



# Insights on strange quark hadronization in small collision systems with ALICE: multiple strange hadrons and $\Sigma^\pm$ baryons

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on behalf of the ALICE Collaboration



ALICE

QM 2023 - Houston  
03-09 Sept.

1. Università degli Studi di Torino

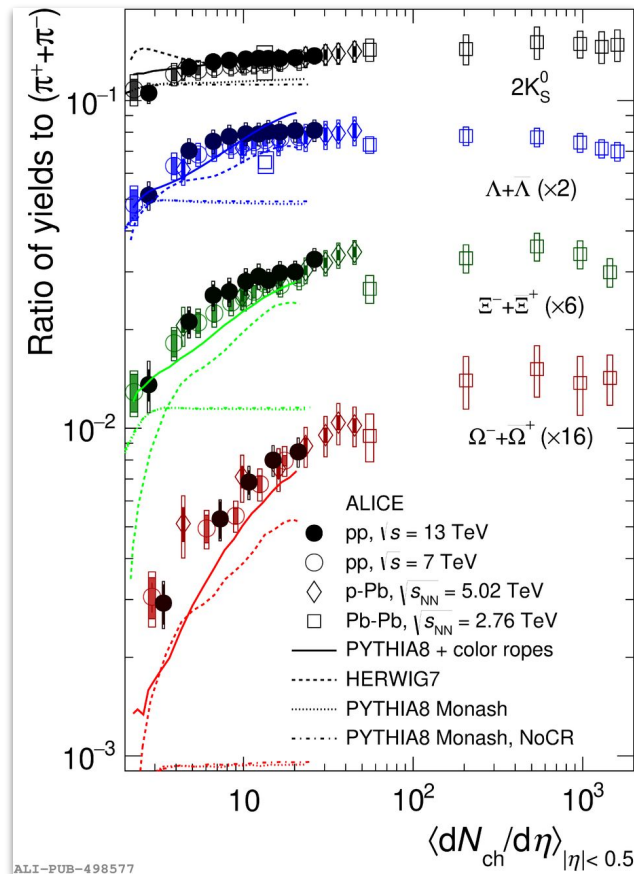
2. INFN Torino





## Strangeness Enhancement (SE):

- $S/\pi$  increases as a function of multiplicity identically across different energies and collision systems.
- Enhancement proportional to the strangeness content in the hadron



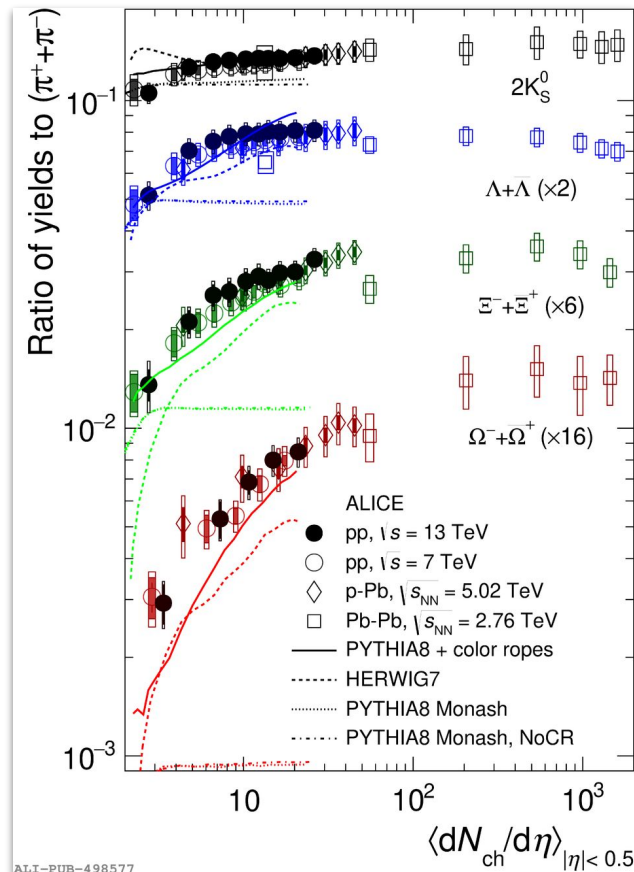


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## New in this talk:

- **(anti) $\Sigma^\pm$  baryons**  $\rightarrow$  same strangeness content as  $\Lambda$
- **Probability Density Function (PDF)** for  $K_S^0$ ,  $\Lambda$ ,  $\Xi$ ,  $\Omega$ 
  - extend beyond the average of the distribution
  - **unique opportunity to test the connection between charged and strange particle multiplicity production**





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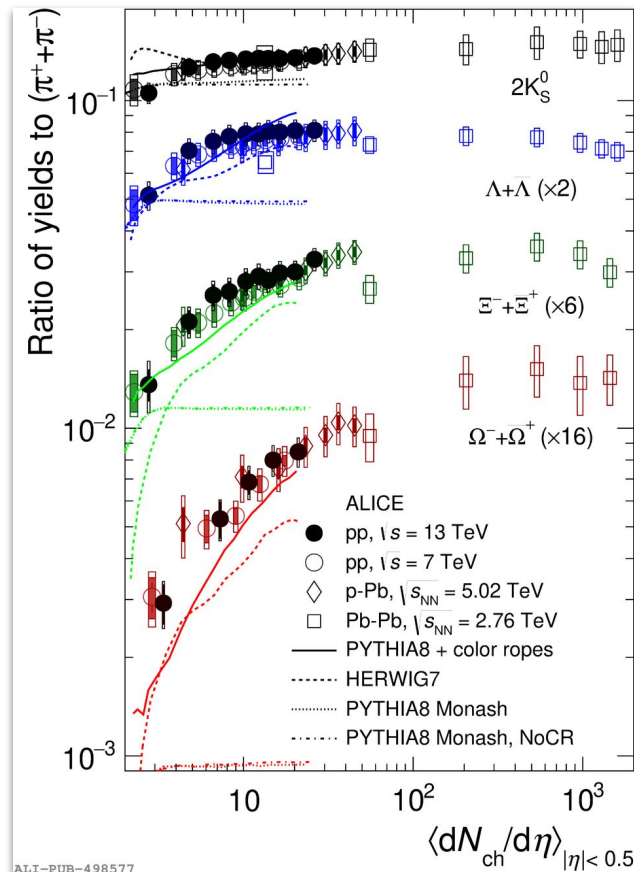
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## Comparison to models:

- in-vacuum hadronization (e.g. [Pythia](#), [AMPT](#), ...)
- thermal production: Statistical Hadronization Model (e.g. [GSI-Heidelberg](#), [FIST](#), ...)
- two-component models (e.g. [EPOS](#), [DCCI](#), ...)





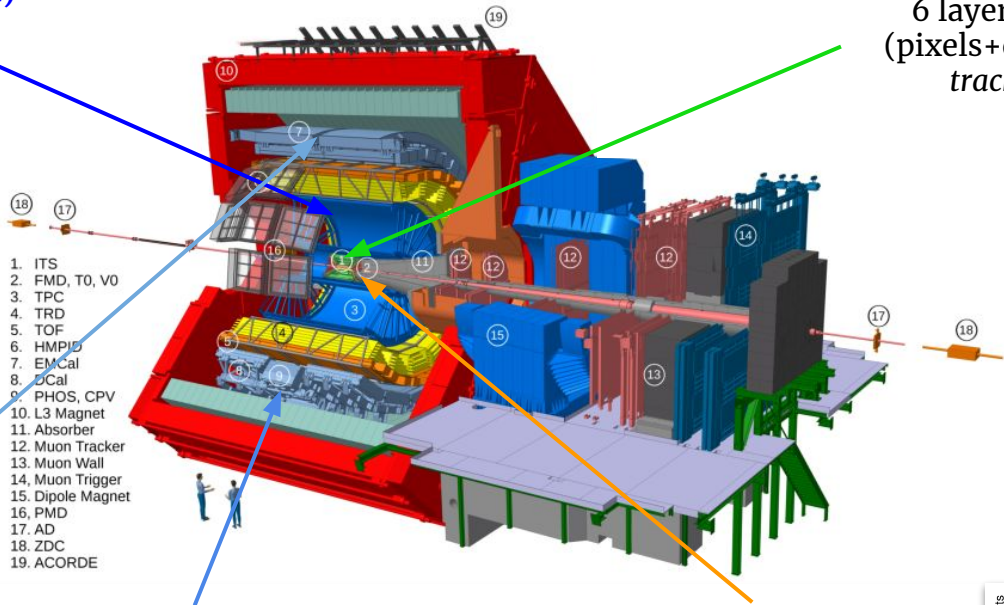


## Time Projection Chamber (TPC)

Gas filled detector  
tracking, PID ( $dE/dx$ )

## Inner Tracking System (ITS)

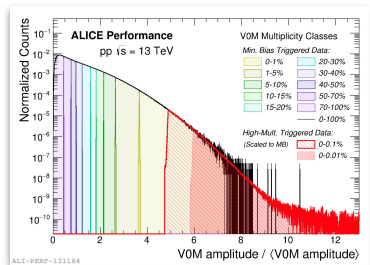
6 layers of silicon detectors  
(pixels+drift+strips) triggering,  
tracking, vertexing, PID



ElectroMagnetic  
CALorimeter (EMCal)  
Pb-scintillator  
sampling calorimeter

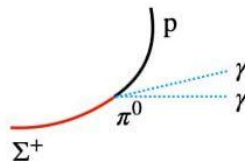
PHOTon Spectrometer (PHOS)  
High-granularity  
electromagnetic calorimeter  
( $\text{PbWO}_4$ )

