

# Quarks and gluons in the Lund plane(s)

Gregory Soyez, with Frederic Dreyer, Andrew Lifson, Gavin Salam and Adam Takacs  
based on arXiv:1807.04758, arXiv:2007.06578 and arXiv:2112.09140

IPhT, CNRS, CEA Saclay

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Understanding high-energy collisions requires  
a description of physics across a wide range of scales (from  $\mathcal{O}(\Lambda_{\text{QCD}})$  to  $\mathcal{O}(\text{TeV})$ )

## This talk

- 1 Lund diagrams as a (historical) conceptual tool for parton showers and resummations
- 2 promoting to a practical tool for jet physics
- 3 (Brief) overview of the wide range of applications
- 4 More extensive description of quark/gluon discrimination

## Basic observation

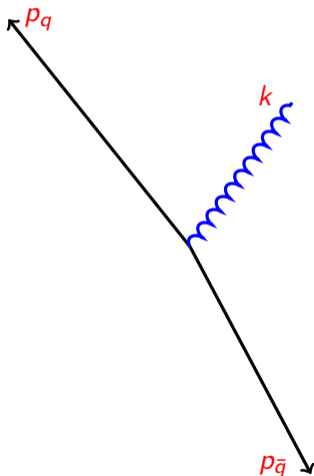
Exploring widely different scales  $\leftrightarrow$  exposing the soft and collinear divergences of QCD  
Obvious connections with parton shower and resummations

## Warmup: Lund diagrams

A useful representation of radiation in a jet

# Basic features of QCD radiations

Take a gluon emission from a  $(q\bar{q})$  dipole



Emission:

$$k^\mu \equiv z_q p_q^\mu + z_{\bar{q}} p_{\bar{q}}^\mu + k_\perp^\mu$$

3 degrees of freedom:

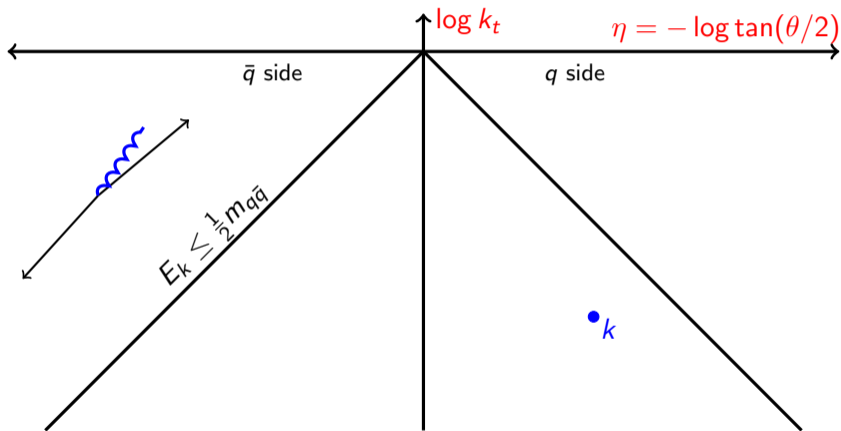
- Rapidity:  $\eta = \frac{1}{2} \log \frac{z_q}{z_{\bar{q}}}$
- Transverse momentum:  $k_\perp$
- Azimuth  $\phi$

In the soft-collinear approximation

$$d\mathcal{P} = \frac{\alpha_s(k_\perp) C_F}{\pi^2} d\eta \frac{dk_\perp}{k_\perp} d\phi$$

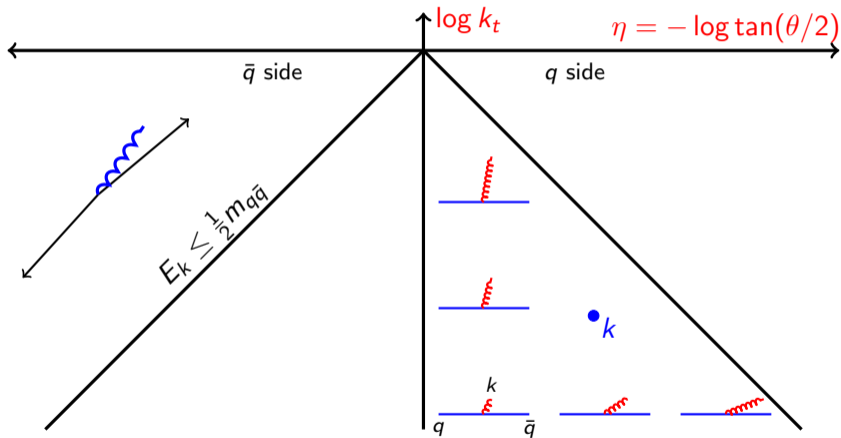
# Basic features of QCD radiations: the Lund plane

Lund plane: natural representation uses the 2 “log” variables  $\eta$  and  $\log k_{\perp}$



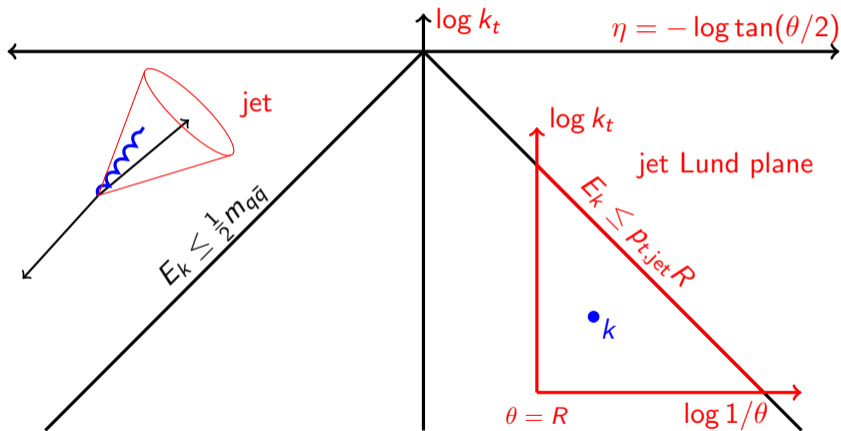
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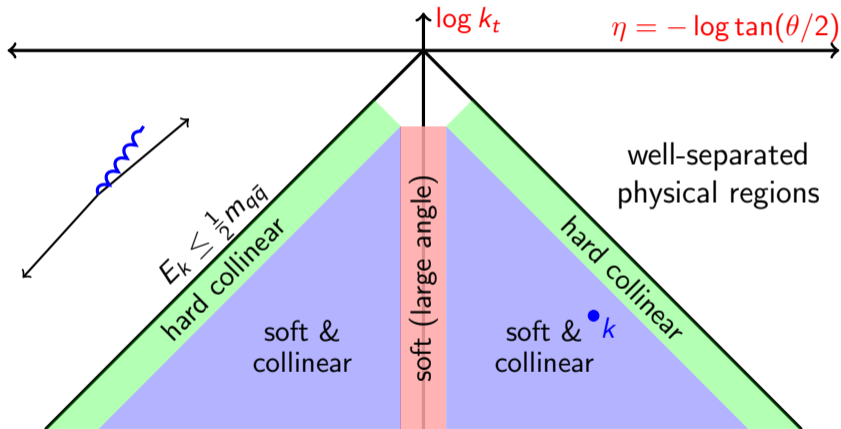
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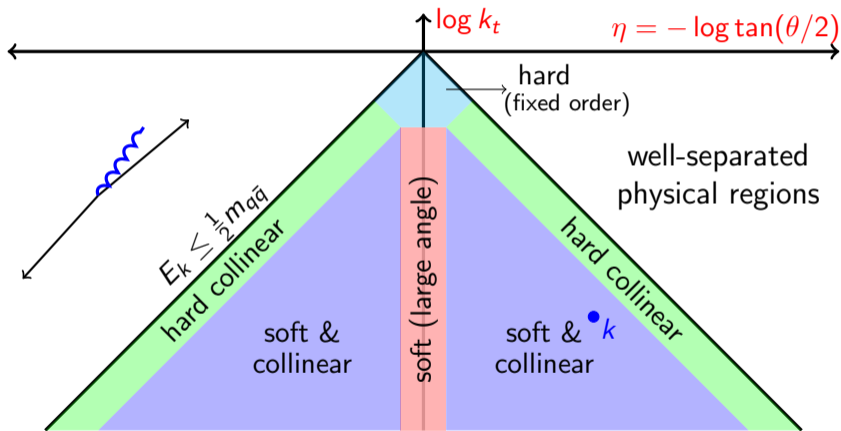
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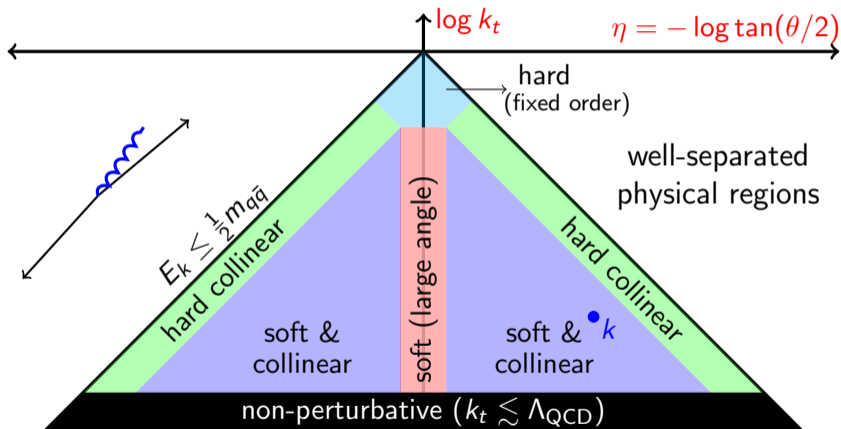
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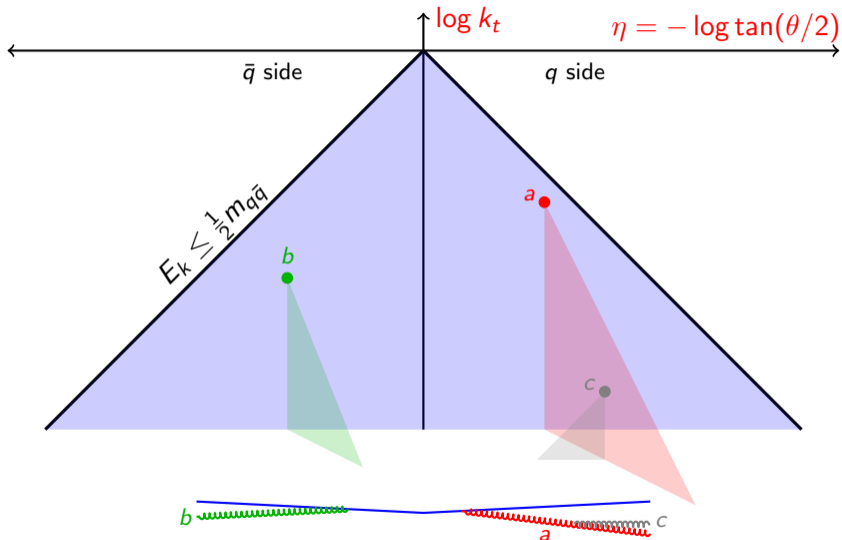


