

Heavy quarks: a (the) golden probe of QGP

E. Scomparin
INFN Torino (Italy)

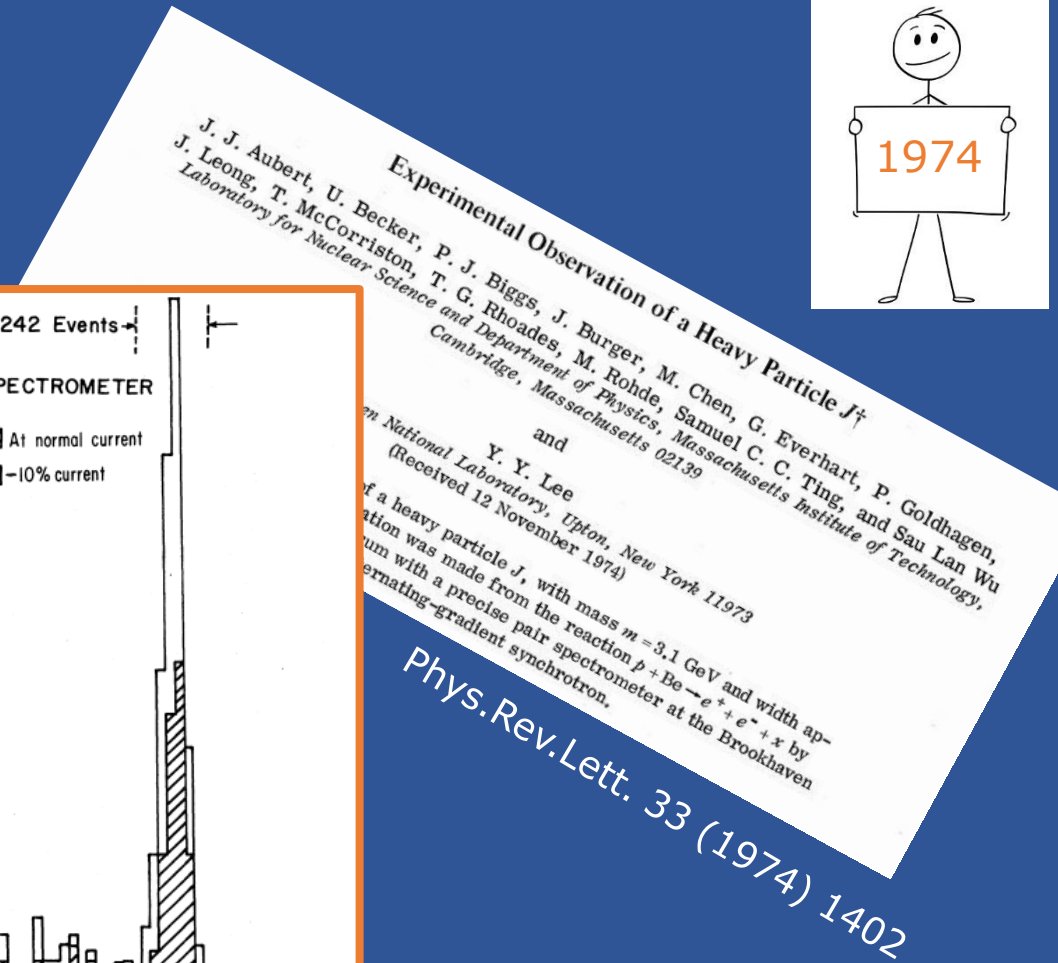
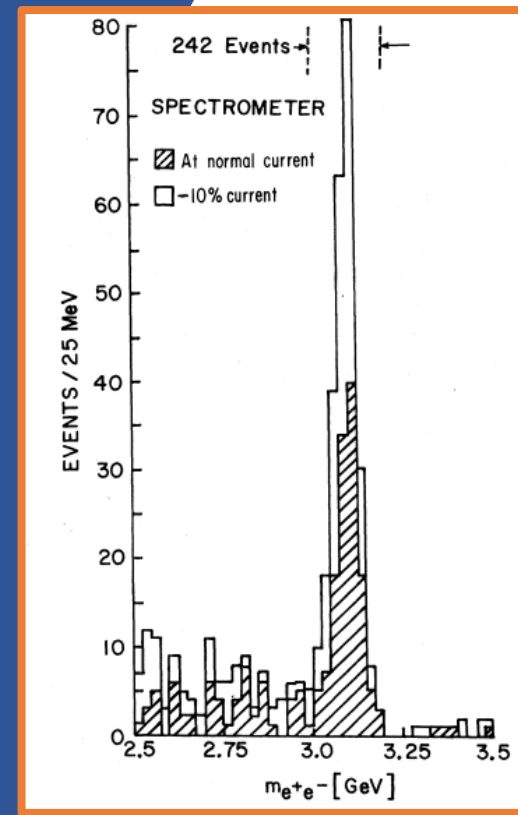
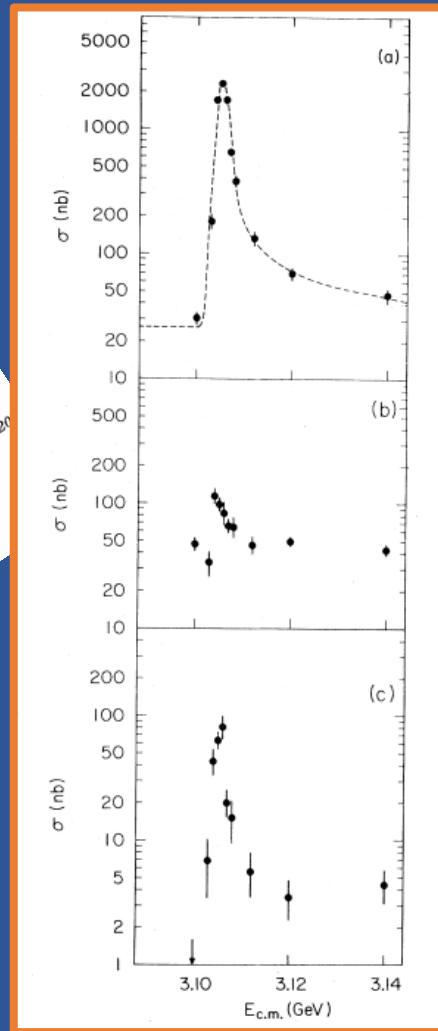
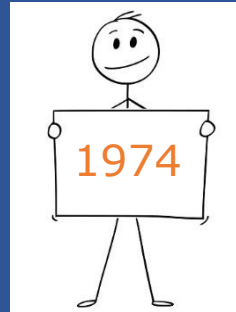


Quark Matter 2023 Student Day



Heavy quarks: a (the) golden probe of QGP

From the November revolution...



Heavy quarks: a (the) golden probe of QGP

From the November
revolution...

... to the discovery
of the J/ψ suppression...

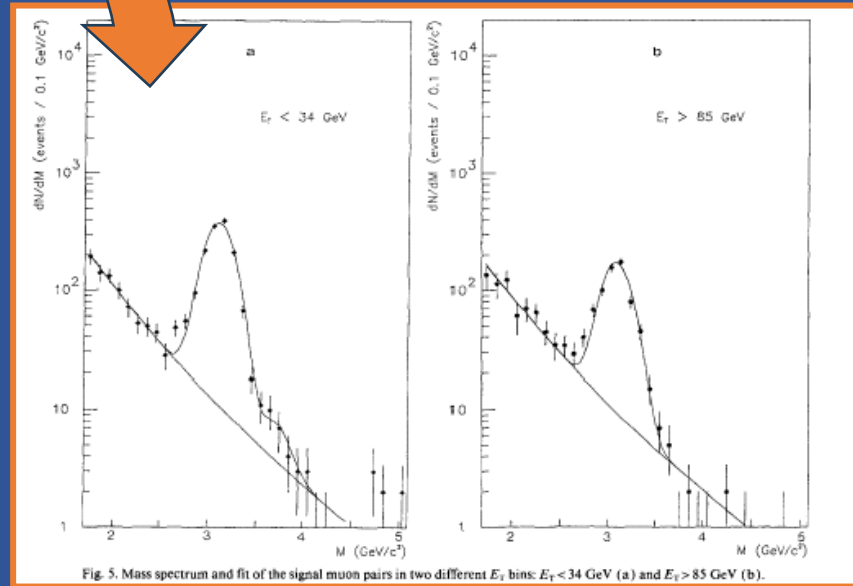
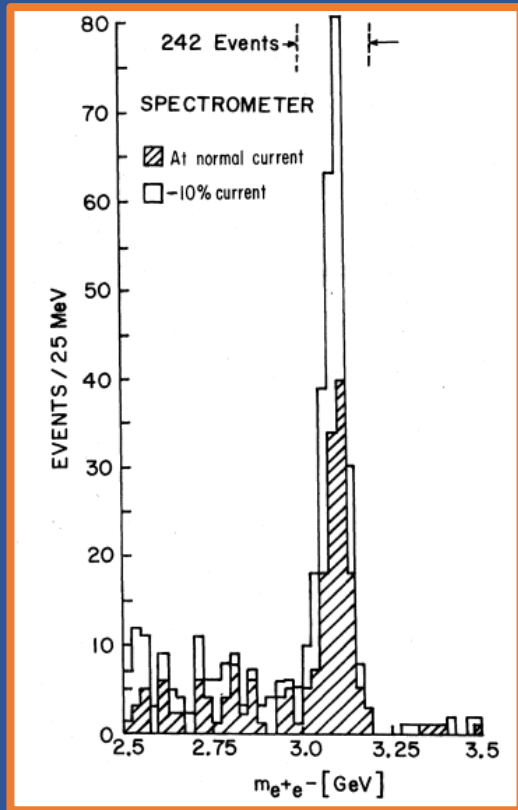
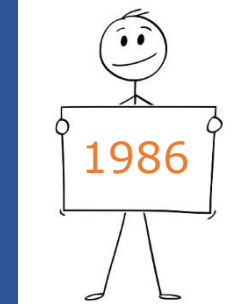


Fig. 5. Mass spectrum and fit of the signal muon pairs in two different E_T bins: $E_T < 34 \text{ GeV}$ (a) and $E_T > 85 \text{ GeV}$ (b).



Abstract

The study of oxygen-uranium reactions at 200 GeV/nucleon shows a significant transverse energy dependence of the yield of J/ψ 's relative to muon pairs produced in the mass continuum. This feature, observed for the first time, is in agreement with predictions from quark-gluon plasma formation, although alternative explanations by hadronic effects cannot be ruled out at this stage.

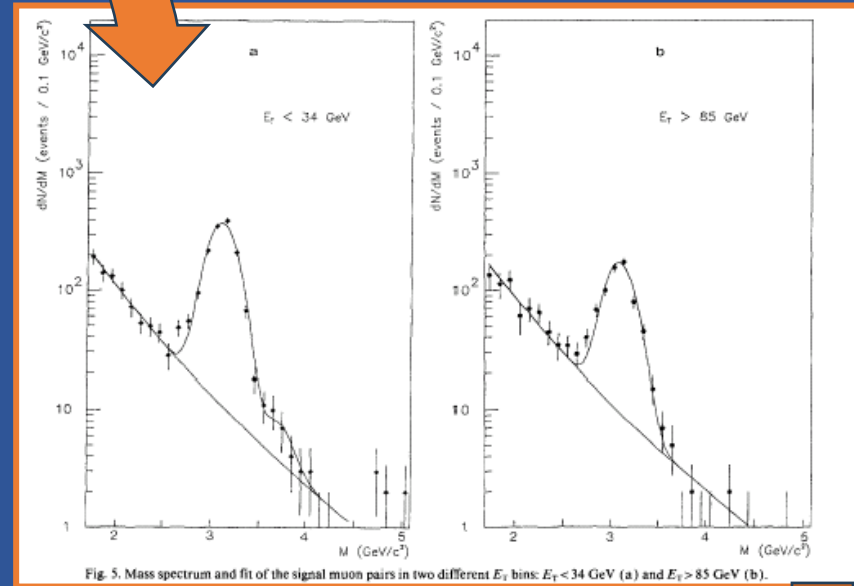
NA38, PLB220 (1989) 471

NA38, Z. Phys. 38(1988) 117

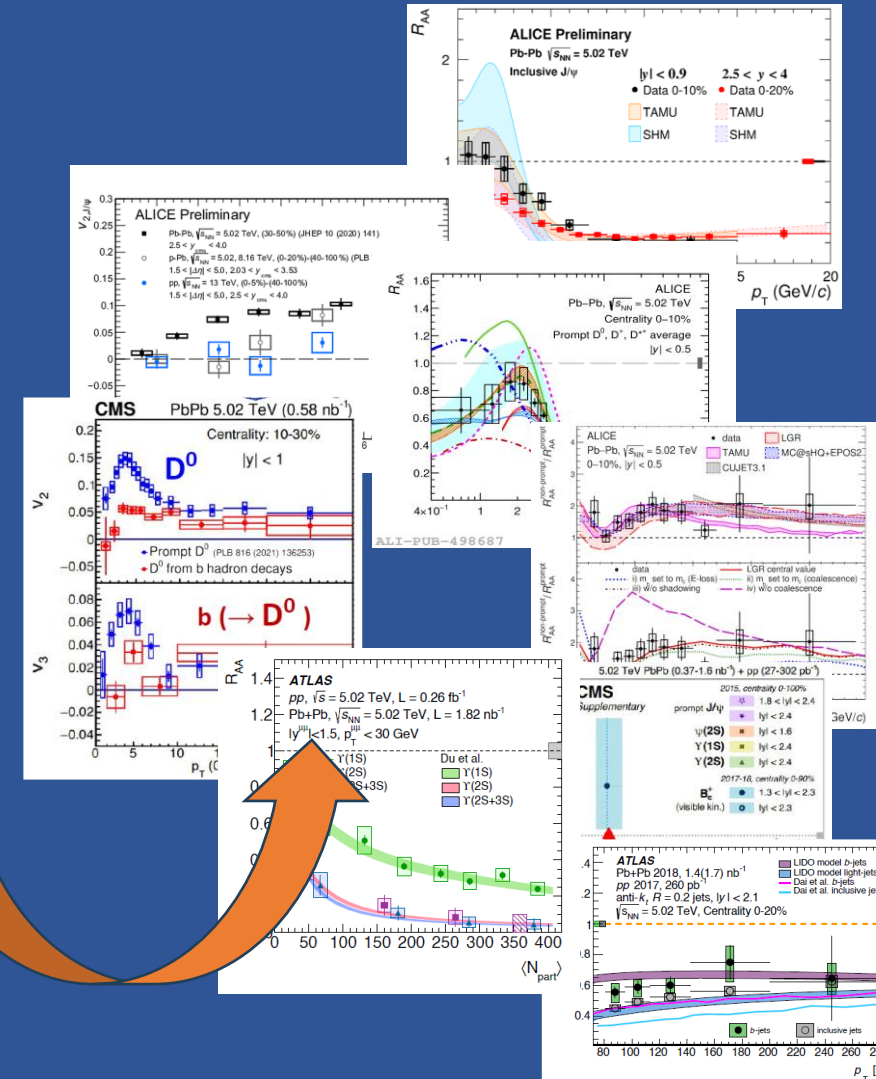
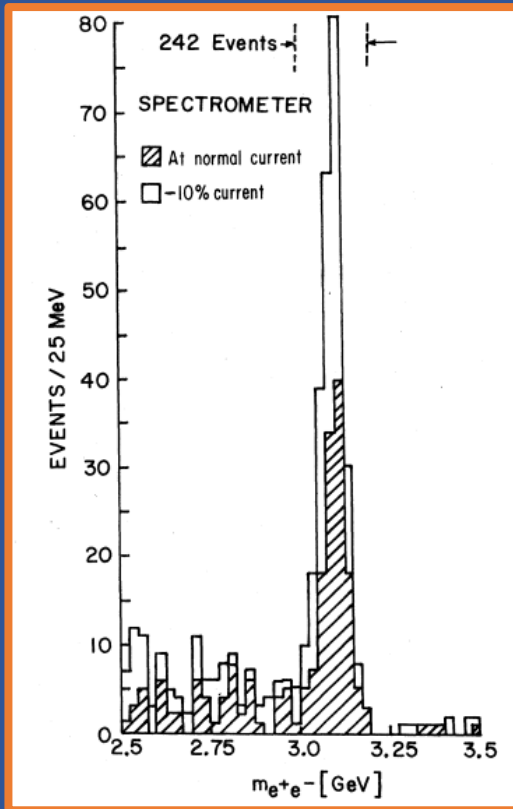
Heavy quarks: a (the) golden probe of QGP

From the November revolution...

... to the discovery of the J/ψ suppression...



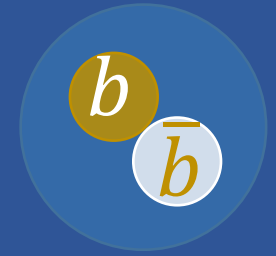
...to today's precision measurements at collider energies!



Heavy quarks: versatile observables

- (Heavy) Quarkonia: bound states of $c\bar{c}$ and $b\bar{b}$ quarks (no $t\bar{t}$!)
 - Rich spectroscopy
 - Wide range of binding energy
 - Several angular momentum states

$\Upsilon, \chi_b, \eta_b, \dots$



$J/\psi, \chi_c, \eta_c, \dots$

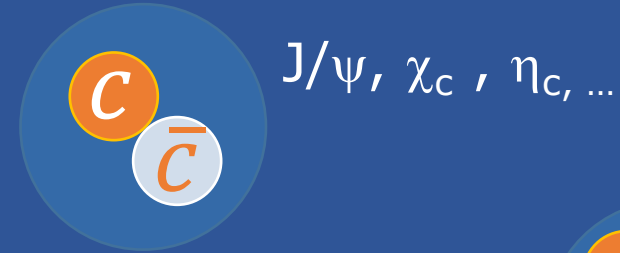
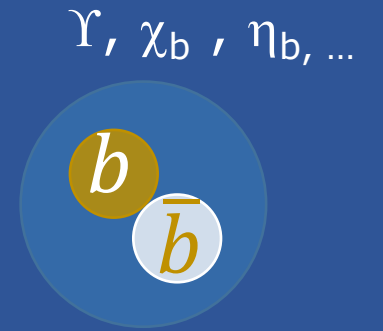


N.B.: “color” is here used to identify “flavour”!

Heavy quarks: versatile observables

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- Rich spectroscopy
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- Open heavy flavour hadrons

- Heavy quark can bind with a light or strange quark(s)
- Large variety of mesons AND baryons



D^0, D^+, \dots



D_s, \dots

$\Lambda_c, \Sigma_c, \dots$

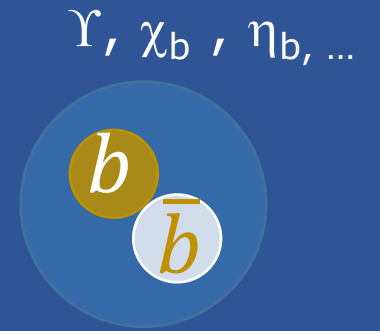
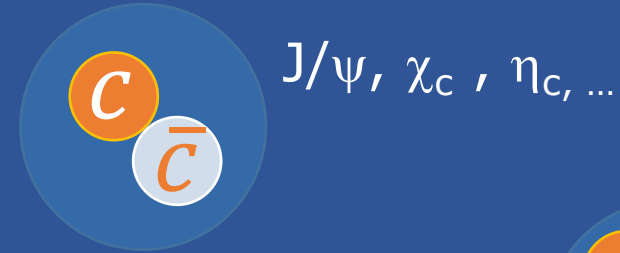


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Heavy quarks: versatile observables

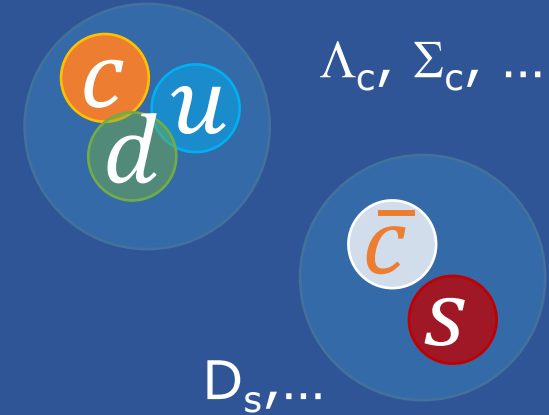
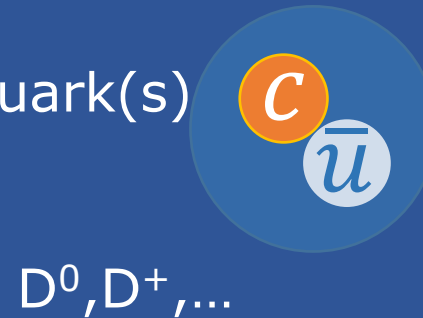
- (Heavy) Quarkonia: bound states of $c\bar{c}$ and $b\bar{b}$ quarks (no $t\bar{t}$!)

- Rich spectroscopy
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- Open heavy flavour hadrons

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- Exotica

- Beyond qqq and $q\bar{q}$
- Large number of states (discoveries continue)



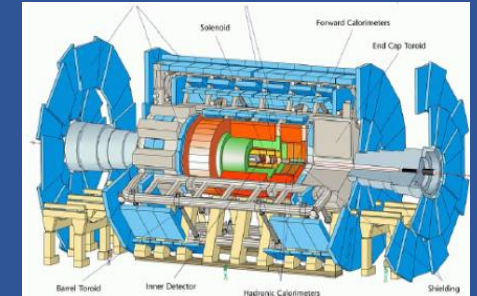
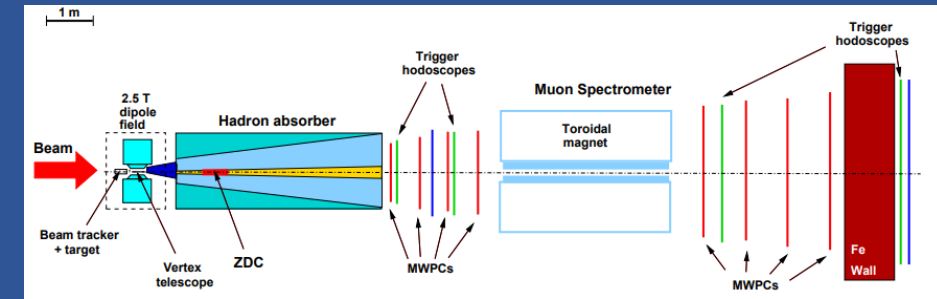
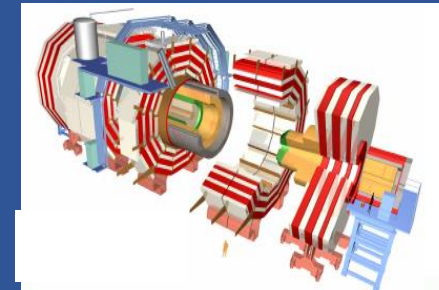
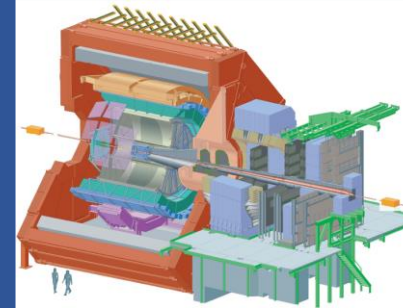
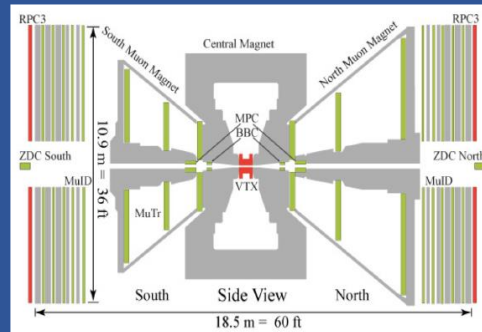
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Quarkonia

- How were they established as a signature of QGP ?
- Do they retain interest until today ?
- Is there space for further progress ?

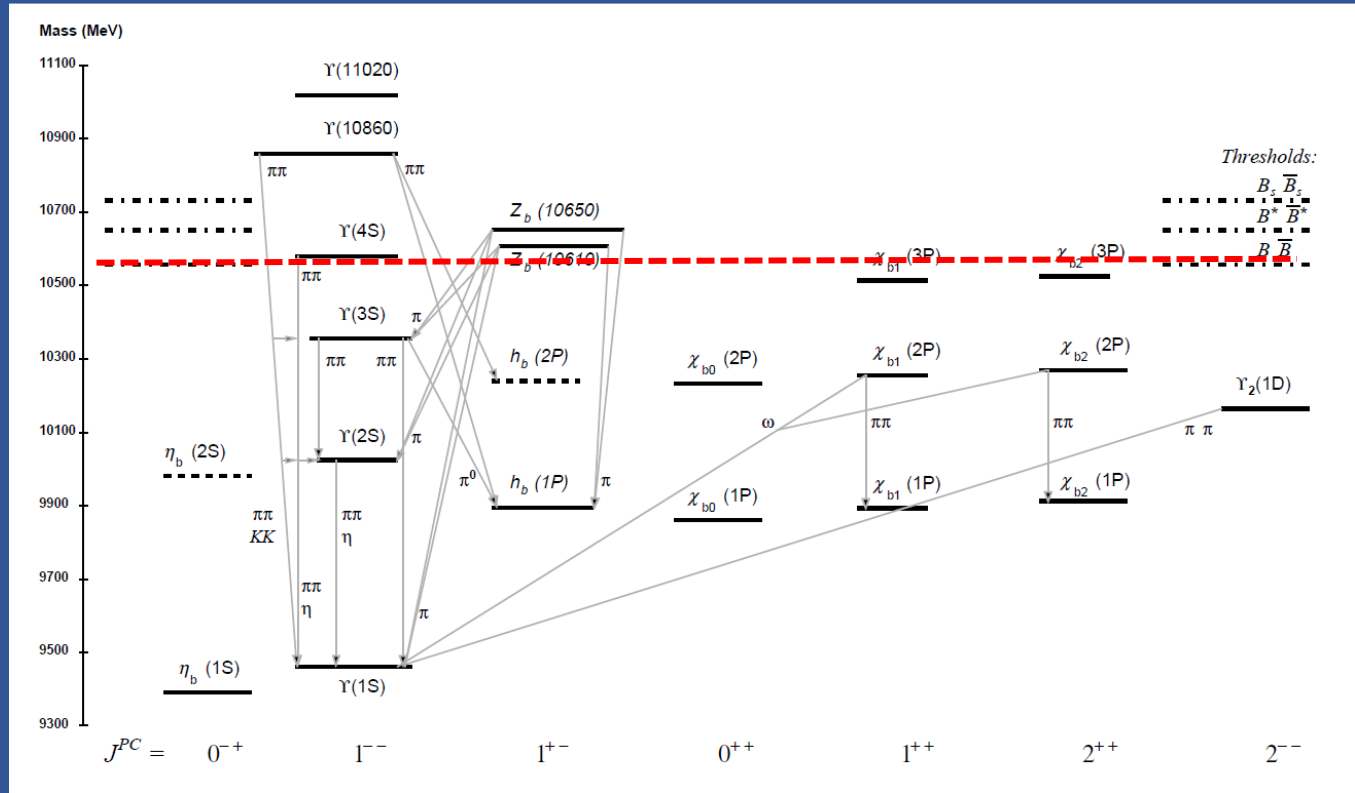
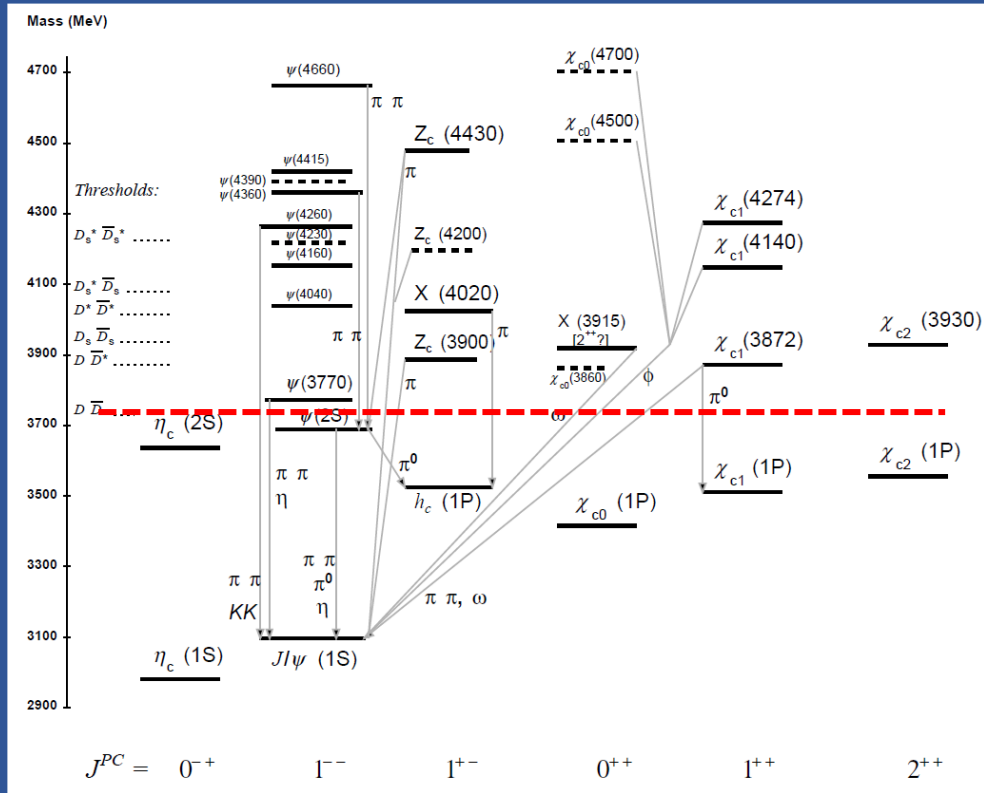
Mostly studied via their **dilepton decay**
(muon/electron spectrometers)

Quarkonia



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Quarkonium spectroscopy



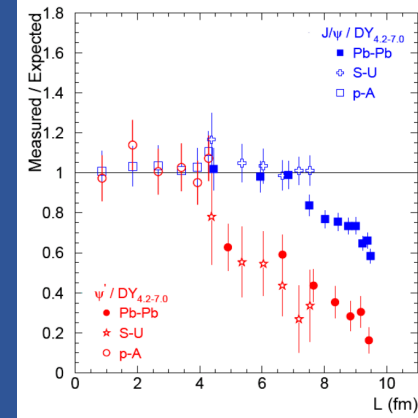
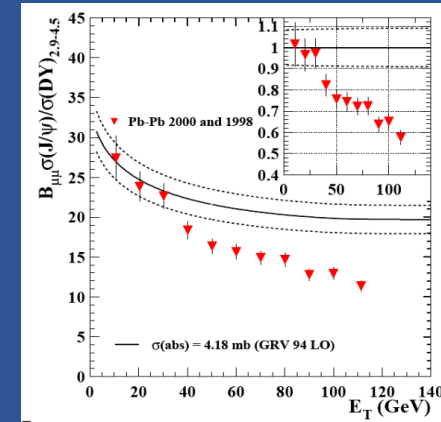
- Below open heavy flavour threshold
 - Extremely **narrow widths** ($\Gamma < 1$ MeV for most vector states)
 - Binding energy can be > 1 GeV ($\Upsilon(1S)$)
 - Non-relativistic system \rightarrow static properties can be explored via Schrodinger equation
 - High-T properties \rightarrow investigated via lattice QCD

Several discoveries in almost 40 years...

Discovery of **anomalous J/ψ suppression**

Discovery of **sequential charmonium suppression**

} SPS



Several discoveries in almost 40 years...

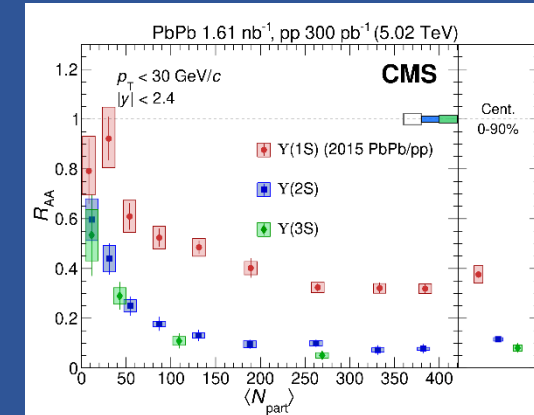
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Discovery of **sequential bottomonium suppression**

SPS

LHC



Several discoveries in almost 40 years...

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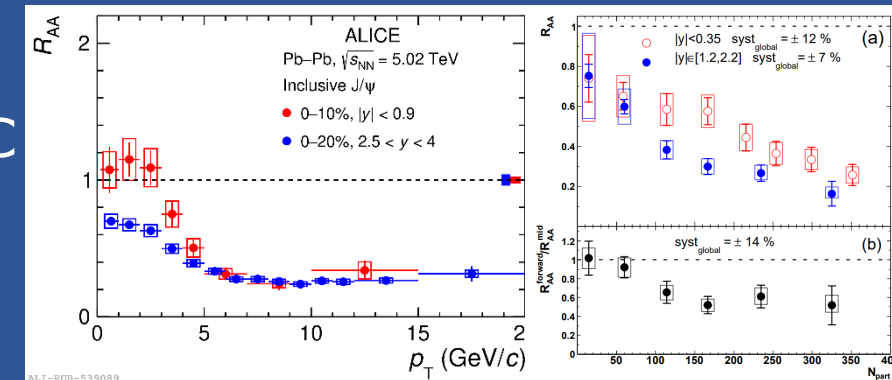
} SPS

Discovery of **sequential bottomonium suppression**

} LHC

Discovery of **J/ψ regeneration**

} RHIC
LHC



Several discoveries in almost 40 years...

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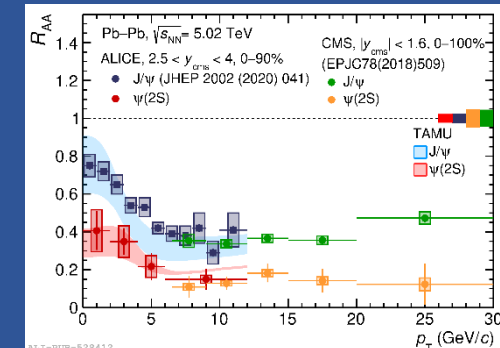
} LHC

Discovery of **J/ψ regeneration**

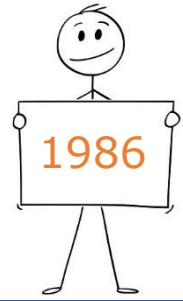
} RHIC
LHC

Observation of **sequential charmonium regeneration**

} LHC



The starting point



受入
86-9-102
高工研図書室

PHYS. LETT. B, in press

BROOKHAVEN NATIONAL LABORATORY

June 1986 BNL-38344

J/ψ SUPPRESSION BY QUARK-GLUON PLASMA FORMATION

T. Matsui
Center for Theoretical Physics
Laboratory for Nuclear Science
Massachusetts Institute of Technology
Cambridge, MA 02139, USA

and

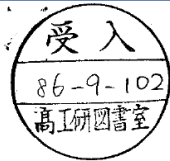
H. Satz
Fakultät für Physik
Universität Bielefeld, D-48 Bielefeld, F.R. Germany
and
Physics Department
Brookhaven National Laboratory, Upton, NY 11973, USA

ABSTRACT

If high energy heavy ion collisions lead to the formation of a hot quark-gluon plasma, then colour screening prevents $c\bar{c}$ binding in the deconfined interior of the interaction region. To study this effect, we compare the temperature dependence of the screening radius, as obtained from lattice QCD, with the J/ψ radius calculated in charmonium models. The feasibility to detect this effect clearly in the dilepton mass spectrum is examined. We conclude that J/ψ suppression in nuclear collisions should provide an unambiguous signature of quark-gluon plasma formation.

From **theory...**

The starting point



PHYS. LETT. B, in press

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If high energy heavy ion collisions lead to the formation of a hot quark-gluon plasma, then colour screening prevents $c\bar{c}$ binding in the deconfined interior of the interaction region. To study this effect, we compare the temperature dependence of the screening radius, as obtained from lattice QCD, with the J/ψ radius calculated in charmonium models. The feasibility to detect this effect clearly in the dilepton mass spectrum is examined. We conclude that J/ψ suppression in nuclear collisions should provide an unambiguous signature of quark-gluon plasma formation.

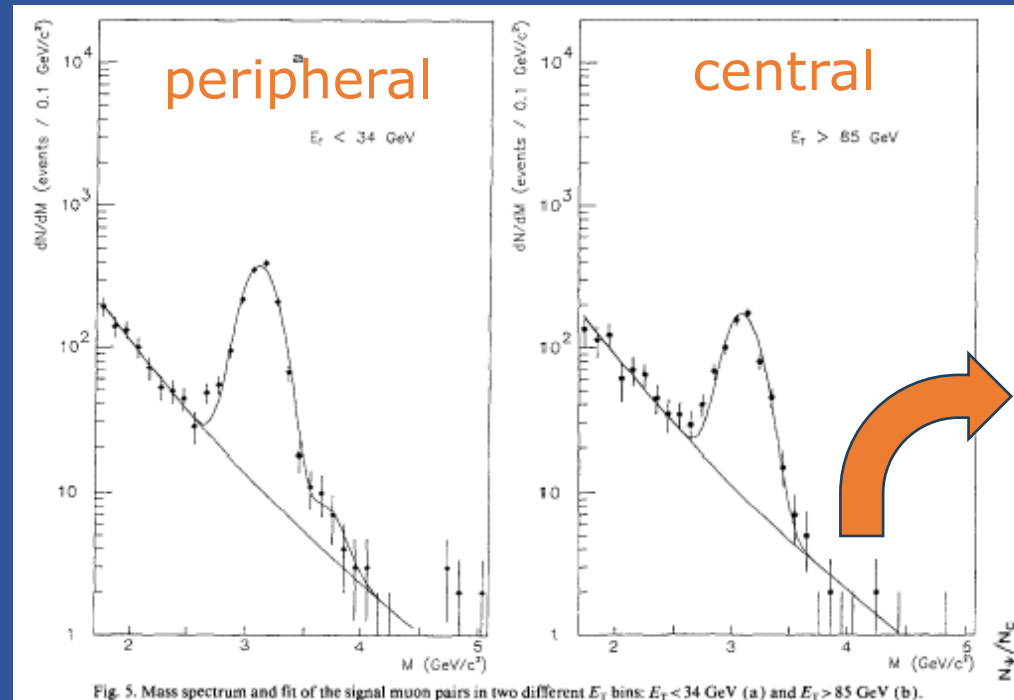
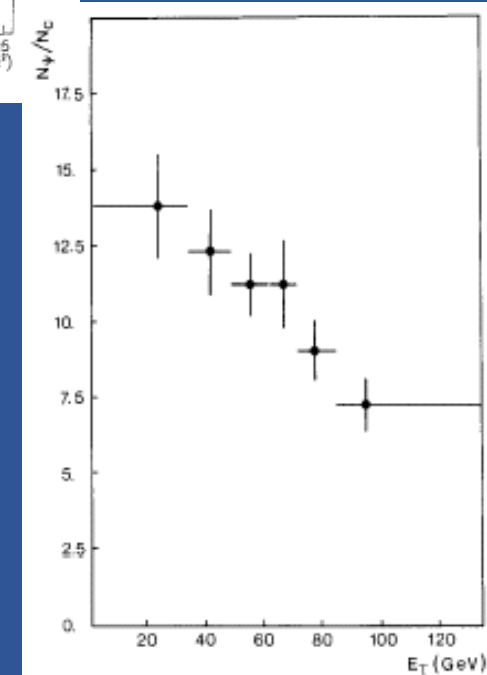


Fig. 5. Mass spectrum and fit of the signal muon pairs in two different E_T bins: $E_T < 34$ GeV (a) and $E_T > 85$ GeV (b).

NA38, Z. Phys. 38(1988) 117

O-U collisions at $\sqrt{s_{NN}} \sim 20$ GeV
 N_ψ/N_c decreases by a factor ~ 2
from peripheral to central collisions

Is also $\psi(2S)$ suppressed ?



