

# Symmetries and the nature of the Higgs.

G.G.Ross, CERN, May 2011



# SUSY extensions of the Standard Model

$$\begin{aligned} W = & h^E L H_d \bar{E} + h^D Q H_d \bar{D} + h^U Q H_u \bar{U} + \mu H_d H_u \\ & + \lambda L L \bar{E} + \lambda' L Q \bar{D} + \kappa L H_u + \lambda'' \bar{U} \bar{D} \bar{D} \\ & + \frac{1}{M} (Q Q Q L + Q Q Q H_d + Q \bar{U} \bar{E} H_d + \dots (\cancel{L})) \end{aligned}$$

R-parity:       $Z_2$        $H_u, H_d$   $+1$       SUSY states odd  
                         $L, \bar{E}, Q, \bar{D}, \bar{U}, \theta$   $-1$       Weinberg, Sakai

# SUSY extensions of the Standard Model

$$\begin{aligned} W = & h^E L H_d \bar{E} + h^D Q H_d \bar{D} + h^U Q H_u \bar{U} + \mu H_d H_u \\ & + \lambda L \bar{L} \bar{E} + \lambda' L \bar{Q} \bar{D} + \kappa L H_u + \lambda'' \bar{U} \bar{D} \bar{D} \\ & + \frac{1}{M} (QQQL + QQQH_d + Q\bar{U}\bar{E}H_d + \dots(\mathcal{L})) \end{aligned}$$

R-parity:  $Z_2$  SUSY states odd

Weinberg, Sakai

Baryon "parity":  $Z_3$  LSP unstable

$$\begin{array}{ll} \frac{Q}{D}, H_u & 1 \alpha \\ L, \bar{E}, \bar{U}, H_d & \alpha^2 \end{array}$$

Discrete gauge symmetry  
-anomaly free

Ibanez, GGR

# SUSY extensions of the Standard Model

$$W = h^E L H_d \bar{E} + h^D Q H_d \bar{D} + h^U Q H_u \bar{U} + \mu H_d H_u \\ + \lambda L L \bar{E} + \lambda' L Q \bar{D} + \kappa L H_u + \lambda'' \bar{U} \bar{D} \bar{D} \\ + \frac{1}{M} (Q Q Q L + Q Q Q H_d + Q \bar{U} \bar{E} H_d + \dots (\mathcal{L}))$$

$\mu$  term,  
GUTs?

R-parity:  $Z_2$  SUSY states odd

Baryon "parity":  $Z_3$  LSP unstable

Proton hexality:  $Z_6 = Z_2^R \times Z_3^B$  LSP stable

$$\frac{1}{M} L L H_u H_u$$

Dreiner, Luhn, Thormeier

# SUSY extensions of the Standard Model

$$\begin{aligned} W = & h^E L H_d \bar{E} + h^D Q H_d \bar{D} + h^U Q H_u \bar{U} + \mu H_d H_u \\ & + \lambda L L \bar{E} + \lambda' L Q \bar{D} + \kappa L H_u + \lambda'' \bar{U} \bar{D} \bar{D} \\ & + \frac{1}{M} (Q Q Q L + Q Q Q H_d + Q \bar{U} \bar{E} H_d + \dots (\mathcal{L})) \end{aligned}$$

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$Z_N^R$  R-symmetry N=4,6,8,12,24 LSP stable  
 $\frac{1}{M} L H_u H_u$

# A unique solution : discrete R symmetry

c.f.Hyun Min Lee's talk

MSSM spectrum

No perturbative  $\mu$  term

Commutes with  $SO(10)$

Anomaly cancellation

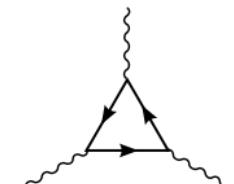
$$A_{G-G-\mathbb{Z}_N} = \rho \bmod \eta$$

Green Schwarz term  
 N, N odd  
 N/2, N even

$$A_{SU(3)-SU(3)-\mathbb{Z}_N} = \frac{1}{2} \sum_i [3 \cdot q_{10_i} + q_{\bar{5}_i} - 4R] + 3R$$

$$A_{SU(2)-SU(2)-\mathbb{Z}_N} = \frac{1}{2} \sum_i [3 \cdot q_{10_i} + q_{\bar{5}_i} - 4R] + 2R + \frac{1}{2} (q_H + q_{\bar{H}} - 2R)$$

$$A_{U(1)_Y-U(1)_Y-\mathbb{Z}_N^R} = \frac{1}{2} \sum_{g=1}^3 (3q_{10}^g + q_{\bar{5}}^g) + \frac{3}{5} \left[ \frac{1}{2} (q_{H_u} + q_{H_d}) - 11 \right] \quad (R=1)$$