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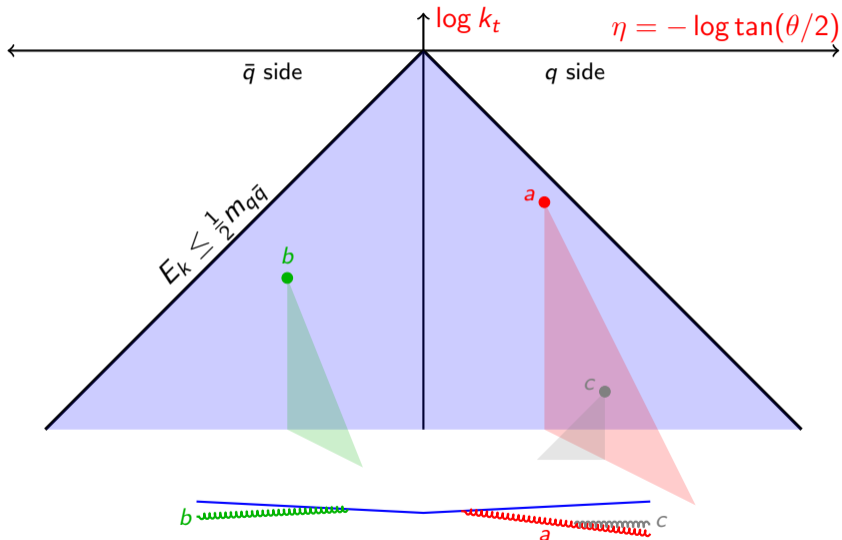


# Multiple emissions in the Lund plane



Each emission spawns  
its own plane  
 $a, b$  primary  
 $c$  secondary  
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**Respects angular  
ordering**  
( $\theta_c < \theta_a$ )

## Lund diagrams represent (multiple) radiation across scales

Set of nice properties:

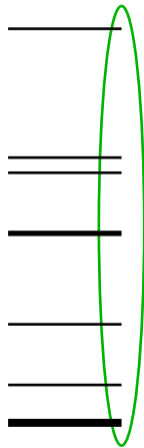
- natural for thinking about resummations and parton showers
- different physical regions (soft, collinear, hard, non-perturbative) well separated
- organised in planes respecting angular ordering

# Lund planes: promoting Lund diagrams to a practical tool

For simplicity, take a high- $p_t$  LHC jet (similar for full  $e^+e^-$  events)

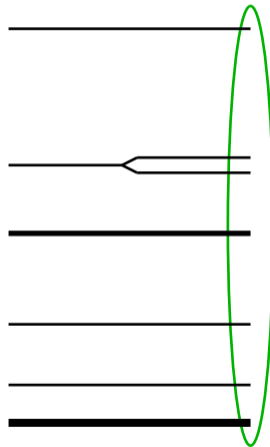
# The Lund plane(s) representation (1/3)

use **Cambridge/Aachen** to iteratively recombine the closest pair



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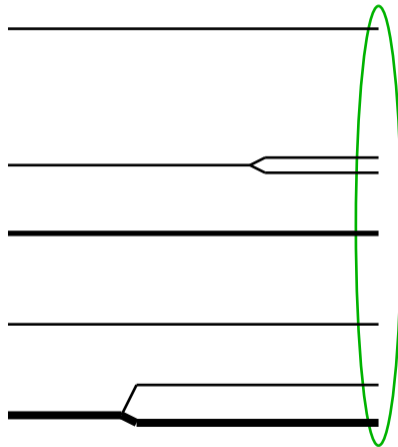
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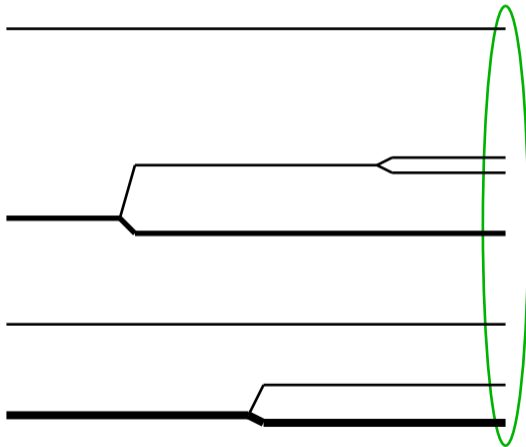
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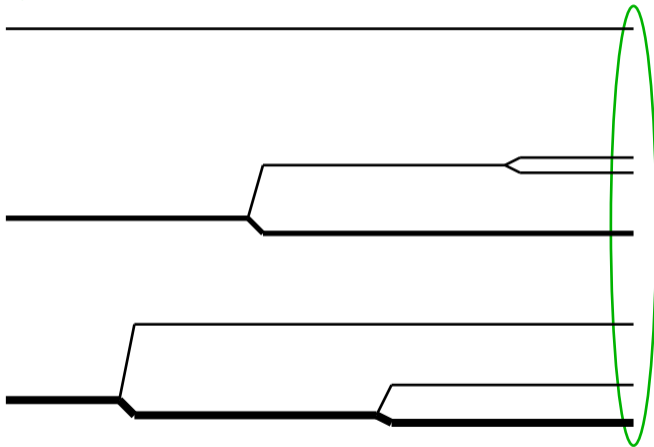
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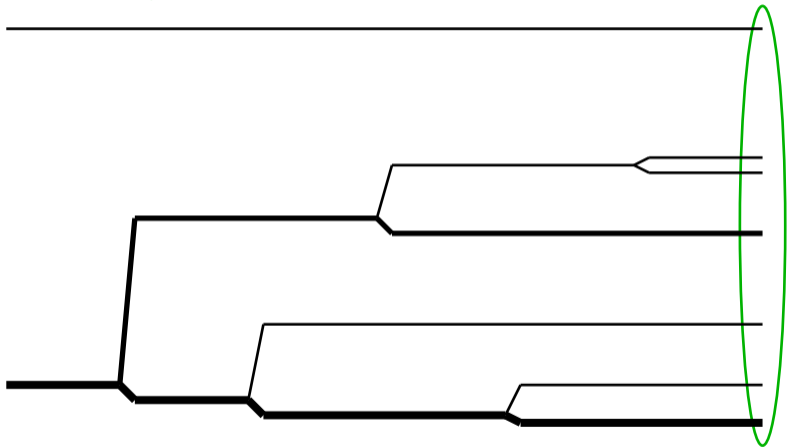
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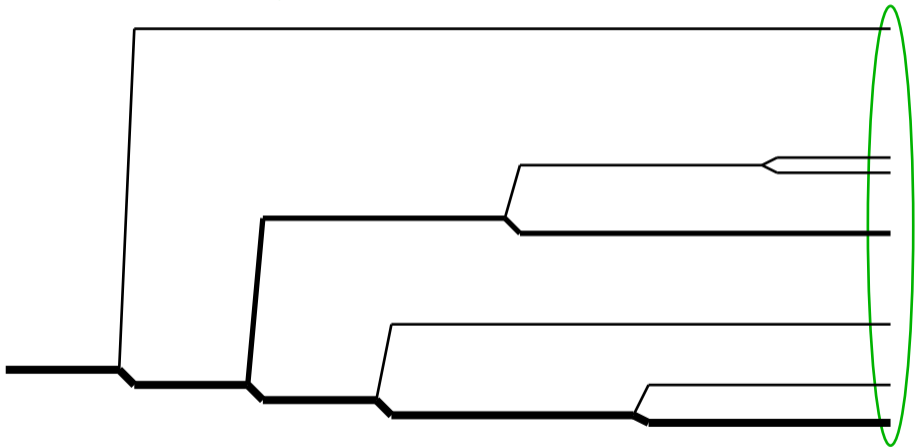
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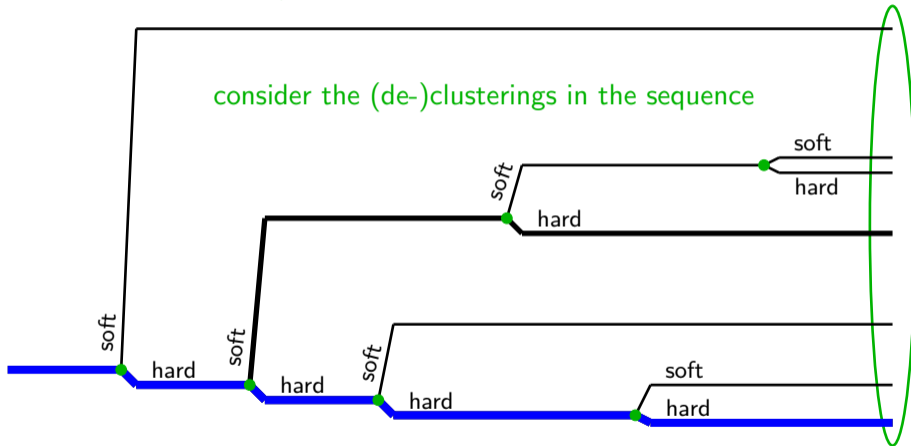
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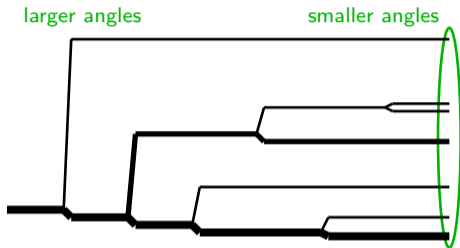
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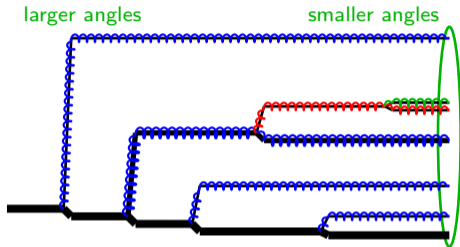
Note: conceptually the largest-energy ( $p_t$  or  $z$ ) branch  $\equiv$  emissions from the “leading parton”

