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Prompt photon physics with the ALICE FoCal

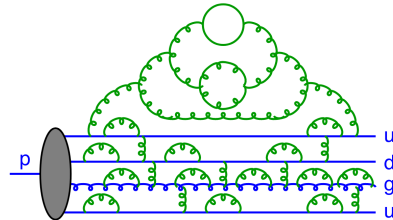
Florian Jonas

Forward QCD workshop, 24.05.2022

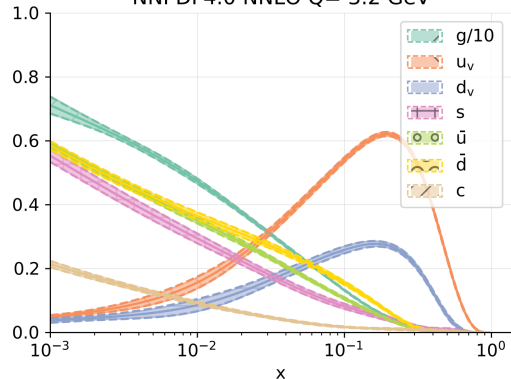


PDFs: What does it look like inside a proton?

- protons are a complex structure of quarks bound to each other via gluons and a constant creation/annihilation of virtual quarks and gluon
- can not be described from first principles using perturbative QCD (pQCD) → use universal parametrizations determined by experimental data
- Parton Distribution Function (PDF) gives probability density to find parton i carrying fraction x of protons momentum at scale Q
- only 50% is carried by quarks; the rest by gluons!
- PDFs evolve with scale Q → described by QCD evolution equations (DGLAP etc.)
- low PDF uncertainties → low prediction uncertainties!

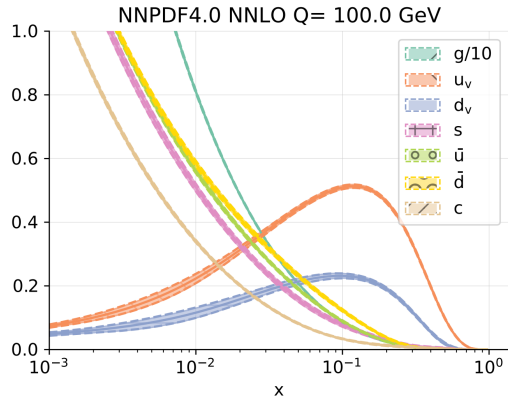
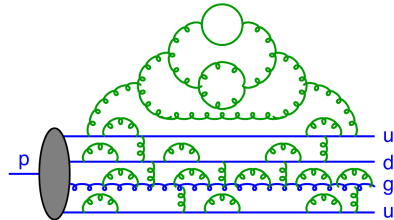


NNPDF4.0 NNLO $Q = 3.2$ GeV

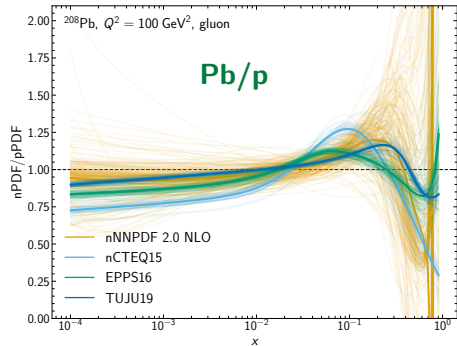
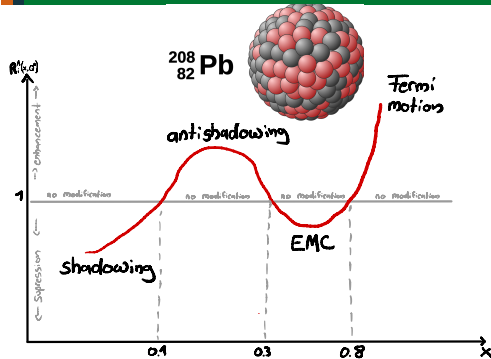


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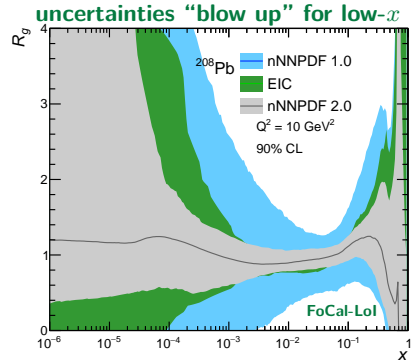
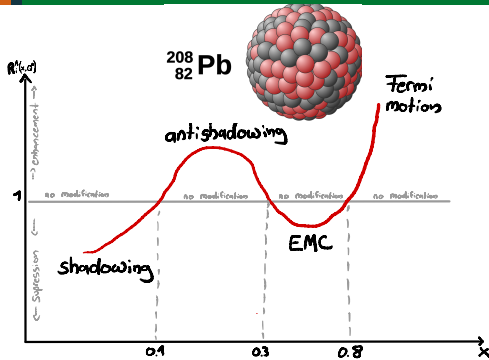


The initial state: Nuclear modification



- the initial state of **proton in nucleus** (e.g. Pb) is **modified** with respect to the free proton case!
 - shadowing**: outer nucleons "shield" inner nucleons
 - EMC effect**: explanations vary, some change in scale involved (increase in quark confinement size, nucleon size, nuclear treatment without quarks ...)
 - fermi motion**: nucleons themselves move inside nucleus
 - energy loss of incoming partons **before** hard scattering?
 - better understanding of nuclear effects in the initial state is crucial for disentangling from final-state effects (e.g. QGP)
- ⇒ still a lot to understand and learn!

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Constraining PDFs and nPDFs → improve predictions

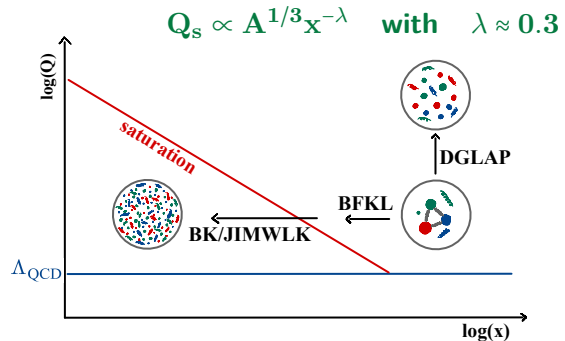
- gluon PDFs are one of the least constrained PDFs
- gluons drive significant fraction of scatterings at the LHC → precise knowledge improves predictions (prominent example $gg \rightarrow H$) (and enter a variety of measurements)

Quantifying nuclear effects in the initial state (IS)

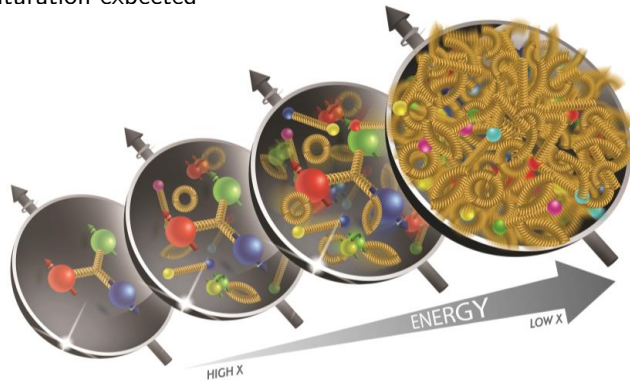
- how strong is gluon shadowing at low- x ?

