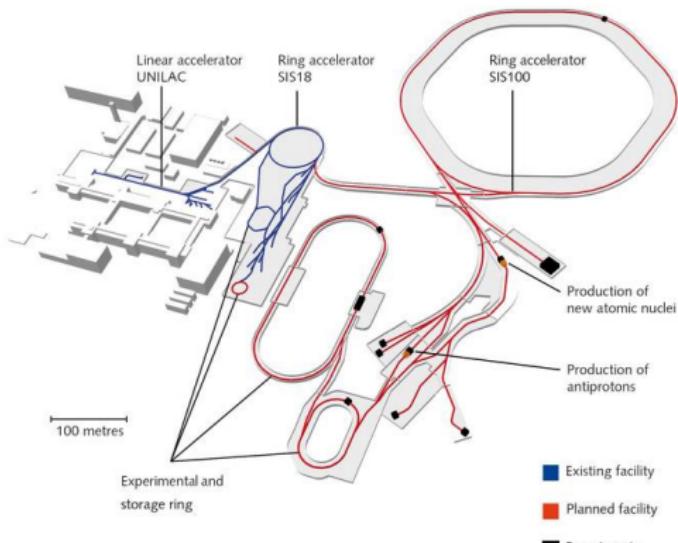


Space Charge Limits and Possible Mitigation Approaches in the FAIR Synchrotrons

Adrian Oeftiger

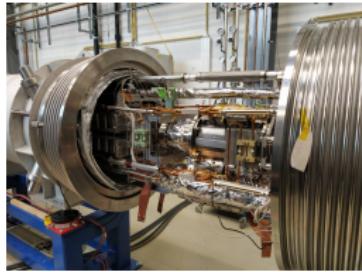
HB2023, Geneva, Switzerland

9 October 2023



Timelapse 2018 - 2023

drone video of construction site ↗



- string test of full SIS100 arc cell installed
- first SIS100 accelerator section to be installed in January 2024,
IPAC'23 paper on SIS100 status ↗

Motivation

Facility for Antiproton and Ion Research

- SIS100: deliver high-intensity hadron beams

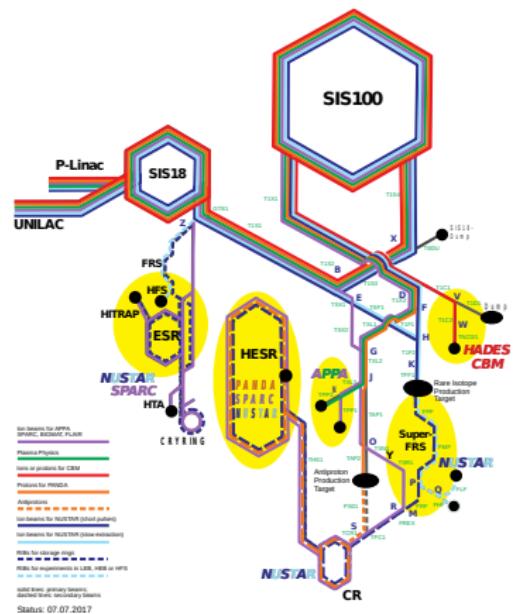


Figure: FAIR complex

Motivation

Facility for Antiproton and Ion Research

- SIS100: deliver high-intensity hadron beams
- crucial for performance: maintain beam quality during 1-sec injection plateau
- reference case: uranium U^{28+} beam
 - largest beam size vs. transverse aperture
 - space charge induced losses
 - ~~ important: dynamic vacuum stability
 - ==> low-loss operation <5%!

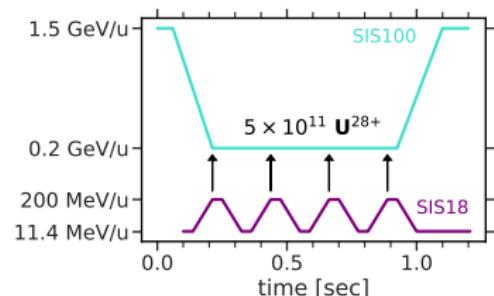


Figure: SIS18 to SIS100 transfer

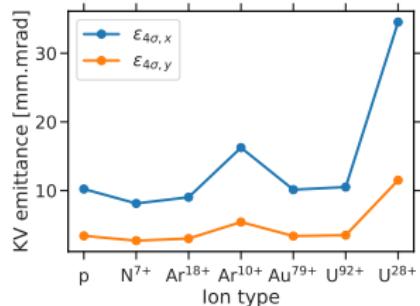


Figure: scaled beam sizes at 18 Tm

Motivation

Facility for Antiproton and Ion Research

- SIS100: deliver high-intensity hadron beams
- crucial for performance: maintain beam quality during 1-sec injection plateau
- reference case: uranium U^{28+} beam
 - largest beam size vs. transverse aperture
 - space charge induced losses
 - ~~ important: dynamic vacuum stability
 - ==> low-loss operation < 5%!

key questions

- What is the maximum tolerable intensity at the space charge limit?
- How to increase the space charge limit?

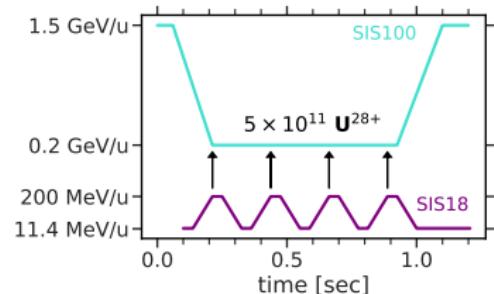


Figure: SIS18 to SIS100 transfer

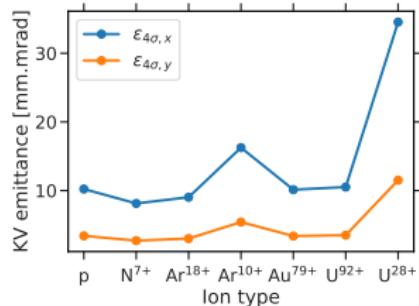


Figure: scaled beam sizes at 18 Tm

Structure:

A. The Model

B. Betatron Resonances:

- Intrinsic from Space Charge
- External from Field Errors

C. Space Charge Limit

D. Mitigation Measures: Conventional & Novel

- β -beat Compensation
- Bunch Flattening
- Pulsed Electron Lenses

A. The Model

Simulation model:

- track macro-particles (m.p.) through accelerator lattice & space charge kicks

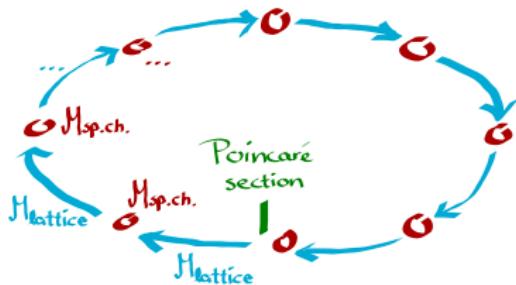


Figure: sketch of simulation model

Simulation model:

- track macro-particles (m.p.) through accelerator lattice & space charge kicks
- nonlinear 3D space charge (SC) models:
 - self-consistent **PIC**: particle-in-cell for open-boundary Poisson equation
 - fixed frozen (**FFSC**): constant field map independent of m.p. dynamics
 - (adaptive frozen (**AFSC**): field map scaled with m.p. distribution momenta)

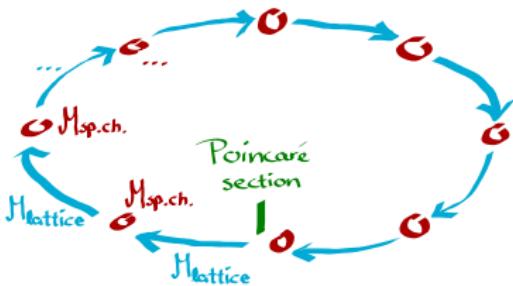


Figure: sketch of simulation model

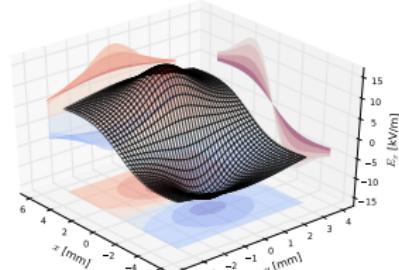


Figure: horizontal space charge field

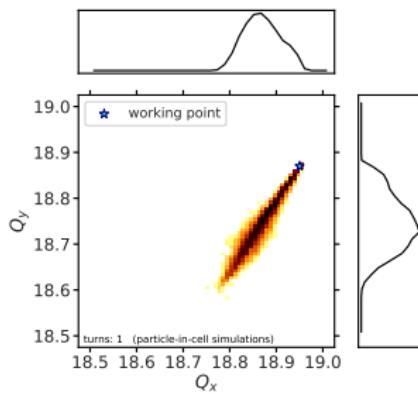
Maximum SC Tune Shift

$$\Delta Q_y^{\text{SC}} = -\frac{r_c \lambda_{\max}}{\beta_0^2 \gamma_0^3} \oint \frac{ds}{2\pi} \frac{\beta_y(s)}{\sigma_y(s)(\sigma_x(s) + \sigma_y(s))}$$

r_c : classical ion radius

λ_{\max} : maximum line density

β_0 : speed in [c] γ_0 : Lorentz factor $\sigma_{x,y}$: local rms beam size



Hor. norm. rms emittance ϵ_x	5.9 mm mrad
Vert. norm. rms emittance ϵ_y	2.5 mm mrad
Rms bunch length σ_z	13.2 m
Bunch intensity N_0 of U^{28+} ions	6.25×10^{10}
Max. space charge ΔQ_y^{SC}	-0.30
Rms chromatic $Q'_{x,y} \cdot \sigma_{\Delta p/p_0}$	0.01
Synchrotron tune Q_s	4.5×10^{-3}
Kinetic energy	$E_{\text{kin}} = 200 \text{ MeV/u}$
Relativistic β factor	0.568
Revolution frequency f_{rev}	157 kHz

