



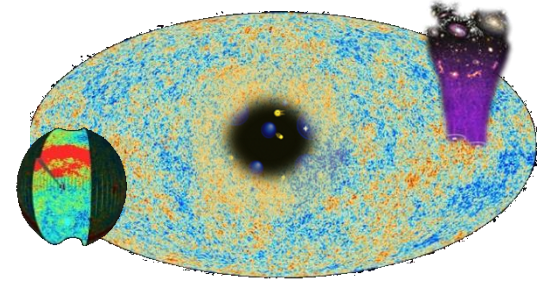
Some principles for the Hidden Sector explorer

Two possibilities for Beyond Standard Model with light particles

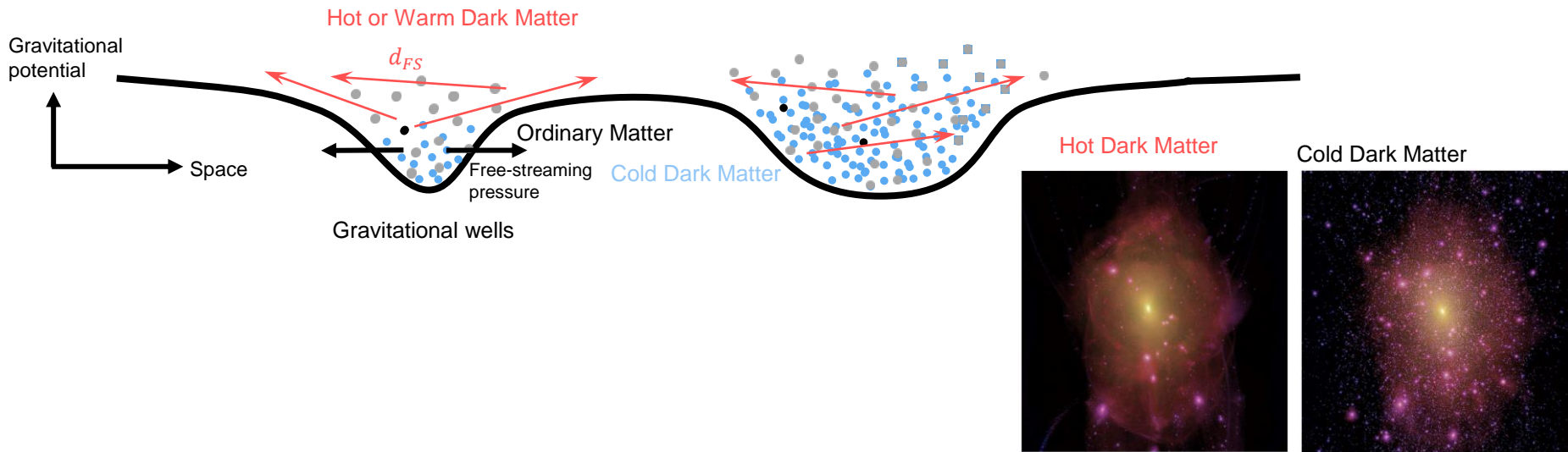
1. Wider theory exist at new high energy scale (SUSY, extra dim., etc) with degrees of freedom that stay relevant at low energies. Particles may be intrinsically light or be light by dynamic effects
2. SM + Hidden Sector with light messengers is all there is up to Planck scale – no new visible scale
3. or both...

E.g. Structure formation and Dark Matter

- At CMB $\delta\rho/\rho \sim 10^{-5}$
 - $\delta\rho/\rho$ grow with \sim scale a during matter domination
 - $a_{today}/a_{dec} = 1 + z_{dec} \sim 10^3$
 - Not enough!



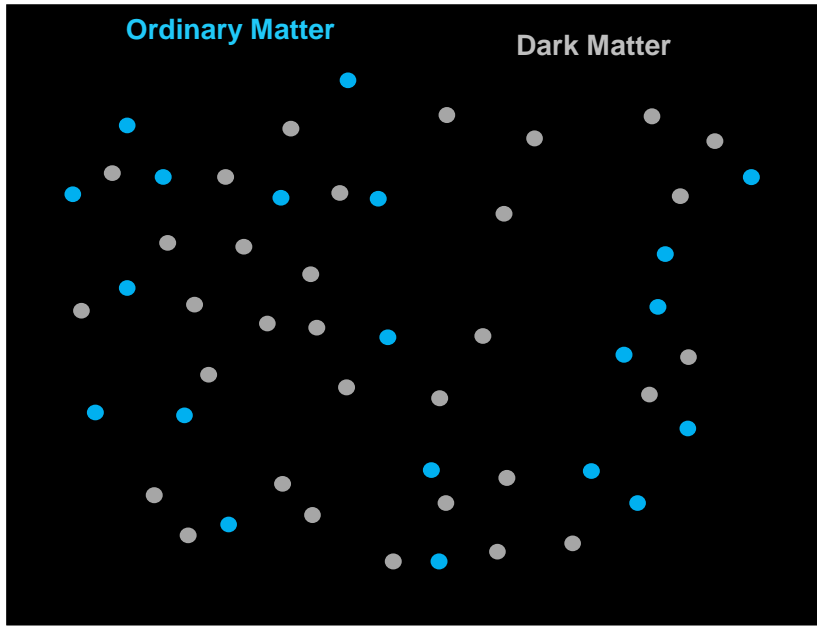
- DM can contribute in two ways:
 - Increasing mass density (increasing Jeans' scale for gravitational collapse $\propto \sqrt{T/m\rho}$)
 - Damping clustering of (too) small structures due to free-streaming $d_{FS} \propto v/\sqrt{\rho}$



- DM could produce a drop-off in the power spectrum of structures as a function of the scale
 - Wash out of structures with sizes in the range $10^6 - 10^8$ solar masses
 - How to measure?
 - Structures of $<10^8$ solar masses are very unlikely to have formed stars!

Multi-million dollar question

A sample of space and matter



$$\mathcal{L} = \boxed{i\bar{\psi}\gamma^\mu \partial_\mu \psi} - \boxed{m\bar{\psi}\psi} -$$

Free fermion

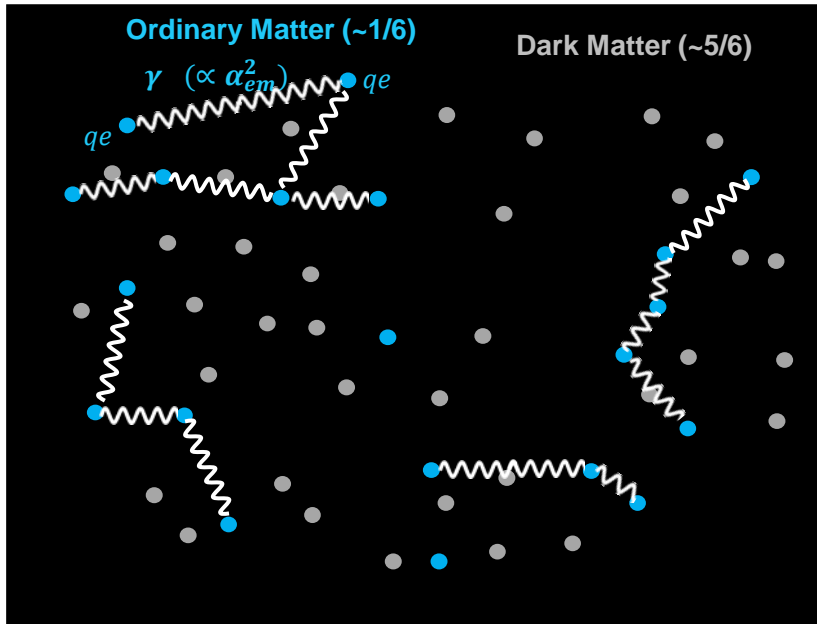
Fermion mass

$$+ \boxed{i\bar{\chi}\gamma^\mu \partial_\mu \chi} - \boxed{m\bar{\chi}\chi} -$$

Free DM

DM mass

Multi-million dollar question



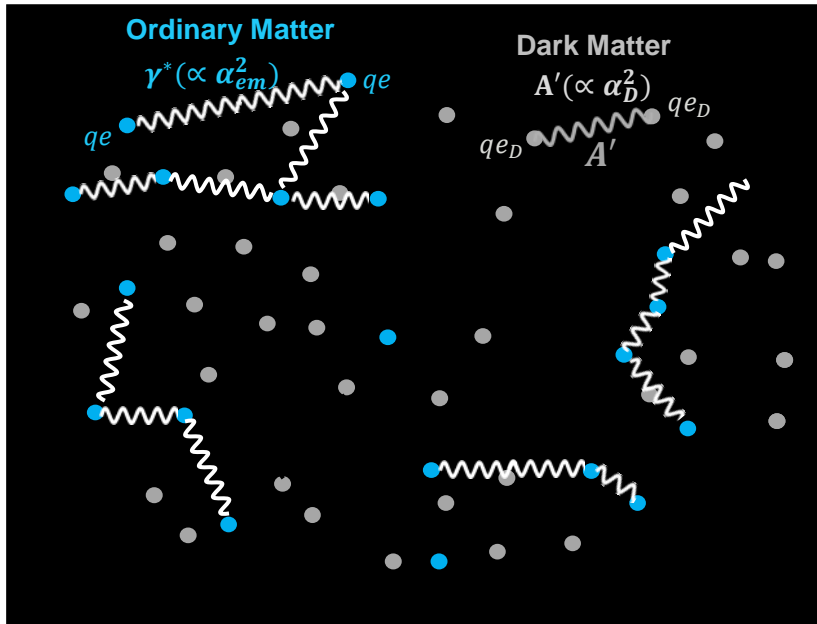
$$\mathcal{L} = \boxed{i\bar{\psi}\gamma^\mu \partial_\mu \psi} - \boxed{m\bar{\psi}\psi} - \boxed{\frac{1}{4}F_{\mu\nu}F^{\mu\nu}} + \boxed{e\bar{\psi}\gamma^\mu A_\mu \psi}$$

Free fermion Fermion mass Free photon Interaction fermion-photon γ

$$+ \boxed{i\bar{\chi}\gamma^\mu \partial_\mu \chi} - \boxed{m\bar{\chi}\chi}$$

Free DM DM mass

Multi-million dollar question



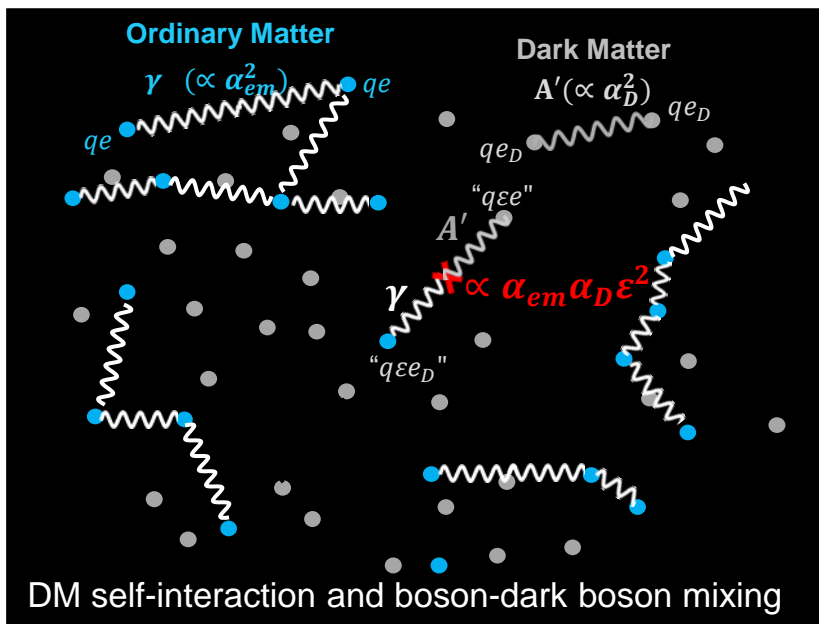
$$\mathcal{L} = \underbrace{i\bar{\psi}\gamma^\mu \partial_\mu \psi}_{\text{Free fermion}} - \underbrace{m\bar{\psi}\psi}_{\text{Fermion mass}} - \underbrace{\frac{1}{4}F_{\mu\nu}F^{\mu\nu}}_{\text{Free photon}} + \underbrace{e\bar{\psi}\gamma^\mu A_\mu \psi}_{\text{Interaction fermion-photon}}$$

$$+ \underbrace{i\bar{\chi}\gamma^\mu \partial_\mu \chi}_{\text{Free DM}} - \underbrace{m\bar{\chi}\chi}_{\text{DM mass}} - \underbrace{\frac{1}{4}V_{\mu\nu}V^{\mu\nu}}_{\text{Free dark photon}} + \underbrace{\epsilon\bar{\chi}\gamma^\mu A'_\mu \chi}_{\text{Interaction dark fermion - dark photon DM self-interaction...}} + \dots$$

