

Multiparton interactions in pp collisions using charged-particle flattenicity

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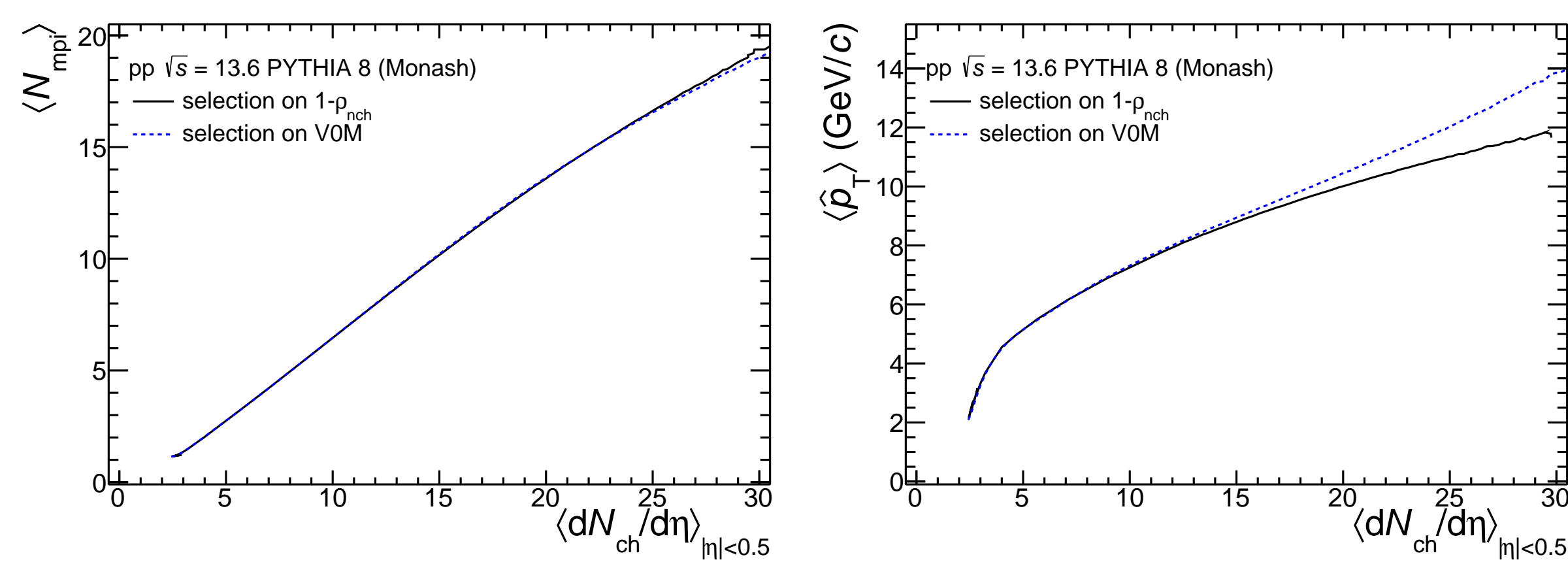


Introduction

The study of pp collisions is relevant given the discovery of collective-like effects and strangeness enhancement in high-multiplicity events. The origin of the effects is not fully understood yet because of several reasons:

- Medium-induced jet modifications have not been observed.
- The existing multiplicity estimators strongly bias the sample towards multijet final states.

The present work discusses that combining charged particle multiplicity (N_{ch}) and flattenicity (ρ), both measured at forward pseudorapidity (η), it is possible to control the bias [1]. The figure below shows that they are correlated with multiparton interactions (MPI: several parton-parton scatterings involving momentum transfers of a few GeV/c occurring in the same pp collision); however, the pp collisions tagged with N_{ch} are harder than those selected with ρ .