B. Betatron Resonances

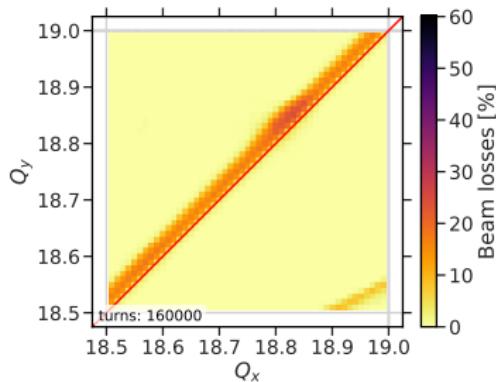


Figure: tune diagram of beam loss

Symmetric error-free SIS100 lattice:

- perfect dipole and quadrupole magnets
 - exact symmetry of $S = 6$
 - space charge → only source for resonances
 - simulated for 160'000 turns = 1 second
- ⇒ mainly Montague resonance visible
- ⇒ absence of low-order structure resonances!

Only Space Charge

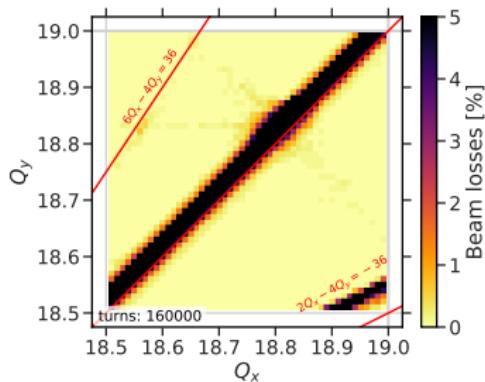


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- simulated for 160'000 turns = 1 second
- \Rightarrow mainly Montague resonance visible
- \Rightarrow absence of low-order structure resonances!

Montague Resonance

Montague resonance $2Q_x - 2Q_y = 0$:

- 4th-order resonance
 - intrinsically driven by space charge
 - transverse emittance exchange for anisotropic beams
- ⇒ stopband always present around $Q_x \approx Q_y$ for SIS100 beams

Space charge model predictions:

- bad: “adaptive frozen” resolves full exchange but predicts too large stopband extent
- + good: “fixed frozen” reproduces stopband edges well!

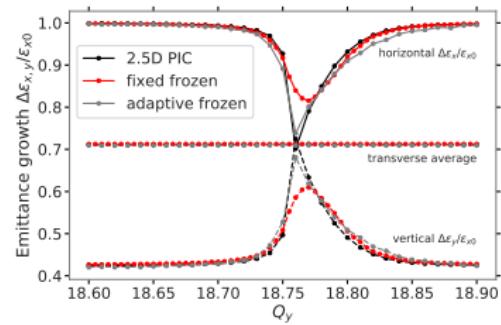


Figure: emittance exchange

Montague resonance $2Q_x - 2Q_y = 0$:

- 4th-order resonance
- intrinsically driven by space charge
- transverse emittance exchange for anisotropic beams

⇒ stopband

for “Fixed frozen” model much better suited than “adaptive frozen” to approximate realistic PIC when **identifying loss-free conditions!**

Space charge model predictions:

- bad: “adaptive frozen” resolves full exchange but predicts too large stopband extent
- + good: “fixed frozen” reproduces stopband edges well!



Figure: emittance exchange

Warm Quadrupoles

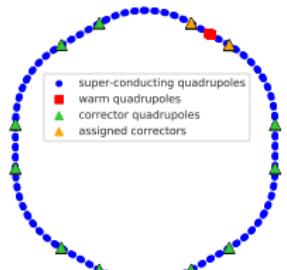


Figure: SIS100 quadrupole survey

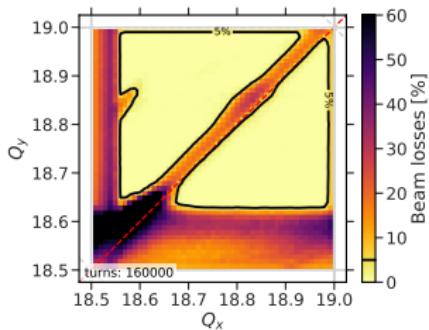
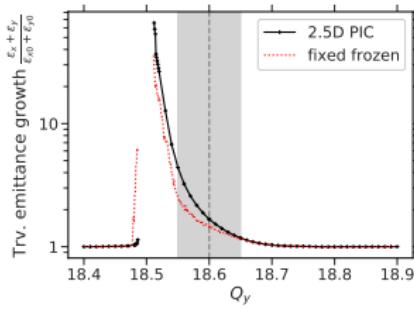


Figure: corrected warm quadrupoles

Real SIS100 lattice:

- 2 cold quadrupoles replaced by warm / normalconducting quadrupoles (radiation hardened, required in extraction region)
- breaking of $S = 6$ symmetry by gradient error
 - externally driven half-integer resonance
 - can be minimised by quadrupole correctors
- FFSC reproduces PIC stopband edges!



Field Error Model

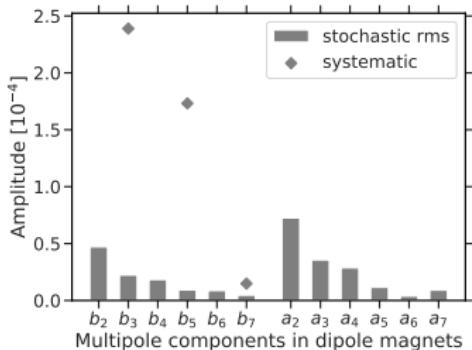


Figure: dipole magnets

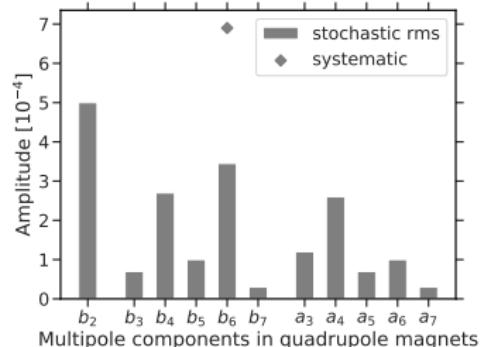


Figure: quadrupole magnets

Field error model extracted from cold bench measurements of magnet units:

- stochastic amplitudes drive non-systematic resonances
- random number sequence → multipole errors for every dipole and quadrupole magnet

quadrupole model displayed here corresponds to PRAB paper version (based on stamped FoS),
see GSI-2021-00450 report / for model based on series production and its comparison

Field Error Model from Beam

HB'23 talk on Wednesday, C. Caliari on “Deep Lie Map Network” ↗:

- machine learning approach:
train linear & nonlinear field errors on kick turn-by-turn data
 - start from model lattice, learn error multipoles $k \cdot L$
- ➡ effective lattice to better reproduce machine behaviour

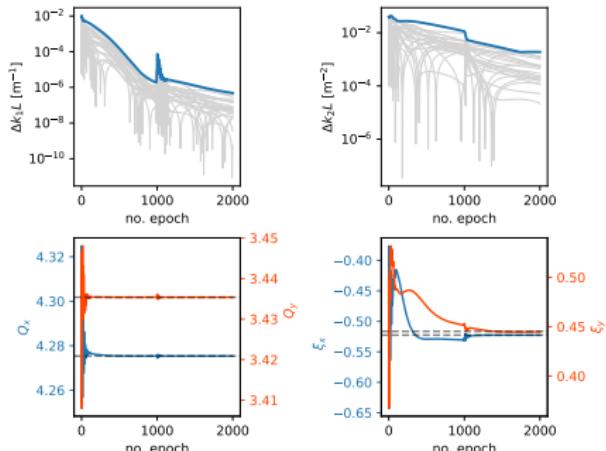
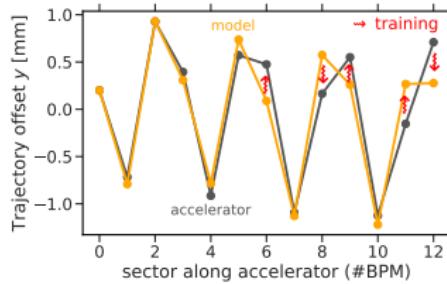


Figure: learning of quadrupole & sextupole errors

Full Model with Space Charge

Linear and nonlinear resonances driven by magnet field errors. Resonance condition without space charge:

$$mQ_x + nQ_y = p \quad \text{for } m, n, p \in \mathbb{Z}$$

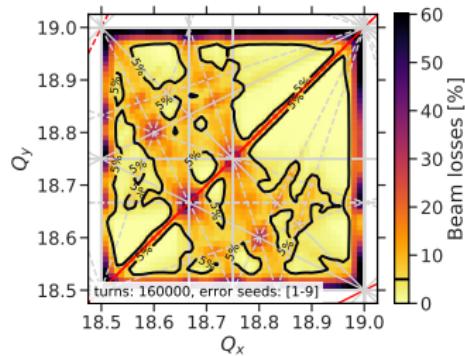


Figure: no space charge

Full Model with Space Charge

Linear and nonlinear resonances driven by magnet field errors. Resonance condition without space charge:

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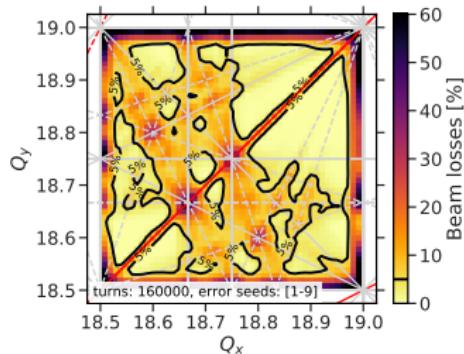


Figure: no space charge

include
SC

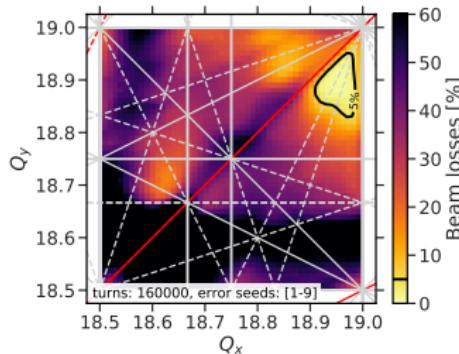


Figure: with fixed frozen space charge

- SC broadens existing resonance stopbands
- optimal working point area around $(Q_x, Q_y) = (18.95, 18.87)$

Validation with Self-consistent PIC

Self-consistent PIC simulations:

- ✓ validated Montague resonance
- ✓ validated half-integer resonance
- now validate full error model FFSC predictions for beam loss

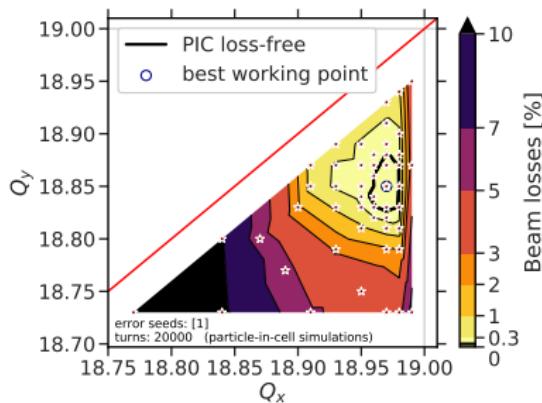


Figure: self-consistent PIC simulations

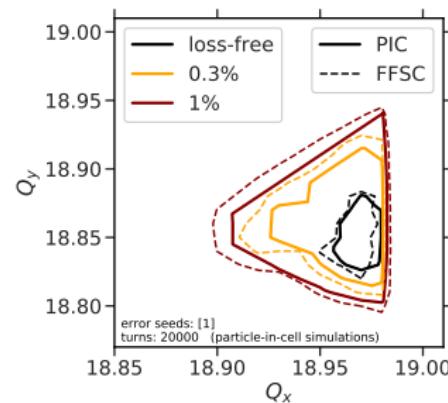


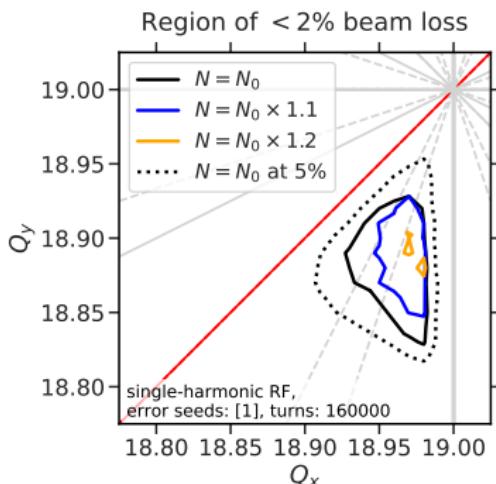
Figure: comparison between SC models

note: PIC simulations take 2 days (on NVIDIA V100 GPU) vs. FFSC simulations with 7 min (on 16 CPU cores, HPC AMD)

C. Space Charge Limit

dynamic definition of space charge limit

reached when loss-free working point area vanishes



Keeping all beam parameters identical,
increasing N :

⇒ U^{28+} space charge limit at **120%** of
nominal bunch intensity N_0 :

$$\max |\Delta Q_y^{\text{SC}}| = 0.36$$

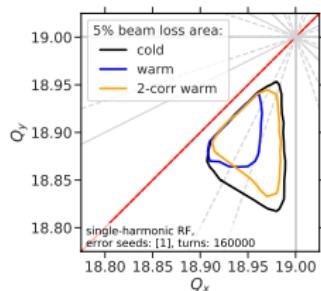
Figure: low-loss area for increasing N

D. Mitigation Measures: Conventional & Novel

Correction of β -beat

Two sources of β -beat (gradient error):

- warm quadrupoles: uncorrected = 2%



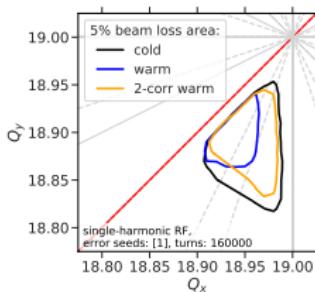
(a) low-loss area with warm quads

Correction of β -beat

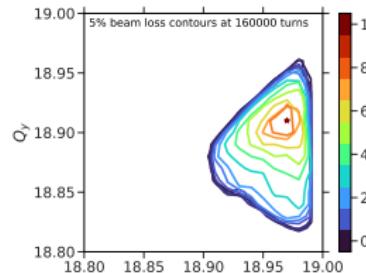
Two sources of β -beat (gradient error):

- warm quadrupoles: uncorrected = 2%
- distributed b_2 : $\approx 0.5\%$

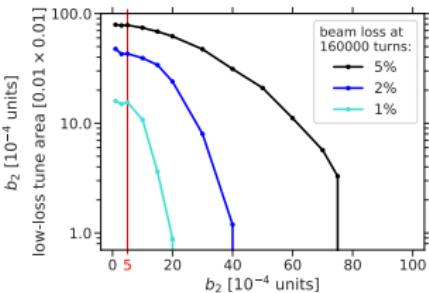
⇒ below $b_2 = 10$ units: no significant effect on low-loss area size



(a) low-loss area with warm quads



(b) low-loss area with b_2



(c) size of low-loss area vs. b_2

A Word on Coherent Stability

For SIS100, space charge parameter $q = \Delta Q_y^{SC} / 2Q_s \approx 33$ at design $N = N_0$:

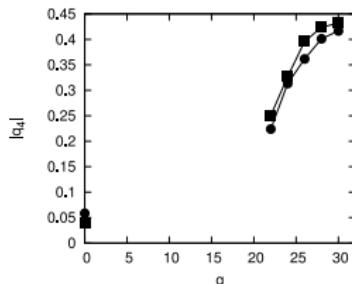
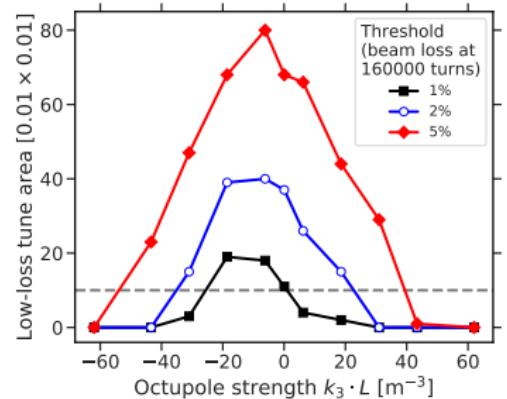


Figure 4: Results of the simulations scans for the $k = 1$ mode: stability thresholds of the octupole power in a dependency from the space-charge parameter q . The circles are for the octupole polarity $q_4 > 0$, the squares are for the octupole polarity $q_4 < 0$.

(a) Required octupole strength for stabilisation of single-bunch resistive-wall instability
[V. Kornilov, IPAC'23]



(b) Incoherently tolerable octupole strengths

- octupole current of $k_3 L = 35 \text{ m}^{-3}$ corresponds to $q_4 = 0.55$.
- ⇒ single-bunch stability through Landau damping from octupoles, transverse feedback system required for coupled-bunch stability

Double-harmonic RF

Add $h=20$ harmonic in bunch lengthening mode:

$$V_{h=20} = V_{h=10}/2$$

⇒ obtain flattened bunches with reduced line density at 80% of nominal λ_{\max} .