V0 detectors (V0A, V0C)  
Forward-rapidity arrays  
of scintillators  
triggering, particle  
multiplicity estimation

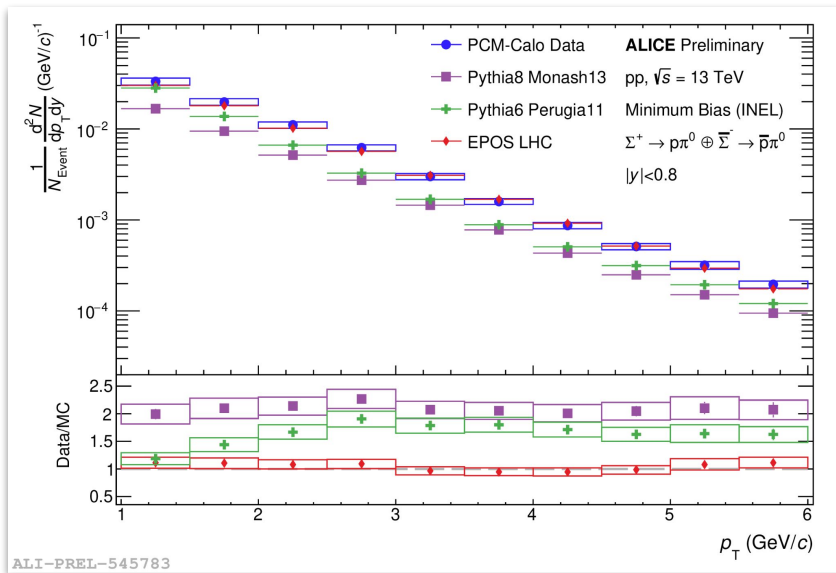




Photons reconstructed in the calorimeters  
or via the photon conversion method (PCM).



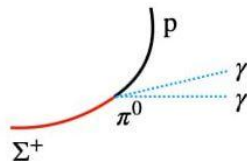
New



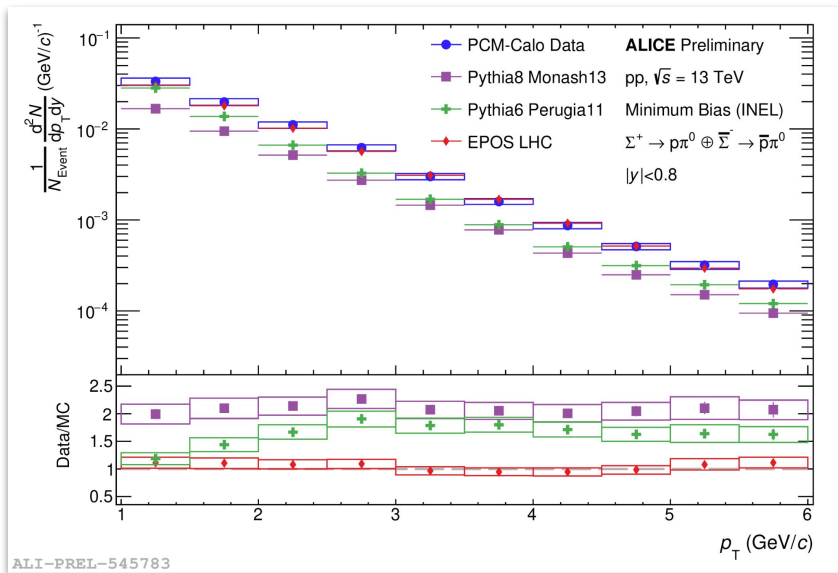
$p_T$ -shape well reproduced by EPOS LHC and Pythia8,  
but the latter underestimates the yield



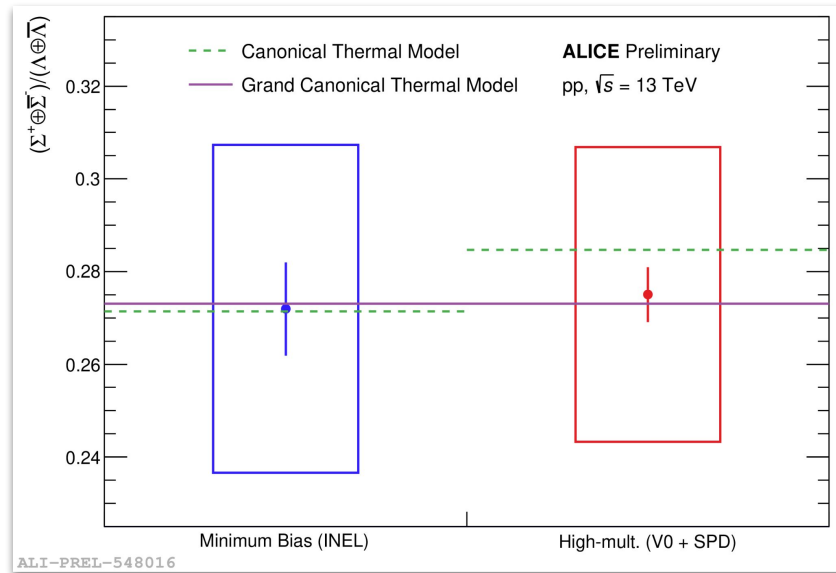
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New



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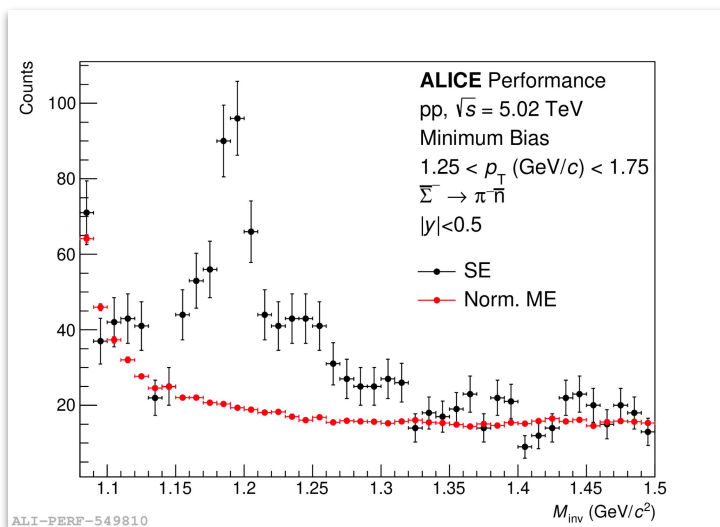
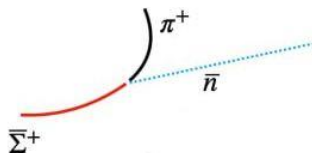


$\Sigma/\Lambda$  does not depend on multiplicity (large uncertainties)  
and is reproduced by the thermal model



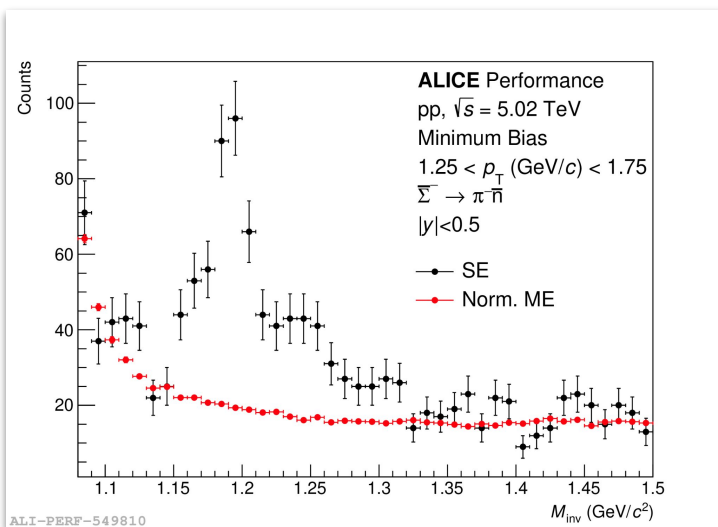
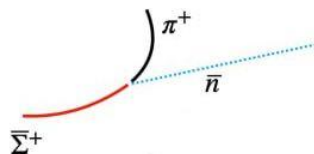
New

$\bar{n}$  detection from signal in the PHOS calorimeter +  
momentum reconstructed using time-of-flight



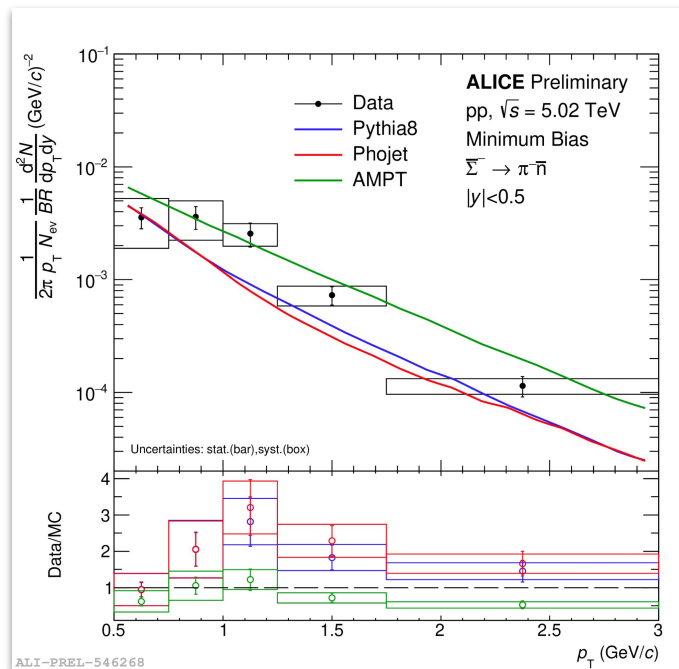


$\bar{n}$  detection from signal in the PHOS calorimeter + momentum reconstructed using time-of-flight