From **theory...**

... to **experiment**



The basic argument

- At $T=0$ the cc interaction can be described by the potential

$$V(r) = -\frac{\alpha}{r} + kr$$

↙ Coulomb term
 $\alpha \sim 0.5$

↘ Confinement term

The basic argument

- At $T=0$ the cc interaction can be described by the potential
- At deconfinement the confinement term vanishes and the Coulomb term is screened

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$$V(r) = -\frac{\alpha}{r} e^{-r/\lambda_D(T)}$$

↘ Debye length

The basic argument

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- The energy of the system can be written as

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Coulomb term
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Debye length

$$E = \frac{\vec{p}^2}{2\mu} - \frac{\alpha}{r} e^{-r/\lambda_D}$$

Reduced mass = $m_{cc}/2 \sim 1.5 \text{ GeV}$

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- Bound state condition can be obtained by minimizing E ; one easily gets

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$$x(1+x)e^{-x} = \frac{1}{\alpha\mu\lambda_D}$$

$x = r/\lambda_D$

The basic argument

➤ At $T=0$ the cc interaction can be described by the potential

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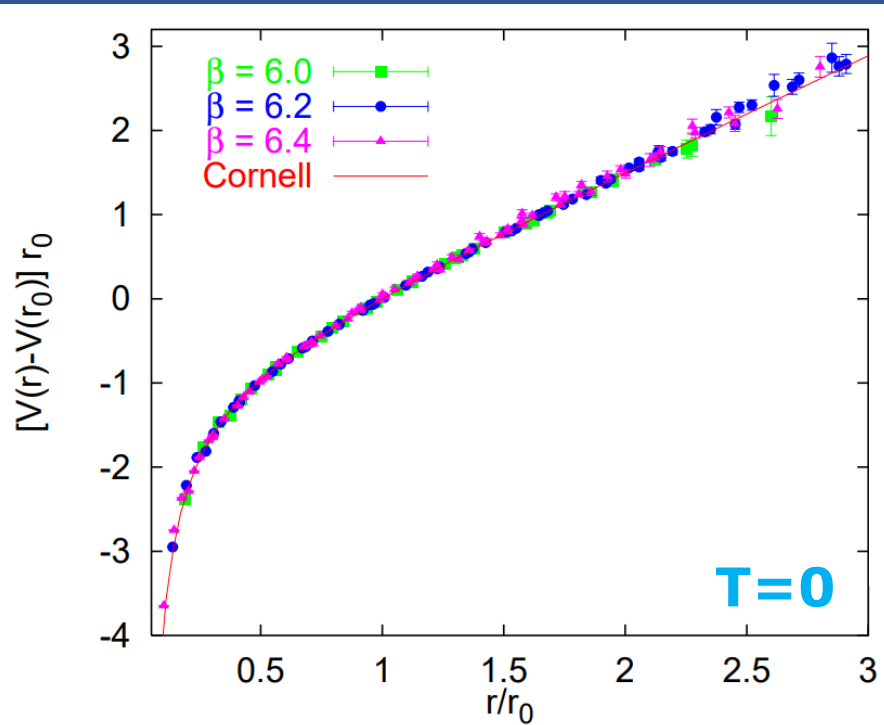
➤ The above relation does **not** admit solutions if $\frac{1}{0.84\alpha\mu} > \lambda_D$ i.e. $\lambda_D \lesssim 0.6 \text{ fm}$

Indications already in the '80s that $\lambda_D \propto T^{-1}$ and **melting conditions** met at $T \sim 200 \text{ MeV}$

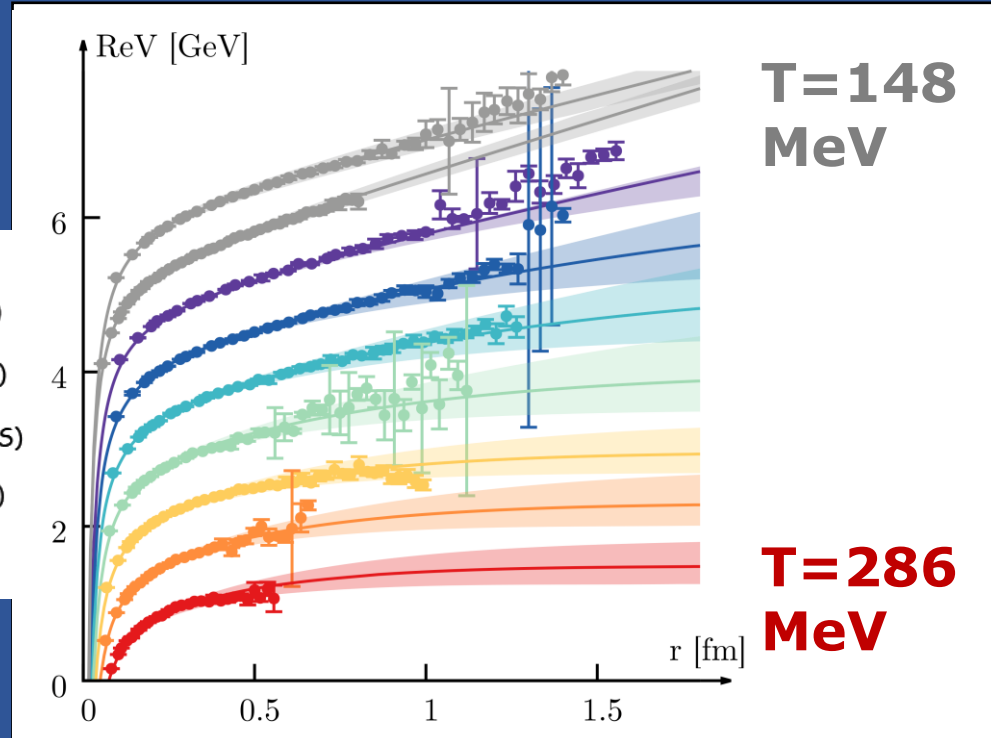
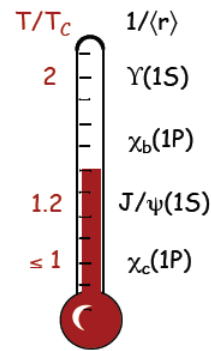
actual λ_D value to be obtained from lattice calculations

Lattice calculations and potentials

G.S. Bali, Phys. Rep. 343 (2001) 1-136



Lafferty and Rothkopf, Phys. Rev. D 101 (2020) 056010



Potential models provide a faithful reproduction of available lattice data

- Gradual transition **from a Cornell to a Debye-screened behaviour** for the (real part of) the potential → **color screening** in a deconfined medium
- Potential also has a finite imaginary part (not shown)
→ decaying of quark-antiquark correlation due to gluonic damping in the plasma

Hot debates

From the QM87 summary talk

The most provocative observation, reported by NA 38 [13], was that J/ψ production seems to be suppressed by $\sim 30\%$ in high E_T events. The second provocative

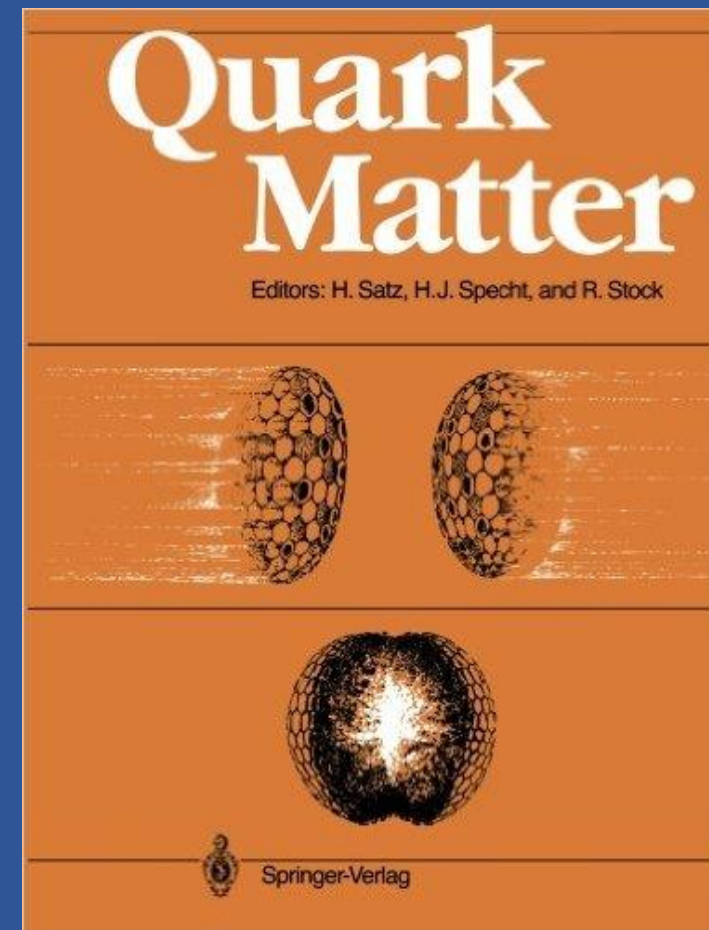
3 Puzzles

3.1 J/ψ suppression

$$N_{\psi}/N_c = \begin{cases} 9.3 \pm 0.6 & \text{for } E_T < 28 \text{ GeV} \\ 5.9 \pm 0.4 & \text{for } E_T > 50 \text{ GeV.} \end{cases} \quad (10)$$

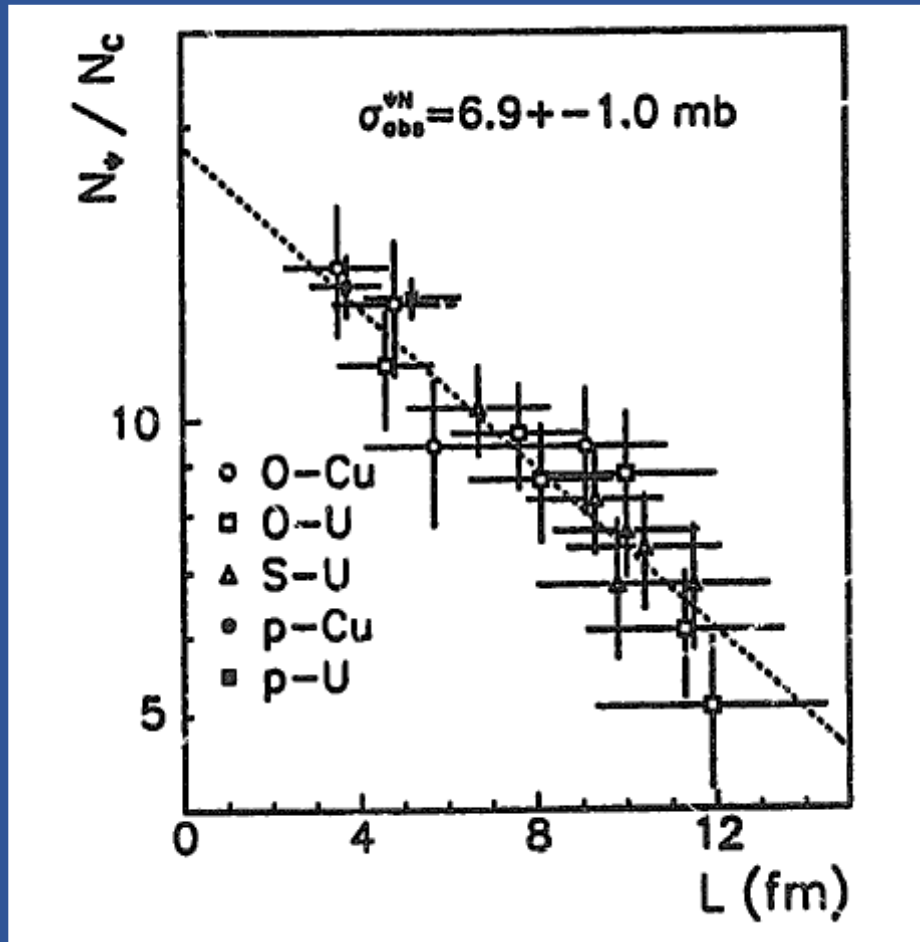
This 30% reduction of ψ production caused the most controversy at Quark Matter '87.

There are naturally several caveats that need further consideration. First, there is the problem of prov-



- Can **competing sources** of J/ψ dissociation involving hadronic interactions with cold nuclear matter and/or hadronic medium reproduce the observations ?

Where is the Quark-Gluon Plasma ?



C. Gerschel et al., PLB207 (1988)253

Nuclear Physics A544 (1992) 513c-516c
North-Holland, Amsterdam

NUCLEAR
PHYSICS A

**Comparison of J/ψ -Suppression in Photon,
Hadron and Nucleus-Nucleus Collisions :
Where is the Quark-Gluon Plasma ?**

C. Gerschel^a and J. Hüfner^b

p-A collision results
imply **significant
dissociation cross
sections in cold
nuclear matter (CNM)**

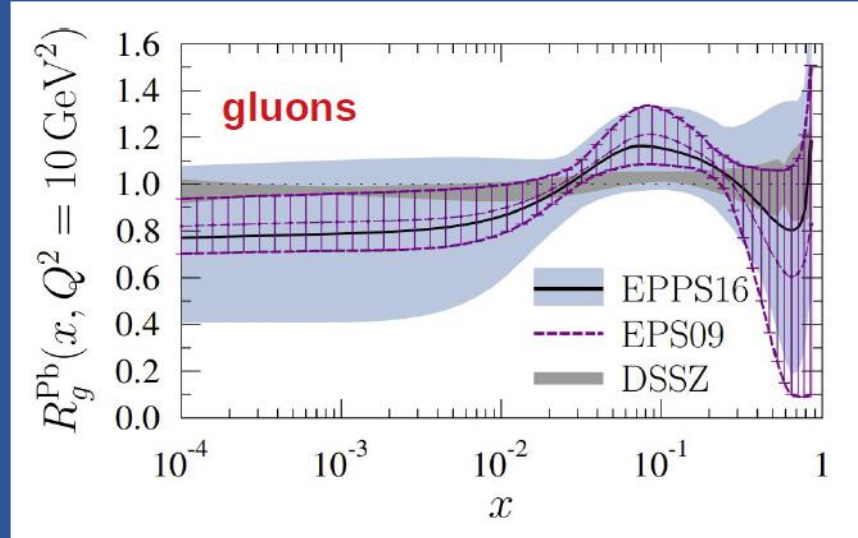
- Crucial ingredient in the interpretation of the data
- Stimulated an intense experimental program of p-A studies at both CERN and FNAL

Quantifying non-QGP effects

- BOTH initial and final state non-QGP effects may lead to a decreased charmonium production
- The relative size depends quite a lot on collision energy (keep in mind for later)

SPS

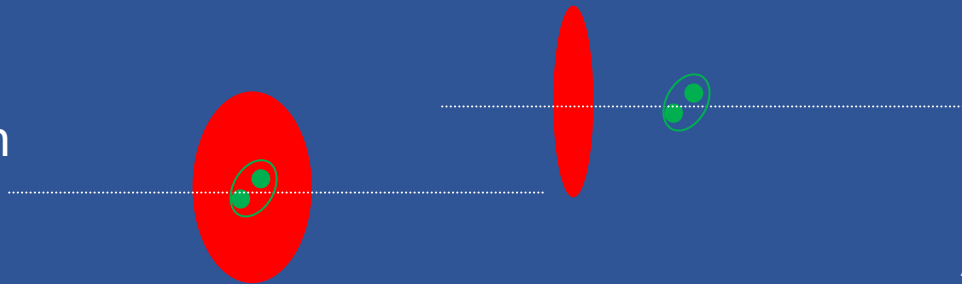
Initial state effects:
moderate anti-shadowing
 $x \sim 10^{-1}$ ($y=0$)



LHC

Initial state effects:
shadowing
 $x \sim 10^{-5}$ ($y \sim 3$)
 $x \sim 10^{-3}$ ($y=0$)
 $x \sim 10^{-2}$ ($y \sim -3$)

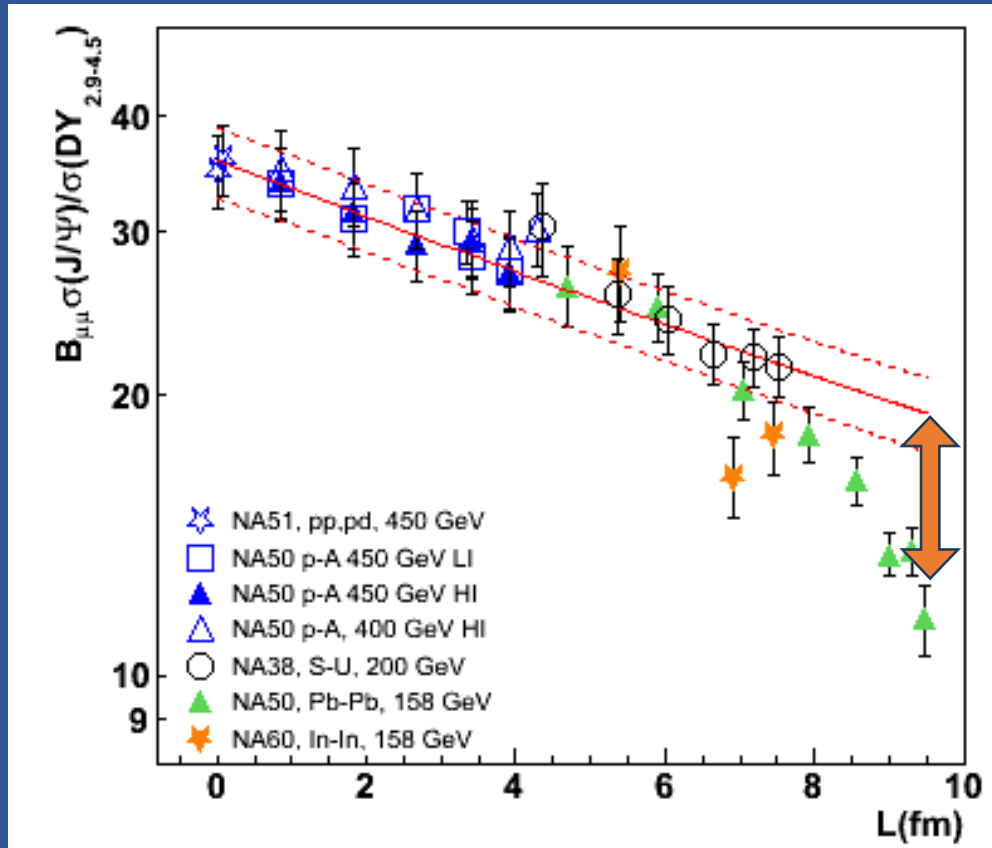
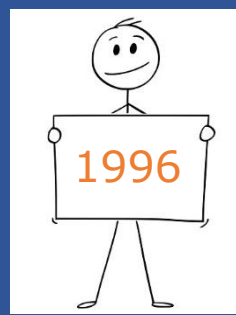
(Final state) CNM effects:
break-up in nuclear matter can
be sizeable
 $\tau = L/(\beta_z \gamma) \sim 0.5 \text{ fm/c}$ ($y=0$)



(Final state) CNM effects:
negligible, extremely short
crossing time
 $\tau = L/(\beta_z \gamma) \sim 7 \cdot 10^{-5} \text{ fm/c}$ ($y \sim 3$)
 $\tau = L/(\beta_z \gamma) \sim 4 \cdot 10^{-2} \text{ fm/c}$ ($y \sim -3$)

The “anomalous” suppression

NA50, Phys.Lett.B 410 (1997) 337-343



Central Pb-Pb collisions clearly depart from the behaviour of p-A and S-U data, where only CNM is assumed to be present

J/ψ suppression in Pb Pb collisions: a hint of quark-gluon plasma production?

J.P. Blaizot

Anomalous suppression

Therefore, it is interesting to speculate whether the new phase of strong absorption may be the quark-gluon plasma.

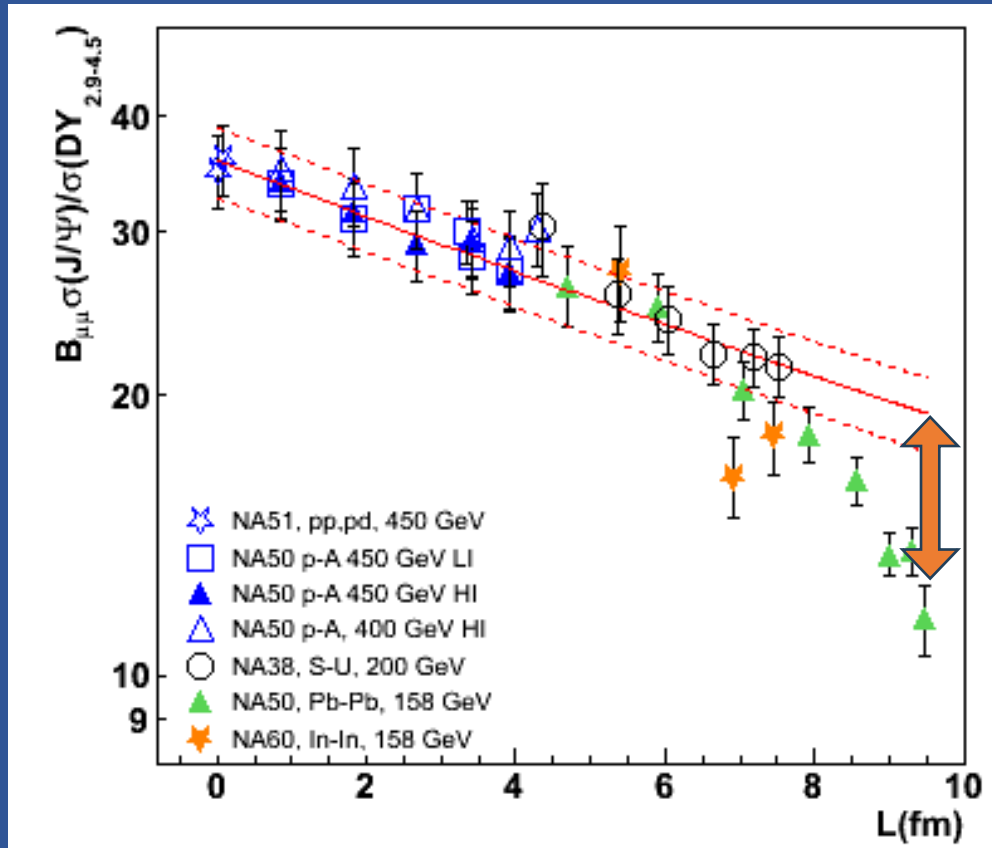
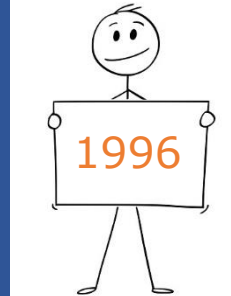
C.Y. Wong

Is this suppression really anomalous? Can we conclude that the quark-gluon plasma is already discovered?

D. Kharzeev

The “anomalous” suppression

NA50, Phys.Lett.B 410 (1997) 337-343

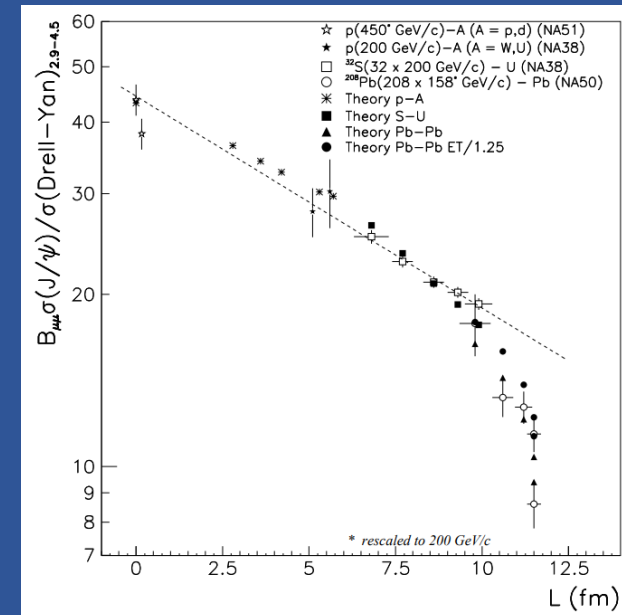


Central Pb-Pb collisions clearly depart from the behaviour of p-A and S-U data, where only CNM is assumed to be present

Anomalous suppression

In conclusion, combining nuclear absorption and final state interaction with co-moving hadrons, we have obtained a reasonable description of the J/ψ and ψ' data.

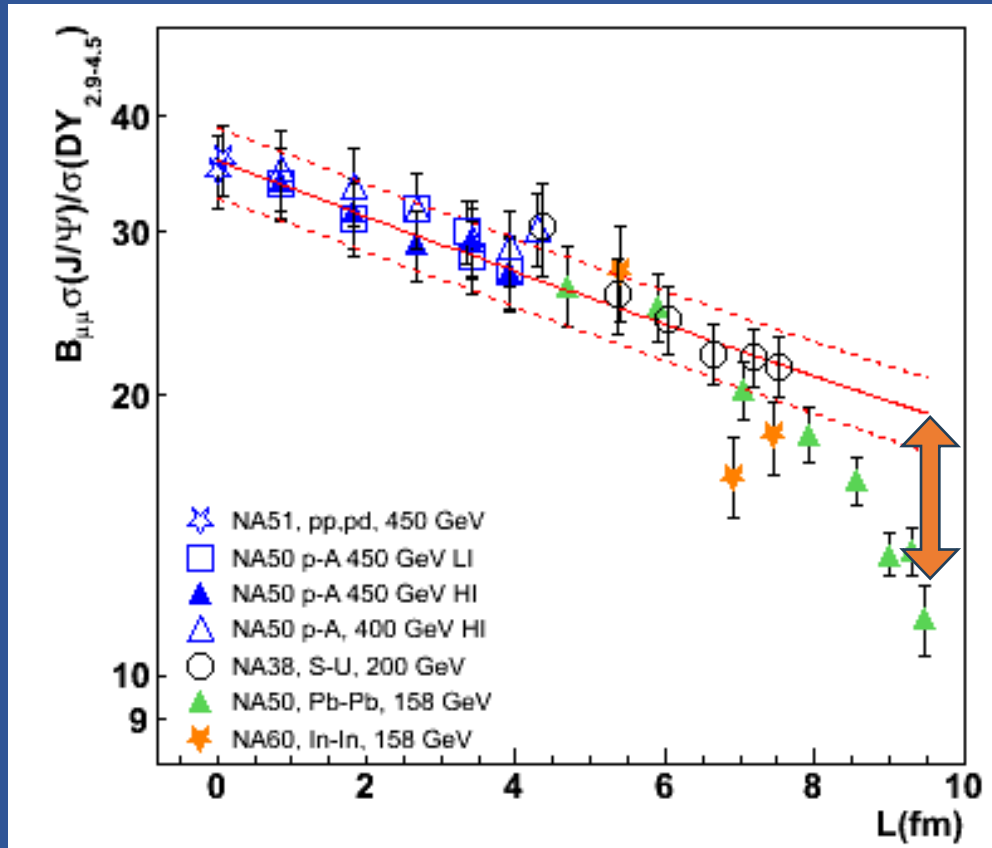
A. Capella



Assuming
 $\sigma_{\psi-co} = 0.6 \text{ mb}$

The “anomalous” suppression

NA50, Phys.Lett.B 410 (1997) 337-343

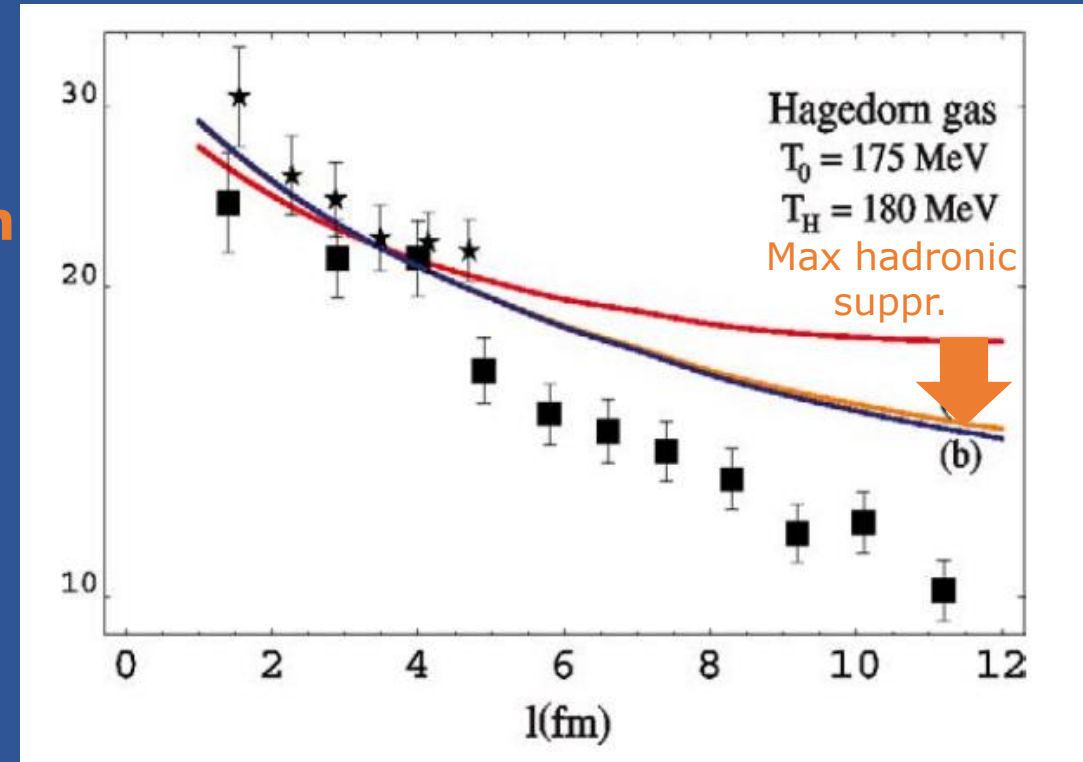


Central Pb-Pb collisions clearly depart from the behaviour of p-A and S-U data, where only CNM is assumed to be present

Anomalous suppression

The dissociation curve cannot become harder, and the prediction falls short from explaining the drop observed by NA50

L. Maiani

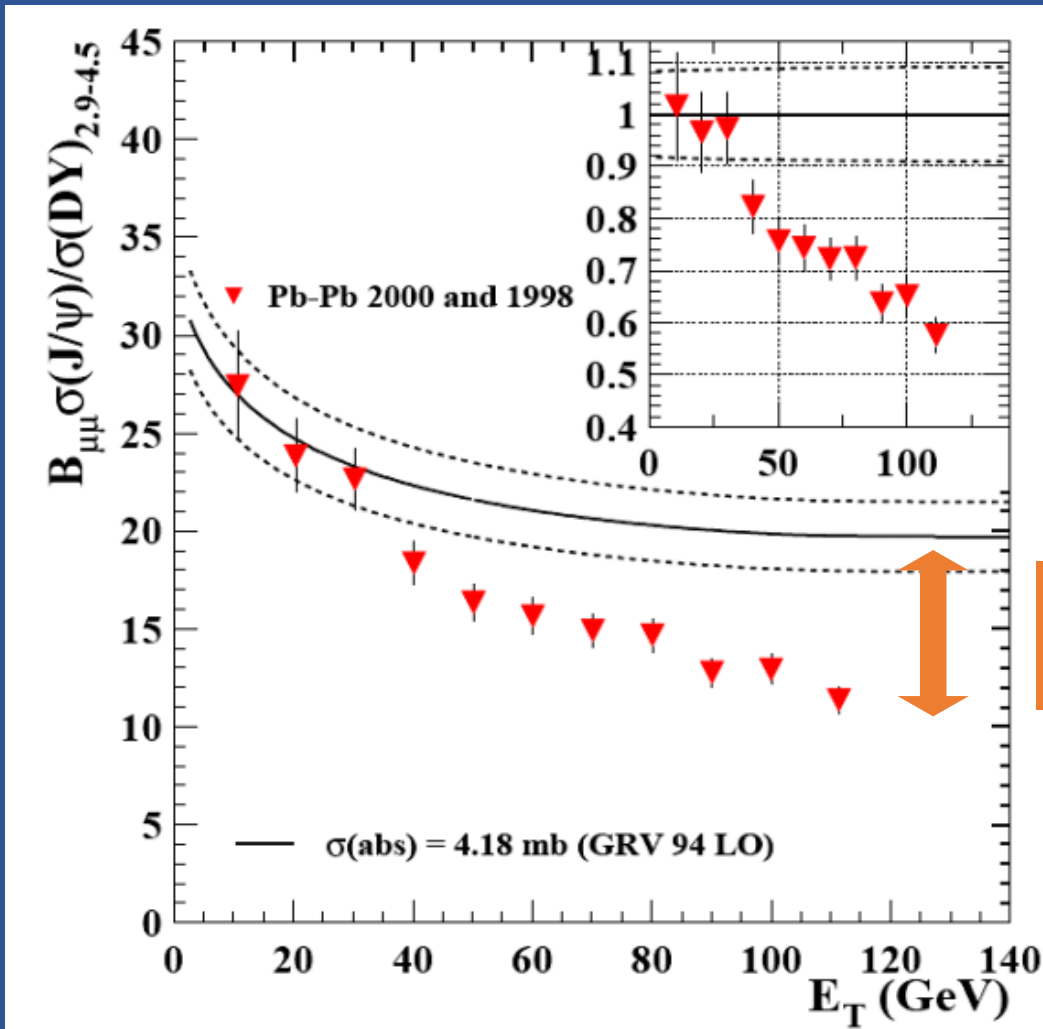


Maiani et al., NP A 748 (2005) 209-225



Feed-down

J.P. Lansberg,
Phys.Rept. 889 (2020) 1

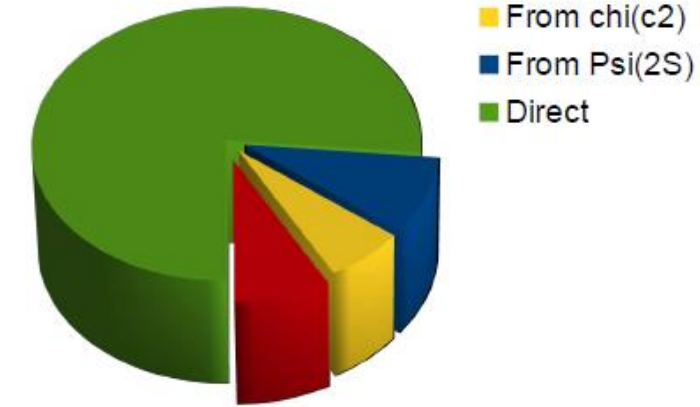


Qualitatively consistent
with the contribution
of feed-down decays
from $\psi(2S)$ and χ_c

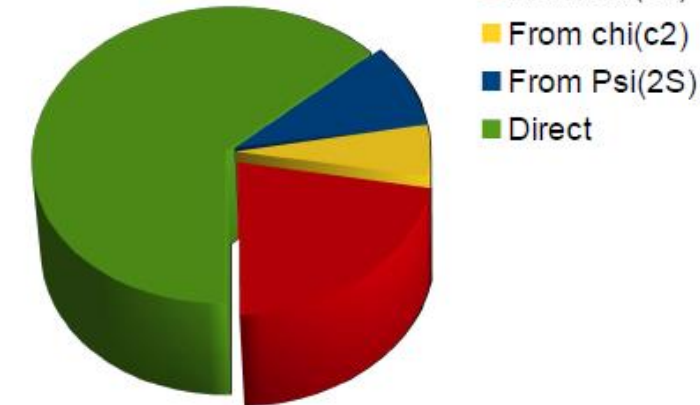


a 30-40% 'anomalous'
suppression effect

Low transverse momentum



High transverse momentum



$\Gamma_{10} \quad J/\psi(1S) \text{ anything}$

$(61.4 \pm 0.6) \%$

← $\psi(2S)$

$\Gamma_{93} \quad \gamma J/\psi(1S)$

$(1.40 \pm 0.05) \%$

← χ_{c0}

$\Gamma_{93} \quad \gamma J/\psi(1S)$

$(34.3 \pm 1.0) \%$

← χ_{c1}

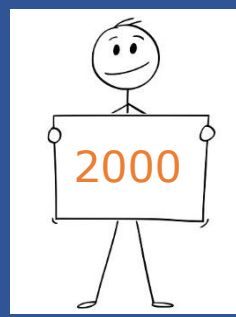
$\Gamma_{93} \quad \gamma J/\psi(1S)$

$(19.0 \pm 0.5) \%$

← χ_{c2}

NA50, Eur.Phys.J.C 39 (2005) 335

Enter RHIC, big expectations

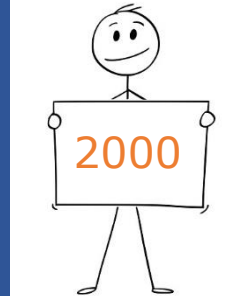


J/ψ

- New mechanism for "nuclear absorption" (J. Qiu)
- Role of transverse energy fluctuations (M. Dinh, A. Capella)
- J/ψ enhancement at RHIC! (R.L. Thews)

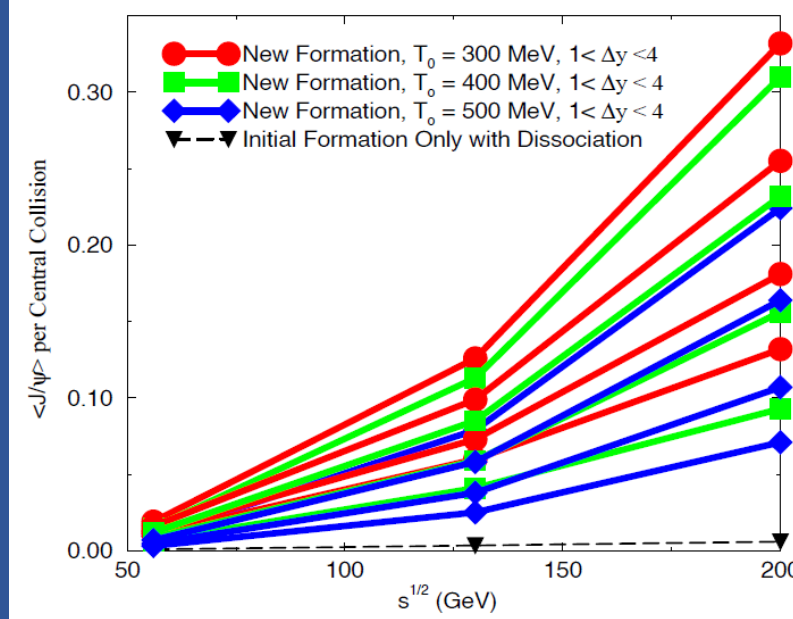
J.P. Blaizot, **QM2001** summary talk

Enter RHIC, big expectations



J/ψ

- New mechanism for "nuclear absorption" (J. Qiu)
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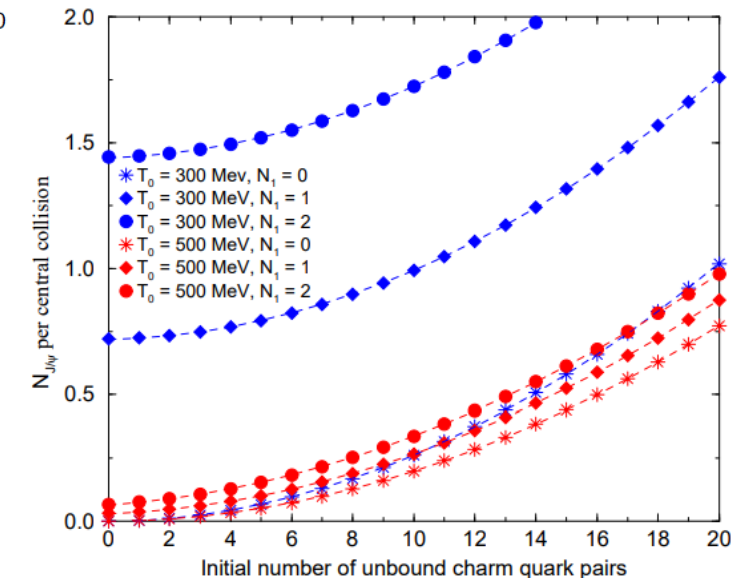


Strong dependence on collision energy

Quadratic dependence on the initial formation rate

PHYSICAL REVIEW C, VOLUME 63, 054905
Enhanced J/ψ production in deconfined quark matter
Robert L. Thews, Martin Schroedter, and Johann Rafelski
Department of Physics, University of Arizona, Tucson, Arizona 85721
(Received 29 August 2000; published 23 April 2001)

Experimental observation of such enhanced production would provide evidence for deconfinement unlikely to be compatible with competing scenarios.



J.P. Blaizot, **QM2001** summary talk

Experiment follows theory

T. Frawley, **QM2002**

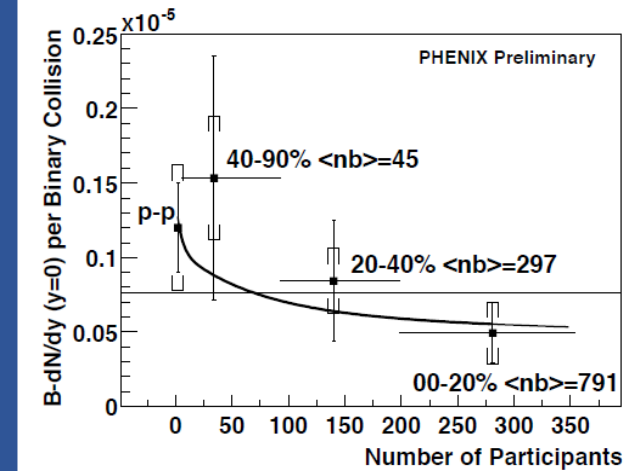
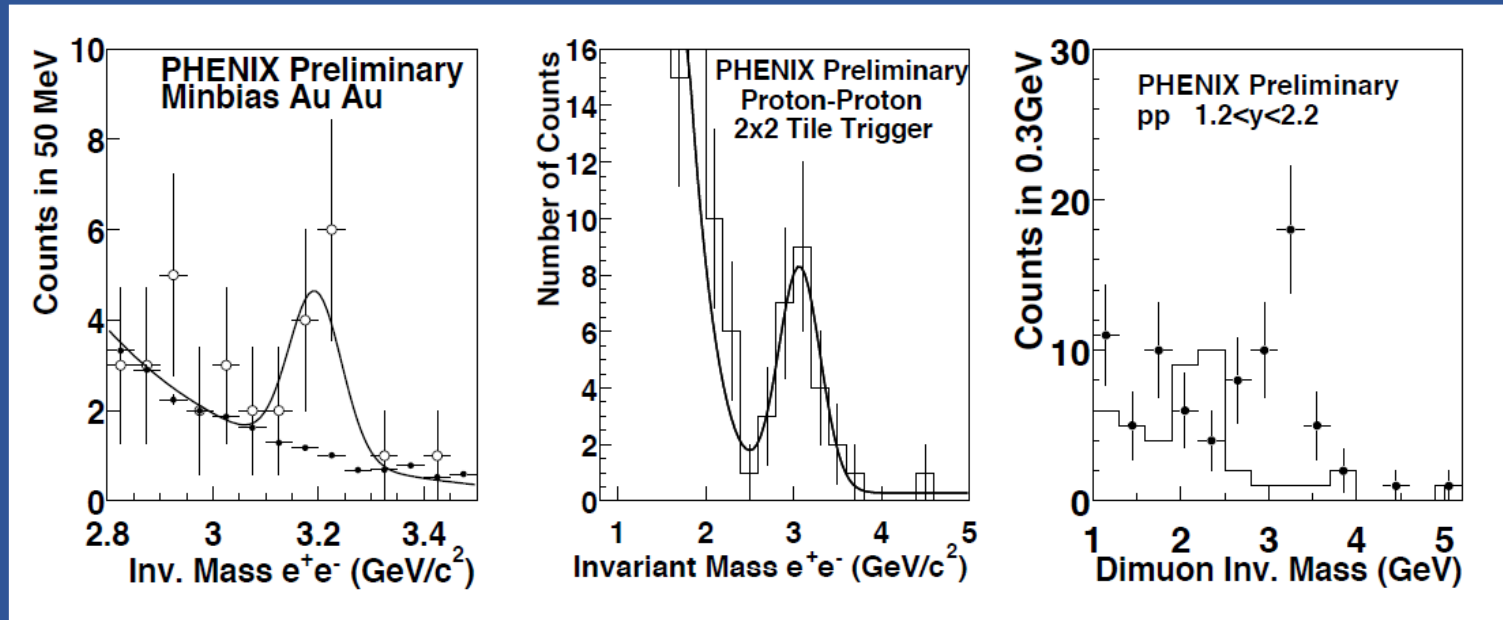
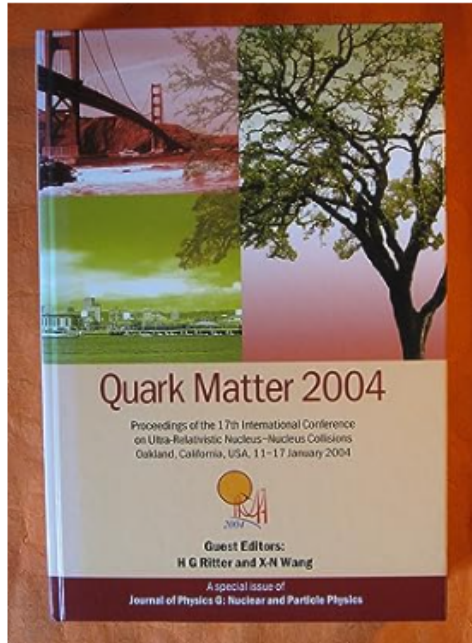


Figure 5. $J/\psi \rightarrow ee$ branching ratio times dN/dy scaled by N_{binary} . The flat line is the best fit binary scaling value. The curve is a normal nuclear absorption model calculation [8].

