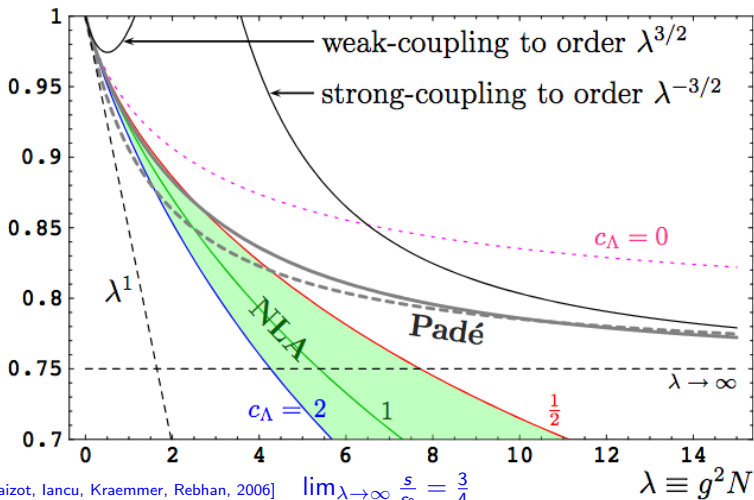


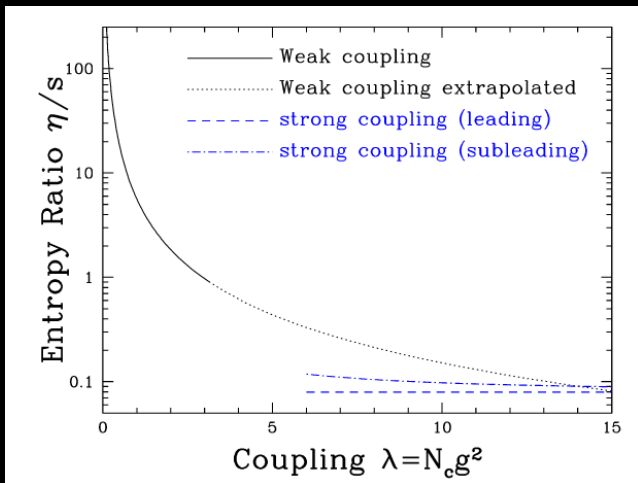
What if ϕ^4 theory in 4d is
non-trivial in the continuum?

Paul Romatschke, CU Boulder

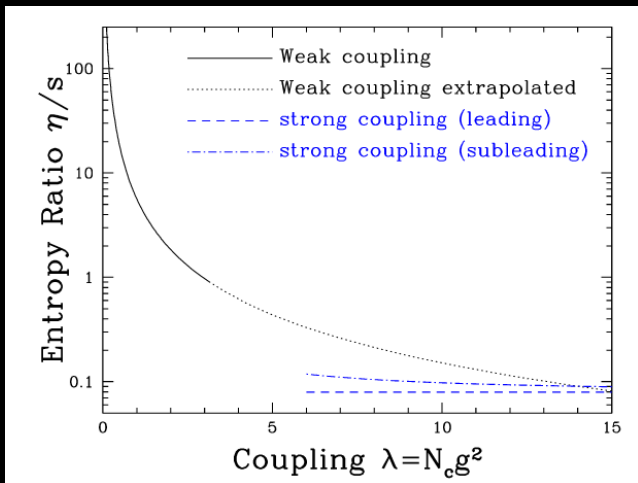
$\mathcal{S}/\mathcal{S}_0$ $\mathcal{N} = 4$ super-Yang-Mills



[Blaizot, Iancu, Kraemmer, Rebhan, 2006]



[Huot, Jeon, Moore, 2006]



[Huot, Jeon, Moore, 2006]

[Policastro, Son, Starinets, 2001] $\lim_{\lambda \rightarrow \infty} \frac{\eta}{s} = \frac{1}{4\pi}$

Can't do QCD

Can't do QCD

Can't really do large N $\mathcal{N} = 4$ SYM

Can't do QCD

Can't really do large N $\mathcal{N} = 4$ SYM

Maybe we can do some other large N theory?

O(N) Model

Defined as

$$Z = \int \mathcal{D}\vec{\phi} e^{-S_E}, \quad S_E = \int_{\mathbf{x}} \left[\frac{1}{2} \partial_{\mu} \vec{\phi} \cdot \partial_{\mu} \vec{\phi} + \frac{\lambda}{N} \left(\vec{\phi}^2 \right)^2 \right],$$

with $\vec{\phi} = (\phi_1, \phi_2, \dots, \phi_N)$.