Production yield extracted in several  $p_T$  bins and compared to phenomenological models

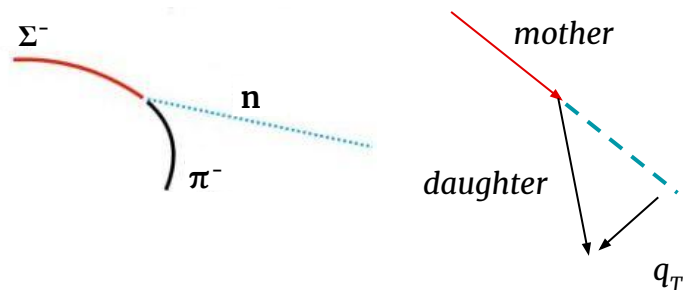
New



Very promising technique.  
Uncertainties will shrink with Run 3

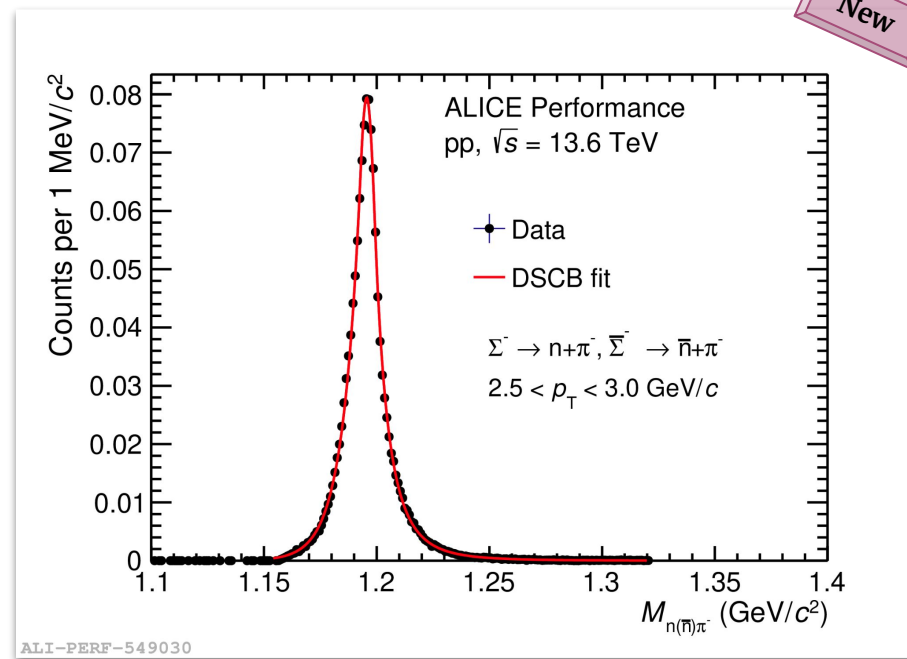


→ B. Heybeck's poster (409)



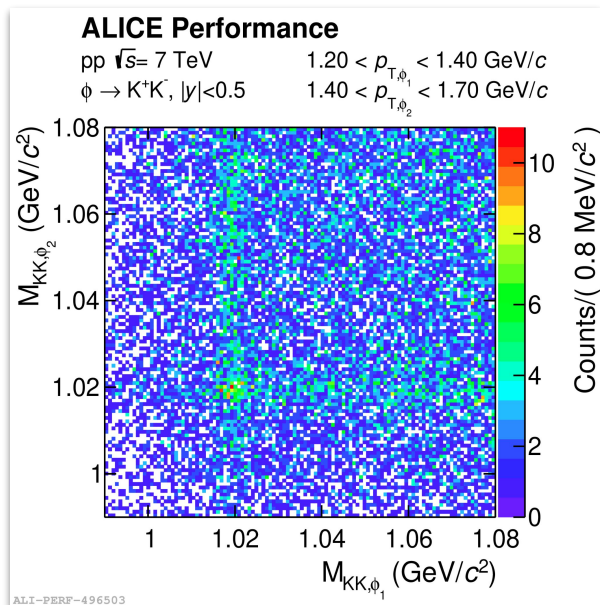
## Kink-topology using the upgraded ITS2 detector

- combine same charge tracks
- minimize the distance between mother and daughter + apply several topological cuts
- $q_T$  = momentum of the neutral daughter = difference between mother and charged daughter momentum  
→ invariant mass of the mother can be evaluated





Measurement of the  $\phi$  meson pair yield at mid-rapidity in MB pp collisions at  $\sqrt{s} = 7$  TeV

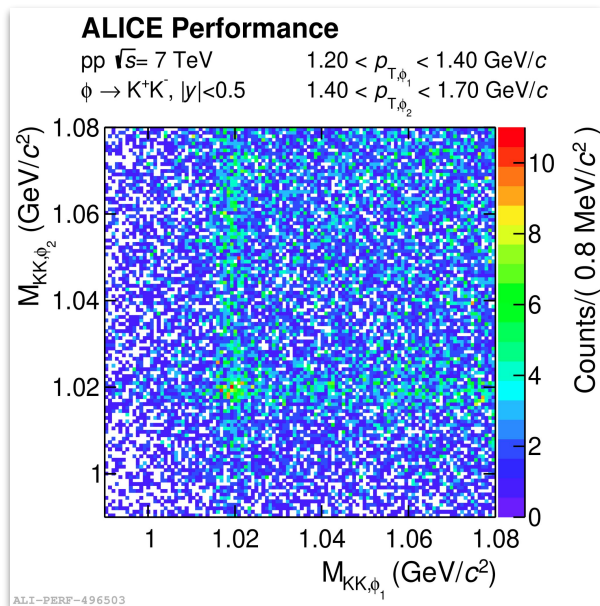


- 2D invariant mass technique in order to extract the signal
- Yield and  $p_T$  distribution extraction

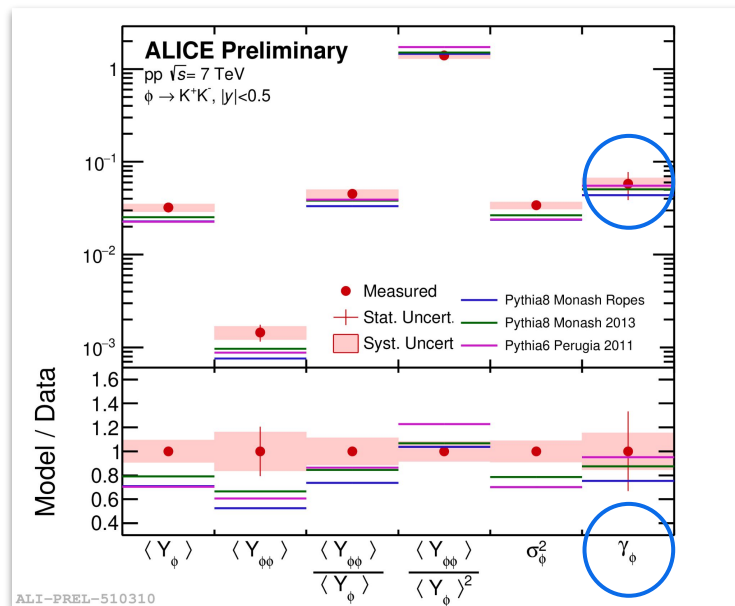




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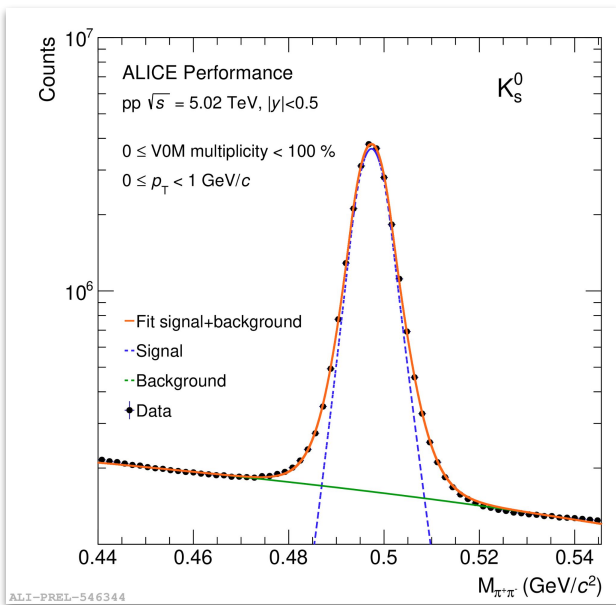
$$\gamma_\phi = \frac{\sigma_\phi^2}{\mu_\phi} - 1 = 2 \frac{\langle Y_{\phi\phi} \rangle}{\langle Y_\phi \rangle} - \langle Y_\phi \rangle$$

- $\gamma_\phi = 0$  for poisson distribution
- observed  $\gamma_\phi > 0 \rightarrow$  not poissonian



New

Analysis based on counting the number of strange particles event-by-event in pp collisions at  $\sqrt{s} = 5.02$  TeV

