

Hadron physics with polarized photons at LEPS/LEPS2

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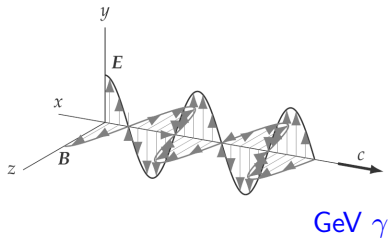
New Vistas in Photon Physics
in Heavy-Ion Collisions

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Institute of Nuclear Physics Polish Academy of Sciences & AGH University of Science and Technology

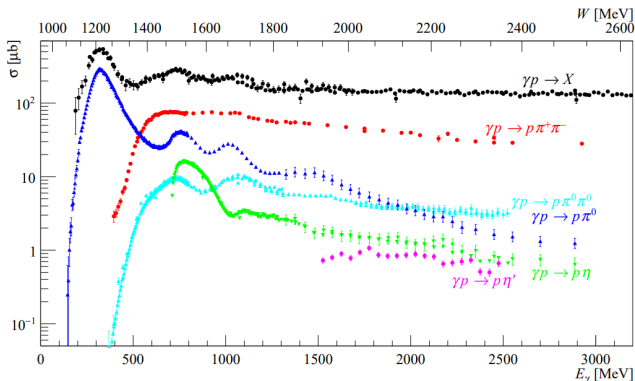
Outline

1. Compton-backscattered photon beam
2. Highlights from the **LEPS experiments**: photoproduction of (π, η, η', K^+) mesons and (ω, ϕ, K^*) mesons.
3. Current status of the **LEPS2 photoproduction experiments**: photoproduction of η and ω mesons with BGOegg detector.
4. Search for Θ^+ , P_s and K^-pp with the LEPS2 detector.



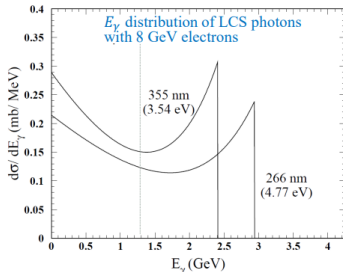
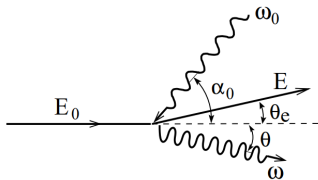
Photon Beam

- Because of a small attenuation of the photon inside the nucleon the photon is an ideal probe to explore the interior of the hadrons.
- The largest γp total cross section is $\sigma \approx 520\mu\text{b}$ at $E_\gamma \approx 0.32$ GeV, which corresponds to a mean free path for a photon in nuclear matter much larger than the nuclear radius ($\lambda \approx 100$ fm).
- At low masses, the $\Delta(1232)$ resonance, the first excited state, can be clearly observed as a peak in the cross section for $p\pi^0$. All other structures are not composites of a single resonant state, but arise due to several overlapping states. ^a



^aA. Thiel, F. Afzal, and Y. Wunderlich, Prog. Part. Nucl. Phys. 125, 103949 (2022).

Backward Compton Scattering



- Kinematics of the backward Compton scattering $e(p_0) + \gamma_0(k_0) \rightarrow e(p) + \gamma(k)$ is characterized by two dimensionless variables x and y :^a

$$x = \frac{2p_0 k_0}{m^2 c^2} \approx \frac{4E_0 \omega_0}{m^2 c^4} \cos^2 \frac{\alpha_0}{2}, \quad y = \frac{k k_0}{p_0} \approx \frac{\omega}{E_0}$$

- The maximum energy of the scattered photon ω_m and the maximum value of the parameter y are:

$$\omega \leq \omega_m = \frac{x}{x+1} E_0, \quad y \leq y_m = \frac{x}{x+1} = \frac{\omega_m}{E_0}$$

- The energy spectrum of Compton backscattered photons is given by:

$$\frac{d\sigma}{dy} \approx \frac{2\sigma_0}{x} \left[\frac{1}{1-y} + 1 - y - 4r(1-r) \right], \quad \text{where } r = \frac{y}{x(1-y)} \leq 1.$$

^aV.G. Serbo, Acta Phys. Polon. B 37, 1333 (2006)

Backward Compton Scattering

- The energy range of Compton backscattered photons spans $E_\gamma = 1.4\text{--}2.9$ GeV at SPring-8, which fits the extensive study of baryon resonances near 2 GeV. Note that photons from UPC are produced at much higher energies ($E_\gamma = \gamma\hbar c/R_A \approx 3$ GeV @ Au+Au at $\sqrt{s_{NN}} = 200$ GeV and $E_\gamma = \gamma\hbar c/R_A \approx 75$ GeV @ Pb+Pb at $\sqrt{s_{NN}} = 5.02$ TeV).^a
- The energy of a Compton backscattered photon depends on its emission angle θ :

$$\omega = \frac{\omega_m}{1 + (\theta/\theta_0)^2}, \quad \theta_0 = \frac{mc^2}{E_0} \sqrt{x+1}$$

- If the laser light is polarized, then high-energy photons are polarized in the same direction, which is characterized by:

$$x = \frac{4E_0\omega_0}{m^2c^4}, \quad a = \frac{1}{1+x}, \quad \rho = \frac{\omega}{\omega_m}, \quad k = \frac{\rho^2(1-a)^2}{1-\rho(1-a)}, \quad \cos\theta_0 = \frac{1-\rho(1+a)}{1-\rho(1-a)}$$

At the Compton edge ($\omega = \omega_m$), $\rho = 1$, $\cos\theta_0 = -1$, and $k = (1-a)^2/a$.