- Collider vs fixed-target: integrating the same luminosity can be tough
 - Fixed target SPS (**2000**): 35 days, $1.2 \cdot 10^7$ Pb/s on a 4mm Pb target $\rightarrow L_{int} \sim 83 \text{ nb}^{-1}$
 - **Total** RHIC Au-Au $\sqrt{s_{NN}}=200 \text{ GeV} \rightarrow L_{int} \sim 130 \text{ nb}^{-1}$

Intermezzo: our literature can be valuable



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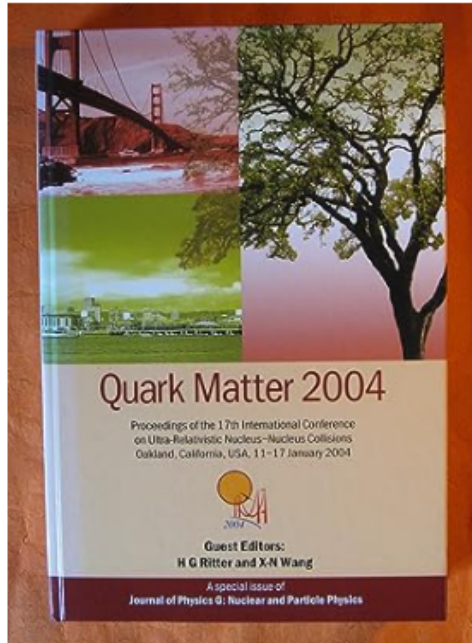
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Intermezzo: our literature can be valuable



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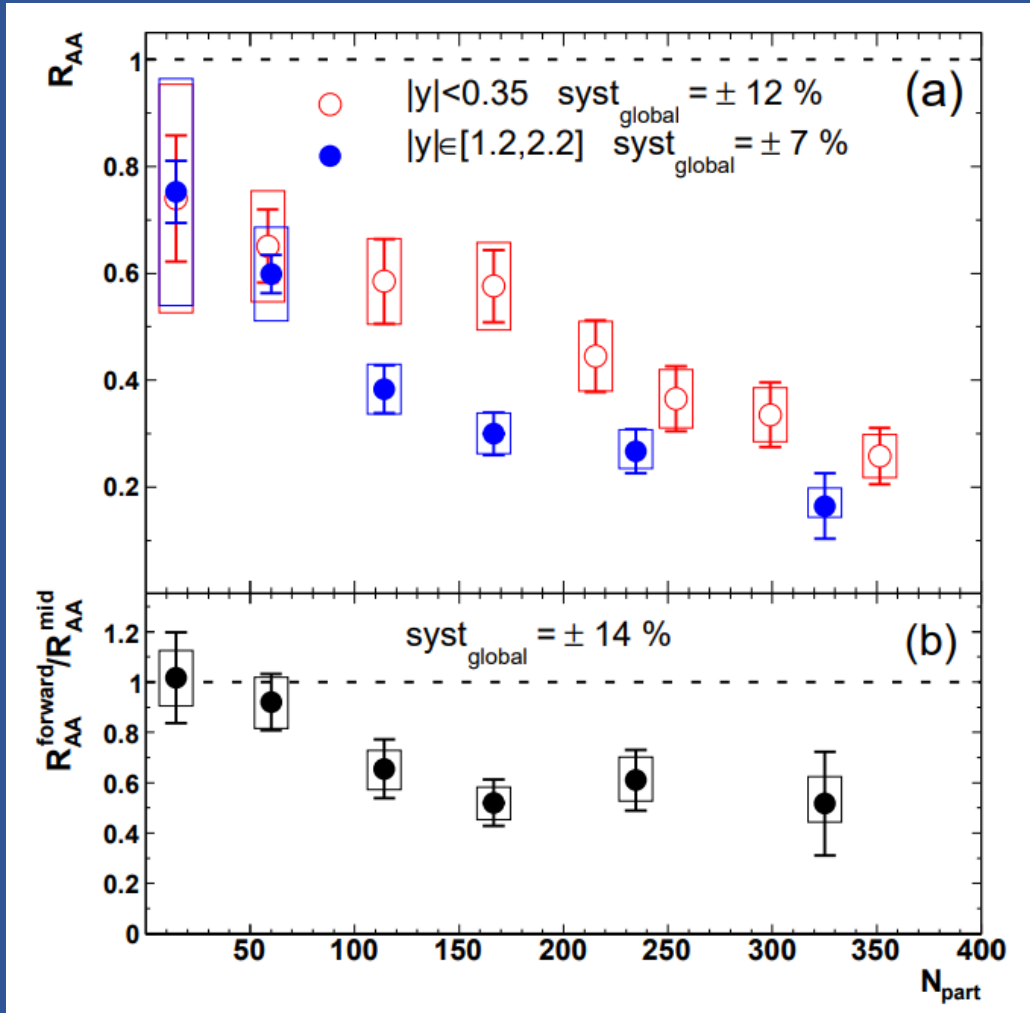
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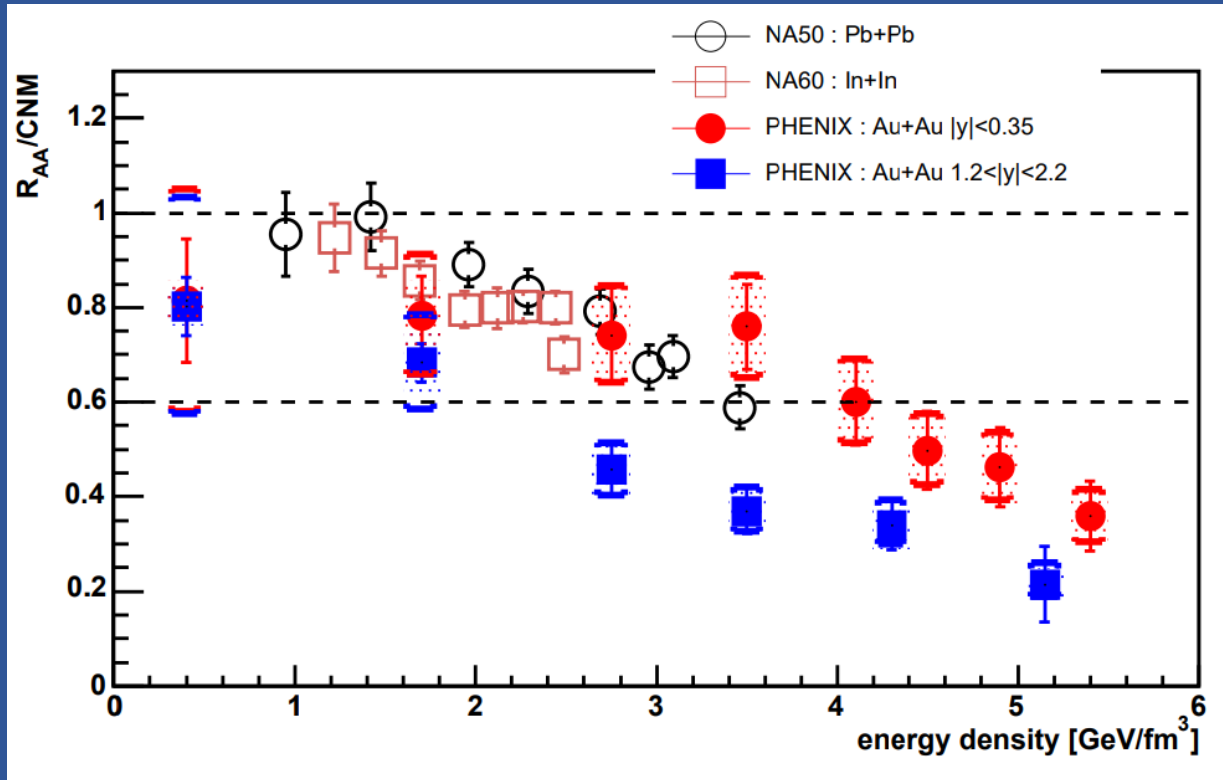
A surprising result



PHENIX, Phys.Rev.Lett.98:232002,2007

- For the first time measurements were performed in **two distinct rapidity ranges**
 - central vs forward y
- Drell-Yan reference becomes unpractical, as high mass dileptons dominated by open HF
 - Use **R_{AA}**
- Clear hierarchy: J/ψ more suppressed at forward y
 - Difficult to reconcile with the “traditional” suppression scenario
 - Can regeneration effects explain the observations ?

Comparing RHIC and SPS results



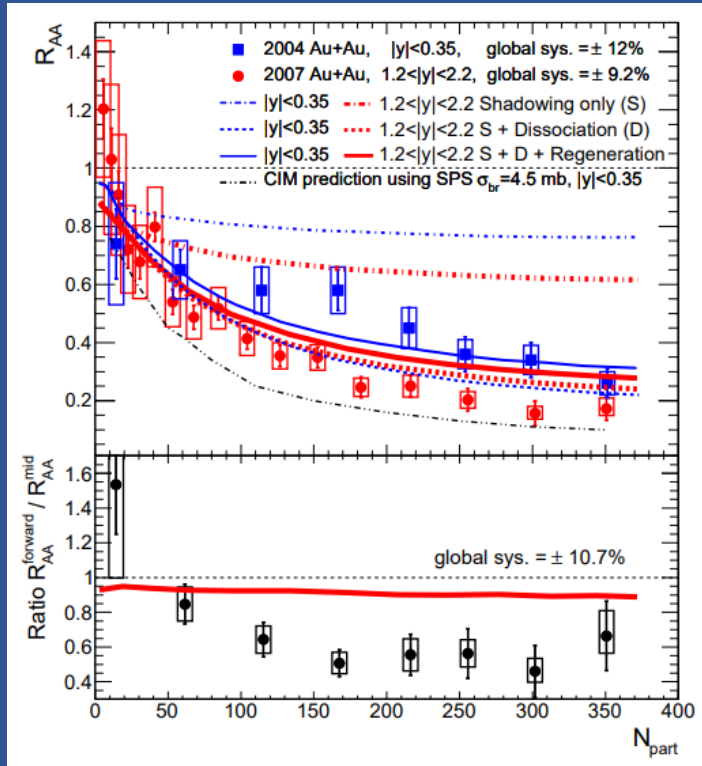
PHENIX, J.Phys.G34:S749-752,2007

- Caveat(s)
 - Assumptions in energy density calculations
 - Approximations in evaluation of CNM effects

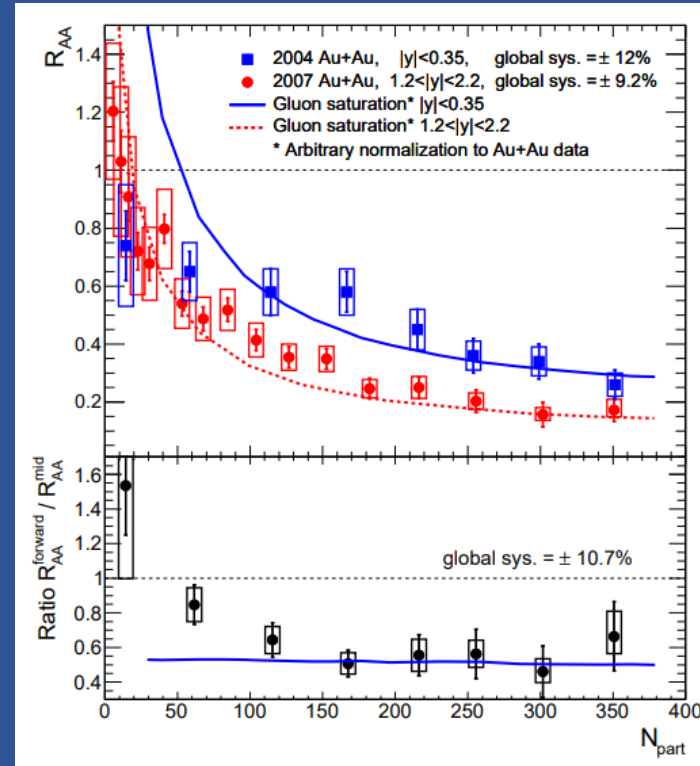
- RHIC forward-y suppression larger than SPS midrapidity
- RHIC and SPS midrapidity suppression quite similar
- Is J/ψ suppression stronger at RHIC but regeneration becomes important at midrapidity ?
- Is the similarity of RHIC and SPS midrapidity results a coincidence ?

Detailed theory comparisons needed!

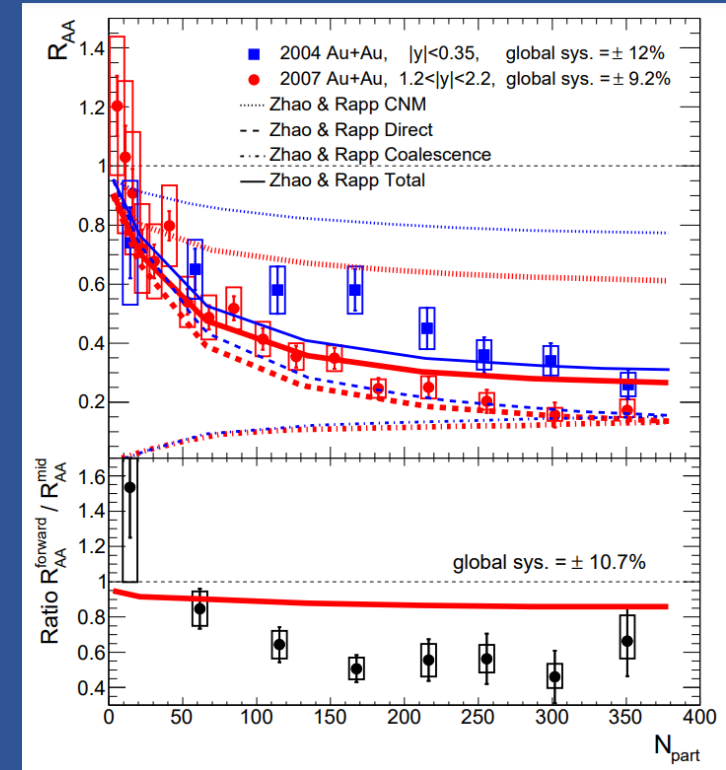
Suppression vs regeneration



Comover interaction model



Gluon saturation model



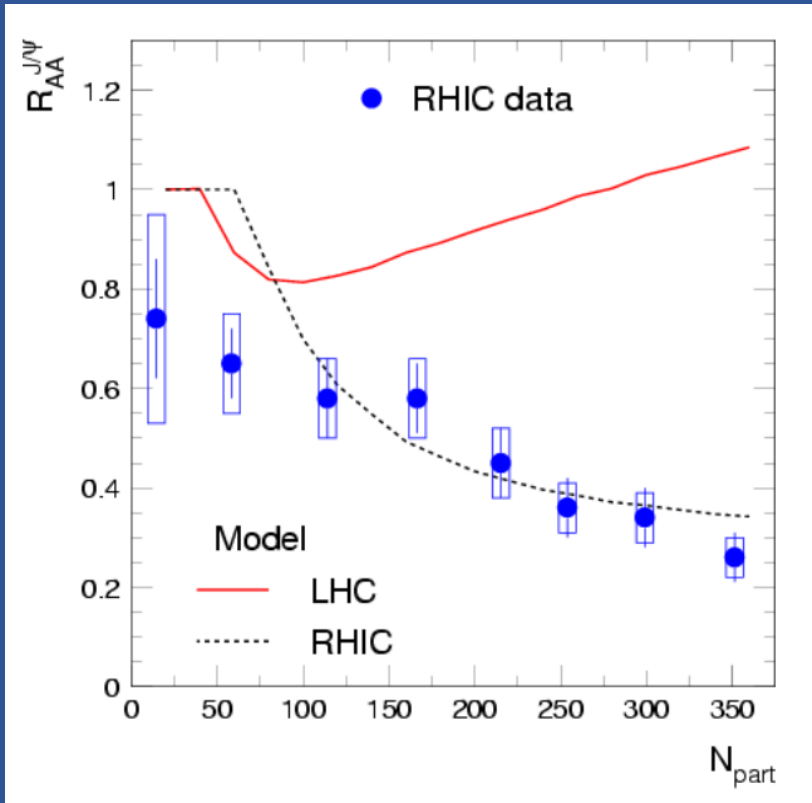
QGP/Hadron gas model

Combination of CNM and hot matter effects represents a tough task for phenomenology

We have presented new and more precise measurements of J/ψ nuclear modification at forward rapidity in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV. The results confirm our earlier published findings of a larger suppression at forward compared with midrapidity. This, combined with the similar suppression of J/ψ at midrapidity between RHIC and lower energy measurements, remains an outstanding puzzle in terms of a full theoretical description.

Moving to LHC energy

- Regeneration is an evidence of deconfinement, indications from RHIC experiments, but LHC represents the ideal testing ground → large charm quark multiplicity

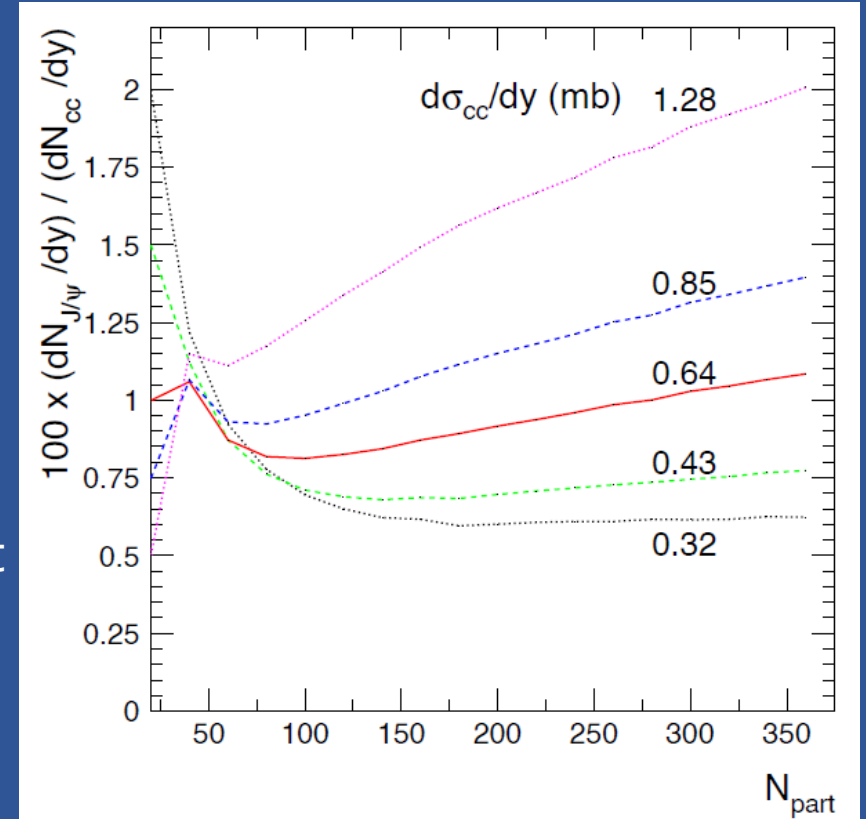


A. Andronic et al., Nucl. Phys. A789:334-356, 2007
P. Braun-Munzinger and J. Stachel, PLB490 (2000) 196

Predictions based on the statistical hadronization model



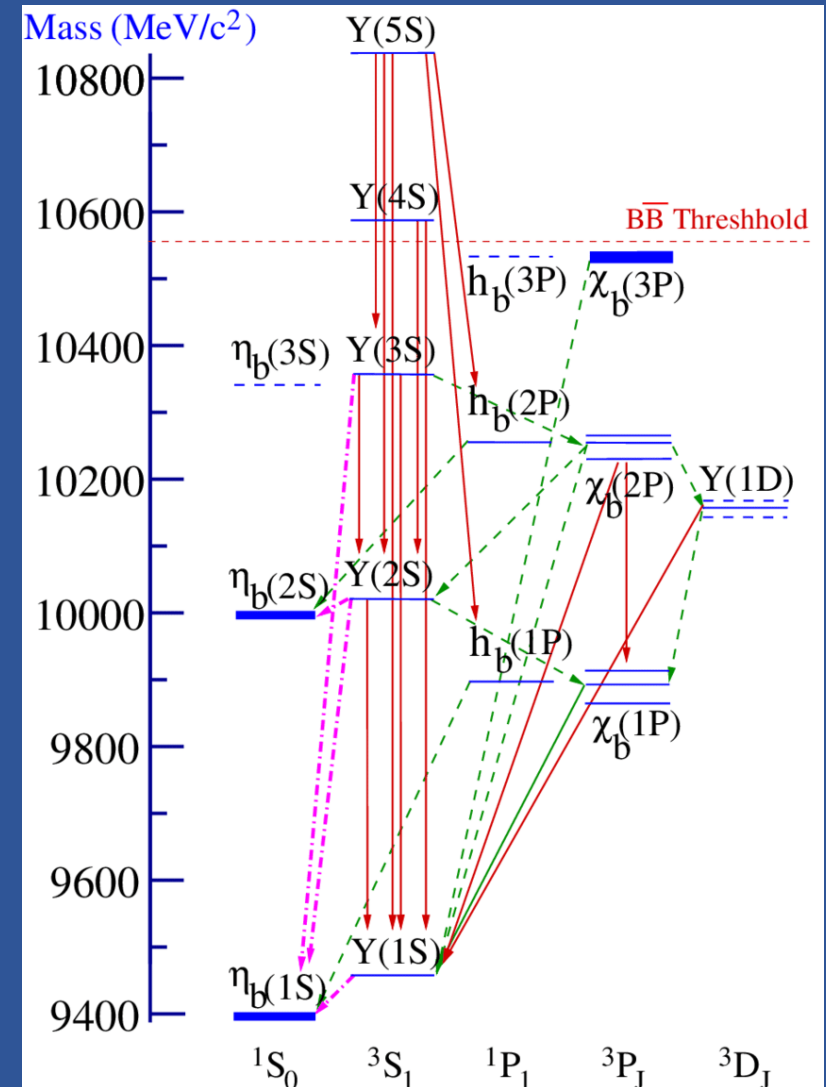
- Charm produced in primary hard collisions
- Total number stays constant
- Thermal equilibration of charm in the QGP
- Quarkonia are **generated** (not regenerated) at T_c



Very strong dependence
on charm cross section!

...but let's first examine bottomonium

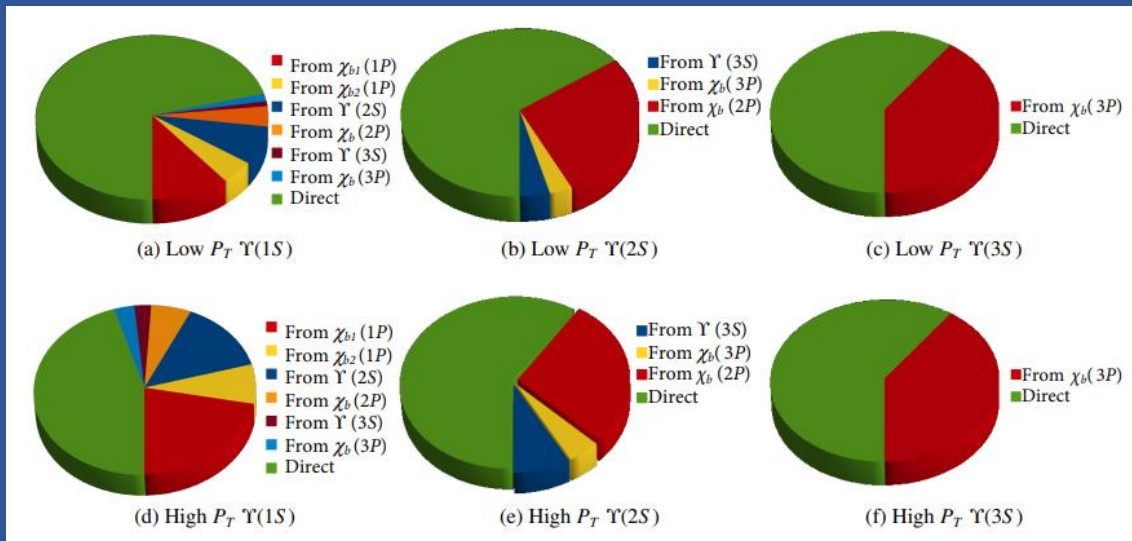
- Below open beauty threshold
 - **Three vector states**, with small but non negligible BR to dileptons ($\sim 2\%$) with extremely different binding energies



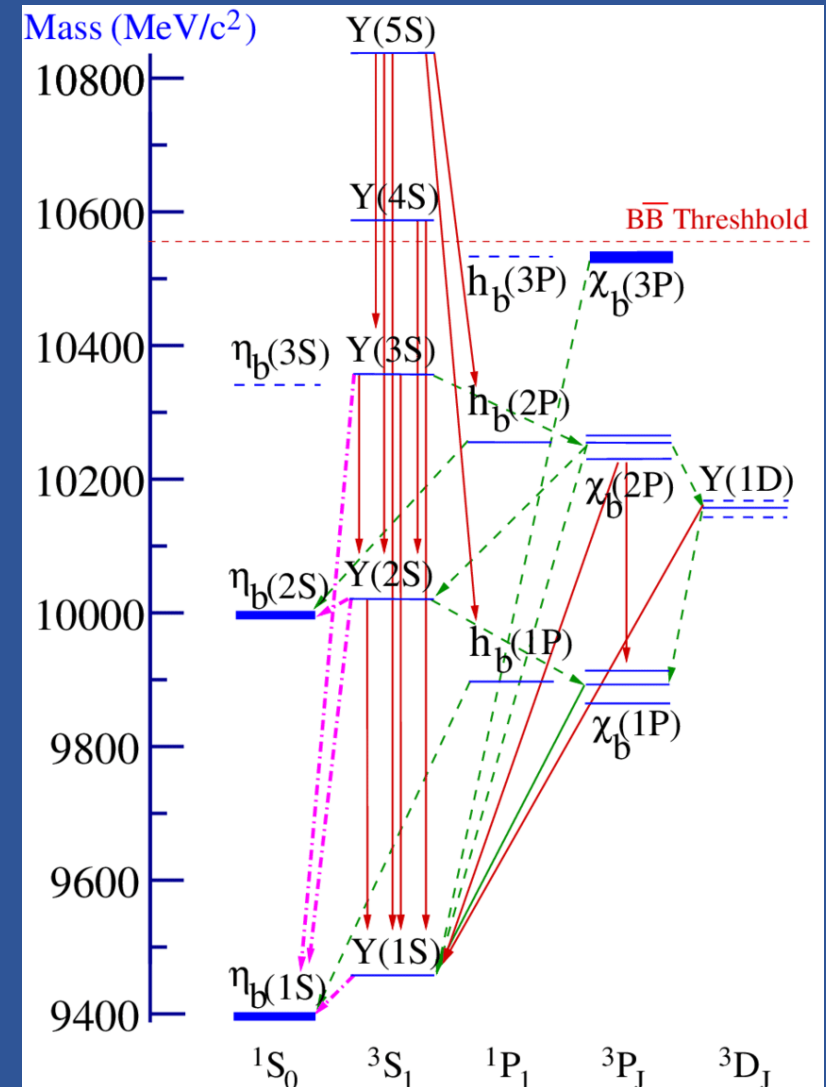
from J.P. Lansberg, Phys. Rep. 889(2020) 1

...but let's first examine bottomonium

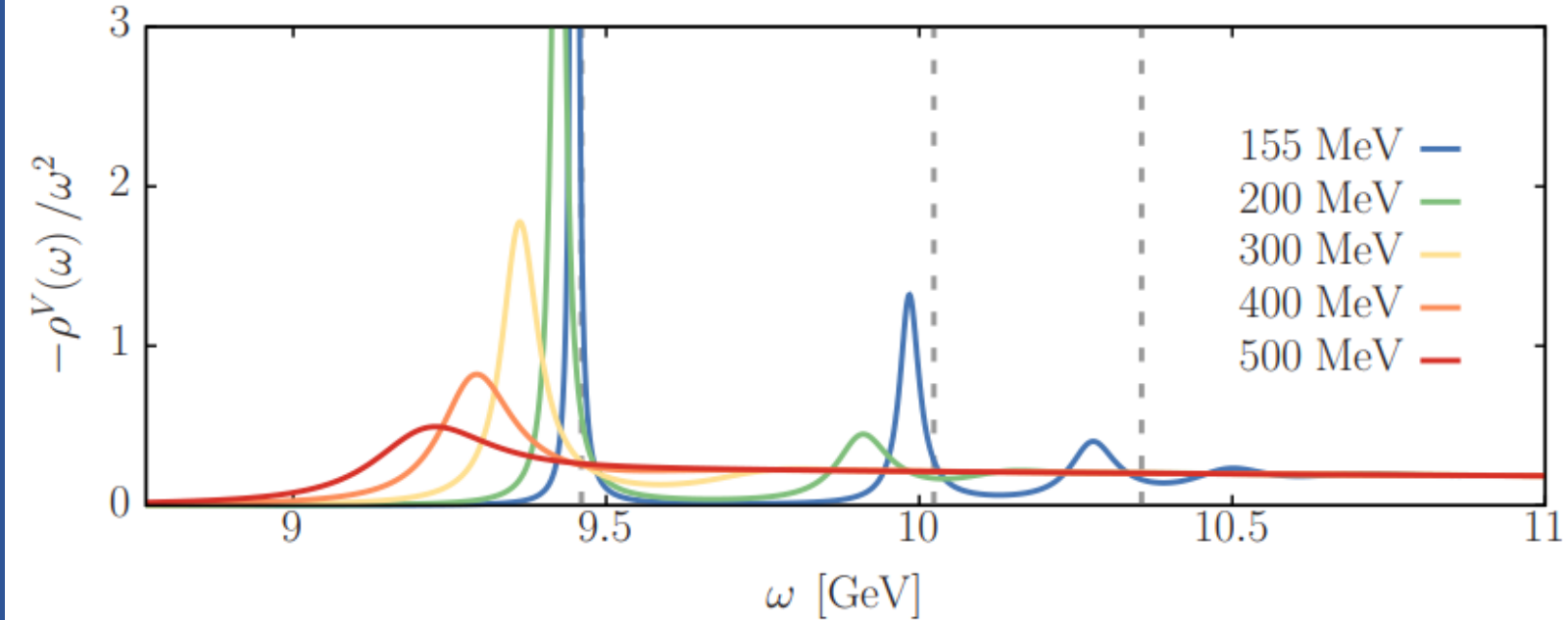
- Below open beauty threshold
 - **Three vector states**, with small but non negligible BR to dileptons ($\sim 2\%$) with extremely different binding energies
 - Rather complicate feed-down structure, to be considered in the interpretation of the results



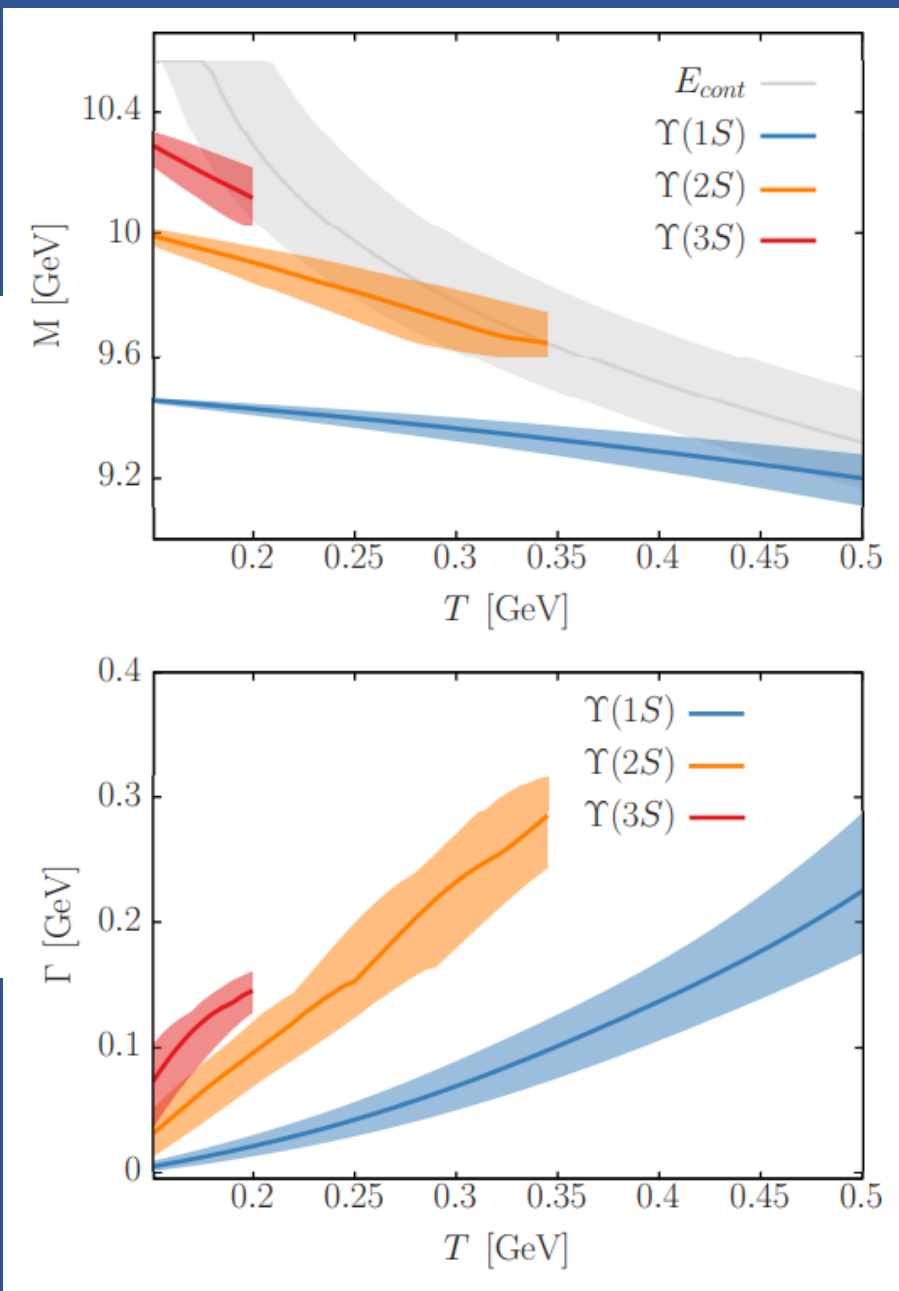
from J.P. Lansberg, Phys. Rep. 889(2020) 1



Modifications of bottomonium properties at $T > 0$

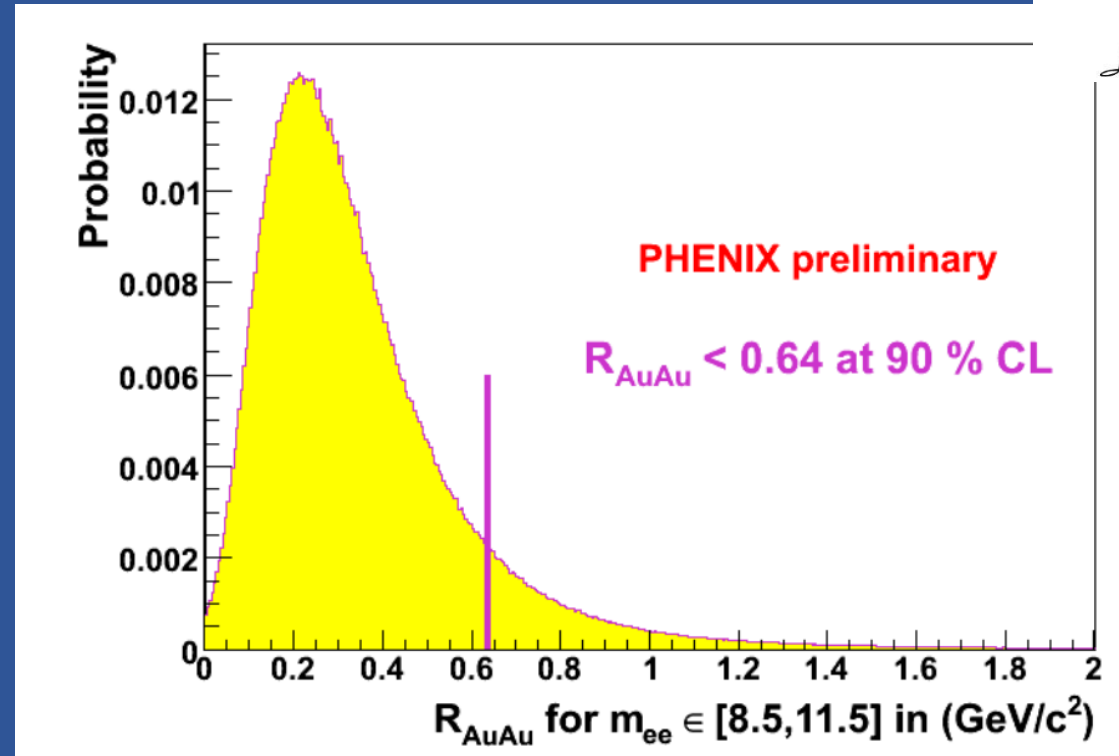
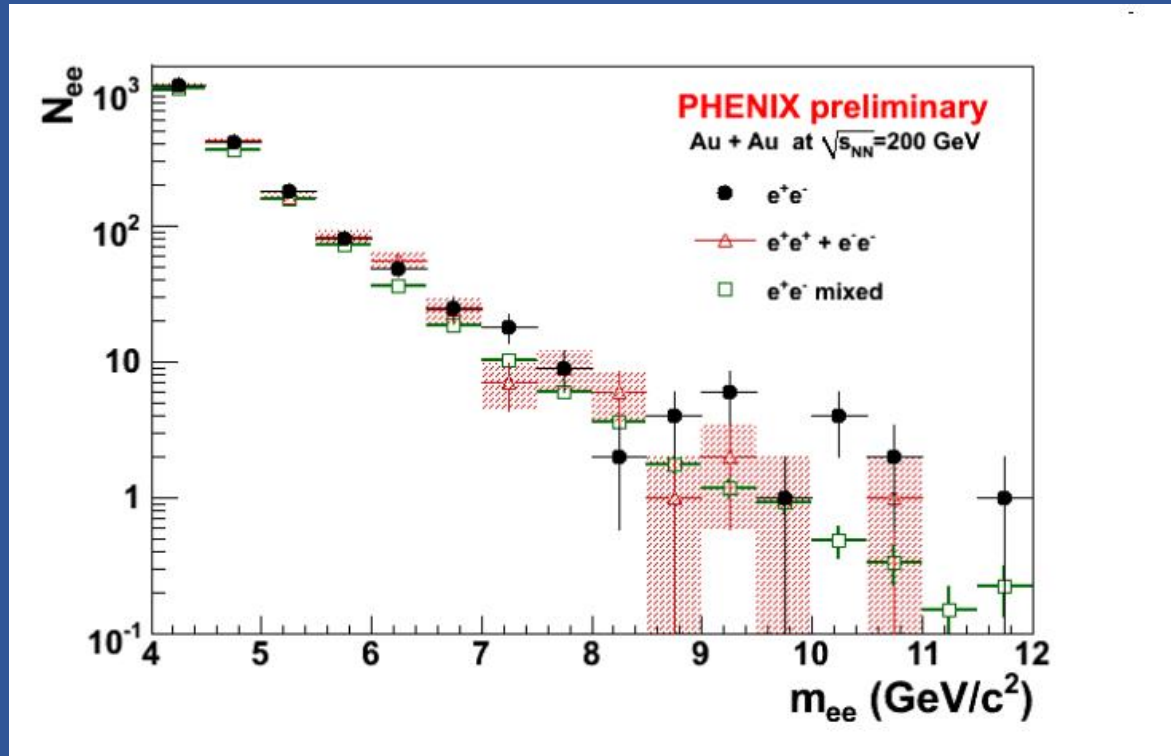
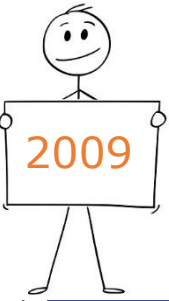


- In-medium peak remnants become “washed out” at **largely different temperatures** (gray line denotes the onset of the continuum)
→ Expect corresponding phenomenology in data



Lafferty and Rothkopf, Phys. Rev. D 101, 056010 (2020)

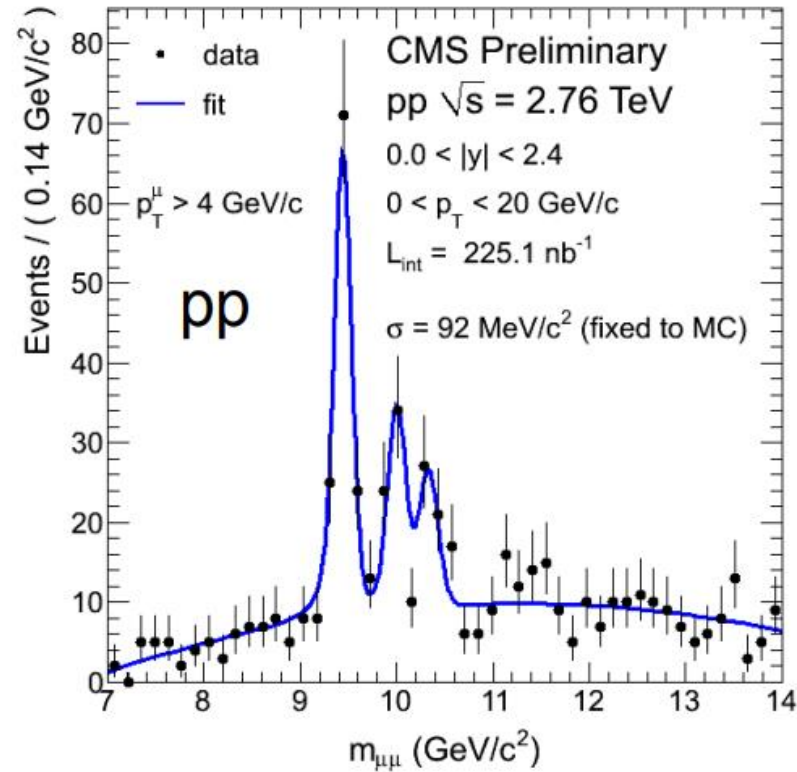
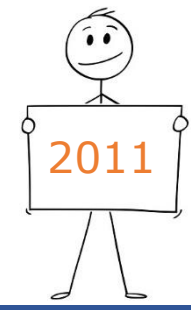
First attempts before LHC



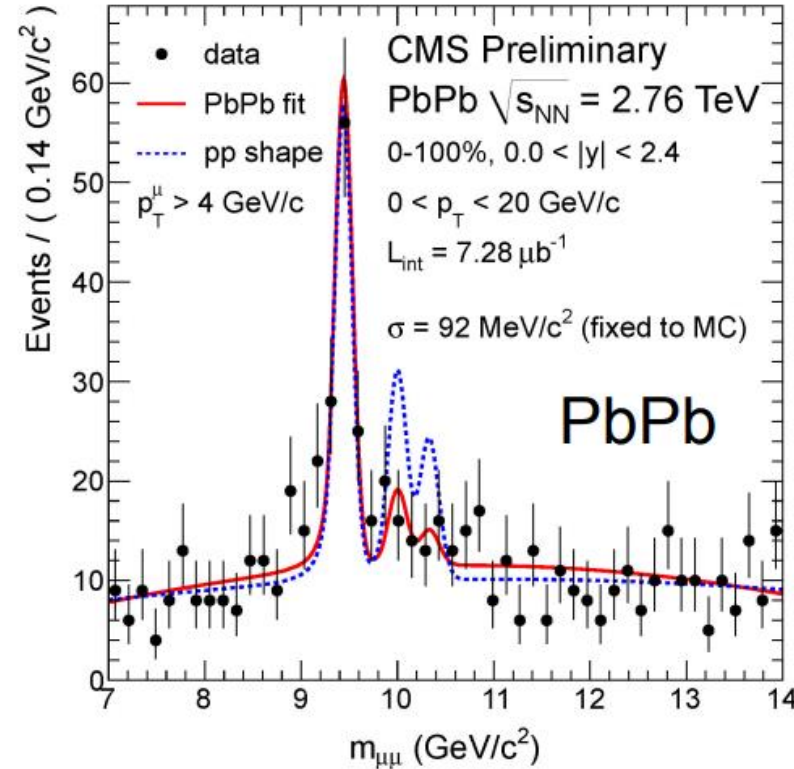
PHENIX, Nucl.Phys.A830:331c-334c,2009

- Large integrated luminosity is needed!
- More accurate results from **STAR** and **PHENIX** came after the LHC start
- Now waiting for **sPHENIX**...

