

# Dynamics of causal hydrodynamic fluctuations in an expanding system



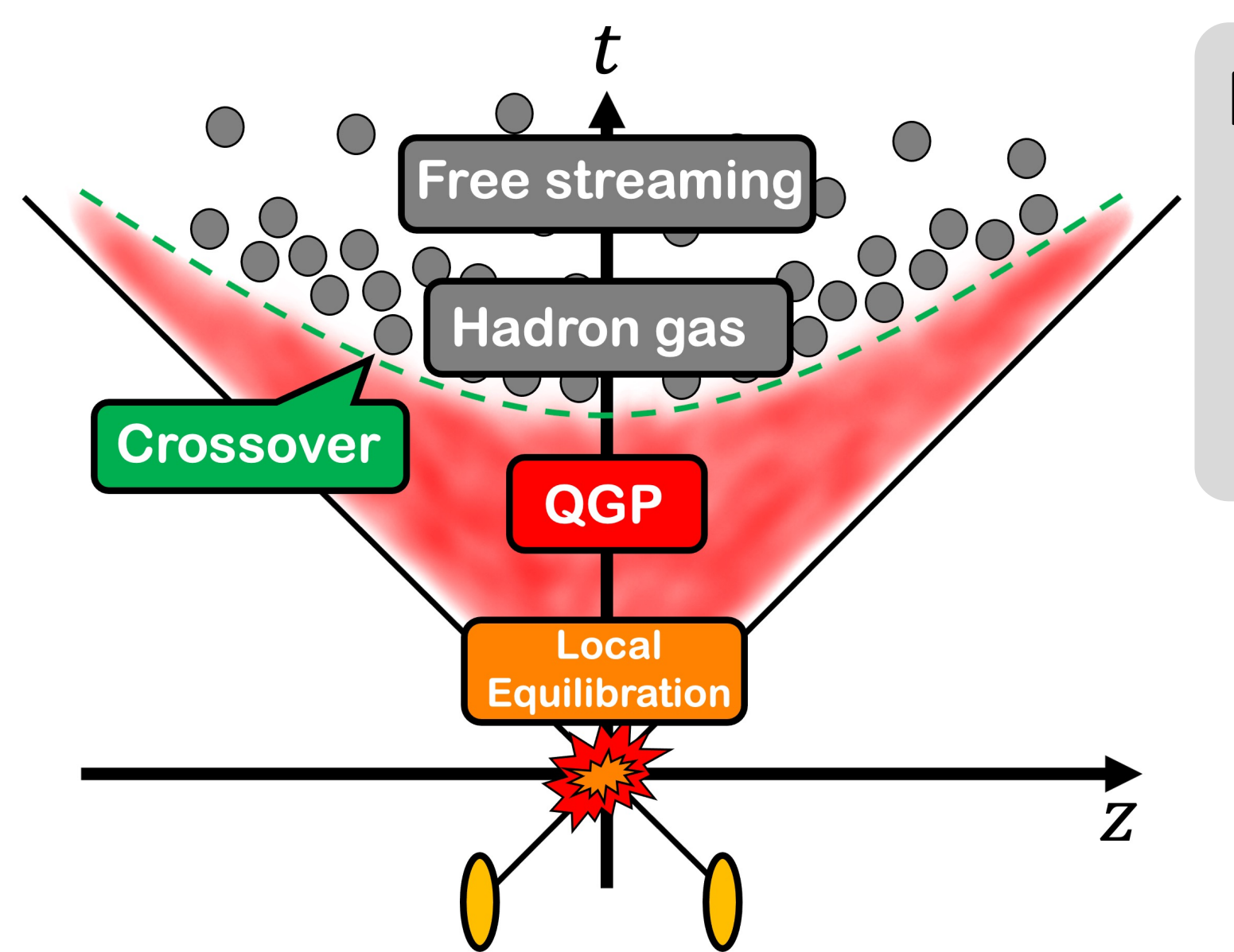
SOPHIA  
HADRON  
PHYSICS  
GROUP

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## 1. Introduction



Space-time evolution of Heavy Ion Collisions

<https://www.bnl.gov/newsroom/news.php?a=110303>

Press release

Discovery of QGP's  
perfect fluid behavior

2005/04/18

Description of  
space-time evolution by  
"Relativistic hydrodynamics"