- For linear polarization,

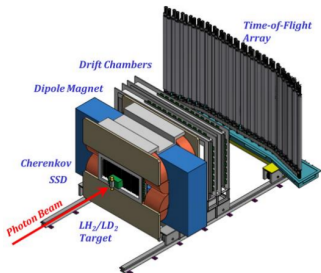
$$\frac{P_\gamma}{P_{\text{laser}}} = \frac{(1 - \cos\theta_0)^2}{2(k + 1 + \cos^2\theta_0)} \underbrace{\rightarrow}_{\text{Compton-edge}} \frac{2}{2+k} < 1$$

- For circular polarization,

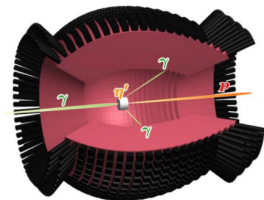
$$\frac{P_\gamma}{P_{\text{laser}}} = \frac{(k+2)\cos\theta_0}{(k+1+\cos^2\theta_0)} \underbrace{\rightarrow}_{\text{Compton-edge}} -1$$

^aMariusz Przybycien, "Overview of UPC Snowmass Lol physics cases", RHIC Science Programs, 2021

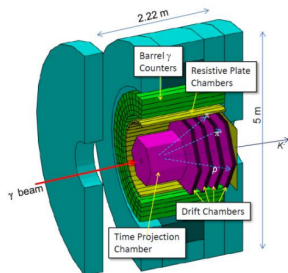
LEPS and LEPS2 at SPring-8



LEPS spectrometer



BGOegg calorimeter



Solenoid spectrometer

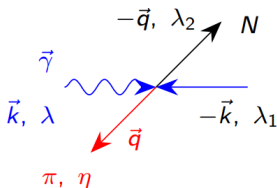
	LEPS (2000 – 2020)	LEPS2 (2013 –)
Tagged γ Energy	$1.5 < E_\gamma < 2.4$ GeV (UV laser) $1.5 < E_\gamma < 2.9$ GeV (DUV laser)	$1.3 < E_\gamma < 2.4$ GeV (UV laser) $1.3 < E_\gamma < 2.9$ GeV (DUV laser)
Photon Beam Intensity	Two laser injection 2×10^6 cps (UV laser) 2×10^5 cps (UV laser)	Four laser injection $< 10^7$ cps (UV laser) $< 10^6$ cps (DUV laser)
Detector	Forward Dipole Spectrometer	BGOegg EM Calorimeter Solenoid Spectrometer

Photoproduction of Pseudoscalar Mesons

- Photoproduction of pseudoscalar mesons (π, η, η', K) is the simplest process to analyze.

$$\gamma + N \rightarrow \begin{bmatrix} \pi \\ \eta \\ \eta' \end{bmatrix} + N$$

$$\lambda = \pm 1 \quad m_s = \pm \frac{1}{2} \rightarrow 0 \quad m_s = \pm \frac{1}{2}$$



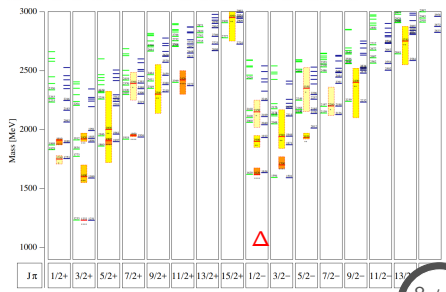
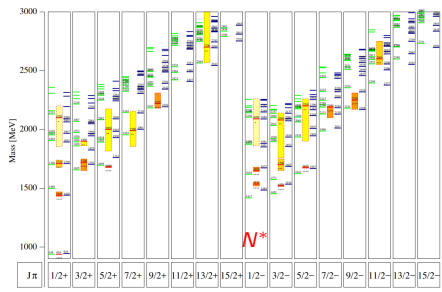
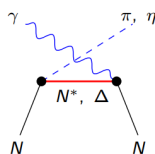
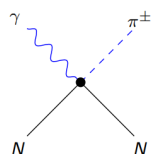
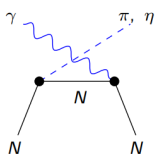
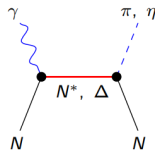
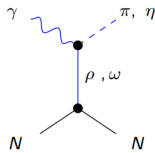
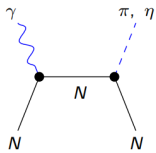
- Four spin degrees of freedom with four complex amplitudes, usually given as CGLN, helicity, or transversity amplitudes:

$$\tilde{\mathcal{F}} = \chi_{m_{s_f}}^\dagger F_{\text{CGLN}} \chi_{m_{s_i}}$$

$$F_{\text{CGLN}} = i\vec{\sigma} \cdot \hat{\epsilon} F_1 + \vec{\sigma} \cdot \hat{q} \vec{\sigma} \cdot \hat{k} \times \hat{\epsilon} F_2 + i\vec{\sigma} \cdot \hat{k} \hat{q} \cdot \hat{\epsilon} F_3 + i\vec{\sigma} \cdot \hat{q} \hat{q} \cdot \hat{\epsilon} F_4$$