Dark electric charge
Coupling extremely weak

Multi-million dollar question - SIDM

A sample of space



$$\mathcal{L} = \boxed{i\bar{\psi}\gamma^\mu \partial_\mu \psi} - \boxed{m\bar{\psi}\psi} - \boxed{\frac{1}{4}F_{\mu\nu}F^{\mu\nu}} + \boxed{e\bar{\psi}\gamma^\mu A_\mu \psi}$$

Free fermion Fermion mass Free photon Interaction fermion-photon

$$+ \boxed{i\bar{\chi}\gamma^\mu \partial_\mu \chi} - \boxed{m\bar{\chi}\chi} - \boxed{\frac{1}{4}V_{\mu\nu}V^{\mu\nu}} + \boxed{\epsilon\bar{\chi}\gamma^\mu A'_\mu \chi} + \dots$$

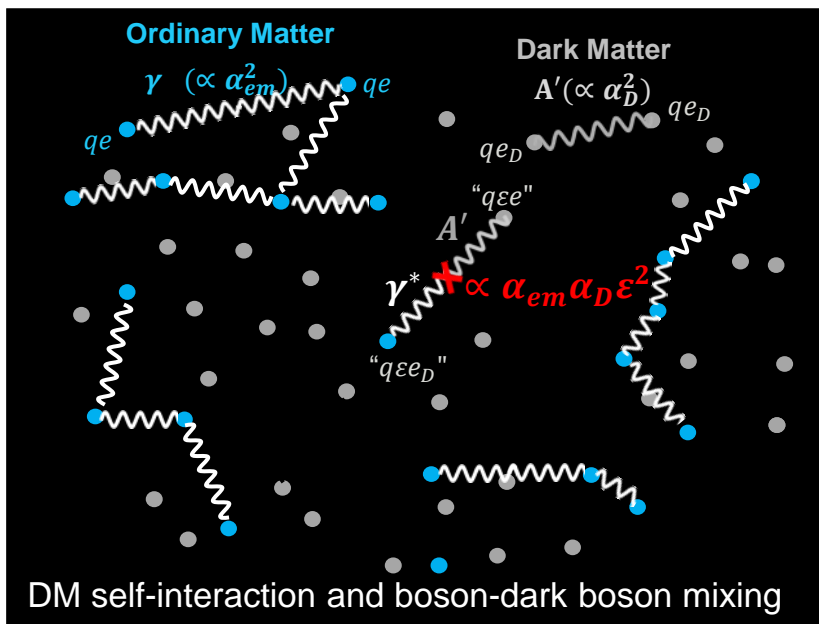
Free DM DM mass Free dark photon Interaction dark fermion – dark photon DM self-interaction...

$$\boxed{-\frac{1}{4}\epsilon F_{\mu\nu}V^{\mu\nu}}$$

Kinetic mixing between ordinary photon and dark photon!
"Vector portal"
→ Coupling extremely weak

Multi-million dollar question – SIDM?

A sample of space



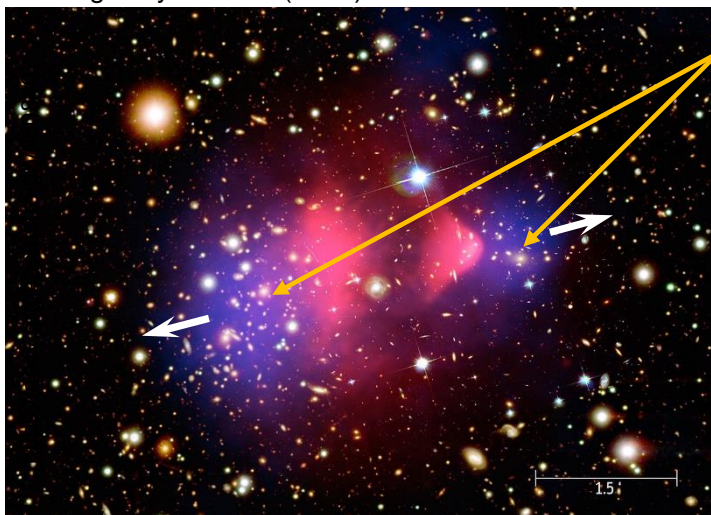
$$\mathcal{L} = \underbrace{i\bar{\psi}\gamma^\mu \partial_\mu \psi}_{\text{Free fermion}} - \underbrace{m\bar{\psi}\psi}_{\text{Fermion mass}} - \underbrace{\frac{1}{4}F_{\mu\nu}F^{\mu\nu}}_{\text{Free photon}} + \underbrace{e\bar{\psi}\gamma^\mu A_\mu \psi}_{\text{Interaction fermion-photon}}$$

$$+ \underbrace{i\bar{\chi}\gamma^\mu \partial_\mu \chi}_{\text{Free DM}} - \underbrace{m\bar{\chi}\chi}_{\text{DM mass}} - \underbrace{\frac{1}{4}V_{\mu\nu}V^{\mu\nu}}_{\text{Free dark photon}} + \underbrace{\epsilon\bar{\chi}\gamma^\mu A'_\mu \chi}_{\text{Interaction dark fermion - dark photon DM self-interaction...}} + \dots$$

$$\underbrace{-\frac{1}{4}\epsilon F_{\mu\nu}V^{\mu\nu}}_{\text{Kinetic mixing between ordinary photon and dark photon! "Vector portal"}}$$

Kinetic mixing between ordinary photon and dark photon!
"Vector portal"
→ Coupling extremely weak

Bullet galaxy clusters (2003)



Collision between two galaxy clusters

- In red, X-ray emitting plasma = dominant baryonic mass (5-15%)
- In blue, reconstruction of total mass distribution from lensing
→ Trace out Dark Matter distribution

Dark Matter remains around galaxy clusters seemingly undisturbed

→ Almost(?) collisionless

→ Dark Matter is, or is just about, non-self-interacting $\sigma/m \lesssim 1 \text{ cm}^2 \text{ g}^{-1}$

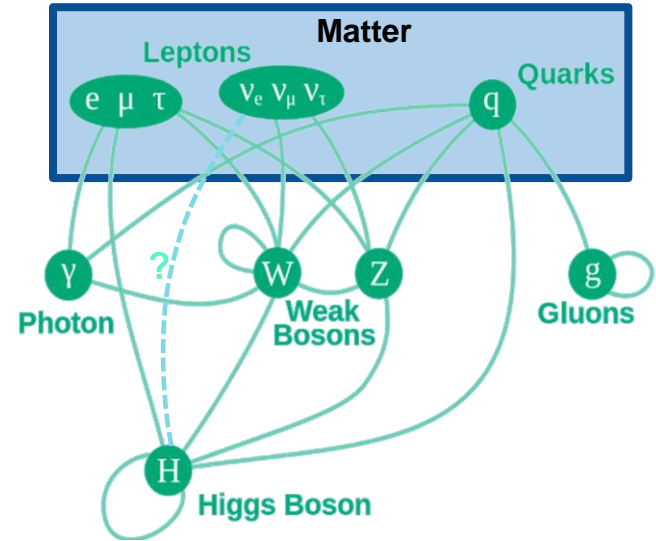
→ Currently we observe >70 colliding galaxy clusters (arXiv:1610.05327)

“Standard model” of Hidden Sector

- Standard Model has taught us successful formalism to implement particles, interaction and mediators
 - SM not only successful, we discovered what it predicted
 - Gives us plausible tools to implement Dark Sector with well-defined phenomenology:

Options for “portals”

- fermionic/scalar DM
- Dark Photons
- Dark Scalars
- Heavy Neutral Leptons
- Axion-Like Particles



A model should be used to teach us something!

- Dynamics of Hidden Sector may drive dynamics and ‘anomalies’ of Visible Sector!
 - Neutrino oscillations and mass, baryon asymmetry, Higgs mass, Dark Matter (abundance, distribution and “behavior”), structure formation, inflation, ...

- Also some Super-Symmetric DM candidates and “portals” etc
 - Sgoldstino, Neutralino in R-Parity Violating SUSY, Hidden Photinos, axinos and saxions....



Search guidelines are not always so clear...

Weinberg (1967)

VOLUME 19, NUMBER 21 PHYSICAL REVIEW LETTERS 20 NOVEMBER 1967

¹¹In obtaining the expression (11) the mass difference between the charged and neutral has been ignored.
¹²M. Ademollo and R. Gatto, *Nuovo Cimento* **44A**, 282 (1966); see also J. Pasupathy and R. E. Marshak, *Phys. Rev. Letters* **17**, 888 (1966).
¹³The predicted ratio [eq. (12)] from the current algebra is slightly larger than that (0.23%) obtained from the ρ -dominance model of Ref. 2. This seems to be true also in the other case of the ratio $\Gamma(\eta \rightarrow \pi^+\pi^-\gamma)/\Gamma(\gamma\gamma)$ calculated in Refs. 12 and 14.
¹⁴L. M. Brown and P. Singer, *Phys. Rev. Letters* **8**, 460 (1962).

A MODEL OF LEPTONS*

Steven Weinberg[†]
 Laboratory for Nuclear Science and Physics Department,
 Massachusetts Institute of Technology, Cambridge, Massachusetts
 (Received 17 October 1967)

Leptons interact only with photons, and with the intermediate bosons that presumably mediate weak interactions. What could be more natural than to unite¹ these spin-one bosons into a multiplet of gauge fields? Standing in

and on a right-handed singlet

$$R \equiv [\frac{1}{2}(1-\gamma_5)]e. \quad (2)$$

The largest group that leaves invariant the kin-

The coupling of ϕ_1 to muons is stronger by a factor M_μ/M_e , but still very weak. Note also that (14) gives g and g' larger than e , so (16) tells us that $M_W > 40$ BeV, while (12) gives $M_Z > M_W$ and $M_Z > 80$ BeV.
 The only unequivocal new predictions made

If $g \gg e$ then $g \gg g'$, and this is just the usual $e-\nu$ scattering matrix element times an extra factor $\frac{1}{2}$. If $g \approx e$ then $g \ll g'$, and the vector interaction is multiplied by a factor $-\frac{1}{2}$ rather than $\frac{1}{2}$. Of course our model has too many arbitrary features for these predictions to be

1265

taken very seriously, but it is worth keeping in mind that the standard calculation³ of the electron-neutrino cross section may well be wrong.

mi, *Z. Physik* **88**, 161 (1934). A model similar to ours was discussed by S. Glashow, *Nucl. Phys.* **22**, 579 (1961); the chief difference is that Glashow introduces symmetry-breaking terms into the Lagrangian, and therefore into the definition of the predictions.

1973: Discovery of neutral currents at CERN

1979: Nobel price Glashow, Salam, Weinberg

Nucl. Phys. B106 (1976)

A PHENOMENOLOGICAL PROFILE OF THE HIGGS BOSON

John ELLIS, Mary K. GAILLARD * and D.V. NANOPOULOS **
CERN, Geneva

Received 7 November 1975

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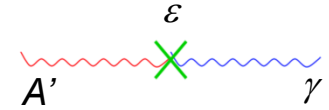
J. Ellis et al. / Higgs boson

We should perhaps finish with an apology and a caution. We apologize to experimentalists for having no idea what is the mass of the Higgs boson, unlike the case with charm [3,4] and for not being sure of its couplings to other particles, except that they are probably all very small. For these reasons we do not want to encourage big experimental searches for the Higgs boson, but we do feel that people performing experiments vulnerable to the Higgs boson should know how it may turn up.

➔ Hidden Sector: Room for progress in theory and expect guidance from the cosmic frontier

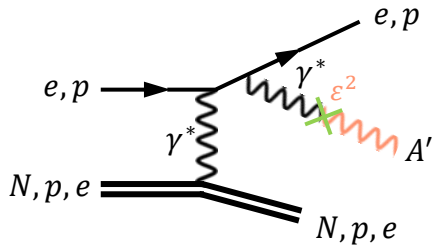
Dark photon production

- Kinetic mixing with massive dark/secluded/paraphoton A' : $\frac{1}{2} \varepsilon F_{\mu\nu}^{SM} F_{HS}^{\mu\nu}$

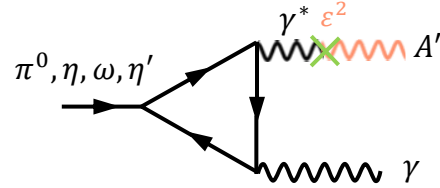


→ Motivated in part by idea of “mirror world” restoring L/R symmetry, dark matter (AMS e^+ excess), $g-2$ anomaly, ...