Examples:

O(N) Model

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Examples:

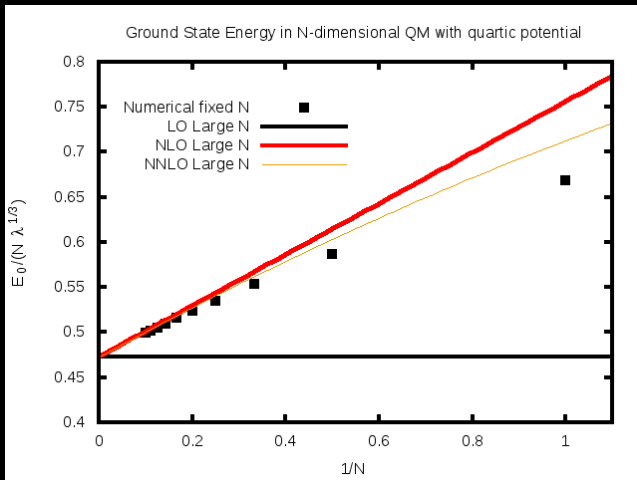
- 0+1d is quantum mechanics in N dimensions (any N)
- 2+1d: conjectured AdS₄ gravity dual for $N \rightarrow \infty$ [hep-th/0210114]
- 3+1d: N=4 is Higgs case

Large N Limit

- $O(N)$ model provides solvable interacting field theory for $N \gg 1$
See review By Moshe and Zinn-Justin [hep-th/0306133]
- Fully non-perturbative, not a weak-coupling expansion!

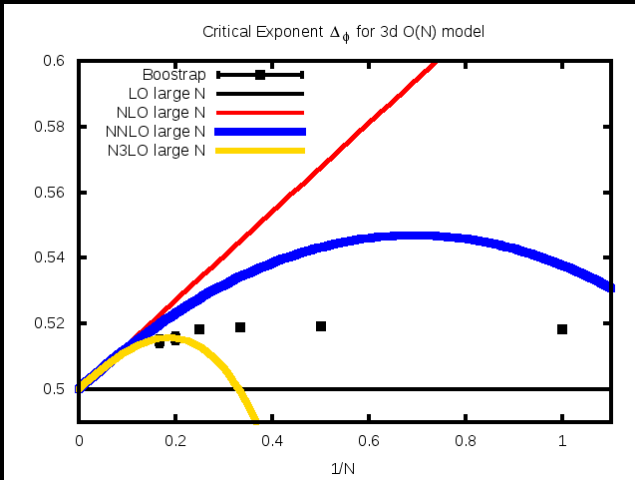
Quantitative tests of large N method in less than 4d

0+1d: Quantum Mechanics



Large N passes QM test even down to $N=1$!

$2+1d$: superrenormalizable QFT

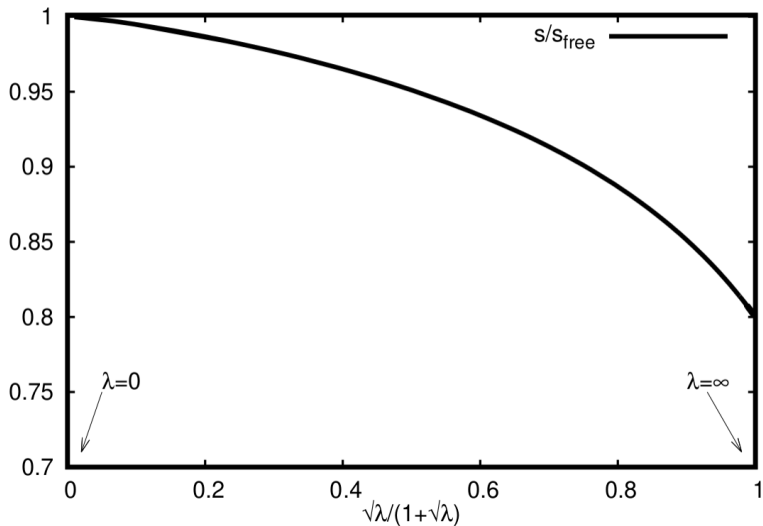


adapted from Kos et al., [1307.6856]