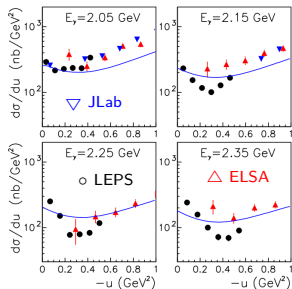
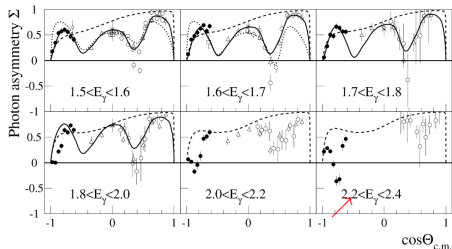
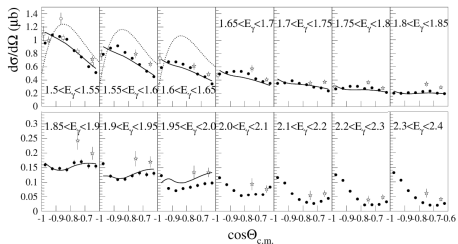
- From four complex amplitudes, one can construct 16 polarization observables, which can be measured with linearly or circularly polarized photon beams, polarized targets, and recoil polarization detection: single polarization observables ($\sigma_0 = d\sigma/d\Omega, \Sigma, T, P$), beam-target observables (E, F, G, H), beam-recoil observables ($C_{x'}, C_{z'}, O_{x'}, O_{z'}$), and target-recoil observables ($T_{x'}, T_{z'}, L_{x'}, L_{z'}$)

Light Baryon Spectroscopy



Backward π^0 Photoproduction

- Differential cross sections and photon beam asymmetries for π^0 photoproduction have been measured at $E_\gamma = 1.5\text{--}2.4$ GeV and at the π^0 scattering angles, $-1 < \cos\theta_{c.m.} < -0.6$.



- Above 2.0 GeV, a dip structure around $\cos\theta_{c.m.} = -0.8$ requires new high-mass resonances combined with the u -channel diagrams. ^a

$$P_\gamma \Sigma \cos 2\Phi = \frac{nN_V - N_h}{nN_V + N_h}$$

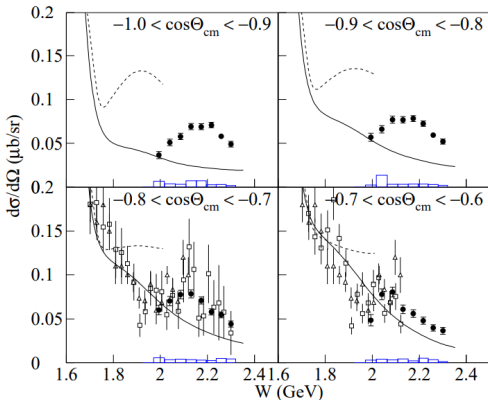
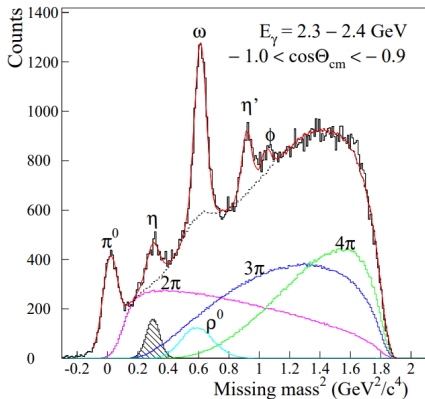
- Regge contributions should come to play and interferences will occur. ^b

^aM. Sumihama *et al.* (LEPS Collab), Phys. Lett. B 657, 32 (2007).

^bA. Sibirtsev *et al.*, Eur. Phys. J A 40, 65 (2009).

Backward η Photoproduction

- Backward η photoproduction off protons has been measured at $E_\gamma = 1.6\text{--}2.4$ GeV at the SPring-8/LEPS facility, by detecting protons scattered at forward angles.

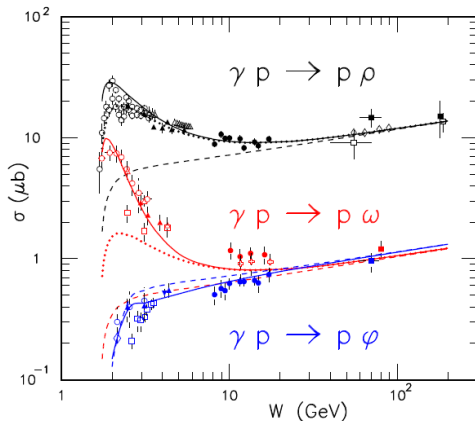
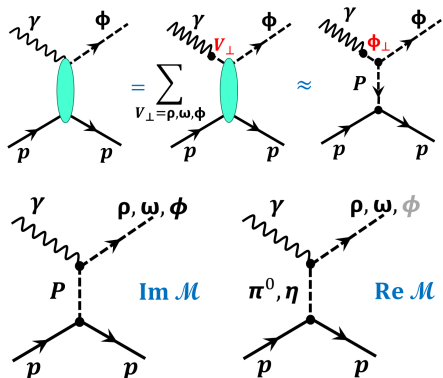


- This work provides unambiguous evidence for a bump structure above $W = 2.0$ GeV. No such structure is seen in η' , ω and π^0 photoproductions, which supports that this unique structure in η photoproduction is due to a baryon resonance with a large $s\bar{s}$ component. ^a

^aM. Sumihama *et al.* (LEPS Collab), Phys. Rev. C 80, 052201 (2009).

