

Pythia 8: Physics and usage

Saariselkä Midsummer School 2024

Ilkka Helenius

June 28, 2024



Outline

Lecture 1:

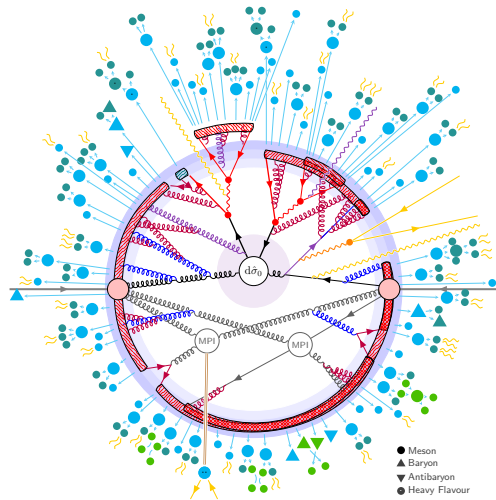
- History of Pythia
- Monte Carlo techniques
- Hard-process sampling

Lecture 2:

- Multiparton interactions
- Parton showers

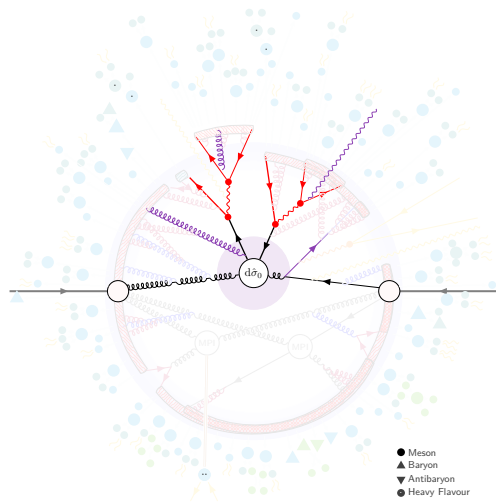
Lecture 3:

- Hadronization
- Beam configurations



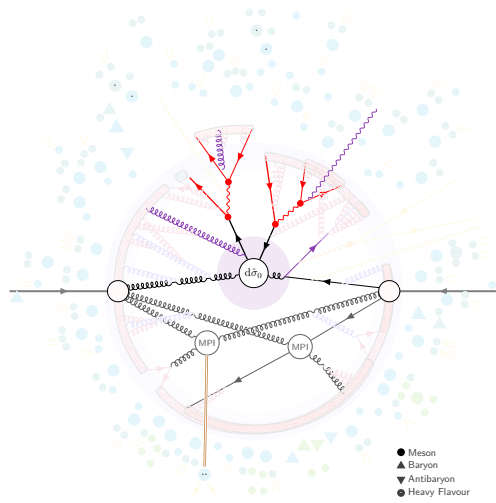
[figure by P. Skands]

Lecture 2:



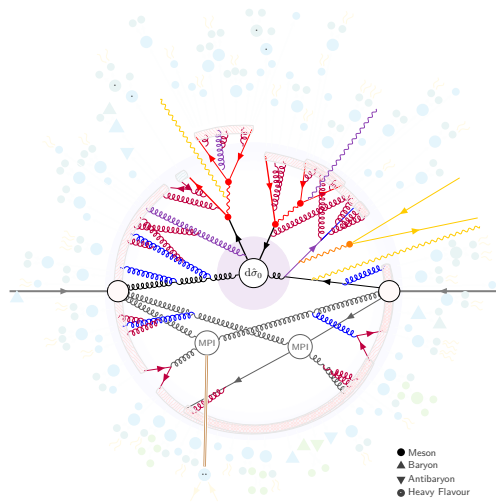
Lecture 2:

- Multiparton interactions



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- Parton showers



Multiparton interactions

Multiparton interactions (MPIs)

- Integrated cross section for QCD $2 \rightarrow 2$ processes

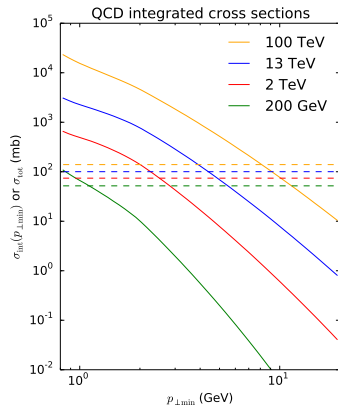
$$\sigma_{\text{int}}(p_{T,\text{min}}) = \int_{p_{T,\text{min}}}^{\sqrt{s}/2} dp_T \frac{d\sigma^{2 \rightarrow 2}}{dp_T^2}$$

