## Far-Forward Detectors at the Electron-Ion Collider

**Michael Murray, Forward QCD Lawrence 24 May 2022** 



# The Electron-Ion Collider (EIC)

Polarized

Electron

Source

- Two interaction regions (IRs) for possible detector locations.
- Only one, IP6, in DOE project scope.

Electron niection

Possible

**Detector** 

Electron **Injector (RCS** 

(Polarized) **Ion Source** 

**Possibl** On-energy **Ion Injector** 

> Electron Cooler



 $\triangleright$  Reference detector based on the 1.5T BaBar solenoid and ECCE reference design.

 $\triangleright$  Contains detectors for tracking, PID, and calorimetry.

## The Electron-Ion Collider (EIC)



• In addition to the central detector  $\rightarrow$ detectors integrated into the beamline on both the hadron-going (**far-forward**) and electron-going (**far-backward**) direction.

**Detectors have to be very tightly integrated with machine. The large crossing angle and short bunch crossing time cause many challenges.** 

### **Great Fun**

## Far-Forward Physics at the EIC



## Far-Forward Physics at the EIC



5

## Far-Forward Physics at the EIC

- $\overline{\phantom{a}}$ charged had  $\overline{\phantom{a}}$ • Physics channels require tagging of **charged hadrons** (protons, proton or neutron tagging1 **Short Saturation**  $r$ apidities ( $\eta > 4.5$ ). artici<del>c</del>s<br>). **Structure** pions) or **neutral particles** (neutrons, photons) at **very-forward**
- ferent 1 • Different final states  $\rightarrow$  tailored detector subsystems.
- UCV, CIP, CIU, CINU, CLC.J. • Various collision systems and energies (h: 41, 100-275 GeV, e: 5-18 GeV; e+p, e+d, e+Au, etc.).
- re uniqual • Placing of far-forward detectors uniquely challenging due to **Short-Bree Function**<br> **Short-Breed in ELC Vallow**
- Y. Furletova, O. Hen, D. W. Higinbotham, C. Hyde, V. BOOIT AND CONCEDIU C **105**, 034001 (2022) [5] AJ, Z. Tu, and C. Weiss, Phys. Rev. C **104,** 065205 and in the ATHENA, ECCE, and CORE EIC detector proposals. **A** EIC Y negrat<br>Jetails s • Details studied in EIC Yellow Report and Conceptual Design Report,

## The Far-Forward Detectors



### B0 Detectors

![](_page_7_Figure_1.jpeg)

### B0 Detectors

 $(5.5 < \theta < 20.0 \text{ mrad})$ 

- $\triangleright$  Charged particle reconstruction and photon tagging.
- $\triangleright$  Precise tracking (~10um spatial resolution).
- $\triangleright$  Fast timing for background rejection and to remove crab smearing (~35ps).
- $\triangleright$  Photon detection (tagging or full reco).

![](_page_8_Picture_6.jpeg)

![](_page_8_Figure_7.jpeg)

**Hadrons**

Alex Jentsch - DIS 2022 - Santiago de Compostela, Spain - May 2nd to 6th, 2022

### B0 Detectors

 $(5.5 < \theta < 20.0 \text{ mrad})$ 

![](_page_9_Figure_2.jpeg)

![](_page_9_Figure_3.jpeg)

- Higher granularity silicon (e.g. MAPS) required.
- Tagging photons important in differentiating between coherent and incoherent heavy-nuclear scattering, and for reconstructing  $\pi^0 \to \gamma \gamma$ .
- **Space is a major concern here – an EMCAL is highly preferred, but we may only have space for a preshower.**

## B0-detectors (calorimetry)

![](_page_10_Figure_1.jpeg)

- For studies of *u*-Channel (Backwardangle) exclusive electroproduction, need capability to reconstruct photons from  $\pi^0$  decays.
	- Physics beyond the EIC white paper!
- Would require full EMCAL with high granularity and energy resolution.  $\triangleright$  PbWO4 used in ECCE studies.
- Longitudinal space in B0pf magnet limited.
	- Would be a great candidate for an upgrade or for IP8 complementarity!

