

Soft QCD and the Production of Secondary Cosmic Rays

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Outline

● EAS : Longitudinal distributions

- ➔ Heitler model: X_{\max}
- ➔ Longitudinal Profile
- ➔ Energy Deposit

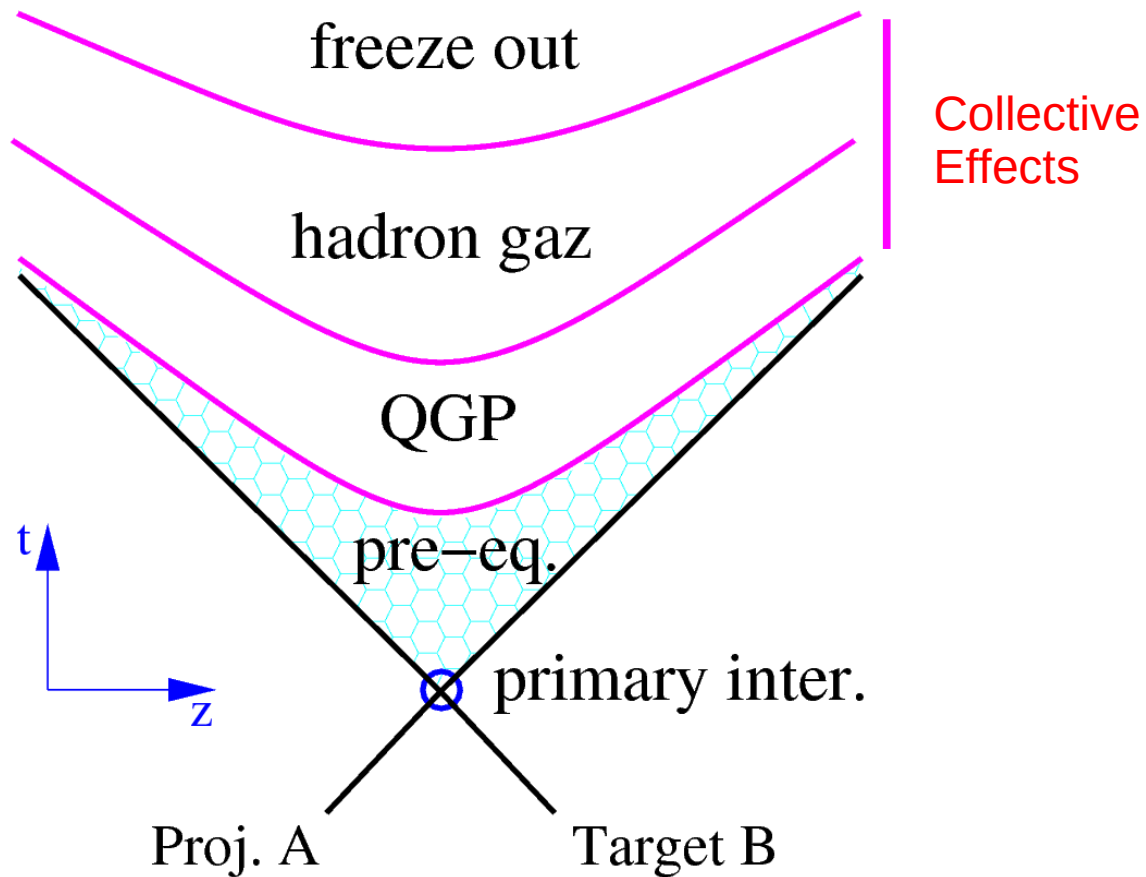
● EAS : Particles at ground

- ➔ Heitler model: N_{μ}
- ➔ Particles at ground
- ➔ Muon puzzle
- ➔ Muon production depth

● EAS : Link with hadronic interactions

- ➔ Hadronic Observables
- ➔ [Hadronic interaction models for Cosmic Rays](#)

Goal of Today's Lecture...



The complete description of hadronic interactions is universal ...
my biased view !

What can be used as projectile/target ?

- Any charged particle with sufficiently long life time can be **accelerated** (electrons (e⁻), positrons (e⁺), protons (p), nuclei (A), hadrons (h), leptons (l) ...) :

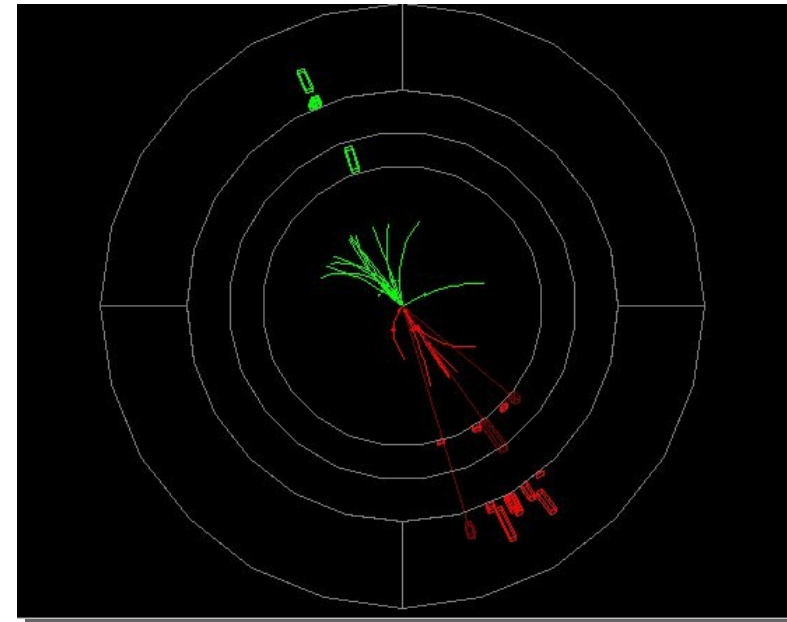
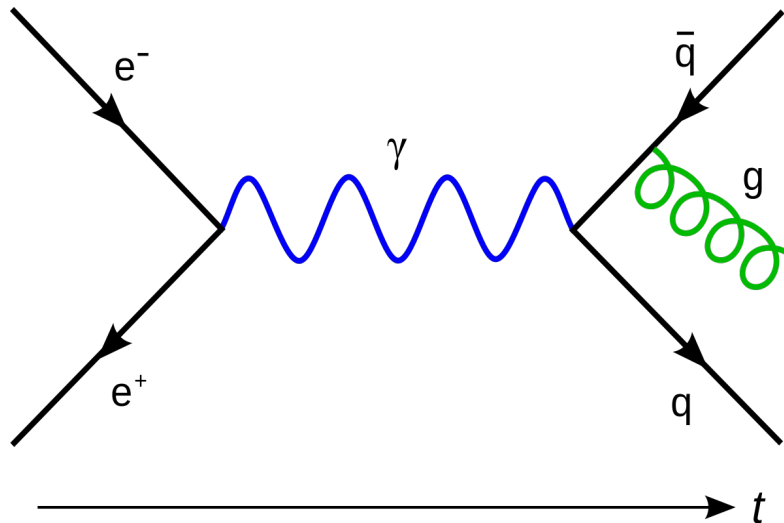
e⁺ vs e⁻ → l vs p (A) → p vs p → h vs A → A vs A

What can be used as projectile/target ?

- Any charged particle with sufficiently long life time can be **accelerated** (electrons (e^-), positrons (e^+), protons (p), nuclei (A), hadrons (h), leptons (l) ...) :



- fundamental particles
- study of new particles
- “clean”
- difficult to go to very high energies ($\sim 100 \text{ GeV}/c$)

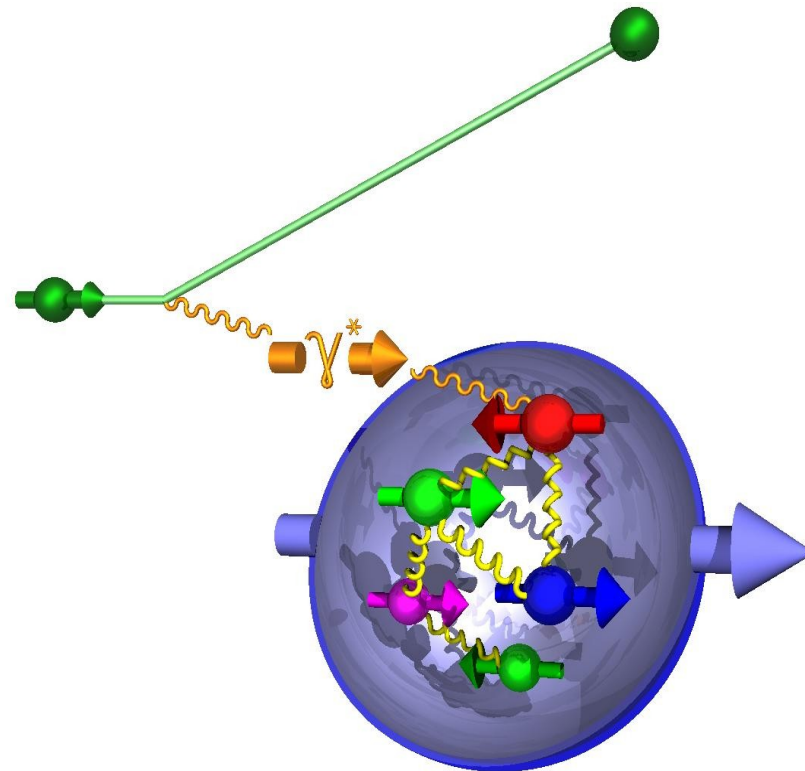


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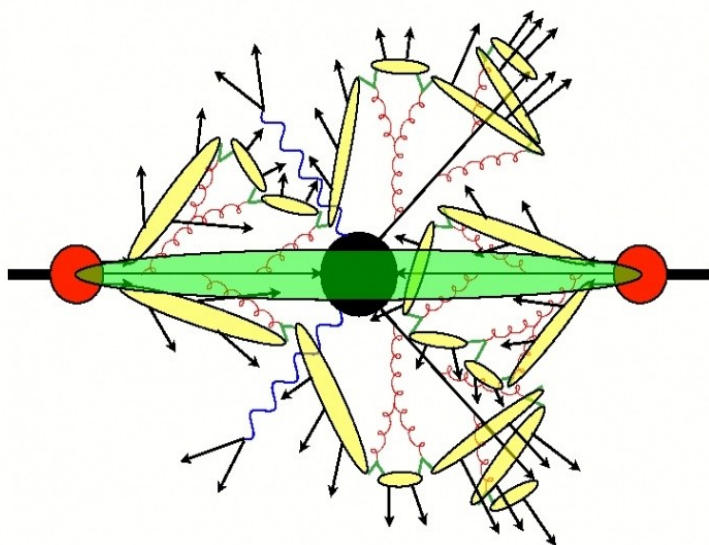
- ◆ one particle with inner structure
- ◆ study proton structure
- ◆ asymmetric detectors



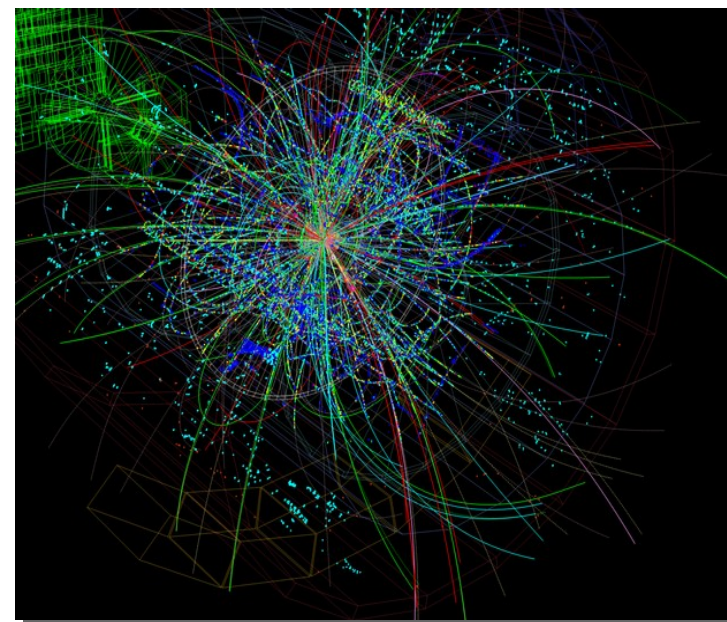
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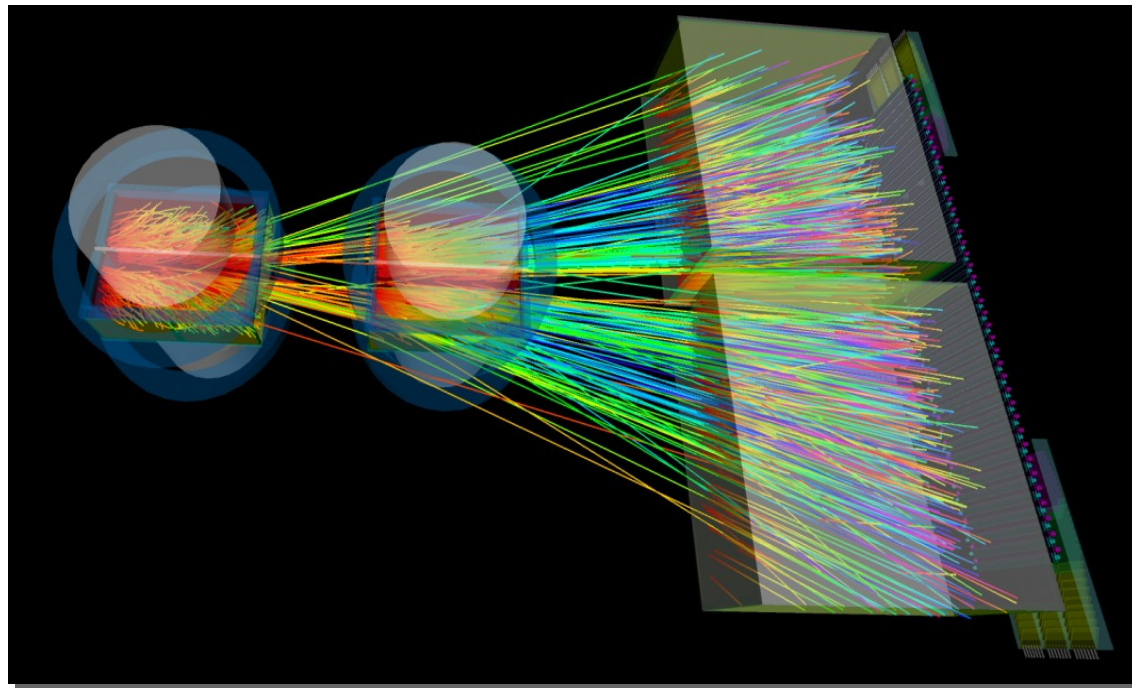
- ◆ complex particles
- ◆ study of new particles
- ◆ multiple interactions of inner structure
- ◆ “dirty”
- ◆ very high energy possible (~10 TeV/c)



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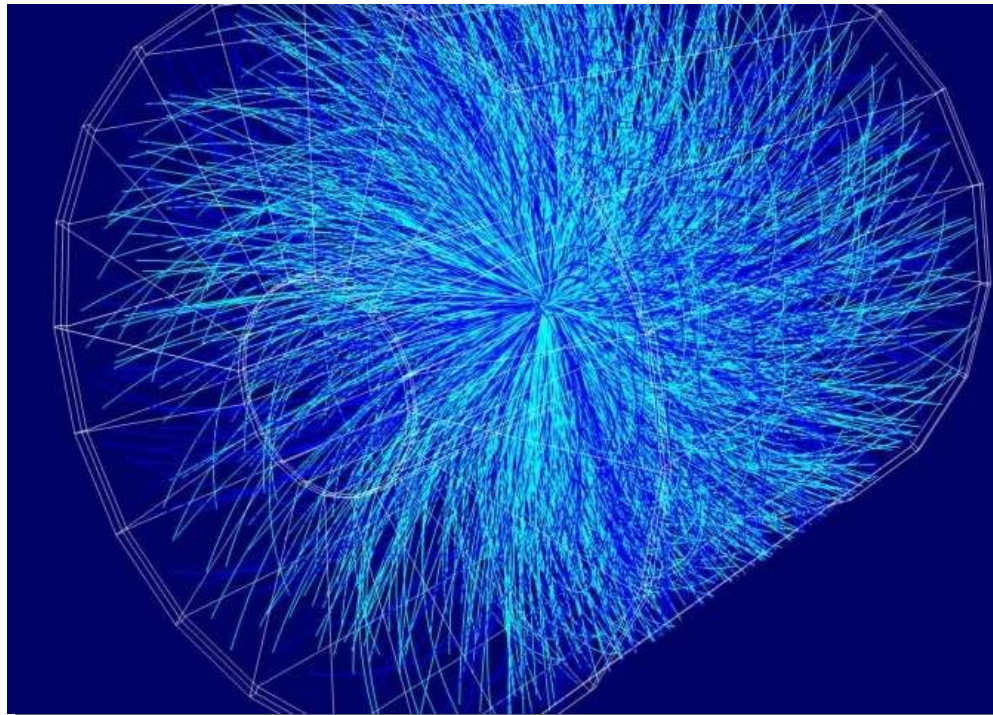
NA61

- ◆ complex particles with short life time
- ◆ important to understand particle cascade
- ◆ fixed target only: “forward” physics
- ◆ only low energy possible (~10 GeV/c)

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ALICE

- ◆ very complex particles
- ◆ study the very early state of the Universe
- ◆ very “dirty”
- ◆ very high energy densities

What can be used as projectile/target ?

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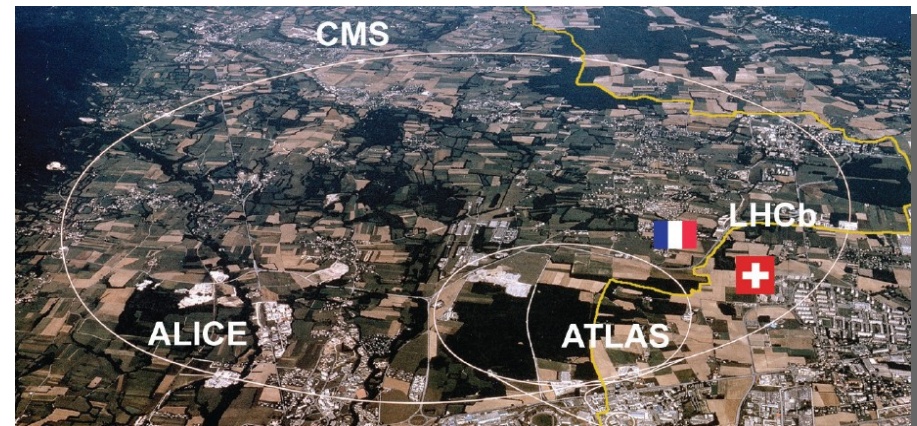
- Different particles = different structure of matter

What can be used as projectile/target (2) ?

- Any charged particle with sufficiently long life time can be **accelerated** (electrons (e⁻), positrons (e⁺), protons (p), nuclei (A), hadrons (h), leptons (l) ...) :

e⁺ vs e⁻ → l vs p (A) → p vs p → h vs A → A vs A

- | | | | | |
|--------------|----------|-------------------|--------------|--------------|
| ◆ LEP (CERN) | ◆ HERA | ◆ LHC (CERN) | ◆ SPS (CERN) | ◆ LHC (CERN) |
| ◆ SLC/PEP | (DESY) | ◆ Tevatron (FNAL) | ◆ (FNAL) | ◆ RHIC (BNL) |
| (SLAC) | ◆ (SLAC) | ◆ RHIC (BNL) | | |
| ◆ KEKB (KEK) | ◆ (CERN) | | | |
| | ◆ (FNAL) | | | |



- Different particles = different detectors

What can be used as projectile/target (3) ?

- Any charged particle with sufficiently long life time can be **accelerated** (electrons (e⁻), positrons (e⁺), protons (p), nuclei (A), hadrons (h), leptons (l) ...) :

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- ◆ Pythia
- ◆ HERWIG
- ◆ Phojet
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- ◆ Hydro models
- ◆ EPOS

- Different particles = different physics = different models

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High Energy Physics models

Cosmic Ray models

Both

- Different particles = **building blocks for a complete approach**

What is measured by detectors ?

