

Status of theory calculations for Z-pole observables

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$\lim_{n \rightarrow \infty} \left\{ \frac{1}{n} \right\} \rho(x, x')$
 $\exists M \in \mathbb{R} \forall n \in \mathbb{N} \forall x, x' \in \mathbb{R} \rho(f(x), f(x')) \leq q \cdot \rho(x, x')$
 $\forall n \in \mathbb{N} x_n \leq y_n \leq z_n$
 $\lim_{n \rightarrow \infty} y_n = g$
 $\lim_{n \rightarrow \infty} \sqrt[n]{1 + e^n + \pi^n + 13^n} \leq q \cdot \lim_{n \rightarrow \infty} x_n = q$
 $\lim_{n \rightarrow \infty} z_n = g$
 $g \in \mathbb{R} - (-\infty, \infty)$

$\{x_n\} \subset \mathbb{R}$
 $\{y_n\} \neq 0 \Leftrightarrow y_n \neq 0$
 $\{z_n\} \subset \mathbb{R}$
 $\lim_{n \rightarrow \infty} \left(1 + \frac{\pi}{1+n^2}\right) \{x_n\} \subset \mathbb{R}$

$N \rightarrow \mathbb{R} x: \rho$
 $\forall n \in \mathbb{N}$, to $\left\{ \frac{x_n}{y_n} \right\} \stackrel{\text{df}}{=} \left\{ \frac{x_n}{y_n} \right\}$
 $n \in \mathbb{N}, A > 0, \Rightarrow \lim_{n \rightarrow \infty} \sqrt[n]{A} = 1$

$\sqrt[n]{5^n} \left\{ \frac{1}{n} \right\} A_y$
 $\sqrt[n]{4^n \cdot \cos 2n}$
 $\left(\frac{n^2 + n - 1}{n^2 - 2n + 3} \right)^5$
 $\forall n \in \mathbb{N} x_n \leq y_n \leq z_n$

ρ
 $\{1 + \frac{1}{n}\}$
 $x_n + y_n$
 $N \rightarrow \mathbb{R} n \geq n_0: (x_n - g) < \epsilon$
 $\text{lokal. max; } \{x_n\}: x_n = \frac{1}{n}; \{y_n\}: y_n = 1 + \frac{1}{n}$
 $\rho(f(x), f(x')) \leq q$

$f(x) \Leftrightarrow \exists q \in [0, 1]: \forall x, x' \in X$
 $(x_n - g) < \epsilon \quad n \geq n_0: (x_n - g) < \epsilon$

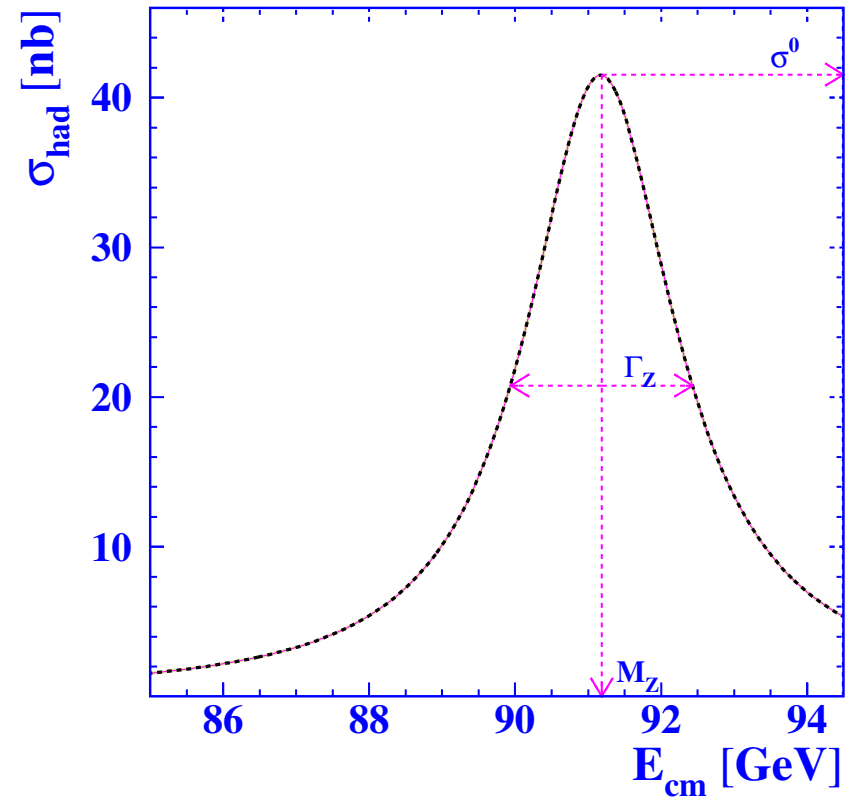
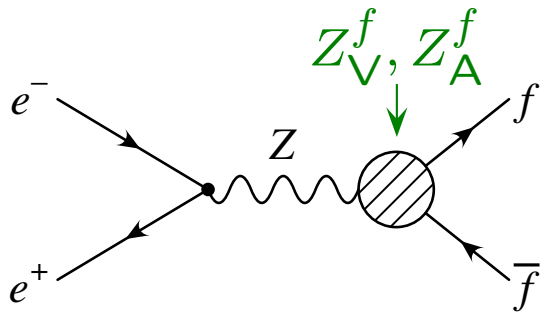
$\left\{ \frac{1}{n} \right\} = \left\{ \frac{1}{n} \right\}$
 $x_n: N \rightarrow \mathbb{R}$
 $\lim_{n \rightarrow \infty} \min$
 lok. min
 $\sqrt[n]{4!} \cdot \sqrt[n]{13^n} \cdot \sqrt[n]{13^n}$
 $\sqrt[n]{0+0+0+13^n} \leq \sqrt[n]{1^n + e^n + \pi^n + 13^n} \leq \sqrt[n]{13^n + 13^n + 13^n + 13^n}$
 $\{x_n\} \subset \mathbb{R}$
 $\{y_n\} \stackrel{\text{df}}{=} \{x_n + y_n\}$
 $\{x_n\} \cdot \{y_n\} \stackrel{\text{df}}{=} \{x_n \cdot y_n\}$

$g \in [0, 1)$
 $\left\{ \frac{1}{n} \right\} \left\{ \frac{1}{n} \right\}$
 g
 g
 $\sqrt[n]{5^n} = 5$
 $f: X \rightarrow x$

$e^+e^- \rightarrow f\bar{f}$ for $E_{\text{CM}} \sim M_Z$:

- Mass M_Z
- Width $\Gamma_Z = \sum_f \Gamma_{ff}$
- Branching ratio $R_f = \Gamma_{ff}/\Gamma_Z$
- $\sigma^0 \approx \frac{12\pi \Gamma_{ee} \Gamma_{ff}}{(s-M_Z^2)^2 + M_Z^2 \Gamma_Z^2} = \frac{12\pi}{M_Z^2} R_e R_f$

$$\Gamma_{ff} = C \left[(Z_V^f)^2 + (Z_A^f)^2 \right]$$

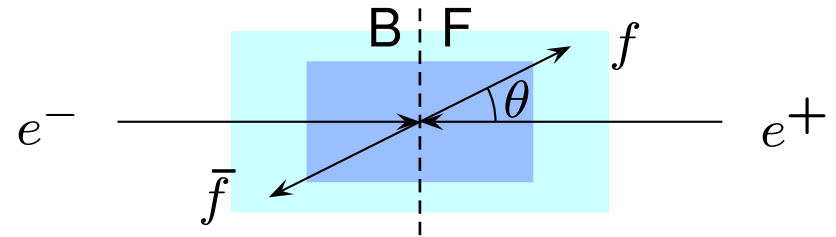


Forward-backward asymmetry:

$$A_{FB} \equiv \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B} = \frac{3}{4} A_e A_f$$

$$A_f = \frac{2(1 - 4\sin^2 \theta_{\text{eff}}^f)}{1 + (1 - 4\sin^2 \theta_{\text{eff}}^f)^2}$$

$$\sin^2 \theta_{\text{eff}}^f = \frac{1}{4|Q_f|} \left[1 - \frac{Z_V^f}{Z_A^f} \right]$$



Left-right asymmetry:

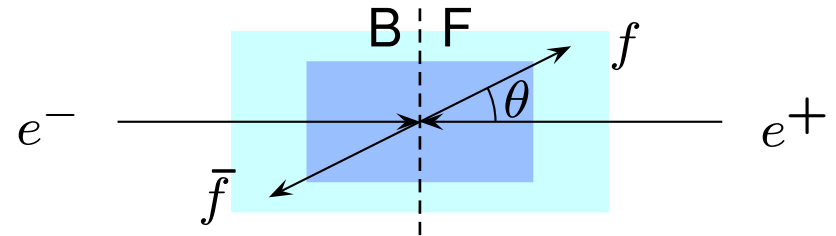
With polarized e^- beam: $A_{LR} \equiv \frac{\sigma_L - \sigma_R}{\sigma_L + \sigma_R} = A_e$

Polarization asymmetry:

Average τ pol. in $e^+e^- \rightarrow \tau^+\tau^-$: $\langle \mathcal{P}_\tau \rangle = -A_\tau$