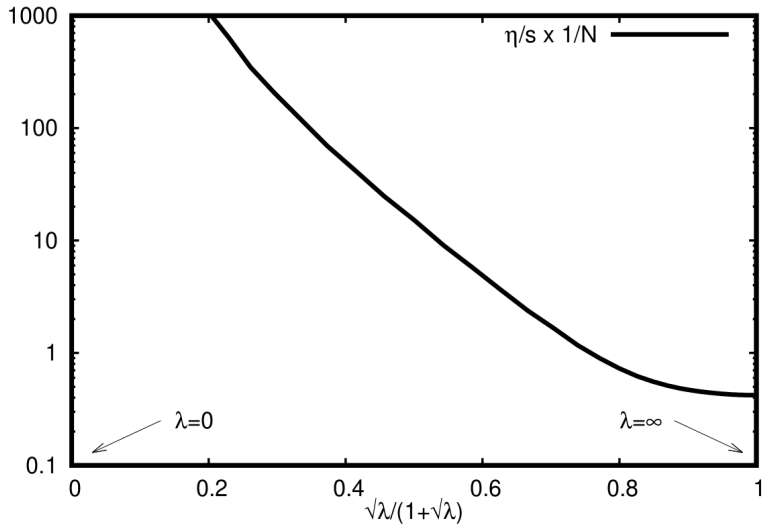
Large N passes test down to $N \sim 4$!

2+1d finite temperature: thermodynamics and transport

Entropy density in O(N) model in 2+1d



Shear viscosity in $O(N)$ model in 2+1d



scalars in less than 4d:

- Large N math is solid
- Large N is quantitatively reliable down to $N \sim \mathcal{O}(1)$
- Large N results for $c_s^2, \frac{\eta}{s}$ for ALL $\lambda \in [0, \infty]$

scalars in $3+1d$: the mental roadblock

Scalar ϕ^4 theory is trivial in four dimensions:

Scalar ϕ^4 theory is trivial in four dimensions:

$$\lim_{\Lambda_{UV} \rightarrow 0} \lambda_{IR} = 0$$

physics arguments/proofs for $d = 4 + \varepsilon$: [Wilson; Fröhlich; Lüscher;
Weisz;...]

math proof for $d = 4$: [Aizenman, Duminil-Copin 2019]

“Loopholes” in proofs:

- assume one or two component scalar fields
- assume $\lambda_{UV} > 0$

Here's what actually happens for 4d $O(N \gg 1)$:

$O(N \gg 1)$ model in 4d in continuum

- $O(N)$ model renormalization is non-perturbative
- In the **continuum limit** $\Lambda_{UV} \rightarrow \infty$, running coupling is

$$\lambda_R(\mu) = \frac{(2\pi)^2}{\ln \frac{\Lambda_{\overline{MS}}^2}{\mu^2}}$$

- Non-vanishing coupling in the continuum. Theory is non-trivial!
- Trick to avoiding triviality:

$$\lim_{\mu \rightarrow \infty} \lambda_R(\mu) = 0^-,$$

$O(N \gg 1)$ model in 4d in continuum

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- Trick to avoiding triviality:

$$\lim_{\mu \rightarrow \infty} \lambda_R(\mu) = 0^-,$$

- Bare coupling is **negative**; situation explicitly excluded in all previous math proofs and physics assumptions!

$O(N \gg 1)$ model in 4d in continuum

Raises new questions:

- How does one make sense of QFT with unbounded potential (e.g. $V(x) = -gx^4$?)

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[A: non-Hermitian or PT-symmetric construction; see extensive work on this by C. Bender!]

$O(N \gg 1)$ model in 4d in continuum

Raises new questions:

- How does one make sense of QFT with unbounded potential (e.g. $V(x) = -gx^4$?)

[A: non-Hermitian or PT-symmetric construction; see extensive work on this by C. Bender!]

- If we set aside problems with 'intuitive/traditional' QFT interpretation, are there actual problems with this theory for any observable?

