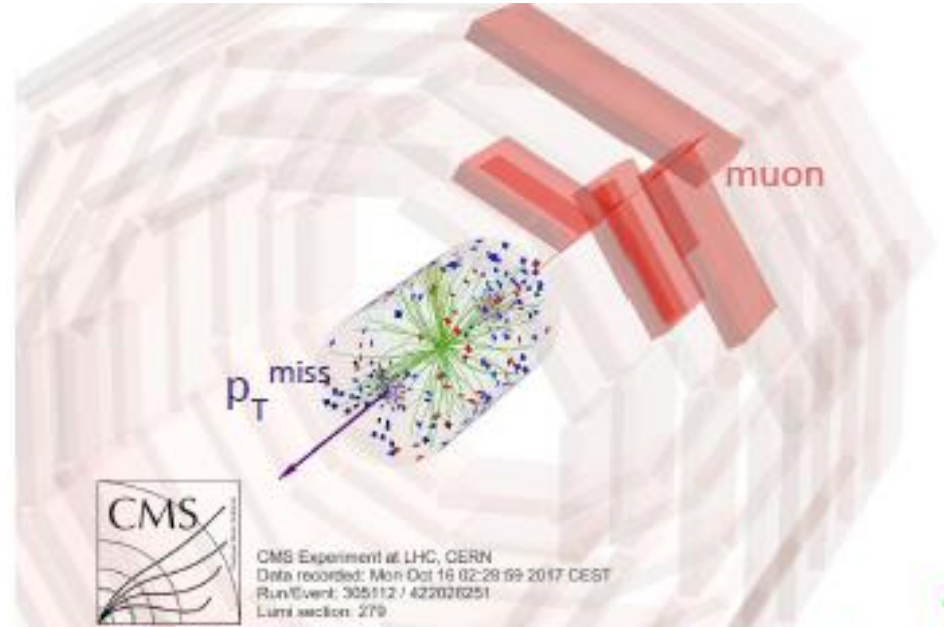


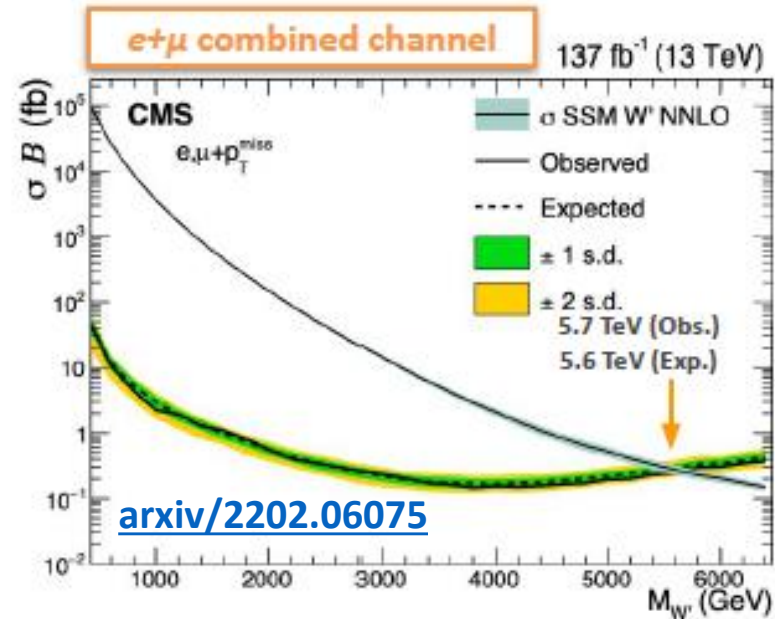
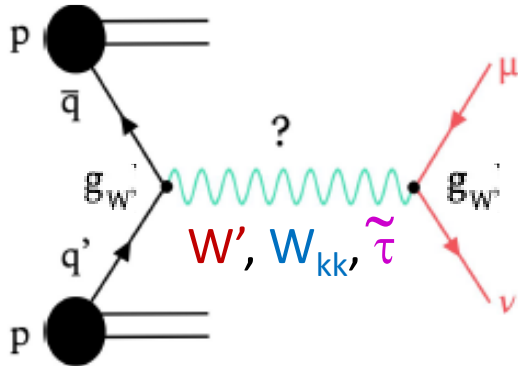
Direct & Indirect constraints on new physics from $l+p_t^{\text{miss}}$ (l^+l^-) final states in CMS

B. de la Cruz, I. Bachiller, J. Alcaraz
(CIEMAT)
+ CMS Collaboration

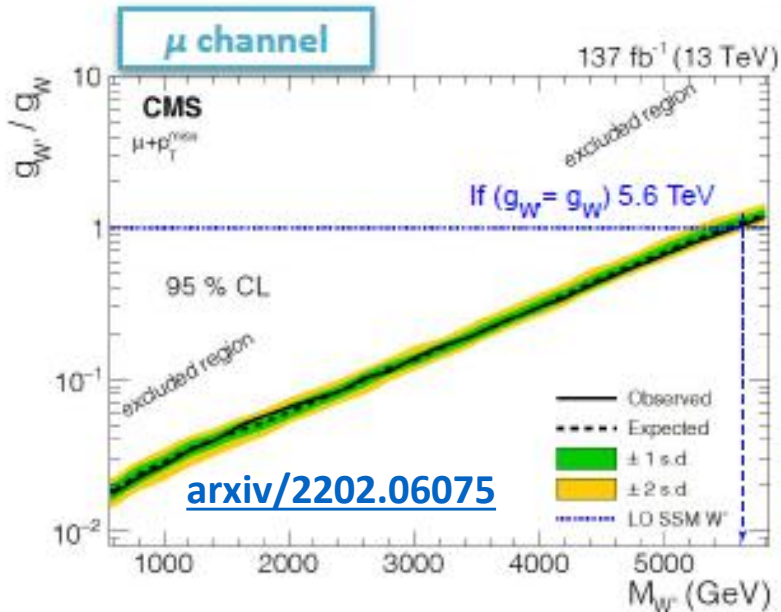


Motivation and context

137 fb⁻¹ pp collision data from CMS during LHC Run2 used to search for new physics



M_{W'} > 5.7 (5.6) TeV Obs. (Exp.)