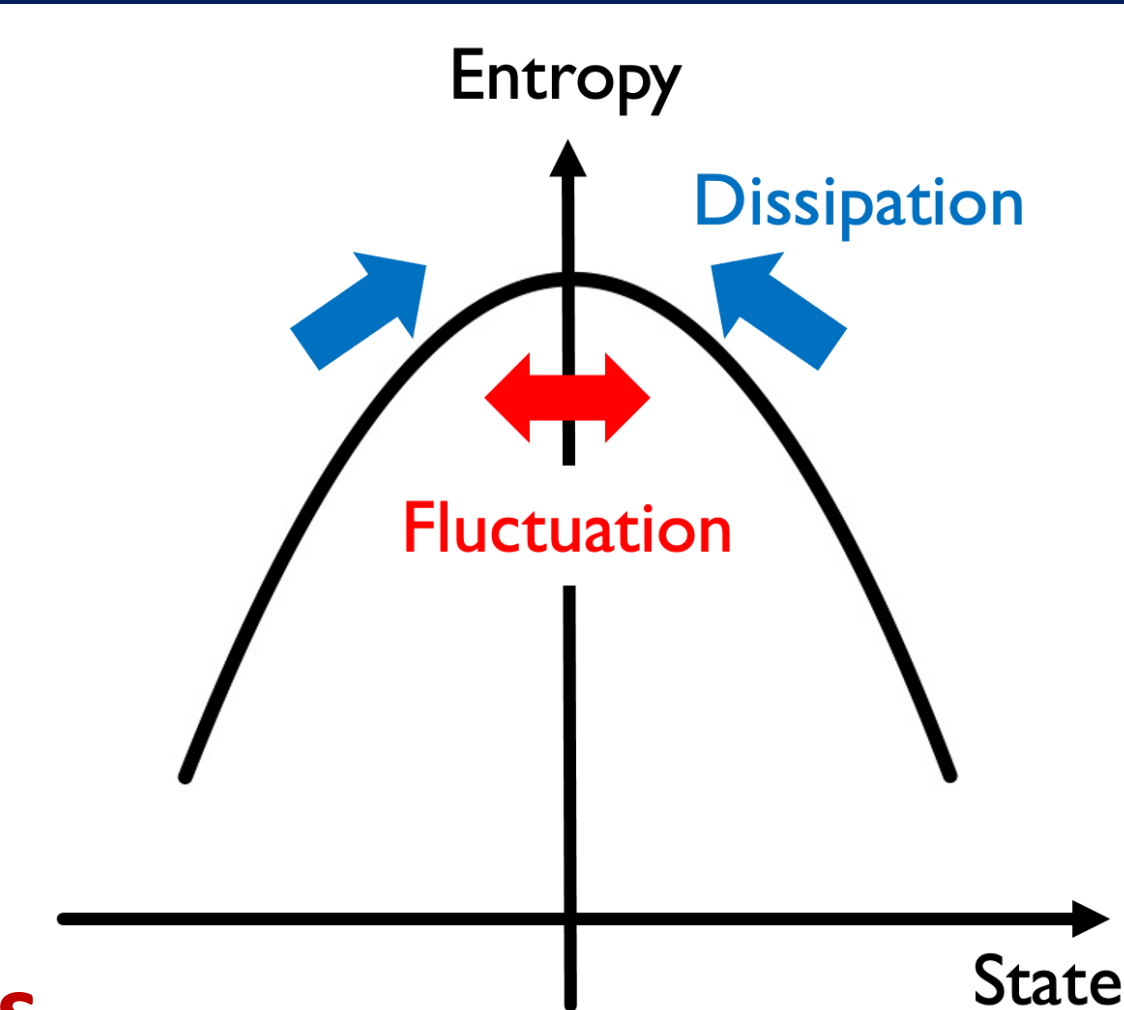
## Fluctuation Dissipation Relation (FDR)

Dissipation (viscosity) and fluctuations  
are always accompanied.

ex.) Bulk pressure

$$\langle \xi_{\Pi}(x) \xi_{\Pi}(x') \rangle = 2T\zeta \delta^{(4)}(x - x')$$

$\xi_{\Pi}$ : fluctuation  $\zeta$ : viscosity



## Evolution of relativistic hydrodynamic models

+ Finite viscosity

+ FDR

Ideal hydro  
around 2000

Dissipative hydro  
around 2010

Fluctuating hydro  
around 2015

Hot Topic

Research on QGP  
properties using  
fluctuations

## 2. Formalisms

J. D. Bjorken, Phys. Rev. D **27**, 140 (1983)

### Perturbative expansion around the Bjorken's solution

$$u_{\text{Bj}}^{\mu} = (\cosh \eta_s, 0, 0, \sinh \eta_s) \quad \eta_s = \frac{1}{2} \ln \left( \frac{t+z}{t-z} \right): \text{coordinate rapidity}$$

### Small deviations

$$u^{\mu} \rightarrow (\cosh(\eta_s + \delta y(\tau, \eta_s)), 0, 0, \sinh(\eta_s + \delta y(\tau, \eta_s)))$$

$$e \rightarrow e_0(\tau) + \delta e(\tau, \eta_s) \text{ etc.} \quad e: \text{energy density, } \tau = \sqrt{t^2 - z^2}: \text{proper time}$$

### Energy-momentum conservation

Balance equation for Background (0th order perturbation)

Balance equation for Fluctuation (1st order perturbation)

for arbitrary constitutive equations for  $\pi$  and  $\Pi$   $\pi \equiv \pi^{00} - \pi^{33}$ : shear stress  
 $\Pi$ : bulk pressure

### Background

$$\frac{d}{d\tau} e_0 + \frac{1}{\tau} (w_0 + \Pi_0 - \pi_0) = 0 \quad (\text{Bjorken equation})$$

$w = e + p$ : enthalpy density  
 $p$ : hydrostatic pressure

### Fluctuation

$$\frac{\partial}{\partial \tau} \left( \frac{\delta e}{\delta y(w_0 + \Pi_0 - \pi_0)} \right) + \frac{1}{\tau} \frac{\partial}{\partial \eta_s} \left( \frac{\delta y(w_0 + \Pi_0 - \pi_0)}{\delta p + \delta \Pi - \delta \pi} \right) + \frac{1}{\tau} \left( \frac{\delta w + \delta \Pi - \delta \pi}{2\delta y(w_0 + \Pi_0 - \pi_0)} \right) = 0$$