### Roman Pots @ the EIC

![](_page_11_Figure_1.jpeg)

## Roman "Pots" @ the EIC

![](_page_12_Figure_1.jpeg)

![](_page_12_Figure_2.jpeg)

- Putting silicondirectly into machine vacuum maximizes geometric coverage
- Need space for detector insertion tooling and support structure.
- Cooling is vital

## Roman "Pots" @ the EIC

![](_page_13_Figure_1.jpeg)

**DD4HEP Simulation**

![](_page_13_Figure_3.jpeg)

• **Two main options**

- AC-LGAD sensor provides both fine pixilation (~140um spatial resolution), and fast timing (~35ps).
- MAPS + LYSO timing layer.
- "Potless" design concept with thin RF foils surrounding detector components.

## Roman "Pots" @ the EIC

![](_page_14_Figure_1.jpeg)

**DD4HEP Simulation**

 $\sigma(z)$  is the Gaussian width of the beam,  $\beta(z)$  is the RMS transverse beam size.

 $\varepsilon$  is the beam emittance.

$$
\sigma(z) = \sqrt{\varepsilon \cdot \beta(z))}
$$

![](_page_14_Figure_6.jpeg)

- **Low-pT cutoff determined by beam optics.**
	- The safe distance is  $\sim 10\sigma$  from the beam center.
	- $\mathbf{1}\sigma \sim \mathbf{1}$ mm
- **Optics choices change with energy, but be changed within a single energy to maximize** *either acceptance at the RP, or the luminosity.*

## Off-Momentum Detectors

![](_page_15_Figure_1.jpeg)

**EICROOT GEANT4 simulation.**

### Summary of Detector Performance (Trackers)

![](_page_16_Figure_1.jpeg)

## Zero-Degree Calorimeter

![](_page_17_Figure_1.jpeg)

 $E_{\gamma}$  [GeV]

Neutron Energy [GeV]

# Summary and Takeaways

- All FF detector acceptances and detector performance well-understood with currently available information.
	- Numerous impact studies done!
	- Some final choices on technology underway  $\rightarrow$  also important for IP8 complementarity.
	- Full effort benefitted from three (ECCE, ATHENA, CORE) proposals to identify multiple technology solutions!
- More realistic engineering considerations need to be added to simulations as design of IR vacuum system and magnets progresses toward CD-2/3a.
	- Lots of experience in performing these simulations, so this work will progress rapidly as engineering design matures.
	- Already well-established line of communication between detector and physics parties and the EIC machine/IR development group <sup>⇒</sup> Crucial for success!!!

![](_page_18_Figure_8.jpeg)

![](_page_19_Picture_0.jpeg)

### Summary of Detector Performance (Trackers)

![](_page_20_Figure_1.jpeg)

- Includes realistic considerations for pixel sizes and materials
	- More work needed on support structure and associated impacts.
- Roman Pots and Off-Momentum detectors suffer from additional smearing due to improper transfer matrix reconstruction.
	- This problem is close to being solved!

### Digression: Machine Optics

### 275 GeV DVCS Proton Acceptance

![](_page_21_Figure_2.jpeg)

![](_page_21_Figure_3.jpeg)

![](_page_21_Figure_4.jpeg)

**High Divergence:** smaller  $\beta^*$  at IP, but bigger  $\beta(z = 30m) \rightarrow$ higher lumi., larger beam at RP

### Digression: Machine Optics

### 275 GeV DVCS Proton Acceptance

![](_page_22_Figure_2.jpeg)

![](_page_22_Figure_3.jpeg)

![](_page_22_Figure_4.jpeg)

**High Divergence:** smaller  $\beta^*$  at IP, but bigger  $\beta(z = 30m) \rightarrow$ higher lumi., larger beam at RP

**High Acceptance:** larger  $\beta^*$  at IP, smaller  $\beta(z = 30m) \rightarrow$ lower lumi., smaller beam at RP

### Digression: Machine Optics

### 275 GeV DVCS Proton Acceptance

![](_page_23_Figure_2.jpeg)

**30**

## So how does the FF system perform for measurements (non-exhaustive)?

![](_page_25_Figure_0.jpeg)

## Off-Momentum Detectors

• Off-momentum protons  $\rightarrow$  smaller magnetic rigidity → **greater bending in dipole fields.**

![](_page_26_Figure_2.jpeg)

![](_page_26_Figure_3.jpeg)

# Digression: particle beams

### • Angular divergence

- Angular "spread" of the beam away from the central trajectory.
- Gives some small initial transverse momentum to the beam particles.
- Crab cavity rotation
	- Can perform rotations of the beam bunches in 2D.
	- Used to account for the luminosity drop due to the crossing angle – allows for head-on collisions to still take place.

