

# Molecular dynamics analysis of particle number fluctuations from a first-order phase transition

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## Motivation

Proton number fluctuations in heavy-ion collisions are a prime observable in the search for the QCD critical point and the associated first-order phase transition. This is motivated by the equilibrium expectations of large critical fluctuations of baryon density near the critical point. However, it remains a non-trivial task to understand the development of critical point signatures. To shed light on this question, we explore the behavior of critical point fluctuations in a framework of molecular dynamics with Lennard-Jones (LJ) potential.

## Molecular dynamics of the Lennard-Jones fluid

We solve the classical  $N$ -body problem by performing molecular dynamics simulations of the LJ fluid in a box with periodic boundary conditions. The LJ interaction potential reads

$$V_{\text{LJ}}(\tilde{r})/\varepsilon = 4 \left[ \left( \frac{\sigma}{\tilde{r}} \right)^{12} - \left( \frac{\sigma}{\tilde{r}} \right)^6 \right] = 4(\tilde{r}^{-12} - \tilde{r}^{-6}), \quad (1)$$

where the first and second terms correspond to short-range repulsion and intermediate-range attraction, respectively. We use reduced variables ( $\mathbf{r}^* = \mathbf{r}/\sigma$ ,  $T^* = T/(\varepsilon)$ ,  $n^* = n\sigma^3$ ,  $p^* = p\sigma^3/\varepsilon$ ,  $\tilde{t} = t\sqrt{\varepsilon/(m\sigma^2)}$ ) as well as normalized relative to the critical point location ( $\tilde{\mathbf{r}} = \mathbf{r}^*/r_c^*$ ,  $\tilde{T} = T^*/T_c^*$ ,  $\tilde{n} = n^*/n_c^*$ ,  $\tilde{p} = p^*/p_c^*$ ).

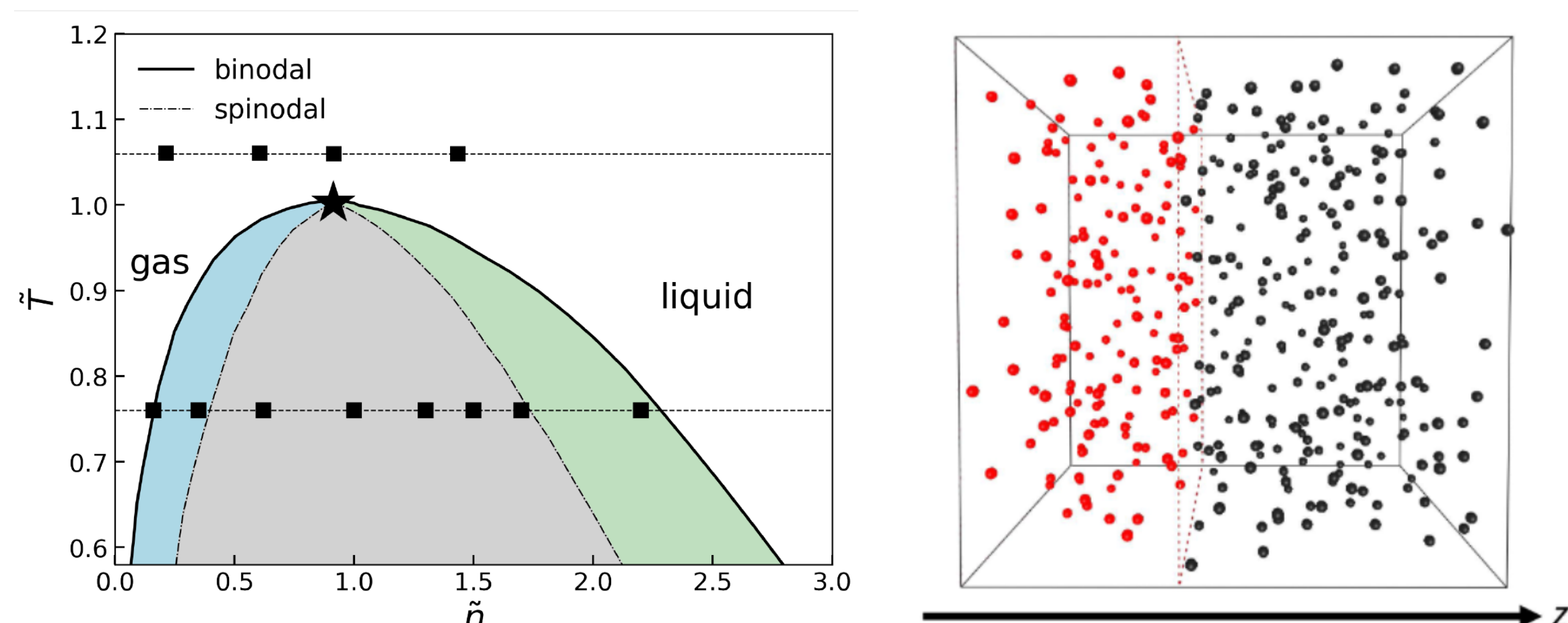


Figure 1: (Left):  $(T, n)$  phase diagram in liquid-gas region. (Right): subensemble in simulation.

