

# Design of a fixed field accelerating ring for high power applications

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HB2023 at CERN

9 October 2023



ISIS Neutron and Muon Source

# Outline

- A bit of history
  - Midwestern Universities Research Association (MURA)
  - ASPUN at Argonne National Lab
  - Initial ESS project
- Why Fixed Field Accelerator?
  - High power by high repetition
  - Sustainable option
- Toward high power Fixed Field Accelerator
  - Constant tune
  - FD (DF) spiral sector
  - Superperiodicity
  - Physical and dynamic aperture
  - Collimation
  - Correction by trim coils
- Beam stacking experiment
- (Modelling space charge effects)
- Summary



# Fixed Field Accelerating Ring

What is Fixed Field Accelerating Ring?

- Cyclotron
- Synchrocyclotron
- Fixed-Field Alternating-Gradient (FFA, used to be FFAG)

Accelerators I am going to discuss have

- Main lattice with DC magnets
- Alternating gradient focusing
- Non isochronous, RF frequency is modulated
- Zero chromaticity, transverse tune is constant



I call it “FFA”

# A bit of history



ISIS Neutron and Muon Source

 [www.isis.stfc.ac.uk](http://www.isis.stfc.ac.uk)

  [@isisneutronmuon](https://www.instagram.com/isisneutronmuon)

 [uk.linkedin.com/showcase/isis-neutron-and-muon-source](https://uk.linkedin.com/showcase/isis-neutron-and-muon-source)

## O CAMELOT !

### A MEMOIR OF THE MURA YEARS\*

*F.T.Cole*

April 1, 1994

### 16.3 A New Single-Beam Proposal

We put together a new proposal with no colliding beams at all. We chose a proton energy of 10 GeV to be high enough above the antiproton production threshold to make usable intensities, but were constrained from going higher by concern about the total cost. We claimed we would reach a time-average intensity of 30 microamperes or  $2 \times 10^{14}$  protons per second, three orders of magnitude above what the synchrotrons were then doing (of course their higher energy took away some of that advantage in antiproton production). It was a spiral-sector ring

# A bit of history

## High current beam studies at MURA

Proceedings of the 2003 Particle Accelerator Conference

**MURA DAYS**

**Invited talk at PAC2003**

Keith R. Symon, University of Wisconsin-Madison, Madison, WI 53706, USA

**YEARS\***

- (i) beam stacking,
- (ii) Hamiltonian theory of longitudinal motion,
- (iii) useful colliding beams (the idea itself is quite old),
- (iv) storage rings (independently invented by O'Neill),
- (v) spiral-sector geometry used in isochronous cyclotrons,
- (vi) lattices with zero-dispersion and low- $\beta$  sections for colliding beams,
- (vii) multiturn injection into a strong-focusing lattice,
- (viii) first calculations of the effects of nonlinear forces in accelerators,
- (ix) first space-charge calculations including effects of the beam surroundings,
- (x) first experimental measurement of space-charge effects,
- (xi) theory of negative-mass and other collective instabilities and correction systems,
- (xii) the use of digital computation in design of orbits, magnets, and rf structures,
- (xiii) proof of the existence of chaos in digital computation, and
- (xiv) synchrotron-radiation rings

# A bit of history

## Spallation neutron source proposal

T. K. Khoe and R. L. Kustom  
 Physics Division  
 Argonne National Laboratory  
 9700 S. Cass Avenue  
 Argonne, IL 60439

## ASPUN, ANL

Keith R. Symon, University of W

### PAC1983

### 16.3 A New Single

- (i) beam stack
- (ii) Hamiltonia
- (iii) useful coll
- (iv) storage rin
- (v) spiral-secto
- (vi) lattices wit
- (vii) multium
- (viii) first ca
- (ix) first space
- (x) first experi
- (xi) theory of
- (xii) the use of

- (xiii) proof of t-Field Alternating
- (xiv) synchro

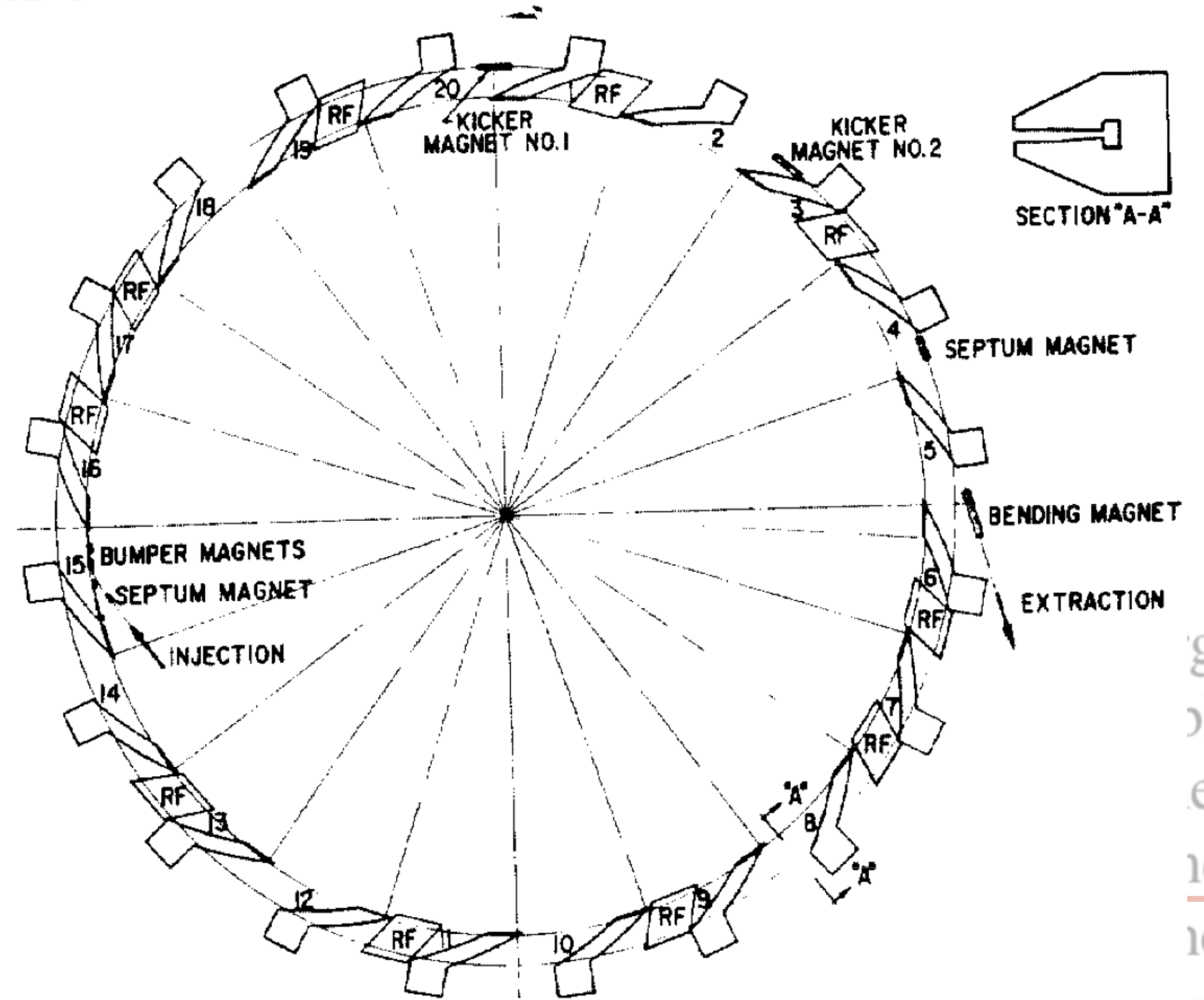


Fig. 1. Schematic View of FFAG Ring.

Table 1. FFAG Accelerator Characteristics

# A bit of history

## One of two options at (old) ESS Proposal

EPAC1992

T. V. Khoe and R. L. Kustom  
Physics Division  
Argonne National Laboratory  
9700 S. Cass Avenue

Proceedings of the 2003 Particle Accelerator Conference

### Accelerator Design Parameters for a European Pulsed Spallation Neutron Source.

Report from workshop for a European Spallation Source.

S Martin ( KFA, Jülich, Germany) and C W Planner ( RAL, UK)

Linac+Compressor ring

Linac+FFAG accelerator

Keith R

ASPUN, ANL at PAC1983

PAC2003

16.3 A

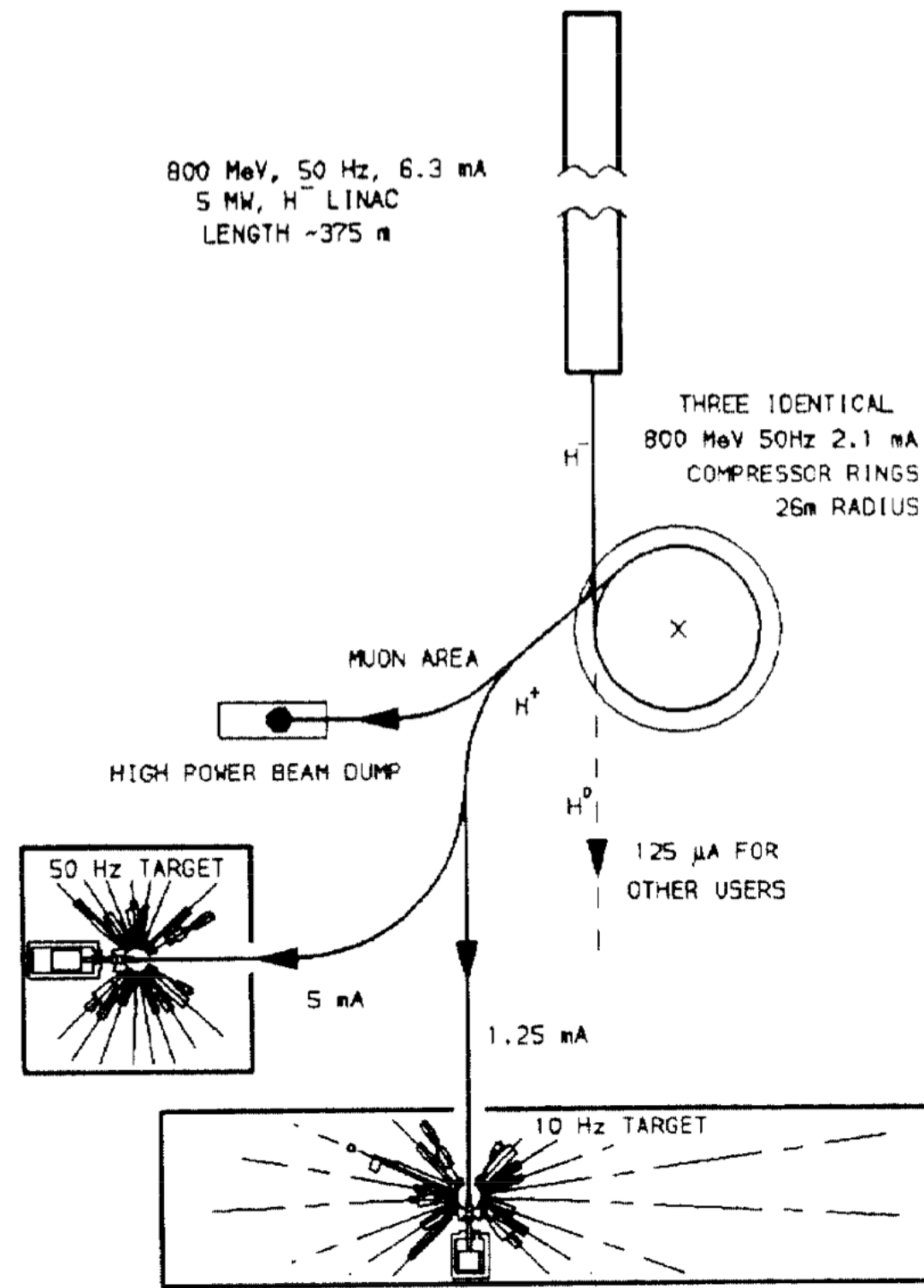


Figure 1. Linac - Compressor Rings Proposal.

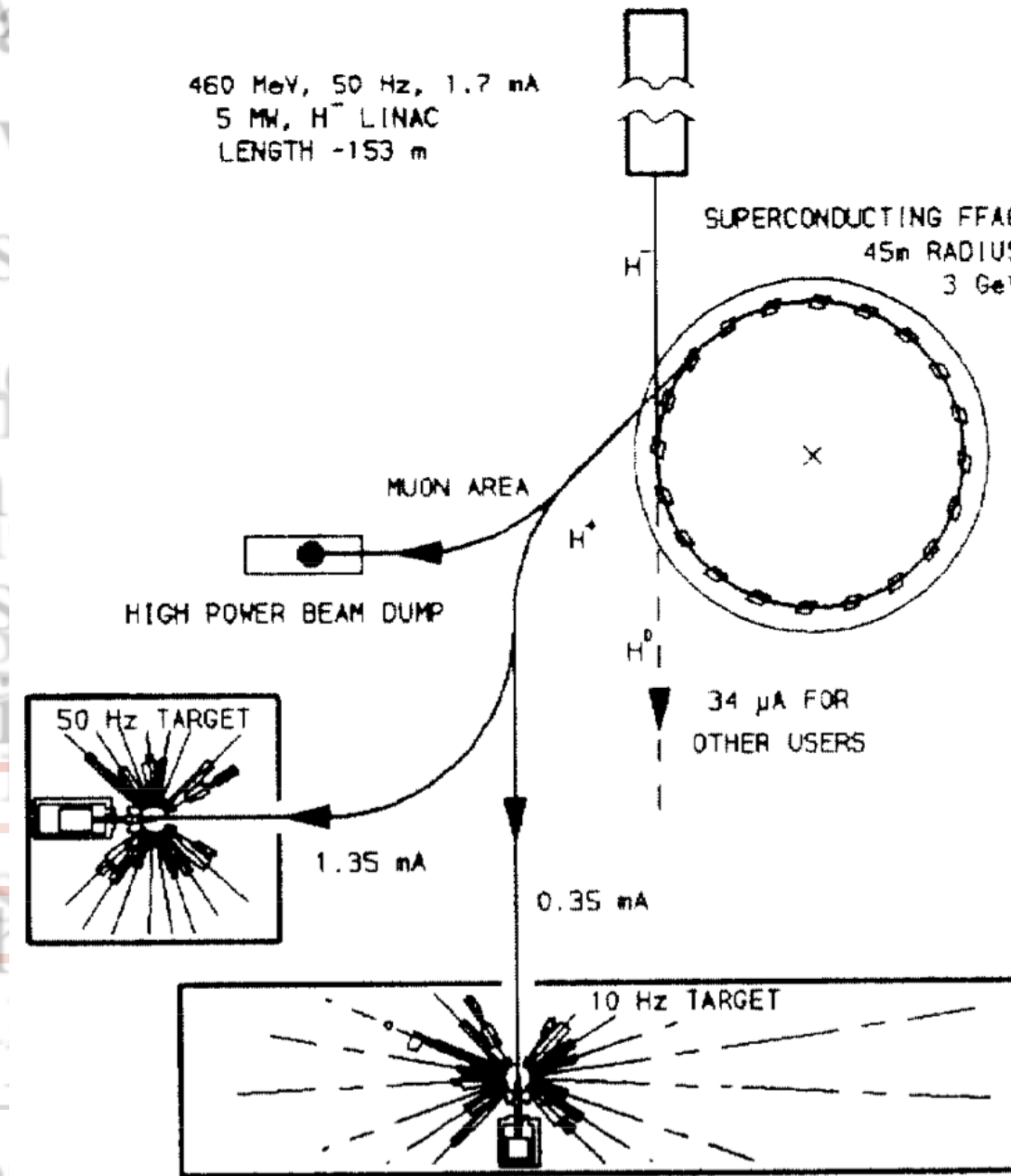
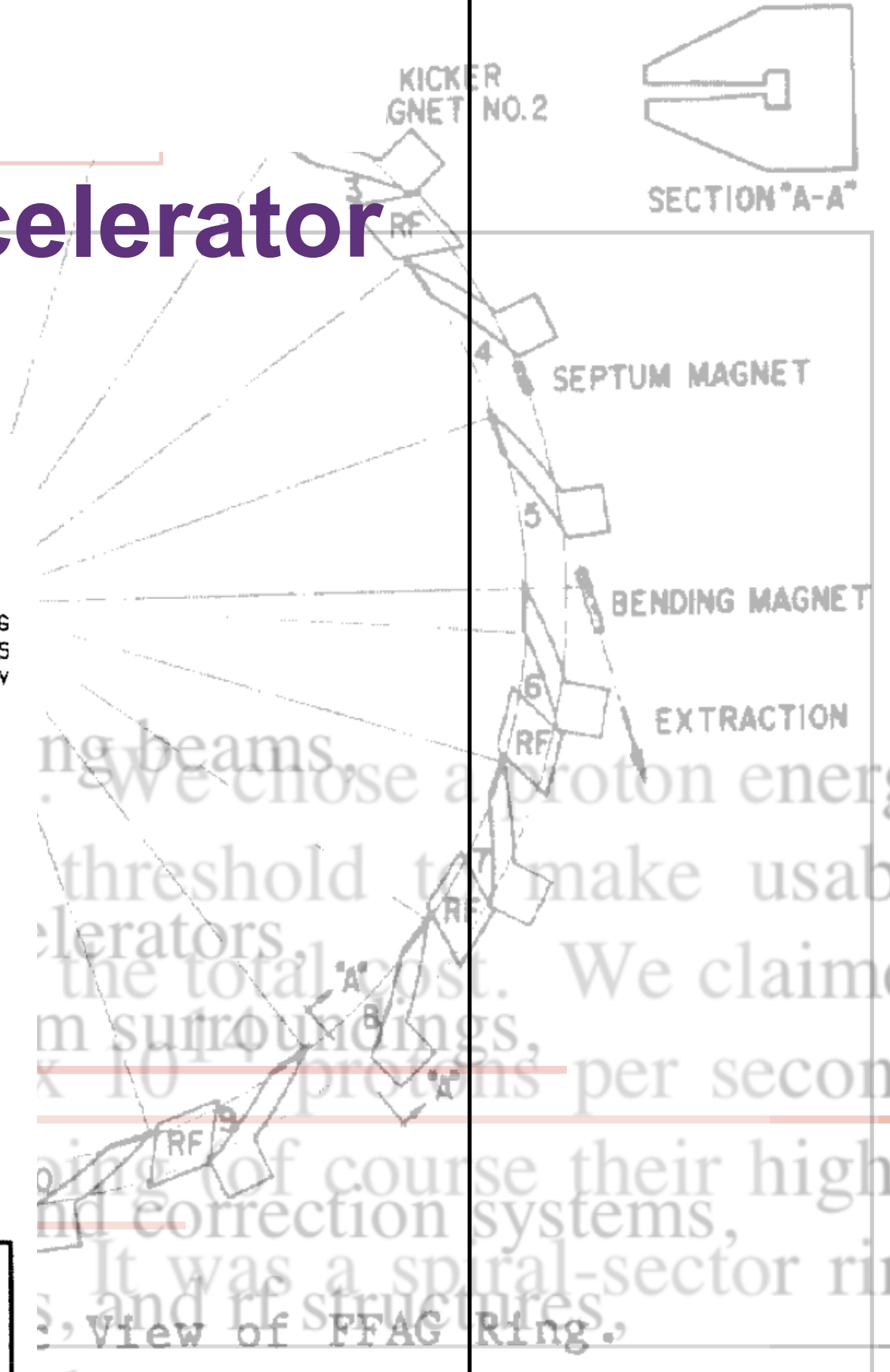


Figure 2. Linac - FFAG Proposal.





# Why FFA for high power applications?

## Why we are considering now?

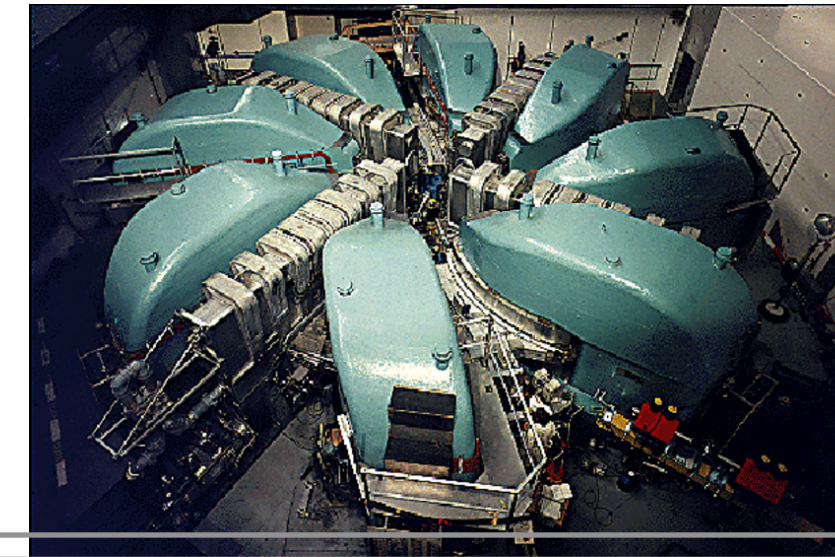


# Why FFA for high power applications?

*High power by high repetition, but not necessarily low peak current*

**Continuous acceleration like a cyclotron is the best way to increase the average beam power.**

- When synchrotron was invented, people had to accept the huge reduction of the beam current.



PSI cyclotron

ISIS synchrotron



By giving up the isochronous condition, **the accelerator lattice or magnet size can be reduced.**

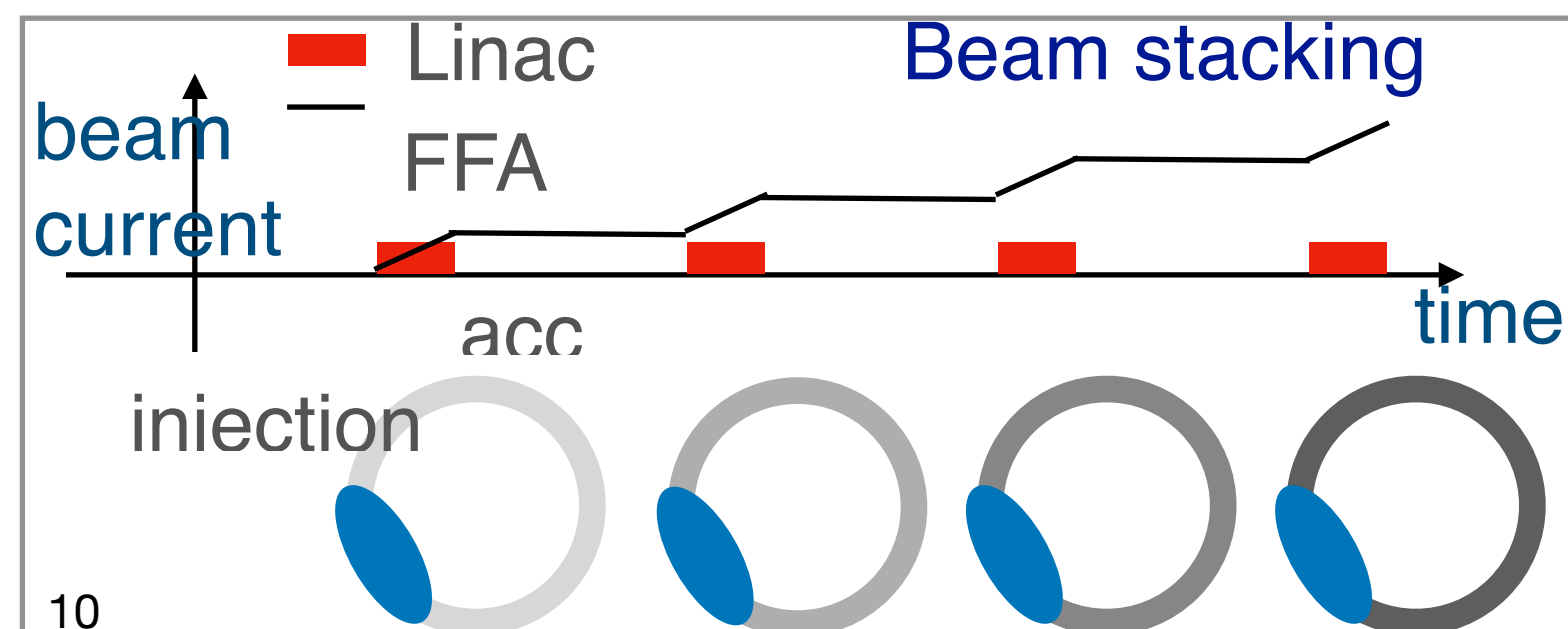
- This is the reason why a synchrotron took over as a high energy accelerator.

Reduction of the beam current can be **compensated by increasing the repetition rate.** That is possible by fixed field magnets.

- Some high power applications prefer pulsed beam to CW beam.



Kyushu U. FFA



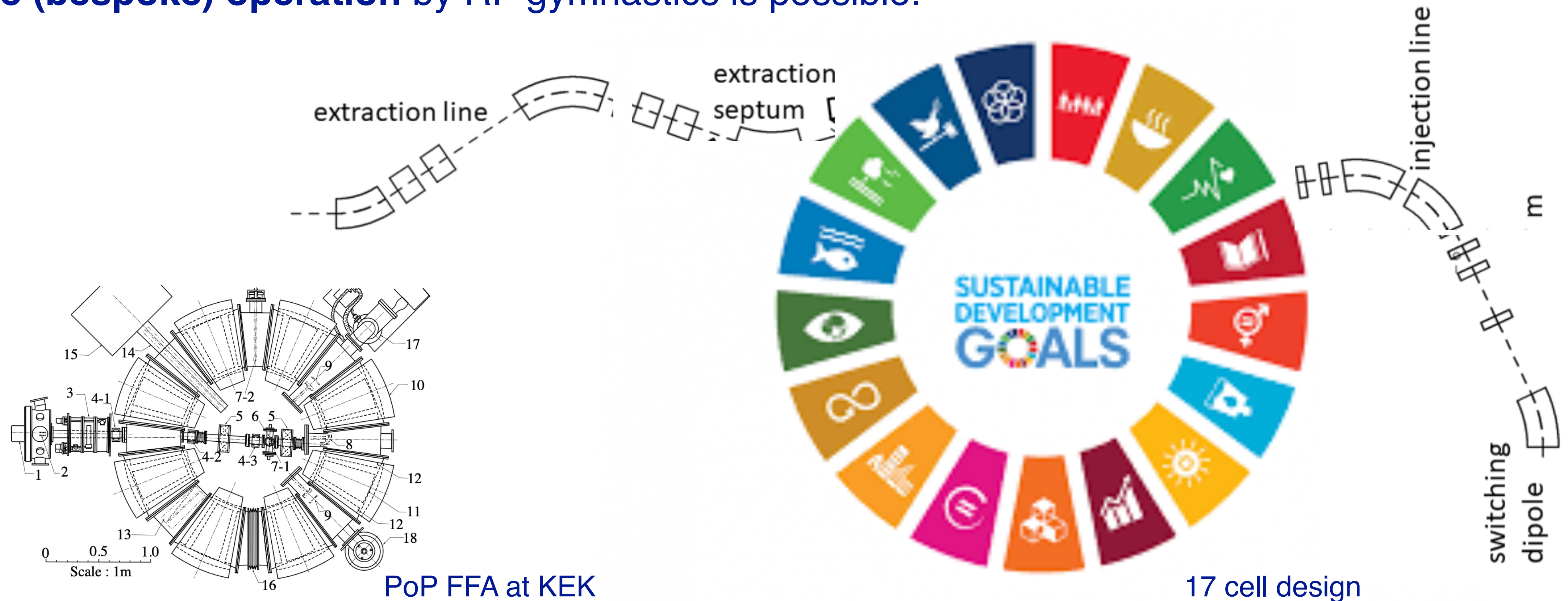
**By beam stacking, the pulsed peak current can be increased** keeping the average power with a lower repetition rate ( $\sim 10$  Hz).

- As a proton drive for a muon collider, spallation neutron source, etc.

# Why FFA for high power applications?

## *Other advantage with DC magnets*

- Required wall power is less to produce the same magnetic field compared with AC magnets for RCS.
  - **“Sustainability”** is the important keyword for future facilities.
- Main magnet can be superconducting and permanent magnet.
- **Reliability** increases without switched power supply.
- **Flexible (bespoke) operation** by RF gymnastics is possible.

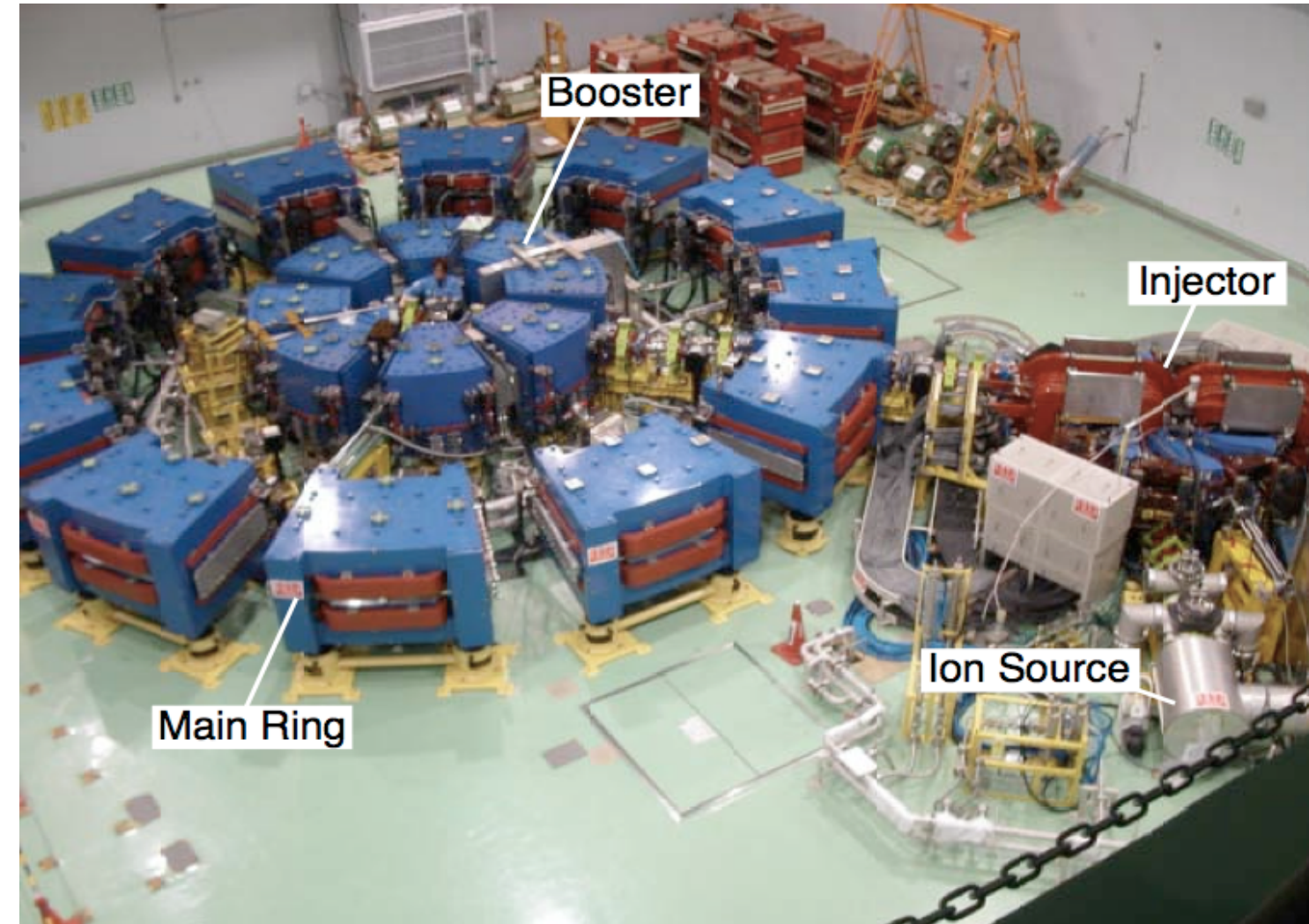
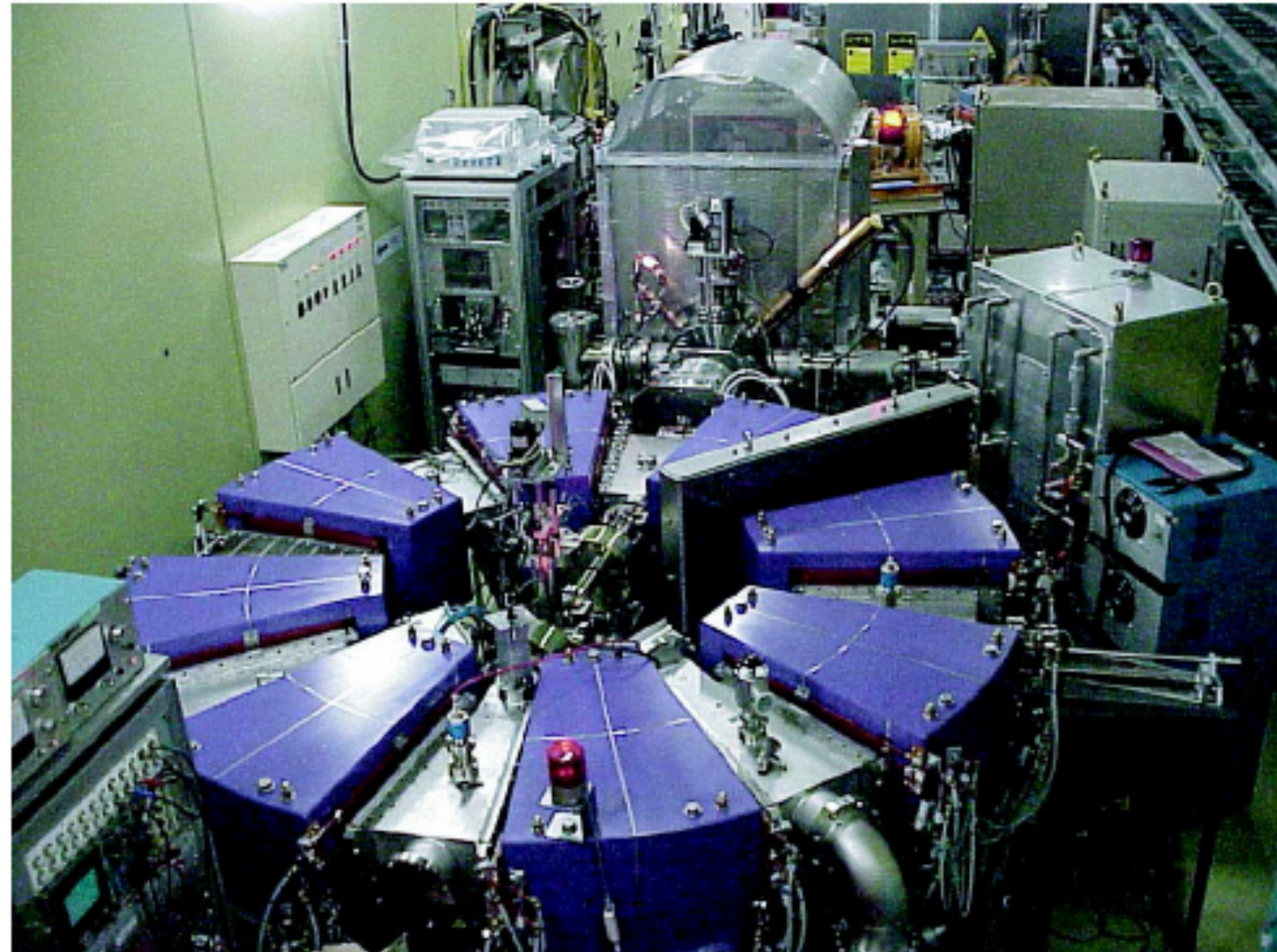


# Why we are considering now?

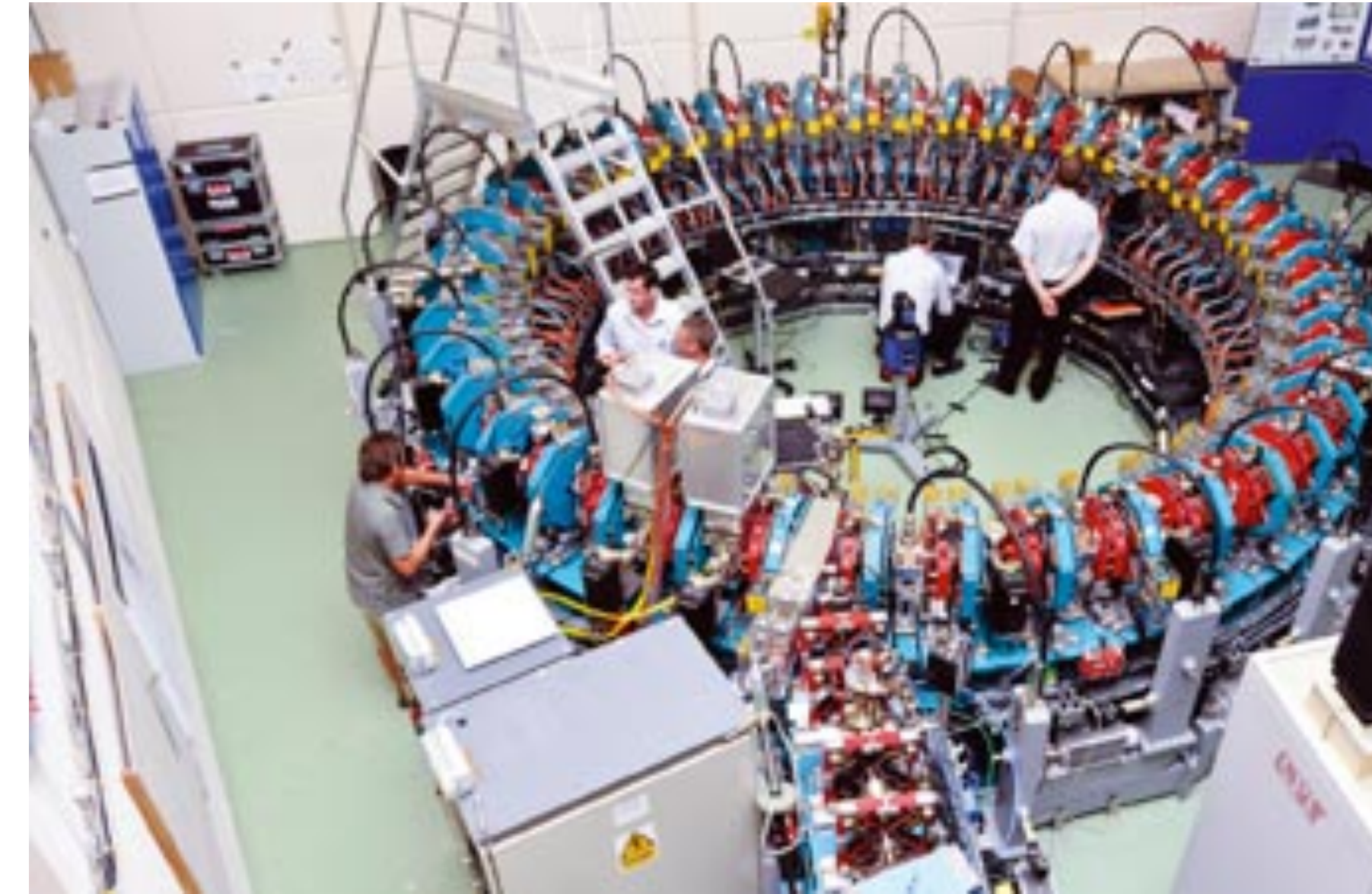
## Rebirth of an FFA

- High energy acceleration to 150 MeV.
- Cascade FFAs.

- Acc from 50 to 500 keV of protons.
- Repetition of 1 kHz operation.



- EMMA: monscaling FFA.
- Serpentine channel acceleration.



**No high power FFA yet.**



- CBETA: multipass arc of ERL.
- Permanent magnet lattice.