## Causal constitutive equations

W. Israel and J. M. Stewart, Annals Phys. **118**, 341 (1979)

$$\text{Israel-Stewart equation + noise} \quad (1 + \tau_{\pi} D) \pi = \frac{4\eta}{3} \theta + \xi_{\pi}$$

$$D \equiv u^{\mu} \partial_{\mu}$$

$$\theta \equiv \partial_{\mu} u^{\mu}$$

$\tau_{\pi}$ : relaxation time

$\xi_{\pi}$ : noise

### Perturbative expansion

### Background

$$\left(1 + \tau_{\pi 0} \frac{d}{d\tau}\right) \pi_0 = \frac{4\eta_0}{3\tau}$$

### Fluctuation

$$\left(1 + \tau_{\pi 0} \frac{\partial}{\partial \tau}\right) \delta \pi = -\frac{\delta \tau_{\pi}}{\tau_{\pi 0}} \left(\frac{4\eta_0}{3\tau} - \pi_0\right) + \frac{4\eta_0}{3\tau} \frac{\partial}{\partial \eta_s} \delta y + \frac{4\delta \eta}{3\tau} + \xi_{\pi}$$

### Fluctuations

FDR for shear stress

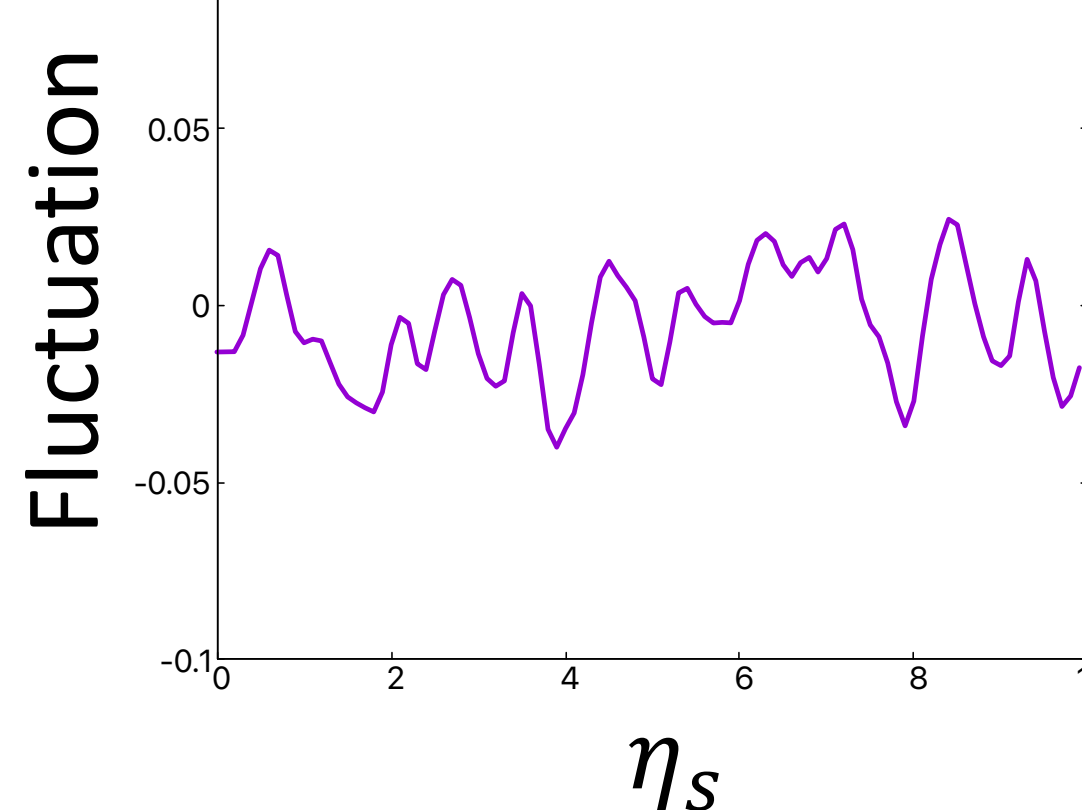
$$\langle \xi_{\pi}(\tau, \eta_s) \xi_{\pi}(\tau', \eta_s') \rangle = \frac{8\eta_0 T_0}{3\tau \Delta x \Delta y} G(\eta_s - \eta_s') \delta(\tau - \tau')$$

$\eta_0$ : shear viscosity (background)

$T_0$ : temperature (background)

$\Delta x = \Delta y = 2 \text{ fm}$ ,  $\langle \xi_{\pi}(\tau, \eta_s) \rangle = 0$ ,  $\sigma_{\eta} = 0.5$

$$G(\eta_s - \eta_s') = \frac{1}{\sqrt{2\pi\sigma_{\eta}^2}} \exp\left(-\frac{(\eta_s - \eta_s')^2}{2\sigma_{\eta}^2}\right)$$



### Standpoint

Representative research in 1-dimensional expanding system

	Hydro	0th (Ideal hydro)	1st (Acausal)	2nd (Causal)
Perturbation				
0th		Bjorken (1983)	Hosoya et al. Danielewicz et al. (1985) etc.	Muronga (2002) etc.
1st			Kouno et al. (1990)	Denicol et al. (2008)
			N/A	Kapusta et al. (2012)
				Our study