$O(N \gg 1)$ model in 4d – Results

- Symmetric Mass Generation in $O(N)$ model: radiative corrections generate a VEV for $\vec{\phi}^2$ (Coleman-Weinberg)
- Acts like a mass term for vector field, but NO $O(N)$ symmetry breaking
- At large N , vector (e.g. Higgs for $N=4$) mass

$$m = \sqrt{e} \Lambda_{\overline{\text{MS}}},$$

is prediction of theory (not a parameter)

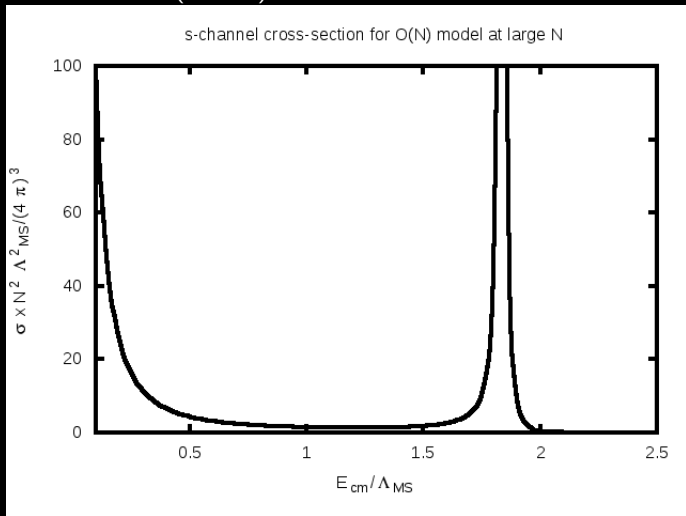
$O(N \gg 1)$ model in 4d – Results

- Finite vacuum free energy density

$$\mathcal{F} = -\frac{N\Lambda_{\overline{\text{MS}}}^4}{64\pi^2}$$

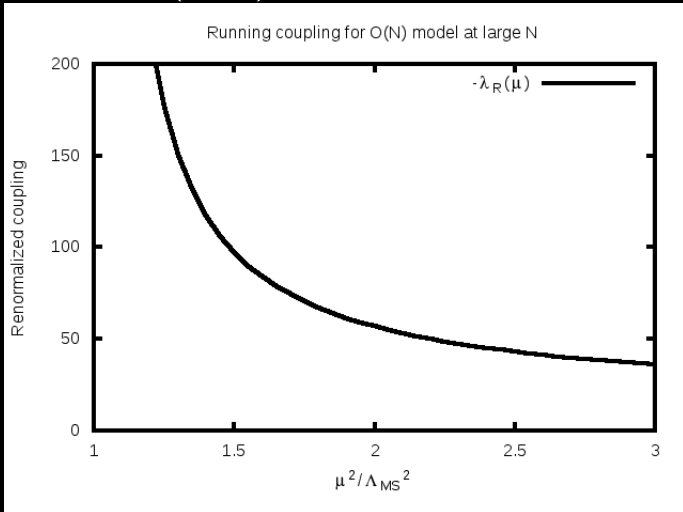
- dominates over perturbative vacuum ($\mathcal{F} = 0$)

$O(N \gg 1)$ model in 4d – Results



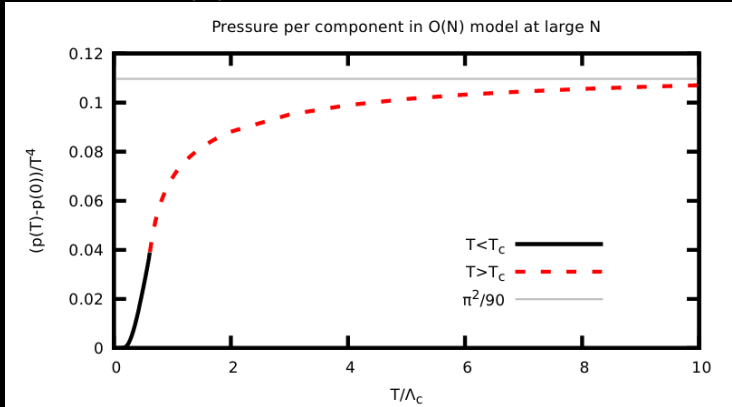
Well behaved scattering cross-section for any CM energy; prediction for scalar bound state at $m \simeq 3\Lambda_{\overline{\text{MS}}}$

$O(N \gg 1)$ model in 4d – Results



Asymptotic freedom in terms of $-\lambda_R(\bar{\mu})$

O(N) model at finite temperature



low T and high T phase separated by 2nd order phase transition

Talk Summary

ϕ^4 theory in 4 dimensions IS non-trivial in the continuum if

- we allow for negative bare coupling $\lambda_0 < 0$
- we stop repeating (false) beliefs about triviality proofs
- we stop equating perturbative QFT with QFT

Stop using QFT 'intuition'!
Calculate observables and check!