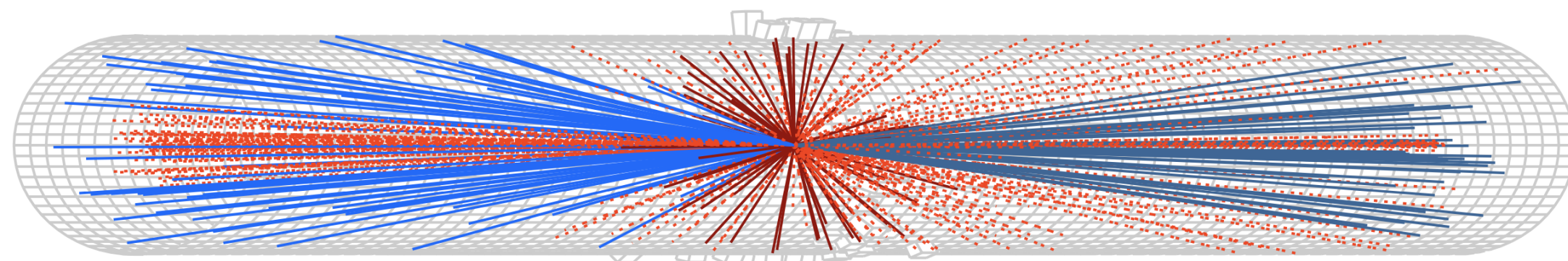
What is flattenicity?

Flattenicity, ρ :

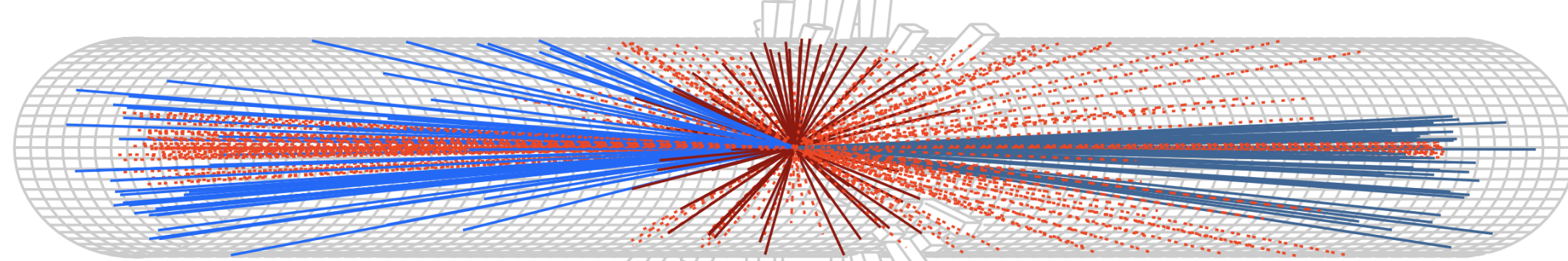
$$\rho = \sqrt{\frac{\sum_i (N_{ch}^{cell,i} - \langle N_{ch}^{cell} \rangle)^2 / N_{cell}^2}{\langle N_{ch}^{cell} \rangle}}, \quad (1)$$

is measured in the V0 detector of the Run 2 ALICE configuration, which consists of V0A ($2.8 < \eta < 5.1$) and V0C ($-3.6 < \eta < -1.7$) covering the full azimuth. Each subdetector has four rings along η and eight equidistant sectors in φ . This leads to a grid with 64 cells ($N_{cell} = 64$). In Eq. 1, $N_{ch}^{cell,i}$ is the multiplicity in the i -th cell and $\langle N_{ch}^{cell} \rangle$ is the event-by-event average multiplicity in the V0 cells. Figure 1 shows two event displays for high-multiplicity pp collisions produced with PYTHIA 8 [2] event generator.