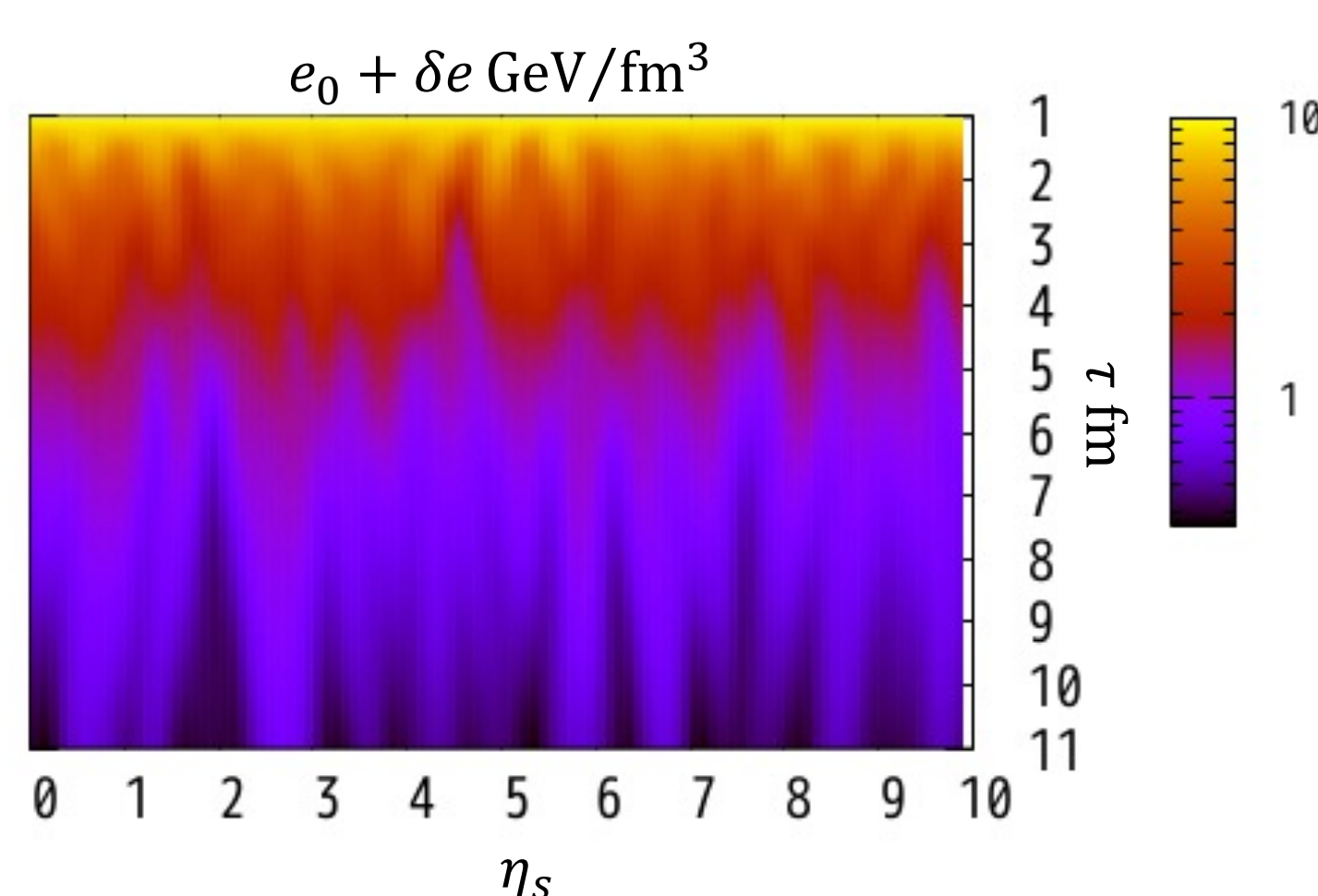
2nd order hydro  
+ 1st order  
perturbation  
+ Fluctuation

## 3. Results

Lattice EoS: A. Bazavov et al., Phys. Rev. D **90**, 094503 (2014)

### Space-time evolution of energy density

Initial condition:  $e_0(\tau = 1 \text{ fm}) = 10 \text{ GeV/fm}^3$ ,  $\delta e(\tau = 1 \text{ fm}) = 0 \text{ GeV/fm}^3$ ,  $\pi_0(\tau = 1 \text{ fm}) = \frac{4\eta}{3\tau} \text{ GeV/fm}^3$