● 2 types of scatterings

→ **elastic** : only momentum transfer between projectile and target : **particle shape**

→ **inelastic** : new particles are produced : **standard model**

● Only particles with long enough life time can be directly observed :

→ proton, neutron, charged pions, charged kaons, electrons, muons and photons

→ easier to measure charged particles

● Hadronic models necessary to compare theory processes with data by a proper **hadronization** of more fundamental particles

→ hadronization can not be calculated from first principles (non-perturbative QCD process)

→ phenomenological approach compared to data to fix (tune) parameters

→ importance of the different “building blocks”

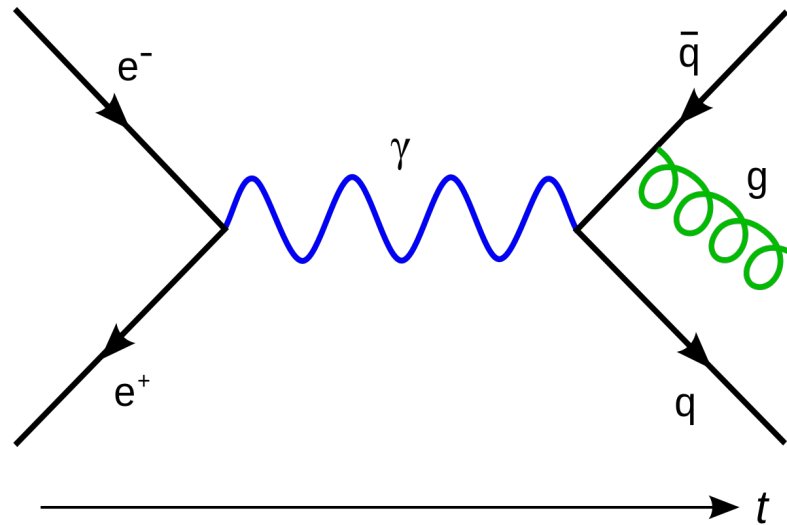
→ different type of particle scattering

“soft” Low Momentum Particle Production

- **All high energy physics analysis** (top, Higgs, Electroweak, supersymmetry, ...) **rely on the hadronic interaction models** :
 - ➔ directly : pQCD, hadronization of top jets, particle decay, missing energy ...
 - ➔ indirectly : detector simulations, background, underlying events ...
- **Different type of hadronic models**
 - ➔ all based on **Monte Carlo** methods : intensive use of random numbers
 - ➔ High Energy Physics models (HEP) : event build around a selected **hard process (can be used as minimum bias event generator)**
 - fast and precise
 - need data to fine tune the parameters (low predictive power for “soft”)
 - ➔ Cosmic Ray models (CR) : **minimum bias event generator only**
 - slow and do not include very rare hard processes
 - one set of parameter to reproduce all relevant data

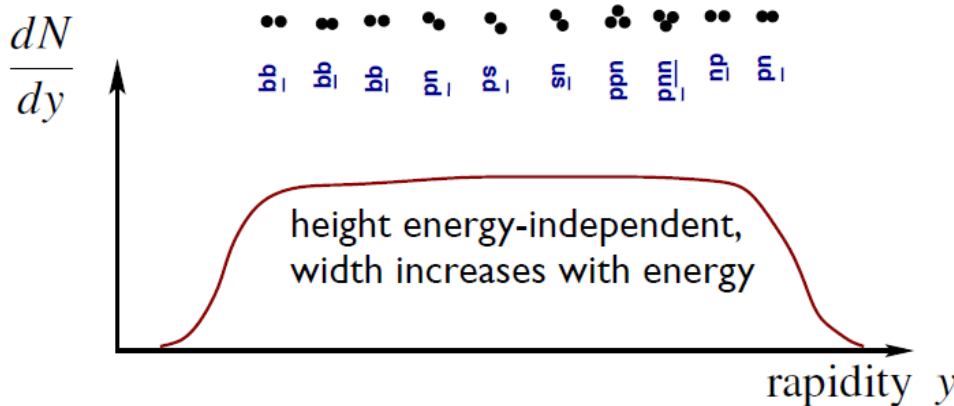
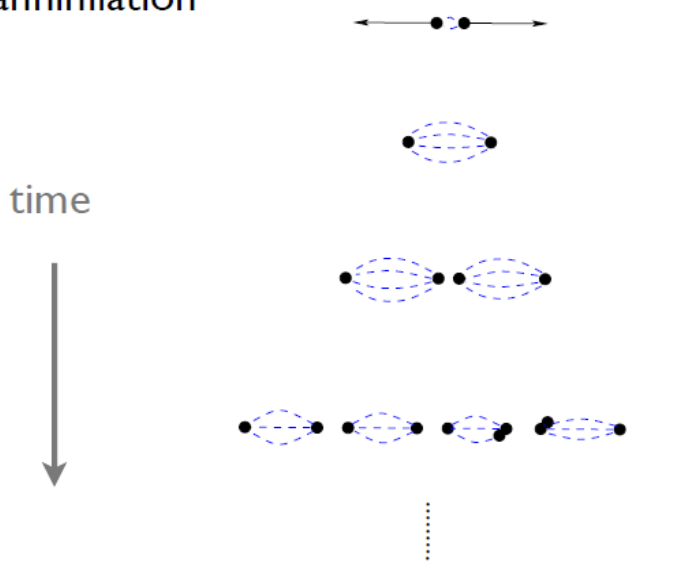
Hadronization of Quarks

e⁺ vs e⁻ → l vs p (A) → p vs p → h vs A → A vs A



String fragmentation and rapidity

Example:
q-qbar pair produced
in e⁺e⁻ annihilation



Rapidity

$$y = \frac{1}{2} \ln \frac{E + p_{\parallel}}{E - p_{\parallel}}$$

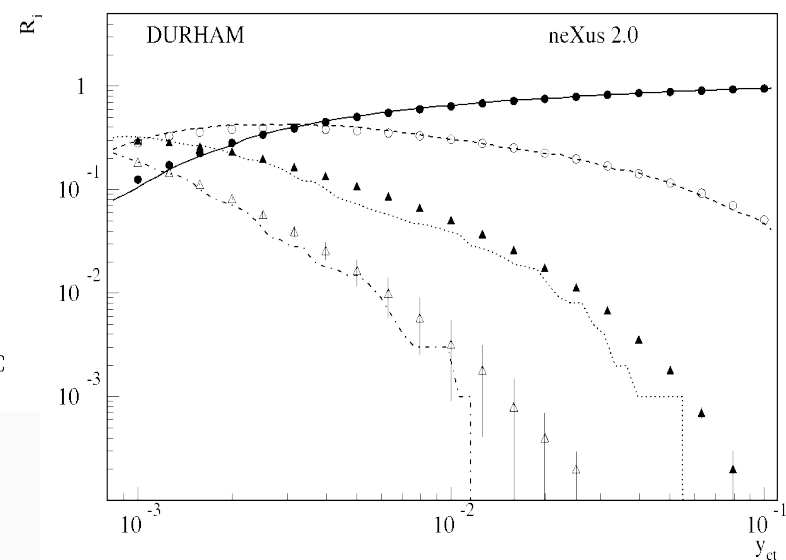
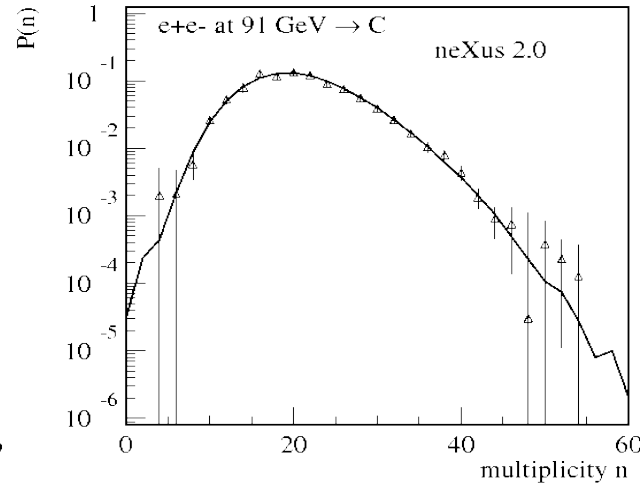
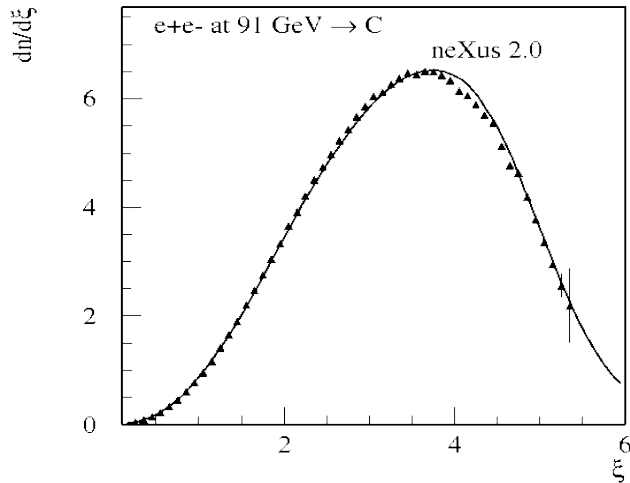
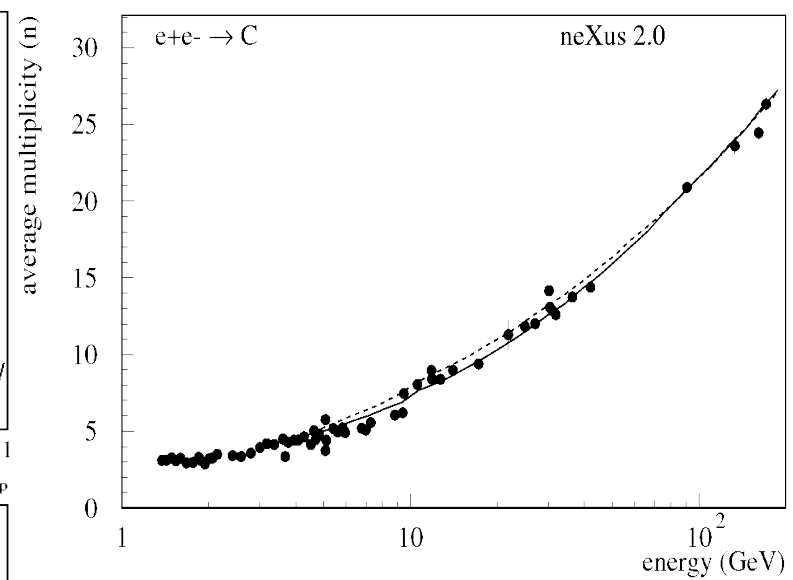
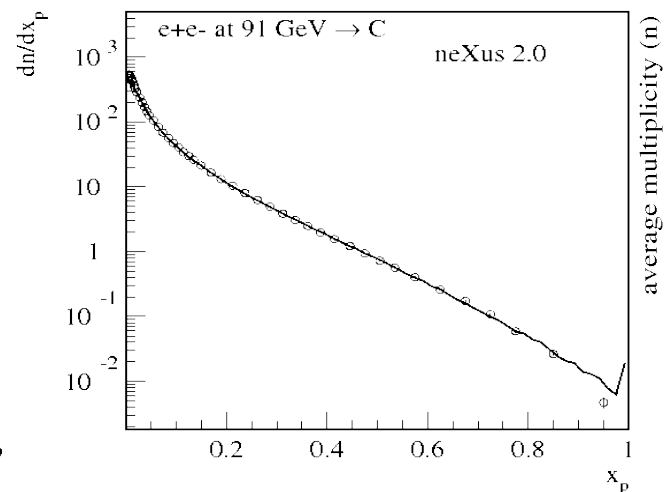
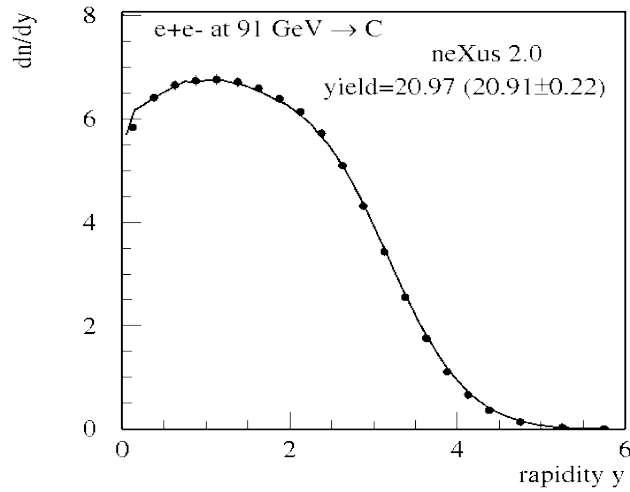
Rapidity of massless particles

$$y = \frac{1}{2} \ln \frac{1 + \cos \theta}{1 - \cos \theta} = -\ln \tan \frac{\theta}{2}$$

Pseudorapidity (all particles)

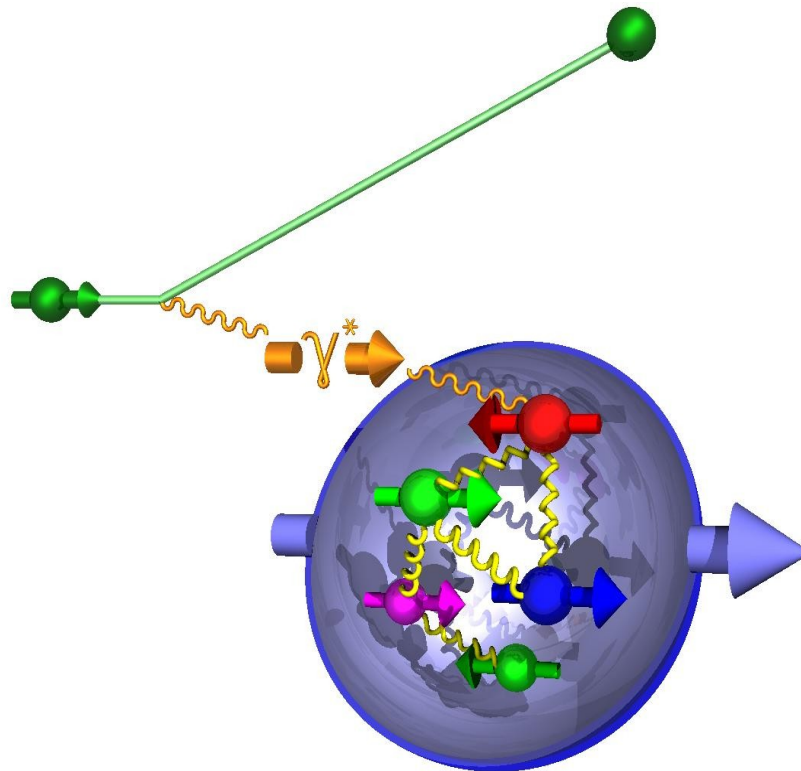
$$\eta = -\ln \tan \frac{\theta}{2}$$

Test at LEP

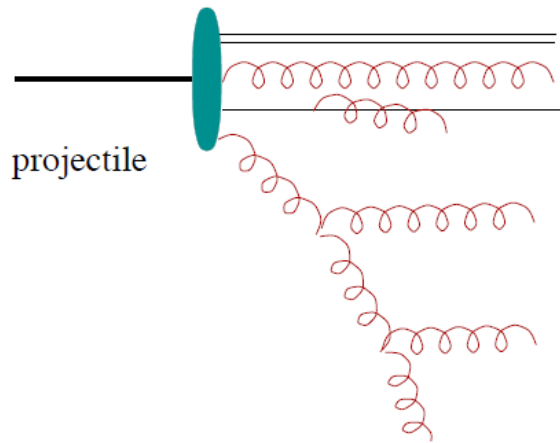


Parton Distributions

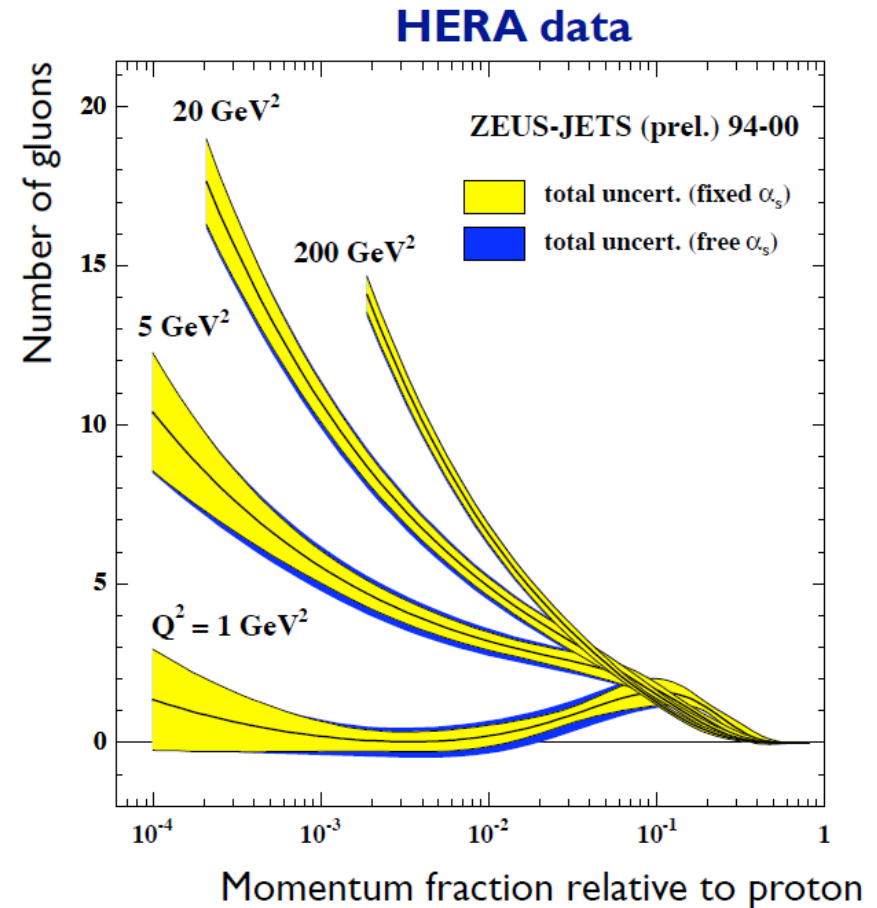
e⁺ vs e⁻ → l vs p (A) → p vs p → h vs A → A vs A



Perturbative QCD predictions for parton densities



Evolution of parton number
given by DGLAP equation
(and non-linear versions of it)



$$\frac{df_i(x, Q^2)}{d \log Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dy}{y} \sum_j f_j(y, Q^2) P_{j \rightarrow i} \left(\frac{x}{y} \right)$$

Prediction of
perturbative QCD

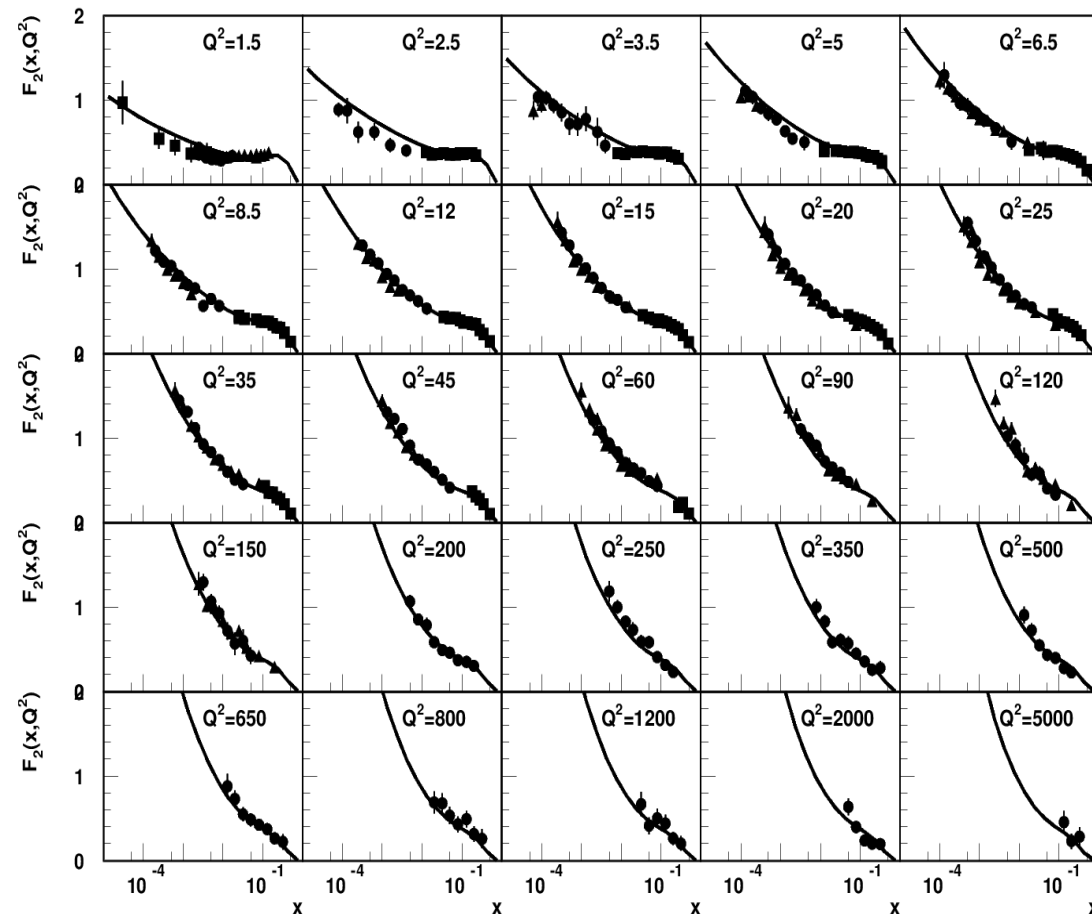
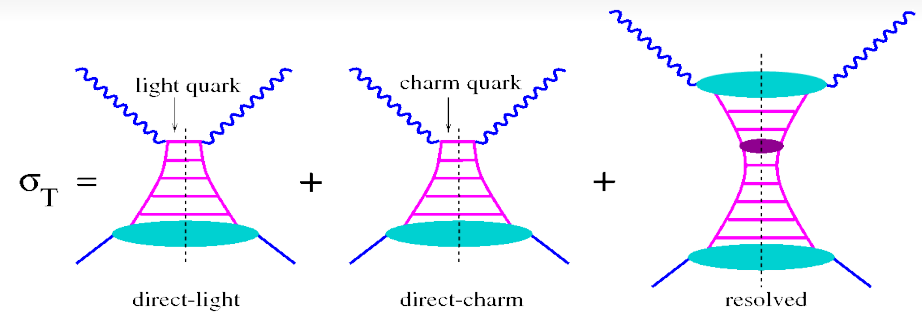
Test of parton interactions using deep inelastic scattering

Theory based parton distribution function (pdf)

- pQCD based for small x and large Q^2
- Regge “soft” parametrization at small x and low Q^2
- explicit fit of data only for valence quarks

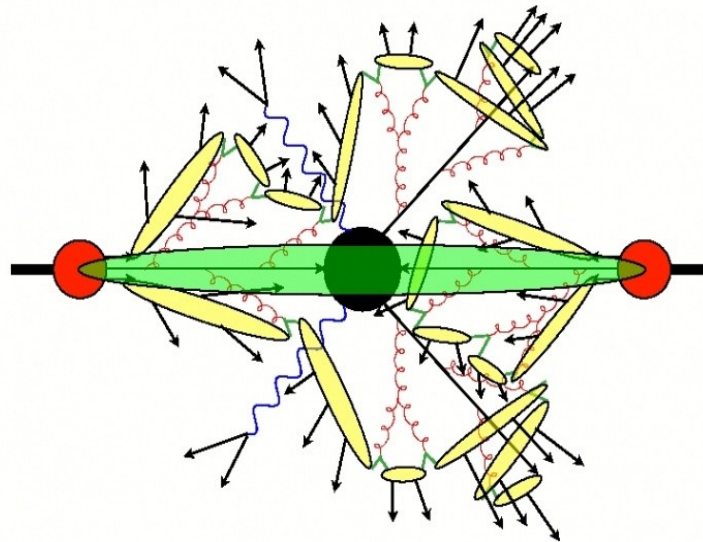
External pdf

- HEP models
- fit of data (including uncertainty) for low Q^2
- pQCD based evolution for large Q^2

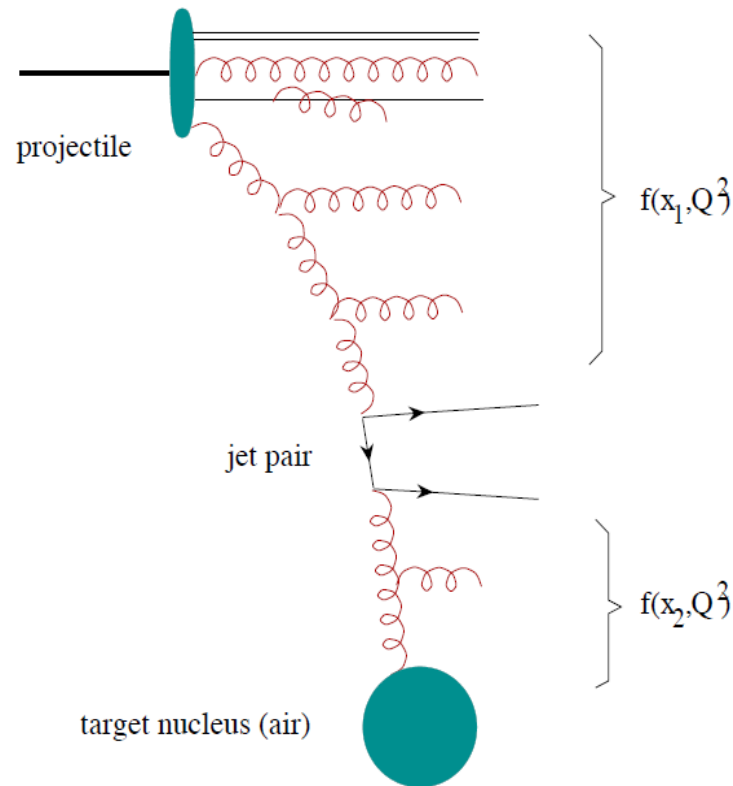


Parton Interactions

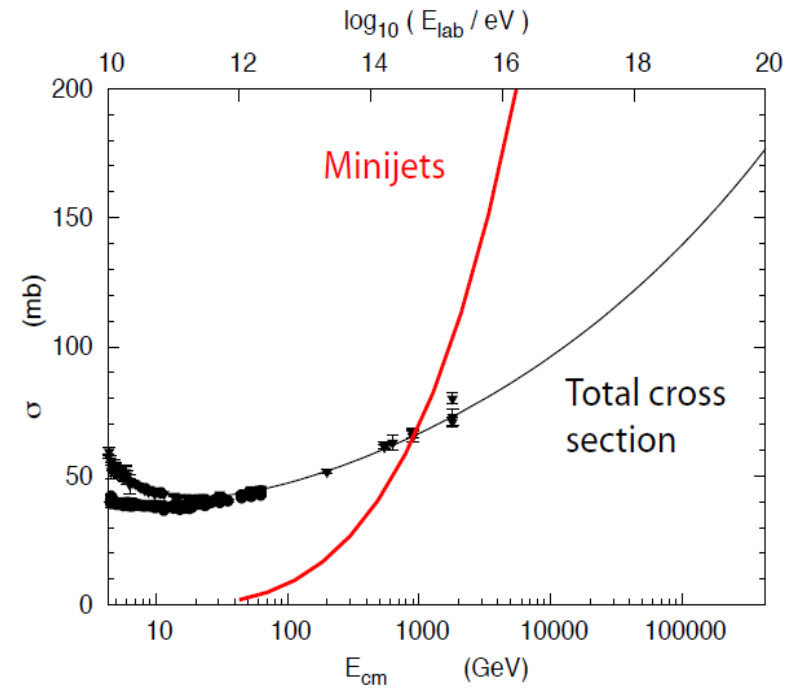
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QCD parton model: minijets



