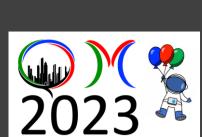
# Understanding effect of event-by-event fluctuations on Light-Nuclei Yield Ratio



Koichi Murase (YITP, Kyoto Univ), Shanjin Wu (Lanzhou Univ)





**Abstract:** To understand how the *event-by-event fluctuations* of the final-state nucleon distributions f(r, p) affect the yield ratio of light nuclei  $N_p N_t / N_d^2$  qualitatively, we model the ensemble of f(r, p) by the overlap of random n-Gaussian hot spots and obtain analytical formulae based on the coalescence model, which is used to establish qualitative understanding.

H. Liu *et al*, PLB **805**, 135452 (2020)

[Recap. of NA49, STAR]

Collision Energy  $\sqrt{s_{NN}}$  (GeV)

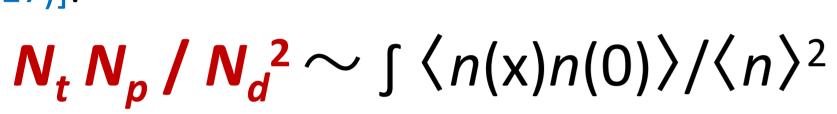
# 1. Background and Motivation

## **■** Critical-point search in heavy-ion collisions

The critical point (CP) of quantum chromodynamics (QCD) is predicted in the finite- $\mu_B$  (baryon chemical potential) region of the phase diagram based on effective theories. Heavy-ion collision experiments at lower collision energies, such as RHIC BES programs, FAIR CBM, etc. aim at finding the CP signal, yet its clear experimental evidence has not been established.

## ■ Signals of critical point

The collision energy  $Vs_{NN}$  roughly corresponds to  $\mu_B$  through baryon stopping, which suggests us to see non-monotonic  $Vs_{NN}$  dependence of observables. Among many observables suggested for the CP signals, we focus on *the yield ratio of light nuclei* [K.-J. Sun *et al* (2017)]:



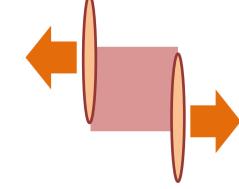
 $N_t$ : triton yield,  $N_p$ : proton yield,  $N_d$ : deuteron yield  $\leftarrow$  Calculated by the *coalescence model* from the *nucleon distribution*  $f(\mathbf{r}, \mathbf{p})$ .

- ✓ Sensitive to local spatial correlations, i.e., critical correlations
- Cancels unwanted volume effect, etc. in the ratio
- ✓ Non-monotonic behaviors found in experimental data

## Dynamical effects in realistic collisions

Previous analyses are based on *ideal setups* with Gaussian source or boost

invariance. In realistic collisions, there are various dynamical effects:



- Expansion → coordinate-momentum correlation
- Initial fluctuations → non-critical correlations
- Non-flow effects  $\rightarrow f(x, p)$  not in equilibrium form
- Finite size and time (Kibble-Zurek) → correlation development ltd.
   cf S. Wu, K. Murase, S. Tang and H. Song, Phys. Rev. C 106, 034905 (2022)

Question: How does the dynamical effects, especially the event by event fluctuations, affect  $N_t N_p / N_d^2$ ?

Event-by-event

$$f^{(1)}(r, p), f^{(2)}(r, p), f^{(3)}(r, p), ...$$

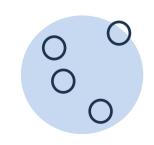
Event average  $N_A^{(1)}, N_A^{(2)}, N_A^{(2)}, N_A^{(3)}, ... \rightarrow \langle N_A^{(i)} \rangle$ 

# 2. Model: fluctuating *n*-Gaussian hot spots

## **■** Coalescence model

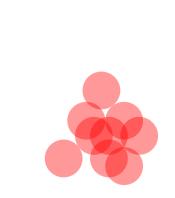
Nucleus of mass number A is formed by nucleons that are close in phase space.

$$N_A = g_A \int \left[\prod_i^A d^3m{r}_i d^3m{p}_i f_i(m{r}_i,m{p}_i)
ight] W_A(\{m{r}_i,m{p}_i\}_{i=1}^A)$$
 sudden frzout



