# Vector meson photoproduction in heavy ions at LHC

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# OUTLINE

- Ultraperipheral collisions
- Vector meson photoproduction
- Heavy quarkonia
- Light vector mesons
- Summary



### Ultrapheripheral collisions (UPCs)

- When ions don't "touch" each other: Impact parameter  $b > R_1 + R_2$ :
- These processes dominate nucleusnucleus cross section.
- Electromagnetic (photon) interactions dominate.
- Photons are quasi-real:  $Q^2 \lesssim 1/R_A^2$
- Large photon flux  $\sim Z^2$  (Fermi/ Weizsacker-Williams). In the case of Pb, probability 82<sup>4</sup> larger than proton!





### Ultrapheripheral collisions (UPCs)





### **Experimentally very clean events !**







# Types of processes in UPCs

### **Photon - photon:**



• BSM physics

QED precision tests

### **Photon - target:**



**Gluons from nucleus (target)** 

- Sensitive to saturation gluon saturation
- Test nuclear PDFs

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# Types of processes in UPCs

### **Photon - photon:**



- BSM physics
- QED precision tests



. . .



**Gluons from nucleus (target)** 

Sensitive to saturation gluon saturation Test nuclear PDFs





### **UPCs and saturation**

- energies.
- recombinations.
- $Q_{\rm S} \propto A^{1/3}$



Photon resolution power (Q)-





# Vector meson photoproduction

- Vector mesons: particles with spin 1 and odd parity.
- Photon quantum numbers ( $J^{PC} = 1^{--}$ ) same for
- VM  $\rightarrow$  highly likely a photon will fluctuate into a VM !
- VM photoproduction cross section  $\propto (xg(x, Q^2))^2$  at LO  $\rightarrow$  sensitive to saturation
- Photoproduced VM cross section at small x can test gluon density
- In particular, heavy quarkonia (J/ $\psi$ , Y), sets an enough large scale for pQCD.





### Vector meson photoproduction kinematics

• A given  $y \rightarrow$  Fixes  $\omega, x, W$ 



• 
$$\omega = \frac{M_{VM}}{2} e^{\pm y}$$

- y: Rapidity of the VM
- $\omega$ : Photon energy
- $M_{VM}$ : Mass of the VM

• 
$$x = \left(\frac{M_{VM}}{\sqrt{s_{NN}}}\right) e^{\mp y}$$

• 
$$W^2 = M_{VM} \sqrt{s_{NN}} \cdot e^{\pm y}$$

 W: Centre-of-mass energy of the photontarget system



### **Coherent and Incoherent VM production**

### Coherent:

- Photon couples with the nuclei as a whole
- Nuclei usually remains intact
- Incoherent:
  - Photon couples to single nucleons
  - Nuclei usually breaks
- In both cases we can have additional photon exchanges:
  - Nuclei collective excitation
  - Neutron emission as de-excitation processes.









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### Photoproduction in pp/ep vs. AA



- Advantage:
- Wavelength  $\lambda \propto R_{\rm proton}$ , then higher photon virtualities (better "resolution power".).
- Disadvantage:
- Low photon flux



- Disadvantage:
- Wavelength  $\lambda \propto R_{Ion}$ , then lower photon virtualities (worse "resolution power").
- Advantages:
- High photon flux (enhanced by  $Z^2$ ).
- Saturation scale easier to be accessed (scale by  $A^{1/3}$ ).



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# Looking inside protons...

 $\gamma + p \rightarrow J/\psi + p$ 



- Investigated with ep, pPb and pp collisions by HERA and LHC
- Consistent results between
  HERA and LHC data.
- Data follow a power-law trend, consistent with the rapidly increasing gluon density.
- No evidence for saturation !

(dn) (d+ψ/L <

a(γ+p



EPJC 79 (2019) 402

### Looking inside Pb...

 $\gamma + Pb \rightarrow J/\psi + Pb$ 



### • Gluons inside Pb:

•  $\sigma(J/\psi) < \text{Impulse}$ 

Approximation(simplest model only putting protons and neutrons together)

 $\rightarrow$  strong nuclear modification in nuclei.

 No theory model can describe data at both forward and central rapidities.