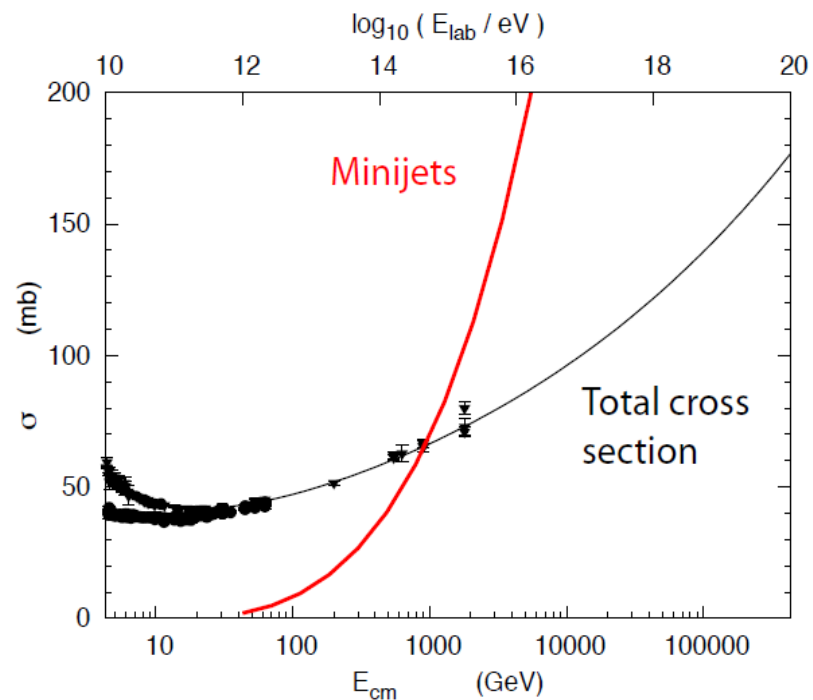
Proton-proton cross section



$$\sigma_{\text{QCD}} = \sum_{i,j,k,l} \frac{1}{1 + \delta_{kl}} \int dx_1 dx_2 \int_{p_{\perp}^{\text{cutoff}}} dp_{\perp}^2 f_i(x_1, Q^2) f_j(x_2, Q^2) \frac{d\sigma_{i,j \rightarrow k,l}}{dp_{\perp}}$$

Solution: Multiple parton-parton interactions (MPI)

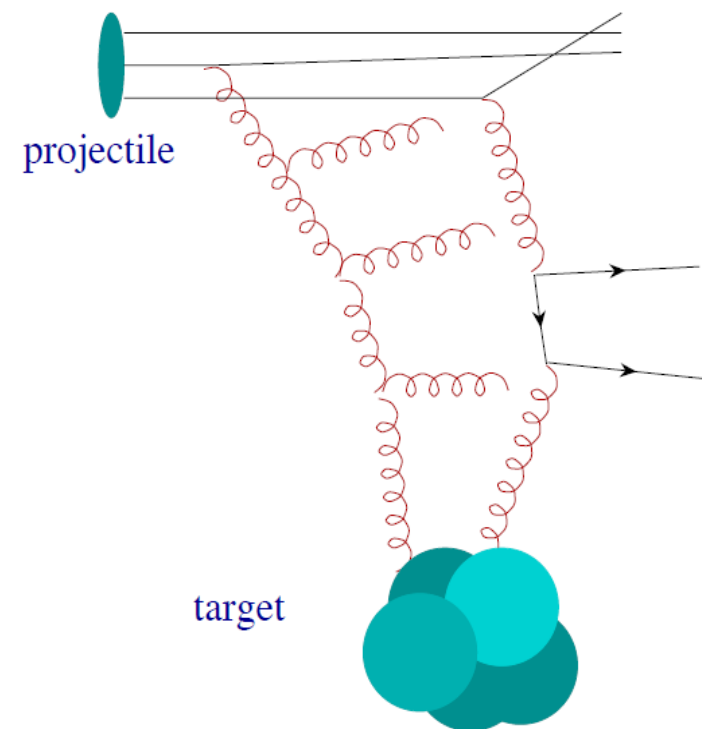
Proton-proton cross section



Average number
of minijet pairs

$$\langle n_{jet} \rangle = \frac{\sigma_{QCD}}{\sigma_{ine}}$$

QCD prediction:
inclusive cross section



MPI depends on model

● High Energy Physics models

- $\langle n_{\text{jet}} \rangle$ and cross-section (fit) are independent
- no soft multiple scattering
- no constrain from total cross-section to have independent access of inclusive class of events (Higgs, electroweak, etc ...)

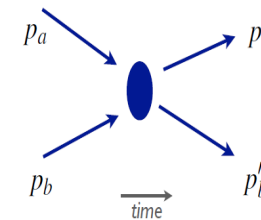
● Cosmic Ray models

- Gribov Regge Theory (GRT) used to compute total cross-section
- Parton model approach
 - fix σ_{hard} (pQCD) and σ_{tot} (data)
 - GRT using $\langle n_{\text{jet}} \rangle$ as final goal to reach
- or Pomeron approach
 - first built the Pomeron from soft and hard component
 - then add corrections to the bare amplitude to fit the total cross-section using GRT
 - $\langle n \rangle$ is a consequence of the Pomeron choice and the cross-section.

Pomeron Definition

$$A(s, t) = \eta(\alpha(t)) \beta(t) \left(\frac{s}{s_0} \right)^{\alpha(t)}$$

Amplitude of elementary parton-parton interaction=mini-jet

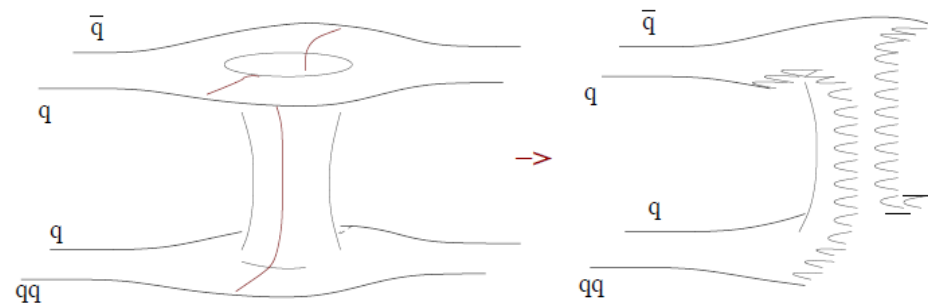


Lorentz-invariant description with Mandelstam variables

$$s = (p_a + p_b)^2$$

$$t = (p'_a - p_a)^2$$

- Quasi-particle that effectively accounts for all exchanged hadronic states
- Amplitude exhibits power-law dependence on energy
- Regge trajectory of Pomeron: exchanged particles might be glue balls ($\alpha(t)$)
- Pomeron trajectory only phenomenologically known
- Large N_c - n_f approximation of QCD: Pomeron corresponds to cylinder topology
- Final state configuration: leading contribution is two chains (strings) of hadrons
- Gluon-gluon scattering in pQCD corresponds to 'hard' contribution to Pomeron

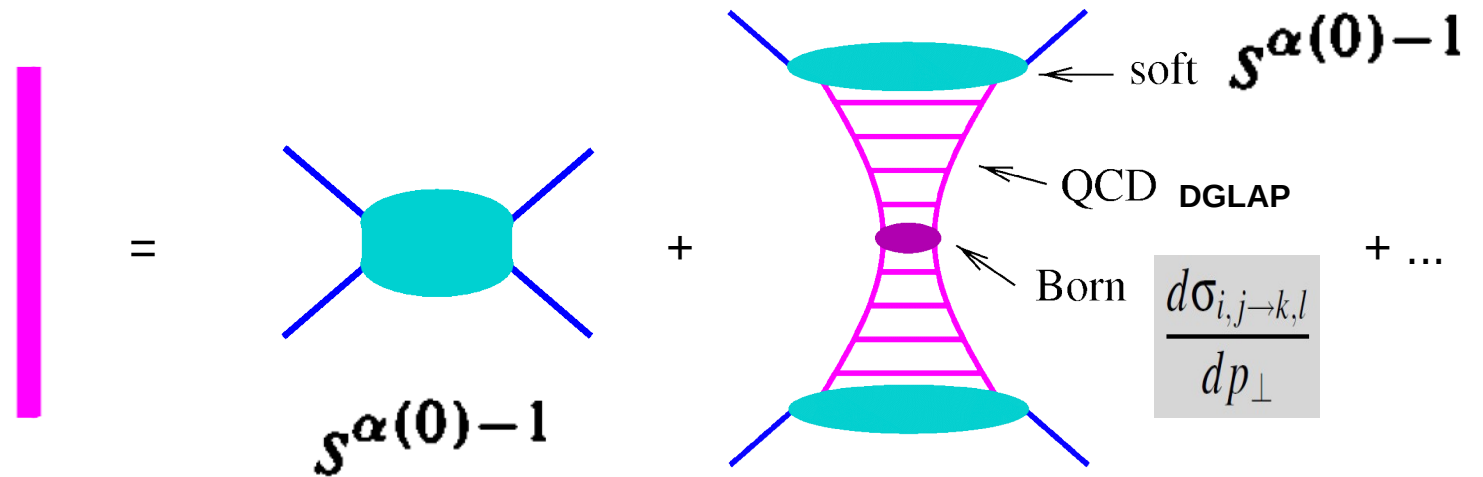


Multiple exchanges & interaction of quasi-particles (Pomerons):

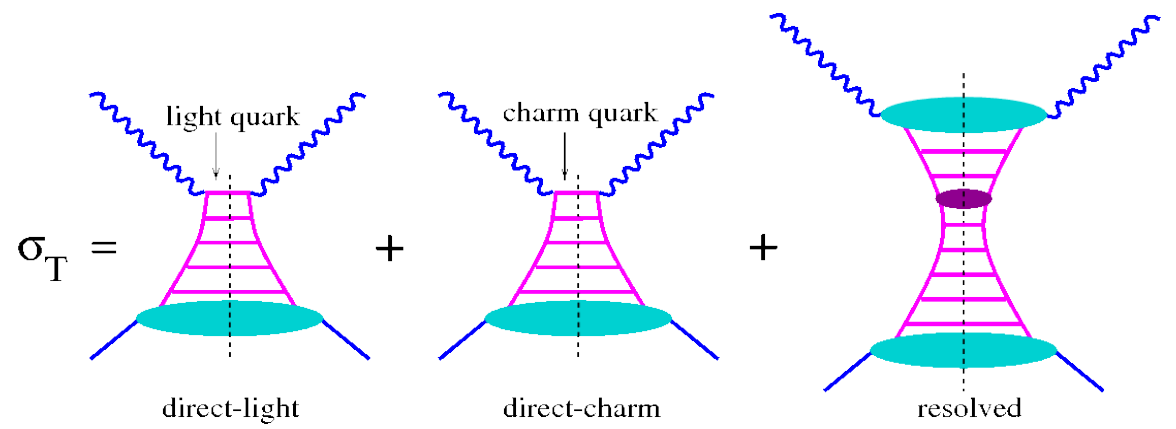
Gribov's Reggeon Field Theory

Pomeron in Models

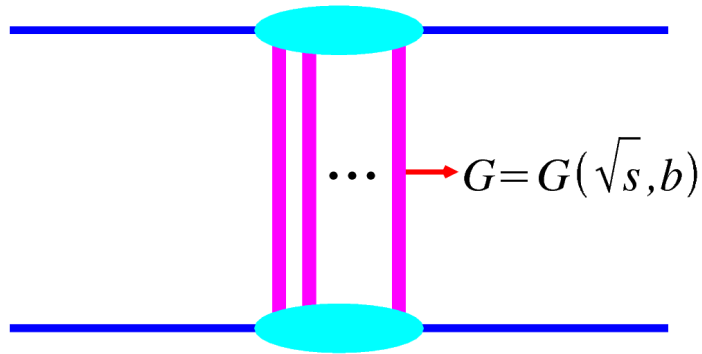
Semi-hard Pomeron :



Test of semi-hard Pomeron with Deep Inelastic Scattering: (Parton Distribution Function from HERA)



Gribov-Regge Based Models



Multiple elementary scattering

- ➔ Using Gribov-Regge (GR) : cross section from optical theorem :

$$\sigma_{ine}(\sqrt{s}) = \int d^2 b (1 - \exp(-G(\sqrt{s}, b)))$$

where $G(\text{energy}, \text{impact parameter}) = \text{Pomeron amplitude}$

- ➔ Probability for the number of elementary interactions (Pomeron) per event (Poisson)

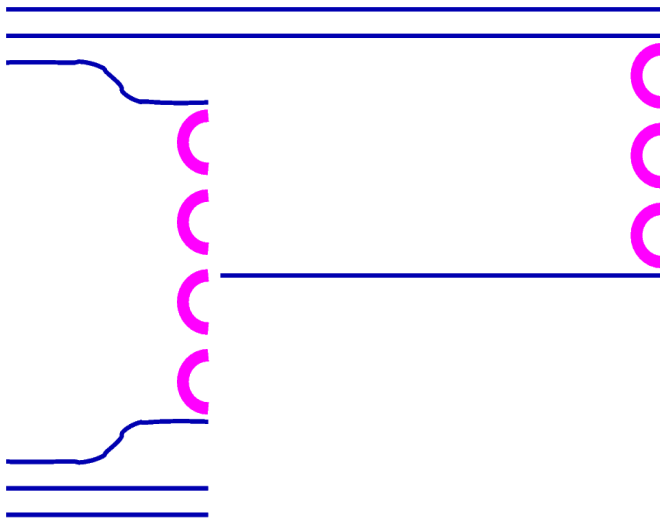
Successful description of hadronic cross-sections

But

Energy conservation NOT considered between the elementary interactions G

No possibility to deduce directly particle production !

Particle Production in GR based Models

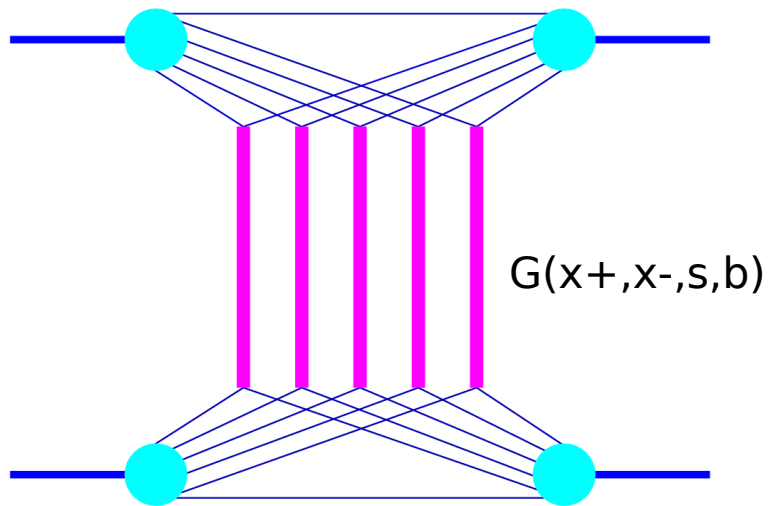


- **Number of strings from GR**
 - ➔ No energy conservation
- **Energy sharing**
 - ➔ Not consistent with cross-section
- **String fragmentation**
 - ➔ Proper energy conservation

**Link between cross-section and particle production
not consistent !**

Parton-Based Gribov-Regge Theory* (PBGRT) developed to solve the problem :
same formalisme for cross section and particle production
 used first in NEXUS and now in EPOS

Cross Section Calculation with Energy Conservation



- ➔ Gribov-Regge but with energy sharing at parton level
(Parton Based Gribov Regge Theory: H.J. Drescher et al., Phys. Rept. 350 (2001) 93)
- ➔ amplitude parameters fixed from QCD and pp cross section (semi-hard Pomeron)
- ➔ cross section calculation take into account interference term

$$\sigma_{\text{ine}}(s) = \int d^2b (1 - \Phi_{\text{pp}}(1, 1, s, b))$$

$$\begin{aligned} \Phi_{\text{pp}}(x^+, x^-, s, b) &= \sum_{l=0}^{\infty} \int dx_1^+ dx_1^- \dots dx_l^+ dx_l^- \left\{ \frac{1}{l!} \prod_{\lambda=1}^l -G(x_\lambda^+, x_\lambda^-, s, b) \right\} \\ &\times F_{\text{proj}}\left(x^+ - \sum x_\lambda^+\right) F_{\text{targ}}\left(x^- - \sum x_\lambda^-\right). \end{aligned}$$

Particle Production in EPOS

m number of exchanged elementary interaction per event fixed from elastic amplitude taking into account energy sharing :

➔ m cut Pomerons from :

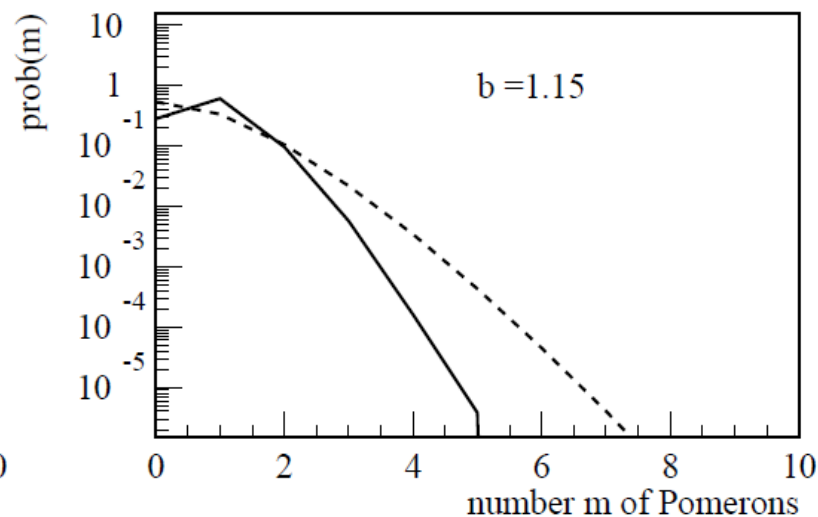
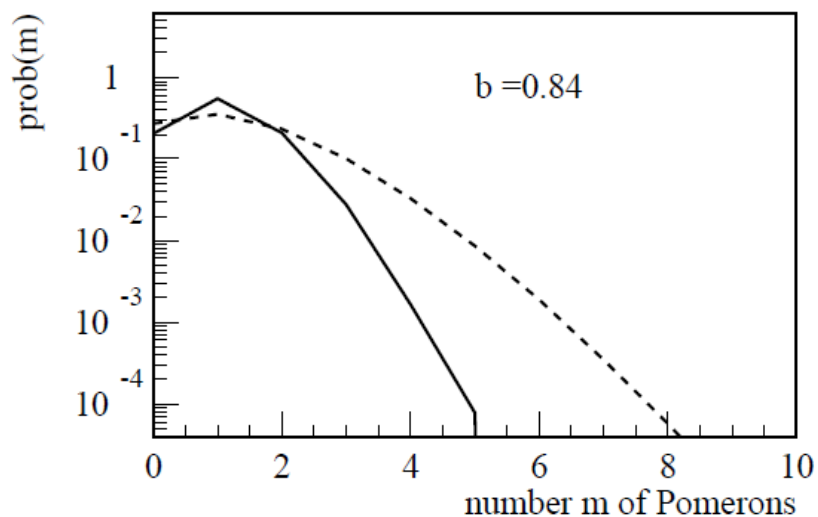
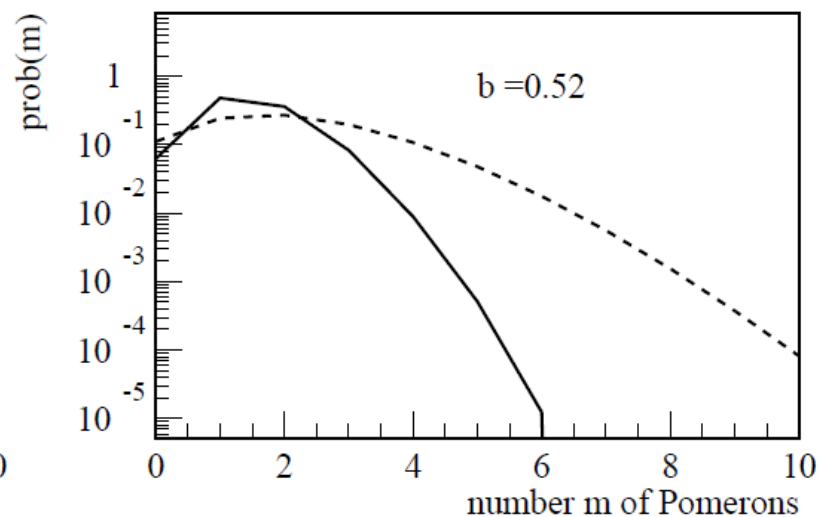
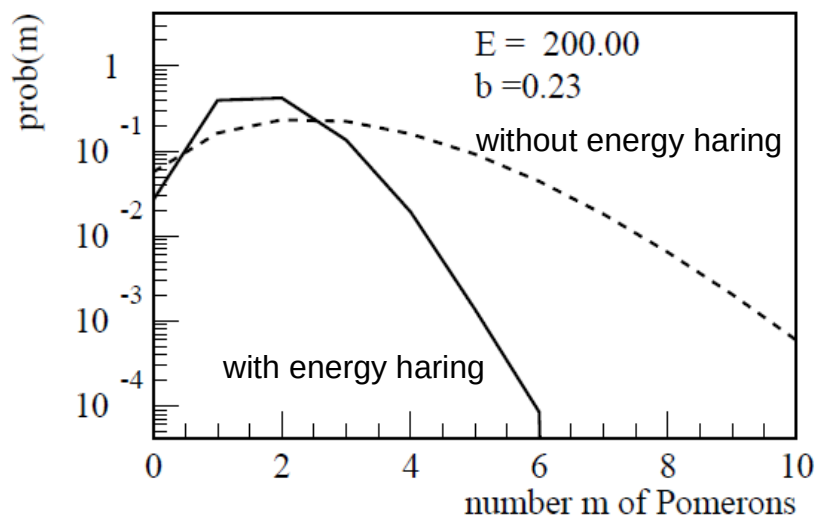
$$\Omega_{AB}^{(s,b)}(m, X^+, X^-) = \prod_{k=1}^{AB} \left\{ \frac{1}{m_k!} \prod_{\mu=1}^{m_k} G(x_{k,\mu}^+, x_{k,\mu}^-, s, b_k) \right\} \Phi_{AB}(x^{\text{proj}}, x^{\text{targ}}, s, b)$$

- m and X fixed together by a complex Metropolis (Markov chain)
- ➔ 2m strings formed from the m elementary interactions
- **energy conservation** : energy fraction of the 2m strings given by X
- ➔ consistent scheme : energy sharing reduce the probability to have large m

Consistent treatment of cross section and particle production:
number AND distribution of cut Pomerons depend on cross section