- σ_{int} exceeds σ_{tot} when p_T^2 small
⇒ Multiple partonic interactions per event
- Partonic cross section diverges at $p_T \rightarrow 0$
⇒ Introduce a screening parameter p_{T0}

$$\frac{d\sigma^{2 \rightarrow 2}}{dp_T^2} \propto \frac{\alpha_s(p_T^2)}{p_T^4} \rightarrow \frac{\alpha_s(p_{T0}^2 + p_T^2)}{(p_{T0}^2 + p_T^2)^2}$$

- Energy-dependent parametrization:

$$p_{T0}(\sqrt{s}) = p_{T0}^{\text{ref}}(\sqrt{s}/\sqrt{s_{\text{ref}}})^\alpha$$



Multiparton interactions (MPIs)

- Now a finite QCD cross section

$$\sigma_{\text{int}}(p_{T,0}) = \int_0^{\sqrt{s}/2} dp_T \frac{d\sigma^{2 \rightarrow 2}(p_{T,0})}{dp_T^2}$$

$\sigma_{\text{nd}}(\sqrt{s})$ is the non-diffractive cross section

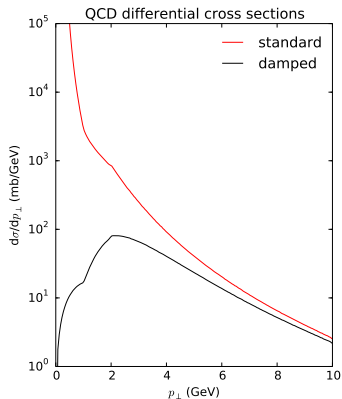
- Number of interactions: $\langle n \rangle = \sigma_{\text{int}}(p_{T,0}) / \sigma_{\text{nd}}$

1. Sample n independent interactions

2. Order interactions in p_T

- Energy conservation by rescaling PDFs
- Sudakov factor to account for the p_T -ordering

$$\mathcal{P}_{\text{MPI}}(p_T) = \frac{1}{\sigma_{\text{nd}}(\sqrt{s})} \frac{d\sigma^{2 \rightarrow 2}(p_{T,0})}{dp_T^2} \exp \left[- \int_{p_T}^{p_{T,\text{max}}} dp_T'^2 \frac{1}{\sigma_{\text{nd}}(\sqrt{s})} \frac{d\sigma^{2 \rightarrow 2}(p_{T,0})}{dp_T'^2} \right]$$



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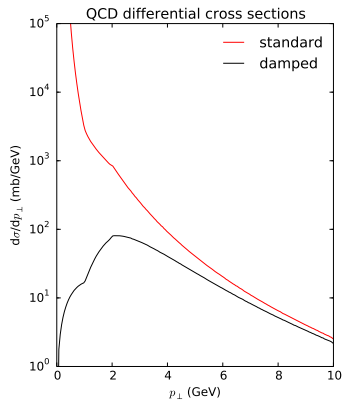
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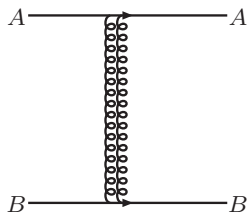
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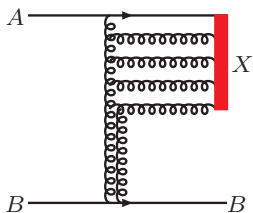
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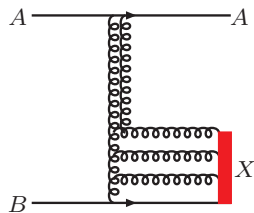
Soft QCD processes



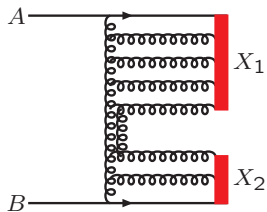
elastic



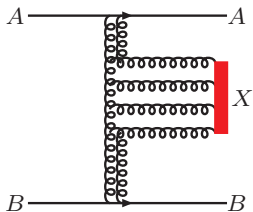
single diffractive (XB)