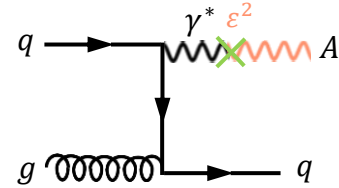
Dark photons: need loads of photons!



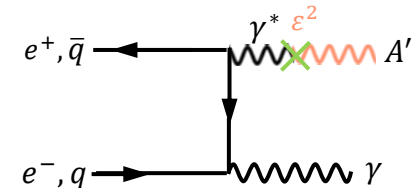
Bremsstrahlung



Light neutral meson decays



Quark/gluon scattering



Annihilation

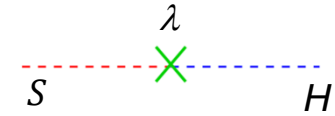
→ Proton and electron beam dumps, pp and e^+e^- colliders

→ Signature: decay or missing energy

	Physics model	Final state
HSDS	HNL, SUSY neutralino	$\ell^\pm \pi^\mp, \ell^\pm K^\mp, \ell^\pm \rho^\mp (\rho^\mp \rightarrow \pi^\mp \pi^0)$
	DP, DS, ALP (fermion coupling), SUSY sgoldstino	$\ell^+ \ell^-$
	DP, DS, ALP (gluon coupling), SUSY sgoldstino	$\pi^+ \pi^-, K^+ K^-$
	HNL, SUSY neutralino, axino	$\ell^+ \ell^- \nu$
	ALP (photon coupling), SUSY sgoldstino	$\gamma\gamma$
	SUSY sgoldstino	$\pi^0 \pi^0$

Dark scalar (“Higgs”) production

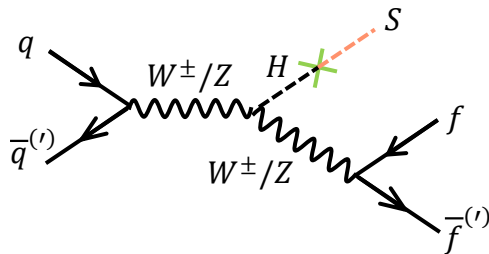
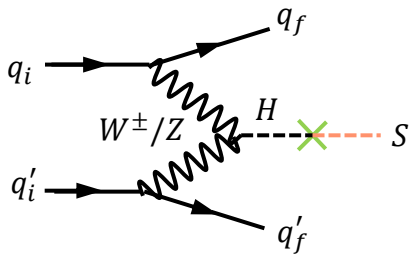
- Mass mixing with dark singlet scalar S : $(gS + \lambda S^2)H^\dagger H$



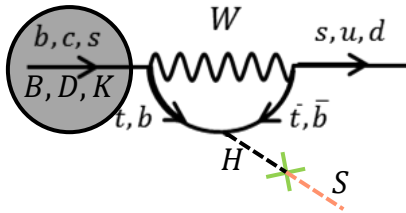
→ Mass to Higgs boson and mass generation in dark sector, inflaton, dark phase transitions BAU, dark matter,...

→ Production of Dark Scalars: Loads of Higgses (real or virtual)!

pp colliders



pp colliders or p beam dumps

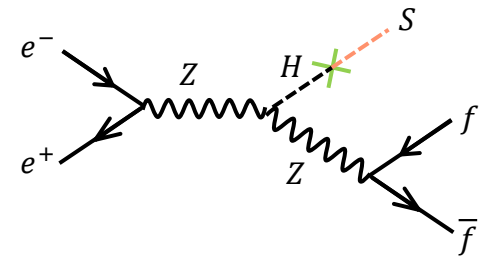


$$\Gamma(K \rightarrow \pi S) \propto (m_t^2 |V_{ts}^* V_{td}|)^2$$

$$\Gamma(D \rightarrow \pi S) \propto (m_b^2 |V_{cb}^* V_{ub}|)^2$$

$$\Gamma(B \rightarrow \pi S) \propto (m_t^2 |V_{tb}^* V_{ts}|)^2$$

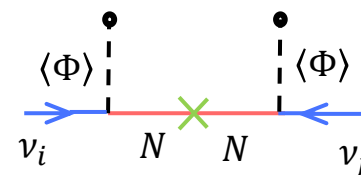
Higgs factory



Sterile neutrino/heavy neutral lepton production

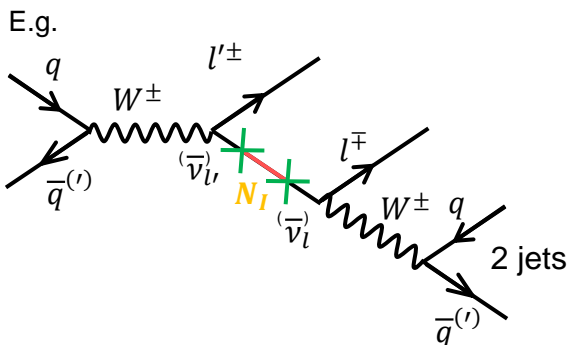
◉ Mixing with right-handed neutrino N (Heavy Neutral Lepton): $Y_{I\ell} H^\dagger \bar{N}_I L_\ell$

→ Neutrino oscillation, baryon asymmetry, dark matter



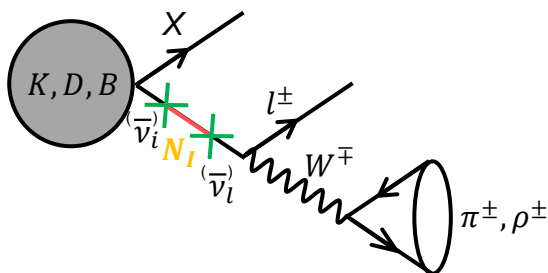
→ Production of Heavy Neutral Leptons: loads of neutrinos!

pp colliders

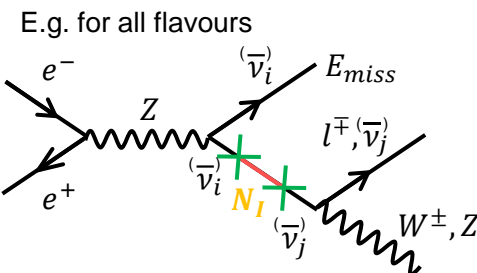


Like-sign leptons: Lepton Number Violation

pp colliders or p beam dumps

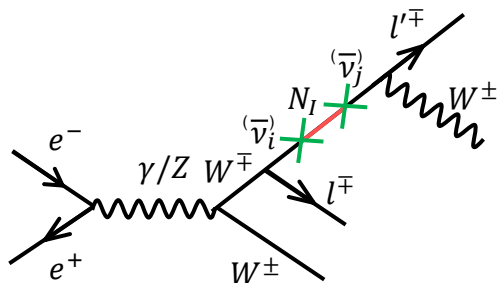


e^+e^- colliders (W,Z,H)
(also $\mu\mu$ -colliders)

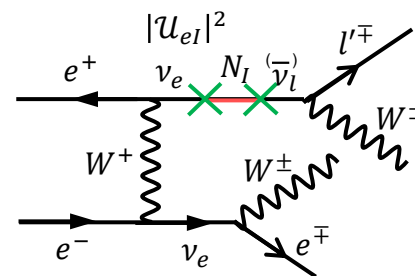
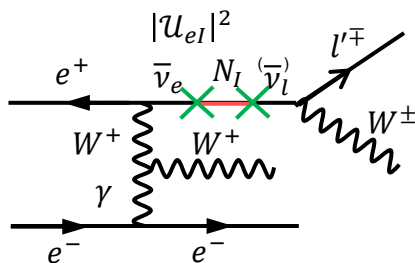


Most promising channel:
 $e^+e^- \rightarrow N(\rightarrow l^\mp W^\pm) \nu_l : l^\mp + 2j + E_{miss}$

e^+e^- colliders (W,ZH) continued
(also $\mu\mu$ -, $\gamma\gamma$ -, $e\gamma$ -colliders)



Alternative: $e^+e^- \rightarrow N(\rightarrow l^\mp W^\pm) l^\mp W^\pm : dilepton + 4j$

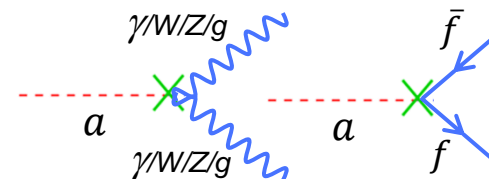


Axion-like particle production

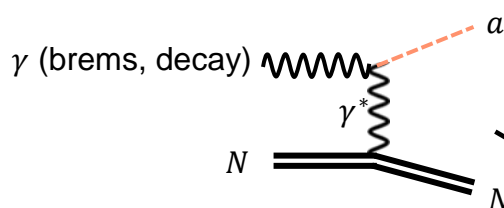
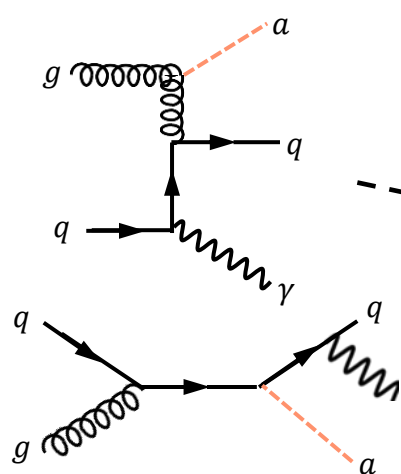
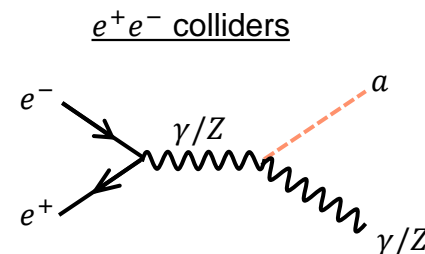
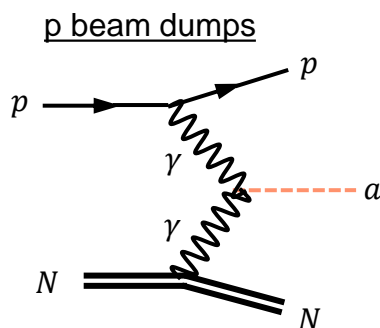
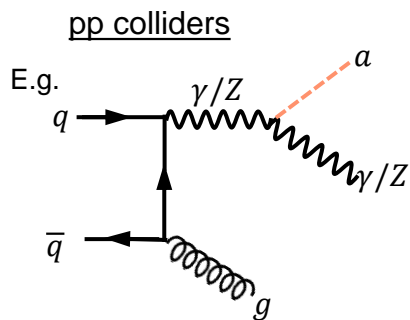
- Mixing with Axion-Like Particles (ALPs) a with coupling to photons, fermions, gluons :

$$\frac{a}{F} G_{\mu\nu} \tilde{G}^{\mu\nu}, \frac{\partial_\mu a}{F} \bar{\psi} \gamma_\mu \gamma_5 \psi, \text{ etc}$$

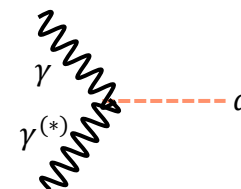
- Extended Higgs, SUSY breaking, dark matter, possibility of inflaton,...



- Production of ALPs: Couplings to photons, gluons and fermions, loads of interactions!**



Gamma factory



And more...

Examples with mass $\sim \mathcal{O}(GeV)$ and production branching ratio $\sim \mathcal{O}(10^{-10})$

→ Light super-goldstinos [Gorbunov, 2001]

“Axion- and dilaton-like”

$$\rightarrow D \rightarrow \pi X, X \rightarrow \pi^+ \pi^-, \pi^0 \pi^0, l^+ l^-$$

→ R-parity violating neutralinos in SUSY [Dedes et al., 2001]

“Heavy-neutrino like”

$$\rightarrow D \rightarrow l \tilde{\chi}, \tilde{\chi} \rightarrow l^+ l^- \nu$$

◎ Overlap in decay signatures!