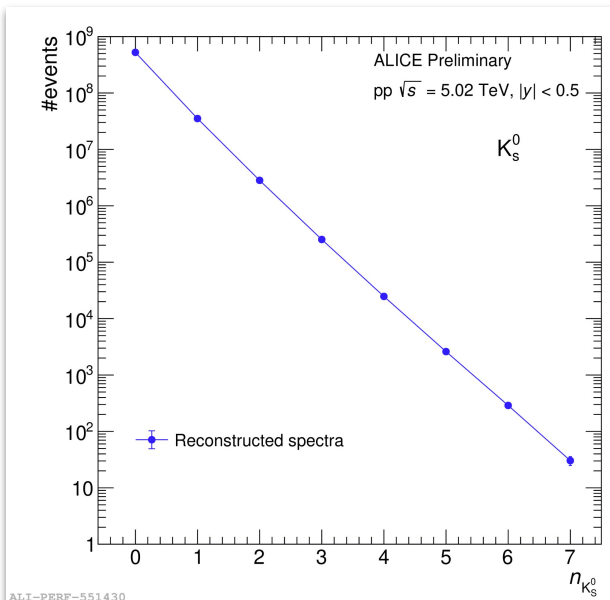
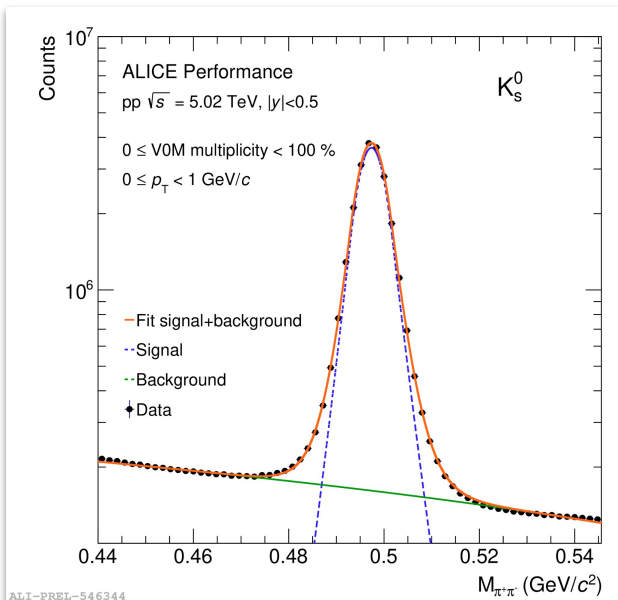


Each candidate **weighted by  $P(S)$**   
or  **$P(B)$**  estimated by **1D** invariant  
mass **fit** in  $p_T$ /multiplicity bins



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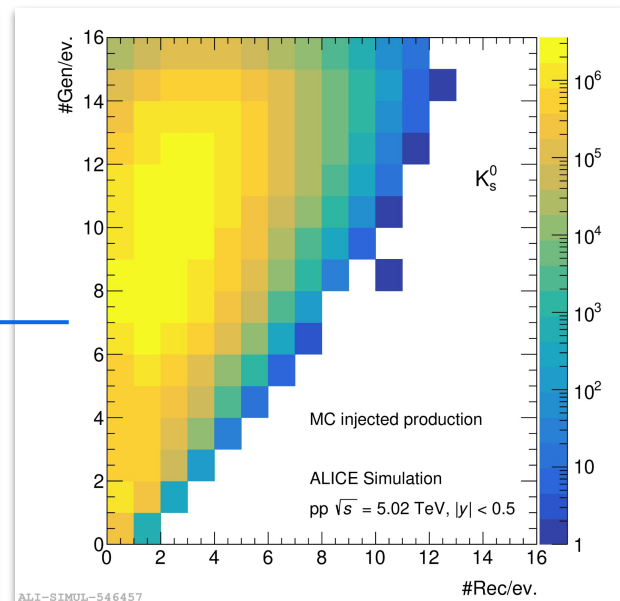
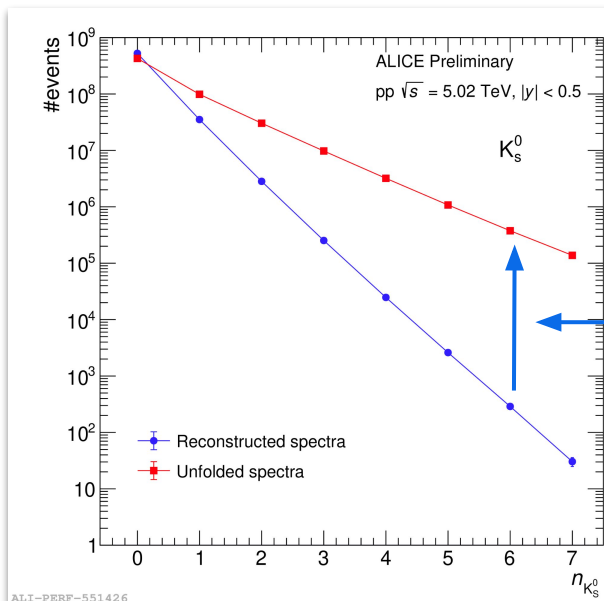
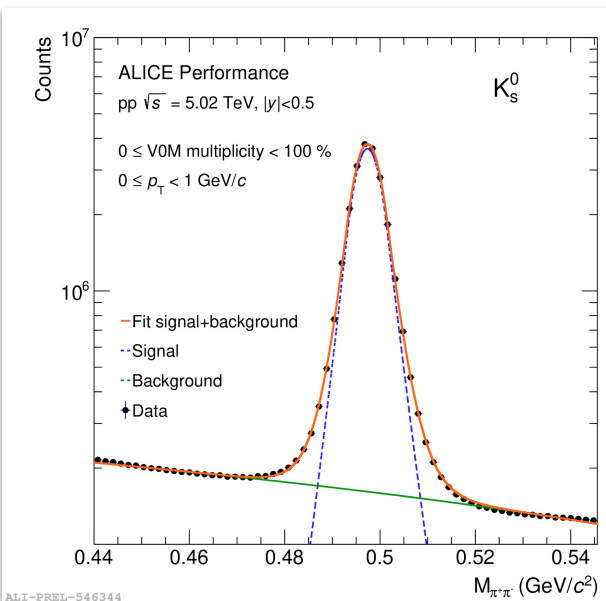
**Weights** associated to each of the  $N$  candidates **combined** to obtain:  
 $P(\text{all-sig}), \dots, P(\text{all-bkg})$

For each event: full probability spectrum spanning from 0 to  $N$



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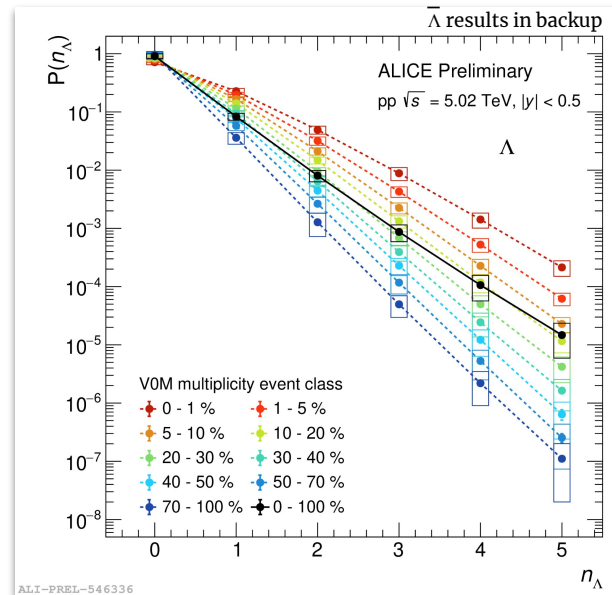
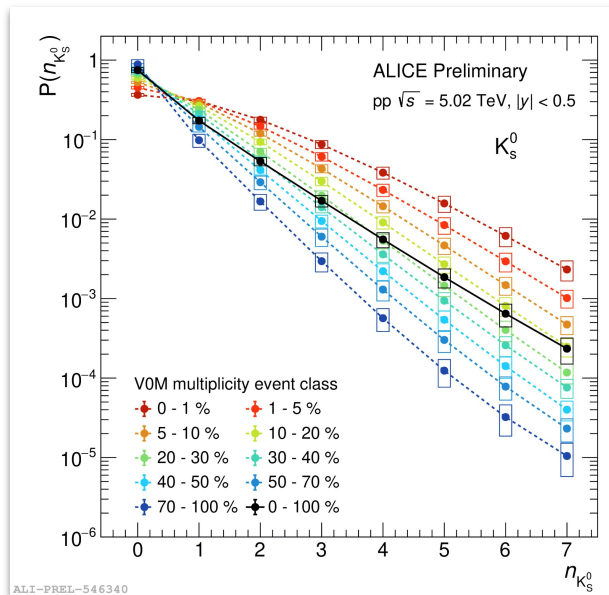
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Correction for detector response (MC production: measured  $p_T$  distribution)