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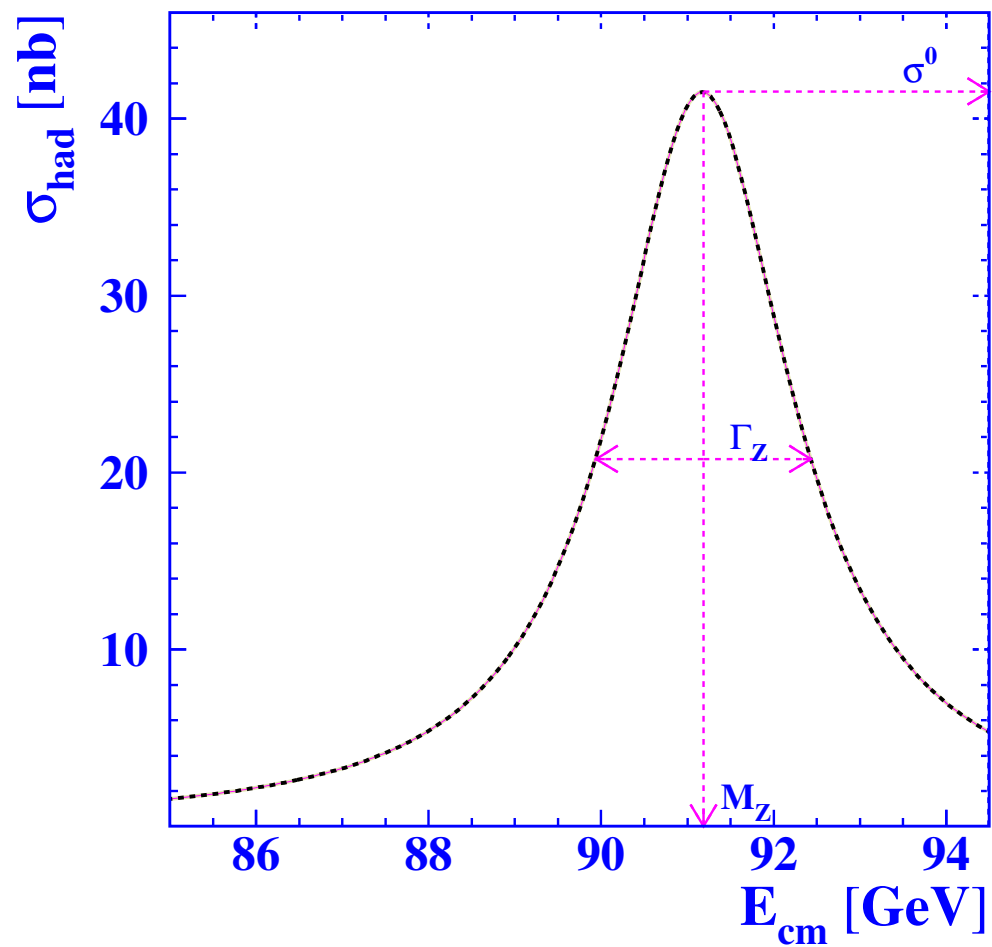
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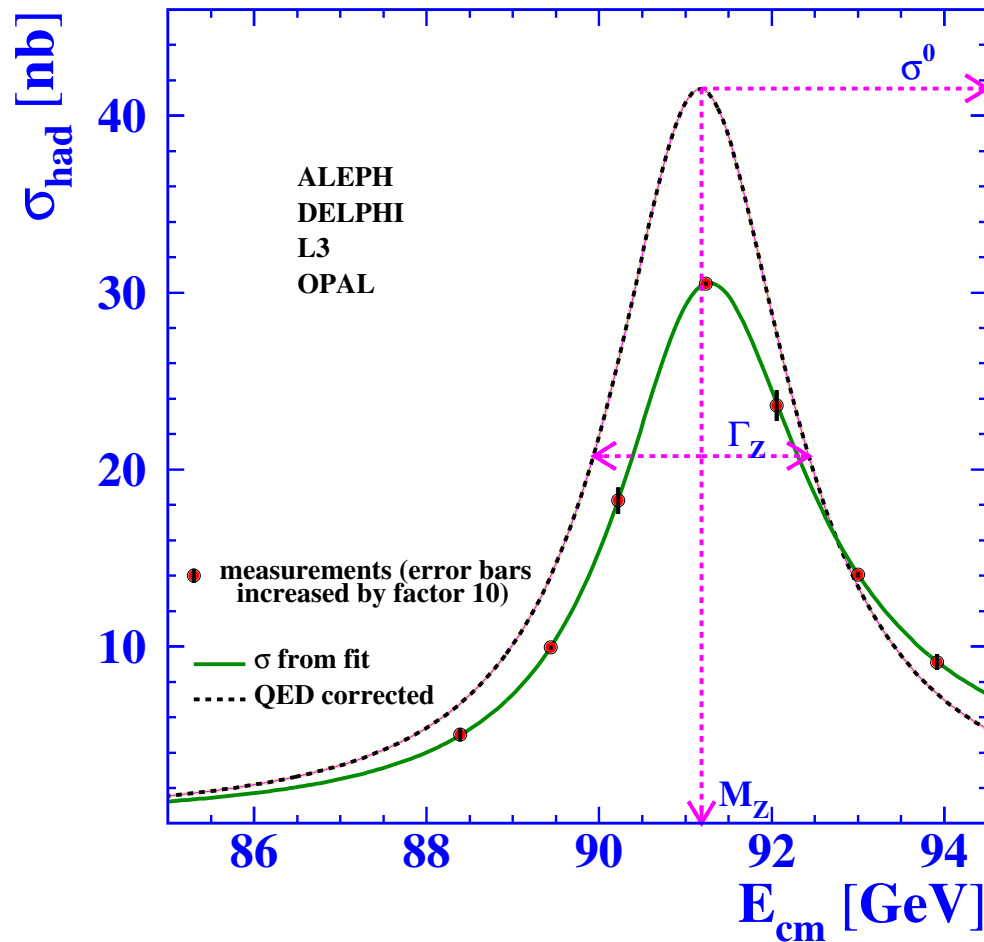
Average τ pol. in $e^+e^- \rightarrow \tau^+\tau^-$:
$$\langle \mathcal{P}_\tau \rangle = -A_\tau$$

Decay widths in terms of $\sin^2 \theta_{\text{eff}}^f$:

$$\Gamma_{ff} = C [F_V^f + F_A^f] = C [F_A^f (1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f)^2 + F_A^f]$$

- Comparison of EWPOs with SM to **probe new physics**
→ multi-loop corrections in full SM
- Extraction of EWPOs (**pseudo-observables**) from **real observables**
→ backgrounds (in full SM), QED/QCD, MC tools
- “Other” electroweak parameters (“**input**” parameters)
→ m_t , α_s , etc. extracted from other processes





LEP EWWG '05

- Large effects from initial-state QED radiation
- Theory input necessary to extract relevant EWPOs (“pseudo-observables”)

■ Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{ini}(s, s') \otimes \sigma_{hard}(s')$$

- Kureav, Fadin '85
- Berends, Burgers, v. Neerven '88
- Kniehl, Krawczyk, Kühn, Stuart '88
- Beenakker, Berends, v. Neerven '89
- Bardin et al. '91; Skrzypek '92
- Montagna, Nicrosini, Piccinini '97

Soft photons (resummed) + collinear photons

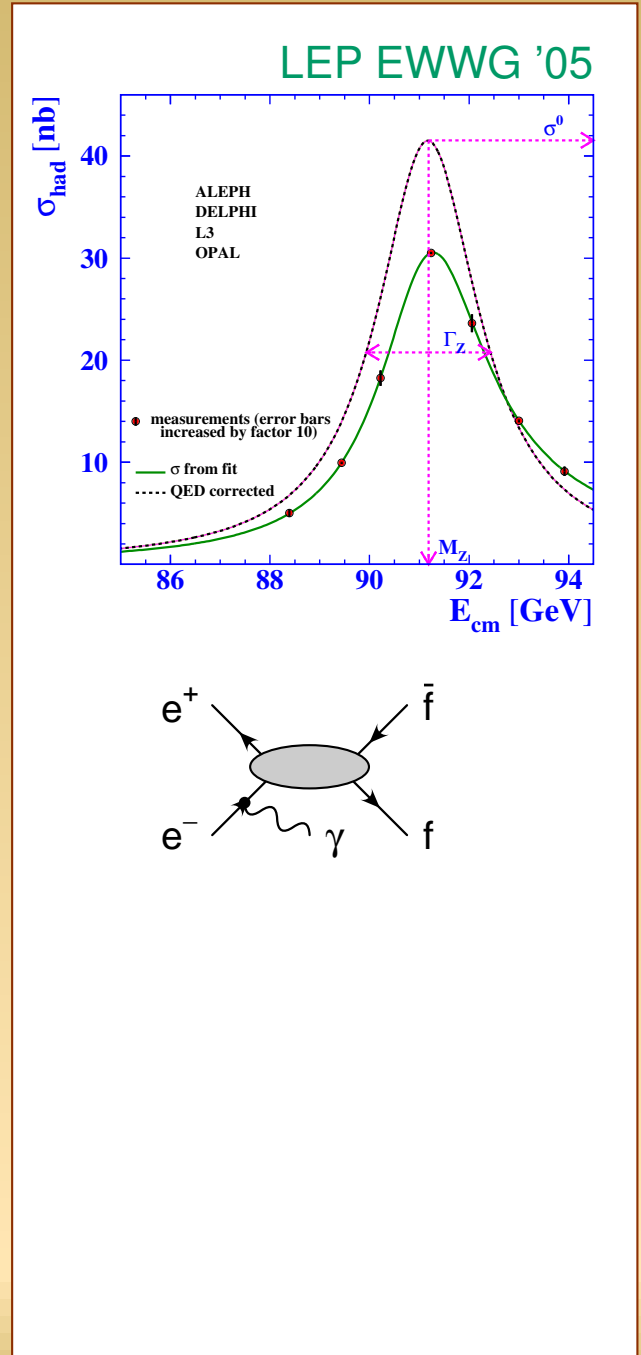
$$\mathcal{R}_{ini} = \sum_n \left(\frac{\alpha}{\pi}\right)^n \sum_{m=0}^n h_{nm} \ln^m\left(\frac{s}{m_e^2}\right)$$