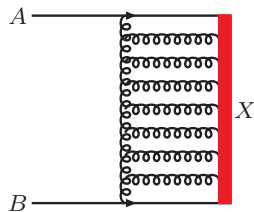
single diffractive (AX)



double diffractive



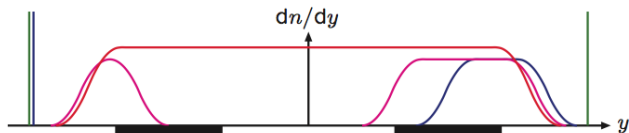
central diffractive



nondiffractive

[figure by T. Sjöstrand]

Particle production in minimum bias collisions



[figure by T. Sjöstrand]

- Diffractive processes relevant at large rapidities
- Non-diffractive at mid-rapidities

Soft-QCD non-diffractive processes

- Need to sample QCD processes without any phase-space cuts
 - Standard cross section diverges when $p_T \rightarrow 0$
- ⇒ Use the regulated cross sections from the MPI framework
 - Contain the same QCD sub-processes as hard QCD event class
 - ⇒ Could also be used for jet production (with infinite statistics)

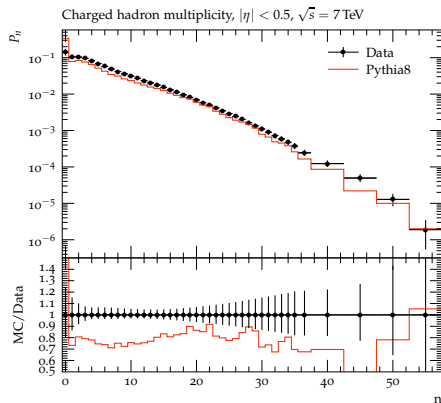
Exercise III: Charged-particle multiplicities

CMS analysis: CMS_2011_S8884919

- Number of charged particles in p+p collisions at 7 TeV
- Non-single-diffractive (NSD) events

Exercise IIIa:

- Go to Pythia online manual <https://pythia.org//latest-manual/Welcome.html>
- Find how to enable processes for the NSD trigger (Process Selection -> Soft QCD Processes)
- Make a `.cmd` file and run with `pythia8-main93`



[CMS: JHEP 01 (2011) 079]

Exercise III: Charged-particle multiplicities

Exercise IIIb:

1. Vary $p_{T,0}^{\text{ref}}$ parameter

```
MultipartonInteractions:  
  pT0Ref = 2.28
```

- Run with `pythia8-main93`

2. Turn off MPIs completely with

```
PartonLevel:MPI=off
```

- Run with `pythia8-main93`
- Compare results

Exercise III: Charged-particle multiplicities

Exercise IIIb:

1. Vary $p_{T,0}^{\text{ref}}$ parameter

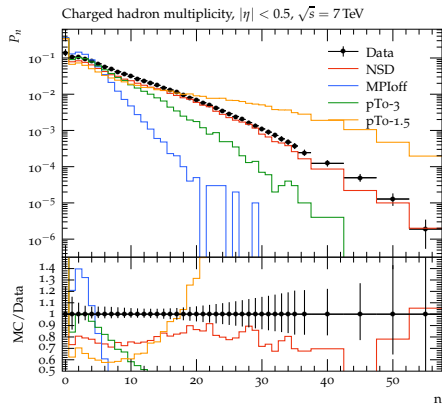
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[CMS: JHEP 01 (2011) 079]

Parton showers

Final state radiation

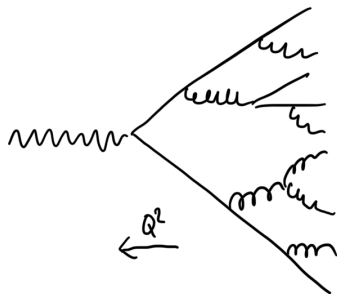
- Splitting probabilities from DGLAP

$$d\mathcal{P}_{a \rightarrow bc} = \frac{dQ^2}{Q^2} \frac{\alpha_s}{2\pi} P_{a \rightarrow bc}(z) dz$$