For M_{W'} < 1 TeV, limit g_{W'}/g_W > 0.03 excluded

Direct searches, limited by √s

Motivation and context

- **SM** constitutes an effective theory, applicable up to energies not exceeding a certain **scale Λ** .
- **Extension of SM** at higher energies imply including **higher-dimensional operators** which appear in the extended-SM Lagrangian as **suppressed by powers of Λ** .

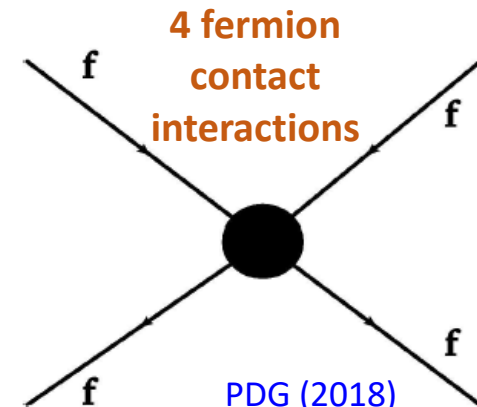
$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{i=1}^{2499} \frac{C_i}{\Lambda^2} \mathcal{O}_i, \quad \text{With } \Lambda > \sqrt{s}$$

Motivation and context

A. Wulzer et al
[10.1016/j.physletb.2017.06.043](https://doi.org/10.1016/j.physletb.2017.06.043)

- **SM** constitutes an effective theory, applicable up to energies not exceeding a certain **scale Λ** .
- **Extension of SM** at higher energies imply including **higher-dimensional operators** which appear in the extended-SM Lagrangian as **suppressed by powers of Λ** .
- Using the **formalism of the SM effective field theories (SMEFT)**, truncating the effective expansion at dimension 6 operators, universal new physics effects in **4 fermion contact interactions** type appear, which can be conveniently* described by the well-known “**oblique parameters**” **S , T , W** and **Y** .

$$\Delta\mathcal{L}_{\text{Universal}} = \frac{S}{16\pi v^2} \mathcal{O}_{WB} - \frac{2\alpha T}{v^2} \mathcal{O}_{HD} - \frac{Y}{2M_W^2} \mathcal{O}_{2B} - \frac{W}{2M_W^2} \mathcal{O}_{2W}$$



DOI: [10.22323/1.314.0467
/hep-ph/0405040](https://doi.org/10.22323/1.314.0467/hep-ph/0405040)

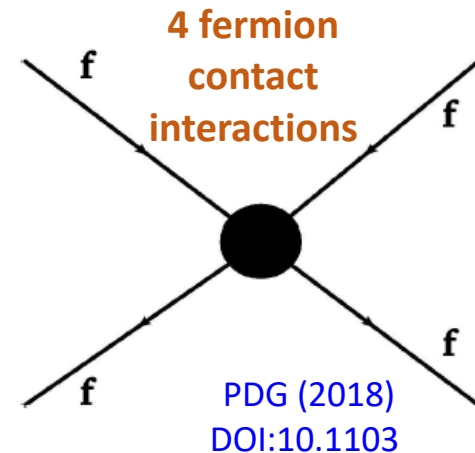
Phys.Rev. D46 (1992) 381

* Involving operators of W , Z , γ , Higgs

Motivation and context

- **SM** constitutes an effective theory, applicable up to energies not exceeding a certain **scale Λ** .
- **Extension of SM** at higher energies imply including **higher-dimensional operators** which appear in the extended-SM Lagrangian as **suppressed by powers of Λ** .
- Using the **formalism of the SM effective field theories (SMEFT)**, truncating the effective expansion at dimension 6 operators, universal new physics effects in **4 fermion contact interactions** type appear, which can be conveniently* described by the well-known “**oblique parameters**” **S , T , W** and **Y** .

$$\Delta\mathcal{L}_{\text{Universal}} = \frac{S}{16\pi v^2} \mathcal{O}_{WB} - \frac{2\alpha T}{v^2} \mathcal{O}_{HD} - \frac{Y}{2M_W^2} \mathcal{O}_{2B} - \frac{W}{2M_W^2} \mathcal{O}_{2W}$$



* Involving operators of W , Z , γ , Higgs

S , T : induce effects constant with \sqrt{s}
current limits coming from fit to EW Precision Observables

W , Y : induce effects that grow with \sqrt{s} → high-energy collider is ideal to test new physics → Highest energies at LHC compensate worse experimental precision

DOI: [10.22323/1.314.0467/hep-ph/0405040](https://doi.org/10.22323/1.314.0467/hep-ph/0405040)
Phys.Rev. D46 (1992) 381

W, Y parameters

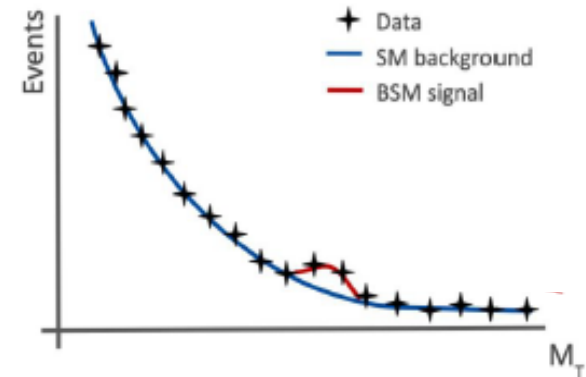
W, Y: leading constraints on W, Y from off-pole measurements at LEP2, but 8 TeV LHC results are more stringent → expected to be improved with LHC @13 TeV and HL-LHC in the future

$$\begin{pmatrix} S \\ T \\ W \\ Y \end{pmatrix} = \begin{pmatrix} -0.10 \pm 0.13 \\ 0.02 \pm 0.08 \\ (-0.1 \pm 0.6) \times 10^{-3} \\ (-1.2 \pm 0.9) \times 10^{-3} \end{pmatrix}, \quad \rho = \begin{pmatrix} 1 & & & \\ 0.86 & 1 & & \\ -0.12 & -0.06 & 1 & \\ 0.70 & 0.39 & -0.49 & 1 \end{pmatrix}$$

<https://arxiv.org/pdf/1706.03783.pdf>

Useful Variables: Differential cross sections vs angle (η), invariant mass, M_T

Dilepton final states studied at LHC: $qq \rightarrow l+l-, lv$ ($l = e, \mu$)



W, Y parameters

W, Y: leading constraints on W, Y from off-pole measurements at LEP2, but 8 TeV LHC results are more stringent → expected to be improved with LHC @13 TeV and HL-LHC in the future