# The Lund plane(s) representation (2/3)



- closely follows our beloved angular ordering

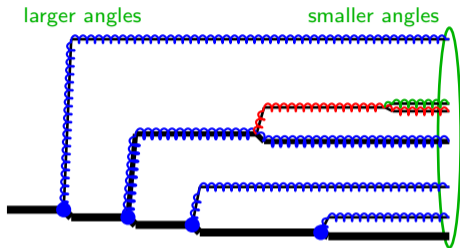
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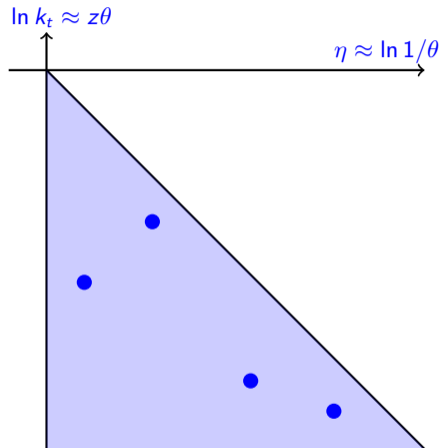
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- i.e. mimics partonic cascade



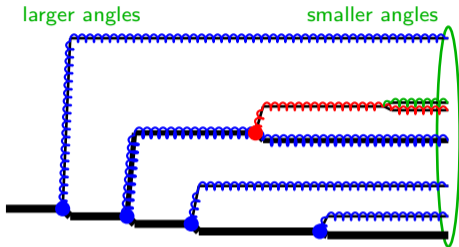
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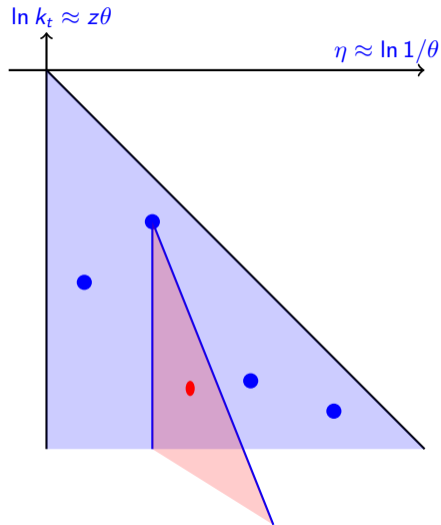
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- can be organised in **Lund planes**
  - primary



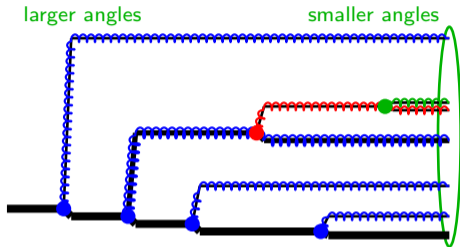
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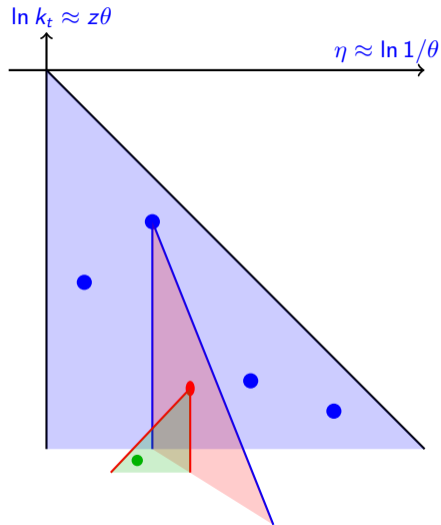
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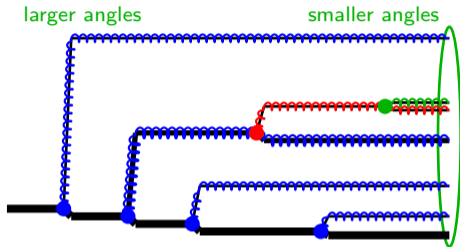
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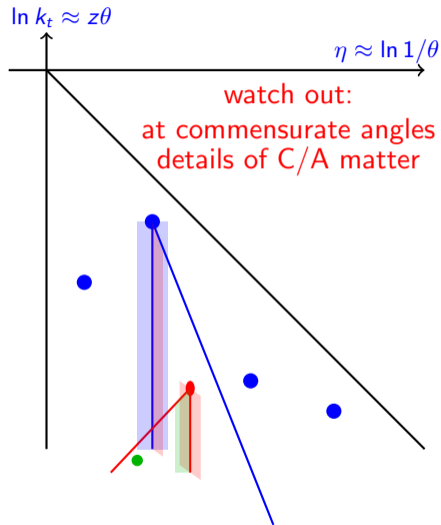
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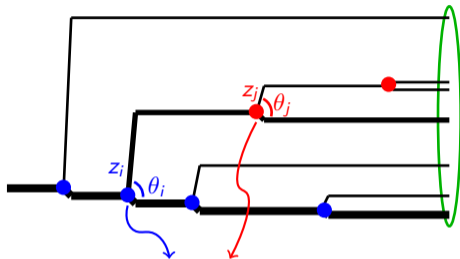
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# The Lund plane(s) representation (3/3)



$$\mathcal{T}_i \equiv \{ \theta_i, k_{t,i}, z_i, \psi_i, m_i, \dots \}$$

for jets in  $pp$ : (similar for  $ee$  events)

$$\eta = -\ln \Delta R$$

$$k_t = p_{t,\text{soft}} \Delta R$$

$\psi \equiv$  azimuthal angle

$$z = \frac{p_{t,\text{soft}}}{p_{t,\text{parent}}}$$

Two different Lund ( $\mathcal{L}$ ) structures

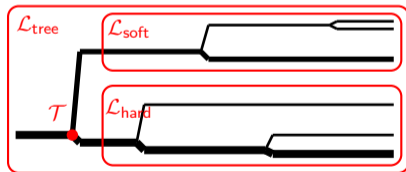
“primary plane”  
(follow hard branch)

OR

full (de-)clustering tree

$$\mathcal{L}_{\text{prim}} \equiv \{ \mathcal{T}_i \}$$

$$\mathcal{L}_{\text{tree}} \equiv \{ \mathcal{T}, \mathcal{L}_{\text{hard}}, \mathcal{L}_{\text{soft}} \}$$



Note: branchings with  $k_t > t_{t,\text{min}} \Rightarrow$  perturbative

## Central observation

For a jet (or a  $ee$  event) one can **construct** a structure that captures the properties of Lund diagrams

The rest of this talk covers several applications:

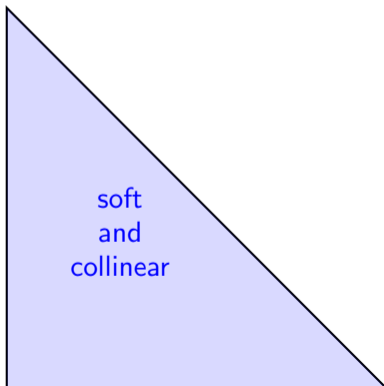
- ✓ Calculations (and measurements)
- ✓ Monte-Carlo developments
- ✓ Tagging (incl. machine learning and **quark/gluon discrimination**)

# Application #1: QCD calculations

# Primary Lund plane multiplicity

Average number of emission at given  $k_t$ ,  $\Delta R$ :

$$\rho = \frac{1}{N_{\text{jets}}} \frac{d^2 N}{d \ln \Delta R d \ln k_t}$$



[A. Lifson, G. Salam, GS, arXiv:2007.06578]

- Double-logarithmic behaviour:

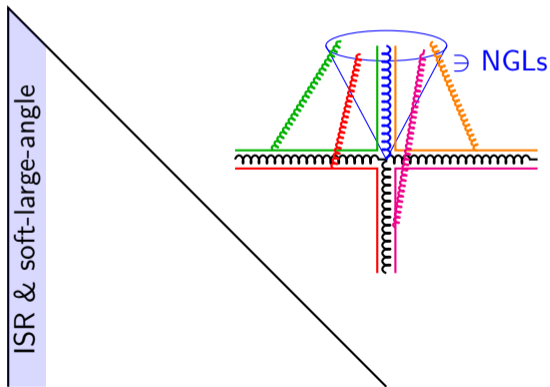
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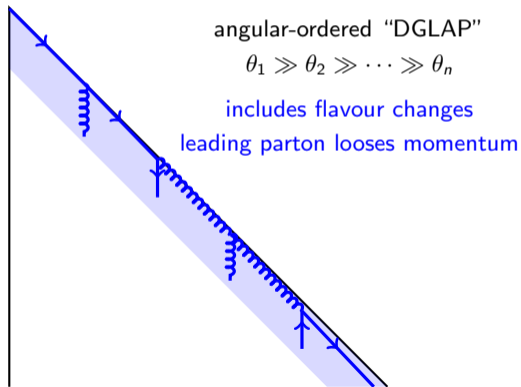
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  - ✓ Running-coupling (trivial)
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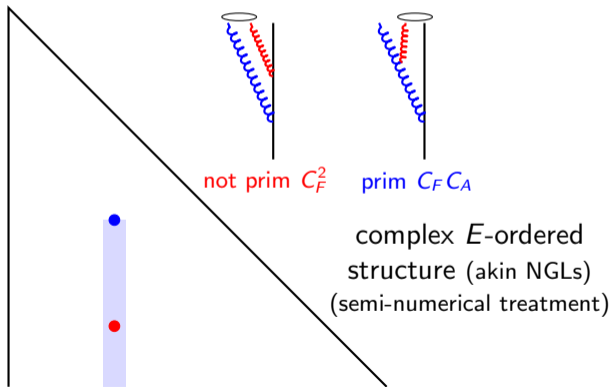
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from NLOJet++  
(some non-trivial details)

2 → 3 at NNLO would  
greatly help!

