Modern calculation techniques for multi-scale loop amplitudes

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Precision calculations for future e^+e^- colliders: targets and tools, \mbox{CERN}







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- 1. Introduction
- 2. Loops & legs: state of the art

3. Recent developments on the multi-scale frontier

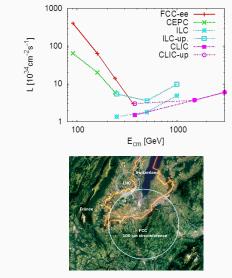
4. Towards $e^+e^- \rightarrow 4j$ @ NNLO QCD

5. Discussion

Introduction

Precision at future e^+e^- colliders

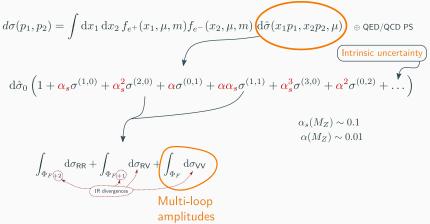
[1910.11775]



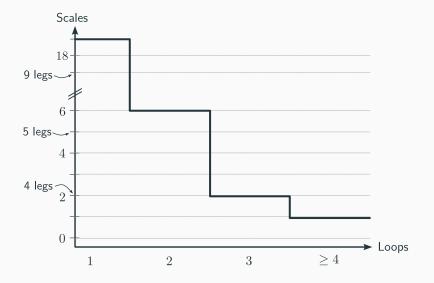
- Measure EW & Higgs observables, α_s to unprecedented precision
- Discovery via precision: search anomalous deviations from SM
- Sub-percent uncertainties
- Theoretical input crucial

Fixed order partonic cross sections

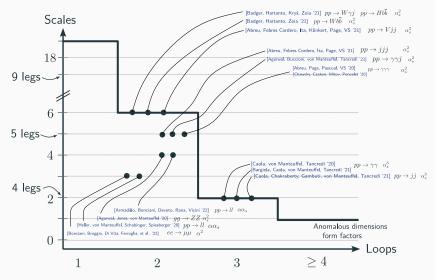
Collinear factorization:



Loops & legs: state of the art

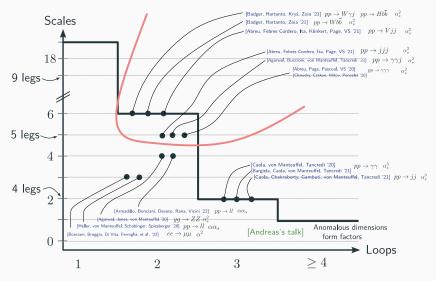


Loops&legs: state of the art



Warning: a biased selection of references!

Loops&legs: state of the art



Warning: a biased selection of references!

Dynamic scales

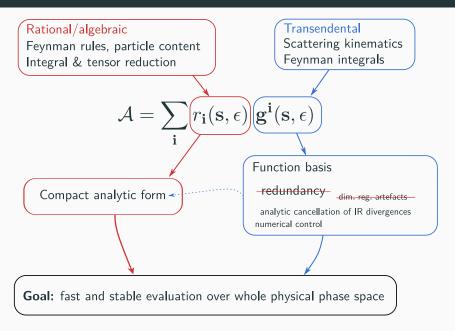
- Mandelstam invariants s_{ij} , off-shell legs p_i^2
- Monte Carlo integrals over phase space $\int d\Phi_n \left(s_{ij}, p_i^2\right) |\mathcal{A}_{2 \to n}(s_{ij}, p_i^2)|^2$
- Need fast and robust numerical evaluation of $\mathcal{A}_{2 \rightarrow n}$ over phase space

Fixed scales

- Particle (complex) masses, e.g. m_t, m_W
- Mathematical complexity can escalate very quickly [see Stefan's talk]
- With few dynamic scales can profit the most from numerical methods [see talks by Vitalii, Janusz, Martijn, Xiao]

Recent developments on the multi-scale frontier

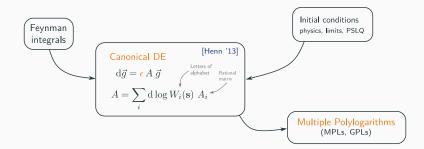
Structure of analytic loop amplitudes



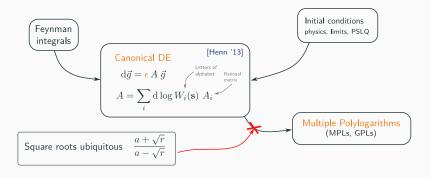
Recent developments on the multi-scale frontier