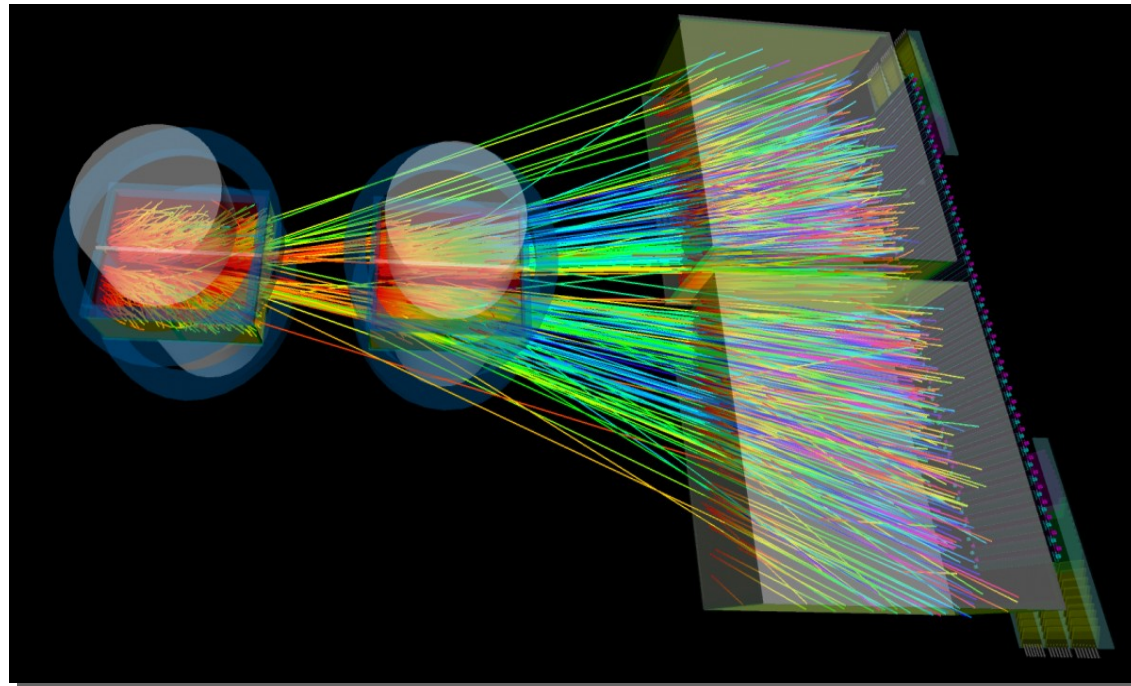
Number of cut Pomerons

Fluctuations reduced by energy sharing (mean can be changed by parameters)

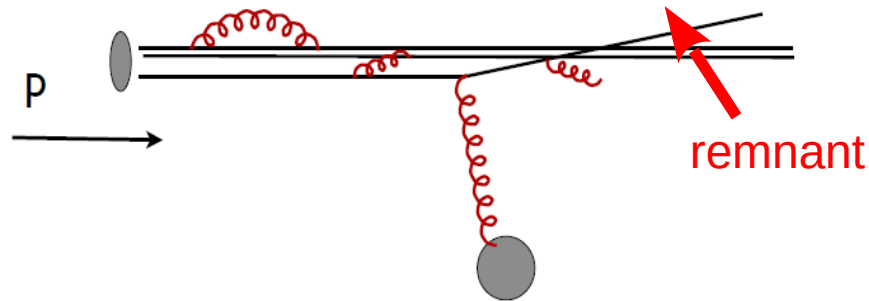


Hadron Interactions

$e^+ \text{ vs } e^-$ \longrightarrow $l \text{ vs } p (A)$ \longrightarrow $p \text{ vs } p$ \longrightarrow $h \text{ vs } A$ \longrightarrow $A \text{ vs } A$



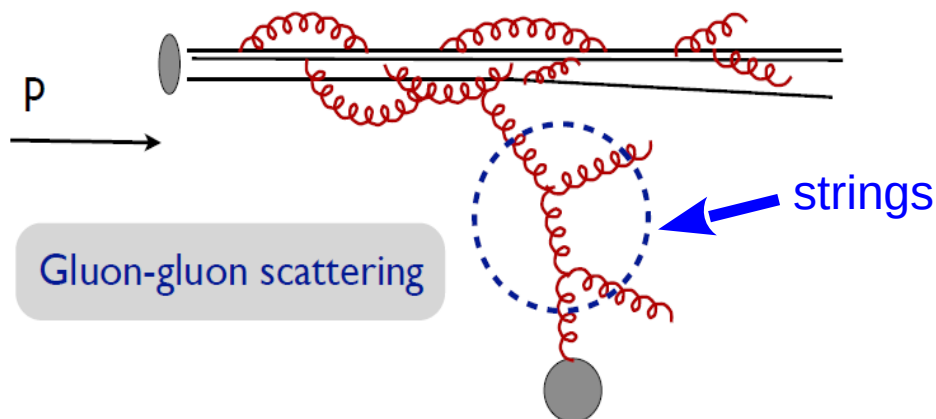
Transition from intermediate to high energy



Intermediate energy:

- $E_{\text{lab}} < 1,500 \text{ GeV}$
- $E_{\text{cm}} < 50 \text{ GeV}$
- dominated by valence quarks

Increase of effective nucleon radius due to small x partons

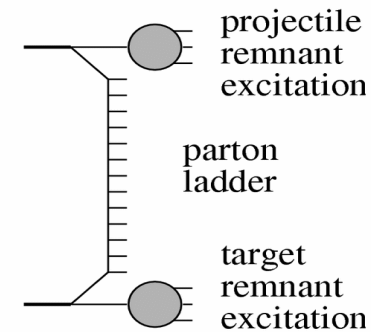
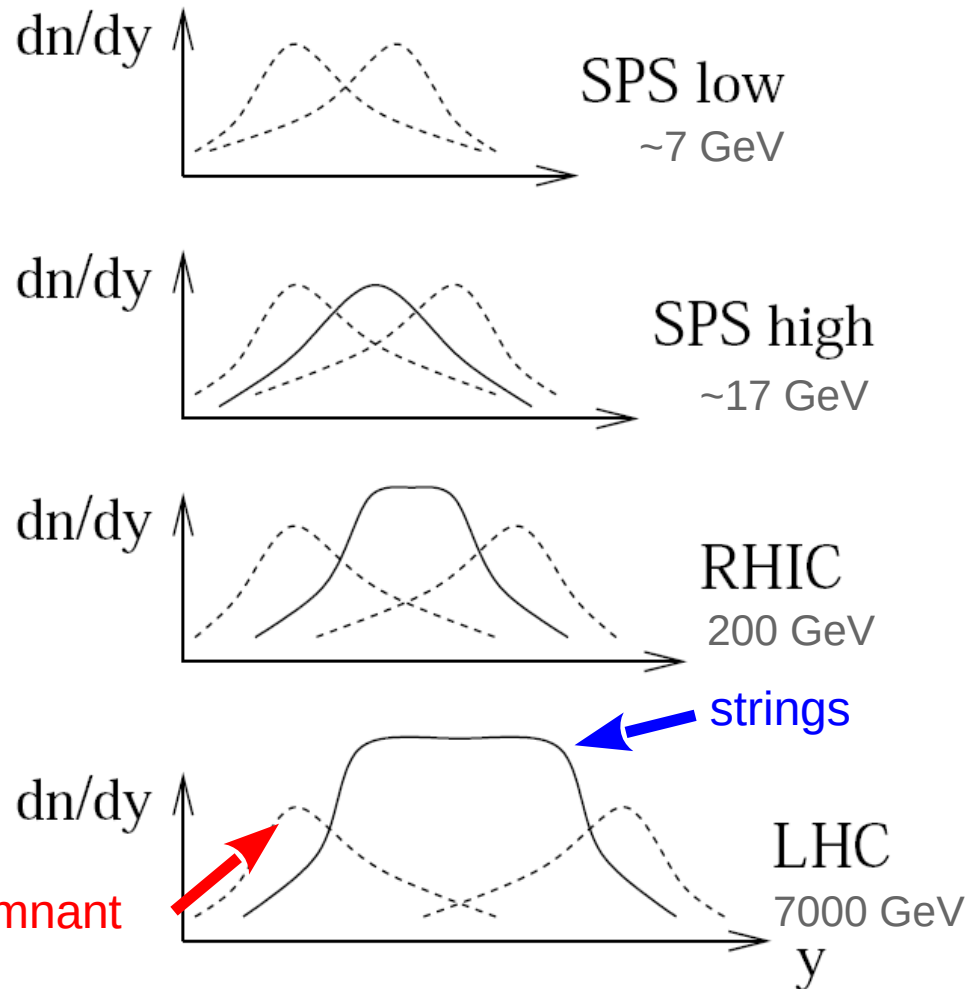


High energy regime:

- $E_{\text{lab}} > 21,000 \text{ GeV}$
- $E_{\text{cm}} > 200 \text{ GeV}$
- dominated by gluons and sea quarks

Remnants

Forward particles mainly from projectile remnant



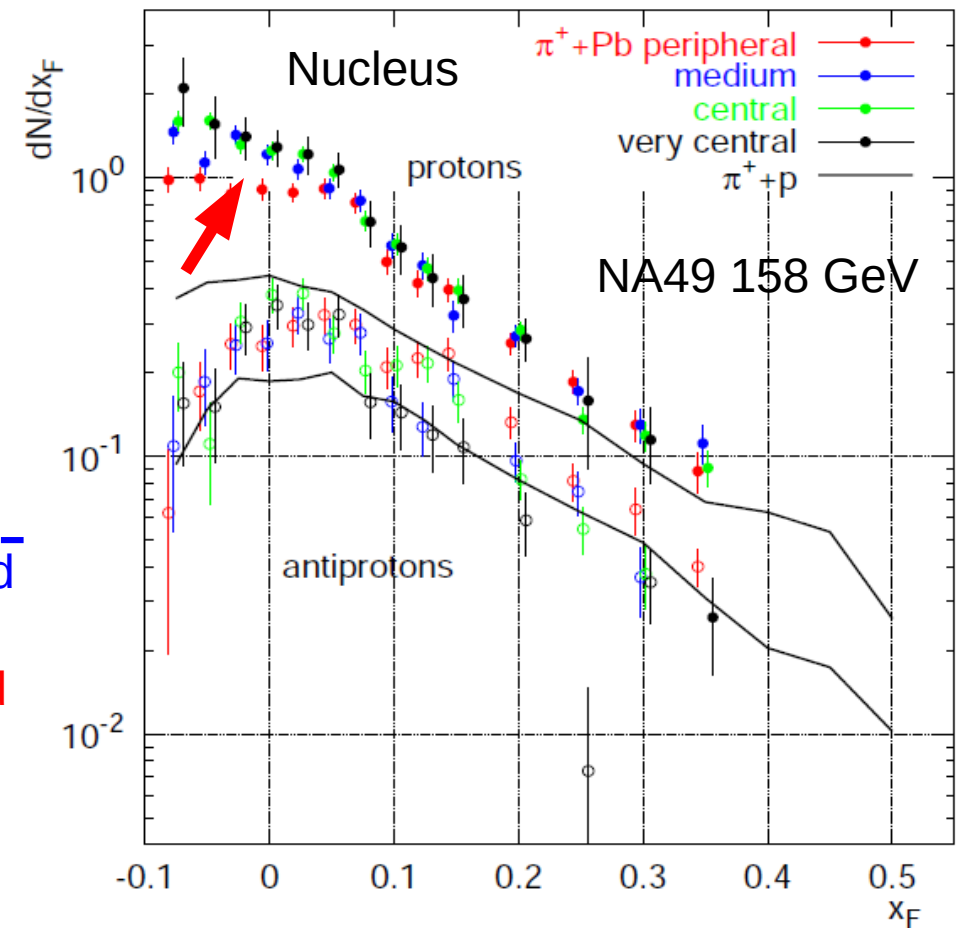
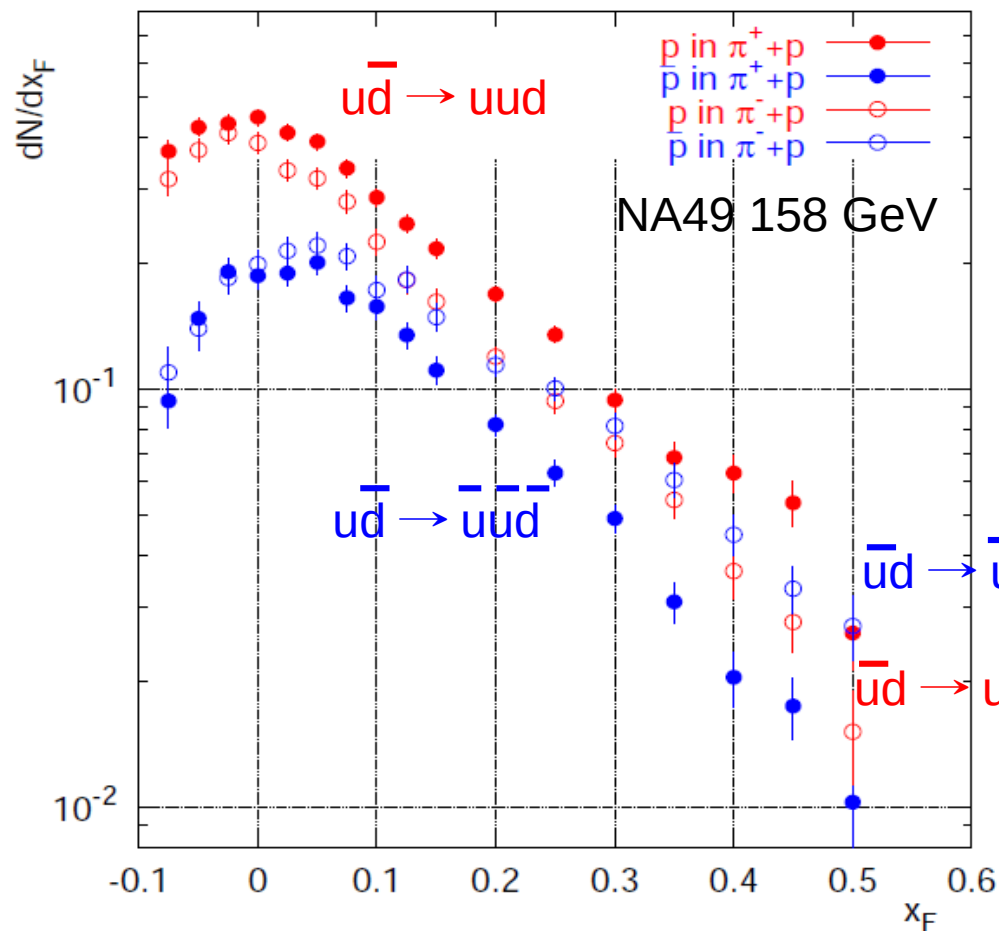
- ➔ At very low energy only particles from remnants
- ➔ At low energy (fixed target experiments) (SPS) strong mixing
- ➔ At intermediate energy (RHIC) mainly string contribution at mid-rapidity with tail of remnants.
- ➔ At high energy (LHC) only strings at mid-rapidity (baryon free)

Different contributions of particle production at different energies or rapidities

Leading Particle Effect

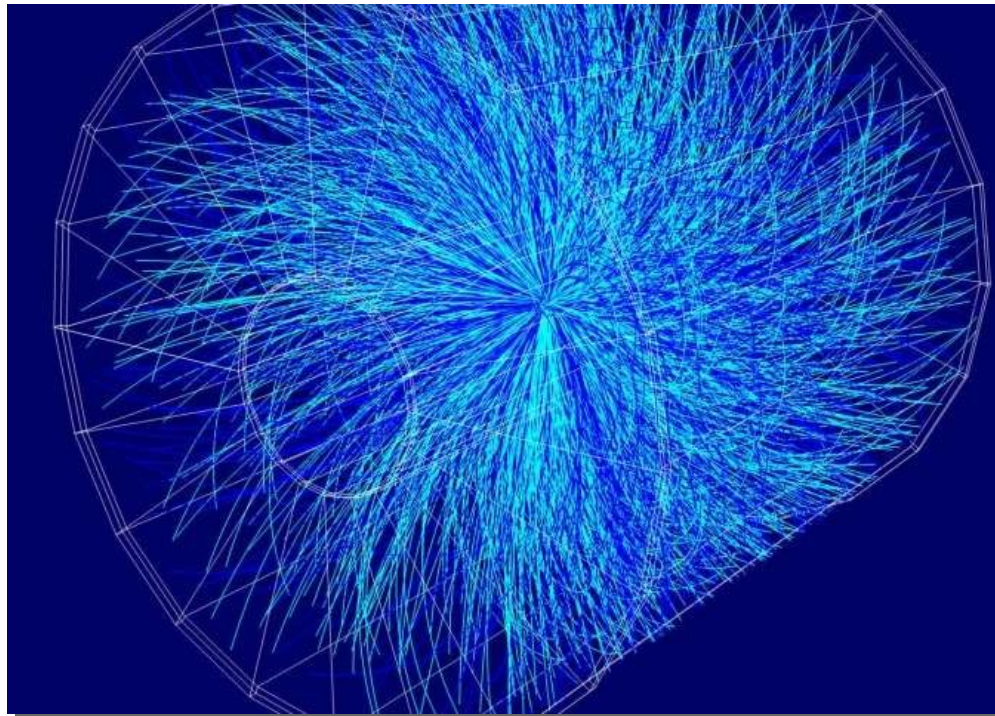
Remnant (leading particle) effect

- ➔ Different forward production depending on the projectile
- ➔ Low energy : even the target hadronization can change “forward” distributions

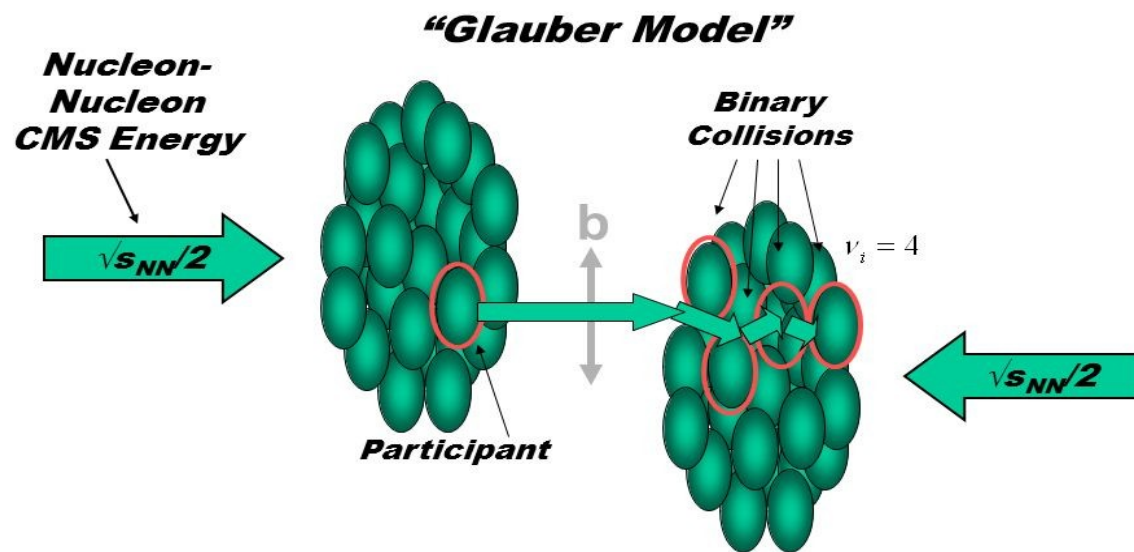


Nucleus Interactions

e⁺ vs e⁻ → l vs p (A) → p vs p → h vs A → A vs A



Nucleus Interactions

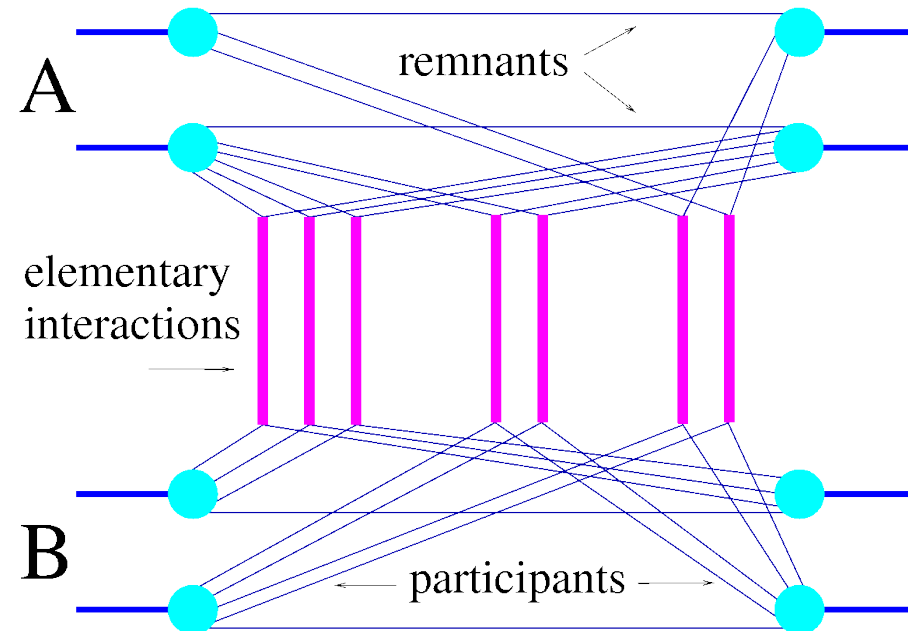


● Gribov Regge extension to nuclei

- ➔ coherent approach
- ➔ energy conservation can be taken into account
- ➔ good for cross-section and multiplicity
- ➔ special care for hard processes

● Glauber model

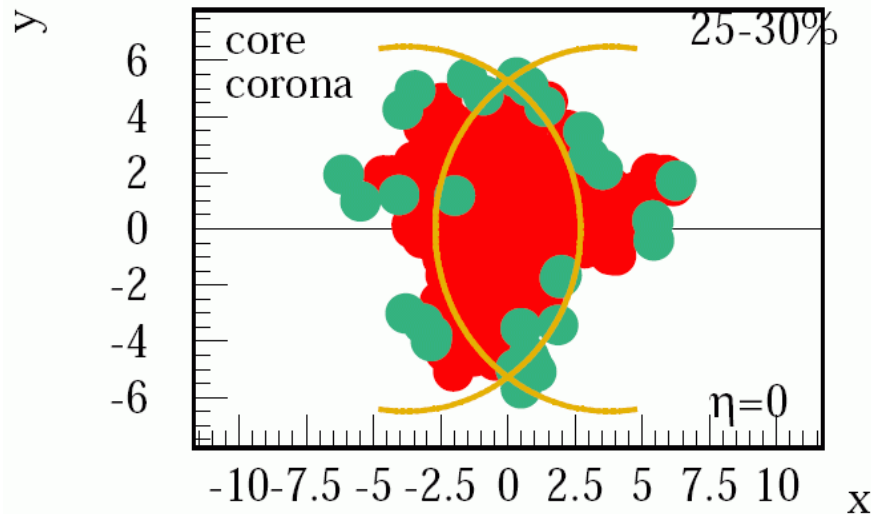
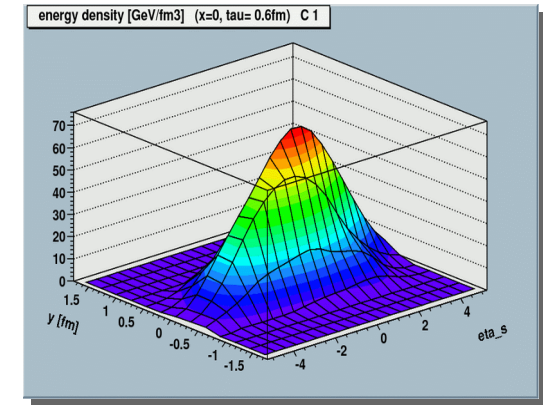
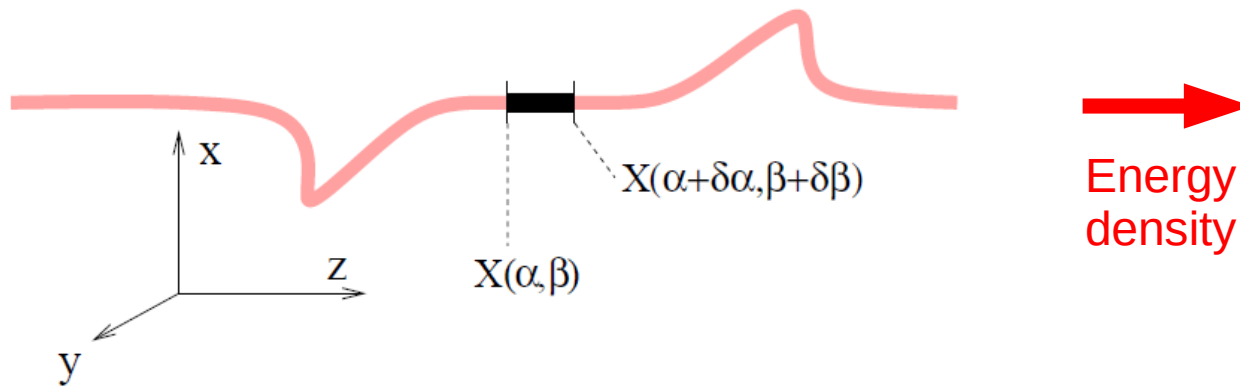
- ➔ Simple non-coherent approach
- ➔ all collisions taken independently
- ➔ works well for rare hard processes which will happen only once per participant
- ➔ cross-section calculation with some corrections



