## **Modification of heavy quark hadronization in highmultiplicity collisions at LHCb**

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**Abstract.** The ratio of heavy flavor hadrons is very sensitive to the hadronization mechanism. This proceeding will present recent LHCb results on the crosssection ratios of  $D_s^+/D^+$ ,  $\Xi_c^+/ \Lambda_c^+$  and  $\Lambda_b^0/B^0$  in different collision systems. The<br>significantly enhanced production ratios  $D^+/D^+$  and  $\Lambda^0/B^0$  with the increase significantly enhanced production ratios  $D_s^+ / D^+$  and  $\Lambda_b^0 / B^0$  with the increase<br>of multiplicity may imply that hadronization mechanisms are modified in highof multiplicity may imply that hadronization mechanisms are modified in highmultiplicity events.

### **1 Introduction**

In the context of hadron colliders, heavy quarks primarily originate from hard parton-parton interactions in the initial stages of the collisions. These interactions are well described by perturbative QCD calculations, which rely on the factorization theorem. According to this theorem, the cross-sections of heavy-flavour hadrons depend on several key factors: the parton distribution functions within the incoming nucleons, the cross-section of hard partonparton scattering, and the fragmentation functions. For different types of heavy-flavored hadrons, the contributions from the first two items are similar, and only the hadronization process makes the difference. Traditionally, assuming that hadronization of heavy quarks is a universal process independent of colliding systems, these fragmentation functions are parameterized using data collected from *ee* or *ep* collisions. Some recent measurements from the LHCb experiment on the  $D_s^+/D^+$  and  $\Lambda_b^0/B^0$  ratios have revealed a significant enhance-<br>ment from low-multiplicity collisions to high-multiplicity collisions. These results indicate ment from low-multiplicity collisions to high-multiplicity collisions. These results indicate the existence of other hadronization mechanisms that are dependent on the collision size.

### 2 The cross-sections ratio  $D_s^+/D^+$  versus multiplicity in  $pPb$ <br>collisions **collisions**

In *p*Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV, LHCb conducted measurement of prompt  $D_s^+$  and  $D^+$ <br>production and studied relative production ratios [1]. The  $D^+ / D^+$  ratio observed in backward production, and studied relative production ratios [1]. The  $D_s^+/D^+$  ratio observed in backward<br>rapidity is slightly higher than that in forward rapidity and *nn* collisions [2]. Previous findings rapidity is slightly higher than that in forward rapidity and *pp* collisions [2]. Previous findings from LHCb [3] suggest that backward rapidity result in a higher yield of charged particles compared to forward rapidity within the symmetric kinematic interval. These results suggest a potential enhancement in the  $D_s^+/D^+$  ratio with increasing multiplicity.<br>Figure 1 illustrates the cross-section ratio  $D^+/D^+$  as a function of no

Figure 1 illustrates the cross-section ratio  $D_s^+ / D^+$  as a function of normalized event mul-<br>city in aPb collisions at  $\sqrt{s_{\text{DM}}}$  = 8.16 TeV [4]. The event multiplicity is denoted by tiplicity in *pPb* collisions at  $\sqrt{s_{NN}}$  = 8.16 TeV [4]. The event multiplicity is denoted by

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 $N_{\text{tracks}}^{\text{PV}}$ , representing the number of tracks used for reconstructing the primary vertex (PV). The charged particle multiplicity is normalized to the mean value observed in minimumbias events. All panels exhibit a significant dependence on multiplicity, with a pronounced enhancement observed, particularly in the backward rapidity. This enhancement can be attributed to the combined effects of the coalescence mechanism and strangeness enhancement. However, this enhancement shows no tendency to weaken in the high transverse momentum  $(6 < p<sub>T</sub> < 12 \text{ GeV}/c)$  ranges.



**Figure 1.** Cross-section ratio  $\sigma_{D_s^+}/\sigma_{D^+}$  versus the normalised event multiplicity in different  $p_T$  ranges for (top) forward and (bottom) backward rapidity [4] for (top) forward and (bottom) backward rapidity [4].

# **3 Measurement of**  $\Xi_c^+$  production in  $p{\rm Pb}$  collisions at  $\sqrt{s_{\rm NN}} = 8.16$ <br>TeV **TeV**

Figure 2 shows production ratio of strange baryons to non-strange baryons,  $\Xi_c^+/\Lambda_c^+$ , and the production ratio of strange baryons to non-strange mesons  $\Xi^+/\Lambda_c^0$  in *n*Ph collisions at the production ratio of strange baryons to non-strange mesons,  $\Xi_c^+/D^0$ , in *p*Pb collisions at  $\Lambda_{\text{Cov}} = 8.16 \text{ TeV}$  [5]. Both the  $\Xi^+/A^+$  and  $\Xi^+/D^0$  ratios show no significant *n*<sub>n</sub> dependence  $\sqrt{s_{NN}}$  = 8.16 TeV [5]. Both the  $\Xi_c^+/Λ_c^+$  and  $\Xi_c^+/D^0$  ratios show no significant *p*<sub>T</sub> dependence.<br>The  $\Xi_c^+/Λ_c^+$  and  $\Xi_c^+/D^0$  ratios are consistent within the error for forward and backward rapid-The  $\Xi_c^+/\Lambda_c^+$  and  $\Xi_c^+/D^0$  ratios are consistent within the error for forward and backward rapid-<br>ity but the  $\Xi^+/\Lambda^+$  ratio in the backward rapidity is slightly higher than that in the forward ity, but the  $\Xi_c^+/\Lambda_c^+$  ratio in the backward rapidity is slightly higher than that in the forward rapidity rapidity.