PYTHIA 8.303 (Monash 2013), pp $\sqrt{s} = 13$ TeV, $N_{mp}=16$, $N_{ch}=318$, primary charged particles, flattenicity=0.11 (ALICE V0 acceptance)



PYTHIA 8.303 (Monash 2013), pp $\sqrt{s} = 13$ TeV, $N_{mp}=11$, $N_{ch}=301$, primary charged particles, flattenicity=0.17 (ALICE V0 acceptance)



— $|\eta| < 0.8$ (ALICE TPC acceptance)
— $2.8 < \eta < 5.1$ (ALICE V0A acceptance)
— $-3.6 < \eta < -1.7$ (ALICE V0C acceptance)

Figure 1: Event displays for “low ρ ” (top) and “high ρ ” (bottom) pp collisions simulated with PYTHIA 8.

For similar multiplicities, pp collisions are dominated by:

- Multijet-final states: high ρ (“hard” pp collisions)
- Minijets (MPI): low ρ (“soft” pp collisions)

Analysis details

- The p_T spectra of charged particles, pions, kaons and (anti)protons are measured as a function of forward multiplicity (V0M) and flattenicity in pp collisions at $\sqrt{s} = 13$ TeV.
- For p_T below 2-3 GeV/c, particle identification (PID) is done exploiting the dE/dx provided by the Time Projection Chamber, and the particle velocity measured with the Time-Of-Flight detector. For higher p_T , the PID is performed deconvoluting the dE/dx spectrum measured in the relativistic rise regime of the Bethe-Bloch curve.
- The analysis follows the well established methods described in several ALICE publications [3].

Results

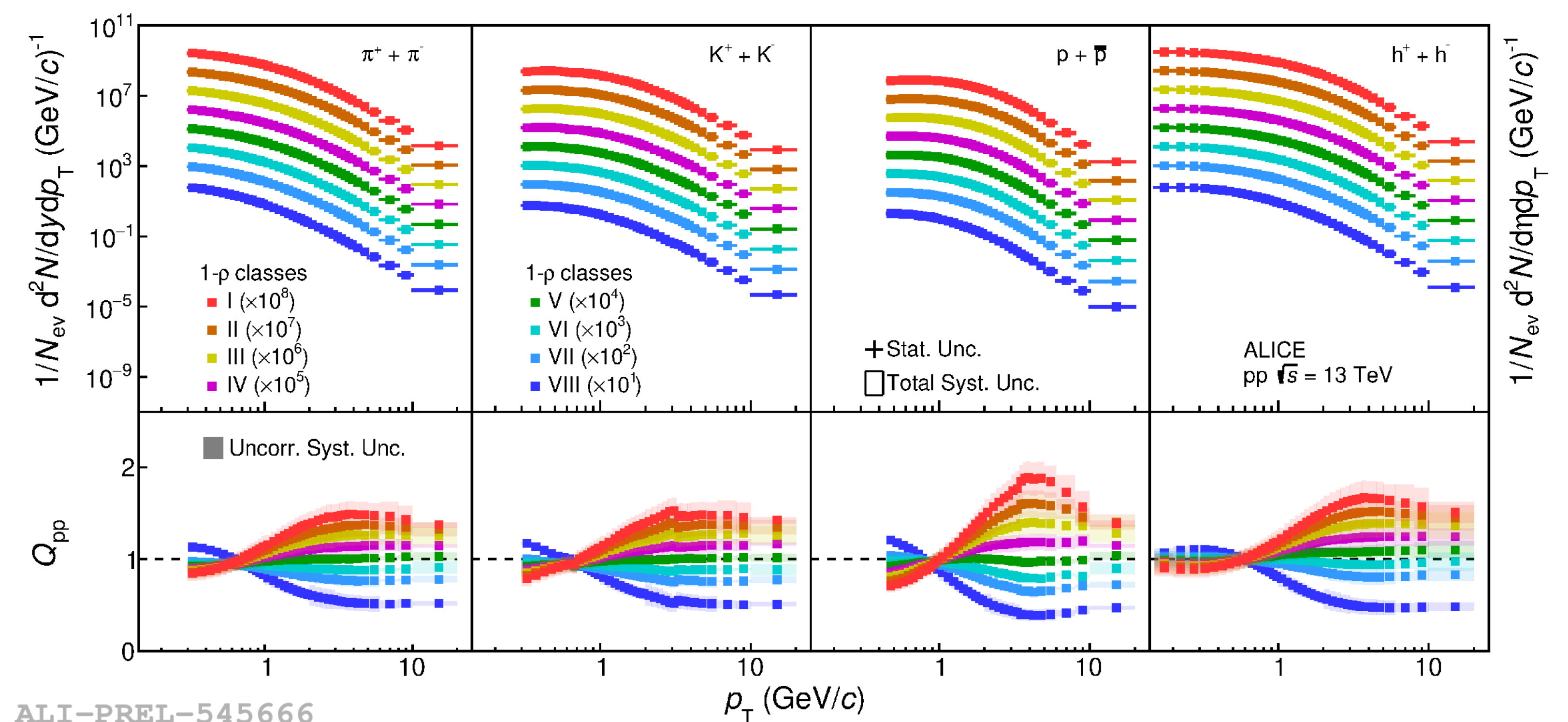
The sample is divided into event classes based on flattenicity. The table below lists the average charged-particle pseudorapidity densities $\langle dN_{ch}/d\eta \rangle$ within $|\eta| < 0.8$ for the different flattenicity classes, it shows the implicit multiplicity dependence of flattenicity.