- Missing energy or decays:

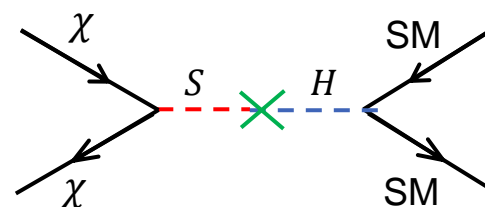
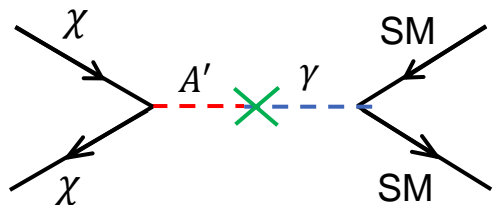
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HSDS	DP, DS, ALP (gluon coupling), SUSY sgoldstino	$\pi^+ \pi^-, K^+ K^-$
	HNL, SUSY neutralino, axino	$l^+ l^- \nu$
	ALP (photon coupling), SUSY sgoldstino	$\gamma\gamma$
	SUSY sgoldstino	$\pi^0 \pi^0$

Also DP, DS, HNL, ALP,... to jets

(L)DM searches

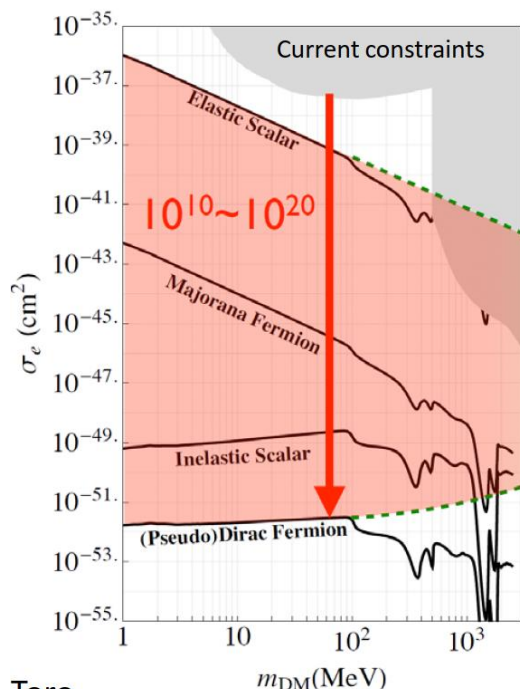
DM search

- Interpretation of invisible energy accompanied by SM signature (scattering, decay...) and assumption on DM-dark boson coupling

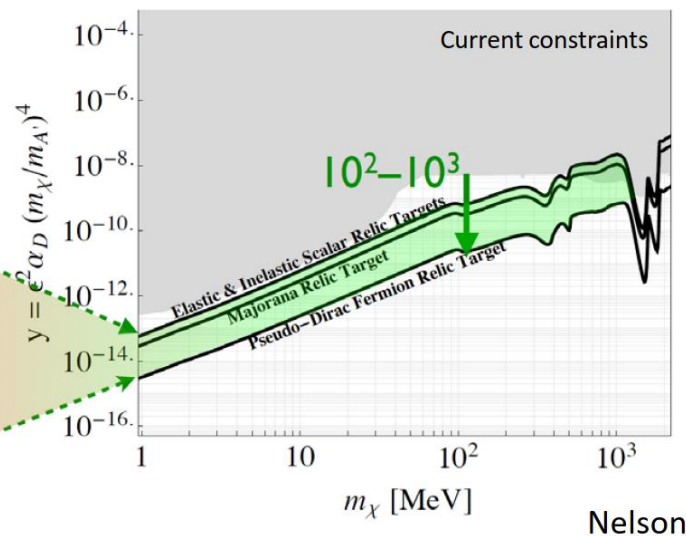


- Scattering against atomic electrons and nuclei

Detection by scattering relic DM



Detection by scattering accelerator-produced DM



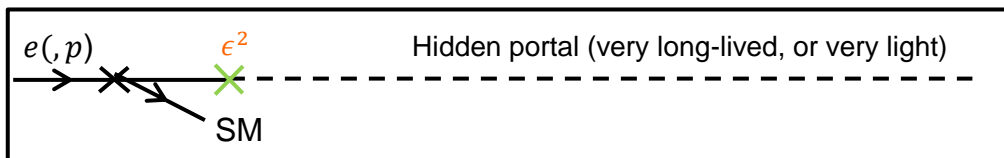
Relativistic production (high- q^2) at accelerators almost insensitive to spin and mass

Long-lived search techniques

Note : $\epsilon \ll 1$

Production

Detection



Escape detector
Missing energy/momentum/mass
Probability $\propto \epsilon^2$

Indirect (invisible)

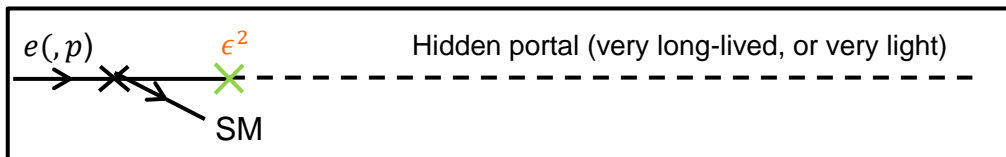
Long-lived search techniques

Note : $\epsilon \ll 1$

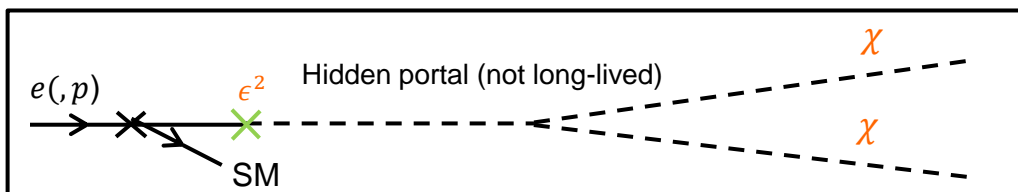
Indirect (invisible)

Production

Detection



Escape detector
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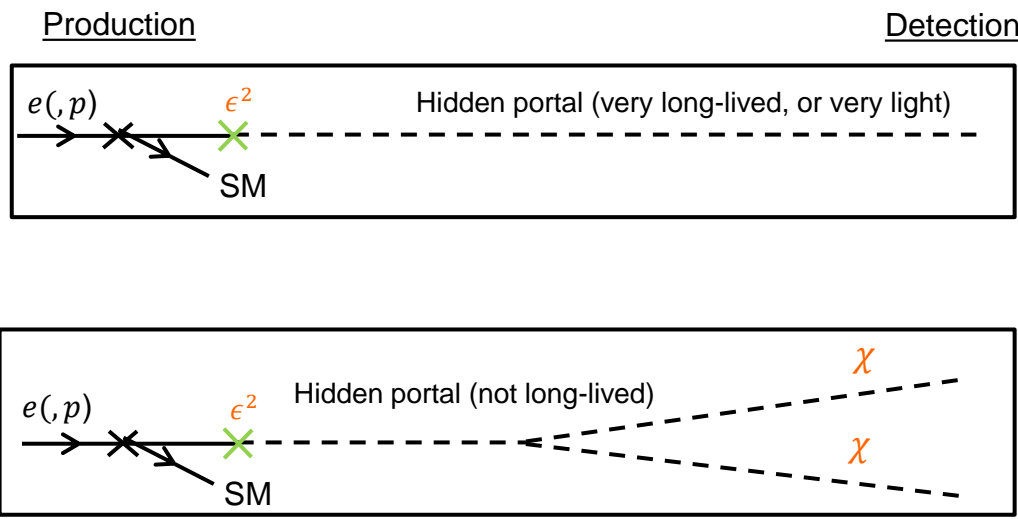


If $m_{HP} > 2m_{DM}$
Escape detector
Missing energy/momentum/mass
Probability $\propto \epsilon^2$

Long-lived search techniques

Note : $\epsilon \ll 1$

Indirect



Escape detector
Missing energy/momentum/mass
Probability $\propto \epsilon^2$

If $m_{HP} > 2m_{DM}$
Escape detector
Missing energy/momentum/mass
Probability $\propto \epsilon^2$

- ATLAS
- CMS
- LHCb
- NA64 (++)
- LDMX
- ...

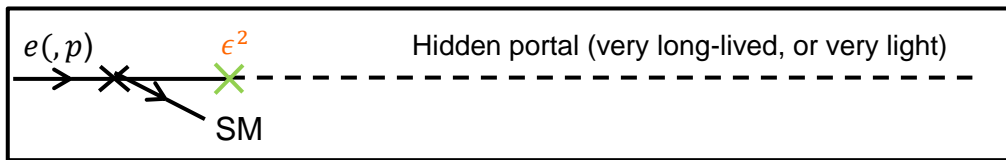
Long-lived search techniques

Note : $\epsilon \ll 1$

Indirect (invisible)

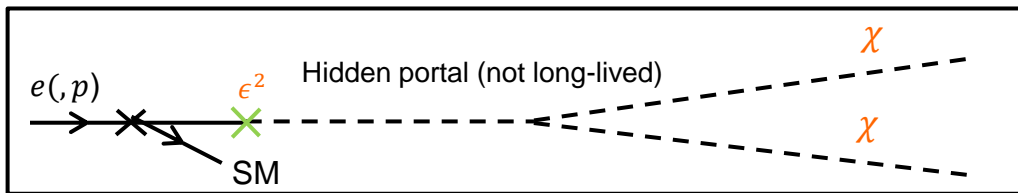
Production

Detection



Escape detector
Missing energy/momentum/mass
Probability $\propto \epsilon^2$

Cannot be directly distinguished
→ Model-dependent interpretation



If $m_{HP} > 2m_{DM}$
Escape detector
Missing energy/momentum/mass
Probability $\propto \epsilon^2$

Use of electrons limits search to Dark Photon mediator (&ALPs)

ATLAS
CMS
LHCb
NA64 (++)
LDMX

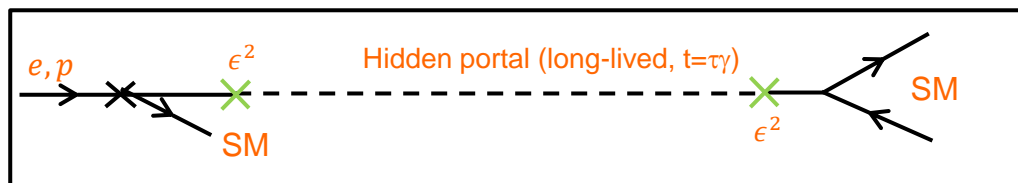
Long-lived search techniques

Production

Detection

Note : $\epsilon \ll 1$

“Direct” (indirect visible)



Decay signature (“displaced vertex”)
 Probability $\propto \epsilon^4$
 Model independent

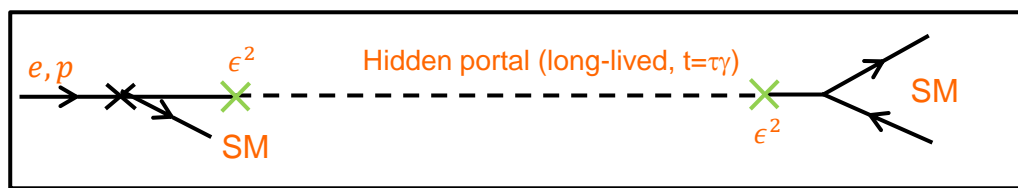
Long-lived search techniques

Production

Detection

Note : $\epsilon \ll 1$

“Direct” (indirect visible)



Decay signature (“displaced vertex”)
 Probability $\propto \epsilon^4$
 Model independent

MATHUSLA
 ANUBIS
 CODEX-b

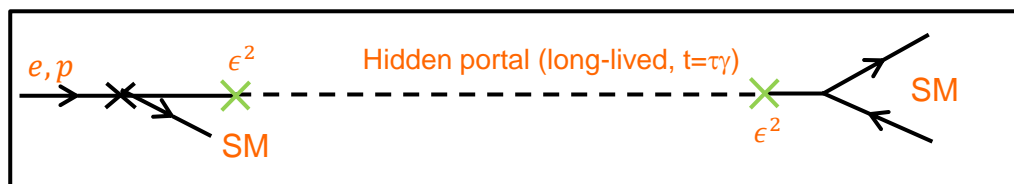
Long-lived search techniques

Production

Detection

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“Direct” (indirect visible)



Decay signature (“displaced vertex”)
 Probability $\propto \epsilon^4$
 Model independent
 +
 Reconstruction of decay: mass, PID
 → Distinguish models
 → Measurement of properties