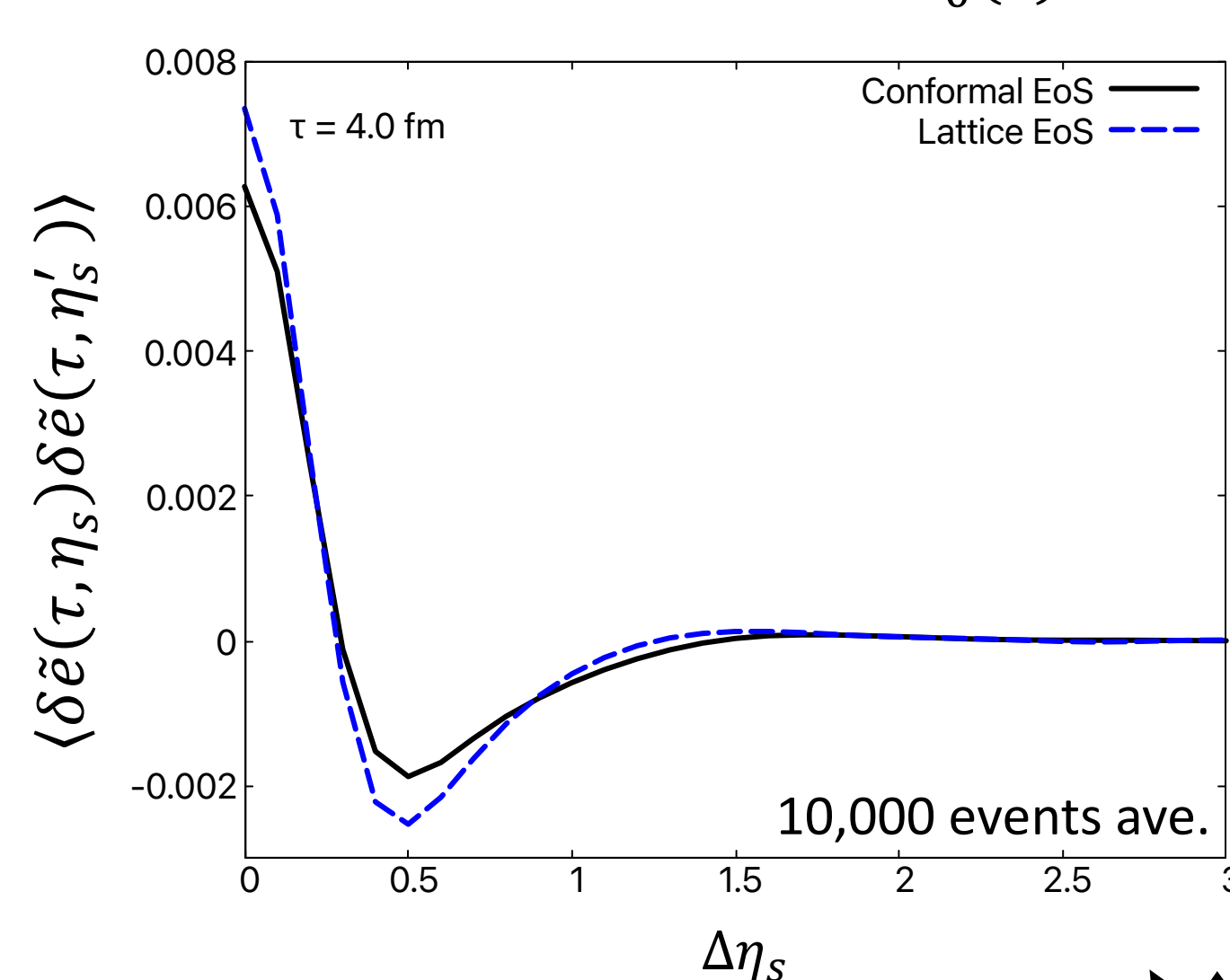
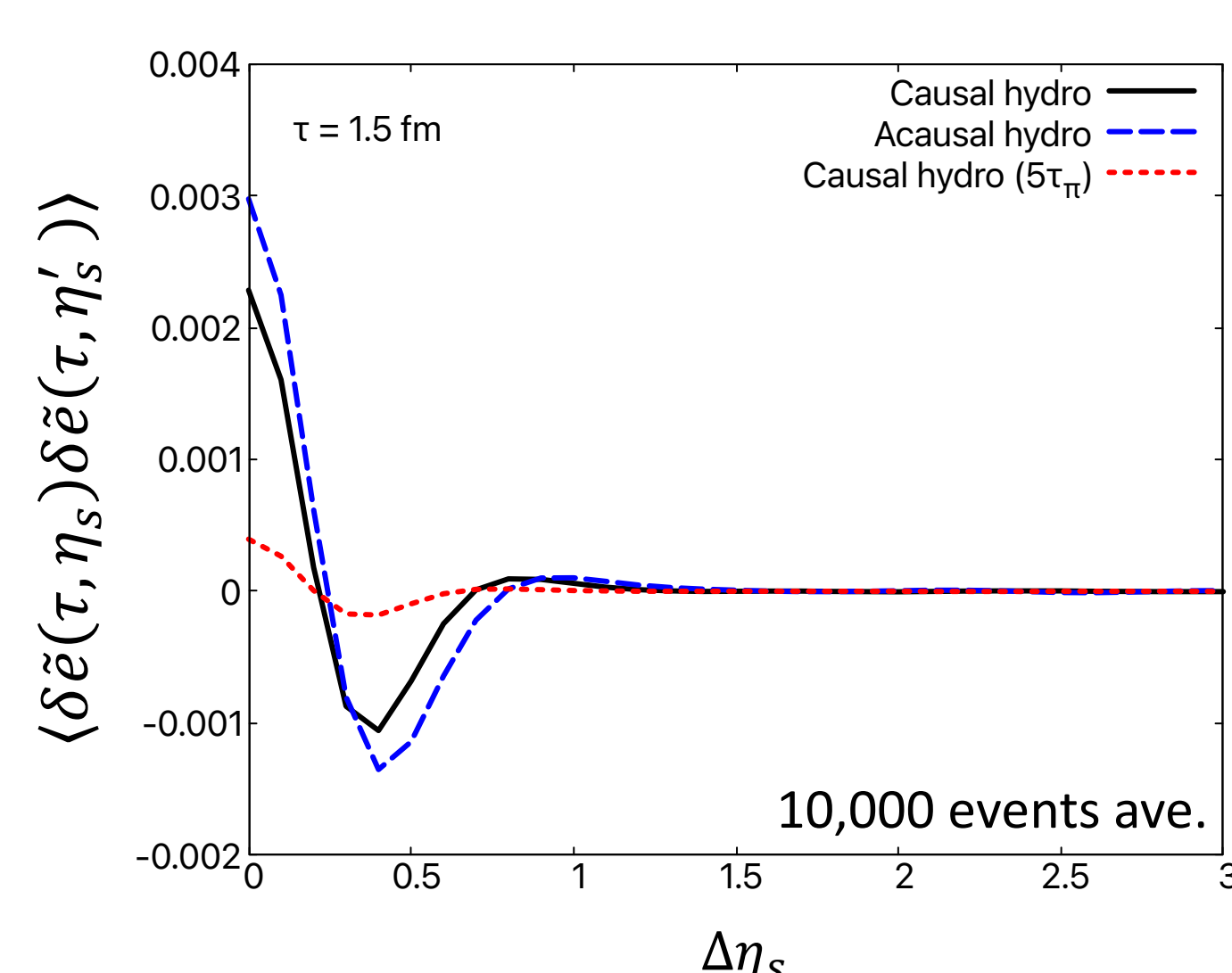


Streak like structure appears

Distribution of energy density  
is frozen and could carry the  
information of the early stage.