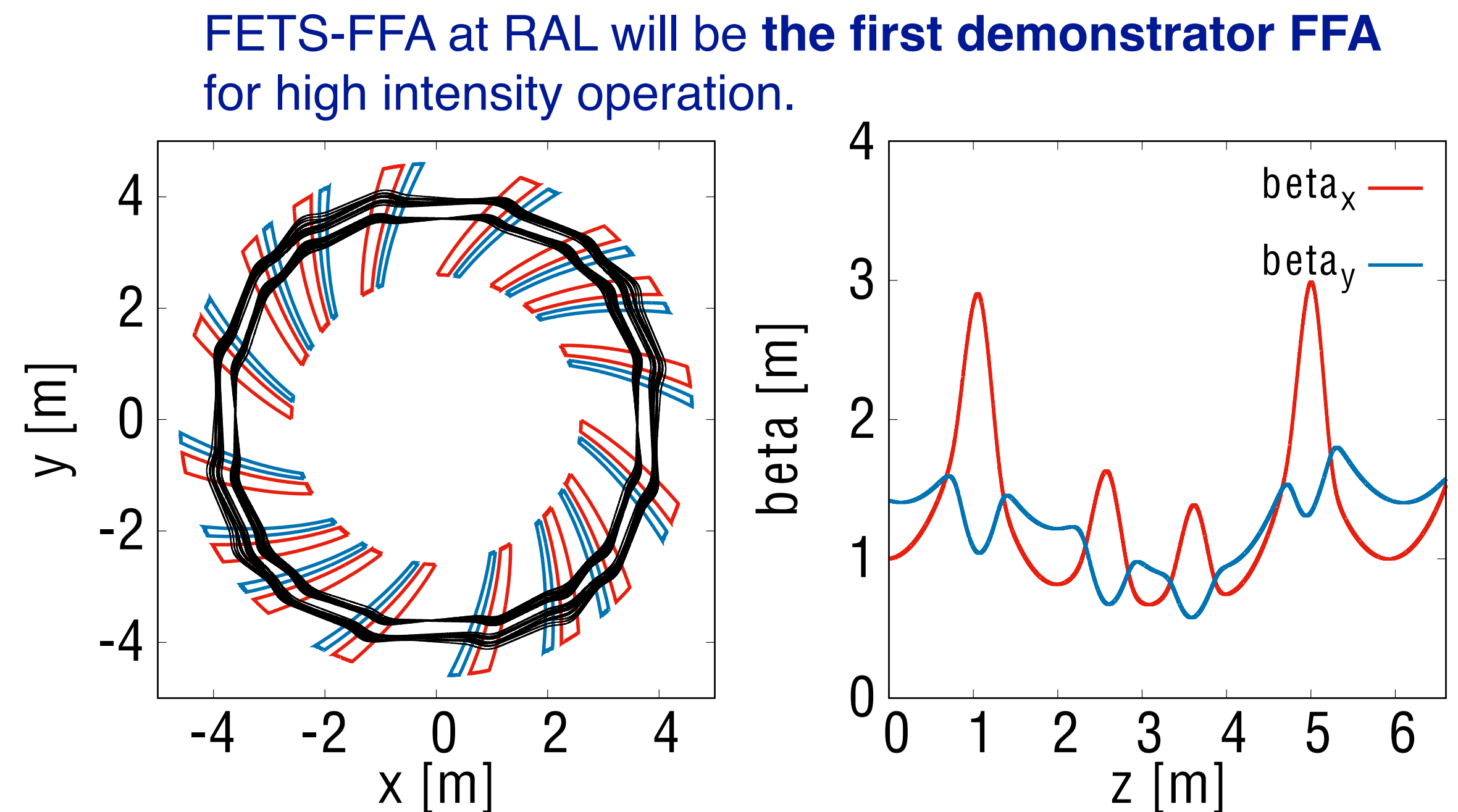
# Toward high power FFA



# Before we start ...

- In the following, I tried to keep the discussion as in general as possible.
- However, I occasionally use specific design parameters where the discussion points become clearer.
- The parameters are based on a demonstrator of high power FFAs which we plan to built at RAL.

Energy	3 - 12 MeV
Minimum radius	3.6 m
Particle	Proton
Maximum intensity	$3 \times 10^{11}$
Emittance (nor.)	10 pi mm mrad
Space charge tune shift	-0.3
Repetition	100 Hz (50 pps)
Average beam power	~ 50 W

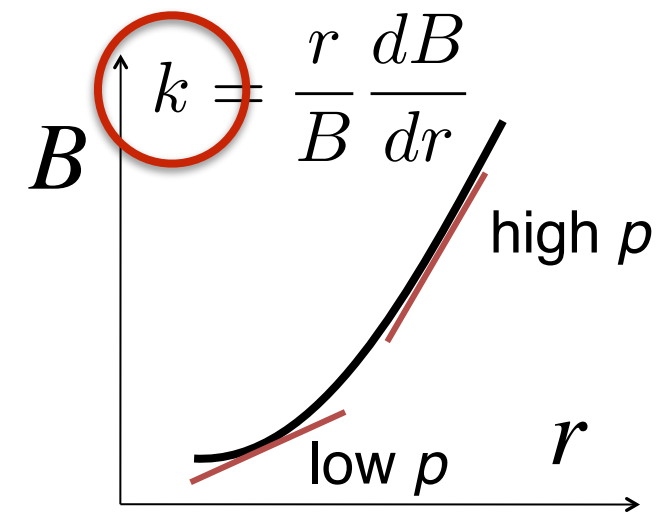
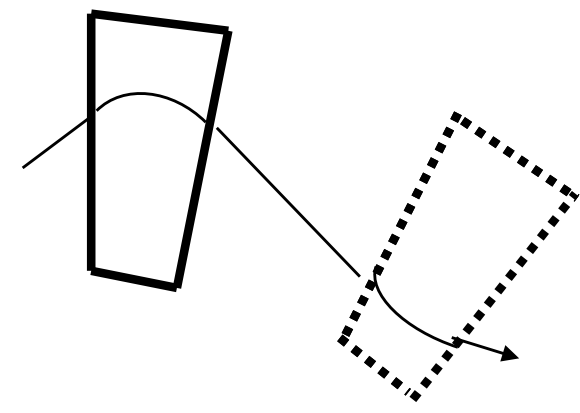


# DF (FD) spiral sector

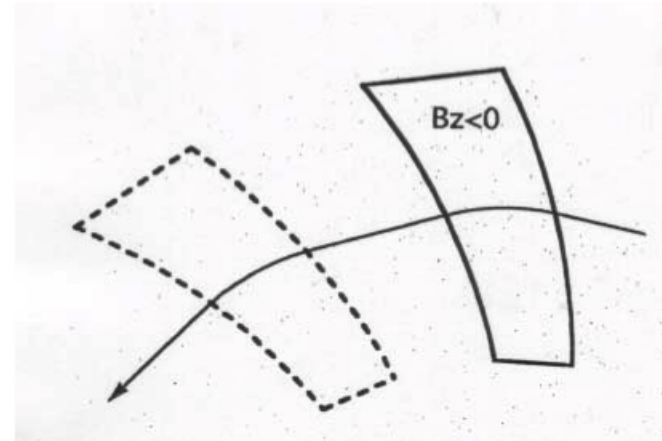
## Combination of radial and spiral

Flexibility of operating point (transverse tune) is essential for high intensity operation ( $Q_h \sim Q_v$ ).

radial sector

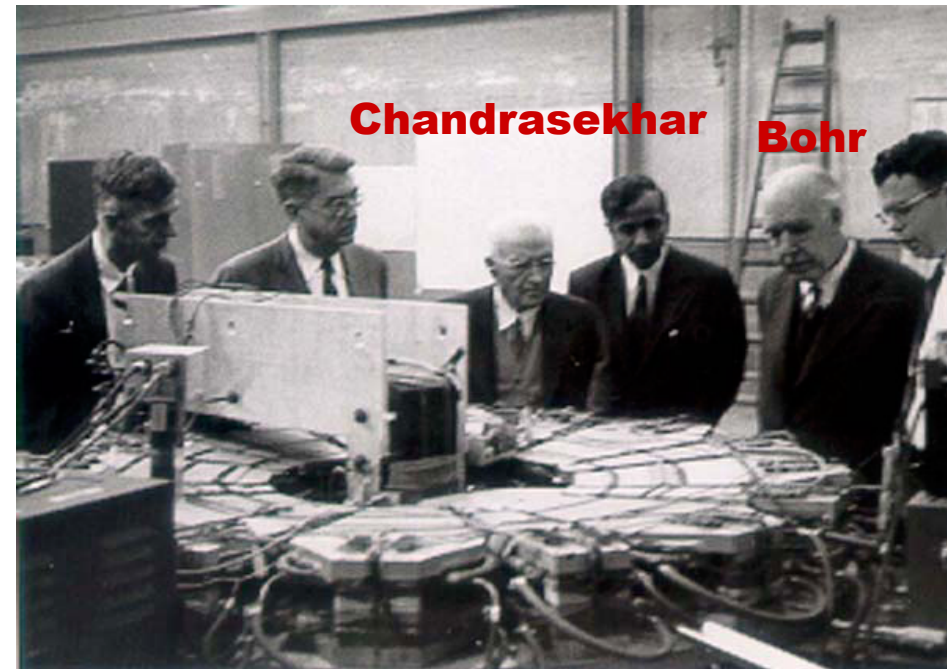


spiral sector

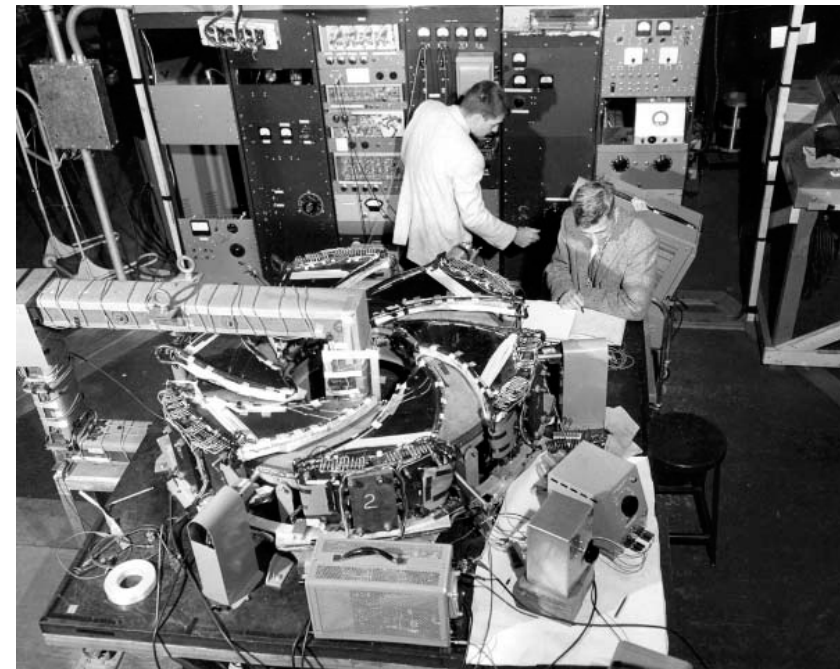


Alternating gradient focusing by focusing (normal bend) and defocusing (**reserve bend**)

Alternating gradient focusing by focusing (normal bend) and defocusing (**edge angle**)



400 keV radial sector  
Science and Technology Facilities Council



180 keV spiral sector

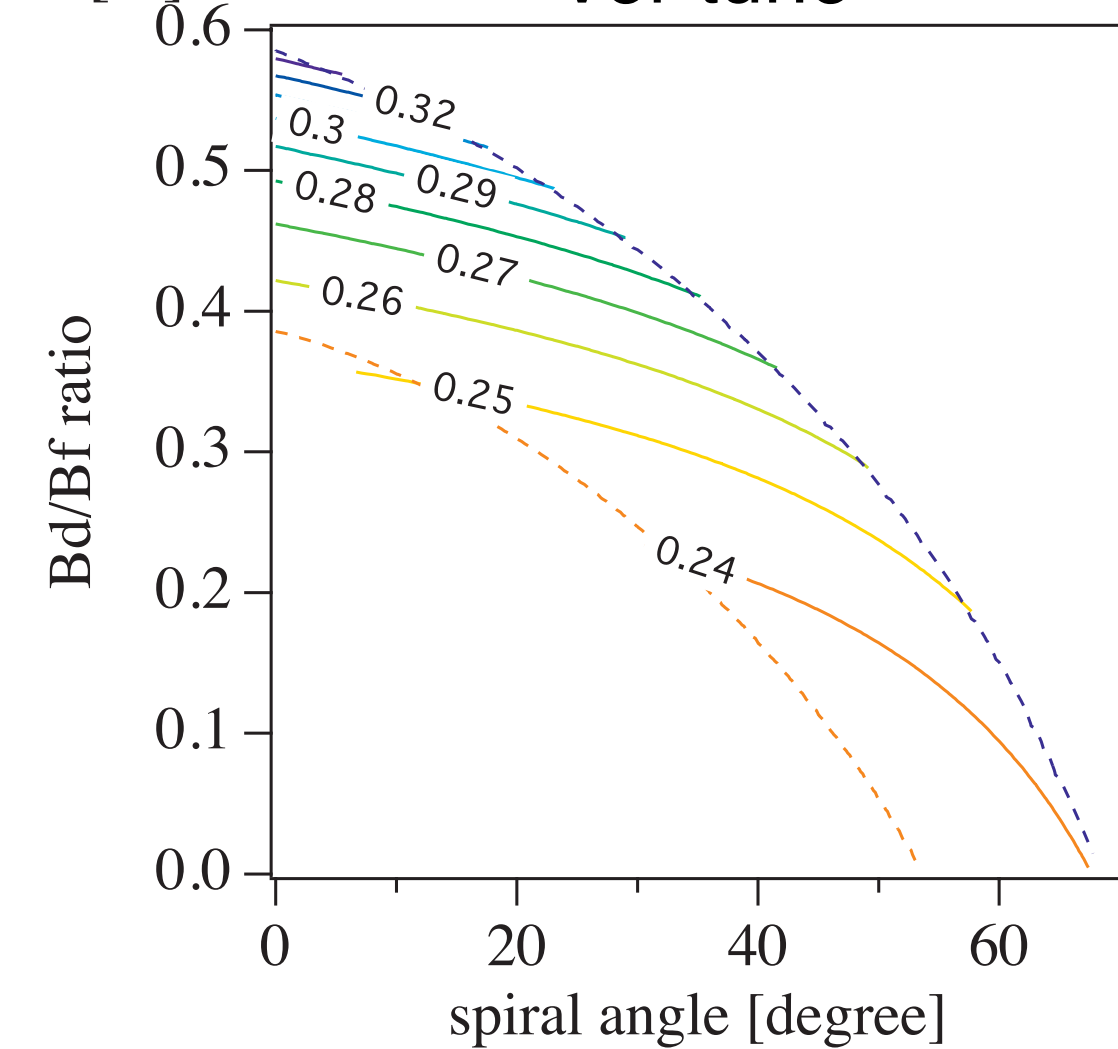
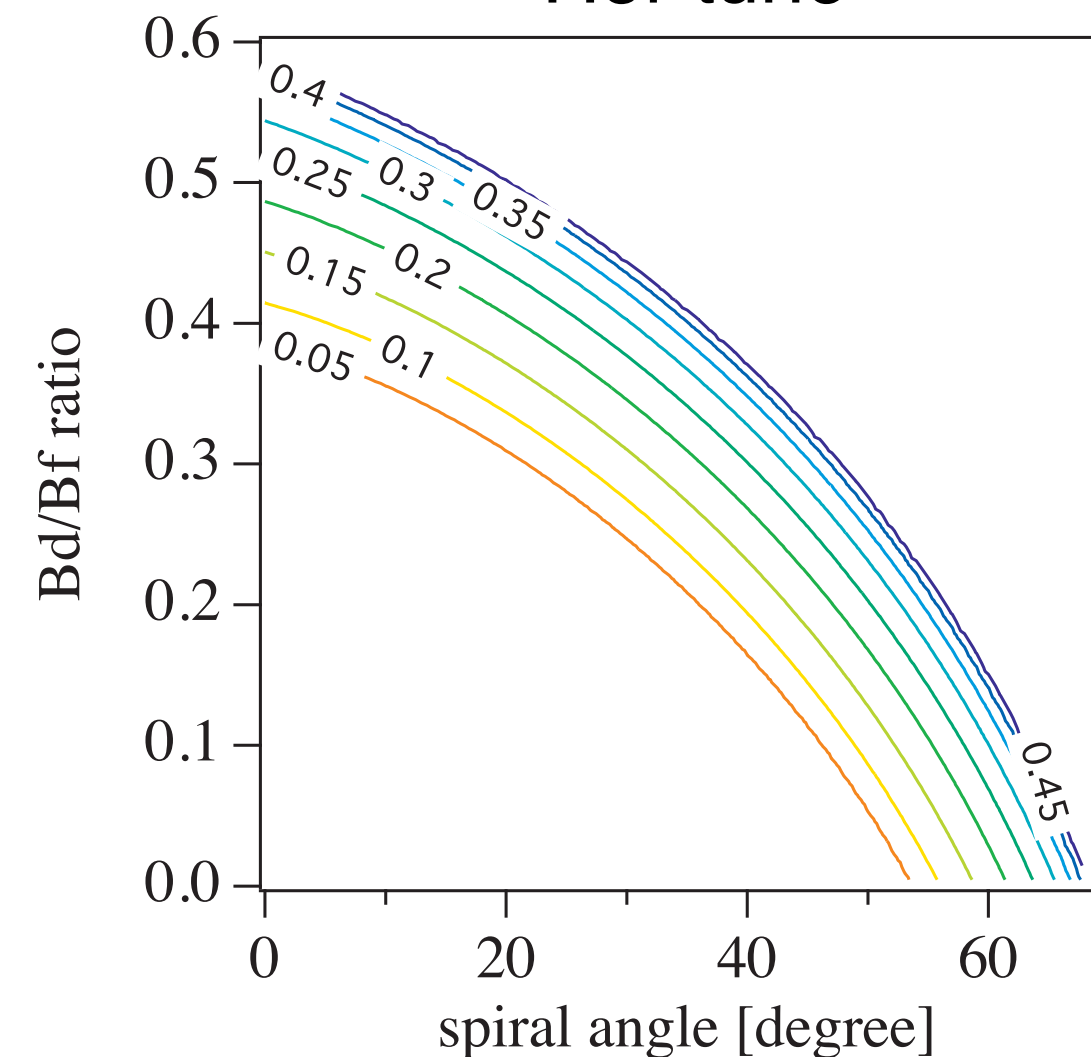
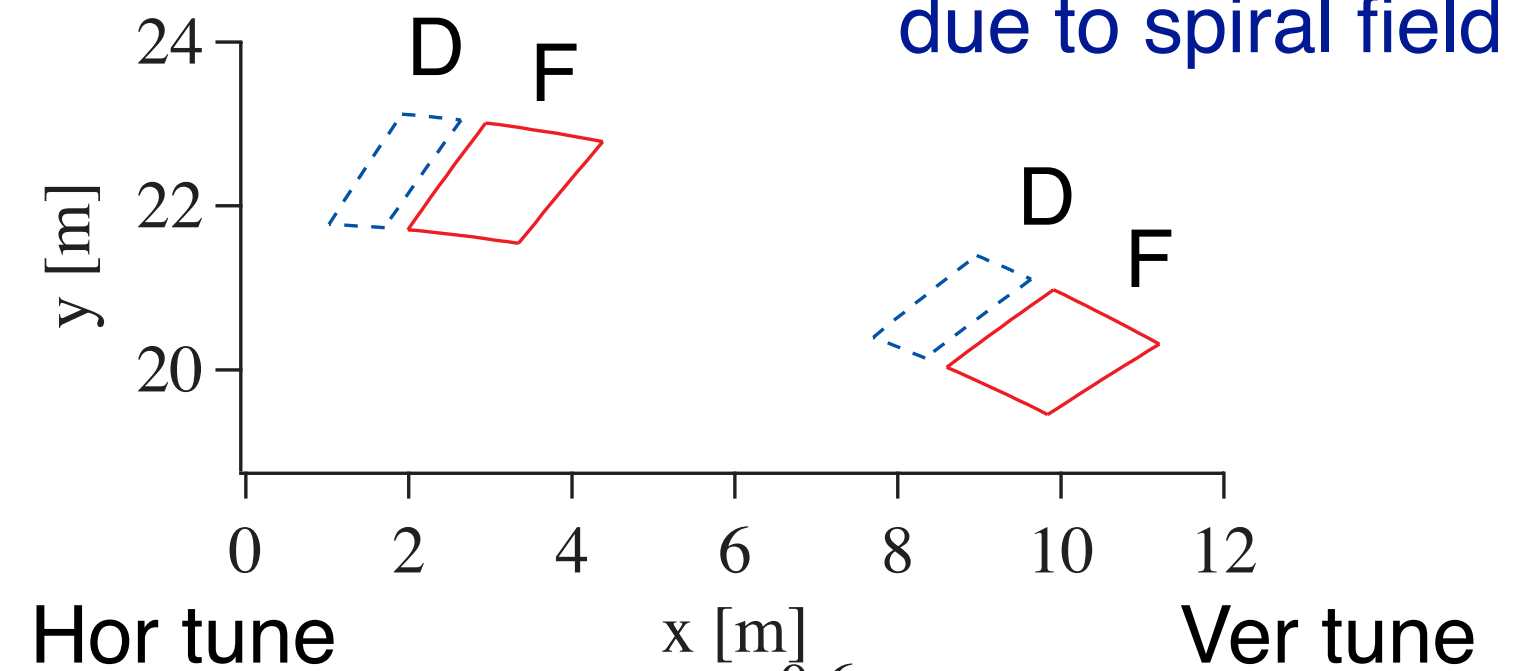
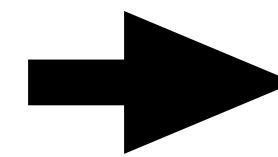
Strong focusing produced by the gradient variation with azimuth arising from the undulation of the orbit.

$$Q_x^2 \approx 1 + k + \frac{k^2 S^2}{N^2 b_0^2}, \quad (4)$$

$$Q_z^2 \approx -k + \frac{k^2 S^2}{N^2 b_0^2} + \frac{\Phi^2}{b_0^2} (1 + 2 \tan^2 \delta), \quad (5)$$

Field gradient averaged over the azimuth.

Specific strong focusing due to spiral field shape.

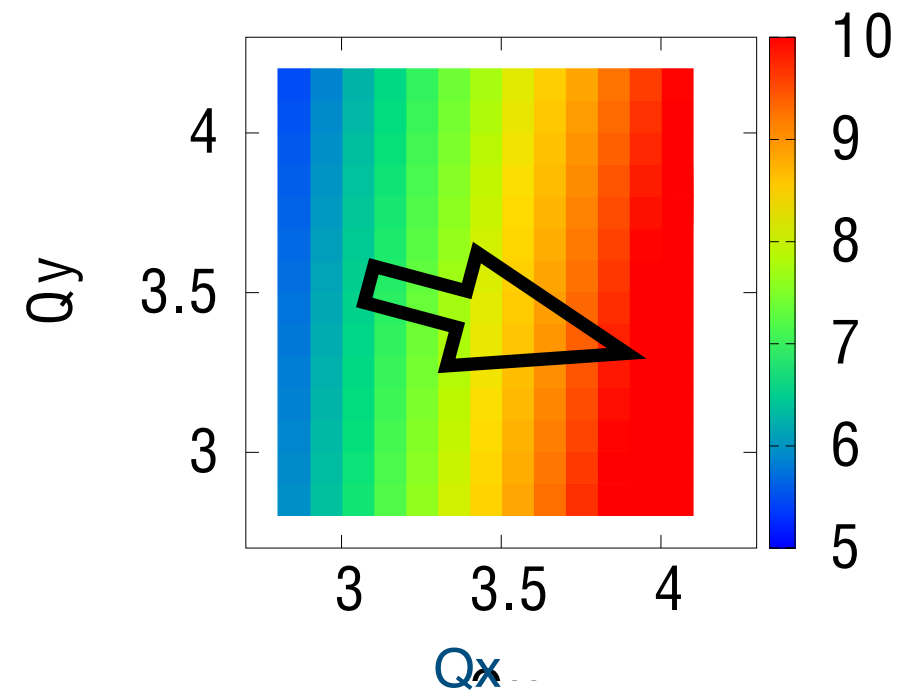


Machida, PRL 119, 064802 (2017)

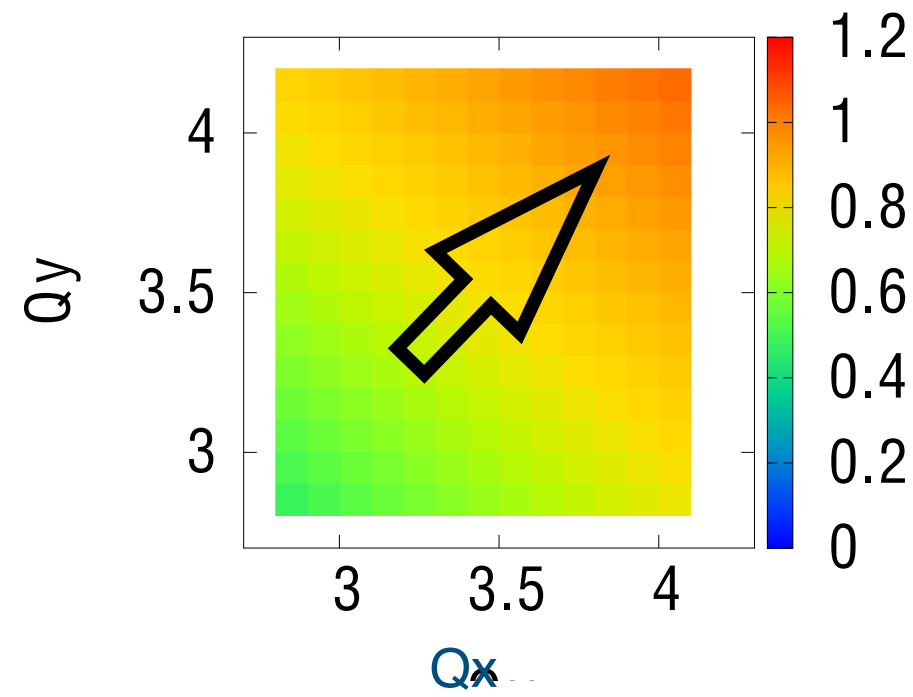
# DF (FD) spiral sector

*Explore wide operating point*

k-value



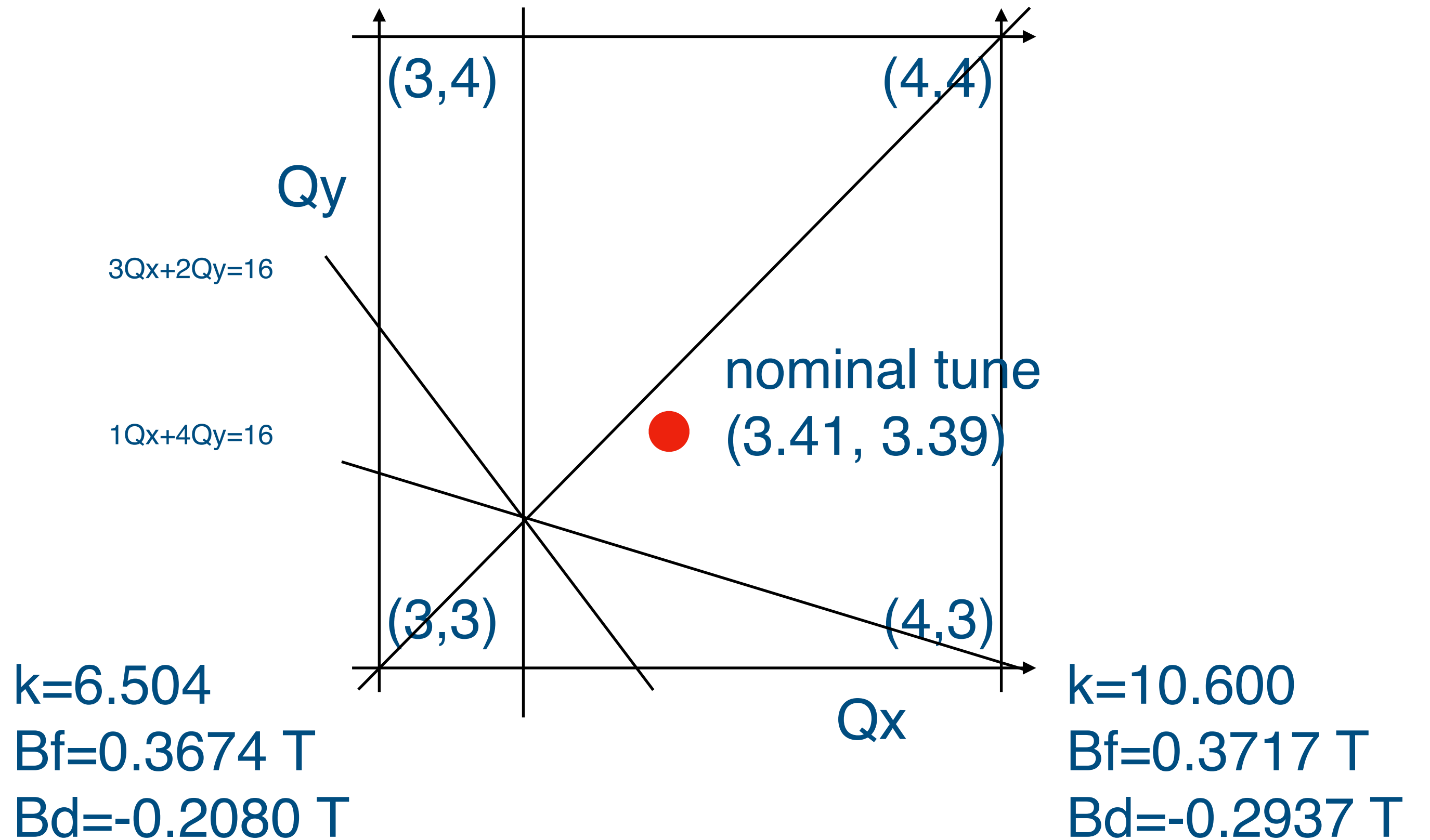
Bd/Bf ratio



k-value and Bd/Bf strength ratio are two parameters to adjust tune Qx and Qy.

k=6.102  
Bf=0.4231 T  
Bd=-0.3462 T

k=9.891  
Bf=0.4153 T  
Bd=-0.4135 T



k=6.504  
Bf=0.3674 T  
Bd=-0.2080 T

k=10.600  
Bf=0.3717 T  
Bd=-0.2937 T

**Tune space can be explored without much depending on a reverse bend.**

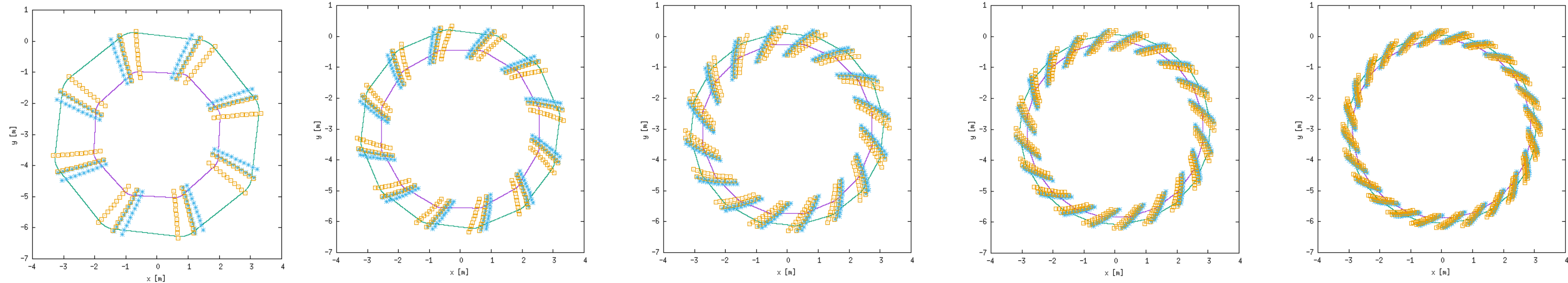


# Superperiodicity

## Orbit excursion vs number of cells

$$B_z(r, \theta) = B_{z0} \left( \frac{r}{r_0} \right)^k F(\theta) \quad k = \frac{r}{B} \frac{dB}{dr}$$

- Increasing the number of cells
  - > higher field index  $k$  -> small **orbit excursion** (good).
  - > shorter **straight section** (bad).



- Let us keep reasonable number of cells, but allocate straight sections unevenly.

Introduction of **superperiod** by exciting  $m$  not equal to the number of cell.

**Long straight section is essential for proper handling of the high intensity beams.**

for example, phase space painting with charge exchange injection.

$$F(\theta) = \sum_m f_m \exp(im\theta)$$

# Superperiodicity

## Long straight section with zero chromaticity

$$B_z(r, \theta) = B_{z0} \left( \frac{r}{r_0} \right)^k F(\theta)$$

### 16-fold symmetry

$$F(\theta) = f_{16} \exp(i16\theta)$$

Straight length: **0.95 m**

Dynamic aperture: 110 pi mm mrad

Field index k: 8.00

Spiral angle: 45 degree

**Magnet families: 2**

### 4-fold symmetry

$$F(\theta) = f_4 \exp(i4\theta) + f_{16} \exp(i16\theta)$$

Straight length: **1.55 m, 0.90 m, 0.45 m**

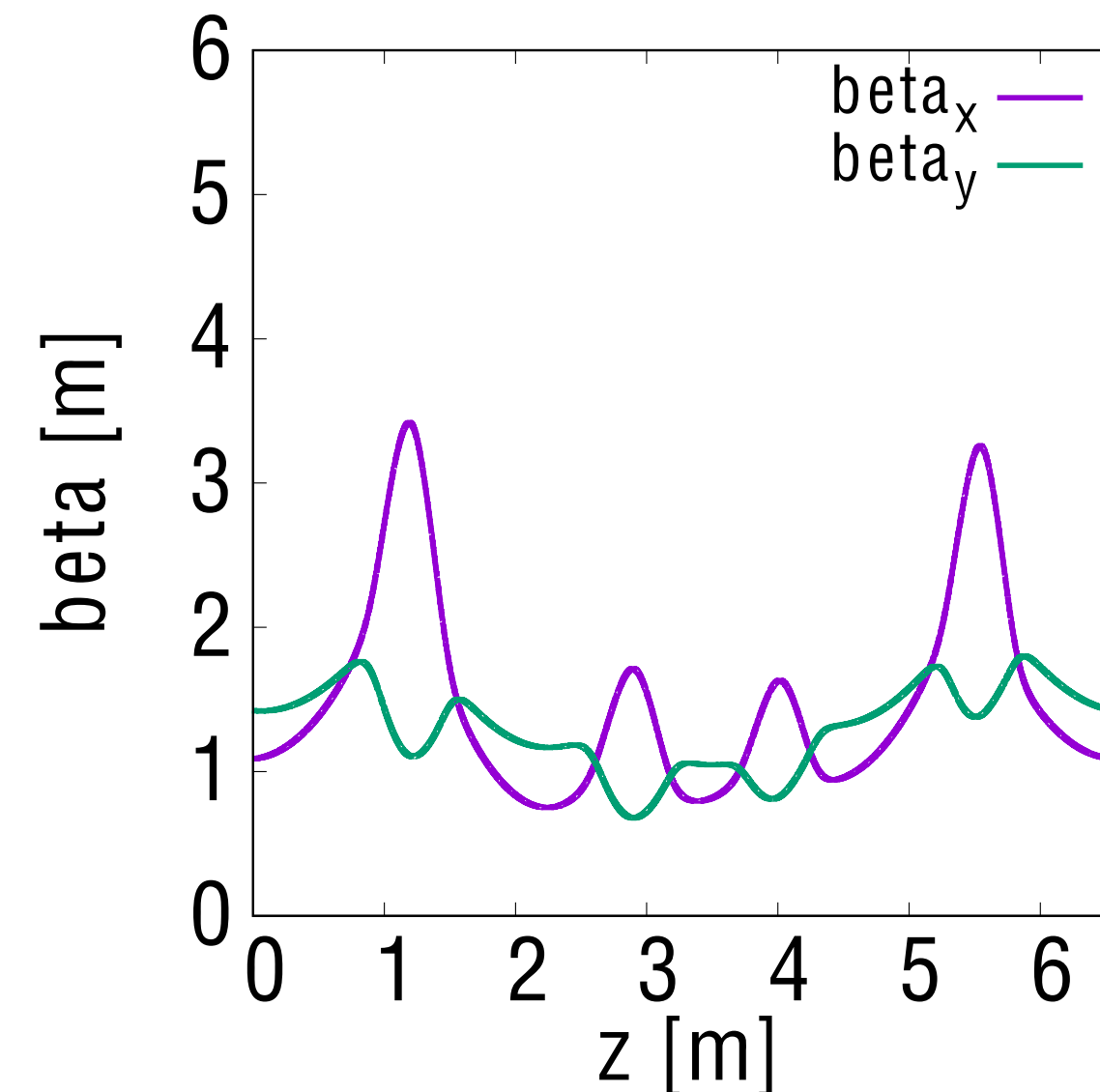
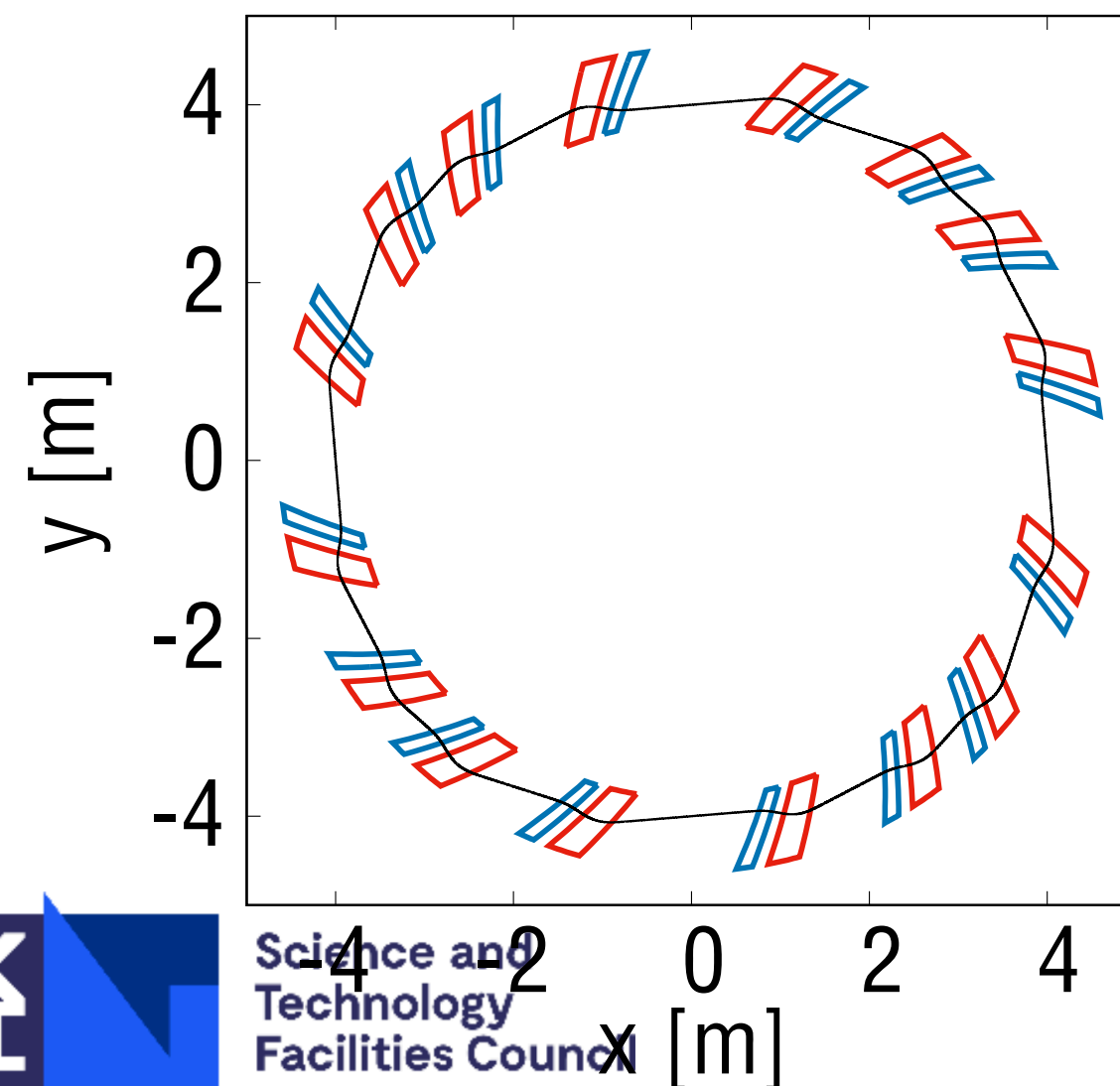
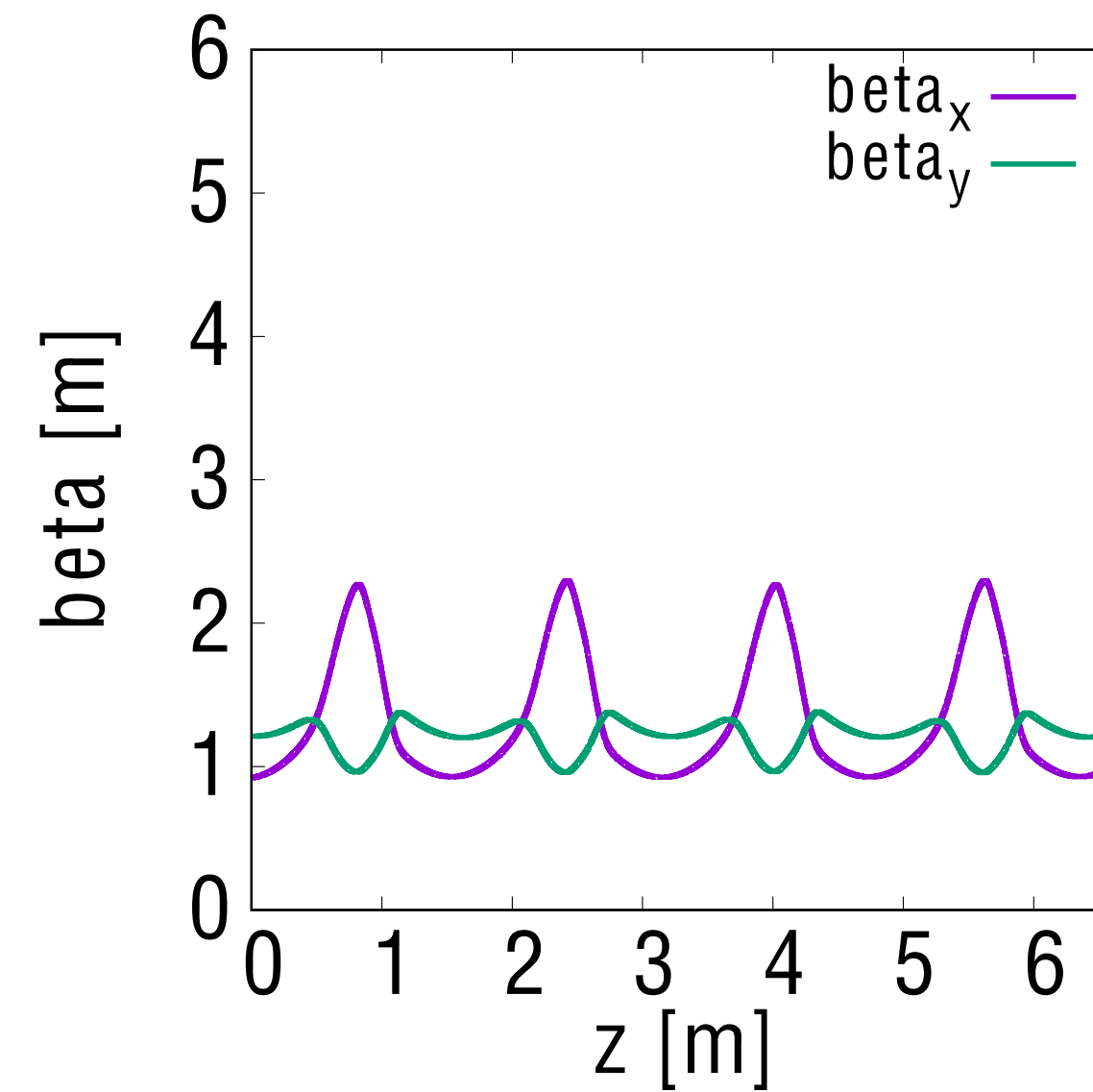
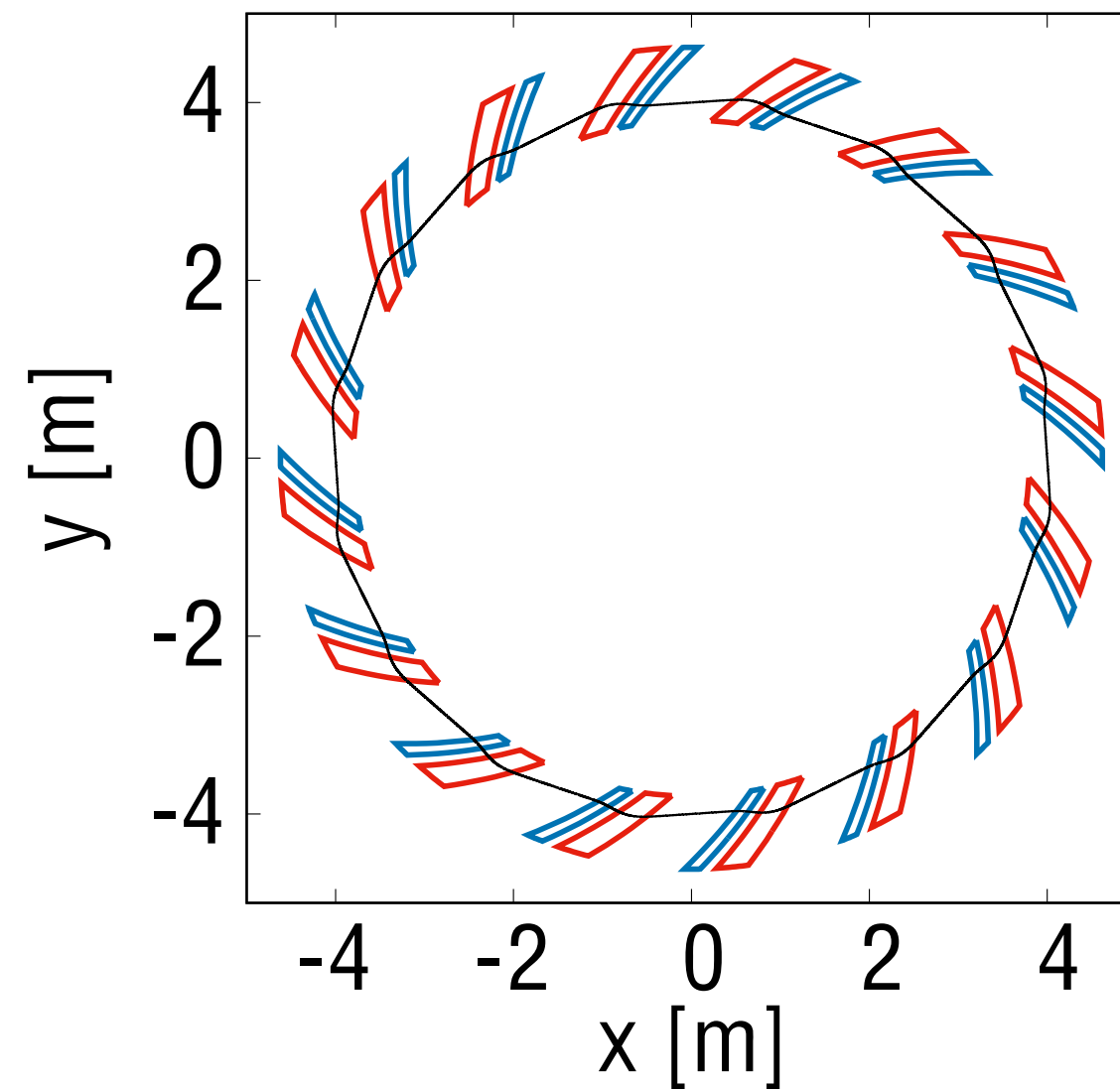
Dynamic aperture: 80 pi mm mrad