**Bayesian unfolding** procedure applied

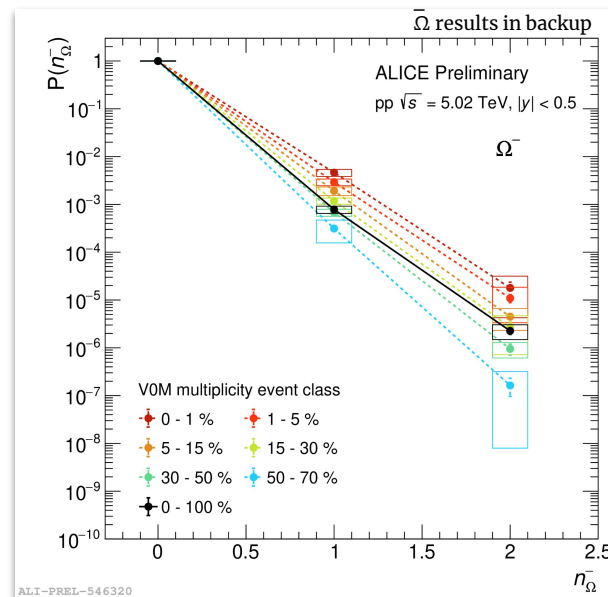
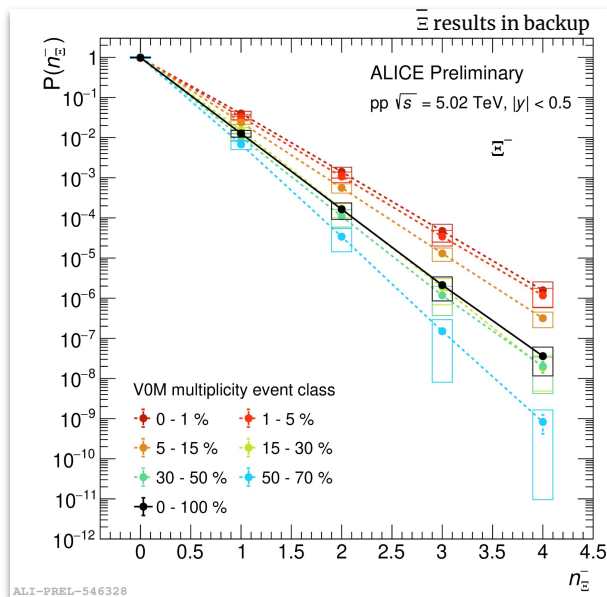


Probability to produce  $n$  particle ( $n$  up to 7 for  $K_s^0$ , 5 for  $\Lambda$ ) of a given species per event

Spanning across large ranges of strange/multiplicity variations, all the way to very “extreme” situations  
(e.g. 7  $K_s^0$  at low average multiplicity, 0  $K_s^0$  at high multiplicity)

Unique opportunity to test the connection between  
charged and strange particle multiplicity production

NOTE: in each VOM bin multiplicity can fluctuate and  $\langle dN_{ch}/d\eta \rangle$   
can significantly change for events with small/large  $n_s$



Probability to produce  $n$  particle ( $n$  up to 4 for  $\Xi$ , 2 for  $\Omega$ ) of a given species per event

Spanning across large ranges of strange/multiplicity variations, all the way to very “extreme” situations  
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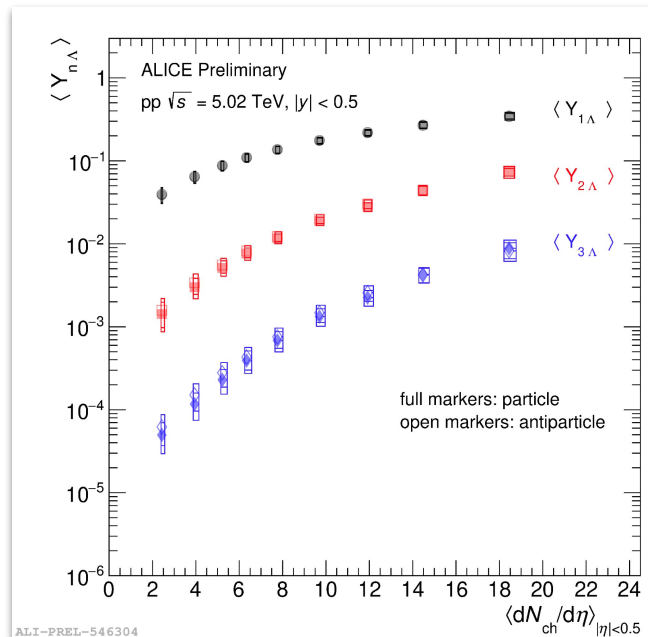
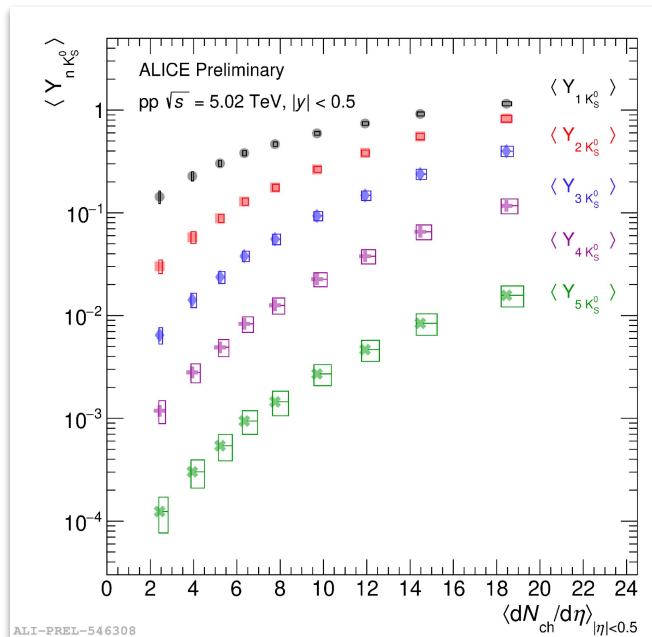
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The PDF allows to calculate the **production yield of 1, 2, 3, ... particles/event**:

$$\langle Y_{k-part} \rangle = \sum_{n=k}^{\infty} \frac{n!}{k!(n-k)!} P(n)$$



The increase with multiplicity of the probability for multiple strange hadrons is more than linear

NOTE: very good agreement between  $\langle Y_{1-part} \rangle$  and previous results ([1],[2])

[1] ALICE, [Nature Physics v13, pages 535–539](#) (1617)

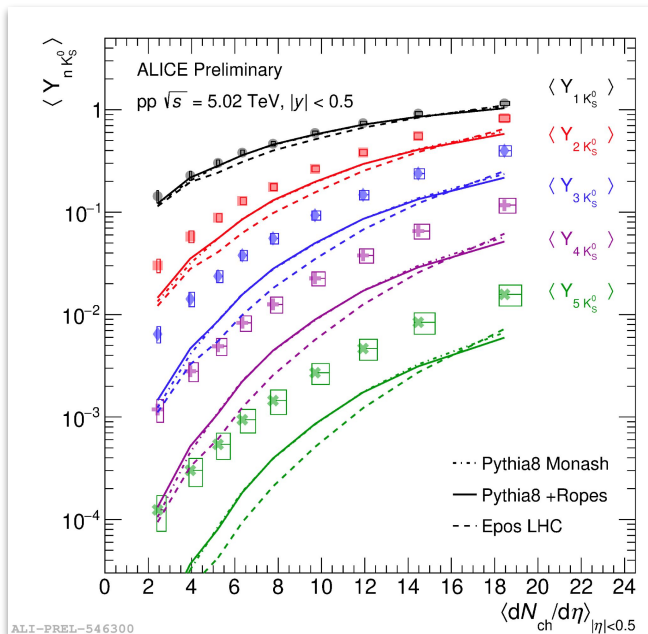
[2] ALICE, [Eur. Phys. J. C 80, 167](#) (1616)



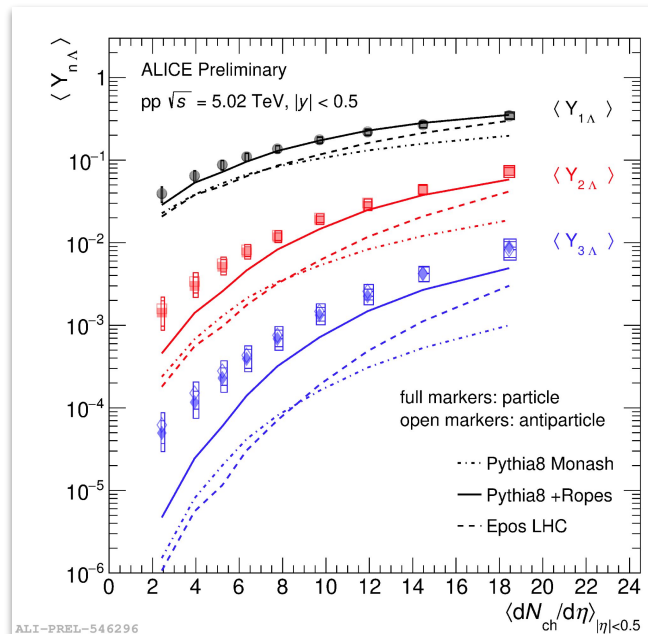


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No difference between Pythia8 Monash and Ropes for  $K^0$ : Pythia8 + Ropes (with QCD-CR) tends to increase baryons

