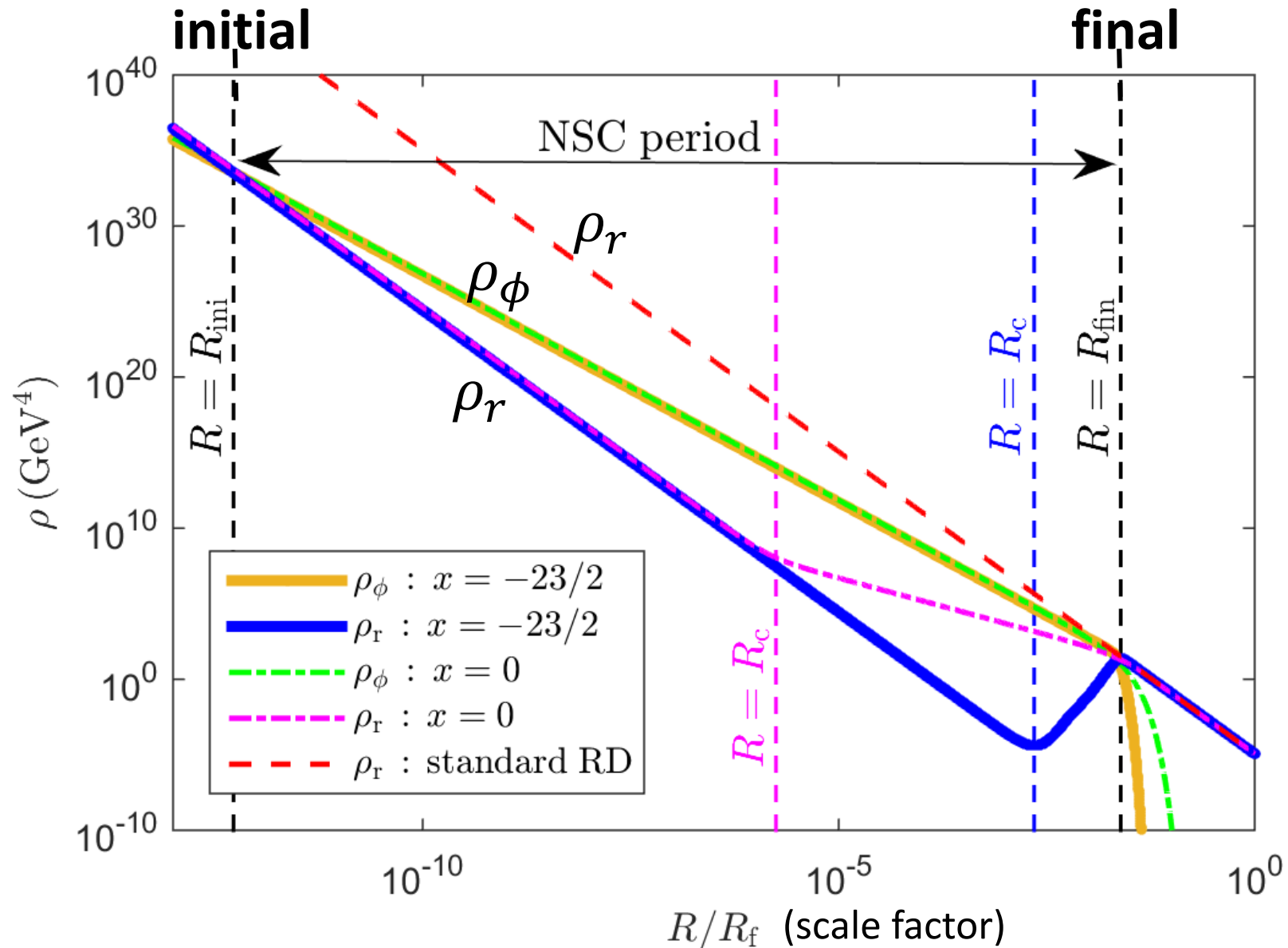
$$\Gamma_\phi \propto R^{-x} \qquad T \propto R^{\frac{-3-2x}{8}}$$