High Density Core Formation

● Heavy ion (HI) collisions :

- ➔ the usual procedure has to be modified, since the density of strings will be so high that they cannot possibly decay independently : **core**



- ➔ Each string splitted into a sequence of string segments, corresponding to widths $\delta\alpha$ and $\delta\beta$ in the string parameter space
- ➔ If energy density from segments high enough
 - ◆ segments fused into core
 - ➔ hydrodynamical-evolution
 - ➔ statistical hadronization
- ➔ If low density (corona)
 - ◆ segments remain hadrons

Core Hadronization

● 2 types of hadronization

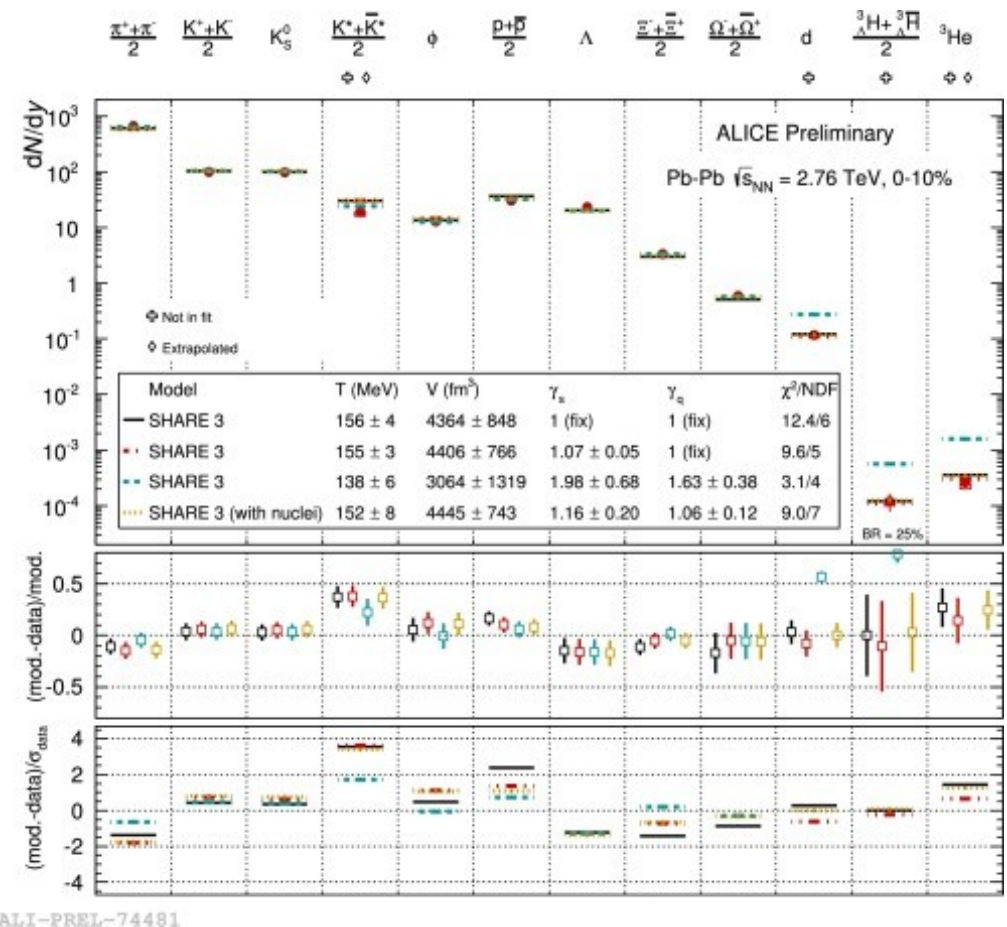
- ➔ Corona (low density = pp ?) : standard string hadronization (in vacuum)
- ➔ Core (high density = HI ?) : collective (thermal) hadronization (in medium)

● Thermal hadronization

- ➔ Good description of all particle yields in a central heavy ion collision with 2 parameters (temperature and chemical potential)
- ➔ Apply to extended source and based on conservation laws only.

● Energy/system/centrality evolution fixed by core/corona ratio !

- ➔ No need to change hadronization of a given system (color reconnection ?)



Collective Hadronization

- ➔ One decade of RHIC experiments (heavy ion, pp, and dAu scattering, up to 200 GeV)

heavy ion collisions produce matter which expands as an almost ideal fluid

- ➔ mainly because azimuthal anisotropies (particle correlations) can be explained on the basis of ideal hydrodynamics (mass splitting, ridge, etc ...)

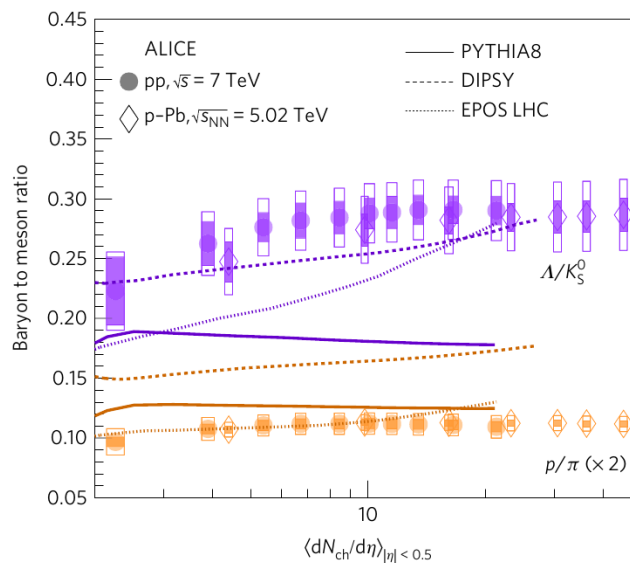
LHC pp results: first signs for collective behavior as well ... ?

Global Approach

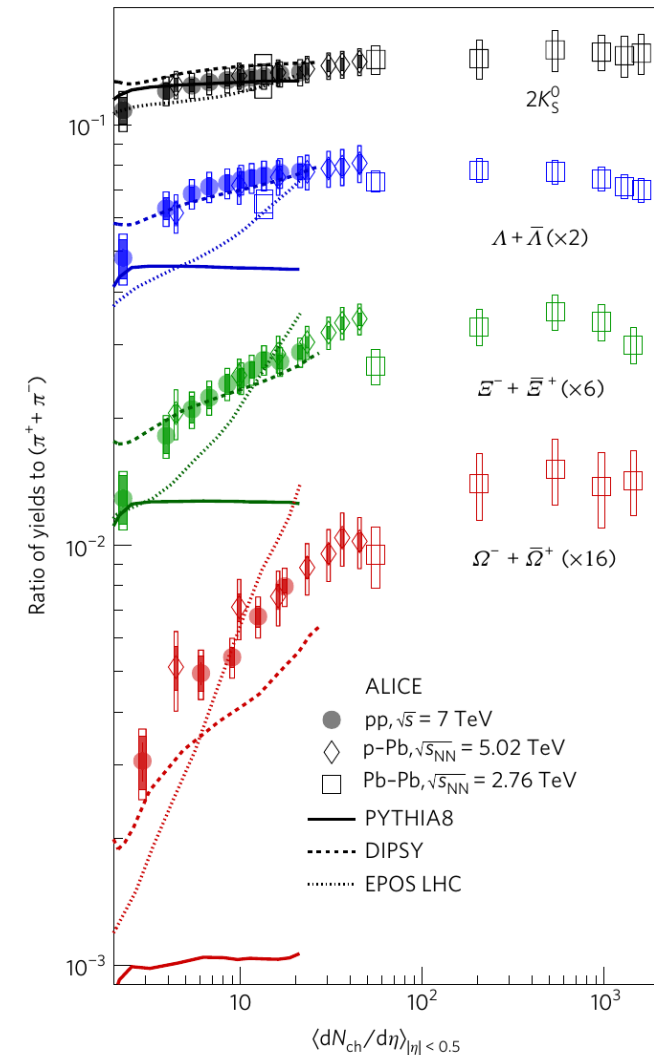
Moving from projectile/target to energy density (particle number)

- ➔ New point of view on particle production
- ➔ Continuity in particle rasion evolution between pp, pPb and PbPb
- ➔ No model can reproduce every thing

- ➔ PYTHIA : string only
- ➔ DIPSY : overlapping strings with modified parameters
- ➔ EPOS LHC : core/corona



Nature Letters DOI: 10.1038/NPHYS4111



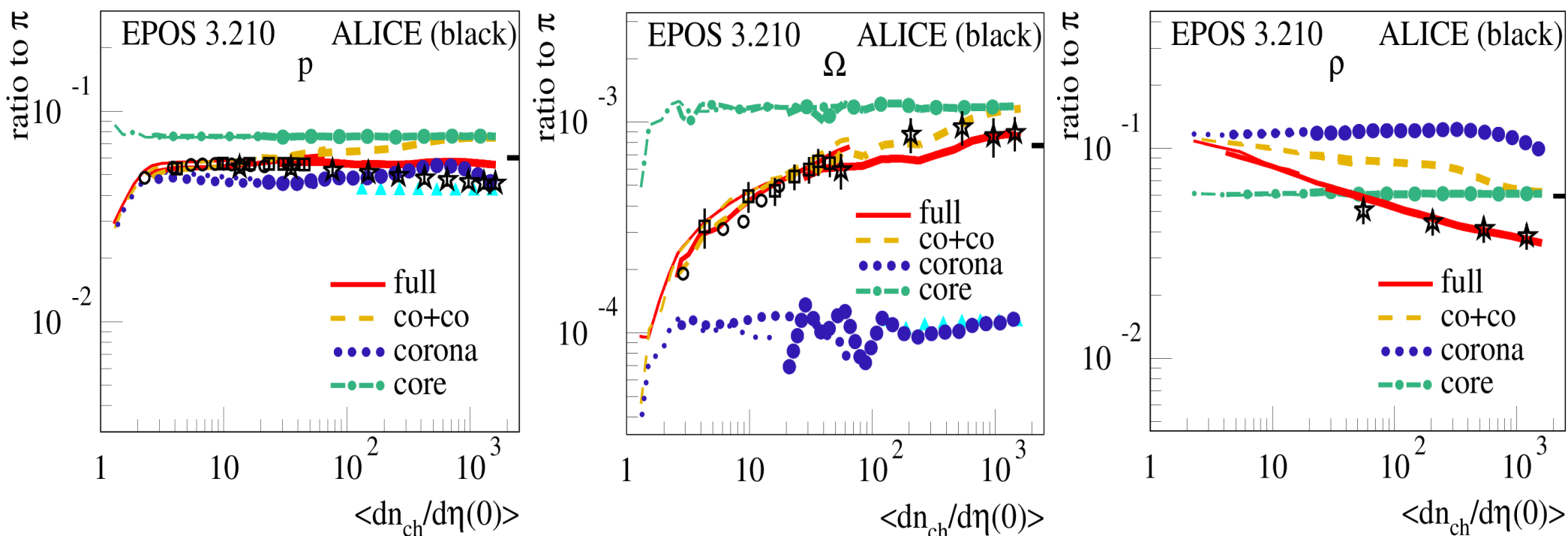
Global Approach : Core+Corona

Moving from projectile/target to energy density (particle number)

➔ New point of view on particle production

➔ Good description achieved with core/corona + hadron reinteractions (hadron gas)

➔ Same physics for all systems but different mixtures of the 2 components (low + high energy density hadronization)



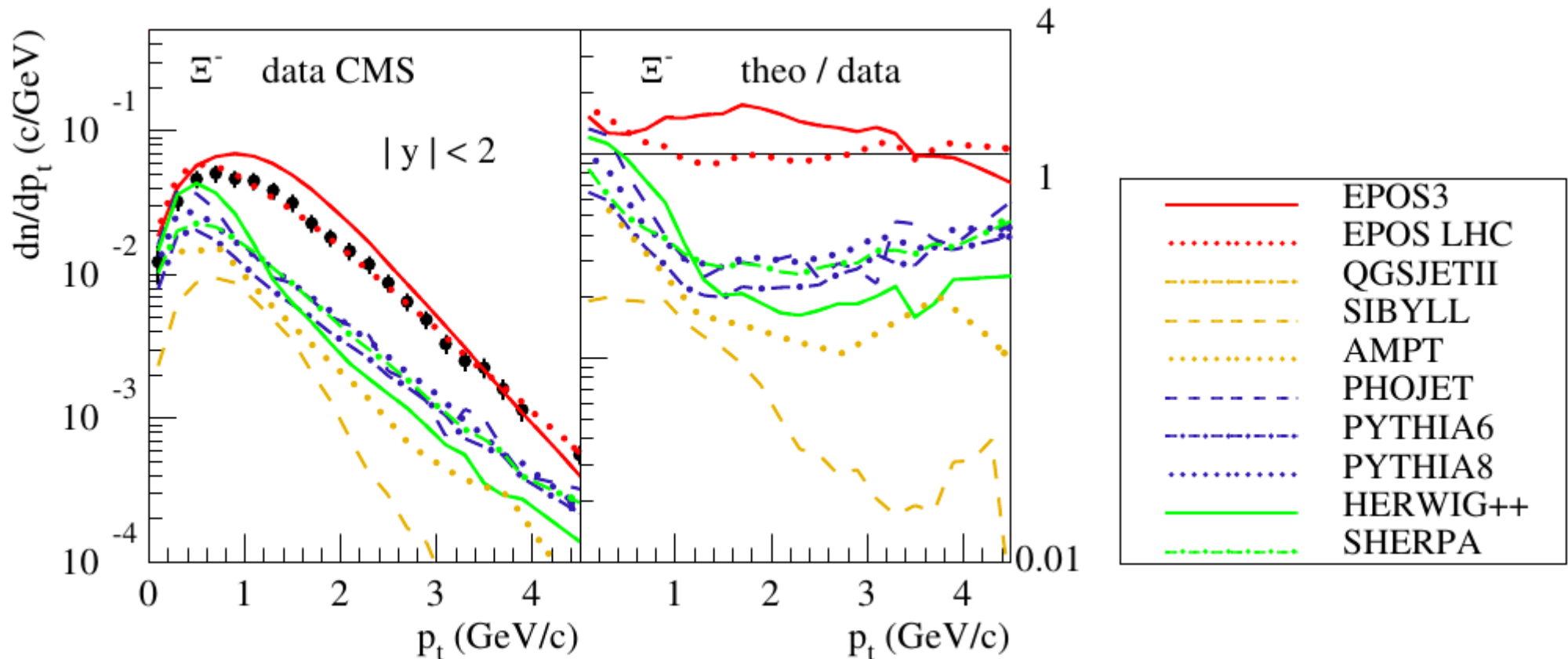
Global Approach : EPOS

Detailed description can be achieved

➔ identified spectra

➔ p_t behavior driven by collective effects (statistical hadronization + flow)

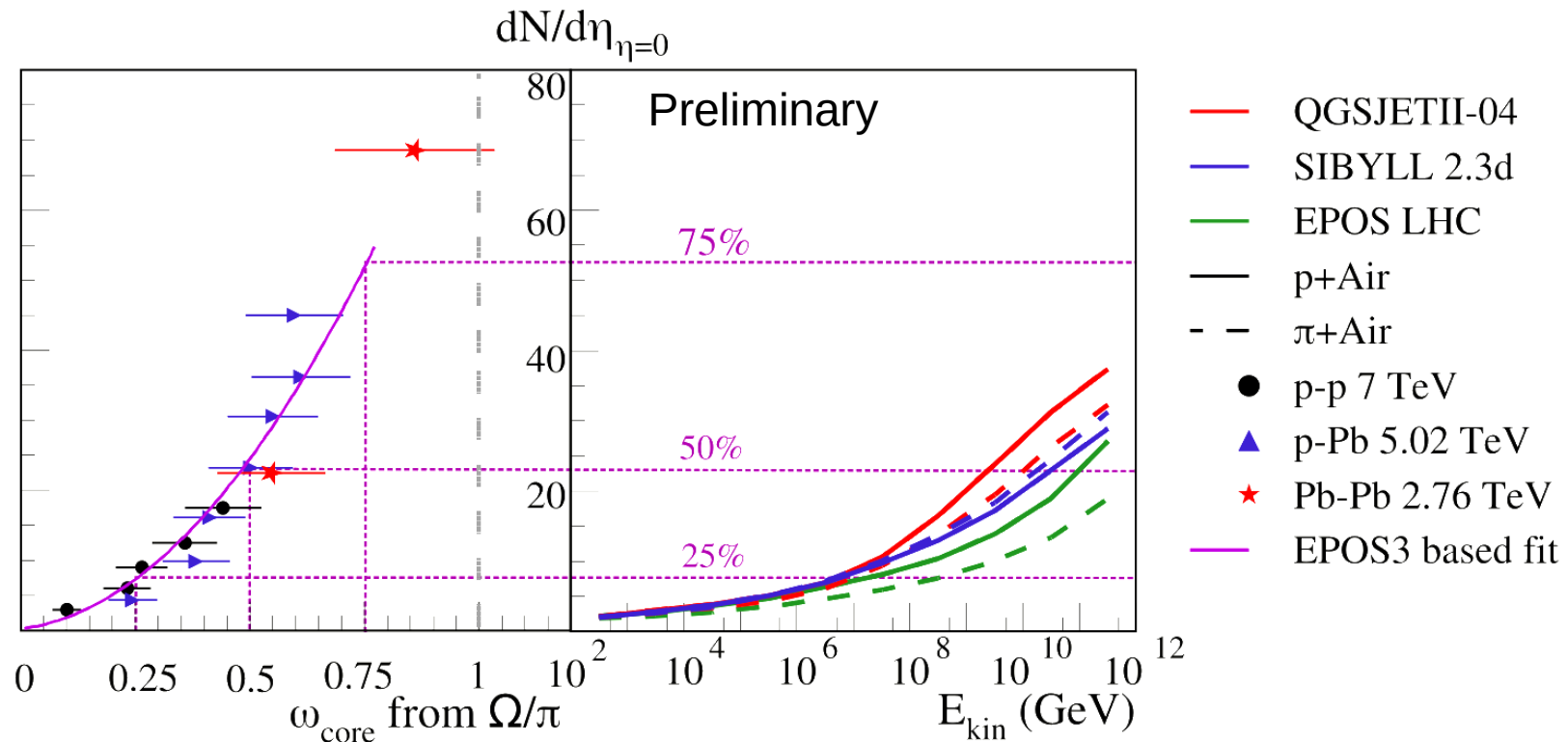
➔ large effect for multi-strange baryons (yield AND $\langle p_t \rangle$)



Particle Densities in Air Showers

Is particle density in air shower high enough to expect core formation ?

- ➔ Core formation start quite early according to ALICE data
- ➔ Cosmic ray primary interaction likely to have 50% core at mid-rapidity ... but forward can be different !



Core-Corona approach and EAS

To test if a QGP like hadronization can account for the missing muon production in EAS simulations a core-corona approach can be artificially apply to any model

- ➔ Particle ratios from statistical model are known (tuned to PbPb) and fixed : core
- ➔ Initial particle ratios given by individual hadronic interaction models : corona
- ➔ Using CONEX, EAS can be simulated mixing corona hadronization with an arbitrary fraction ω_{core} of core hadronization:

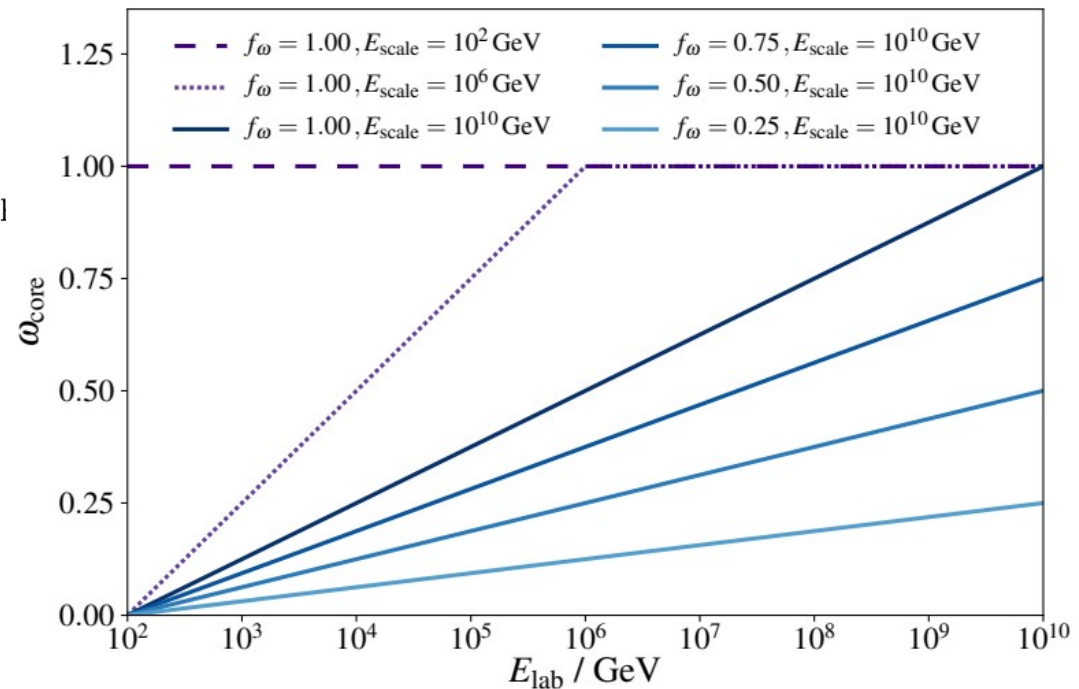
$$N_i = \omega_{\text{core}} N_i^{\text{core}} + (1 - \omega_{\text{core}}) N_i^{\text{corona}}$$

$$\omega_{\text{core}}(E_{\text{lab}}) = f_{\omega} \underbrace{F(E_{\text{lab}}; E_{\text{th}}, E_{\text{scale}})}_{\frac{\log_{10}(E_{\text{lab}}/E_{\text{th}})}{\log_{10}(E_{\text{scale}}/E_{\text{th}})} \text{ for } E_{\text{lab}} > E_{\text{th}}}$$

$$E_{\text{th}} = 100 \text{ GeV}$$

Different scenarii can be studied playing with f_{ω} and E_{scale} .