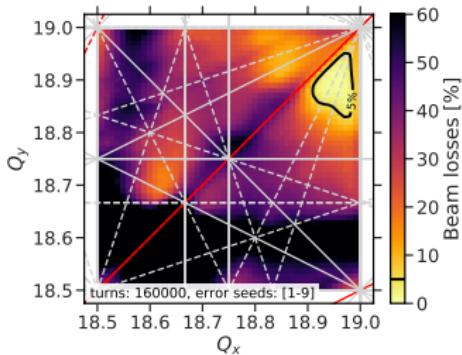


Figure: single-harmonic RF

flatten
⇒
bunch

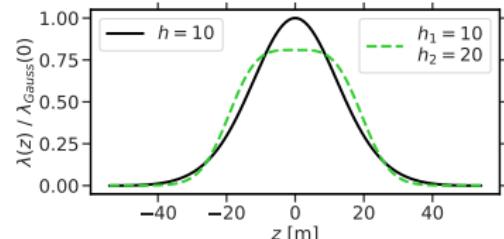


Figure: rms-equivalent line densities

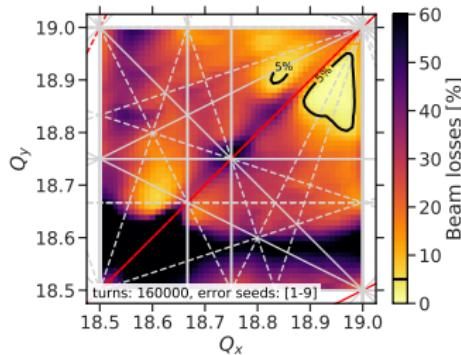


Figure: double-harmonic RF

SC Limit with Double-harmonic RF

Increasing N for double-harmonic RF:

- find space charge limit at **150%** of nominal intensity N_0

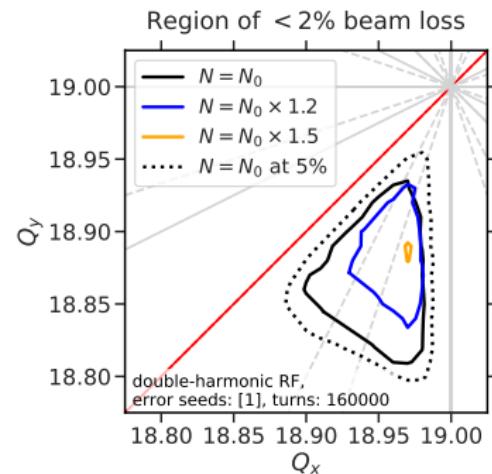


Figure: low-loss area for increasing N

Novel: Pulsed Electron Lenses

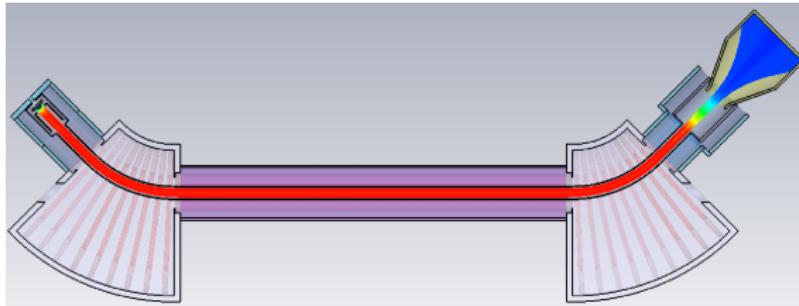


Figure: e-lens model for SIS18 [K. Schulte-Urlich et al., IPAC'22] ↗



Figure: Modulation grid.

Short insertion (here $L = 3.36\text{ m}$) with co-propagating electron beam:

- transversely homogeneous distribution
- longitudinally modulated to match ion bunch profile
- compensate longitudinal dependency of space charge
- ⇒ **suppress periodic resonance crossing**
- additionally provide strong Landau damping for head-tail modes:
V. Gubaidulin et al., PRAB 25, 084401 (2022) ↗ [tbc with strong SC]

Tune Footprint vs. E-Lens Compensation

Some n_{el} e-lenses with I_e current and rms beam size σ_e provide tune shift:

$$\Delta Q_y^e = \frac{1}{4\pi} \sum_{k=1}^{n_{\text{el}}} \beta_y(s_k) \frac{r_c}{Ze} \frac{I_e}{\sigma_e^2 \gamma_0} \frac{1 - \beta_e \beta_0}{\beta_e} \frac{L}{\beta_0 c}$$

Define linear compensation degree (for Gaussian bunches $\Delta Q^{\text{KV}} = \Delta Q^{\text{SC}}/2$):

$$\alpha \doteq \frac{\Delta Q^e}{|\Delta Q^{\text{KV}}|}$$

Remarks:

- dipole tune increases with

$$\Delta Q_{\text{dip}} = \alpha \cdot \Delta Q^e$$

- without chroma, $\alpha = 0.5$ yields smallest tune spread!

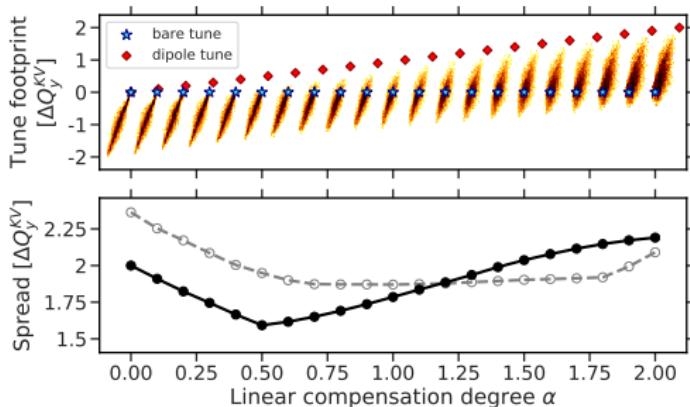
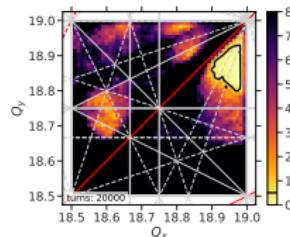


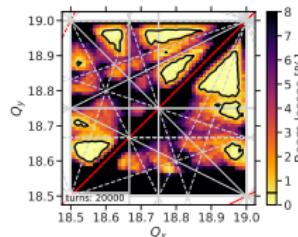
Figure: Gaussian bunch, tune footprint vs. e-lens strength (black: $\Delta p/p_0 = 0$, grey: with natural chromatic detuning)

Optimal E-Lens Configuration

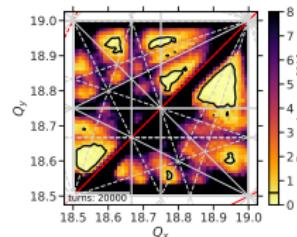
In SIS100 with natural chromaticity:



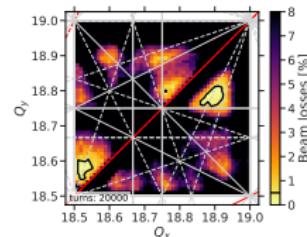
(a) $\alpha = 0$



(b) $\alpha = 0.5$



(c) $\alpha = 0.7$



(d) $\alpha = 1$

Figure: FAIR design intensity $N = N_0$ with $n_{\text{el}} = 3$ pulsed e-lenses.

- optimal choice of α depends on nearby resonances
 - ➡ depends on particularities of synchrotron
- SIS100: at low $n_{\text{el}} \leq 6$, $\alpha = 0.5$ optimal vs. high $n_{\text{el}} > 6$, $\alpha = 0.7$ better

SC Limit with Pulsed E-Lenses

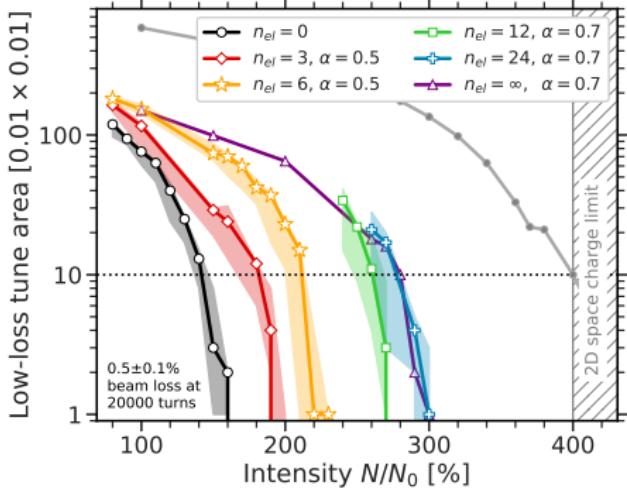


Figure: low-loss area for increasing N

Table: SC limit with electron lenses.