- Temperature increases for  $x < -3/2$

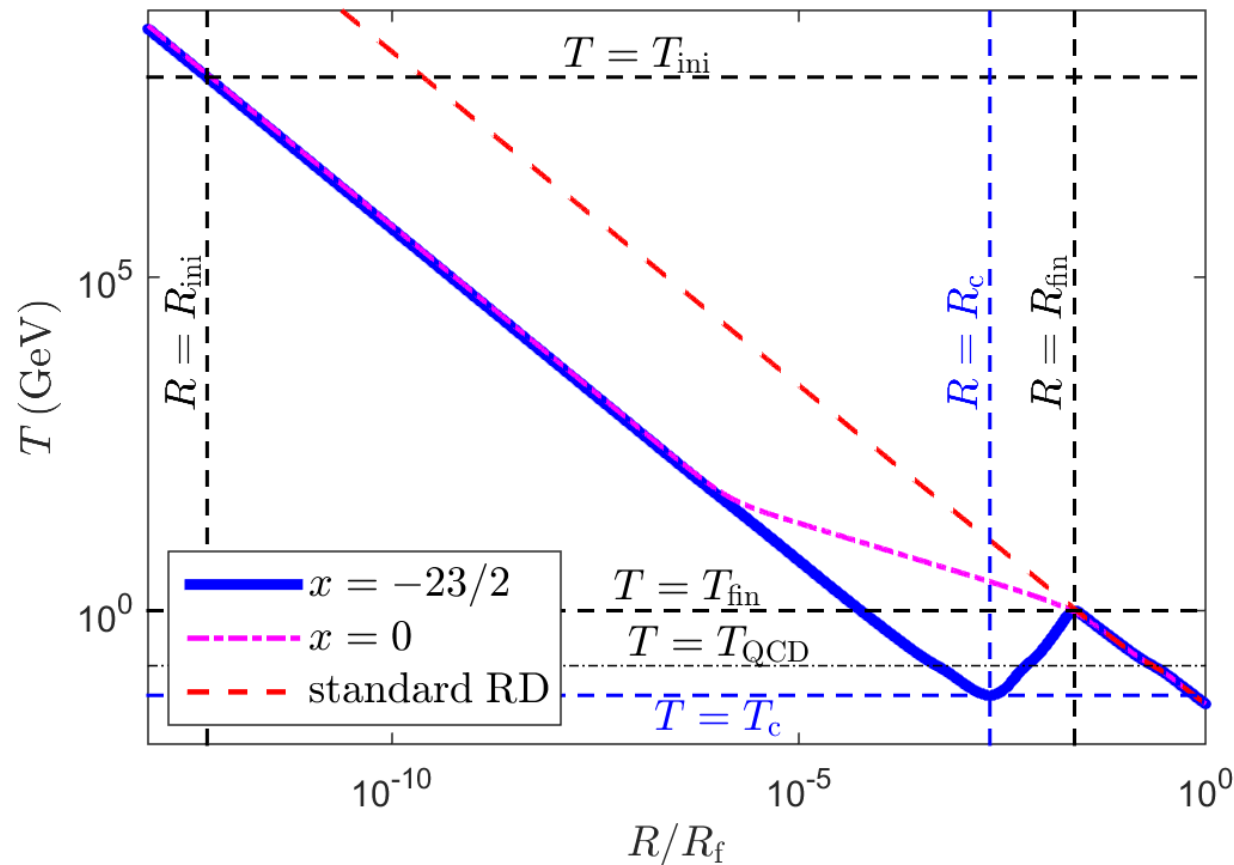


# Background energy density evolution

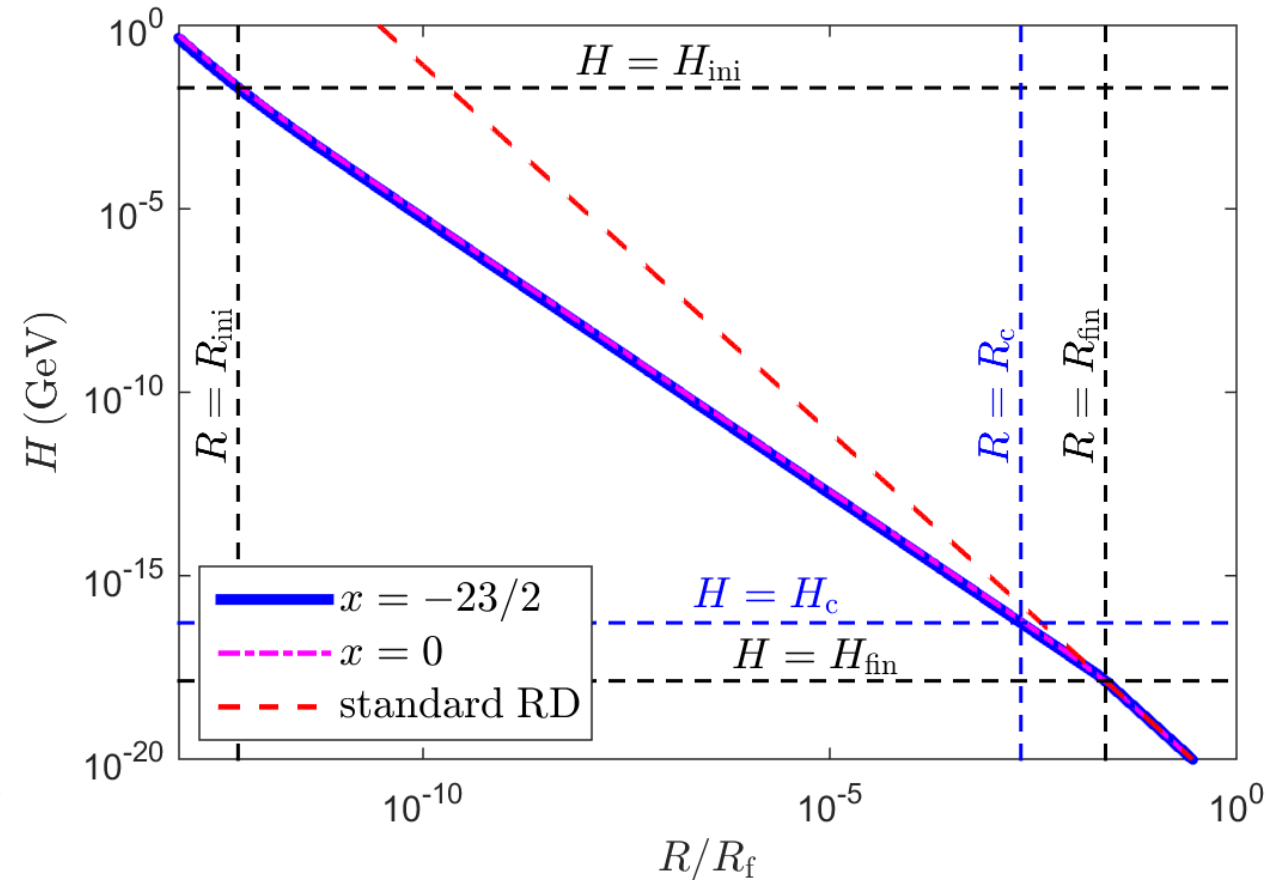
- Three example histories:
  - Standard RD ( - - - )
  - NSC  $x = 0$  ( - - - )
  - NSC  $x = -23/2$  ( — — )
- Adiabatic-nonadiabatic transition in NSC at  $R = R_c$  as decay of  $\phi$  affects radiation evolution