Transcendental part

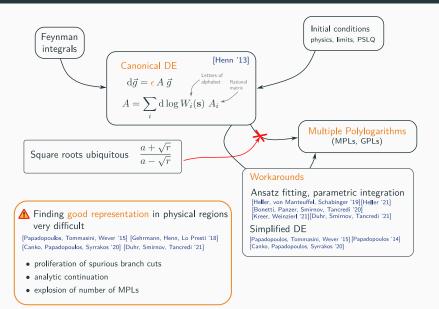
Feynman integrals: the canonical way



Feynman integrals: the canonical way



Feynman integrals: the canonical way



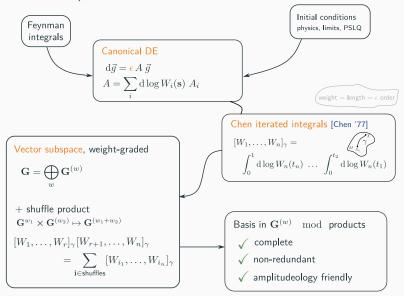
Special function basis construction

[Chicherin, VS, Zoia '21] (see also [Chicherin, VS '20] [Badger, Hartanto, Zoia '21])



Special function basis construction

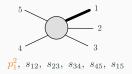
[Chicherin, VS, Zoia '21] (see also [Chicherin, VS '20] [Badger, Hartanto, Zoia '21])



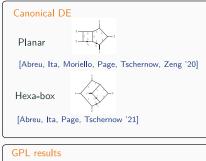
Example: two-loop five-point one-mass integrals

One-mass kinematics

e.g. $pp \rightarrow Vjj, \, e^+e^- \rightarrow 4j$



Example: two-loop five-point one-mass integrals



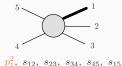
[Papadopoulos, Tommasini, Wever '15] [Canko, Papadopoulos, Syrrakos '20] [Syrrakos '20] [Kardos, Papadopoulos, Smirnov, Syrrakos, Wever '22]

Function basis (planar)

[Badger, Hartanto, Zoia '21] color-ordered, numerical evaluation

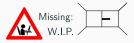
[Chicherin, VS, Zoia '21]

One-mass kinematics e.g. $pp \rightarrow Vjj$, $e^+e^- \rightarrow 4j$

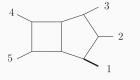


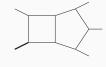
Semi-numerical DE solution

DiffExp [Moriello '19] [Hidding '20] AMFLow [Liu, Ma, Wang '17] [Liu, Ma '21] Initial values, validation, small scale sampling

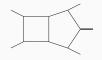


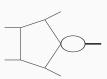
Example: planar function basis

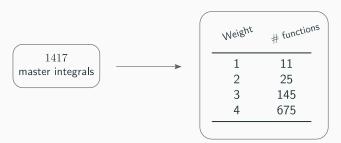




+ 24 permutations $\{p_2, p_3, p_4, p_5\}$

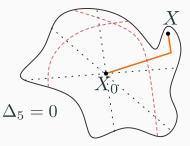






Weights 1 and 2

Well-defined combinations of $\log,\,\mathrm{Li}_2$ functions

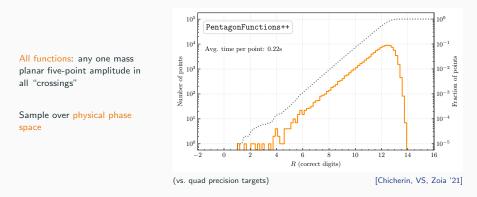


Weights 3 and 4

- Numerical one-fold integration [Caron-Huot, Henn '14] of analytic integrands

 exponential convergence [Takahasi, Mori '73]
- No crossing of physical thresholds \implies no analytic continuation needed
- Dedicated series expansions around spurious singularities

Numerical performance



Available as a C++ library PentagonFunctions++

https://gitlab.com/pentagon-functions/PentagonFunctions-cpp

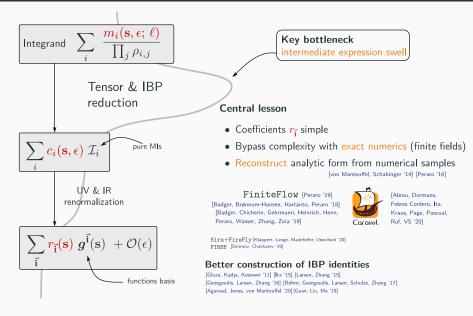
(also Mathematica interface)

So far the only method to get to $2 \rightarrow 3$ cross sections!

