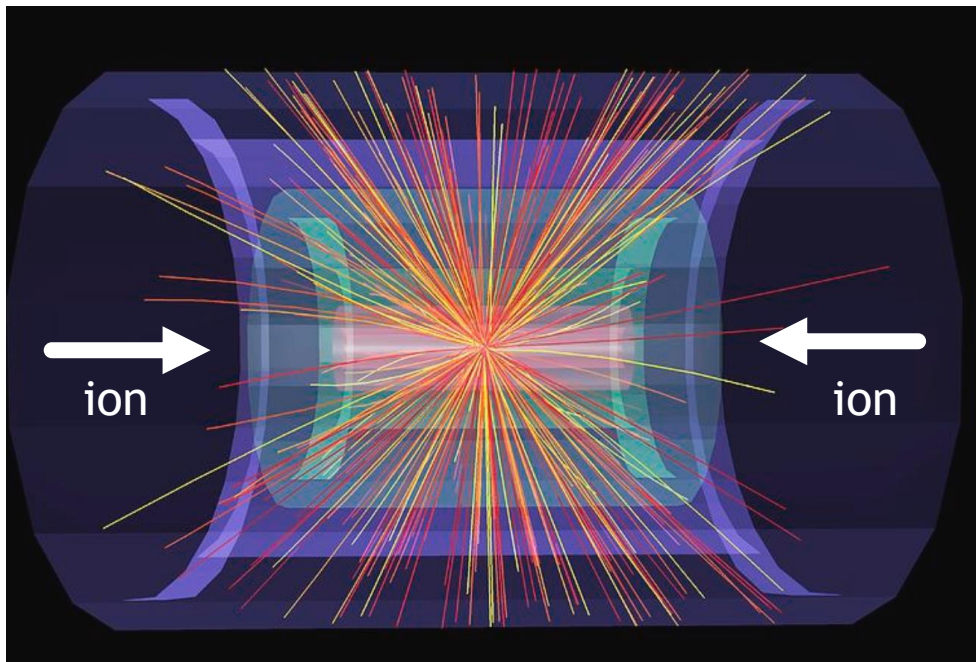


Detectors

TK Hemmick

How do you want to do your physics?



► Outline of the presentation:

- Measuring E with Calorimeters
- Measuring \vec{p} with Tracking
- Measuring PID

- Beam Species; Beam Energy; (E, p_x, p_y, p_z)
 - Seems logical to use real 4-vectors.
 - Close match to textbook calculations
 - Does this mean $(\text{Calorimeter}, \overrightarrow{\text{Tracker}})$ is ideal?
 - Not exactly...
- All measured particles are “on shell”.
 - $E^2 = p^2 + m^2$
 - Past generations have measured on shell masses well.
 - Your precision is enhanced using this knowledge.
 1. Calorimeter-centric $(ID, E\hat{p})$
 2. Spectrometer-centric (ID, \vec{p})
 3. Best = properly weighted average of 1 & 2

Particles must interact

pion →
proton →
muon →
electron →

Infinite Block of Material

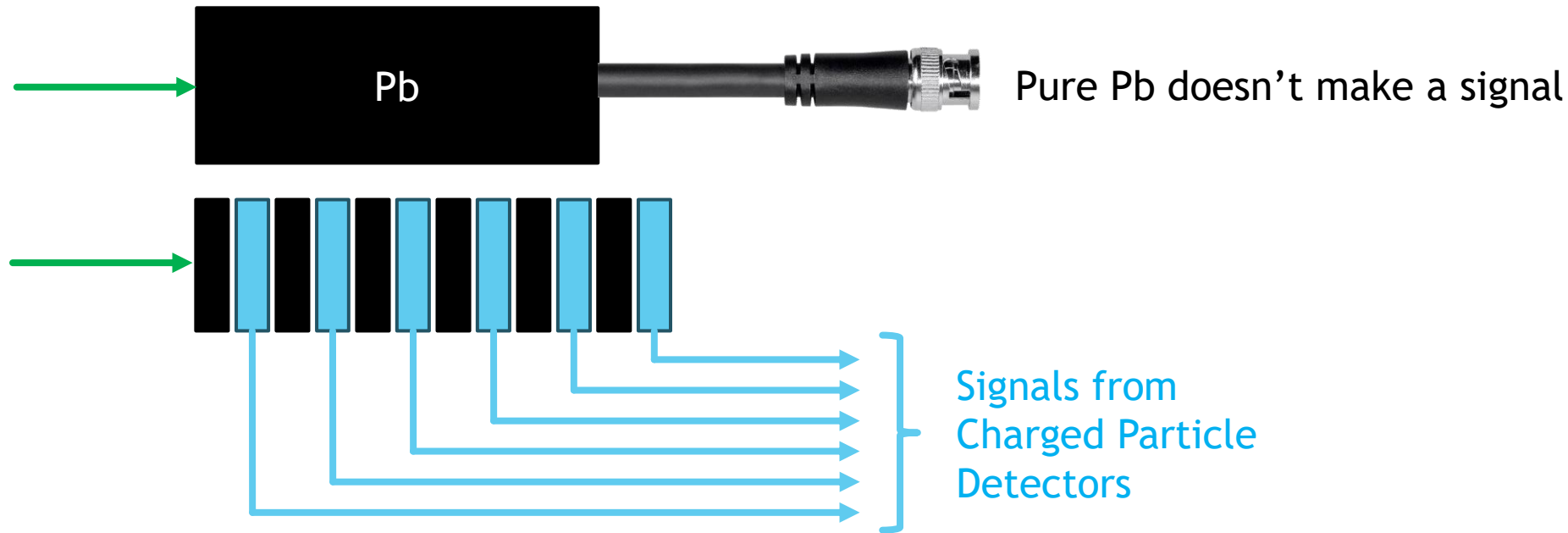
▶ Four fundamental forces: Gravity, Weak, E&M, strong

▶ Participants

▶ Pion	Strong, E&M	}	Medium Range
▶ Proton	Strong, E&M		
▶ Muon	E&M	→	Long Range
▶ Electron	E&M	→	Short Range

- ▶ Bremsstrahlung sets in for the electron.
- ▶ Allows short detector
- ▶ Allows high precision

Instrument a Pb Brick

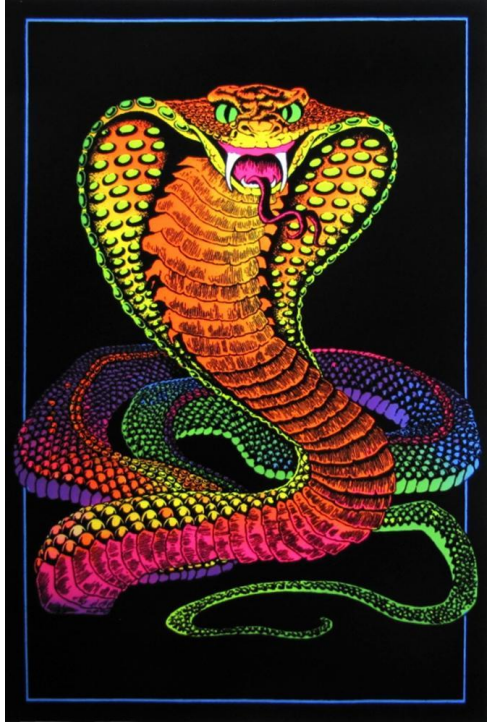


- ▶ A piece of Pb does not produce a signal
 - ▶ CsI scintillates
 - ▶ PbWO_4 scintillates
- ▶ Slice the Pb brick like bread and insert CP detectors
 - ▶ “Sampling Calorimeter”
 - ▶ The big question is how the signal gets out

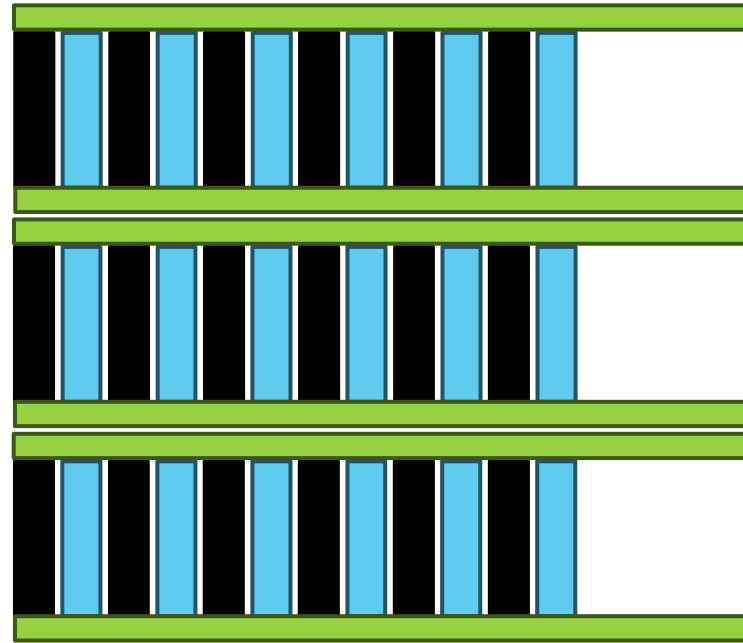
Crystal Calorimeters

- Best Precision
- Highest Cost

Fluorescence is a Cheap Redirect

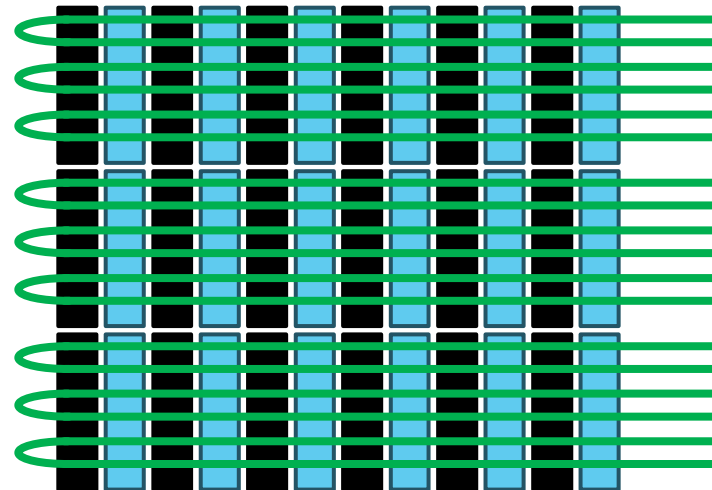


Drug Paraphernalia
from the 1970's



Wavelength Shift Bars

Fluorescent "Wavelength Shifter"