- MATHUSLA
- ANUBIS
- CODEX-b
- ATLAS
- CMS
- LHCb
- NA62 (++)
- NA64 (++)
- FASER
- SHiP

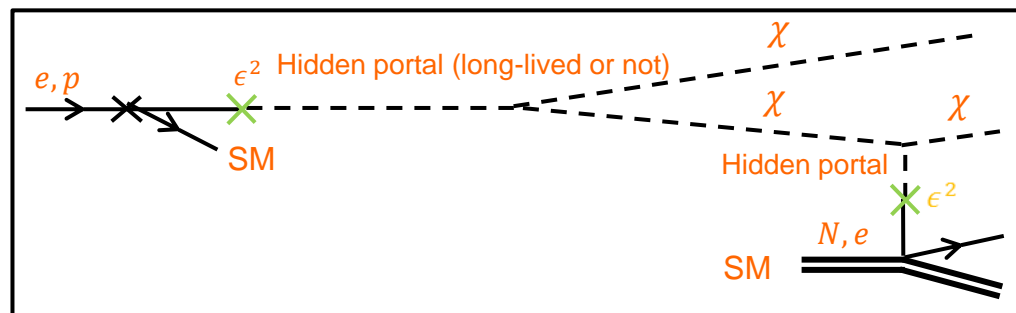
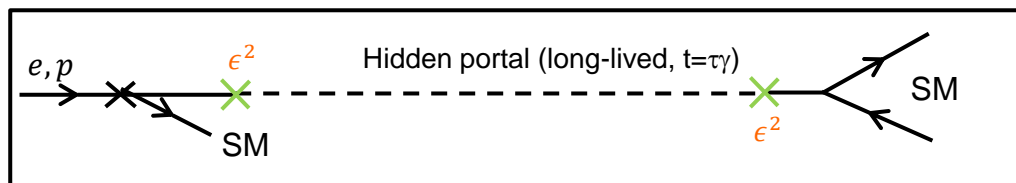
Long-lived search techniques

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Dark matter scattering
 Recoil against electron or nuclei
 Probability $\propto \epsilon^4$
 Model-dependent interpretation

- MATHUSLA
- ANUBIS
- CODEX-b
- +
- ATLAS
- CMS
- LHCb
- NA62 (++)
- NA64 (++)
- FASER
- SHIP

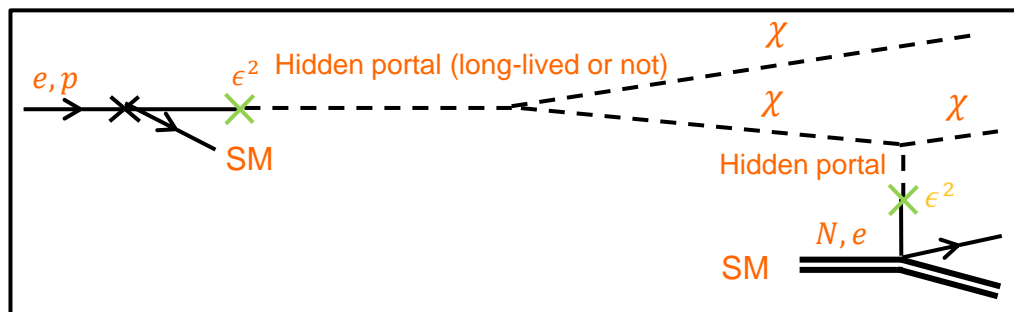
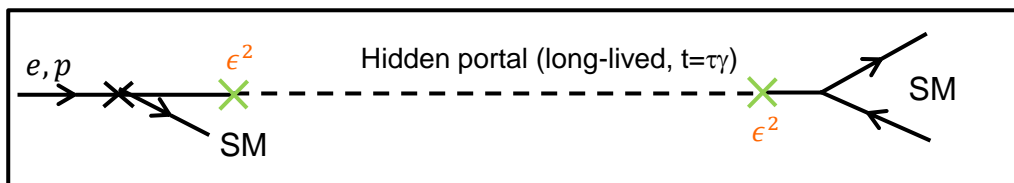
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- SHiP
- +
- SHiP
- SND@LHC

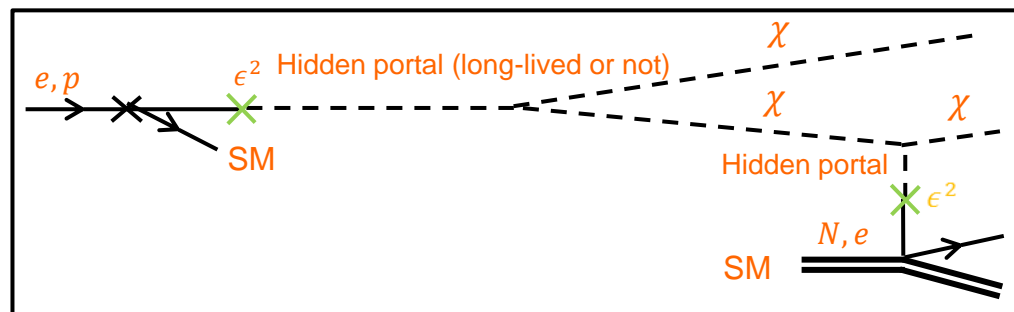
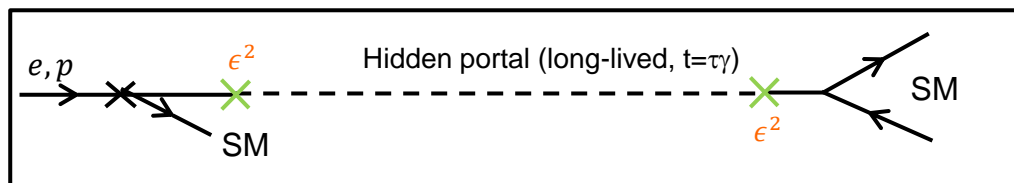
Long-lived search techniques

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 Model-dependent interpretation

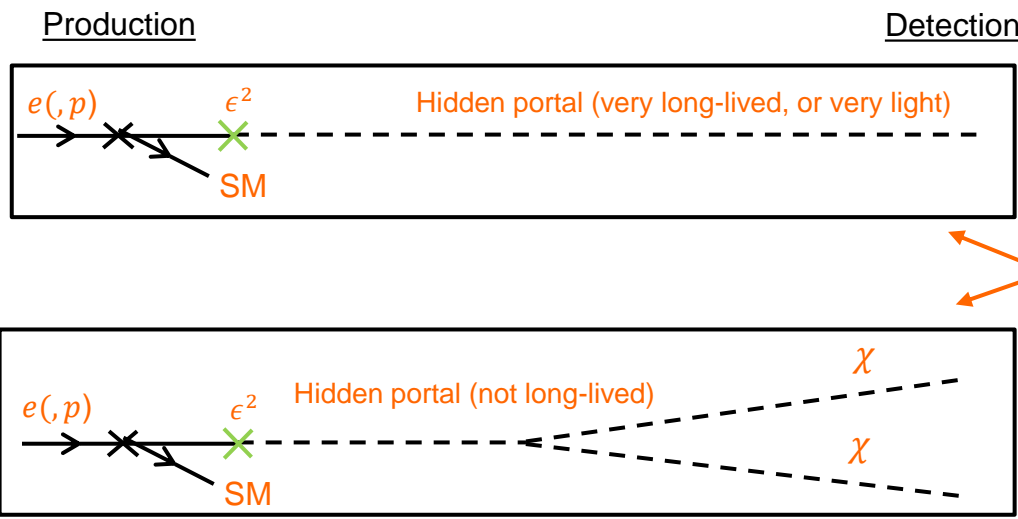
ϵ^4 dependence requires higher luminosity

- MATHUSLA
- ANUBIS
- CODEX-b
- +
- ATLAS
- CMS
- LHCb
- NA62 (++)
- NA64 (++)
- FASER
- SHIP
- +
- SHIP
- SND@LHC

Long-lived search techniques

Note : $\epsilon \ll 1$

Indirect (invisible)

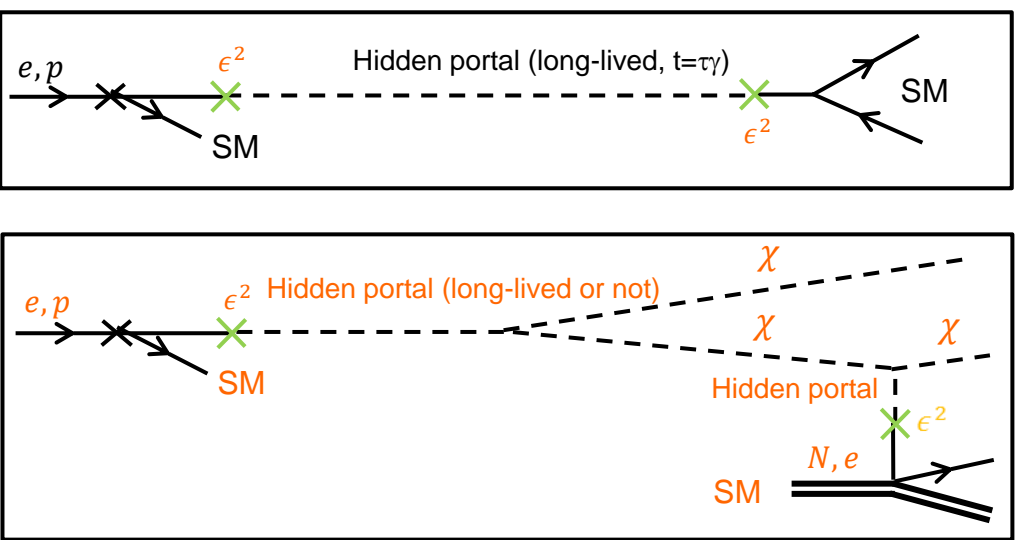


Escape detector
Missing energy/momentum/mass
Probability $\propto \epsilon^2$

Cannot be directly distinguished
→ Model-dependent interpretation

ATLAS
CMS
LHCb
NA64
LDMX

"Direct" (indirect visible)



Decay signature ("displaced vertex")
Probability $\propto \epsilon^4$
Model independent

+
Reconstruction of decay: mass, PID
→ Distinguish models
→ Measurement of properties

MATHUSLA
ANUBIS
CODEX-b

ATLAS
CMS
LHCb
NA62 (++)
NA64 (++)
FASER
SHiP

Dark matter scattering
Recoil against electron or nuclei
Probability $\propto \epsilon^4$
Model-dependent interpretation

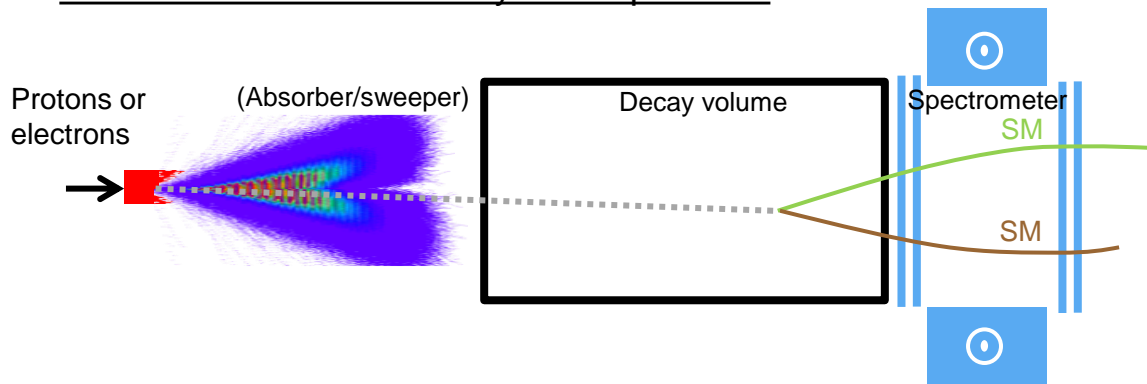
SHiP
SND@LHC

ϵ^4 dependence requires higher luminosity

→ Background situation very different in the different techniques!