## Temperature vs scale factor



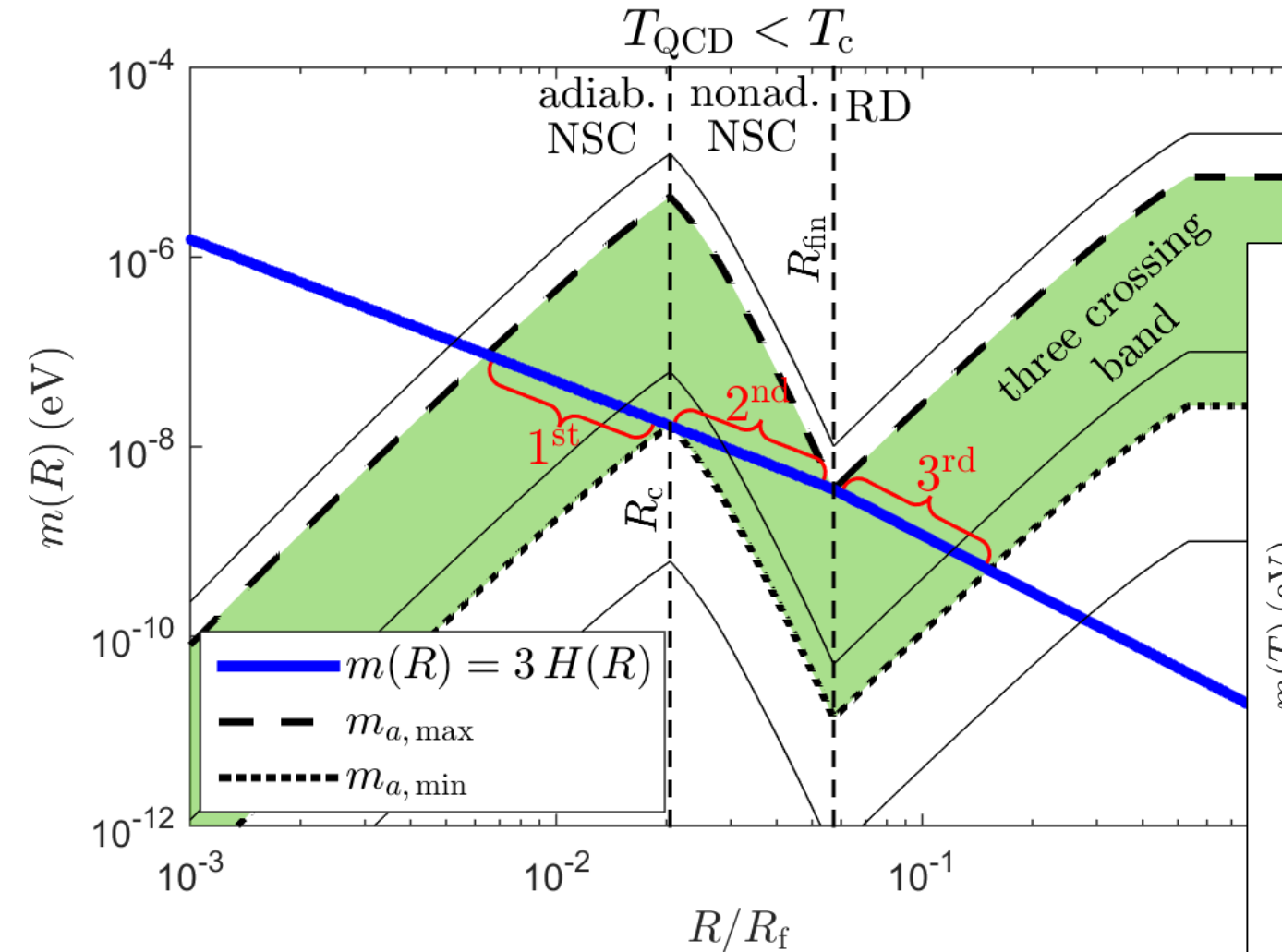
## Hubble rate vs scale factor



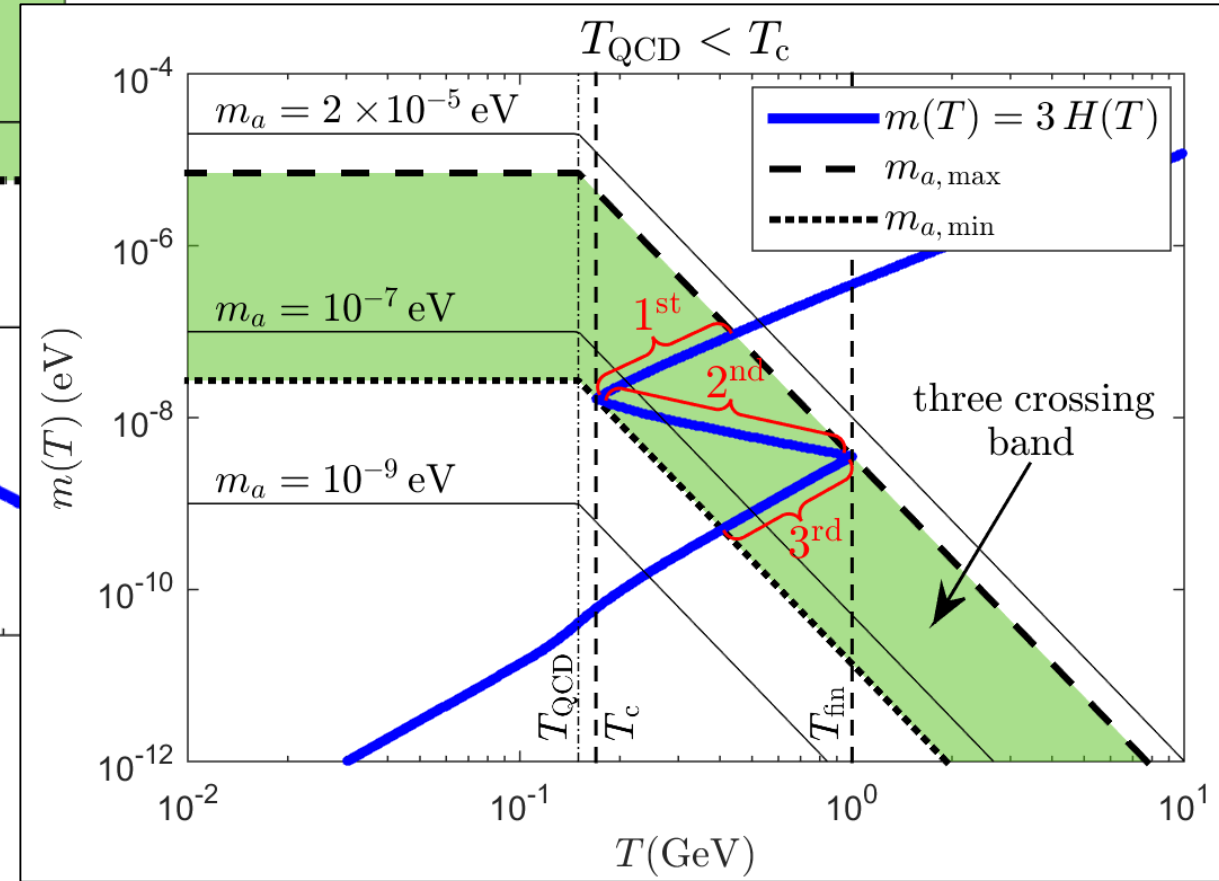
- Lowest temperature reached during NSC is now  $T_c$  rather than reheat temperature  $T_{fin}$   $\rightarrow$  require  $T_c > T_{\text{BBN}} \sim \text{MeV}$
- Same temperature can occur multiple times