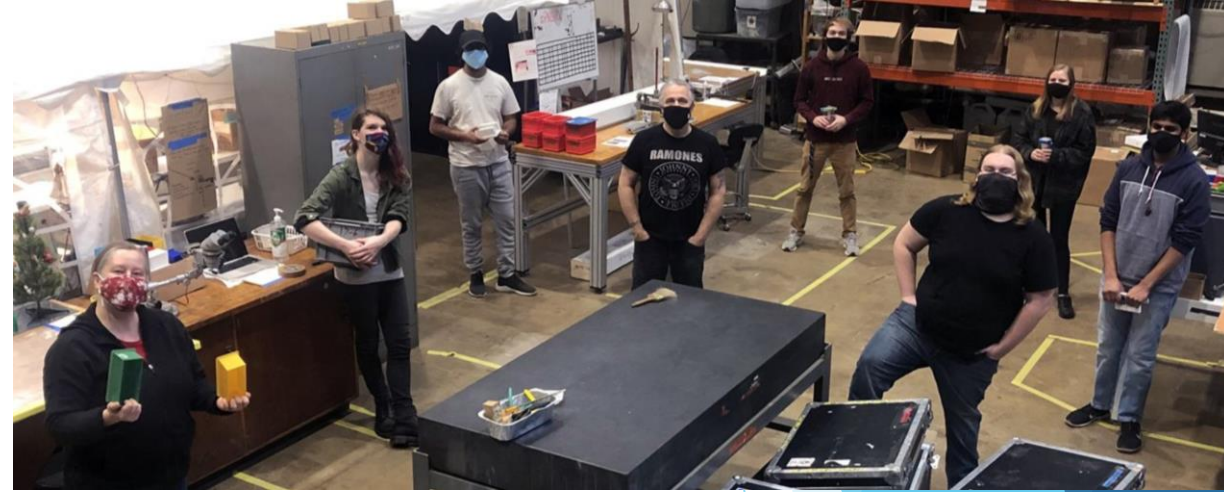
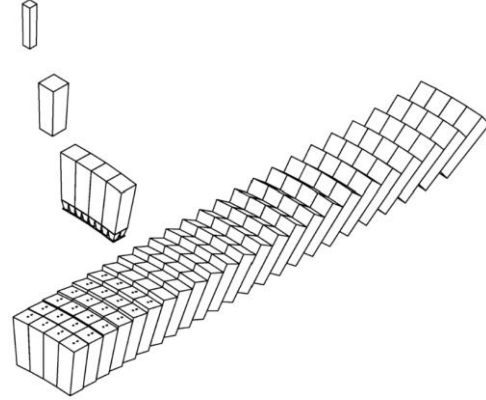


Wavelength Shift Fibers

- ▶ Fluorescence absorbs light at shorter wavelength and re-emits it at a longer wavelength.
- ▶ "Blacklight Posters" used inks containing Fluorescent Dye
- ▶ Plastic can be purchased with Fluorescent Dye in bars and in fibers.
- ▶ Fluorescence "redirects" the light, much of which is totally internally reflected to the back.
- ▶ Shashlik or Shish Kebab is best.

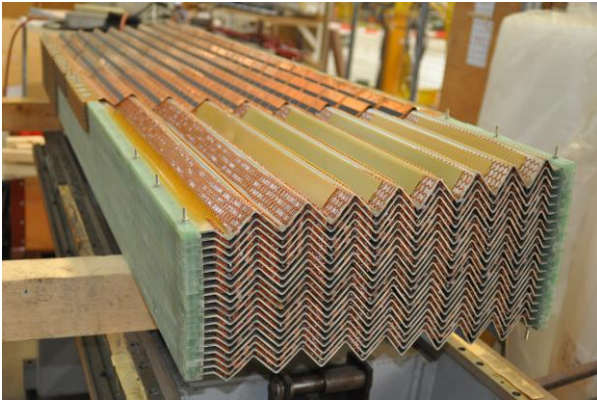
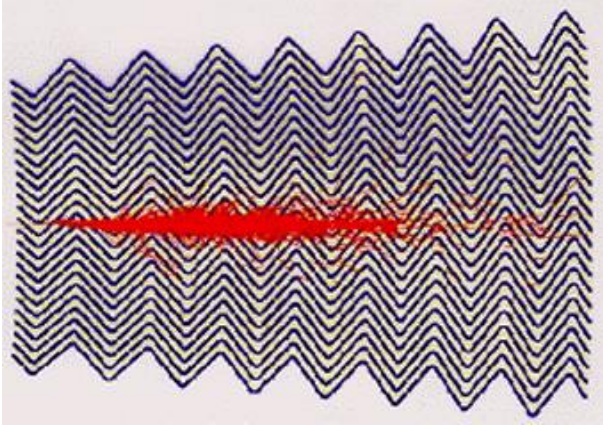


SciFi (aka Scintillating Fibers)

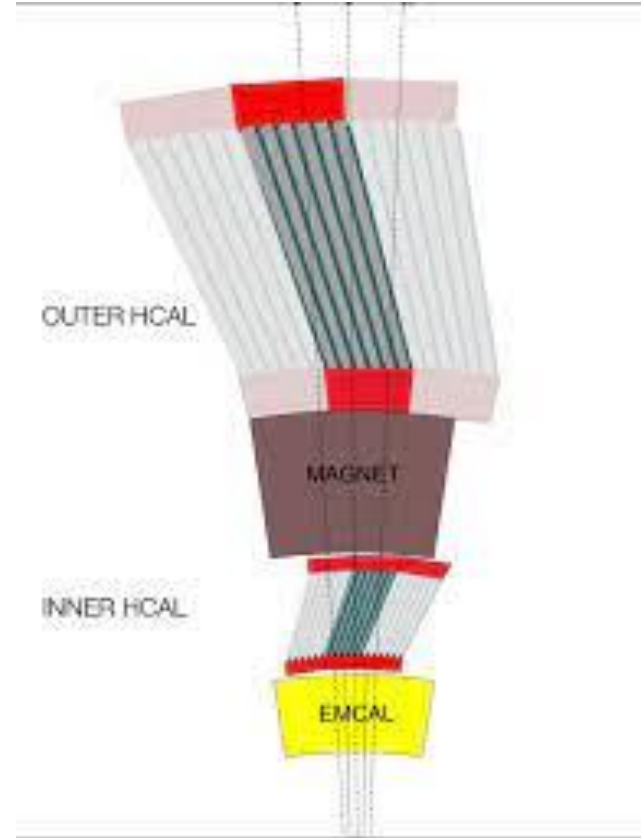


- ▶ Scintillating fibers with close spacing direct light out the back without shifters
- ▶ Here an sPHENIX EMCAL block is shown.
 - ▶ The surrounding material is W powder mixed with a minimal amount of epoxy that is been “injection-molded” around a tight array of scintillating fiber.

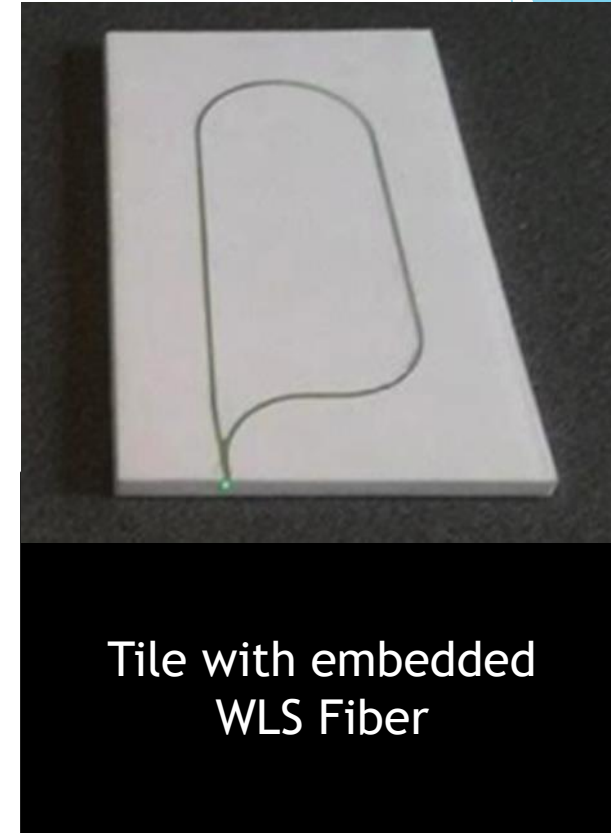
Accordion



ATLAS Liquid Argon



sPHENIX HCAL



Tile with embedded
WLS Fiber

- ▶ ATLAS Liquid Argon Calorimeter pioneered the “accordion” to channel ionization charge from within the Argon along the accordion folds.
- ▶ sPHENIX mimicked this with a simple “tilt” of each tile and concentrated the light using a wavelength shifting fiber.

Calorimeter Energy Resolution

- ▶ Fundamentally the calorimeter “counts” something (charge particles, ionization, photons, ...) that is proportional to the energy deposit.

▶ Statistical Term:

- ▶ $N \propto E$
- ▶ $\partial N = \sqrt{N} \propto \sqrt{E}$
- ▶ $\frac{\partial E}{E} \propto \frac{\partial N}{N} \propto \frac{1}{\sqrt{N}} \propto \frac{1}{\sqrt{E}}$

▶ Scale/Calibration Term:

- ▶ $E = aN$
- ▶ $\partial E = \partial a N = \text{const} * E$
- ▶ $\frac{\partial E}{E} = \text{const}$

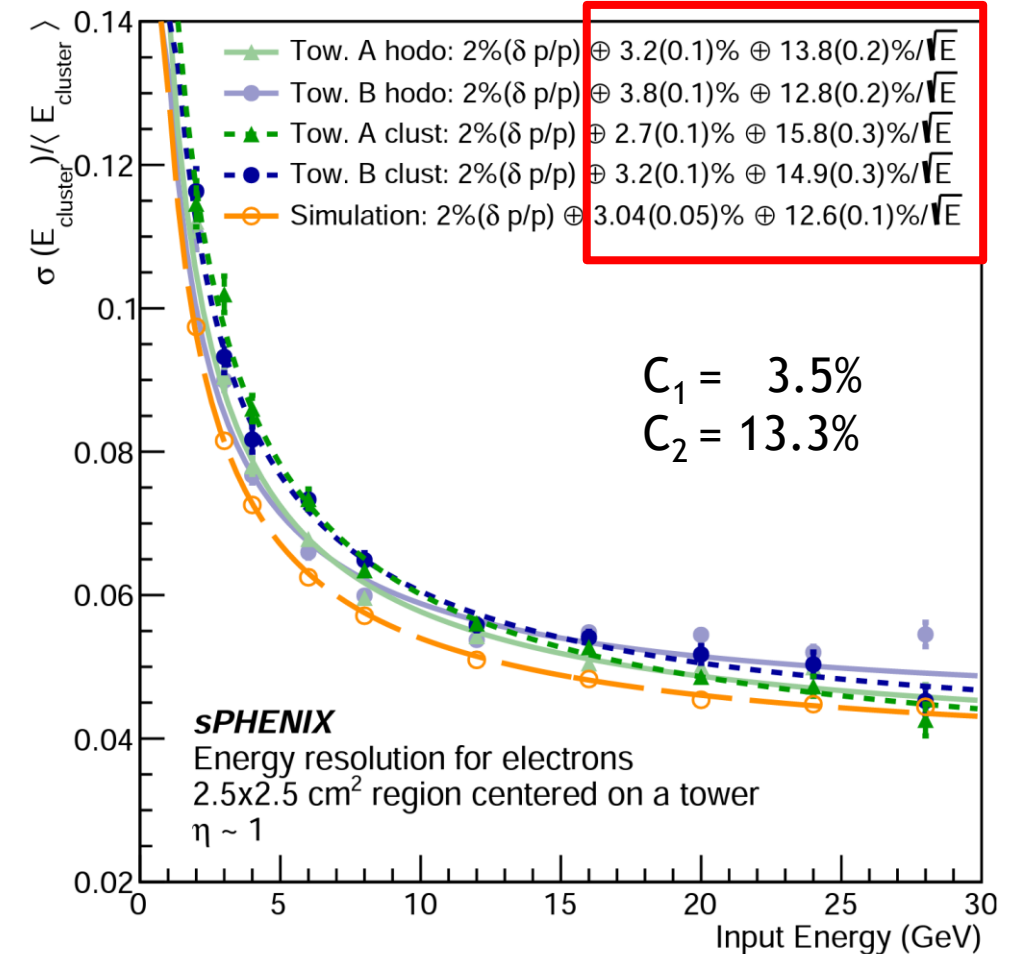
▶ Noise Term:

- ▶ $\partial E = \text{const}$
- ▶ $\frac{\partial E}{E} = \frac{\text{const}}{E}$

▶ Full Resolution

$$\frac{\partial E}{E} = \sqrt{c_1^2 + \left(\frac{c_2}{\sqrt{E}}\right)^2 + \left(\frac{c_3}{E}\right)^2} = c_1 \oplus \frac{c_2}{\sqrt{E}} \oplus \frac{c_3}{E}$$

Often negligible



Calorimeters vs Spectrometers

► Calorimeters

- Measures energy
- Best energy resolution at highest momentum/energy
- DESTRUCTIVE to the particle (cannot measure again)

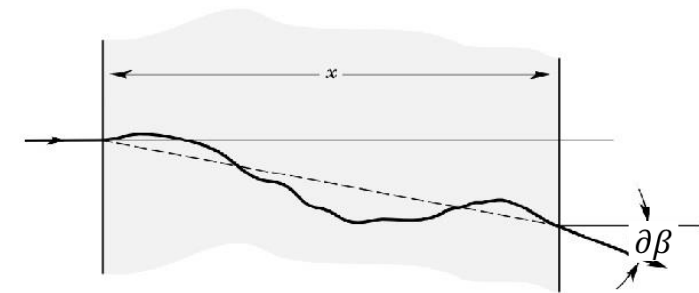
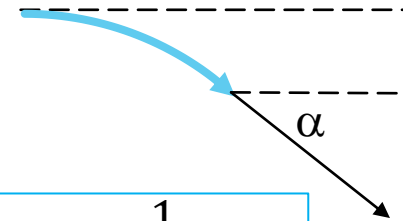
► Magnetic spectrometers

- Measures momentum
- Best momentum resolution at lower momentum
- Position Resolution: $\frac{\partial p}{p} = c_2 p$
- Multiple Scattering: $\frac{\partial p}{p} = c_1$
- Total: $\frac{\partial p}{p} = \sqrt{c_1^2 + (c_2 p)^2} = c_1 \oplus c_2 p$
- Non-destructive to the particle

$$\beta \propto \frac{1}{p}$$

$$\partial\beta = \text{const} \propto \frac{\partial p}{p^2}$$

$$\frac{\partial p}{p} \propto \text{const} \times p$$



$$\partial\beta = \frac{13.6}{\beta c p} z \sqrt{x/X_0} [1 + 0.038 \ln(x/X_0)]$$

$$\partial\beta = \frac{\text{const}}{p} \propto \frac{\partial p}{p^2}$$

$$\frac{\partial p}{p} = \text{const}$$

$$X_0 = \frac{716.4A}{Z(Z+1) \ln(287/\sqrt{Z})}$$

Rad.len.	$dE/dx _{\min}$	Density
X_0	{ MeV }	{ g cm ⁻³ }
{ g cm ⁻² }	g ⁻¹ cm ²	({ gℓ ⁻¹ })

Electron Motion in a Gas

▶ Low electric fields

- ▶ Collisions of electrons with gas molecules provide an effect backward “Force” on the electron proportional to velocity:

$$F_B = -bv.$$

- ▶ Terminal velocity when $qE = bv$; $v_T = \frac{q}{b}E$

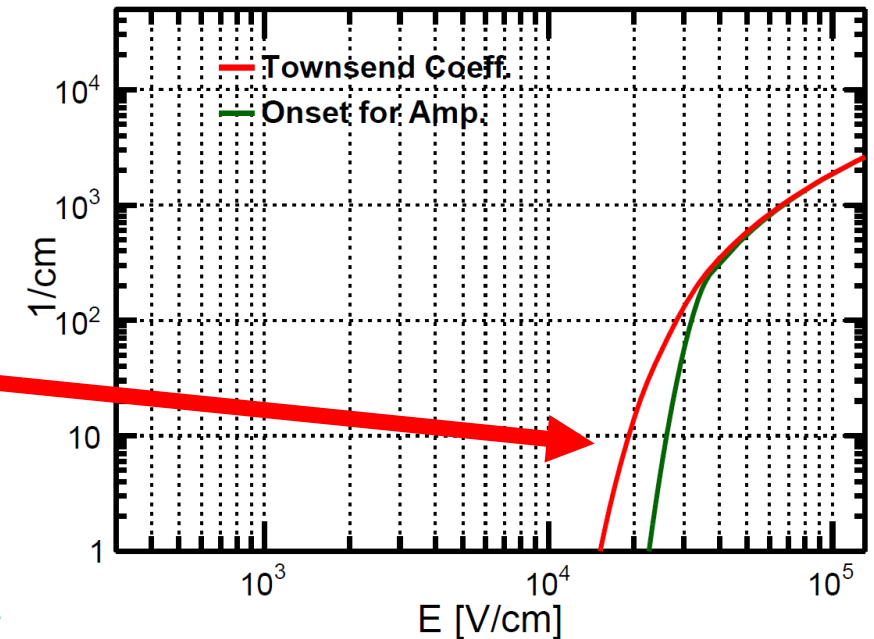
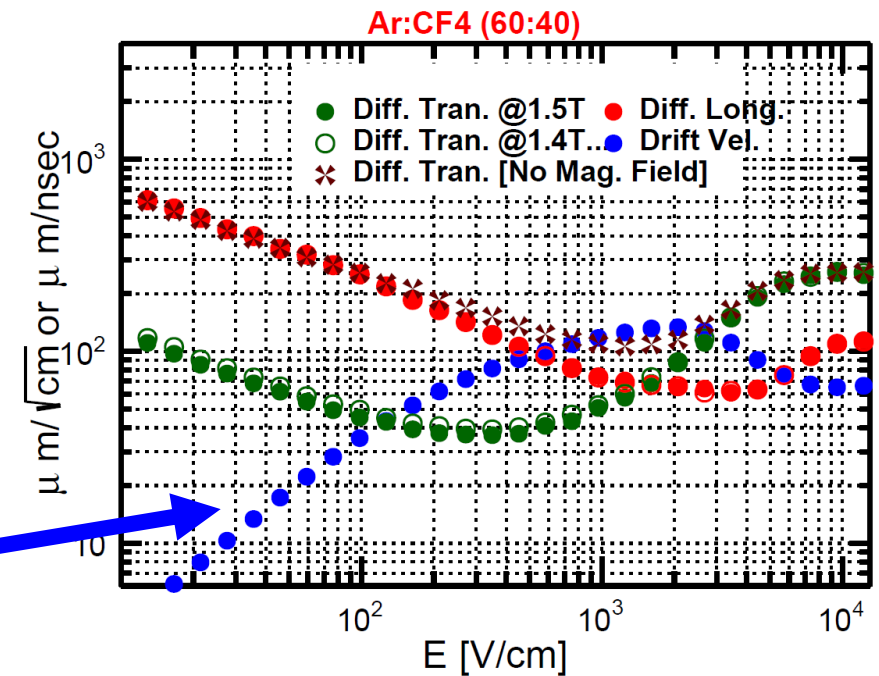
- ▶ Ions are simple (“b” is effectively constant and $v_T \propto E$)
- ▶ Electrons are more complicated as “b” depends upon v

▶ High Electric Fields

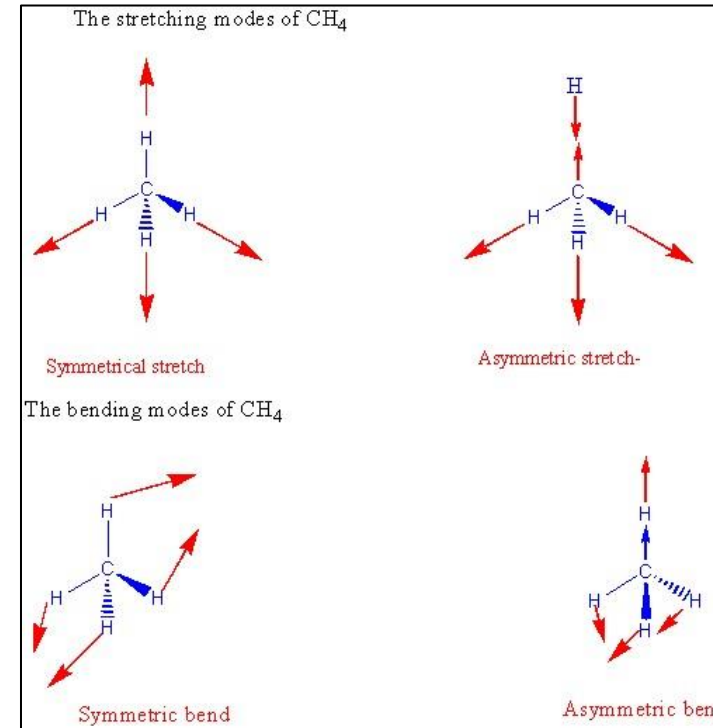
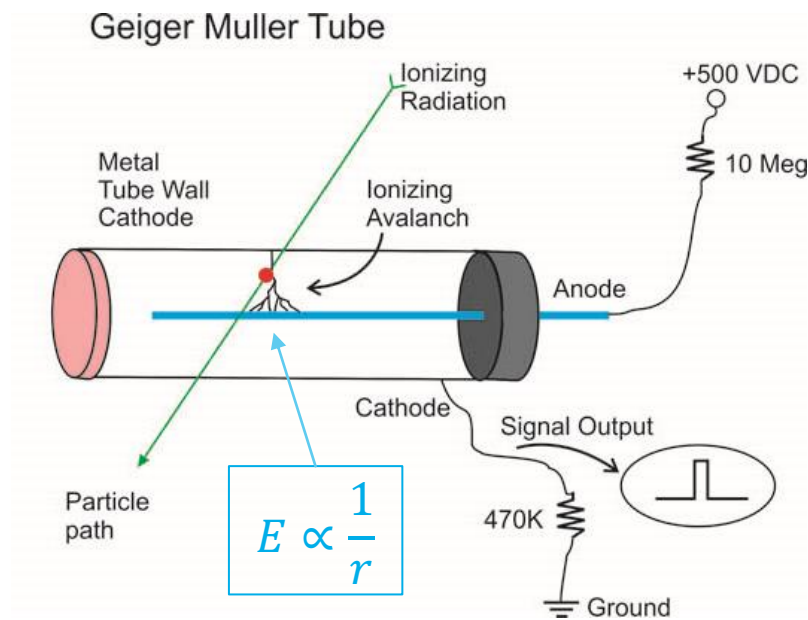
- ▶ Electrons gain enough energy in a mean free path that they are capable of ionizing the gas.
- ▶ Exponential growth in electron current:

$$\frac{I}{I_0} = e^{\alpha x}$$

- ▶ Growth controlled by the Townsend Coefficient, α .
- ▶ Virtually all gas detectors use a low electric field to drift electrons into an avalanche region for detection.