Experimental setups

○ Direct search: visible decay to SM particles

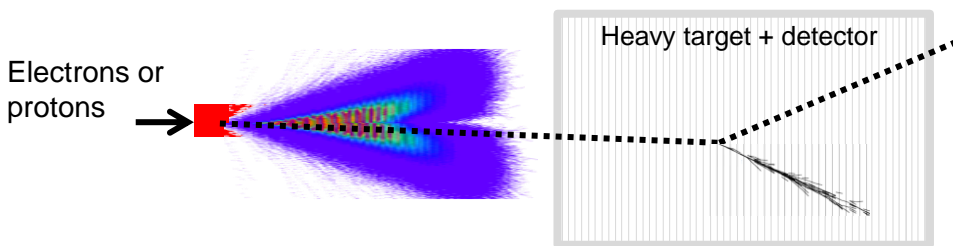


“Fixed target mode setups”:

NA62++@CERN (p@400, 10^{18})
 HPS, APEX, DarkLight@JLAB (e@1-10)
 SHiP@CERN (p@400, 2×10^{20}),
 SeaQuest@FNAL (p@120, 10^{18} – 10^{20})
 (LBNF@FNAL)

ATLAS, CMS, LHCb @LHC (no absorbers)
 BELLE2@sKEKB (no absorber)
 FASER@LHC
 MATHUSLA@LHC (no spectrometer)

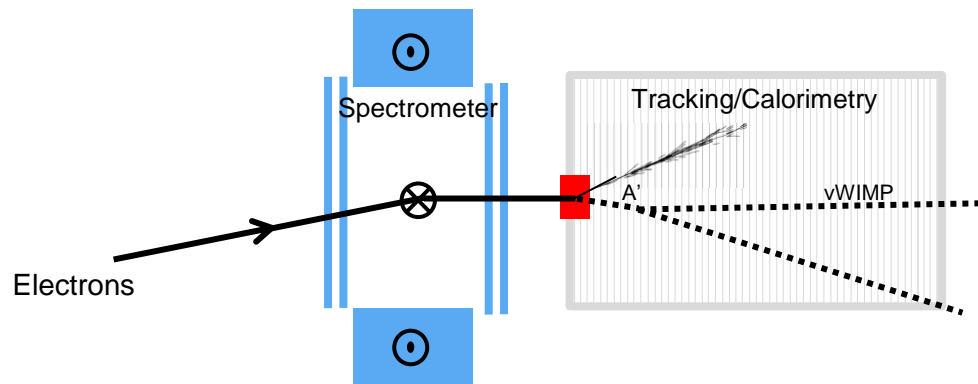
○ Direct search: Scattering off atomic electrons and nuclei



“Fixed target mode setups”:

BDX@JLAB (e@11, 10^{22}),
 MiniBooNE@FNAL (p@8.9, 10^{20}),
 SHiP@CERN (p@400, 2×10^{20})
 (interest for BDX-like
 experiments at
 LNF, Mainz (MESA),
 SLAC, Cornell...)

○ Indirect search: Missing energy/momentum (slow extraction/electron association)

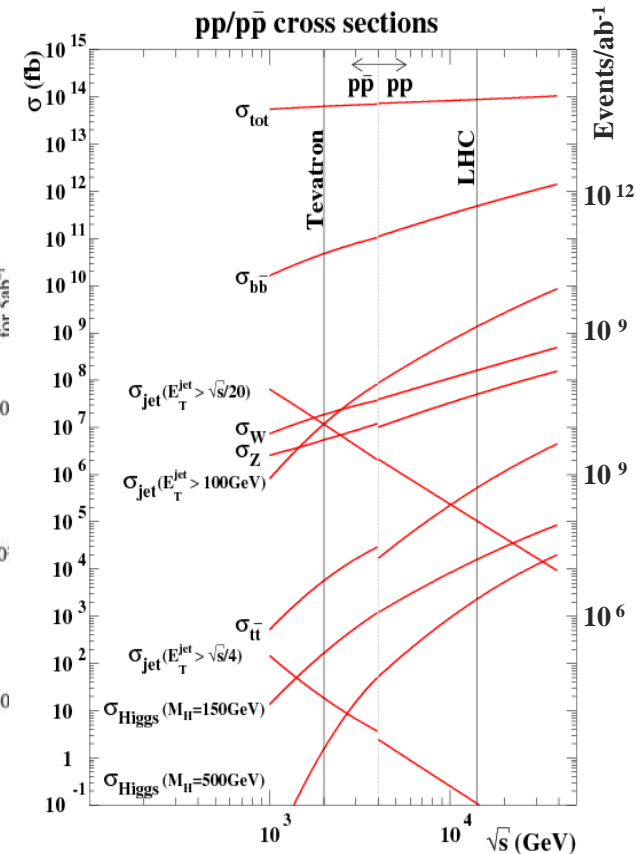
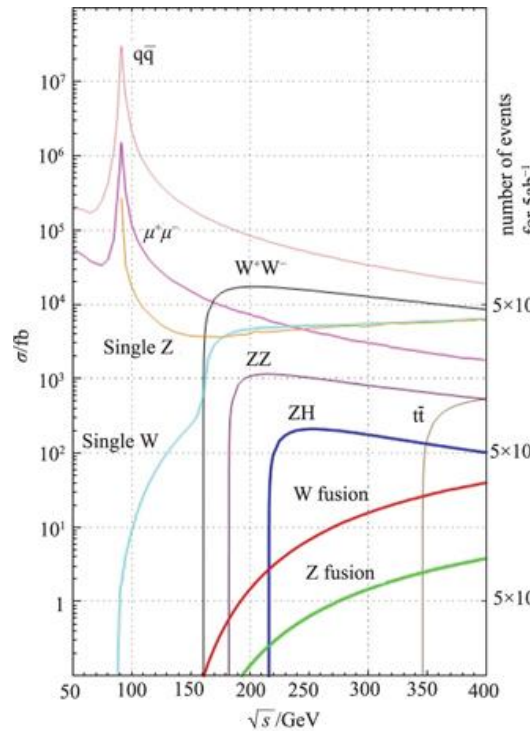
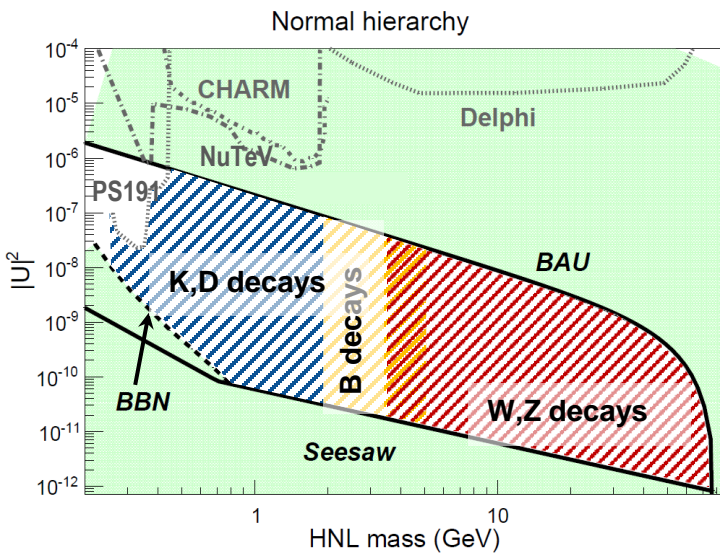


NA64/NA64++ @CERN (e@100, 10^{12})
 LDMX@SLAC/CERN (e@4-8/16, 10^{14} - 10^{16})

Putting it all together

○ We need massive production of $\gamma, g/q, b, W, Z, H$!

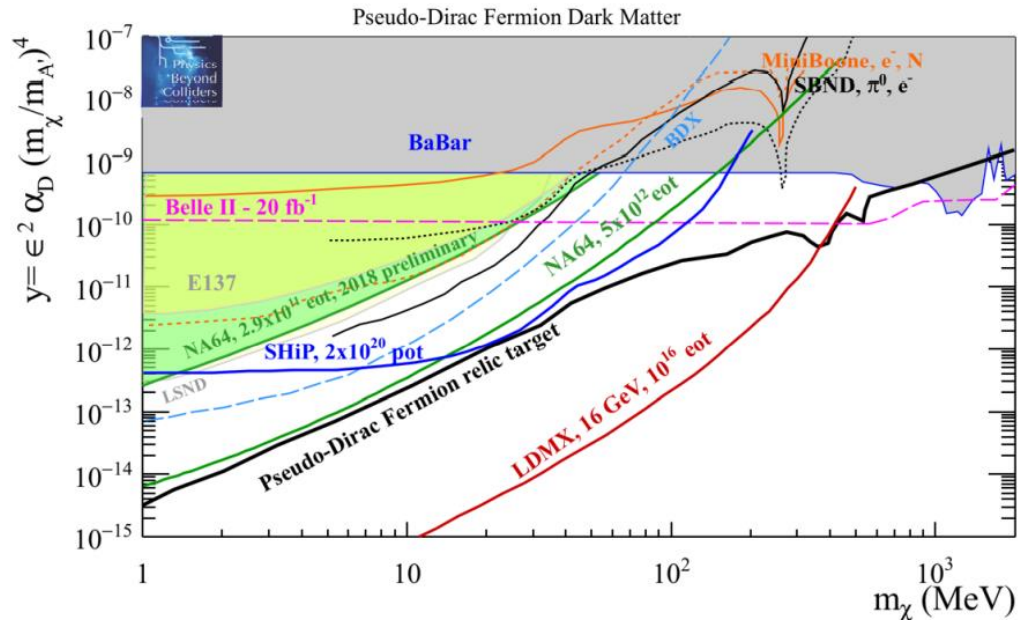
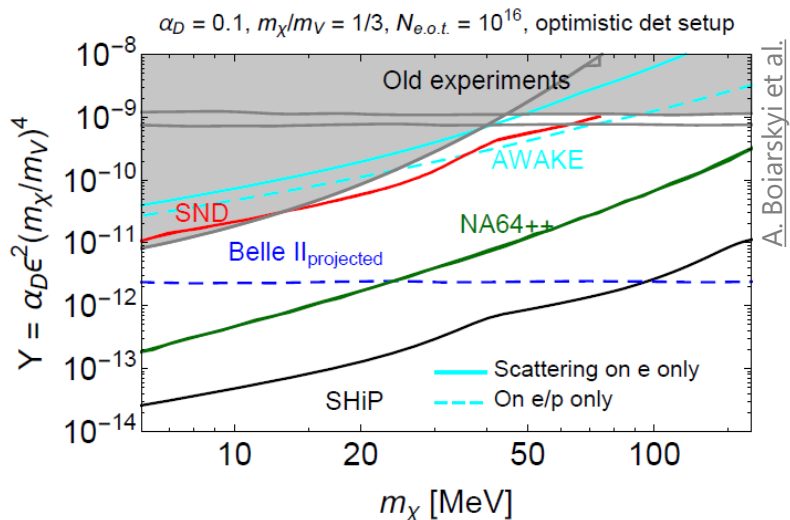
- ➔ General purpose machines: ee, pp colliders and proton injectors for beam dumps
- ➔ Higher energy is not the driver (cross-sections at lepton colliders, background at hadron colliders)



Putting it all together...

Other machines...