Note : the leading particle is NOT modified (projectile remnant)

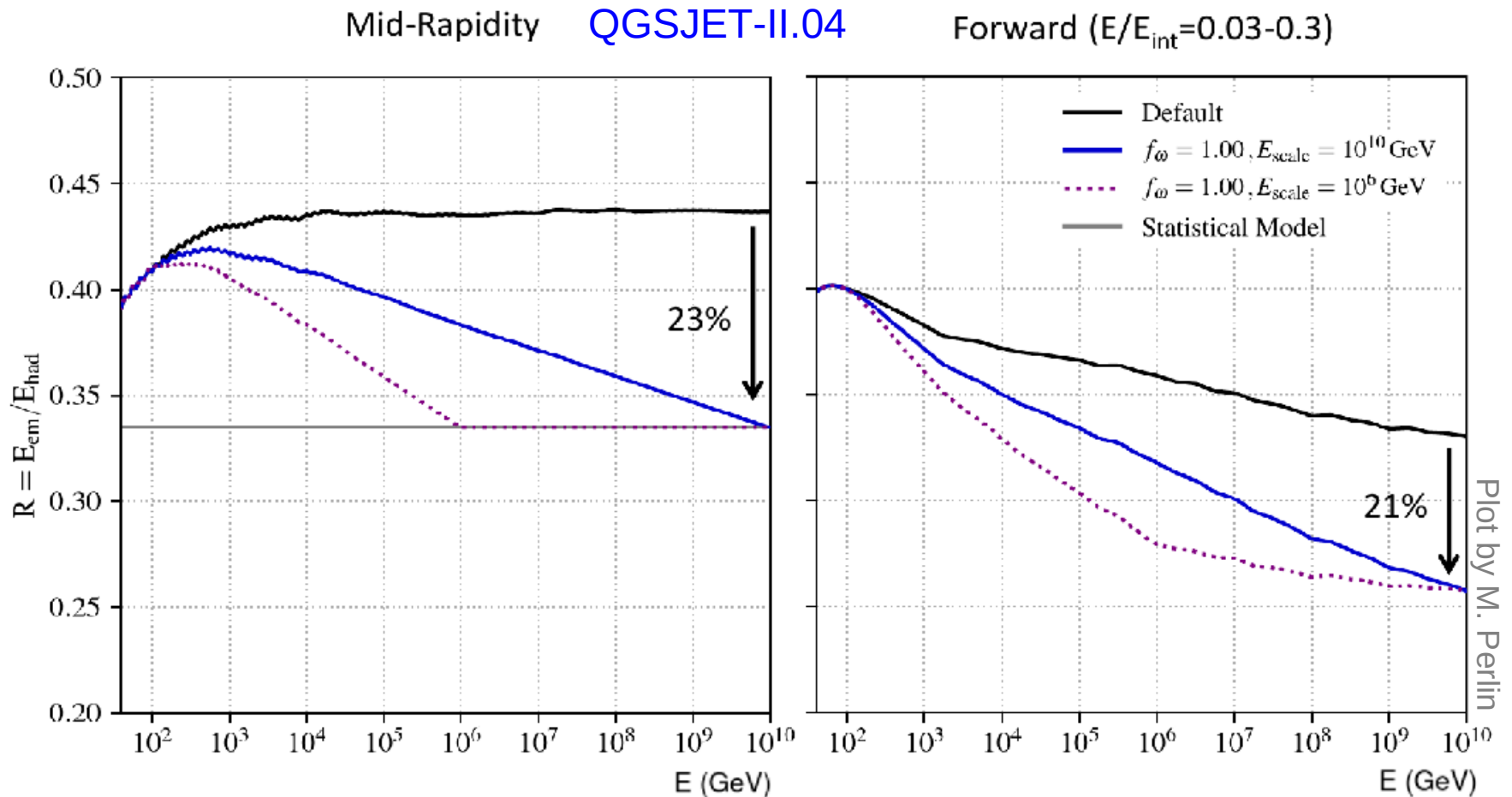


Evolution of hadronization from core to corona

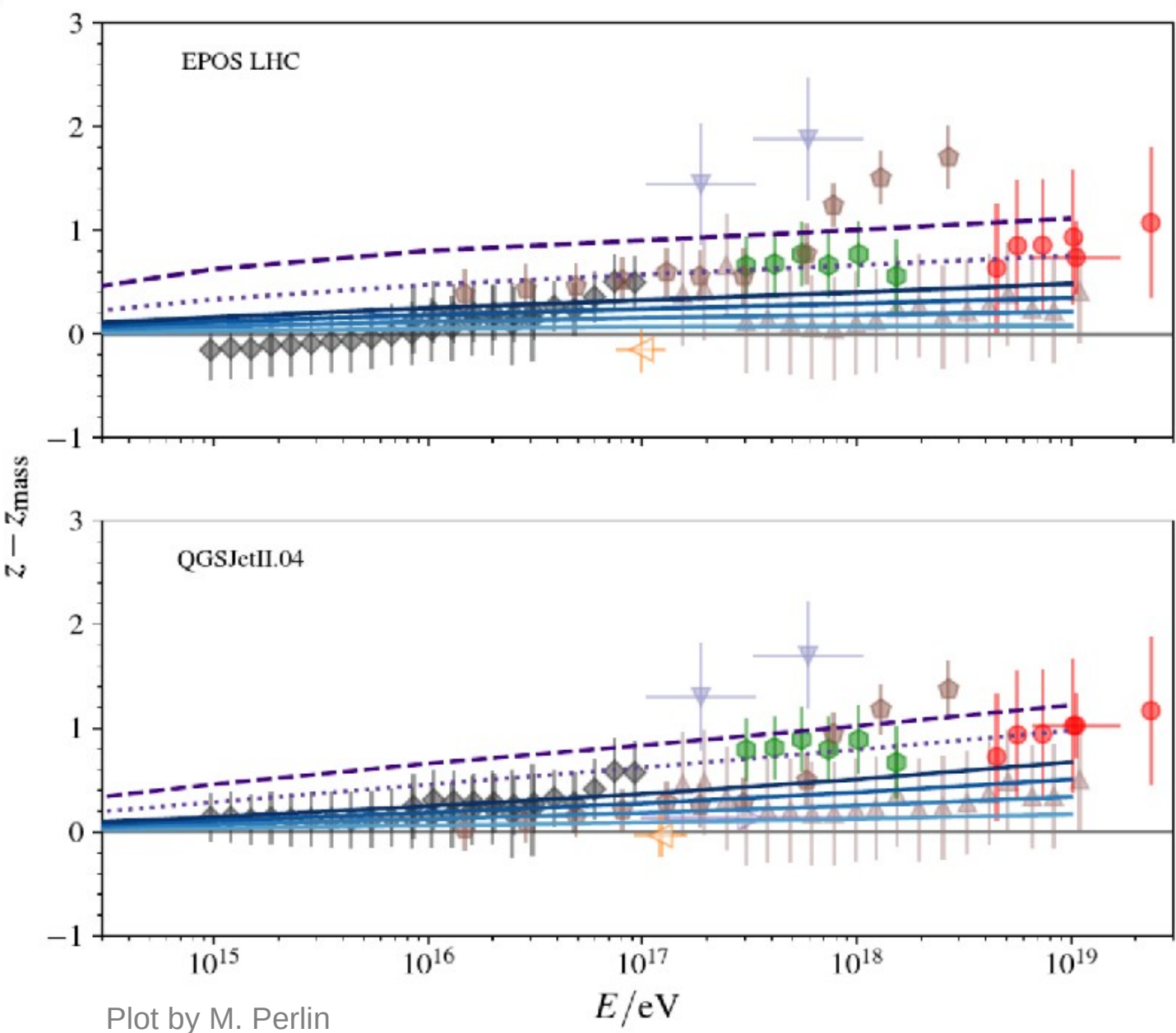
The relative fraction of π^0 depends on the hadronization scheme

→ Change of ω_{core} with energy changes $c = \frac{N_{\pi^0}}{N_{\text{mult}}}$ or $R(\eta) = \frac{\langle dE_{\text{em}}/d\eta \rangle}{\langle dE_{\text{had}}/d\eta \rangle}$

which define the muon production in air showers.



Results for z-scale



$$z = \frac{\ln N_{\mu}^{\text{det}} - \ln N_{\mu,p}^{\text{det}}}{\ln N_{\mu,\text{Fe}}^{\text{det}} - \ln N_{\mu,p}^{\text{det}}}$$

- $f_{\omega} = 1.00, E_{\text{scale}} = 10^2 \text{ GeV}$
- ⋯ $f_{\omega} = 1.00, E_{\text{scale}} = 10^6 \text{ GeV}$
- $f_{\omega} = 1.00, E_{\text{scale}} = 10^{10} \text{ GeV}$
- $f_{\omega} = 0.75, E_{\text{scale}} = 10^{10} \text{ GeV}$
- $f_{\omega} = 0.50, E_{\text{scale}} = 10^{10} \text{ GeV}$
- $f_{\omega} = 0.25, E_{\text{scale}} = 10^{10} \text{ GeV}$
- $f_{\omega} = 0$ (Default model)

- Pierre Auger MD+SD [Preliminary]
- ◆ IceCube [Preliminary]
- ◆ NEVOD-DECOR
- Pierre Auger FD+SD
- ▽ SUGAR
- ▲ Yakutsk [Preliminary]
- ▽ EAS-MSU
- ◁ KASCADE-Grande

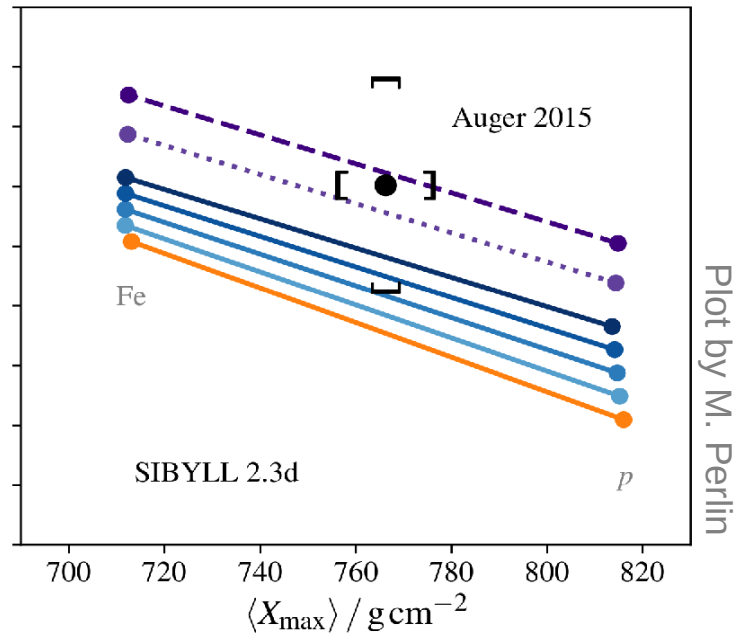
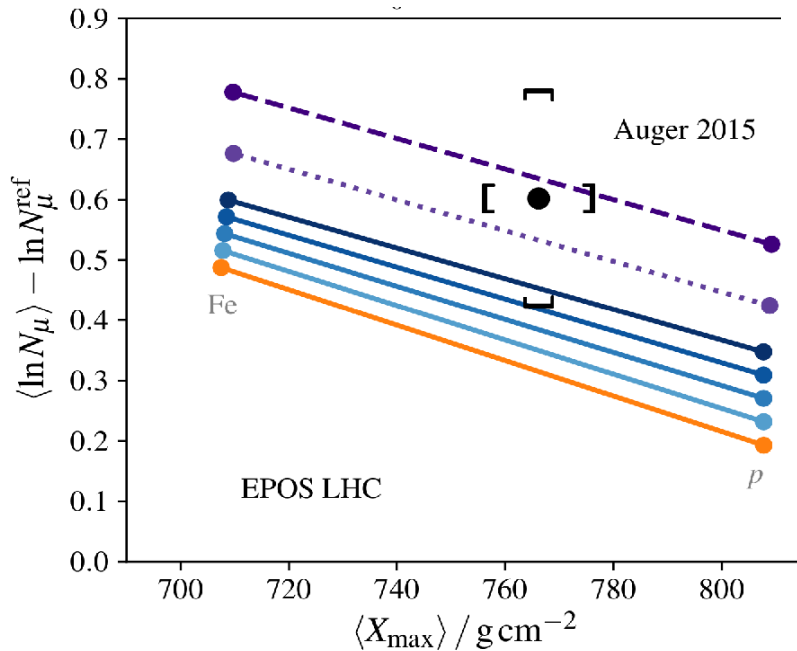
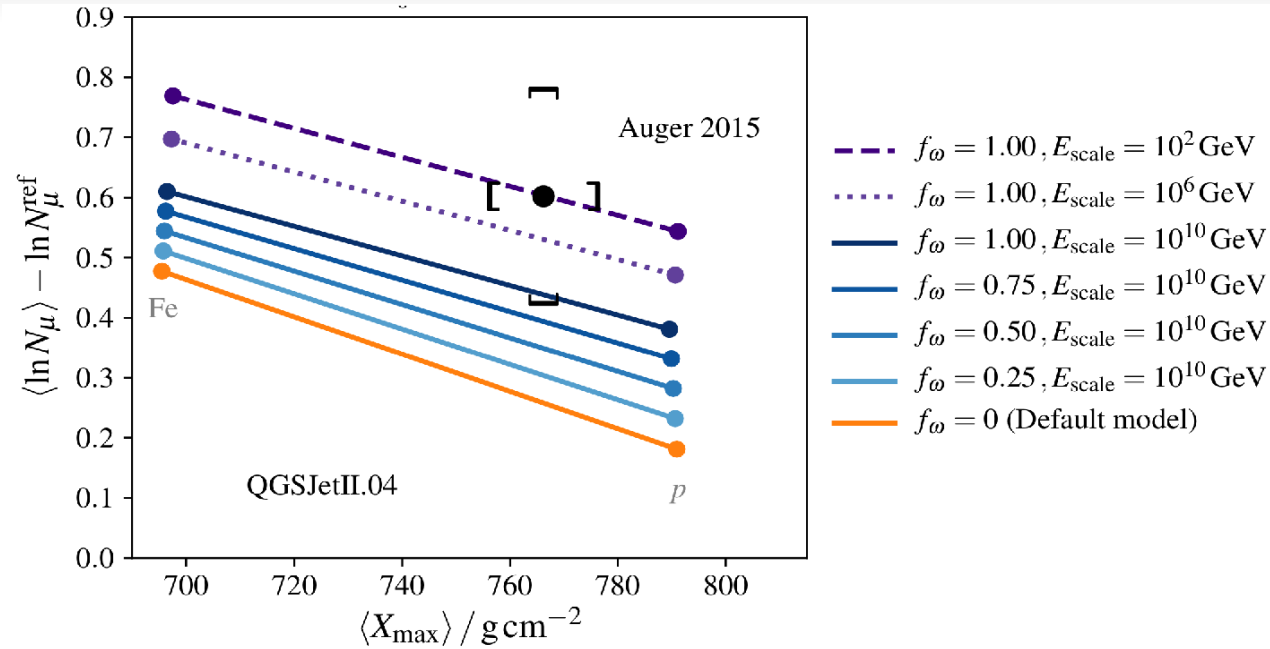
$$z_{\text{mass}} = \frac{\langle \ln A \rangle}{\ln 56}$$

Plot by M. Perlin

Complete Picture and Muon Puzzle

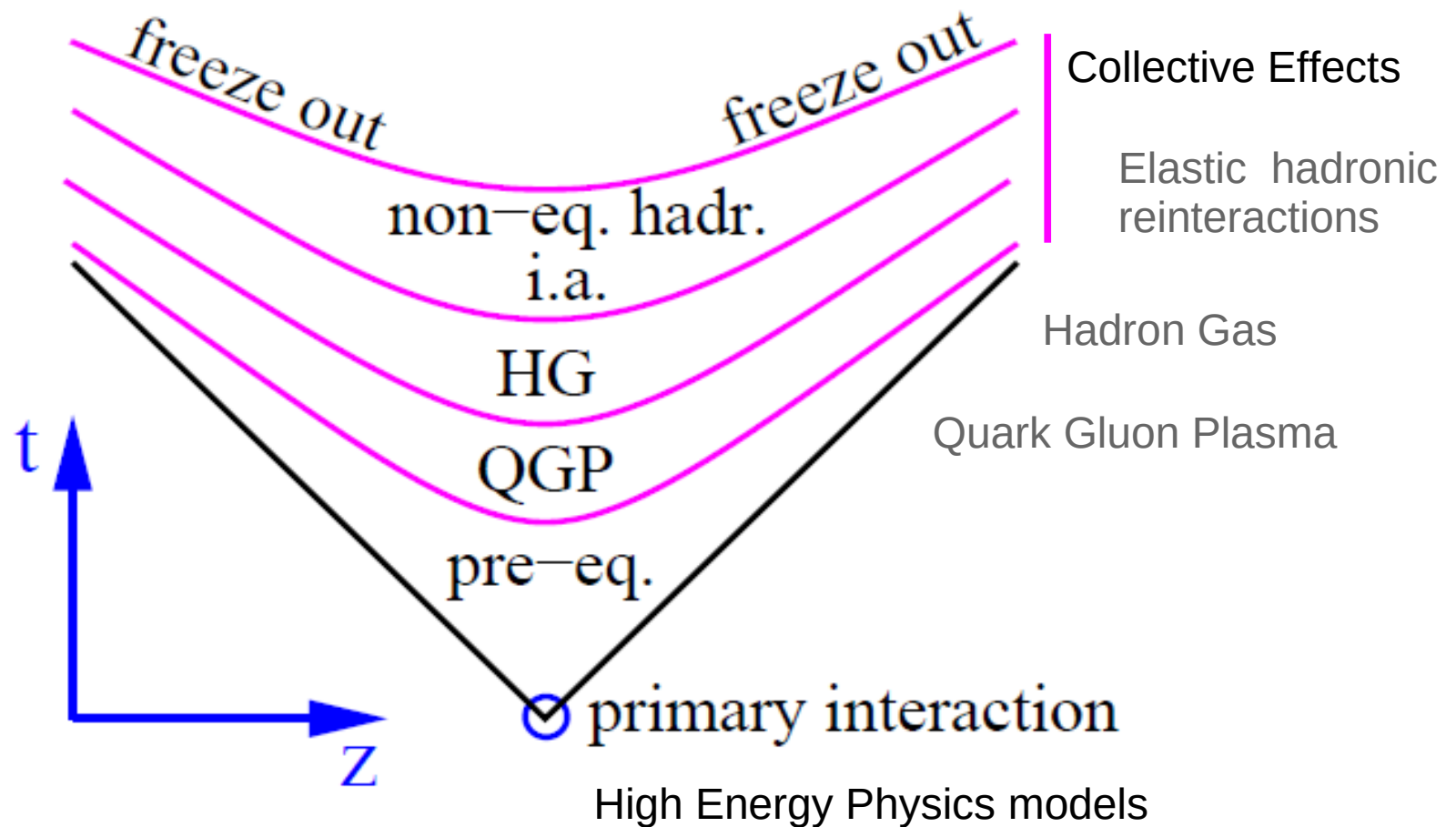
Significant effect on EAS observed if core-corona introduced in models

- ➔ No change in X_{\max}
- ➔ Needs a large part of core hadronization at maximum energy to reach Auger point
- ➔ Sibyll with higher mass (deep X_{\max}) need less



Plot by M. Perlin

High Energy Hadronic Interactions



References :

Ralph Engel PhD Thesis (1997)

K. Werner, Phys.Rept. 232, Nos. 2-5 (1993) 87-299

K. Werner, arXiv:hep-ph/0206111v1

References for High Energy Models (new)

● DPMJET III

- ➔ J. Ranft, R. Engel, S. Roesler, Nucl. Phys. B Proc. Suppl.122, 392 (2003)
- ➔ A. Fedynitch, Cascade equations and hadronic interactions at very high energies. PhD thesis, KIT, Karlsruhe, Dept.Phys. (2015)
- ➔ A. Fedynitch, R. Engel, 14th Inter-national Conference on Nuclear Reaction Mechanisms:Varenna, Italy, p. 291. CERN, Geneva (2015)

● EPOS LHC

- ➔ T. Pierog et al., Phys. Rev. **C92** (2015) 034906 arXiv:1306.0121 [hep-ph]

● QGSJETII-04

- ➔ S. Ostapchenko, Phys. Rev. **D83** (2011) 014018
- ➔ S. Ostapchenko, Phys. Rev. **D89** (2014) 074009

● SIBYLL 2.3d

- ➔ F. Riehn et al. Phys. Rev. **D102** (2020), 063002 arXiv:1912.03300 [hep-ph]

References for High Energy Models (old)

- **DPMJET II**

- ➔ J. Ranft, Phys. Rev. **D51**, 64 (1995)