No D=4  $\cancel{\beta}, \cancel{\gamma}$

$$A_3 - A_2 \Rightarrow (q_H + q_{\bar{H}}) = 4R \bmod 2\eta$$

$$A_3 - A_1 \Rightarrow \frac{3}{5} \left[ \frac{1}{2} (q_{H_u} + q_{H_d}) - 6 \right] = 0 \bmod \eta$$

$$\Rightarrow N = 3, 4, 6, 8, 12, 24$$

$M$	$q_{10}$	$q_{\bar{5}}$	$q_{H_u}$	$q_{H_d}$	$q_{H_u}^{sh}$	$q_{H_d}^{sh}$	$\rho$
4	1	1	0	0	16	16	1
6	5	3	4	0	28	24	0
8	1	5	0	4	24	28	1
12	5	9	4	0	28	24	3
24	5	9	16	12	88	84	9

$Z_{4R}$ 

# Phenomenology

MSSM spectrum

No perturbative  $\mu$  termCommutes with  $SO(10)$ 

Anomaly cancellation

$M$	$q_{10}$	$q_{\bar{5}}$	$q_{H_u}$	$q_{H_d}$	$q_{H_u}^{\text{sh}}$	$q_{H_d}^{\text{sh}}$	$\rho$
4	1	1	0	0	16	16	1
6	5	3	4	0	28	24	0
8	1	5	0	4	24	28	1
12	5	9	4	0	28	24	3
24	5	9	16	12	88	84	9

D=5 operators

up and down Yukawas allowed

$$3q_{10} + q_{\bar{5}} + q_{H_u} + q_{H_d} = 4 \pmod{M} \Rightarrow 3q_{10} + q_{\bar{5}} = 0 \pmod{M} \Rightarrow \frac{1}{M} Q \cancel{QQQL} - \frac{1}{M} LLH_u H_u$$

Weinberg operator

SUSY breaking $\langle W \rangle, \langle \lambda \lambda \rangle$  R=2 non-perturbative breaking $Z_{4R} \rightarrow Z_2^R$  R-parity

Domain walls safe

 $\mu \sim m_{3/2}, O(\frac{m_{3/2}}{M^2} QQQL)$  $M_{\text{higgs}} \approx M_{\text{SUSY}}$ 

Residual fine tuning?

 $\mu, \beta, \mathcal{L}$

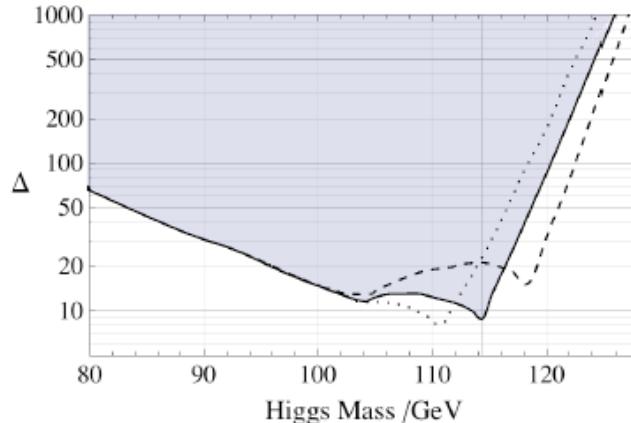
# The CMSSM

$$\lambda v^2 = -\mu^2 + 0.34(m_{\tilde{Q}}^2 + m_{\tilde{U}}^2) + 1.86\tilde{M}_3^2 - 0.22\tilde{M}_2^2 + \dots$$

SPS1a

$$\Delta \equiv \max |\Delta_p|_{p=\{\mu_0^2, m_0^2, m_{1/2}^2, A_0^2, B_0^2\}}, \quad \Delta_p \equiv \frac{\partial \ln v^2}{\partial \ln p}$$

Relic density unrestricted



SUSY particle masses

$$3.20 < 10^4 \text{ Br}(b \rightarrow s\gamma) < 3.84$$

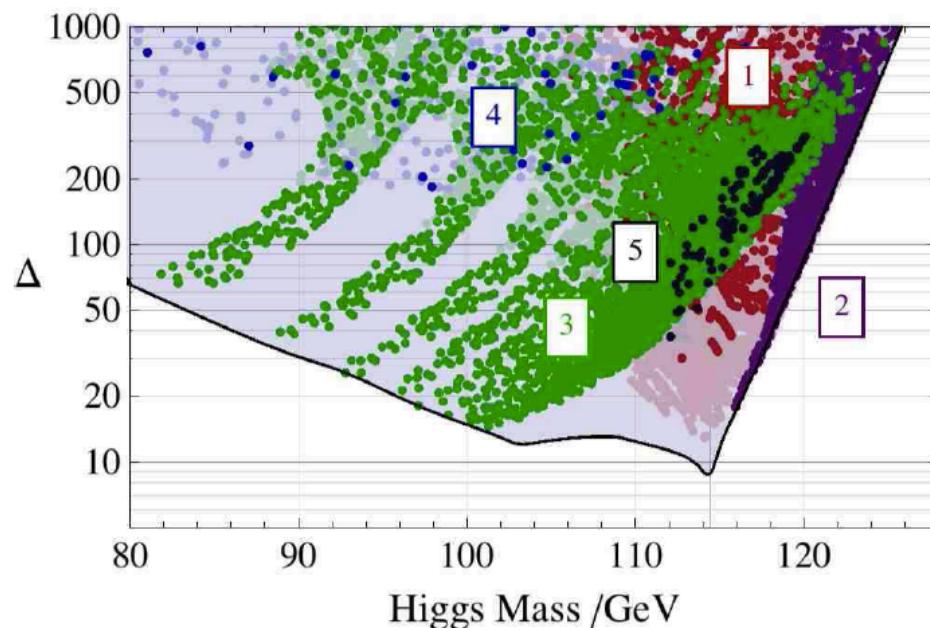
$$\text{Br}(b \rightarrow \mu\mu) < 1.8 \times 10^{-8}$$