The measurements are compared with the EPPS16 calculation [6], but both  $\Xi_c^+/\Lambda_c^+$  and  $D^0$  are overestimated. The calculations from Pythia 8.3 with color reconnection [7] and  $\Xi_c^+/D^0$  are overestimated. The calculations from Pythia 8.3 with color reconnection [7] and<br>EPOS4HO [8, 9] are also shown in Figure 2, both of which are based on results from an EPOS4HQ [8, 9] are also shown in Figure 2, both of which are based on results from *pp* collisions.

#### **4 Enhanced production of**  $\Lambda^0_h$ Enhanced production of  $\Lambda_b^0$  in high-multiplicity  $pp$  collisions at  $\sqrt{s}$  = 13 **TeV**

Figure 3 presents cross-sections ratio  $\Lambda_b^0/B^0$  as a function of  $p_T$  in different multiplicity bins<br>in an collisions at  $\sqrt{s} = 13$  TeV [10]. The event multiplicity is characterized by N<sup>VELO</sup> Figure 5 presents cross-sections ratio  $A_b/B$  as a function of  $p_T$  in unteresting in the proposition at  $\sqrt{s}$  = 13 TeV [10]. The event multiplicity is characterized by N<sup>VELO</sup> and  $N_{\text{tracks}}^{\text{back}}$ .  $N_{\text{tracks}}^{\text{VELO}}$  denotes the total number of charged tracks reconstructed in the VELO detector, while  $\overline{N}_{\text{tracks}}^{\text{back}}$  represents the subset of  $N_{\text{tracks}}^{\text{VELO}}$  pointing away from the LHCb detector. Both plots show a significant dependence on multiplicity, with clear enhancements observed,



**Figure 2.** Production ratios, (left)  $\Xi_c^+/\Lambda_c^+$  and (right)  $\Xi_c^+/D^0$ , as a function of  $p_T$  in the *pPb* (red triangles) and *Phn* (blue triangles) data samples [5] triangles) and Pb*p* (blue triangles) data samples [5].

particularly evident when utilizing VELO tracks. At low  $p_T$ , the  $\Lambda_b^0/B^0$  ratio is significantly higher than the value observed in  $e^+e^-$  collisions. However, as the *p<sub>p</sub>* increases the  $\Lambda^0/B^0$ higher than the value observed in  $e^+e^-$  collisions. However, as the  $p_T$  increases, the  $\Lambda_b^0/B^0$ <br>ratio tends to align with the results obtained from  $e^+e^-$  collisions ratio tends to align with the results obtained from  $e^+e^-$  collisions.



**Figure 3.** Cross-section ratio  $\sigma_{\Lambda_b^0}/\sigma_{B^0}$  as a function of  $p_T$  in different multiplicity bins.

Figure 4 shows cross-sections ratio  $\Lambda_b^0/B^0$  as a function of  $p_T$ . The data are compared<br>revious *nn* measurement [11] and *n*Ph measurement [12] and generally consistent with to previous *pp* measurement [11] and *p*Pb measurement [12], and generally consistent with them within uncertainties. Additionally, two curves from *b* quarks statistical hadronization model [13] are also shown in Figure 4. The green solid curve considers feeddown contributions from *b* baryons which have been collected by Particle Data Group [14]. The black dashed curve takes into account feeddown contributions from an expanded set of *b* baryons predicted by the Relativistic Quark Model [15].

### **5 Summary**

In recent LHCb studies, cross-section ratios,  $D_s^+/D^+$  and  $\Lambda_b^0/B^0$ , were measured in both low<br>and bigh multiplicity collisions.  $\Xi^+/\Lambda^+$  was measured in both *nPh* and Phn collisions. The and high multiplicity collisions.  $\Xi_c^+ / \Lambda_c^+$  was measured in both *pPb* and *Pbp* collisions. The



**Figure 4.** Cross-section ratio  $\sigma_{\Lambda_b^0}/\sigma_{B^0}$  as a function of  $p_T$ .

 $\frac{\Sigma_c^+}{\Lambda_c^+}$  shows little variation between forward and backward rapidity within uncertainties.<br>The  $D^+ / D^+$  and  $\Lambda^0 / R^0$  show a significant enhancement in high-multiplicity collisions com-The  $D_s^+/D^+$  and  $\Lambda_b^0/B^0$  show a significant enhancement in high-multiplicity collisions com-<br>pared to low-multiplicity collisions. The  $\Lambda^0/B^0$  ratio decreases with  $p_T$  and converges to pared to low-multiplicity collisions. The  $\Lambda_b^0/B^0$  ratio decreases with  $p_T$  and converges to the  $e^+e^-$  result at high  $p_T$ . These suggest the potential presence of other hadronization the  $e^+e^-$  result at high  $p_T$ . These suggest the potential presence of other hadronization mechanisms in high-multiplicity collisions. Furthermore, it is also possible that in highmultiplicity collisions, the more contribution from excited state feeddown could lead to the observed enhancement in cross-section ratios. It is worth noting that this enhanced feeddown of SHM+RQM relative to SHM+PDG does not weaken as  $p<sub>T</sub>$  increases, as predicted in Figure 4.

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