Thank You

Bonus Material

References & Hyperlinks

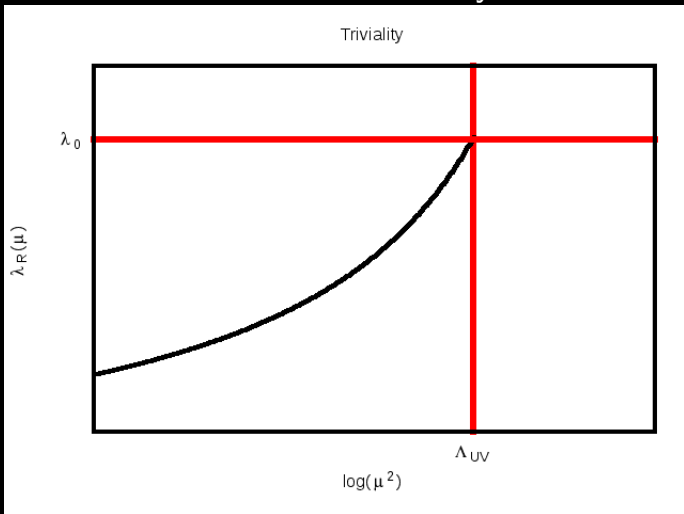
Continuum large N scalar field theory

- PR, “A solvable quantum field theory with asymptotic freedom in 3+1 dimensions”, hyperlink: [\[2211.15683\]](#)
- PR, “Life at the Landau pole”, [\[2212.03254\]](#)
- Grable and Weiner, “A Fully Solvable Model of Fermionic Interaction in 3+1d3+1d”, [\[2302.08603\]](#)
- PR, “What if ϕ^4 theory in 4 dimensions is non-trivial in the continuum?”, [\[2305.05678\]](#)

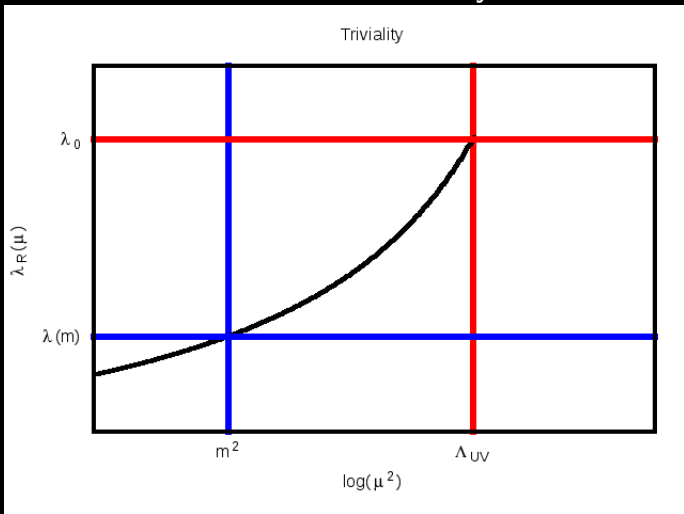
\mathcal{PT} -symmetric Quantum Mechanics and QFT relations

- Bender and Böttcher, “Real Spectra in Non-Hermitian Hamiltonians Having PT Symmetry”, [\[physics/9712001\]](#)
- Ai, Bender and Sarkar, “PT-symmetric -g ϕ^4 theory”, [\[2209.07897\]](#)
- Lawrence, Peterson, PR and Weller, “Instantons, analytic continuation, and PT-symmetric field theory”, [\[2303.01470\]](#)