 $W_A()$  ... Wigner fn. for nucleus wave fn. (assumed to be harmonic oscillator)  $f_i()$  ... phase space distribution of ith nucleon  $g_A$  ... statistical factor from spin

■ Single-event f(r, p) by n-Gaussian hot spots



$$f(\boldsymbol{z}) = \frac{1}{n} \sum_{h=1}^{n} c_h G(\boldsymbol{z} - \boldsymbol{a}_h; \mathcal{C}_2^{\text{hs}})$$

G: Gaussian profile

n Number of hot spots

 $z \sim \sqrt{2}(r/\sigma_A, p\sigma_A)$ 

Center of hot spot h

 $\mathcal{C}_2^{\mathrm{hs}}$  Covariance; hot-spot size/shape in  $\mathbb{R}^6$ 

 $\sigma_A$ : nucleus size

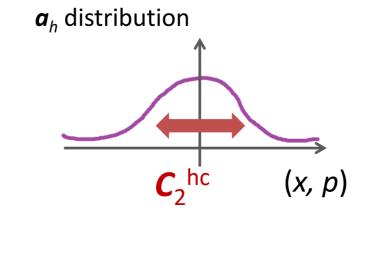
 $C_h$  Magnitude of hot spot h

## ■ Ensemble of $\{f(r, p)\}$ by randomizing $(a_h, c_h)$

Gaussian distribution of hot-spot centers  $\boldsymbol{a}_h$ 

$$\Pr(\{c_h, \boldsymbol{a}_h\}_{h=1}^n) = \prod_{h=1}^n p(c_h) \frac{e^{-\frac{1}{2}\boldsymbol{a}_h^{\mathrm{T}}(\mathcal{C}_2^{\mathrm{hc}})^{-1}\boldsymbol{a}_h}}{\sqrt{\det(2\pi\mathcal{C}_2^{\mathrm{hc}})}}$$

 $h{=}1$   $V^{\det(Z)}$   $\mathcal{C}_2^{\mathrm{hc}}$  Covariance of hot-spot centers

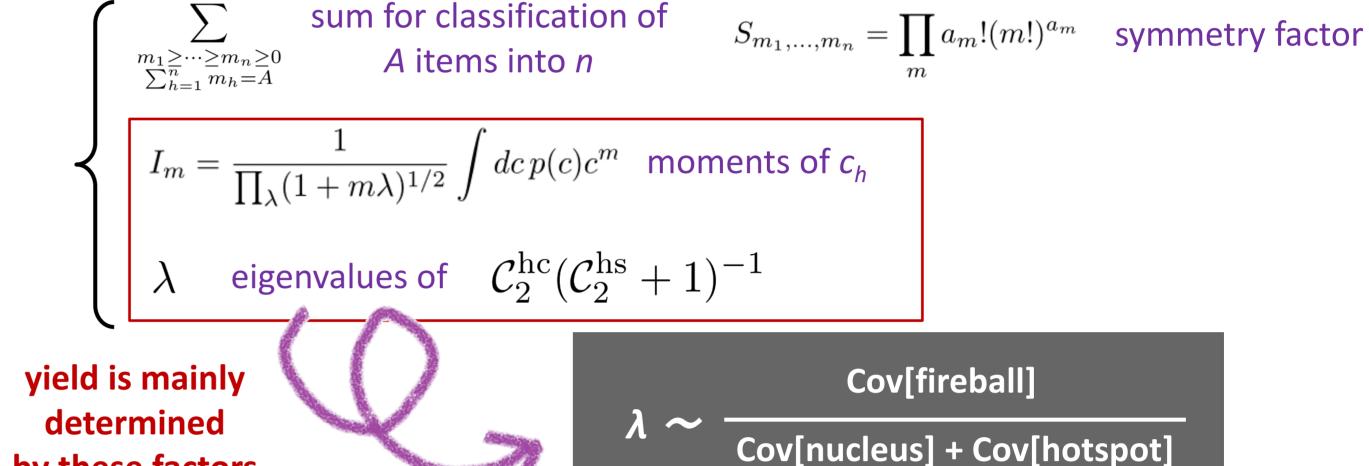


# 3. Results of *n*-Gaussian hot-spot model

■ Analytic result of yield of light nuclei (event averaged)