Field index k: 7.40

Spiral angle: 30 degree

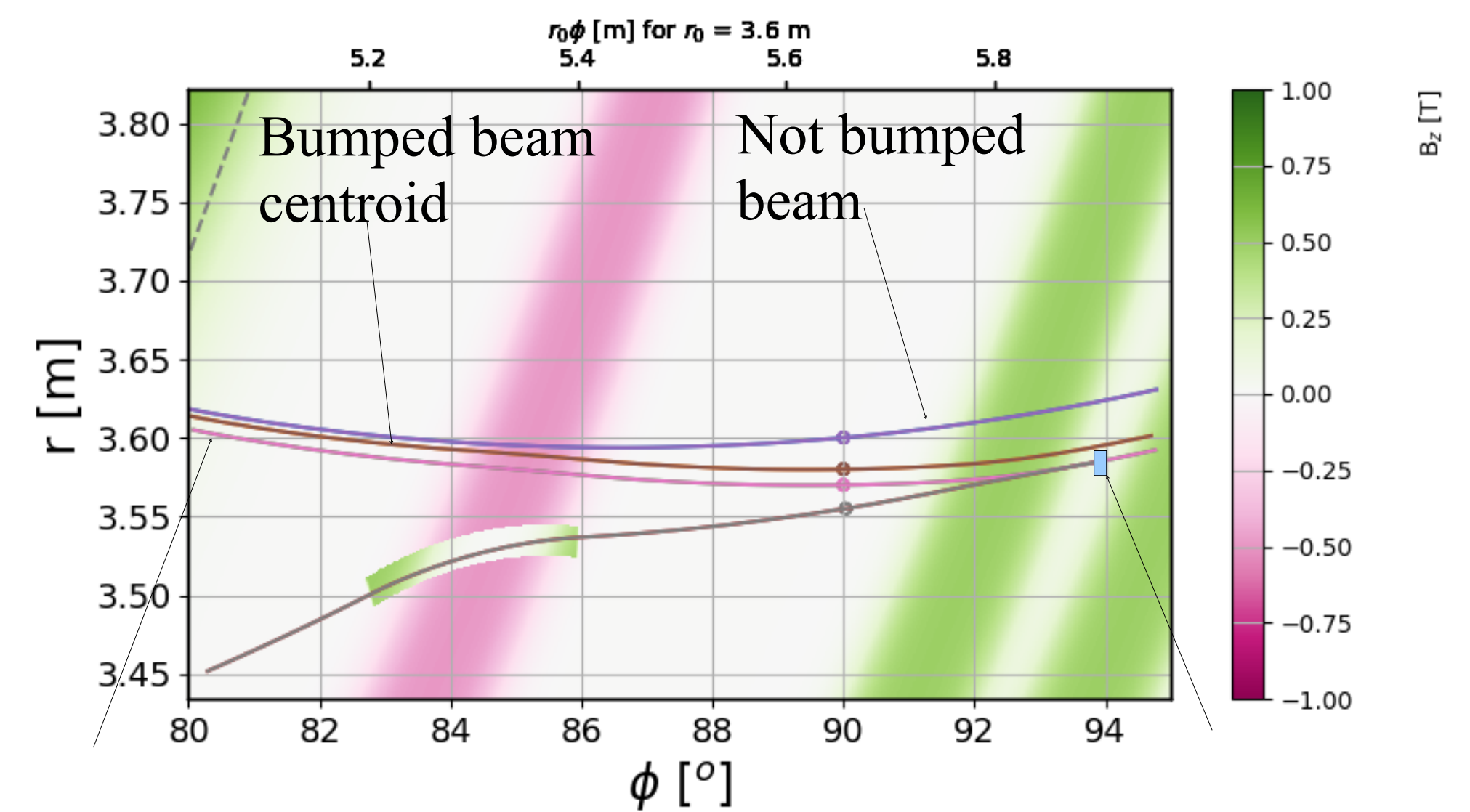
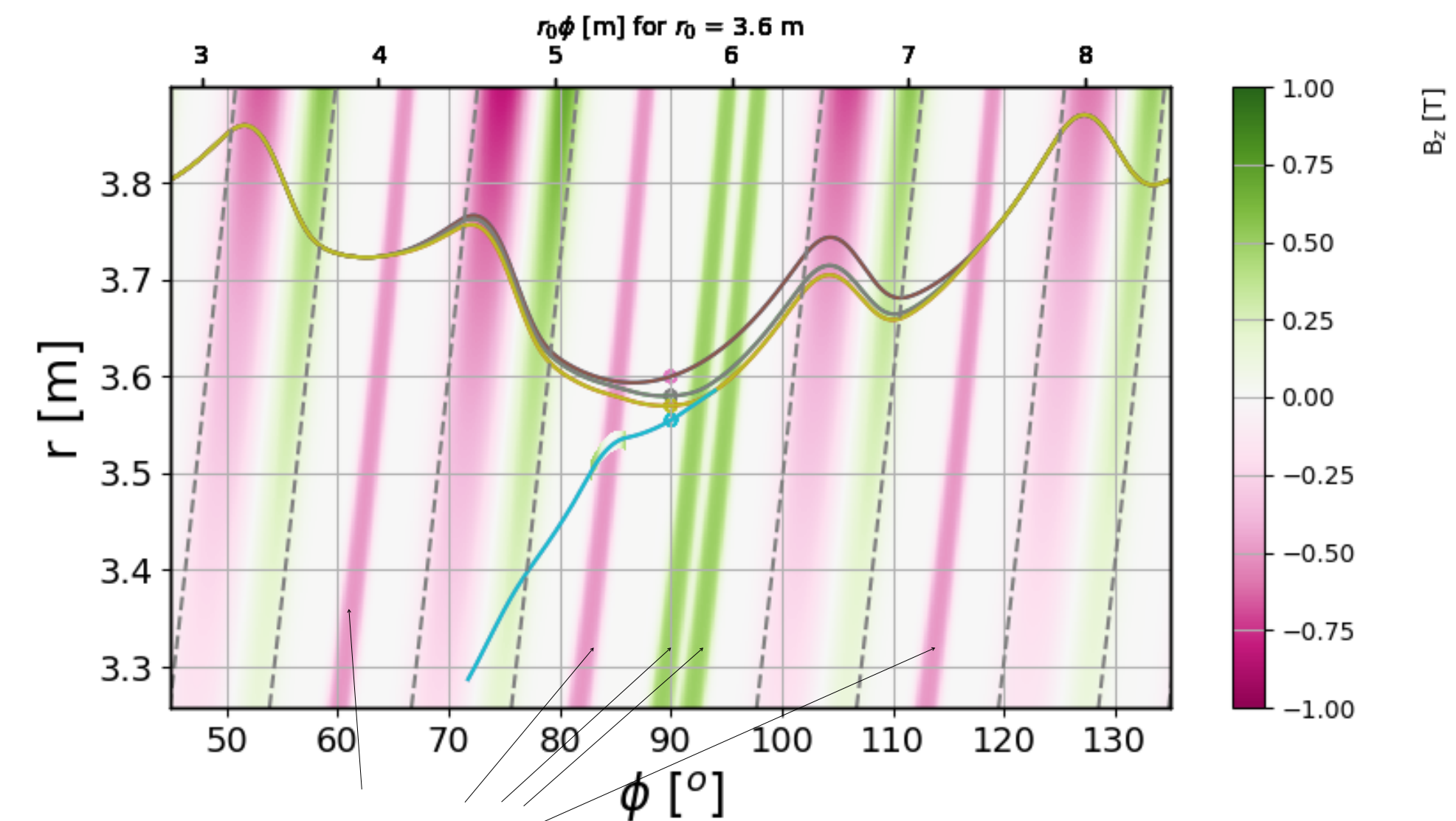
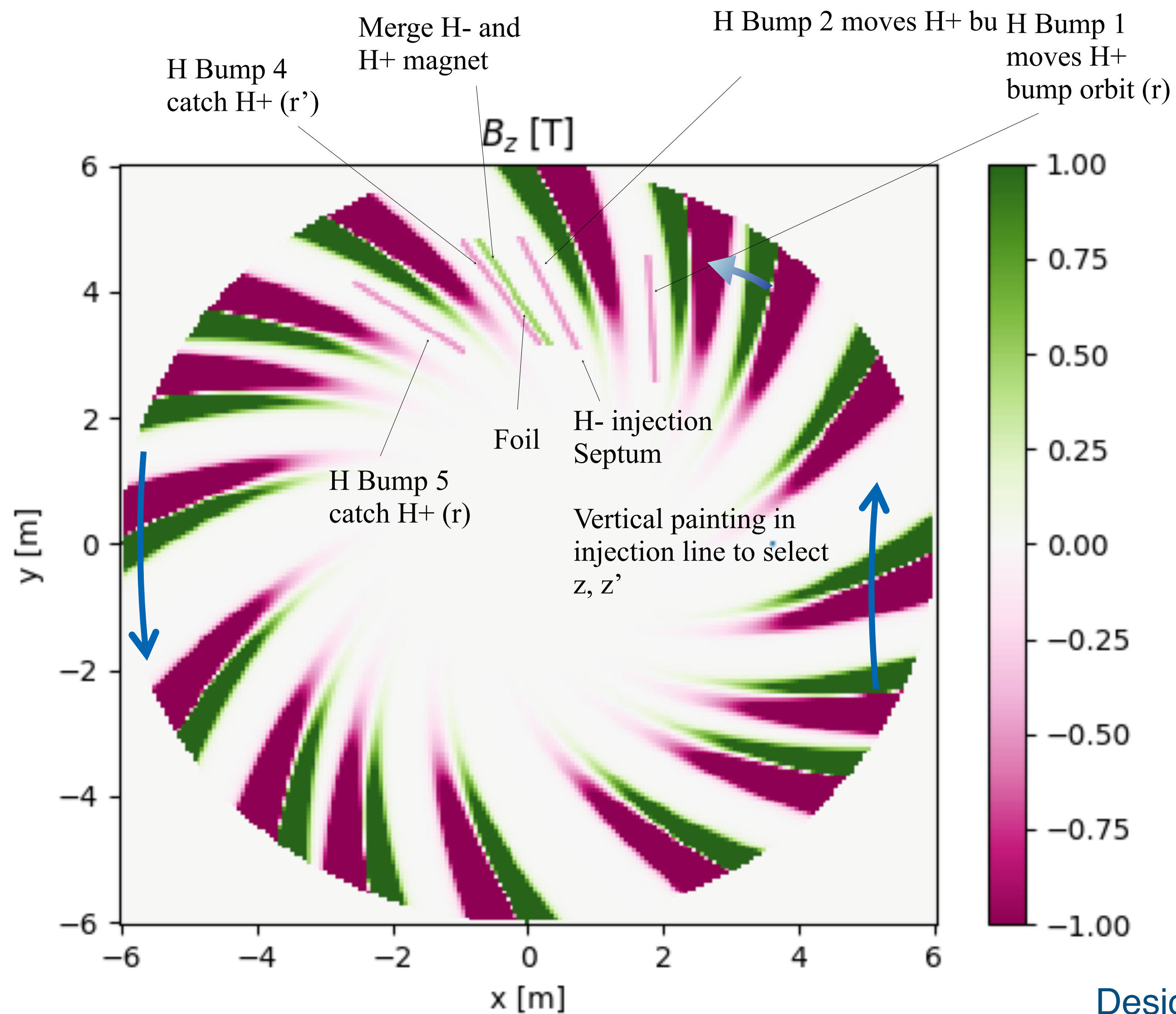
**Magnet families: 8**

Note horizontal beam size is larger.



# Injection

H- charge exchange injection with 5 bump magnets (3 bumps in a long straight).



Design by Chris Rogers

# Physical and dynamic aperture

## Large acceptance pays off

Optimise nonlinearity to obtain the same acceptance of SNS/JPARC

	Normalised emittance	Geometrical acceptance	Vertical beam size [mm]
Beam core	10 [pi mm mrad]	125 [pi mm mrad]	+/- 16 mm
Collimator acceptance	20	250	+/- 22 mm
Vacuum chamber size	40	500 Same as SNS/JPARC	+/- 32 mm

At 3 MeV, uniform beam of 10 pi mm mrad (100%, normalised)

$$\Delta Q = -\frac{r_p n_t}{2\pi\beta\gamma^2\epsilon_n B_f} = -0.12 \quad \text{per } 10^{11} \text{ protons.}$$

FETS injector will reduce both emittance and peak intensity by more than one order of magnitude.

0.25 pi mm mrad, 60 mA

-> 0.02 pi mm mrad, 1 mA (50 turns for  $3 \times 10^{11}$ )

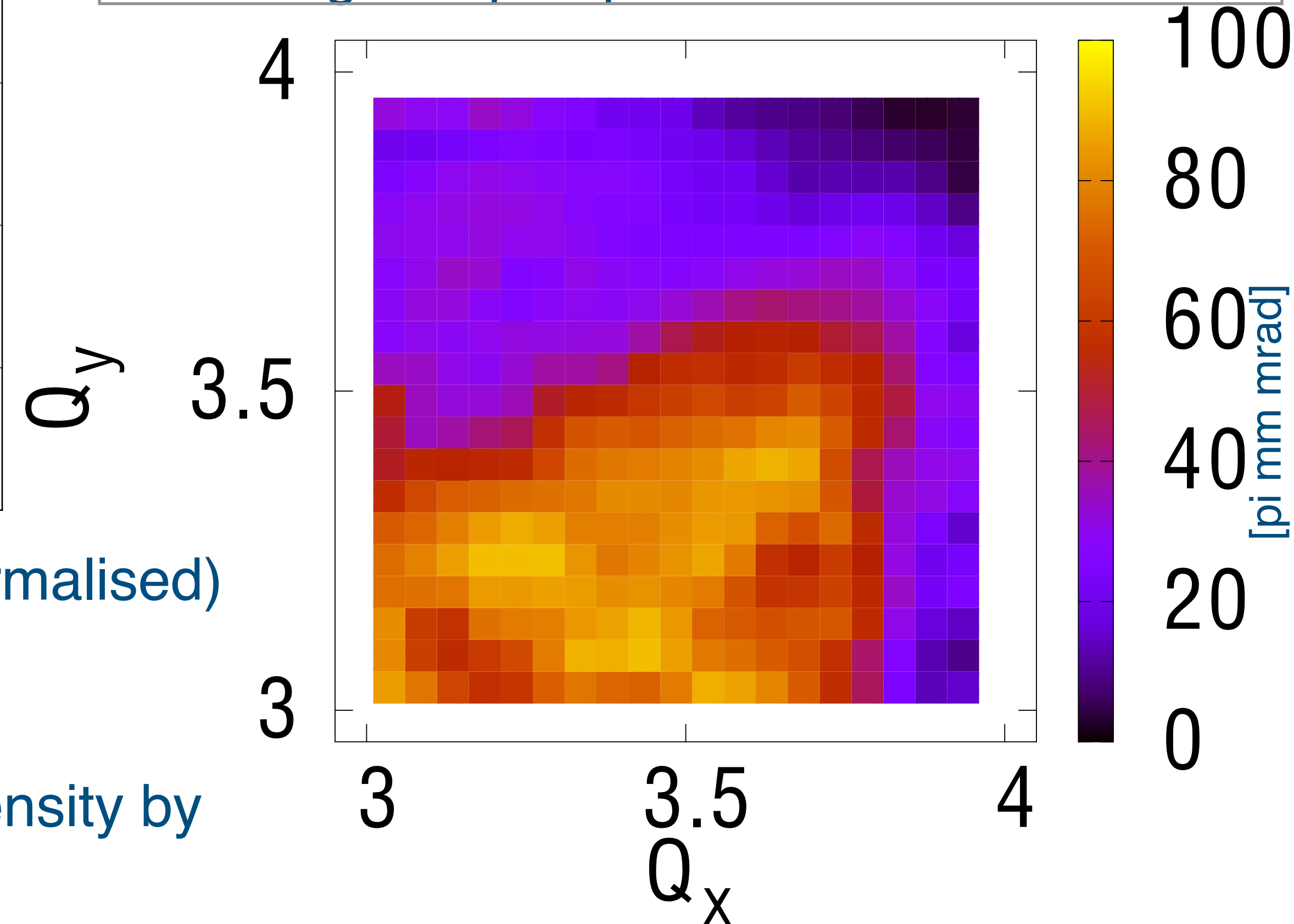
Aiba, et al

$$B_{\text{Loct}} = -\frac{1}{3!} (E^{(1)} k^2 + E^{(3)}) \frac{B_0}{r_0^3} z^3 \sin \beta$$

$$\cong -\frac{E^{(1)} k^2}{3!} \frac{B_0}{r_0^3} z^3 \sin \beta = O(s) z^3 \quad (11)$$

- E(1) is the first derivative of fringe field extent with azimuthal direction.
- Strong octupole at fringe fields.

Large amp dependent tune shift in V

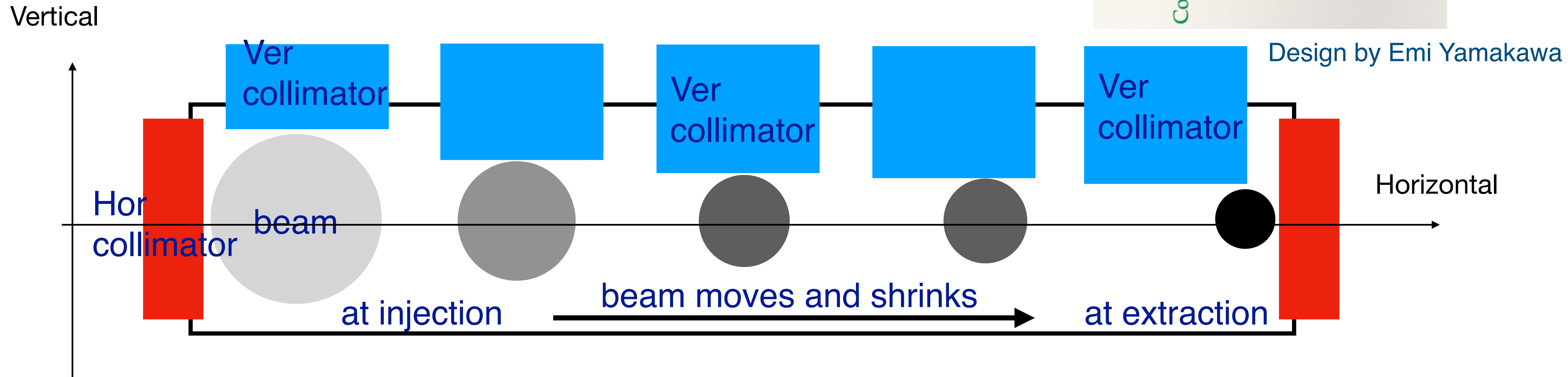
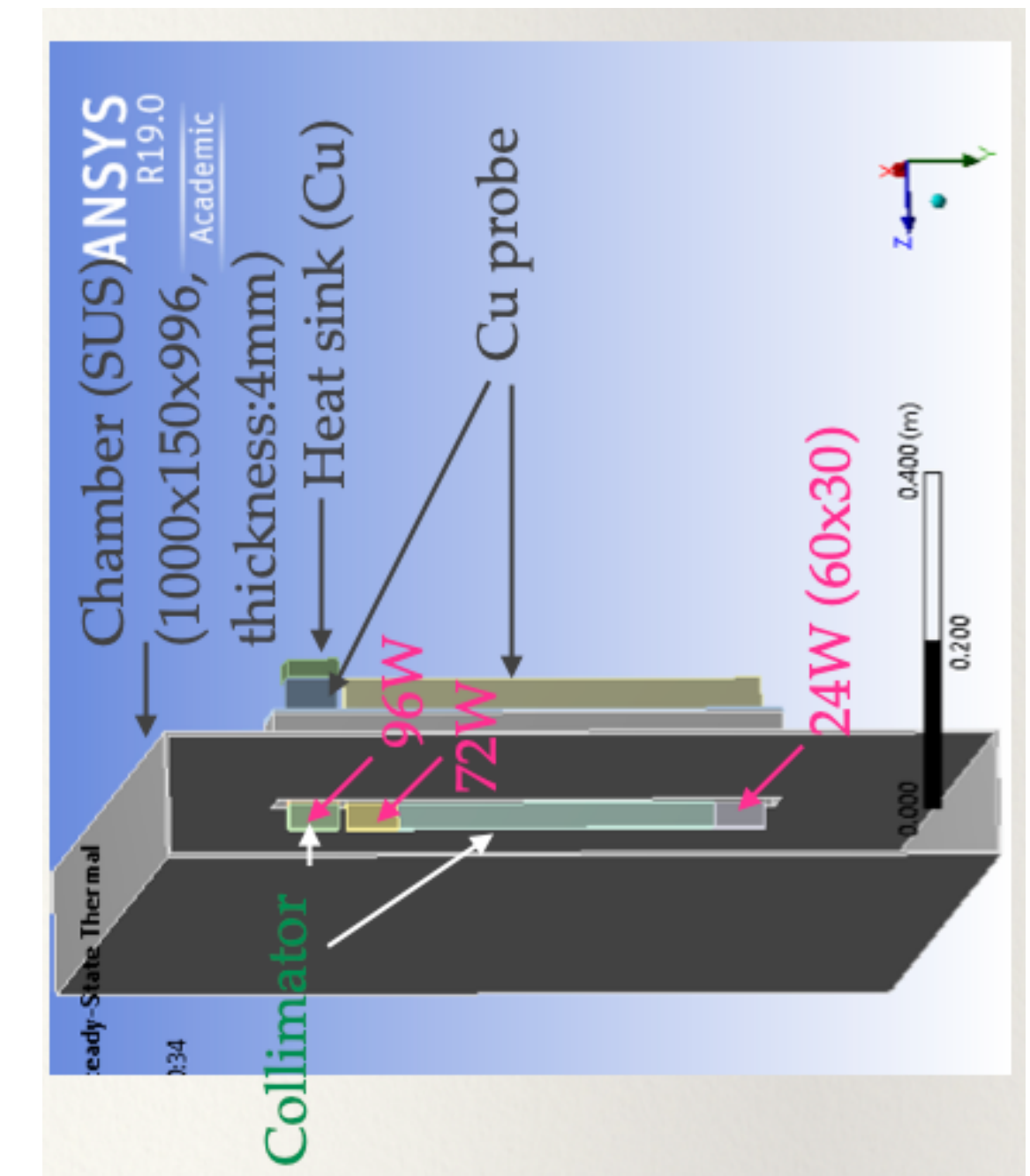


dynamic aperture at 3 MeV (normalised)  
4-fold symmetric lattice

# Collimation

## For all momentum

- In horizontal
  - Collimator at injection (inner side of the aperture) and at extraction (outer side)
- In vertical
  - Collimator continuously or stepwise for all the momentum



Collimate beam halos as soon as developed. No need to wait until extraction.

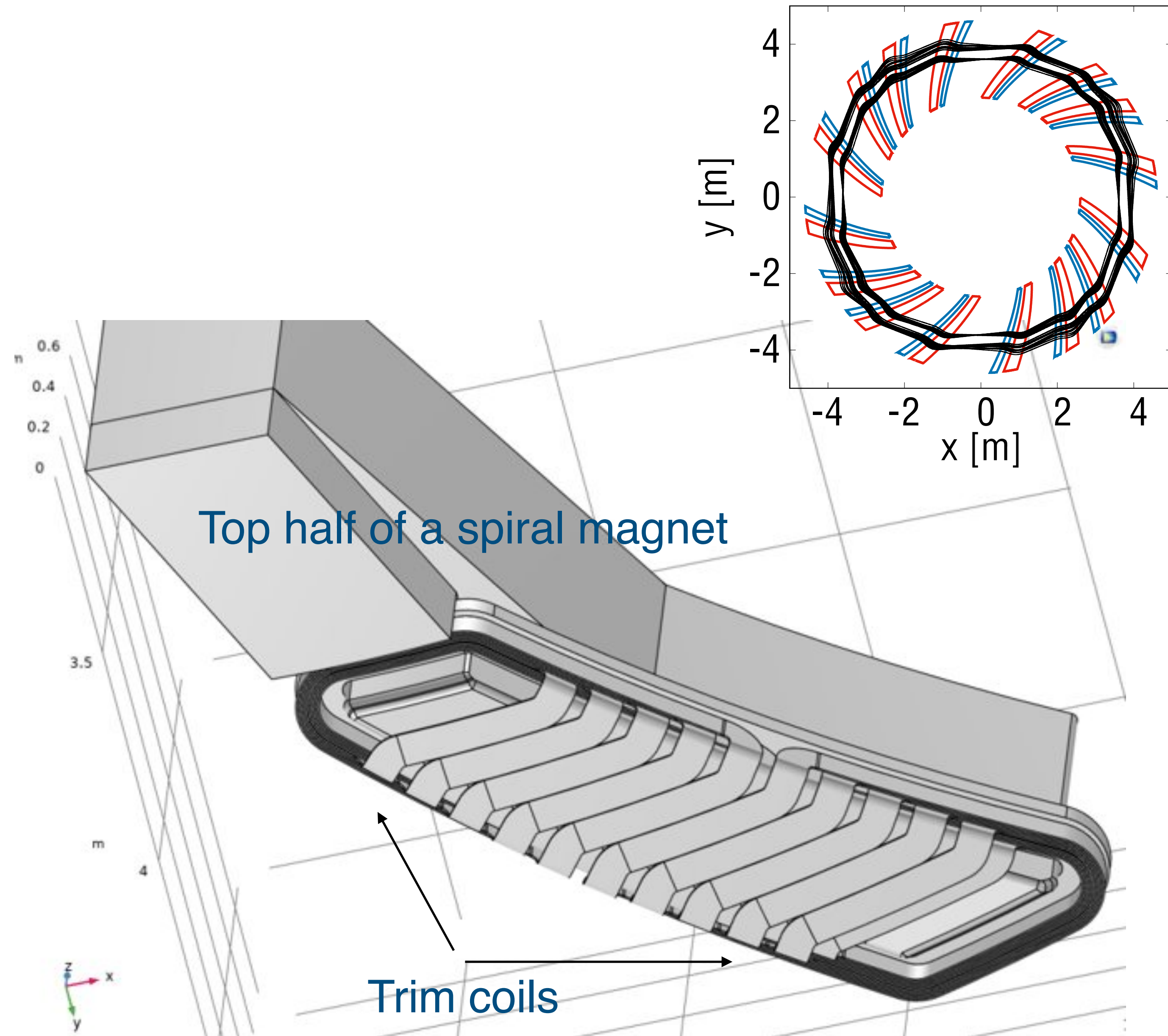
# Correction

## *Orbit, optics, nonlinear harmonic*

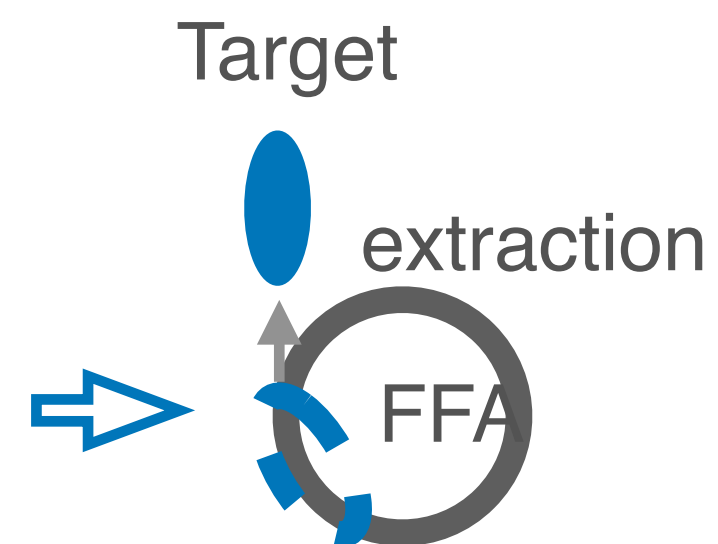
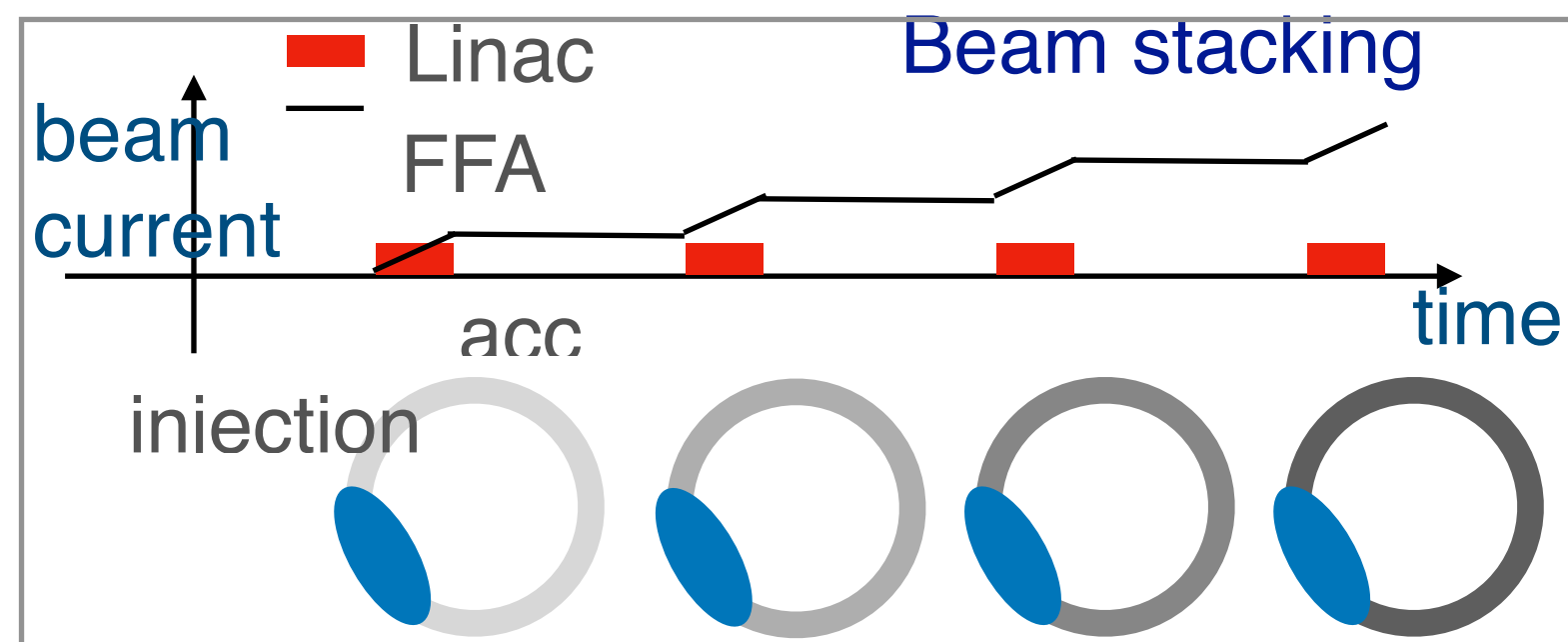
- Each magnet has ~10 trim coil winding on the flat pole.
- **Power supply for each trim coil is independent.**
- Primarily coils are adjusted to make the ideal field.

$$B_z(r, \theta) = B_{z0} \left( \frac{r}{r_0} \right)^k F(\theta)$$

- Additionally small adjustments to excite a harmonic component to correct orbit, optics and nonlinear harmonics.
- How accurately ~10 coils can correct is still a question.



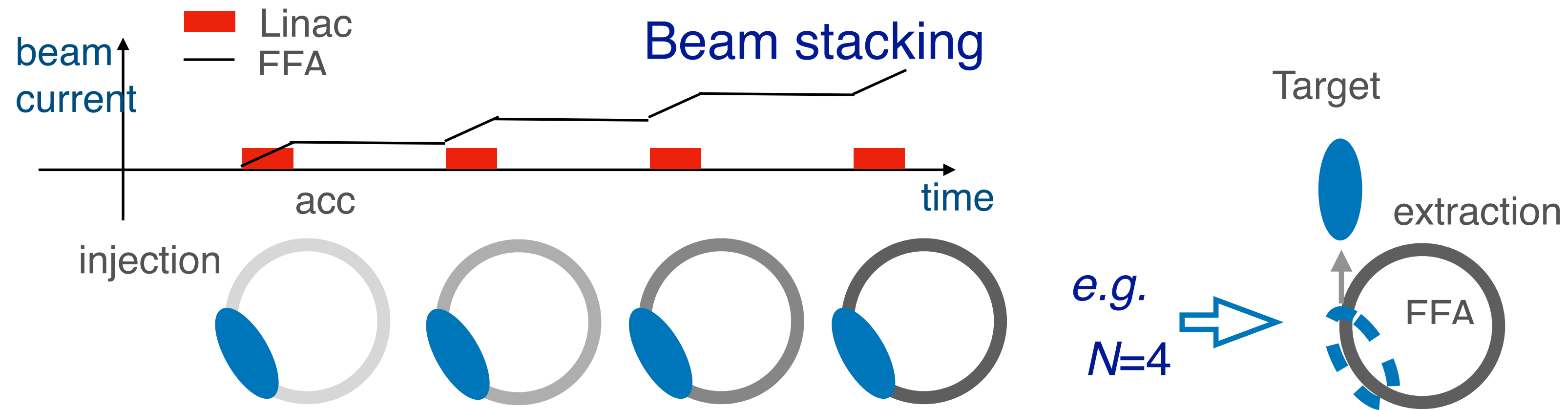
# Beam stacking experiment



**By beam stacking, the pulsed peak current can be increased** keeping the average power with a lower repetition rate ( $\sim 10$  Hz).