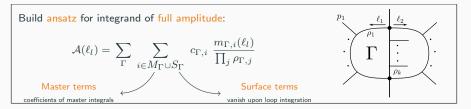
[Czakon, Mitov, Poncelet '21] [Chen, Gehrmann, Glover, Huss, Marcoli '21] [Chawdry, Czakon, Mitov, Poncelet '21] [Kallweit, VS, Wiesemann '20] [Badger, Gehrmann, Marcoli, Moodie '21] [Hartanto, Poncelet, Popescu, Zoia '22] Recent developments on the multi-scale frontier

Rational coefficients

Rational coefficients: algebraic complexity



Numerical unitarity: universal ansatz



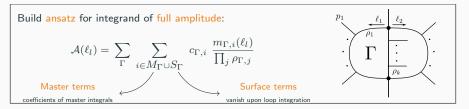
Generalization of one-loop unitarity methods:

[Bern, Dixon, Kosower, Dunbar '94, '95] [Britto, Feng, Cachazo '05] [Ossola, Papadopoulos, Pittau '07] [Ellis, Giele, Kunszt '08] [Giele, Kunszt, Melnikov '08]

Related work

[Badger, Frellesvig, Zhang '12] [Zhang '12] [Mastrolia, Mirabella, Ossola, Peraro '13]
 [Ita '15] [Mastrolia, Peraro, Primo '16] [Abreu, Febres Cordero, Ita, Jaquier, Page, Zeng '17]

Numerical unitarity: universal ansatz



Surface terms produced from unitarity-compatible IBPs [Gluza, Kadja, Kosower '11]:

$$\int \left(\prod_{l} \mathrm{d}^{D} \ell_{l}\right) \sum_{l} \frac{\partial}{\partial \ell_{l}^{\nu}} \left(\frac{\boldsymbol{u}_{l}^{\nu} m(\ell_{l})}{\prod_{j \in P_{\Gamma}} \rho_{j}}\right) = 0, \qquad \boldsymbol{u}_{l}^{\nu} \frac{\partial}{\partial \ell_{l}^{\nu}} \rho_{j} = f_{j} \rho_{j}$$

Generating vectors u_l^{ν} from computational algebraic geometry (e.g. with Singular) [Ita '15] [Larsen, Zhang '15] [Georgoudis, Larsen, Zhang '16] [Abreu, Febres Cordero, Ita, Page, Zeng '17] [Böhm, Georgoudis, Larsen, Schulze, Zhang '17]

- Targeted set of identities for each Γ
- Eliminate linear dependencies with on-shell conditions $\rho_i = 0, \ j \in P_{\Gamma}$ imposed
- No need to invert IBP systems at this stage

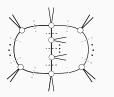
Numerical unitarity: cut equations

Obtain coefficients $c_{\Gamma,i}$ from linear systems of cut equations:

$$\lim_{\substack{\ell_l \to \ell_l^{\Gamma}}} \left(\mathcal{A}(\ell_l) \prod_j \rho_j \right) = \sum_i c_{\Gamma,i} m_{\Gamma,i} + \frac{\text{topologies with more propagators}}{(\text{from previous steps})}$$

on-shell limit for $\boldsymbol{\Gamma}$





Solve cut equations numerically

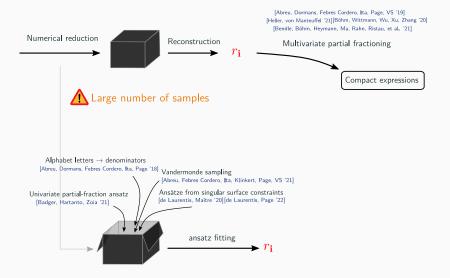
- ✓ analytic integrand or individual Feynman diagrams not needed
- ✓ numerical IBP reduction included
- $\checkmark\,$ suitable for floating point and finite fields