![](_page_27_Figure_7.jpeg)

These effects introduce smearing in our momentum reconstruction.

# Spectator Tagging in Light Nuclei

EIC enables use of deuteron beams  $\rightarrow$  the next best thing to a beam of neutrons!

![](_page_28_Figure_2.jpeg)

- Measurements on unpolarized deuterons<sup>1</sup> (or polarized He-3)<sup>2</sup> at the EIC.
- Spectator proton momentum  $\rightarrow$  enables selection of nuclear (p/n) configurations.
	- Extract free neutron structure function<sup>3</sup>  $\rightarrow$  Not possible elsewhere!
	- Study nuclear modifications of both nucleons in the deuteron (study in progress).

[1] Z. Tu, A. Jentsch, et al., Physics Letters B, (2020) [2] I. Friscic, D. Nguyen, J. R. Pybus, A. Jentsch, *et al.,* Phys. Lett. B, **Volume 823**, 136726 (2021)

[3] A. Jentsch, Z. Tu, and C. Weiss, Phys. Rev. C **104,** 065205, (2021) **(Editor's Suggestion)**

![](_page_29_Figure_0.jpeg)

#### 30

![](_page_30_Figure_0.jpeg)

e'

#### **Proton spectator case.**

Particular process in BeAGLE: incoherent diffractive J/psi

**Short-range correlations!**

![](_page_30_Picture_5.jpeg)

18x110GeV

- Spectator kinematic variables reconstructed over a broad range.
- All beam/detector effects included.
- Bin migration is observed due to smearing in the reconstruction.

Ø **In the proton spectator case, essentially all spectators tagged.**

Ø **Active neutrons only tagged up to 4.5 mrad.**

Alex Jentsch - DIS 2022 - Santiago de Compostela, Spain - May 2nd to 6th, 2022

![](_page_31_Figure_0.jpeg)

![](_page_32_Figure_0.jpeg)

 $\gamma^*$ 

n

d  $\left(\begin{matrix} p \end{matrix}\right)$ 

**t-reconstruction** using **doubletagging (both proton and neutron).** Takes advantage of combined B0 + off-momentum detector coverage. Better coverage in the neutron spectator case.

Ø **Spectator information is the "dial" for the SRC region.**

 $p'$ 

 $t=(p'-p)^2$ 

 $t'=(n'-d)^2-M_p$ 

 $J\!/\mathrm{u}$ 

n'

## Free Neutron  $F<sub>2</sub>$  Extraction

**(Active nucleon reduced cross section)**  $\sim$   $\mathbf{F}_2$ 

![](_page_33_Figure_2.jpeg)

![](_page_33_Picture_3.jpeg)

- Cross-section as a function of the proton spectator kinematics  $\rightarrow$  dial to select nuclear configuration  $\rightarrow$  allows extrapolation to "free" neutron region.
- Enables measurement of **free** neutron structure function!

 $P_{pT}^2 = p_{px}^2 + p_{py}^2$  $\sigma_{red,n} \sim F_{2,n}$  (cross section)

A. Jentsch, Z. Tu, and C. Weiss, Phys. Rev. C **104,** 065205, (2021) **(Editor's Suggestion)**

## Neutron Spin Structure in He3

- Studies of neutron structure with a *polarized* neutron.
- More challenging final state tagging since *both* protons must be tagged in the FF region.
- MC events generated with CLASDIS in fixed-target frame, and then boosted to collider frame.

![](_page_34_Figure_4.jpeg)

![](_page_34_Figure_5.jpeg)

![](_page_35_Figure_0.jpeg)

Alex Jentsch - DIS 2022 - Santiago de Compostela, Spain - May 2nd to 6th, 2022

### Roman Pots

- Active sensor area very large (26cm x 13cm).
- "Potless" design could make better use of space.
- With AC-LGADS + ALTIROC ASIC, current estimates of power dissipation around 400-500 watts for entire subsystem, so roughly 100 watts/layer.
	- With potless design, leveraging experience from LHCb VELO for cooling would allow for cooling of the electronics within the vacuum.
- Support structure only to be placed between hadron pipe and wall to avoid interference with the ZDC.

### Roman Pots

• Updated layout with current design for AC-LGAD sensor + ASIC.

![](_page_37_Figure_2.jpeg)

• Current R&D aimed at customizing ASIC readout chip (ALTIROC) for use with AC-LGADs.