Number n_{el}	SC limit	Gain
0	$1.4 \cdot N_0$	100%
3	$1.8 \cdot N_0$	130%
6	$2.1 \cdot N_0$	150%
12	$2.6 \cdot N_0$	185%
$24, \infty$	$2.8 \cdot N_0$	200%

Remarks:

- SC limit scales well
- $n_{el} = 24$ case saturates gain
- theoretical 2D limit ($Q_s = 0$, no e-lenses) = by construction no periodic resonance crossing
⇒ reached after $n_{el} = 24, \infty$

Involved Publications



PHYSICAL REVIEW ACCELERATORS AND BEAMS **25**, 054402 (2022)

Simulation study of the space charge limit in heavy-ion synchrotrons

Adrian Oeftiger^{a,*}, Oliver Boine-Frankenheim^{a,1,2}, Vera Chetvertkova^{a,1},
Vladimir Kornilov^{a,1}, Dmitrii Rabusov^{a,1}, and Stefan Sorge^{a,1}

^aGSI Helmholtzzentrum für Schwerionenforschung

²Technische Universität Darmstadt

(Received 5 November 21)

Pulsed Electron Lenses for Space Charge Mitigation

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(Dated: October 5, 2023)



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ELECTRON LENSES FOR MODERN AND FUTURE ACCELERATORS

Pulsed electron lenses for space charge compensation in the FAIR synchrotrons

S. Artikova,^a O. Boine-Frankenheim,^{a,b,*} O. Meusel,^c A. Oeftiger,^a D. Ondreka,^a
K. Schulte-Urichs^c and P. Spiller^c



Nuclear Inst. and Methods in Physics Research, A 1040 (2022) 167290

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Nuclear Inst. and Methods in Physics Research, A

journal homepage: www.elsevier.com/locate/nima



Characterization and minimization of the half-integer stop band with space charge in a hadron synchrotron

Dmitrii Rabusov^{b,*}, Adrian Oeftiger^b, Oliver Boine-Frankenheim^{a,b}

^aTechnische Universität Darmstadt, Schlossgartenstr. 8, 64289 Darmstadt, Germany

^bGSI Helmholtzzentrum für Schwerionenforschung GmbH, Planckstr. 1, 64291 Darmstadt, Germany



Conclusion

Summary:

- identified **optimal tune area** in SIS100 around $(Q_x, Q_y) = (18.95, 18.87)$
- explored **space charge limit**: $\max |\Delta Q_y^{\text{SC}}| = 0.36$
 - nominal SIS100: +20% intensity
 - double-harmonic RF: +50% intensity
 - 3 pulsed electron lenses: +70..80% intensity
- FAIR start planned in 2028 with “Early Science” programme

Conclusion

Summary:

- identified **optimal tune area** in SIS100 around $(Q_x, Q_y) = (18.95, 18.87)$
- explored **space charge limit**: $\max |\Delta Q_y^{\text{SC}}| = 0.36$
 - nominal SIS100: +20% intensity
 - double-harmonic RF: +50% intensity
 - 3 pulsed electron lenses: +70..80% intensity
- FAIR start planned in 2028 with “Early Science” programme

take-home messages

- fixed frozen SC model fast & validated tool to identify resonance-free tunes
- dynamic space charge limit: find based on tolerable loss & emittance growth
- pulsed electron lenses: optimum configuration for space charge mitigation

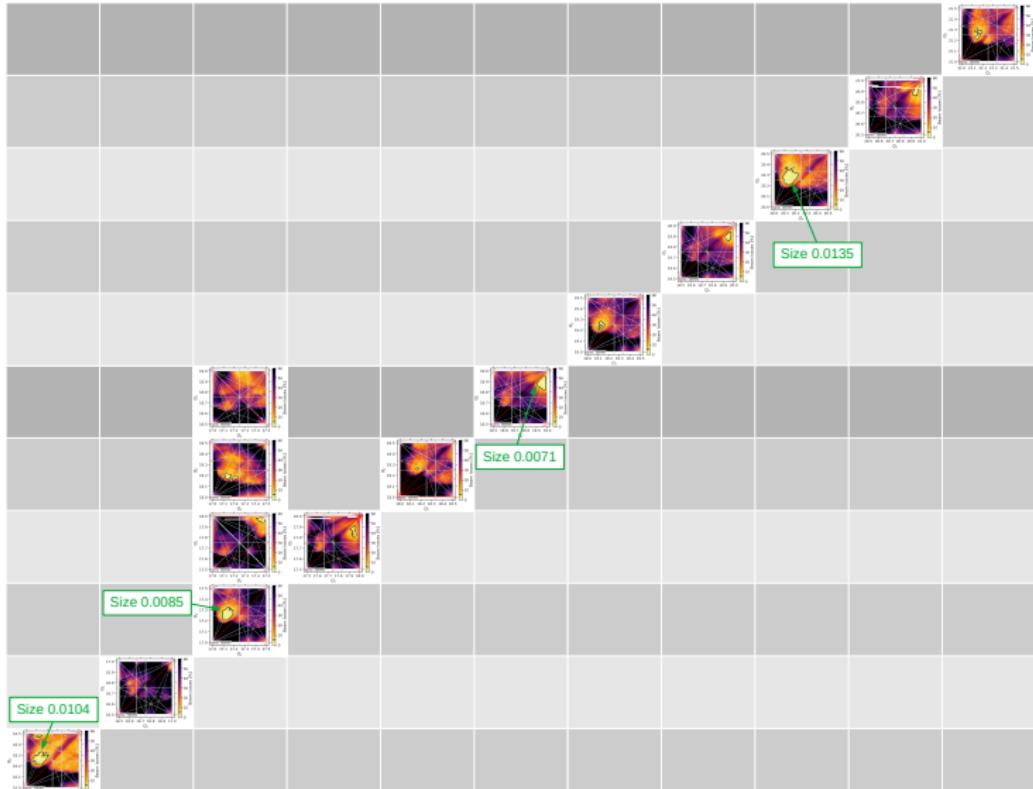
Thank you for your attention!

Acknowledgements:

GSI: O. Boine-Frankenheim, V. Chetvertkova, V. Kornilov, D. Rabusov, S. Sorge,
D. Ondreka, A. Bleile, V. Maroussov, C. Roux, K. Sugita

CERN: R. de Maria, G. Iadarola, M. Schwinzerl

Grand Overview Tune Diagrams



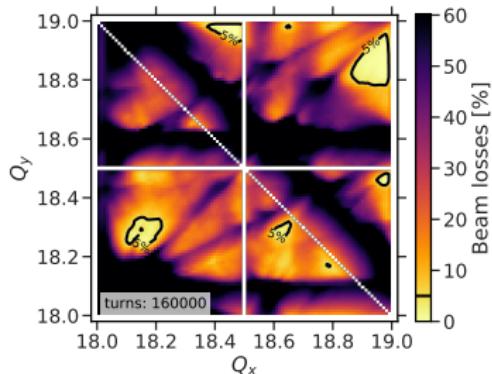


Figure: GPU simulation results for latest magnet field error model

Thanks to GSI's new high-performance GPU cluster in Green Cube:

- 400 GPU cards of today's most performant model (AMD Radeon Instinct MI100)
- even faster simulations, larger tune scans in shorter times
- ➡ following up magnet series production and doublet assembly

Why Homogeneous E-lenses?

O. Boine-Frankenheim and W. Stem, NIM A 896 (2018) 122–128 ↗:

- transversely Gaussian distributed pulsed e-lenses drive **strong systematic nonlinear resonances**
- ➡ configuration with low number of Gaussian e-lenses not feasible

Their 2D simulation results with self-consistent PIC:

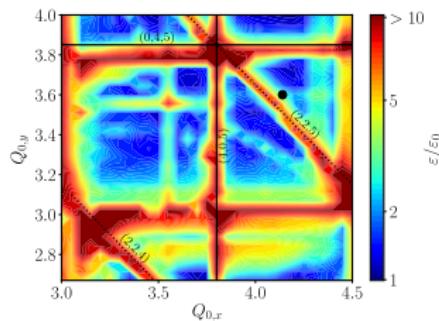


Fig. 8. 2D PIC tune scan for $N_e = 3$, $\Delta Q_{x,y} = -0.2/0.4$ and $\alpha = 0.5$ for a Gaussian electron beam.

(a) Gaussian pulsed e-lens

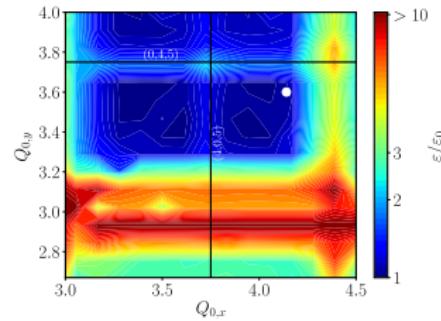


Fig. 9. 2D PIC tune scan for $N_e = 3$, $\Delta Q_{x,y} = -0.2/0.4$ and $\alpha = 0.5$ for a homogeneous electron beam.

(b) Homogeneous pulsed e-lens

Location of E-lenses

Electron lenses are placed in locations

- of equal $\beta_x = \beta_y$ (then electron pulse aspect ratio = ion beam for balanced compensation in both planes)
 - in the straight section (avoid dispersion!)
- ⇒ 4 possible locations per each sector (out of 6)

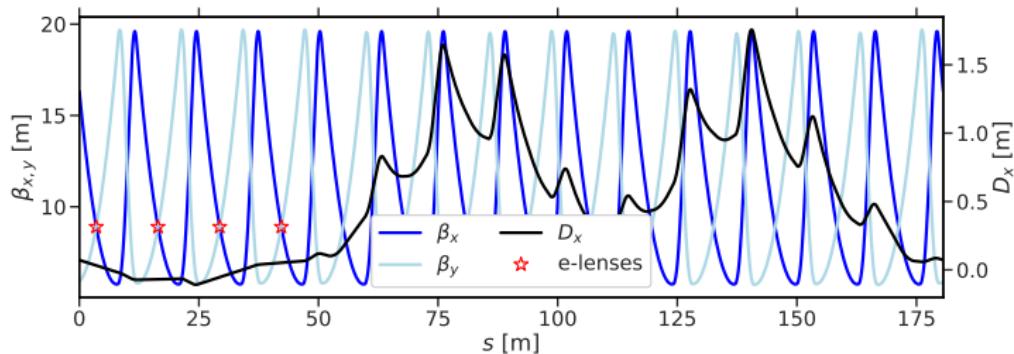
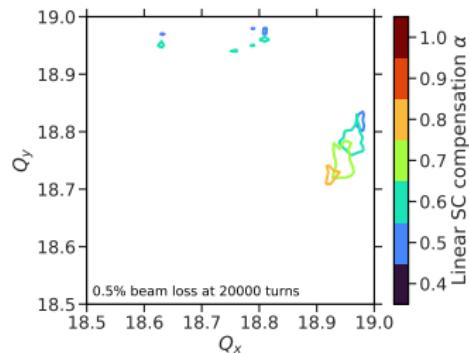