- **EPOS 1.99**

- ➔ K. Werner et al., Phys. Rev. **C74** (2006) 044902

- ➔ T. Pierog et al., ICRC 2009 Proceedings

- **QGSJET 01**

- ➔ N.N. Kalmykov et al., Nucl. Phys. B (Proc. Suppl.) 52B (1997) 17

- **QGSJETII-03**

- ➔ S. Ostapchenko, Phys. Rev. **D74** (2006) 014026

- ➔ Nucl. Phys. B (Proc. Suppl.) 151 (2006) 143 & 147

- **SIBYLL 2.1**

- ➔ R. Engel et al., Proc. 26th ICRC (Salt Lake City) 1 (1999) 415

- ➔ E.-J. Ahn et al., Phys. Rev. **D80** (2009) 094003

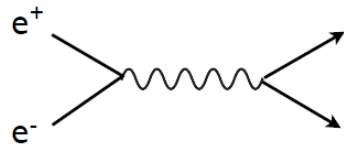
Hadronization Models

2 models well established for 2 extreme cases

➔ String Fragmentation models)

vs Collective hadronization (statistical

Annihilation at high energy

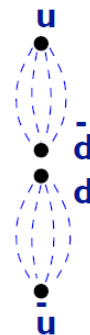


Quarks together are color-neutral system

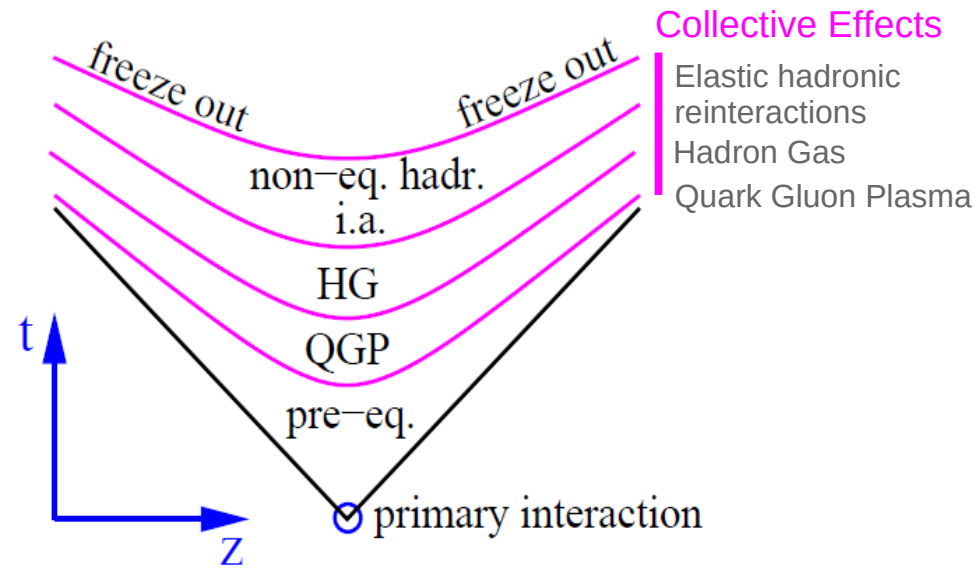


color field

time →



In dilute systems... CORONA
→ "high" π^0 fraction

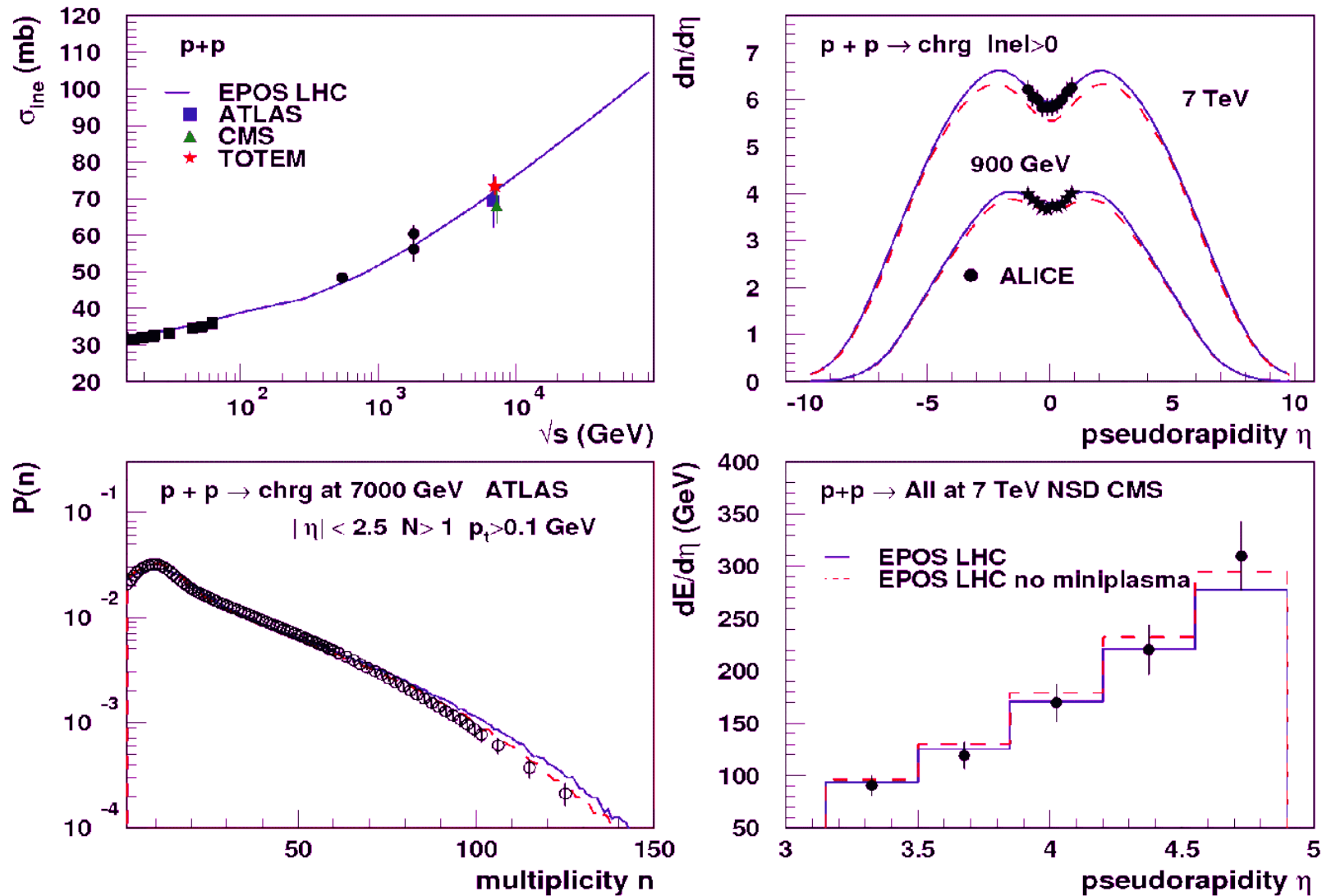


In dense systems... CORE
→ "low" π^0 fraction

➔ Core-corona = mixing of the two for proton-proton, hadron-Air, ..

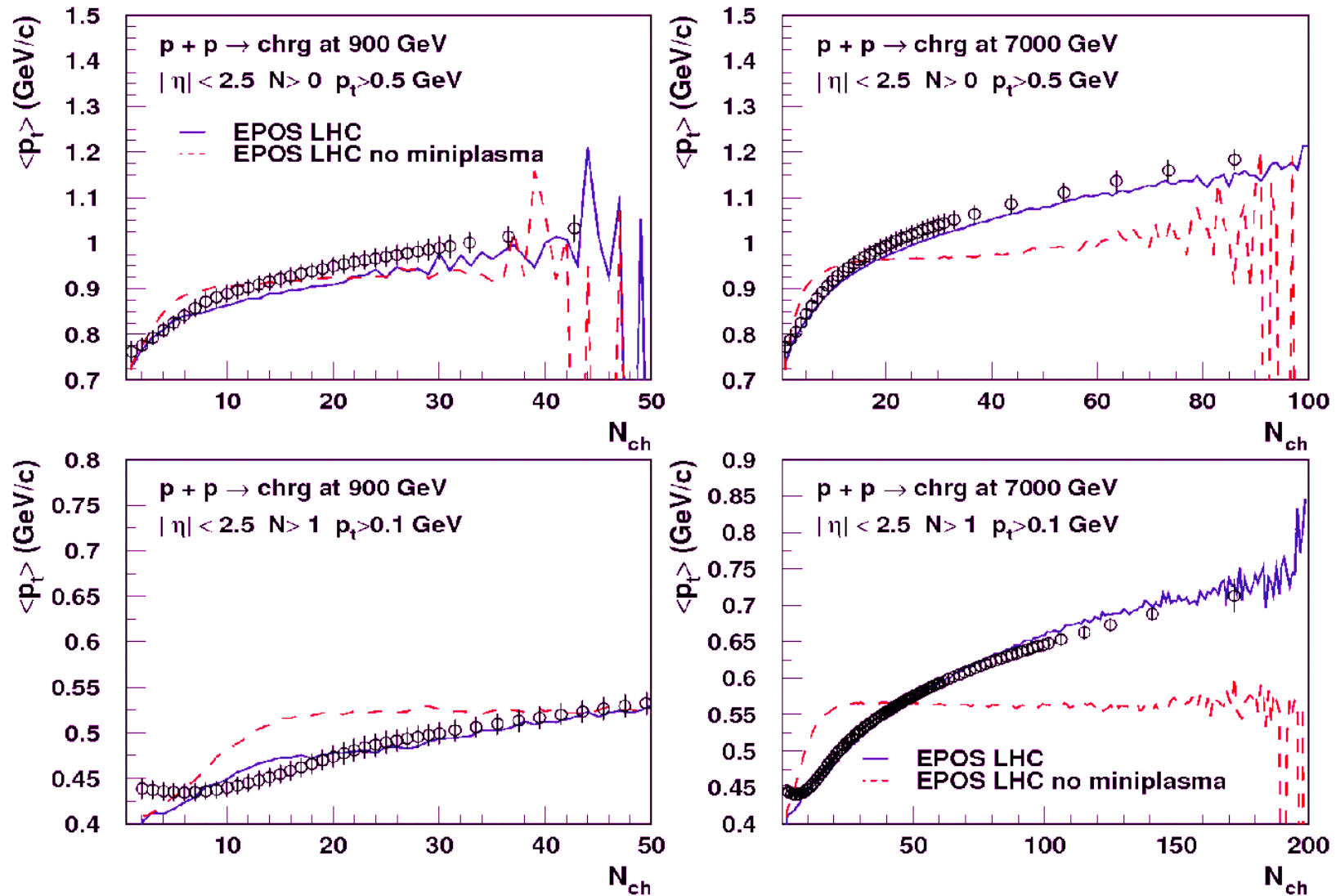
EPOS LHC

● Effective flow treatment



EPOS LHC

● Effective flow treatment

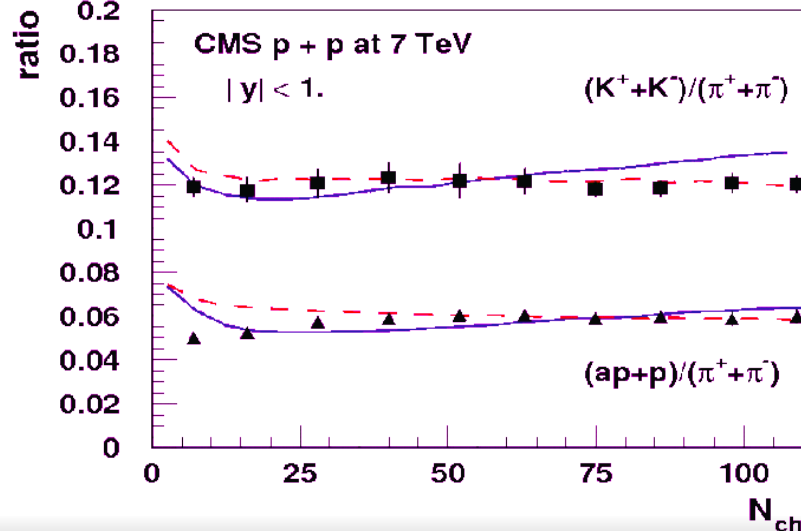
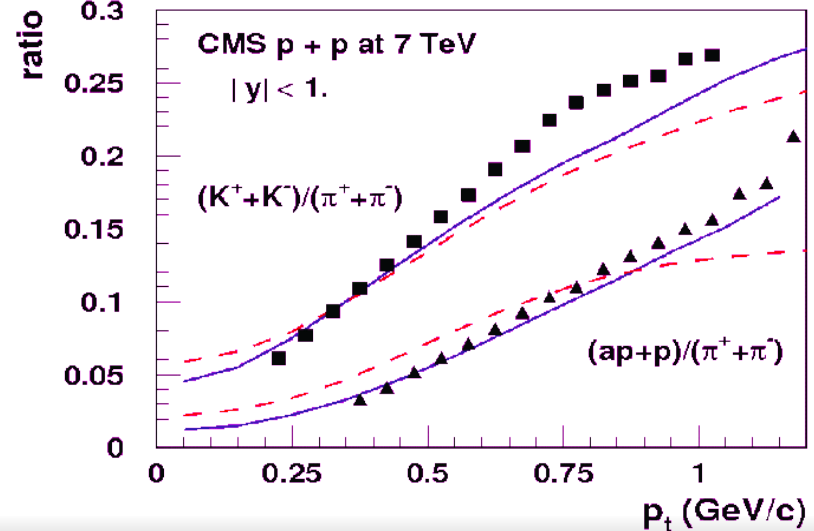
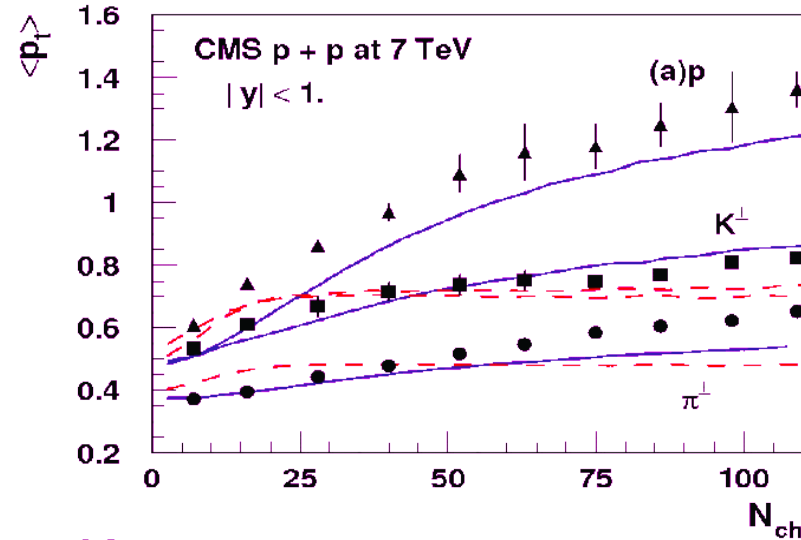
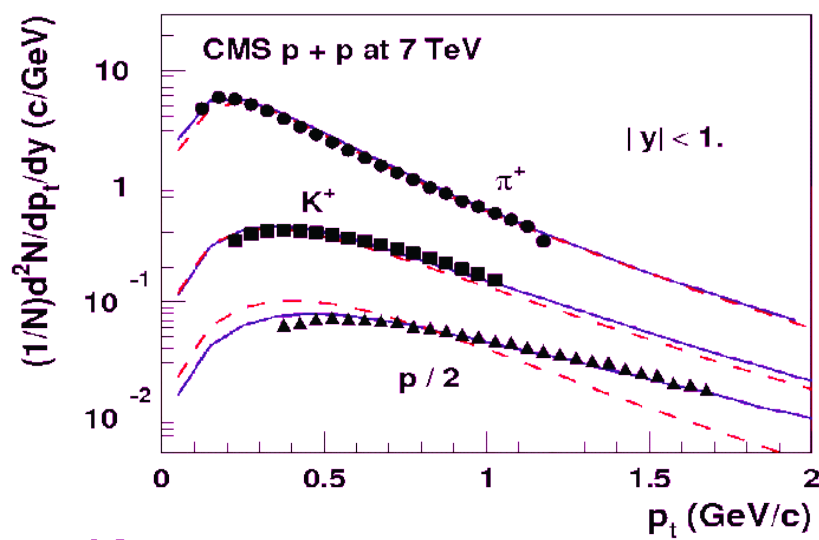


EPOS LHC

● Detailed description can be achieved

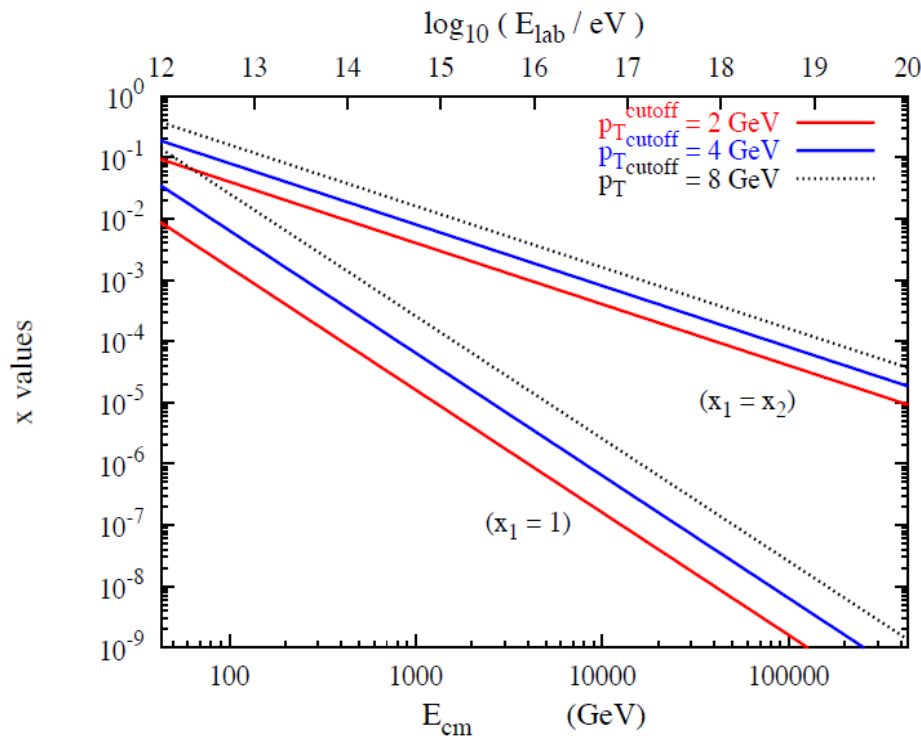
➔ identified spectra

➔ p_t behavior driven by collective effects (statistical hadronization + flow)

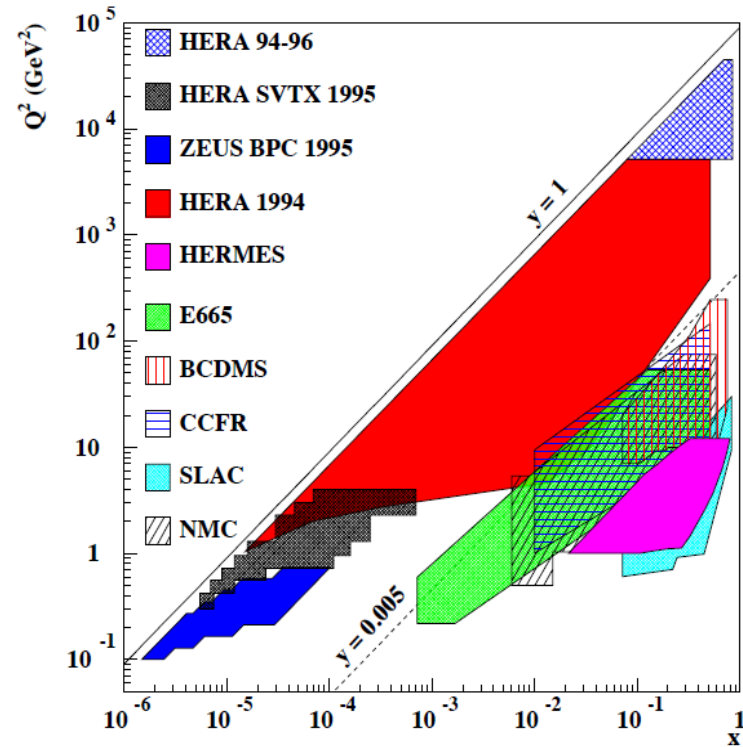


Parton densities not known at very low x

Range of x values (momentum fractions) needed in calculation

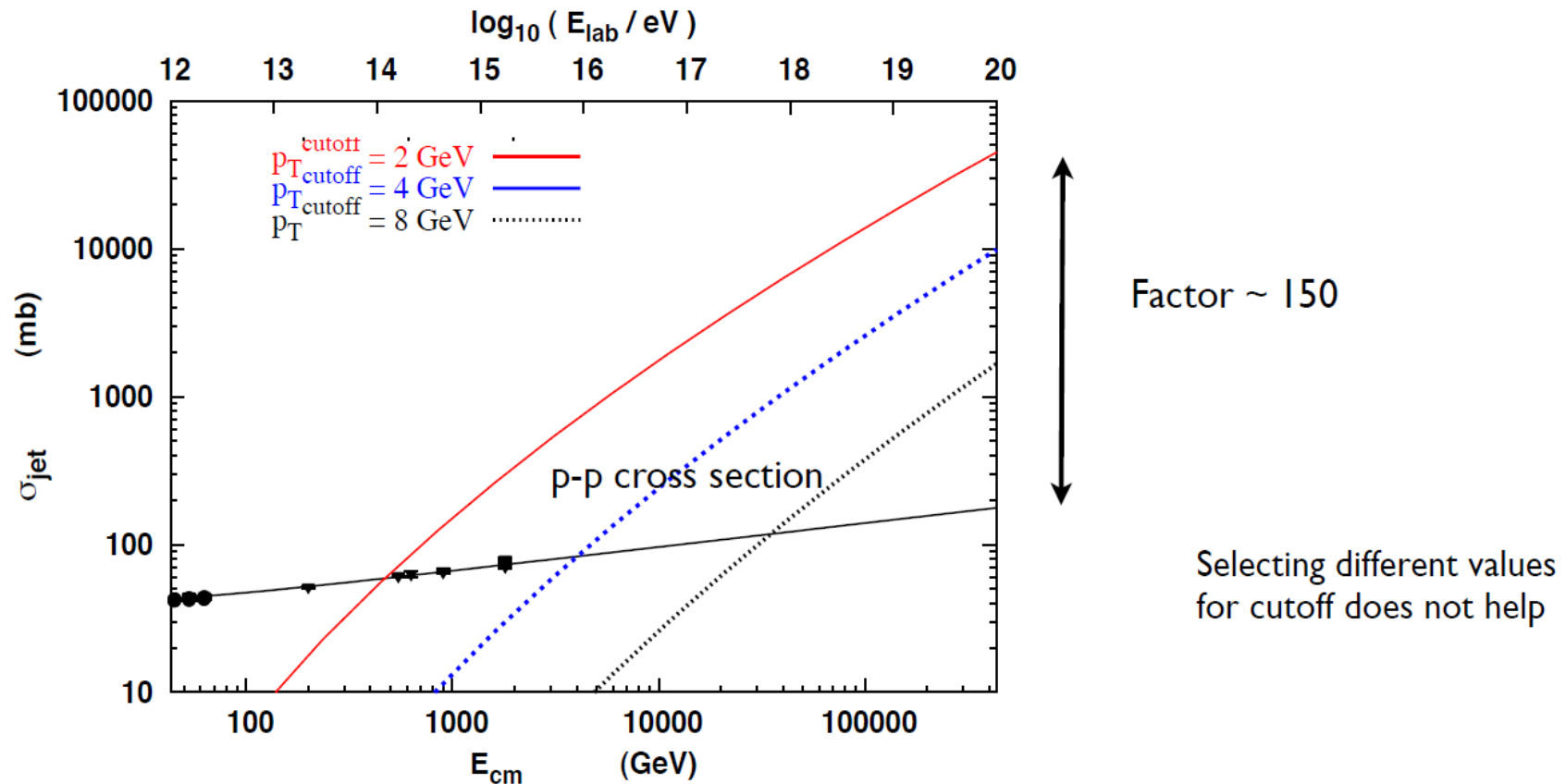


HERA measurement range



$$\hat{s} = x_1 x_2 s \geq 4p_{\perp}^2$$

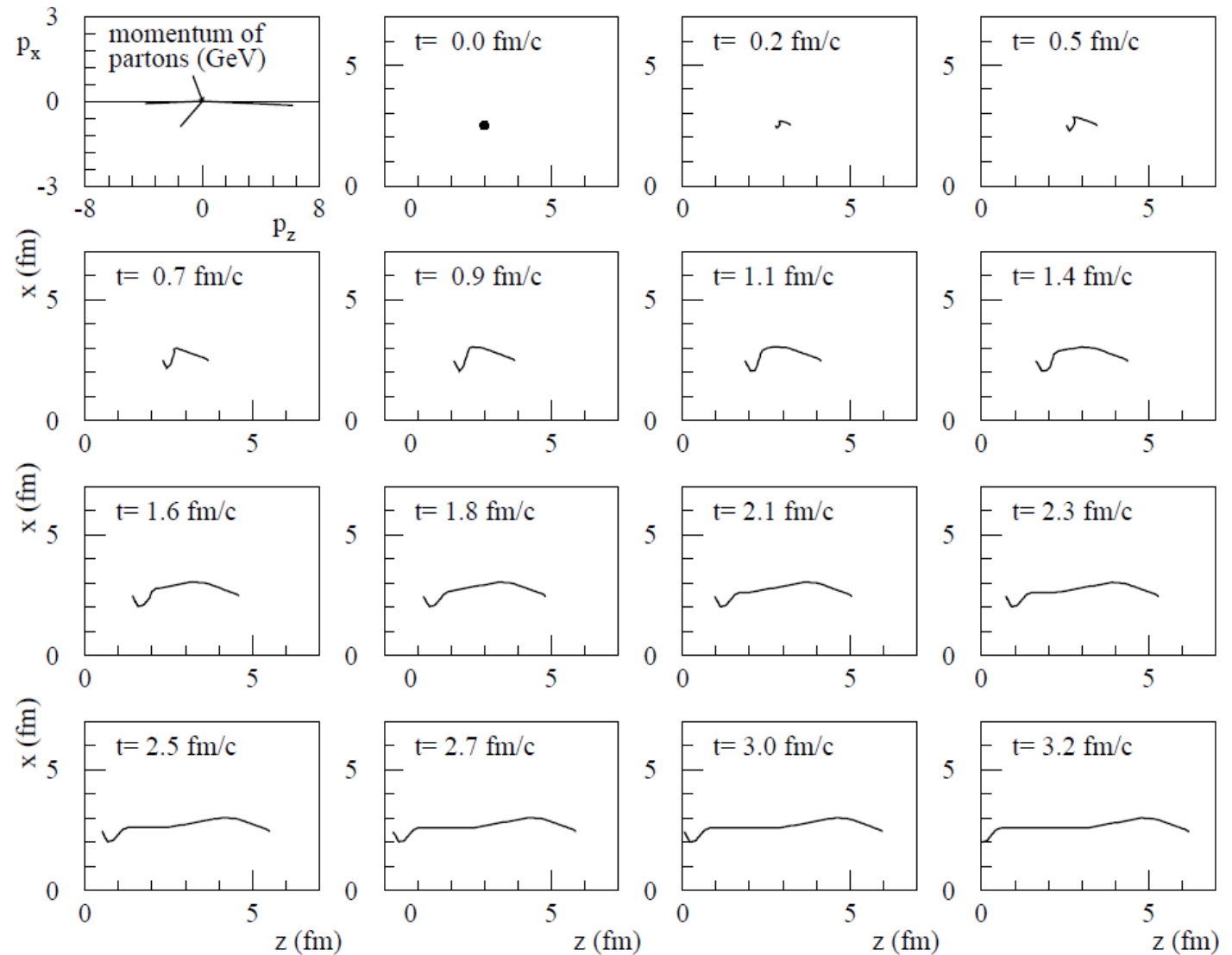
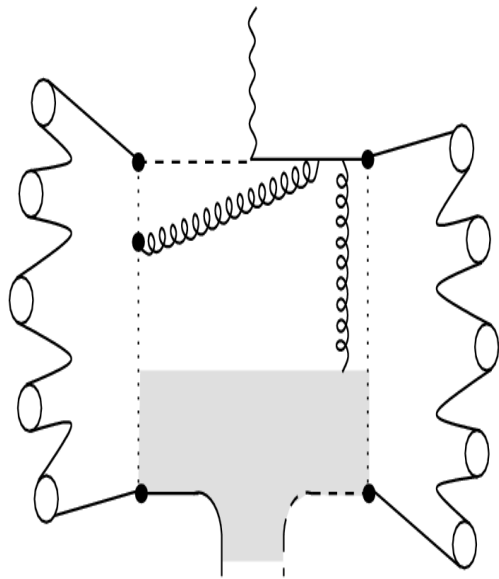
Dependence on transverse momentum cutoff



Cross section of new interaction process of minijet production cannot simply be added to cross section from soft interactions

From cut Pomerons to hadrons: string fragmentation

planar representation
of a cylindrical semi-
hard Pomeron



Are fragmentation parameters universal ?