A breakthrough



$$\Upsilon(2S + 3S)/\Upsilon(1S)|_{pp} = 0.78^{+0.16}_{-0.14} \pm 0.02$$



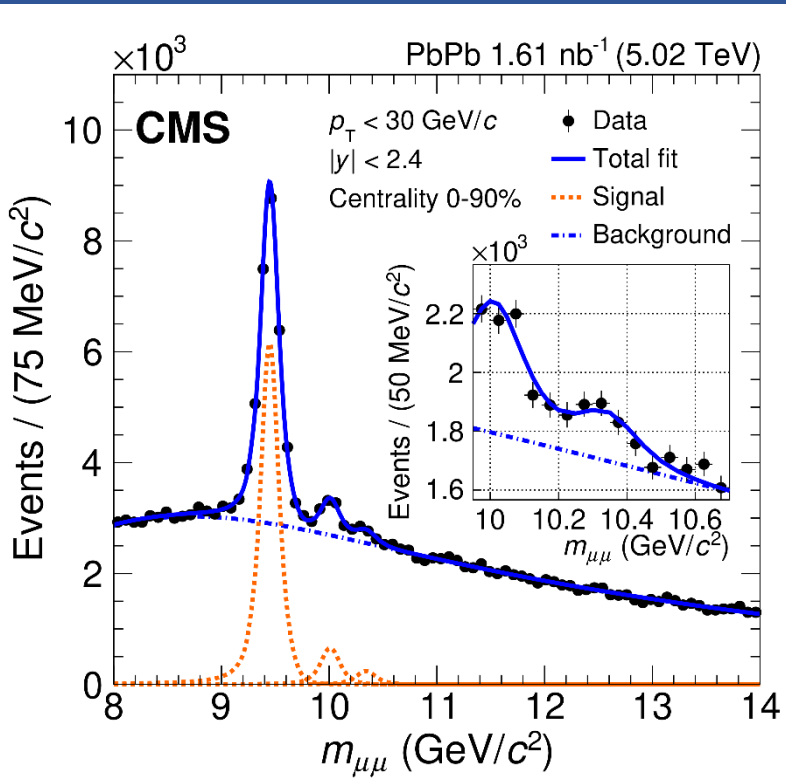
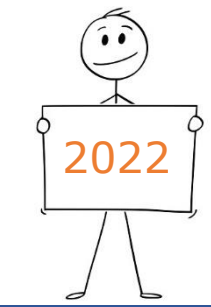
$$\Upsilon(2S + 3S)/\Upsilon(1S)|_{PbPb} = 0.24^{+0.13}_{-0.12} \pm 0.02$$

CMS, J. Phys. G38(2011) 124071
CMS, PRL 107 (2011) 052302

First evidence for a **sequential suppression**
in the bottomonium sector!

After 10 years...

CMS, arXiv:2303.17026

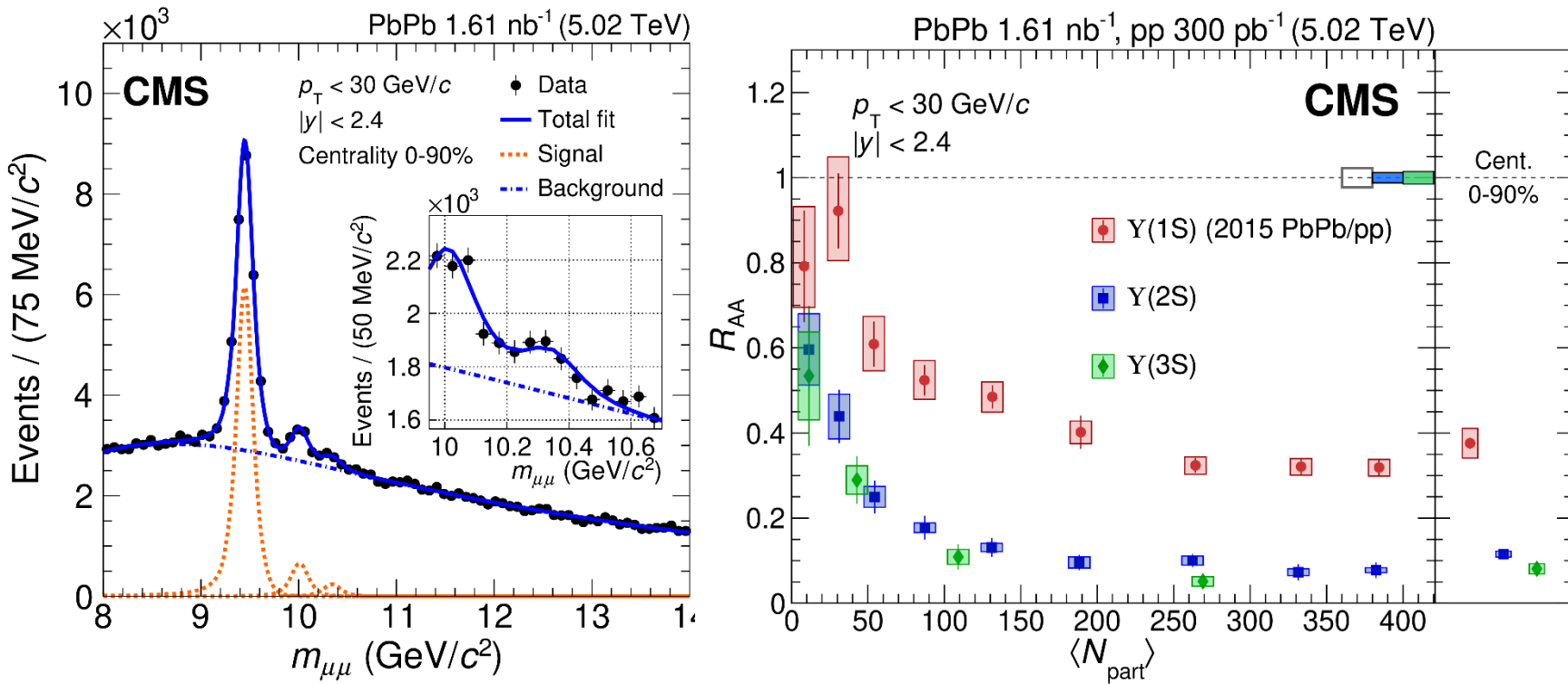


$\Upsilon(3S)$ measured for the first
time in Pb-Pb collisions

After 10 years...

CMS, arXiv:2303.17026

N.B. also the **tightly bound 1S state** is strongly suppressed



$\Upsilon(3S)$ measured for the first time in Pb-Pb collisions

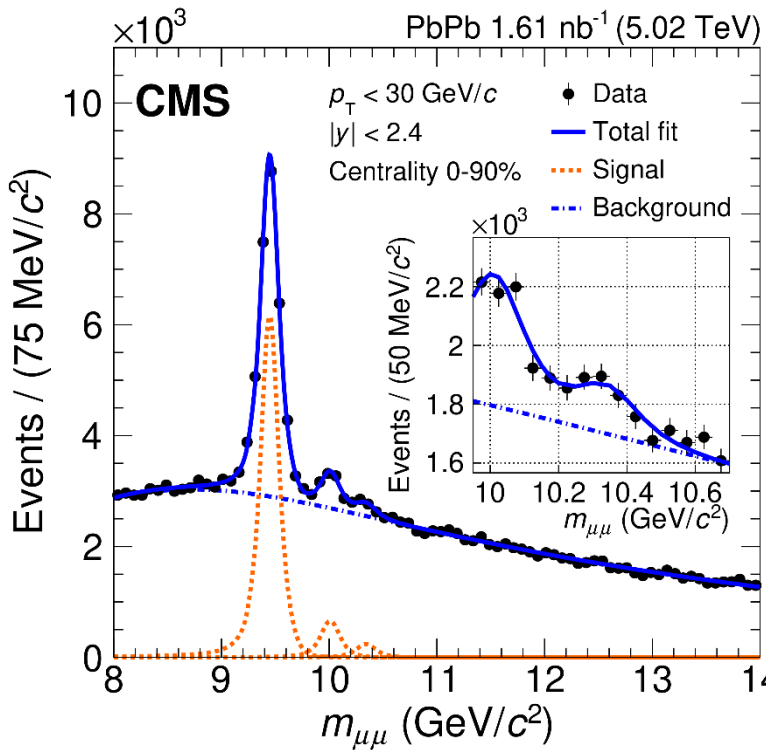
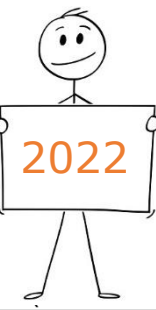


Hierarchy of suppression for the 1S, 2S, 3S states

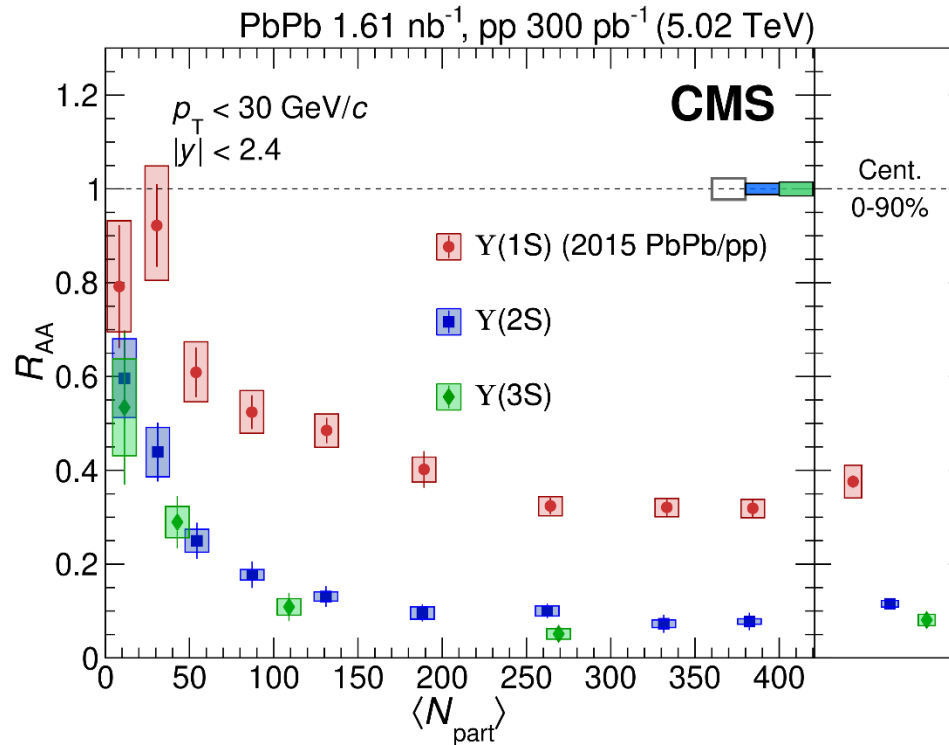
After 10 years...

CMS, arXiv:2303.17026

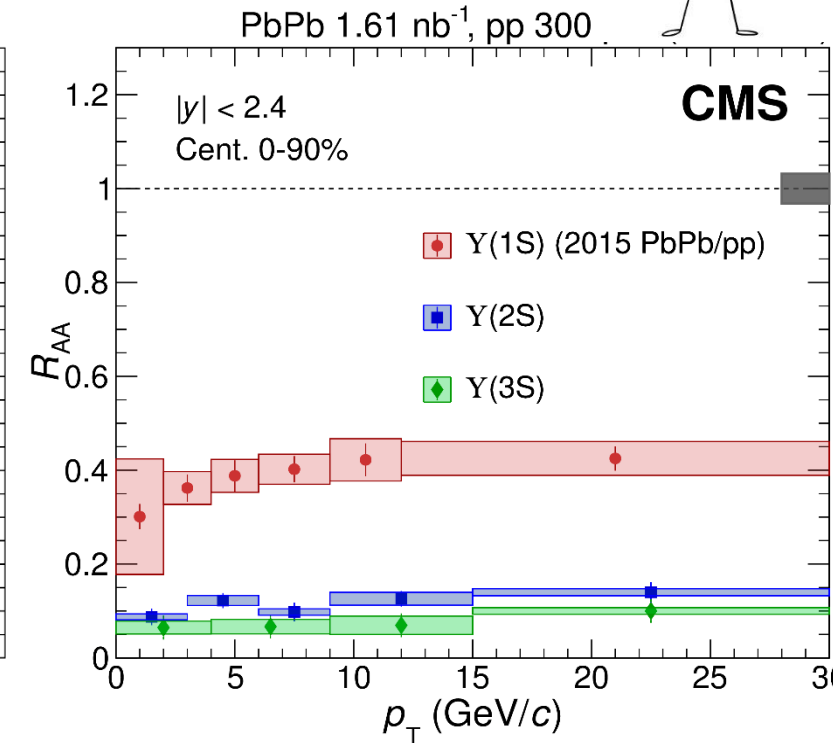
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$\Upsilon(3S)$ measured for the first time in Pb-Pb collisions



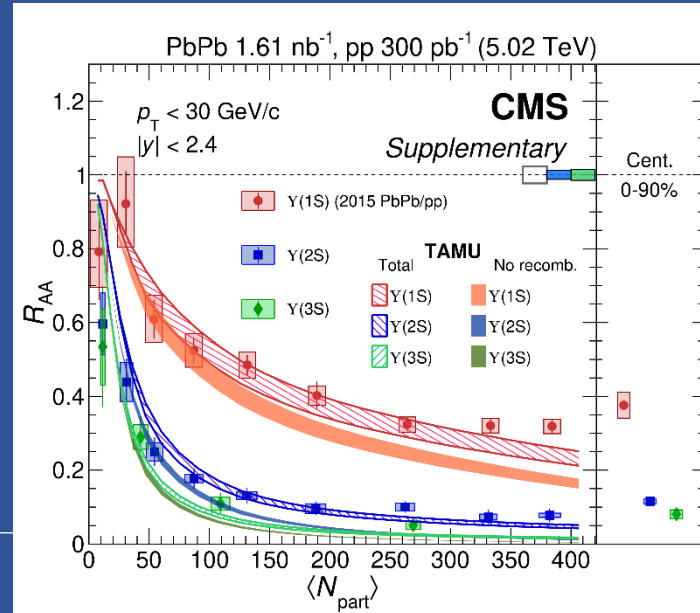
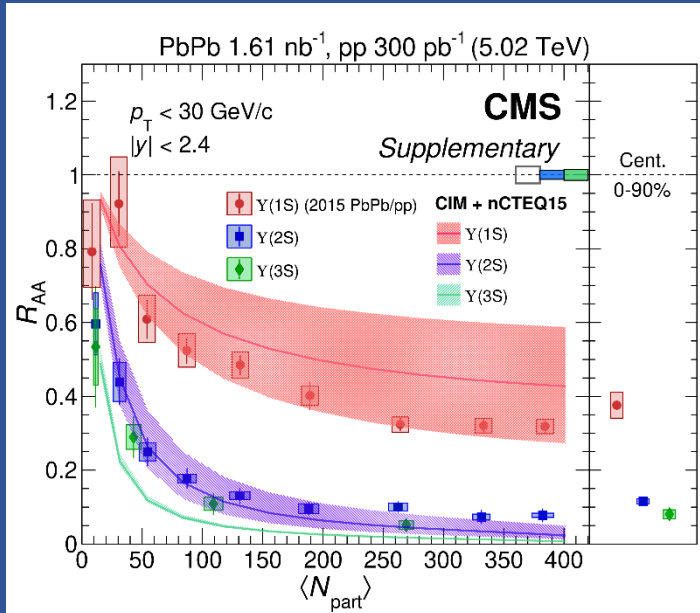
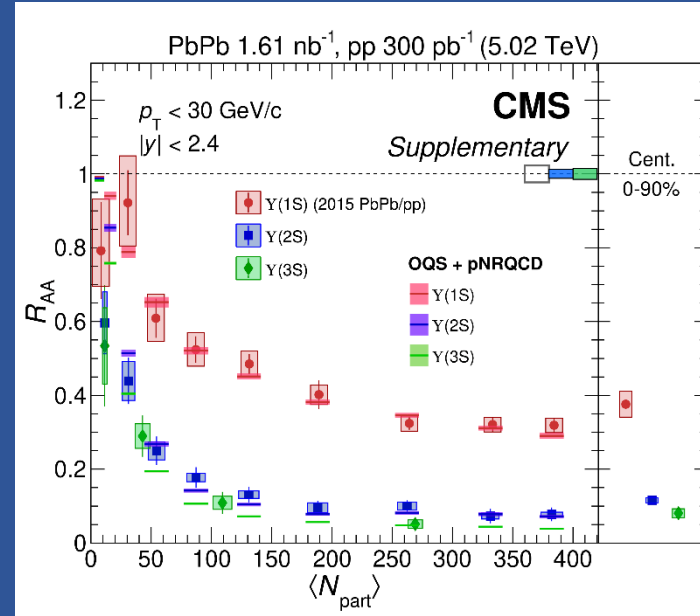
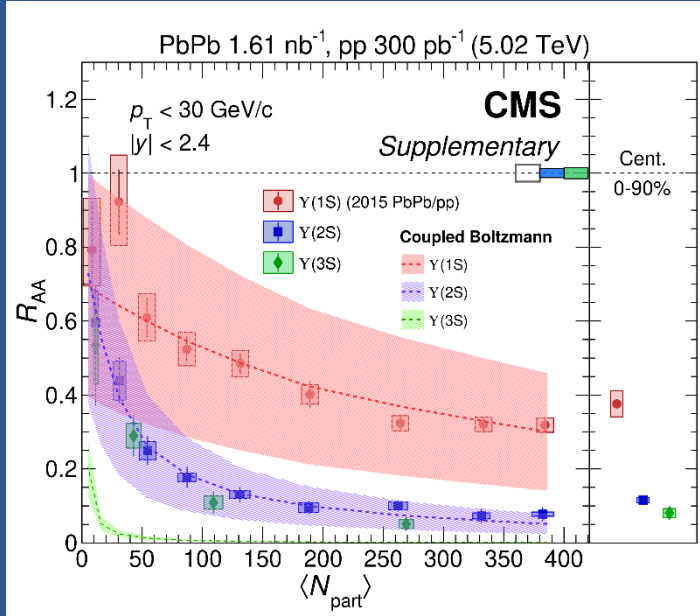
Hierarchy of suppression for the 1S,2S,3S states



No appreciable p_T dependence of R_{AA}

Comparisons with theory models

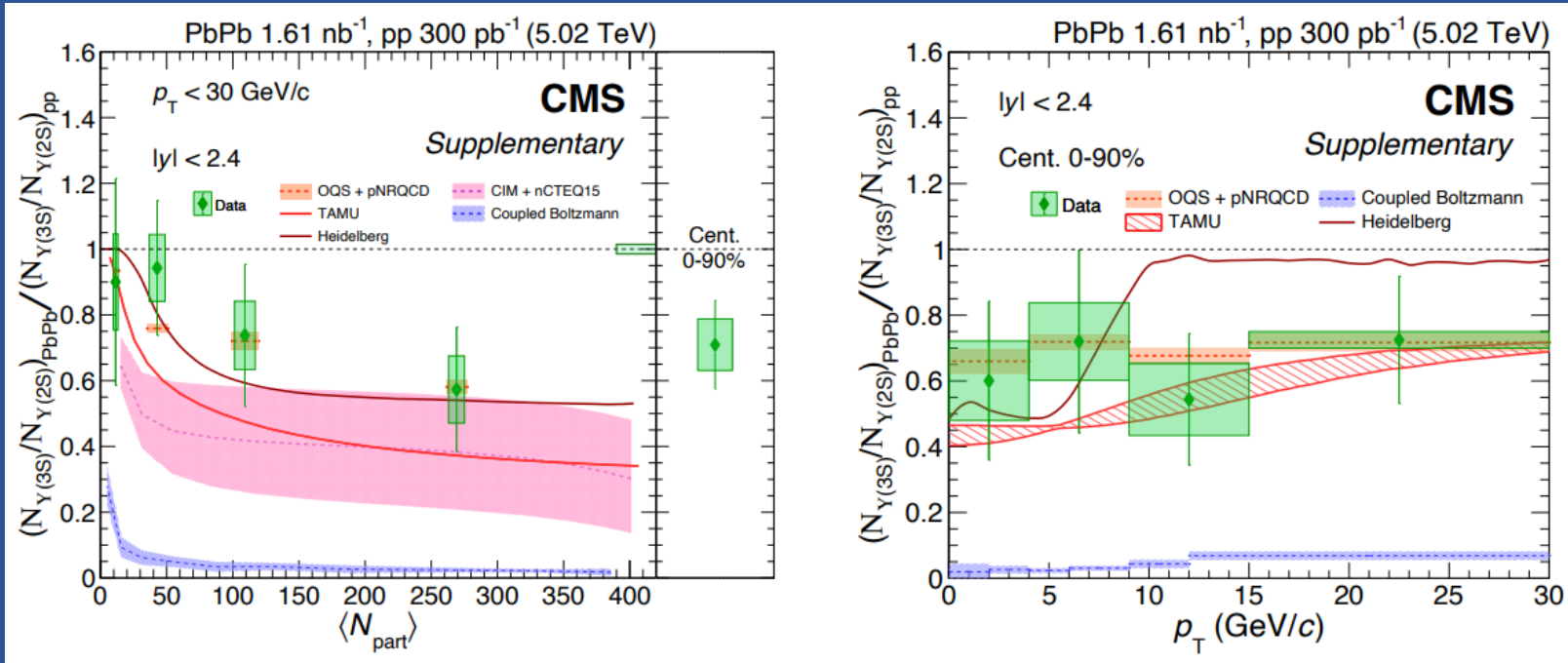
CMS, arXiv:2303.17026



Several approaches are able to reproduce semi-quantitatively the experimental observations! (also the p_T dependence)

Look in more detail at the excited states...

Comparisons with theory models

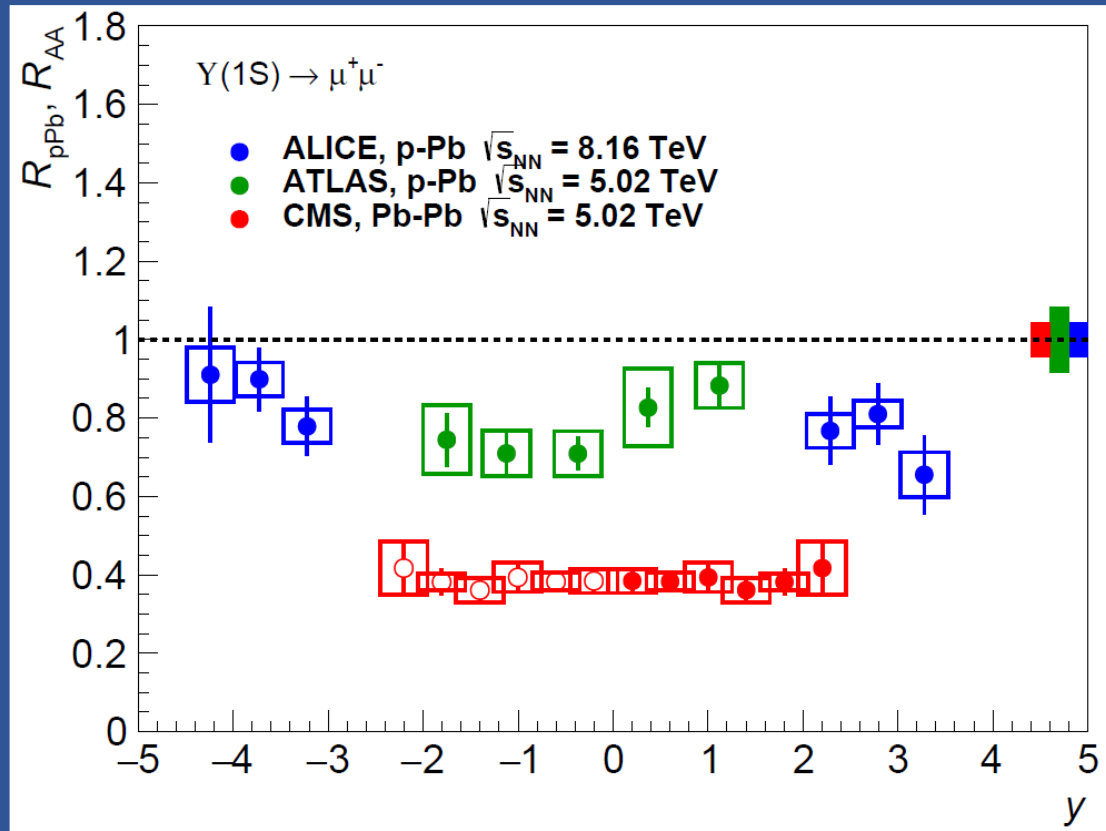


Several approaches are able to reproduce semi-quantitatively the experimental observations! (also the p_T dependence)

Look in more detail at the excited states...

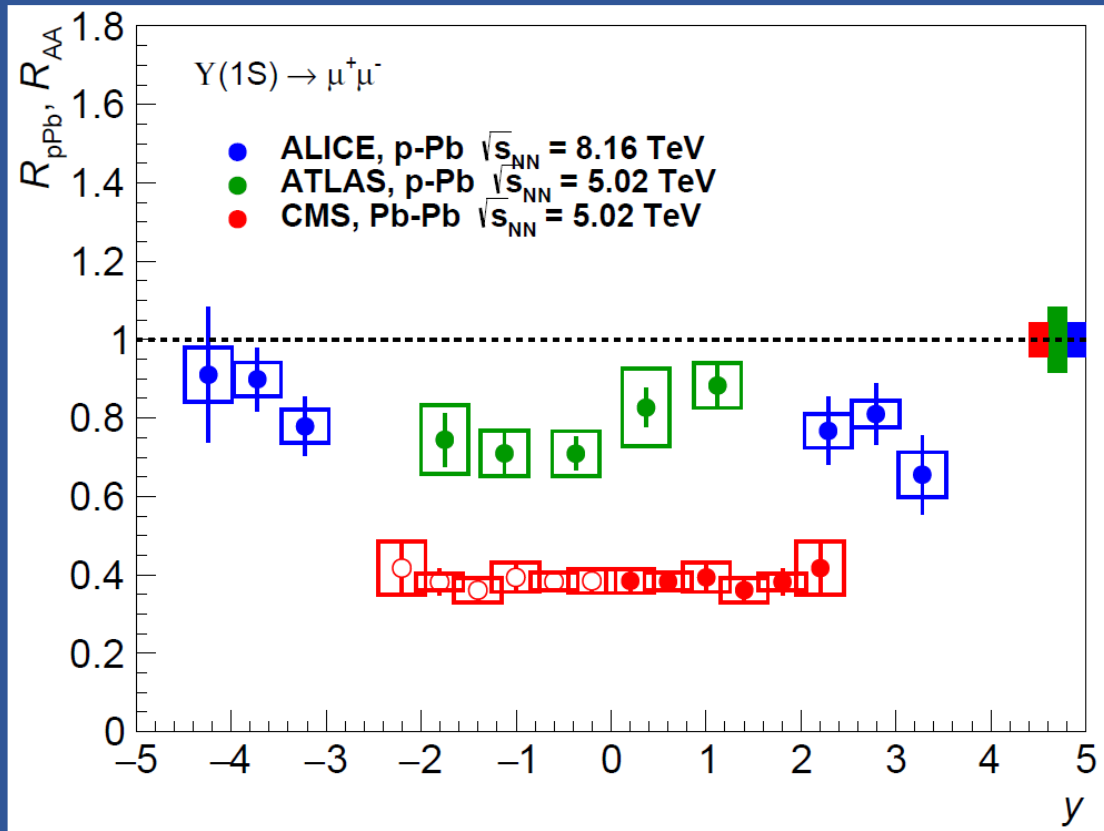
- $Y(3S)/Y(2S)$ double ratio → indication of **stronger suppression for $Y(3S)$** , particularly for central events
- Significant differences among model calculations
- This set of data, in spite of the relatively large uncertainties, poses **strong constraints to the models**

What about CNM effects ?



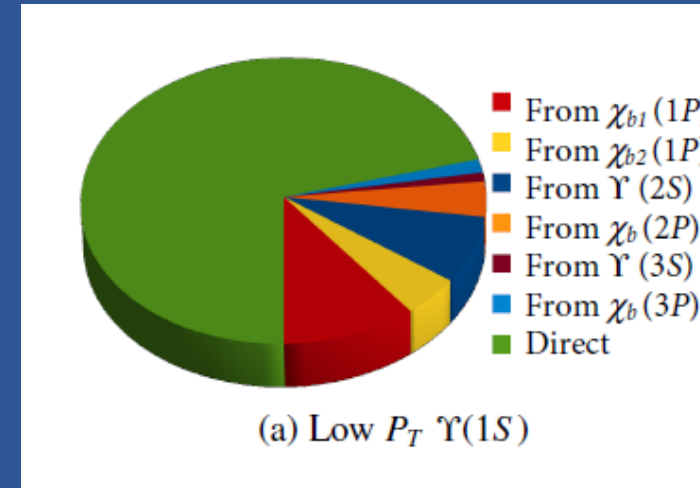
- Is the strong suppression of $\Upsilon(1S)$ “compatible” with its extremely large binding energy (>1 GeV) ?
- Can a fraction of $\Upsilon(1S)$ suppression be due to **non-QGP effects?**

Is $\Upsilon(1S)$ really suppressed ?



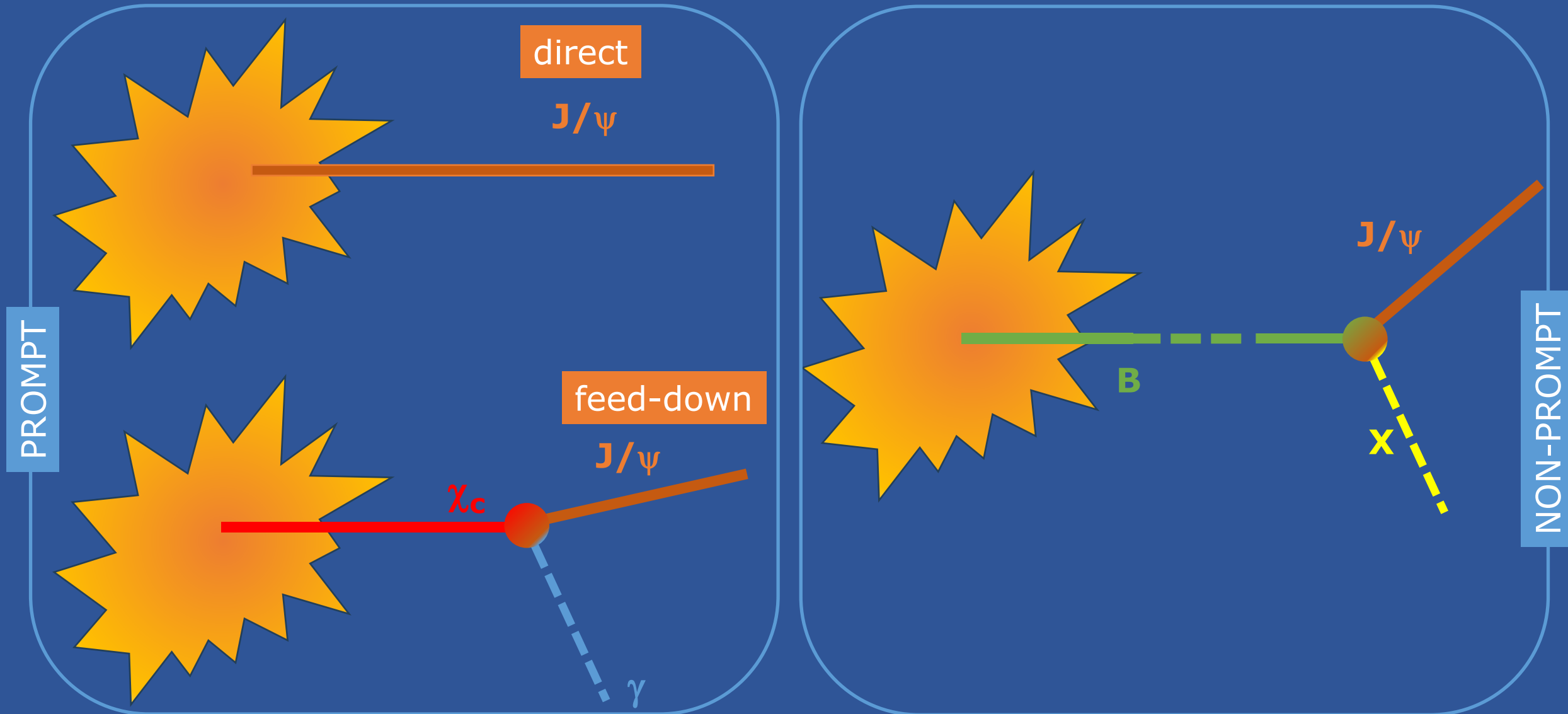
CNM effect on $\Upsilon(1S)$
are shown to be non-negligible

Also considering feed-down effects on $\Upsilon(1S)$ from $\Upsilon(2S,3S)$ and χ_b states



Could the observed **$\Upsilon(1S)$ suppression** be compatible with a **combination of feed-down + CNM effects** ?

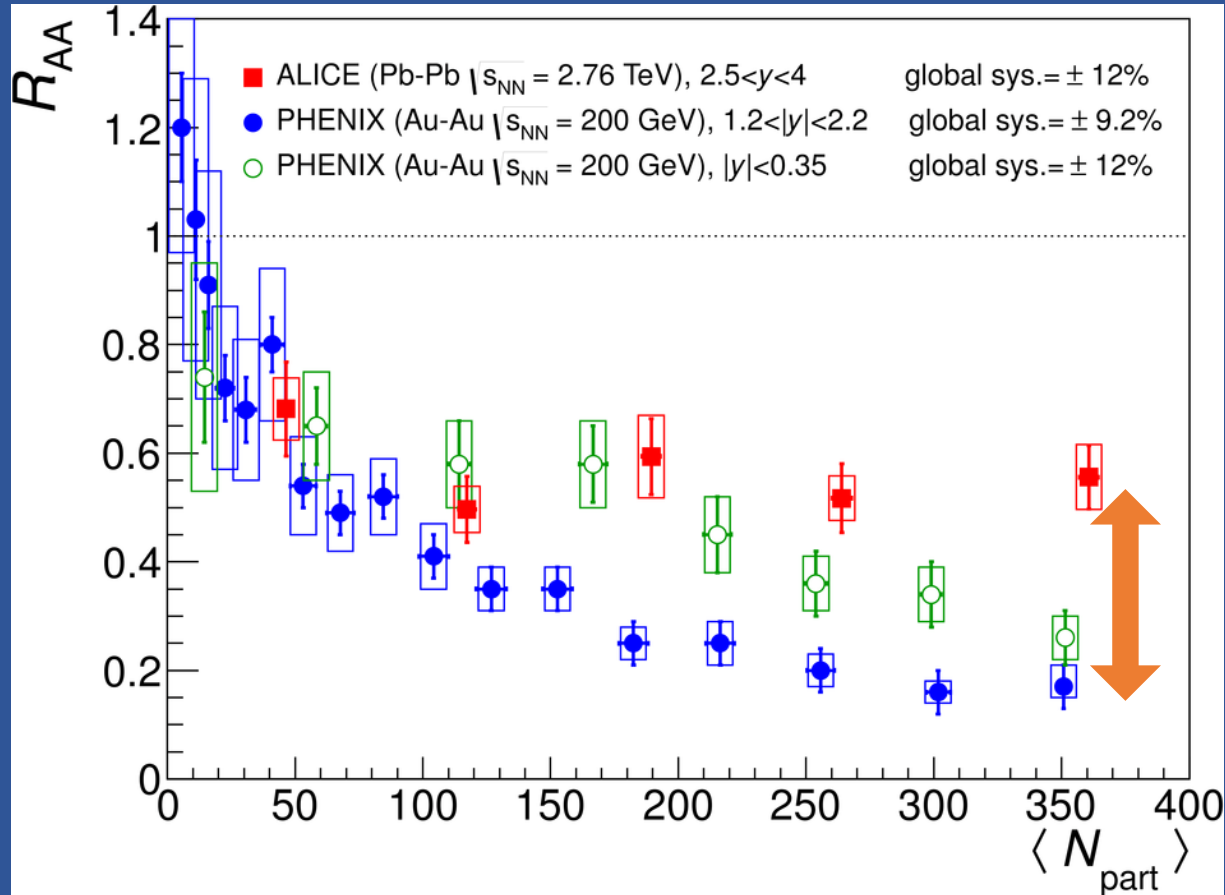
Before coming back to charmonia...



Let's come back to charmonia at LHC



- RHIC to LHC: factor >10 in $\sqrt{s_{NN}}$: a **decisive test** of mechanisms at play



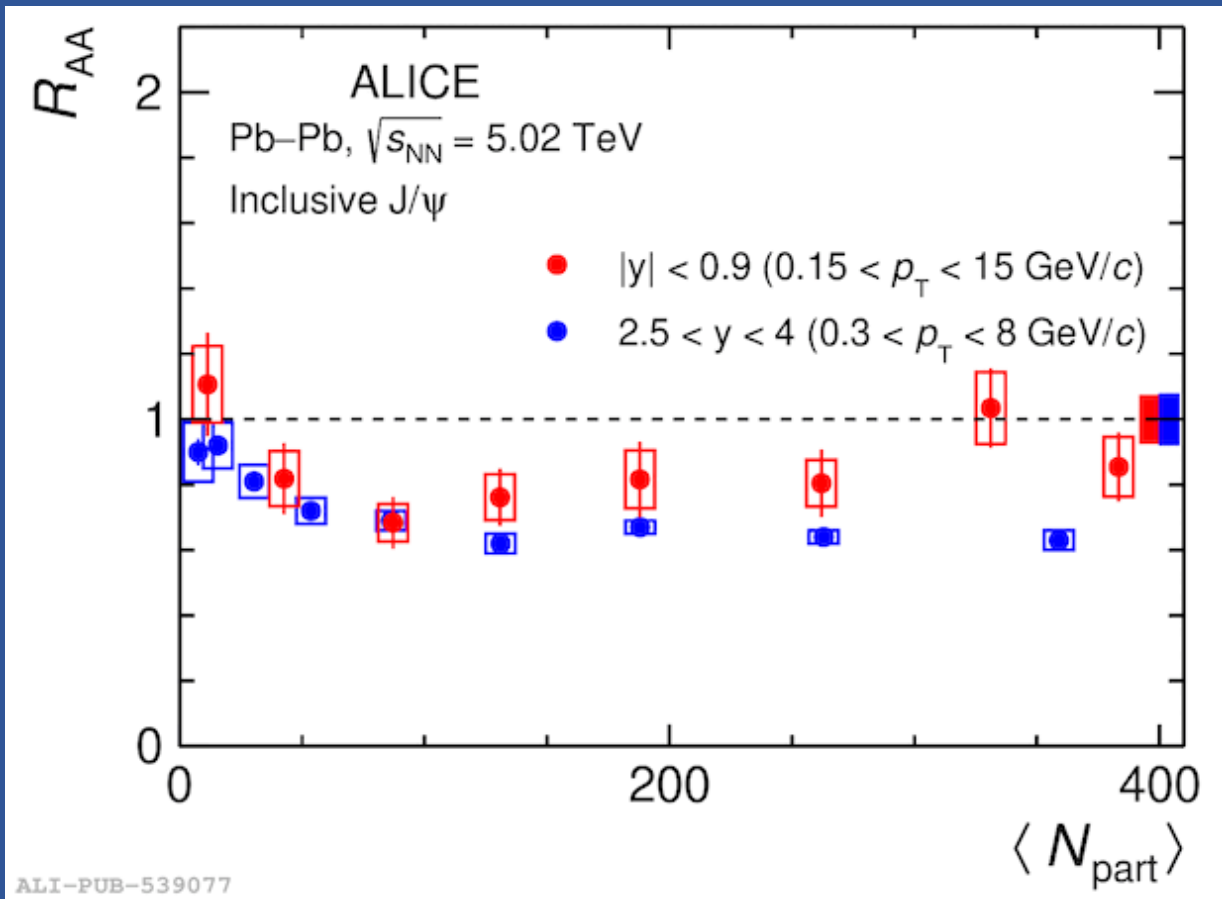
- Central Pb-Pb collisions
 - R_{AA} is **larger by a factor >2** at LHC energy with respect to RHIC
- Forward-y results (muon trigger!)

Smoking gun for **(re)generation effects ?**

ALICE, PRL 109 (2012) 072301

After 10 years...