Universal ($m=n$) logs known to $n = 6$,
also some sub-leading terms

Ablinger, Blümlein, De Freitas, Schönwald '20

Exclusive description: MC tools

→ talk by S. Frixione

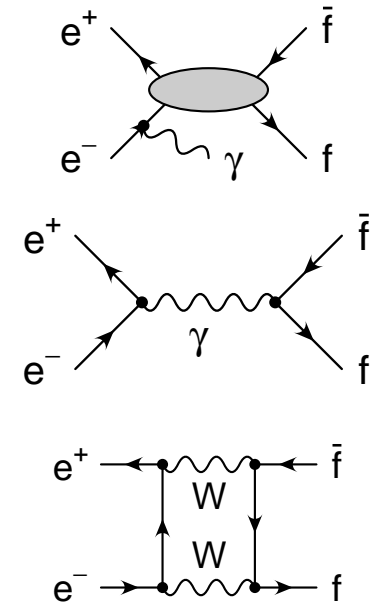
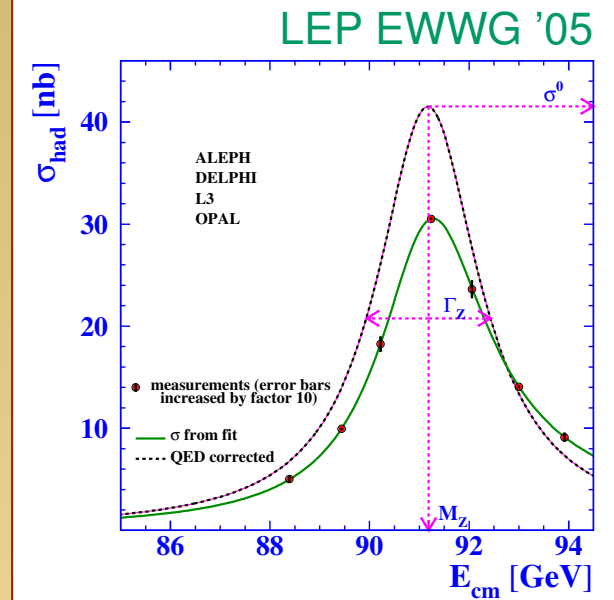


- Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

- Subtraction of γ -exchange, γ - Z interference, box contributions:

$$\sigma_{\text{hard}} = \sigma_Z + \sigma_\gamma + \sigma_{\gamma Z} + \sigma_{\text{box}}$$



- Deconvolution of initial-state QED radiation:

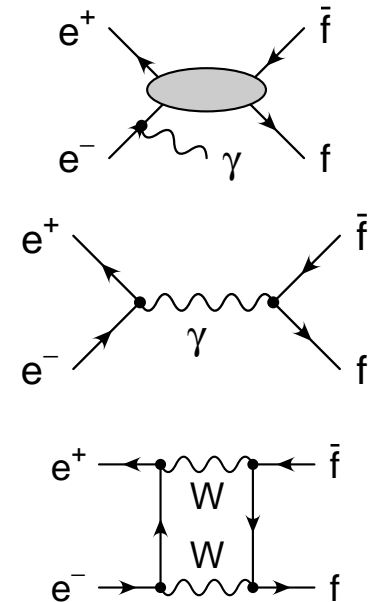
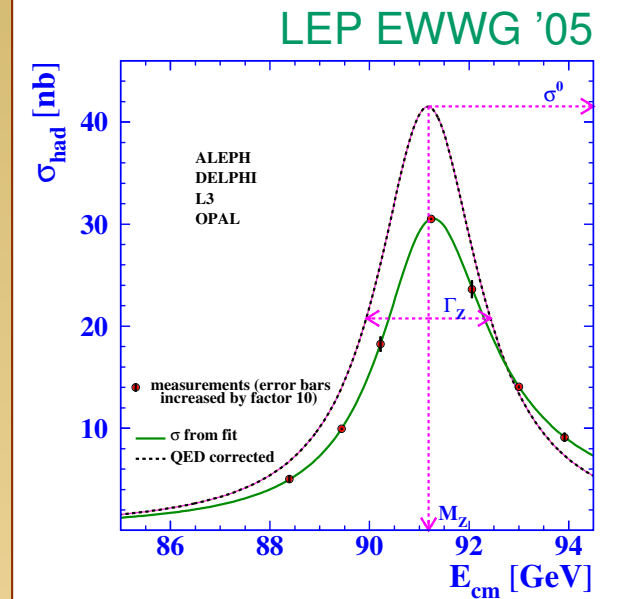
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$$\sigma_{\text{hard}} = \sigma_Z + \sigma_\gamma + \sigma_{\gamma Z} + \sigma_{\text{box}}$$

- Z-pole contribution:

$$\sigma_Z = \frac{R}{(s - \overline{M}_Z^2)^2 + \overline{M}_Z^2 \overline{\Gamma}_Z^2} + \sigma_{\text{non-res}}$$



- Deconvolution of initial-state QED radiation:

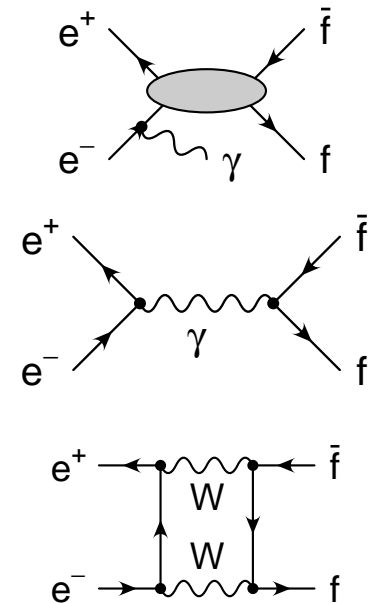
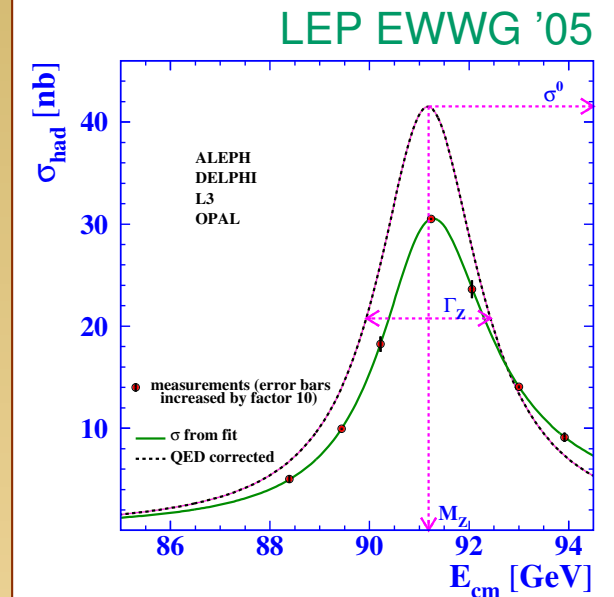
$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

- Subtraction of γ -exchange, γ -Z interference, box contributions:

$$\sigma_{\text{hard}} = \sigma_Z + \underbrace{\sigma_\gamma + \sigma_{\gamma Z} + \sigma_{\text{box}}}_{\text{computed in SM}}$$

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- Deconvolution of initial-state QED radiation:

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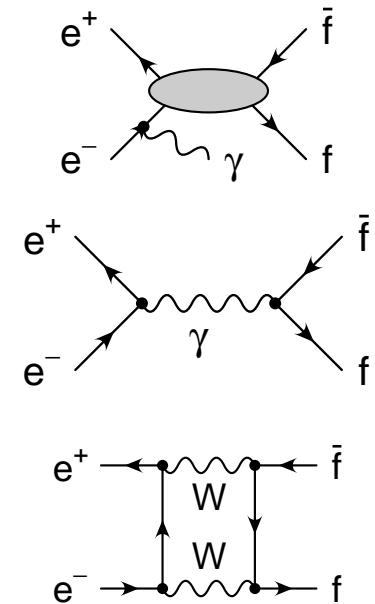
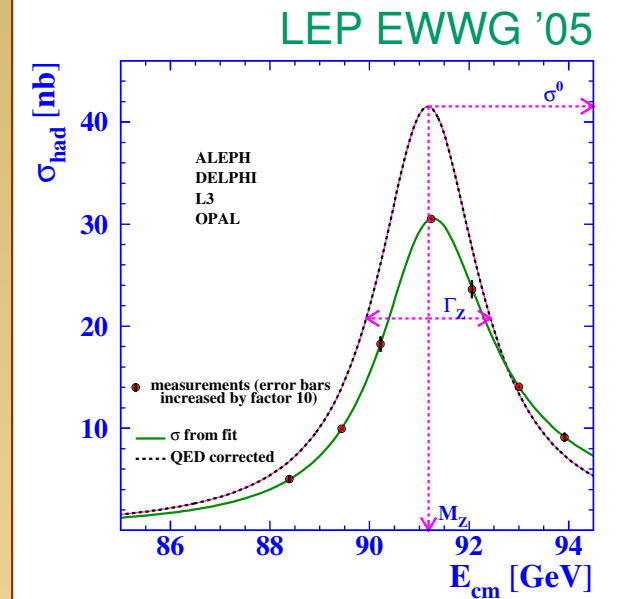
$$\sigma_Z = \frac{R}{(s - \overline{M}_Z^2)^2 + \overline{M}_Z^2 \overline{\Gamma}_Z^2} + \sigma_{\text{non-res}}$$

- In experimental analyses:

$$\sigma \sim \frac{1}{(s - M_Z^2)^2 + s^2 \Gamma_Z^2 / M_Z^2}$$

$$\overline{M}_Z = M_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx M_Z - 34 \text{ MeV}$$

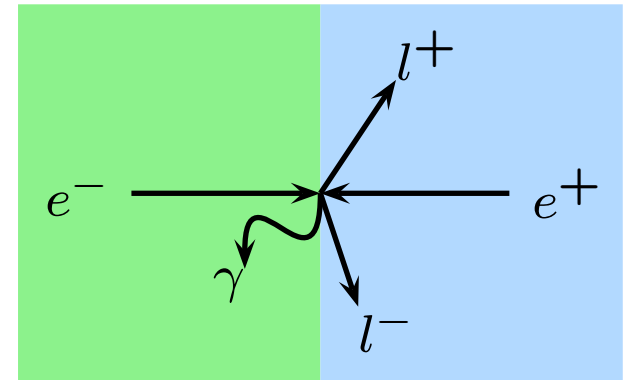
$$\overline{\Gamma}_Z = \Gamma_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx \Gamma_Z - 0.9 \text{ MeV}$$