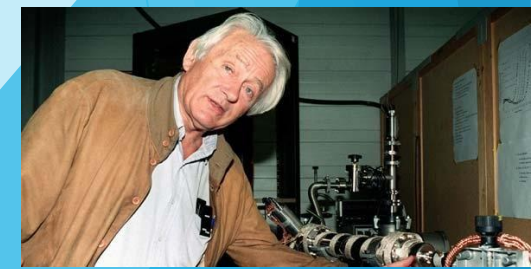


Geiger-Muller → Charpak (Nobel 1993)

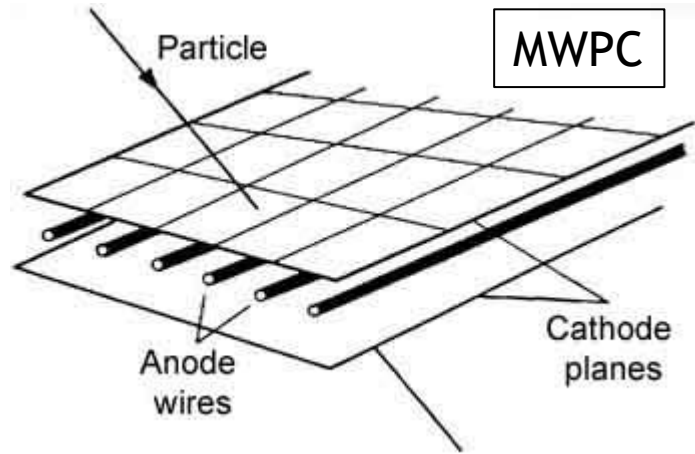


- ▶ With a pure noble gas, photons from the avalanche create ionization in the outer region.
- ▶ The avalanche “runs away”
 - ▶ Good news: BIG signal
 - ▶ Bad news:
 - ▶ Dead time to recharge wire
 - ▶ Insensitive to energy deposited

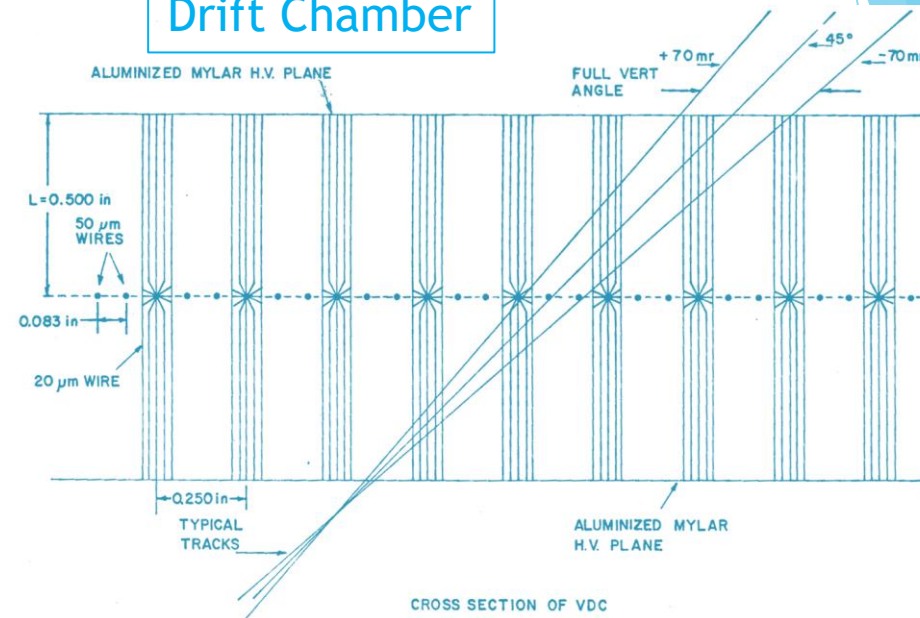
- ▶ Complex molecules have the ability to absorb energy as vibrations and rotations.
- ▶ Adding this “quench” to the gas breaks the feedback loop.
- ▶ “Proportional Counter”



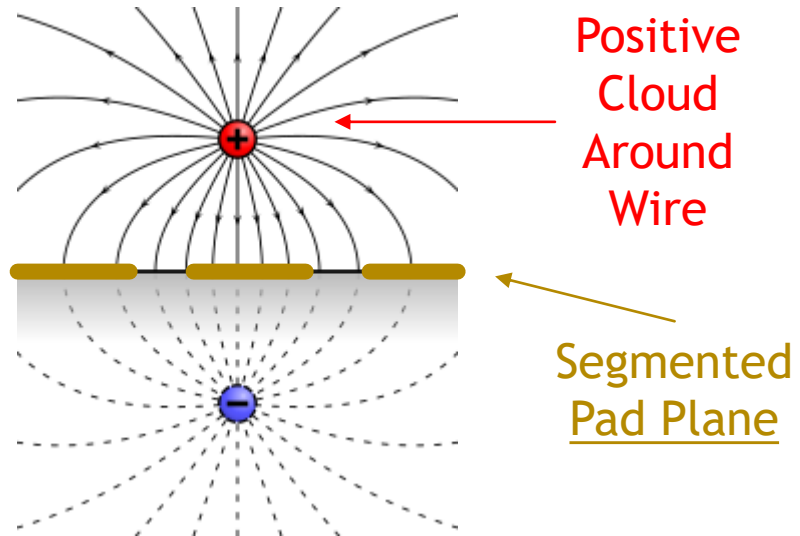
Gas Chamber Evolution



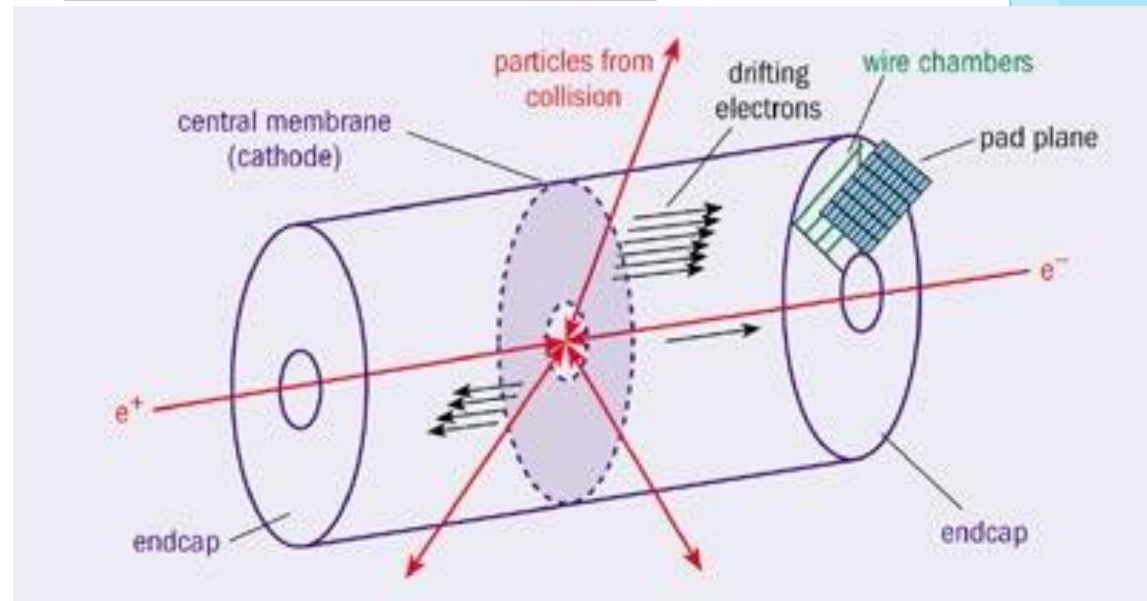
Drift Chamber



Pad Plane

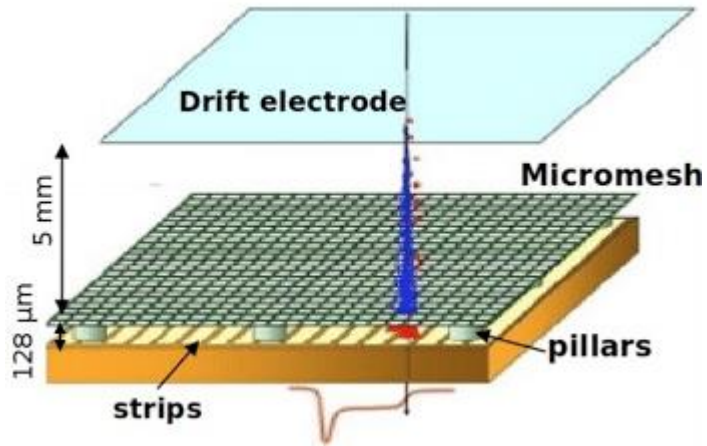


Time Projection Chamber



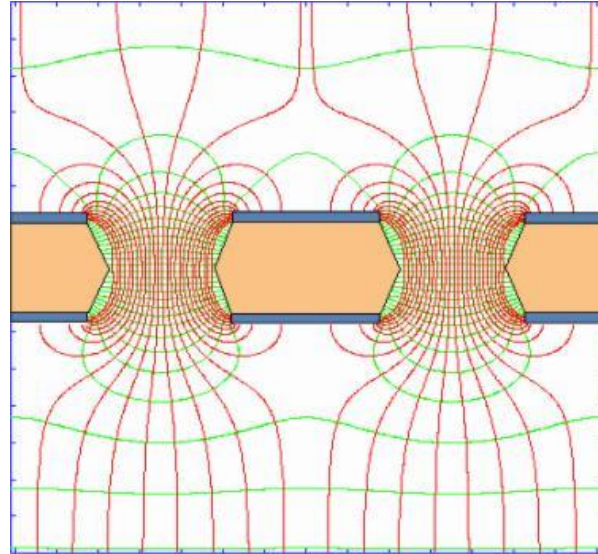
Avalanche Stage Evolution (MPGD)

microMEGAS



Ioannas Giomataris

Gas Electron Multipliers GEMs

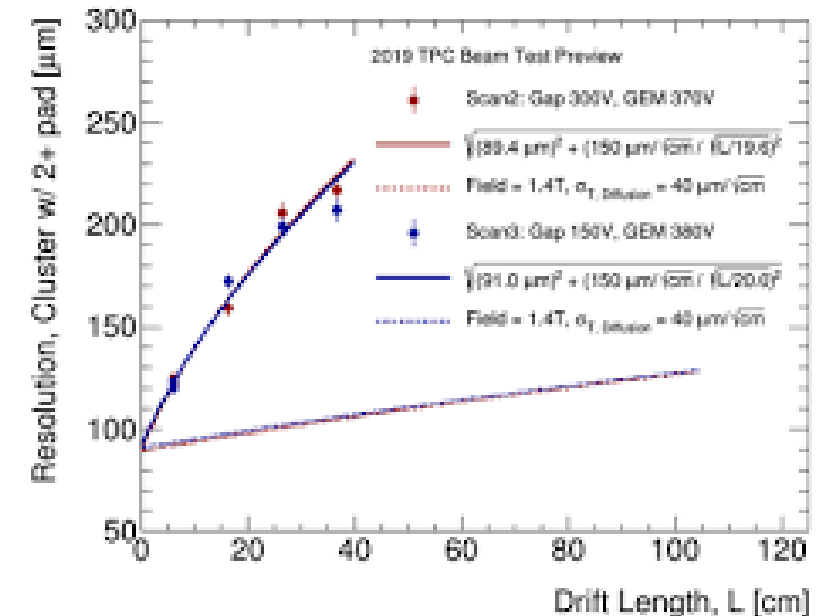
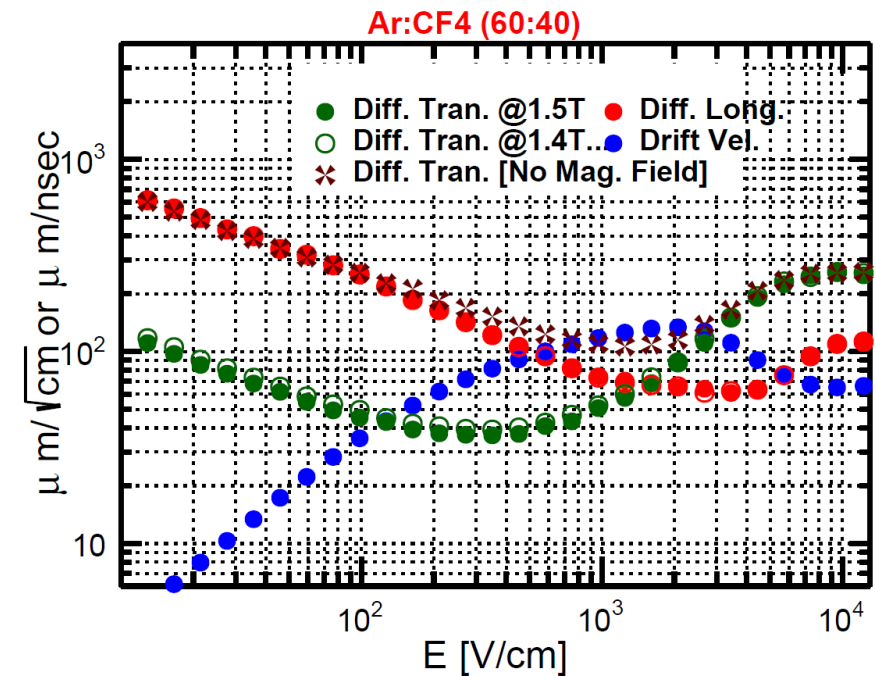