Figure: Optics functions for 1 sector in SIS100

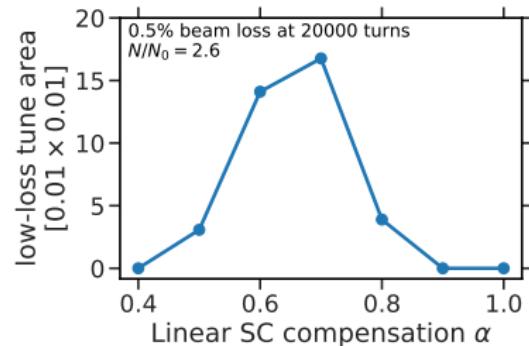
SC-Limit Dependency on α

For $\alpha = 0.5$, SC limit with 24 e-lenses reached at $N = 2.6 \cdot N_0$.

→ can we do better for other α ?



(a) loss boundaries vs. α



(b) low-loss tune area vs. α

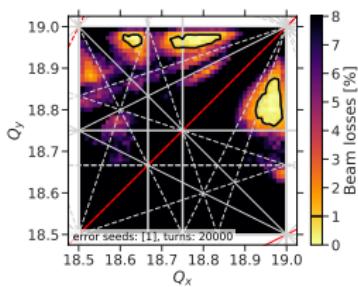
Figure: Finding optimal α at $N = 2.6 \cdot N_0$

→ yes! $\alpha = 0.7$ looks much better at higher $n_{el} > 6!$

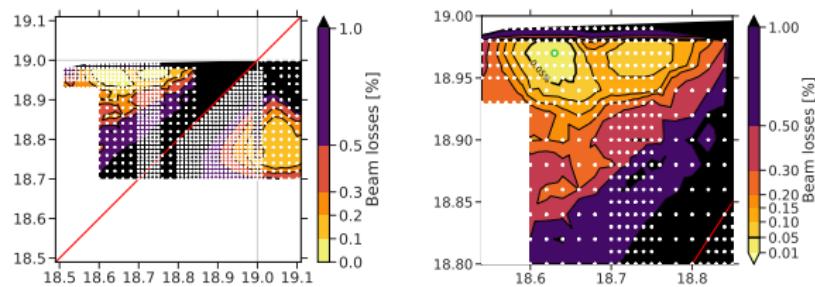
Cross-check with PIC

For 6 e-lenses at $N = 2 \cdot N_0$ and $\alpha = 0.5$, compare FFSC results to PIC:

FFSC (20k turns)



PIC (1k turns)



- no surprises (coherent resonances etc.), low-loss tune areas match
- FFSC results for SC limit validated!

Relevant Field Error Orders

Major resonances confining low-loss area:

- top left: Montague resonance
- right: integer resonance $Q_x = 19$
- bottom: higher-order resonances

Simulations with reduced field error model:

- identify sextupole and octupole orders
- $n = 3, 4$ as main limitation towards low Q_y

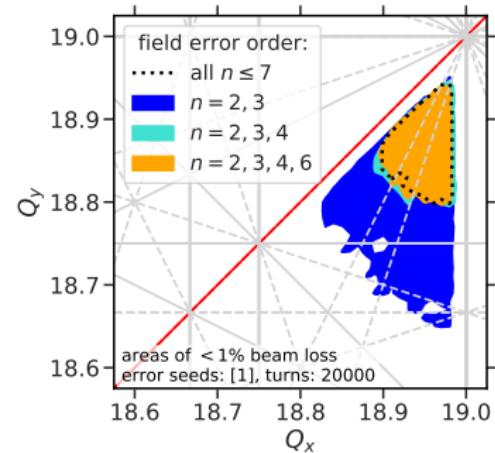


Figure: low-loss tune areas vs. multipole order

Emittance Growth at Low-loss Area

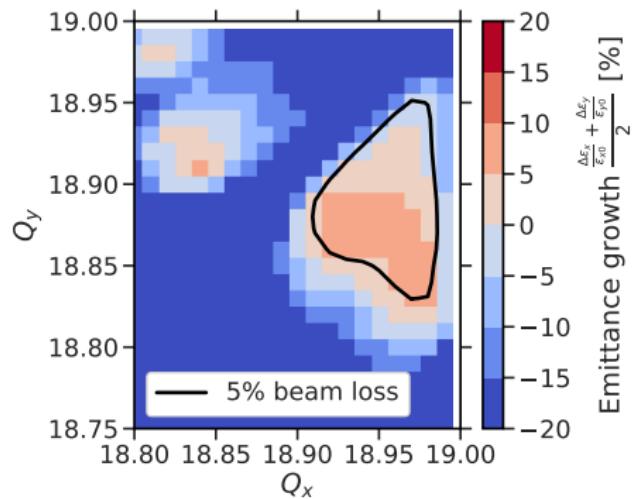


Figure: with space charge, emittance growth

PIC Results for Best Working Point

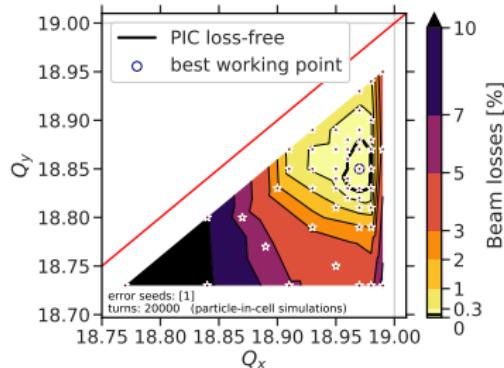


Figure: tune diagram with self-consistent PIC simulations

→ high-resolution¹ PIC results indeed show $\approx 2.5\% < 5\%$ beam loss in identified area

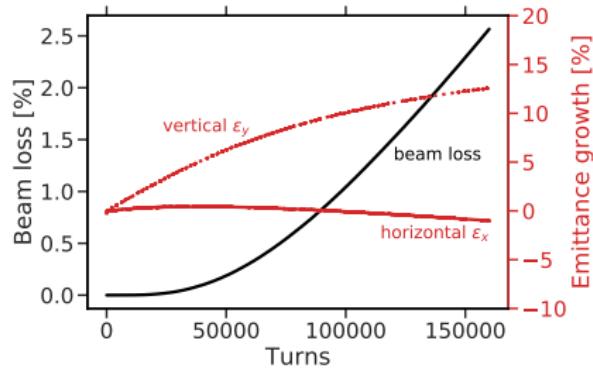


Figure: best working point $(Q_x, Q_y) = (18.97, 18.85)$

¹using 20×10^6 macro-particles in simulation such that resonance-free ideal lattice gives 0% beam loss, 160000 turns compute for \approx weeks on NVIDIA V100 GPU

Comparison 2.5D to 3D PIC

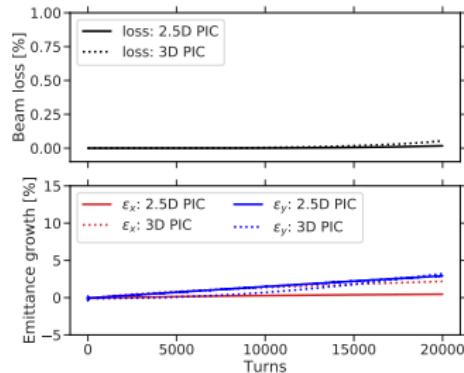


Figure: good working point (Q_x, Q_y) = (18.97, 18.85)

⇒ 3D confirms 2.5D PIC results

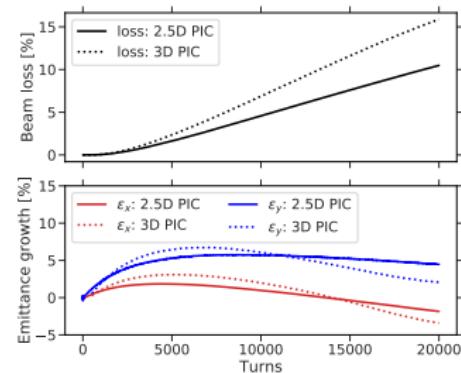
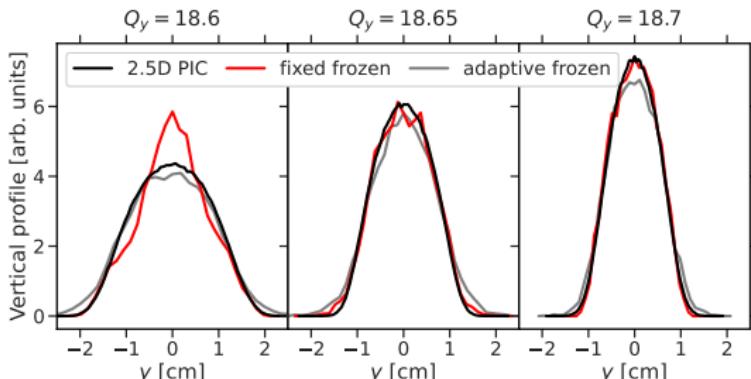


Figure: lossy working point (Q_x, Q_y) = (18.84, 18.73)

Adaptive Frozen SC (AFSC) Model



Observations from vertical beam profiles throughout half-integer stopband:

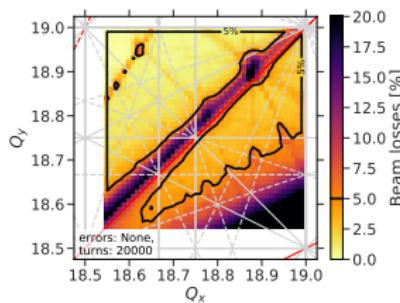
- black PIC = “truth”
 - left two panels: AFSC \approx PIC, only for strongly resonance-affected WPs
 - BUT right panel: FFSC \approx PIC, at edge of stopband
- ⇒ FFSC better suited than AFSC to reproduce stopband edges!