Non-linear QCD evolution

- interesting nonlinear physics at low- x and Q where gluon density is very high:
 - gluon splitting and gluon fusion start to balance each other, leading to gluon saturation
 - most prominent model to describe regime of gluon saturation is the Color Glass Condensate (CGC) model
- nuclear environment → higher saturation scale Q_s

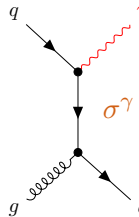


- we would like to know about gluons inside the proton
- very soft gluons (low- x) are especially interesting because ...
 - PDFs are not well constrained (large uncertainties)
 - shadowing of gluons expected
 - gluon saturation expected

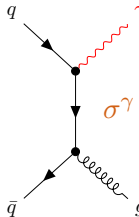


How can we access this regime?

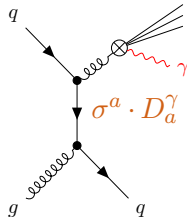
Compton



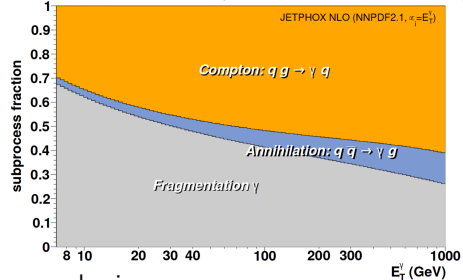
annihilation



fragmentation



10.1016/j.nuclphysb.2012.03.003

LHC, pp $\rightarrow \gamma + X$ @ $\sqrt{s}=14$ TeV, $y=0$ 

- prompt photons produced in hard scattering via two mechanisms:

① **direct production** (compton scattering, annihilation, ...)

→ direct access to incoming parton (e.g. gluon)

② **fragmentation** of outgoing parton

→ relationship to incoming particle complicated by fragmentation function $D_a^\gamma(z; \mu_f)$

Prompt photon cross section

$$\sigma(p_\gamma) = \sum_{a=q, \bar{q}, g} \int_0^1 \frac{dz}{z} \hat{\sigma}^a\left(\frac{p_\gamma}{z}; \mu_r, \mu, \mu_f\right) \cdot D_a^\gamma(z; \mu_f) + \hat{\sigma}^\gamma(p_\gamma; \mu_r, \mu, \mu_f)$$

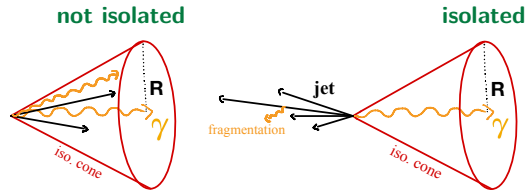
Problem:

- ① (theory+exp) fragmentation prompt photons increase complexity of calculation & relation with initial state
- ② (exp) most photons produced are decay photons (low S/B ratio)

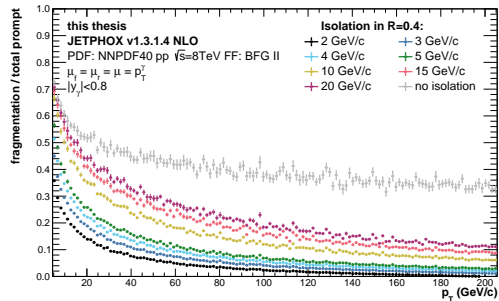
Solution: Isolation

- isolation := applying a restriction on activity in vicinity of photon
- fixed-cone isolation: study $\sum E_T$ in radius R around photon and cut
- smooth-cone isolation (Frixione): reject more activity the closer one gets to photon

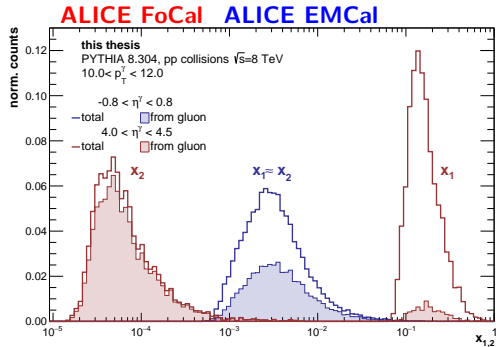
- suppresses fragmentation due to collinear fragmentation
- suppresses decay photons (often come with other particles from hadronization)



fixed-cone isolation



What are we probing with prompt photons?

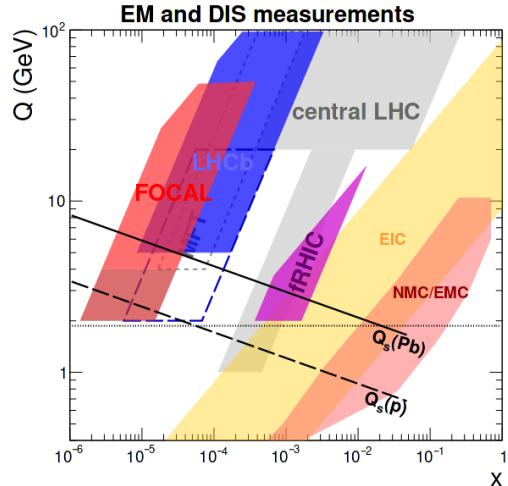


- direct access to gluon density (at LO)
- test for pQCD calculations
- prompt photons allow to probe a wide phasespace in x and Q :

Probed momentum fraction (LO)