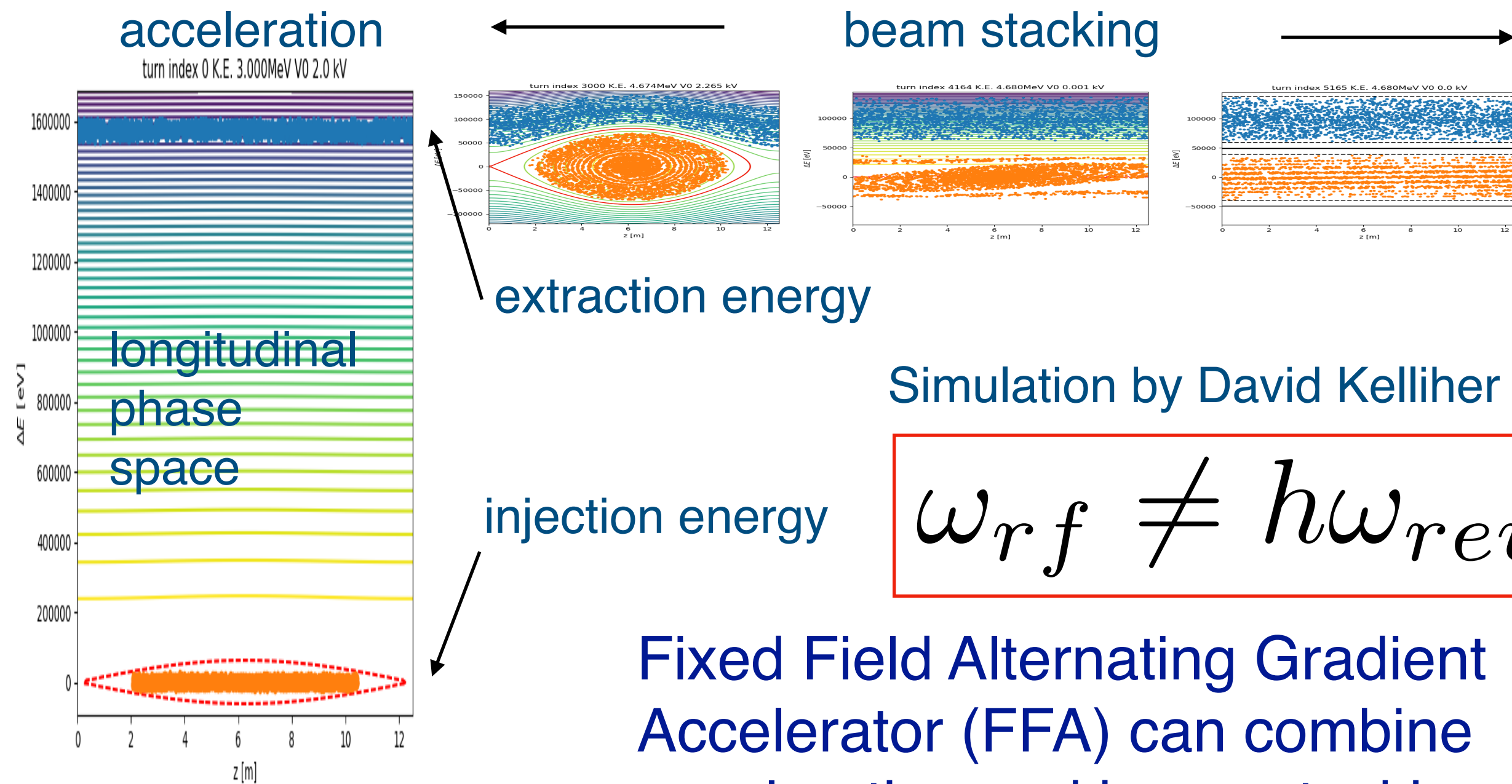
- As a proton drive for a muon collider, spallation neutron source, etc.

# Beam stacking



## Benefits

- Bottleneck to achieve high beam power exists at injection energy.
- By beam stacking, beam power is not limited at injection.
- Repetition rate of an accelerator (120 Hz) can be different from that users will see (30 Hz).
- **Longitudinal emittance is proportional to # of stacking (or larger).**



**Beam stacking is a way to make high peak power with low repetition.**

- No other ring accelerators can make that peak power.
- Good as a proton driver for a muon collider, spallation neutron source, etc.



# Beam stacking

*First experiment at MURA in 1960s*

A beam is injected.

A beam is captured and accelerated.  
Some of particles are not captured.

Repeat 4 times. Momentum spread is larger.

High energy

Low energy

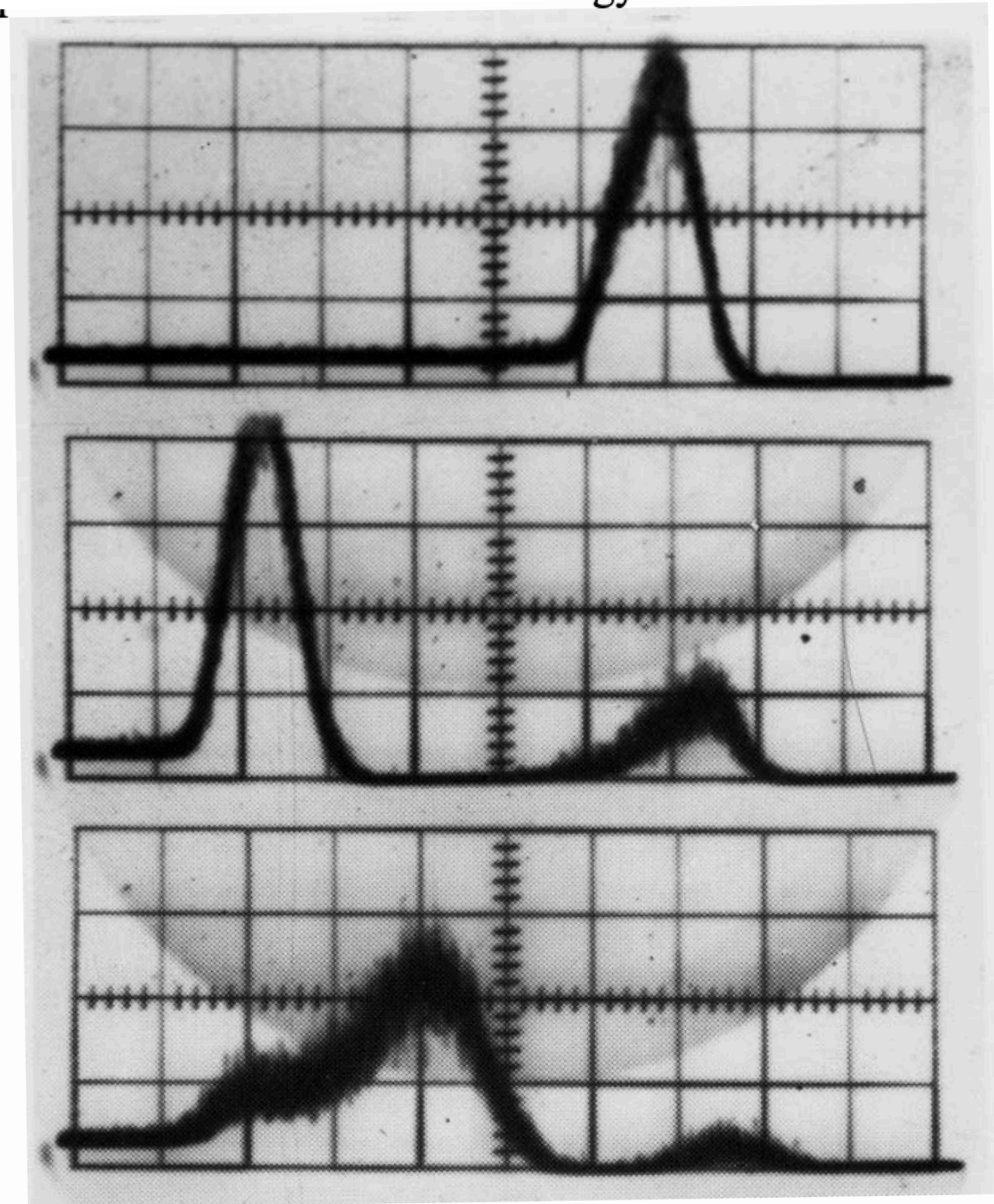


Figure 6: Beam Stacking Experiment

# Beam stacking

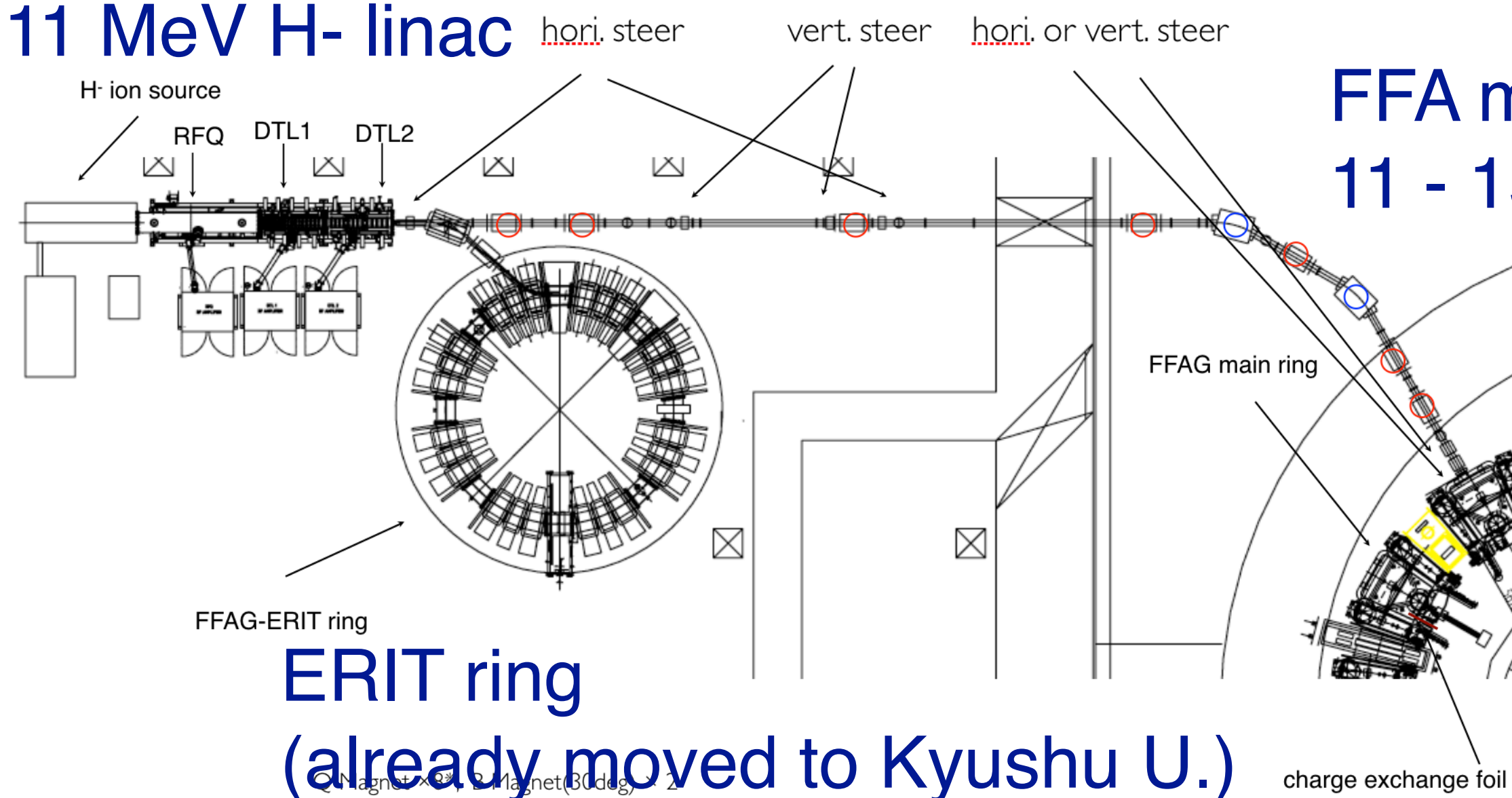
## *New experiment this year*

Experimental demonstration (of 2 beams)

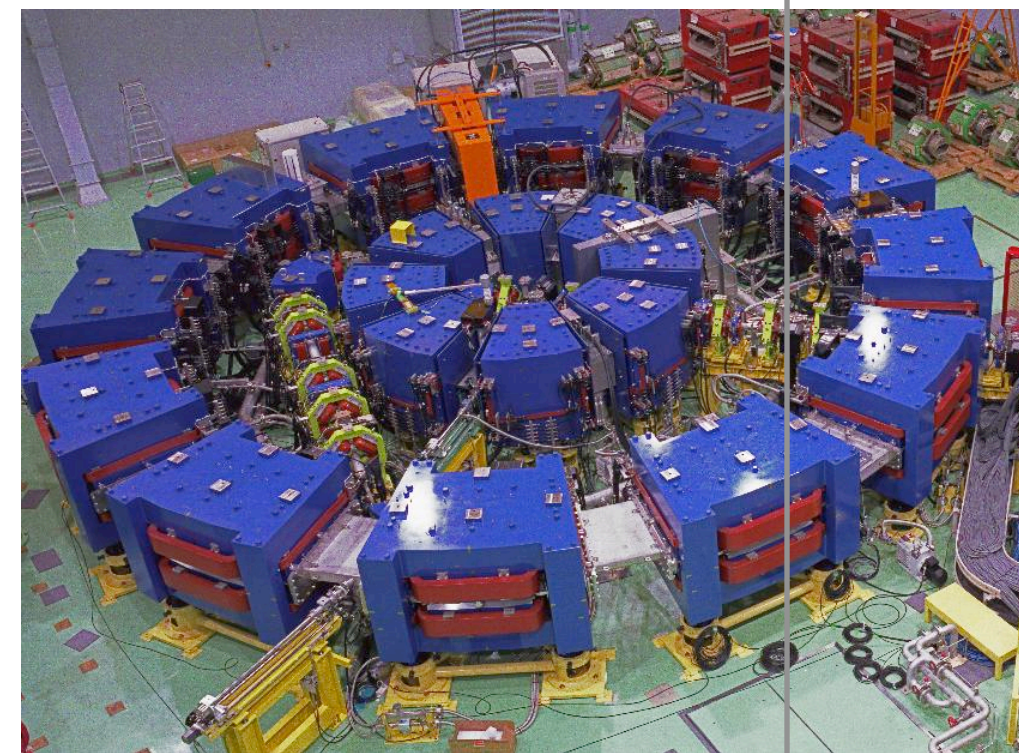
- Is the total **momentum spread  $dp/p$**  2 x each beam?
- Is the total **number of particles** is 2 x each beam?

## KURNS FFA accelerator complex at Kyoto Univ.

### 11 MeV H- linac

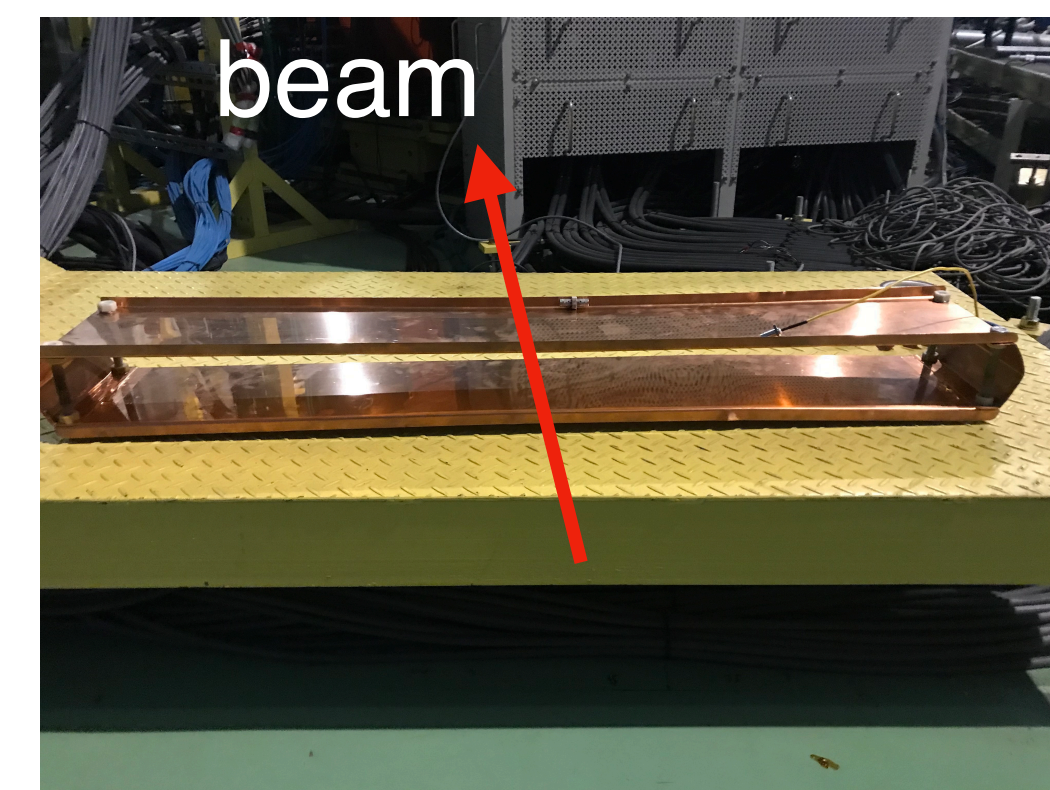
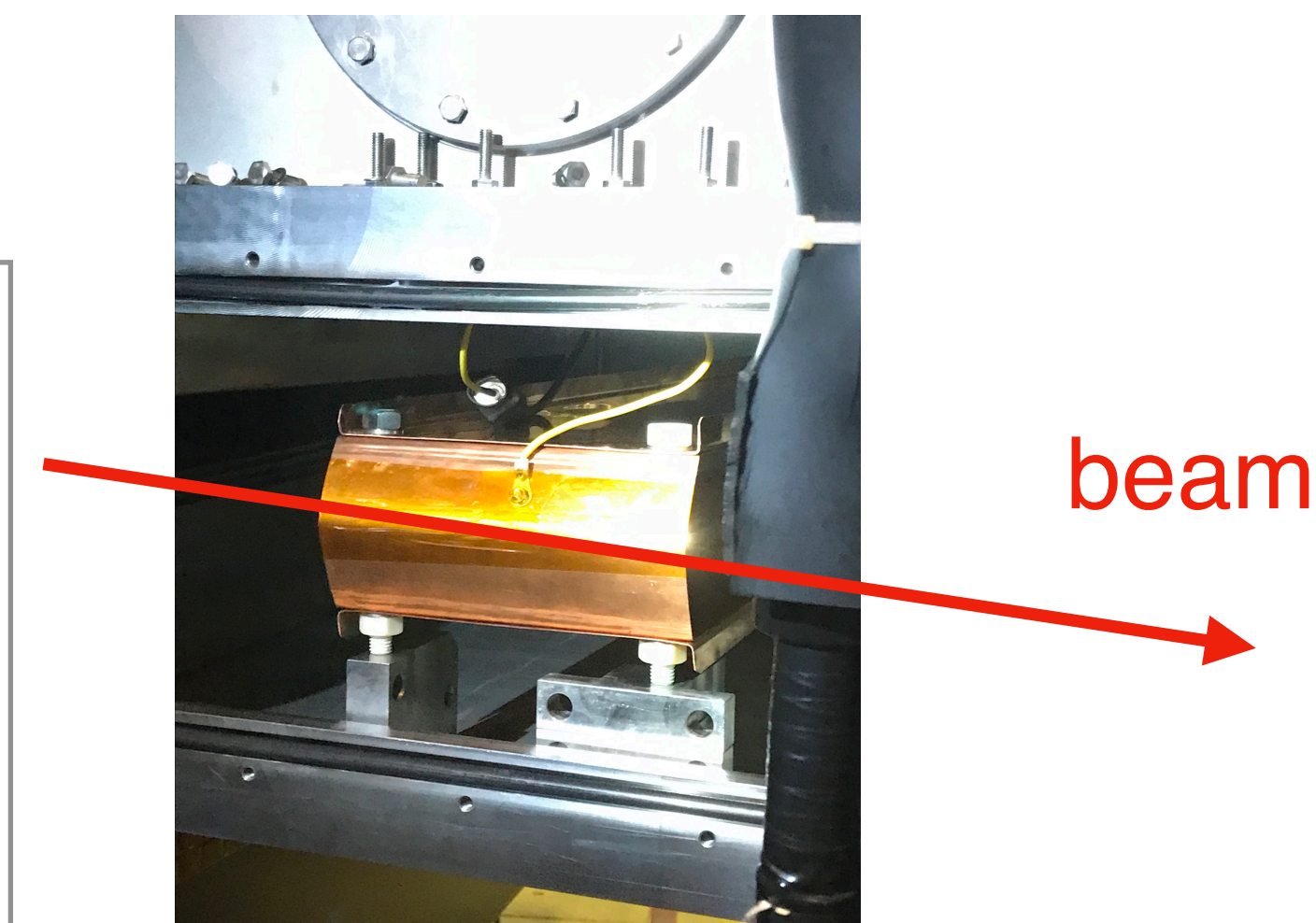


### FFA main ring 11 - 150 MeV



## Full Aperture Bunch (FAB) monitor

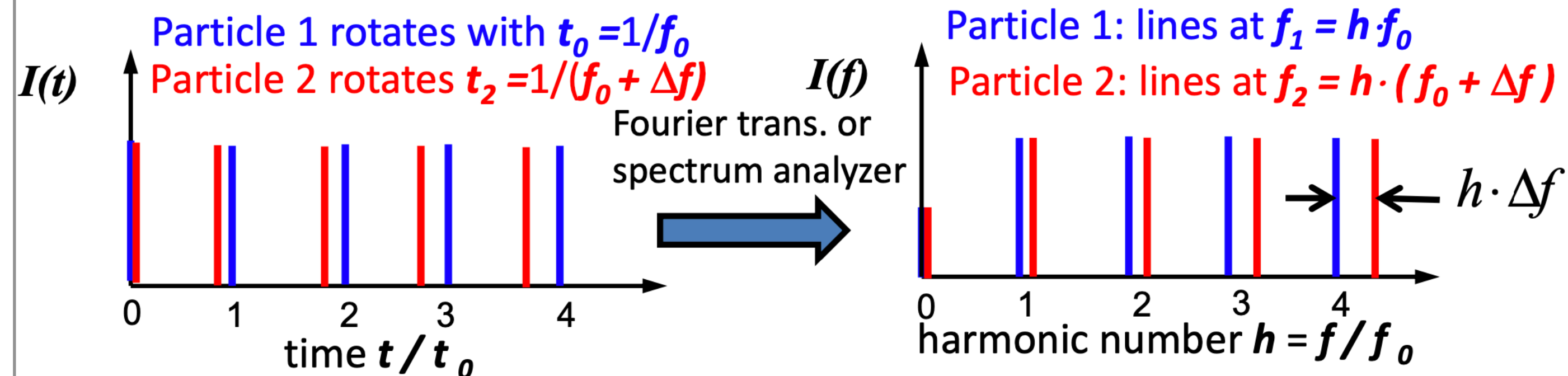
- Pickup bunch structure.
- Signal is amplified to the scope.



# Schottky signal analysis

A Beam consists of finite number of particles

IBIC 2017, Peter Forck



- Momentum spread is seen by different revolution time.
- Spread of frequency spectrum at each harmonic  $h$  can be measured

$$\frac{dp}{p} = \frac{1}{h\eta} \frac{df}{f} \quad \eta : \text{slippage factor}$$

- This is an incoherent signal.
- Sum of frequency spectrum (more precisely PSD) is proportional to the number of particles

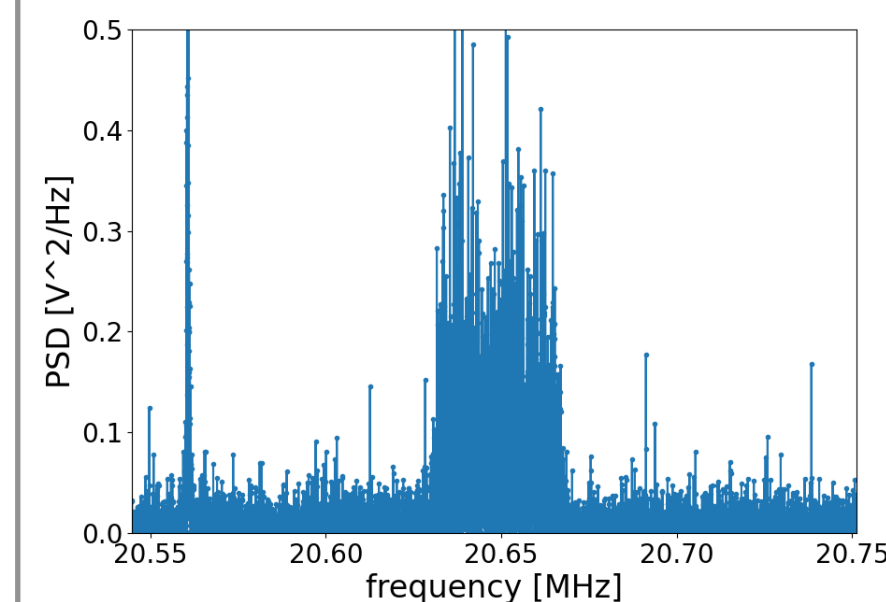
$$\int \left( \frac{dP}{df} \right) df = 2Z_t e^2 f_0^2 \int \left( \frac{dN}{df} \right) df$$

$Z_t$ : transfer impedance



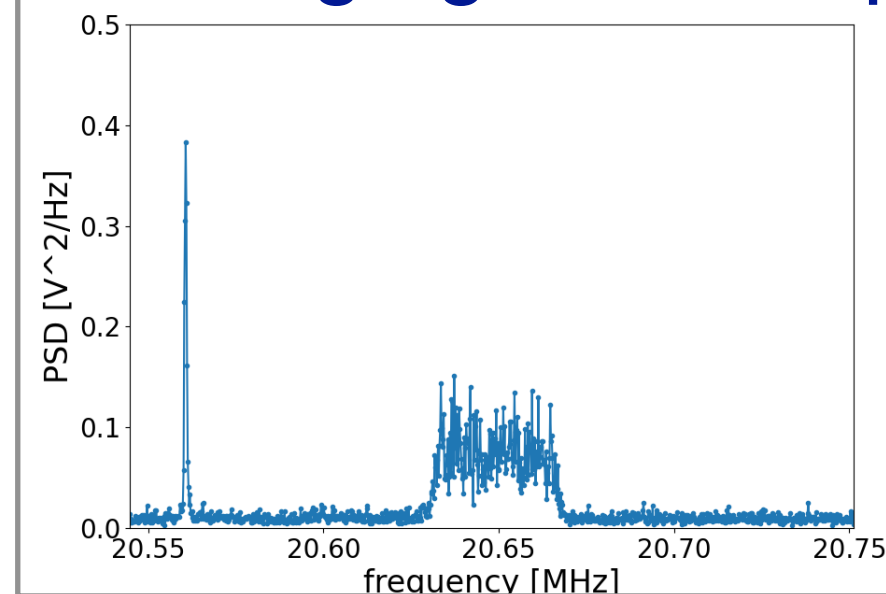
Schottky signal and PSD as an output tells 1) beam intensity and 2) momentum spread

## Power Spectrum Density (PSD)

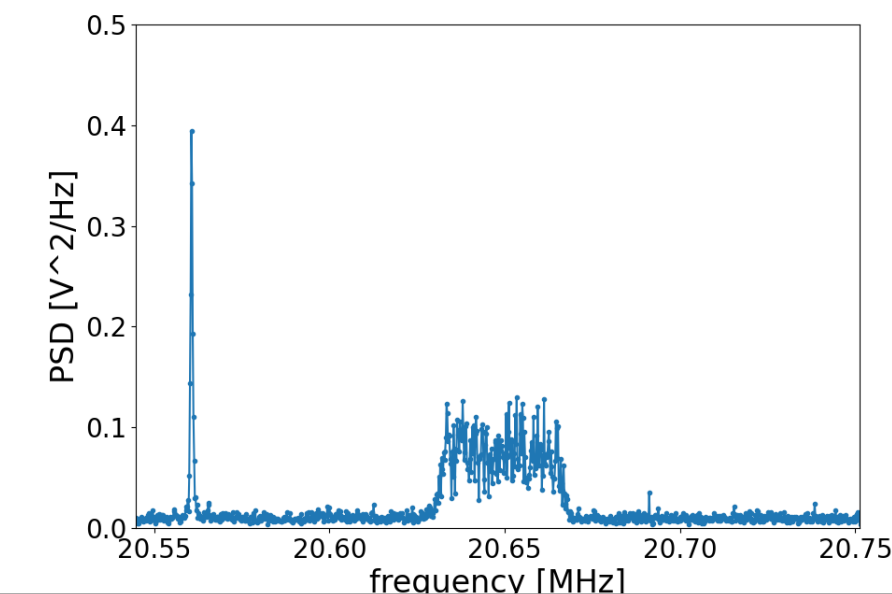


FFA spectrum  
(Vertical axis is power  $V^2$ )

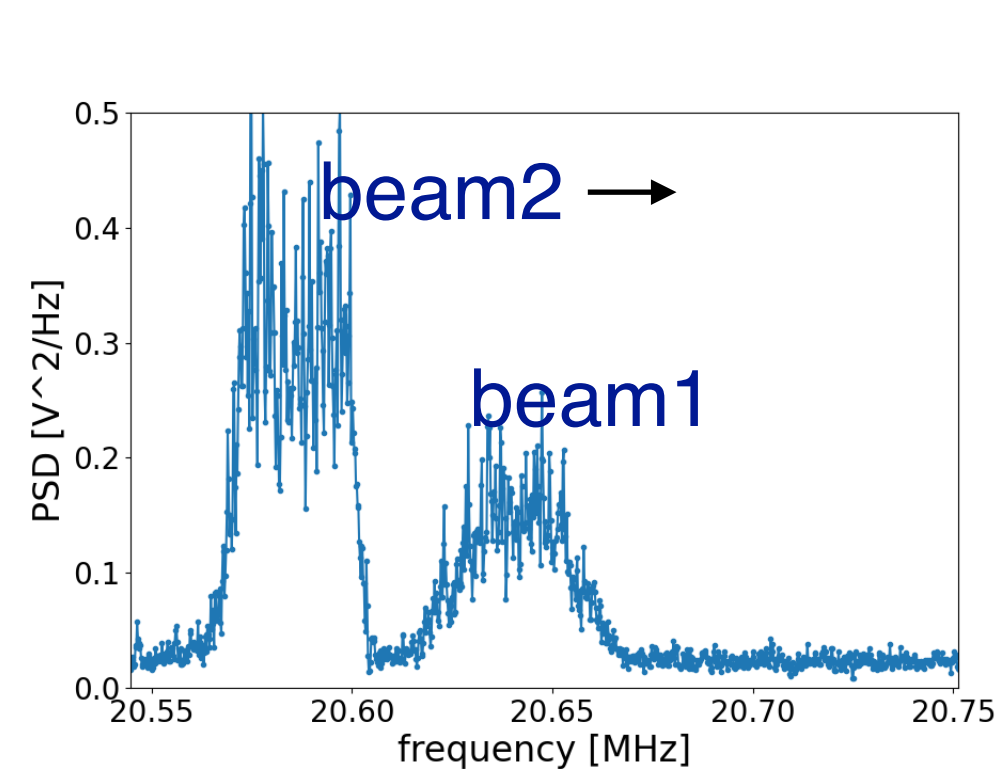
averaging over freq



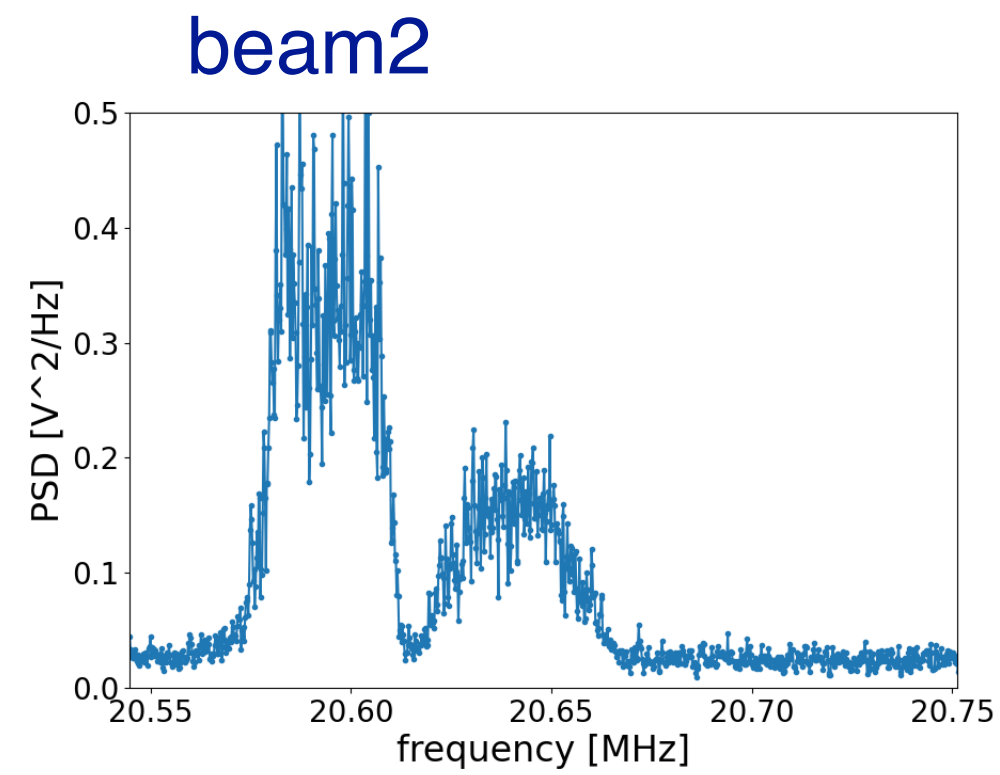
averaging over time  
"Bartlett" or "Welch"



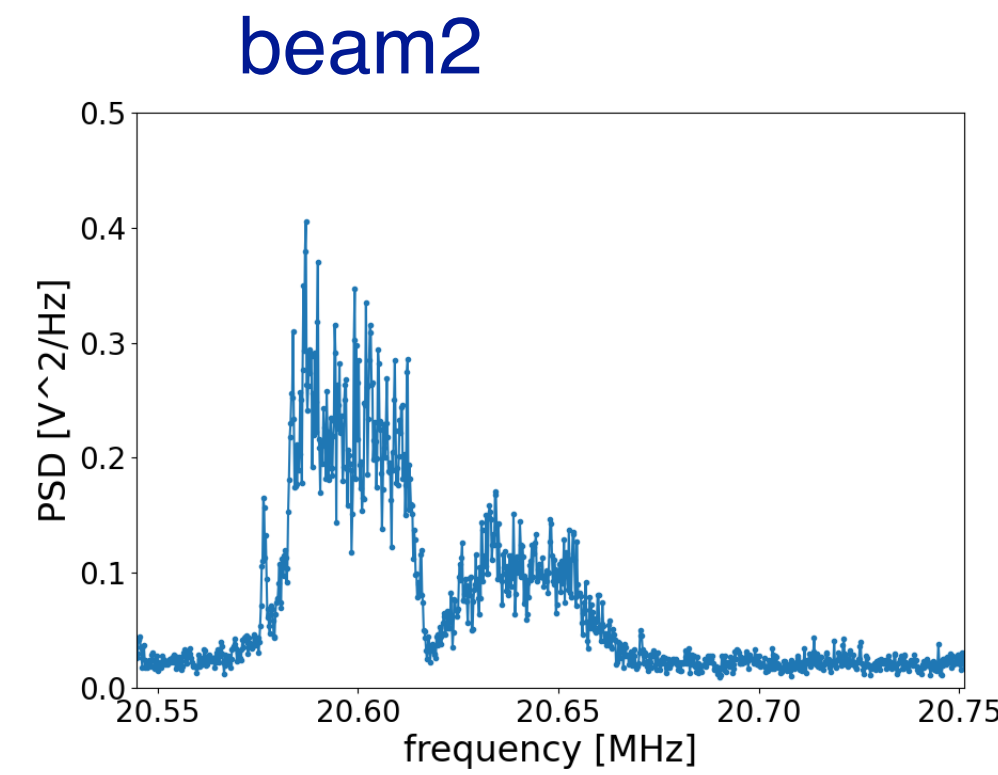
# Schottky signal as a function of the final energy of beam 2.