[S.Abreu,F.Febres Cordero,H.Ita,  
B.Page,V.Sotnikov,2102.13609]

[M.Czakon,A.Mitov,R.Poncelet,2106.05331]

[A. Lifson, G. Salam, GS, arXiv:2007.06578]

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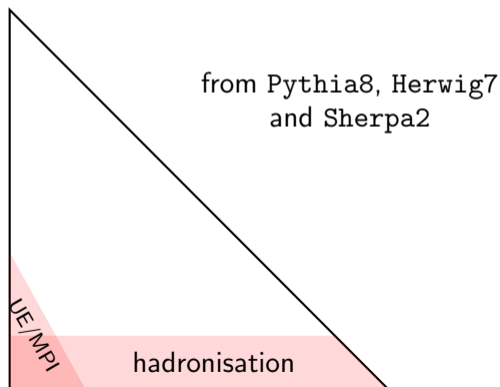
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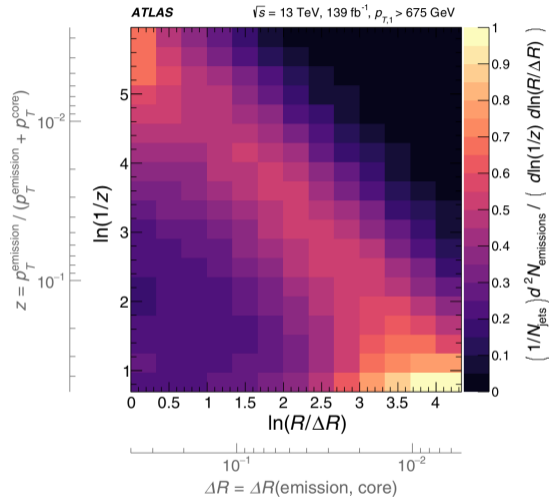


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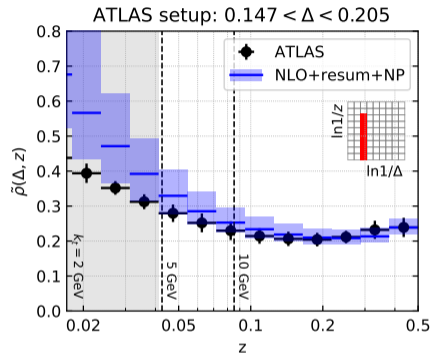
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- + Matching to NLO ( $\sim$  top)
- + NP corrections ( $\sim$  bottom)



[ATLAS, 2004.03540]



- good agreement (particularly for  $k_t \gtrsim 5 \text{ GeV}$ )
- commensurate exp.&th. uncert.
- Can we get  $\alpha_s$  from this?

# Lund multiplicity (1/2)

## Lund multiplicity

count the (average) number of Lund declusterings  
(in the full tree) with  $k_t \geq k_{t,cut}$

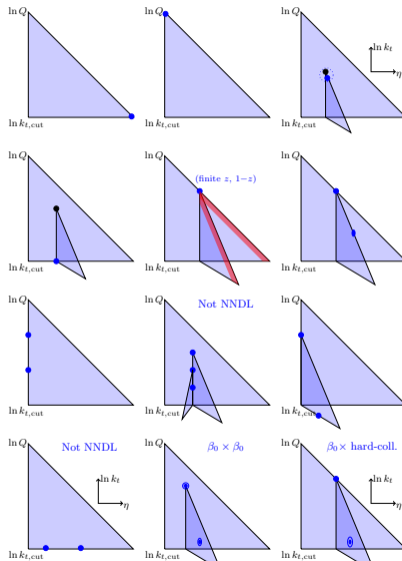
All-order structure ( $L = \ln \frac{Q}{k_{t,cut}}$ ):

$$\langle N^{LP}(L, \alpha_s) \rangle = \underbrace{h_1(\alpha_s L^2)}_{\text{Since 1992}} + \underbrace{\sqrt{\alpha_s} h_2(\alpha_s L^2)}_{\text{New NNDL!!}} + \underbrace{\alpha_s h_3(\alpha_s L^2)}_{\text{New NNDL!!}} + \dots$$

[R. Medves, A. Soto, GS, 2205.02861]

$$\begin{aligned} 2\pi h_3^{(q)} = & D_{end}^{q \rightarrow qg} + (D_{end}^{g \rightarrow gg} + D_{end}^{g \rightarrow q\bar{q}}) \frac{C_F}{C_A} (\cosh \nu - 1) + D_{hmc}^{q\bar{q}g} \cosh \nu + \frac{C_F}{C_A} [(1 - c_s) D_{pair}^{q\bar{q}} (\cosh \nu - 1) + (K + D_{pair}^{gg} + c_s D_{pair}^{q\bar{q}}) \frac{\nu}{2} \sinh \nu] \\ & + C_F \left[ (\cosh \nu - 1 - \frac{1-c_s}{4} \nu^2) D_{clust}^{(prim)} + (\cosh \nu - 1) D_{clust}^{(sec)} \right] + \frac{C_F}{C_A} \left[ D_{e-loss}^g \frac{\nu}{2} \sinh \nu + (D_{e-loss}^g - D_{e-loss}^g) (\cosh \nu - 1) \right] \\ & + \frac{C_F}{C_A} \frac{\pi^2 \beta_0^2}{8C_A} [3\nu(2\nu^2 - 1) \sinh \nu + (\nu^4 + 3\nu^2) \cosh \nu] + \frac{C_F}{2} \{ (B_{gg} + c_s B_{gq})^2 \nu^2 \cosh \nu + 8 [2c_s B_{gg} - 2c_s B_q - (1 - 3c_s^2) B_{gq}] B_{gq} \cosh \nu \\ & + [4B_q(B_{gg} + (2c_s + 1)B_{gq}) - (B_{gg} + c_s B_{gq})(B_{gg} + 9c_s B_{gq})] \nu \sinh \nu + 4(1 - c_s^2) B_{gq}^2 \nu^2 + 8 [2c_s B_q - 2c_s B_{gg} + (1 - 3c_s^2) B_{gq}] B_{gq} \} \\ & + \frac{C_F}{C_A} \frac{\pi \beta_0}{2} \{ (B_{gg} + c_s B_{gq}) \nu^3 \sinh \nu + [2B_q - 2B_{gg} + (6 - 8c_s) B_{gq}] \nu \sinh \nu + 2(B_q + B_{gg} + B_{gq}) \nu^2 \cosh \nu - 4(1 - c_s) B_{gq} (2 \cosh \nu - 2 + \nu^2) \} \end{aligned}$$

Side product: NNDL Cambridge multiplicity for  $y_{cut} = k_{t,cut}^2$



# Lund multiplicity (1/2)

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(in the full tree) with  $k_t \geq k_{t,cut}$

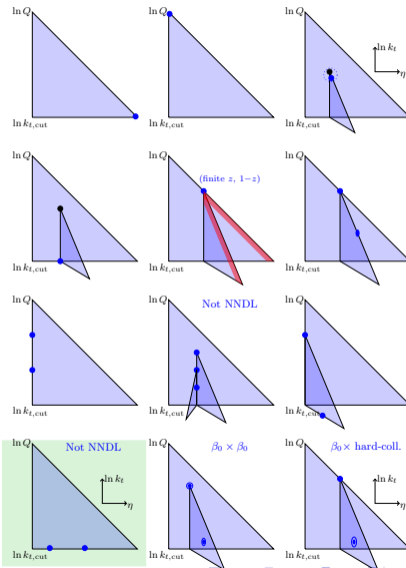
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No “long-distance effect”  $\Rightarrow$  simpler than  $k_t$





# Lund multiplicity (2/2)

[R. Medves, A. Soto, GS,

2205.02861]

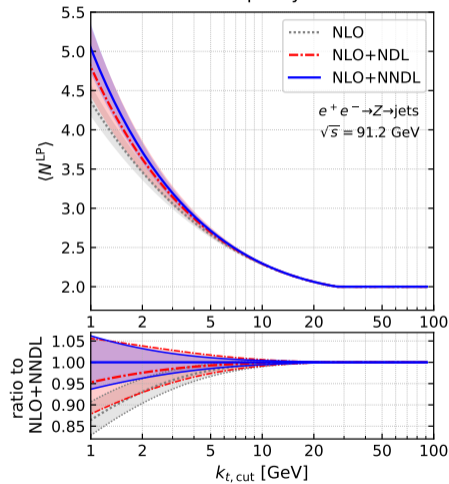
## NNDL Matched to NLO

- Clear effect of resummation
- Clear effect compared to NDL (incl. uncert)

## Several questions

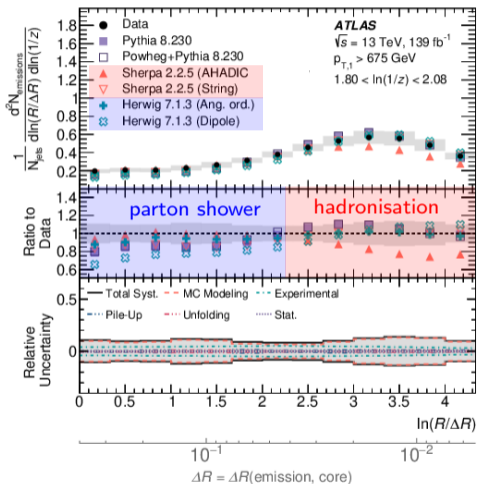
- LEP (ALEPH) measurement?  
see. e.g. Y.Chen *et al.* 2111.09914
- Upgrade to LHC jets?
- Can it lead to an  $\alpha_s$  measurement?
- NNLO? N<sup>3</sup>DL?