$$x \approx \frac{2p_T}{\sqrt{s}} \exp(-y)$$

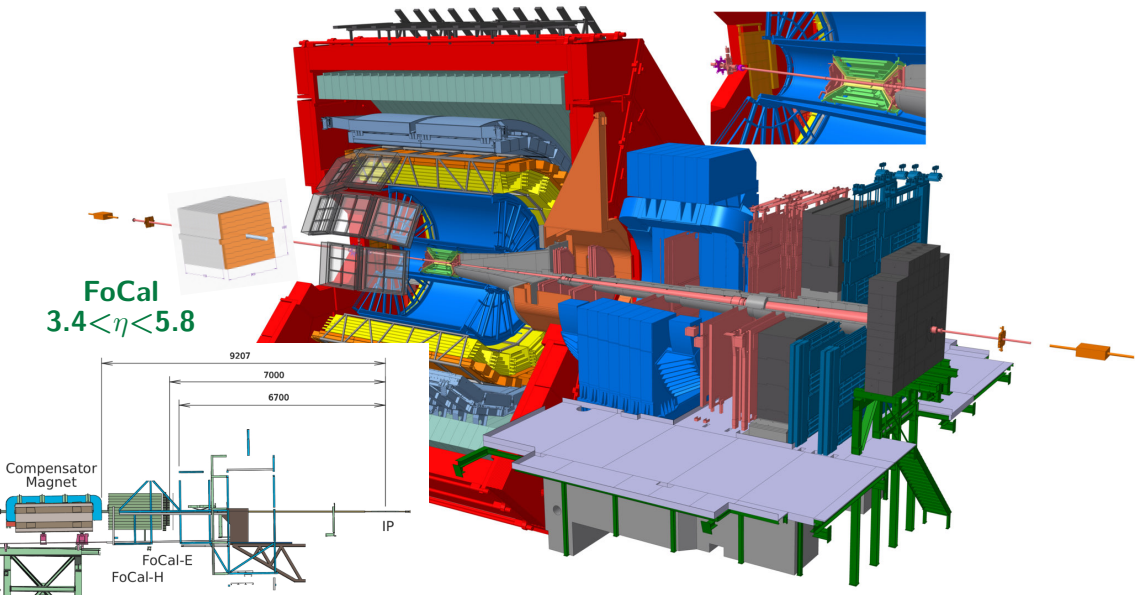
x : momentum fraction carried by parton in proton



the FoCal allows access to saturation regime by measuring at large rapidities!

The Forward Calorimeter (FoCal)

FoCal
 $3.4 < \eta < 5.8$





The Forward Calorimeter (FoCal)

Operational in Run4 (2029)

$$3.4 < \eta < 5.8$$

General:

- very forward calorimeter consisting of two parts (FoCal-E and FoCal-H) located ≈ 7 m from IP
- main physics goal: explore non-linear QCD evolution

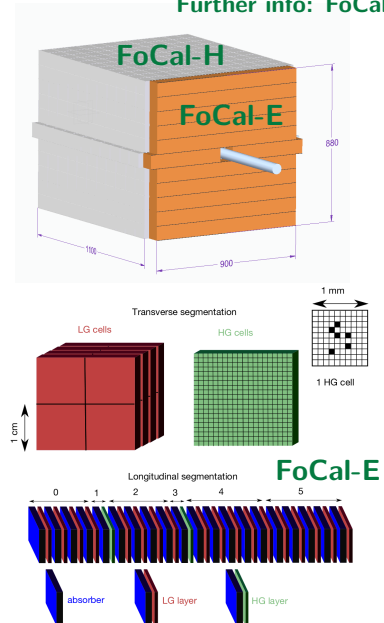
FoCal-E (electromagnetic):

- high-granularity Si-W sampling calorimeter combining two readout granularities ($1 \times 1 \text{ cm}^2$ and $30 \times 30 \mu\text{m}^2$)
- ability to “track” longitudinal component of shower!
- used to measure **photons** and π^0

FoCal-H (hadronic):

- conventional metal-scintillator **hadronic calorimeter** behind FoCal-E
- design using scintillation fibres tested 2021 at SPS test beam
- used to measure **photon isolation**, **jet energy** etc.

Further info: FoCal-Lol

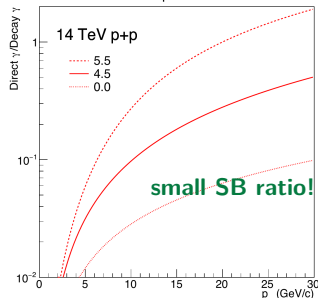
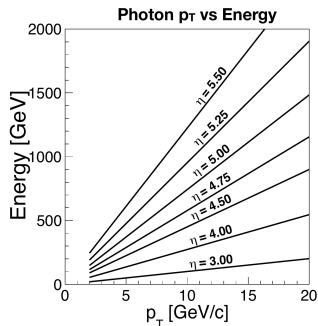


Challenges:

- small signal/background ratio \rightarrow smallest at low- p_T
– the most interesting part for us
- reliable identification of signal and background
- forward rapidity \rightarrow showers with very high energies

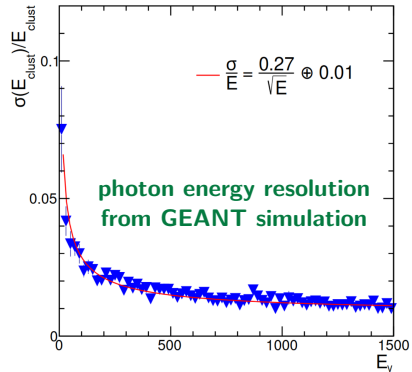
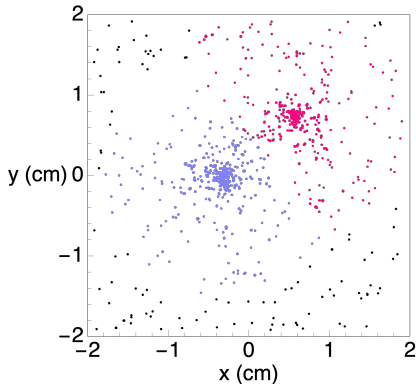
FoCal analysis strategy:

- perform measurement in pp at $\sqrt{s} = 14$ TeV and p-Pb at $\sqrt{s} = 8.8$ TeV to constrain PDFs and access gluon saturation
- measure photons using the FoCal-E
- obtain isolation in cone combining response in FoCal-E and FoCal-H
- further improve S/B by “decay photon tagging”
 - reject pairs of photons that give π^0 mass

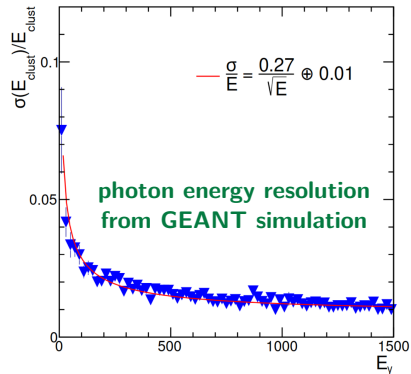
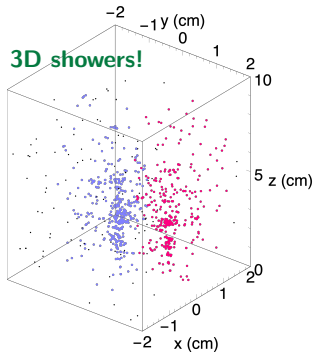