### Correlations of energy density fluctuations

$$\delta \tilde{e}(\tau, \eta_s) \equiv \frac{\delta e(\tau, \eta_s)}{e_0(\tau)}$$

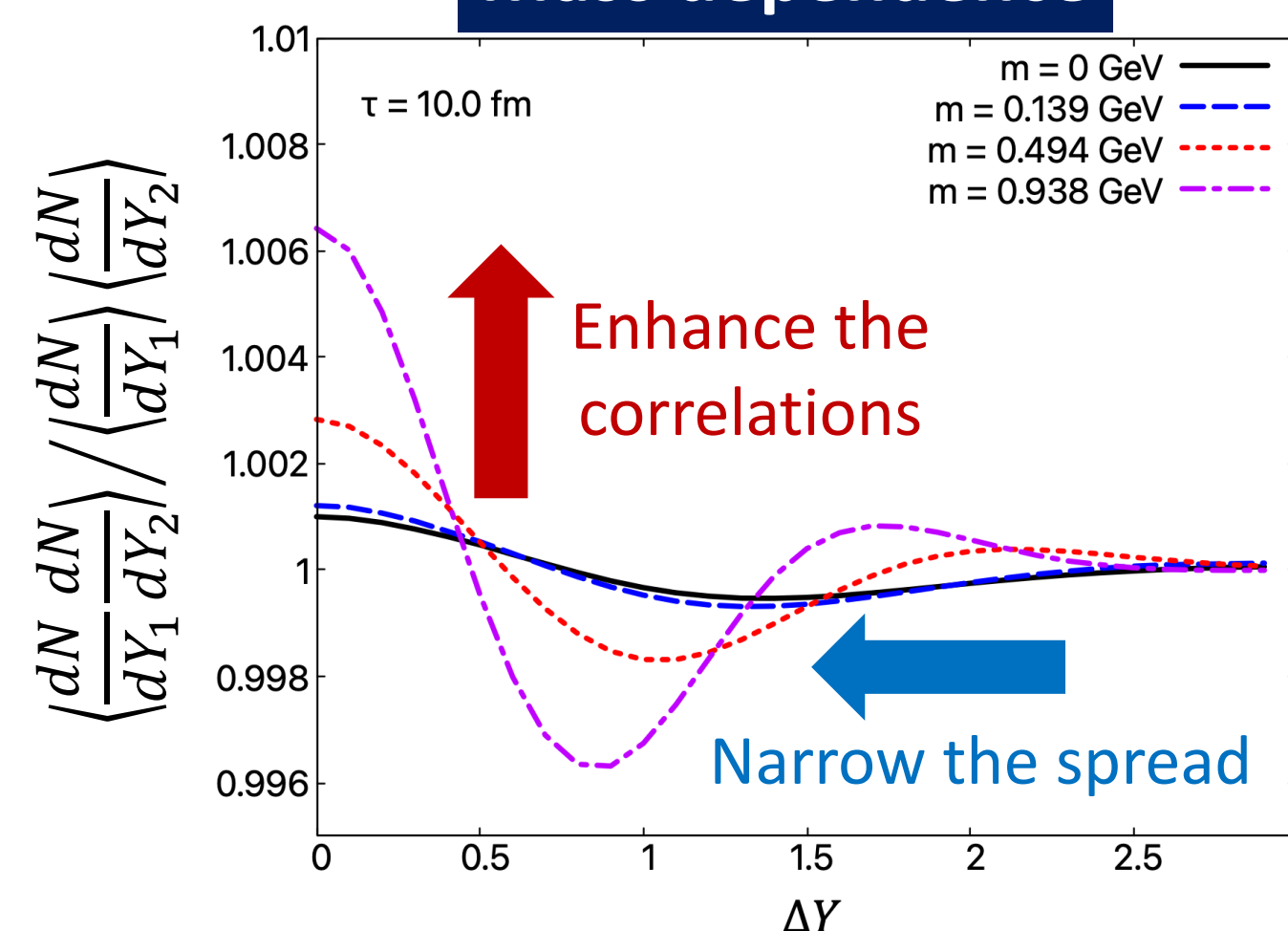


## 2-particle correlations

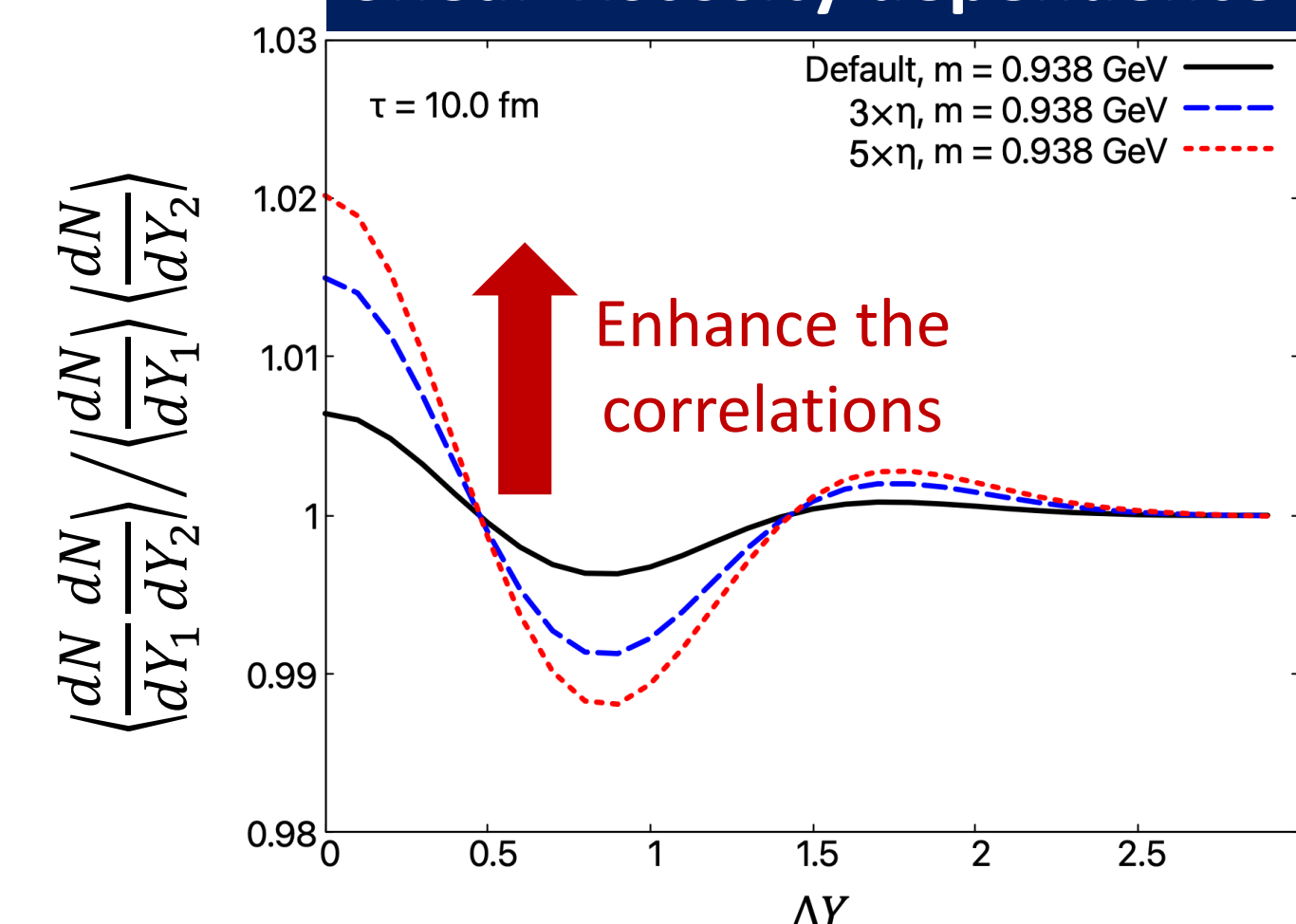
$Y_1$ : rapidity of particle 1

$Y_2$ : rapidity of particle 2

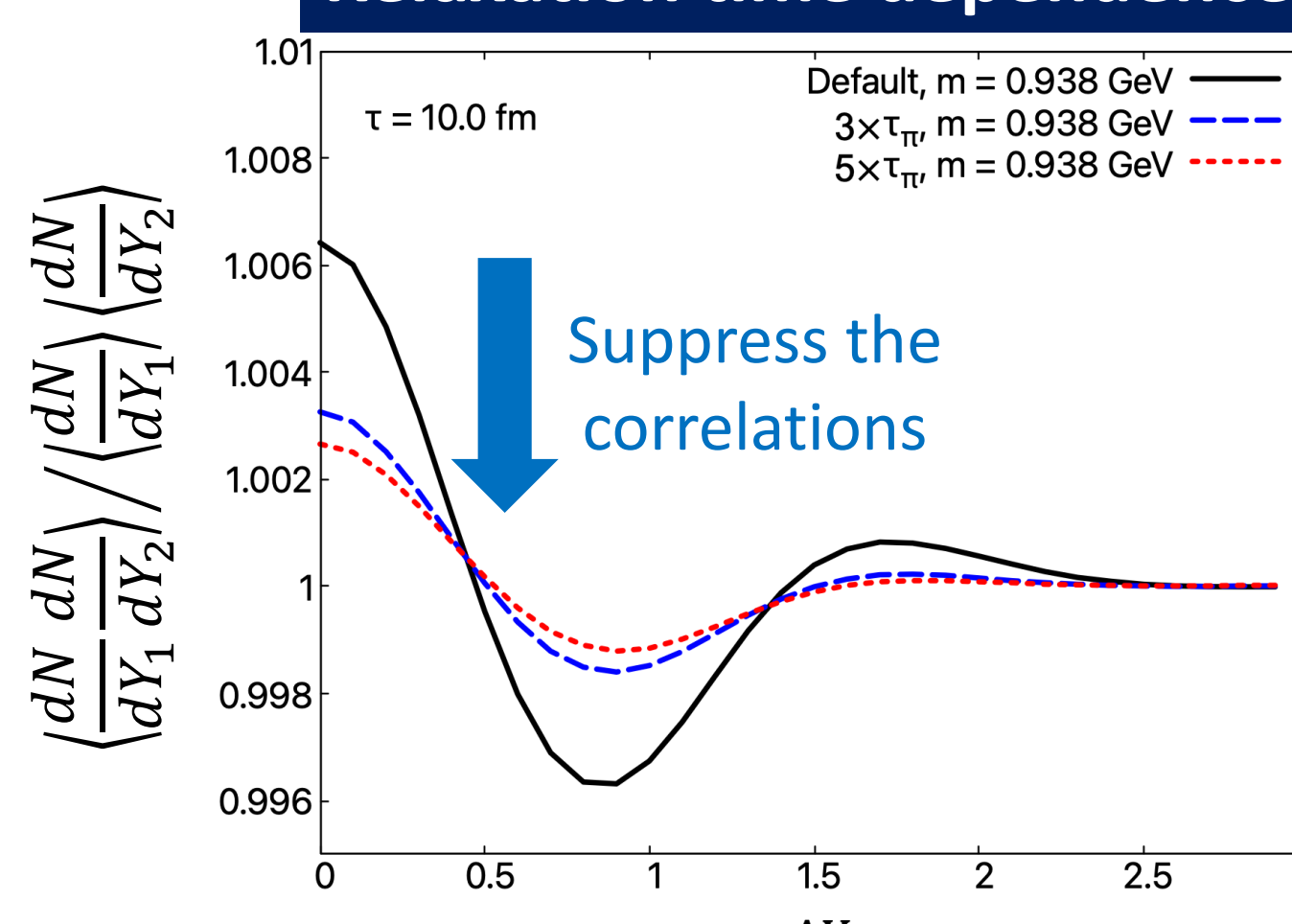