Lund multiplicity at LEP



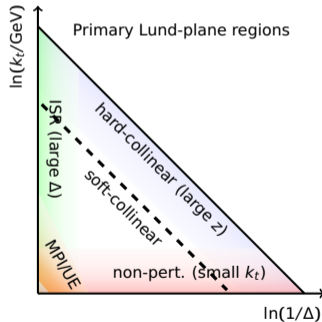
## Application #2: MC development

# Obvious comparisons



“standard” data vs. Monte Carlo comparison

Recall that different Lund regions are sensitive to different physics:



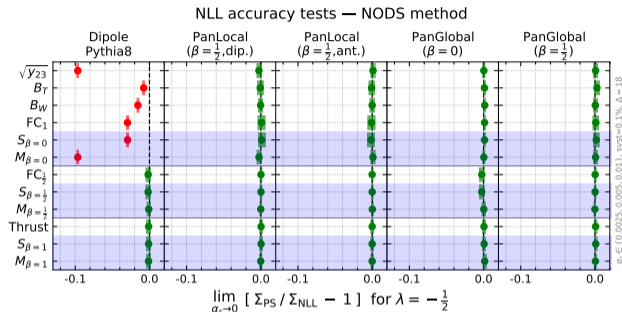
# Revisiting “standard” substructure observables [skip]

- Equivalent to angularities/EECs:

$$S_\beta = \sum_{i \in \mathcal{L}} E_i e^{-\beta \eta_i}$$

$$M_\beta = \max_{i \in \mathcal{L}} E_i e^{-\beta \eta_i}$$

- ✓ *sum* allows for the use of “max”
- ✓ *sum*  $\neq$  *max* at NLL
- ✓ can be defined in *pp*



[M. Dasgupta, F. Dreyer, K. Hamilton, P. Monni, G. Salam, GS, 2002.11114]

[K. Hamilton, R. Medves, G. Salam, L. Scyboz, GS, 2011.10054]

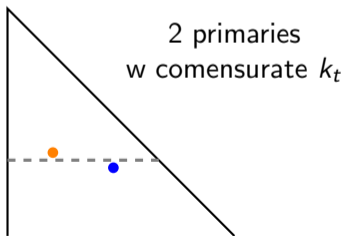
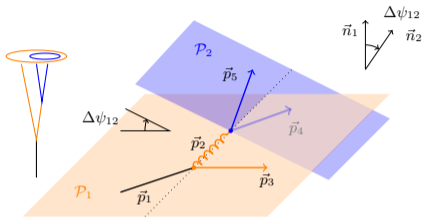
- N*-subjettiness-like: sum excluding the *N* largest

$$\tau_N^{\beta, \text{Lund}} = \sum_{i \in A_N} E_i e^{-\beta \eta_i} \quad \text{with} \quad A_N = \text{argmin}_{X \subset \mathcal{L}, |\mathcal{L} \setminus X| = N-1}$$

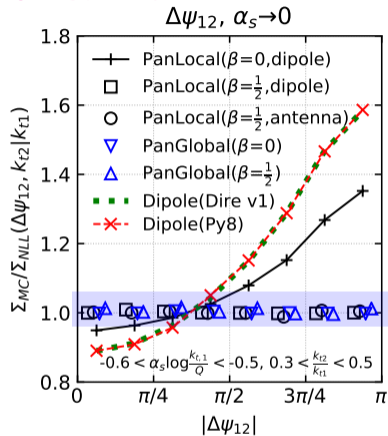
- ✓ Could replace sum by max (likely gaining a simpler resummation structure)
- ✓ Could be defined on the primary plane only

# Crafted observables: example $\Delta\Psi_{12}$

Azimuth between 1<sup>st</sup> and 2<sup>nd</sup> prim. declust.



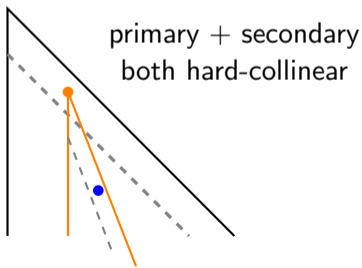
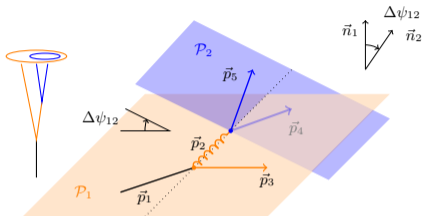
[M.Dasgupta,F.Dreyer,K.Hamilton,P.Monni,G.Salam,GS,2002.11114]



NLL failures for “standard” showers  
“New” PanScales shower OK at NLL

# Crafted observables: example $\Delta\psi_{12}$

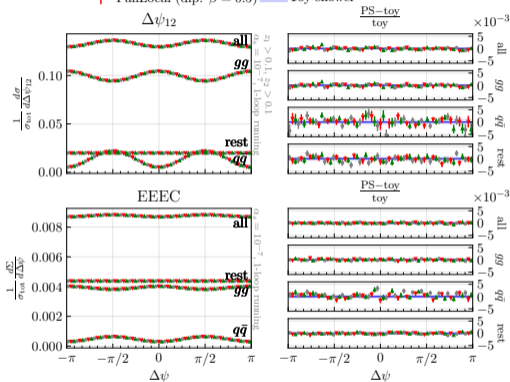
Azimuth between 1<sup>st</sup> and 2<sup>nd</sup> prim. declust.



[A.Karlberg,G.Salam,L.Scyboz,R.Verheyen,2103.16526]

All-order  $\gamma^* \rightarrow q\bar{q}$ ,  $\lambda = -0.5$

↓ PanGlobal ( $\beta = 0$ )     ↓ PanLocal (ant.  $\beta = 0.5$ )  
↓ PanLocal (dip.  $\beta = 0.5$ )     — Toy shower

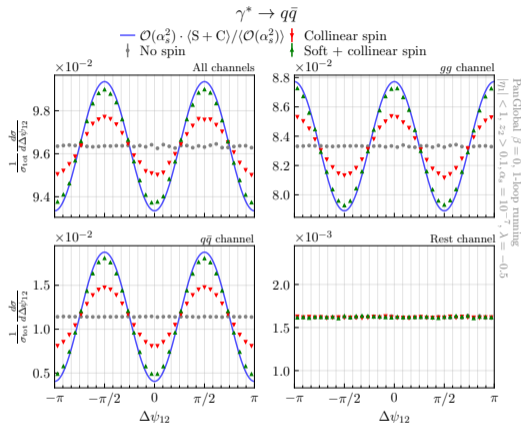
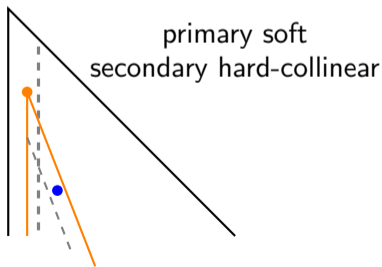
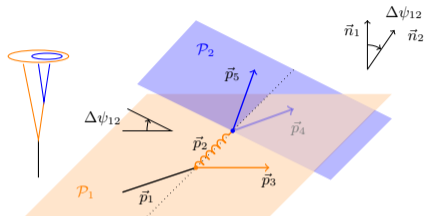


Sensitive to (collinear) spin  
 “New” PanScales shower have spin at NLL  
 agrees w EEEC from 2011.02492 (EEEC less sensitive)

# Crafted observables: example $\Delta\Psi_{12}$

[K.Hamilton,A.Karlberg,G.Salam,L.Scyboz,R.Verheyen,2111.01161]

Azimuth between 1<sup>st</sup> and 2<sup>nd</sup> prim. declust.



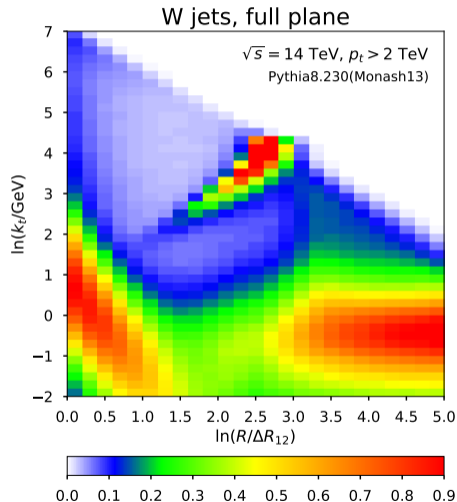
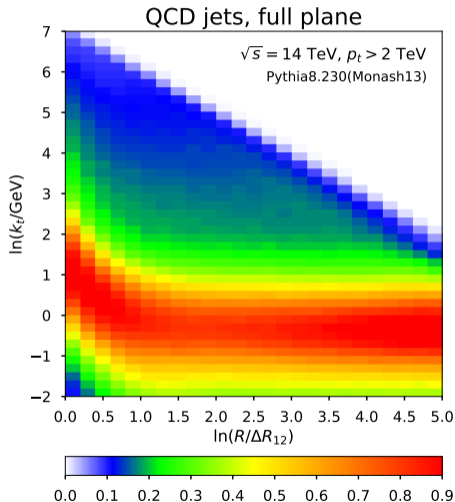
Sensitive to (soft) spin  
“New” PanScales shower have spin at NLL  
first all-order result

# Application #3: Boosted object tagging (mostly illustrative)

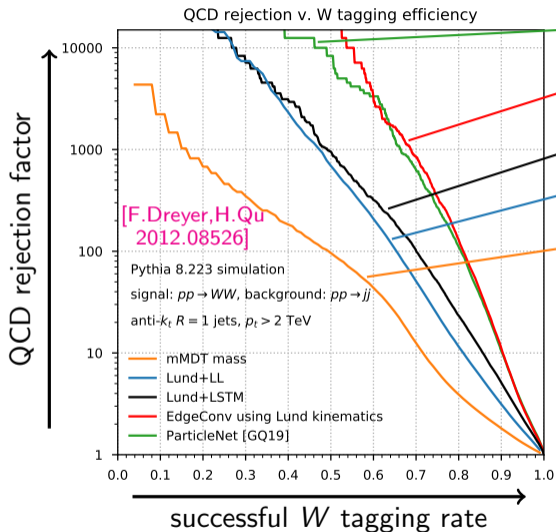


# Tagging boosted $W$ bosons (v. QCD jets) [1/2]

Clear potential on a simple image (also: many basic features recognised)



# Tagging boosted $W$ bosons (v. QCD jets) [2/2]



[graph network using 4-vector (more complex)]