$$\delta a_\mu < 292 \times 10^{-11}$$

$$-0.0007 < \delta\rho < 0.0012$$

$$\Delta_{Min} = 9, \quad m_h = 114 \pm 2 \text{ GeV}$$

Relic density restricted



1  $h^0$  resonant annihilation

2  $\tilde{h}$  t-channel exchange

3  $\tilde{\tau}$  co-annihilation

4  $\tilde{t}$  co-annihilation

• 5  $A^0 / H^0$  resonant annihilation

Within  $3\sigma$  WMAP:

$$\Delta_{Min} = 15, \quad m_h = 114.7 \pm 2 \text{ GeV}$$

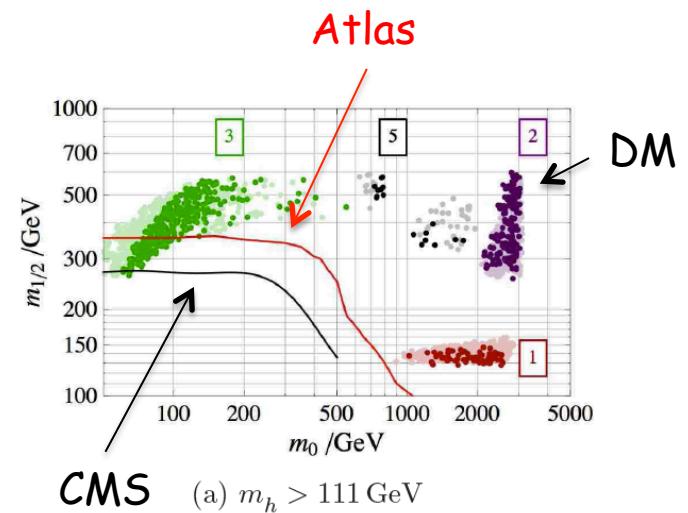
<  $3\sigma$  WMAP:

$$\Delta_{Min} = 18, \quad m_h = 115.9 \pm 2 \text{ GeV}$$

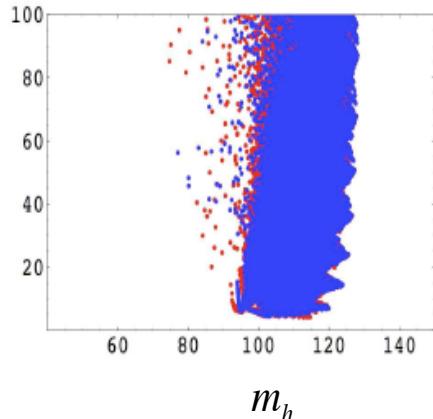
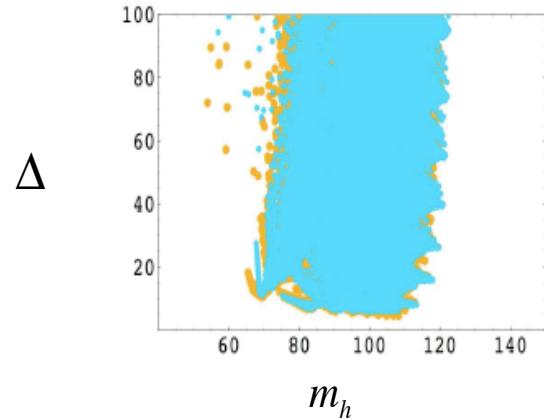
Cassel, Ghilencea, Kraml, Lessa, GGR

# LHC - Regions of low fine tuning $\Delta < 100$ :

	SUG0	SUG1	SUG2	SUG3	SUG5
$m_0$	1455	1508	2270	113	725
$m_{1/2}$	160	135	329	383	535
$A_0$	238	1492	30	-220	1138
$\tan \beta$	22.5	22.5	35	15	50
$\mu$	191	433	187	529	581
$m_{\tilde{g}}$	482	414	900	898	1252
$m_{\tilde{u}_L}$	1469	1509	2331	826	1315
$m_{\tilde{t}_1}$	876	831	1423	602	1000
$m_{\tilde{\chi}_1^+}$	106	104	168	293	416
$m_{\tilde{\chi}_2^0}$	108	104	181	293	416
$m_{\tilde{\chi}_1^0}$	60	53	123	155	222
$\Delta$	9	50	45	68	84
$\Omega_{\tilde{\chi}_1^0} h^2$	0.41	0.13	0.10	0.13	0.10
$\text{BR}(b \rightarrow s\gamma) \times 10^4$	3.4	3.7	3.4	3.2	3.2
$\text{BR}(B_s \rightarrow \mu^+ \mu^-) \times 10^9$	3.0	2.9	2.9	3.4	1.7
$\delta a_\mu \times 10^{10}$	4.5	3.2	3.2	22.5	16.6
$\sigma_{\chi p}^{\text{SI}} (\text{pb}) \times 10^{10}$	108	5	432	24	101
$\sigma^{(LO)}(7 \text{ TeV}) (\text{pb})$	8	12	0.9	0.4	0.02
$\sigma^{(LO)}(14 \text{ TeV}) (\text{pb})$	40	75	3	5	0.4