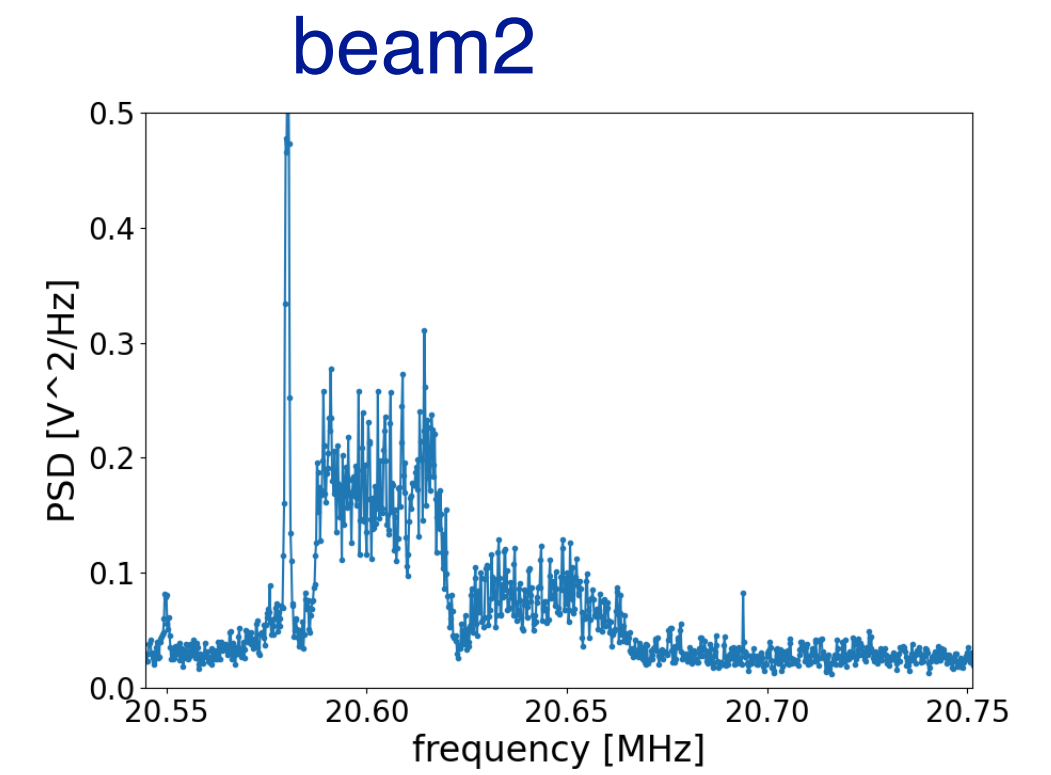
-235.4 keV



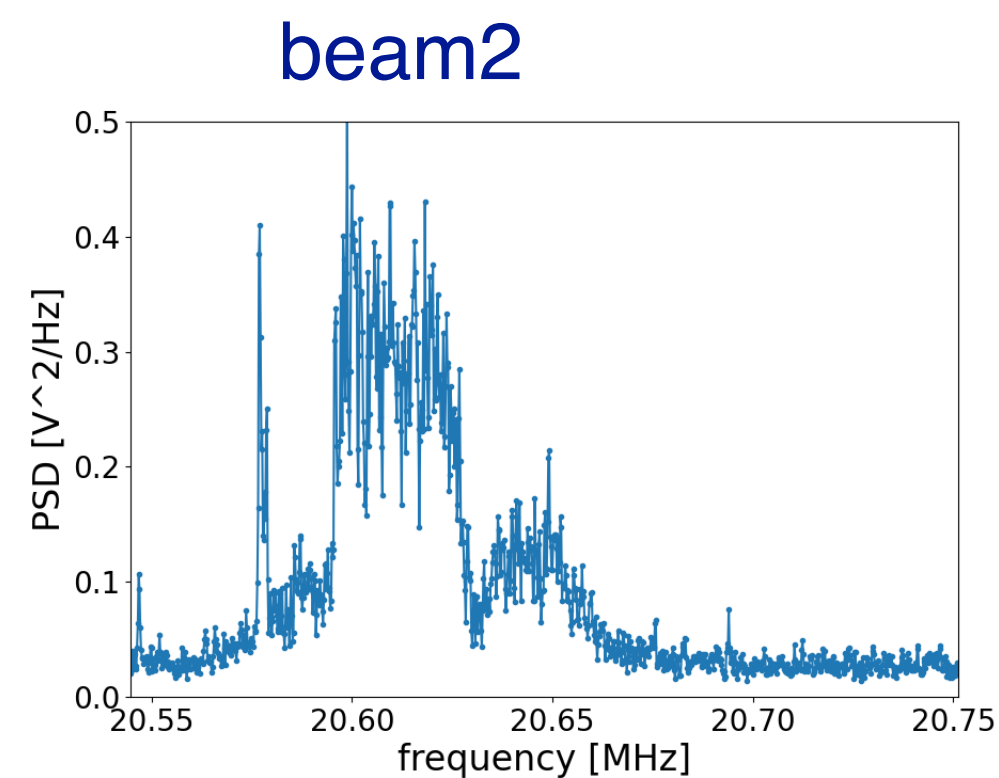
-200.6 keV



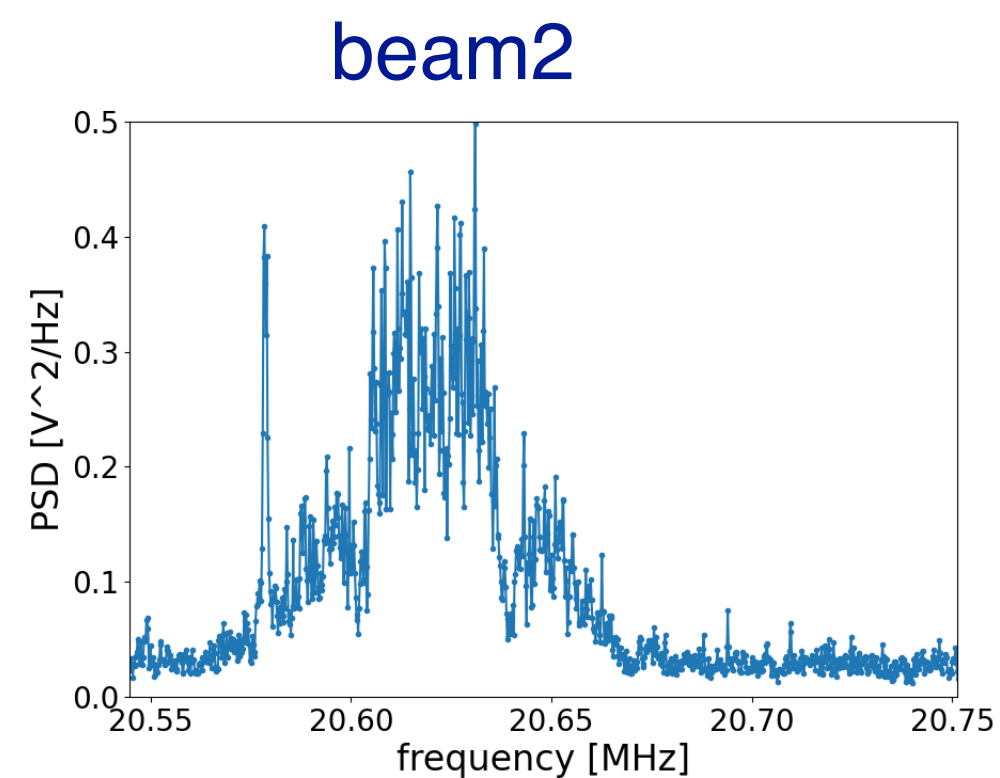
-182.5 keV



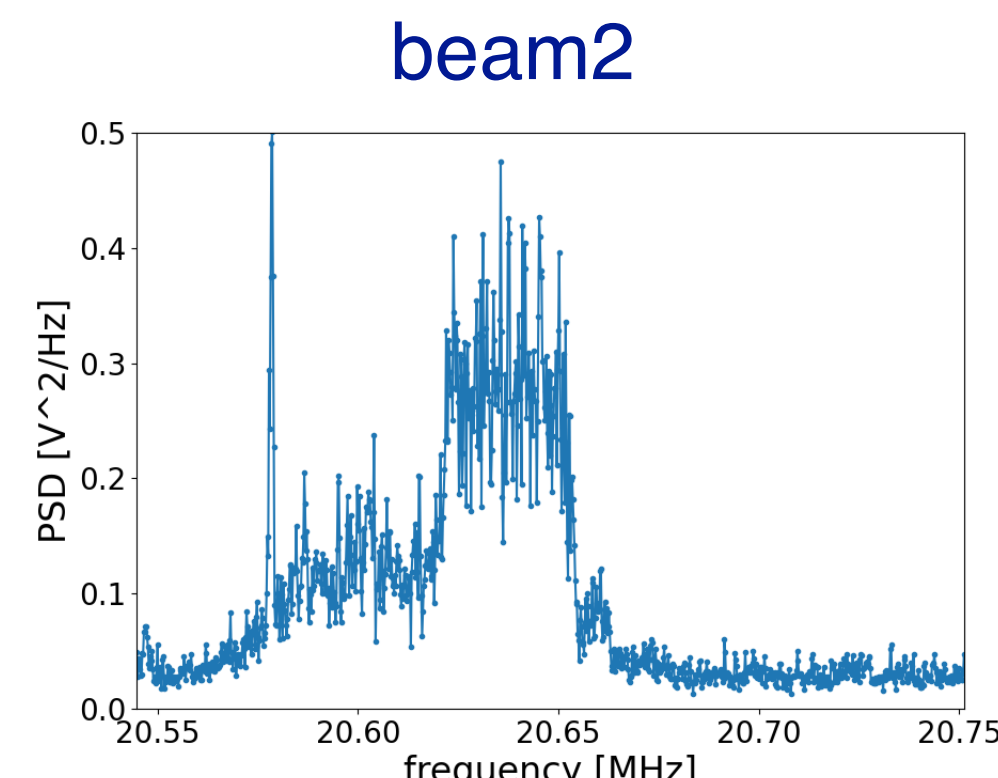
-165.0 keV



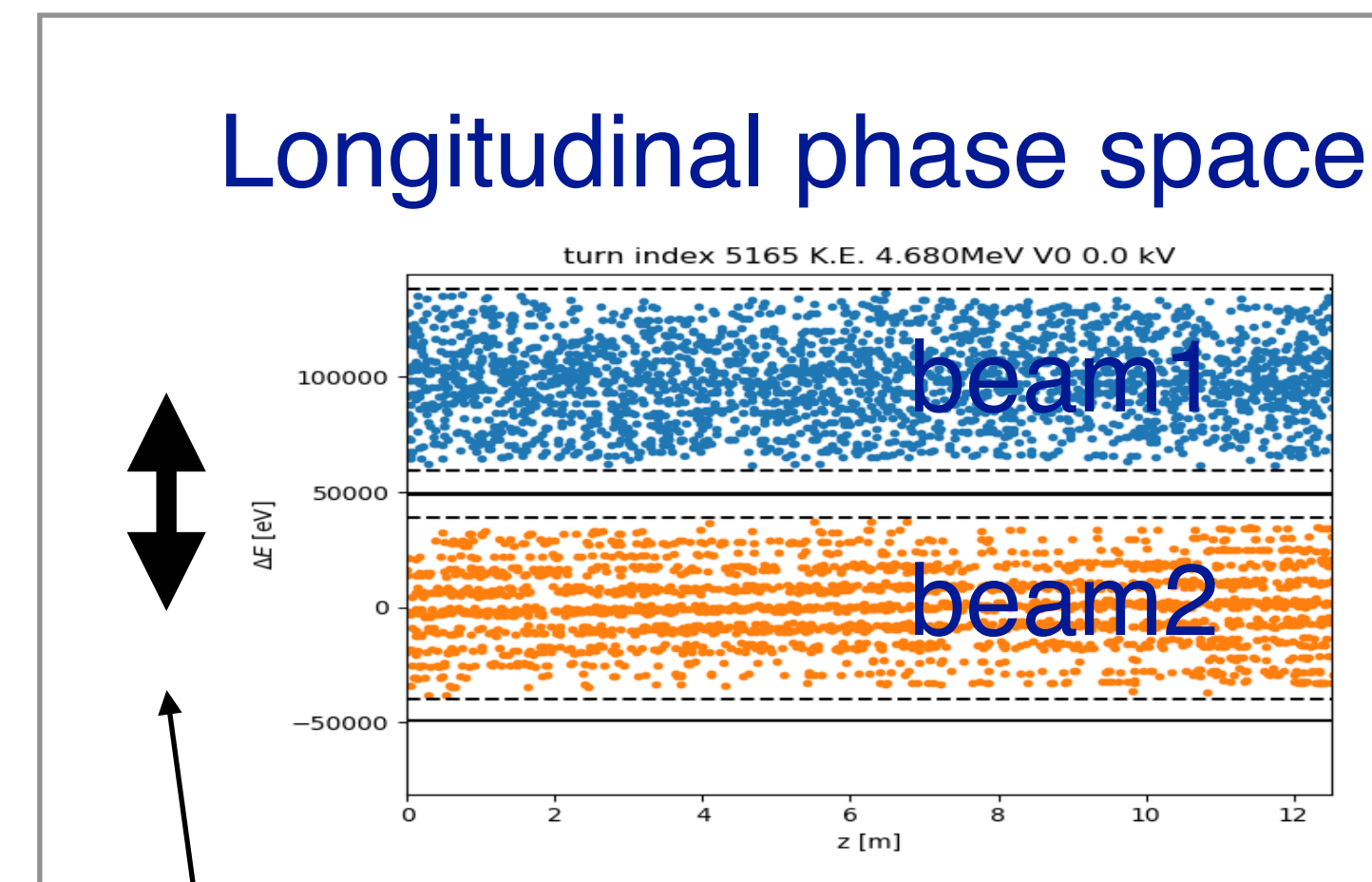
-129.7 keV



-94.4 keV

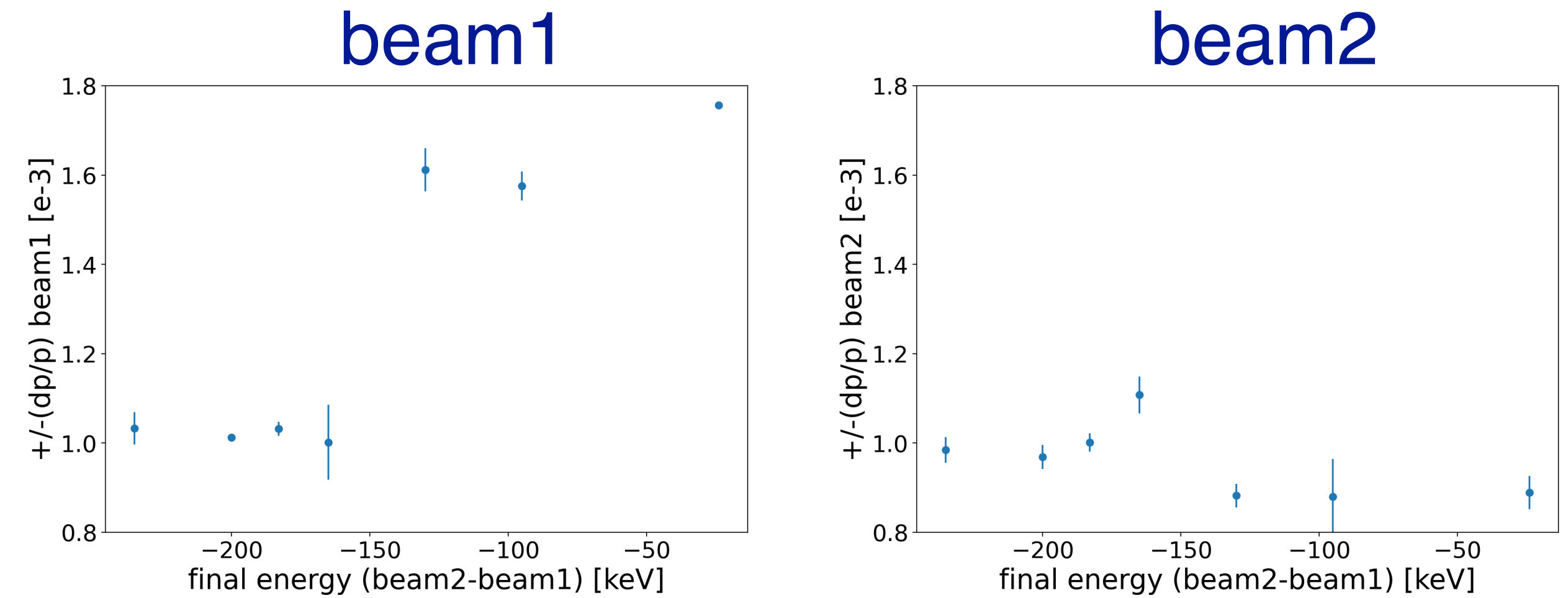
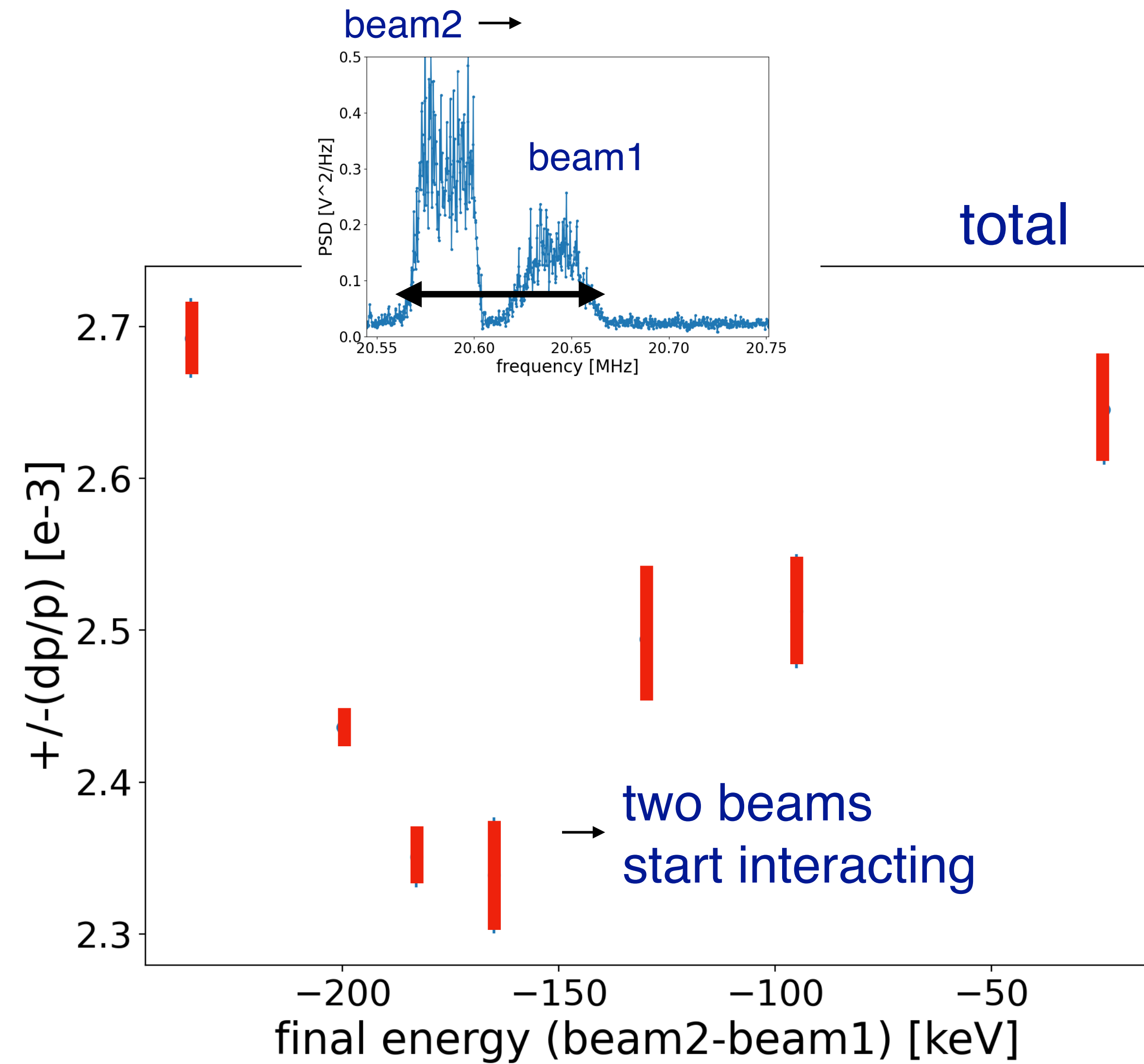


-23.8 keV



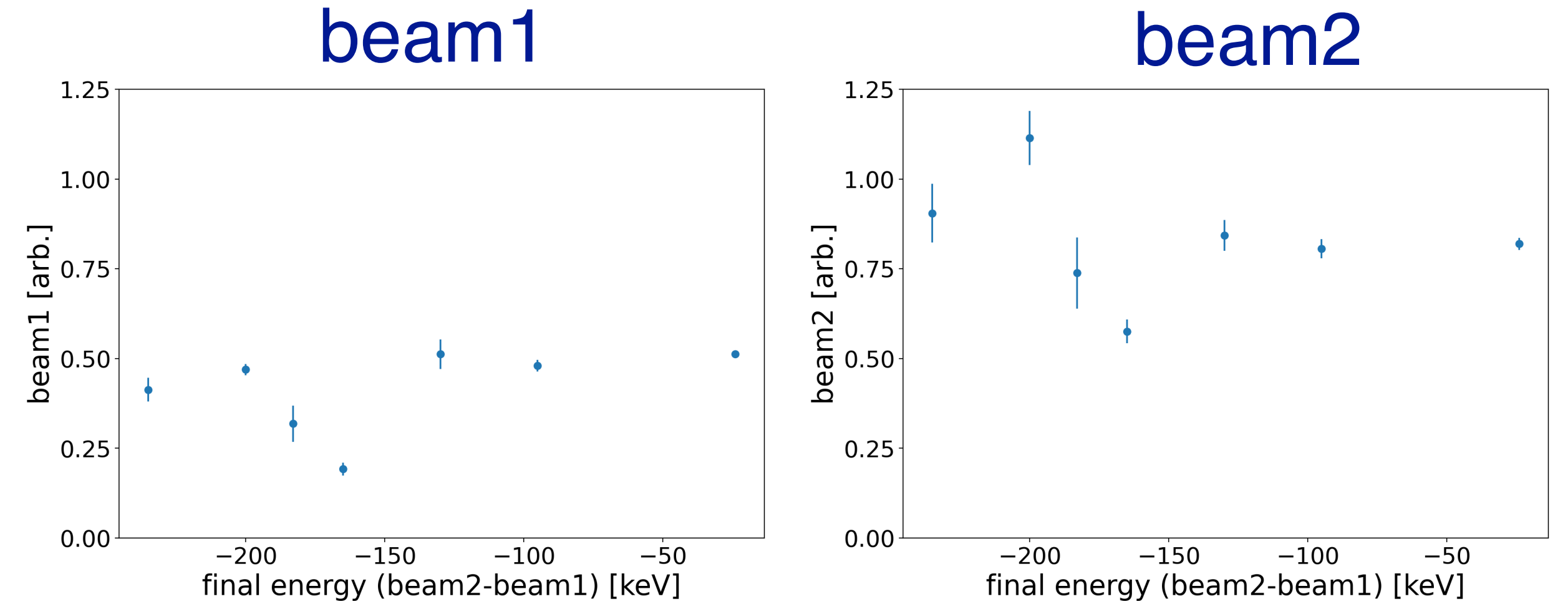
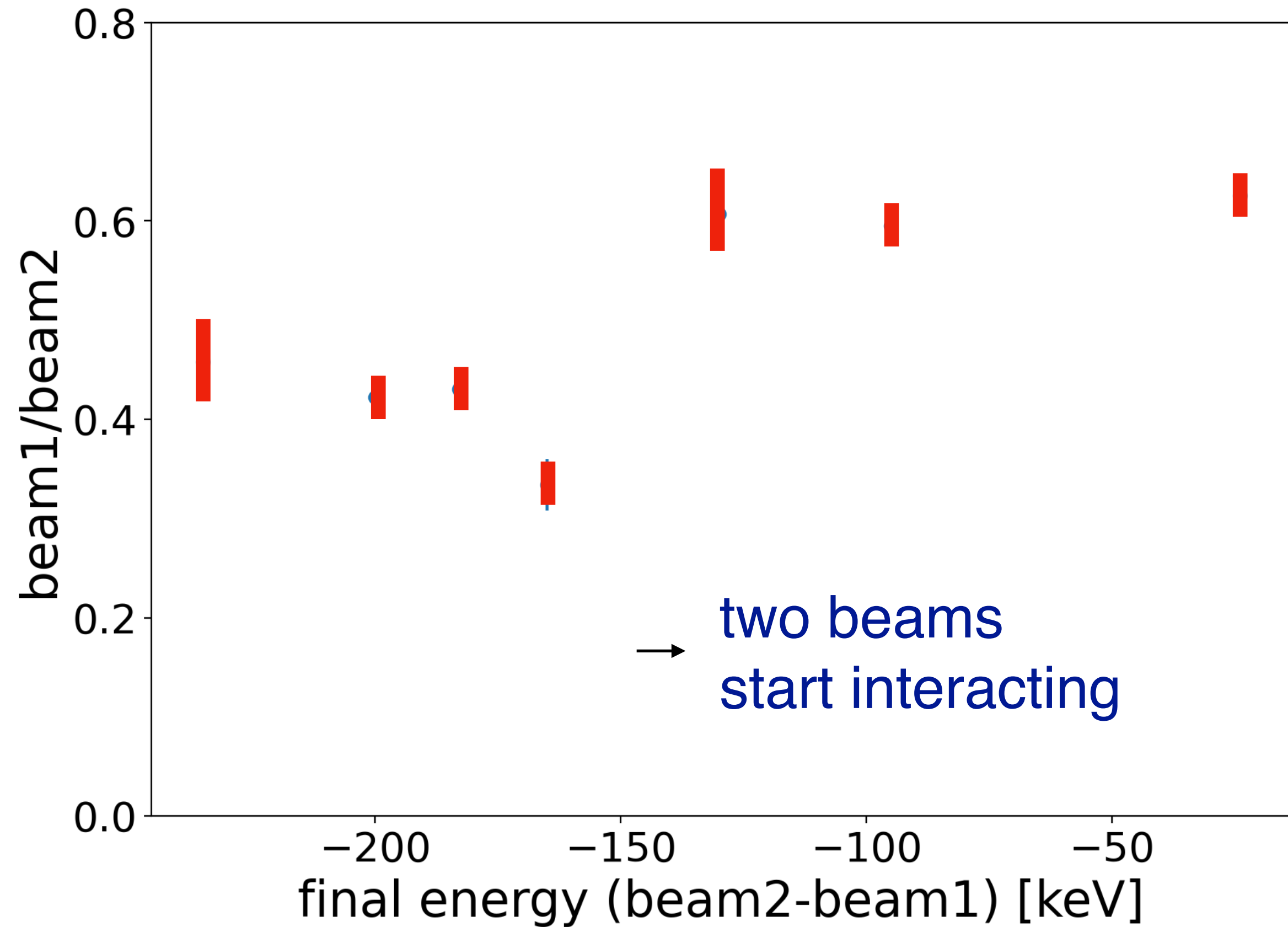
Energy separation (beam2-beam1)

# Result 1: Momentum spread $dp/p$



- **Total  $dp/p$  becomes minimum** just at the point where two beams start interacting.
- Once two beams interact each other, total  $dp/p$  is larger than twice of  $dp/p$  of each beam.
- $dp/p$  of each beam is unchanged until two beams start interacting.

# Result 2: Beam intensity



- Until two beams start interacting, **the ratio of beam 1 and beam 2 is about 40%** independent of final energy of beam 2.
- Ratio of beam 1 and beam 2 looks higher with interaction, but it depends on the definition of beam 1 intensity. In this analysis, beam 1 includes intensity of both sides of beam 2.

**Intensity of beam1 (waiting at top energy) is significantly reduced.**

# RF knock out

## Similar to synchro-beta resonance

When the RF cavity is located at the finite dispersion point  $D_x$ , energy gain induces horizontal displacement.

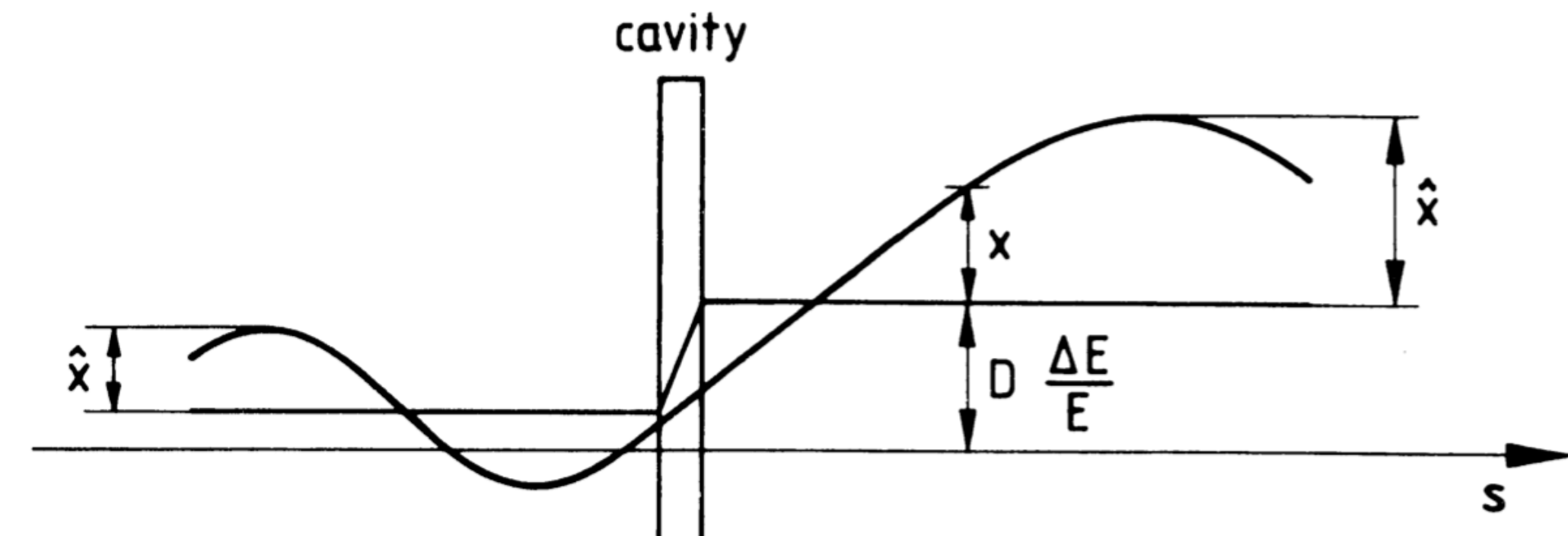
In a bunched beam, energy gain or induced horizontal displacement has a frequency of synchrotron oscillation and its higher harmonics.

$$\delta x = -D_x \frac{dp}{p} = -\frac{D_x}{2} \frac{dT}{T} = -\frac{\pi D_x V_0 a_s}{T \lambda} \cos(\omega_s t)$$

For the stacked (coasting) beam,

$$\delta x = -D_x \frac{dp}{p} = -\frac{D_x}{2} \frac{dT}{T} = -\frac{D_x V_0}{T} \cos(\omega_{rev} - \omega_{rf}) t$$

When it becomes the same frequency of (horizontal) betatron oscillations, **resonance occurs**.

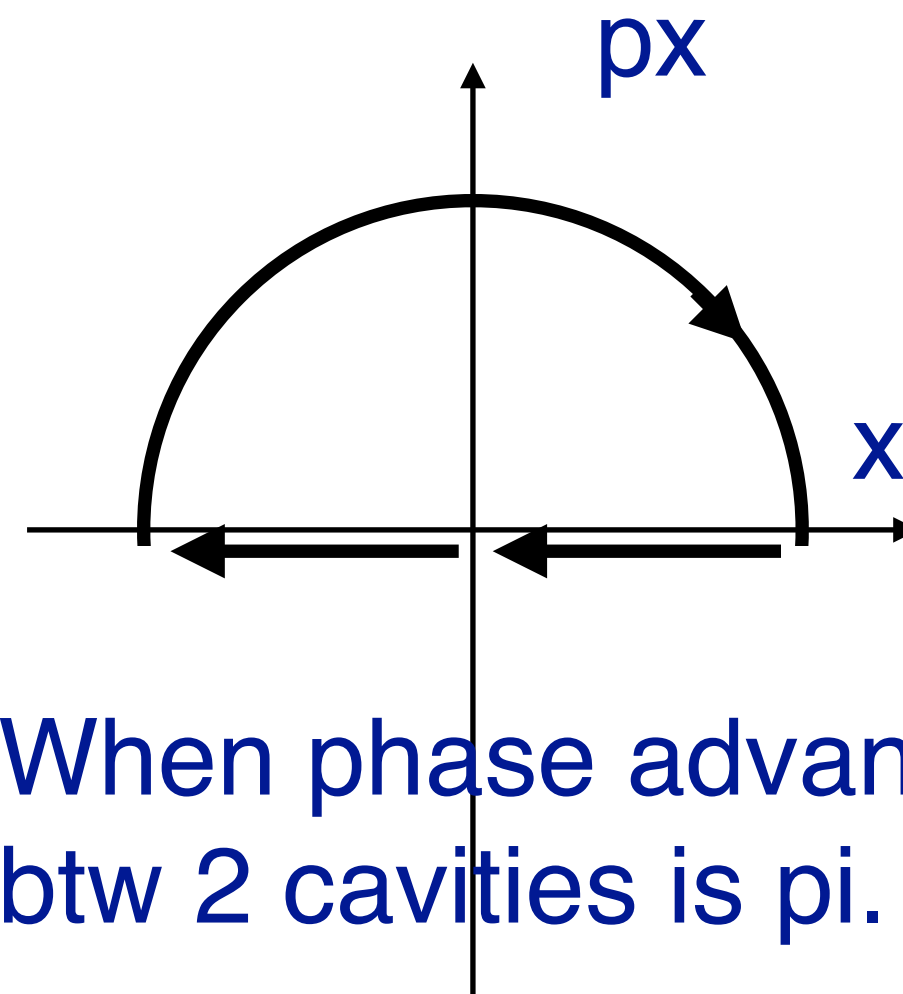


from CERN-87-03

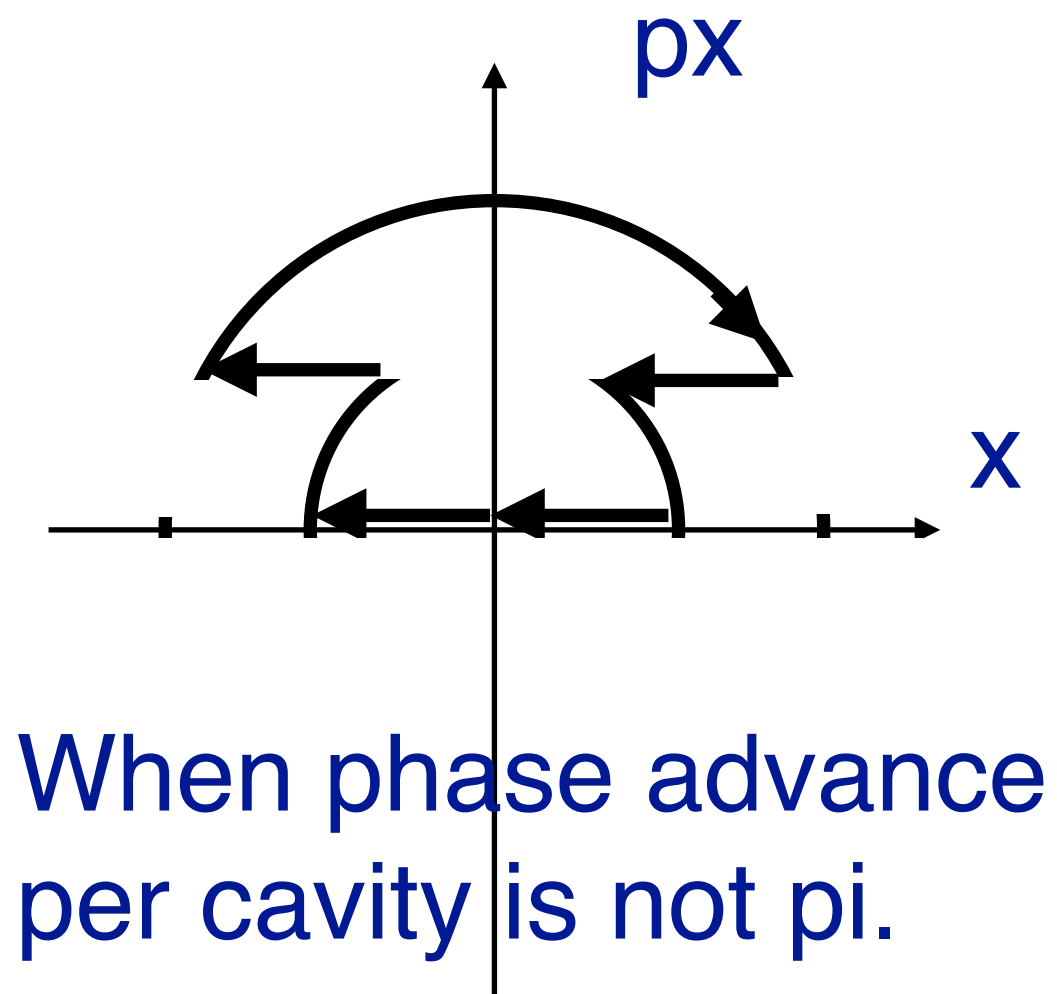
$$\frac{\omega_{rev} - \omega_{rf}}{\omega_{rev}} = \frac{\omega_{\beta,h}}{\omega_{rev}} \quad \text{or} \quad 1 - \frac{\omega_{\beta,h}}{\omega_{rev}} \quad \text{where} \quad \frac{\omega_{\beta,h}}{\omega_{rev}} = Q_{\beta,h}$$

# Proposed mitigation methods (from MURA papers)

- For a ring with single RF cavity
  - Reduce voltage around resonance
  - Control betatron phase around resonance by changing tune for short time (like a jump around transition energy crossing).
- For a ring with two RF cavities
  - **Choose a proper betatron phase advance between two cavities**
  - Tipped RF cavities to cancel transverse fields
- For a ring with multiple RF cavities
  - Place cavities with equal spacing.
  - Place cavities with proper phase.



When phase advance  
btw 2 cavities is  $\pi$ .



When phase advance  
per cavity is not  $\pi$ .



# Before summary

## *Many studies to be done*

- **Impedance** calculation of a wide window shape vacuum chamber, even irregular shape at some points.
- **Instability and its mitigation**
  - Acceleration is fast.
  - Beam stacking requires the stability of high current coasting beams for long time.
- Zero chromaticity for the entire energy to keep the tune constant.
- Chromaticity is not a knobs to control instabilities.
- Full of systematic multiples and large amplitude dependent tune shift.
- Does it help?



# Summary

- High intensity is the primary goal of the Fixed Field Accelerator development at the start.
- Many ideas and proposals existed, but hardware was not ready until recently.
- Now time to revisit the initial idea with the state of the art equipment and new technique.
- It has a potential to give the highest peak power without sacrificing the average current.
- Prototype construction of a high intensity Fixed Field Accelerator is the next step.
  - First, beam loss handling with the same space charge level of SNS/JPARC.
  - Second, study beam instability and its mitigation.



# Thank you for your attention



ISIS Neutron and Muon Source

 [www.isis.stfc.ac.uk](http://www.isis.stfc.ac.uk)

  [@isisneutronmuon](https://www.instagram.com/isisneutronmuon)

 [uk.linkedin.com/showcase/isis-neutron-and-muon-source](https://uk.linkedin.com/showcase/isis-neutron-and-muon-source)



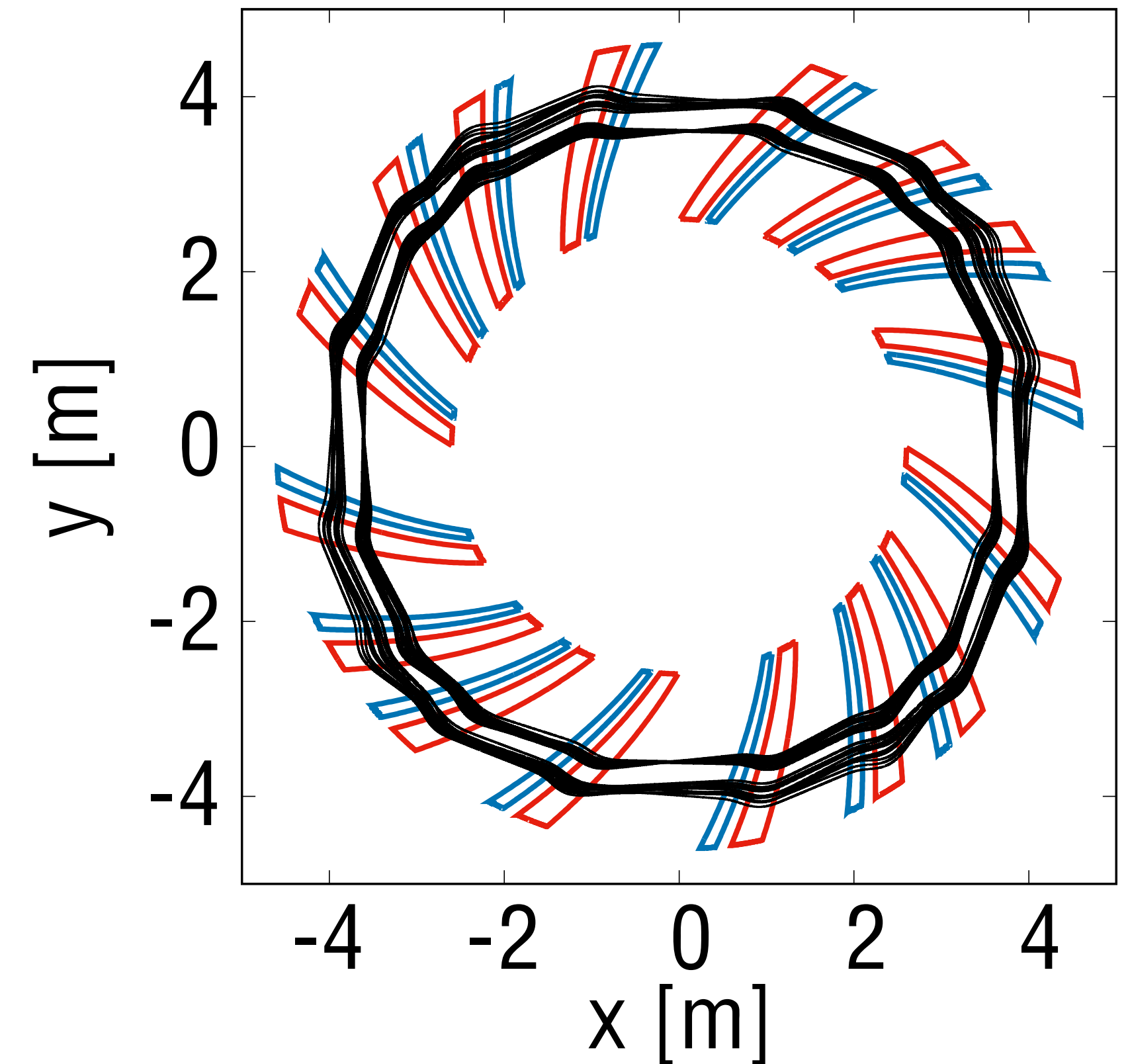
# Backup

# Modelling space charge effects



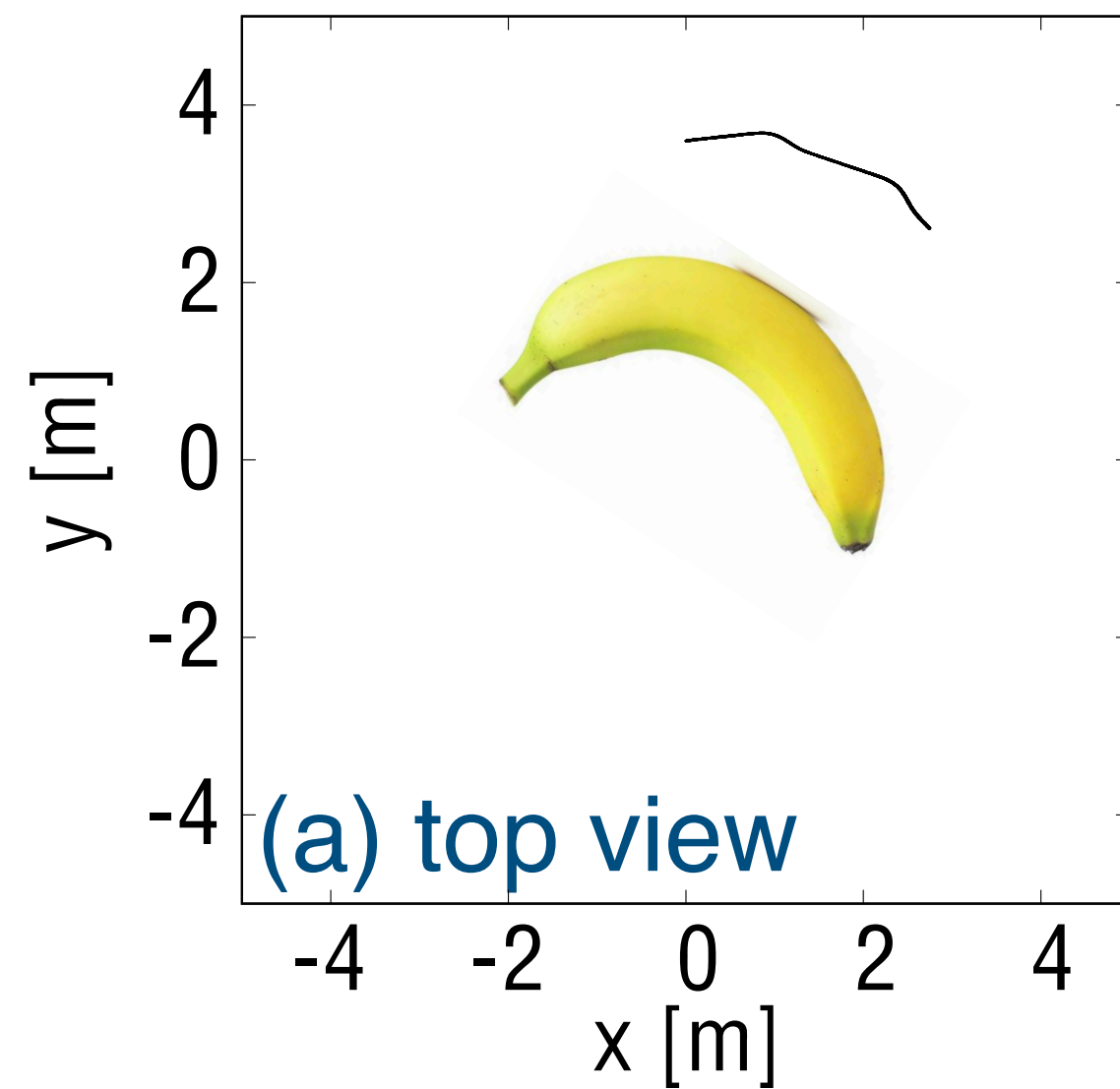
# Space charge modelling in an FFA

- Equilibrium orbit is **a function of time** (momentum) and **operating point**.
  - Equilibrium orbit is fixed in a synchrotron.
- Important to know where the centre of charge distribution in order to calculate space charge effects.
  - Perturbation to betatron oscillation frequency matters.
- A bunch occupies **the large fraction of the circumference**,  $1/2 \sim 1/4$ . The longitudinal size is much larger than the transverse.
  - A beam size has a similar aspect ratio in 3D in a cyclotron.