Graph Net trained on full Lund tree

Deep-learning (LSTM) using Lund primaries

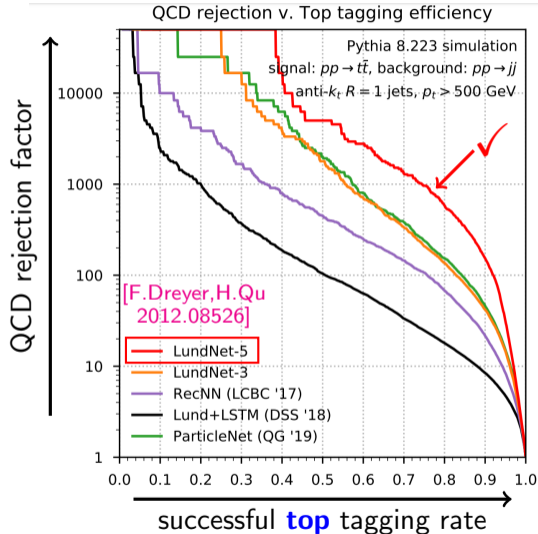
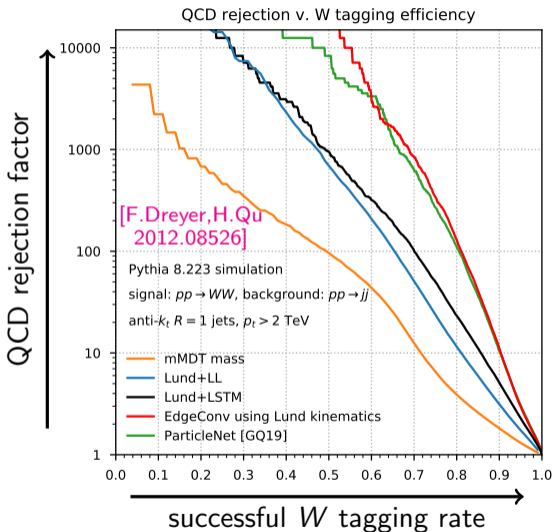
Likelihood ratio based on prim. Lund images

Historical mMDT/SoftDrop

## Main messages

- Large gain from info in the primary plane
- Yet another gain from the full Lund tree
- non-negligible amount of info for  $k_t \lesssim 1$  GeV
- non-negligible differences between generators or parton/hadron level

# Tagging boosted $W$ bosons (v. QCD jets) [2/2]



## Lund plane variables helpful in all areas of jet substructure

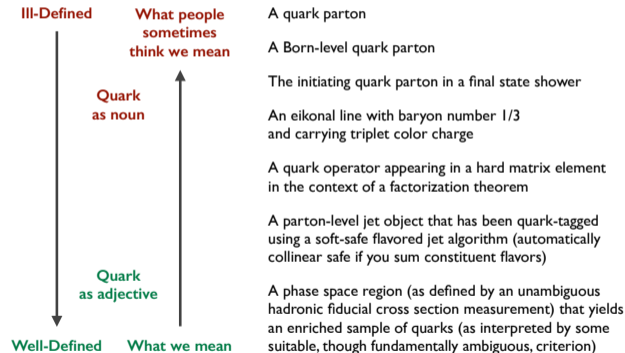
- Variables to test/develop Monte-Carlo generators
- New calculations in (p)QCD and comparisons to data
- Efficient input to Deep-Learning boosted taggers
- Possibilities to craft new observables for a specific purpose  
(Interesting also in heavy-ion collisions)

## Quark/gluon discrimination

**Goal: using the Lund declustering info (primary or full-tree)  
can we say if a jet is quark- or gluon-initiated?**

## What is a Quark Jet?

From lunch/dinner discussions



### pedestrian summary

- there is no such thing as a “quark” or a “gluon” jet
- well-defined: tagging process **A** (“quark-enriched”<sup>(\*)</sup>) against process **B** (“gluon-enriched”<sup>(\*)</sup>)

(\*) ambiguous

### Our approach(es)

- discuss process-independent aspects (at least analytically)
- probe changes for different processes

# Quark v. gluon jets: I. approach

Optimal discriminant (Neyman–Pearson lemma)

$$\mathbb{L}_{\text{prim,tree}} = \frac{p_g(\mathcal{L}_{\text{prim,tree}})}{p_q(\mathcal{L}_{\text{prim,tree}})}$$

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Approach #1

Deep-learn  $\mathbb{L}_{\text{prim,tree}}$   
LSTM with  $\mathcal{L}_{\text{prim}}$  or Lund-Net with  $\mathcal{L}_{\text{tree}}$



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LSTM with  $\mathcal{L}_{\text{prim}}$  or Lund-Net with  $\mathcal{L}_{\text{tree}}$

Approach #2

Use pQCD to calculate  $p_{q,g}(\mathcal{L}_{\text{prim,tree}})$

- Consider  $k_t \geq k_{t,\text{cut}}$  to stay perturbative
- Resum logs to all orders in  $\alpha_s$ , up to **double** logs
  - ▶ Each primary radiation comes with a factor  $\frac{2\alpha_s(k_t)C_R}{\pi}$
  - ▶ Each subsidiary radiation comes with a factor  $\frac{2\alpha_s(k_t)C_A}{\pi}$
- Probabilities:  $p_{q,g} = \prod_{i \in \text{prim}} \frac{2\alpha_s(k_{ti})C_{F,A}}{\pi} \prod_{i \in \text{others}} \frac{2\alpha_s(k_{ti})C_A}{\pi}$  (up to a negligible Sudakov)
- The ratio largely cancels:  $\mathbb{L}_{\text{prim,tree}} = \left(\frac{C_F}{C_A}\right)^{n_{\text{prim}}}$  [C.Frye,A.Larkoski,J.Thaler,1704.06266]
- **The optimal discriminant is the primary multiplicity i.e. the Iterated SoftDrop multiplicity**

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Optimal discriminant (Neyman–Pearson lemma)

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Use pQCD to calculate  $p_{q,g}(\mathcal{L}_{\text{prim,tree}})$

- Consider  $k_t \geq k_{t,\text{cut}}$  to stay perturbative
- Resum logs to all orders in  $\alpha_s$ , up to **single** logs
  - ▶ **single logs** from “DGLAP” collinear splittings

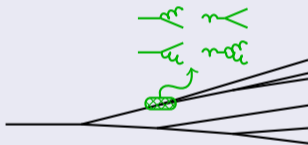
$$P_q(\mathcal{L}_{\text{parent}}) = S_q(\Delta_{\text{prev}}, \Delta) \left[ \tilde{P}_{qq}(z) p_q(\mathcal{L}_{\text{hard}}) p_g(\mathcal{L}_{\text{soft}}) + \tilde{P}_{gq}(z) p_g(\mathcal{L}_{\text{hard}}) p_q(\mathcal{L}_{\text{soft}}) \right]$$

$$p_g(\mathcal{L}_{\text{parent}}) = S_g(\Delta_{\text{prev}}, \Delta) \left[ \tilde{P}_{gg}(z) p_g(\mathcal{L}_{\text{hard}}) p_g(\mathcal{L}_{\text{soft}}) + \tilde{P}_{qg}(z) p_q(\mathcal{L}_{\text{hard}}) p_q(\mathcal{L}_{\text{soft}}) \right]$$

- ▶ some **single logs** for emissions at commensurate angles

Note: all-order not tractable analytically; we resum any *pair* of commensurate-angle emissions

- ▶ **running coupling** (in the Sudakov)

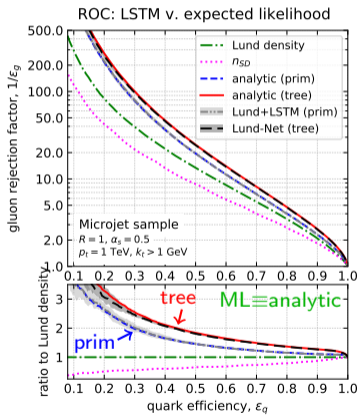


# Quark v. gluon jets: II. ML validation

our analytic discriminant is exact/optimal in the dominant collinear limit  $\theta_1 \gg \theta_2 \gg \dots \gg \theta_n$   
 $\Rightarrow$  ML expected to give the same performance

# Quark v. gluon jets: II. ML validation

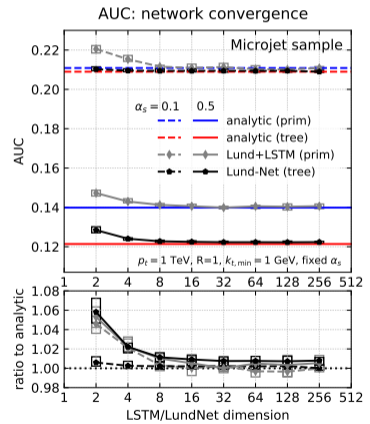
our analytic discriminant is exact/optimal in the dominant collinear limit  $\theta_1 \gg \theta_2 \gg \dots \gg \theta_n$   
 $\Rightarrow$  ML expected to give the same performance



Microjet  
 $\equiv$   
 exact  
 pure-collinear

[M.Dasgupta,F.Dreyer  
 G.P.Salam,G.Soyez,  
 1411.5182]

ROC curves agree

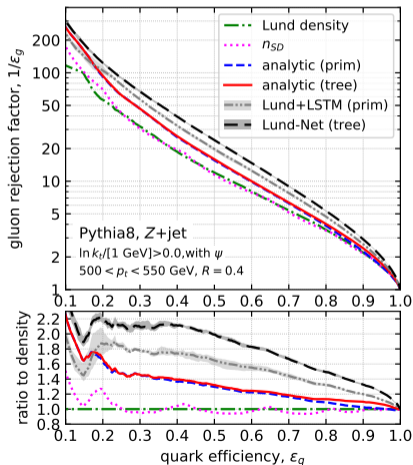


Converges for large-enough networks

# Quark v. gluon jets: III. performance

$pp \rightarrow Zq$  v.  $pp \rightarrow Zg$  ( $p_t \sim 500$  GeV,  $R = 0.4$ )

ROC: Pythia sample

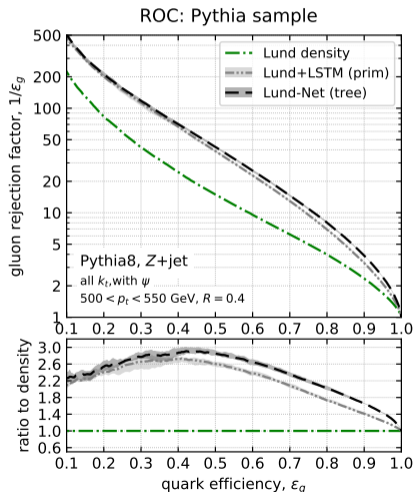


- clear performance ordering:

- 1 Lund+ML > Lund analytic > ISD
- 2 tree > prim

# Quark v. gluon jets: III. performance

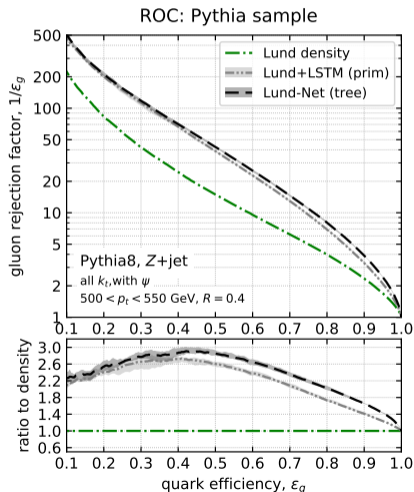
$pp \rightarrow Zq$  v.  $pp \rightarrow Zg$  ( $p_t \sim 500$  GeV,  $R = 0.4$ )



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- larger gains with no  $k_t$  cut

# Quark v. gluon jets: III. performance

$pp \rightarrow Zq$  v.  $pp \rightarrow Zg$  ( $p_t \sim 500$  GeV,  $R = 0.4$ )



- clear performance ordering:
  - 1 Lund+ML > Lund analytic > ISD
  - 2 tree > prim
- larger gains with no  $k_t$  cut
- Interesting questions:
  - ▶ Analytic approach to NP?
  - ▶ Apply analytics to other systems ( $W/Z/H$ , top)

Question: is your tagger resilient to uncontrolled effects?

