

**Quarkonium production: probing QCD at the LHC,
Vienna, April 17-21, 2011**

Quarkonium, QCD and AdS/CFT

D. Kharzeev



Outline

- **The beginnings**

From atoms (V century B.C.) to quarks;
atoms, **void and geometry**

- **The successes**

Symmetry; QCD as the gauge theory of strong interactions;
asymptotic freedom; hard processes; strong color fields

- **The challenges**

Confinement; chiral symmetry breaking; extreme QCD:
temperature, density, energy; the spin; **void and geometry**

- **The future**

**From particles to fields: collective phenomena
as the essence of QCD; quarkonium as a probe**

The beginnings

Everything consists of

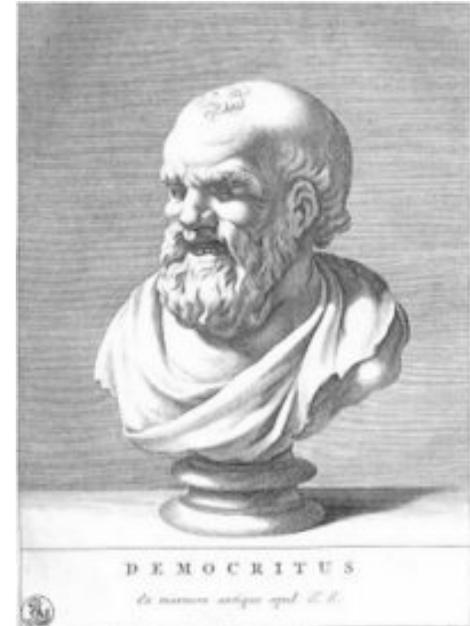
Atoms and Void

ἄτομος un - cuttable, indivisible

α- τέμνω



Λεύκιππος
Leucippus, V B.C.



Δημόκριτος
Democritus, ca 460 -370 B.C.

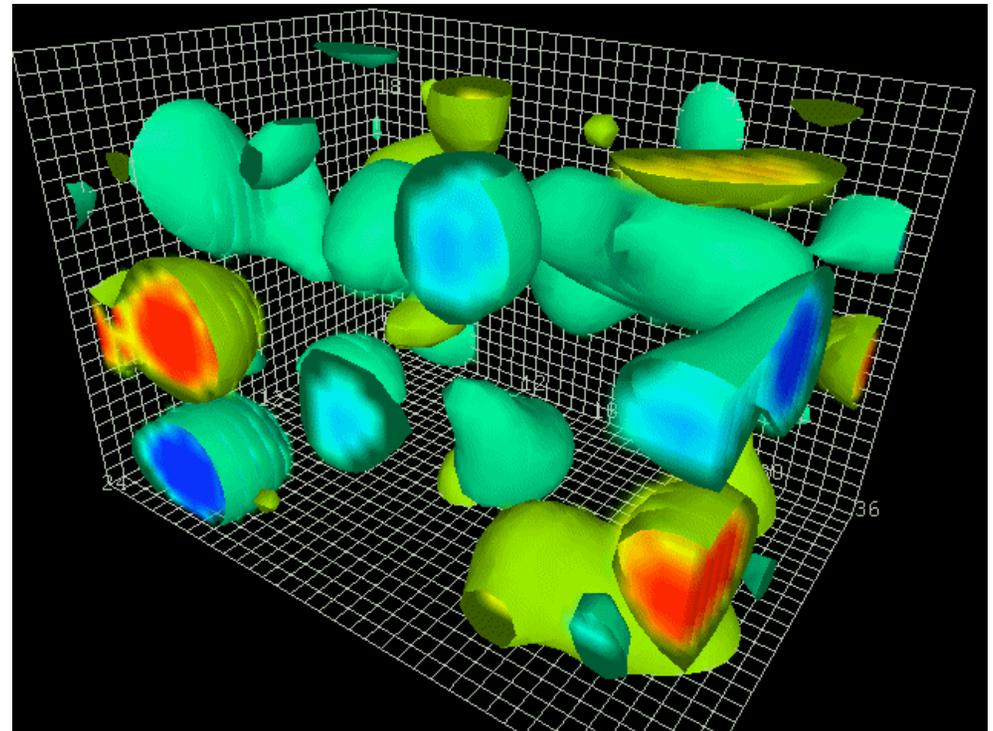
Everything consists of

Atoms and Void

“Atoms”,
ca 2011 A.D.