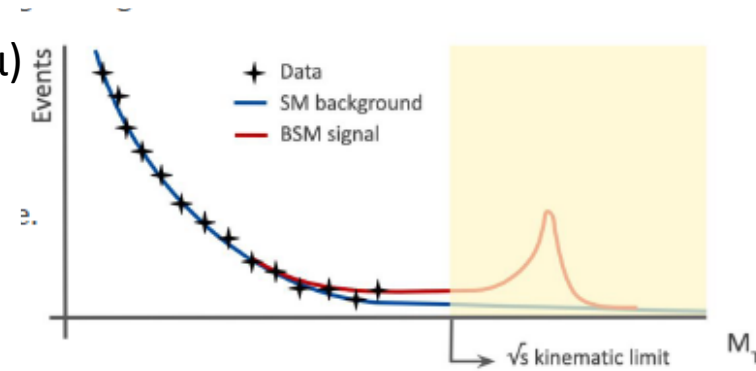
$$\begin{pmatrix} S \\ T \\ W \\ Y \end{pmatrix} = \begin{pmatrix} -0.10 \pm 0.13 \\ 0.02 \pm 0.08 \\ (-0.1 \pm 0.6) \times 10^{-3} \\ (-1.2 \pm 0.9) \times 10^{-3} \end{pmatrix}, \quad \rho = \begin{pmatrix} 1 & & & \\ 0.86 & 1 & & \\ -0.12 & -0.06 & 1 & \\ 0.70 & 0.39 & -0.49 & 1 \end{pmatrix}$$

<https://arxiv.org/pdf/1706.03783.pdf>

Useful Variables: Differential cross sections vs angle (η), invariant mass, m_T

Dilepton final states studied at LHC: $qq \rightarrow l+l-, lv$ ($l = e, \mu$)

Effect of new terms in 4-fermion contact interactions is the **modification of SM propagators** of neutral (γ, Z) and charged (W) EWK gauge bosons, $P(W, Y)$.

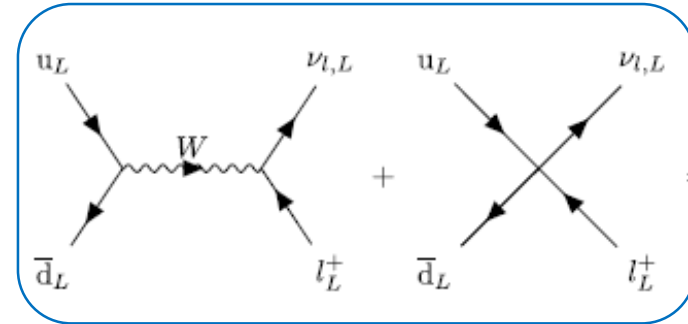


- relative deviations wrt SM given by the weight $|P(W, Y)/P(W=Y=0)|^2$
- a **reweighting** on evt-by-evt basis at generator level is possible
- $\frac{g^2 W}{m_W^2}, \frac{g'^2 Y}{m_W^2} \rightarrow \frac{\mathbf{C}}{\text{Scale}^2}$ Terms in the interference between SM and New Physics, related to the ratio **coupling/scale** of NP

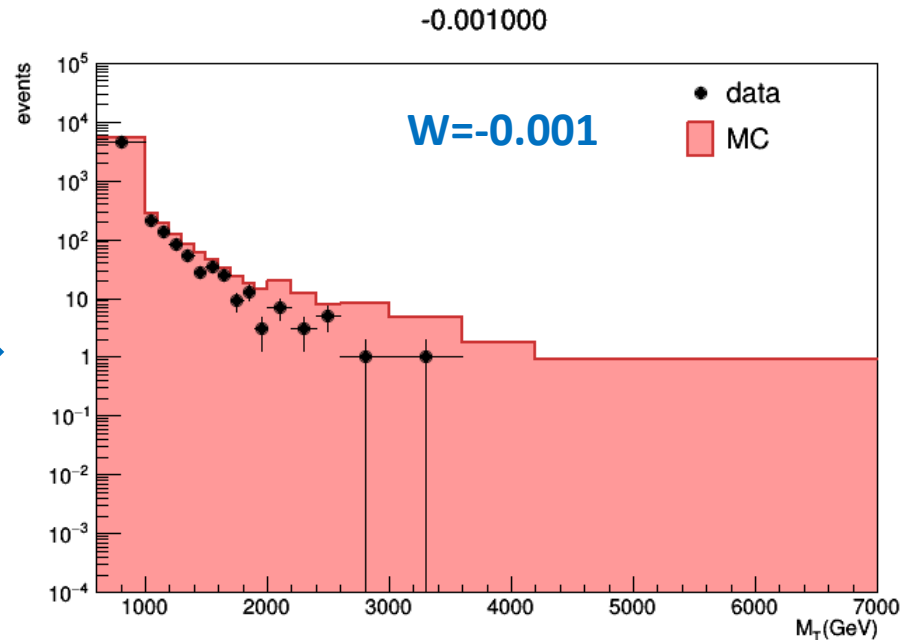
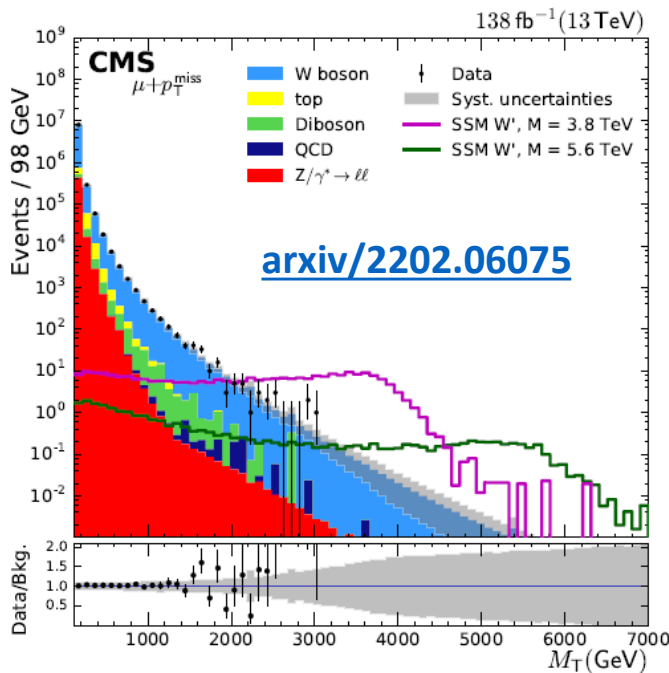
$qq \rightarrow lv \ (l = e, \mu)$

$$\left| \frac{P(W, Y)}{P(W=Y=0)} \right|^2 = \left(1 + \frac{(2t^2 - 1)W}{1 - t^2} + \frac{t^2 Y}{1 - t^2} - \frac{W(q^2 - m_W^2)}{m_W^2} \right)$$