Why no Chromaticity Correction?

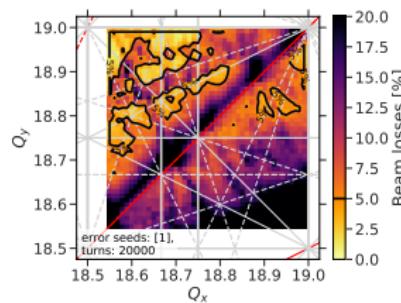
SIS100 lattice: optimised for dynamic vacuum stability with intermediately-charged heavy ions

- suppress dispersion where possible
 - cryo-catchers where dispersion \Rightarrow removal of charge-stripped ions
- require strong sextupole current for full chromaticity correction $Q'_{x,y} = 0$
- at injection: large beam sizes vs. restricted dynamic aperture

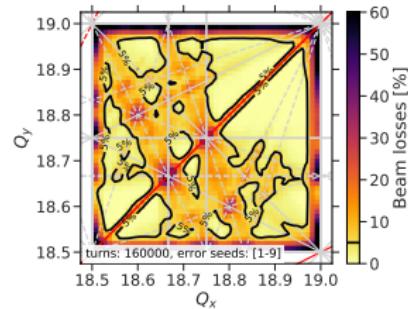
Simulation results without space charge:



(a) just sextupoles, no field errors



(b) sextupoles + field errors



(c) no sextupoles, just field errors

Beta-Beating from Warm Quadrupoles



Beta-beating with and without adjacent quadrupole correctors at $Q_{x,y} = (18.84, 18.73)$:

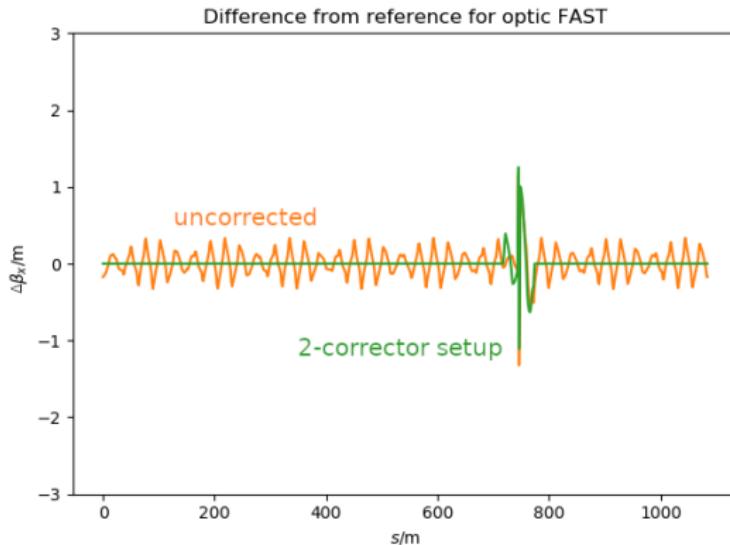
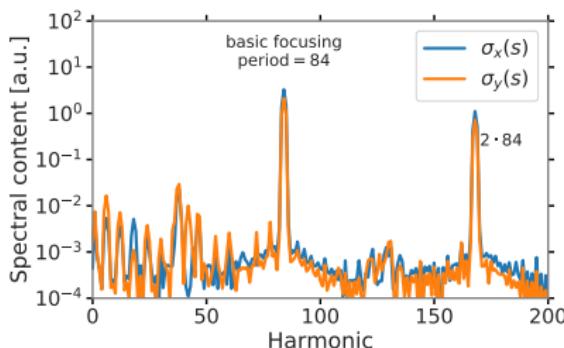


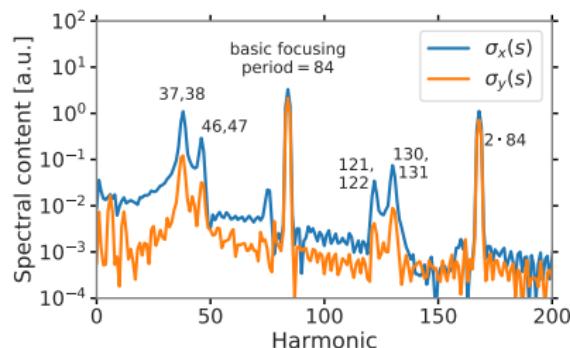
Figure: β -beat around SIS100 [courtesy D. Ondreka]

Harmonics with 2 Warm Quadrupoles

Fourier spectrum of single-turn envelope motion with space charge:



(a) Only cold quadrupoles



(b) With 2 warm quadrupoles

Observations:

- strong harmonics emerge at $2 \cdot Q_{x,y} \approx 38$ for half-integer resonance

Half-integer Resonance vs. SC



Characterization and minimization of the half-integer stop band with space charge in a hadron synchrotron

Dmitrii Rabusov ^{a,*}, Adrian Oeftiger ^b, Oliver Boine-Frankenheim ^{a,b}

^a Technische Universität Darmstadt, Schlossgartenstr. 8, 64289 Darmstadt, Germany

^b GSI Helmholtzzentrum für Schwerionenforschung GmbH, Planckstr. 1, 64291 Darmstadt, Germany

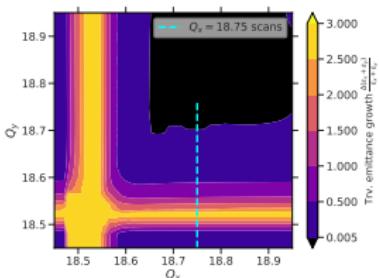


Figure: emittance growth in 2D tune diagram

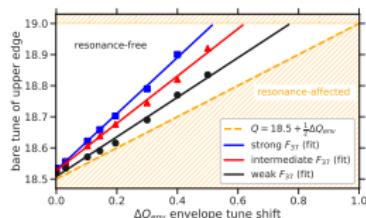


Figure: space charge limit for Gaussian bunches

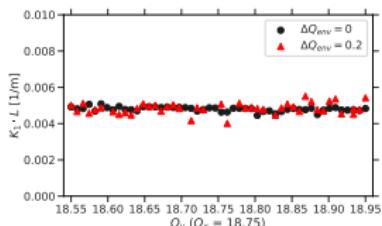


Figure: correction independent of space charge

Half-integer Resonance vs. SC



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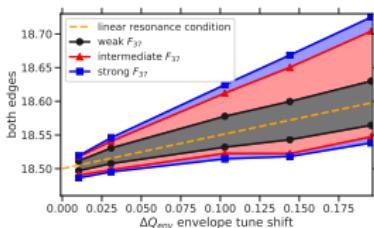


Figure: coasting beam stopband edges vs. space charge

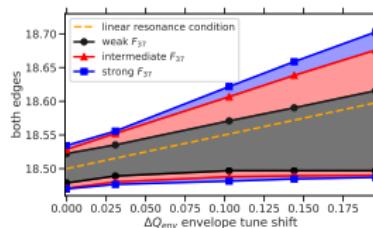


Figure: bunched beam stopband edges vs. space charge

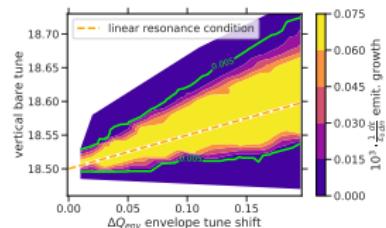


Figure: choose a threshold

Half-integer Resonance

Vertical half-integer stopband:

- no space charge, without $\Delta p/p_0$: $\delta Q_{\text{stopband}} = 0.023$
 - no space charge, with $\Delta p/p_0$: $\delta Q_{\text{stopband}} \sim 0.1$
 - incl. space charge: $\delta Q_{\text{stopband}} \sim 0.25$
- ⇒ fixed frozen SC model reproduces stopband edges from PIC