$$N_{A} = g_{A} 8^{A-1} \det(\mathcal{C}_{2}^{\text{hs}} + 1)^{-(A-1)/2}$$

$$\times \sum_{\substack{m_{1} \geq \dots \geq m_{n} \geq 0 \\ \sum_{h=1}^{n} m_{h} = A}} \frac{n! A!}{n^{A} S_{m_{1}, \dots, m_{n}}} \frac{I_{m_{1}} \cdots I_{m_{n}}}{\prod_{\lambda} \left(\sum_{h=1}^{n} \frac{m_{h}}{A} \frac{1}{1 + m_{h} \lambda}\right)^{1/2}},$$

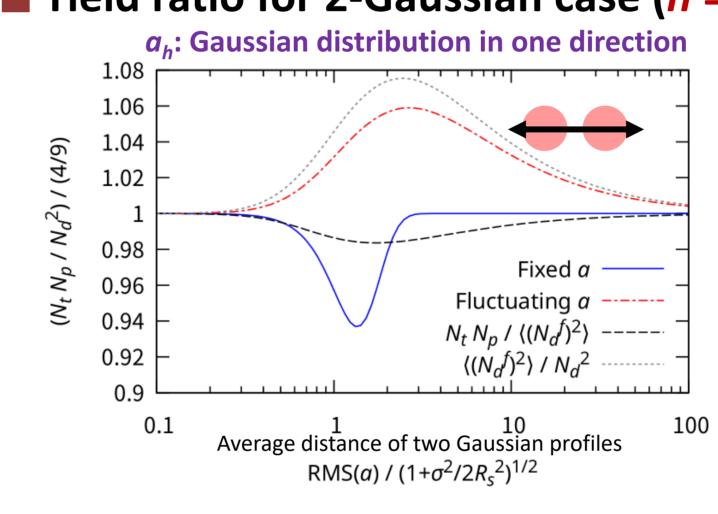


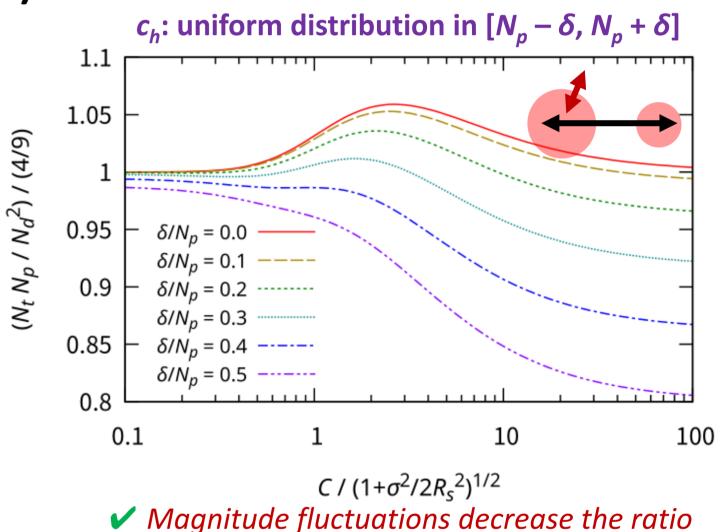
■ Yield ratio in limits n = 1 and  $n \rightarrow \infty$ 

by these factors

n=1: single Gaussian, Ratio = 4/9 (ideal value)  $n=\infty$ : infinite number of hot spots Ratio = 4/9 (fluctuations smeared out)

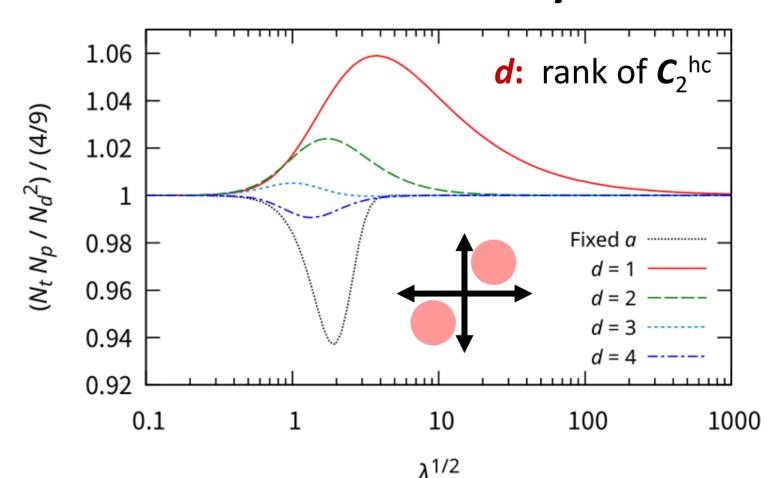
■ Yield ratio for 2-Gaussian case (n = 2)