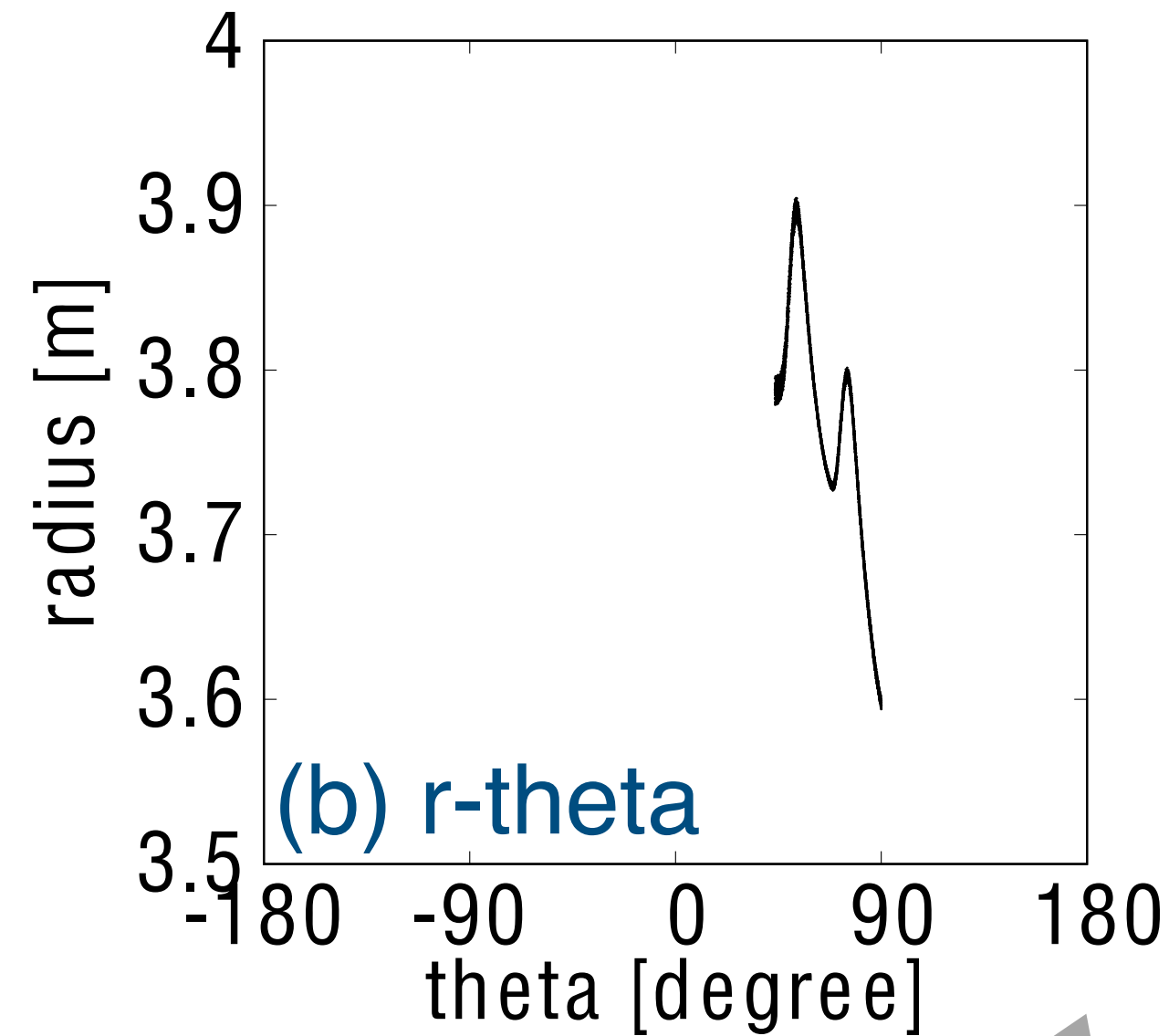


# Scode's way of modelling

- Beam in arbitrary position.

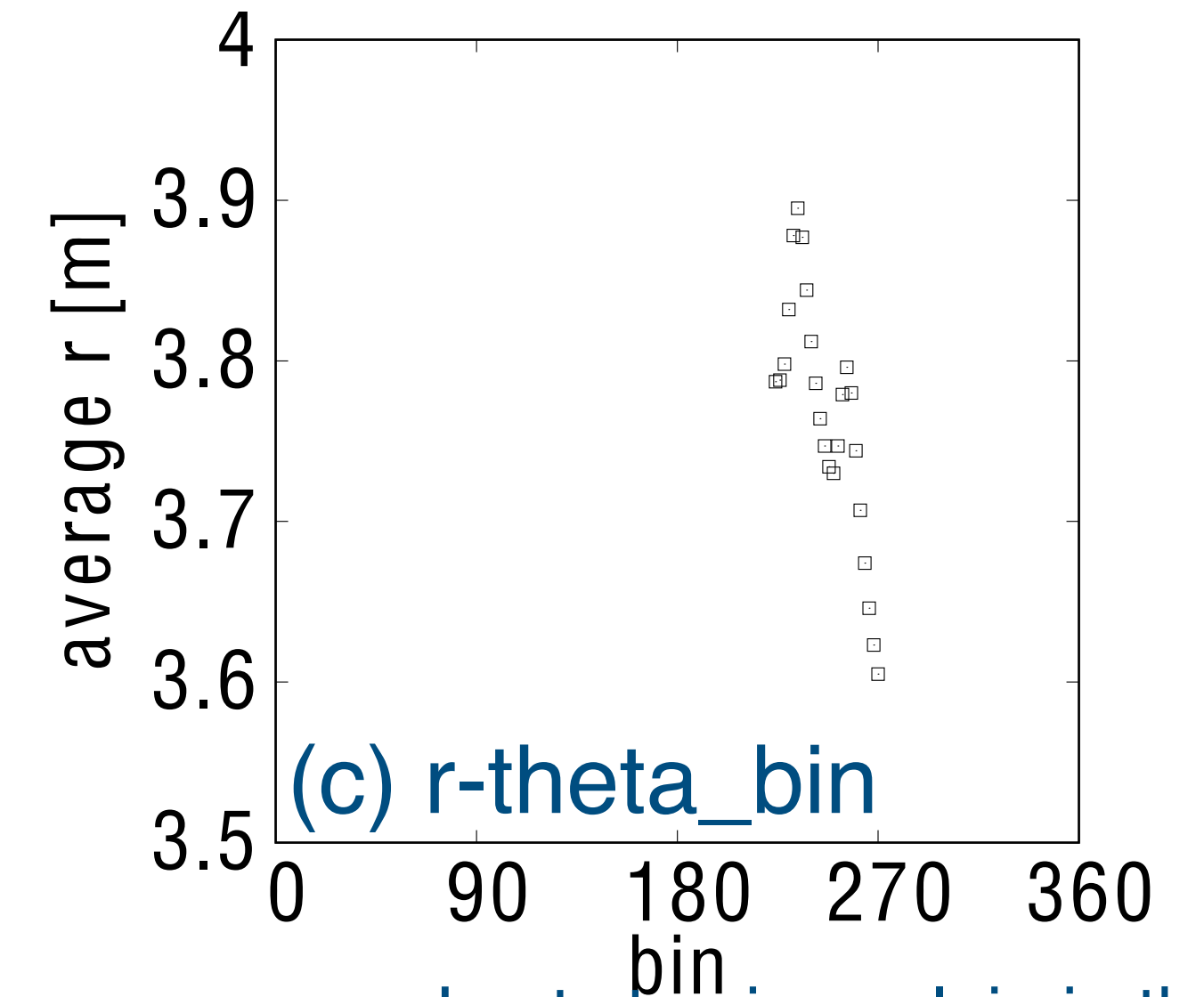
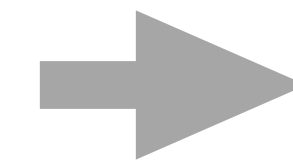


- Convert coordinates to r-theta



- Binning in theta direction
- Calculate average r in each bin

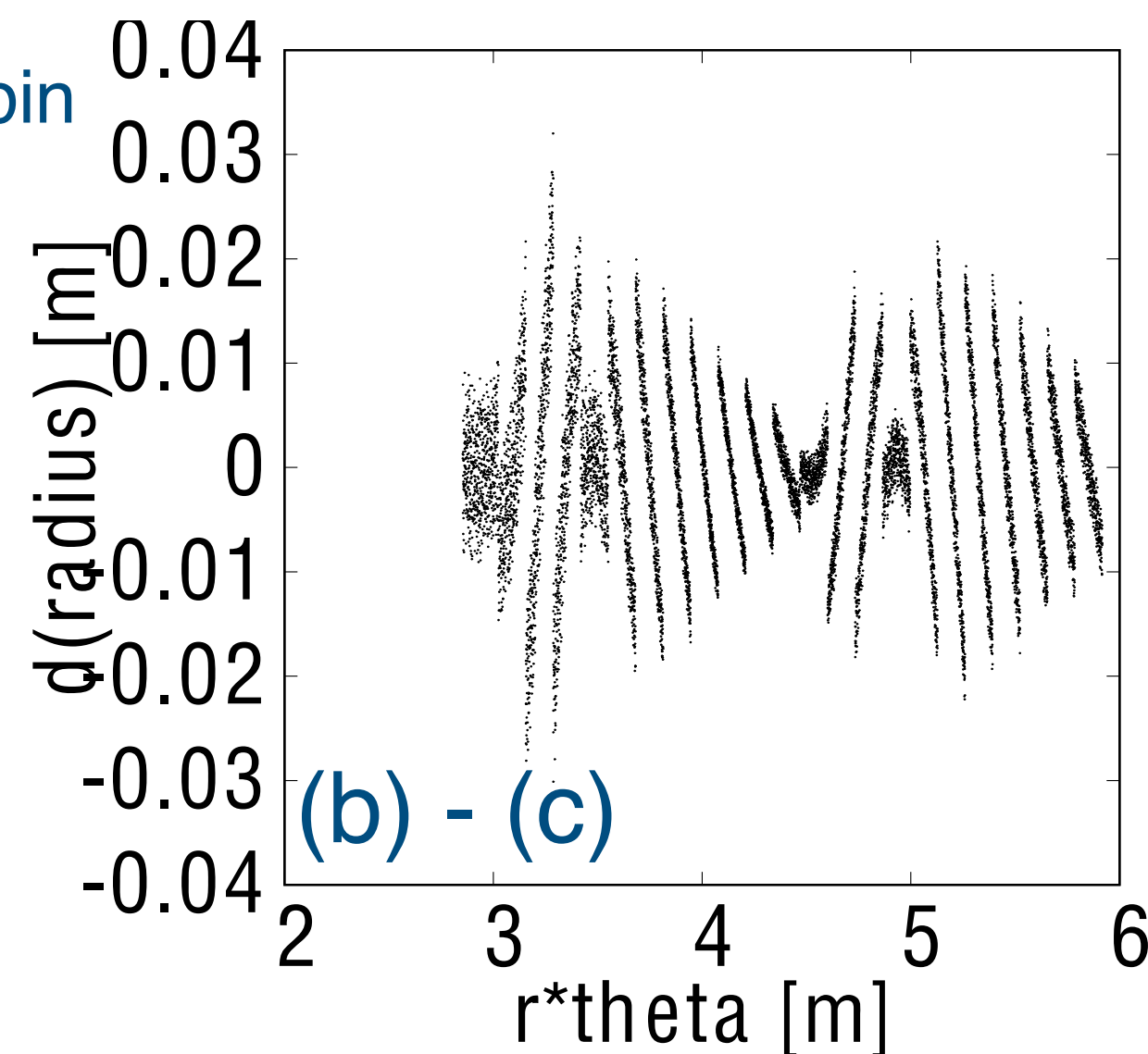
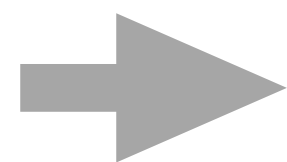
$$\left( \frac{\sum x_i}{n}, \frac{\sum y_i}{n} \right)$$



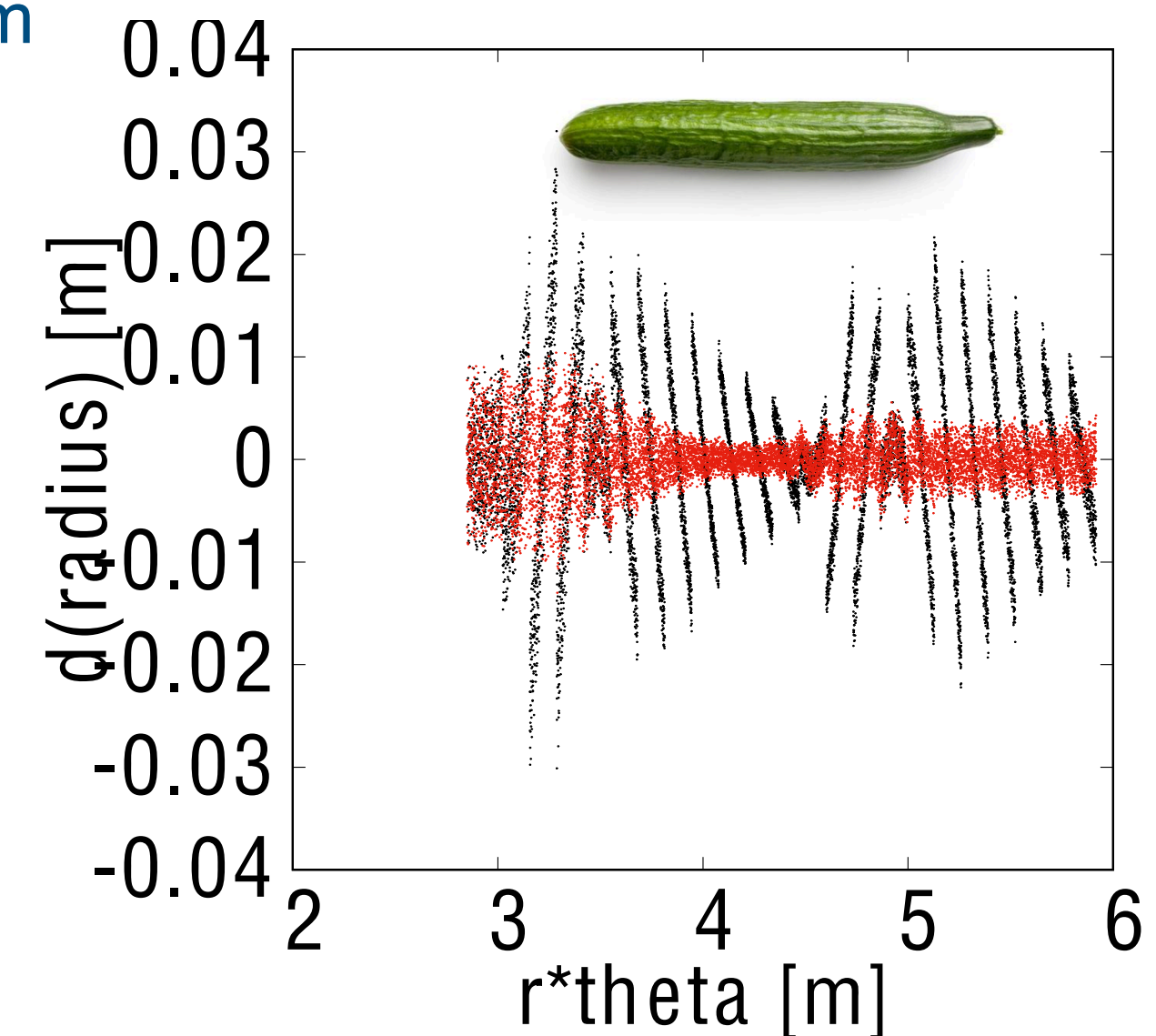
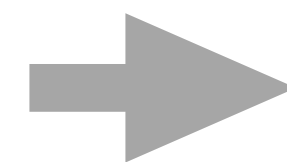
- Last step is explain in the next page.

- Same binning in theta direction
- Calculate tilt angle psi in each bin

$$\tan(\psi) = \frac{\sum p_{y,i}}{\sum p_{x,i}}$$



- Finally a straight beam along beam path.



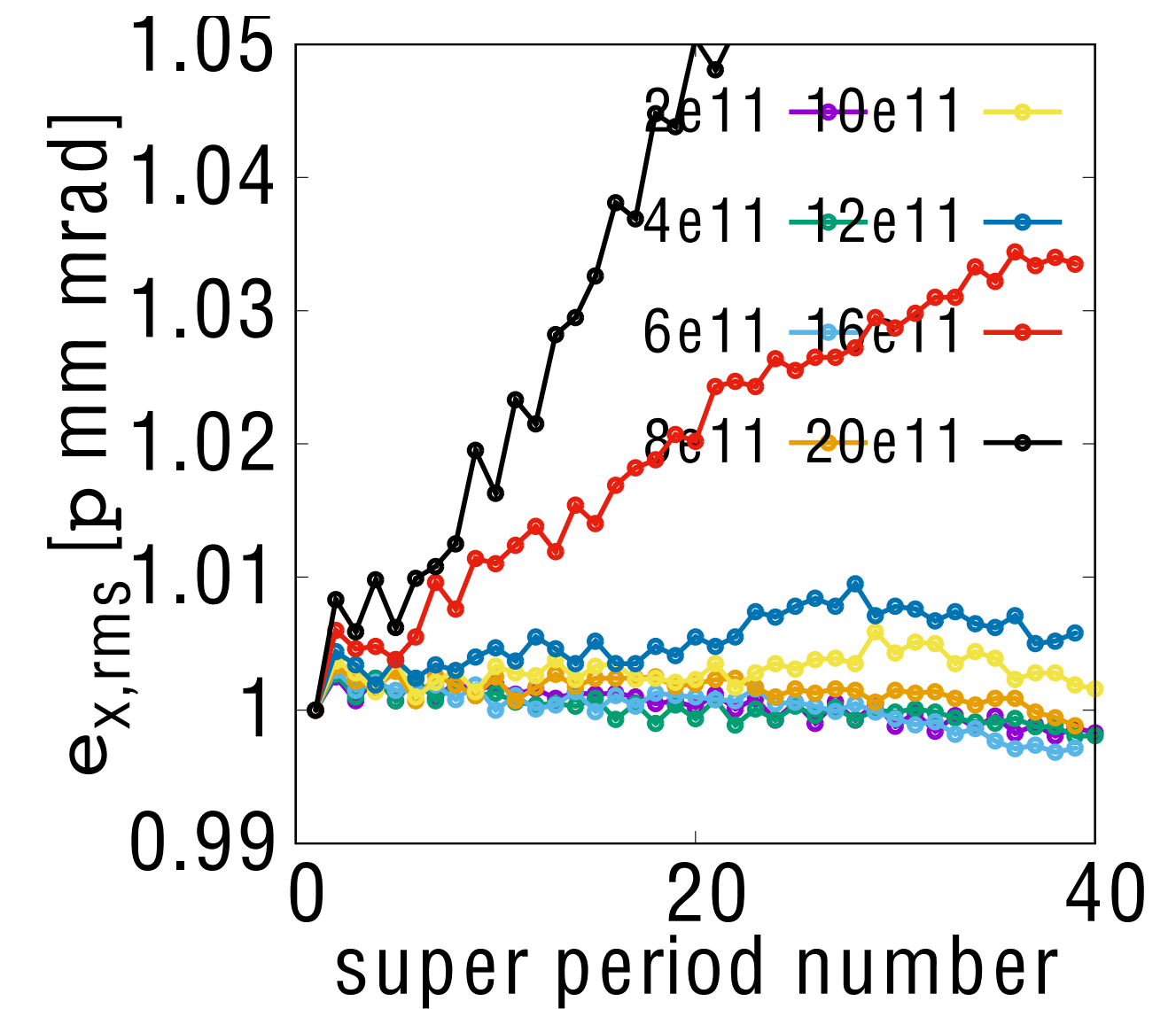
# Simulation result (preliminary)

Lattice	FETS-FFA
Circumference	~ 23 m
Energy	3 MeV
Longitudinal distribution	Coasting
Transverse distribution	KV
Emittance (100%)	10 pi mm mrad, normalised
Injection	Single turn
Operating point	(3.41, 3.39)
Longitudinal bin	360 / ring
# of macro particle	10000

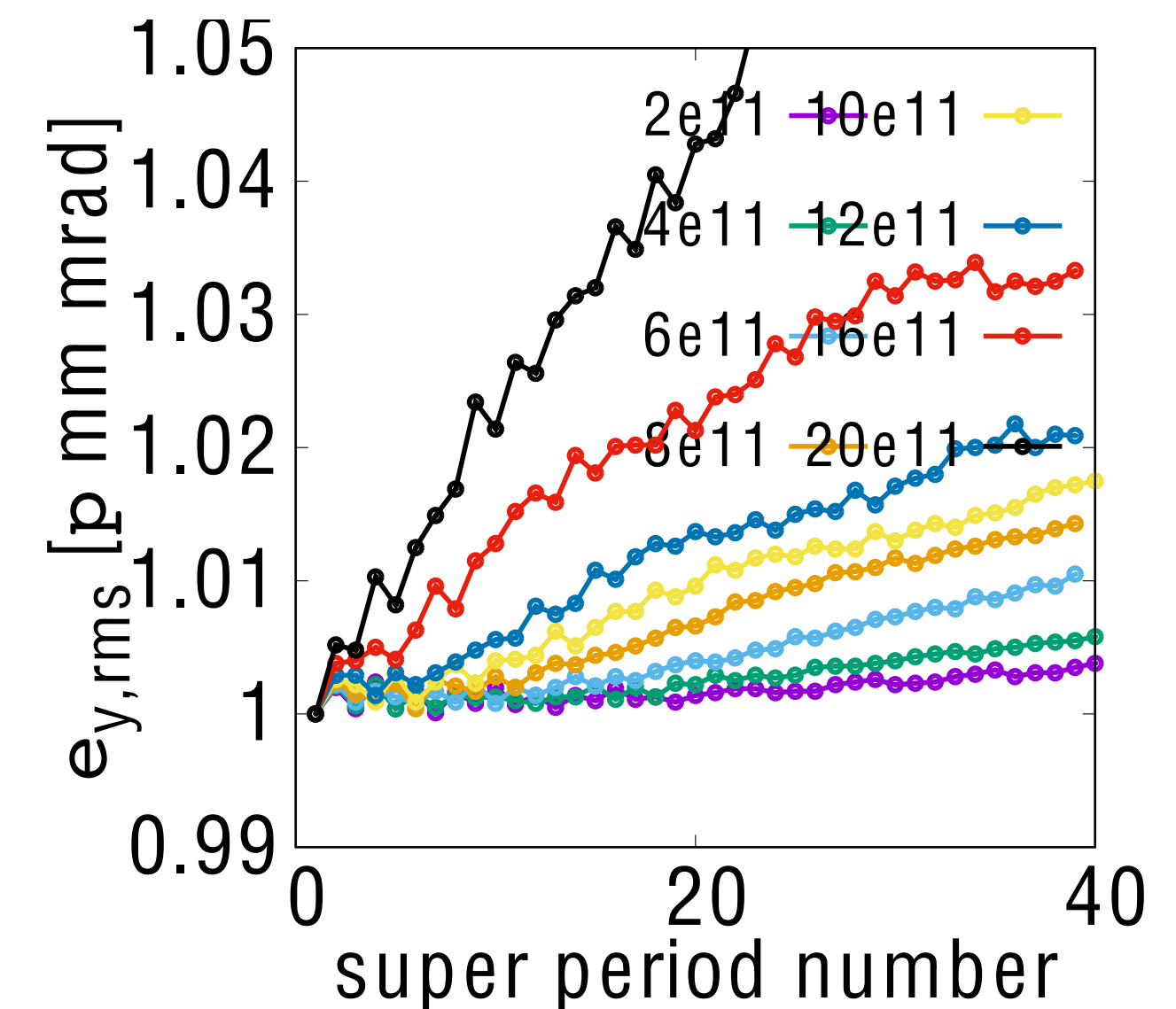
Emittance growth starts happening at  $12 \times 10^{11}$  and significant one above the intensity of  $16 \times 10^{11}$ .

(Partially it is due to mismatch.)

Hor



Ver





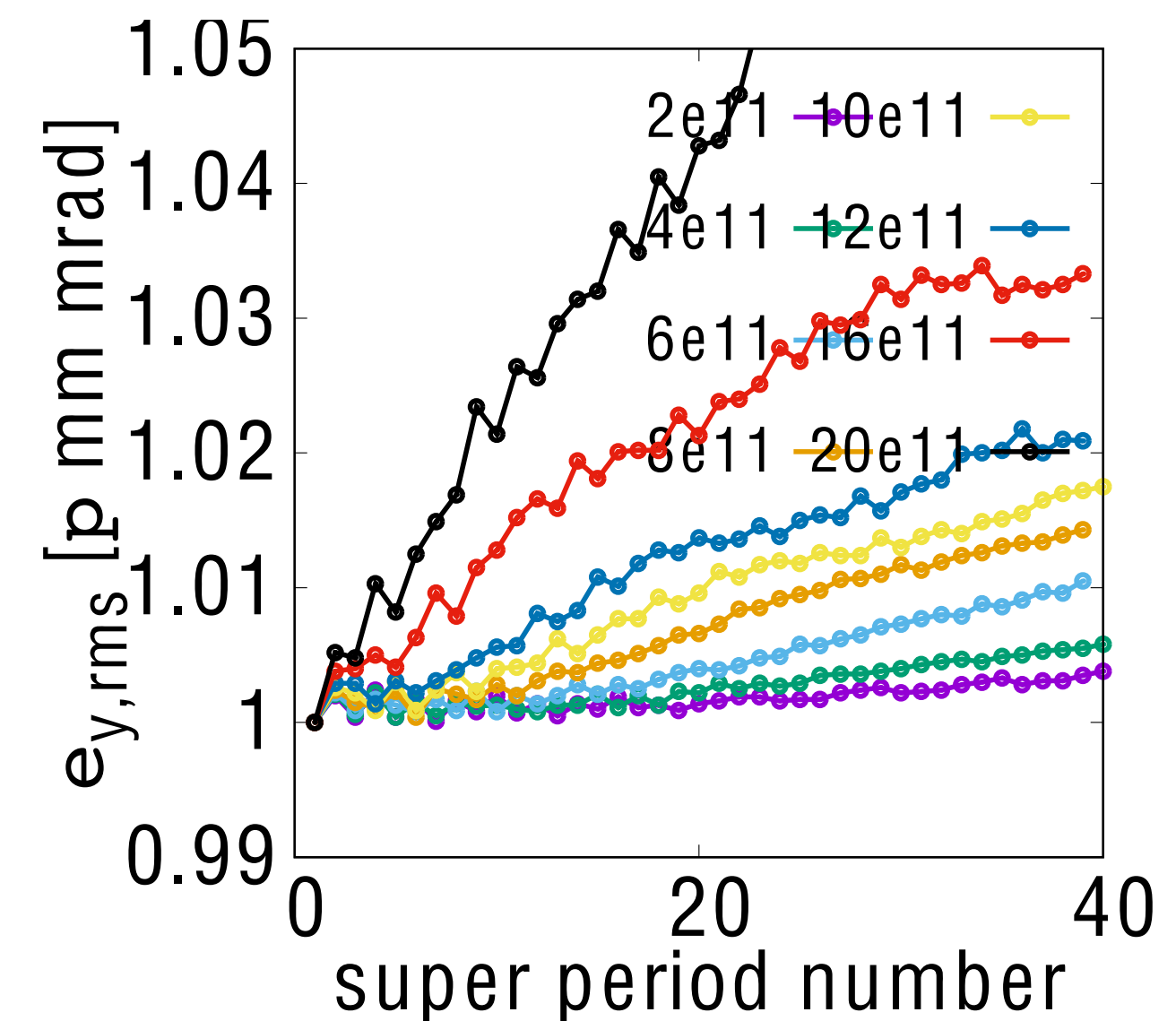
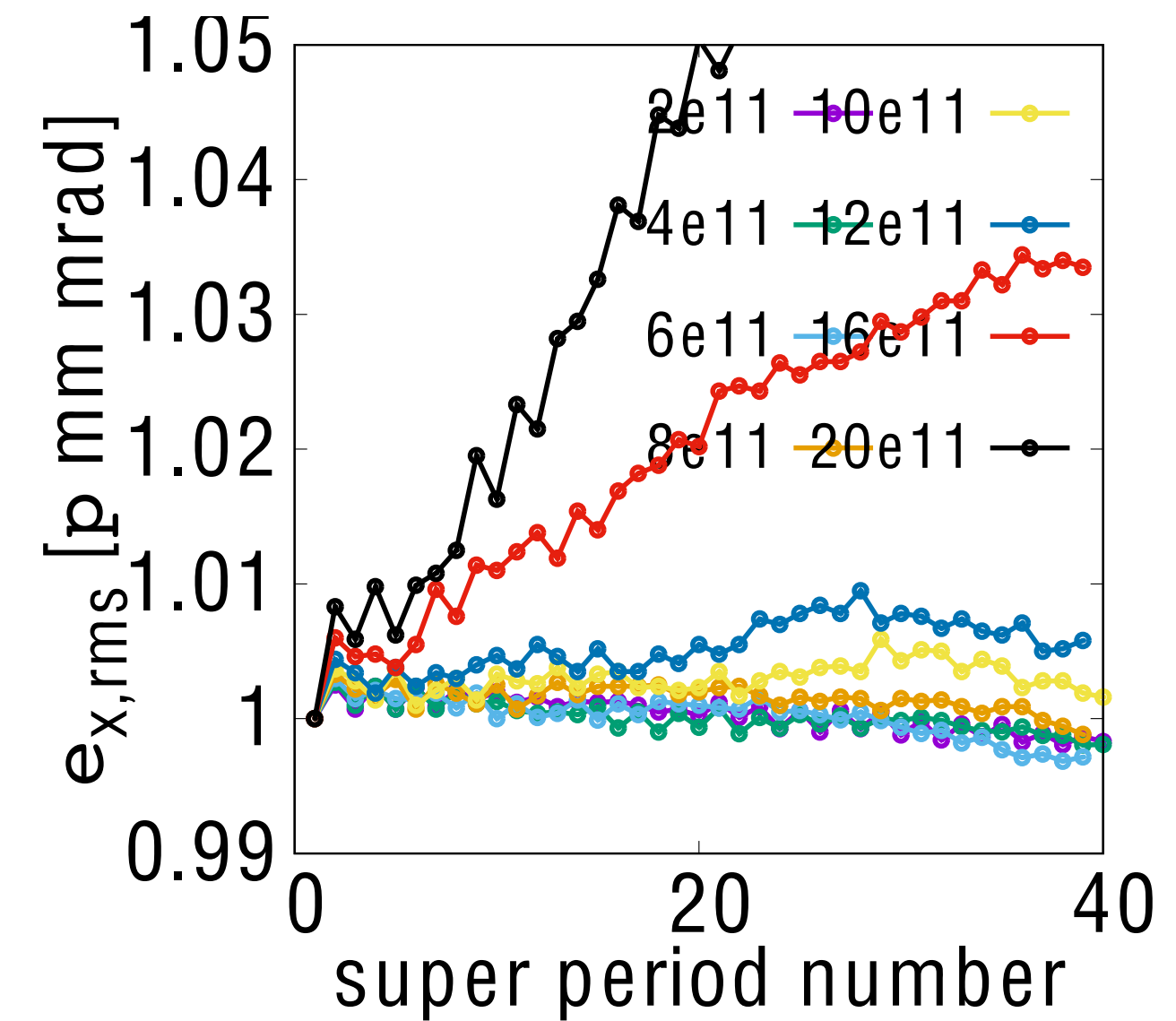
# Simulation result (preliminary)

Space charge incoherent tune shift

$$\Delta Q_v = - \frac{n_t r_p}{\pi \epsilon_v (1 + \sqrt{\epsilon_h / \epsilon_v}) \beta^2 \gamma^3} \frac{1}{B_f}$$

	Maximum inc. tune	RMS inc. tune shift	Coherent tune shift
$10 \times 10^{11}$	-0.304	-0.304	-0.228
$12 \times 10^{11}$	-0.365	-0.365	-0.274
$16 \times 10^{11}$	-0.484	-0.484	-0.365
$20 \times 10^{11}$	-0.608	-0.608	-0.456

- Distance between operating point (3.41, 3.39) and nearby resonance (3.00, 3.00) is (0.41, 0.39).



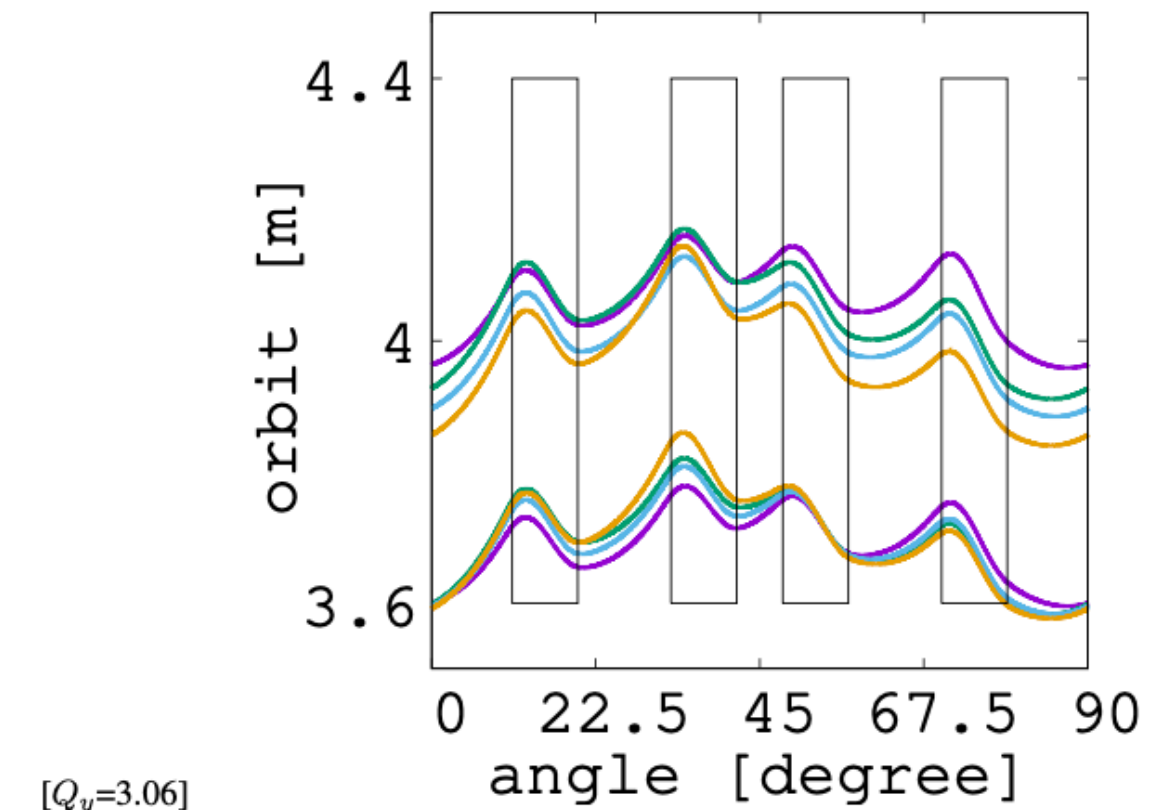
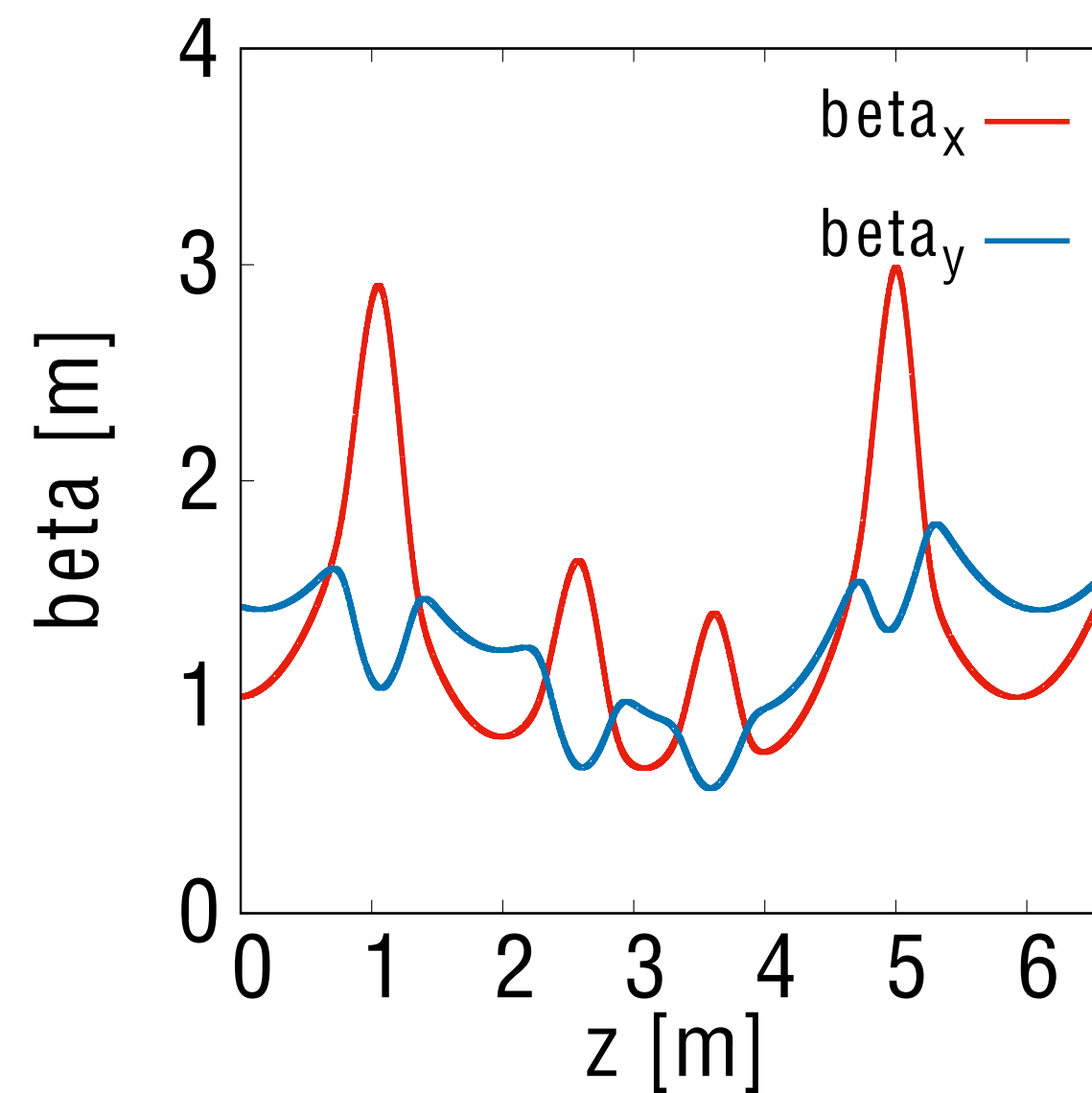
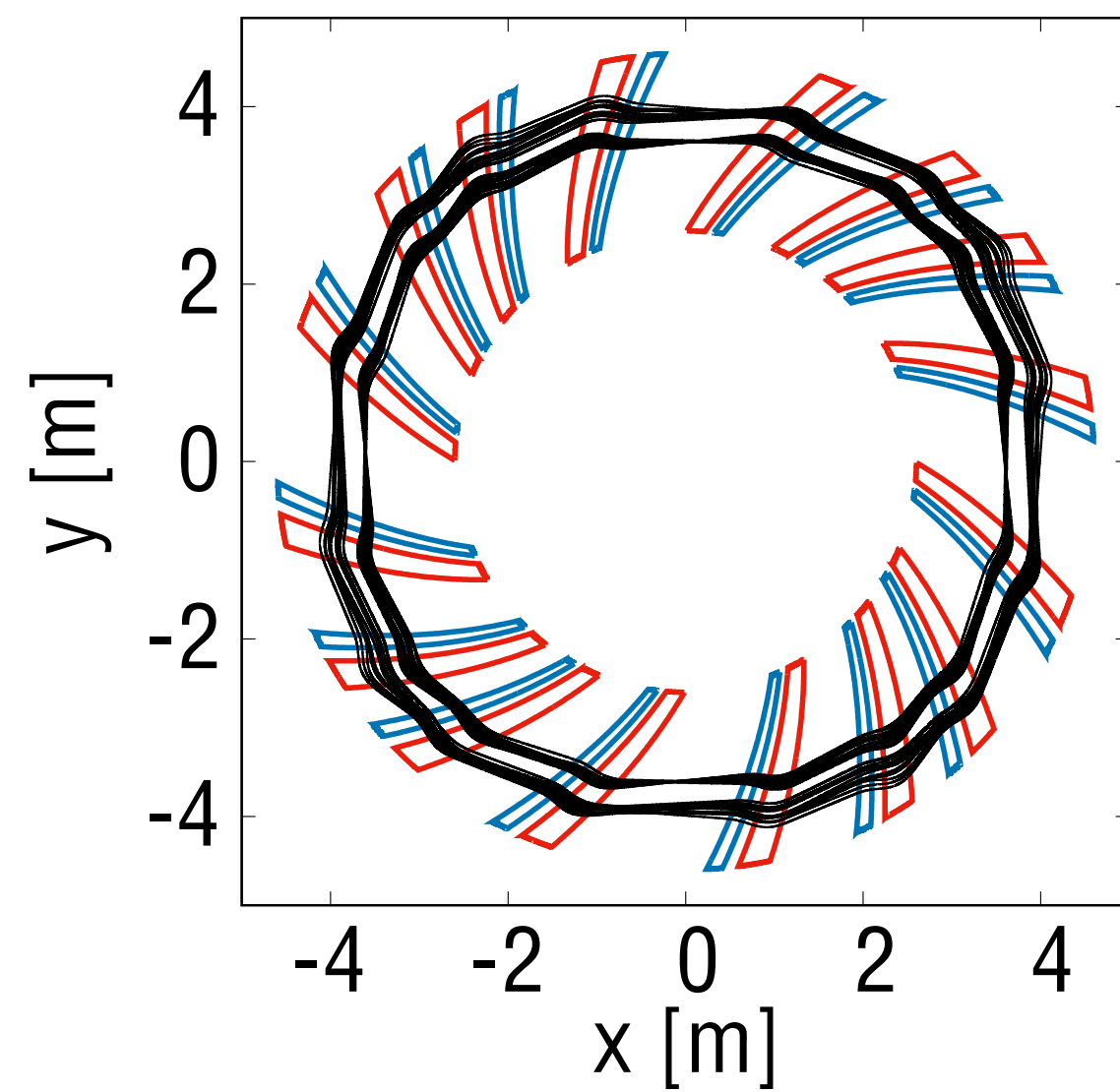
- Emittance growth starting around  $12 \sim 16 \times 10^{11}$  is reasonable (**no surprise!**).