$(t^2 = \text{tangent squared of SM weak mixing angle} \sim 0.3)$



- Relevant term at high energy is the last one, depending on q^2
- lv final state is sensitive to W , independent of Y at high energy.
- Only info needed at generator level is q^2 , i.e. lv invariant mass

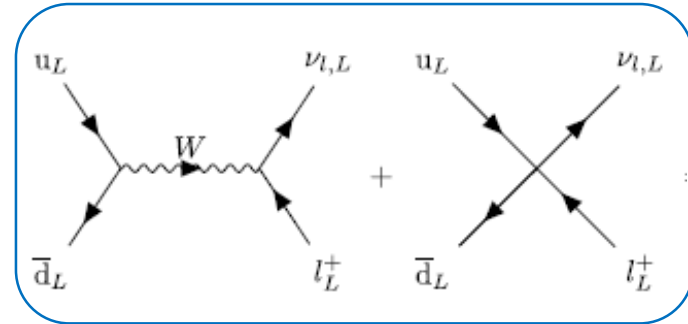


$ev + \mu\nu, L = 101 \text{ fb}^{-1}$

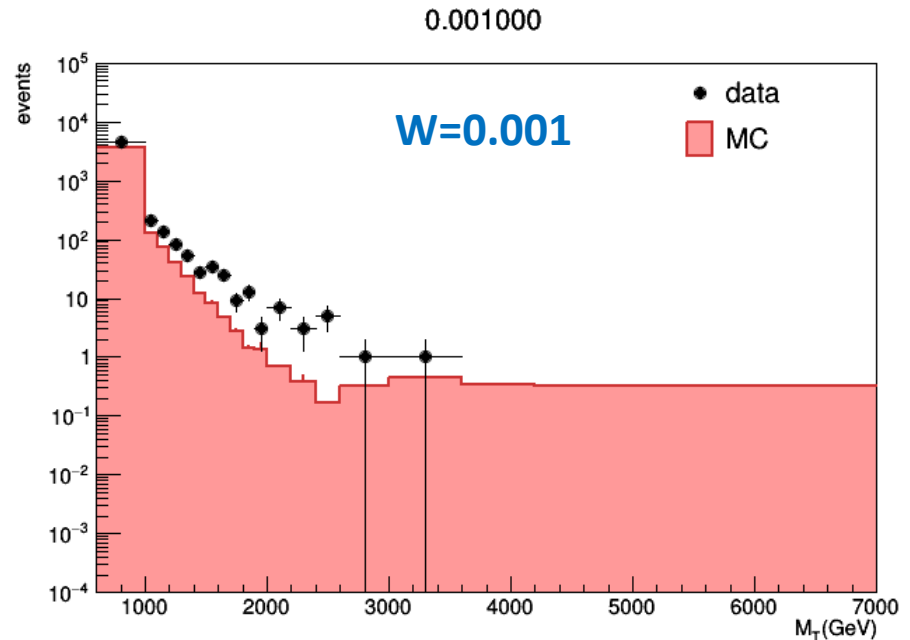
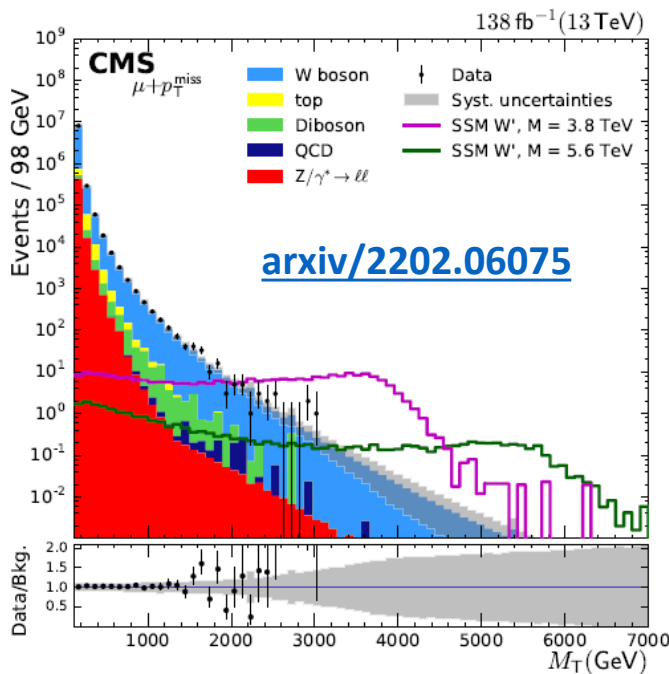
$qq \rightarrow lv \ (l = e, \mu)$

$$\left| \frac{P(W, Y)}{P(W=Y=0)} \right|^2 = \left(1 + \frac{(2t^2 - 1)W}{1 - t^2} + \frac{t^2 Y}{1 - t^2} - \frac{W(q^2 - m_W^2)}{m_W^2} \right)$$

$(t^2 = \text{tangent squared of SM weak mixing angle} \sim 0.3)$



- Relevant term at high energy is the last one, depending on q^2
- lv final state is sensitive to W , independent of Y at high energy.
- Only info needed at generator level is q^2 , i.e. lv invariant mass