QED radiation in principle cancels in asymmetries, e.g. $A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B}$

Some effects from detector acceptance and cuts

Typical influence $< 10^{-3}$



Implementation of QED effects:

a) Analytical formulae, e.g. ZFITTER

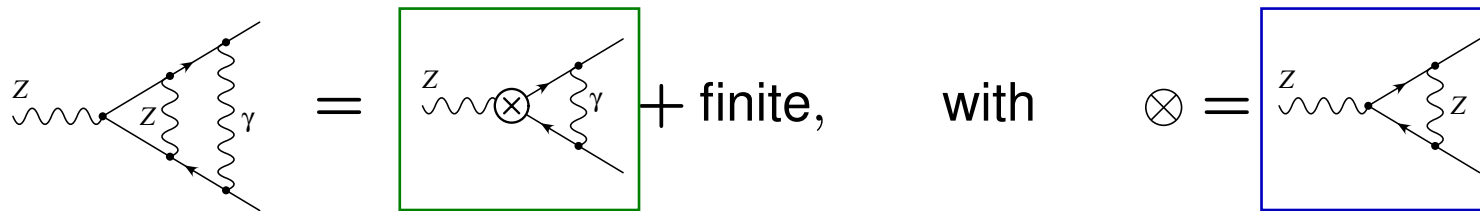
Arbuzov, Bardin, Christova, Kalinovskaya, Riemann, Riemann, ...

b) Monte Carlo event generator, e.g. KORALZ, KKMC

Jadach, Ward, ...

Factorization of massive and QED/QCD FSR:

$$\bar{\Gamma}_f \approx \frac{N_c \bar{M}_Z}{12\pi} \left[\mathcal{R}_V^f F_V^f + \mathcal{R}_A^f F_A^f \right]_{s=\bar{M}_Z^2}$$

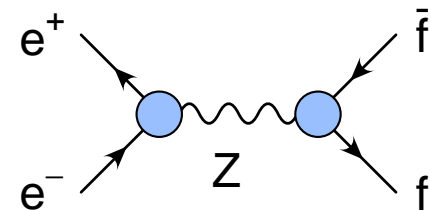


$\mathcal{R}_V^f, \mathcal{R}_A^f$: Final-state QED/QCD radiation;
 known inclusively to $\mathcal{O}(\alpha_s^4), \mathcal{O}(\alpha^2), \mathcal{O}(\alpha\alpha_s)$ Kataev '92
 Chetyrkin, Kühn, Kwiatkowski '96
 Baikov, Chetyrkin, Kühn, Ritinger '12

→ talk by S. Kluth

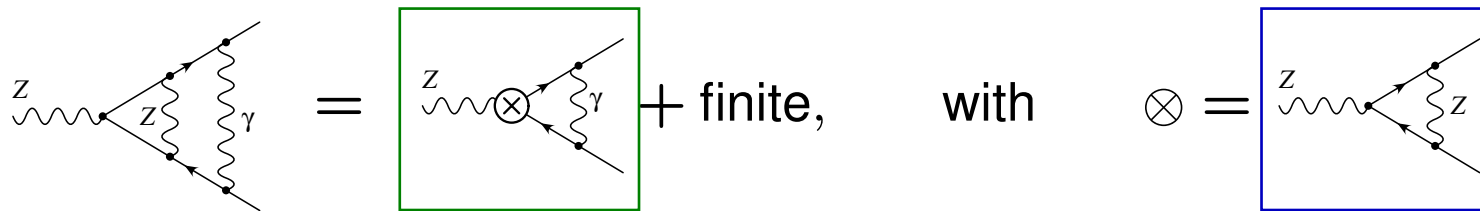
or compute exclusively using MC methods

F_V^f, F_A^f : Electroweak corrections

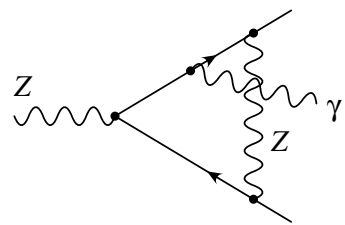


Factorization of massive and QED/QCD FSR:

$$\bar{\Gamma}_f \approx \frac{N_c \bar{M}_Z}{12\pi} \left[\mathcal{R}_V^f F_V^f + \mathcal{R}_A^f F_A^f \right]_{s=\bar{M}_Z^2}$$



Additional non-factorizable contributions, e.g.



→ Known at $\mathcal{O}(\alpha\alpha_s)$ Czarnecki, Kühn '96
Harlander, Seidensticker, Steinhauser '98

→ Currently not known at $\mathcal{O}(\alpha^2)$ and beyond

→ $\mathcal{O}(0.01\%)$ uncertainty on Γ_Z, σ_Z , maybe larger for A_b

→ How to account for in MC simulations?

- Implementation in MC program to evaluate exp. efficiency and particle ID

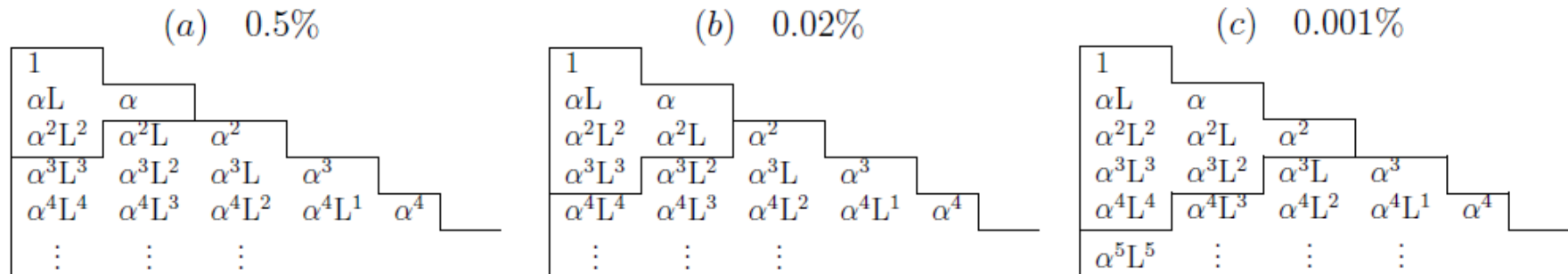
- Current state of art: e.g. KORALZ, KKMC
 $\rightarrow \mathcal{O}(\alpha^2 L)$ accuracy [$L = \ln(s/m_e^2)$]

Jadach, Ward, ...

- One to two orders improvement needed:

Observable	Where from	Present (LEP)	FCC stat.	FCC syst	$\frac{\text{Now}}{\text{FCC}}$
M_Z [MeV]	Z linesh. [28]	$91187.5 \pm 2.1\{0.3\}$	0.005	0.1	3
Γ_Z [MeV]	Z linesh. [28]	$2495.2 \pm 2.1\{0.2\}$	0.008	0.1	2
$R_i^Z = \Gamma_h/\Gamma_l$	$\sigma(M_Z)$ [33]	$20.767 \pm 0.025\{0.012\}$	$6 \cdot 10^{-5}$	$1 \cdot 10^{-3}$	12
σ_{had}^0 [nb]	σ_{had}^0 [28]	$41.541 \pm 0.037\{0.25\}$	$0.1 \cdot 10^{-3}$	$4 \cdot 10^{-3}$	6
N_ν	$\sigma(M_Z)$ [28]	$2.984 \pm 0.008\{0.006\}$	$5 \cdot 10^{-6}$	$1 \cdot 10^{-3}$	6
$\sin^2 \theta_W^{eff} \times 10^5$	$A_{FB}^{lept.}$ [33]	$23099 \pm 53\{28\}$	0.3	0.5	55
$A_{FB,\mu}^{M_Z \pm 3.5\text{GeV}}$	$\frac{d\sigma}{d\cos\theta}$ [28]	$\pm 0.020\{0.001\}$	$1.0 \cdot 10^{-5}$	$0.3 \cdot 10^{-5}$	100

Jadach,
Skrzypek '19



- \rightarrow Need matching of h.o. matrix elements with QED parton shower
 (exclusive in all fs particles)

Expand amplitude for $e^+e^- \rightarrow f\bar{f}$ about **complex pole** $s_0 \equiv \overline{M_Z^2} + i\overline{M_Z}\overline{\Gamma_Z}$:

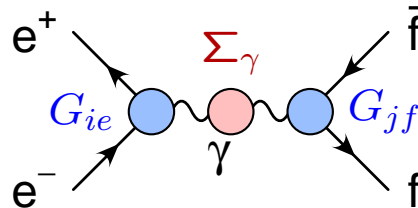
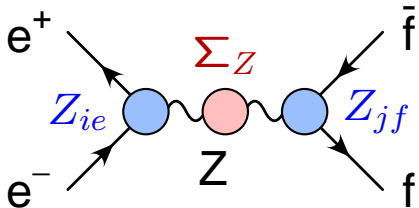
→ All terms are individually gauge-invariant

$$\mathcal{M}_{ij} = \frac{R_{ij}}{s - s_0} + S_{ij} + (s - s_0)S'_{ij} + \dots \quad (i, j = V, A)$$

$$R_{ij} = \left. \frac{Z_{ie}Z_{jf}}{1 + \Sigma'_Z} \right|_{s=s_0}$$

$$S_{ij} = \left[\frac{Z_{ie}Z'_{jf} + Z'_{ie}Z_{jf}}{1 + \Sigma'_Z} - \frac{Z_{ie}Z_{jf}\Sigma''_Z}{2(1 + \Sigma'_Z)^2} + \frac{G_{ie}G_{jf}}{s + \Sigma_\gamma} + B_{ij} \right]_{s=s_0}$$

$$S'_{ij} = \dots$$



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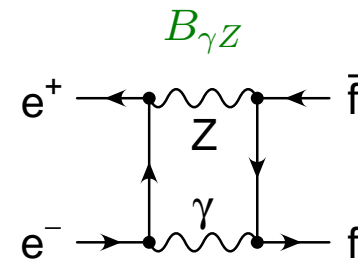
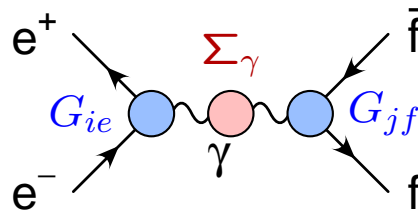
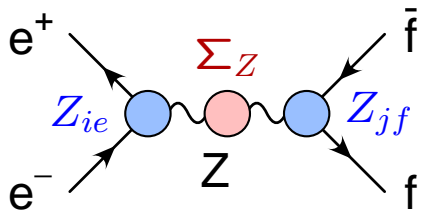
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$$R_{ij} = \left. \frac{Z_{ie}Z_{jf}}{1 + \Sigma'_Z} \right|_{s=s_0} + B_{\gamma Z, ij}^R + B_{\gamma Z, ij}^{RL} \ln\left(1 - \frac{s}{s_0}\right)$$

$$S_{ij} = \left[\frac{Z_{ie}Z'_{jf} + Z'_{ie}Z_{jf}}{1 + \Sigma'_Z} - \frac{Z_{ie}Z_{jf}\Sigma''_Z}{2(1 + \Sigma'_Z)^2} + \frac{G_{ie}G_{jf}}{s + \Sigma_\gamma} + B_{ij} \right]_{s=s_0} + B_{\gamma Z, ij}^S + B_{\gamma Z, ij}^{SL} \ln\left(1 - \frac{s}{s_0}\right)$$

$$S'_{ij} = \dots$$



Express R_{ij} in terms of $\sin^2 \theta_{\text{eff}}^f$ and F_A^f (with NNLO corrections):

$$\begin{aligned}
 R_{ij} = & 4I_e^3 I_f^3 \sqrt{F_A^e F_A^f} \left[Q_i^e Q_j^f \left(1 + i r_{AA}^I - \frac{1}{2} (r_{AA}^I)^2 + \frac{1}{2} \delta \bar{X}_{(2)} \right) \right. \\
 & \left. + (Q_i^e I_{j,f} + Q_j^f I_{i,e}) (i - r_{AA}^I) - I_{i,e} I_{j,f} \right] \\
 & + M_Z \Gamma_Z Z_{ie(0)} Z_{jf(0)} x_{ij}^I,
 \end{aligned}$$

$$Q_V^f = 1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f, \quad Q_A^f = 1$$

$$I_{V,f} = \frac{1}{(a_{f(0)}^Z)^2} \left[a_{f(0)}^Z \text{Im} Z_{Vf(1)} - v_{f(0)}^Z \text{Im} Z_{Af(1)} \right], \quad I_{A,f} = 0$$

$$\delta \bar{X}_{(2)} = -(\text{Im} \Sigma'_{Z(1)})^2 + 2 \bar{b}_{\gamma Z(1)}^R,$$

$$r_{ij}^I = \frac{\text{Im} Z_{ie(1)}}{Z_{ie(0)}} + \frac{\text{Im} Z_{jf(1)}}{Z_{jf(0)}} - \text{Im} \Sigma'_{Z(1)},$$

$$x_{ij}^I = \frac{\text{Im} Z'_{ie(1)}}{Z_{ie(0)}} + \frac{\text{Im} Z'_{jf(1)}}{Z_{jf(0)}} - \frac{1}{2} \text{Im} \Sigma''_{Z(1)},$$

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$$\mathcal{M}_{ij} = \frac{R_{ij}}{s - s_0} + S_{ij} + (s - s_0)S'_{ij} + \dots \quad (i, j = V, A)$$

Current state of art: R @ NNLO + leading higher orders

S @ NLO

S' @ (N)LO

For future ee colliders: (at least) one order more!

→ also matching to Monte-Carlo for QED/QCD ISR/FSR/IFI

- M_Z, Γ_Z : From $\sigma(\sqrt{s})$ lineshape; $\delta M_Z, \delta \Gamma_Z \sim 0.1$ MeV at FCC-ee
→ Main theory uncertainties: QED ISR
- m_t : Current status $\delta m_t \sim 0.3$ GeV at LHC PDG '20
→ Theoretical ambiguity in mass def. of $\mathcal{O}(0.5$ GeV)
Hoang, Plätzer, Samitz '18; Ferrario Ravasio, Nason, Oleari '18
→ Robust det. from $e^+e^- \rightarrow t\bar{t}$ at $\sqrt{s} \sim 350$ GeV → talk by M. Beneke
- m_b, m_c : From quarkonia spectra using Lattice QCD
 $\delta m_b^{\overline{\text{MS}}} \sim 30$ MeV, $\delta m_c^{\overline{\text{MS}}} \sim 25$ MeV LHC HXSWG '16
→ estimated improvements $\delta m_b^{\overline{\text{MS}}} \sim 13$ MeV, $\delta m_c^{\overline{\text{MS}}} \sim 7$ MeV
Lepage, Mackenzie, Peskin '14
- M_H : from kinematic constraint fits $HZ(\ell\ell), H(b\bar{b})Z$
→ $\delta M_H \sim 10\dots 20$ MeV
→ theory errors subdominant

- α_S :

d'Enterria, Skands, et al. '15

- Most precise determination using Lattice QCD:

$\alpha_S = 0.1184 \pm 0.0006$ HPQCD '10

$\alpha_S = 0.1185 \pm 0.0008$ ALPHA '17

$\alpha_S = 0.1179 \pm 0.0015$ Takaura et al. '18

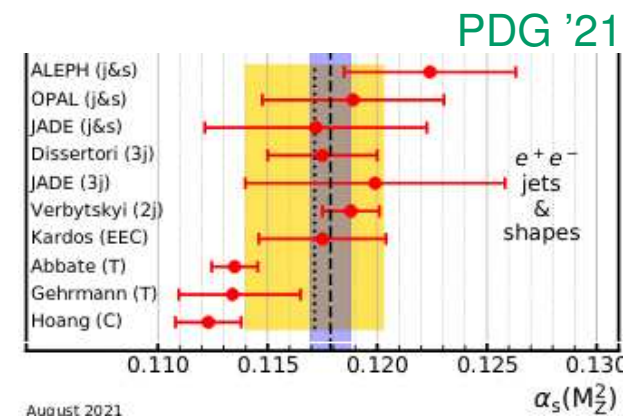
$\alpha_S = 0.1172 \pm 0.0011$ Zafeiropoulos et al. '19

→ Difficulty in evaluating systematics

- e^+e^- event shapes: $\alpha_S \sim 0.113...0.119$

→ Large non-perturbative power corrections

→ Systematic uncertainties?



- Hadronic τ decays: $\alpha_S = 0.119 \pm 0.002$

PDG '18

→ Non-perturbative uncertainties in OPE and from duality violation

Pich '14; Boito et al. '15,18

- α_s :

- Electroweak precision ($R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$):

$$\alpha_s = 0.120 \pm 0.003$$

PDG '18

→ No (negligible) non-perturbative QCD effects

$$\text{FCC-ee: } \delta R_\ell \sim 0.001$$

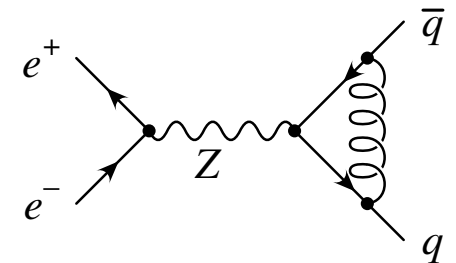
$$\Rightarrow \delta \alpha_s < 0.0001$$

Theory input: **N³LO EW corr. + leading N⁴LO**

to keep $\delta_{\text{th}} R_\ell \lesssim \delta_{\text{exp}} R_\ell$

Caviat: R_ℓ could be affected by new physics

→ talk by S. Kluth



- α_s :

- Electroweak precision ($R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$):

$$\alpha_s = 0.120 \pm 0.003$$

PDG '18

→ No (negligible) non-perturbative QCD effects

FCC-ee: $\delta R_\ell \sim 0.001$

⇒ $\delta \alpha_s < 0.0001$

Caveat: R_ℓ could be affected by new physics

- $R = \frac{\sigma[ee \rightarrow \text{had.}]}{\sigma[ee \rightarrow \mu\mu]}$ at lower \sqrt{s}

e.g. CLEO ($\sqrt{s} \sim 9 \text{ GeV}$): $\alpha_s = 0.110 \pm 0.015$

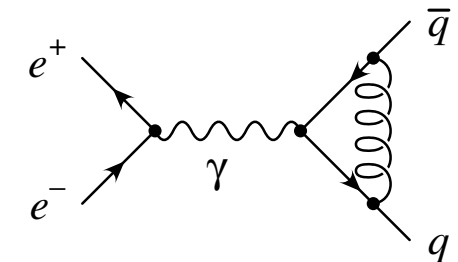
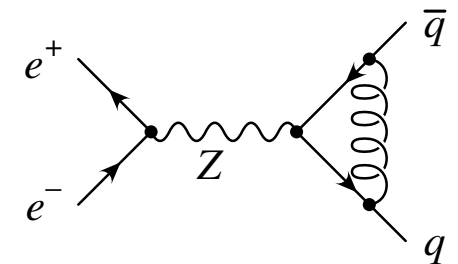
Kühn, Steinhauser, Teubner '07