ϕ Photoproduction near Threshold

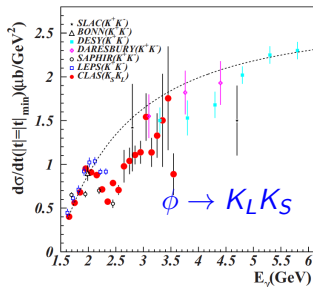
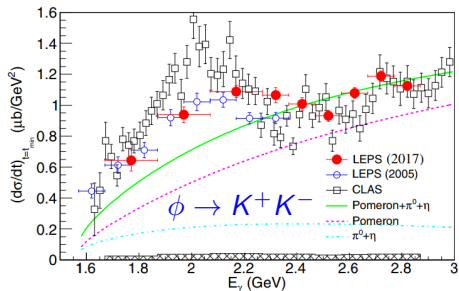
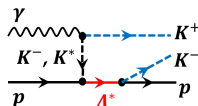
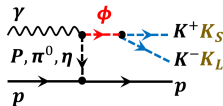
- In the vector meson dominance (VMD) model for photoproduction, a real photon can fluctuate into a virtual vector meson $V = \{\rho, \omega, \phi\}$, which subsequently scatters off the target proton.
- The ϕ meson production has the unique feature within gluon dynamics of being a result of **OZI suppression** due to the dominant $\bar{s}s$ structure.



ϕ Photoproduction near Threshold

- For Pomeron exchanges, the photoproduction cross section can be written in the form

$$\frac{d\sigma}{dt}(\gamma p \rightarrow \phi p) = \left(\frac{d\sigma}{dt} \right)_{t=t_{\min}} \exp(-b_\phi |t - t_{\min}|)$$



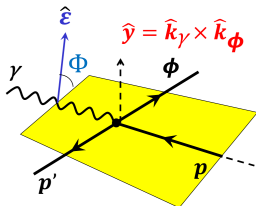
- Differential cross sections for ϕ photoproduction in the reaction $\gamma p \rightarrow p\phi$ followed by $\phi \rightarrow K^+ K^-^a$ and $K_S K_L^b$ show some enhancement in the photon beam energy range near 2.1 GeV.

^aK. Mizutani *et al.* (LEPS Collab), PRC 96, 062201(R)(2017); B. Dey *et al.* (CLAS Collab), PRC 89, 055208(2014)

^bH. Seraydaryan *et al.* (CLAS Collab), PRC 89, 055206(2014).

Decay Angular Distribution and SDMEs

$$\gamma p \rightarrow \phi p$$



$$W(\cos \theta) = \frac{3}{2} \left(\frac{1}{2} (1 - \rho_{00}^0) \sin^2 \theta + \rho_{00}^0 \cos^2 \theta \right)$$

$$W(\phi) = \frac{1}{2\pi} (1 - 2\rho_{1-1}^0 \cos 2\phi)$$

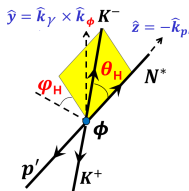
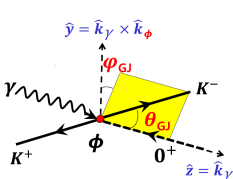
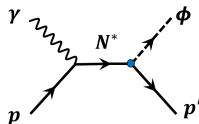
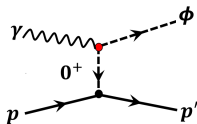
$$W(\phi - \Phi) = \frac{1}{2\pi} (1 + P_\gamma (\rho_{1-1}^1 - \text{Im} \rho_{1-1}^2) \cos 2(\phi - \Phi))$$

$$W(\phi + \Phi) = \frac{1}{2\pi} (1 + P_\gamma (\rho_{1-1}^1 + \text{Im} \rho_{1-1}^2) \cos 2(\phi + \Phi))$$

$$W(\Phi) = 1 - P_\gamma (2\rho_{11}^1 + \rho_{00}^1) \cos 2\Phi$$

$$P_\sigma = \frac{\sigma^N - \sigma^U}{\sigma^N + \sigma^U} = 2\rho_{1-1}^1 - \rho_{00}^1$$

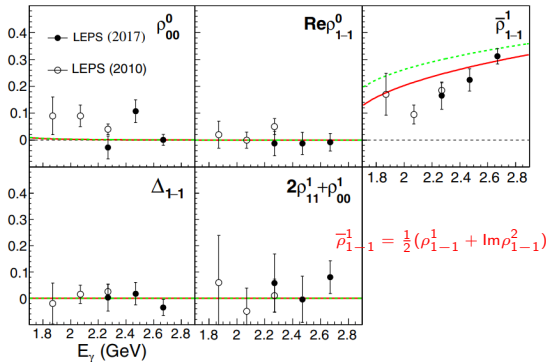
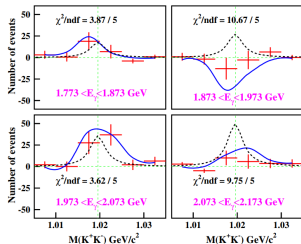
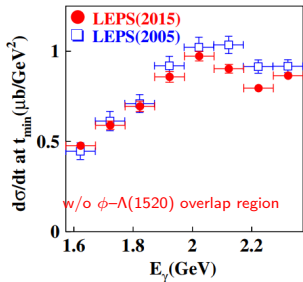
$$\Sigma = \frac{\sigma_{\parallel} - \sigma_{\perp}}{\sigma_{\parallel} + \sigma_{\perp}} = \frac{\rho_{11}^1 + \rho_{1-1}^1}{\rho_{11}^0 + \rho_{1-1}^0}$$



For the helicity-conserving exchanges,

$$\rho_{1-1}^1 = -\text{Im} \rho_{1-1}^2 = \begin{bmatrix} +0.5(N) \\ -0.5(U) \end{bmatrix}$$