*Ergodic hypothesis:* Observables can be calculated as time averages

$$\langle A \rangle = \frac{1}{\tilde{\tau}} \int_{\tilde{t}_{\text{eq}}}^{\tilde{t}_{\text{eq}} + \tilde{\tau}} A(\{\tilde{\mathbf{r}}_i(\tilde{t}), \tilde{\mathbf{v}}_i(\tilde{t})\}) d\tilde{t}, \quad (2)$$

Pressure (virial theorem):

$$\tilde{p} = \tilde{n}\tilde{T} + \frac{\sum_{i=1}^N \sum_{j=i+1}^N \tilde{\mathbf{r}}_{i,j} \cdot \tilde{\mathbf{f}}_{i,j}}{3\tilde{L}^3}. \quad (3)$$

Scaled variance of particle number fluctuations in a subvolume:

$$\tilde{\omega} = (1 - \alpha)^{-1} \frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle}. \quad (4)$$

Here  $\alpha$  is observed volume fraction and the factor  $(1 - \alpha)^{-1}$  is a correction for global conservation [3].

Grand-canonical limit:

$$\omega_{\text{gce}} = \tilde{T} \left( \frac{\partial \tilde{p}}{\partial \tilde{n}} \right)^{-1}_{\tilde{T}}, \quad \omega_{\text{gce}} \rightarrow \infty \text{ at the critical point.} \quad (5)$$

## Results

### Pure phase (crossover region)

We study fluctuations along a supercritical isotherm  $T = 1.06T_c$ , where the system is in a pure phase.

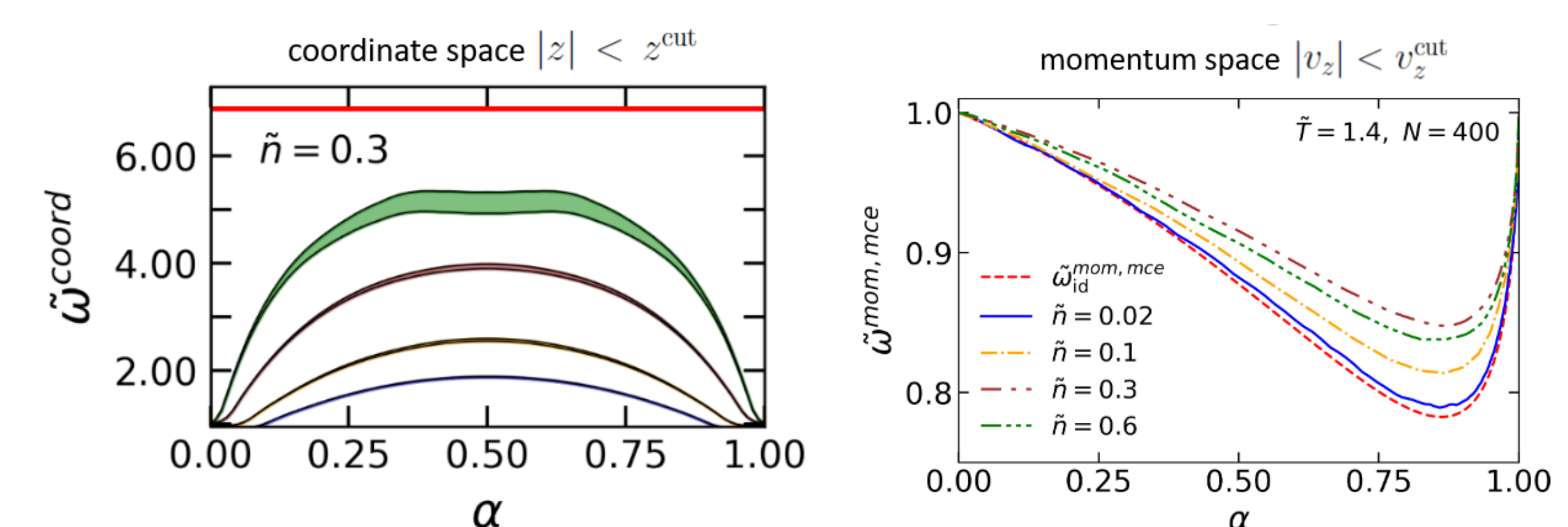


Figure 2: Corrected scaled variance  $\tilde{\omega}$  of particle number fluctuations near the critical point from a crossover side ( $T = 1.06T_c$ ,  $n \approx n_c$ ) in (left) coordinate and (right) momentum space subvolumes.

**Conclusion:** We observe large fluctuations associated with the critical point in coordinate space subvolumes, but in the absence of collective flow, these signals are washed out when momentum cuts are imposed instead.