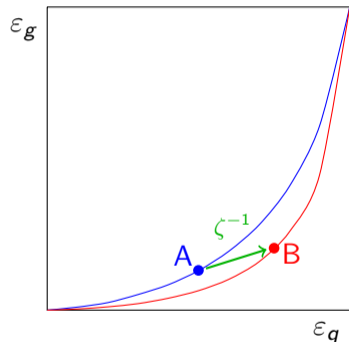
One has:

- a reference sample A  
(e.g. network trained+tested w Pythia)
- an alternate sample B  
(e.g. network tested w Herwig)

We want (for a given working point)

$$\zeta = \left[ \left( \frac{\Delta \varepsilon_q}{\langle \varepsilon_q \rangle} \right)^2 + \left( \frac{\Delta \varepsilon_g}{\langle \varepsilon_g \rangle} \right)^2 \right]^{-1}$$

as small as possible.



(would probably deserve a study on its own)



Question: is your tagger resilient to uncontrolled effects?

One has:

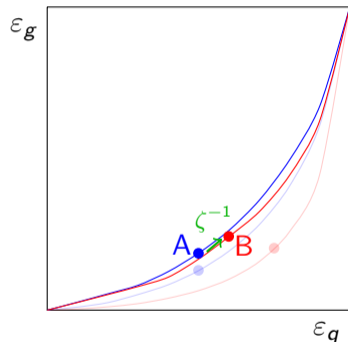
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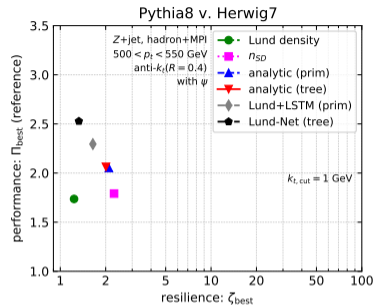
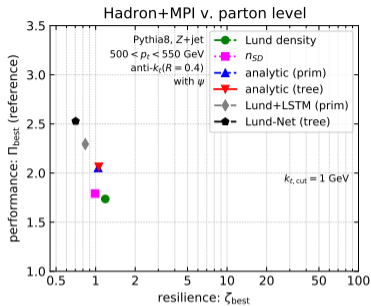
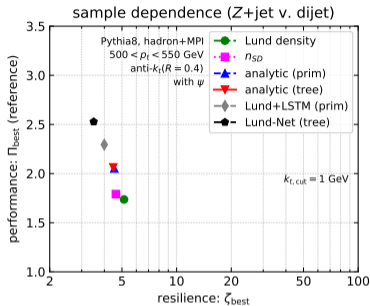
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(would probably deserve a study on its own)



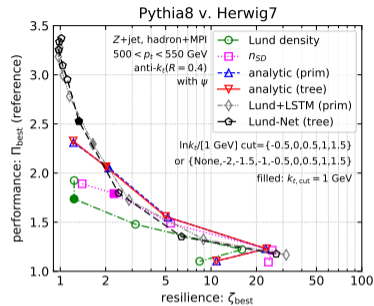
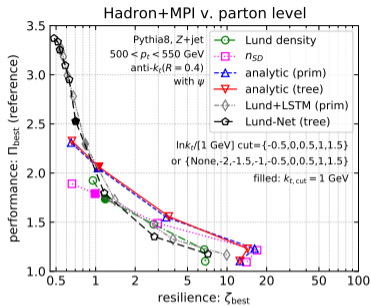
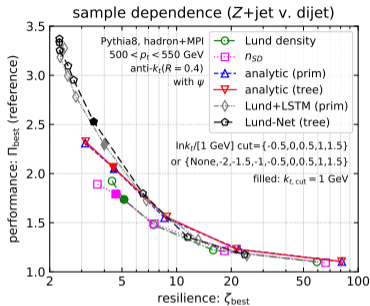
Less performant  
More resilient

# Resilience (2/2)



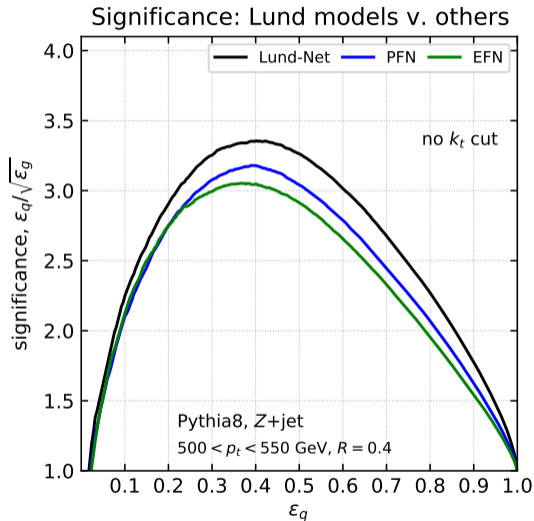
- performance =  $\varepsilon_q / \sqrt{\varepsilon_g}$
- working point:  $k_{t,cut} = 1$  GeV, optimal performance (reference: Pythia, hadron+MPI, Z+jet)
- 3 studies: sample (Z+jet v. dijets), NP effects (hadron v. parton), generator (Pythia v. Herwig)
- performance: same ordering as before
- resilience: network-based < Lund analytics  $\lesssim n_{SD}$

# Resilience (2/2)



- same, varying  $k_{t, \text{cut}}$
- for each curve: “standard” trade-off between performance and resilience
- Overall: better behaviour for the new Lund-based approaches:
  - At “large” resilience: better envelope for the Lund analytic approaches
  - At “small” resilience: ML performance gain pays off

# Comparison to other approaches: ML-based

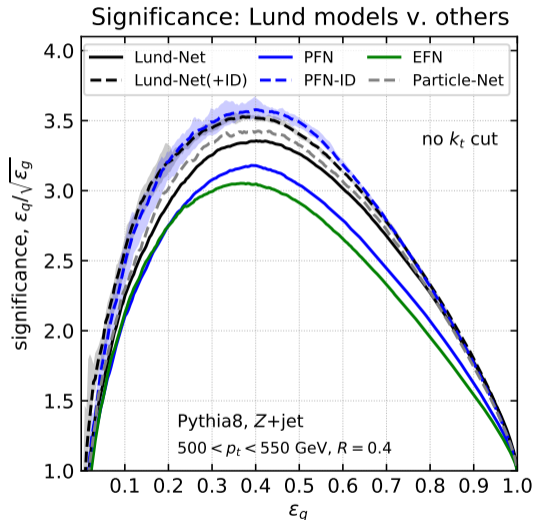


## Approaches:

- Lund-Net (full tree)
- Particle-flow network
- Energy-flow network

- ▶ small performance gain for Lund
- ▶ differences might come from details

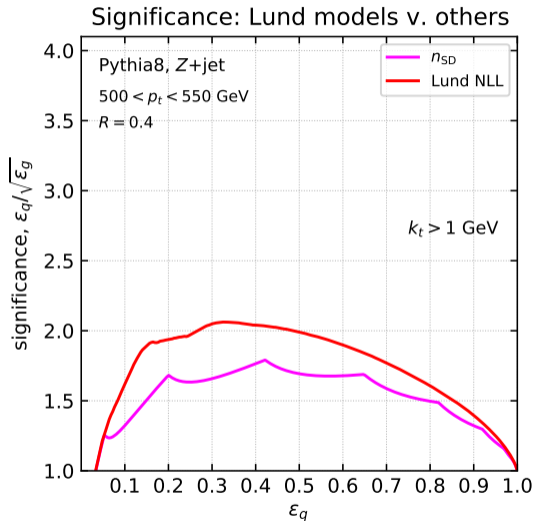
# Comparison to other approaches: ML-based



## Approaches:

- Lund-Net (full tree)
  - Particle-flow network
  - Energy-flow network
  - Dashed: with PDG-ID
  - Particle-Net
- ▶ small performance gain for Lund
- ▶ differences might come from details
- ▶ with PDG-ID:  $\text{PFN} \sim \text{Lund} \gtrsim \text{PNet}$

# Comparison to other approaches: analytics/shapes

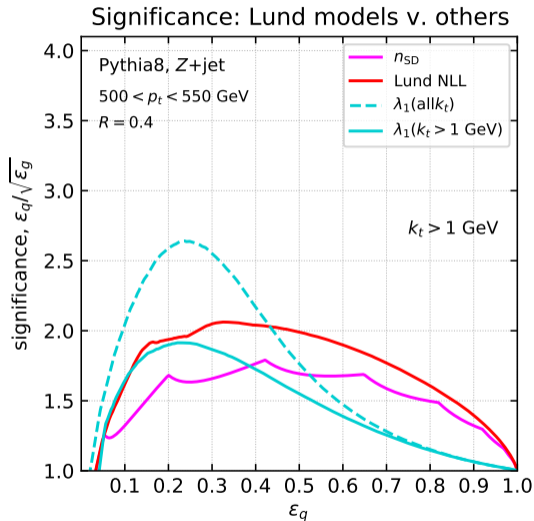


## Approaches:

- ISD mult ( $n_{SD}$ )
- Lund (full tree, analytic)