- Electron beam dumps
 - ➔ Limited to dark photon/ALP (γ -coupling) coupling
 - ➔ ϵ^2 - sensitivity suffers from measuring accurately each incoming electron
 - ➔ ϵ^4 - sensitivity suffers from electron yield



- Muon colliders – higgs and Drell-Yann production
- Neutrino facilities with off axis detector – lower energy (below m_K), light targets i.e. huge neutrino backgrounds...
- Gamma factory – low mass

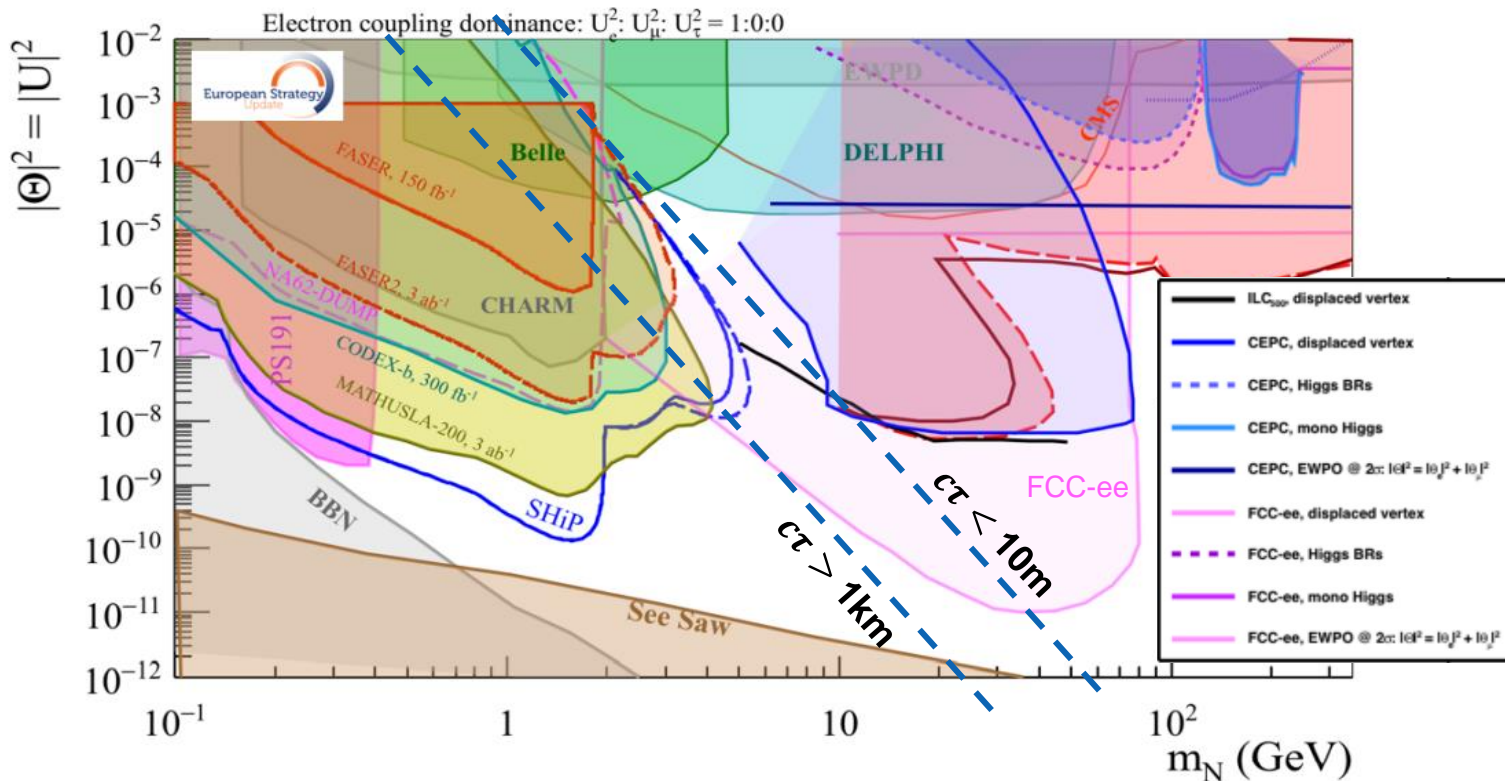
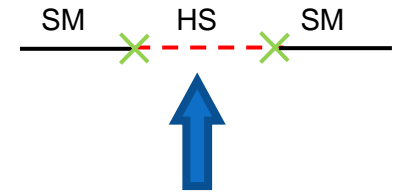


Caught between a rock and a hard place

Acceptance and background are the biggest challenges!

- Dilemma: background/pile-up versus absorbers/sweepers
- New states are typically long-lived, e.g. HNL $\tau_N \sim \frac{96\pi^2 h}{|U|^2 G_F^2 M_N^5}$

→ Lifetime $\otimes \epsilon \times 4\pi$ challenge



Complementarity

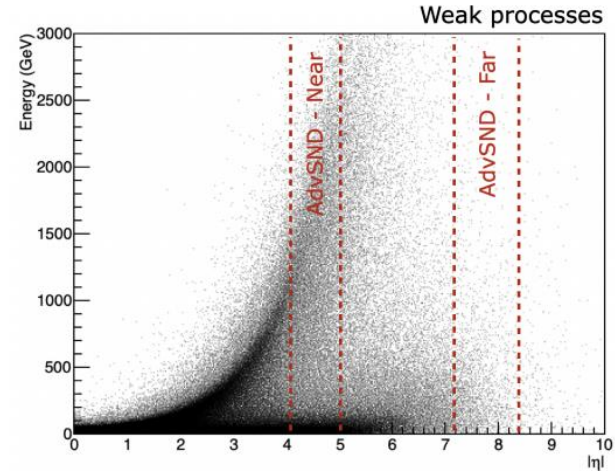
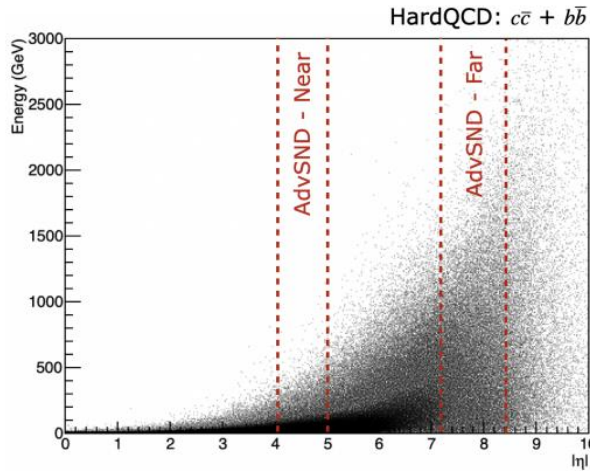
Complementarity between accelerator-based search methods and experimental configurations

- Classical collider detectors:
 - higher particle mass, complete geometric acceptance and short lifetimes
 - Displaced vertex $100\mu m \lesssim c\tau \lesssim 10m$ or missing energy
- Distant voluminous collider detectors
 - higher particle mass, *longer* lifetimes, limited geometric acceptance/ 4π and/or limited detection techniques
- Beam dump experiments
 - High luminosity and geometric acceptance (boost), lower particle mass, long lifetimes
 - E.g. SPS luminosity for a long target (e.g. 1m++ Mo/W) with 4×10^{19} pot/year
 - Maximum atomic number and charge (=short X_0/λ_{int}) to increase cross-section for charm, beauty hadron and radiative processes AND reduce neutrino flux by stopping π^\pm, K^\pm
 - Tuned to $m_K - m_b$ mass range
 - SPS $\mathcal{L}_{int}[yr^{-1}] = 10^6 s \times \int_0^{\infty} \Phi_0 \times \rho_N \times e^{-l/\lambda} dl = \Phi_0 \times \rho_N \times \lambda = \underline{3.6 \times 10^{45} \text{ cm}^{-2}}$ (cascade not incl. $\sim 2.6x$ for charm)
 - HL-LHC $\mathcal{L}_{int}[yr^{-1}] = 10^7 s \times 10^{35} \text{ s}^{-1} \text{ cm}^{-2} = \underline{10^{42} \text{ cm}^{-2}}$
 - LHC ($\sqrt{s} = 14 \text{ TeV}$): with 1 ab^{-1} , i.e. 3-4 years: $\sim 2 \times 10^{16}$ D's in 4π
 - SPS@400 ($\sqrt{s} = 27 \text{ GeV}$) with 2×10^{20} pot, i.e. ~ 5 years: $\sim 2 \times 10^{17}$ D's

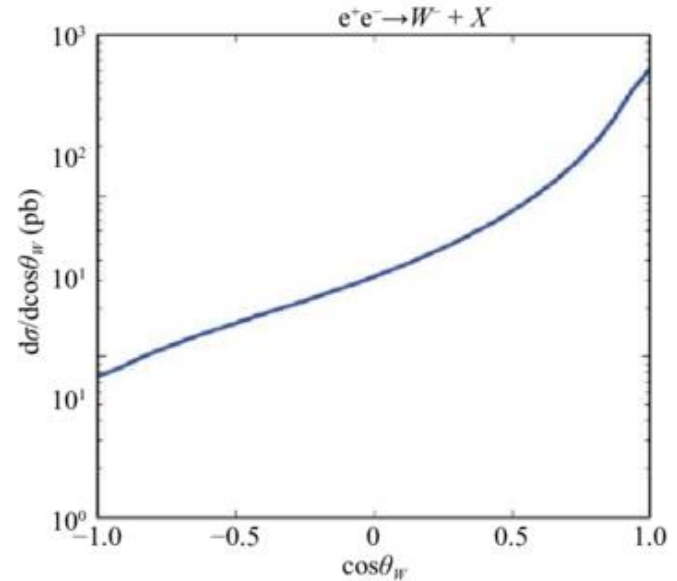
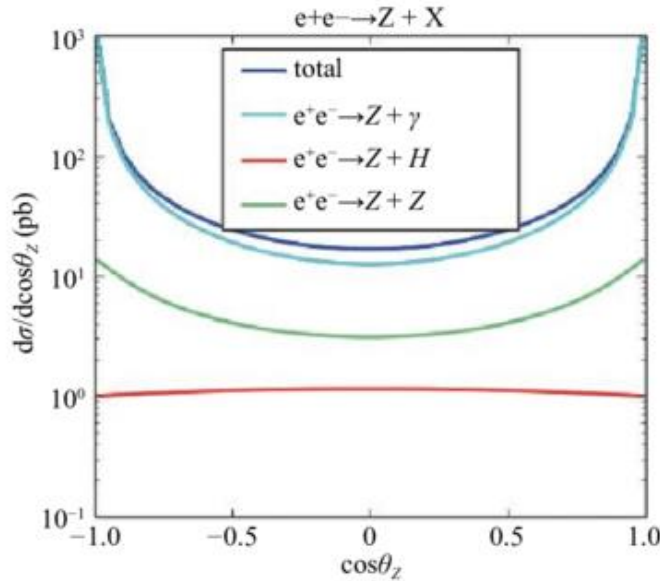
Forward production at colliders

- Beam-dump style detectors at colliders (asymmetric?)

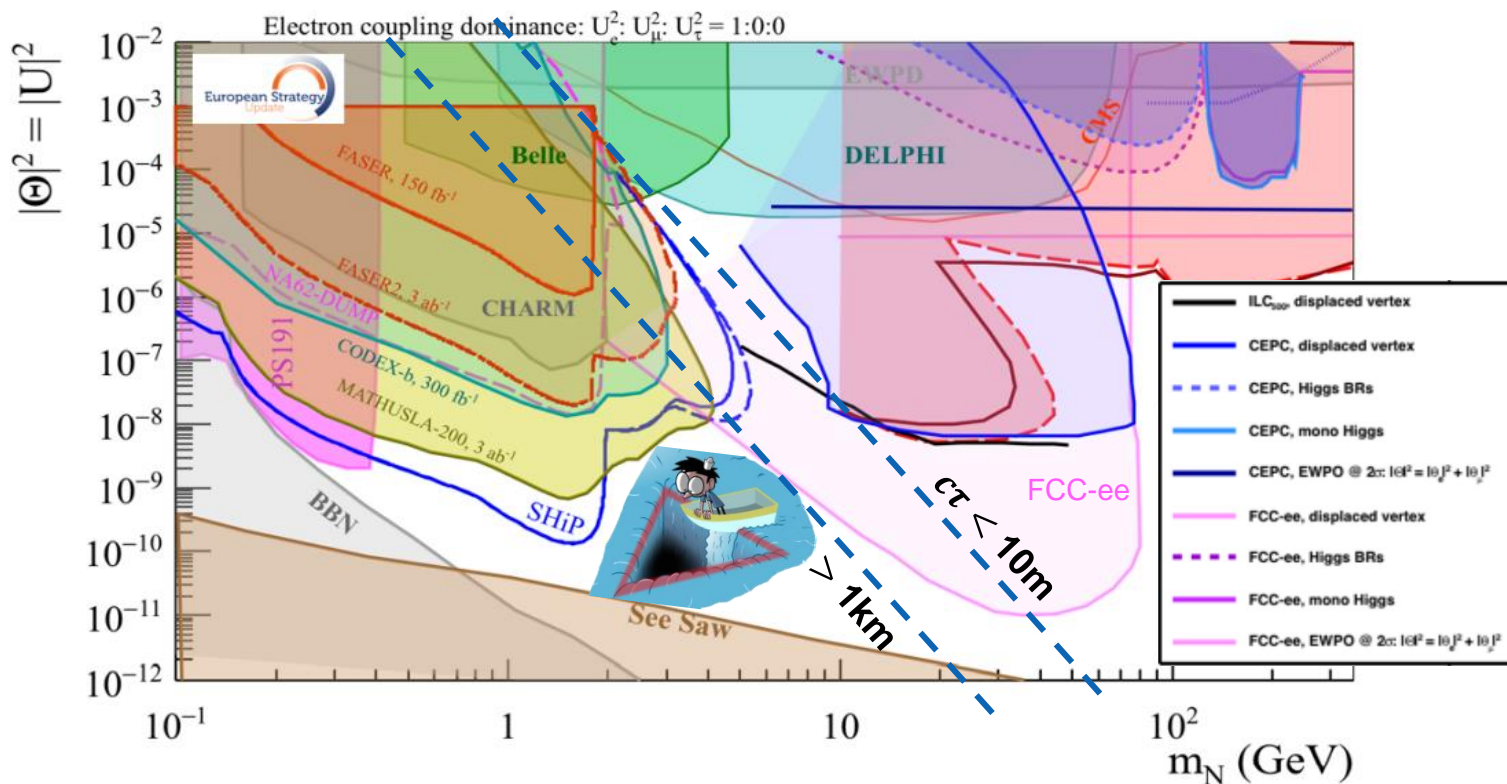
pp-collider



e⁺e⁻-collider



Hidden hidden sectors...