Measuring photons with the FoCal-E

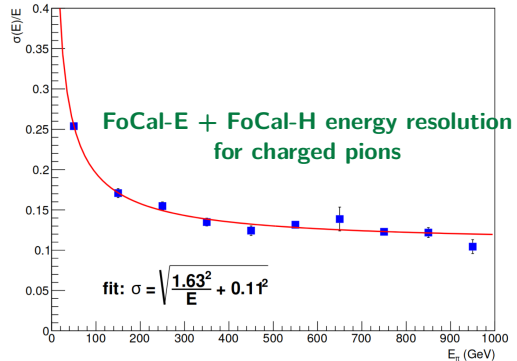
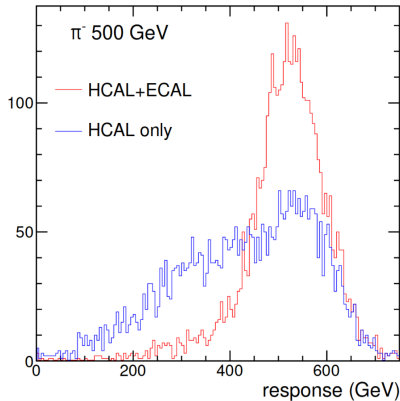
- moderate requirements on energy resolution (constant $< 5\%$ sufficient)
- Tungsten absorber offers low $X_0 \approx 3.5$ mm for a **compact shower** (total depth $\approx 20X_0$)
- **good shower separation** is crucial to identify decay photon background (achieved by the $30 \times 30 \mu\text{m}^2$ pixel layers)
- **photon efficiency (signal)**: ≈ 1 for $E \gtrsim 75$ GeV
- **π^0 efficiency (background)**: over 90 % for $5 \leq p_T \leq 15$ GeV/c



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- to measure isolation of a photon, detector should be capable of measuring charged particles
- main interest in high E hadrons \rightarrow constant term in resolution dominates
- charged pions leave signal in FoCal-E and FoCal-H; response can be combined to single signal

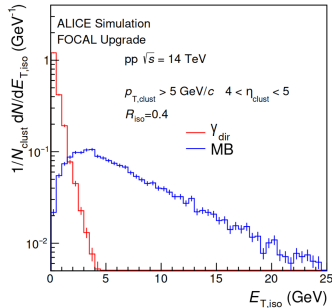


- studies show that already a standard metal-scintillator sampling calorimeter fulfils physics requirements

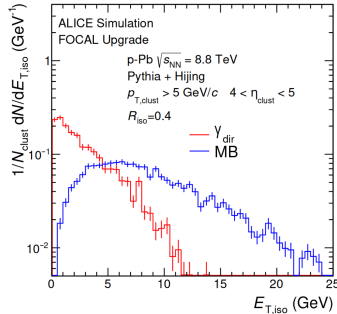
Isolating photons using FoCal-H and FoCal-E

- isolation in cone $R = 0.4$ tested using full MB Pythia simulations and Pythia Gamma-Jet (for p-Pb embedded in Hijing)
- isolation cut $p_T^{\text{iso}} < 5 \text{ GeV}/c$ rejects $\approx 80\% - 90\%$ & selects $\approx 90\%$ signal in pp collisions
- using ECal **and** HCal for isolation improves S/B ratio

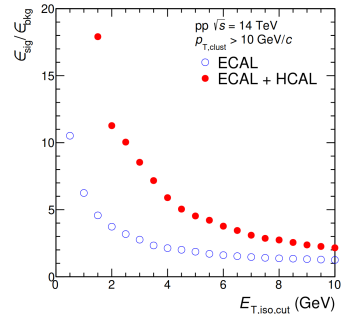
isolation in pp collisions



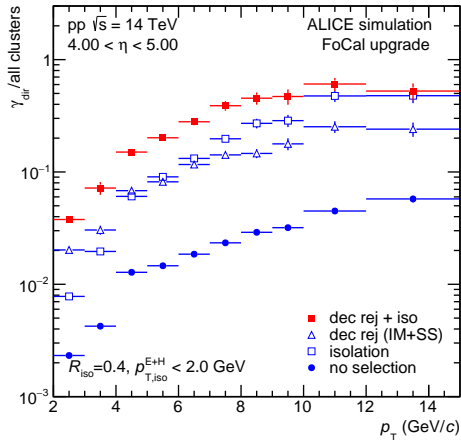
isolation in p-Pb collisions



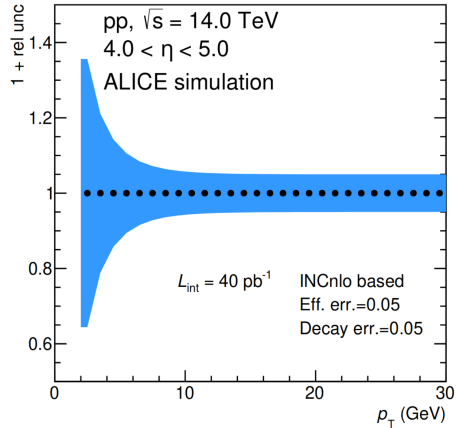
effect of isolation cut



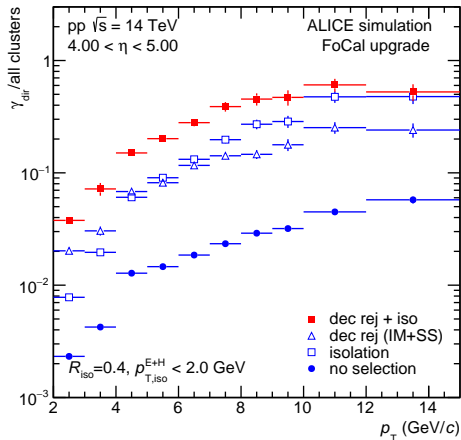
- in addition to isolation **photon tagging** helps to reject photons originating from decays



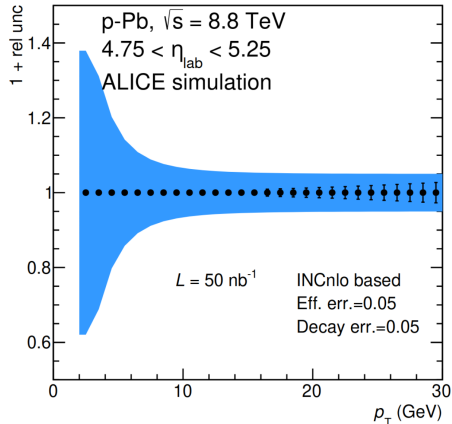
- signal fraction increases significantly by factor 4–8 when applying isolation and decay photon rejection!