# Superperiodicity

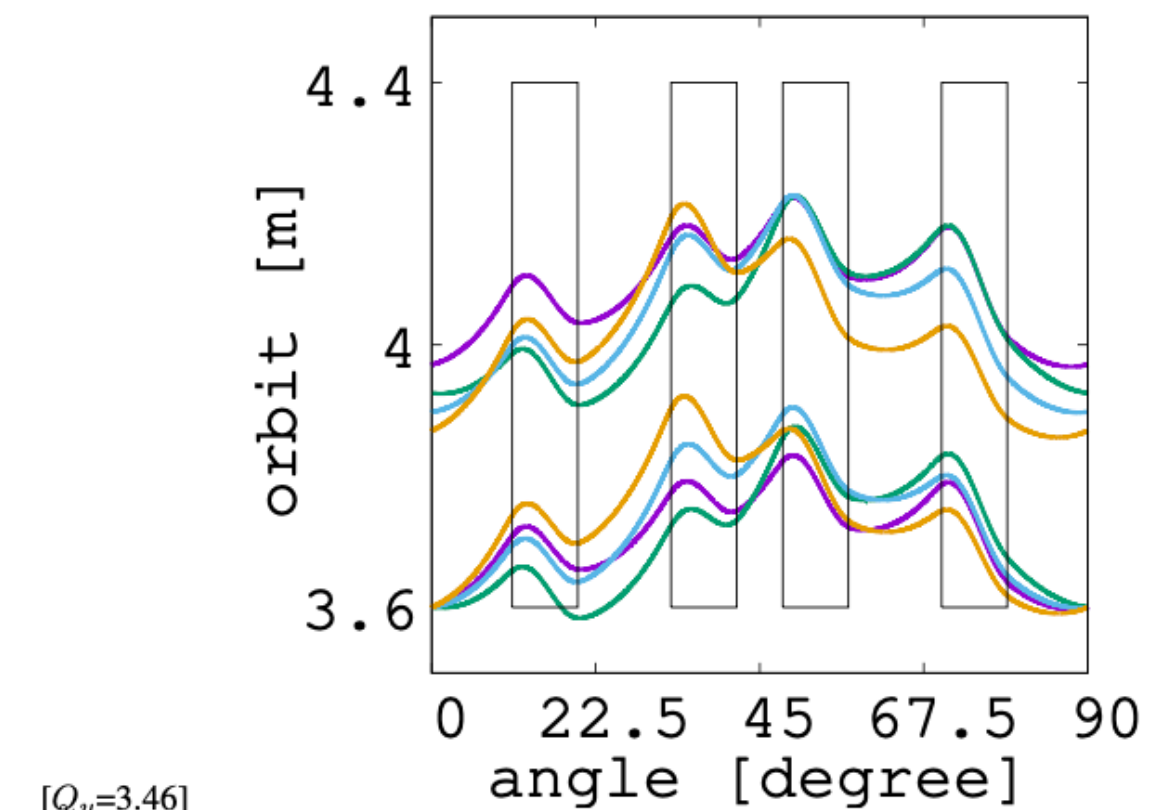
## FETS-FFA proposal at RAL

- 4-fold symmetry lattice with radius of 3.6 m

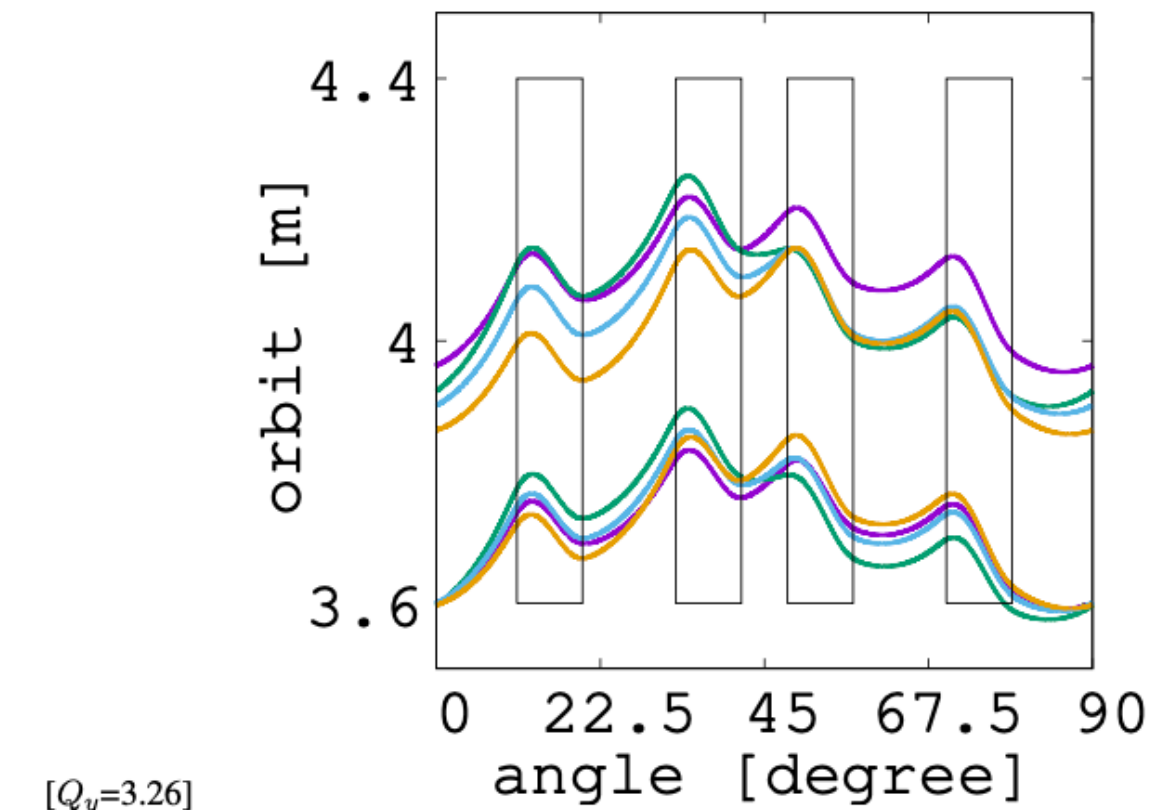
$$B_z(r, \theta) = B_{z0} \left( \frac{r}{r_0} \right)^k F(\theta)$$



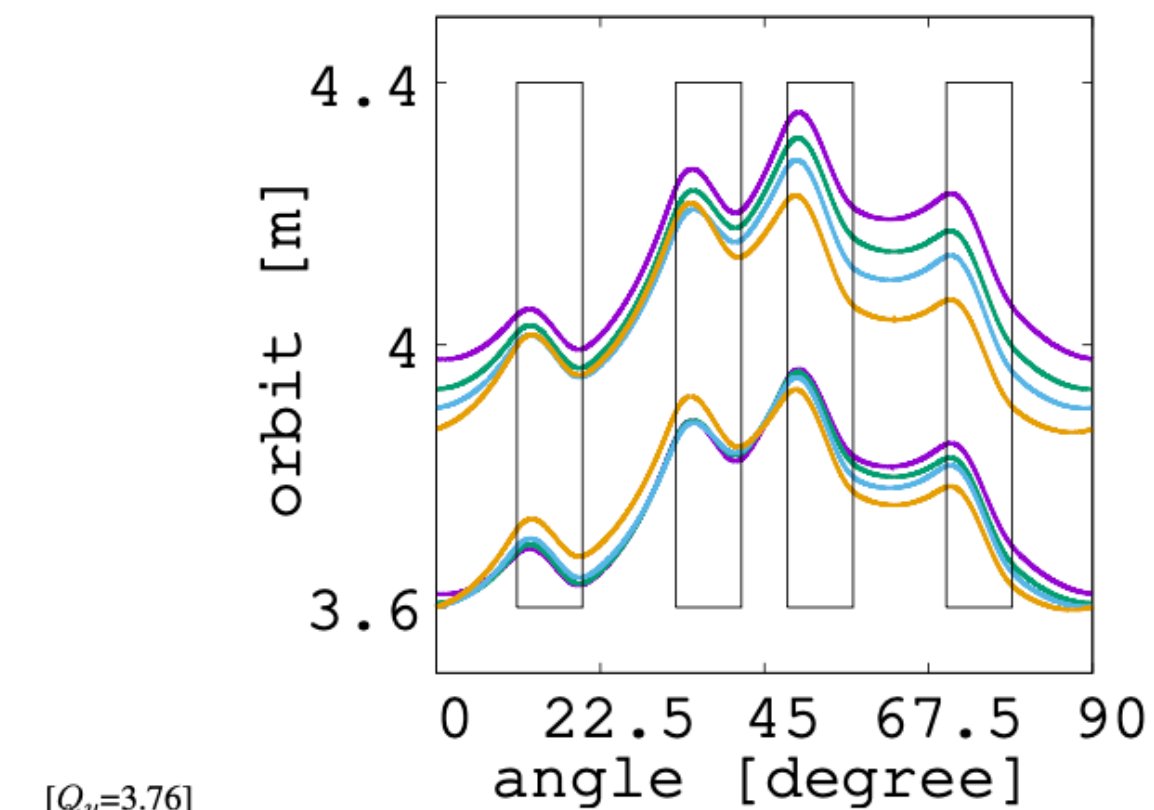
[ $Q_y=3.06$ ]



[ $Q_y=3.46$ ]



[ $Q_y=3.26$ ]



[ $Q_y=3.76$ ]

Figure 2.8: 3 MeV and 12 MeV orbits for 16 operating points.

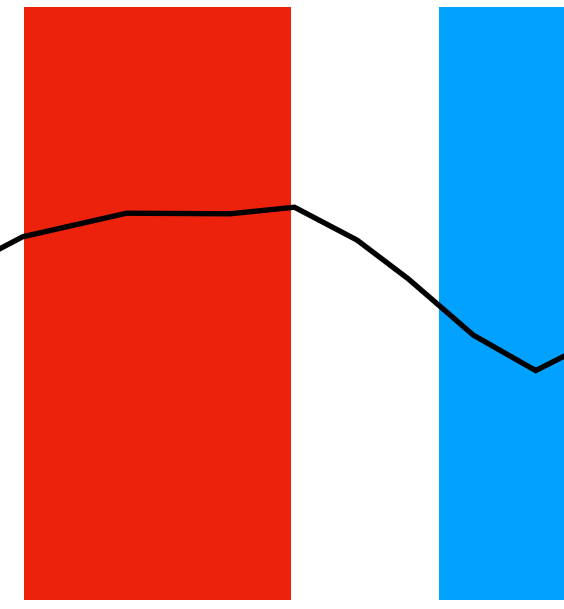
Figures shows injection and extraction orbits which have the momentum ratio of two.

# DF (FD) spiral sector

cell tune is fixed at (0.213125, 0.213125)

## Optics

radial sector



Bf

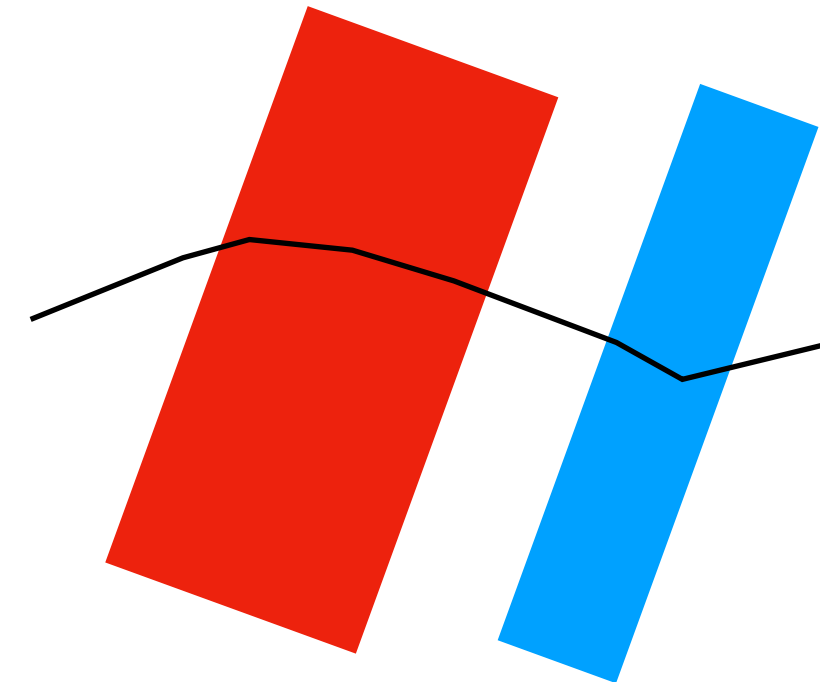
Bd

(DFD)

(DDD)

(entrance, body, exit)

FD spiral

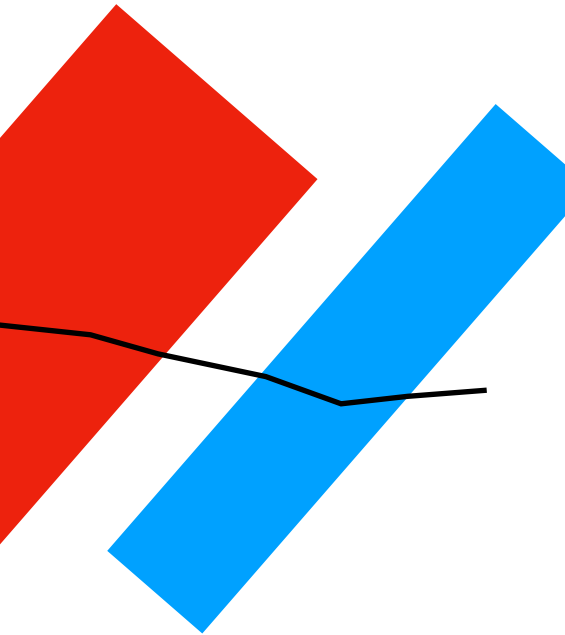


Bf

Bd

(DFO)

(ODD)



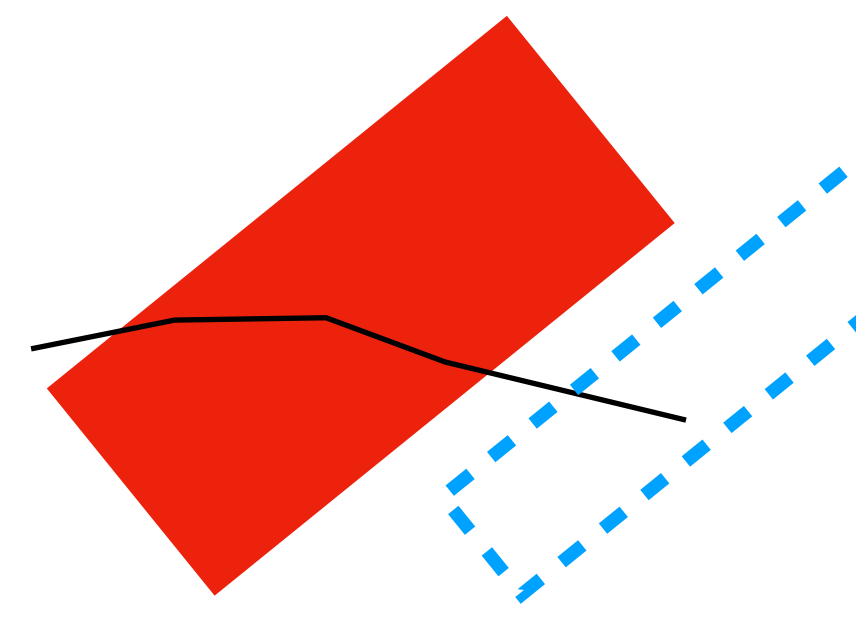
Bf

Bd

(DFF)

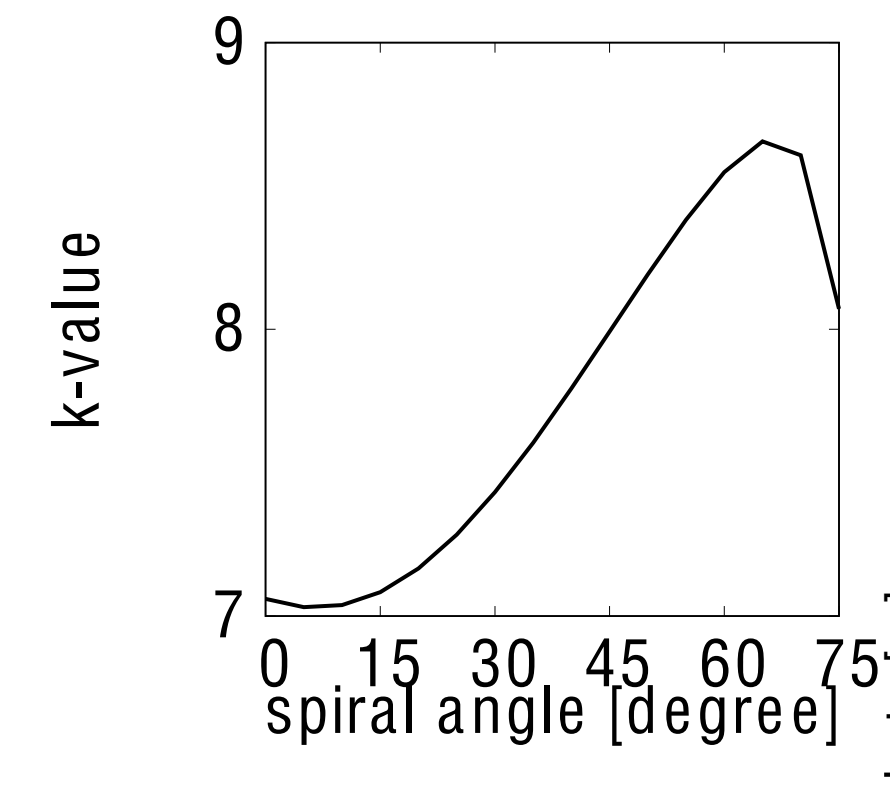
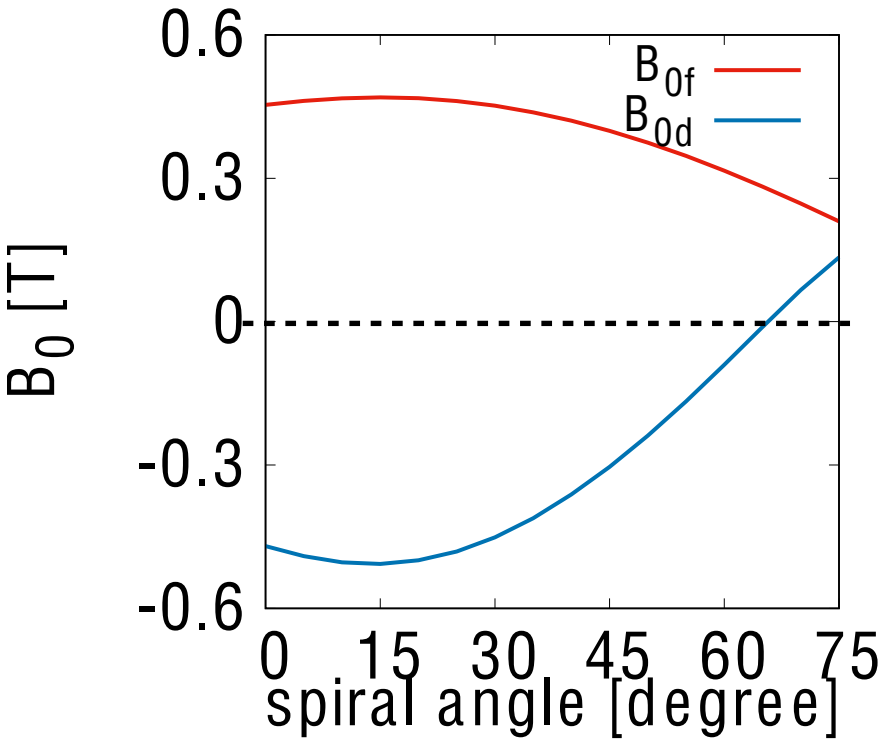
(FDD)

spiral sector



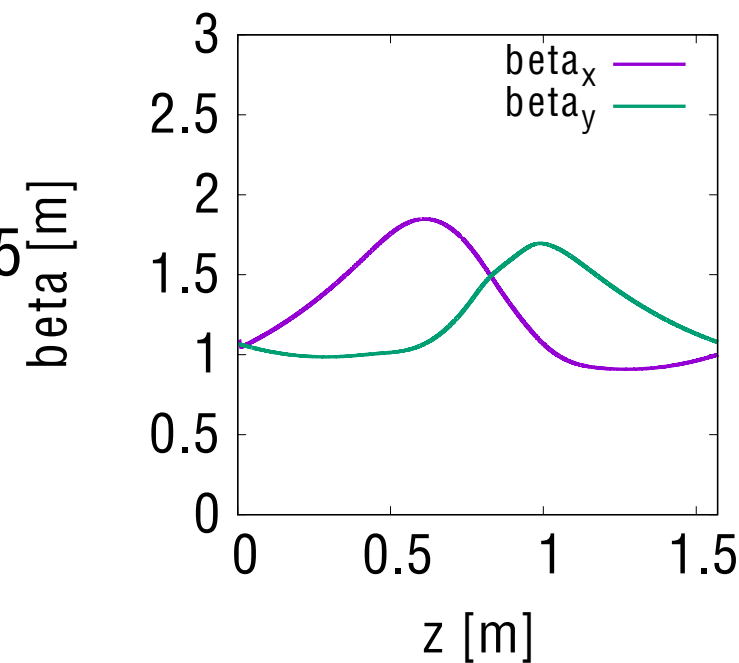
Bf

(DFF)



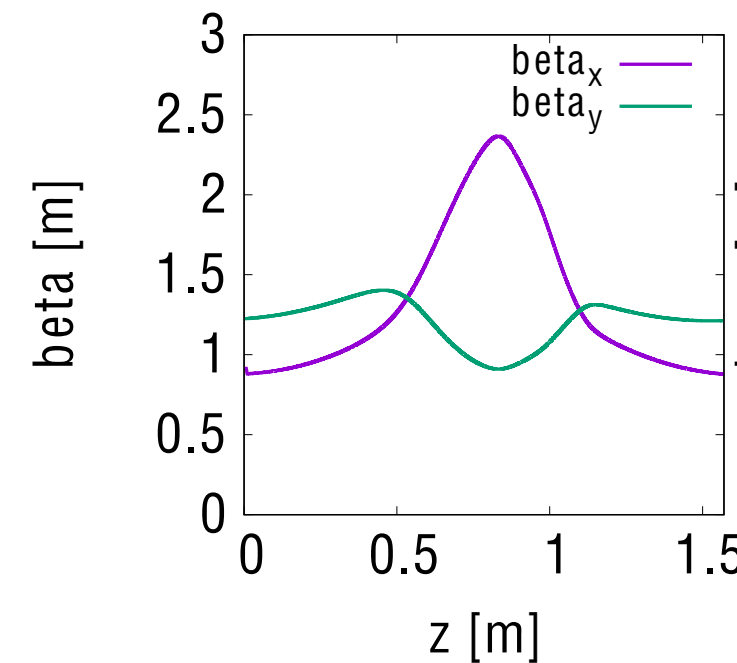
### FD

Spiral angle=0 deg



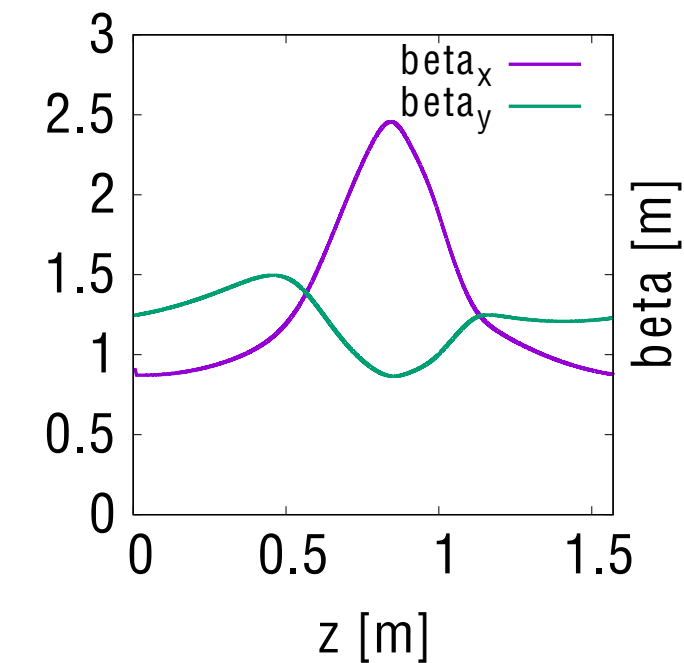
### DFD

35 deg

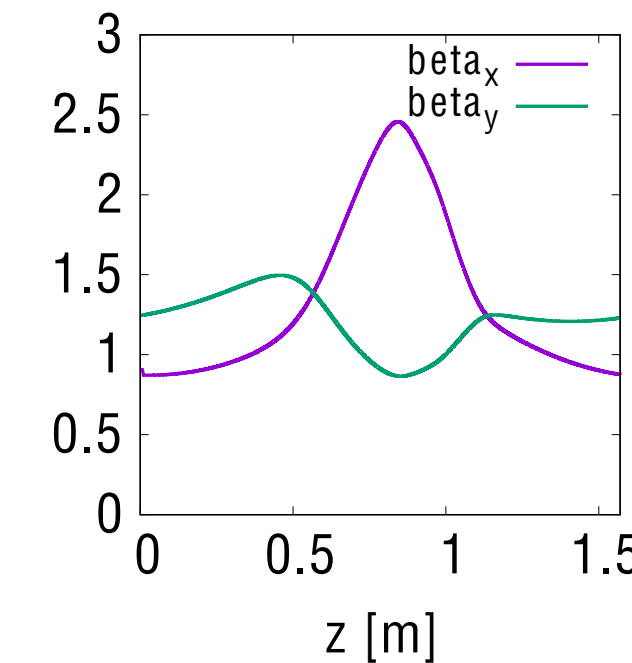


### DFD

40 deg

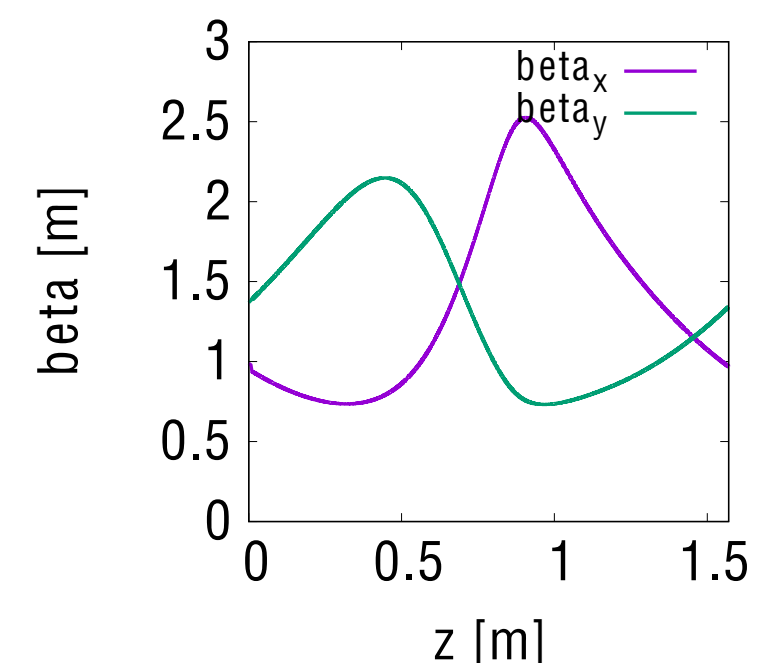


45 deg



### DF

65 deg



Optics looks like triplet, not doublet

# Why it did not go further?

*What we are doing now*

## Magnet R&D

### Demonstrator



FFAG'14 International Workshop on FFAG Accelerators  
September 22 – 27, 2014  
Brookhaven National Laboratory, New York  
(September 23, Tuesday)

### High Temperature Superconductor Magnet for FFAG Accelerators

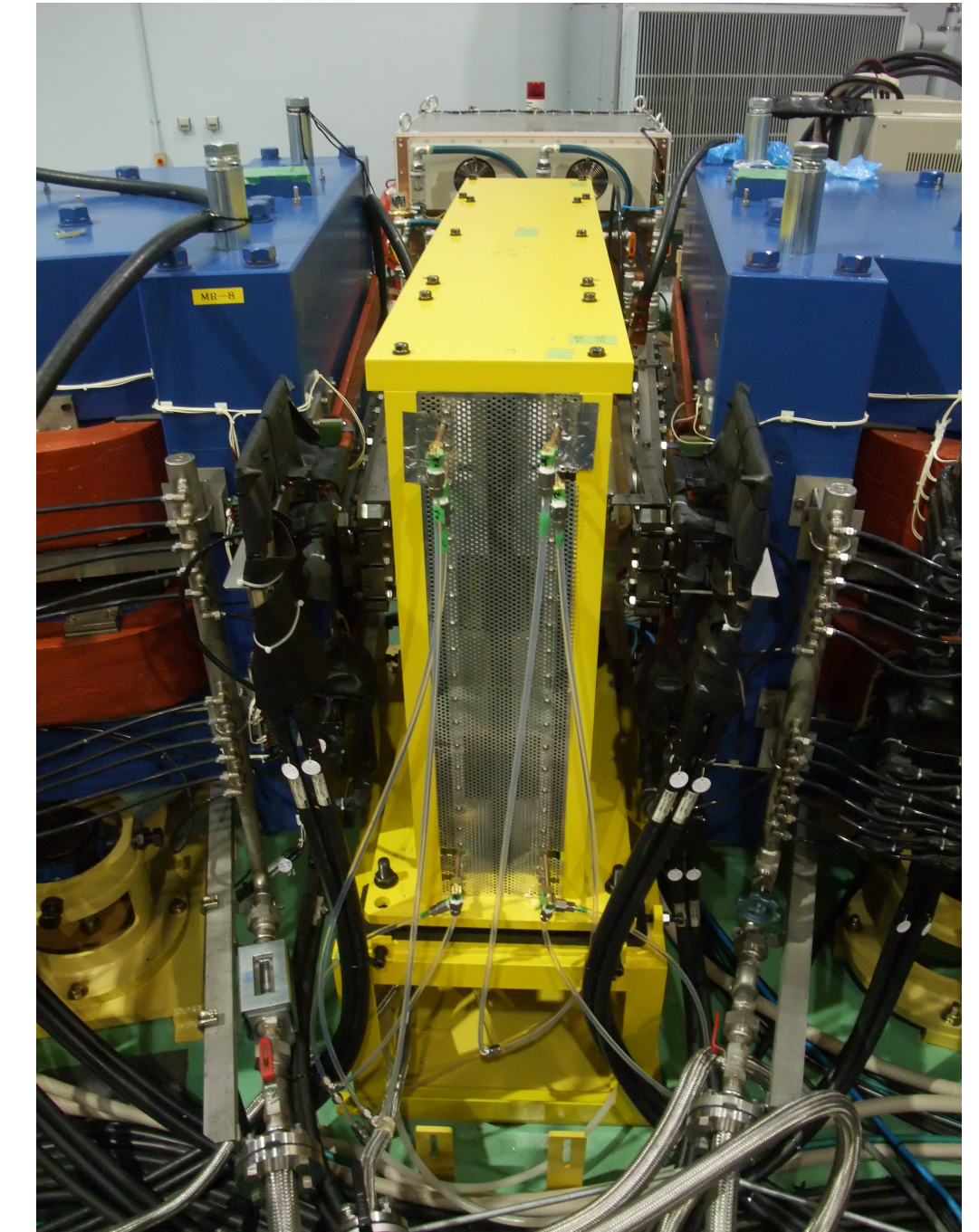
N. Amemiya (Kyoto University), T. Ogitsu (KEK)  
K. Koyanagi, T. Kurusu (Toshiba)  
Y. Mori (Kyoto University)  
Y. Iwata, K. Noda (NIRS), M. Yoshimoto (JAEA)

amemiya.naoyuki.6a@Kyoto-u.ac.jp

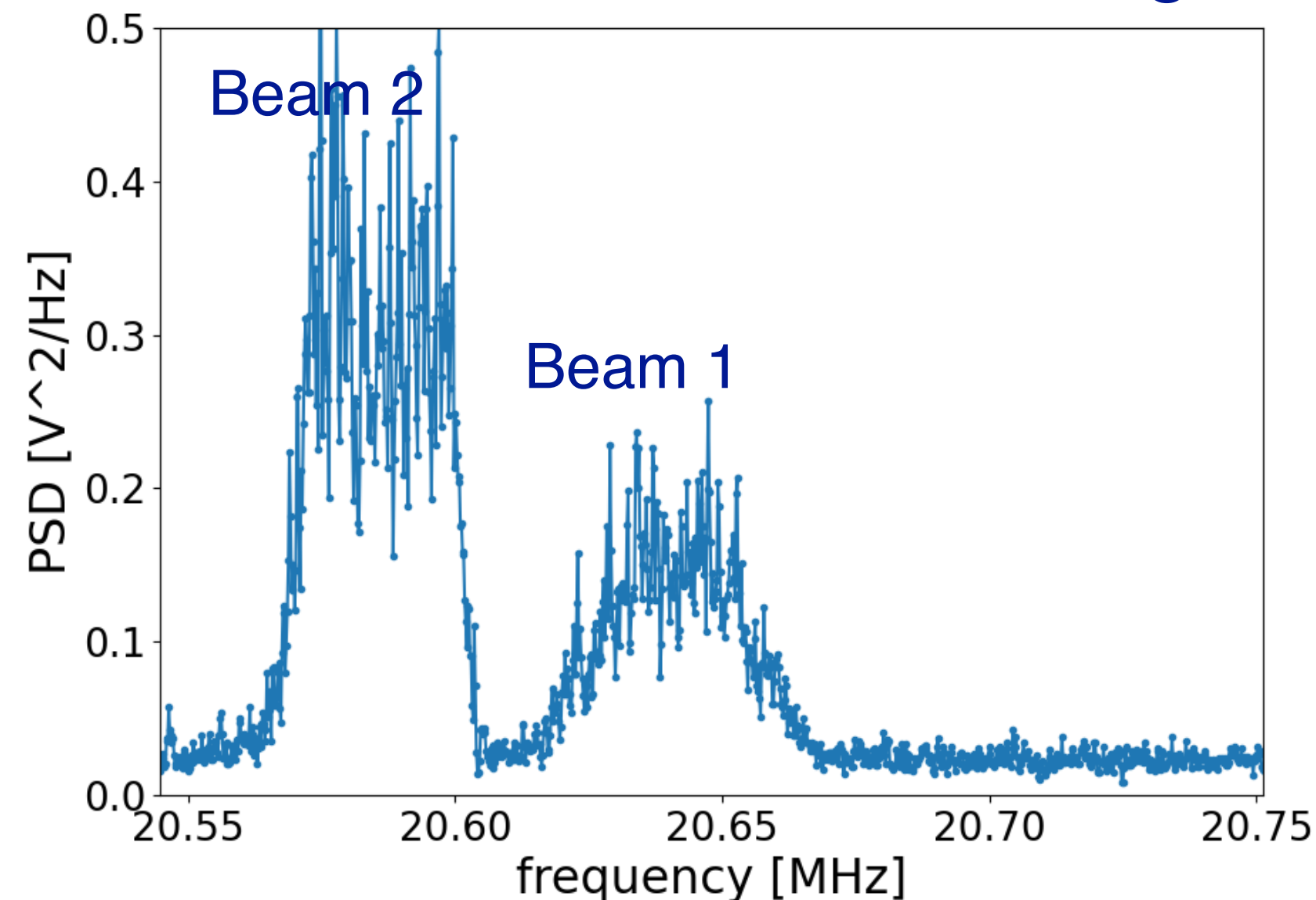
This work was supported by Japan Science and Technology Agency under Strategic Promotion of Innovative Research and Development Program (S-Innovation Program).