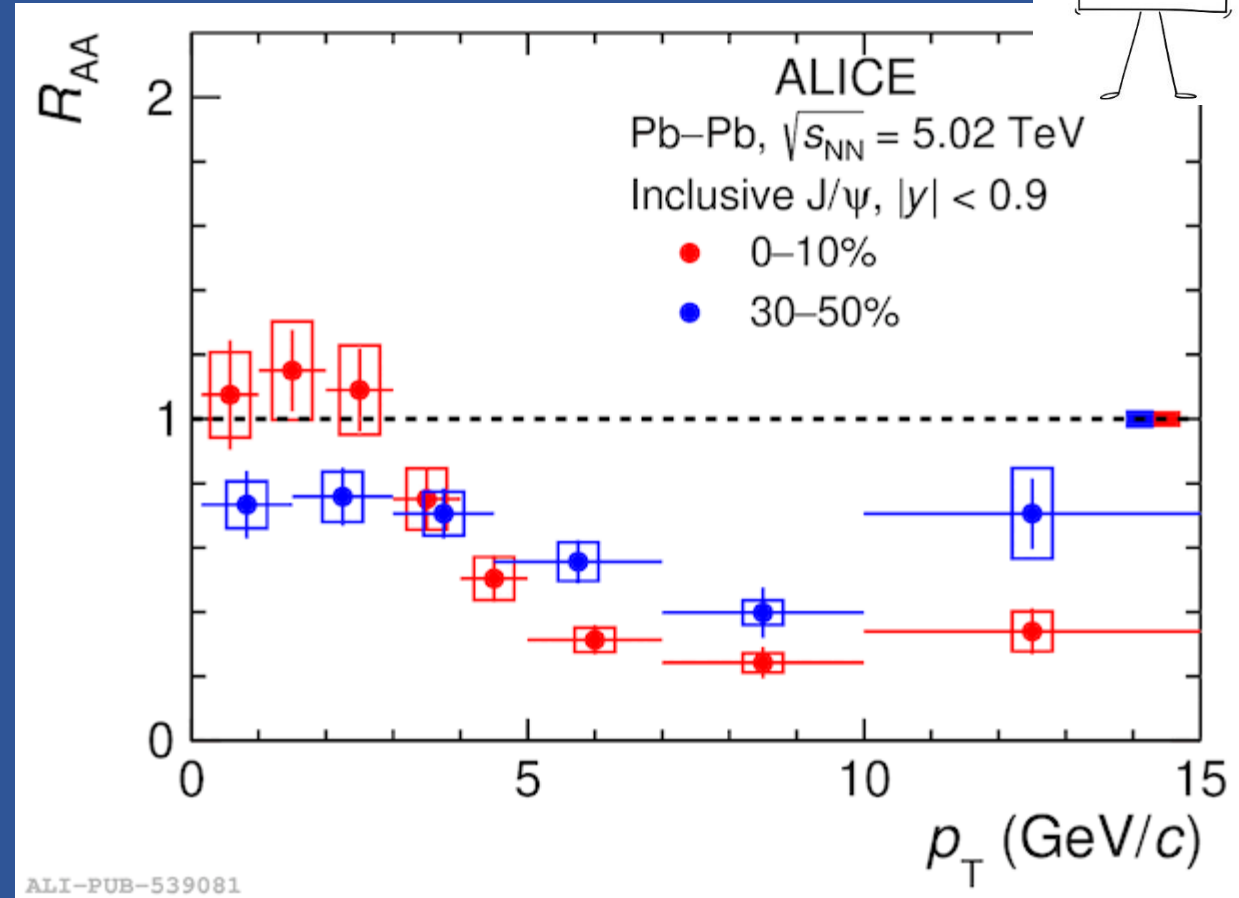
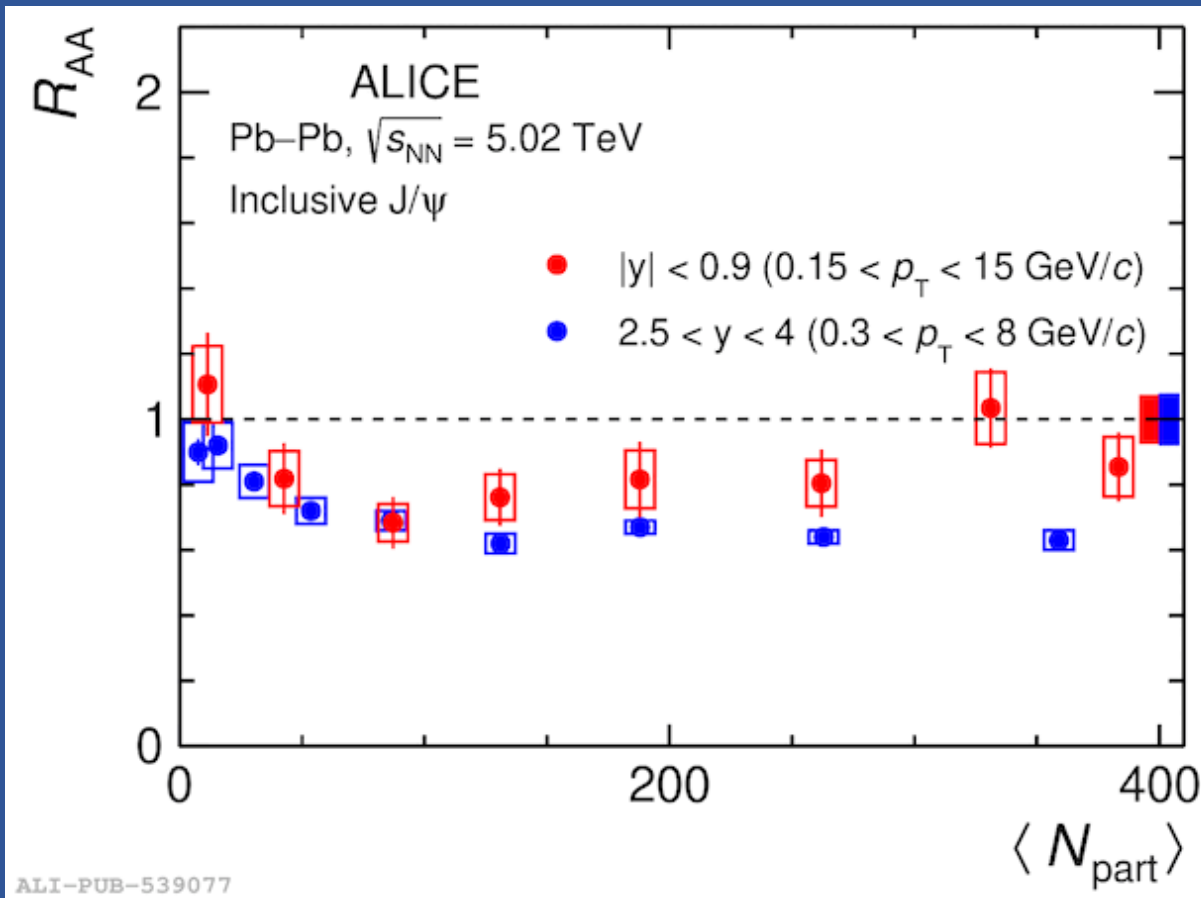
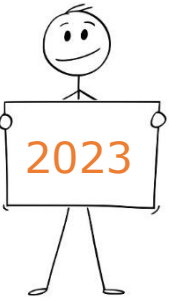
Smoking gun(s) for **(re)generation**



➤ R_{AA} is **larger at midrapidity** ← as expected from a larger charm quark multiplicity

After 10 years...

Smoking gun(s) for **(re)generation**



- R_{AA} is **larger at midrapidity** ← as expected from a larger charm quark multiplicity
- R_{AA} **increases at low p_T** ← as expected from the charm quark p_T distribution

Theory models: generation vs regeneration

Transport

- ❑ Macroscopic rate equation including suppression and regeneration in the QGP
- ❑ Suppression
 - ❑ Calculated starting from modifications of charmonium spectral functions, **constrained by LQCD-validated potentials**
- ❑ Regeneration
 - ❑ Tuned from measured heavy-quark yields

X. Du and R. Rapp,
NPA 943(2015) 14P.7
P. Zhou et al.,
PRC89 (2014) 054911

Statistical hadronization

- ❑ Charmonium **yields determined at chemical freeze-out** according to their statistical weights
- ❑ Charm fugacity factor related to charm conservation and based on experimental data on production cross sections

A. Andronic et al.,
Nature 561 (2018) 321

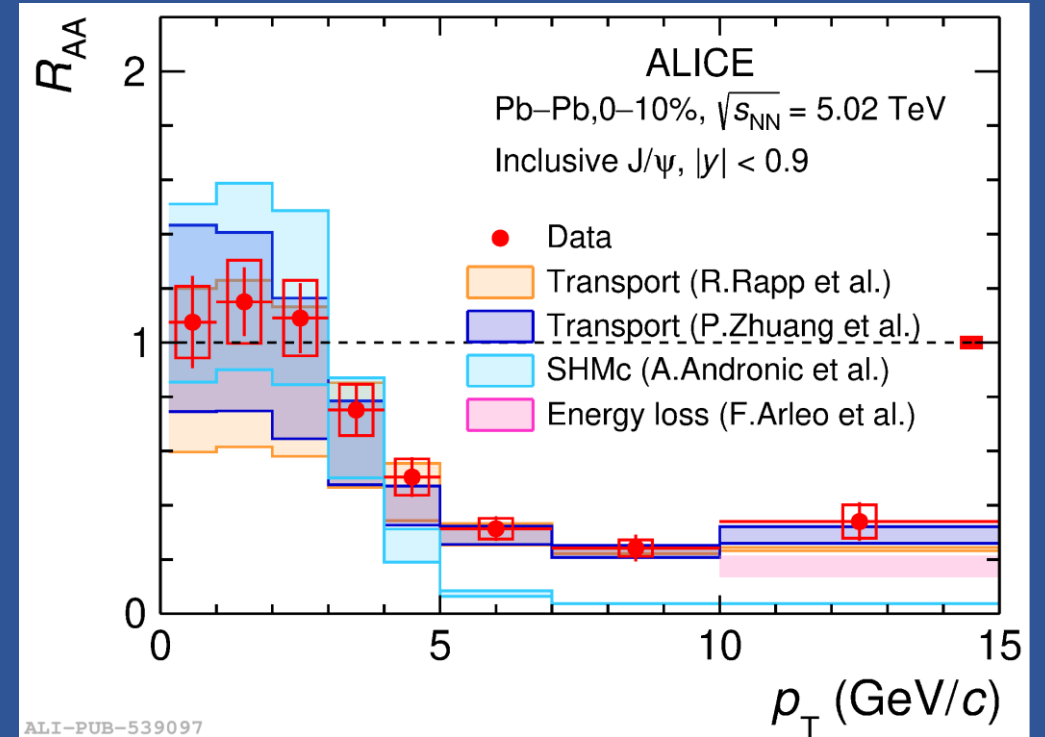
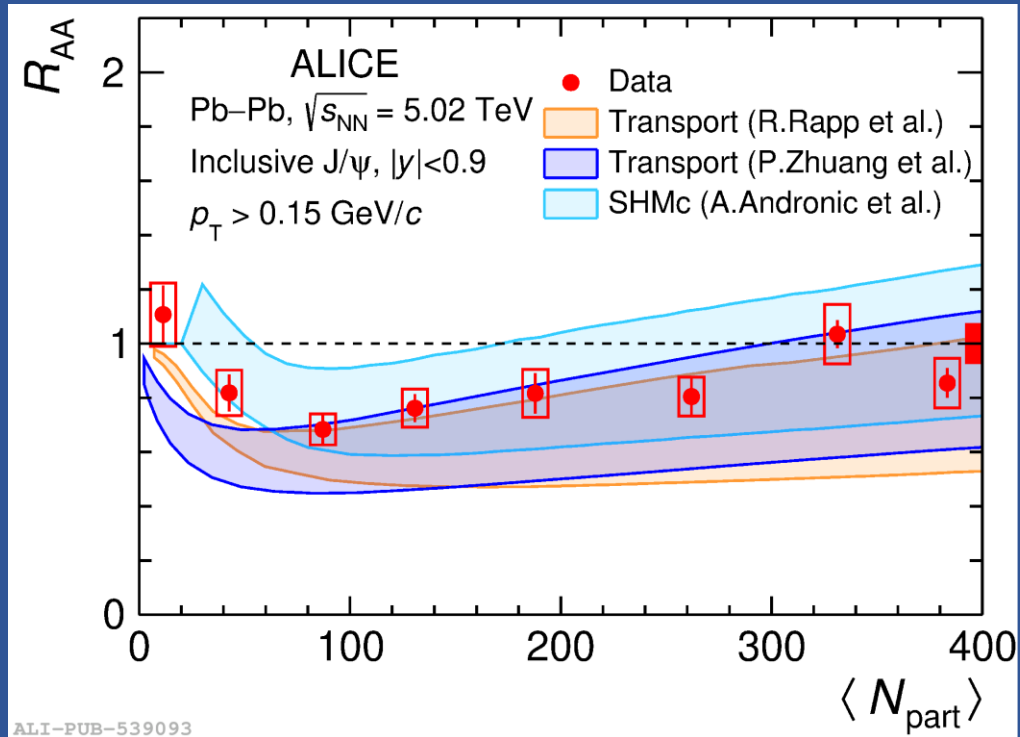
It turns out that both
approaches fairly **reproduce**
LHC experimental results
on the J/ψ

Other approaches exist!
“Comover” models etc.

ALICE, Phys. Lett. B 766 (2017) 212

E. Ferreira,
PLB 731 (2014) 57

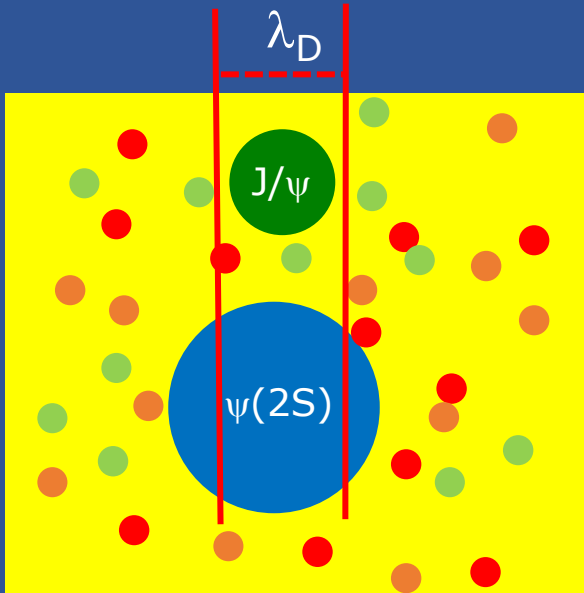
Comparisons with theory models



- Model **uncertainties** dominated by
 - Open charm cross section
 - Initial state effects (shadowing)

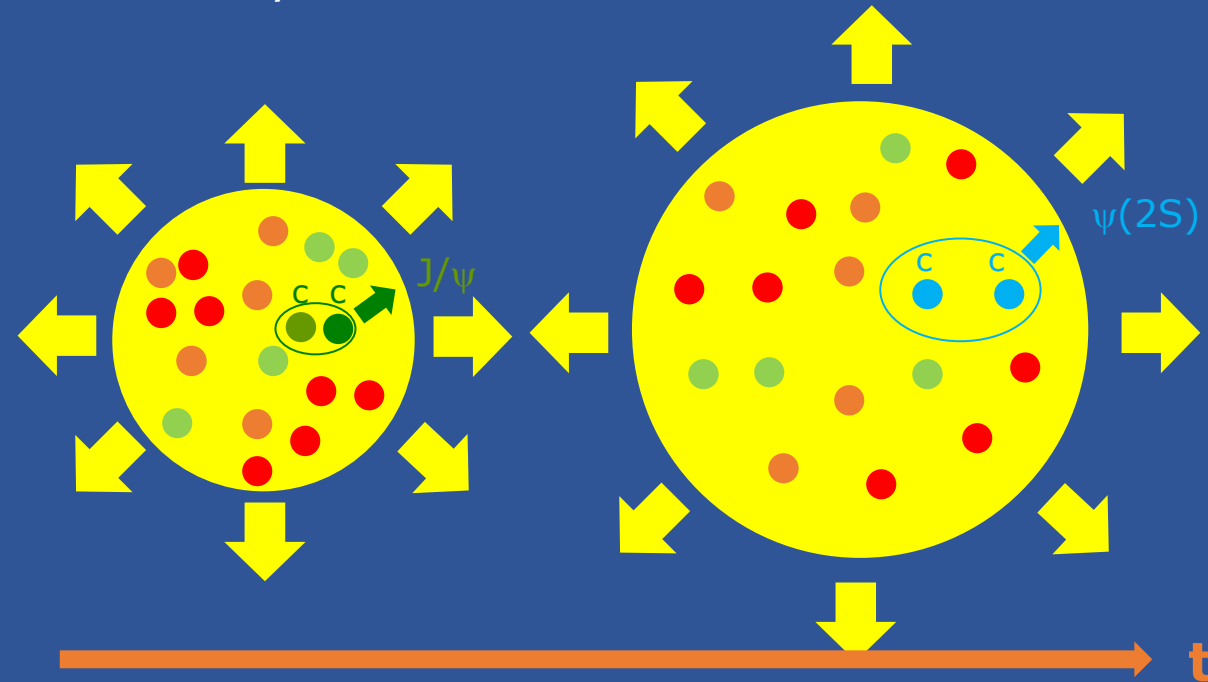
Sequential regeneration for charmonia at LHC ?

- Binding energy $\sim (2m_D - m_\psi) \rightarrow \psi(2S) \sim 60 \text{ MeV}, J/\psi \sim 640 \text{ MeV}$



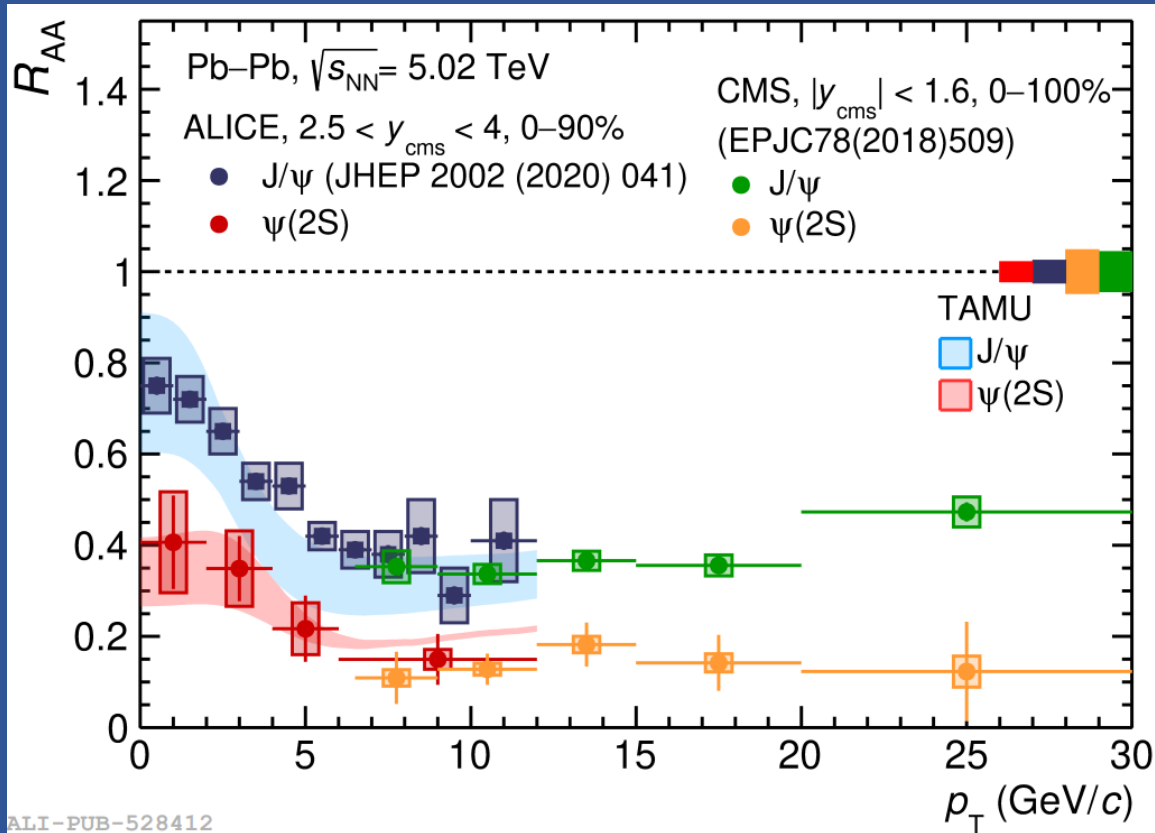
- **Much stronger dissociation effects** for the weakly bound $\psi(2S)$ state already seen at SPS energy
- Effect of re-combination on $\psi(2S)$ could also be important \rightarrow important when the system is **more diluted** ?

Important test
for models!



$\psi(2S)$ at LHC

ALICE, arXiv:2210.08893

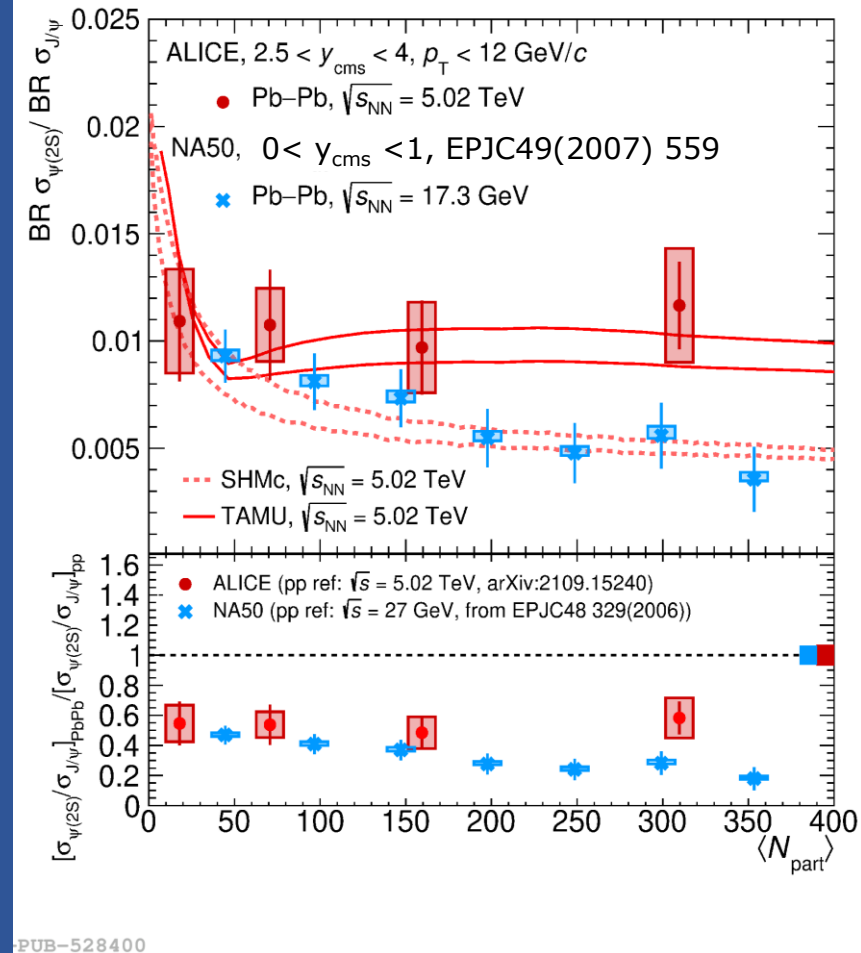
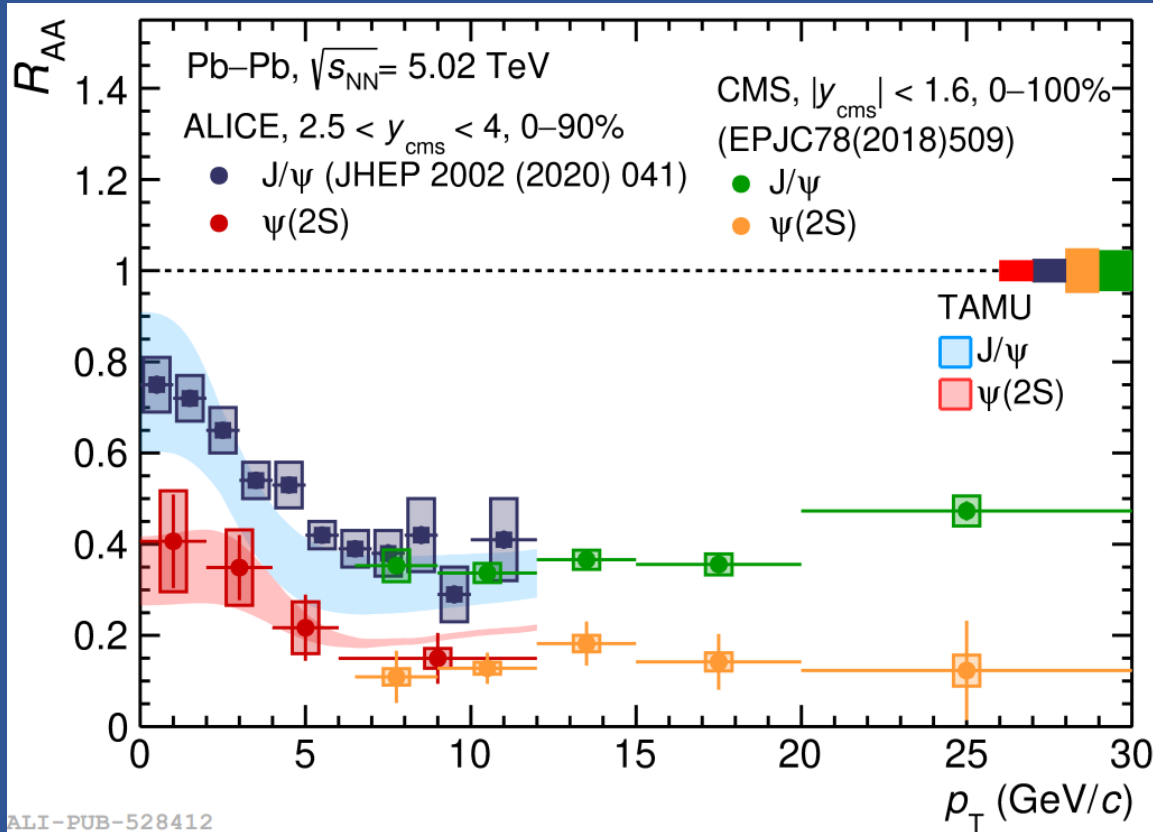


ALI-PUB-528412

- Two main conclusions
- Indication for an **increase of R_{AA} at low p_T** → “sequential” regeneration ?

$\psi(2S)$ at LHC

ALICE, arXiv:2210.08893

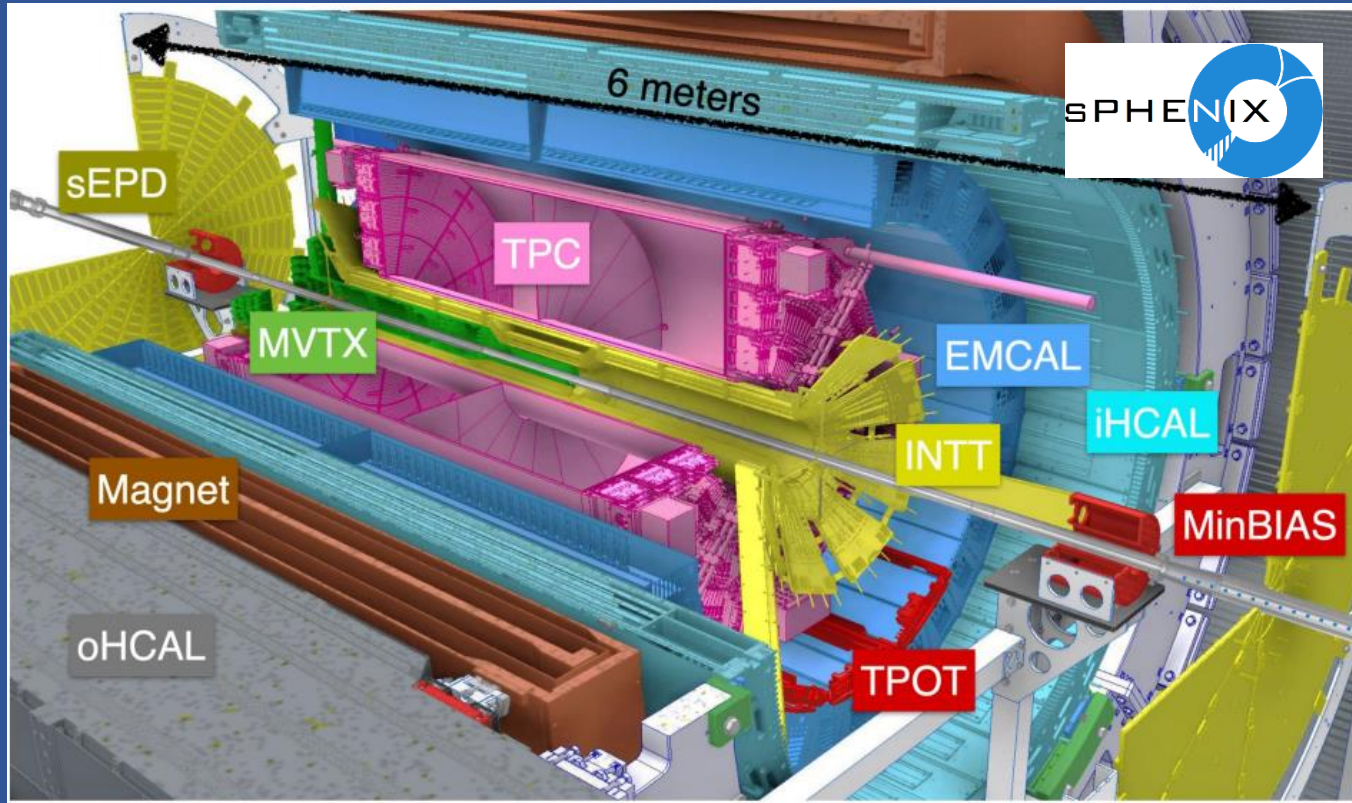


Caveat: **prompt vs inclusive**



- Two main conclusions
- Indication for an **increase of R_{AA} at low p_T** → “sequential” regeneration ?
- First indication for **discrepancy** between transport and statistical approaches → worth investigating!

What remains to be done ? High-energy/low μ_B

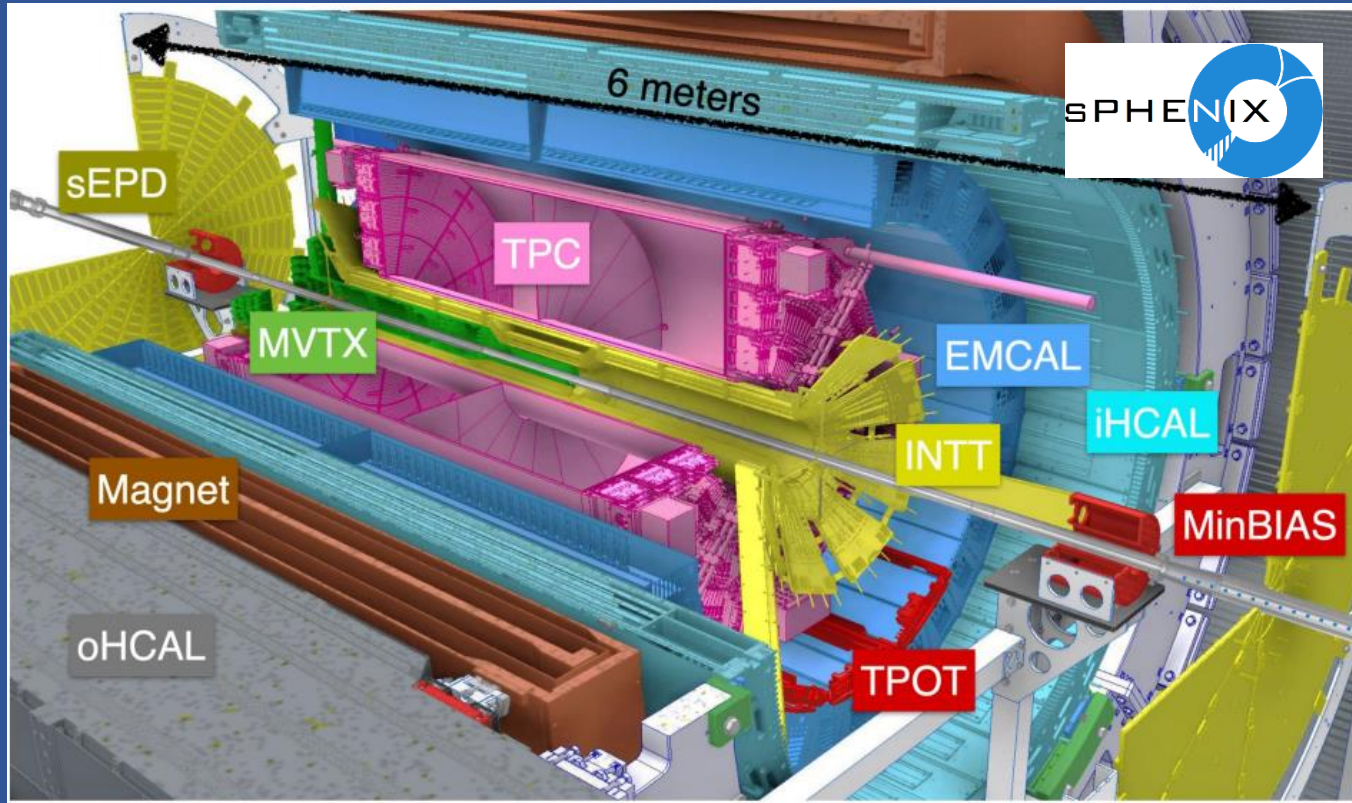


sPHENIX: The first “brand new” heavy-ion experiment since more than a decade

Now taking data!

<https://www.sphenix.bnl.gov/>

What remains to be done ? High-energy/low μ_B

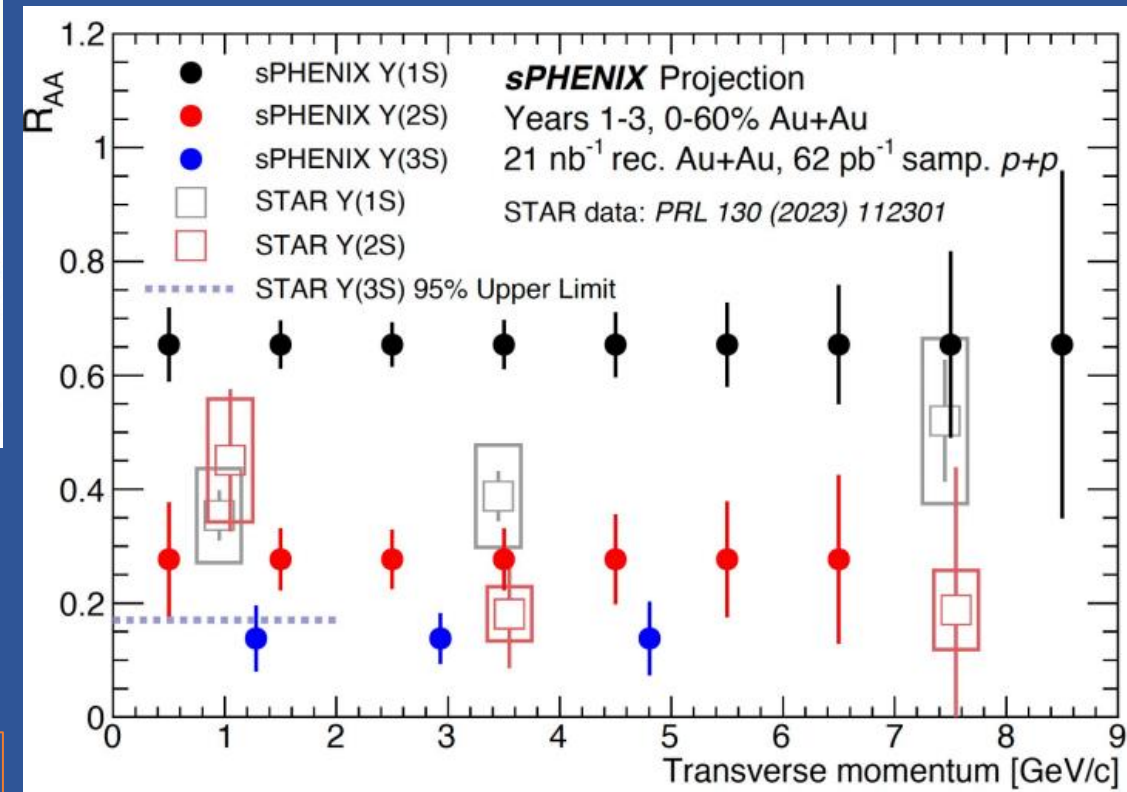


sPHENIX: The first “brand new” heavy-ion experiment since more than a decade

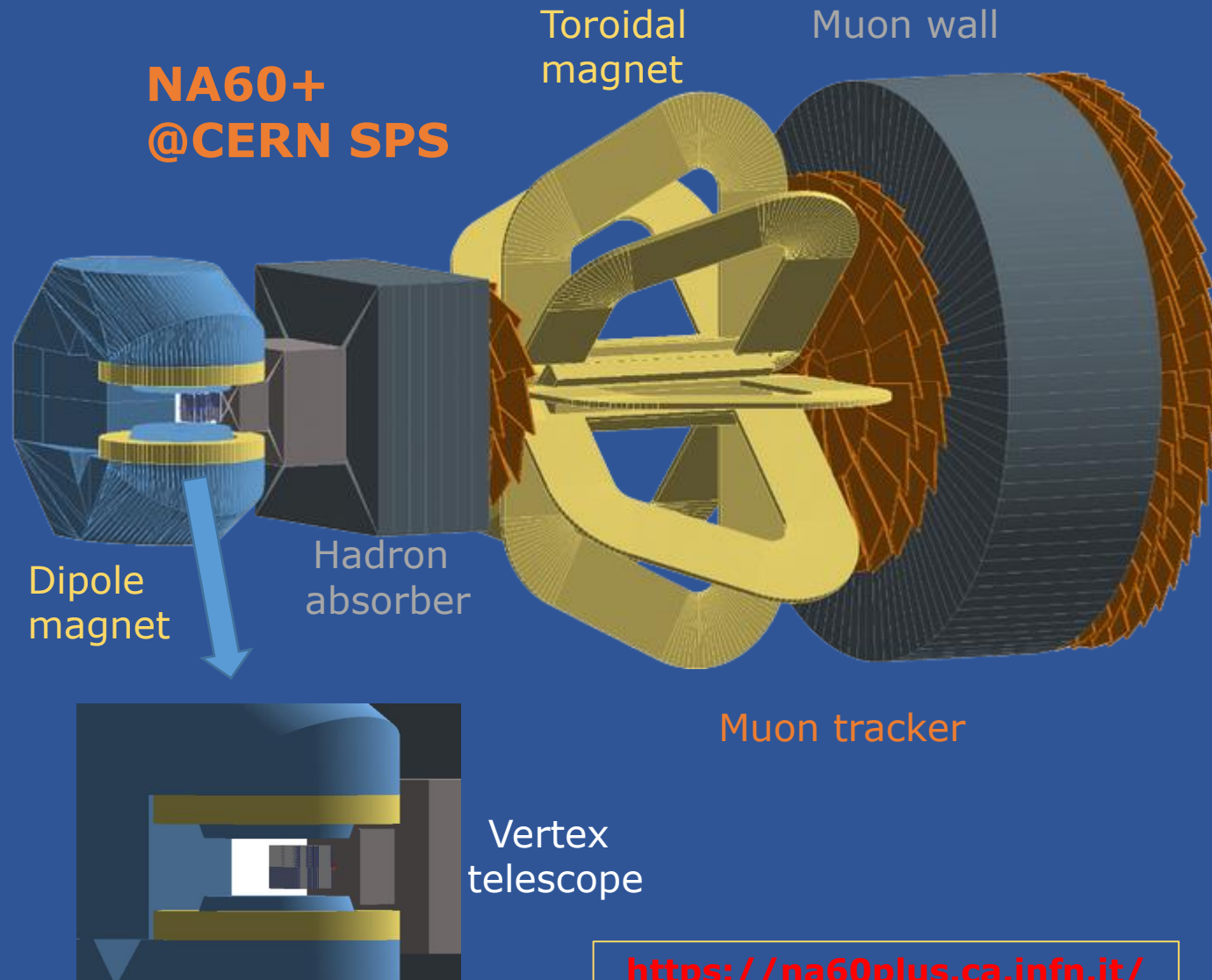
Now taking data!