## Reduced fine tuning - singlet extensions



$$W_{\text{eff}}^{\text{SMSSM}} = (H_u H_d)^2 / \mu_s + \mu H_u H_d$$

$$\xi_1 2 \frac{\mu_0}{M_*} (|h_1|^2 + |h_2|^2) h_1 h_2 \quad \xrightarrow{v^2 = -\frac{m^2}{\lambda}}$$

Cassel, Ghilencea, GGR  
 Dine, Seiberg, Thomas  
 Bastero-Gill, Hugonie, King,  
 Roy, Vempati  
 Delgado, Kolda, Olson, Puente

$$W = W_{\text{Yukawa}} + (\mu + \lambda S) H_u H_d + \frac{\mu_S}{2} S^2 + \frac{\kappa}{3} S^3 + \xi S \quad \text{GNMSSM}$$

$$\text{c.f. } W = W_{\text{Yukawa}} + \lambda S H_u H_d + \frac{\kappa}{3} S^3 \quad \text{NMSSM}$$

GNMSSM (NMSSM) unnatural .....discrete R-symmetry.....

## Singlet extensions with R-symmetry

$$W = W_{NMSSM} + \Delta W$$

$$\begin{aligned}\Delta \mathcal{W}_{\mathbb{Z}_4^R} &= Y + Y^2 N + Y N^2 + Y H_u H_d \\ &\sim m_{3/2} M_P^2 + m_{3/2}^2 N + m_{3/2} N^2 + m_{3/2} H_u H_d\end{aligned}$$

$M$	$q_{\mathbf{10}}$	$q_{\overline{\mathbf{5}}}$	$q_{H_u}$	$q_{H_d}$	$q_N$
4	1	1	0	0	2

$$\begin{aligned}\Delta \mathcal{W}_{\mathbb{Z}_8^R} &= Y + Y^2 (N + Y N^2 + Y H_u H_d) \\ &\sim m_{3/2} M_P^2 + m_{3/2}^2 N + \frac{m_{3/2}^3}{M_P^2} N^2 + \frac{m_{3/2}^3}{M_P^2} H_u H_d\end{aligned}$$

$M$	$q_{\mathbf{10}}$	$q_{\overline{\mathbf{5}}}$	$q_{H_u}$	$q_{H_d}$	$q_N$
8	1	5	0	4	6

Lee, Raby, Ratz, Ross, Schieren, Schmidt-Hoberg, Vaudrevange

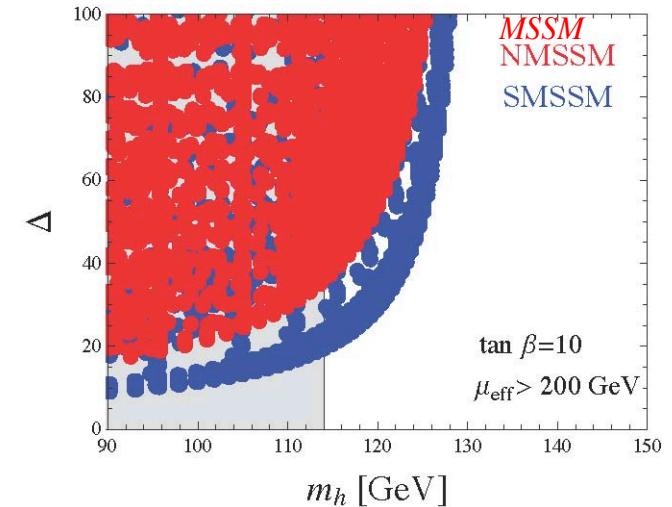
R-symmetry guarantees tadpoles under control

Abel

## General-NMSSM phenomenology

$$W = W_{\text{Yukawa}} + (\mu + \lambda S) H_u H_d + \frac{\mu_S}{2} S^2 + \frac{\kappa}{3} S^3 + \xi S \quad \text{GNMSSM}$$

$$\text{c.f. } W = W_{\text{Yukawa}} + \lambda S H_u H_d + \frac{\kappa}{3} S^3 \quad \text{NMSSM}$$



GGF, Schmidt-Hoberg (preliminary)

