Supernovae From Red Super Giant Stars

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- 1 Supernova Explosions
- 2 Numerical Setup

3 Results

Model HS19.8 Model HS25.5 Model HS14.2

4 Conclusion

- In the last stages of its life, a massive star consists of concentric shells, with a core composed mostly by iron and lighter elements all around.
- The collapse of the inner core of a massive star forming a neutron star (NS) releases a huge amount of gravitational binding energy

$$\Delta E_b \approx \frac{GM_{core}^2}{R} = 3 \times 10^{53} \left(\frac{M_{core}}{M_{\odot}}\right) \left(\frac{R}{10 \, km}\right)^{-1} erg.$$
(1)

• Only a few percent of this energy needs to be trapped to power the outflow of a supernova (SN) explosion.

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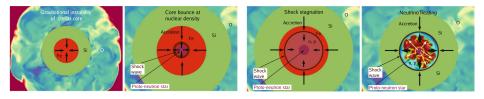
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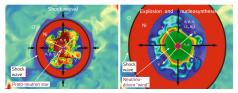
Solution

Neutrinos coming from the cooling proto-neutron star revive the stalled shock (Colgate & White 1966, Wilson 1985).

This idea was confirmed by simulations first (Wilson 1985), and by the observation of SN1987A later.

Phases of the Explosion (Wongwathanarat et al. 2013, Müller et al. 2016)





- PROMETHEUS-HOTB (Janka & Müller 1996, Gabler et al. 2021)
- Models: 7 red super giant (RSG) stars, with $12.5M_{\odot} \le M \le 27M_{\odot}$.
- Grid: $400(r) \times 45(\theta) \times 135(\phi) \times 2$ zones.
- $\Delta \theta = \Delta \phi = 2^{\circ}$.
- rib = 3.5e6 cm.
- gridlx = 4.0e10 cm.
- $t_f = 2.5s$.
- Expansion in the stellar wind of the progenitor star

$$\rho_e = \rho_0 \left(\frac{R_{prog}}{r}\right)^2, \quad T_e = T_0 \left(\frac{R_{prog}}{r}\right)^2.$$
(2)

Compactness (O'Connor & Ott 2011):

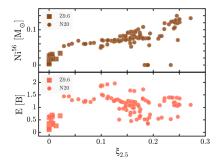
$$\xi_M = \left. \frac{M/M_{\odot}}{R(M_{bary} = M)/1000 \text{ km}} \right|_{t=t_b}$$
(3)

100% accuracy if a two-parameter description is used (Ertl et al. 2015). *Compactness* (O'Connor & Ott 2011):

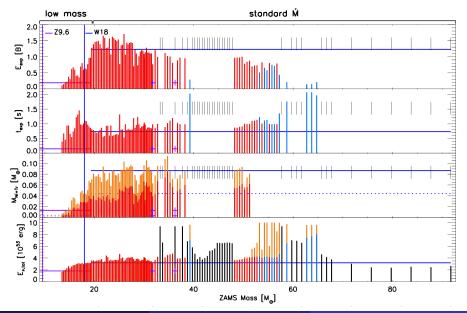
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Figure: Figure 16 from Sukhbold et al. 2016



Explsion Rate (Ertl et al. 2020)

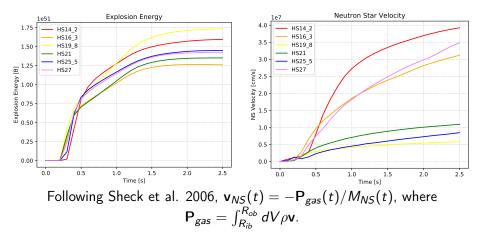


We identified 7 models that can lead to a SN explosion:

| Model | t _{exp} [ms] | E_{exp} [B] | $M_{Ni} [M_{\odot}]$ | $M_{ns} [M_{\odot}]$ | <i>v_{ns}</i> [km/s] |
|--------|-----------------------|---------------|----------------------|----------------------|------------------------------|
| W12.5 | Currently running | | | | |
| HS14.2 | 285 | 1.59 | 0.121 | 1.52 | 392.28 |
| HS16.3 | 207 | 1.26 | 0.097 | 1.35 | 311.72 |
| HS19.8 | 193 | 1.73 | 0.077 | 1.36 | 58.27 |
| HS21 | 190 | 1.35 | 0.093 | 1.35 | 109.14 |
| HS25.5 | 250 | 1.45 | 0.107 | 1.60 | 84.68 |
| HS27 | 203 | 1.42 | 0.110 | 1.46 | 348.26 |