Bermuda triangle?

Production in B and W, Z decays with $c\tau\gamma \sim kmS \dots$

t -channel production from W exchange and forward large detectors
a la *proton-dump* at distance from a collider IP?

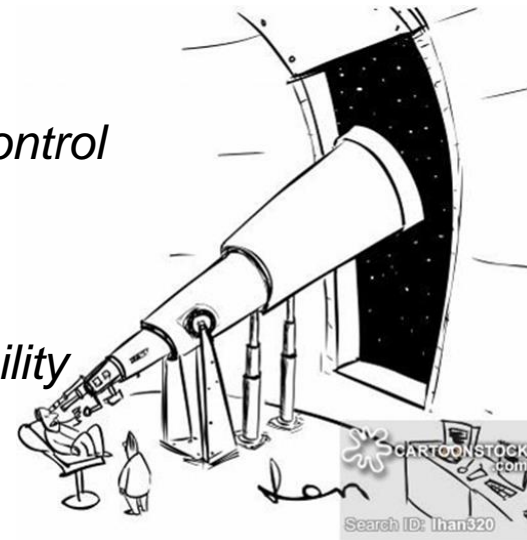
A not so little plea

Do I want to go home with an exclusion, having a doubt that I control the background... ?

or worse...

discovering something with no clue what it is and with no capability to characterise the signal...?

and perhaps still be worrying about my control of the background...?

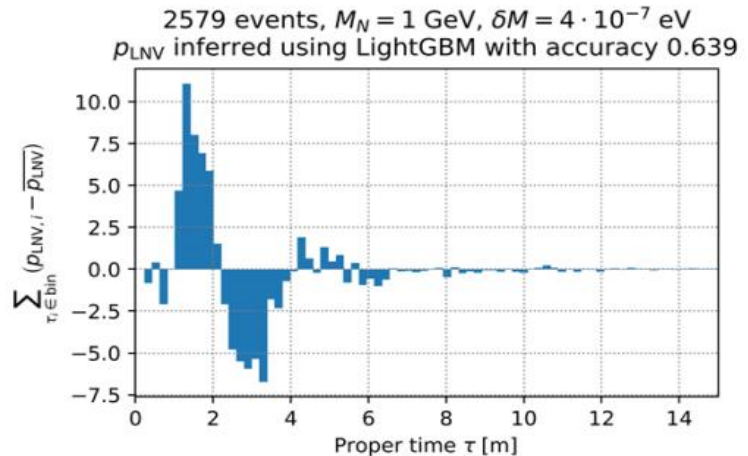
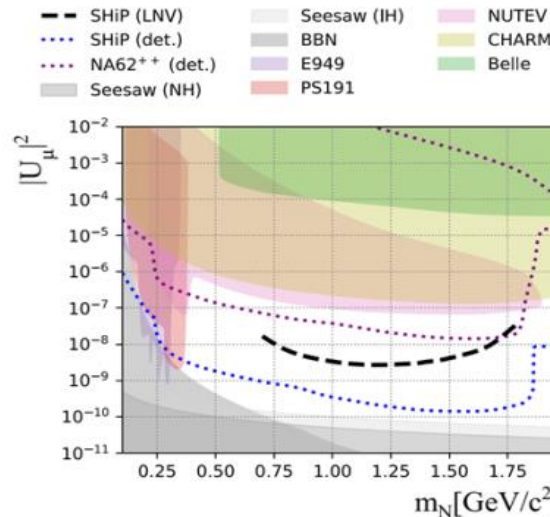
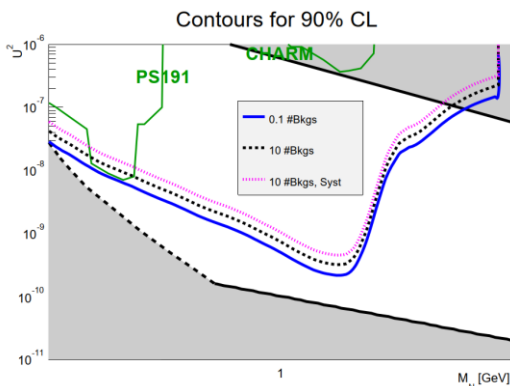


"I've either discovered dark matter, or I've left the lens cap on."

Example of model characterisation by SHiP in case of discovery

→ Sensitivity to oscillations between lepton number conserving and violating rates

[arXiv:1912.05520 \(2020\)](https://arxiv.org/abs/1912.05520)



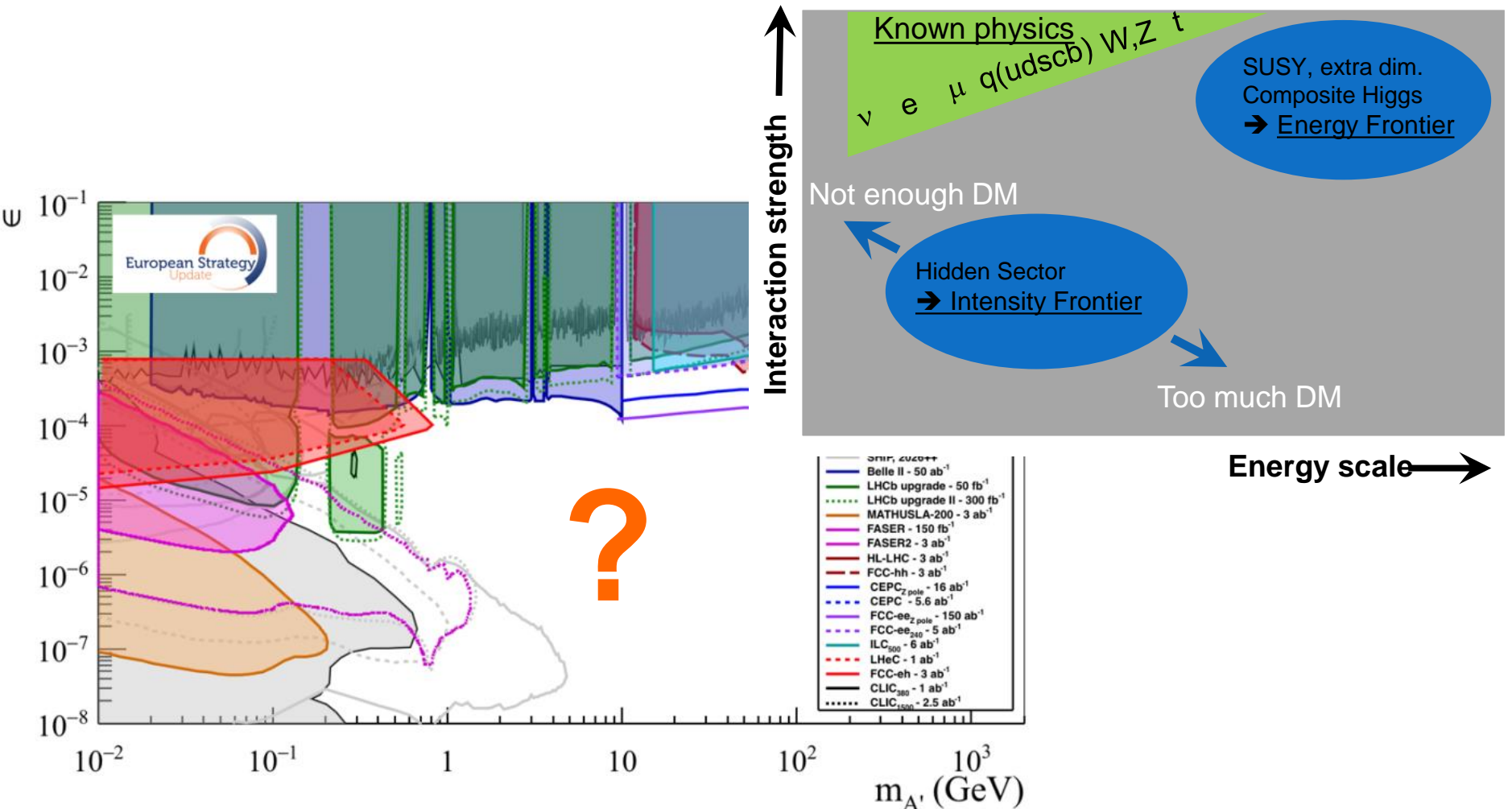
Sensitivity – devil is in the background

“exclusion plots *under the assumption of zero background*”, what does it mean?

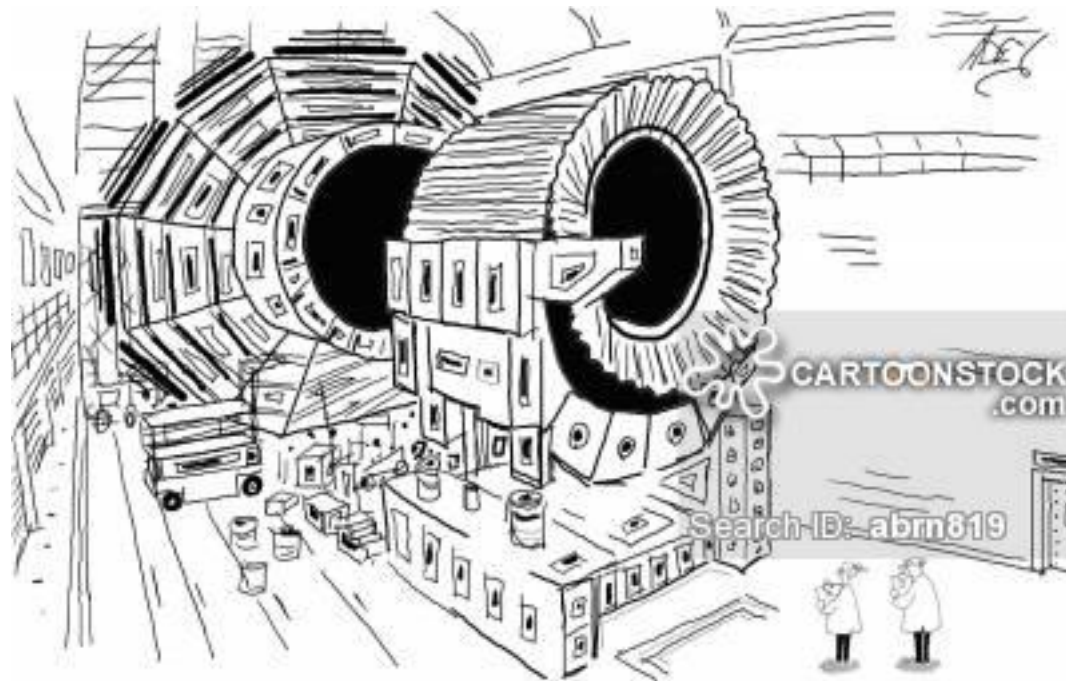
- To exclude or discover (compute the significance of the observation) a hypothetical signal, the experimental data should be tested against the new physics model AND against a background only hypothesis, to see which of the two can be excluded
 - with appropriate well-understood uncertainties on the background
- **Optimizing an analysis for discovery or exclusion may lead to a different setup and selection criteria**
 - Optimising for discovery aim for very high signal-to-background ratio, even at the cost of acceptance and efficiency. Separating events into high signal-to-background ratios and low signal-to-background ratios classes and combining the results gives optimal sensitivity.
 - Optimising for exclusion aim at improving signal acceptance at the cost of letting in more background.
 - **Difference in signal/background of hadron and lepton machines**
- **Control of background essential**
 - **Binning of observables from event reconstruction: kinematics, mass, particle identification**
 - Importance of control regions to constrain uncertainties on background
 - Machine Learning for optimisation relying on multiple quantities
 - **Detailed simulations (several models for event generators and detector simulations), subsidiary measurements, full simulations... define maturity**

Another plea...the log scale deception

- Large number of options and huge parameter spaces
 - All parameter spaces have their “unreachable” regions, even physically attractive regions!
- Theoretical model building and cosmofrontier are essential guides!



Conclusion



“If all else fails - it makes a great frothy latte.”