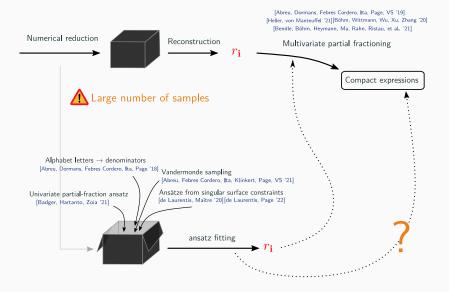
Implemented in C++ library CARAVEL [Abreu, Dormans, Febres Cordero, Ita, Kraus, Page, Pascual, Ruf, VS '20]











Towards $e^+e^- \to 4j$ @ NNLO QCD

Two-loop amplitudes for Vjj production at hadron colliders

$$pp \to Vjj, \quad V = W, Z/\gamma^*$$

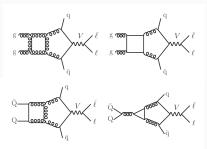
 $\downarrow_{} \ell \bar{\ell}$

Helicity amplitudes

 $\mathcal{M}\left(\bar{q}_{p_{1}}^{+}, g_{p_{2}}^{h_{2}}, g_{p_{3}}^{h_{3}}, q_{p_{4}}^{-}; \bar{\ell}_{p_{5}}^{+}, \ell_{p_{6}}^{-}\right)$ $\mathcal{M}\left(\bar{q}_{p_{1}}^{+}, Q_{p_{2}}^{h}, \bar{Q}_{p_{3}}^{-h}, q_{p_{4}}^{-}; \bar{\ell}_{p_{5}}^{+}, \ell_{p_{6}}^{-}\right)$

• $N_c
ightarrow \infty$, $N_f \sim N_c$, 5FNS, no top loops

[Abreu, Febres Cordero, Ita, Klinkert, Page, VS '21]



$$\mathcal{M}_{\kappa} = \mathcal{M}_{\kappa}^{(0)} \left(1 + \frac{\alpha_s}{2\pi} \tilde{\mathcal{M}}_{\kappa}^{(1)} + \left(\frac{\alpha_s}{2\pi}\right)^2 \tilde{\mathcal{M}}_{\kappa}^{(2)} + \dots \right), \quad \text{for } \kappa = g, Q$$

$$\begin{split} \mathcal{M}_{\mathbf{g}}^{(k)} &\xrightarrow{\text{L.C.}} \left(\frac{N_c}{2}\right)^k \sum_{\sigma \in S_2} \left(T^{a_{\sigma(3)}} T^{a_{\sigma(2)}}\right)_{i_4}^{\tilde{i}_1} \sum_{j=0}^k \left(\frac{N_f}{N_c}\right)^j \mathcal{A}_{\mathbf{g}}^{(k)[j]} \\ \mathcal{M}_q^{(k)} &\xrightarrow{\text{L.C.}} \left(\frac{N_c}{2}\right)^k \delta_{i_2}^{\tilde{i}_1} \delta_{i_4}^{\tilde{i}_3} \sum_{j=0}^k \left(\frac{N_f}{N_c}\right)^j \mathcal{A}_q^{(k)[j]} \end{split}$$

Numerical benchmark [Hartanto, Badger, Brønnum-Hansen, Peraro '19]

Analytic reconstruction

Numerical reduction (in \mathbb{F}_p) to finite remainders expressed in basis of pentagon functions [Chicherin, VS, Zoia '21]

$$\mathcal{R}^{(2)} = \mathcal{A}^{(2)} - \mathbf{I}^{(1)} \mathcal{A}^{(1)} - \mathbf{I}^{(2)} \mathcal{A}^{(0)} = \sum_{\vec{i}} r_{\vec{i}}(\vec{s}) g^{\vec{i}}$$

with CARAVEL.

Black-box reconstruction too hard: would require 10^7 samples, each few minutes!