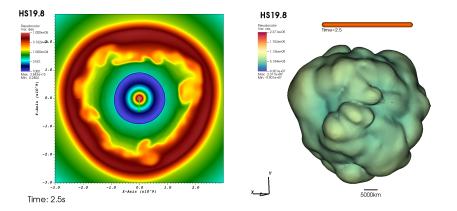
Calibration of E_{exp} and ⁵⁶Ni mass on 1D models by Sukhbold et al. 2016 (calibration models: SN1054 and SN1987A).

Explosion Energy and Remnant Kick

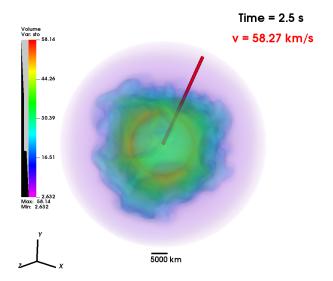


Model HS19.8

Density and ⁵⁶Ni distribution

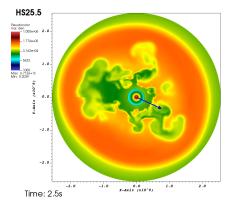


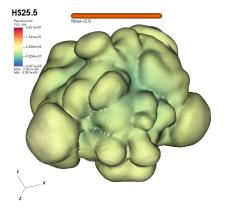
Volume plot of the entropy



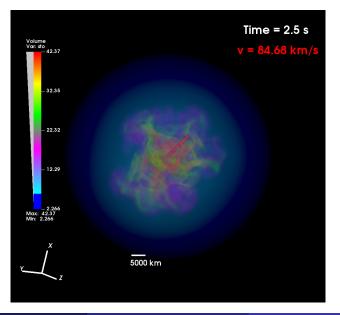
Model HS25.5

Density and ⁵⁶Ni distribution



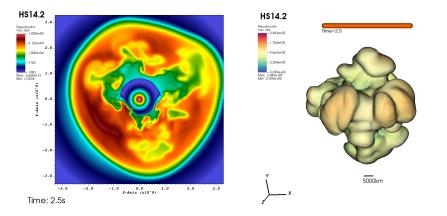


Volume plot of the entropy

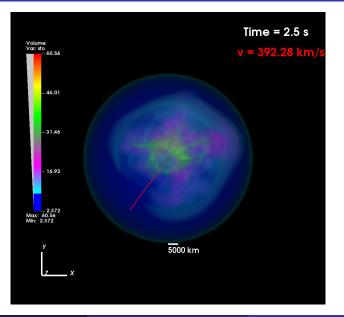


Model HS14.2

Density and ⁵⁶Ni distribution



Volume plot of the entropy



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- Supernova explosions are affected by hydrodynamical instabilities from very early stages, and this cause a deviation from spherical symmetry.
- The HD kick velocity does not contradict observations.

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THANK YOU!!!

Yin-Yang Grid (Kageyama & Sato 2004)

$$\theta = \left[\frac{\pi}{4} - \delta, \frac{3}{4}\pi + \delta\right] \cap \phi = \left[-\frac{3}{4}\pi - \delta, \frac{3}{4}\pi + \delta\right], \quad \delta \equiv \Delta \theta = \Delta \phi$$

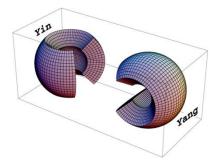


Figure: Figure 3 from Kageyama & Sato 2004.

- Eases the severe restriction of the time step size.
- Easy and straightforward implementation in a code that already uses spherical polar coordinates.
- It eliminates flaws near coordinate symmetry axis of the polar grid → No evidence of preferred radial direction.
- Reduced computational cost.

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