ϕ Photoproduction near Threshold

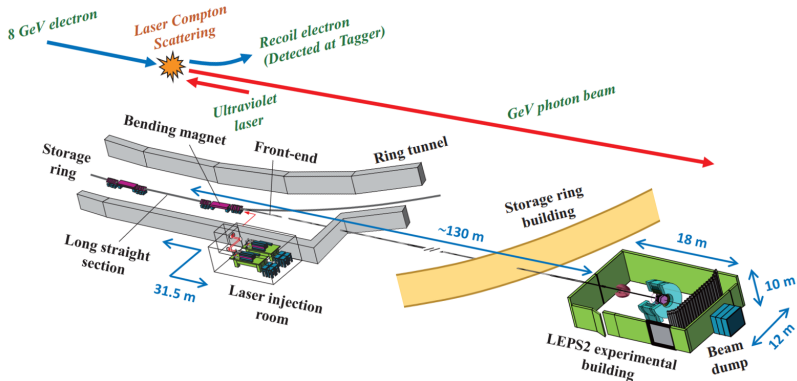
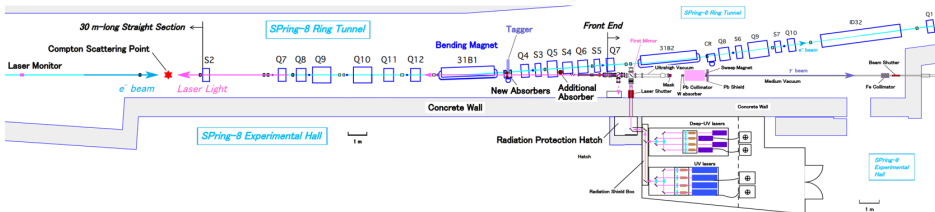


- $\bar{\rho}_{1-1}^0$ deviates largely from the model prediction in $1.97 < E_\gamma < 2.17 \text{ GeV}$, which reflects the contributions of N^* resonances. ^a
- The $\sqrt{s} = 2.1 \text{ GeV}$ bump structure was reconfirmed without the ϕ - $\Lambda(1520)$ interference region. ^b

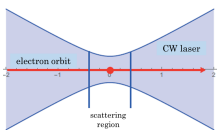
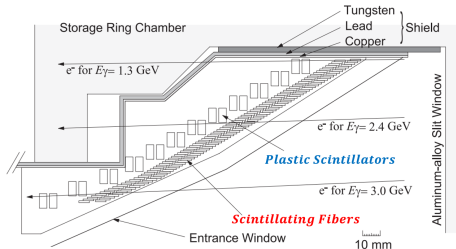
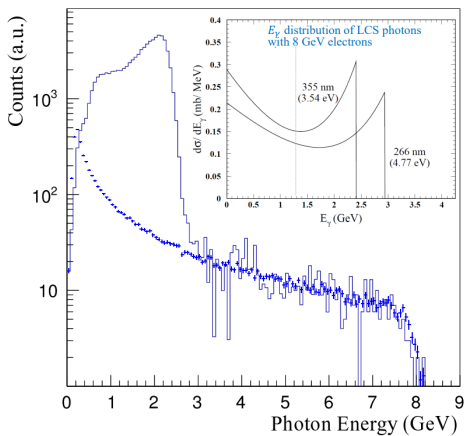
^a K. Mizutani *et al.*, PRC 96, 062201(R)(2017).

^b S.Y. Ryu *et al.* (LEPS Collab), PRL 116, 232001(2016).

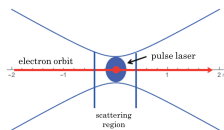
GeV Photon Beam at LEPS2/SPring-8



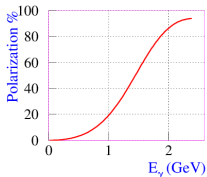
Backward Compton Scattered Photons



CW laser



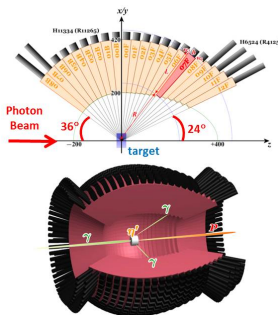
Synchronized pulse laser



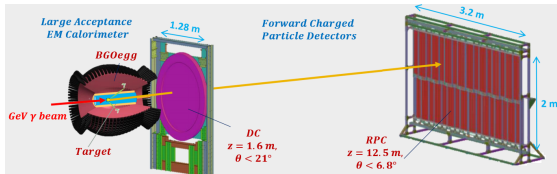
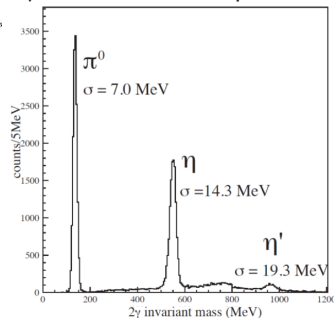
- The tagged photon energy resolution is **12.3 MeV** and the maximum intensity is $2.3 \times 10^6/s$ in the energy from 1.3 to 2.4 GeV. ^a

^a N. Muramatsu *et al.*, NIMA 1033, 166677(2022).