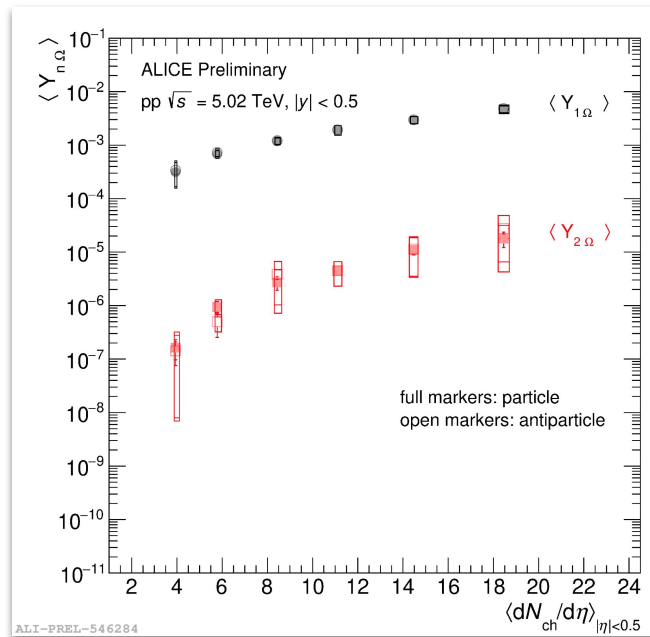
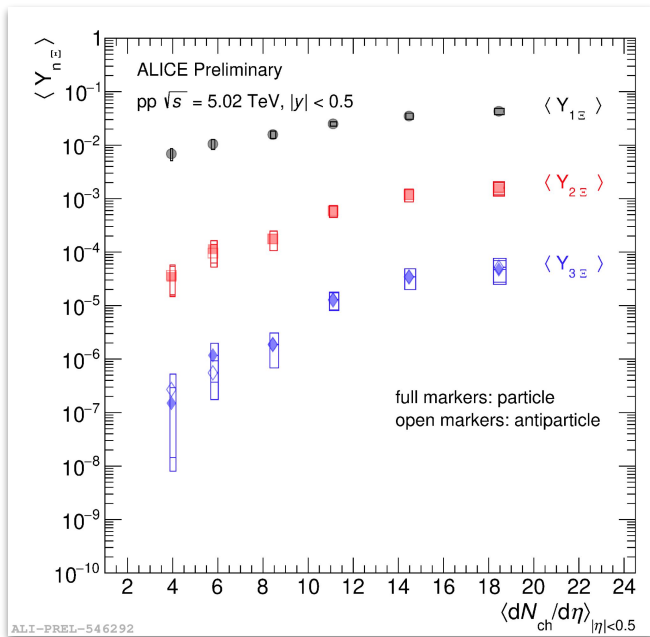


Ropes approaches the data at high multiplicity for  $\Lambda$   
Epos LHC does a rather good job at high multiplicity, but shows larger discrepancy at low multiplicity



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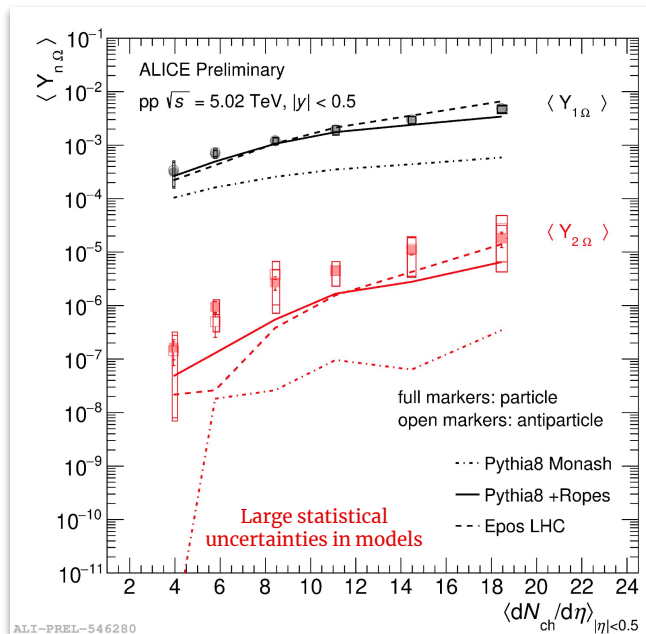
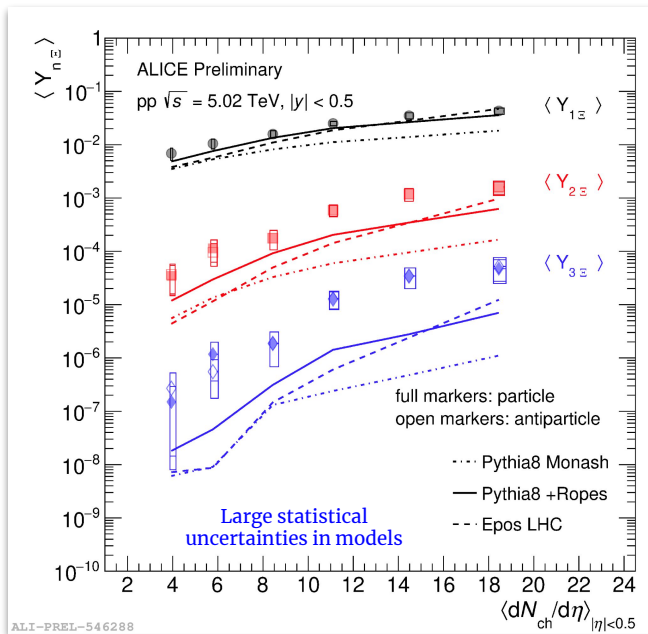
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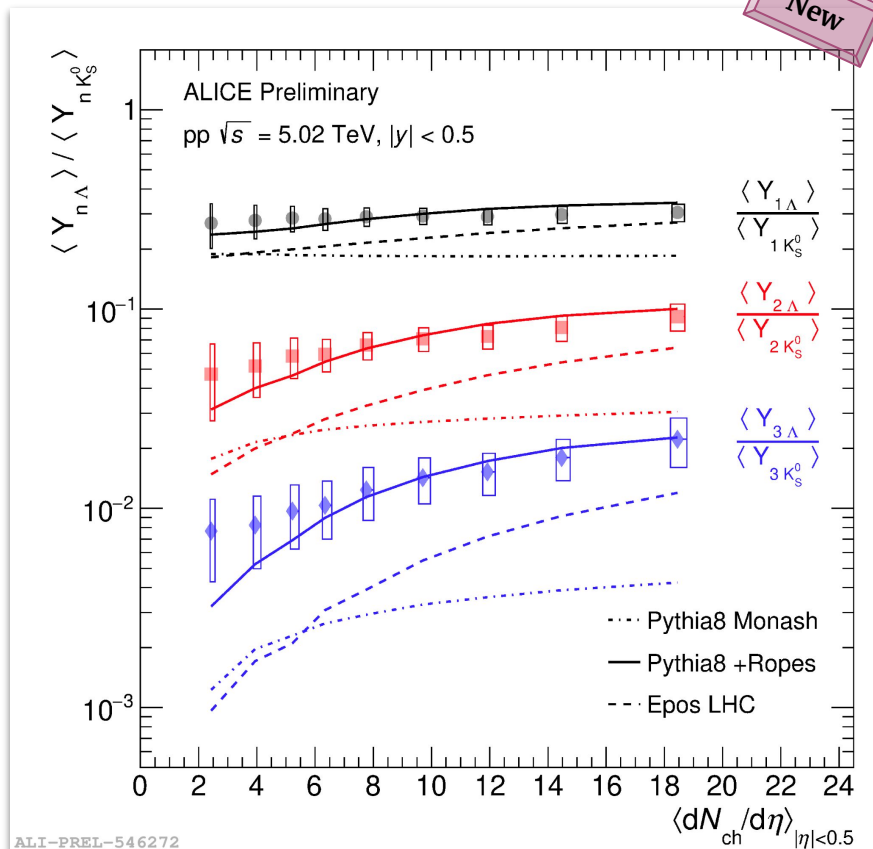
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Larger MC statistics will be produced to compare models in  $\langle Y_{k-part} \rangle$



Very important to disentangle baryon-related from strangeness-related effects!

- Increase of  $\Lambda/K_s^0$  VS multiplicity when looking at multiple production!
- Possibly in all strange-hadron/ $\pi$  VS multiplicity plots we have a strangeness-related AND a baryon-related contribution to the enhancement
- Baryon-related effect **well reproduced by Ropes** (with QCD-CR) at high multiplicity





- **First measurement of (anti) $\Sigma^\pm$  baryons at LHC energies:**

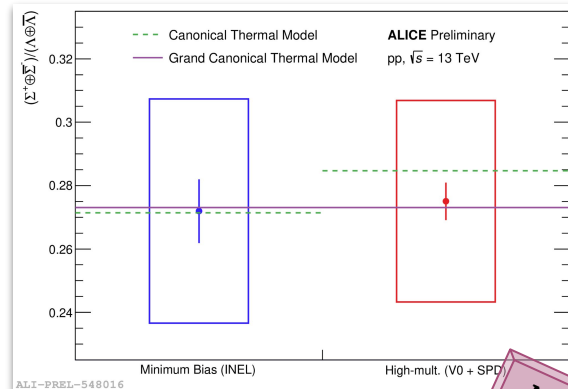
→ 2 different techniques tested with Run 2 data

- $p_T$  distribution and  $\Sigma/\Lambda$  compared with phenomenological models
- $\Sigma/\Lambda$  does not depend on multiplicity

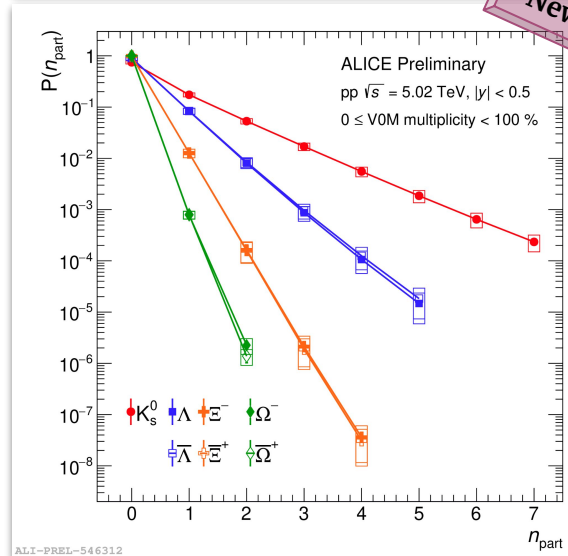
- **Multiple strange hadron production:**