- Iterative structure, emissions ordered in Q^2
- Use Sudakov factor to account for the ordering (A parton can only branch if it had not already branched)

$$d\mathcal{P}_{a \rightarrow bc} = \frac{dQ^2}{Q^2} \frac{\alpha_s}{2\pi} P_{a \rightarrow bc}(z) dz \exp \left[- \sum_{b,c} \int_{Q^2}^{Q_{\max}^2} \frac{dQ'^2}{Q'^2} \int dz' \frac{\alpha_s}{2\pi} P_{a \rightarrow bc}(z') \right]$$

- Several options for the ordering variable, different phase-space mapping
⇒ Pythia uses $p_T = z(1-z)Q^2$



Initial state radiation

- Start from a highly-virtual parton participating to the hard scattering
- Backwards evolution, trace back splittings that have occurred before the hard interactions (\sim undo DGLAP evolution of the PDFs)
- Need to consider conditional probability for the splitting

$$d\mathcal{P}_{a \leftarrow b} = \frac{df_b}{f_b} = \frac{x' f_a(x', Q^2)}{x f_b(x, Q^2)} \frac{dQ^2}{Q^2} \frac{\alpha_s}{2\pi} P_{a \rightarrow bc}(z) dz \quad (x' = x/z)$$



- Similarly need a Sudakov factor to account for the non-emission probability
- Evolution variable in Pythia: $p_T = (1 - z)Q^2$

Parton-level evolution

Common evolution scale (p_T) for FSR, ISR and MPIs

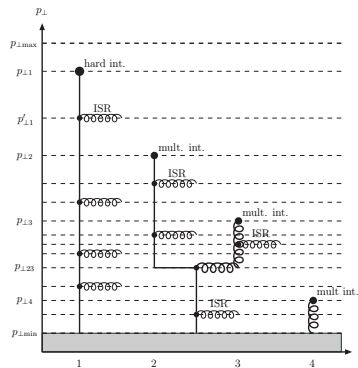
- Probability for something to happen at given p_T

$$\frac{d\mathcal{P}}{dp_T} = \left(\frac{d\mathcal{P}_{\text{MPI}}}{dp_T} + \sum \frac{d\mathcal{P}_{\text{ISR}}}{dp_T} + \sum \frac{d\mathcal{P}_{\text{FSR}}}{dp_T} \right) \times \exp \left[- \int_{p_T}^{p_T^{\text{max}}} dp'_T \left(\frac{d\mathcal{P}_{\text{MPI}}}{dp'_T} + \sum \frac{d\mathcal{P}_{\text{ISR}}}{dp'_T} + \sum \frac{d\mathcal{P}_{\text{FSR}}}{dp'_T} \right) \right]$$

where $\exp[. . .]$ is a Sudakov factor (probability that nothing else has happened before p_T)

Simultaneous partonic evolution

1. Start the evolution from the hard-process scale
2. Sample p_T for each \mathcal{P}_i , pick one with highest p_T
3. Continue until $p_{T\text{min}} \sim \Lambda_{\text{QCD}}$ reached



[T. Sjöstrand, P. Skands:
EPJC 39 (2005) 129-154]

Parton shower options in Pythia: The Simple Shower

The default parton shower implementation

[T. Sjöstrand, P. Skands, EPJC 39 (2005) 129-154]

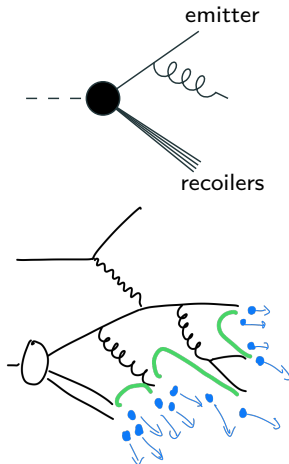
- The only one available until 8.3 release
- Improved Leading-Log (LL) approximation
- Recoil distributed to the whole final state
- Applicable to photoproduction

DipoleRecoil variant

[B. Cabouat, T. Sjöstrand, EPJC 78 (2018 no.3, 226)]

- Restrict recoil to the emitting dipole instead of the whole final state
- Can be applied processes like DIS or Vector-Boson fusion

DGLAP



Parton shower options in Pythia: Vincia and Dire

Vincia antenna shower `PartonShowers:model = 2`

[H. Brooks, C. T. Preuss, P. Skands, JHEP 07 (2020) 032]