I	II	III	IV	V	VI	VII	VIII
0-1%	1-5%	5-10%	10-20%	20-30%	30-40%	40-50%	50-100%
22.2 ± 0.7	18.2 ± 0.6	15.3 ± 0.5	12.6 ± 0.4	10.0 ± 0.3	8.06 ± 0.19	6.47 ± 0.13	3.51 ± 0.04

The p_T spectra as a function of flattenicity are shown in the figure below. From red (0-1% 1- ρ) to blue (50-100% 1- ρ) the spectra get harder for $p_T > 5$ GeV/c. The spectral shape modification is studied with Q_{pp} , which is defined as follows:

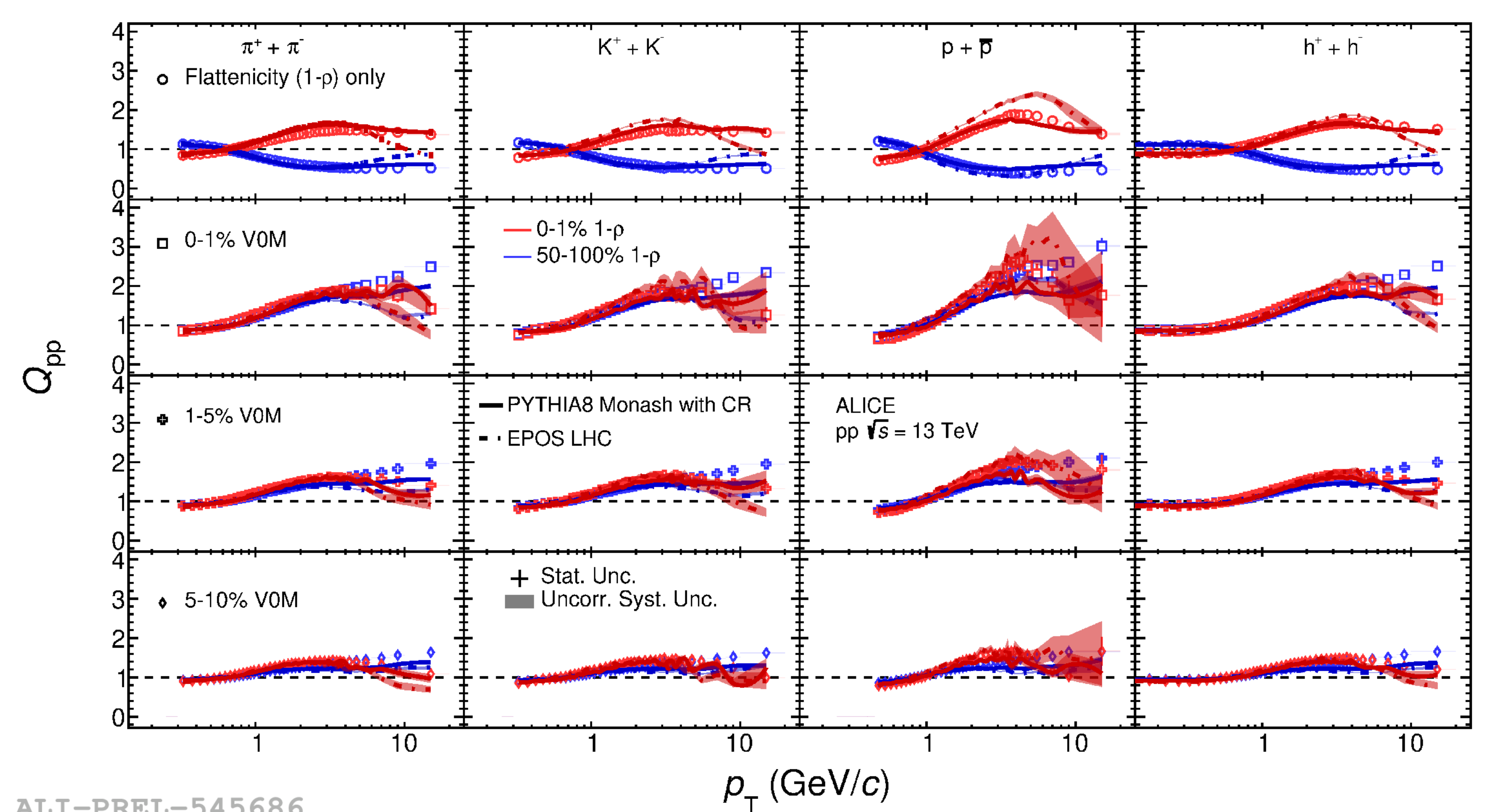
$$Q_{pp} \equiv [d^2N / \langle N_{ch} \rangle dy dp_T]^{\rho_{class}} / [d^2N / \langle N_{ch} \rangle dy dp_T]^{\text{minimum bias}}. \quad (2)$$

For independent parton-parton scatterings occurring in the same pp collision: $Q_{pp} \rightarrow 1$. However, for all flattenicity classes Q_{pp} seems to converge to unity only at high p_T .



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The figure below shows Q_{pp} for pp collisions with fixed $\langle N_{ch} \rangle$ values. In the highest multiplicity class (0-1% V0M), Q_{pp} is flattenicity independent for $p_T < 3$ GeV/c. However, for higher p_T , Q_{pp} keeps rising for the 50-100% 1- ρ event class. Whereas for the 0-1% 1- ρ class, Q_{pp} reduces with increasing p_T developing a bump structure. The other event classes exhibit a similar behaviour. Data are compared with EPOS LHC and PYTHIA 8.



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Conclusions

- In contrast with previous measurements, Q_{pp} exhibits a reduction (bump) at low- (intermediate-) p_T going from low to high $\langle N_{ch} \rangle$ values. This bump is mass dependent. At higher p_T , Q_{pp} seems to approach to unity.
- Overall, data are better described by PYTHIA (with color reconnection).

References

- [1] A. Ortiz, A. Khuntia, O. Vázquez-Rueda, S. Tripathy, G. Bencedi, S. Prasad, and F. Fan, “Unveiling the effects of multiple soft partonic interactions in pp collisions at $\sqrt{s}=13.6$ TeV using a new event classifier,” *Phys. Rev. D* **107** no. 7, (2023) 076012, arXiv:2211.06093 [hep-ph].
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- [3] ALICE Collaboration, S. Acharya *et al.*, “Production of charged pions, kaons, and (anti-)protons in Pb-Pb and inelastic pp collisions at $\sqrt{s_{NN}} = 5.02$ TeV,” *Phys. Rev. C* **101** no. 4, (2020) 044907, arXiv:1910.07678 [nucl-ex].

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