► clear gain from our analytic approach

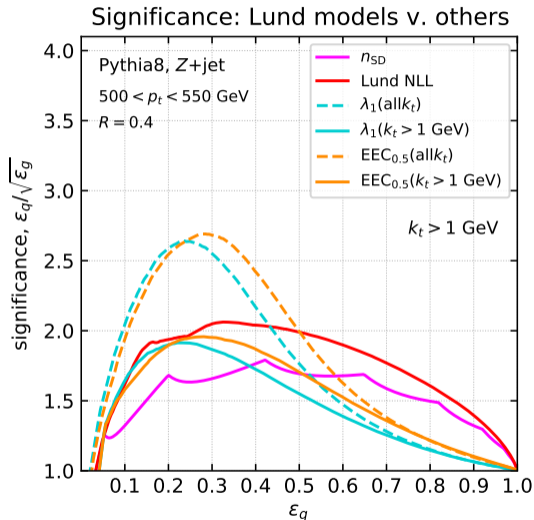
# Comparison to other approaches: analytics/shapes



## Approaches:

- ISD mult ( $n_{SD}$ )
  - Lund (full tree, analytic)
  - width ( $\sum_i p_{ti} \Delta R_i$ )
  - Dashed: use subjets with  $k_t > 1$  GeV
- 
- ▶ clear gain from our analytic approach
  - ▶ Different behaviour for shapes
  - ▶ Lund (expectably) better for same info

# Comparison to other approaches: analytics/shapes



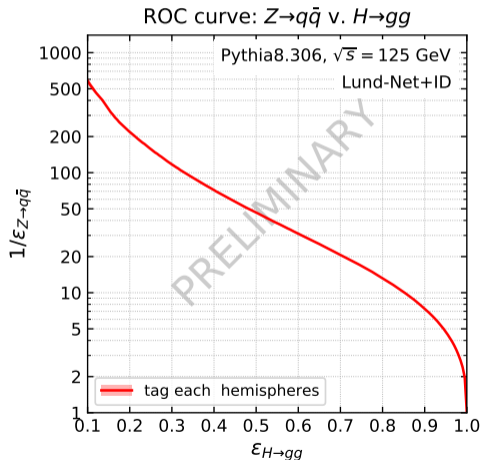
## Approaches:

- ISD mult ( $n_{SD}$ )
- Lund (full tree, analytic)
- width ( $\sum_i p_{ti} \Delta R_i$ )
- EE correlation ( $\sum_{i,j} p_{ti} p_{tj} \Delta R_{ij}^\beta$ )
- Dashed: use subjets with  $k_t > 1$  GeV

- ▶ clear gain from our analytic approach
- ▶ Different behaviour for shapes
- ▶ Lund (expectably) better for same info



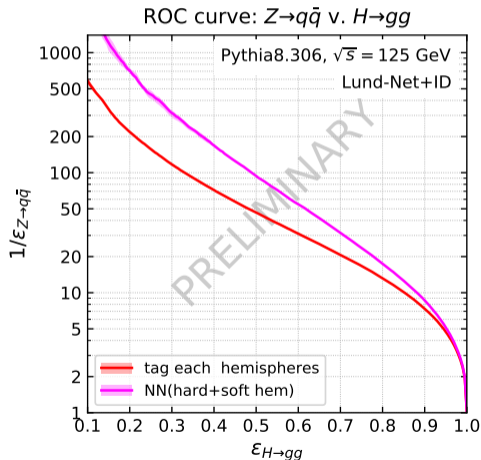
$$e^+e^- \rightarrow Z \rightarrow q\bar{q} \text{ v. } e^+e^- \rightarrow H \rightarrow gg \quad (\sqrt{s} = 125 \text{ GeV, no ISR})$$



observed performance:

- tagging both hemispheres  
i.e. both jets should be tagged
- full event clearly worse than  $(\text{jet})^2$

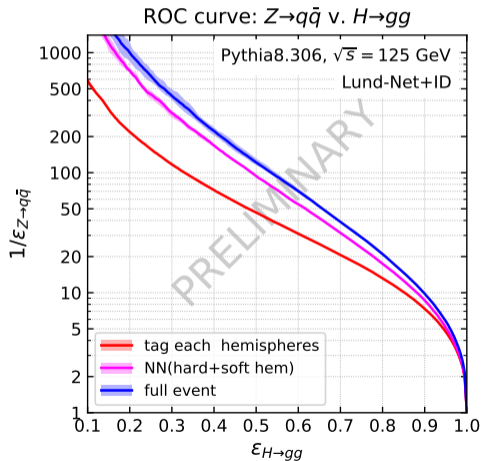
$$e^+e^- \rightarrow Z \rightarrow q\bar{q} \text{ v. } e^+e^- \rightarrow H \rightarrow gg \quad (\sqrt{s} = 125 \text{ GeV, no ISR})$$



observed performance:

- tagging both hemispheres
  - double Lund-Net tag
- train separately on hard & soft hemispheres  
use another NN (or MVA) to combine the two
- clear performance gain

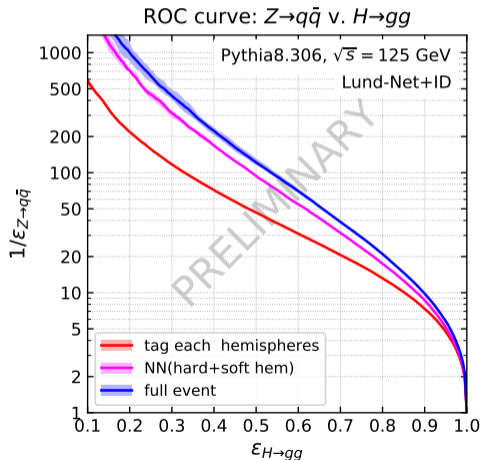
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observed performance:

- tag both hemispheres
  - double Lund-Net tag
  - Lund-Net for the full event
- Another performance gain

$$e^+e^- \rightarrow Z \rightarrow q\bar{q} \text{ v. } e^+e^- \rightarrow H \rightarrow gg \quad (\sqrt{s} = 125 \text{ GeV, no ISR})$$



observed performance:

- tagging both hemispheres
  - double Lund-Net tag
  - Lund-Net for the full event
- Another performance gain

Open questions/work in progress

- How does the analytic do?  
e.g. what gain from full-event tagging?
- Applications to other cases (e.g. at the LHC)?

- ① Lund diagrams have helped thinking about resummation and MCs  
Now they can be reconstructed in practice
  - They provide a view of a jet/event which mimics angular ordering
  - They provide a separation between different physical effects
- ② Broad spectrum of applications:
  - Wide range of possible (p)QCD calculations  
Main limitation: (non-global) clustering logs; can we apply grooming-like techniques?
  - Large scope for crafting new observables ((p)QCD calculations, MC devel/validation)
  - More connections to deep learning, heavy-ion collisions, ...
- ③ Quark-gluon tagging:
  - analytic: single-log gives a systematic improvement over ISD multiplicity
  - deep-learning: Lund-Net shows very good performance (also for  $W$  and top tagging)

# Backup

## Construct the Lund tree in practice: use the Cambridge(/Aachen) algorithm

Main idea: Cambridge(/Aachen) preserves angular ordering

### $e^+e^-$ collisions

① Cluster with Cambridge ( $d_{ij} = 2(1 - \cos \theta_{ij})$ )

② For each (de)-clustering  $j \leftarrow j_1 j_2$ :

$$\eta = -\ln \theta_{12}/2$$

$$k_t = \min(E_1, E_2) \sin \theta_{12}$$

$$z = \frac{\min(E_1, E_2)}{E_1 + E_2}$$

$$\psi \equiv \text{some azimuth, ...}$$

### Jet in $pp$

① Cluster with Cambridge/Aachen ( $d_{ij} = \Delta R_{ij}$ )

② For each (de)-clustering  $j \leftarrow j_1 j_2$ :

$$\eta = -\ln \Delta R_{12}$$

$$k_t = \min(p_{t1}, p_{t2}) \Delta R_{12}$$

$$z = \frac{\min(p_{t1}, p_{t2})}{p_{t1} + p_{t2}}$$

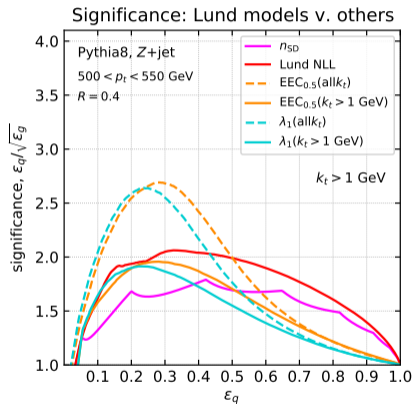
$$\psi \equiv \text{some azimuth, ...}$$

### Primary Lund plane

Starting from the jet, de-cluster following the “hard branch” (largest  $E$  or  $p_t$ )

# Quark v. gluon jets: III. performance v. others

$pp \rightarrow Zq$  v.  $pp \rightarrow Zg$  ( $p_t \sim 500$  GeV,  $R = 0.4$ )

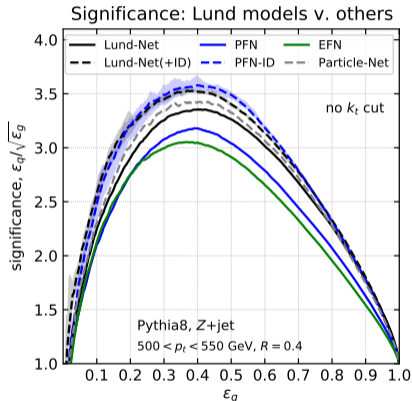


- Analytic approach shows gains for  $k_t > 1$  GeV (shapes improve at small  $\epsilon_q$  by adding smaller  $k_t$ )



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- Analytic approach shows gains for  $k_t > 1$  GeV (shapes improve at small  $\epsilon_q$  by adding smaller  $k_t$ )
- ML performance on par with PFN, slightly better than Particle-Net (treatment of PDG-ID could maybe be improved)