<https://www.sphenix.bnl.gov/>

- Clear distinction of three Υ states
- Probing the QGP with color dipoles at three length scales
- Kinematic range allows for **comparison between RHIC and LHC measurements**

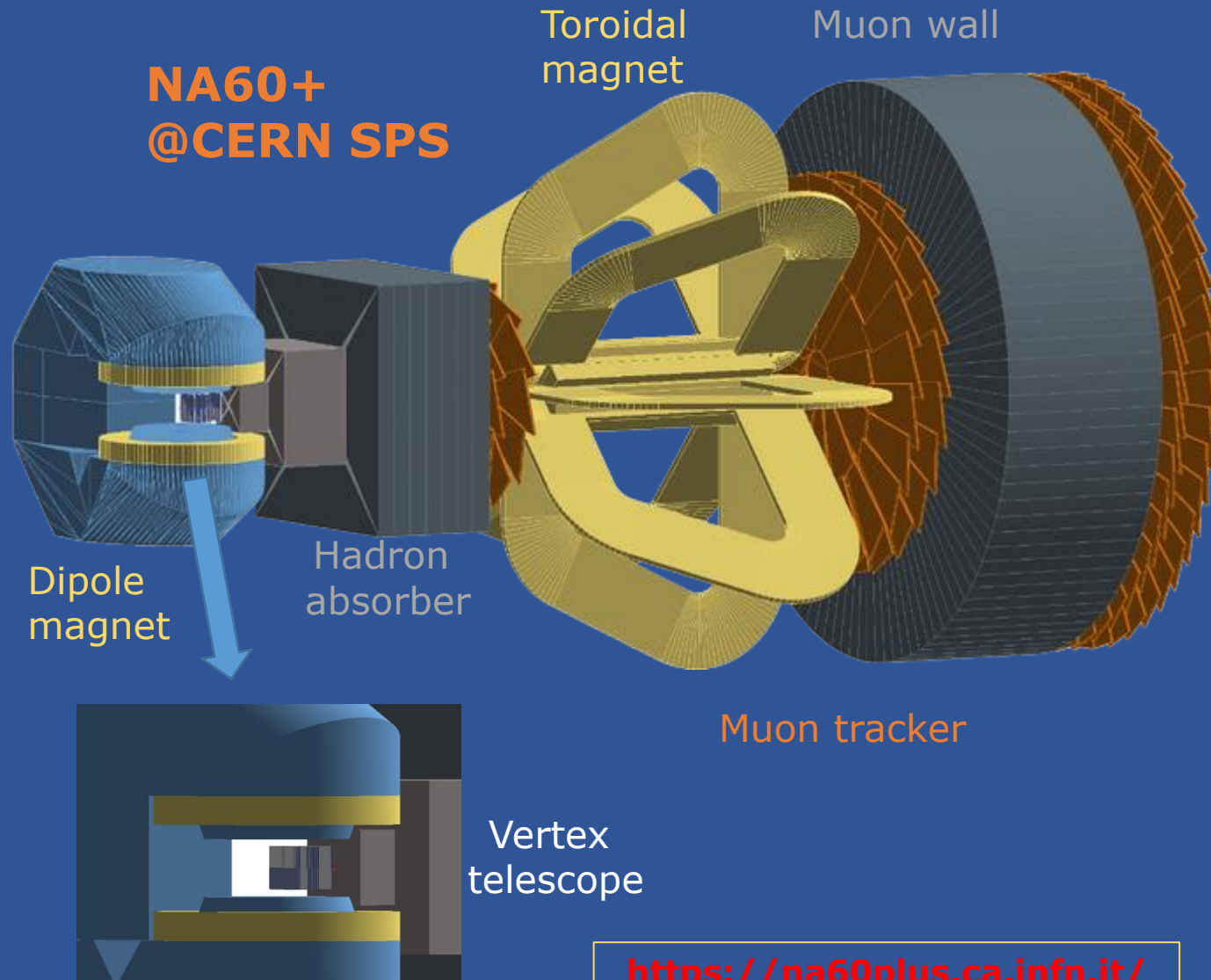


What remains to be done ? Low-energy/high μ_B

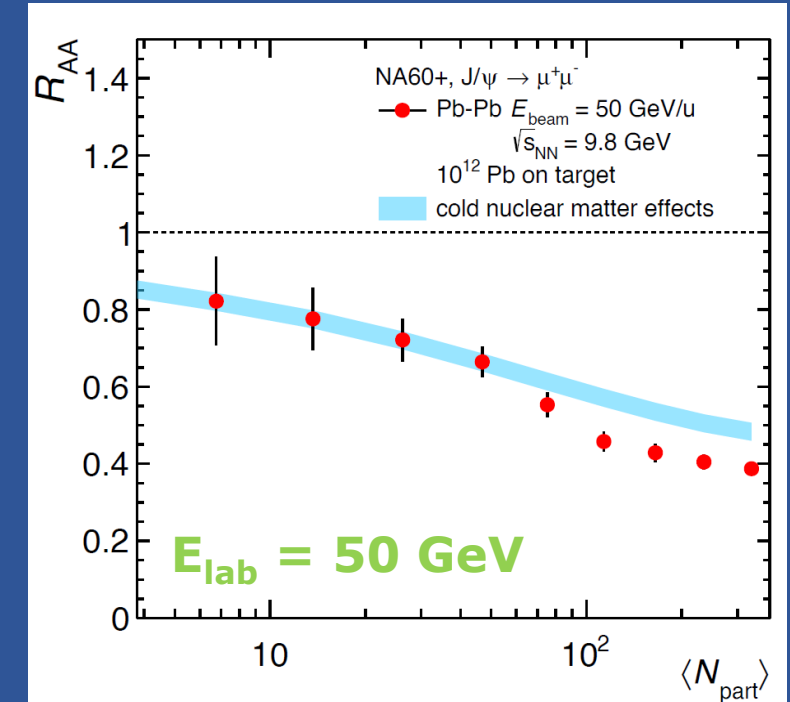


<https://na60plus.ca.infn.it/>

What remains to be done ? Low-energy/high μ_B



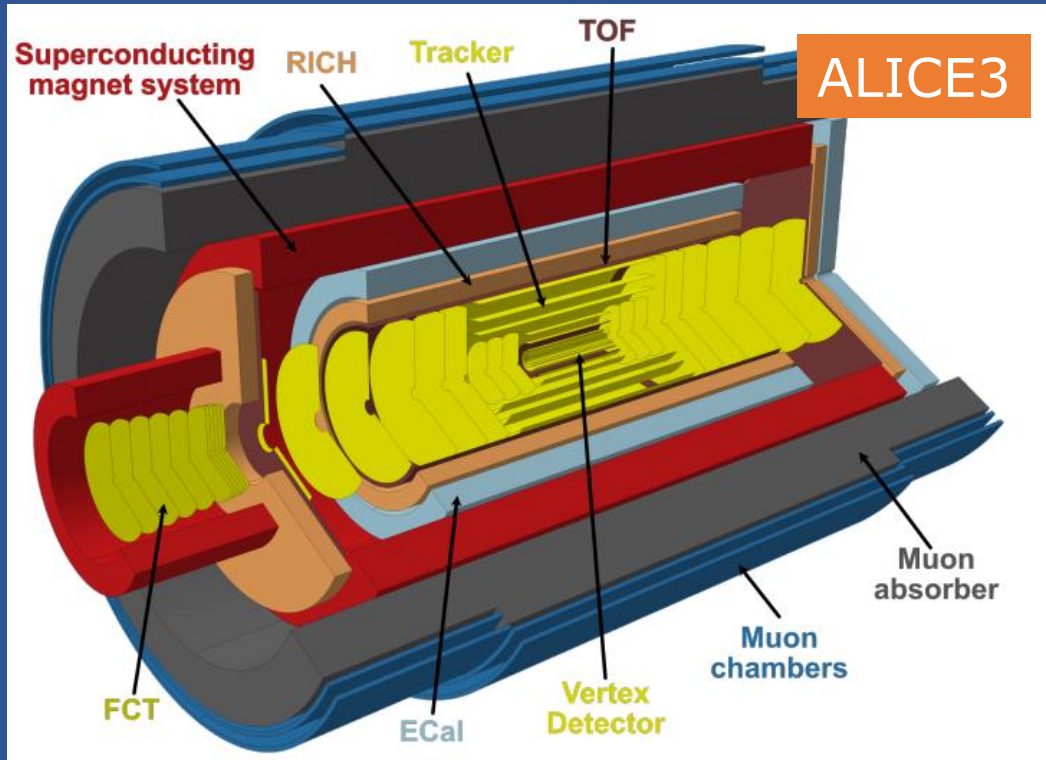
- No charmonium results available below top SPS energy ($\sqrt{s_{NN}}=17$ GeV)



- Study the onset of anomalous suppression and correlate with temperature (thermal dileptons)

<https://na60plus.ca.infn.it/>

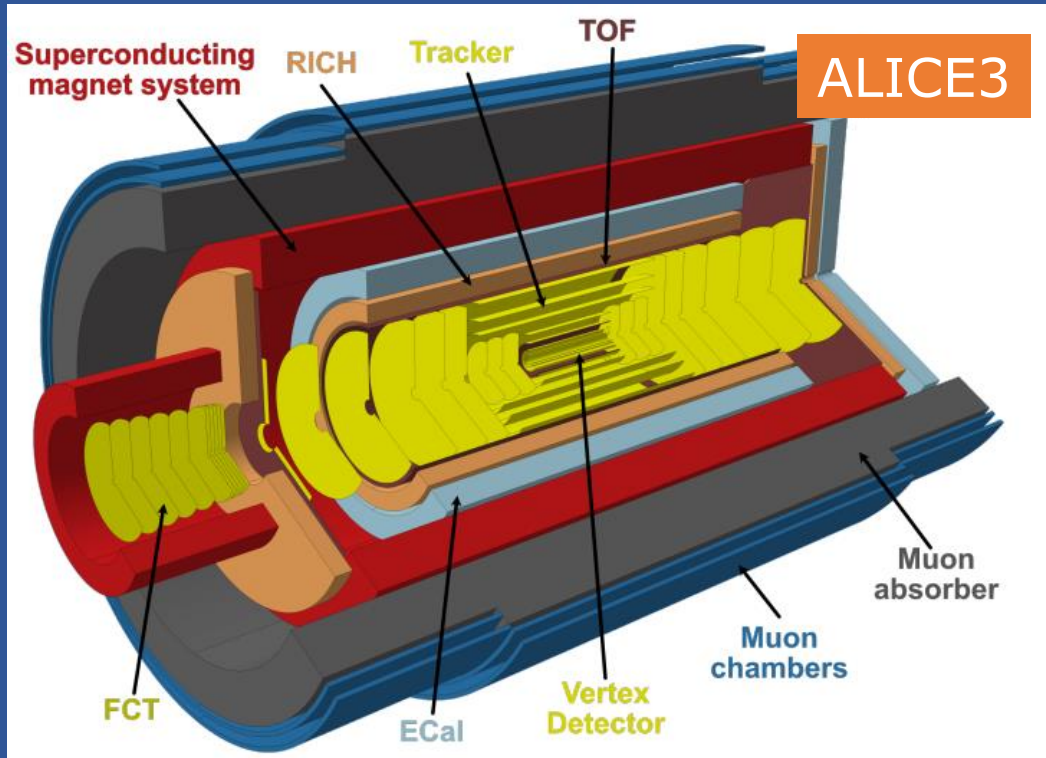
What remains to be done ? High-energy/zero μ_B



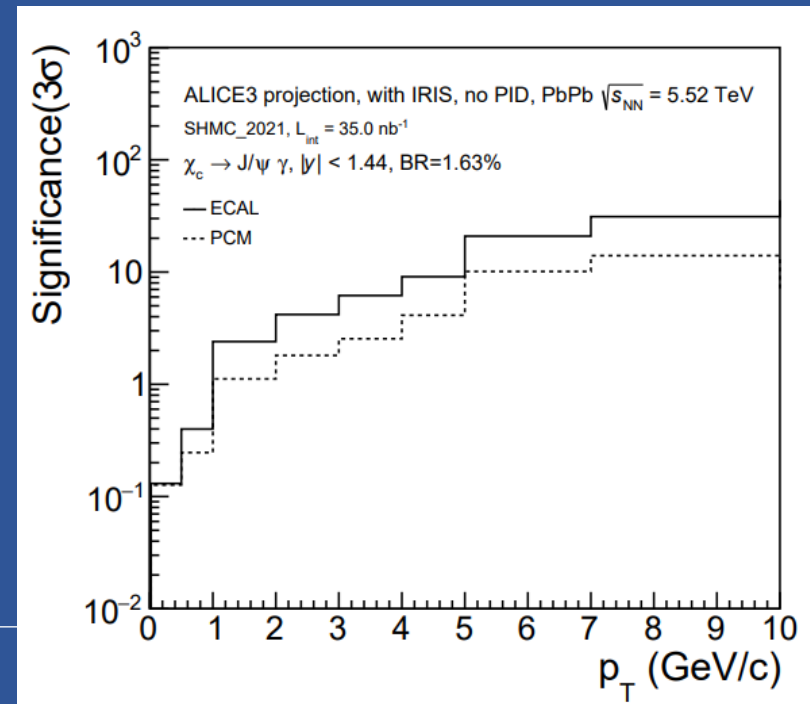
- Unique capabilities for
 - reconstruction of quarkonium states down to $p_T = 0$
 - low energy photons (0.5 GeV and below)

[LoI, arXiv:2211.02491](#)

What remains to be done ? High-energy/zero μ_B



- Aim at performing **quarkonium spectroscopy in the Quark-Gluon Plasma**
- Pseudoscalar (η_c, η_b) and P-wave (χ_c, χ_b) states largely unexplored in heavy-ion collisions
- Access
 - $\chi_c \rightarrow J/\psi \gamma, \chi_b \rightarrow \Upsilon \gamma$
 - $\eta_c \rightarrow p\bar{p}, \eta_c \rightarrow \Lambda\bar{\Lambda}$ (performance under study)



$$L_{\text{int}} = 35 \text{ nb}^{-1}$$

Good significance
for χ_c down to
 $p_T \sim 2 \text{ GeV/c}$

[LoI, arXiv:2211.02491](#)

Exotica

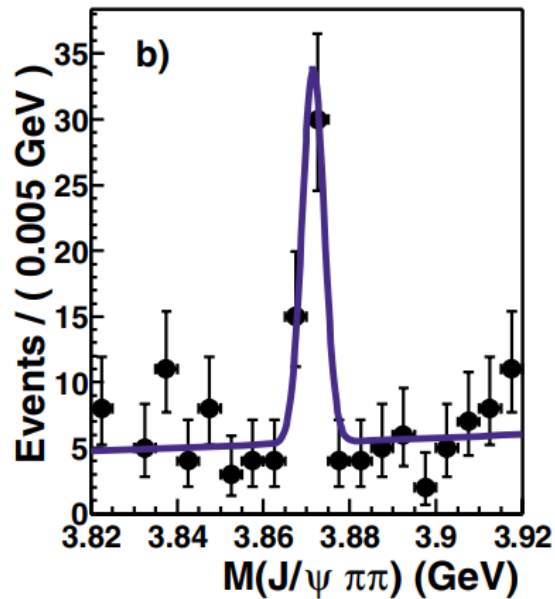
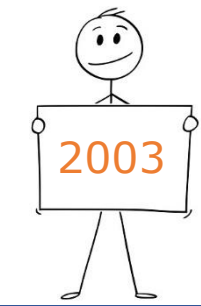
- Can non- $q\bar{q}$ /non qqq states tell us something on QGP properties ?
- Can we learn something on their structure by producing them in HI collisions?

Where all started...

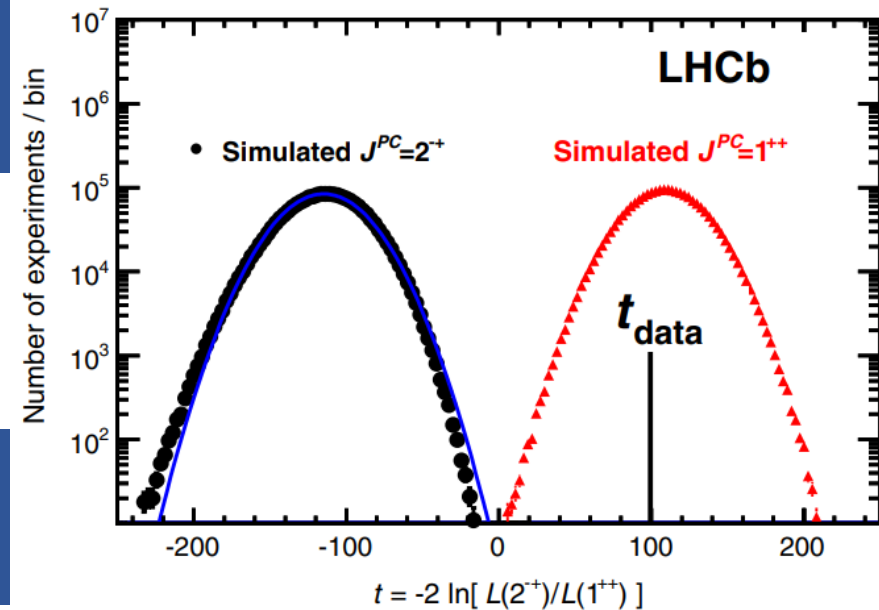
Observation of a Narrow Charmoniumlike State in Exclusive $B^\pm \rightarrow K^\pm \pi^+ \pi^- J/\psi$ Decays

S.-K. Choi *et al.* (Belle Collaboration)

Phys. Rev. Lett. **91**, 262001 – Published 23 December 2003



Among the remaining possibilities are the $\chi_{c1}(2^3P_1)$ charmonium disfavored by the value of the $X(3872)$ mass [34], and unconventional explanations such as a $D^{*0}\bar{D}^0$ molecule [8], tetraquark state [9], or charmonium-molecule mixture [10].



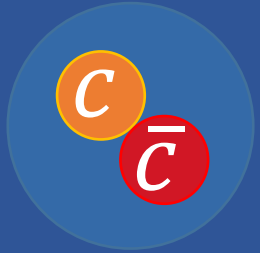
Determination of the $X(3872)$ Meson Quantum Numbers

R. Aaij *et al.* (LHCb Collaboration)

Phys. Rev. Lett. **110**, 222001 – Published 29 May 2013

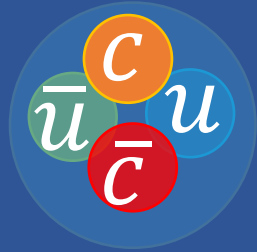
Can HI collisions help us decipher its nature?

charmonium



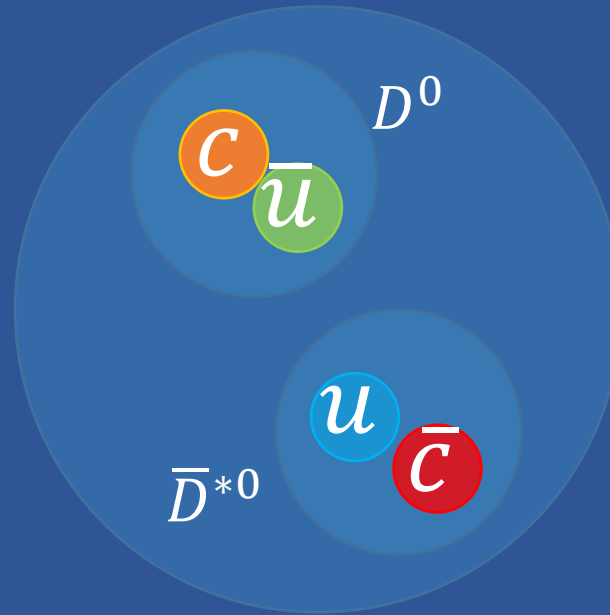
wrong mass
predicted with
 $J^{PC} = 1^{++}$

tetraquark



$r \sim 0.3 \text{ fm}$

$D^0 - \bar{D}^{*0}$ molecule

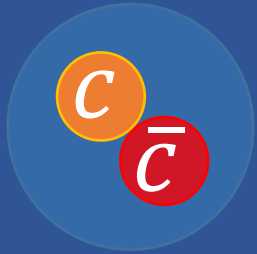


$r > 5 \text{ fm}$, small binding energy

➤ **Production in a QCD medium** might provide insight on its inner structure?

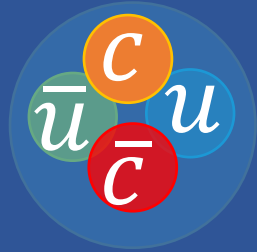
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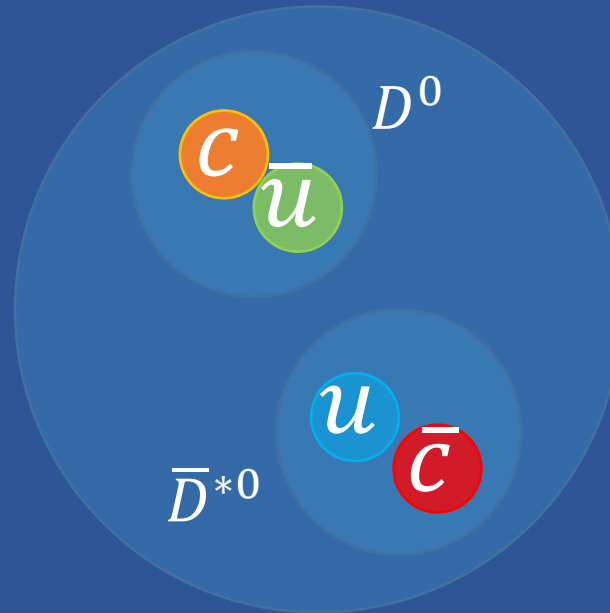
wrong mass
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tetraquark

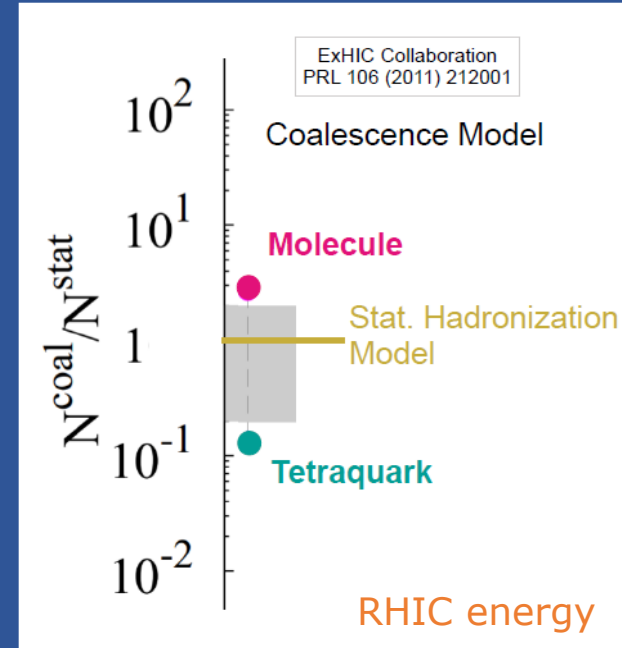


$r \sim 0.3 \text{ fm}$

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$r > 5 \text{ fm}$, small binding energy



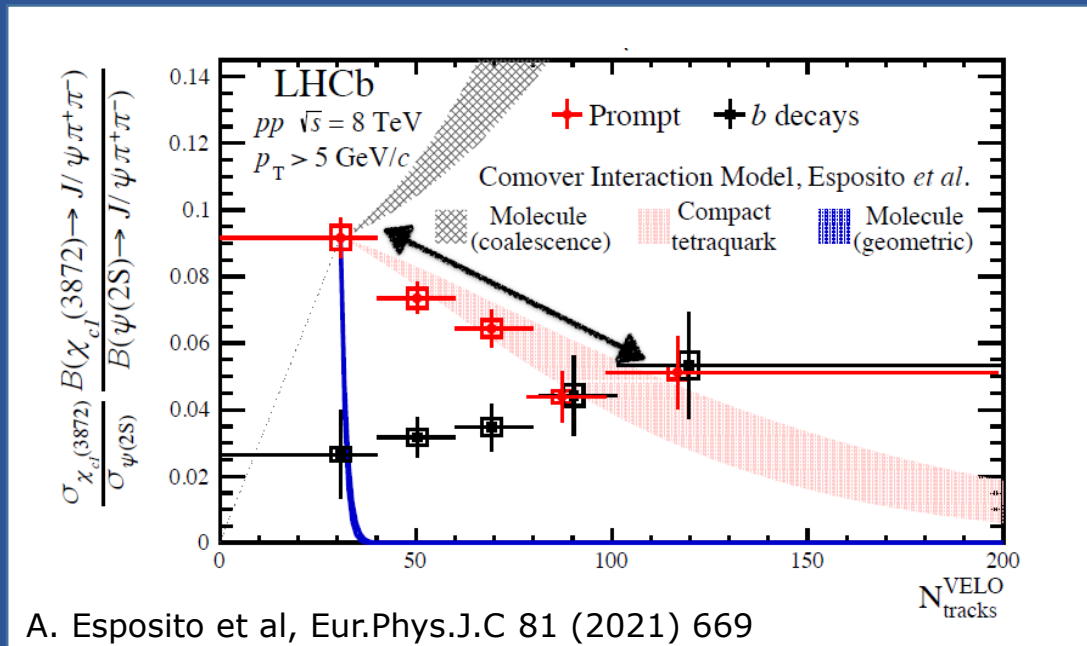
We conclude from the above discussions that the yield of a hadron in relativistic heavy ion collisions reflects its structure and thus can be used as a new method to discriminate the different pictures for the structures of multi-quark hadrons.

- **Production in a QCD medium** might provide insight on its inner structure?
- Early coalescence-based models predicted **lower yields for a compact multiquark state**

X(3872): yield vs multiplicity in pp

- At the LHC, high-multiplicity pp collisions create a dense hadronic environment
- LHCb studied the **ratio $X(3872)/\psi(2S)$ as a function of hadronic multiplicity**

LHCb, PRL 126 (2021) 092001 (2021)

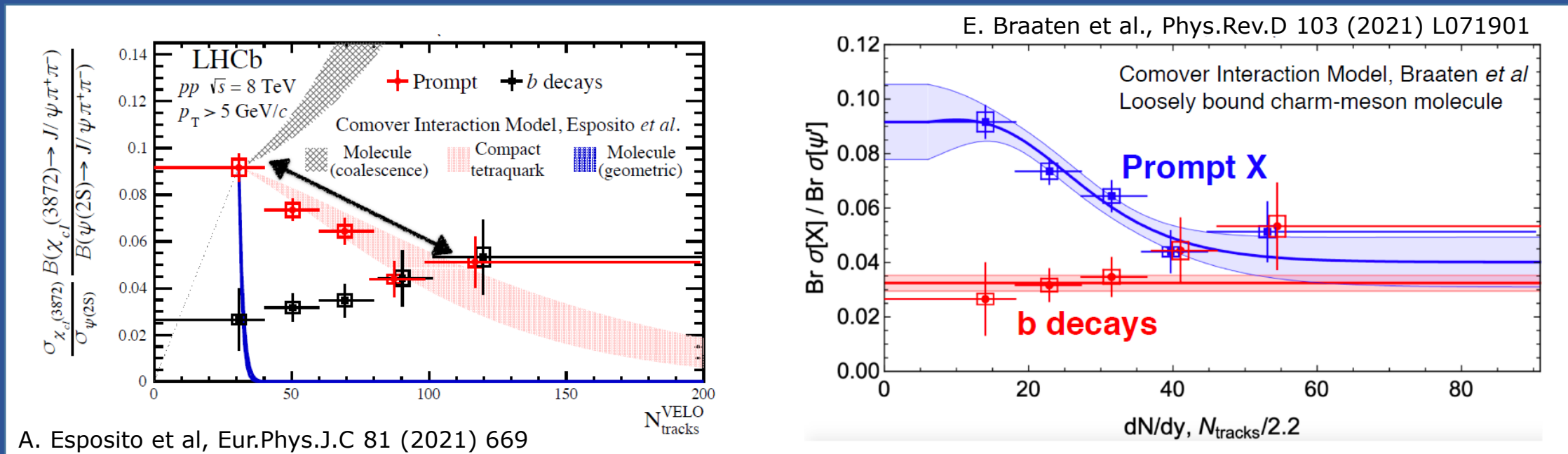


- Data described by comover interaction model assuming X(3872) to be a **tetraquark**
→ breakup reaction rate approximated by the geometric cross section

X(3872): yield vs multiplicity in pp

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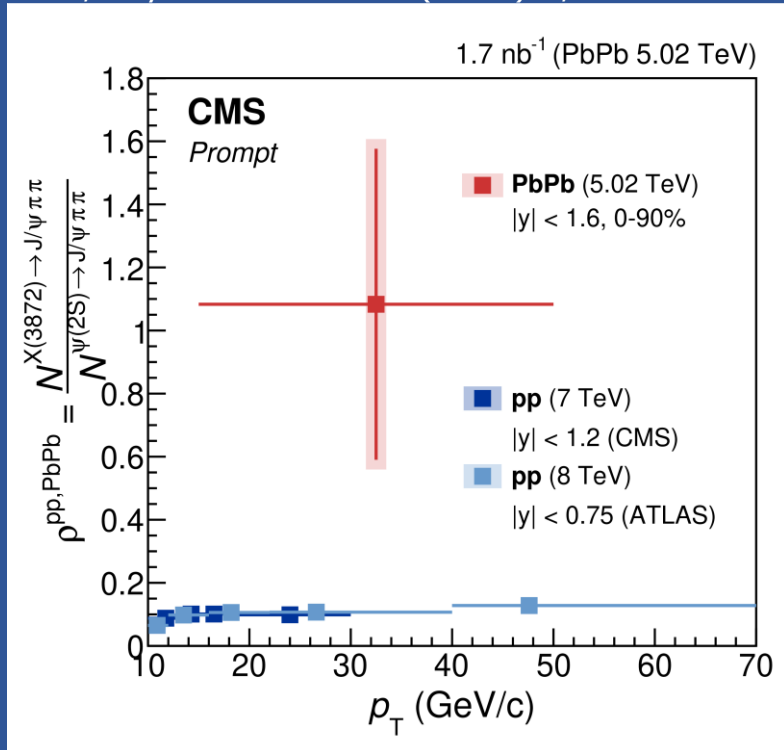
LHCb, PRL 126 (2021) 092001 (2021)



- Data described by comover interaction model assuming X(3872) to be a **tetraquark**
 - breakup reaction rate approximated by the geometric cross section
- However, using a different ansatz for CIM can also favour X(3872) being a **molecule**
 - scattering of comoving pions from the charm-meson constituents of X(3872) (no coalescence effects assumed)

X(3872): first measurement in Pb-Pb

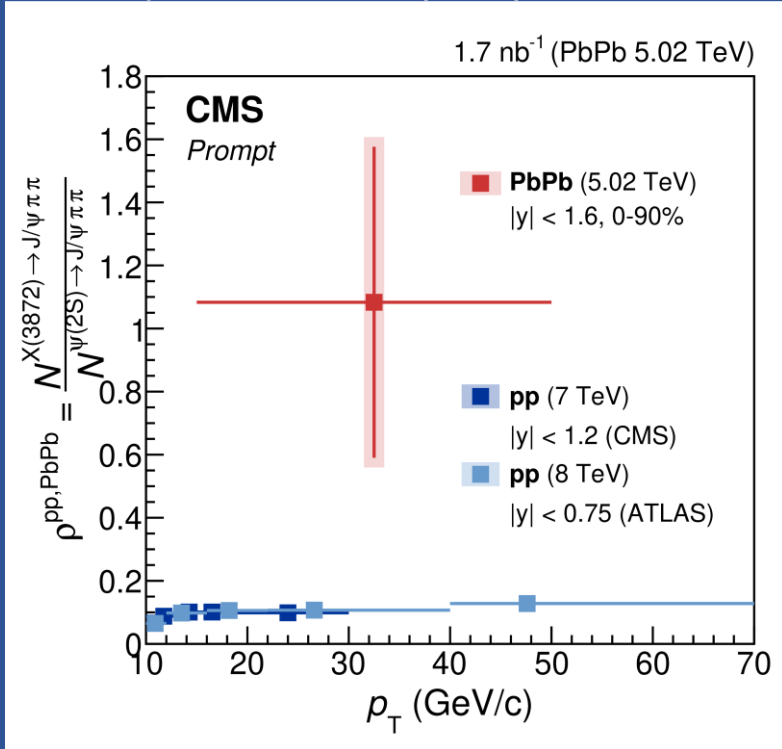
CMS, Phys.Rev.Lett. 128 (2022) 3, 032001



- **Hint of prompt X(3872) to $\psi(2S)$ enhancement in Pb-Pb, at very high p_T ($15 < p_T < 50$ GeV/c)**

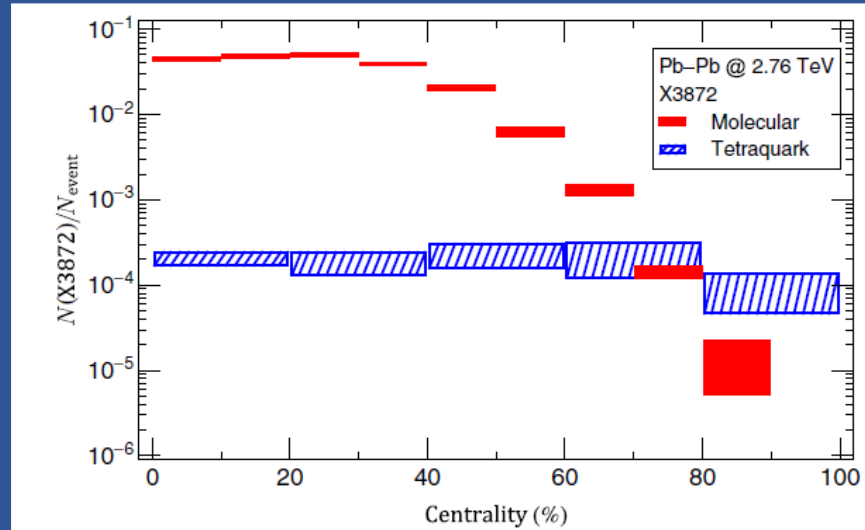
X(3872): first measurement in Pb-Pb

CMS, Phys.Rev.Lett. 128 (2022) 3, 032001



- **Hint of prompt X(3872) to $\psi(2S)$ enhancement in Pb-Pb, at very high p_T ($15 < p_T < 50$ GeV/c)**

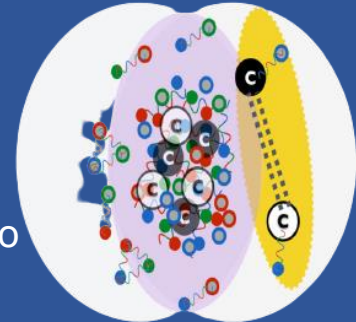
H. Zhang et al., PRL 126(2021) 012301



colored force between a color antitriplet diquark cq and a color triplet antidiquark $\bar{c}\bar{q}$



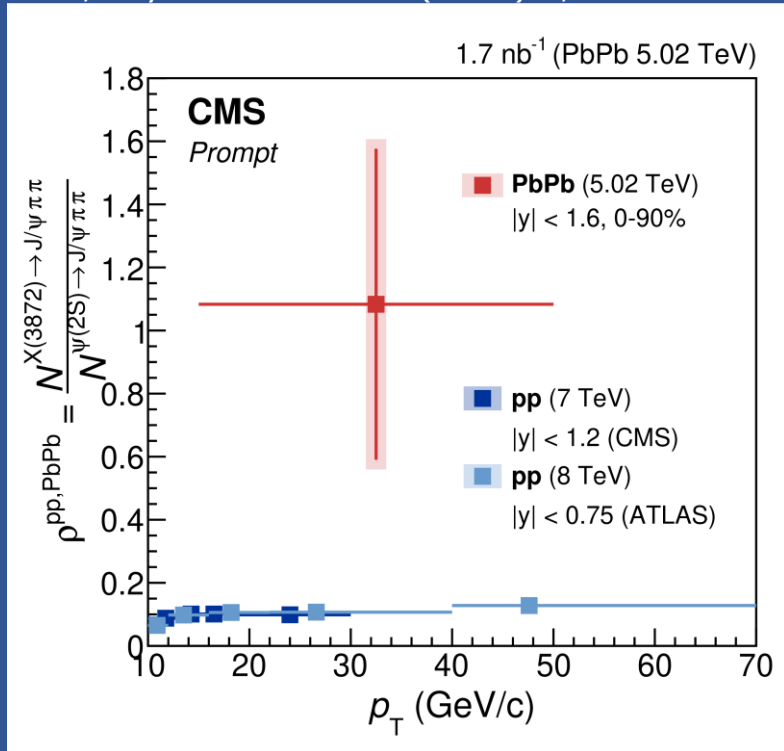
Coalescence of two charmed mesons



- **Coalescence model (AMPT): much larger yields for molecular option, with strong centrality dependence (ccbar more likely separated in space at freeze-out)**

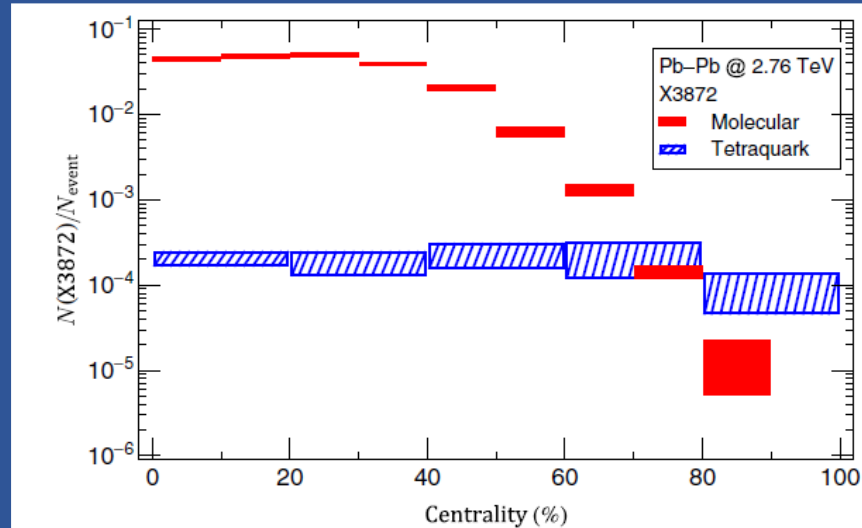
X(3872): first measurement in Pb-Pb

CMS, Phys.Rev.Lett. 128 (2022) 3, 032001

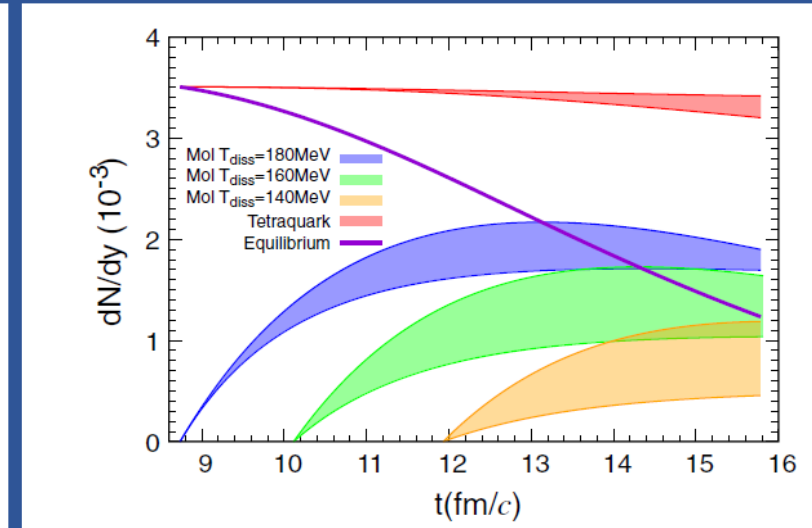


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H. Zhang et al., PRL 126(2021) 012301

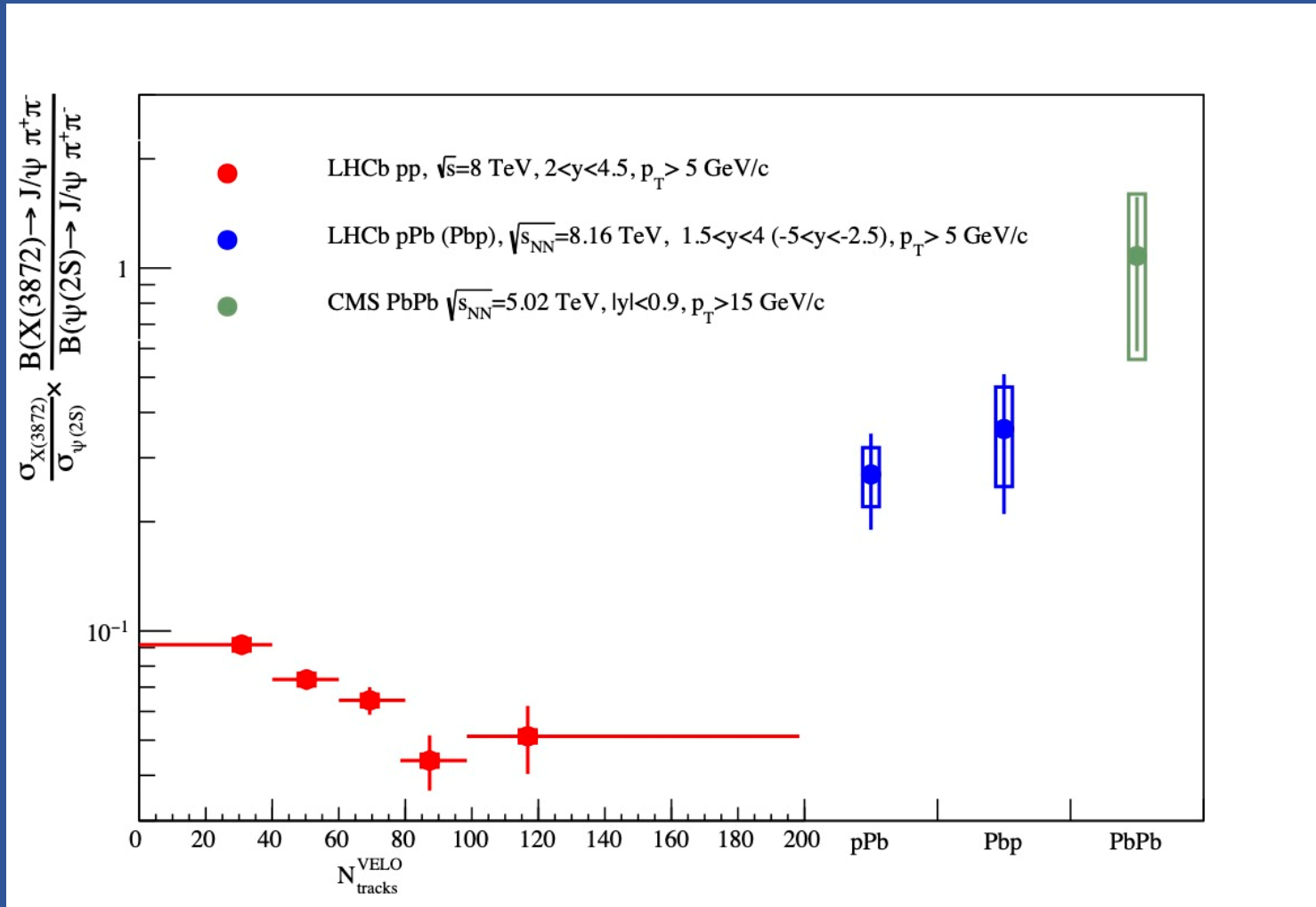


B. Wu et al., EPJA 57(2021) 122



- **Coalescence model (AMPT): much larger yields for molecular option**, with strong centrality dependence (ccbar more likely separated in space at freeze-out)
- **Transport model: moderate difference between yields**, larger reaction rates associated with the loosely bound molecule structure imply that it is formed later in the fireball evolution than the tetraquark and thus its final yields are generally smaller

X(3872): current experimental status

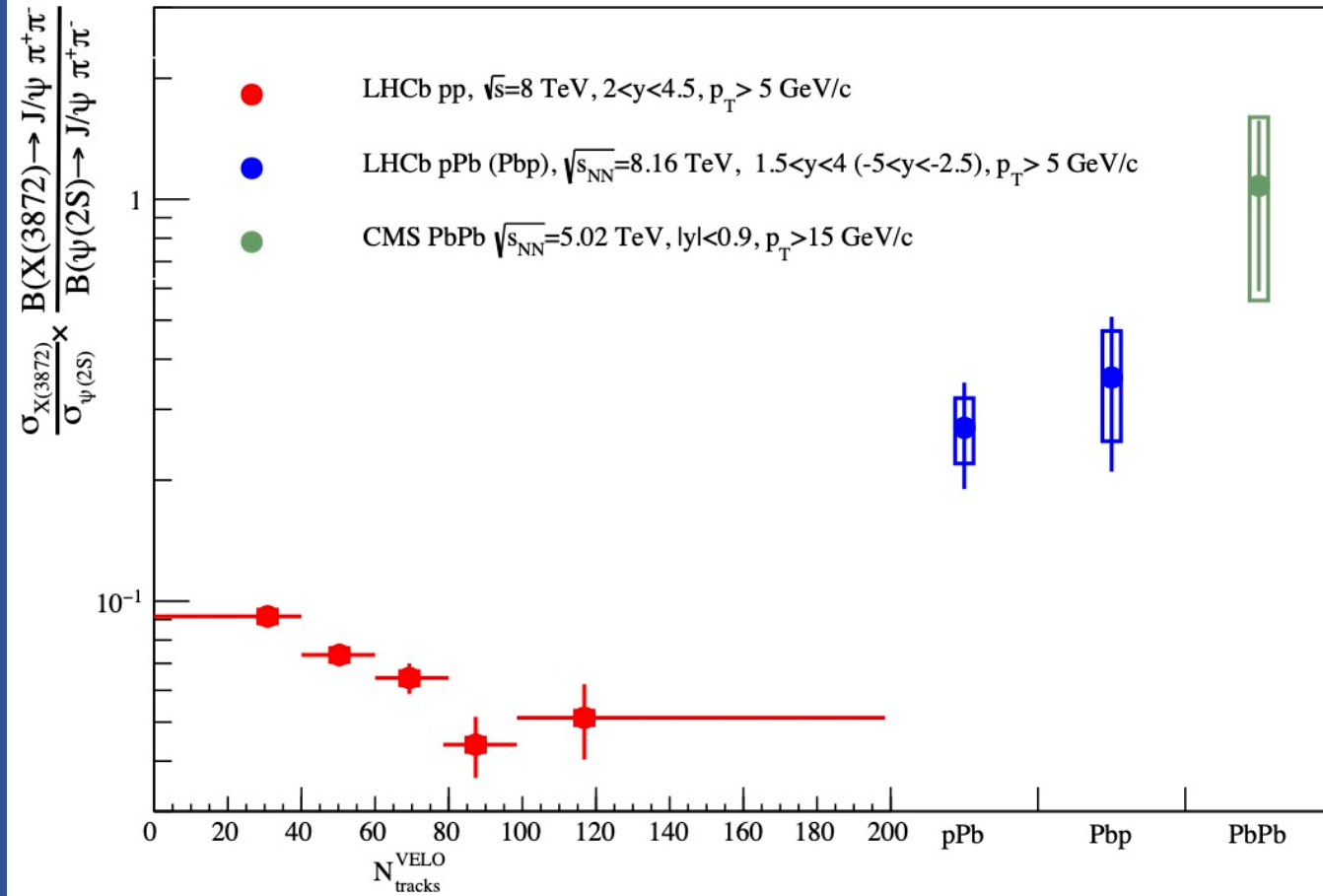


pp \rightarrow p-Pb \rightarrow Pb-Pb

from **suppression** to **enhancement**?