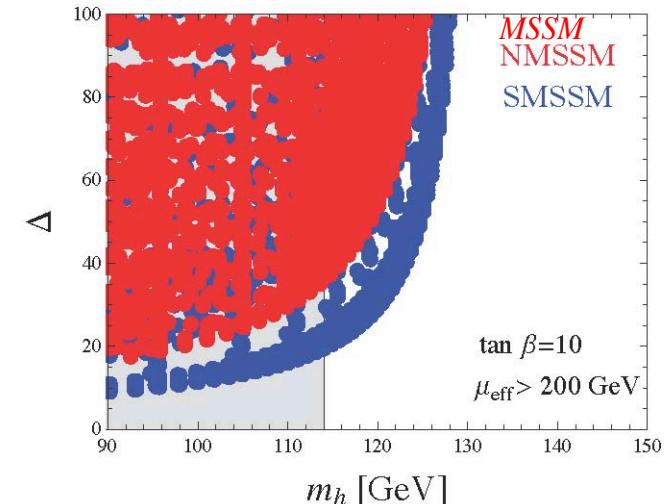
## SUSY phenomenology

$$\mu_s \gg \mu \quad \text{MSSM SUSY structure}$$

# General-NMSSM phenomenology

$$W = W_{\text{Yukawa}} + (\mu + \lambda S) H_u H_d + \frac{\mu_S}{2} S^2 + \frac{\kappa}{3} S^3 + \xi S \quad \text{GNMSSM}$$

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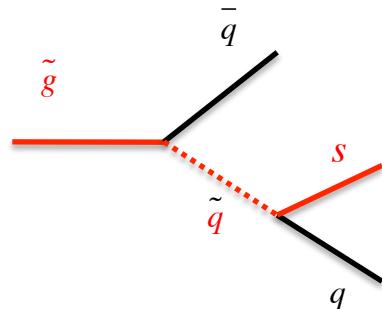


GGR, Schmidt-Hoberg (preliminary)

# SUSY phenomenology

$$\mu_s \gg \mu \quad \text{MSSM SUSY structure}$$

$$\mu_s \sim \mu \quad \tilde{S} \text{ can be LSP}$$



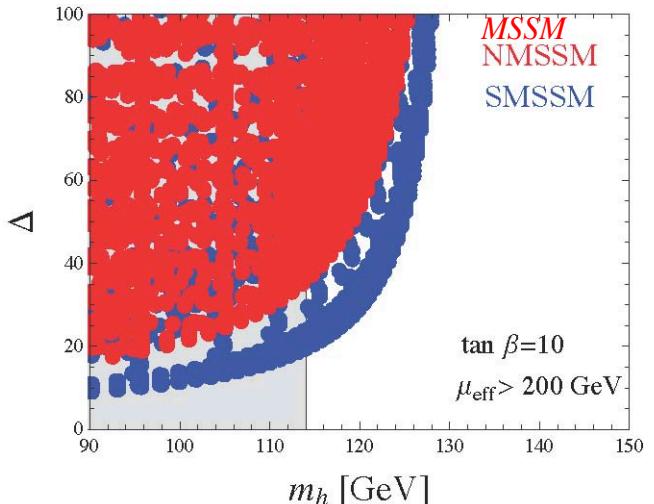
$$\left( \tilde{B}, \widetilde{W}^0, \widetilde{H}_d^0, \widetilde{H}_u^0, \tilde{S} \right)$$

$$M_{\tilde{\chi}^0} = \begin{pmatrix} M_1 & 0 & -M_Z s_W c_\beta & M_Z s_W s_\beta & 0 \\ 0 & M_2 & M_Z c_W c_\beta & -M_Z c_W s_\beta & 0 \\ -M_Z s_W c_\beta & M_Z c_W c_\beta & 0 & -\mu & -\lambda v s_\beta \\ M_Z s_W s_\beta & -M_Z c_W s_\beta & -\mu & 0 & -\lambda v c_\beta \\ 0 & 0 & -\lambda v s_\beta & -\lambda v c_\beta & \mu_s \end{pmatrix}$$

# General-NMSSM phenomenology

$$W = W_{\text{Yukawa}} + (\mu + \lambda S) H_u H_d + \frac{\mu_S}{2} S^2 + \frac{\kappa}{3} S^3 + \xi S$$

$$V_{\text{soft}} = m_s^2 |s|^2 + m_{h_u}^2 |h_u|^2 + m_{h_d}^2 |h_d|^2 + \left( \lambda A_\lambda s h_u h_d + \frac{1}{3} \kappa A_\kappa s^3 + \frac{1}{2} b_s s^2 + h.c. \right)$$