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# Looking inside Pb...

low and high energy photons.



$$\frac{d\sigma_{AA\to AA'J/\psi}}{dy} = N_{\gamma/A}(y) \cdot \sigma_{\gamma A\to J/\psi A'}(y) + N_{\gamma/A}(-y) \cdot \sigma_{\gamma A\to J/\psi A'}(-y)$$



# A solution to the two-way ambiguity puzzle

### **Control impact parameter of UPCs via forward neutron emissions** Ann. Rev. Nucl. Part. Sci. 70 (2020) 323



Additional photon exchanges lead to neutron emissions via EMD.

•  $b_{XnXn} < b_{0nXn} < b_{0n0n}$ 





### Total Coh. $J/\psi$ cross section in neutron categories



arXiv:2303.16984

- ZDC allows to classify events in neutron categories.
- First separation in different neutron categories.



### Coh. $J/\psi$ photo nuclear cross section vs W



- ZDC categories allow us to obtain the W dependence.
- W< 40 GeV: rapidly increasing
- Higher W: slow rise  $\rightarrow$ underlying physics changed!
- First evidence of saturation?
- No models can describe the entire data distribution!







### What about even higher mass?



- different dipole size.
- Important to compare with saturation/shadowing models.



Photoproduction of Y (composed of bb), allows to probe nuclei with a

• The high mass (even higher than J/ $\psi$ ) sets a sufficiently large scale for pQCD.





### Lighter vector mesons

- In Asymmetric systems, such as pPb, photon comes from the Pb nucleus (no emitter/target ambiguity).
- Light vector mesons are more challenging for pQCD.
- LHC data complement HERA extensive program.





### Lighter vector mesons

### In symmetric collisions, we have quantum interference.



### Manifested as angular modulations







### Lighter vector mesons



- It is expected that interference effects vanish at large impact parameter.
- Effect observed in ALICE data.





# Summary

- UPCs are very clean processes to test QED and QCD phenomena.
- Photoproduction of VMs in UPCs are powerful probes to test the gluonic structure of nuclei.
- The topic has brought the attention lately in the LHC community and continues the HERA legacy.
- Many interesting ongoing related studies in CMS, to be unveiled soon !



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10<sup>2</sup>

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10<sup>3</sup>

W (GeV)



# Backup slides



### **Event classification via neutron multiplicity**

•Energy depositions in ZDCs allow to classify events in neutron categories.







### A solution to the two-way ambiguity puzzle

What is measured

**Dominant b ranges of** different neutron classes:

- 0n0n: b > 40 fm٠
- $0nXn: b \sim 20 \text{ fm}$ ٠
- XnXn: b < 15 fm•

$$\frac{d\sigma_{AA\to AAJ/\psi}^{0n0n}}{dy} = N_{\gamma}^{0}$$
$$\frac{d\sigma_{AA\to AA'J/\psi}^{0nXn}}{dy} = N_{\gamma}^{0}$$
$$\frac{d\sigma_{AA\to AA'J/\psi}^{0nXn}}{dy} = N_{\gamma}^{0}$$

 $\rightarrow$  Solve for  $\sigma_{\gamma A \rightarrow J/\gamma}$ 

## What we want Photon flux from theory $\sigma_{\gamma/A}^{0n0n}(y) \cdot \sigma_{\gamma A \to I/\psi A'}(y) + N_{\gamma/A}^{0n0n}(-y) \cdot \sigma_{\gamma A \to I/\psi A'}(-y)$ $\int_{V/A}^{0nXn}(y) \cdot \sigma_{\gamma A \to I/\psi A'}(y) + N_{\gamma/A}^{0nXn}(-y) \cdot \sigma_{\gamma A \to I/\psi A'}(-y)$ $\sum_{\nu/A}^{XnXn}(y) \cdot \sigma_{\nu A \to I/\psi A'}(y) + N_{\nu/A}^{XnXn}(-y) \cdot \sigma_{\nu A \to I/\psi A'}(-y)$

$$\psi_{A'}(y)$$
 and  $\sigma_{\gamma A \to J/\psi A'}(-y)$ , and  $x = \left(\frac{M_{VM}}{\sqrt{s_{NN}}}\right) e^{-\frac{1}{2}}$ 

Entering a new regime of small  $x \sim 10^{-4} - 10^{-5}$  in nuclei!