# Axion mass and Hubble rate

$$m(T) \approx m_a \begin{cases} \left(\frac{T_{\text{QCD}}}{T}\right)^4 & T > T_{\text{QCD}} \\ 1 & T < T_{\text{QCD}} \end{cases}$$

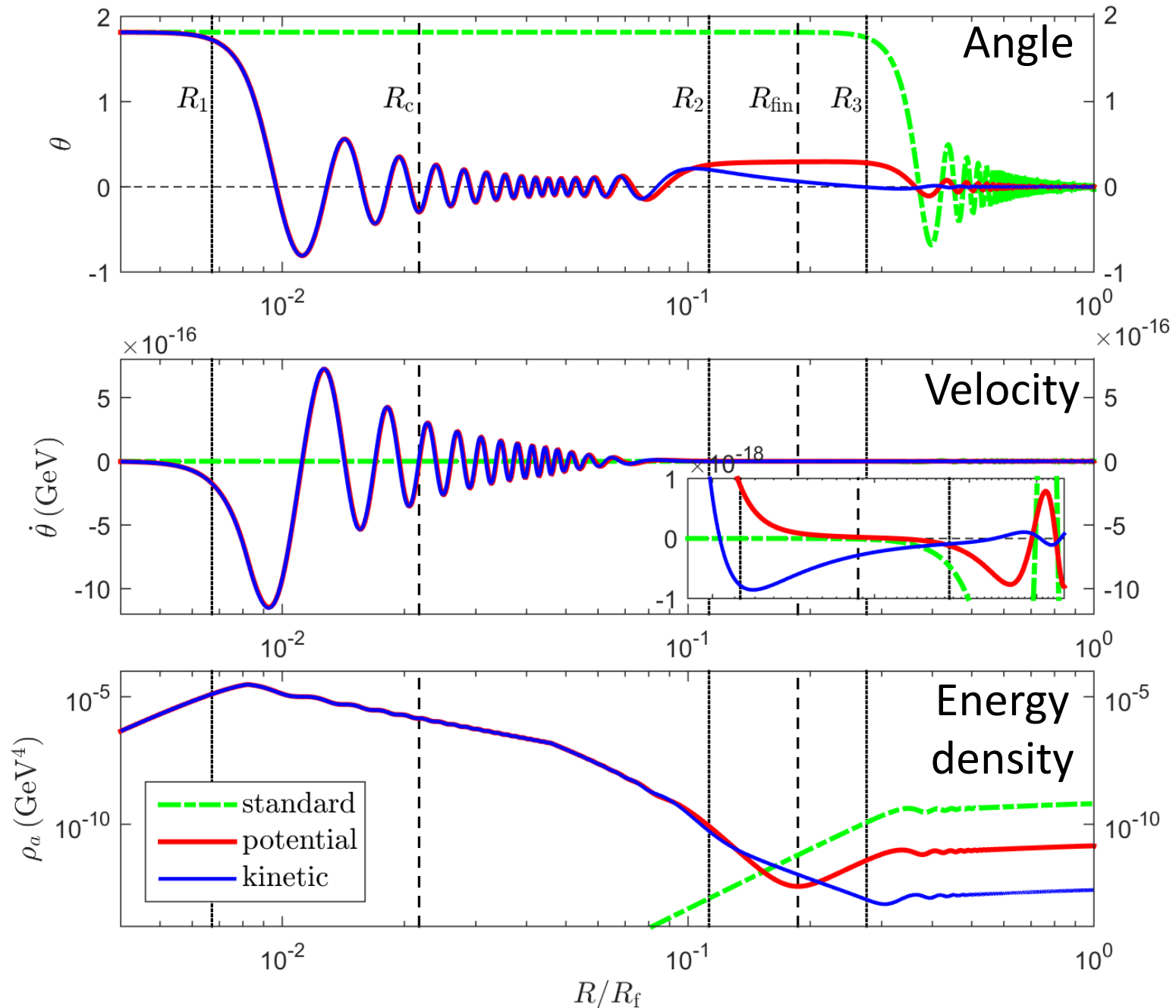


→  $3H \approx m$  can occur up to three times, provided that  $x < -3$



# Misalignment with three crossings

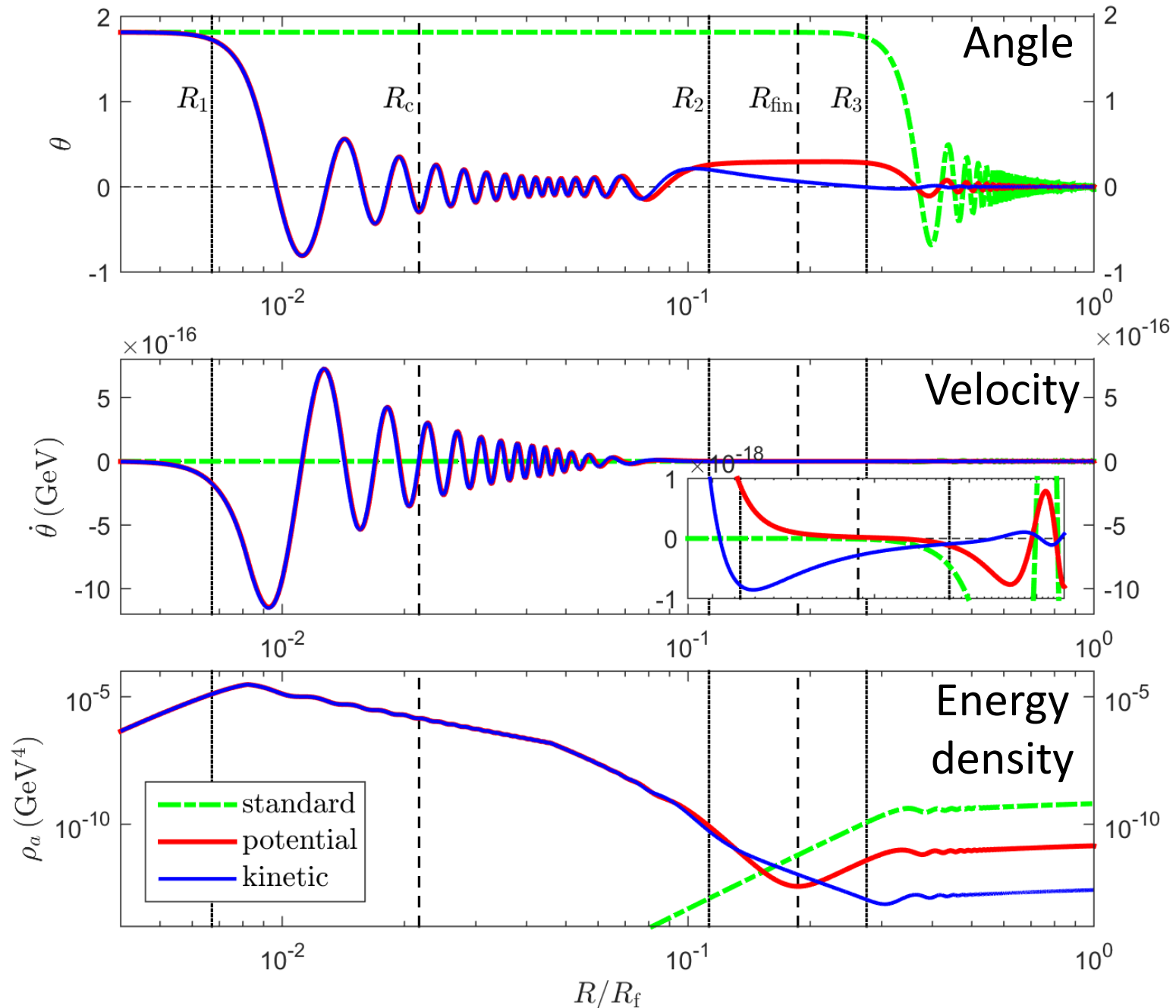
- First crossing ( $R_1$ ): oscillations begin as normal
- Second crossing ( $R_2$ ): Hubble friction restored, axion evolution dependent on kinetic/potential energy
- Third crossing ( $R_3$ ): oscillations resume from a new configuration (smaller angle, velocity tending toward zero)



# Misalignment with three crossings

- Misalignment is altered by restoration of Hubble friction
- Second period of oscillation with new “initial” condition
- Resultant axion energy density is smaller due to entropy injection and smaller amplitude