## Higgs structure

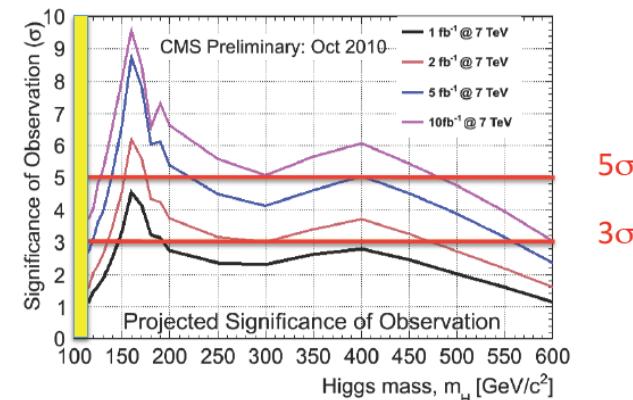
$$(h_u, h_d, s)$$

GGR, Schmidt-Hoberg (preliminary)

$$\mu_s \gg \mu$$

MSSM Higgs structure

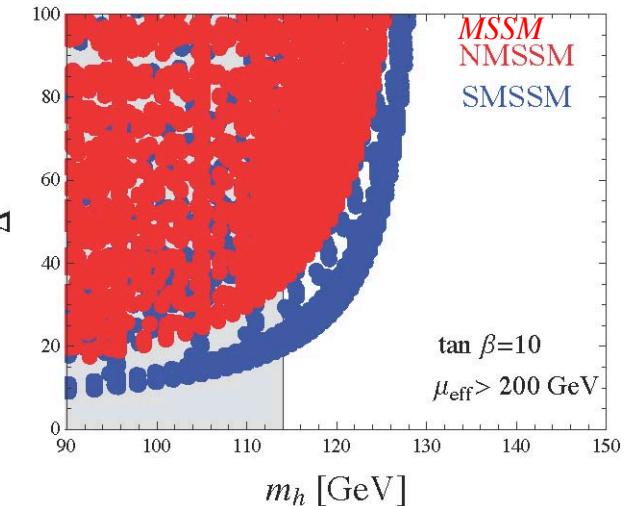
$\Delta$  significantly reduced even for  $\mu_s \sim 5 \text{ TeV}$



# General-NMSSM phenomenology

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## Higgs structure

$$(h_u, h_d, s)$$

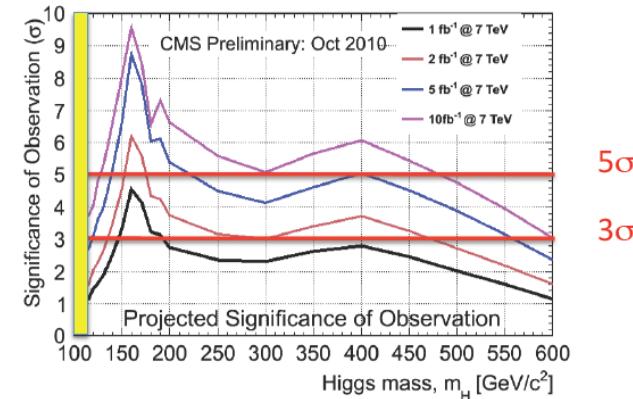
GGR, Schmidt-Hoberg (preliminary)

$$\mu_s \gg \mu$$

MSSM SUSY structure

$$\mu_s, m_s, b_s \sim \mu$$

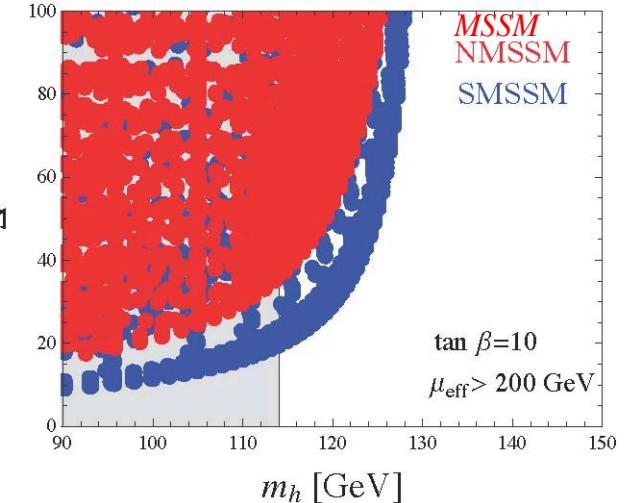
h, s mixing, reduced production cross section



# General-NMSSM phenomenology

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## Higgs structure

$$(h_u, h_d, s)$$

GGR, Schmidt-Hoberg (preliminary)

$$\mu_s \gg \mu$$

MSSM SUSY structure

$$\mu_s, m_s, b_s \sim \mu$$

$$h_1 = H_u + \varepsilon S, \quad h_2 = S - \varepsilon H_u$$

$$M_{h_1} \ll M_{h_2}, \quad \varepsilon \ll 1$$

$$BR\left(\frac{h_1 \rightarrow \tilde{S}\tilde{S}}{h_1 \rightarrow \gamma\gamma, b\bar{b}}\right) \gg 1$$