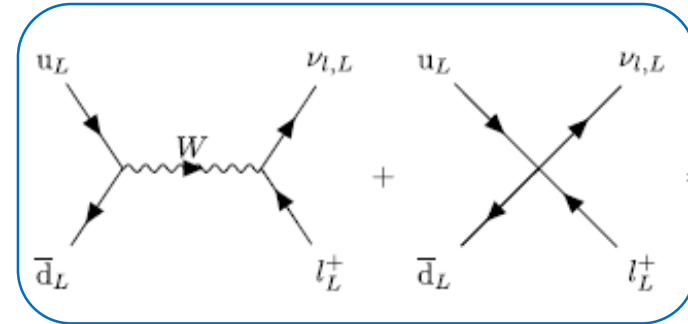


$ev + \mu\nu, L=101 \text{ fb}^{-1}$

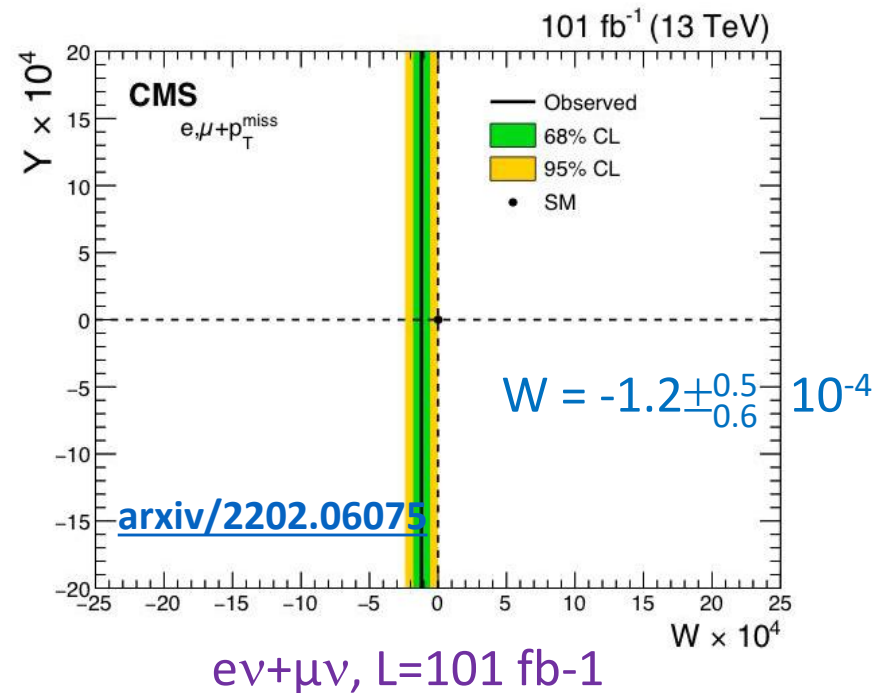
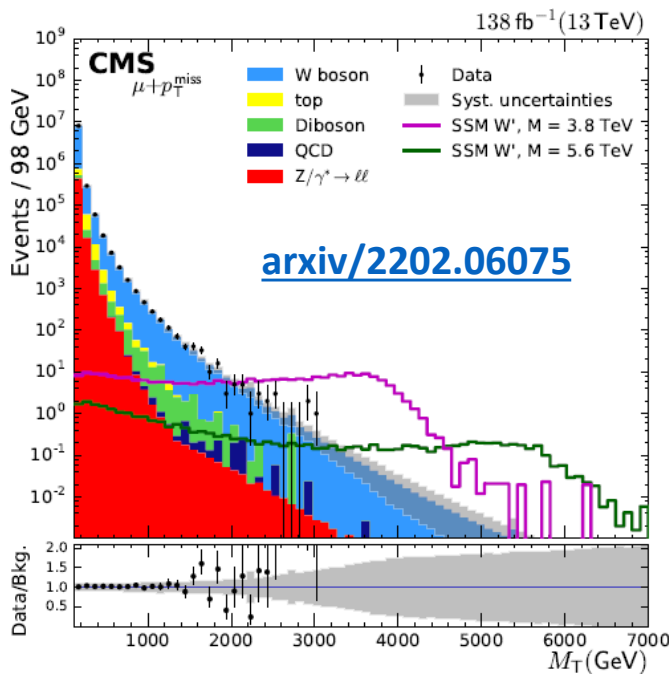
$qq \rightarrow l\nu$ ($l = e, \mu$)

$$\left| \frac{P(W, Y)}{P(W=Y=0)} \right|^2 = \left(1 + \frac{(2t^2 - 1)W}{1 - t^2} + \frac{t^2 Y}{1 - t^2} - \frac{W(q^2 - m_W^2)}{m_W^2} \right)$$

(t^2 =tangent squared of SM weak mixing angle ~ 0.3)



- Relevant term at high energy is the last one, depending on q^2
- $l\nu$ final state is sensitive to W , independent of Y at high energy.
- Only info needed at generator level is q^2 , i.e. $l\nu$ invariant mass

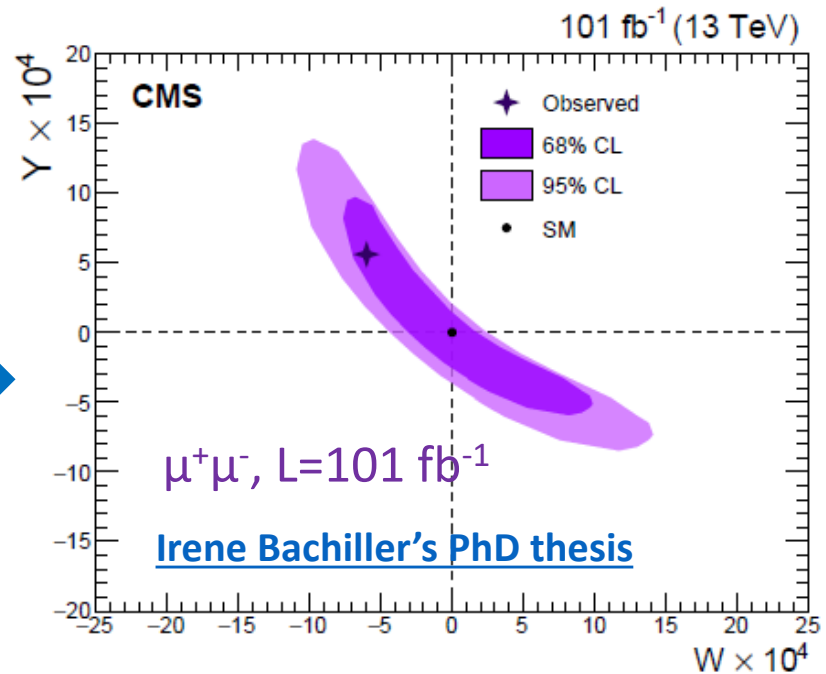
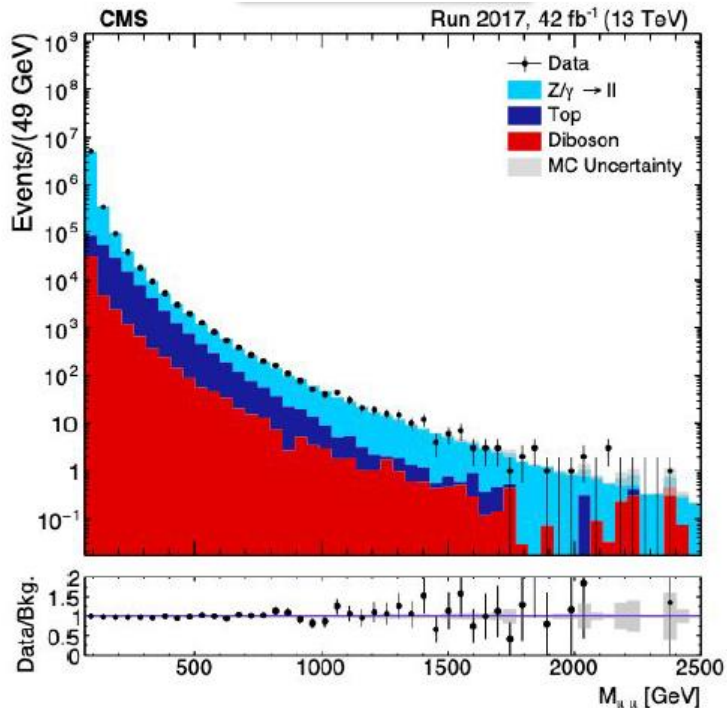