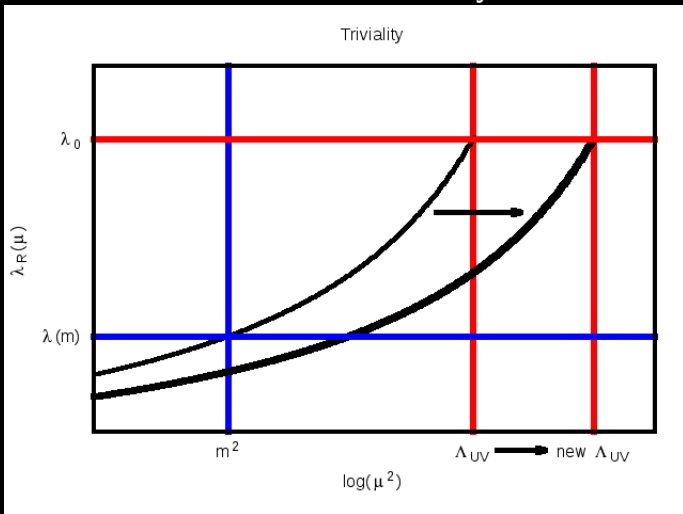
Quantum Triviality



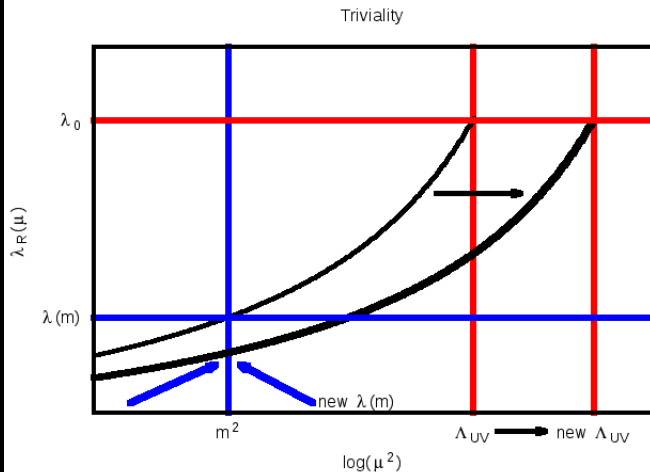
Quantum Triviality



Quantum Triviality



Quantum Triviality



Negative Coupling Field Theory History

A Field Theory with Computable Large-Momenta Behaviour.

K. SYMANZIK

Deutsches Elektronen-Synchrotron DESY - Hamburg

(ricevuto il 12 Dicembre 1972)

In the current extensive discussions (*) of φ^4 theory it is usually taken for granted that the renormalized coupling constant g must be positive. As emphasized previously (³) there is no known reason, axiomatic or otherwise, for $g > 0$ to be required for a physically acceptable theory. The feeling that otherwise the theory cannot have a vacuum and particles of discrete mass is not rigorously founded as discussed near the end of this letter. The interesting feature of the theory with $g < 0$, however, appears worth pointing out: If one assumes the theory to exist, the large-momenta behaviour of its Feynman amplitudes can be computed at generic momenta to arbitrary accuracy. Besides, we find that the imaginary part of the four-point vertex function in φ^4 theory should not change sign in momentum space.

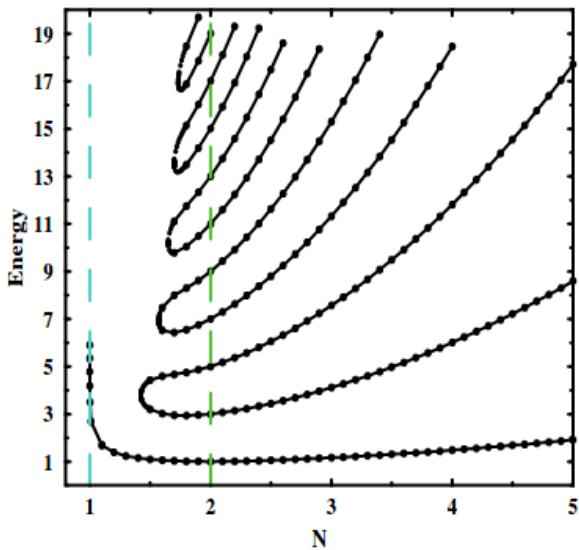
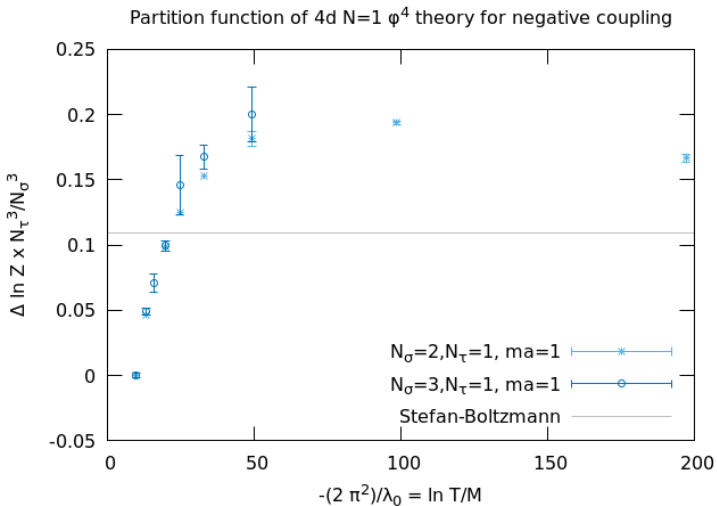