- uncertainties in pp (p–Pb) have been estimated to range from 5 % to ≈ 30 %
- main sources: background subtraction (low p_{T}); energy resolution (high p_{T})

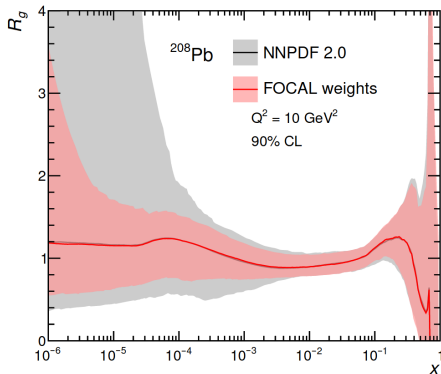
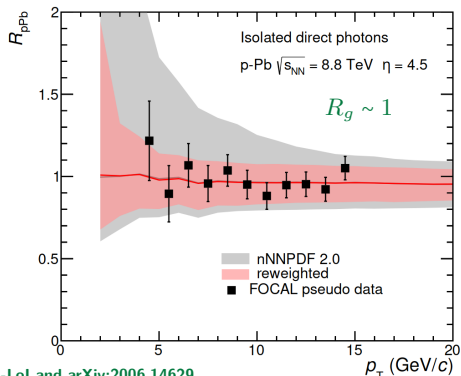
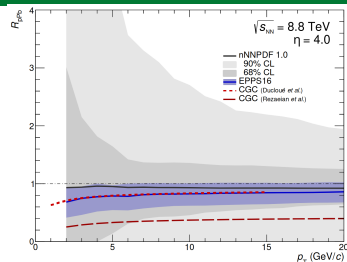


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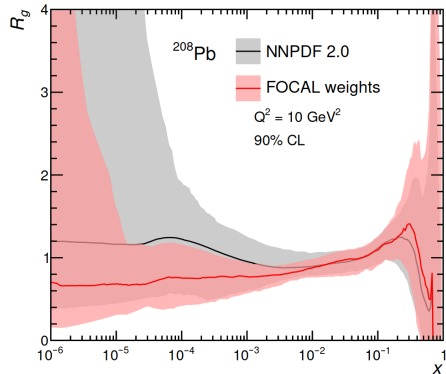
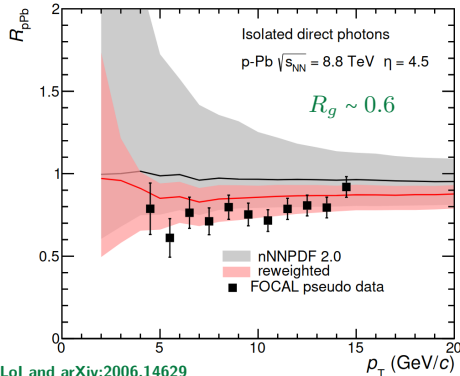
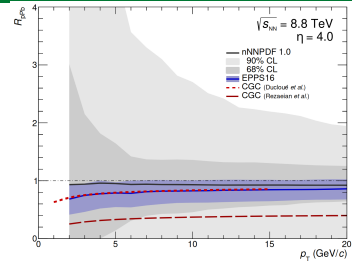


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- pseudo-data projections using INCNLO allow to estimate physics impact
- sys. uncertainties estimated using phys. simulations (dominant low- p_T : purity)
- re-weighting of nNNPDF2.0 using pseudo-data illustrates **potential to significantly reduce nPDF uncertainties at low- x**

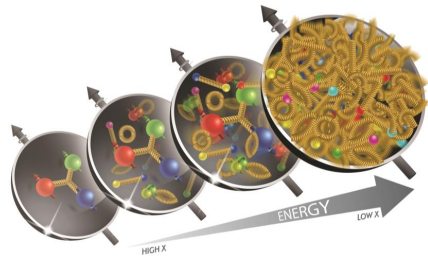
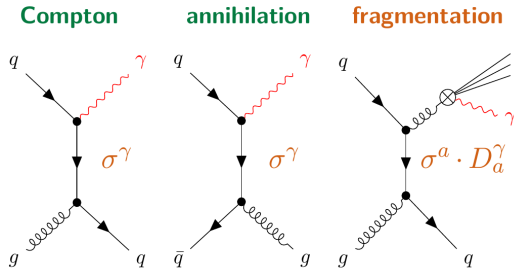


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- isolated photons are a valuable direct probe for the gluon PDF
 - improving uncertainties of PDFs and nPDFs
 - improving knowledge of nuclear effects in the initial state
 - interesting physics at low- x (gluon saturation, CGC)
- the FoCal (planned operation Run4) calorimeter allow to access very low $x \sim 10^{-5}$ by measuring at forward rapidities
- isolated photons are a key ingredients for the physics goal of exploring non-linear QCD evolution

Interested? Comments? Feel free to discuss with me over a coffee and check out the FoCal Lol!



Backup

Mandelstam invariant s

$$\sqrt{s} = \sqrt{(p_1 + p_2)^2} = \sqrt{(p_3 + p_4)^2} \quad (1)$$

Coverage x :

$$\sqrt{s_{12}} = x_1 x_2 \sqrt{s} \quad (2)$$

$$x_1 = \frac{p_T}{\sqrt{s}} (\exp(\eta_3) + \exp(\eta_4)) \quad (3)$$

$$x_2 = \frac{p_T}{\sqrt{s}} (\exp(-\eta_3) + \exp(-\eta_4)) \quad (4)$$

$$x \approx \frac{2p_T}{\sqrt{s}} \exp(-\eta) \quad (5)$$

$$x = \frac{Q^2}{2p_2 \cdot q} \quad (6)$$

$$Q^2 = -q^2 \quad (7)$$

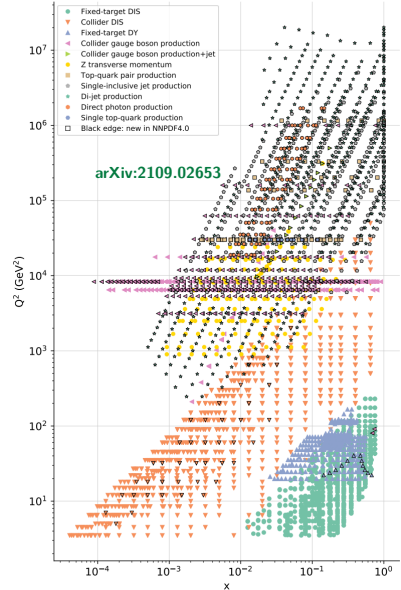
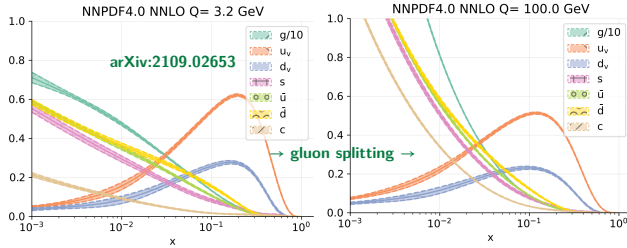
The initial state: Parton Distribution Functions (PDFs)

virtual photon $\lambda \propto 1/Q$

→ DGLAP →

x momentum fraction; Q momentum transfer

Kinematic coverage

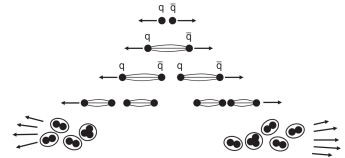
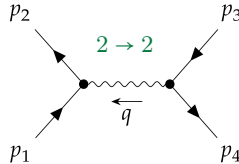
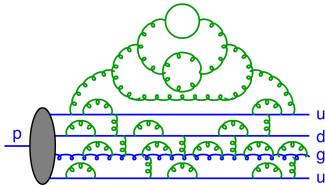


- can not be obtained a priori from calculations*
- PDFs obtained via fitting of vast amount of experimental data, each measurement constraining a different region in x and Q , as well as different parton flavours (a lot of DIS but also collider data from LHC)
- valence quarks peak at $\approx 0.3x$ and as expected $u_v \approx 2d_v$
- only 50% is carried by quarks; the rest by gluons!
- PDFs evolve with scale Q → described by QCD evolution equations (DGLAP etc.)
- low PDF uncertainties → low prediction uncertainties!