Fabio Sauli

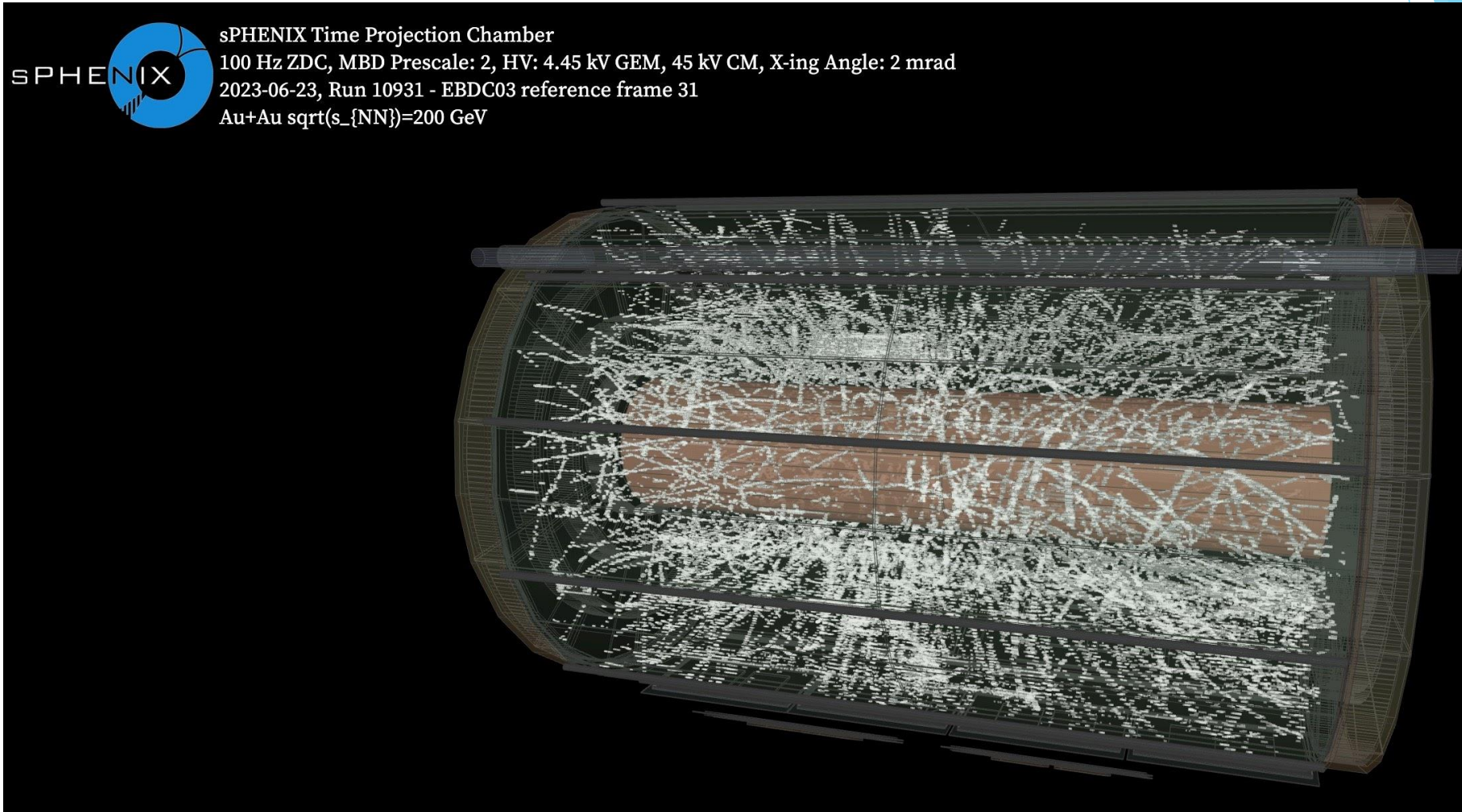
Precision of Gas Measurements

- ▶ Many factors can affect resolution, but two terms affect all detectors.
 - ▶ Readout Intrinsic Resolution
 - ▶ Constant...segmentation, noise, etc...
 - ▶ σ_I
 - ▶ Diffusion
 - ▶ Diffusion is a transverse kick due to electron-molecule collisions
 - ▶ Diffusion follows the mathematics of a random walk.
 - ▶ A single electron “misses” it’s target by $\sigma_1 = D_T \sqrt{L}$
 - ▶ The error on the mean of N electrons revied is $\sigma_T = \frac{\sigma_1}{\sqrt{N}}$
 - ▶ Other effects are either sensitive to the number of electrons in a signal or not. The typical formula becomes: $\sigma = \sigma_i \oplus \frac{D_T}{\sqrt{N_{eff}}} \sqrt{L}$
 - ▶ Notice that transverse diffusion falls STEEPLY with magnetic field.

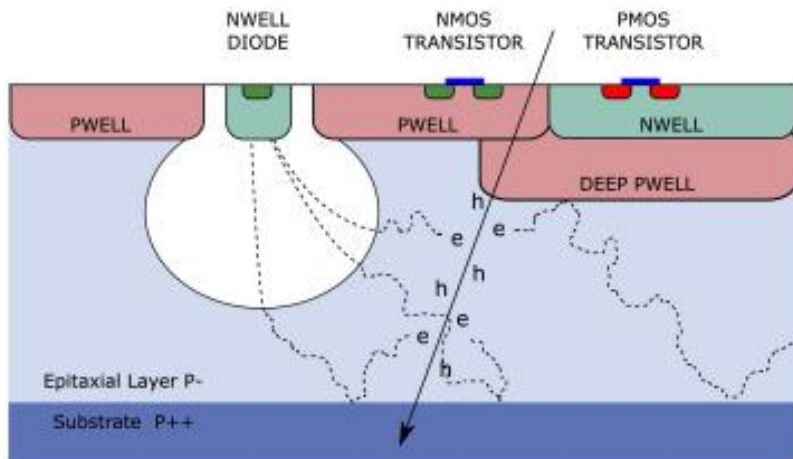
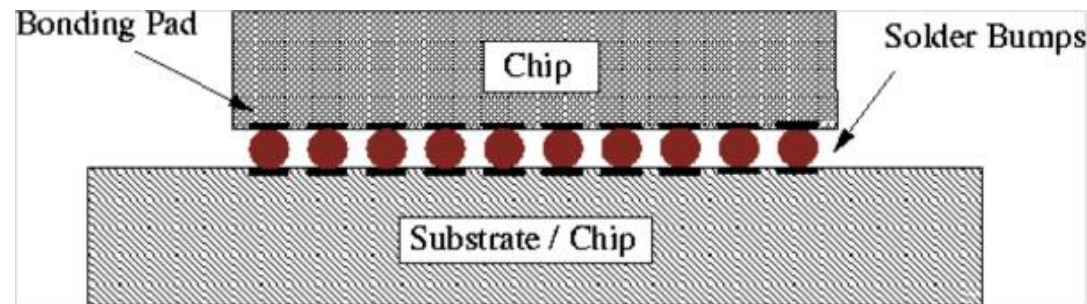
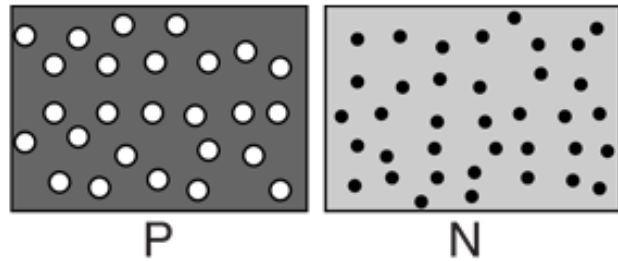
D_T has units of $\frac{\mu m}{\sqrt{cm}}$



sPHENIX TPC Example

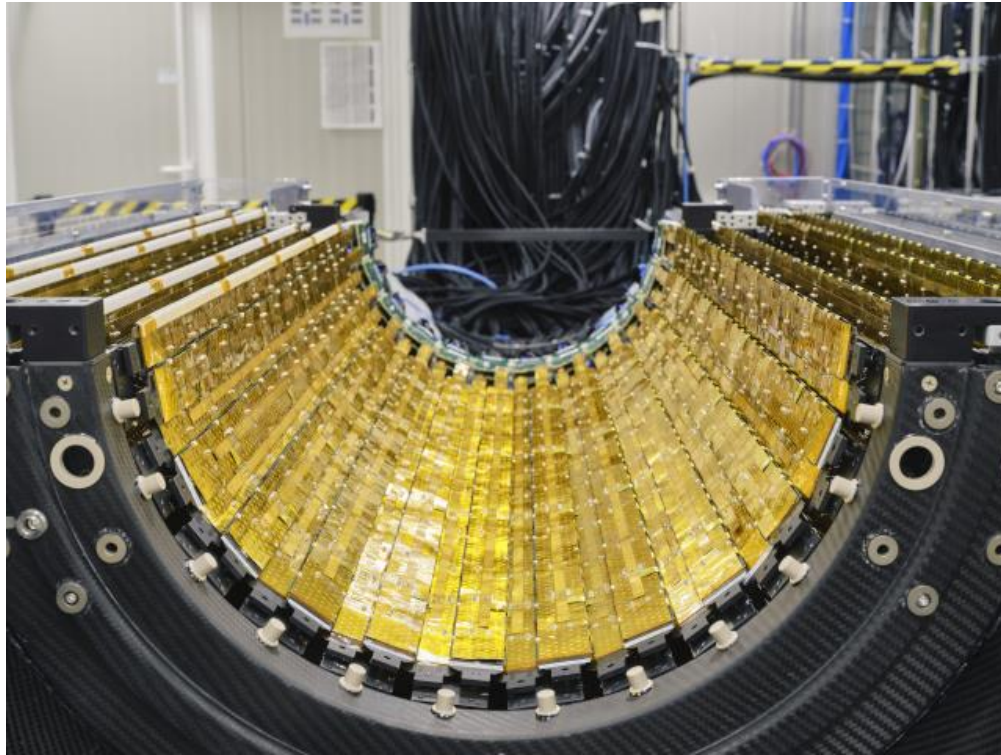


Silicon Detectors



- ▶ Reverse bias diode ideally flows no current and created a depletion region in the middle.
- ▶ A charge particle energy deposit in the depleted region makes current surge.
- ▶ Chip technology allows a pattern of sensors to be printed onto a chip.
- ▶ “Bump bonding” allows individual sensors to be connected to a readout chip.
- ▶ Monolithic Active Pixel Sensor (MAPS)
- ▶ Point-like collection of the signal frees space on chip to allow readout.
- ▶ Point-like collection reduces input capacitance and makes noise low.
- ▶ ALPIDE (shown) is a leading MAPs chip.