→ dominated by s -channel photon, less room for new physics

→ QCD still perturbative

naive scaling to 50 ab^{-1} (BELLE-II): $\delta \alpha_s \sim 0.0001$

→ talk by S. Kluth



- $\Delta\alpha \equiv 1 - \frac{\alpha(0)}{\alpha(M_Z)} \approx 0.059 = 0.0315_{\text{lept}} + 0.0276_{\text{had}}$

a) $\Delta\alpha_{\text{had}}$ from $e^+e^- \rightarrow \text{had.}$
using dispersion relation



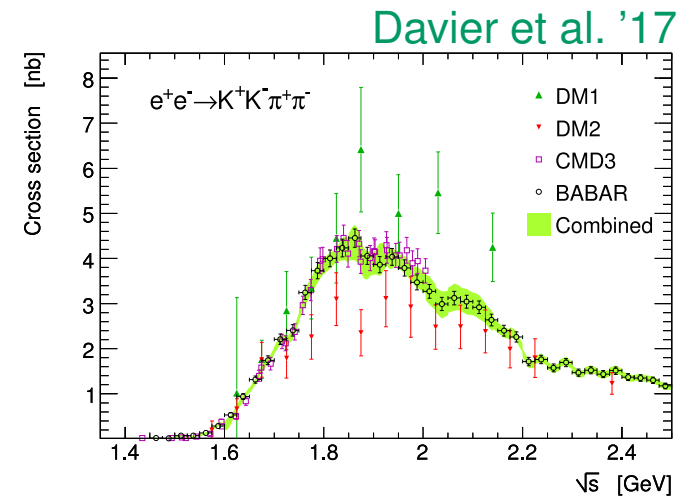
New data from BaBar, VEPP, BES, KLOE

→ Robust precision $\sim 10^{-4}$

Davier et al. '19; Jegerlehner '19; Keshavarzi, Nomura, Teubner '19

Improvement to $\delta(\Delta\alpha_{\text{had}}) \sim 5 \times 10^{-5}$ likely

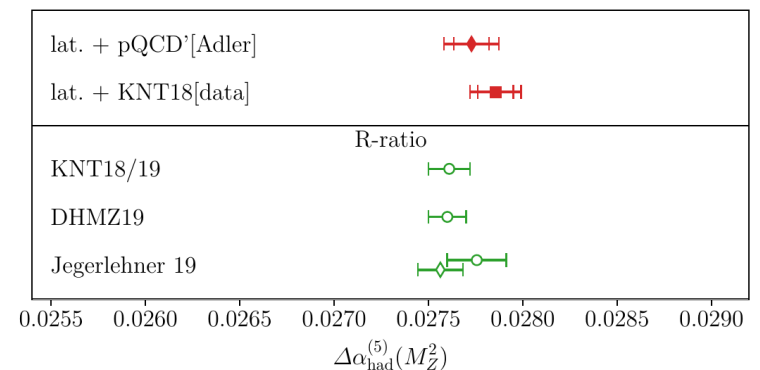
Jegerlehner '19



b) $\Delta\alpha_{\text{had}}$ from Lattice QCD
(challenging but much progress)

Burger et al. '15

Cè et al. '22



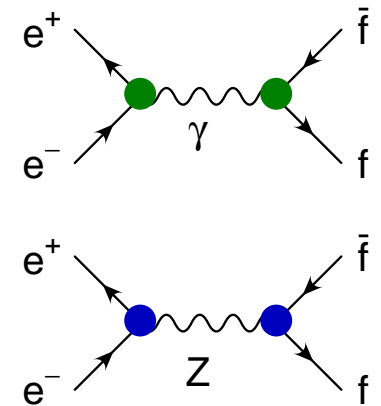
- $\Delta\alpha \equiv 1 - \frac{\alpha(0)}{\alpha(M_Z)} \approx 0.059 = 0.0315_{\text{lept}} + 0.0276_{\text{had}}$

c) Direct det. of $\Delta\alpha_{\text{had}}$ from $e^+e^- \rightarrow \mu^+\mu^-$ off the Z peak

Janot '15

$$|\mathcal{M}_{ij}|^2 \propto |g_i^\ell|^2 |g_j^\ell|^2 + (s - M_Z^2) \alpha(M_Z) |g_{i,j}^\ell|^2 + \dots$$

↑
determined
from Z pole



- $\Delta\alpha \equiv 1 - \frac{\alpha(0)}{\alpha(M_Z)} \approx 0.059 = 0.0315_{\text{lept}} + 0.0276_{\text{had}}$

c) Direct det. of $\Delta\alpha_{\text{had}}$ from $e^+e^- \rightarrow \mu^+\mu^-$ off the Z peak

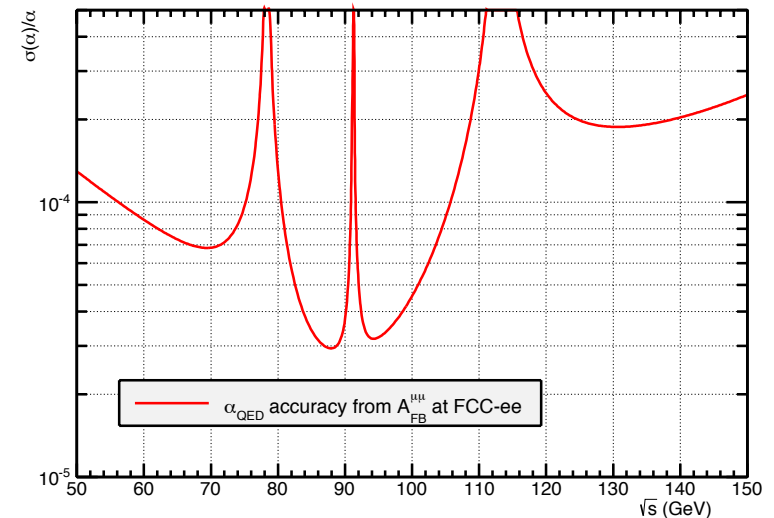
Janot '15

$$|\mathcal{M}_{ij}|^2 \propto |g_i^\ell|^2 |g_j^\ell|^2 + (s - M_Z^2) \alpha(M_Z) |g_{i,j}^\ell|^2 + \dots$$

→ Use $A_{\text{FB}}^{\mu\mu}$ and two cms energies to reduce systematics

→ Sensitivity maximized for $\sqrt{s_1} \sim 88 \text{ GeV}$, $\sqrt{s_2} \sim 95 \text{ GeV}$

→ $\delta(\Delta\alpha_{\text{had}}) \sim 3 \times 10^{-5}$
for $\mathcal{L}_{\text{int}} = 85 \text{ ab}^{-1}$



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Requires theory input:
2-/3-loop corrections for
 $e^+e^- \rightarrow \mu^+\mu^-$

- To probe new physics, compare EWPOs with SM theory predictions

	Current exp.	Current th.	CEPC exp.	FCC-ee exp.
M_W^* [MeV]	15	4	0.5	0.4
Γ_Z [MeV]	2.3	0.4	0.025	0.025
$R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$ [10^{-3}]	25	5	2	1
$R_b = \Gamma_Z^b / \Gamma_Z^{\text{had}}$ [10^{-5}]	66	10	4.3	6
$\sin^2 \theta_{\text{eff}}^\ell$ [10^{-5}]	16	4.5	<1	0.5

* computed from G_μ

- Theory error estimate is not well defined, ideally $\Delta_{\text{th}} \ll \Delta_{\text{exp}}$
- Common methods:
 - Count prefactors (α , N_c , N_f , ...)
 - Extrapolation of perturbative series
 - Renormalization scale dependence
 - Renormalization scheme dependence

- Many seminal works on 1-loop and leading 2-loop corrections

Veltman, Passarino, Sirlin, Marciano, Bardin, Hollik, Riemann, Degrassi, Kniehl, ...

- Full 2-loop results for M_W , Z -pole observables

Freitas, Hollik, Walter, Weiglein '00

Awramik, Czakon '02

Onishchenko, Veretin '02

Awramik, Czakon, Freitas, Weiglein '04

Awramik, Czakon, Freitas '06

Hollik, Meier, Uccirati '05,07

Awramik, Czakon, Freitas, Kniehl '08

Freitas '14

Dubovyk, Freitas, Gluza, Riemann, Usovitsch '16,18

- Approximate 3- and 4-loop results (enhanced by Y_t and/or N_f)

Chetyrkin, Kühn, Steinhauser '95

Faisst, Kühn, Seidensticker, Veretin '03

Boughezal, Tausk, v. d. Bij '05

Schröder, Steinhauser '05

Chetyrkin et al. '06

Boughezal, Czakon '06

Chen, Freitas '20

Freitas, Heinemeyer, et al. '19

	CEPC	FCC-ee	perturb. error with 3-loop [†]	Param. error	
				scen. 1*	scen. 2*
M_W [MeV]	0.5	0.4	1	2.1	0.6
Γ_Z [MeV]	0.025	0.025	0.15	0.15	0.1
R_b [10^{-5}]	4.3	6	5	< 1	< 1
$\sin^2 \theta_{\text{eff}}^\ell$ [10^{-5}]	< 1	0.5	1.5	2	1

[†] **Theory scenario:** $\mathcal{O}(\alpha\alpha_s^2)$, $\mathcal{O}(N_f\alpha^2\alpha_s)$, $\mathcal{O}(N_f^2\alpha^2\alpha_s)$, leading 4-loop
 (N_f^n = at least n closed fermion loops)

Parametric inputs:

***Scenario 1:** $\delta m_t = 600$ MeV, $\delta\alpha_s = 0.0002$, $\delta M_Z = 0.5$ MeV,
 $\delta(\Delta\alpha) = 5 \times 10^{-5}$

***Scenario 2:** $\delta m_t = 50$ MeV, $\delta\alpha_s = 0.0002$, $\delta M_Z = 0.5$ MeV,
 $\delta(\Delta\alpha) = 3 \times 10^{-5}$

Analytical techniques:

- Computational intensive reduction to master integrals (MIs)
- Not fully understood function space of MIs
- Works best for problems with few (no) masses