* but their evolution with scale can be calculated

Things are complicated ...

pQCD=perturbative QCD



The initial state (before coll.)

- initial state, i.e. state before collision, e.g. partons bound in proton
- complicated structure that also evolves with time (virt. gluons, sea quarks)
- **out of reach for pQCD**

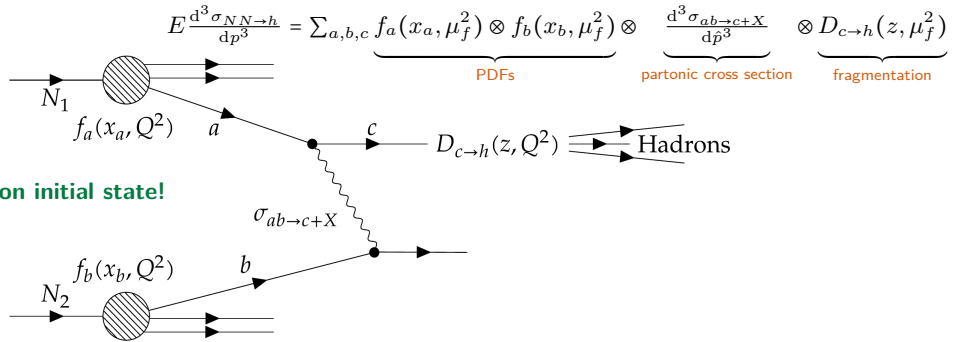
Hard scattering

- process is happening on level of partons e.g. quark-quark \rightarrow quark-quark
- for high enough momentum transfers $-q^2 = Q^2$
 \rightarrow can be treated in pQCD

Final state (after coll.)

- outgoing quarks can not exist as free quarks (confinement)
- process of hadronisation through string breaking
- complicated final state of bound objects and a complicated way to get from quarks to hadrons
 \rightarrow **out of reach for pQCD**

Factorization: A recipe to describe hadronic collisions



Focus today on initial state!

Solution: the problem can be factorized into three components

- factorization possible since fluctuations inside hadron happen on timescale much longer than initial scattering process (scattering “sees” snapshot of structure)
- Parton Distribution Function (PDF)** absorbs non-perturbative physics of initial state in parametrizations that give **probability density** to find a parton i , carrying **momentum fraction x at scale Q** → needs to be determined from experimental data
 - partonic cross section:** describes scattering, can be treated in pQCD for large enough Q
 - Fragmentation function:** relates outgoing quark c to hadron h (absorbs hadronization process)

Sources of photons in hadronic collisions

photons do not interact strongly in final state → carry direct information of production source

Decay photons

- dominant source of photons produced in a collisions
- mainly originating from decays $\pi^0 \rightarrow \gamma\gamma$ and $\eta \rightarrow \gamma\gamma$
- can be used to reconstruct a variety of hadrons that allow to test pQCD, PDF and fragmentation
- interesting rare decays e.g. Higgs $\rightarrow \gamma\gamma$

Thermal photons

- in heavy-ion collisions (e.g. Pb–Pb) energy densities reached are high enough to form QGP
- interactions in partonic- and hadronic phase of QGP produce photons
- slope of spectrum related to effective QGP temperature

relevant at $p_T^\gamma \lesssim 4 \text{ GeV}/c$

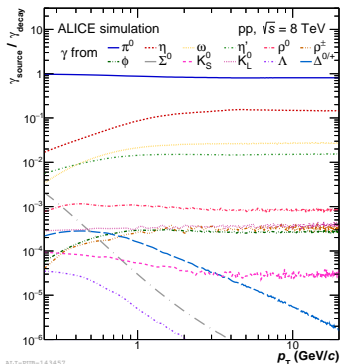
Prompt photons

- photons created in the initial hard scattering
- at LO point-like coupling to partons going into the scattering process → prompt photons contain information about the initial state
- main production mechanism: $qg \rightarrow \gamma q$ (more on next slide)

increasingly relevant at large p_T^γ

photons do not interact strongly in final state \rightarrow carry direct information of production source

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Thermal photons

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- interactions in partonic- and hadronic phase of QGP produce photons
- slope of spectrum related to effective QGP temperature

relevant at $p_T^\gamma \lesssim 4$ GeV/c

Prompt photons

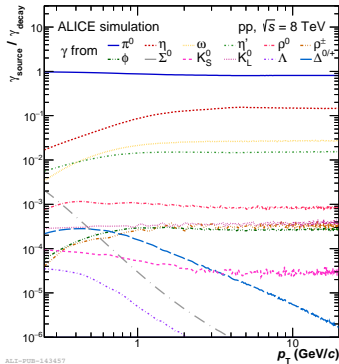
- photons created in the initial hard scattering
- at LO point-like coupling to partons going into the scattering process \rightarrow prompt photons contain information about the initial state
- main production mechanism: $qg \rightarrow \gamma q$ (more on next slide)

increasingly relevant at large p_T^γ

Sources of photons in hadronic collisions

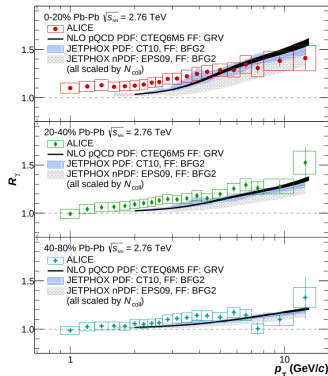
photons do not interact strongly in final state \rightarrow carry direct information of production source

Decay photons



Thermal photons

$$R_\gamma = N_{incl.}^\gamma / N_{decay}^\gamma$$



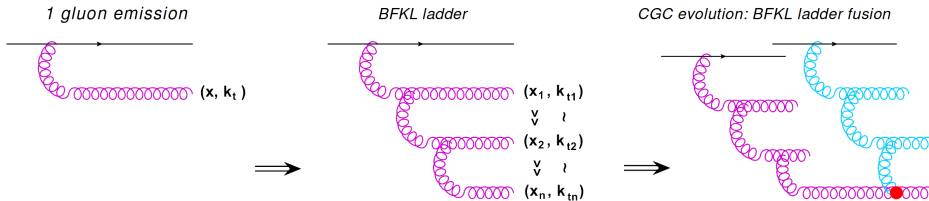
10.1016/j.physletb.2016.01.020

Prompt photons

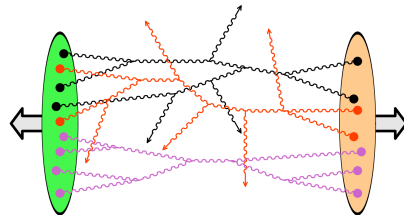
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increasingly relevant at large p_T^γ