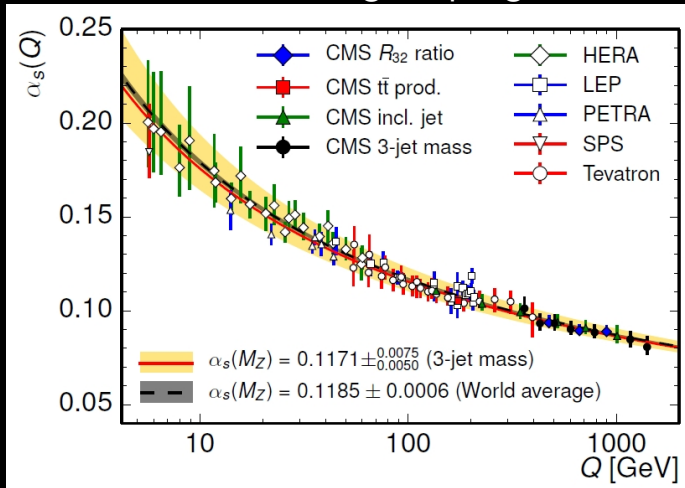
FIG. 1. Energy levels of the Hamiltonian $H = p^2 - (ix)^N$ as a function of the parameter N . There are three regions:

Negative coupling ϕ^4 in 4d on the lattice



adapted from [2305.05678]

QCD running coupling



Somewhat misleading: really a fit of perturbation theory to experimental measurements

QCD at infinite coupling

- In pQCD, $\alpha_s(\bar{\mu})$ does diverge at $\bar{\mu} = \Lambda_{\overline{\text{MS}}} \sim 0.3 \text{ GeV}$
- Usually dismissed as an artifact of perturbation theory
- Non-perturbative extractions (lattice+NRQCD) exist down to $\bar{\mu} = 1.5 \text{ GeV}$ where

$$\alpha_s(1.5\text{GeV}) \simeq 0.336$$

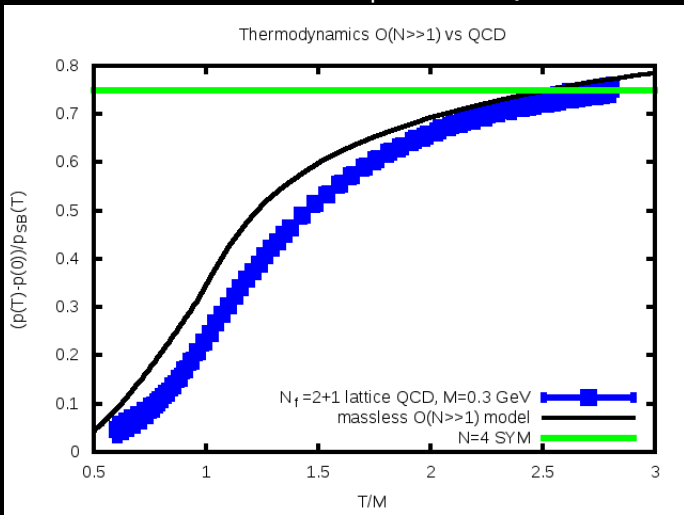
[Bazavov et al, 1407.8437]

- QCD could have a Landau pole at $\Lambda_{\overline{\text{MS}}} \sim 0.3 \text{ GeV}$
- No issues in QCD

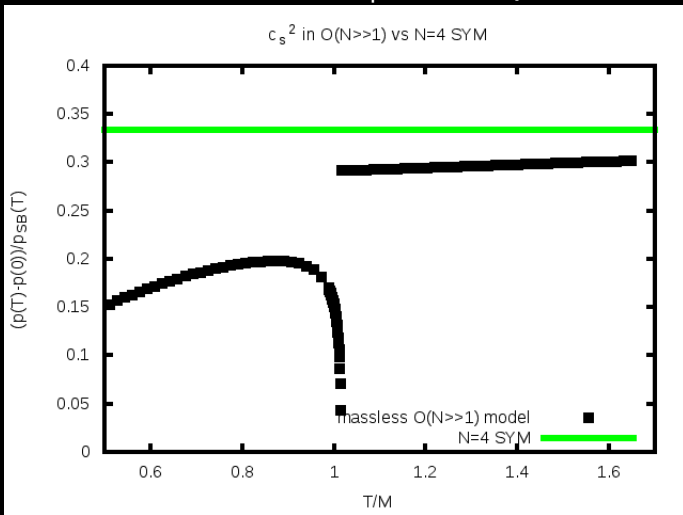
The $O(N \gg 1)$ Model as a Model for QCD

- Only one scale M
- Is M the same as $\Lambda_{\overline{MS}}$ in QCD?
- Let's compare!