## Invisible Higgs decay

# LEP: invisible Higgs searches

$$e^+ e^- \rightarrow h^0 Z$$

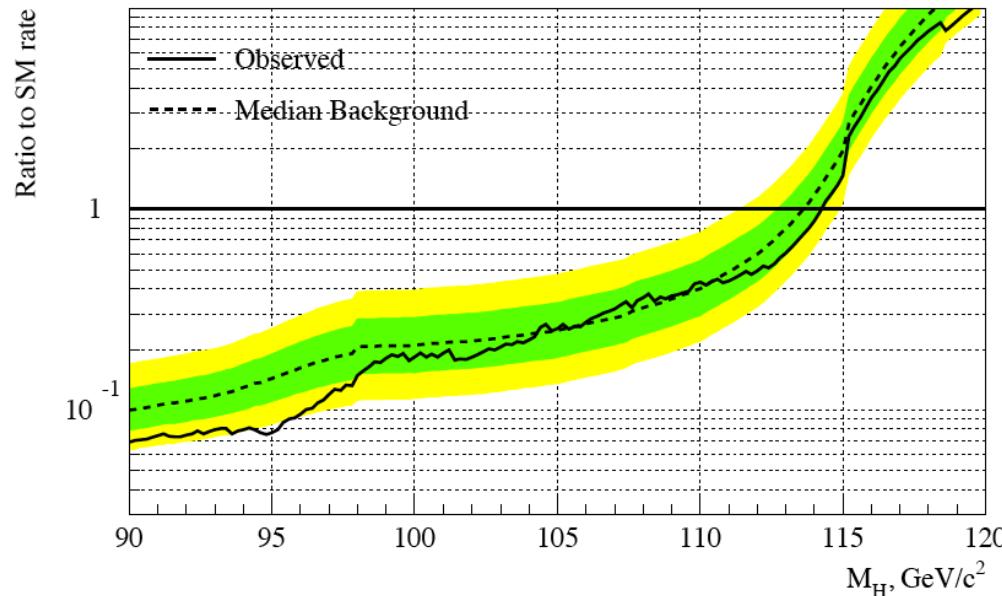
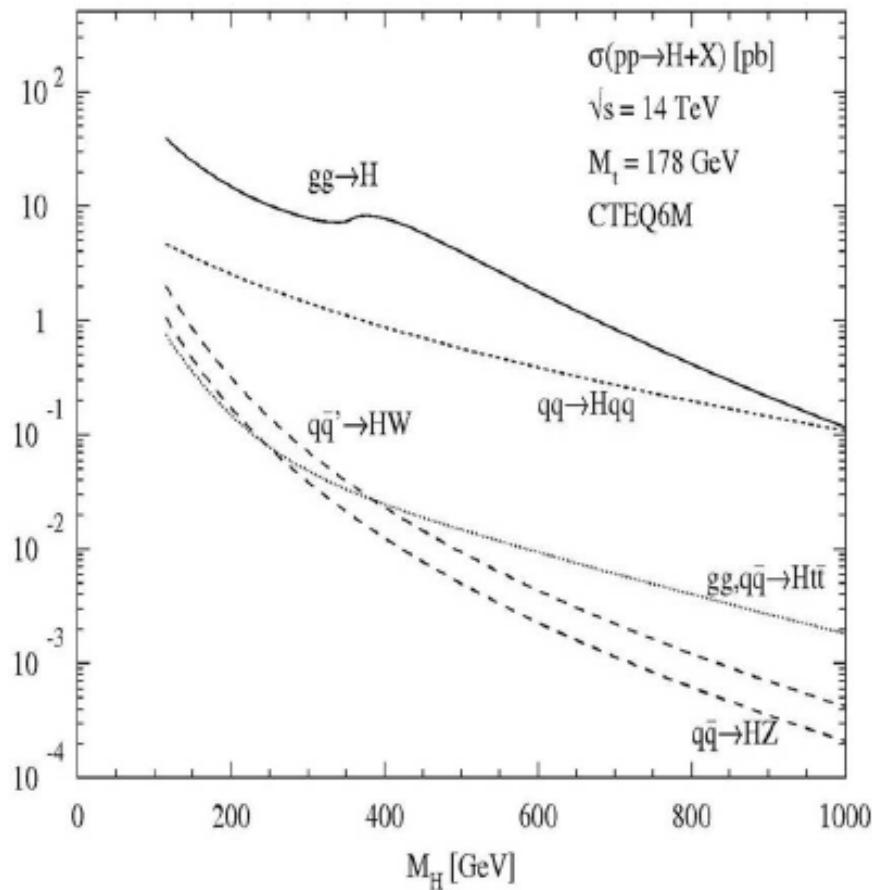
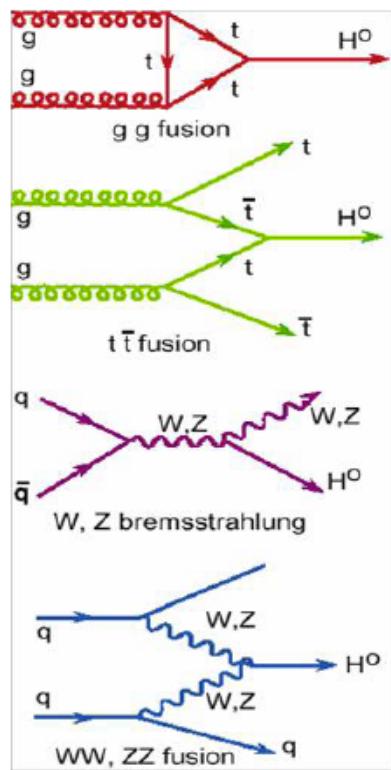


Figure 4: The region excluded by the combined LEP results in the  $h^0 \rightarrow \text{invisible}$  search. The 95% CL upper limit on,  $\xi^2$ , the production rate as a fraction of the Standard Model total rate, is shown, together with the expected range assuming there is no signal.