• Denominators are letters of the alphabet,

$$r_i = \frac{n_i}{\prod_{j=1}^{37} W_j^{q_{ij}}}, \qquad \left(\text{recall:} \quad \mathrm{d}\vec{g} = \epsilon \left(\sum_i \mathrm{d}\log W_i \, A_i \right) \vec{g} \right)$$

 \implies exponents q_{ij} from univariate reconstruction

- Reconstruct in one variable $s_{23} \to$ partial fraction \to dense ansatz for polynomials in remaining variables
- Maximal ansatz size 500k, exploit special structure of Vandermonde matrix to invert in $\mathcal{O}(N^2)$

Analytic results available at http://www.hep.fsu.edu/~ffebres/W4partons

From $pp \rightarrow Vjj$ to $e^+e^- \rightarrow 4j$

Analytic results for 6-point helicity amplitudes in the form

$$\mathcal{R} = \sum_{\vec{i}} r_{\vec{i}}(\vec{s}) \ \boldsymbol{g}^{\vec{i}}(\vec{s})$$

alternatively we know the vector current

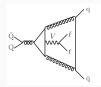
$$\mathcal{R}^{\mu} = \sum_{\vec{\mathbf{i}}} r^{\mu}_{\vec{\mathbf{i}}}(\vec{s}) \ \boldsymbol{g}^{\vec{\mathbf{i}}}(\vec{s}), \qquad \mathcal{R}^{\mu} J_{\mu}, \quad J_{\mu} = \frac{\kappa}{s_{56} - M_v^2 + i \,\Gamma_v M_v} \bar{u}(p_6) \gamma_{\mu} v(p_5)$$

Crossing from scattering to decay kinematics

- Attach appropriate lepton current; permutations, charge/parity conjugation of $r_{\vec{i}}^{\mu}(\vec{s})$ trivial
- Special function basis constructed for scattering kinematics $2 \rightarrow 3$. Need to redo the construction [Chicherin, VS, Zoia '21] for decay kinematics $1 \rightarrow 4$.

Nonplanar contributions

- All subleading in ${\cal N}_c$ not included, likely suppressed
- Contribution proportional to $\sum_{f} Q_{f}$ not small?



Summary

- Loop amplitudes remain major bottleneck in precision of theoretical predictions
- $\bullet\,$ Future e^+e^- colliders will benefit from advances in calculation techniques for the LHC
- Two-loop QCD corrections for $2\to 3$ processes, mixed QCDxEW corrections for $2\to 2$ are becoming a reality

Main lessons

- Multiscale calculations require new techniques and ideas
- $\bullet\,$ Large final-state phase space $\implies\,$ loop amplitudes must evaluate fast
- Good grasp on analytic structure of Feynman integrals and associated function spaces has been essential
- Judicious combination of analytic and numerical methods fruitful

Discussion

Analytic vs numerical methods: rational

Analytic reconstruction from samples in \mathbb{F}_p

- Best bet so far: framework for educated guess work, we hope to get much better at it!
- Also useful for deriving DEs for Feynman integrals
- Can we handle even more scales comfortably?
- Special function basis needed

Fully numerical?

- Numerical (floating point) amplitude reduction is in principle possible right now (e.g. Caravel)
- But: slow due to severe numerical instabilities in integral reduction, cancellations between Feynman integrals \implies high intermediate precision needed
- Can developments by [Lang, Pozzorini, Schär, Zhang, Zoller] (OpenLoops@2loops) help with this? [see Max's talk]

Analytic

- $\checkmark\,$ MPLs successful for $2\rightarrow2,$ also with masses, when possible
- $\checkmark~{\rm For}~2\to3$ up to two loops "pentagon functions" method, when possible
- ? Biggest issue: general class of functions not understood, even mathematically [see Stefan's talk]

(Semi-)numerical

Solutions of DEs by matching local series expansions [see talks by Martijn, Xiao]

- \checkmark Successfully sidestep analytic complexity with few dynamic scales
- $\checkmark\,$ Even very difficult multi-scale integrals integrals can be tackled [Liu, Ma '21]
- ? Still too slow for many dynamic scales?

To the shopping list

Basically there

• $e^+e^- \rightarrow 4j$, α_s^2 (also as a 4f background)

Still difficult, but imaginable with incremental progress

- $e^+e^- \rightarrow 3j$, α_s^3 (do we need $\alpha_s \alpha$?)
- $e^+e^- \rightarrow 5j$, α_s^2
- $e^+e^- \rightarrow QQj$, α_s^2 , massive b

Terrifying

 α^2 corrections to five-point functions

- $e^+e^- \rightarrow \nu\nu H$
- $\bullet ~ e^+e^- \rightarrow l^+l^-H$

A break-through in understanding of elliptic Feynman integrals might be required

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