LEPS2 BGOegg Detector



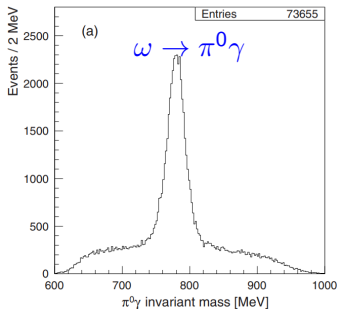
2γ invariant mass spectrum



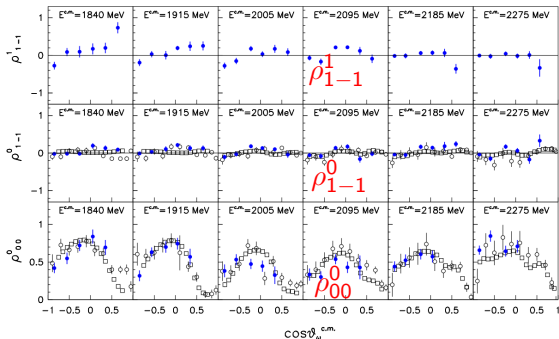
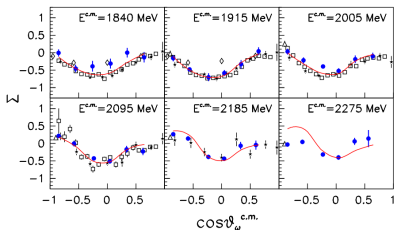
- BGOegg consisting of 1320 BGO crystals ($20 X_0$) views a target in $\theta = 24^\circ - 144^\circ$ with the world's best energy resolution (1.3% at 1 GeV).^a

^aT. Ishikawa *et al.*, NIMA 837, 109(2016).

ω Photoproduction with BGOegg Detector



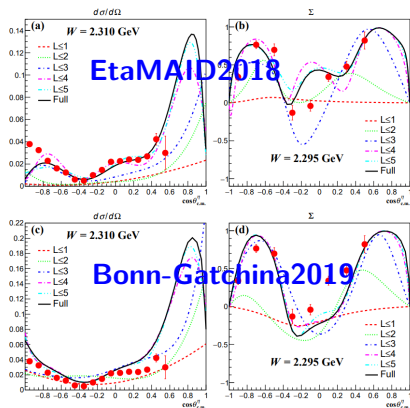
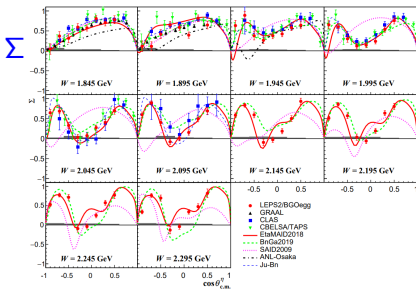
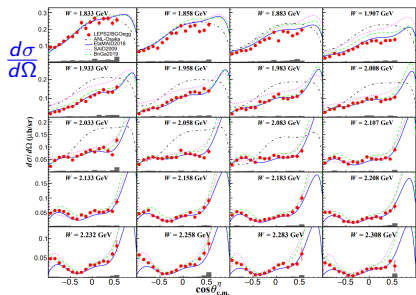
$\gamma p \rightarrow \omega p$



- ρ_{1-1}^1 values decrease at forward angles as a function of total energy.
- The ρ_{00}^0 shows large values at intermediate angles, while an increase toward higher energies is seen at backward angles.
- The observed behaviors of the helicity-flip amplitudes surely reflects the contributions of N^* resonances. ^a

^aN. Muramatsu et al. (LEPS2 Collab), PRC 102, 025201(2020)

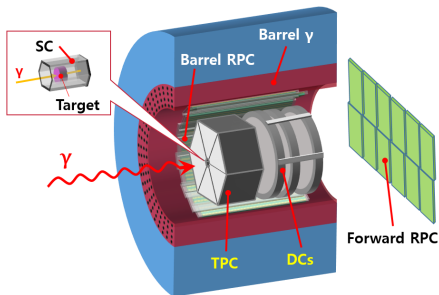
η Photoproduction with BGOegg Detector



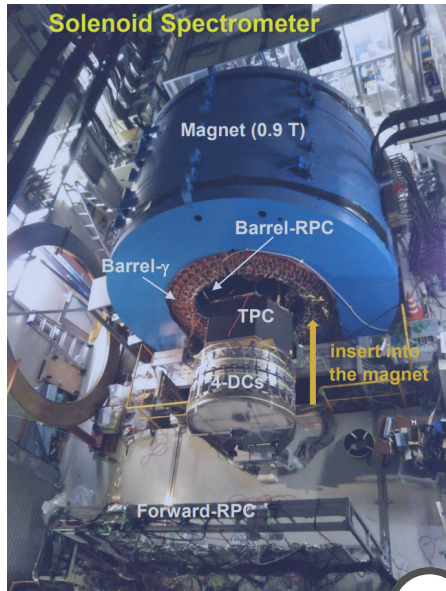
- A bump structure at $W = 2.0\text{--}2.3$ GeV is confirmed at extremely backward η polar angles.
- This bump structure is likely associated with high-spin resonances that couple with $s\bar{s}$ quarks. ^a

^aT. Hashimoto *et al.* (LEPS2/BGOegg Collab), PRC 106, 035201 (2022).

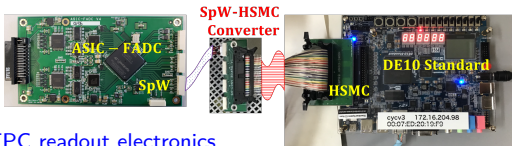
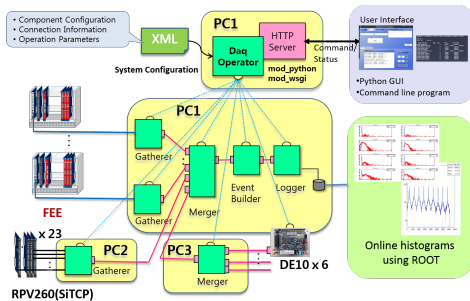
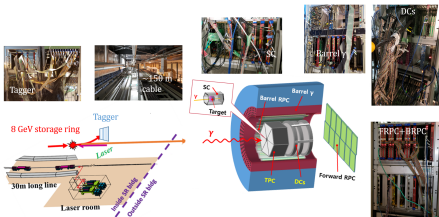
LEPS2 Solenoid Detector



- The LEPS2 solenoid detector comprises TPC, SC, DCs, forward and barrel RPCs, Barrel Pb/Scint calorimeter ($14.3 X_0$), and neutron counters as well as a photon tagger placed approximately 150 m upstream from the detector.