ALPIDE in ALICE and sPHENIX



- ▶ ALPIDE sensors developed by ALICE.
- ▶ Same technology also used in sPHENIX.
- ▶ NOTE: Silicon and Gas resolutions use different units!
 - ▶ Silicon uses pixel size; gas uses sigma (factor of $\sqrt{12}$)
 - ▶ $29 \times 27 \mu\text{m}^2$ means sigmas of $8.4 \times 7.8 \mu\text{m}^2$

PID ~ Velocity

▶ $p = m\gamma\beta$ $E = m\gamma$ velocity(β) measurement yields mass.

▶ Direct measurement:

- ▶ Record signal time at multiple locations, calculate v .
- ▶ “Fast” detector = low transit time spread (most easily achieved at small transit time)

▶ Velocity-dependent interaction(s) with detector:

▶ Specific Ionization (aka $\frac{dE}{dx}$)

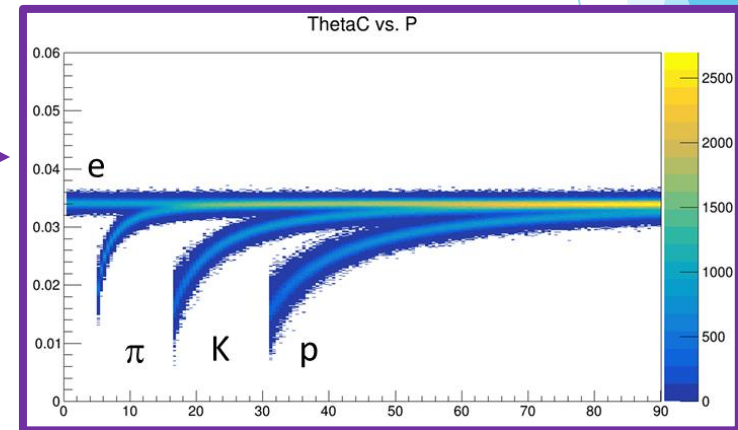
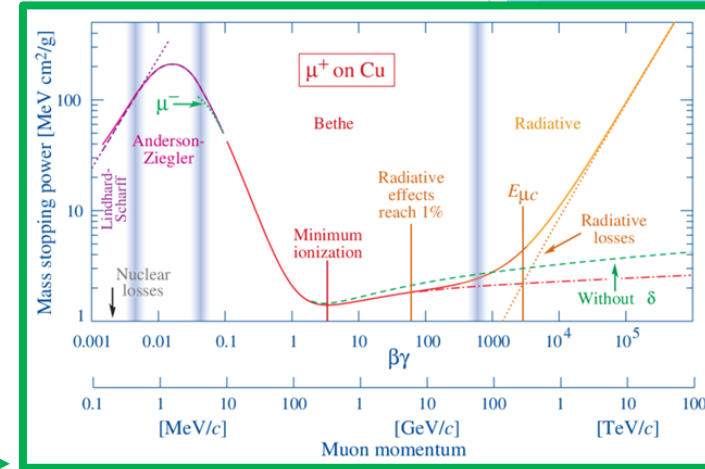
▶ Cherenkov Radiation: $\cos \theta_C = \frac{1}{n\beta}$

- ▶ θ_C measured wrt track direction.
- ▶ Thus dependent upon deliverables from tracking

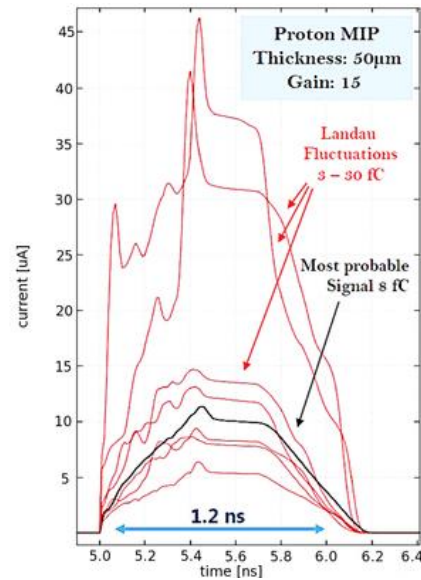
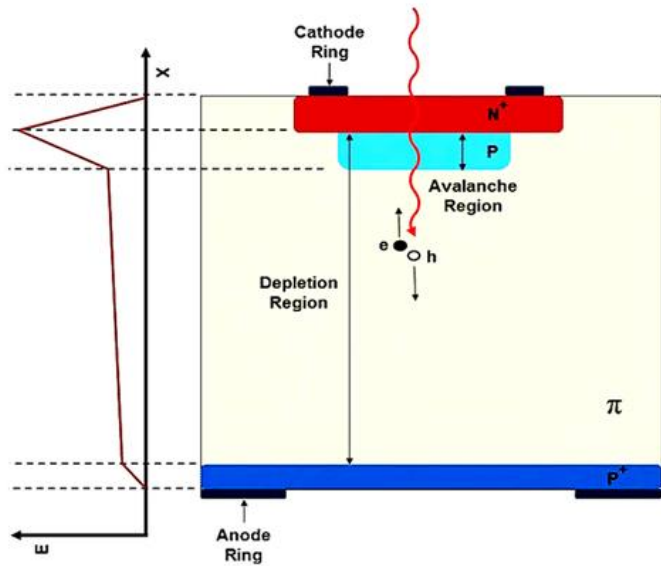
▶ Bremsstrahlung: $P = \frac{q^2\gamma^4}{6\pi\epsilon_0 c} \left(\dot{\beta}^2 + \frac{(\vec{\beta} \cdot \dot{\vec{\beta}})^2}{1-\beta^2} \right)$

▶ Transition Radiation: $I = \frac{Z^2 e^2 \gamma \omega_p}{3c}$

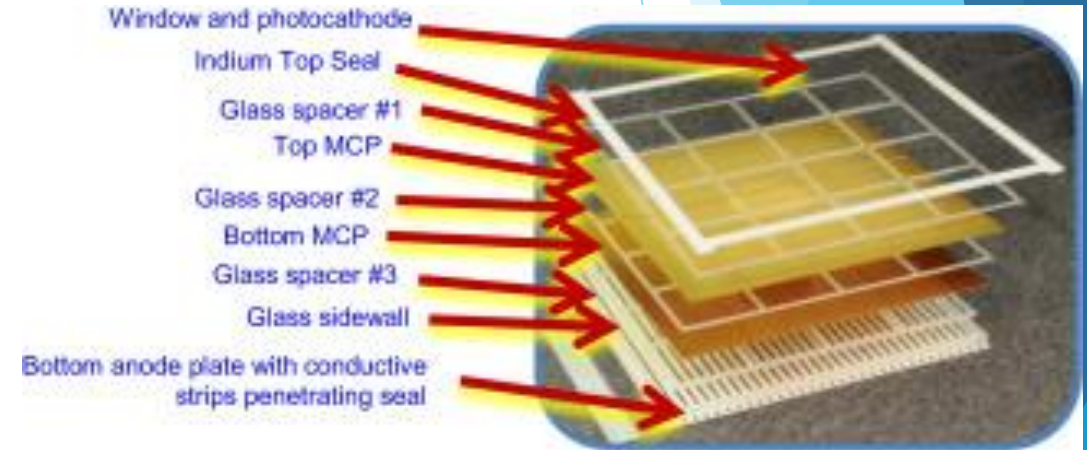
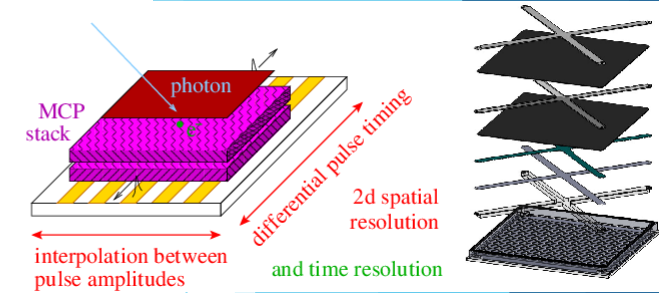
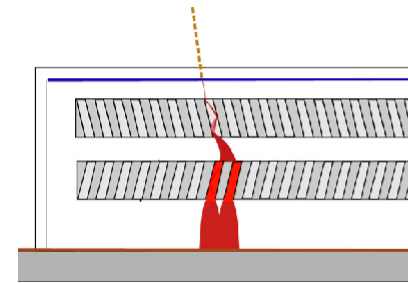
eID mechanisms