→ Smaller mass for axion DM



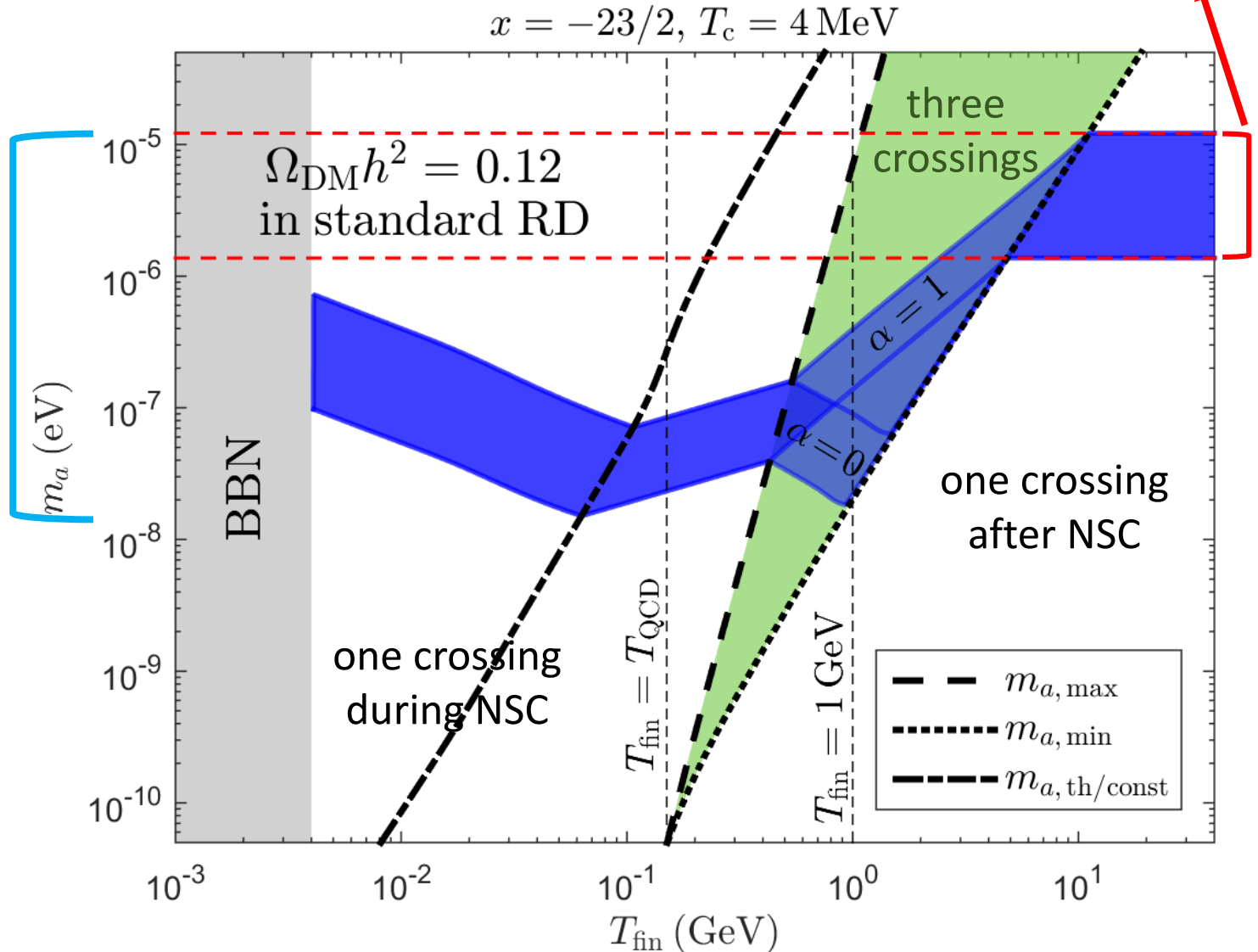
# Extended DM axion mass window

$$\frac{1}{2} < \theta_i < \frac{\pi}{\sqrt{3}}$$

- Mass window for observed DM abundance extended toward smaller values
- Lowest extent determined by  $\theta_i$  and  $x$

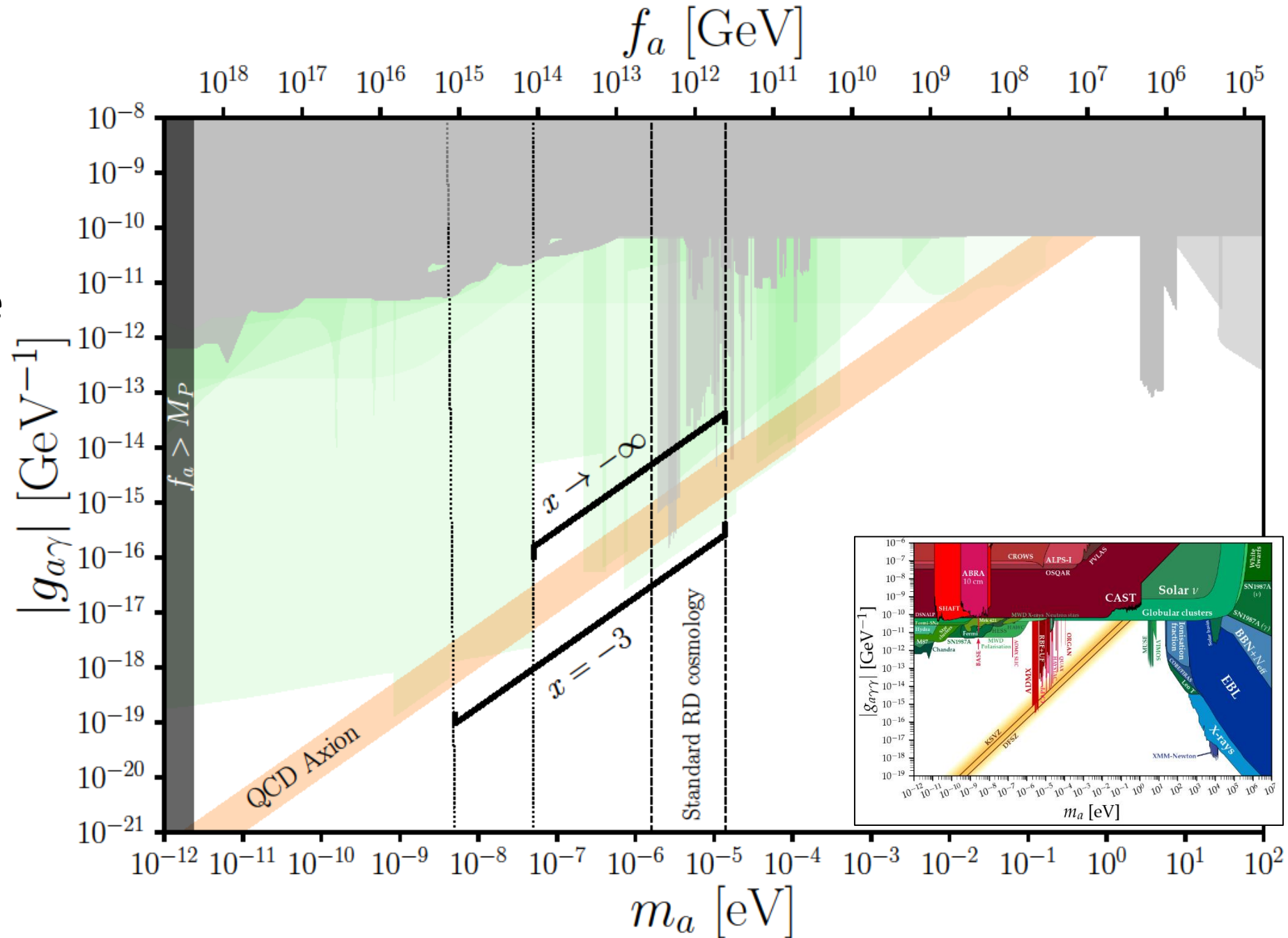
$\alpha \approx 0$ : kinetic

$\alpha \approx 1$ : potential



# Axion photon coupling

- Extended window overlaps with future projections outside of standard range
- Can probe NSC scenarios in coming years
- Adds to motivation to look out of standard window