Aleph, Delphi, L3 and Opal combined

## LHC:



$gg \rightarrow h^0$       forward-backward jets

$q\bar{q} \rightarrow h^0 Z$        $\times 10^{-1}$

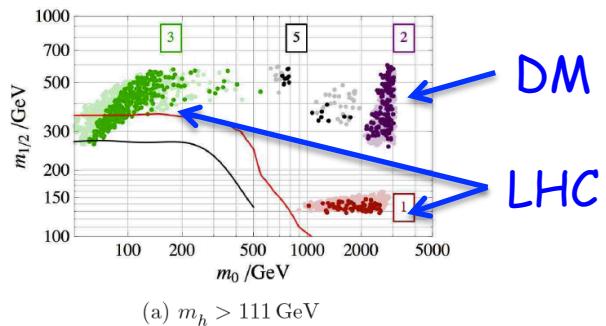
....no invisible Higgs LHC sensitivity reported yet

# Summary

- Hierarchy problem  $\Rightarrow$  Low-energy SUSY + further symmetry

$$Z_{4R} : \mu, \beta, \lambda$$

- CMSSM Complementary DM & LHC searches



$\Delta \leq 100$  Sensitivity  $\times (10 - 100)$

$\Delta \leq 100$   $\sim LHC 7TeV 1fb^{-1}$

(Full region  $LHC 14TeV 10fb^{-1}$ )

- NMSSM Reduced  $\Delta \Rightarrow GMSSM \Rightarrow Z_{4R}, Z_{8R}$   
SUSY states can be heavier

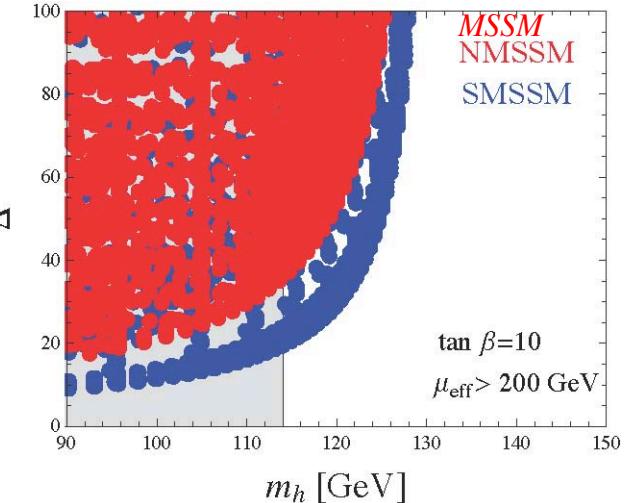
...Light Higgs search **may** provide the first crucial test!



# General-NMSSM phenomenology

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## Higgs structure

$$(h_u, h_d, s)$$

GGR, Schmidt-Hoberg (preliminary)

$$\mu_s \gg \mu$$

MSSM SUSY structure

$$\mu_s, m_s, b_s \sim \mu$$

$$h_1 = H_u + \varepsilon S, \quad h_2 = S - \varepsilon H_u$$

$$M_{h_1} \ll M_{h_2}, \quad \varepsilon \ll 1$$

$$BR\left(\frac{h_1 \rightarrow \tilde{S}\tilde{S}}{h_1 \rightarrow \gamma\gamma, bb}\right) \gg 1$$

## Invisible Higgs decay

$$M_{11}^2 = \cot(\beta)(\lambda(v_s(A_\lambda + v_s \kappa + \mu_s) + \xi) + b\mu) + M_Z^2 \sin^2(\beta)(1 + \delta)$$

$$M_{22}^2 = \tan(\beta)(\lambda(v_s(A_\lambda + v_s \kappa + \mu_s) + \xi) + b\mu) + M_Z^2 \cos^2(\beta)$$

$$M_{33}^2 = v_s \kappa (A_\kappa + 4v_s \kappa + 3\mu_s) + \frac{1}{2v_s} (v^2 \lambda (A_\lambda + \mu_s) \sin(2\beta) - 2v^2 \lambda \mu - 2\mu_s \xi)$$

$$M_{12}^2 = -\lambda(v_s(A_\lambda + v_s \kappa + \mu_s) + \xi) - b\mu + (v^2 \lambda^2 - M_Z^2/2) \sin(2\beta)$$

$$M_{13}^2 = v \lambda (2 \sin(\beta)(v_s \lambda + \mu) - \cos(\beta)(A_\lambda + 2v_s \kappa + \mu_s))$$

$$M_{23}^2 = v \lambda (2(v_s \lambda + \mu) \cos(\beta) - (A_\lambda + 2\kappa v_s + \mu_s) \sin(\beta))$$