### Mixed phase

Here we study fluctuations along a subcritical isotherm  $T = 0.76T_c$ .

## Results

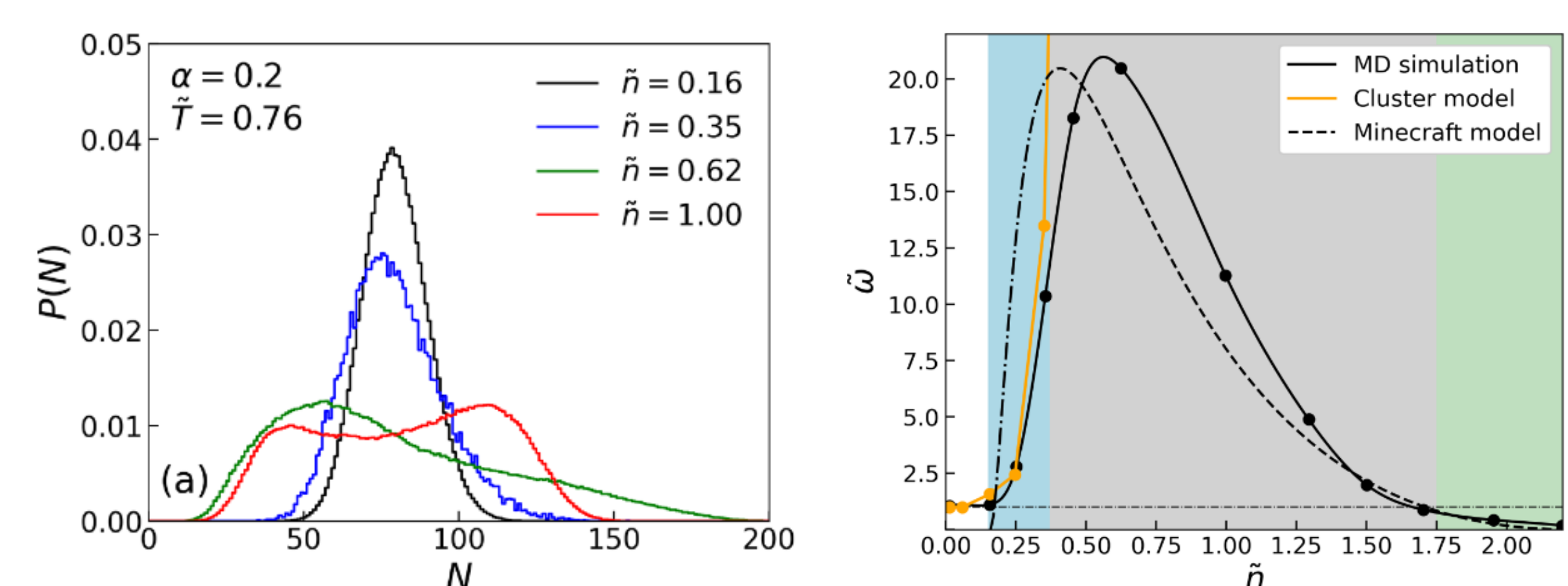


Figure 3: (Left):  $P(N)$  distribution at mixed phase. (Right): Fluctuations at mixed phase. Here  $\alpha = 0.2$ ,  $N = 400$ .

**Conclusion:** We observe large particle number fluctuations in the nucleation region  $\tilde{n} \leq 0.35$ , where *cluster model* can be used,

$$Z_{\text{GCE}} = \prod_{k \geq 1} \exp \left[ V(2\pi k m T)^{3/2} g(k) \exp \left\{ \frac{\mu k}{T} \right\} \right], \quad (6)$$

and in the spinodal decomposition region at  $\tilde{n} > 0.35$  where we use the so-called Minecraft model (see [2] for explicit formulas).

## Outlook

- Fluctuations in expanding systems as appropriate for heavy-ion collisions.
- Precise study of fluctuations in the mixed phase region, dependence on system size and subvolume form, and test of the thermodynamic limit.

## References

- [1] Volodymyr A. Kuznietsov, Oleh Savchuk, Mark I. Gorenstein, Volker Koch **and** Volodymyr Vovchenko. “Critical point particle number fluctuations from molecular dynamics”. *in Phys. Rev. C*: 105 (2022).
- [2] Volodymyr A. Kuznietsov, Oleh Savchuk, Roman V. Poberezhnyuk, Volodymyr Vovchenko, Mark I. Gorenstein **and** Horst Stoecker. “Molecular dynamics analysis of particle number fluctuations in the mixed phase of a first-order phase transition”. *in Phys. Rev. C*: 107 (2023).
- [3] Volodymyr Vovchenko, Oleh Savchuk, Roman V. Poberezhnyuk, Mark I. Gorenstein **and** Volker Koch. “Connecting fluctuation measurements in heavy-ion collisions with the grand-canonical susceptibilities”. *in Physics Letters B*: 811 (2020).