Time-of-Flight

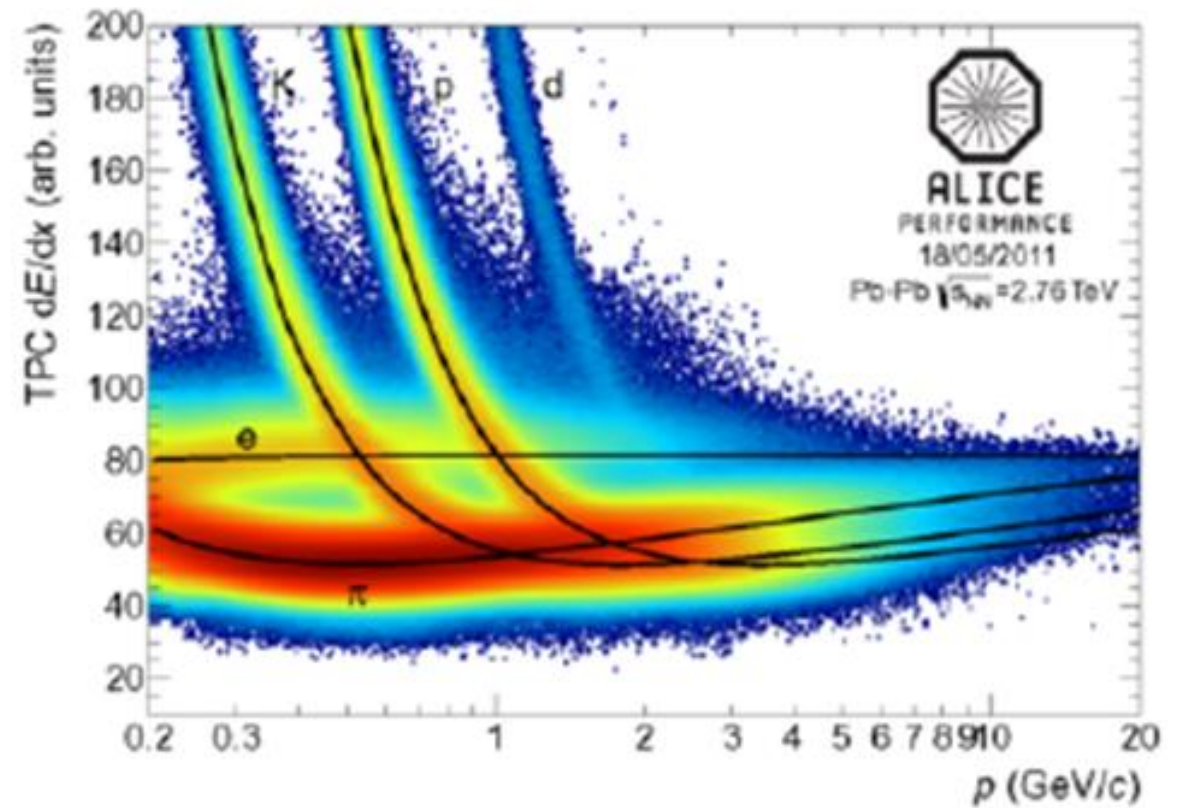
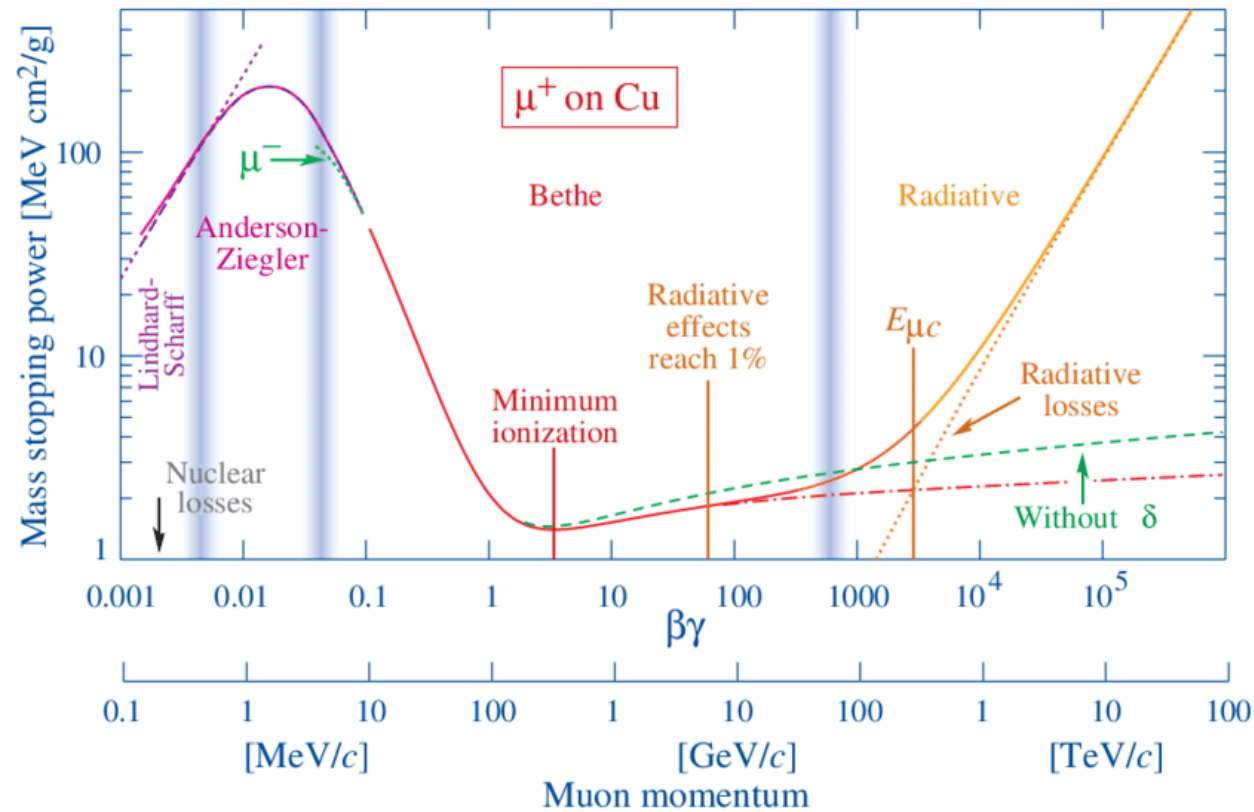


MCP: Micro Channel Plate



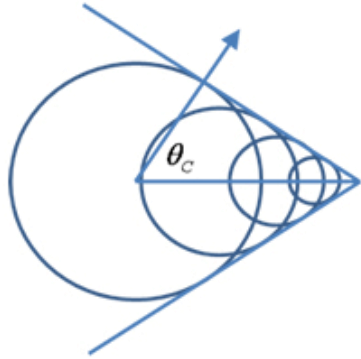
- ▶ Direct measurements of velocity require a fast charge collection device.
- ▶ The time resolution is automatically helped by small devices:
 - ▶ LGAD: Low Gain Avalanche Device. Instantly small.
 - ▶ MCP: Micro Channel Plate → LAPPD (Large Area Picosecond Photon Detector)

dE/dx



- ▶ dE/dx is a velocity-dependent interaction with the detector.
- ▶ Combined with momentum it provides PID.
- ▶ “Band crossings” are at $\beta \sim 0.7$ and can be distinguished many ways.

Cherenkov Light



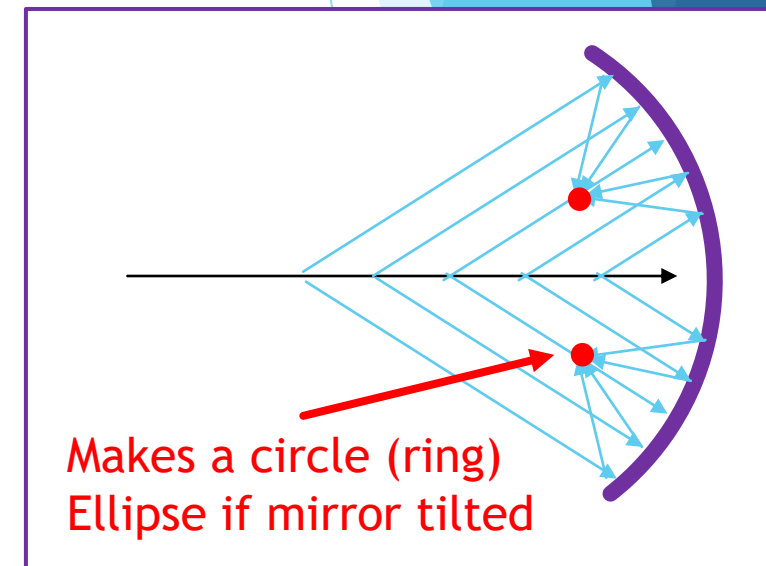
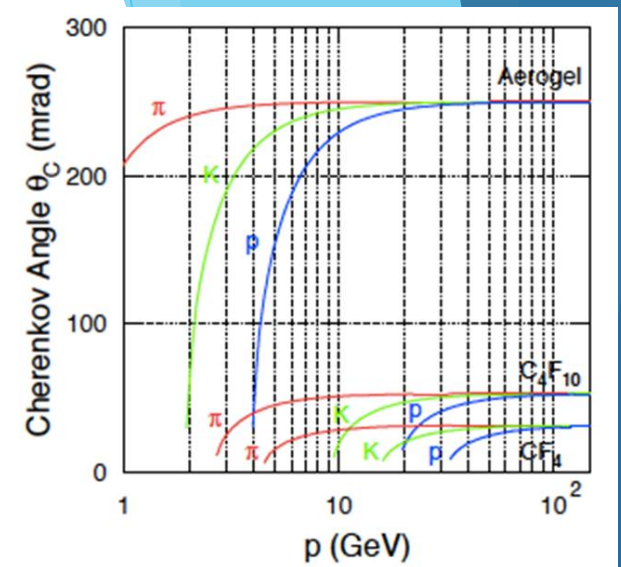
Cherenkov Angle
determines β

$$\theta_c = \cos^{-1} \left(\frac{1}{\beta n} \right)$$

- ❑ Cherenkov Light - “Optical Boom”
- ❑ Cherenkov angle depends upon velocity & index
 - Low n distinguishes particles to high momentum.
 - Low n yields less light!
- ❑ Go where the light is:

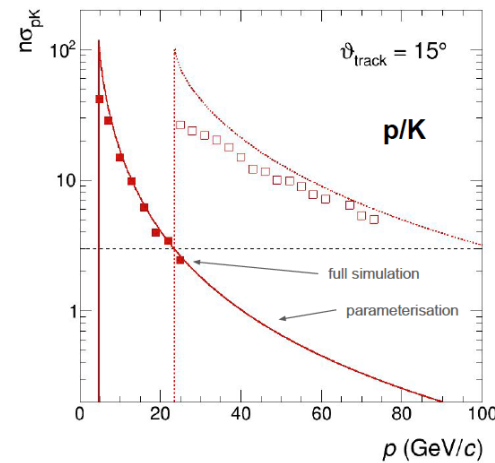
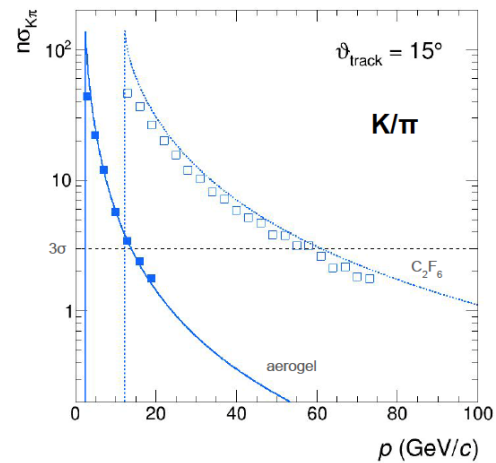
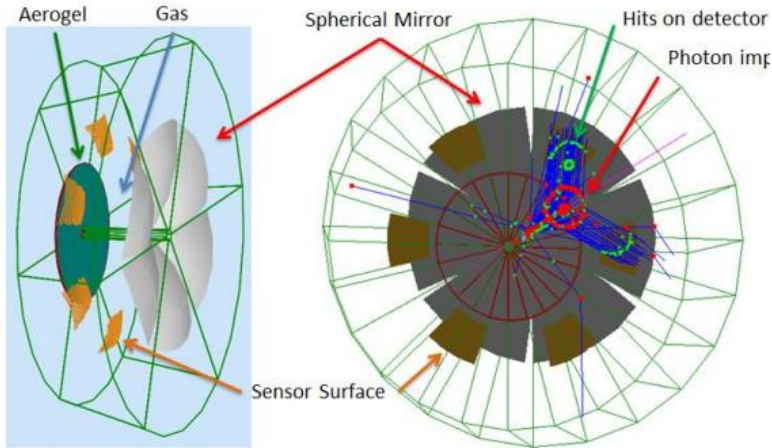
$$\frac{dN}{dL} = 2\pi\alpha_{EM} \sin^2 \theta_c \int_{\lambda_{MIN}}^{\infty} \epsilon(\lambda) QE(\lambda) \frac{1}{\lambda^2} d\lambda$$

- Deep UV Cherenkov yield is highest...how to get there?