LEPS2 DAQ System

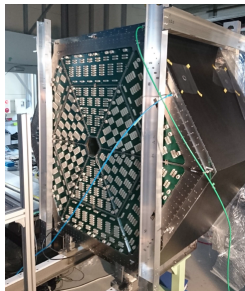
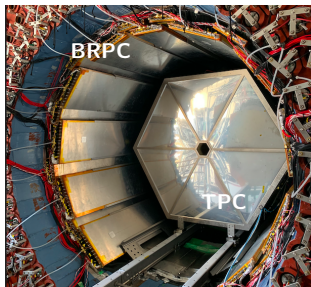


TPC readout electronics

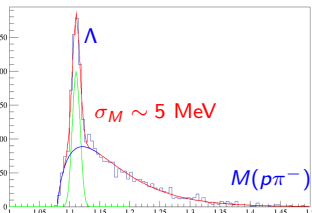
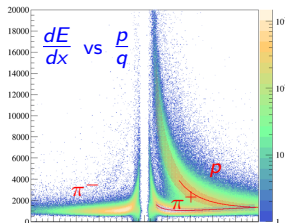
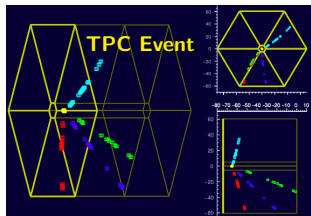
- We have developed a new, network-distributed data acquisition system based on DAQ-middleware framework.^a

^a S. Y. Ryu for LEPS2 Collab., AIP Conf. Proc. 2249, 030024 (2020).

Time Projection Chamber



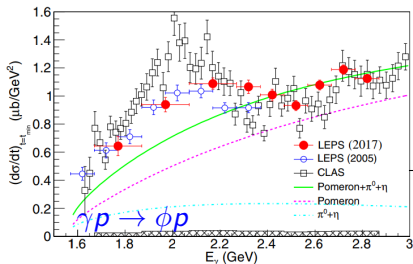
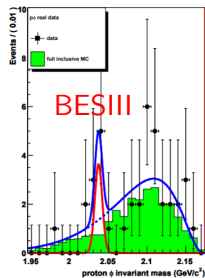
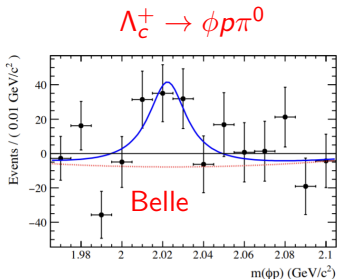
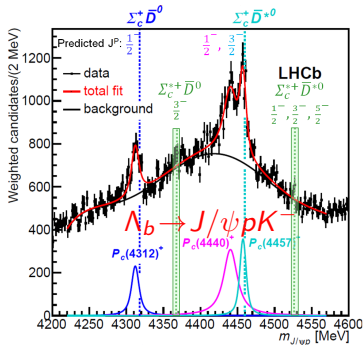
- Hexagonal prism-shaped drift volume : 112 cm (corner-to-corner) \times 71 cm (length).
- 24 straight pad layers grouped into 6 sections forming a hexagonal web-like structure.
- 10,830 pads with single pad size of $4.6 \times 10 \text{ mm}^2$.



- Λ decays are successfully reconstructed with p and π^- tracks in the first dataset.^a

^a R. Kobayakawa for LEPS2 Collab. JPS fall meeting, 2021.

Strange Partner P_s of P_c Pentaquark States

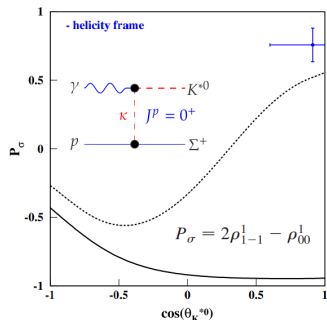
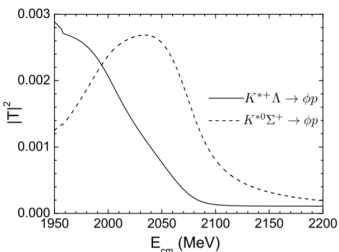


- A strange partner $P_s(uuds\bar{s})$ was searched for in the Cabibbo-suppressed $\Lambda_c^+ \rightarrow \phi p \pi^0$ decay at Belle^a and BESIII.^b
 $\mathcal{B}(\Lambda_c^+ \rightarrow P_s \pi^0) \cdot \mathcal{B}(P_s \rightarrow \phi p) < 8.3 \times 10^{-5}$ at 90% CL)

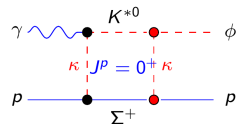
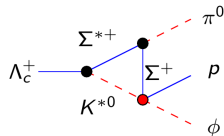
^aB. Pal *et al.* (Belle Collab), PRD 96, 051102 (2017).

^bG. Mezzadri, Ph.D thesis, Ferrara U (2018).