Numerical techniques:

- Large computing time
- Numerical instabilities, in particular for diagrams with physical cuts
- Works best for problems with many masses

New techniques, e.g.:

- Numerical reduction to MIs, numerical MIs via differential equations (DEs)
Mandal, Zhao '18, Czakon, Niggetiedt '20
- DEs with respect to auxiliary parameter, $\frac{1}{k_i^2 - m_i^2 + i\epsilon}$
Liu, Ma, Wang '17
Liu, Ma '18,21,22
- Series solutions of DEs
Moriello '19, Hidding '20
- Dispersion relations + Feynman parameters
Song, Freitas '21

→ next week's talks

Extension of SM by **higher-dimensional operators**:

Wilson '69
Weinberg '79

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_{d=5}^{\infty} \frac{1}{\Lambda^{d-4}} \sum_i c_i \mathcal{O}_i^{(d)}$$

- $\mathcal{O}_i^{(d)}$ depend on all SM fields (including Higgs doublet)
- Operators must satisfy SM gauge invariance
- Valid description for energies $E \ll \Lambda$ ($\Lambda \sim$ mass of heavy particles)
- Operators ranked by suppression power Λ^{4-d}

Review: 1706.08945

HEFT: similar to SMEFT, but Higgs and Goldstone bosons treated independently

↑
gauge singlet ↑
SU(2) triplet

Example:

$$\text{HEFT: } \frac{c_1}{\Lambda} h(\bar{\psi}\not{D}\psi), \quad \frac{c_2}{\Lambda^2} h^2(\bar{\psi}\not{D}\psi), \quad \dots$$

$$\text{SMEFT: } \frac{c_{\phi\psi}}{\Lambda^2} (\phi^\dagger\phi)(\bar{\psi}\not{D}\psi) \quad \rightarrow \quad \frac{c_{\phi\psi}}{2\Lambda^2} (v+h)^2(\bar{\psi}\not{D}\psi)$$

$$\Rightarrow \quad c_1 = c_{\phi\psi}v, \quad c_2 = \frac{1}{2}c_{\phi\psi}$$

→ Fewer operators, less freedom in SMEFT

→ HEFT needed to describe some composite-Higgs theories

Leading dim-6 contribution: $\mathcal{L} = \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \mathcal{O}(\Lambda^{-3}) \quad (\Lambda \gg M_Z)$

$$\mathcal{O}_{\phi 1} = \frac{1}{4} |\Phi^\dagger \overleftrightarrow{D}_\mu \Phi|^2$$

$$\alpha \Delta T = -\frac{v^2}{2} \frac{c_{\phi 1}}{\Lambda^2}$$

$$\mathcal{O}_{\text{BW}} = \Phi^\dagger B_{\mu\nu} W^{\mu\nu} \Phi$$

$$\alpha \Delta S = -e^2 v^2 \frac{c_{\text{BW}}}{\Lambda^2}$$

$$O_{\text{LL}}^{(3)\mu e} = (\bar{L}_L^\mu \sigma^a \gamma_\mu L_L^\mu) (\bar{L}_L^e \sigma^a \gamma^\mu L_L^e) \quad \Delta G_F = -\sqrt{2} \frac{c_{\text{LL}}^{(3)\mu e}}{\Lambda^2}$$

$$O_{\text{R}}^f = i(\Phi^\dagger \overleftrightarrow{D}_\mu \Phi) (\bar{f}_R \gamma^\mu f_R)$$

$$f = e, \mu, \tau, b, lq$$

$$O_{\text{L}}^F = i(\Phi^\dagger \overleftrightarrow{D}_\mu \Phi) (\bar{F}_L \gamma^\mu F_L)$$

$$F = \begin{pmatrix} \nu_e \\ e \end{pmatrix}, \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}, \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}, \begin{pmatrix} u, c \\ d, s \end{pmatrix}, \begin{pmatrix} t \\ b \end{pmatrix}$$

$$O_{\text{L}}^{(3)F} = i(\Phi^\dagger \overleftrightarrow{D}_\mu^a \Phi) (\bar{F}_L \sigma_a \gamma^\mu F_L)$$

More operators than EWPOs \rightarrow need assumptions:

- Family universality, e.g. $c_{\text{R}}^e = c_{\text{R}}^\mu = c_{\text{R}}^\tau$
- $U(2) \times U(1)$ flavor symmetry, e.g. $c_{\text{R}}^e = c_{\text{R}}^\mu \neq c_{\text{R}}^\tau$

Leading Z pole contribution:

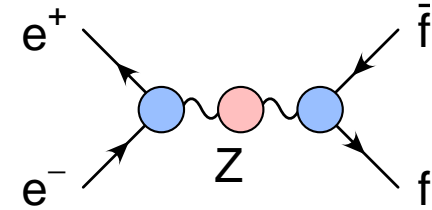
Modified propagators:

e.g. $\mathcal{O}_{\phi 1} = (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi)$

$$\mathcal{O}_{\text{BW}} = \Phi^\dagger B_{\mu\nu} W^{\mu\nu} \Phi$$

Modified Z-fermion couplings:

e.g. $\mathcal{O}^f = i(\Phi^\dagger \overleftrightarrow{D}_\mu \Phi)(\bar{f}\gamma^\mu f) \quad f = e, \mu, \tau, b, \dots$



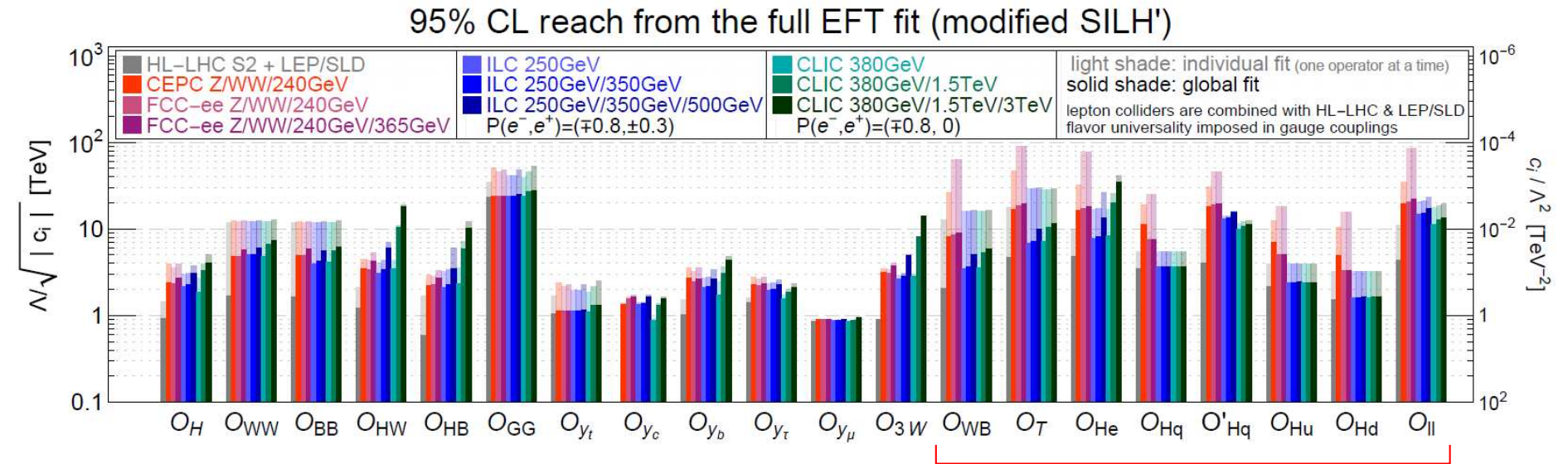
Direct correspondence between pseudo-observables and SMEFT ops.:

$$\frac{\delta \Gamma_{ff}}{\Gamma_{ff}} = 2 \left[\frac{\delta Z_L^f}{Z_L^f} + \frac{\delta Z_R^f}{Z_R^f} \right] \quad \frac{\delta \sin^2 \theta_{\text{eff}}^f}{\sin^2 \theta_{\text{eff}}^f} = \frac{1}{1 - Z_L^f/Z_R^f} \left[\frac{\delta Z_L^f}{Z_L^f} - \frac{\delta Z_R^f}{Z_R^f} \right]$$

$$\delta Z_X^f = \frac{1}{\Lambda^2} \left[I_3^{fX} c_L^{(3)f} - \frac{1}{2} c_X^f - Q_f \frac{sc}{c^2 - s^2} c_{WB} - \left(\frac{1}{16} c_{\phi 1} - \frac{1}{4} c_{LL}^{(3)} + c_L^\ell \right) \left(I_3^{fX} + Q_f \frac{s^2}{c^2 - s^2} \right) \right]$$

$$\frac{\delta M_W}{M_W} = \delta Z_L^\nu - \delta Z_L^e - \frac{1}{4} c_{LL}^{(3)}$$

Project reach of future colliders to constrain SMEFT operators:

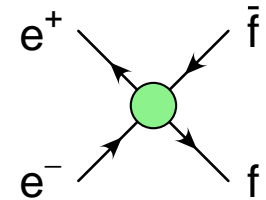


de Blas, Durieux, Grojean, Gu, Paul '19

Beyond leading Z pole (“background”):

Additional 4-fermion operators

e.g. $O_{ff'}^{(1)} = (\bar{f}\gamma_\mu f)(\bar{f}'\gamma^\mu f')$ $f, f' = e, \mu, \tau, b, \dots$



(related via e.o.m. to energy-dependent $Zff/\gamma ff$ couplings $O_f^D \sim \bar{f}\gamma^\mu F_{\mu\nu} D^\nu f$)