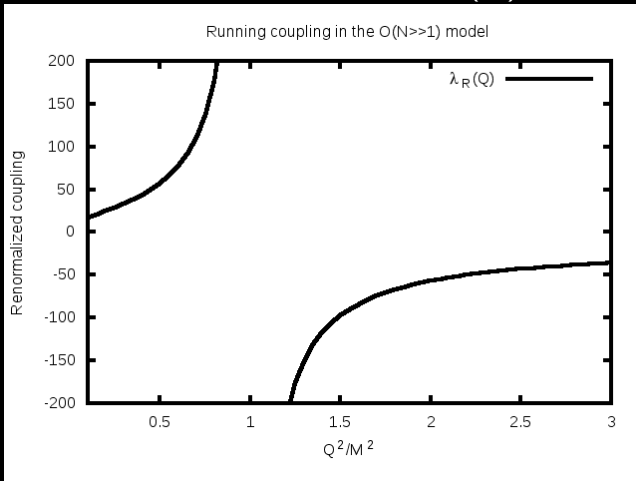
Parameter-free comparison to QCD



Parameter-free comparison to QCD



Exact Running coupling in $O(N)$ Model



Scattering for NR fermions

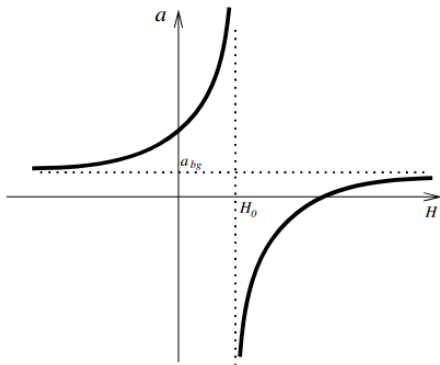


Fig. 2. A schematic of a typical, experimentally observed behavior of an s -wave scattering length $a(H)$ as a function of magnetic field H in a vicinity of a Feshbach resonance.

[Gurarie, Radzihovsky, 2007]

Analytic Continuation for Path Integral

Let's consider 0d field theory defined by

$$Z(\lambda) = \int dx e^{-\lambda x^4}$$

- Well defined for $\text{Re}(\lambda) > 0$,

$$Z(\lambda) = 2\lambda^{\frac{1}{4}} \Gamma\left(\frac{5}{4}\right)$$

- $Z(\lambda)$ is formally divergent for $\text{Re}(\lambda) < 0$
- But it is an analytic function with a well-defined analytic continuation

Analytic Continuation for Path Integral

Let's consider 0d field theory with negative coupling

$$Z_{PT}(g) = \int dx e^{+g x^4}$$

- Particular analytic continuation: “the cone”
- Instead of integrating along real x axis, deform the integration contour into the complex plane so that integral converges

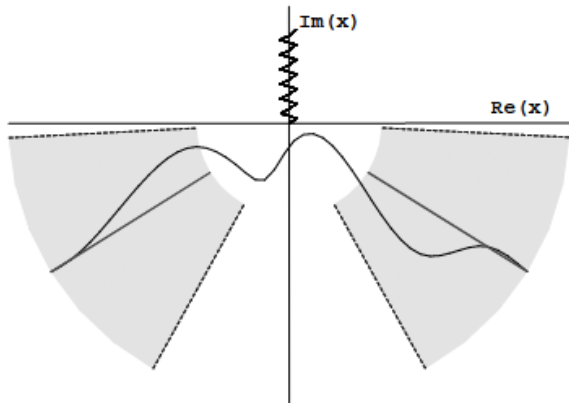


FIG. 2. Wedges in the complex- x plane containing the contour on which the eigenvalue problem for the differential

[Bender & Böttcher, 1997]

Analytic Continuation for Path Integral

$$Z_{\mathcal{PT}}(g) = \int_{\mathcal{C}} dx e^{+gx^4}$$

- Result for $Z_{\mathcal{PT}}(g)$ does not depend on details of \mathcal{C} as long as is inside Bender's wedges
- Result is

$$Z_{\mathcal{PT}}(g) = \sqrt{2} g^{\frac{1}{4}} \Gamma\left(\frac{5}{4}\right) = \text{Re} Z(\lambda \rightarrow -g + i0^+)$$

- Finite and well defined!