→ first measurements of the PDF for (multi-)strange particles

- perfect benchmark to test production models in events spanning from extreme unbalances between charged and strange particle multiplicity
- 2- and 3-  $\Lambda/K_S^0$  yield ratios increase with multiplicity (baryon-related effect)



New





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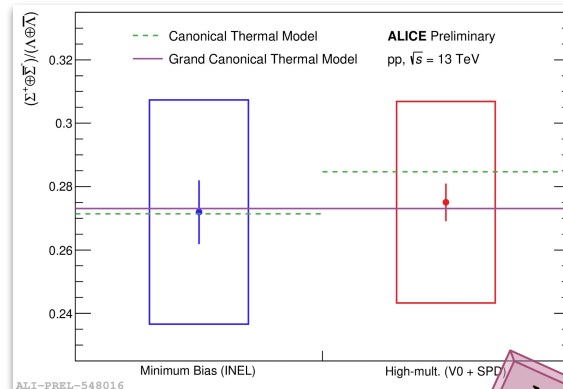
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## Run 3:

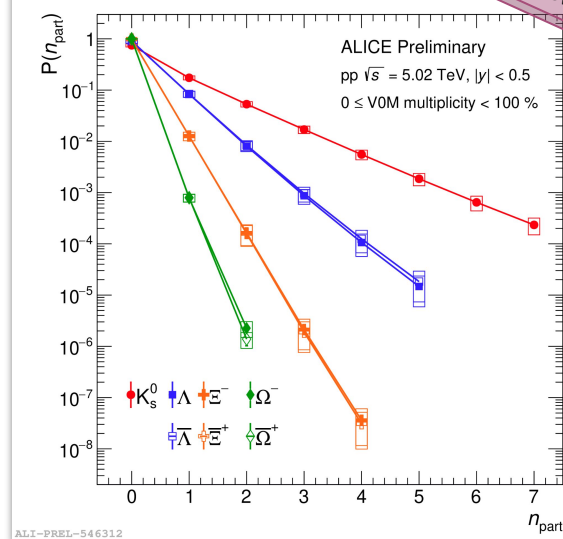
- will allow to apply kink topology to the measurement of (anti) $\Sigma^\pm$  baryons
- larger statistics (3/4 orders of magnitude higher) useful for cascade analyses
- for the future: extended PDF study to higher number of particles/event

→ B. Heybeck's poster (409)

→ F. Ercolessi's talk (Tuesday, h 12.40)  
C. De Martin's poster (150)



New



# Backup slides

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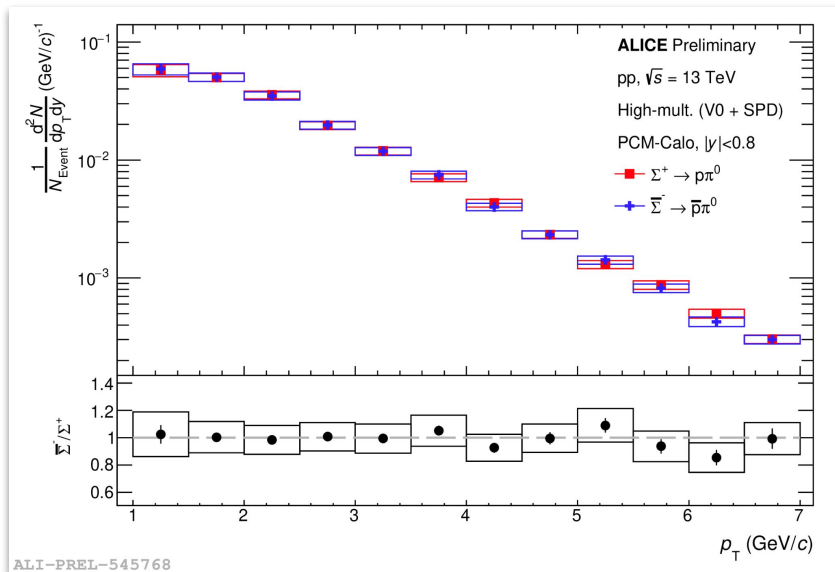


1. Università degli Studi di Torino

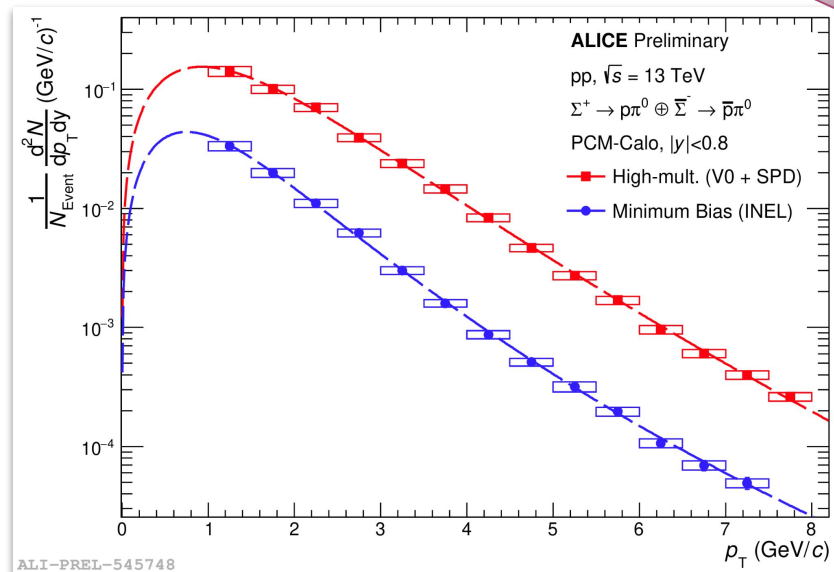
2. INFN Torino







Comparison between particle and antiparticle  $p_T$  spectra  $\rightarrow$  good agreement



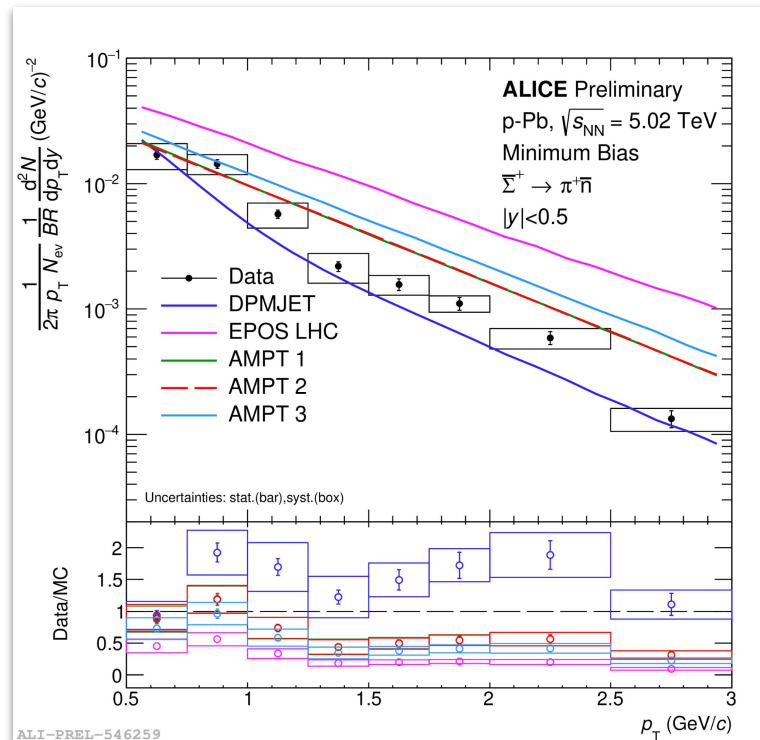
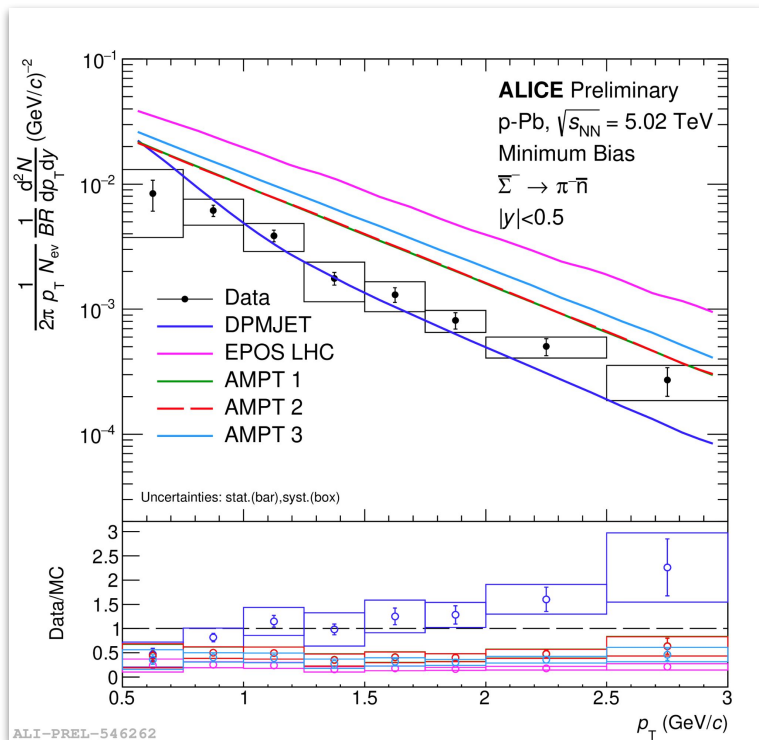
$\rightarrow$  Comparison between particle and antiparticle  $p_T$ -spectra in high multiplicity and minimum bias triggered pp collisions