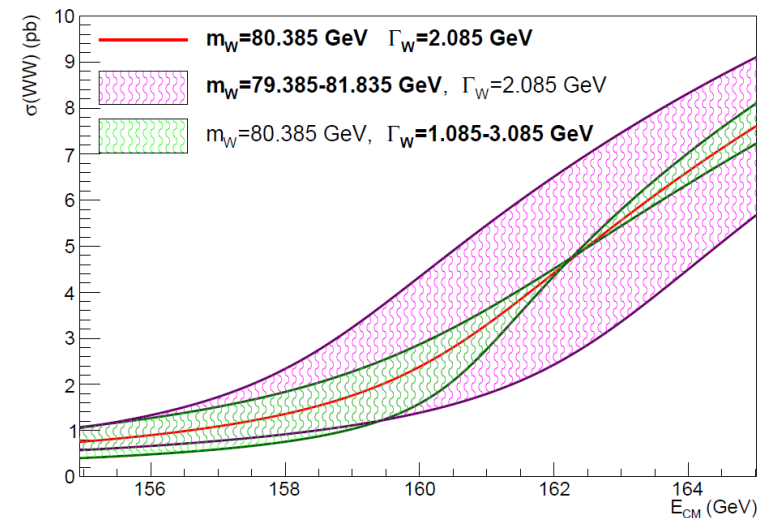
→ Model-independent analysis requires additional parameters (or Wilson coeffs.) in non-resonant terms

- High-precision measurement of M_W from $e^+e^- \rightarrow W^+W^-$ at threshold

- a) Corrections near threshold enhanced by $1/\beta$ and $\ln \beta$

$$\beta \sim \sqrt{1 - 4 \frac{M_W^2 - iM_W \Gamma_W}{s}} \sim \sqrt{\Gamma_W / M_W}$$

- b) Non-resonant contributions are important

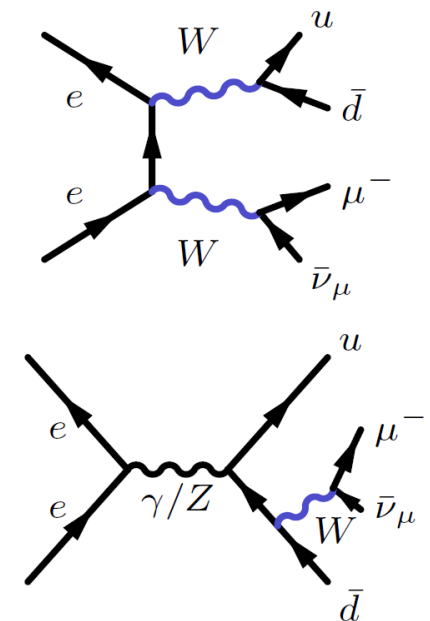


- Full $\mathcal{O}(\alpha)$ calculation of $e^+e^- \rightarrow 4f$
Denner, Dittmaier, Roth, Wieders '05

- EFT expansion in $\alpha \sim \Gamma_W / M_W \sim \beta^2$
Beneke, Falgari, Schwinn, Signer, Zanderighi '07

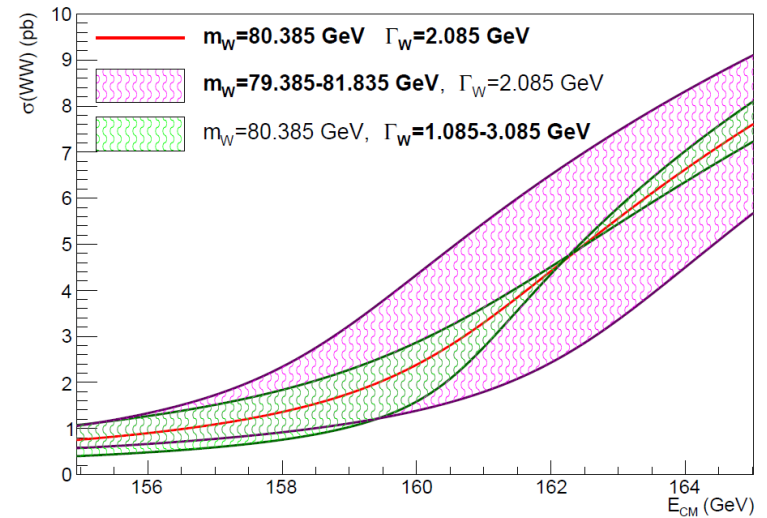
- NLO corrections with NNLO Coulomb correction ($\propto 1/\beta^n$): $\delta_{\text{th}} M_W \sim 3 \text{ MeV}$
Actis, Beneke, Falgari, Schwinn '08

- Adding NNLO corrections to $ee \rightarrow WW$ and $W \rightarrow f\bar{f}$ and NNLO ISR: $\delta_{\text{th}} M_W \lesssim 0.6 \text{ MeV}$

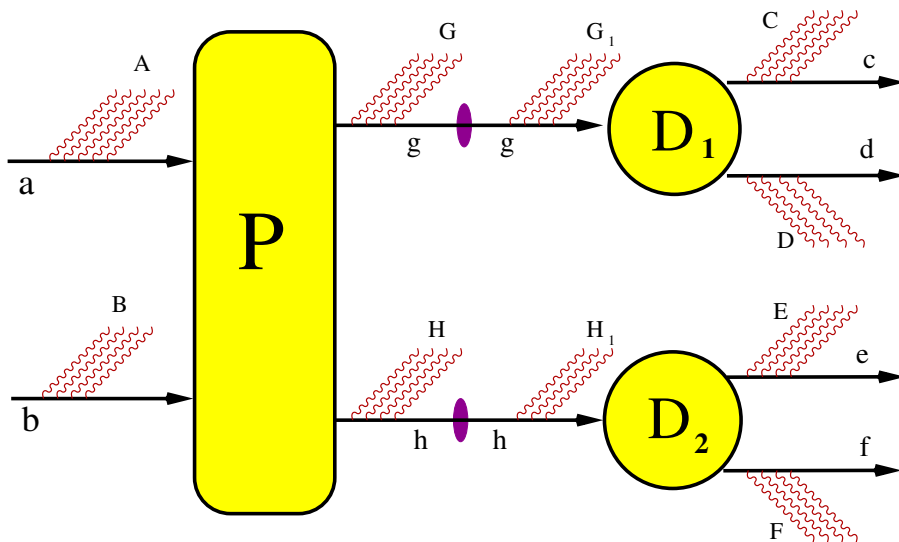


- High-precision measurement of M_W from $e^+e^- \rightarrow W^+W^-$ at threshold
- a) Corrections near threshold enhanced by $1/\beta$ and $\ln \beta$

$$\beta \sim \sqrt{1 - 4 \frac{M_W^2 - iM_W\Gamma_W}{s}} \sim \sqrt{\Gamma_W/M_W}$$
- b) Non-resonant contributions are important



- Resummation of soft photon radiation



Jadach, Płaczek, Skrzypek '19

- Electroweak precision tests require theory input for **measurements of pseudo-observables** (BRs, widths, masses, cross-sections, ...) and their **SM/BSM interpretation**
- **Future e^+e^- colliders** improve precision by 1–2 orders of magnitude
- Theory progress needed both for **fixed-order loop corrections** as well as **MC tools**
- Model-independent description of BSM effects through **EWPOs** (for leading Z-pole term) and **SMEFT operators**
- Accurate description of non-resonant terms (backgrounds) requires a **gauge-invariant framework, SM corrections**, and accounting of **BSM effects**

Backup slides

Theory calculations: Uncertainties

	Experiment	Theory error	Main source
M_W	80.379 ± 0.012 MeV	4 MeV	$\alpha^3, \alpha^2\alpha_s$
Γ_Z	2495.2 ± 2.3 MeV	0.4 MeV	$\alpha^3, \alpha^2\alpha_s, \alpha\alpha_s^2$
R_ℓ	20.767 ± 0.025	0.005	$\alpha^3, \alpha^2\alpha_s$
R_b	0.21629 ± 0.00066	0.0001	$\alpha^3, \alpha^2\alpha_s$
$\sin^2 \theta_{\text{eff}}^\ell$	0.23153 ± 0.00016	4.5×10^{-5}	$\alpha^3, \alpha^2\alpha_s$

- Theory error estimate is not well defined, ideally $\Delta_{\text{th}} \ll \Delta_{\text{exp}}$
- Common methods:
 - Count prefactors (α, N_c, N_f, \dots)
 - Extrapolation of perturbative series
 - Renormalization scale dependence
 - Renormalization scheme dependence

Example: Error estimation for Γ_Z

■ Geometric perturbative series

$$\alpha_t = \alpha m_t^2$$

$$\mathcal{O}(\alpha^3) - \mathcal{O}(\alpha_t^3) \sim \frac{\mathcal{O}(\alpha^2) - \mathcal{O}(\alpha_t^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha^2) \sim 0.26 \text{ MeV}$$

$$\mathcal{O}(\alpha^2 \alpha_s) - \mathcal{O}(\alpha_t^2 \alpha_s) \sim \frac{\mathcal{O}(\alpha^2) - \mathcal{O}(\alpha_t^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s) \sim 0.30 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^2) - \mathcal{O}(\alpha_t \alpha_s^2) \sim \frac{\mathcal{O}(\alpha \alpha_s) - \mathcal{O}(\alpha_t \alpha_s)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s) \sim 0.23 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^3) - \mathcal{O}(\alpha_t \alpha_s^3) \sim \frac{\mathcal{O}(\alpha \alpha_s) - \mathcal{O}(\alpha_t \alpha_s)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s^2) \sim 0.035 \text{ MeV}$$

$$\mathcal{O}(\alpha_{\text{bos}}^2) \sim \mathcal{O}(\alpha_{\text{bos}})^2 \sim 0.1 \text{ MeV}$$

■ Parametric prefactors:

$$\mathcal{O}(\alpha_{\text{bos}}^2) \sim \Gamma_Z \alpha^2 \sim 0.1 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^2) - \mathcal{O}(\alpha_t \alpha_s^2) \sim \frac{\alpha n_{|q}}{\pi} \alpha_s^2 \sim 0.29 \text{ MeV}$$

Total: $\delta\Gamma_Z \approx 0.5 \text{ MeV}$