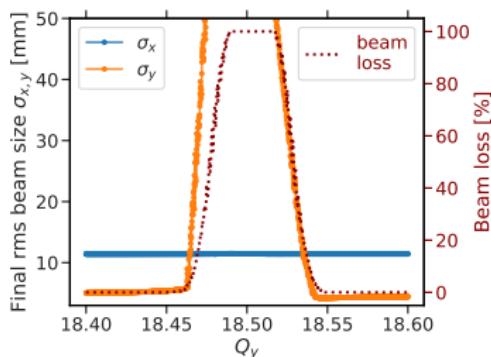


Figure: no space charge

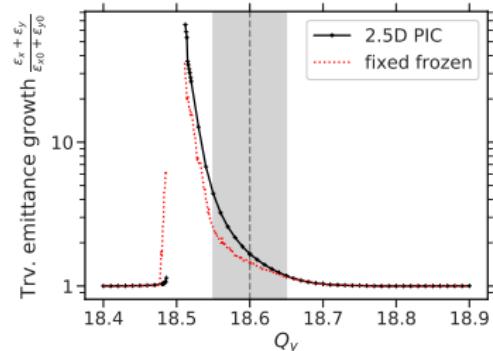
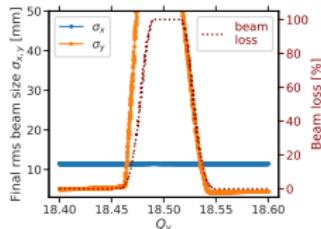


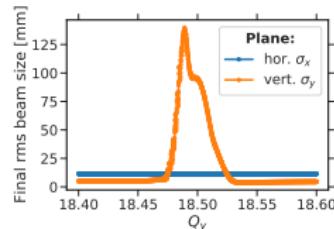
Figure: with space charge

Half-integer without SC but with Chroma

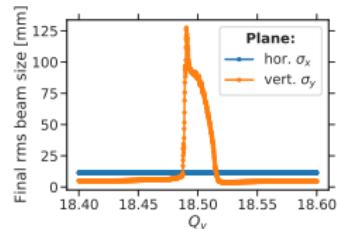
Chromaticity critical for half-integer width! ($Q'_y = -22.5$ and $\sigma_{\Delta p/p_0} = 0.5 \times 10^{-3}$)



(a) full momentum spread



(b) half momentum spread

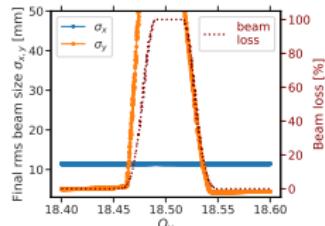


(c) 10% momentum spread

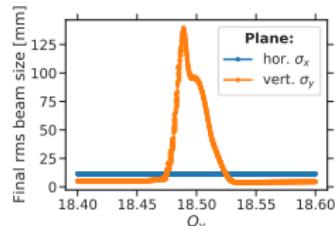
→ for vanishing $\sigma_{\Delta p/p_0} \rightarrow 0$, simulated width matches analytical $\Delta Q_{1/2} = 0.023$

Half-integer without SC but with Chroma

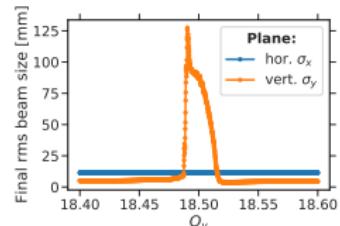
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(c) 10% momentum spread

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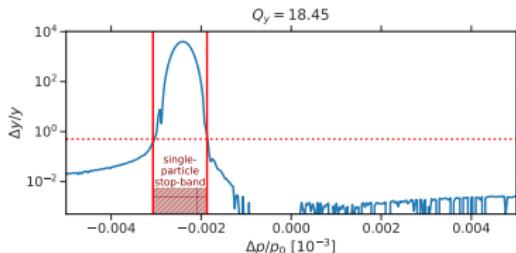
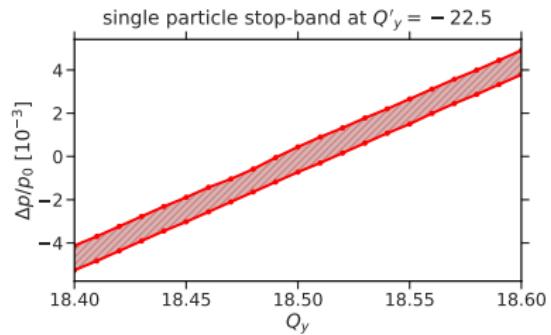
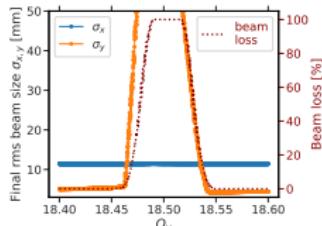


Figure: shifted single-particle stopband

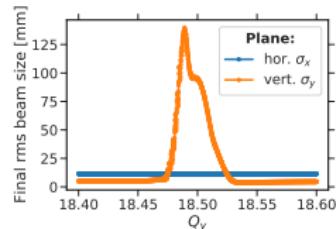


Half-integer without SC but with Chroma

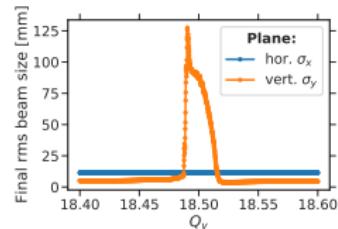
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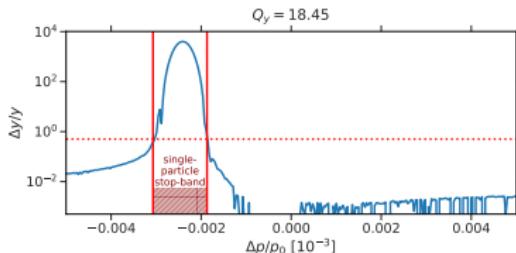
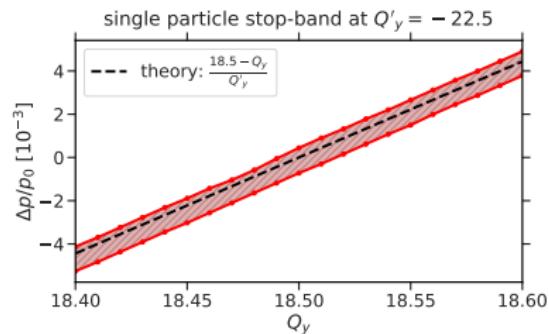
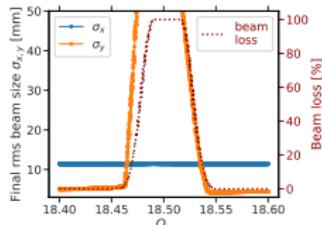


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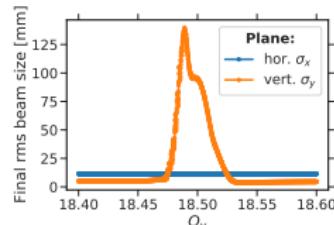


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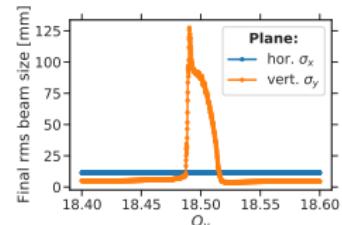
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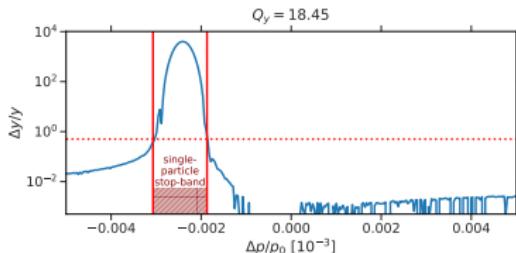
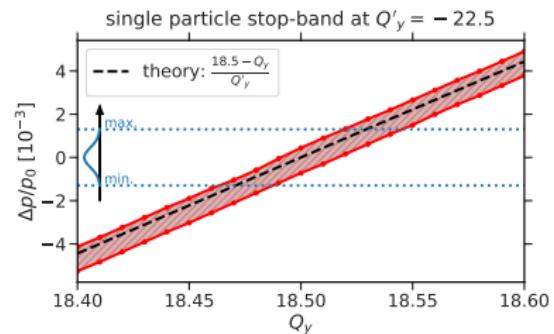
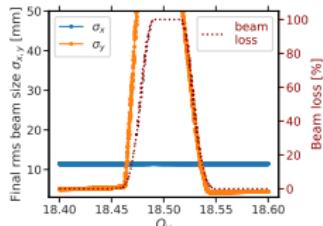


Figure: shifted single-particle stopband

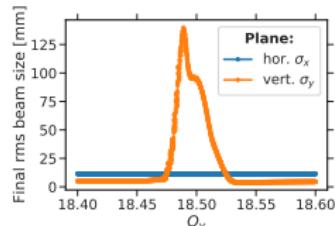


Half-integer without SC but with Chroma

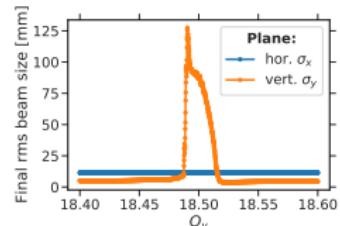
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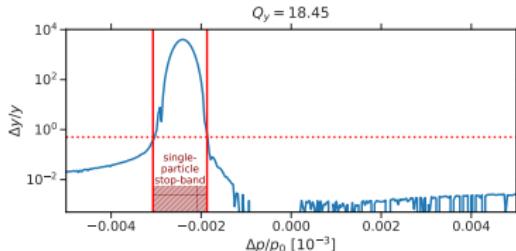


Figure: shifted single-particle stopband