$\rightarrow$  Results have been fitted with a Levy-Tsallis function



New

Production yield extracted in several  $p_T$  bins and compared to phenomenological models





*iterative procedure based on the Bayes' theorem using a picture of causes C ("true values") and effects E ("observed values")*

$$P(C_i | E_j) = \frac{P(E_j | C_i) \cdot \pi(C_i)}{\sum_{i=1}^{n_C} P(E_j | C_i) \cdot \pi(C_i)}$$

$P(E_j | C_i)$  estimated by using Monte Carlo (response matrix)  
 $P(C_i | E_j) \rightarrow$  probability that different  $C_i$  were responsible for the observed effect  $E_j \rightarrow$  GOAL  
 $\pi(C_i) \rightarrow$  prior probabilities (initially arbitrary, but updated on subsequent iterations)

- Choosing a prior distribution in order to apply Bayes' theorem  $\rightarrow$  posterior probability matrix obtained
- Applied to "observed spectra"  $\rightarrow$  1<sup>st</sup> estimation of the corrected spectra
- The corrected spectra obtained in the previous step becomes the prior probability and the correction proceeds as before
- Procedure is re-iterated until stability is achieved (**regularization parameter:  $n_{\text{iter}}$** )

$$\hat{n}(C_i) = \frac{1}{\epsilon_i} \sum_{j=1}^{n_E} n(E_j) \cdot P(C_i | E_j) = \sum_{j=1}^{n_E} M_{ij} \cdot n(E_j)$$

expected number of events in the cause bin  $i$

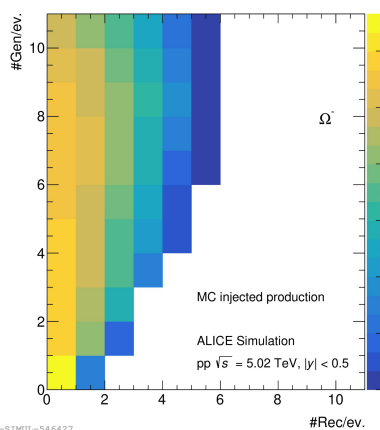
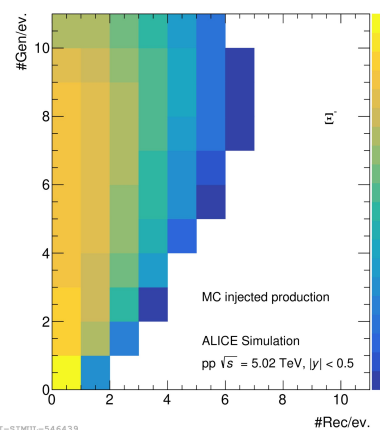
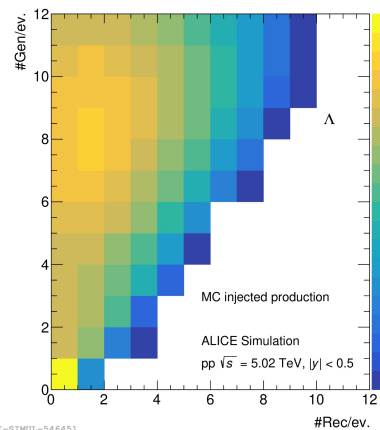
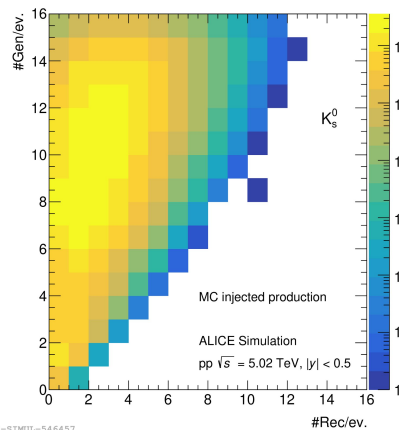
$\rightarrow M_{ij}$  is the unfolding matrix:  $M_{ij} = \frac{P(E_j | C_i) \cdot \pi(C_i)}{\epsilon_i \cdot \sum_{l=1}^{n_C} P(E_j | C_l) \cdot \pi(C_l)}$

$\rightarrow n(E_j)$  measurements (effects)

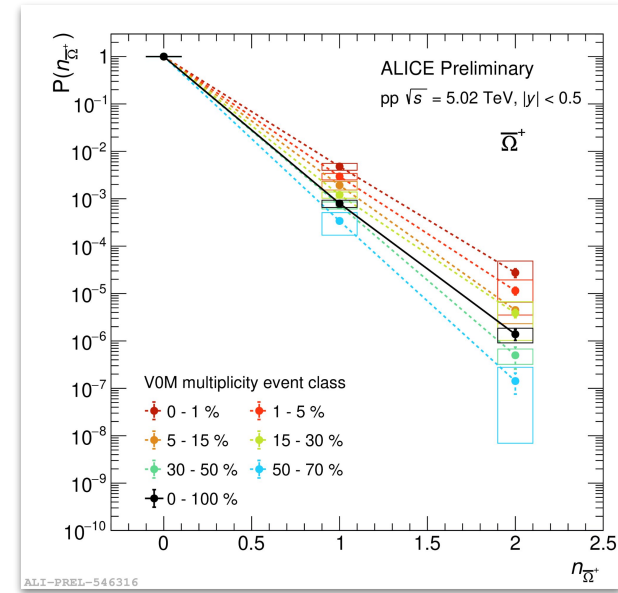
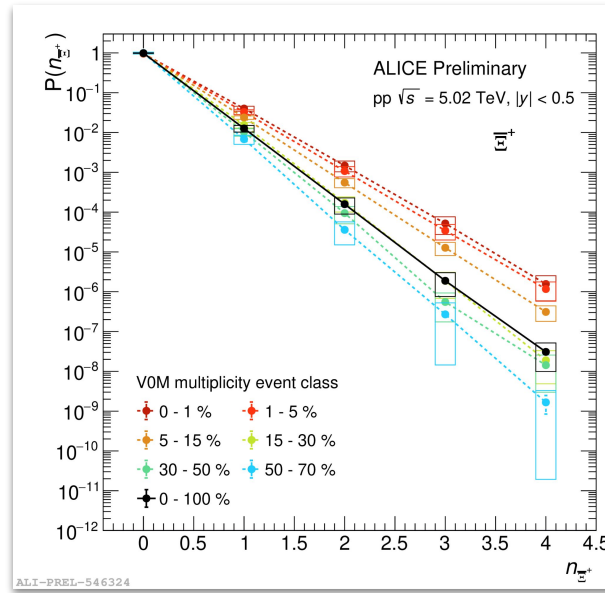
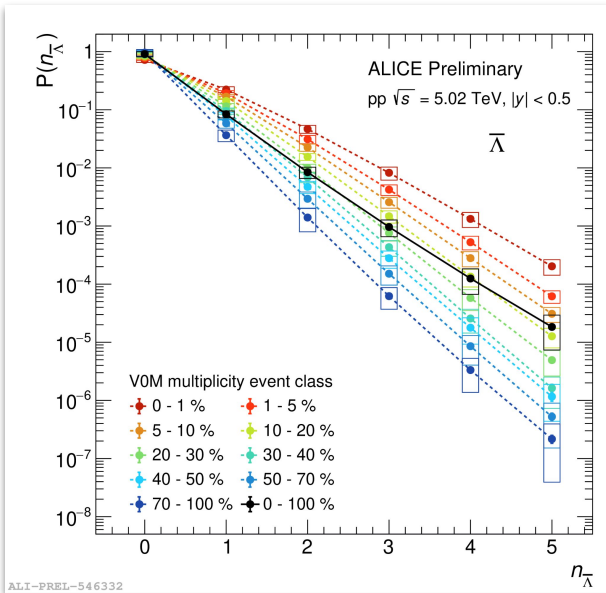
$\rightarrow \epsilon_i$  efficiencies

unfolding errors: covariance matrix

$$V(\hat{n}(C_k), \hat{n}(C_l)) = \sum_{i,j=1}^{n_E} \frac{\partial \hat{n}(C_k)}{\partial n(E_i)} V(n(E_i), n(E_j)) \frac{\partial \hat{n}(C_l)}{\partial n(E_j)}$$



Moving from  $K_s^0$  to  $\Omega$  particle the response matrices are increasingly "squeezed" toward a low number of reconstructed particles/event



Probability to produce  $n$  particle ( $n$  up to 5 for  $\bar{\Lambda}$ , 4 for  $\bar{\Xi}$ , 2 for  $\bar{\Omega}$ ) of a given species per event

Spanning across large ranges of strange/multiplicity variations, all the way to very “extreme” situations

Unique opportunity to test the connection between charged and strange particle multiplicity production

NOTE: in each V0M bin multiplicity can fluctuate and  $\langle dN_{ch}/d\eta \rangle$  can significantly change for events with small/large  $n_s$