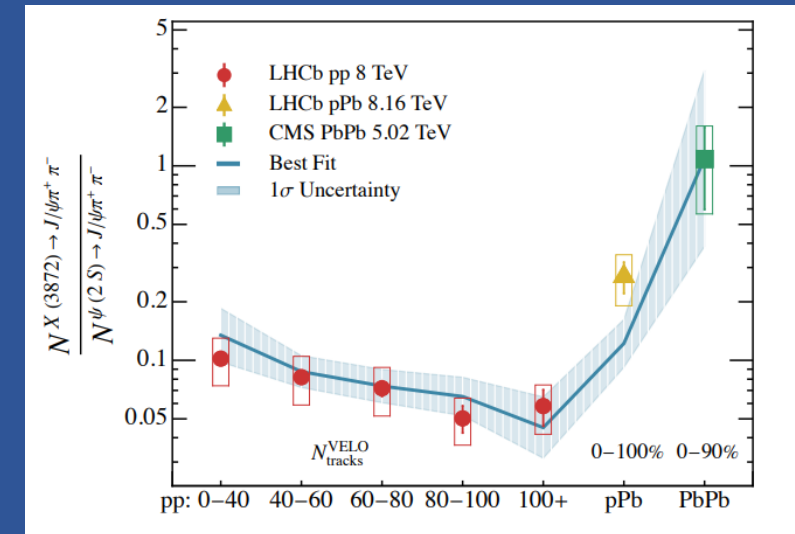
➤ Extension of measurements toward low p_T badly needed \rightarrow **LHC run 3/4**

X(3872): current experimental status



pp \rightarrow p-Pb \rightarrow Pb-Pb

from **suppression** to **enhancement**?



First attempts at a coherent description of yields vs system size

➤ Extension of measurements toward low p_T badly needed \rightarrow **LHC run 3/4**

Conclusions

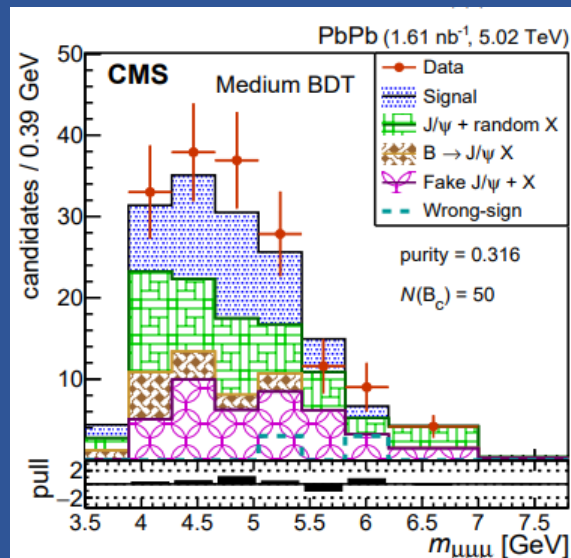
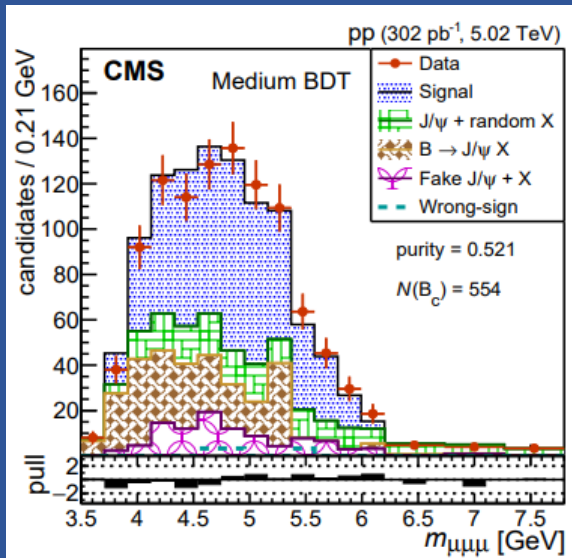
- Heavy quarks: great sensitivity to the medium, spectacular phenomenology!
- Quarkonia
 - Lots of results on vector states
 - P-wave and pseudoscalar states almost completely unexplored
 - B_c needs further investigation
 - Related observables (not covered today): polarization,...
- Exotica
 - Only the surface was scratched
 - Beyond $X(3872)$: $T_{cc}^+(ccud)$, $X(6900)$ ($cccc\bar{c}$),...
- Open heavy-flavors
 - A world on its own... sorry for not discussing them today!



Looking forward to new exciting results at QM2023

A “hybrid” quarkonium state: B_c^+

- Binding energy intermediate between J/ψ and $\Upsilon(1S)$, can be dissociated in the QGP
- **Regeneration** effects **could be important** (small $\sigma_{pp}^{B_c}$, large charm multiplicity in Pb-Pb)
- Energy loss: study mass and color-charge dependence
- First measurement by CMS in Pb-Pb collisions via $B_c^+ \rightarrow (J/\psi \rightarrow \mu\mu) \mu^+ \nu_\mu$ (displaced vertex of 3 muons, with OS pair in the J/ψ region)
- ✓ Needs good understanding of background in $3.2 < M_{\mu\mu} < 6.3$ GeV
→ Use **BDT technique**

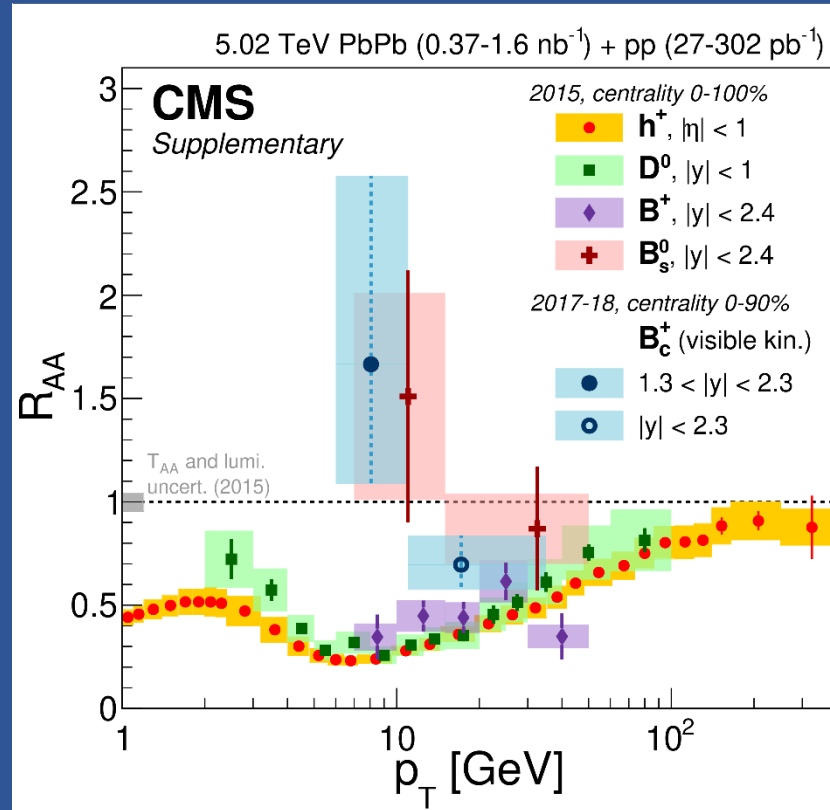
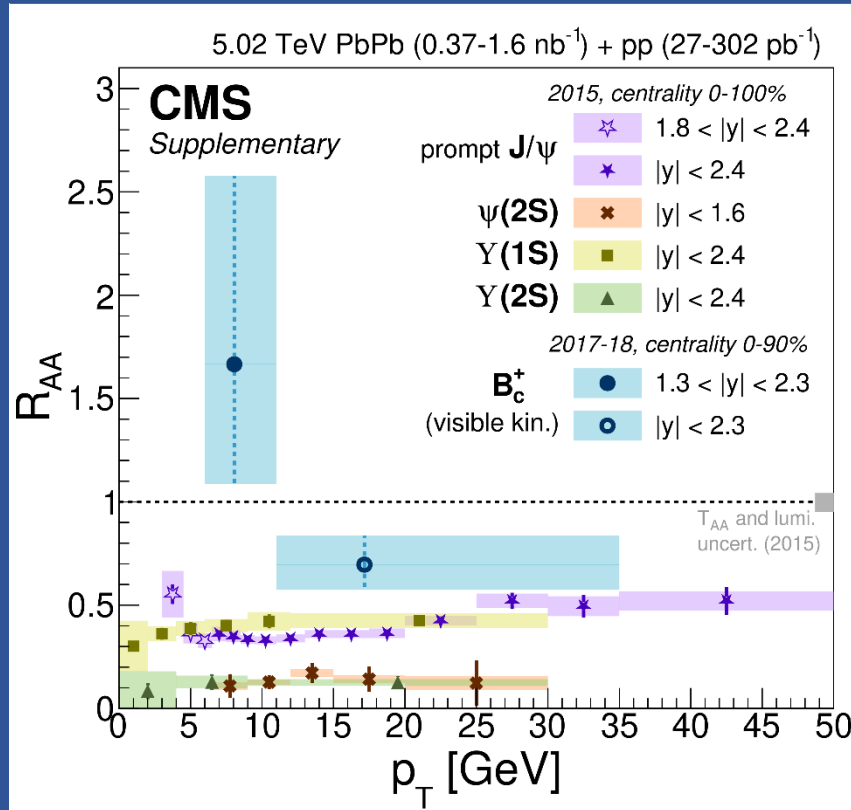


CMS, Phys. Rev. Lett. 128 (2022) 252301

Significance in Pb-Pb
well **above 5σ**

Fake J/ψ : OS muons not coming from J/ψ (sidebands)
B decays: $B \rightarrow J/\psi + \mu$ from same vertex (simulation)
 $J/\psi + \text{random muon}$

A “hybrid” quarkonium state: B_c^+

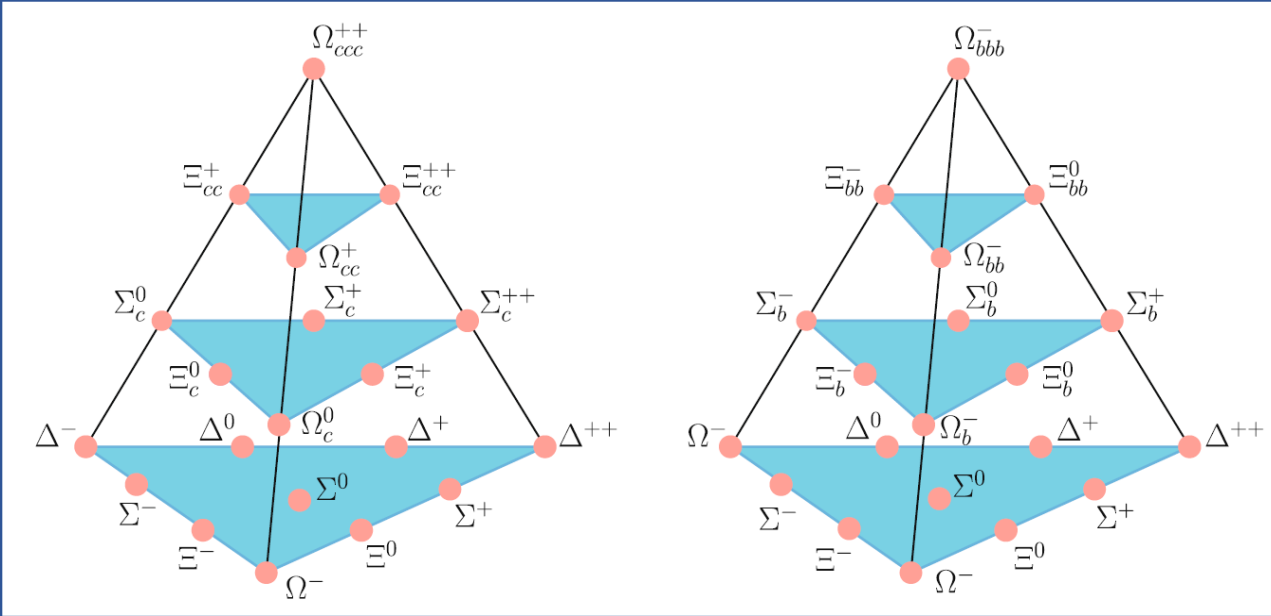


- Reminiscent of J/ψ **behaviour**, but **larger R_{AA}** values
- **High- p_T** region likely sensitive to **energy loss** effects too
- Very promising channel in view of **higher luminosity** data samples

- Hint for a p_T dependence of R_{AA}
→ from **enhancement to suppression** when increasing p_T (1.6 σ effect)

- Other heavy mesons typically show more suppression, may indicate **recombination** as a significant production effect

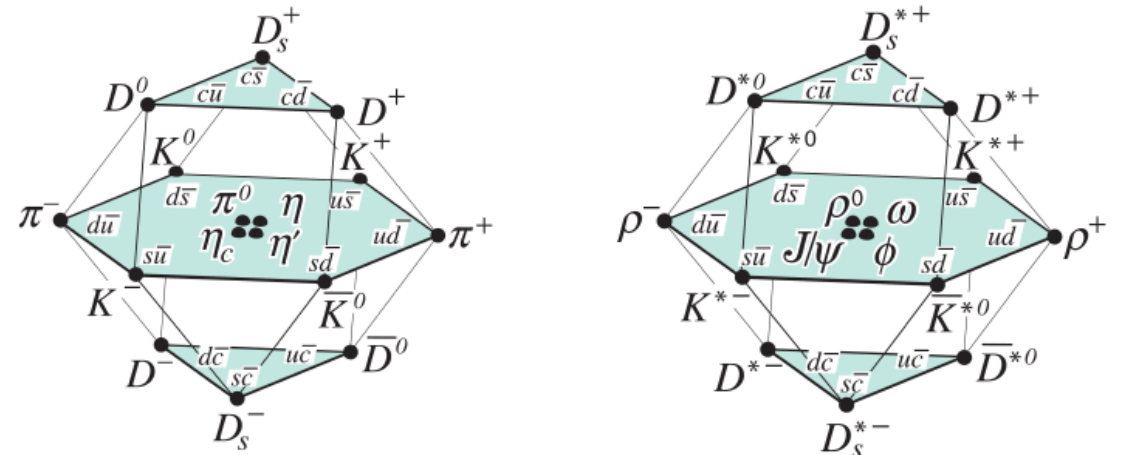
Charm/beauty mesons/baryons



(from Huang et al. EPJC (2021) 81:276)

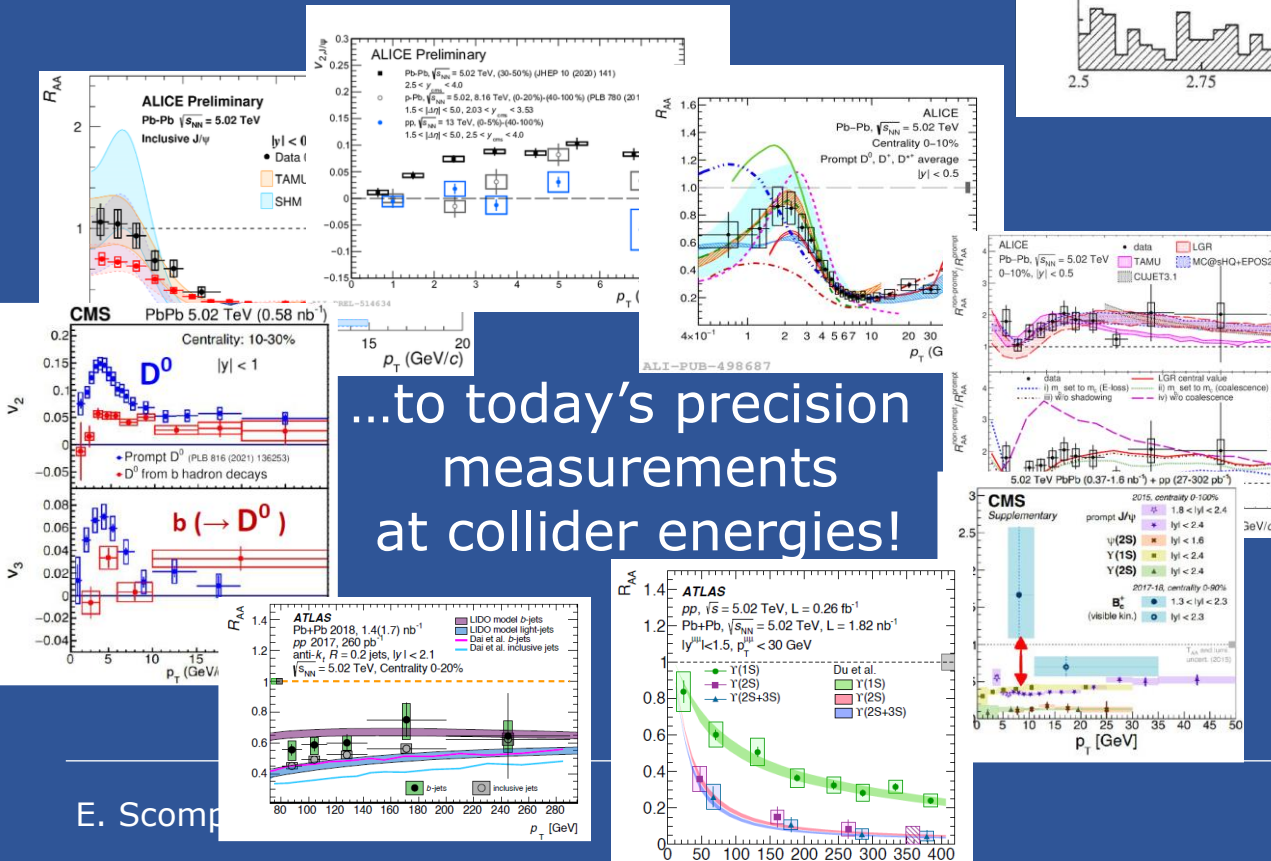
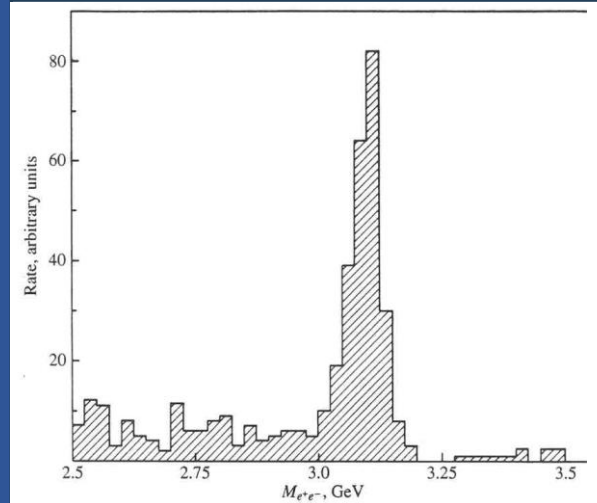
- Give rise to a large variety of particles
- Investigate hadron formation processes
- Baryons vs mesons
- Multi-heavy quark states

- Charm and beauty quarks are created in the early stages and probe the QGP phase
 - Energy loss
 - Thermalization



Heavy quarks: a (the) golden probe of QGP

From the November revolution...



...to today's precision measurements at collider energies!

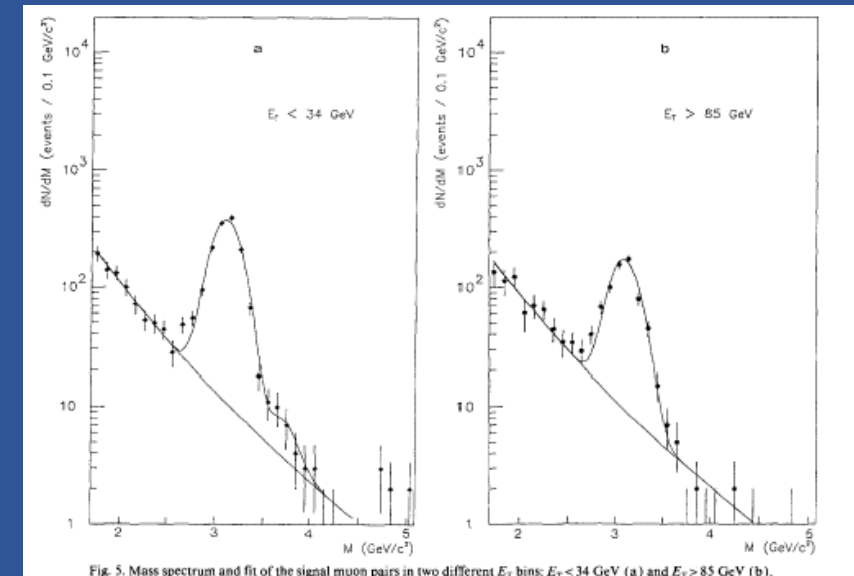
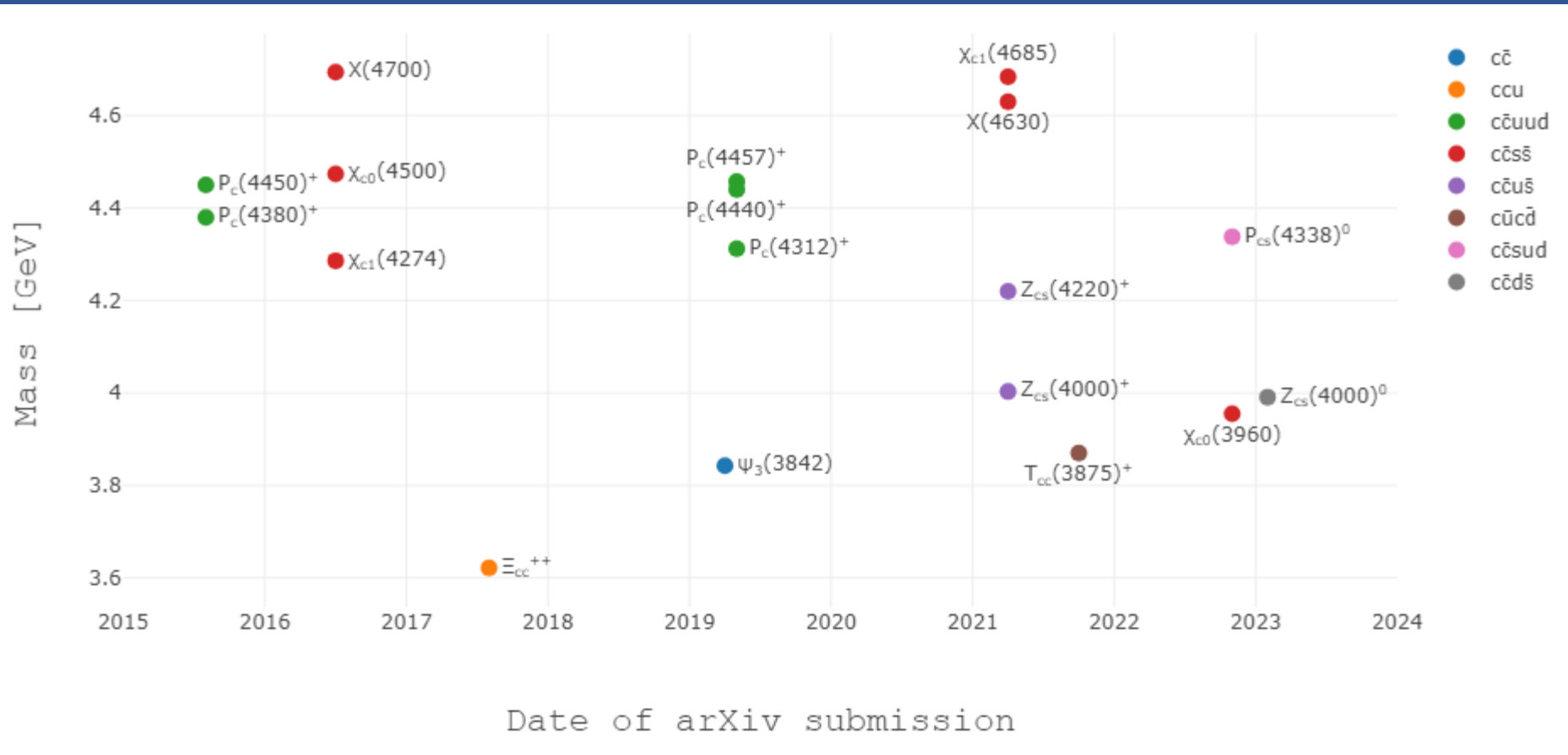


Fig. 5. Mass spectrum and fit of the signal muon pairs in two different E_T bins: $E_T < 34$ GeV (a) and $E_T > 85$ GeV (b).

... to the discovery of the J/ψ suppression...

Beyond $q\bar{q}$ and qqq



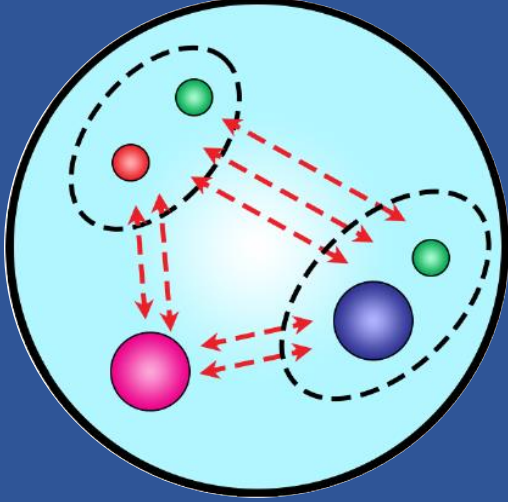
Recent years have seen a flourishing of discoveries of new states that do not fit inside the qqq or $q\bar{q}$ category

Here the states **discovered by LHCb**

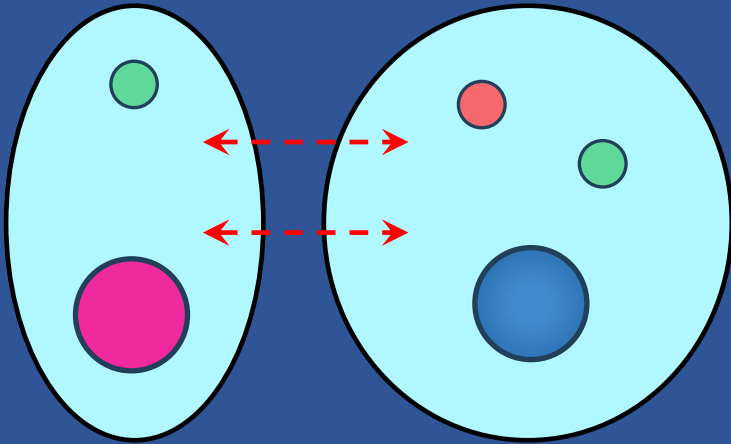
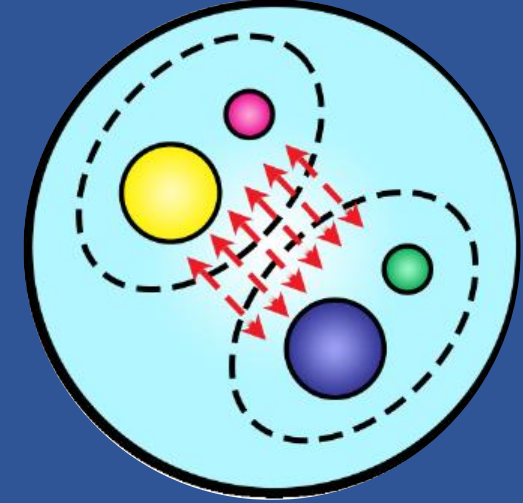
➤ Strong contributions from **CMS, ATLAS, BELLE, BES III** as well

Beyond $q\bar{q}$ and qqq

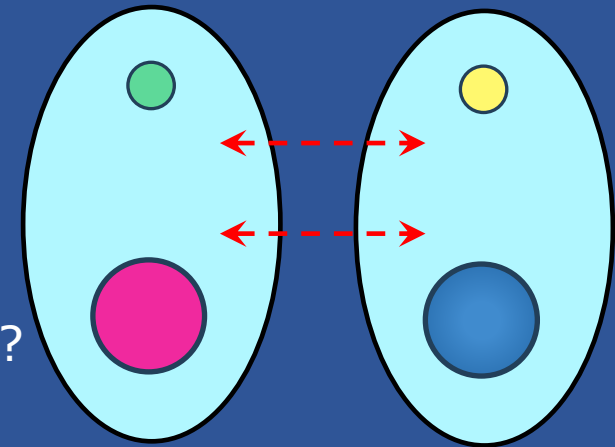
From Chen et al.,
Rept. Prog. Phys. 86 (2023) no.2, 026201



Compact multiquark states are tightly bound by the strong interaction directly, while the **hadronic molecular states** are weakly bound by the residual strong interaction



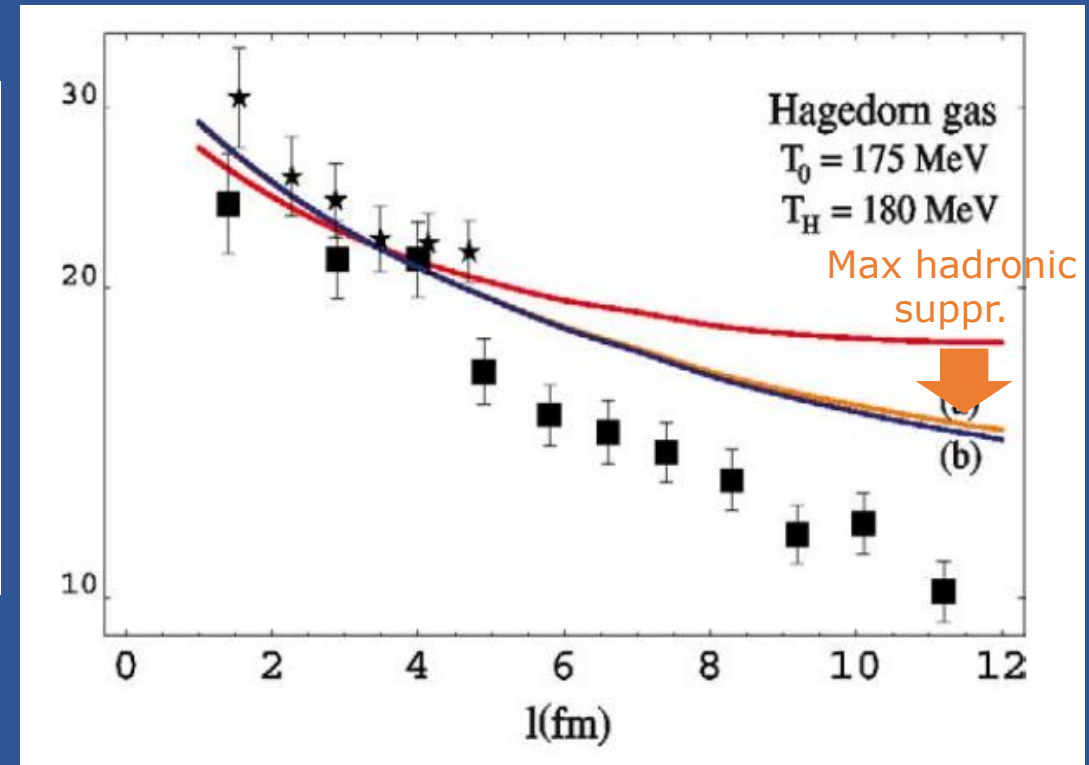
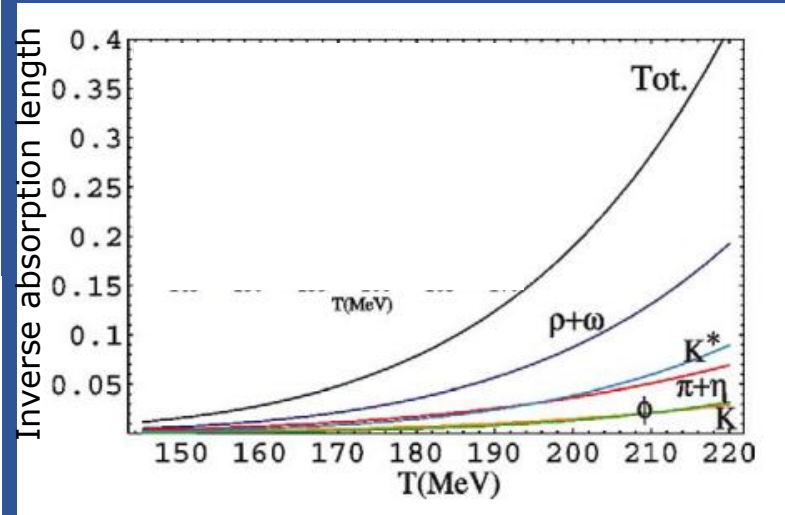
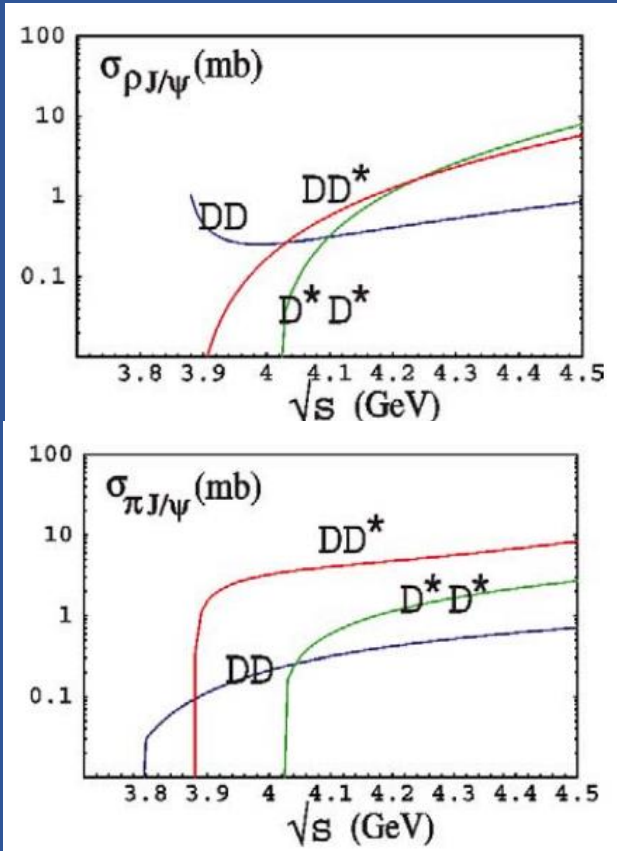
Can we investigate their nature by studying their production in HI collisions?



Can hadronic suppression really explain data ?

Wild discussions at the time...

Maiani et al., NP A 748 (2005) 209–225



The sharp rise of the degrees of freedom due to the vicinity of the Hagedorn temperature makes so that the temperature of the gas practically does not rise at all, the dissociation curve cannot become harder, and the prediction falls short from explaining the drop observed by NA50.

Sequential suppression...

| state | J/ψ | χ_c | ψ' | Υ | χ_b | Υ' | χ'_b | Υ'' |
|------------------|----------|----------|---------|------------|----------|-------------|-----------|--------------|
| mass [GeV] | 3.10 | 3.53 | 3.68 | 9.46 | 9.99 | 10.02 | 10.26 | 10.36 |
| ΔE [GeV] | 0.64 | 0.20 | 0.05 | 1.10 | 0.67 | 0.54 | 0.31 | 0.20 |
| ΔM [GeV] | 0.02 | -0.03 | 0.03 | 0.06 | -0.06 | -0.06 | -0.08 | -0.07 |
| r_0 [fm] | 0.50 | 0.72 | 0.90 | 0.28 | 0.44 | 0.56 | 0.68 | 0.78 |

charmonia

bottomonia



Binding energy

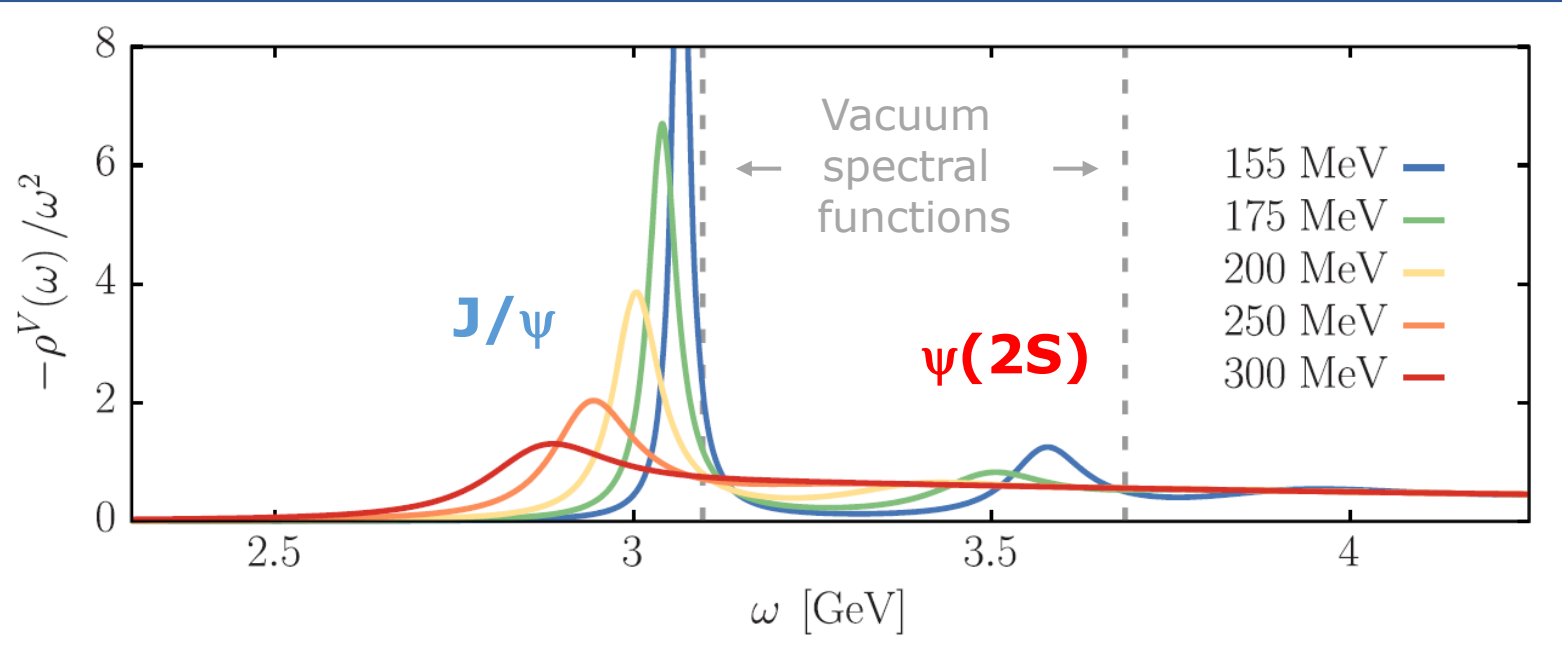


Resonance size

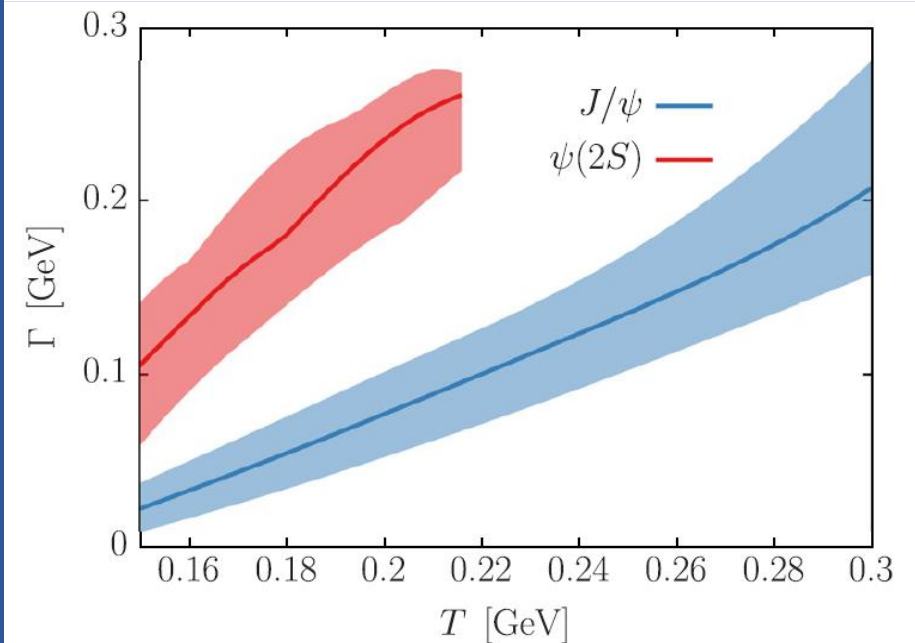
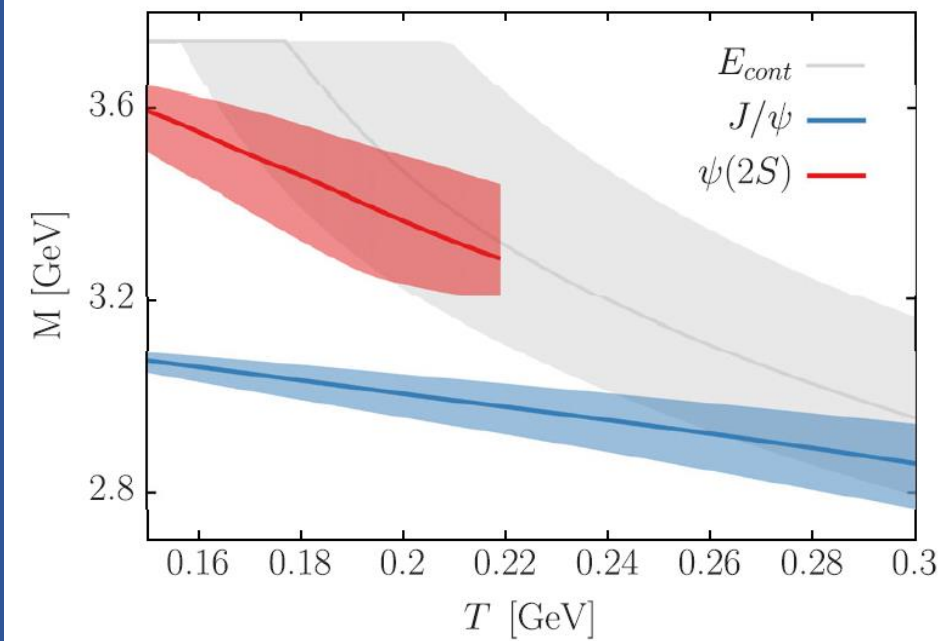
- Expect “ordering” in the suppression with weakly bound states more strongly affected