### Mass dependence



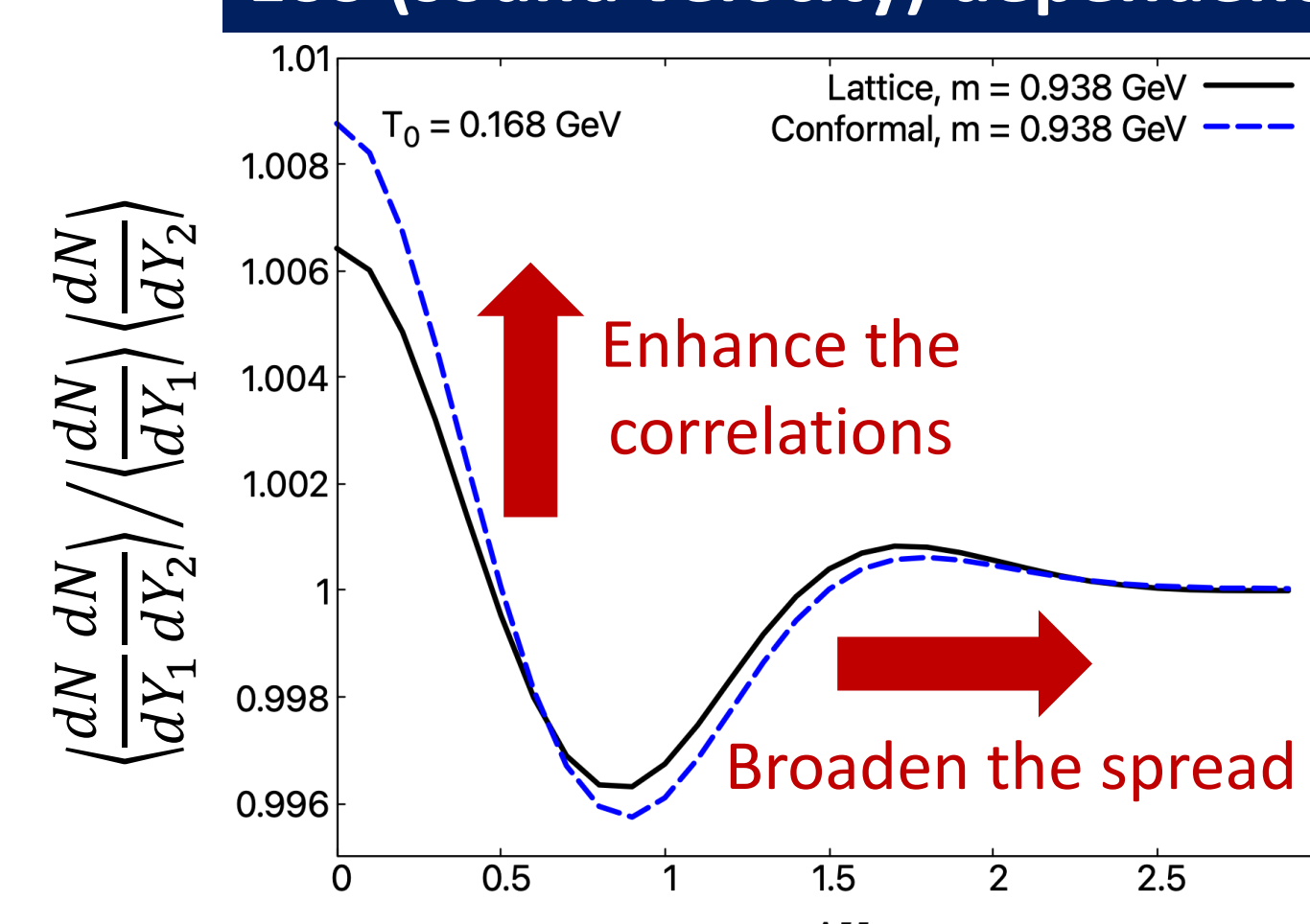
### Shear viscosity dependence



### Relaxation time dependence



### EoS (sound velocity) dependence



- Heavier hadrons are good probes of correlations.
- Shear viscosity and relaxation time work oppositely.
- Correlations include the information of EoS.

Extract the properties of the medium from 2-particle correlations!

## 4. Summary

- We developed a framework which deals with **causal hydrodynamic fluctuations** in 1-dimensional expanding system.
- We observed streak like structure through time evolution of energy density caused by a **freeze of distribution**.
- We found behaviors of **correlations of thermodynamic variables and particles** are closely related with the properties of the medium.