- **In principle yes but ...**

- ➔ Only EPOS (Pythia) uses LEP data

- Without remnant and statistical hadronization of high density regions (Pythia) : NO

- If everything is taken into account : YES (it seems ...)

- ➔ QGSJET and SIBYLL use hadronic interaction data at low energy to fix parameters for string fragmentation

- possible bias due to remnant contribution

- **String fragmentation models**

- ➔ Lund fragmentation function (Pythia)

- ➔ Area law (EPOS)

Cross Section Calculation

Gribov-Regge for multiple scattering :

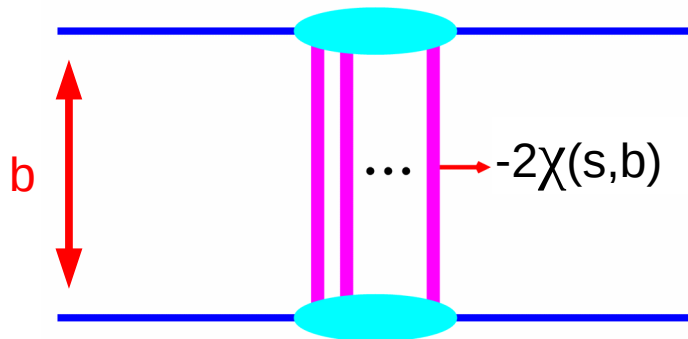
→ elastic amplitude : $-2\chi(s,b)$

→ sum n interactions :

■ Gribov :

$$\frac{(-2\chi)^n}{n!} \rightarrow \exp(-2\chi)$$

$s = (\text{cms energy})^2$
 $b = \text{impact parameter}$



$$\sigma \sim 1 - \exp(-2\chi) \leftarrow \text{Not the same } \chi \text{ in different models}$$

→ $\chi(s,b)$ parameters fixed with pp cross-section

→ pp to pA or AA cross section from GRT (including nuclear effect in QGSJETII)

classical GRT: energy conservation not taken into account at this level

Particle Production from Pomerons

number n of exchanged elementary interaction per event fixed from elastic amplitude (cross section) :

→ n cut Pomerons from :

$$P(n) = \frac{(2\chi)^n}{n!} \cdot \exp(-2\chi)$$

- no energy sharing accounted for (interference term)

→ $2n$ strings formed from the n elementary interactions

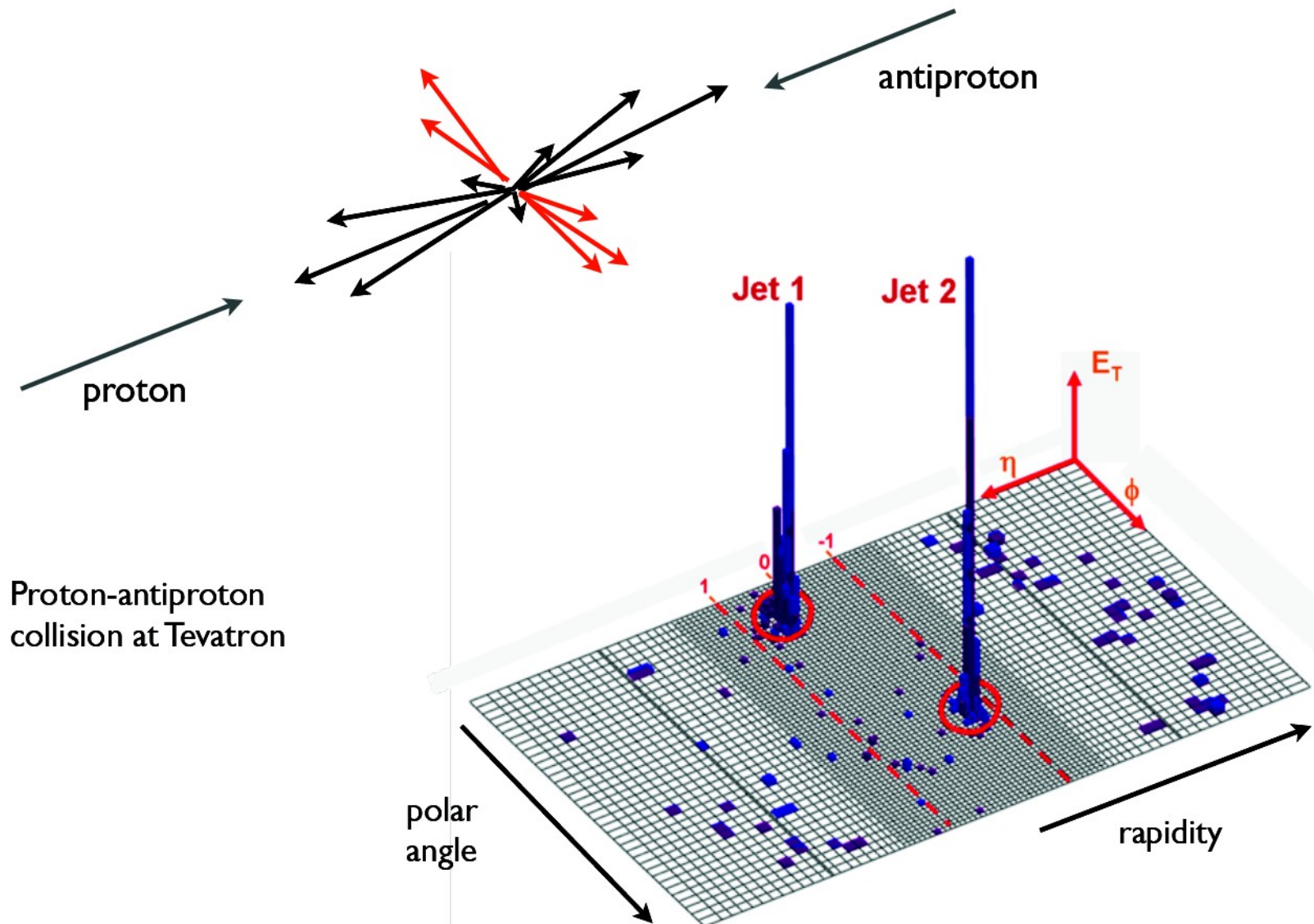
- energy conservation : energy shared between the $2n$ strings

- particles from string fragmentation

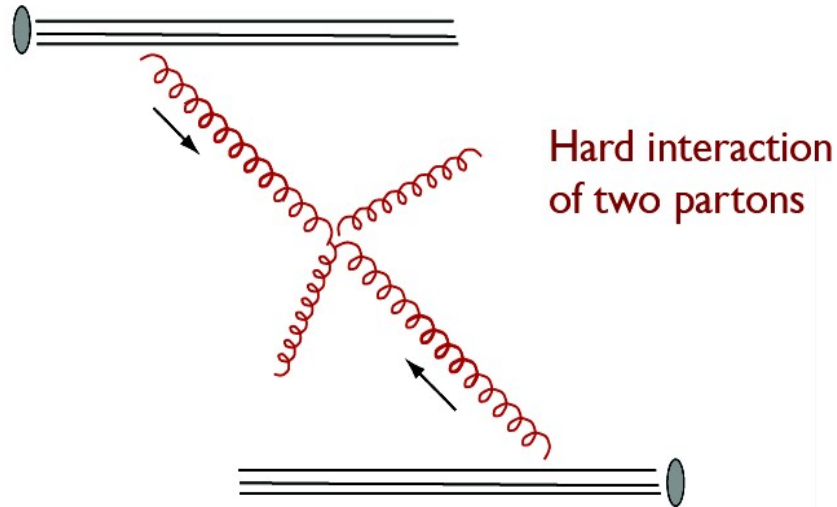
→ **inconsistency** : energy sharing should be taken into account when fixing n

→ alternative approach with energy conservation from the beginning (EPOS)

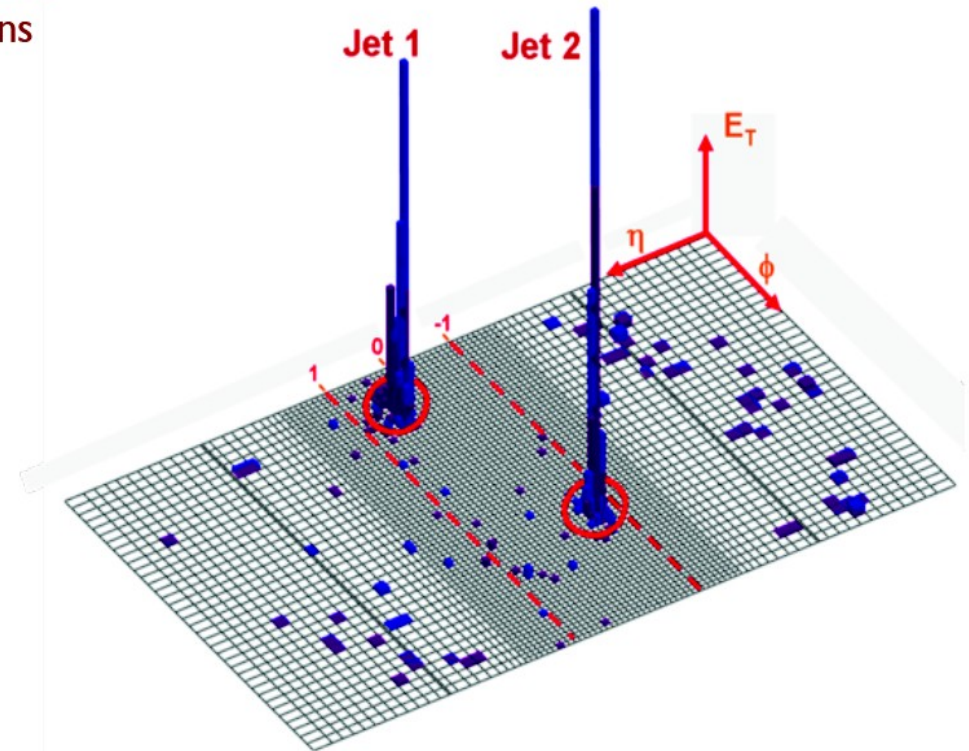
Scattering of quarks and gluons: jet production



Interpretation within perturbative QCD



QCD predictions known
for parton-parton cross sections



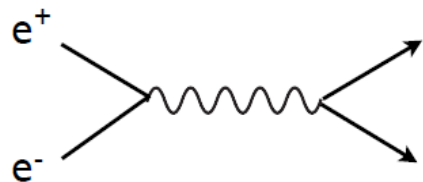
Terminology

Soft interaction: no large momentum transfer

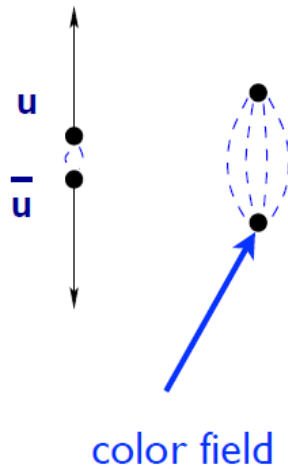
Hard interaction: large momentum transfer ($|t| > 2 \text{ GeV}^2$)

Simplest case: e^+e^- annihilation into quarks

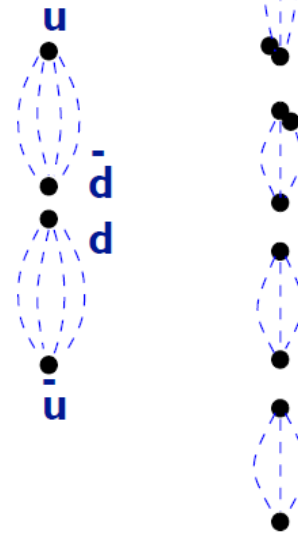
Annihilation at high energy



Quarks together are color-neutral system



time



.....

- $u\bar{d}$
- $d\bar{u}$
- $\bar{u}\bar{u}\bar{d}$
- udd
- $u\bar{s}$
- $s\bar{d}$
- $u\bar{d}$
- $q\bar{q}$
- $q\bar{q}$
- $q\bar{q}$

Chain of hadrons

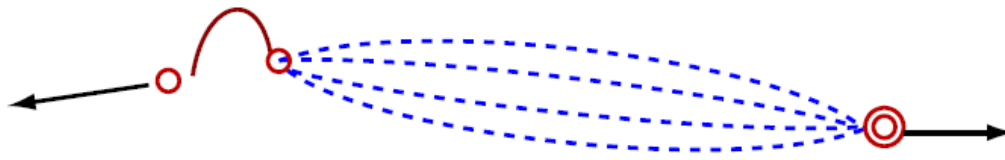
String fragmentation

Fragmentation function (SIBYLL)

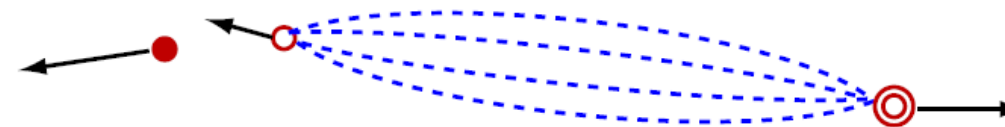
String characterized by momentum fractions of partons at ends



$$E_{\text{str}} = \frac{1}{2} \sqrt{s} (x_1 + x_2)$$



$$p_{\text{str}} = \frac{1}{2} \sqrt{s} (x_1 - x_2)$$



Momentum fraction z
of new meson relative
to quark at string end

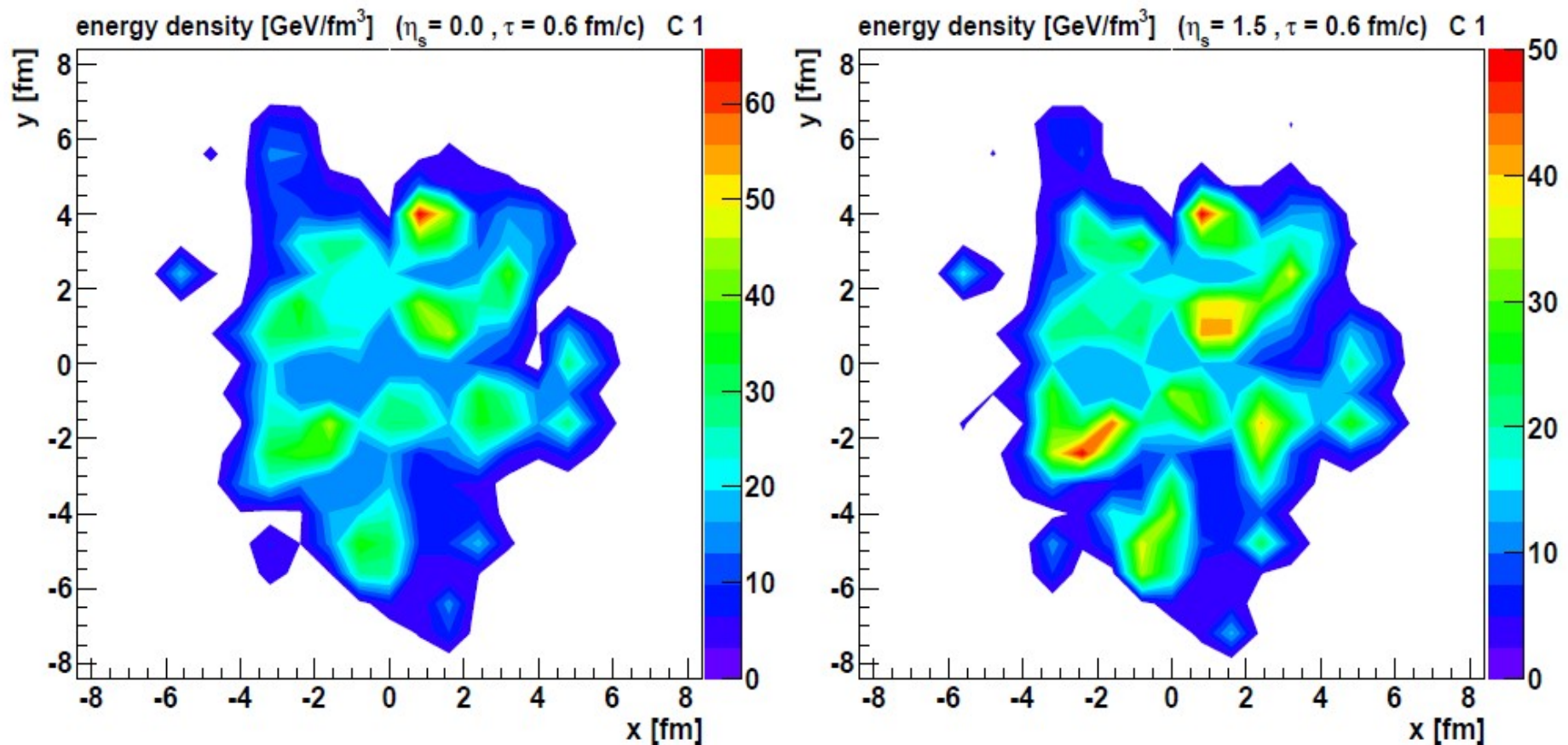
$$m_{\perp}^2 = p_{\perp}^2 + m^2$$

$$f(z) = \frac{(1-z)^a}{z} \exp \left\{ -\frac{bm_{\perp}^2}{z} \right\}$$

Event-by-Event Energy Density : AuAu

➔ Bumpy structure of energy density in transverse plane, but translational invariance

● pseudorapidity extension of flux tubes (strings)

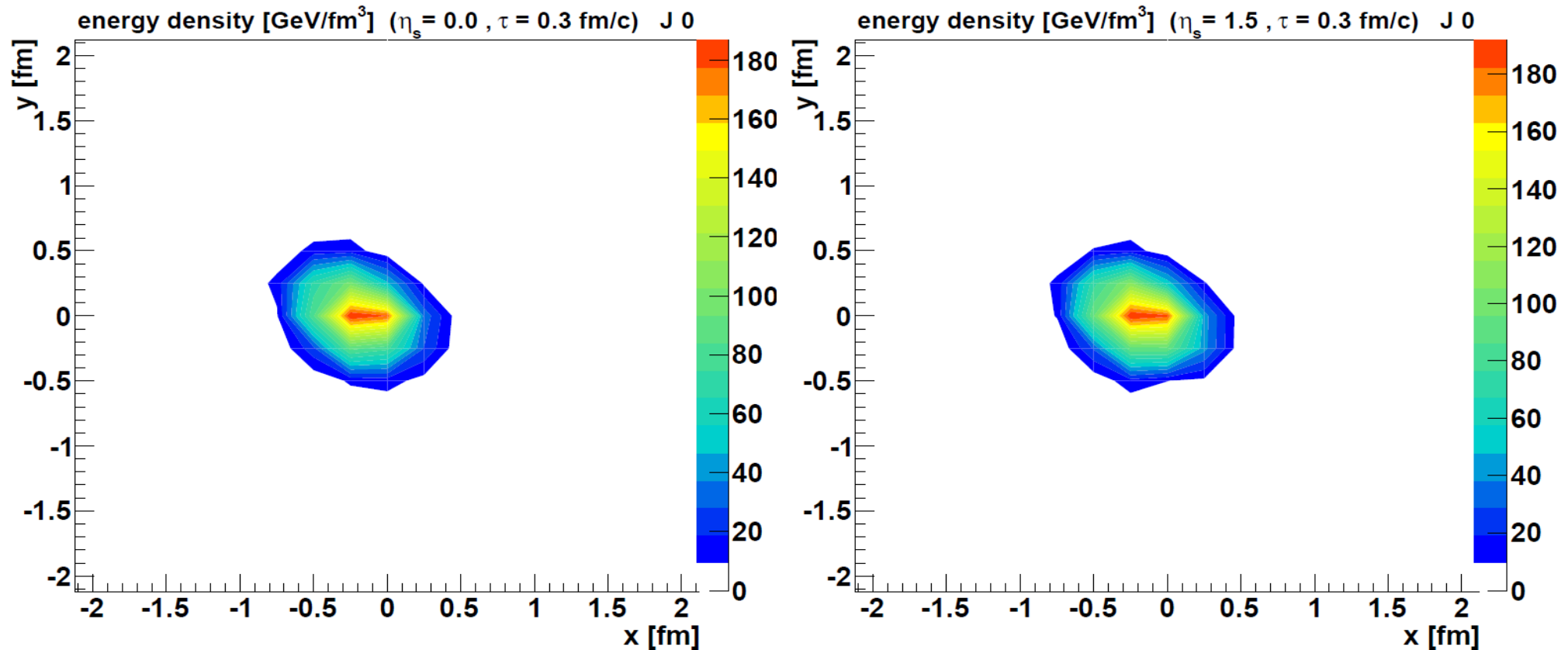


Initial energy density in the transverse plane for two different η_s

Event-by-Event Energy Density : pp

➔ Random azimuthal asymmetries of initial energy density but translational invariance

● pseudorapidity extension of flux tubes (strings)



Initial energy density in the transverse plane for two different η_s