$qq \rightarrow l^+l^- (l = \mu)$

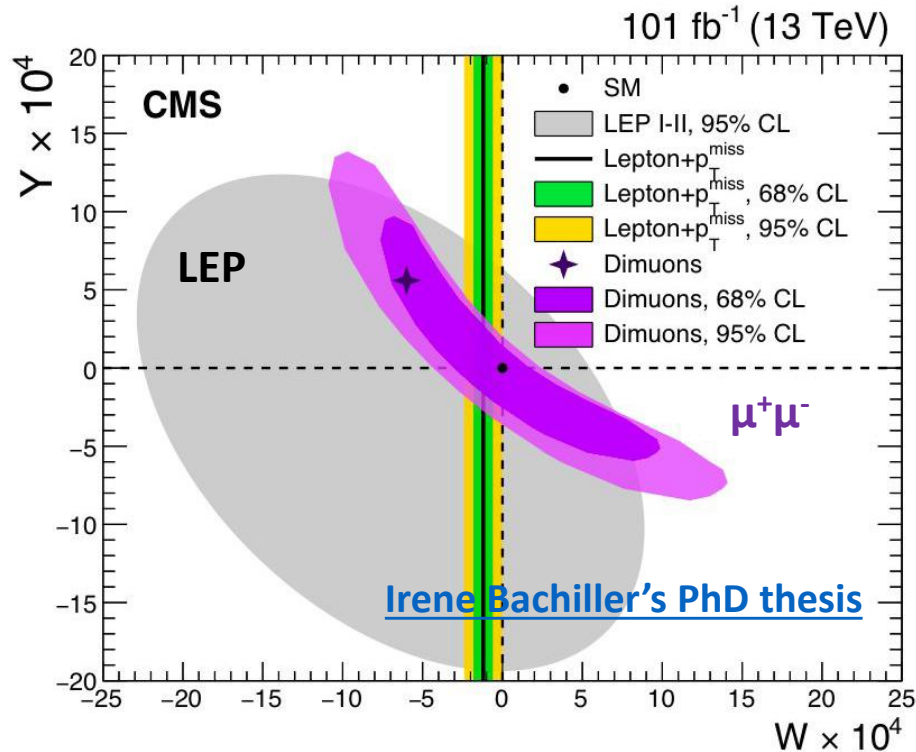
$$\text{Weight}(q^2, z, W, Y) = \frac{\left[\frac{d\sigma}{dz}(q^2, z, W, Y) \right]}{\left[\frac{d\sigma^0}{dz}(q^2, z) \right]}$$

With $z = \cos\theta$, angle of final state l wrt incoming q in $qq\bar{q}$ rest frame

- l^+l^- depends on both $W, Y \rightarrow$ can constrain both parameters
- Info needed at generator level is q^2 (ll inv. mass), $\cos\theta$ (could be replaced by CS angle), flavour of incoming q & $q\bar{q}$.