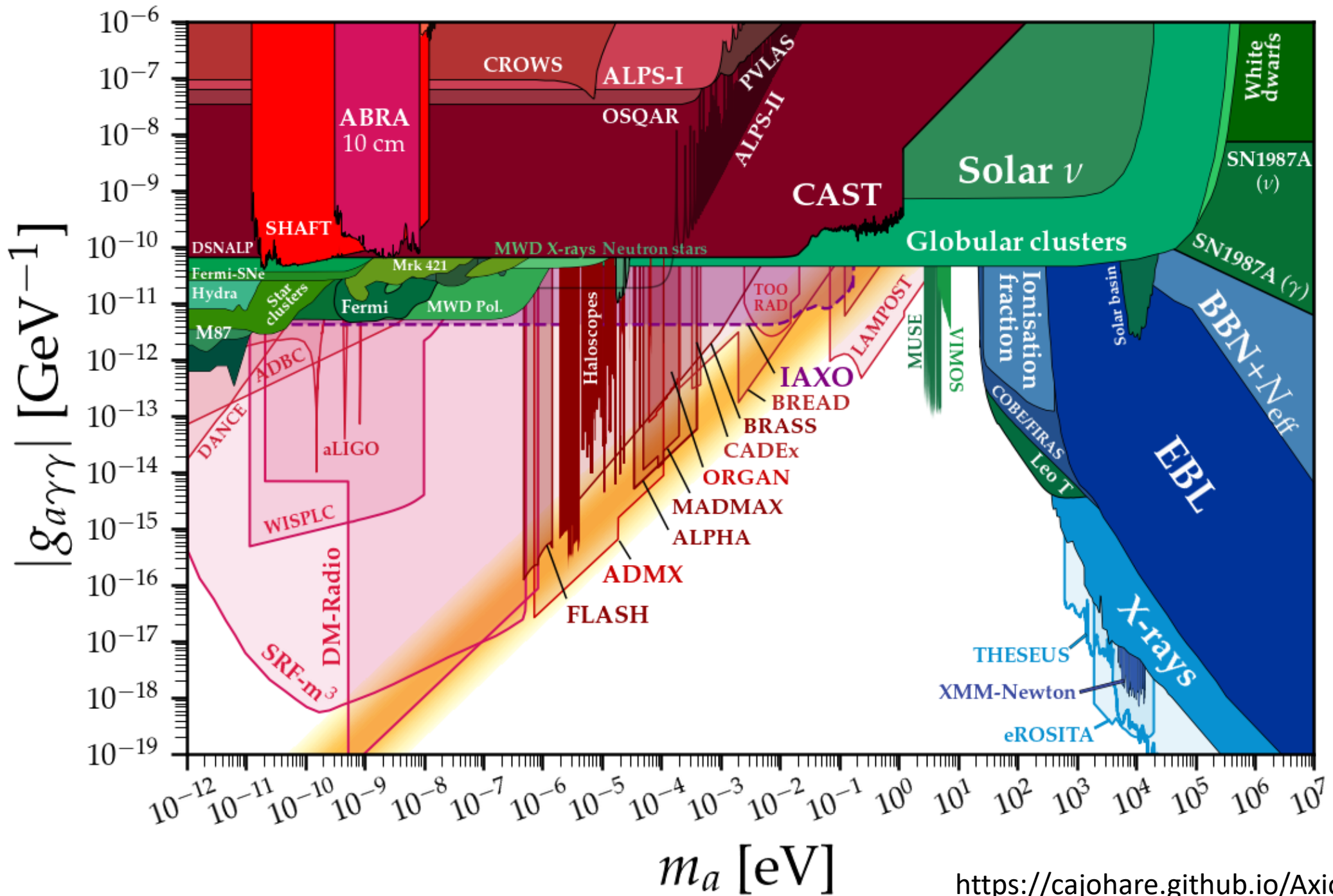


# Summary

- Axion misalignment mechanism depends on the cosmological history through  $m$  and  $H$
- NSC period with increasing temperature can significantly alter axion production  $\rightarrow H(t) \approx m(t)$  up to three times
- Hubble friction can be temporarily restored  $\rightarrow$  second period of oscillation with new initial condition
- Viable axion mass window for observed DM abundance extended to include smaller masses