...checked with modern techniques

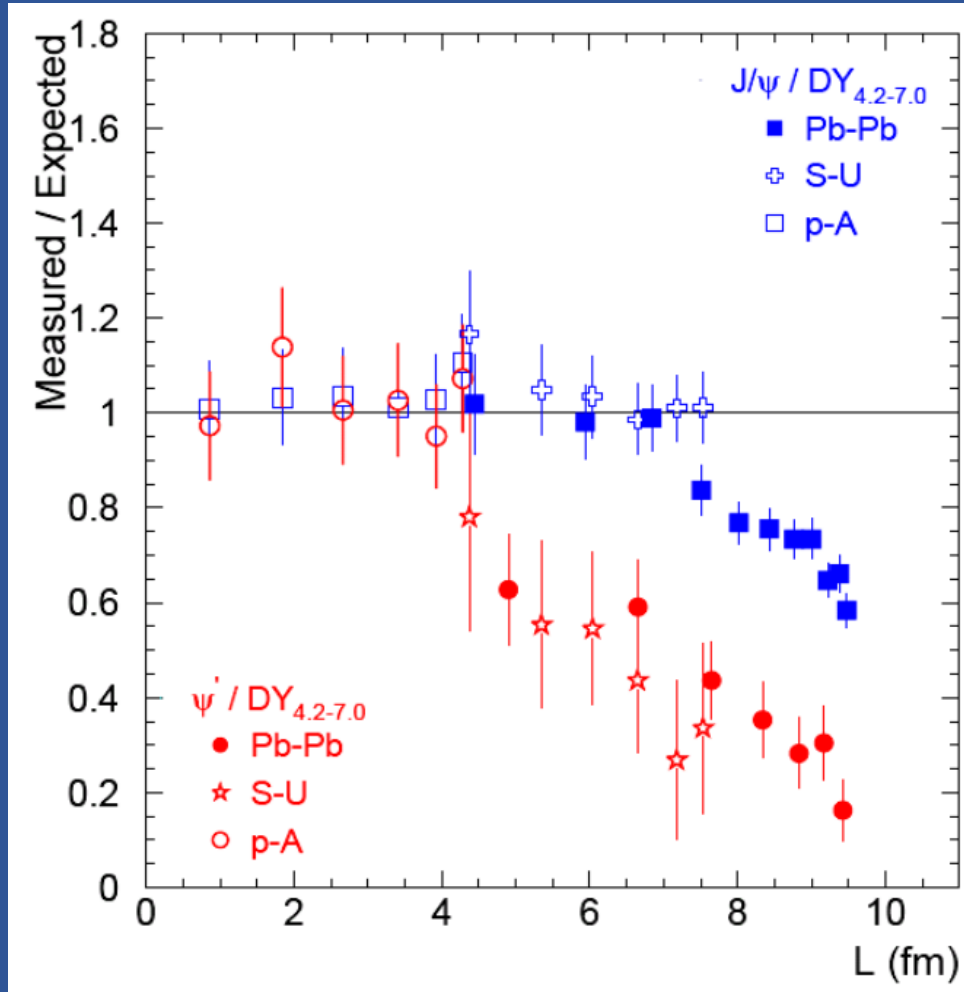
Lafferty and Rothkopf, Phys. Rev. D 101 (2020) 056010



- ❑ Strong effects on the **mass AND width** of the charmonium states, with distinctive **differences between J/ψ and $\psi(2S)$**
- ❑ As intuitively expected, the more deeply the state is bound, the less is susceptible to medium effects



“Sequential” suppression revealed: $\psi(2S)$ vs J/ψ

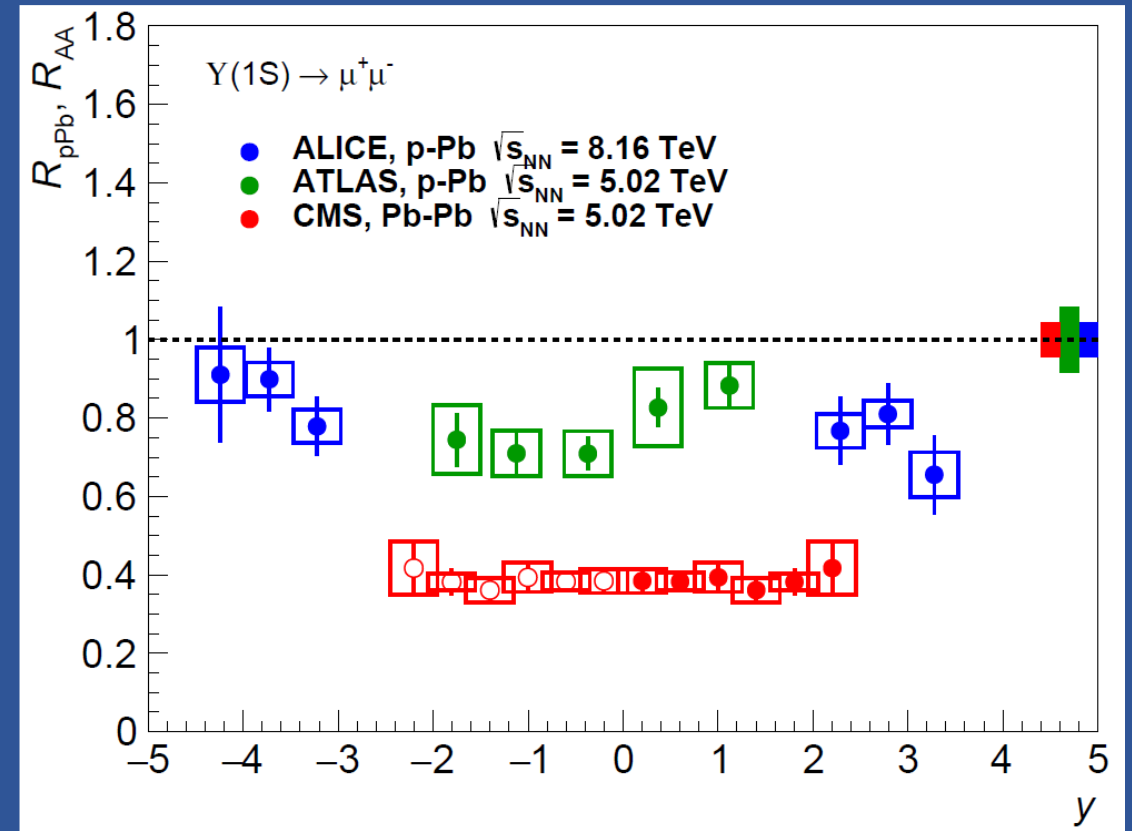
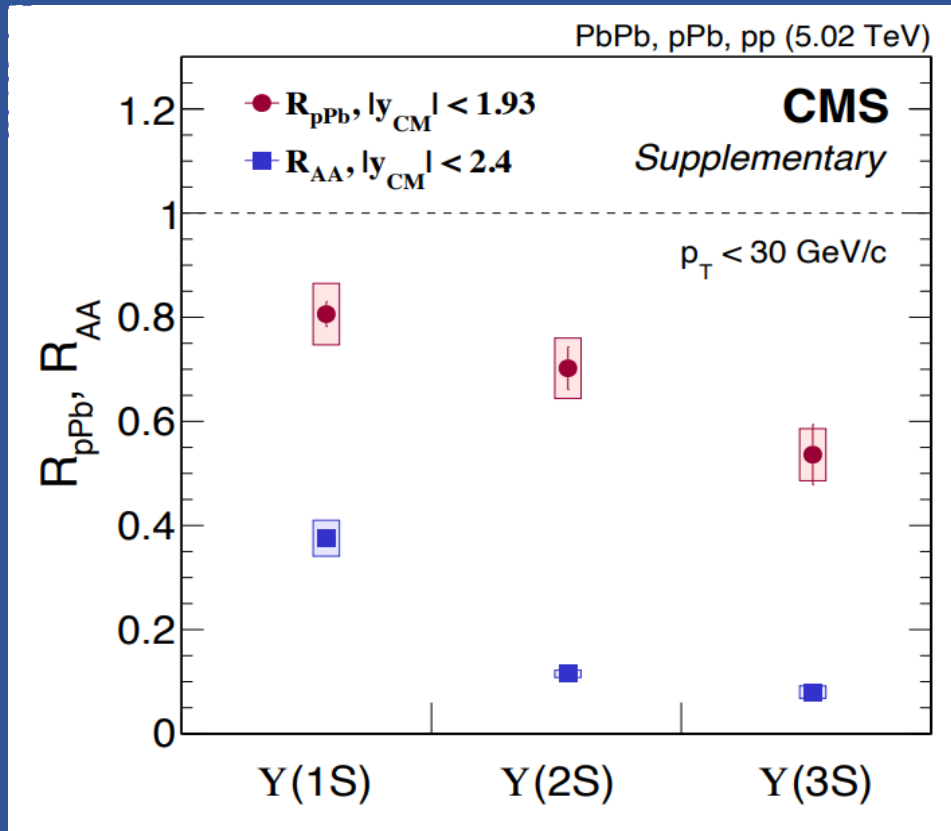


- With respect to the same reference process
- Having corrected for respective CNM effects, calibrated with p-A data



- $\psi(2S)$ suppression effects turn in for more peripheral events in a given collision system
- The effects are much stronger for $\psi(2S)$ at a given centrality

What about CNM effects ?



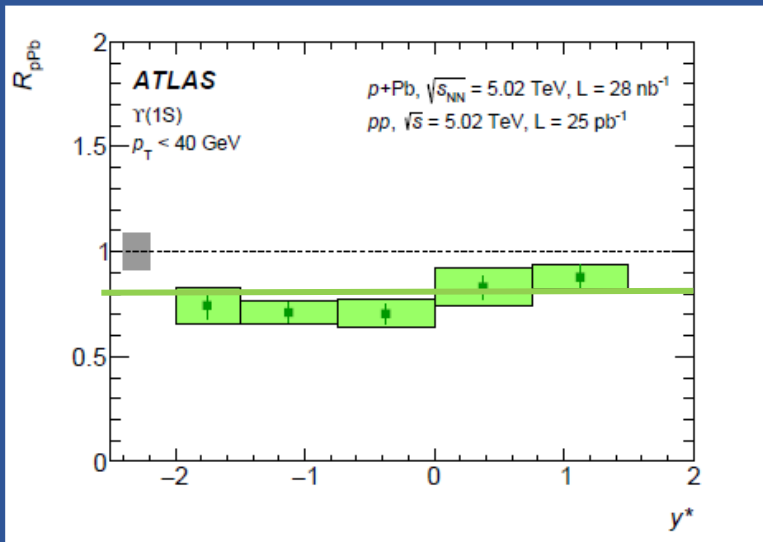
- Shadowing acts similarly for the three states
- Are we seeing final state dissociation effects?
- Nuclear suppression must be negligible
- Is a **dense system created in p-Pb?**

- Is the strong suppression of $Y(1S)$ “compatible” with its extremely large binding energy ($>1 \text{ GeV}$) ?
- Can a fraction of $Y(1S)$ suppression be due to **non-QGP effects?**

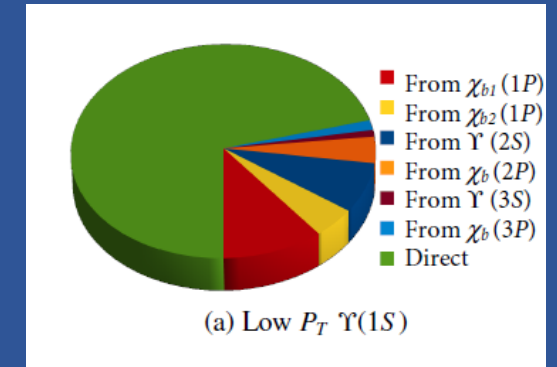
Direct $\Upsilon(1S)$ suppression ?

- Two aspects to be considered
 - Non-negligible **CNM effects**
 - **Feed-down** from S and P bottomonium states, with LHCb results implying a $\sim 30\%$ effect at (fairly) low p_T in pp

ATLAS, arxiv:1709.03089



Lansberg, arXiv:1903.09185



- ❑ Consider the $\Upsilon(1S)$ suppression seen by CMS and assume all the remaining Pb-Pb $\Upsilon(1S)$ to be direct

CMS: $R_{AA}^{\text{incl } \Upsilon(1S)} \sim 0.38$

$R_{AA}^{\text{direct } \Upsilon(1S)} \sim 0.38/0.7 = 0.54$

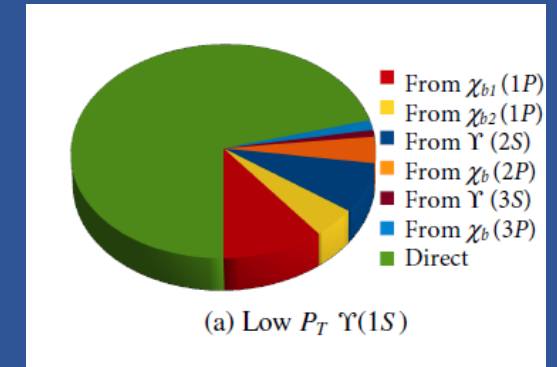
CNM effects: $(R_{pA})^2 \sim 0.7^2 \sim 0.5$

The observed **$\Upsilon(1S)$ suppression** could be compatible with a **combination of (i) 2S,3S feed-down + (ii) CNM effects**

Direct $\Upsilon(1S)$ suppression ?

- Two aspects to be considered
 - Non-negligible **CNM effects**
 - **Feed-down** from S and P bottomonium states, with LHCb results implying a $\sim 30\%$ effect at (fairly) low p_T in pp

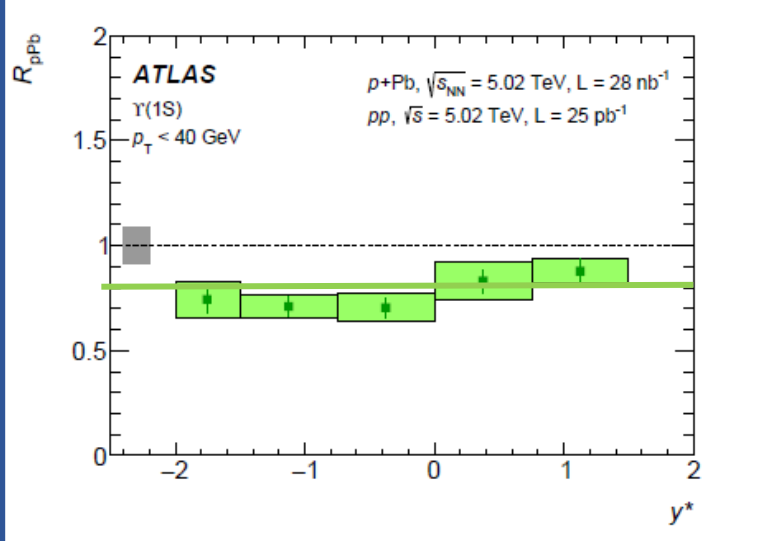
Lansberg, arXiv:1903.09185



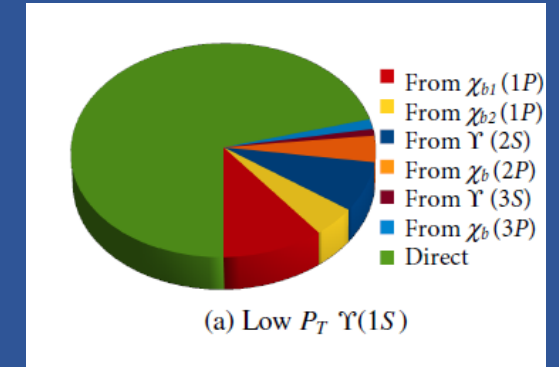
Direct $\Upsilon(1S)$ suppression ?

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ATLAS, arxiv:1709.03089



Lansberg, arXiv:1903.09185



- ❑ Consider the $\Upsilon(1S)$ suppression seen by CMS and assume all the remaining Pb-Pb $\Upsilon(1S)$ to be direct

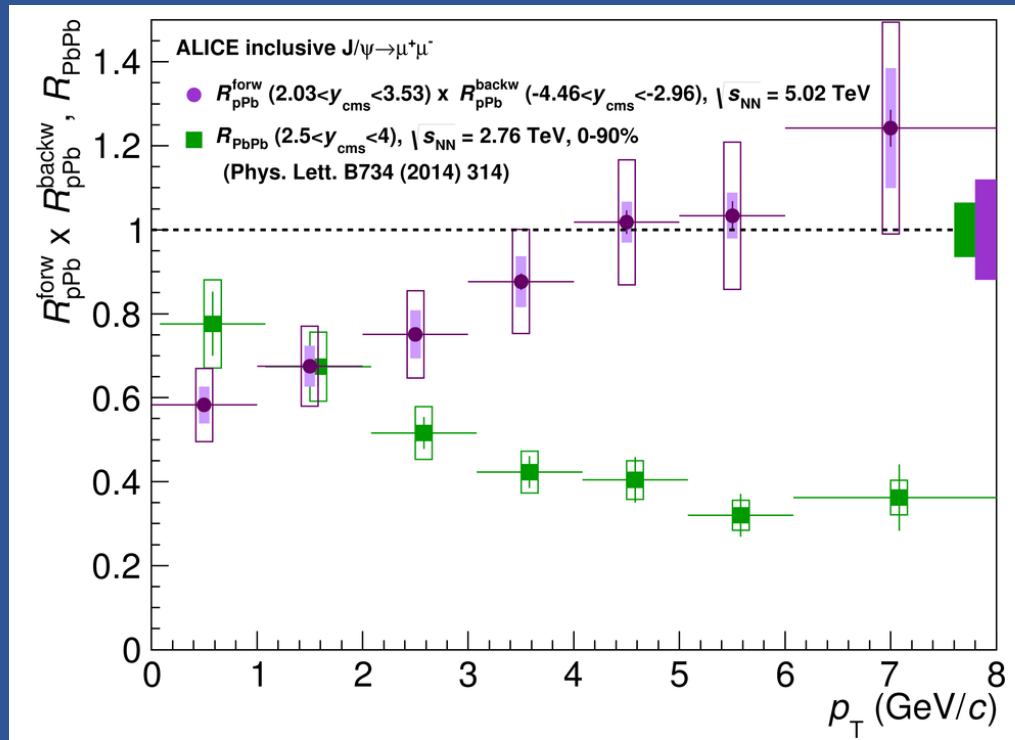
CMS: $R_{AA}^{\text{incl } \Upsilon(1S)} \sim 0.38$

$$R_{AA}^{\text{direct } \Upsilon(1S)} \sim 0.38/0.7 = 0.54$$

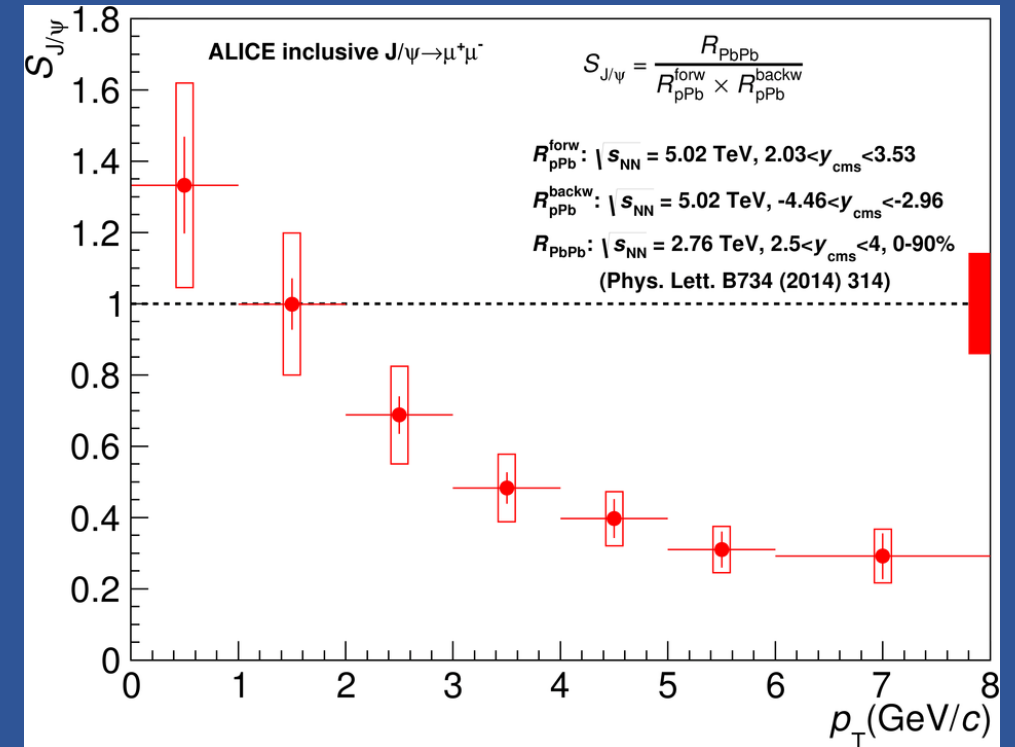
CNM effects: $(R_{pA})^2 \sim 0.7^2 \sim 0.5$

Do non-QGP effects matter ?

- Simple arguments show this is the case



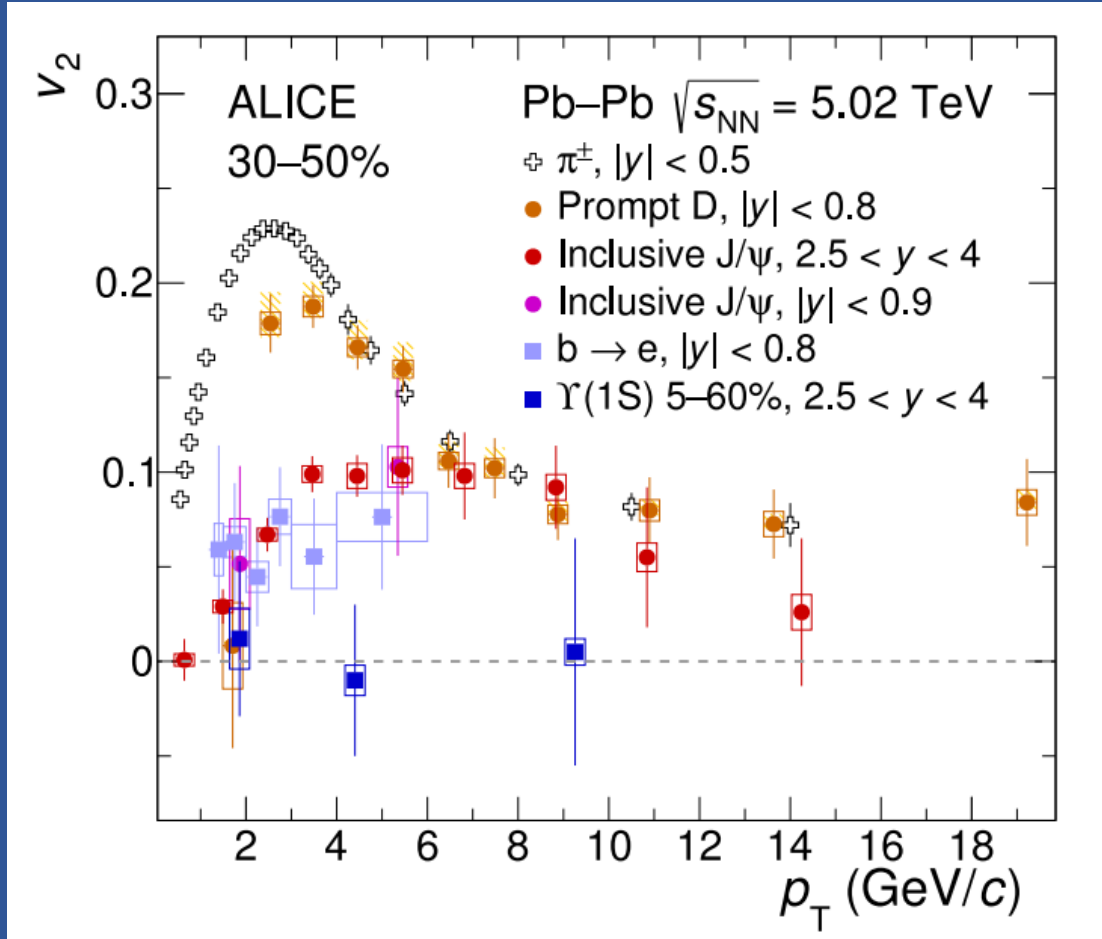
- Compare the measured R_{AA} with the product of backward and forward R_{pPb}



- Strong effect in particular at low p_T
- **J/ψ enhancement!**

Beyond suppression/regeneration \rightarrow flow

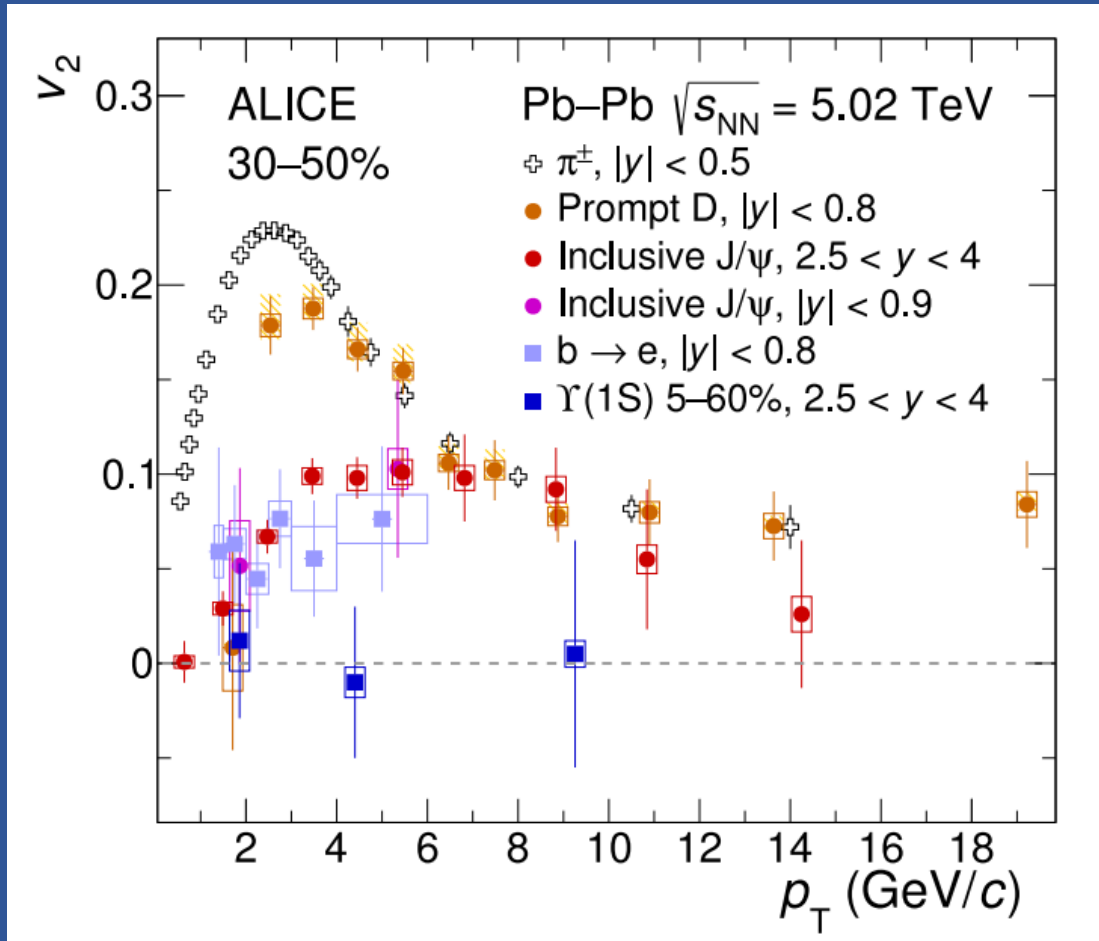
ALICE, arXiv:2211.04384



- **Quark flavour hierarchy** observed in the low- p_T range

Beyond suppression/regeneration \rightarrow flow

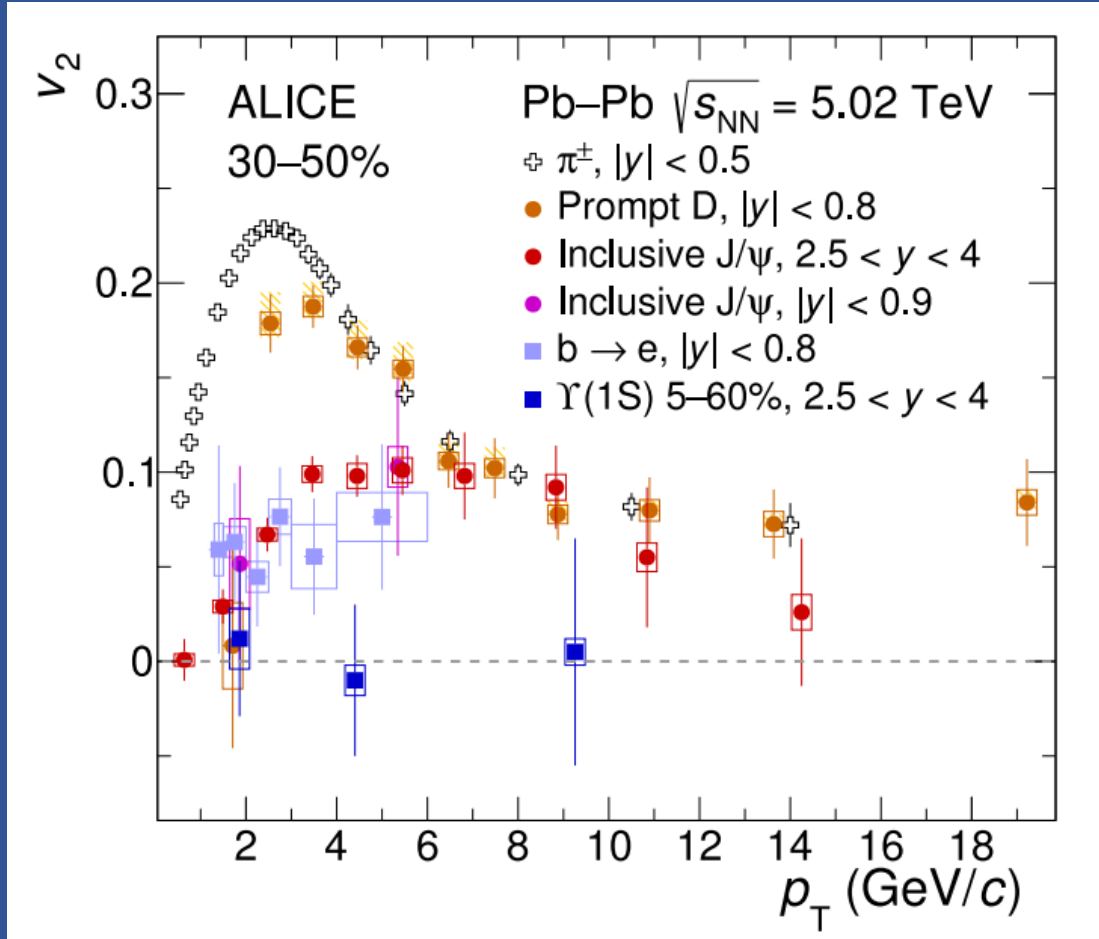
ALICE, arXiv:2211.04384



- **Quark flavour hierarchy** observed in the low- p_T range
- Both open and hidden charm hadrons show a significant amount of anisotropic flow
 \rightarrow **charm quarks are at least partly thermalised in the QGP medium**

Beyond suppression/regeneration \rightarrow flow

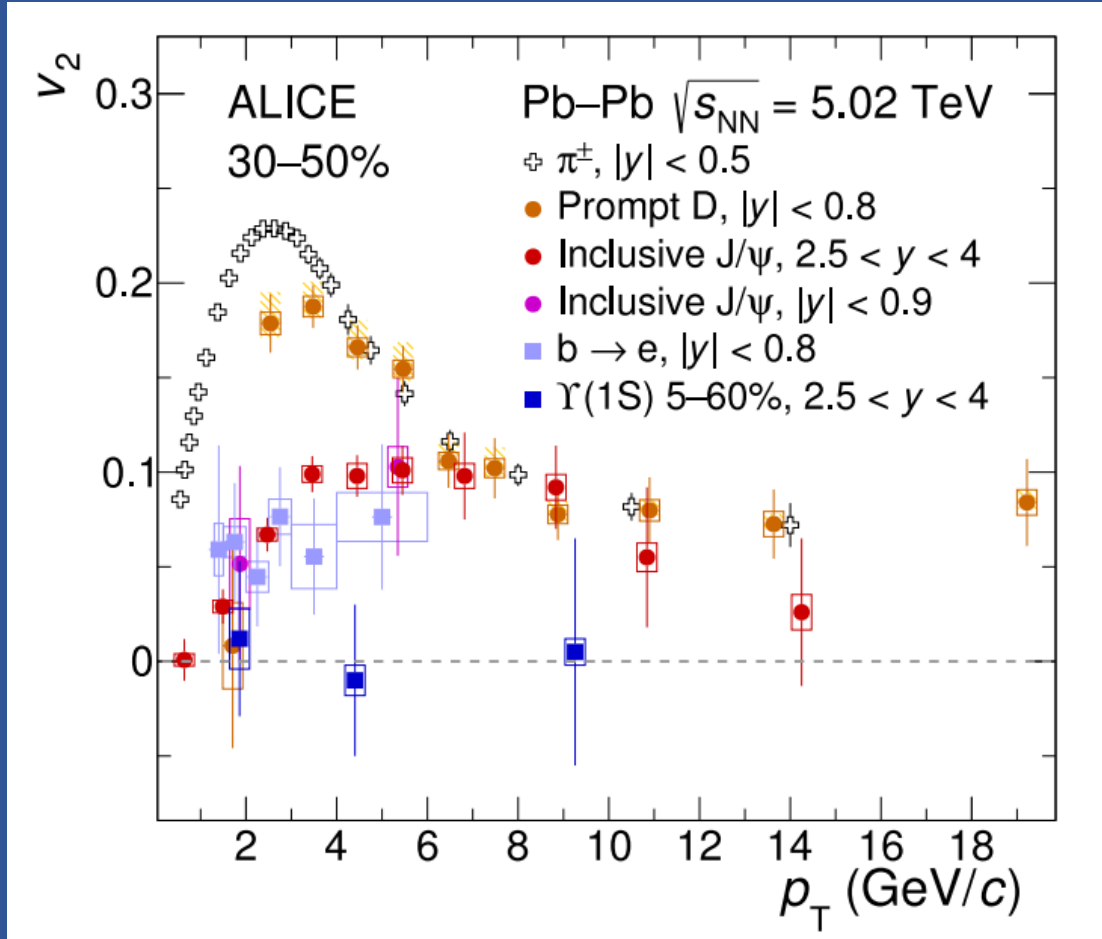
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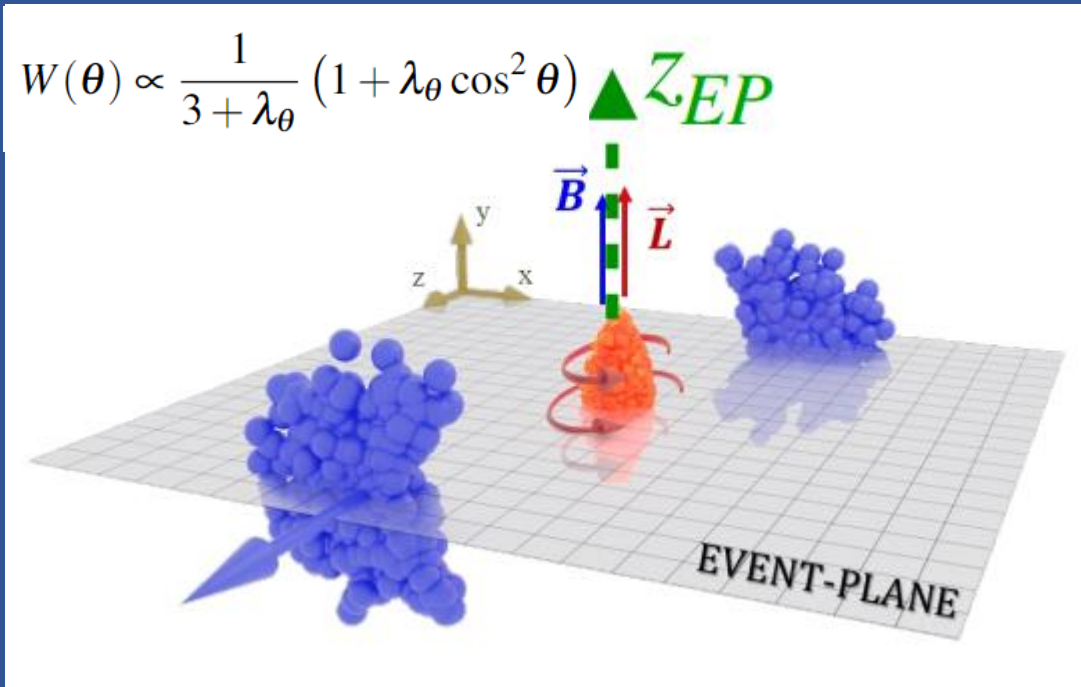
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ALICE, arXiv:2211.04384



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- Inclusive $\Upsilon(1S)$
 - v_2 compatible with **zero** (large uncertainties)
 - Contribution from (re)generation in the beauty sector is small

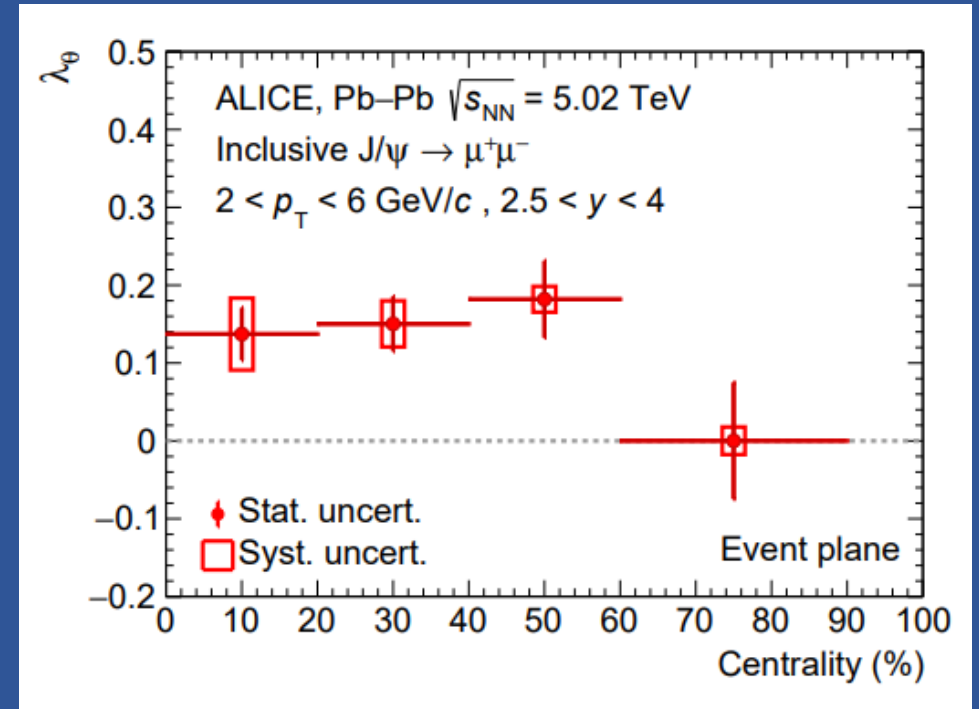
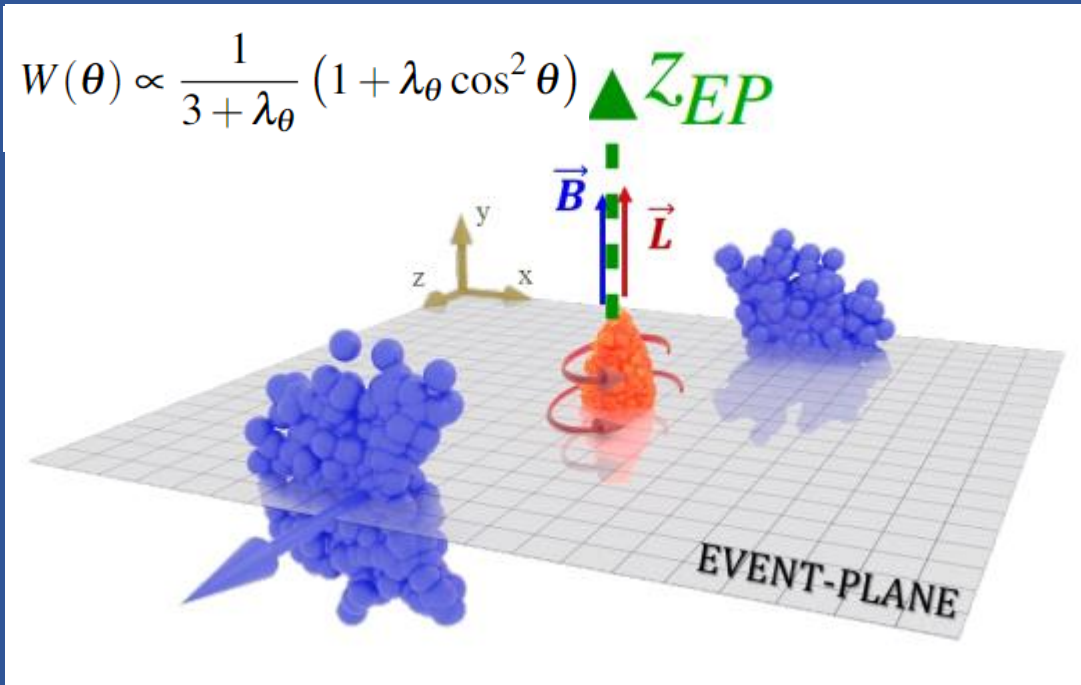
Beyond suppression/regeneration \rightarrow polarization



- Large **magnetic field** ($B \sim 10^{14}$ T, $\tau \sim 1$ fm/c) and **angular momentum** L (up to 10^{22} s $^{-1}$) produced in the QGP formation, perpendicular to the event plane

Beyond suppression/regeneration → polarization

ALICE, Phys.Rev.Lett. 131 (2023) 042303



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- Small but **significant polarization**
- 3σ effect in the 40-60% centrality range
- Most important at low p_T
- Qualitatively consistent with observations for K^* and ϕ

Might be consistent with an effect related to L (B should rather induce a negative polarization)