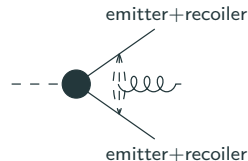
- QCD, QED, EW, interleaved with MPIs
- Interleaved resonance decays
- Iterated LO matrix-element corrections
- Efficient multi-jet merging with sectors

Dire in Pythia `PartonShowers:model = 3`

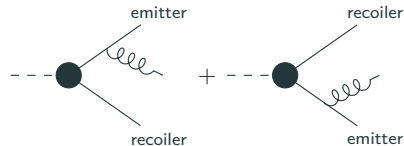
[S. Höche, S. Prestel, EPJC 75 (2015) no.9, 461]

- QCD, QED, \sim EW and dark photons
- Correct soft-gluon interference at lowest order
- Inclusive NLO corrections to collinear splittings
- Recoil given to the non-emitting side of the dipole

Antennae



Dipoles



[figures by C. Preuss] 12

Exercise IV: Z-boson p_T

ATLAS analysis: ATLAS_2011_S9131140

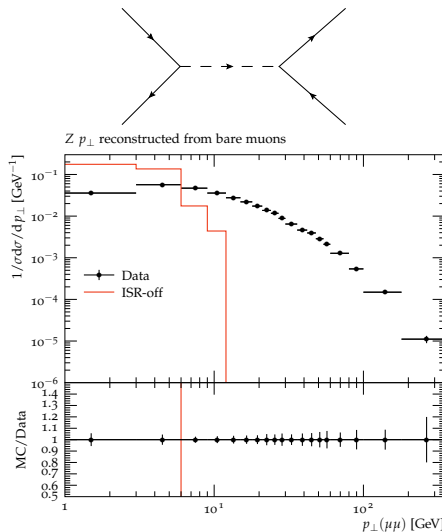
- Transverse momentum of Z-boson
 \Rightarrow LO calculation: $p_T^Z = 0$
- Reconstructed from lepton pair

Exercise IVa:

1. Enable Z production and turn off ISR

```
WeakSingleBoson:ffbar2gmZ = on  
PartonLevel:ISR=off
```

- Run and compare to data
2. Turn ISR on
 - Compare to data and result with ISR



Exercise IV: Z-boson p_T

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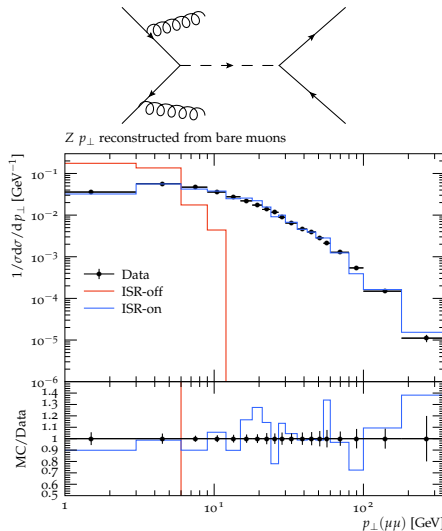
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Exercise IVb:

- Generate events $\mathcal{O}(10k)$ with all three parton-shower models, Default (Simple), Vincia and Dire

```
PartonShowers:model=1,2,3
```

- Run and compare to data

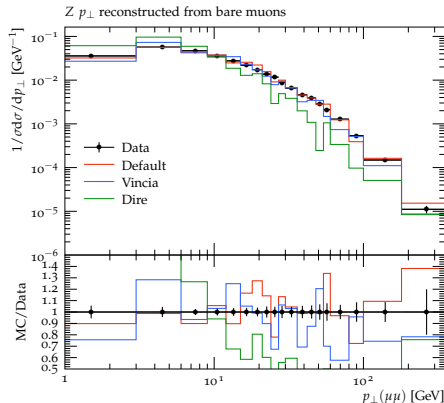
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[ATLAS: PLB 705 (2011) 415-434]

Results with 100k events

Number of MPIs

- Study the online manual how to extract number of MPIs from the event information
 - Calculate the average number of MPIs in non-diffractive (min. bias) events and in events with a Z boson
- ⇒ Are these the same?

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 - Calculate the average number of MPIs in non-diffractive (min. bias) events and in events with a Z boson
- ⇒ Are these the same?
- Pedestal effect: Harder processes with smaller impact parameter
- ⇒ More MPIs and larger $\langle n_{\text{ch}} \rangle$