Strange Partner of P_c Pentaquark States



- The bump structure observed in ϕ photoproduction could be regarded as a ΣK^* molecular state ($J = 3/2^-$). Triangular singularity could also explain the bump structure with $\Sigma K^* \rightarrow \phi p$.^a



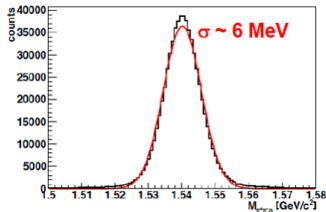
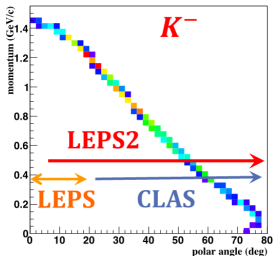
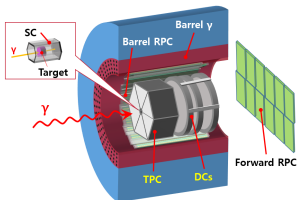
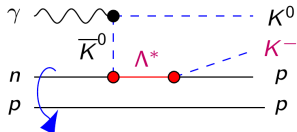
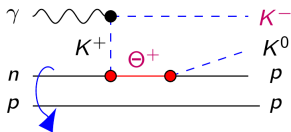
- The measured parity spin asymmetry shows that natural-parity exchange is dominant in $\gamma p \rightarrow K^{*0}\Sigma^+$ reaction.^b
- The nature of the bump structure should be further investigated using circularly polarized photon beam and a polarized (HD) target.

^aJ.-J. Xie and F.-K. Guo, PLB 774, 108 (2017); J. He, PRD 95, 074031(2017).

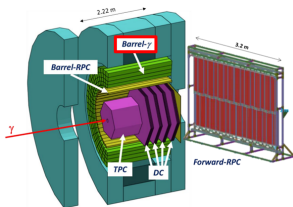
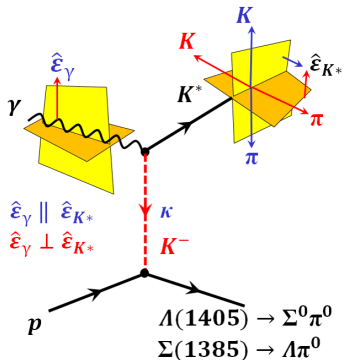
^bS.H. Hwang *et al.* (LEPS Collab.), PRL 108, 092001(2012).

Θ^+ Search at LEPS2

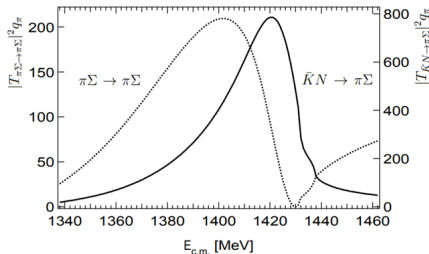
- We search for the Θ^+ ($S = +1$, $uudd\bar{s}$) via $\gamma d \rightarrow K^- K^0 pp$ reaction, followed by $\Theta^+ \rightarrow K^0 p$; $K^0(K_S) \rightarrow \pi^+ \pi^-$.
- All final state particles can be reconstructed using the large-acceptance LEPS2 detector, which facilitate a wider angular coverage for K^- detection.



Photoproduction of $\Lambda(1405)$ with K^{*+}

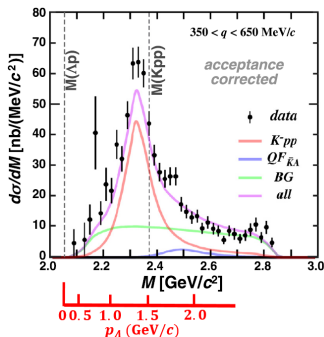
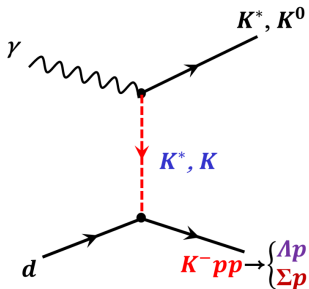


- $K^{*+} (\rightarrow \pi^+ K_S^0)$ decay plane \perp (\parallel) the photon beam polarization ($\hat{\epsilon}$) for **unnatural-parity exchange** (**natural-parity exchange**)
- Unnatural-parity K^- exchange selects $\Lambda(1405)$ strongly coupled to a $K^- p$ pole.^a
- $I = 0$ channel $\Lambda(1405) \rightarrow \Sigma^0 \pi^0$ can be reconstructed using Barrel- γ detector.



^a Jido *et al.*, NPA 725, 181 (2003); T. Hyodo *et al.*, PLB 593, 75 (2004).

Search for K^-pp nuclei at LEPS2



- J-PARC E15 reported a peak in Λp mass of ${}^3\text{He}(K^-, \Lambda p)n$ well below K^-pp mass threshold. ^a
- K^-pp search in $\gamma d \rightarrow K^+\pi^-X$, K_S^0X , and $K^{*0}X$ reactions at LEPS2. facilitating complete kinematics with a detection of K/K^* and decay products from K^-pp .
- The existence of K^-pp nuclei can also be reconfirmed with Λp scattering in the range of 0.5–2.0 GeV/c.

$$\gamma p \rightarrow K^+\Lambda; \quad \Lambda p \rightarrow (K^-pp) \rightarrow \Lambda p$$

^aJ-PARC E15 Collab., Phys. Lett. B 789, 620 (2019).

Summary

- A **LEPS2** facility with BGOegg and solenoid detectors has started its full operation for studying hadron spectroscopy from photoproduction using **high-intensity Compton backscattered photons at SPring-8**.
- Highly polarized photons are very powerful to unveil the nature of hadrons from photoproduction; N^* , Δ , Λ^* , Σ^* , Θ^+ , P_s , K^-pp and so on.
- The LEPS2 solenoid detector has completed **its first phase of physics runs**. Analysis effort of the first dataset is now underway. Please stay tuned.