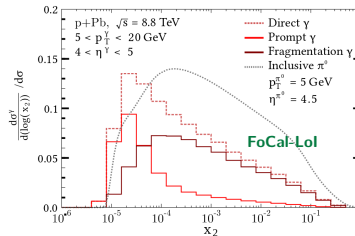
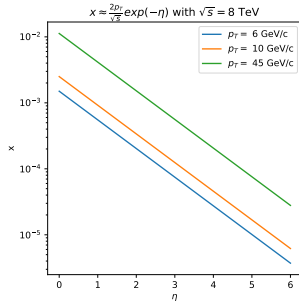
A closer look at saturation



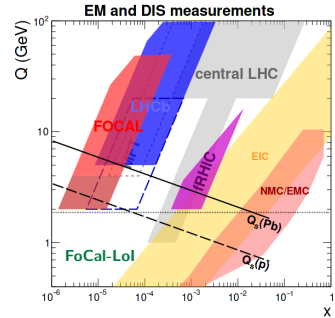
- probability P of Bremsstrahlung emission of a gluon carrying x increases with $P \sim \alpha_s \ln(1/x)$
 → exponential rise of soft gluons at low- x described by BFKL/DGLAP evolution
- when gluon densities get very large, gluon recombination/fusion occurs, probability depends on gluon density
- initial-state of collision is dense gluon color fields described by CGC
- early stage of collision is initial condition for hydro transport models, and affects event bulk properties (multiplicity, flow, ...)



- measuring at very-forward rapidities $3.4 < \eta < 5.8$ allows to access low- x regime
- non-linear evolution becomes accessible below saturation scale Q_s



accessible gluon- $x \sim 10^{-5}$



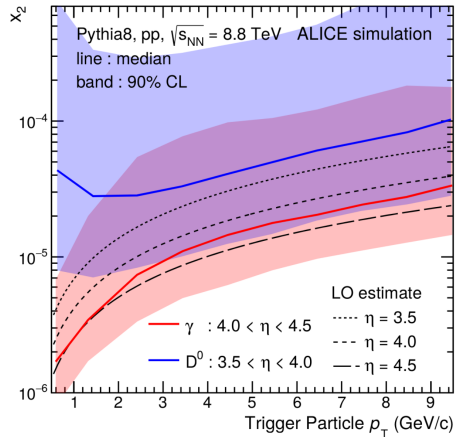
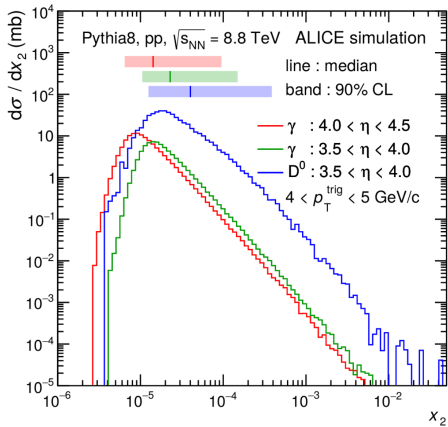
- going forward (large η) exponentially lowers x

- a more detailed look at gluon- x sensitivity using JETPHOX NLO + EPS09

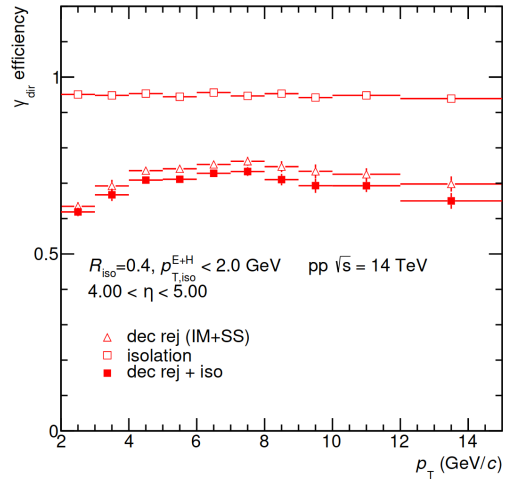
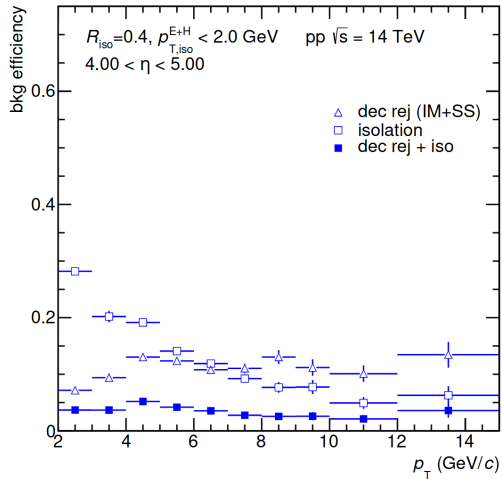
- FoCal competitively placed in (x, Q) phasespace

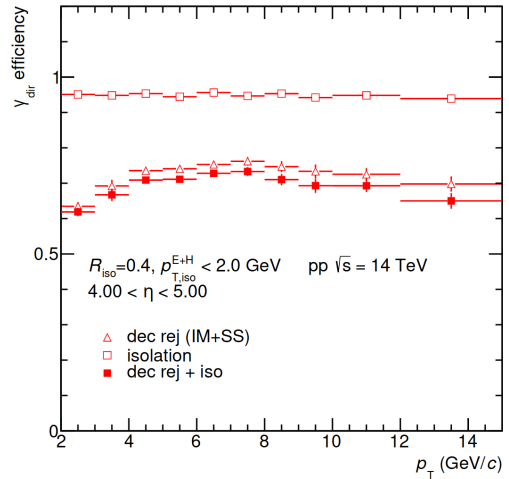
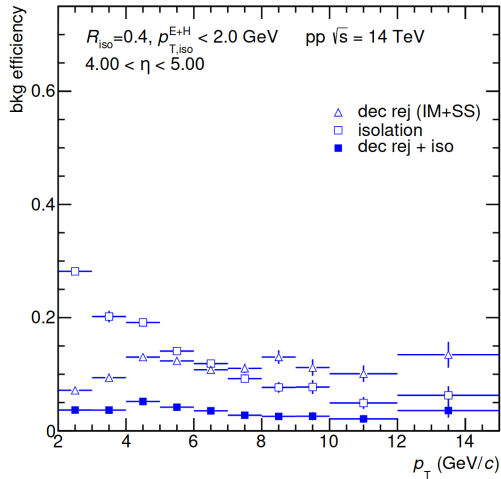
Access to the gluon: Photons vs. D^0