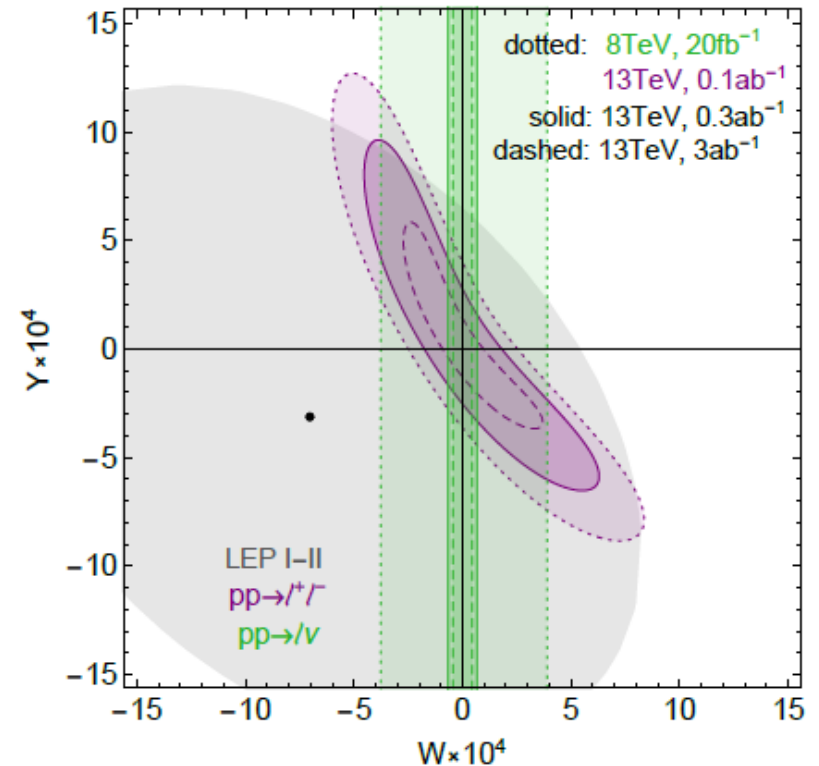
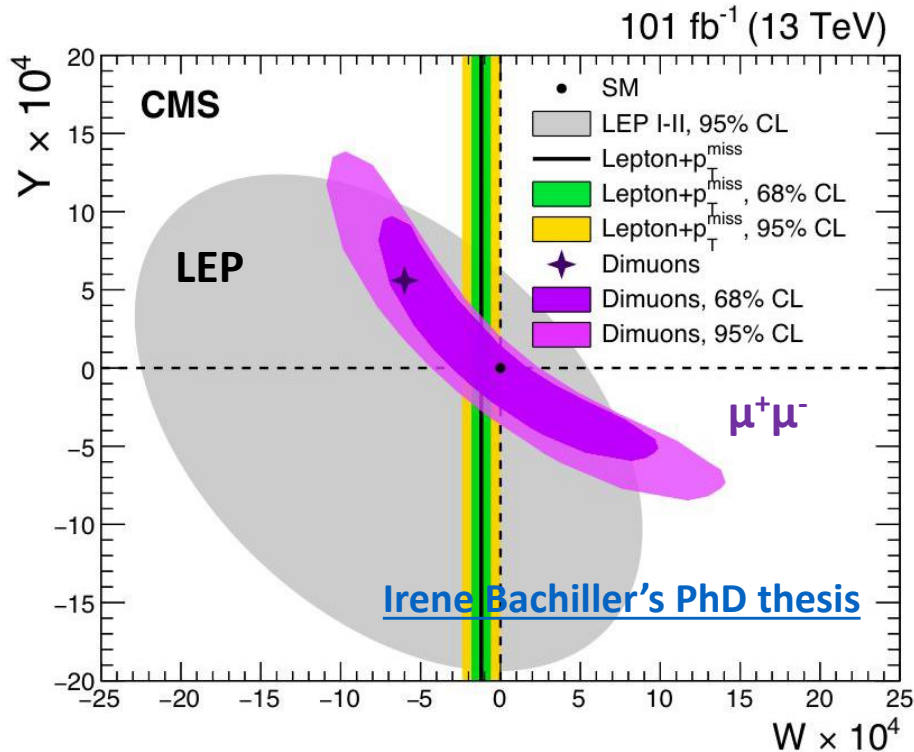
Constraints on W, Y using 101 fb⁻¹ LHC data



W oblique parameter limits improve
~10-fold previous LEP constraints

Large potential of LHC data on Y
oblique parameter from l+l- final
state; less restrictive on W parameter
than lv

Constraints on W, Y using 101 fb⁻¹ LHC data



W oblique parameter limits improve
~10-fold previous LEP constraints

Large potential of LHC data on Y
oblique parameter from l+l- final
state; less restrictive on W parameter
than lv

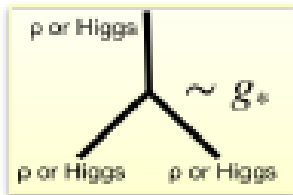
Projection by A. Wulzer et al using
Run 1 & 2 ATLAS and CMS data
<https://arxiv.org/abs/1609.08157>

Higgs Compositeness

Constraints on W & Y oblique parameters can be used to set bounds on several types of new models, among them Higgs Compositeness.

Higgs boson as bound state (eg new strong interaction) → new set of composite resonances (W', Z',...)

- new layer of complexity, m^*
- coupling to SM particles (g^*) or resultant of integrating out the new states (g/g^*)



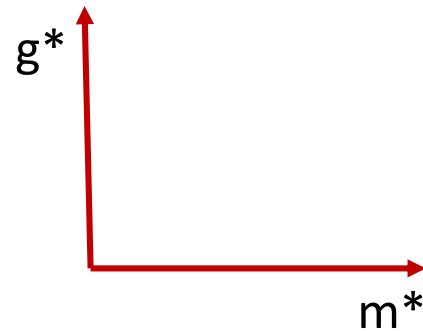
$$\text{gauge } \rho \sim \frac{g_W}{g^*}$$

4f, W'/Z'

$$\frac{c_{2W}}{\Lambda^2} = \frac{1}{g_*^2 m_*^2}$$

H

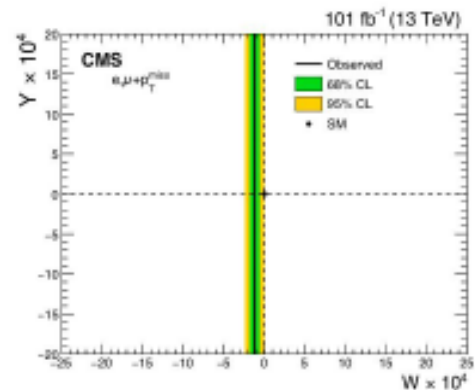
$$\frac{c_\phi}{\Lambda^2} = \frac{g_*^2}{m_*^2}$$