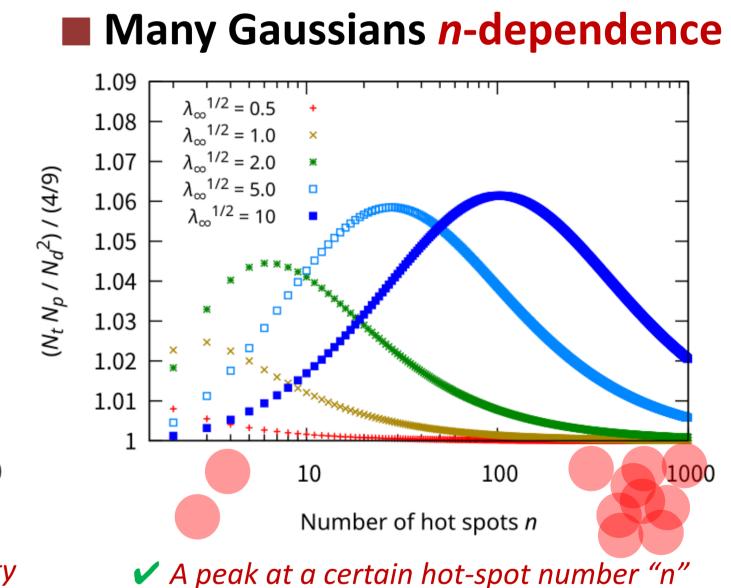




✓ Position fluctuations increase the ratio

## **■** Fluctuation dimensionality





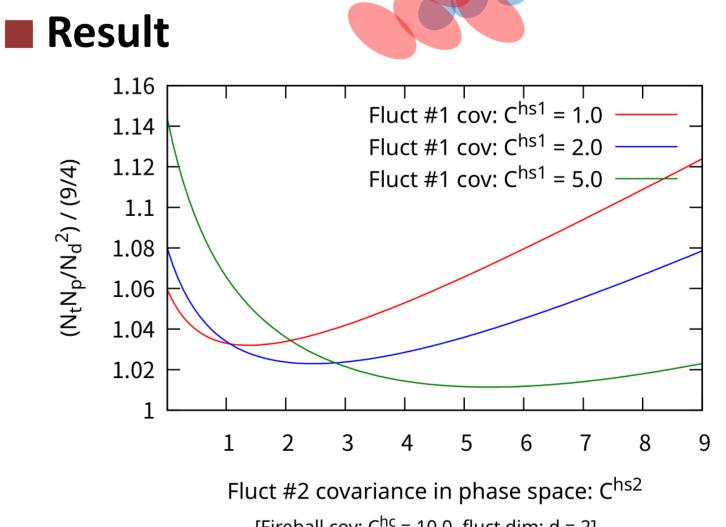
The increase largely depends on the dimensionality

# 4. Two fluctuation sources

## Extended Model

f(r, p): sum of two fluctuations sources with different hot-spot sizes Chs1 & Chs2. (E.g. critical fluctuations vs shorter-scale thermal fluctuations).

Question: When scale separation of two fluctuations are not sufficient in HIC, can we differentiate them?



Fluct #2 covariance in phase space: C<sup>1132</sup>

[Fireball cov: C<sup>hc</sup> = 10.0, fluct dim: d = 2]

The effect becomes larger when two sources have different sizes

# 5. Summary

Introduced "fluctuating n-Gaussian hot spots" for coalescence f(r, p)

- <u>Yield ratio of light nuclei is largely affected by e-by-e fluctuations:</u> Position/magnitude fluctuations increase/decrease the ratio. The dimensionality and # of hot spots also matter.
- Scale separation of two fluctuation sources make effect larger