![](_page_37_Picture_119.jpeg)

## Luminosity Monitor

- Must make measurement in challenging environment.
	- High synchrotron radiation, high bremsstrahlung rates (~10 GHz), etc.
- Need  $\sim$  1% for absolute luminosity measurement,  $\sim$  10<sup>-4</sup> for relative luminosity measurement.
- Can make direct photon measurement, or indirect via pair conversion in exit window, where e+e- pair is steered toward two calorimeters opposite a dipole magnet.
- Direct photon calorimeter includes moveable SR filters/monitors (F1 and F2), and has configurations for high (PCALf) and low (PCALc) luminosity running.

![](_page_38_Figure_6.jpeg)

# Exit window for luminosity monitor

- Part of outgoing electron beam pipe
- Conversion layer for bremsstrahlung photons
- Tilt angle vs. electron (and photon) beam axis against synchrotron radiation

![](_page_39_Figure_4.jpeg)

# Low-Q2 Taggers

- Two taggers for reconstructing electrons from low- $Q^2$  (< 10<sup>-1</sup> GeV<sup>2</sup>) reactions.
- Combination of EM calorimetry for energy reconstruction, and silicon layers (High Resolution Hodoscope – HIHS) for position and angular resolution.

![](_page_40_Figure_3.jpeg)

### Performance for low-Q2 tagger

- Tagger 1 and 2 are placed closer (further) from the IP
- Overlap in Q2 acceptance (< 0.1 GeV^2)
- Complementary in electron energy (higher energies reach Tagger 2)
- Consistent for Pythia6 and quasi-real photoproduction (QR)

![](_page_41_Figure_5.jpeg)

Electron energy E (GeV)

Tagger 1

 $0.9$ 

 $10.8$ 

 $\neg$ 0.7

 $-0.6$ 

 $-0.5$ 

 $-10.4$ 

 $-0.3$ 

### Machine Optics: Roman Pots

![](_page_42_Figure_1.jpeg)

## Momentum Resolution – Timing

For exclusive reactions measured with the Roman Pots we need good timing to resolve the position of the interaction within the proton bunch. But what should the timing be?

![](_page_43_Figure_2.jpeg)

- Because of the rotation, the Roman Pots see the bunch crossing smeared in x.
- Vertex smearing = 12.5mrad (half the crossing angle) \* 10cm = **1.25 mm**
- If the effective vertex smearing was **for a 1cm bunch**, we would have **.125mm** vertex smearing.
- The simulations were done with these two extrema and the results compared.
	- Ø **From these comparisons, reducing the effective vertex smearing to that of the 1cm bunch length reduces the momentum smearing to negligible from this contribution.**
	- Ø **This can be achieved with timing of ~ 35ps (1cm/speed of light).**

### Momentum Resolution – Comparison

• The various contributions add in quadrature (this was checked empirically, measuring each effect independently).

![](_page_44_Figure_2.jpeg)

#### • **Beam angular divergence**

- Beam property, can't correct for it sets the lower bound of smearing.
- Subject to change (i.e. get better) beam parameters not yet set in stone
- **Vertex smearing from crab rotation**
	- Correctable with good timing (~35ps)
- **Finite pixel size on sensor**
	- 500um seems like the best compromise between potential cost and smearing

## Free Neutron F<sub>2</sub> Extraction

![](_page_45_Figure_1.jpeg)

 $\sigma$ 

- Resulting dependence on  $p_{pT}^2$  is very weak and the extrapolation can be performed with a  $1<sup>st</sup>$ -degree polynomial fit.
- Extrapolation only performed for the generator-level distribution.  $R = residue$  of spectral function  $a_T^2$  = position of pole  $S_d(p_{pT}, \alpha_p)[pole]=$  $\overline{R}$  $p_{pT}^2 + a_T^2$ <sup>2</sup>  $R = 2\alpha_p^2 m_N \Gamma^2 (2 - \alpha_p)$  $a_T^2 = m_N^2 - \alpha_p (2 - \alpha_p)$  $M_d^2$  $\begin{array}{cc} 4 & (PpT + \alpha_T) \end{array}$

# What about IP8?

![](_page_47_Figure_0.jpeg)

## Major potential benefit: Secondary Focus

![](_page_48_Picture_1.jpeg)

- Allows for tagging of protons and nuclei at very **high values of xL close to one** (pT ~ 0).
- Complementarity with the IP6 configuration and detector – important for the EIC!

![](_page_49_Figure_0.jpeg)

![](_page_50_Figure_0.jpeg)