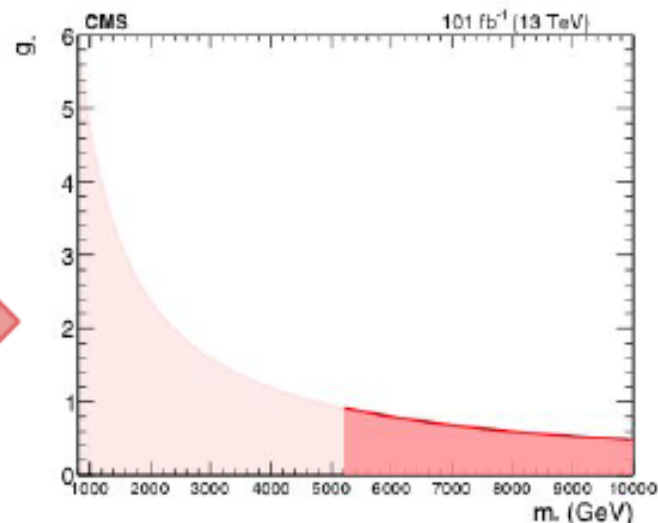
Indirect bound from W oblique parameter

$$g_*^2 = \frac{g_W^2 M_W^2}{W m_*^2}$$

$W = [-1.2 \pm 1.2] \times 10^{-4}$ 95 % CL

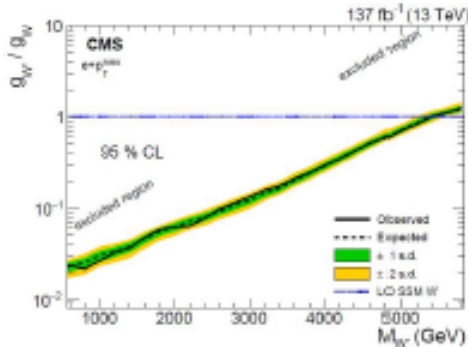


$$g_* < (4770 \text{ GeV}) / m_*$$



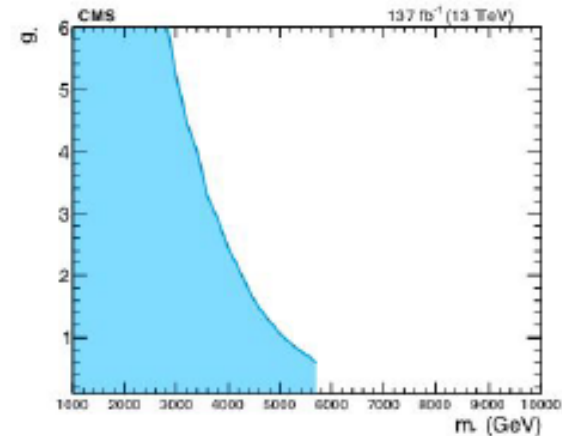
Direct bound from search for W'

SSM: W' decay only to fermions

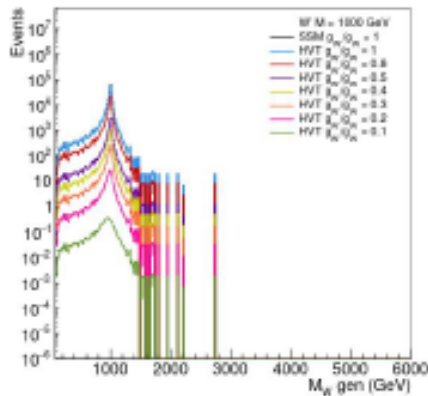


Limit set by the **W' coupling strength limit** with $\Gamma(W')/m(W') < 5\%$ narrow resonance assumption.

$$g_{W'} = \frac{g_W^2}{g_*} \quad g_*^2 < \frac{m_W^2 g^2}{m_*^2 W}$$

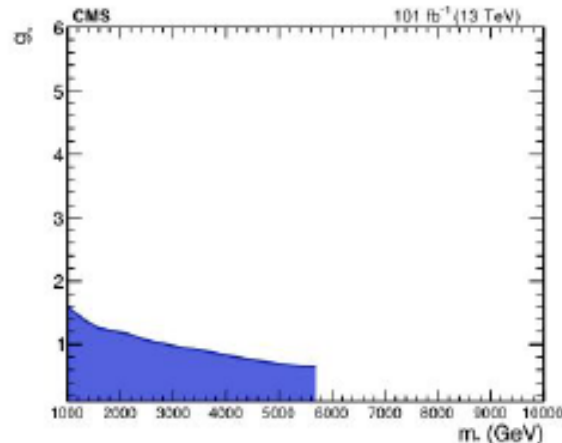


HVT: W' decay to fermions+bosons



Limit set by the W' coupling strength from HVT model where widths and cross sections are consistent with **bosonic and fermionic** couplings.

$$g_{W'} = \frac{g_W^2}{g_*} \quad \Gamma_{HVT} = \Gamma_{SSM} + \frac{M_{W'} g_*^2}{48\pi}$$



Limited by \sqrt{s}

Bounds on Higgs Compositeness

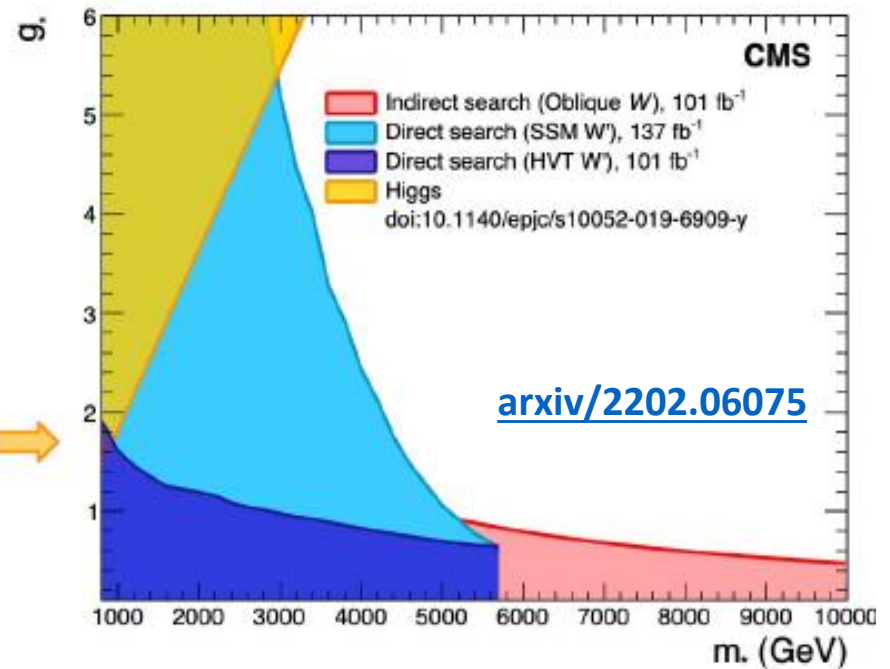
Using Higgs cross section strength, μ , measurement at CMS

$$\Delta\mu = \frac{g_*^2 v^2}{m_*^2}$$

$\mu = \text{CMS measured Higgs boson } \sigma / \text{SM expectation } \sigma = 1.17 \pm 0.10$

doi: 10.1140/epjc/s10052-019-6909-y

$\Delta\mu < 0.20$ at 95 % CL $\Rightarrow g_* < (1.82 \cdot 10^{-3} \text{ GeV}^{-1}) m_* \Rightarrow$



Higgs Compositeness excluded below 1 TeV (HVT) or 3 TeV (SSM)

Thank you