### MA RF cavity

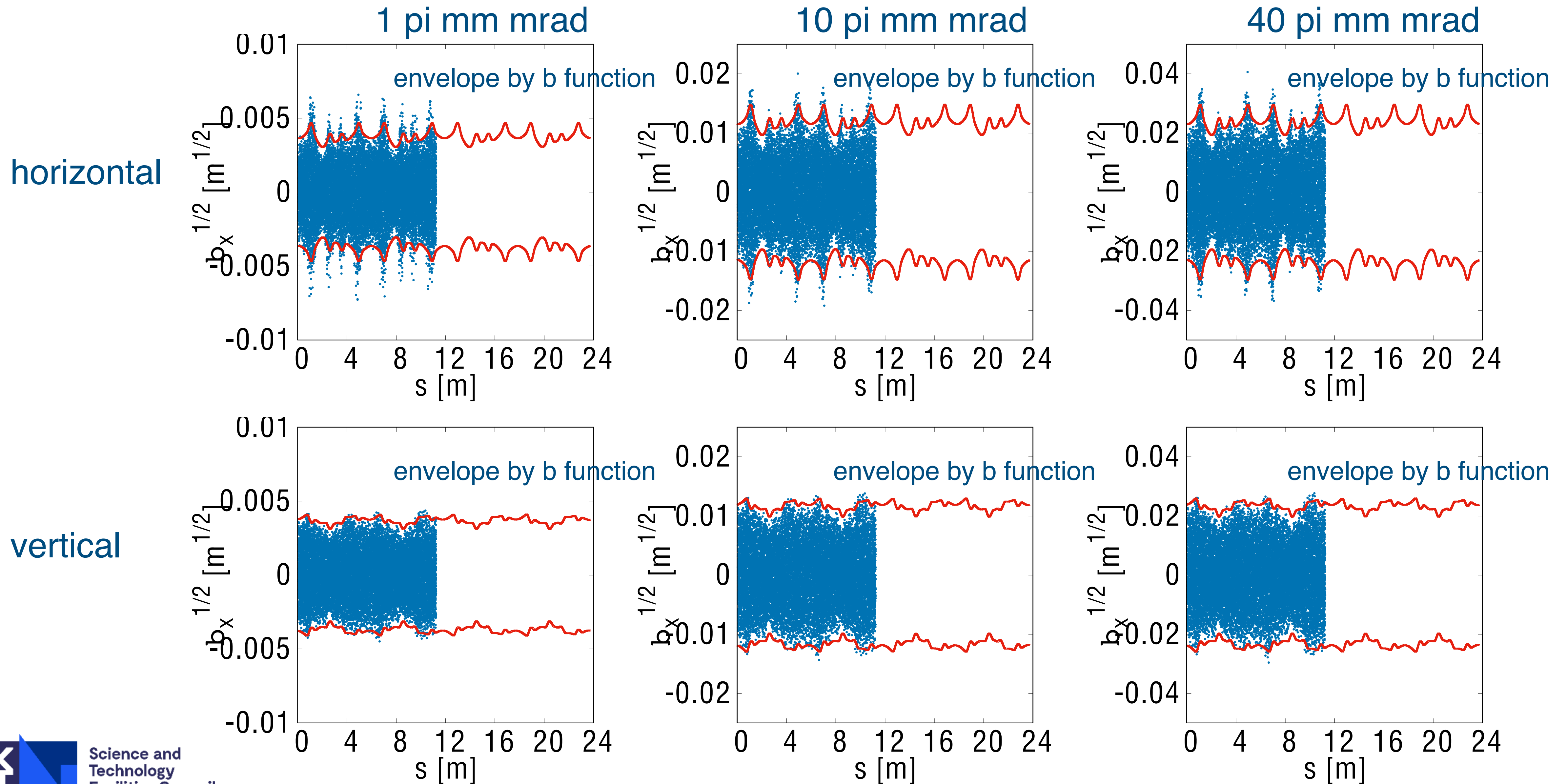


### Beam stacking



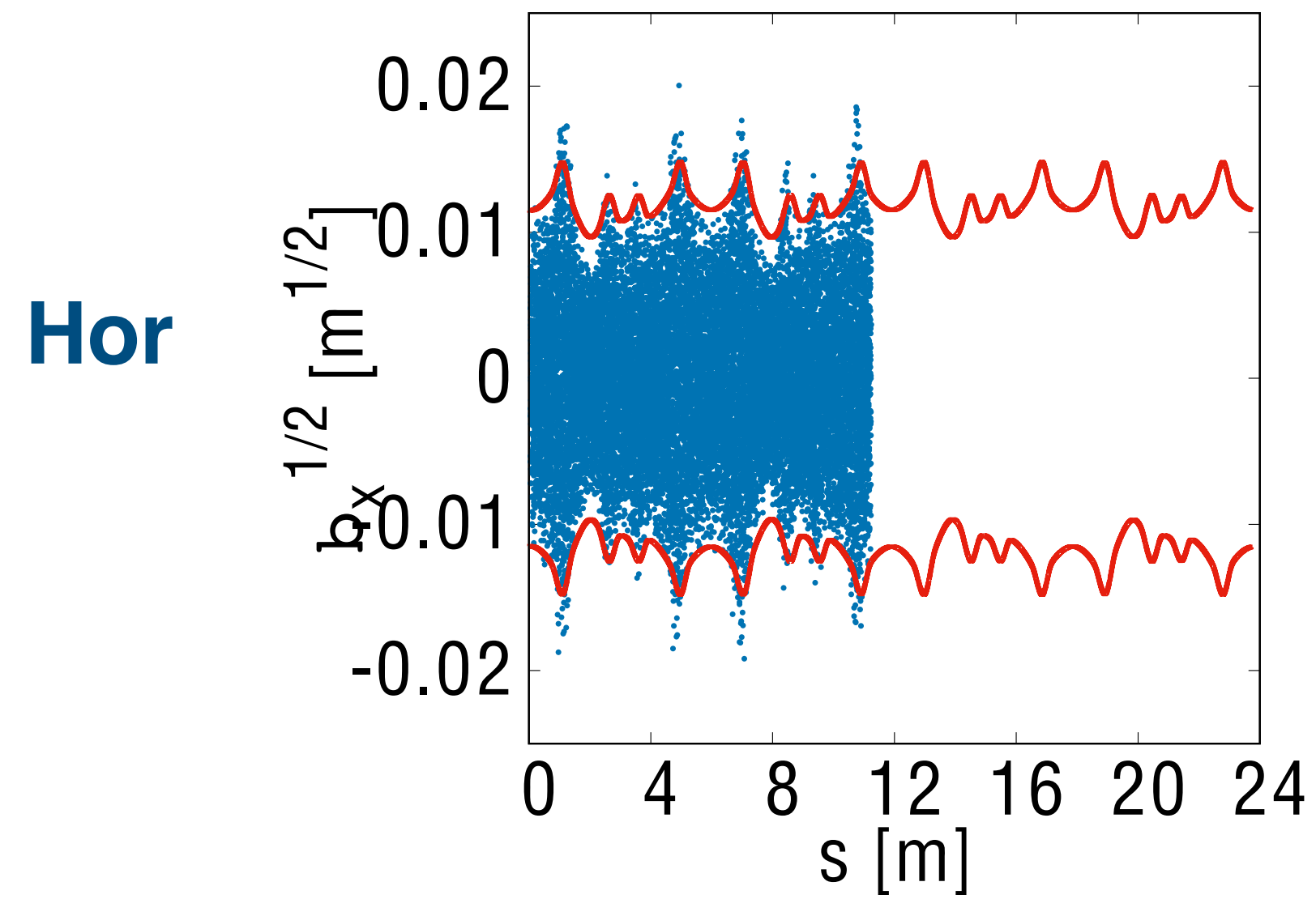
Simply, it was not ready yet.  
(Nothing here existed before.)

# Finally, a curved beam in arbitrary position becomes straight

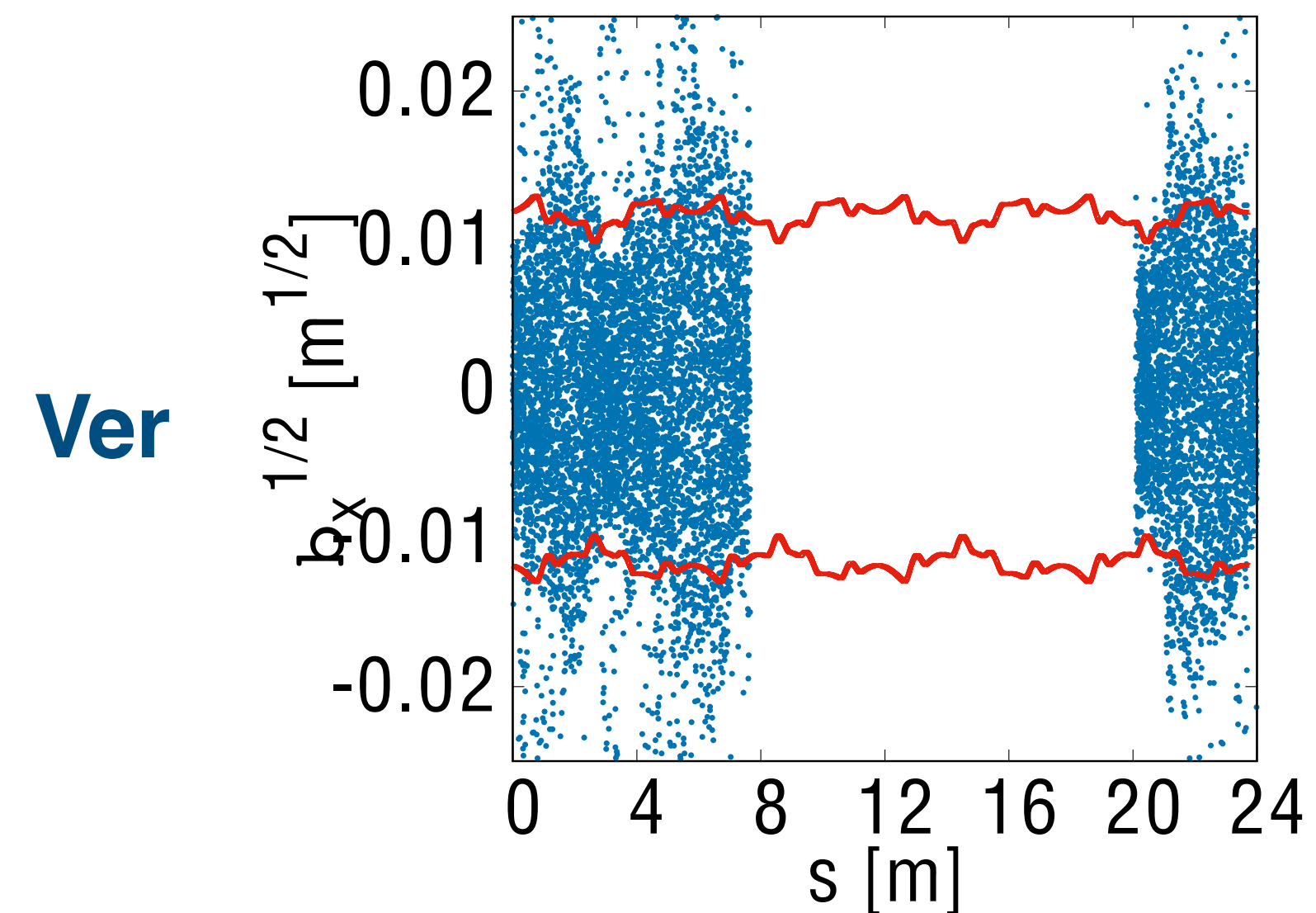
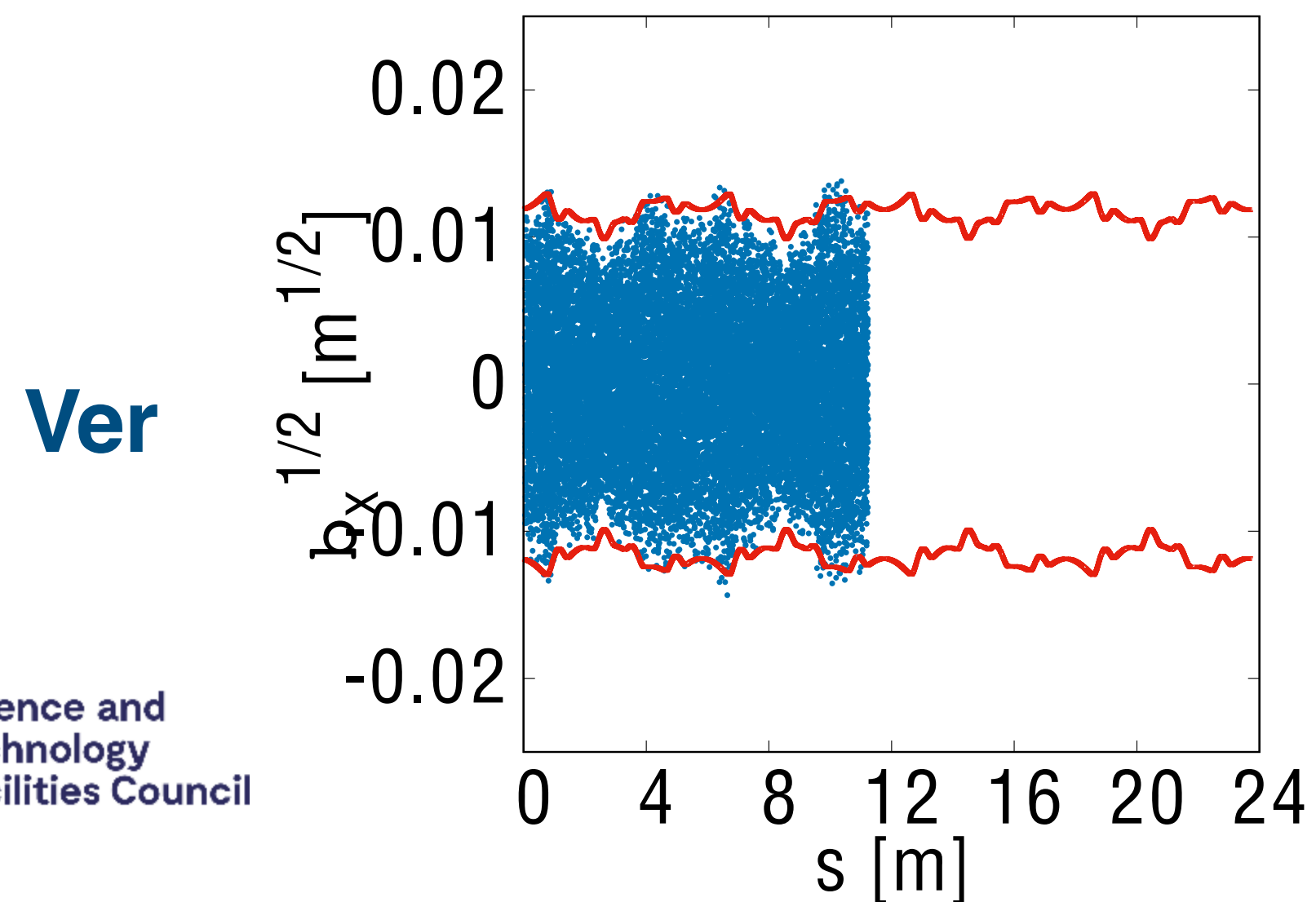
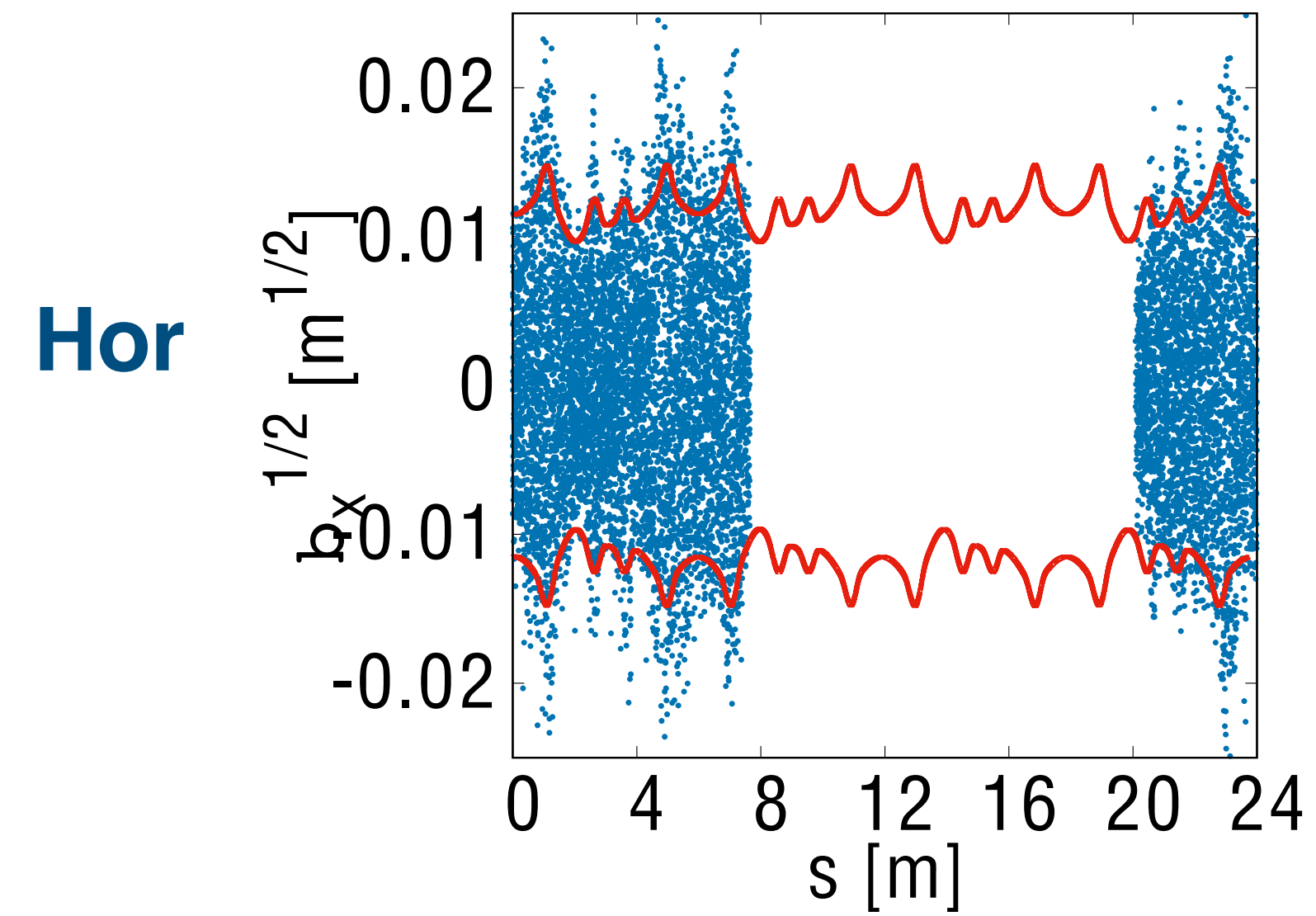


# Simulation result (preliminary)

Initial beam envelope.



Beam envelope at 19th turn with  $20 \times 10^{11}$ .



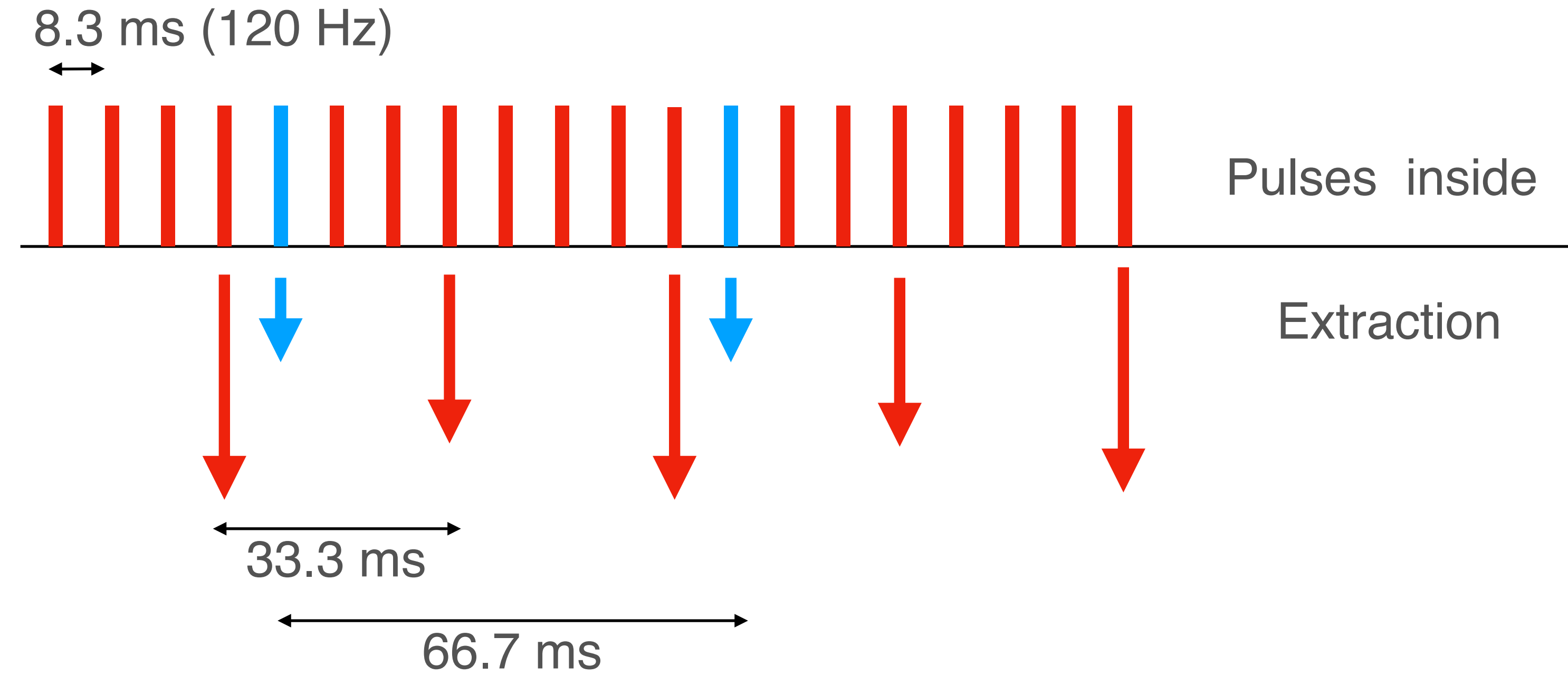
# Space charge tune shift

- Tune shift is inversely proportional to  $\beta^2 \gamma^3$ .
- Space charge effects are strong at injection, but decrease quickly with acceleration.

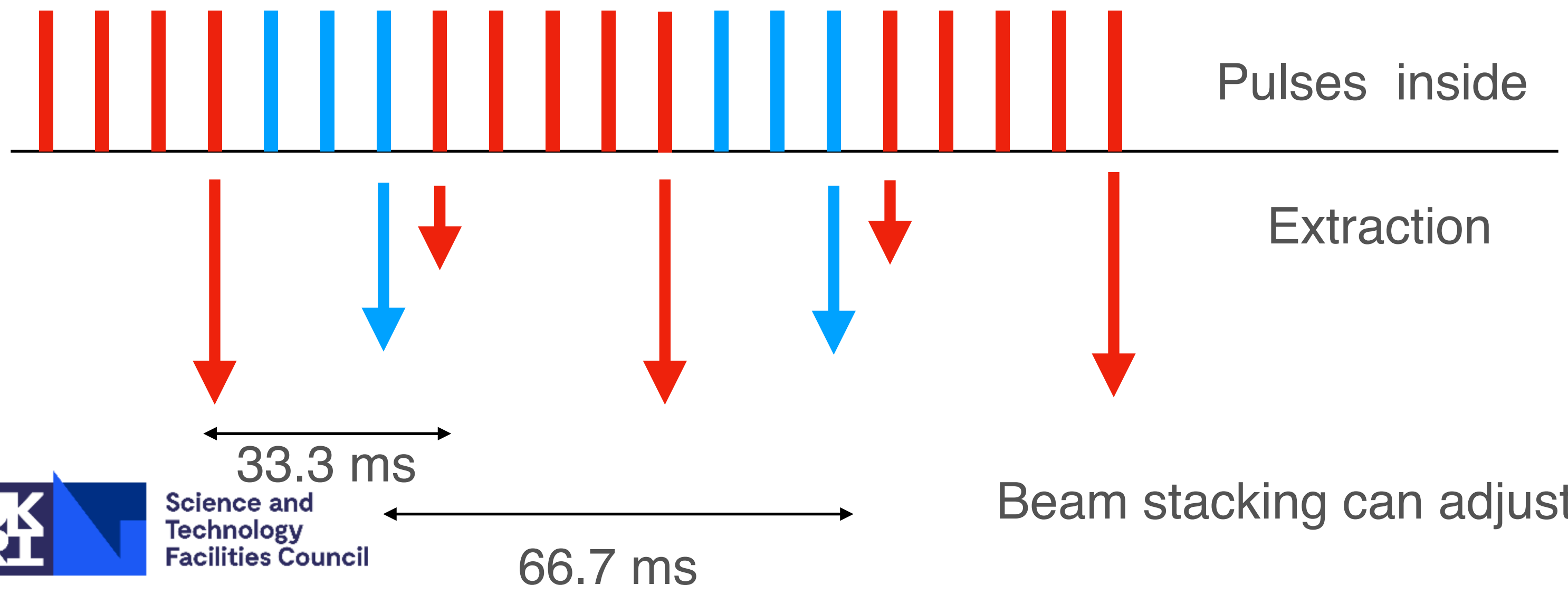
$$\Delta Q_v = - \frac{n_t r_p}{\pi \epsilon_v (1 + \sqrt{\epsilon_h / \epsilon_v}) \beta^2 \gamma^3} \frac{1}{B_f}$$

- If all the particles are injected at the same time, the peak beam power is limited at injection.
- It is possible to combine more number of particles at extraction to increase the peak power.

# 30/15 Hz user cycle by beam stacking from 120 Hz accelerator



	TS1 (red)	TS2 (blue)
Rep. rate	30 Hz	15 Hz
Power	2.1 MW	0.3 MW



	TS1 (red)	TS2 (blue)
Rep. rate	30 Hz	15 Hz
Power	1.5 MW	0.9 MW

Beam stacking can adjust beam power for TS-1 and TS-2.



# Frequency component during beam stacking

The beam sees RF voltage at a cavity location

$$V_{gap} = V_0 \cos \omega_{rf} t \sum_{n=0}^{\infty} \delta(t - nT_{rev})$$

$$= V_0 \sum_{n=0}^{\infty} \cos \omega_{rf} nT_{rev} = V_0 \sum_{n=0}^{\infty} \cos 2\pi n \frac{\omega_{rf}}{\omega_{rev}}$$

$V_{gap}$  (envelope) means the lowest frequency component of RF voltage seen by the beam.

when  $\omega_{rf} \ll \omega_{rev}$

$$V_{gap} \text{ (envelope)} = V_0 \cos \omega_{rf} t$$

Requirement in the longitudinal direction imposes

when  $\omega_{rf} \sim \omega_{rev}$

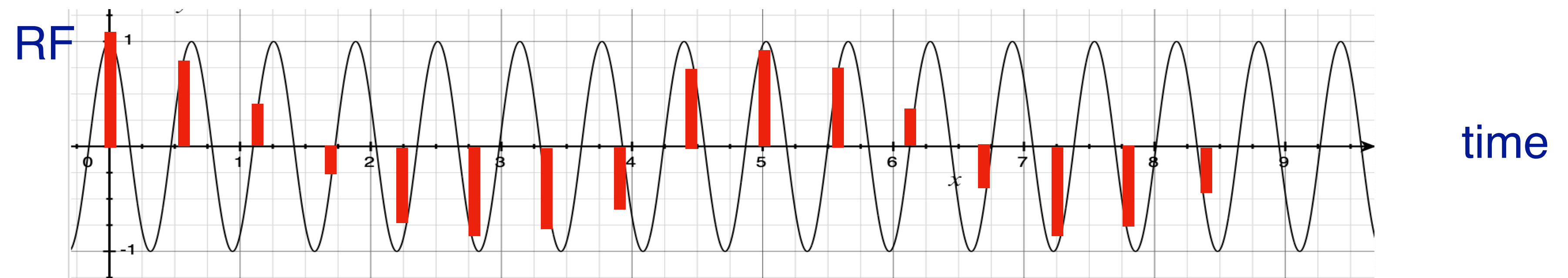
$$V_{gap} \text{ (envelope)} = V_0 \cos (\omega_{rev} - \omega_{rf}) t$$

$$\omega_{rf} > (1/2) \omega_{rev}$$

(aliasing, beat, ...)

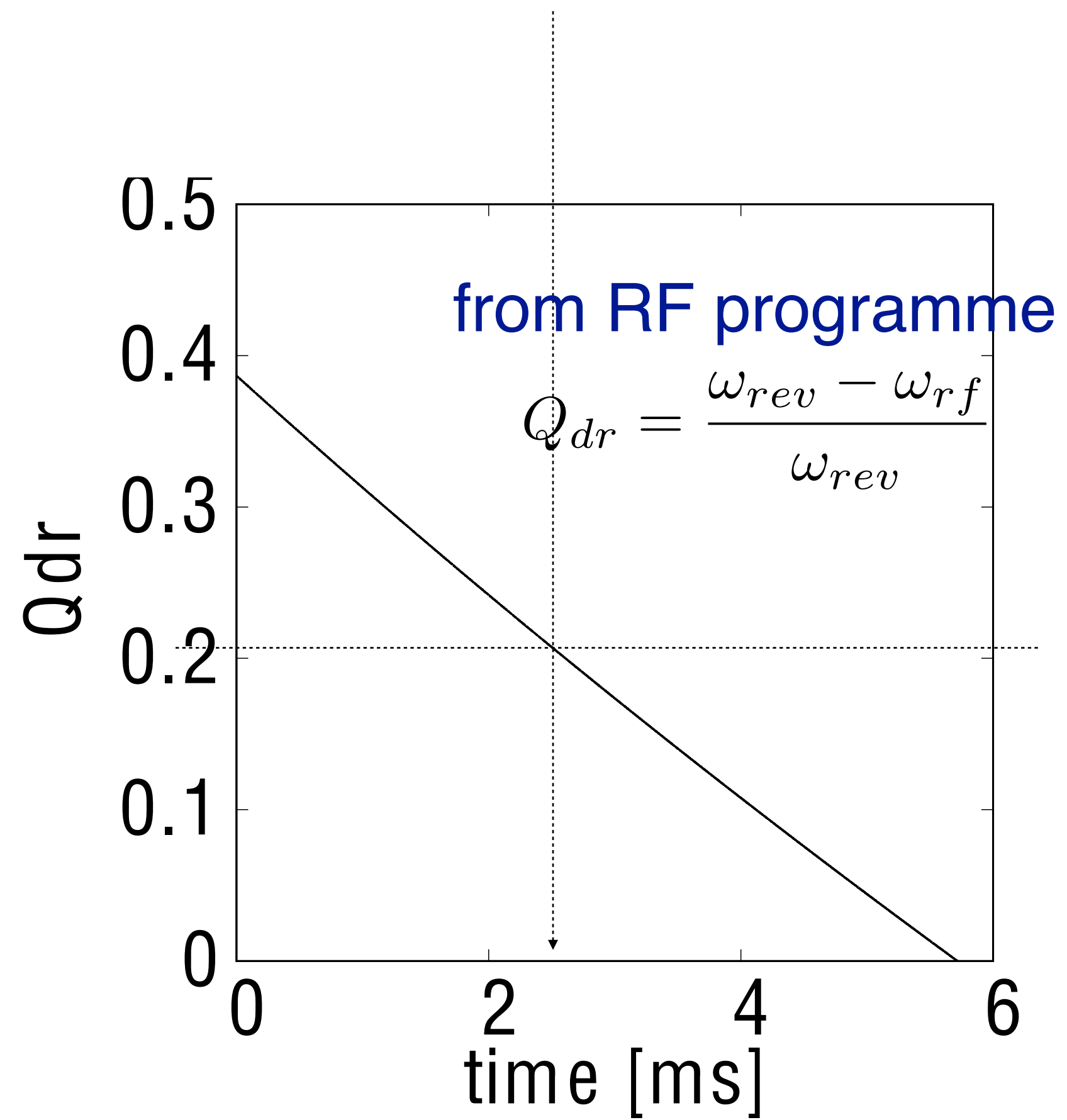
RF voltage seen by the coasting beam

$f_{rf} < f_{rev}$

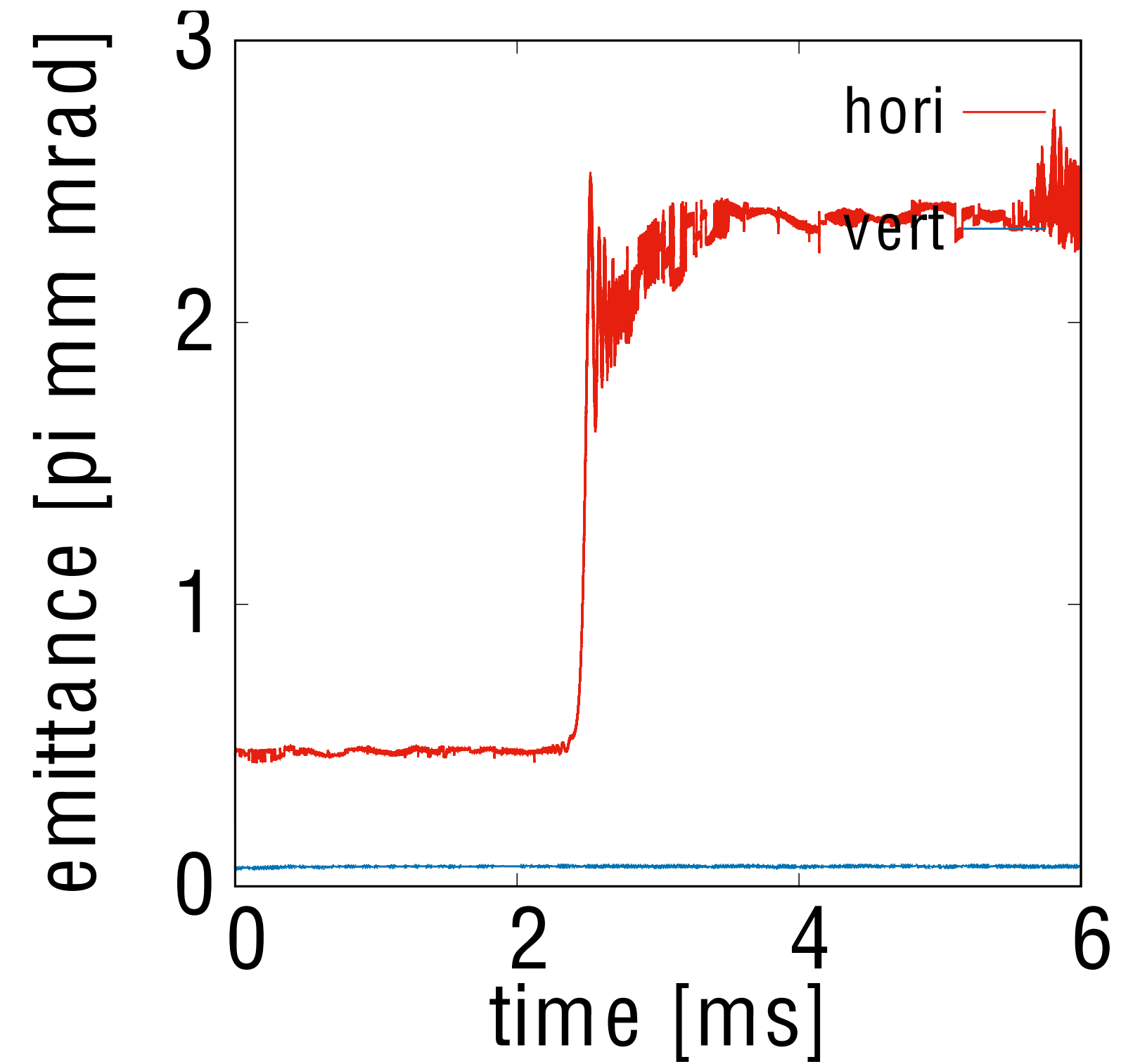


# Simulation with KURNS 3D field map

$$Q_h = 0.79 \quad \text{or} \quad 1 - Q_h = 0.21$$



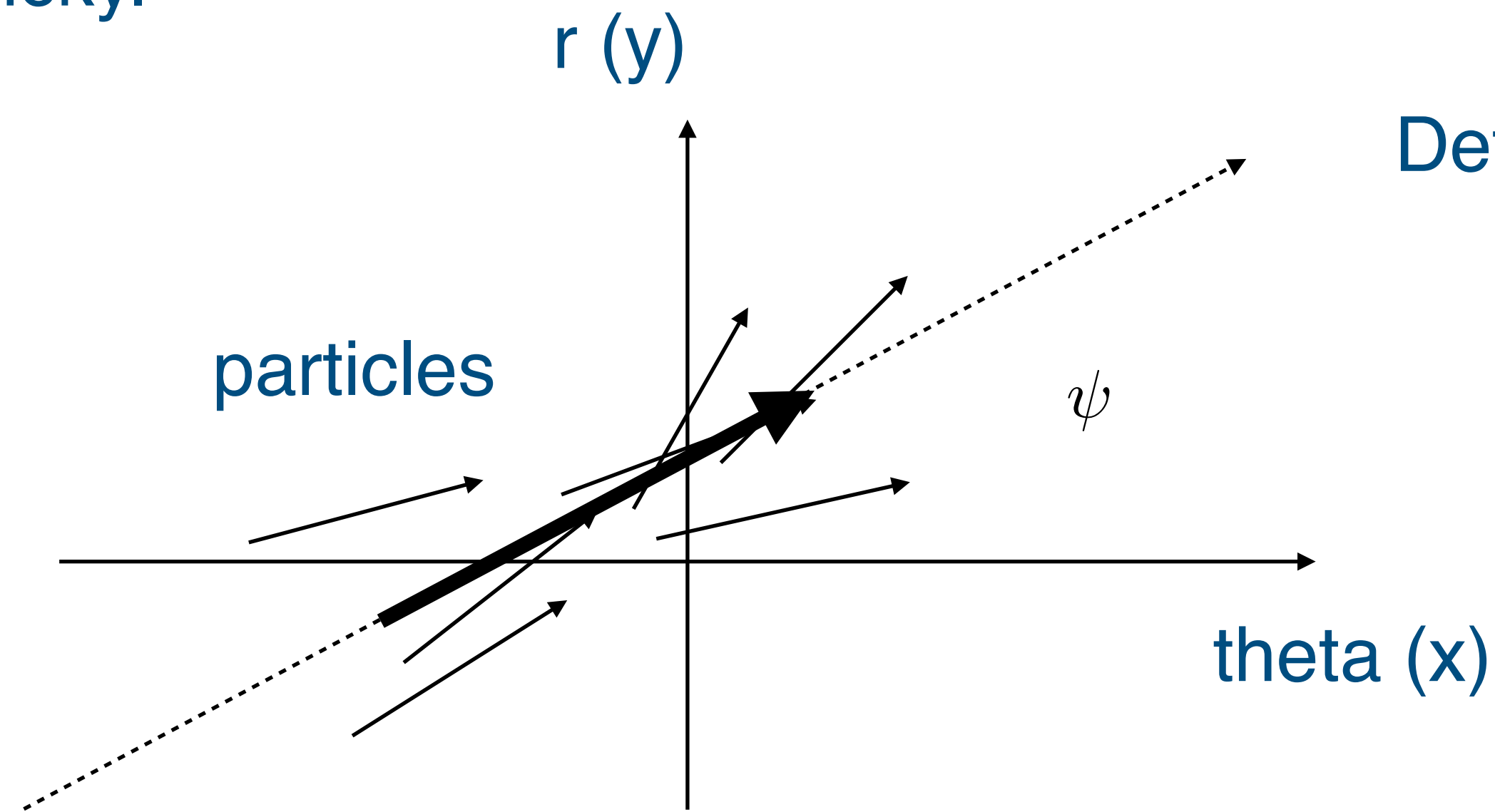
horizontal emittance growth  
at  $\sim 2.5$  ms.



We should see something at  $\sim 2.5$  ms

# Scode's way of modelling

- Ideally, transverse (radial) position is measured from the closed orbit, not from the average position within a bin.
- It may be possible to define the instantaneous closed orbit, but could be tricky.



Define the close orbit within a bin as  
A straight line with a gradient of

$$\tan(\psi) = \frac{\sum p_{y,i}}{\sum p_{x,i}}$$

which goes through the point of

$$\left( \frac{\sum x_i}{n}, \frac{\sum y_i}{n} \right)$$

where n is the number of particle and i is index.