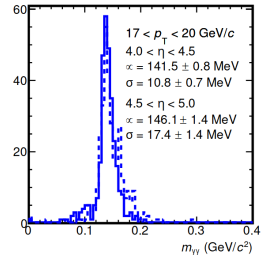
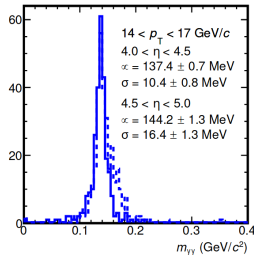
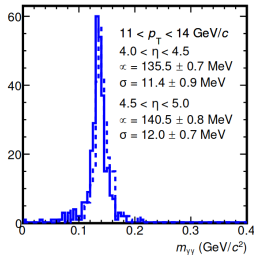
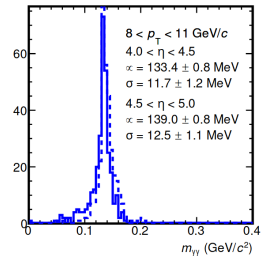
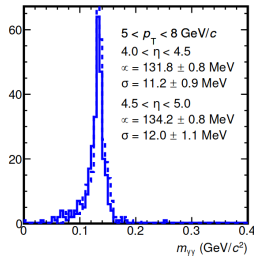
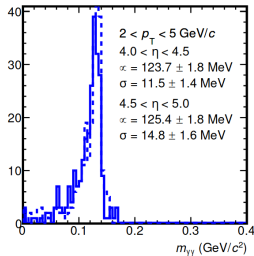
- direct prompt photons are relatively easy to calculation + direct connection to gluons at LO
- fragmentation component highly suppressed when using isolation
- no color-charge \rightarrow no final-state energy loss



- direct photons probe lower x than D^0 (open charm); broadening from fragmentation

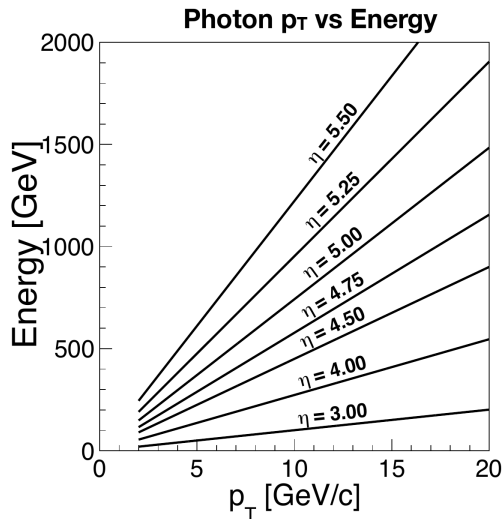
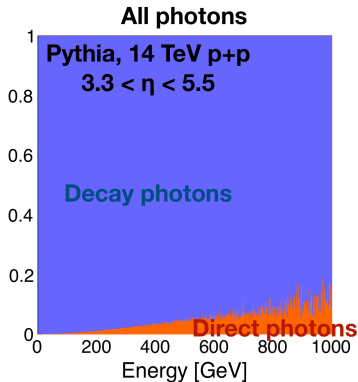


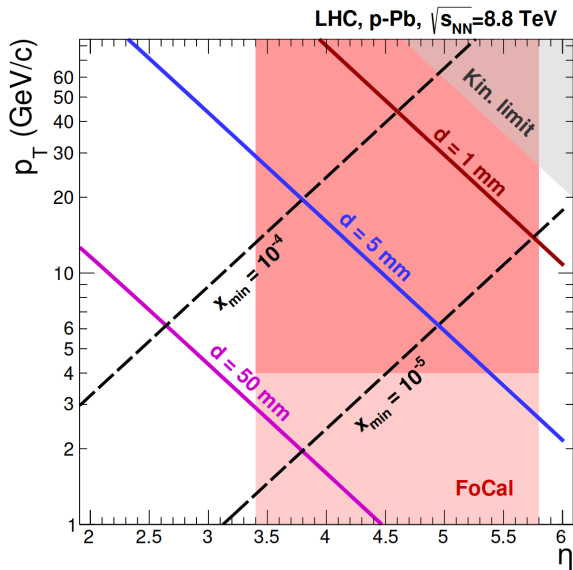


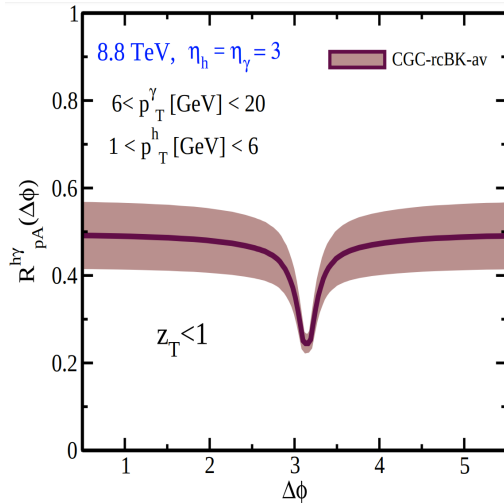
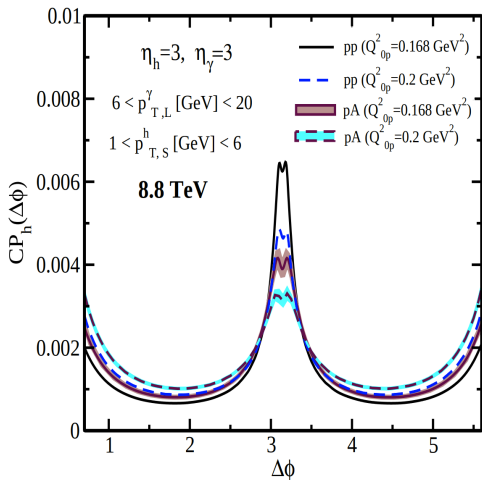


- large forward rapidities \rightarrow large energies!
- measurements in $3.3 < \eta < 5.8$ for $p_T \approx 4 - 20 \text{ GeV}/c$ correspond to 40 to 2500 GeV (!)

\rightarrow detector needs to measure small signal over large energy ranges







- γ -hadron correlations promising measurement to investigate nonlinear QCD effects

- F.Jonas and C. Loizides: Centrality dependence of electroweak boson production in PbPb collisions at the CERN Large Hadron Collider (*Phys. Rev. C* 104, 044905 (2021))
- ALICE Collaboration: Production of ω mesons in pp collisions at $\sqrt{s} = 7$ TeV *The European Phys. Journal C* 80, 1130 (2020)

for more information about myself and contact information, feel free to visit
<https://florianjonas.xyz>