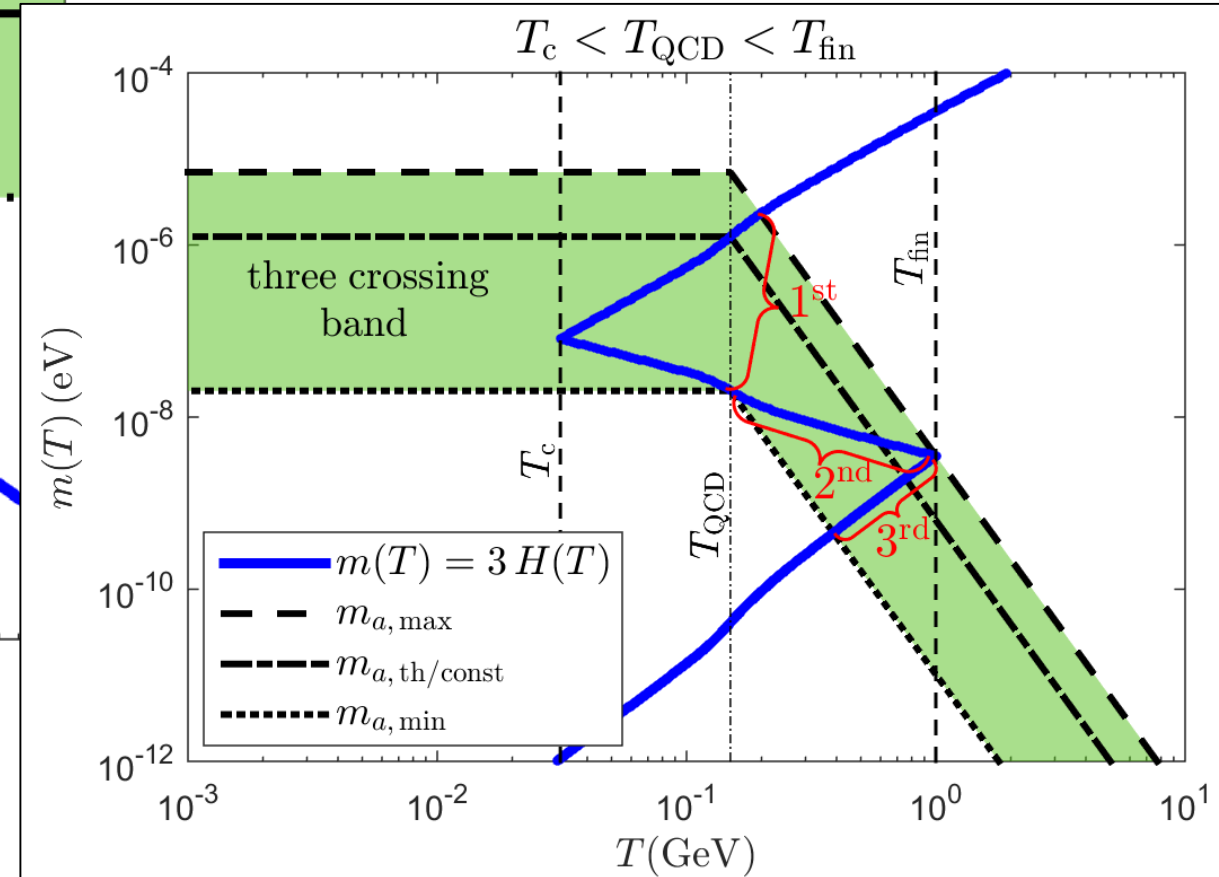
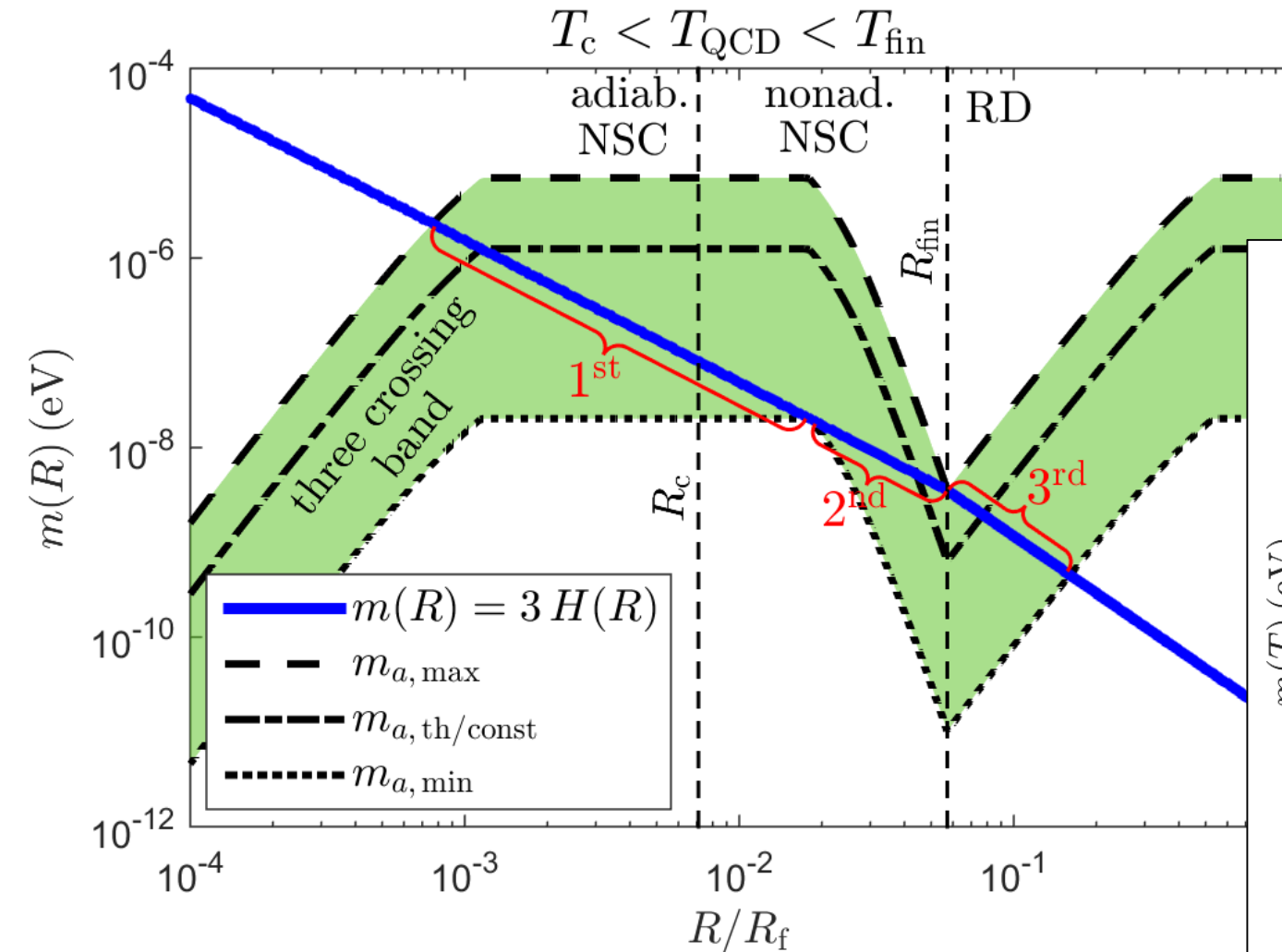


Thank you!



# Axion mass and Hubble rate

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