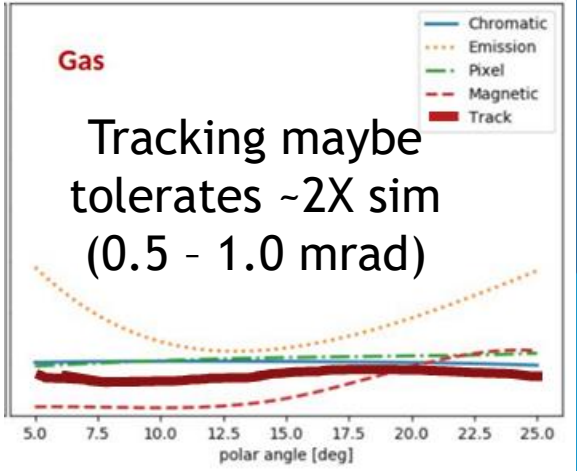
Dual RICH

- Cherenkov radiator
 - Refractive Index
 $n = 1.02$ (aerogel) 1.0008 (C_2F_6)
 - Length of the radiator
 $L = 4$ cm (aerogel) , 160 cm (C_2F_6)
- Mirrors
- Photon Detector
3 mm pixel size; 200-500 nm MAPMT
- Particle Generation
Originate from the vertex



K/π and p/K separation as a function of momentum

- Exquisite detail in simulation.
- AI-based optimization.
- Uses constant external angular resolution assumption.
- Similar external constraint assumed as deduced for separate gas & aerogel RICHes

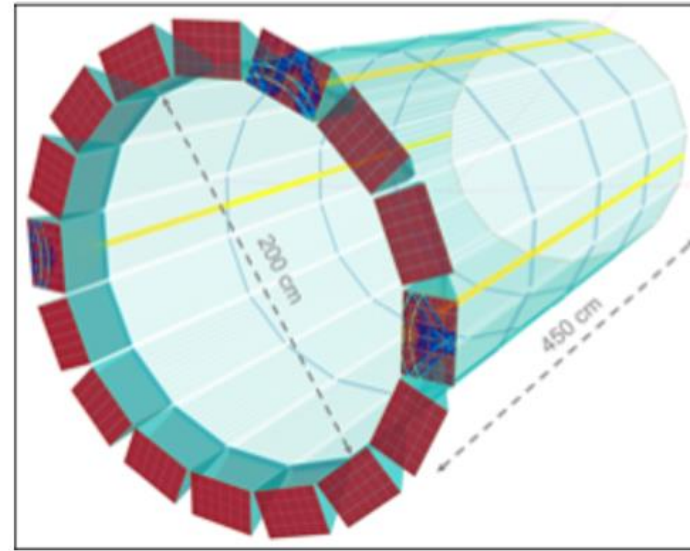
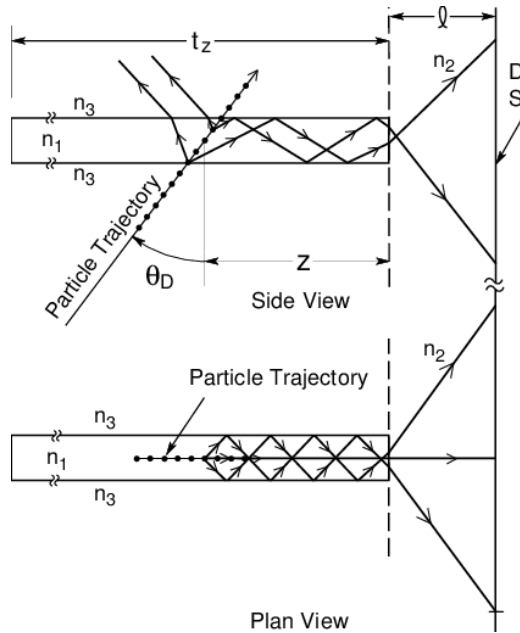


NOTE: PMTs to the side makes optical aberration (“Emission”) dominant

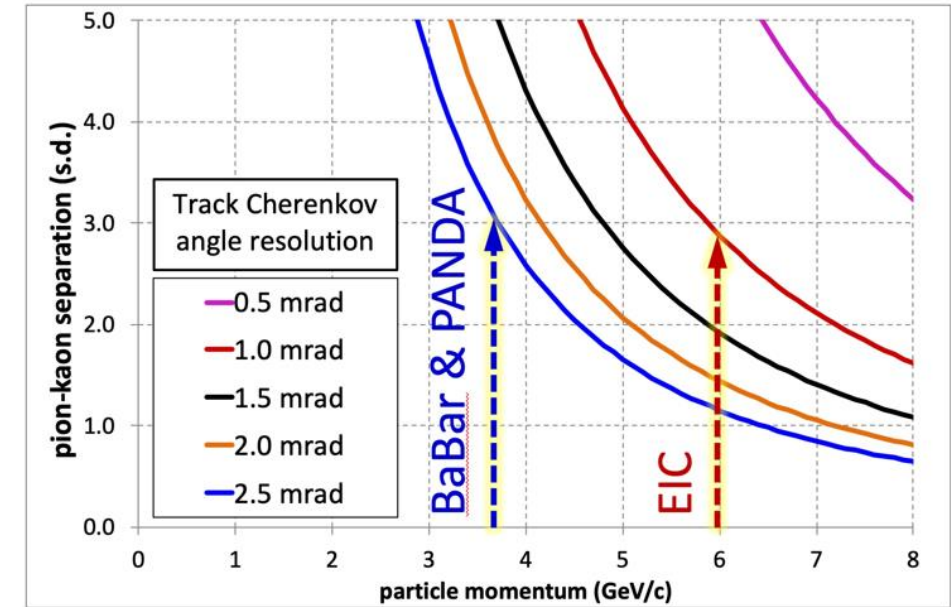
Tracking	
Angular resolution	$\sigma = 0.5$ mrad (1 mm over 2 m) – whole momentum range
Impact point resolution	$\sigma = 0.3$ mm
Momentum resolution dP/P	+/- few percent negligible effects in Cherenkov angle reconstruction
Magnetic Field	
3 Tesla Central Field in JL-MEIC spectrometer	
Space Requirement	
(based on original spectrometer constraints)	
longitudinal length	JLEIC: ≈ 1.6 m, ePHENIX: ≈ 1.0 m
transverse radius	JLEIC: ≈ 2.5 m, ePHENIX: ≈ 2 m
beam pipe radius	< 10 cm
Background	
no direct external background only background produced by the simulated charged particle: Delta rays, Rayleigh scattering ...	

External Assumption

Detection of Internally Reflected Cherenkov: DIRC



hpDIRC



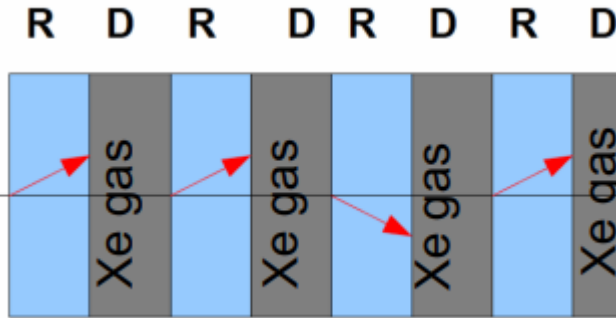
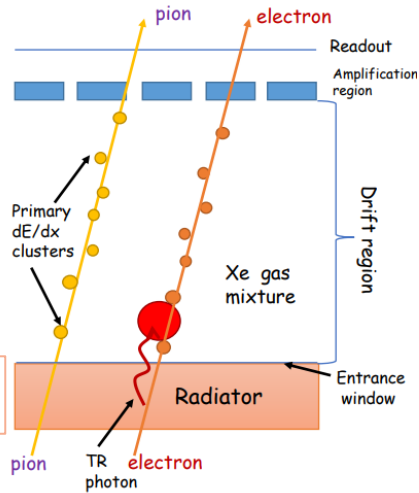
- ▶ At normal incidence, Cherenkov light is internally reflected for $n > \sqrt{2}$
- ▶ Cherenkov angle preserved if bars meticulously flat with 90 degree corners.
 - ▶ If you “back up the photon sensors” this is a proximity focus. (original BaBar)
 - ▶ If you use focusing elements it becomes a high performance hpDIRC
- ▶ Spectacularly compact PID device with photon detection at end.

eID by Transition Radiation (TRD)

1: Make TR photons:

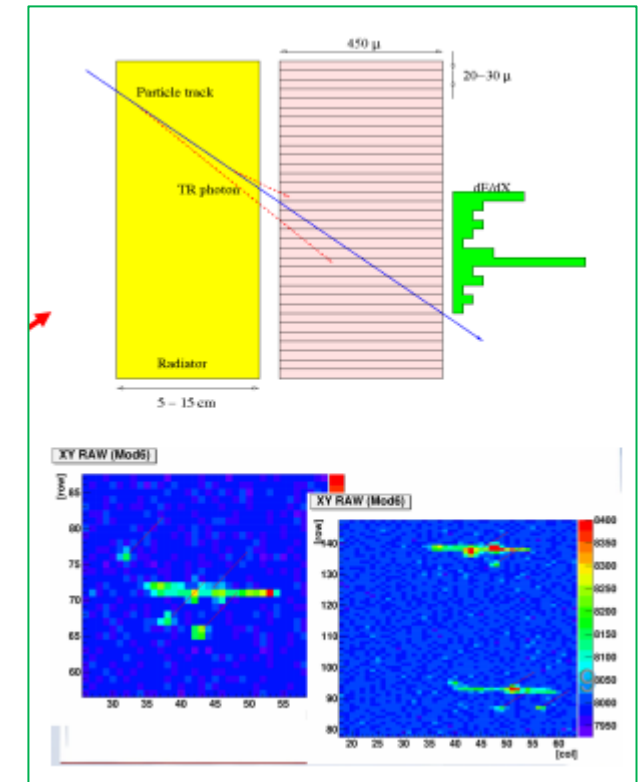
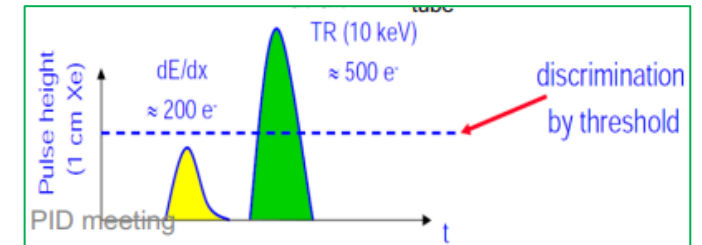
- ▶ Index, n , change
- ▶ X-ray regime

Film, foam, fleece

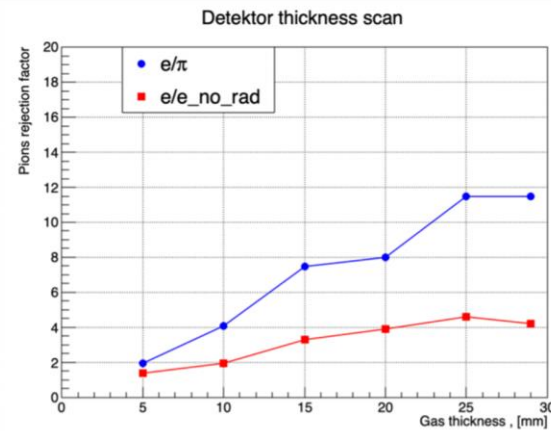
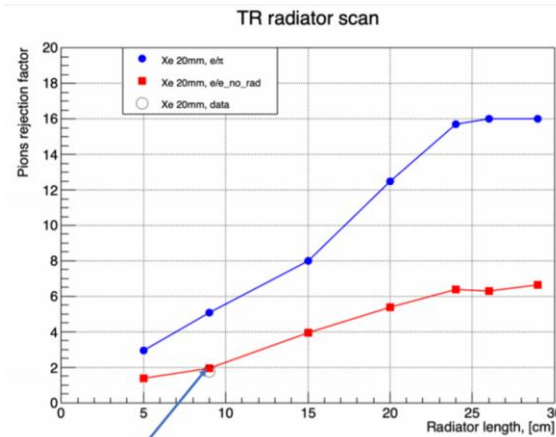


2: Layered to get more photons

3: ID TR photons by energy/pos



4: Reject Pions



- ▶ “Threshold detector” (TR X-rays produced or not)
- ▶ $P_{\min} \sim 1 \text{ GeV}$ (rejection \sim constant at higher momentum)

Balance between X_0 and rejection

Summary



- ▶ Welcome to Nuclear Physics!!!
- ▶ Detector Technology is fun and exciting.
- ▶ No time for so many things SiPM, readout electronics, streaming mode.
- ▶ Present capabilities are excellent and the future is better.
- ▶ You are entering the nuclear physics at an amazing time.
- ▶ Financial investments are excellent, particularly with the future electron-ion collider.