Void,
ca 2011 A.D.

| | | | |
|---|------------------------|-------------------------------|---------------------------------|
| Leptons | ν_e e- Neutrino | ν_μ μ - Neutrino | ν_τ τ - Neutrino |
| | e electron | μ muon | τ tau |
| Quarks | u up | c charm | t top |
| | d down | s strange | b bottom |
| I II III The Generations of Matter | | | |

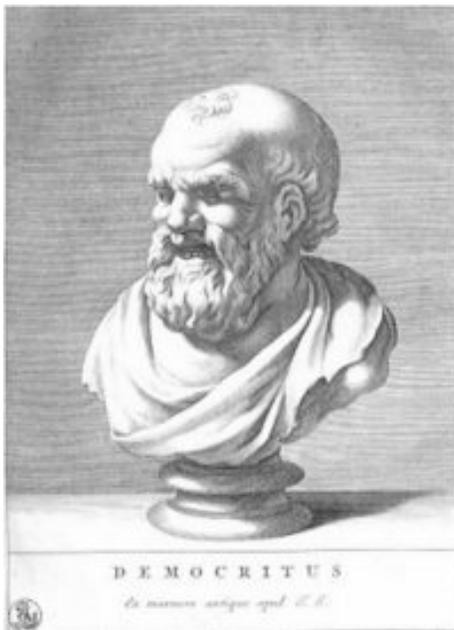




Leucippus, V B.C.

Atomism vs corpuscularianism:

are quarks and leptons
the ultimate indivisible
“atoms” of Nature?



Democritus, ca 460 -370 B.C.



Robert Boyle (1627–1691)

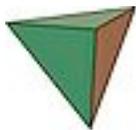
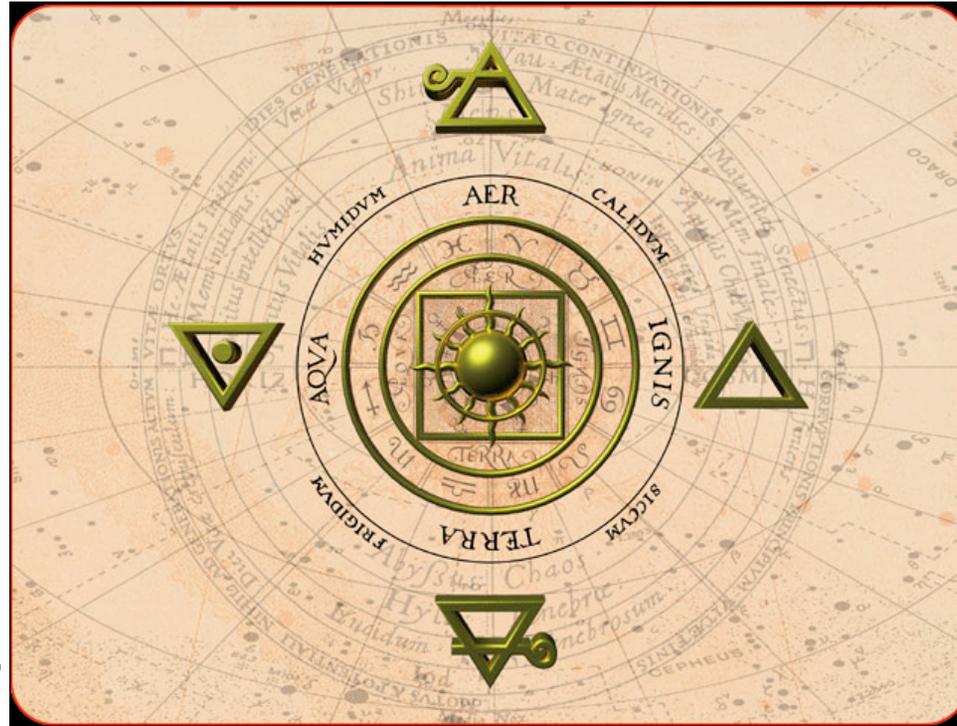
Page from alchemic treatise of
[Ramon Llull](#), 16th century

Atoms and Geometry, IV B.C.



Πλάτων

Plato, 428-348 B.C.



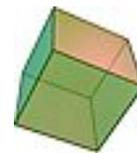
Fire



Air

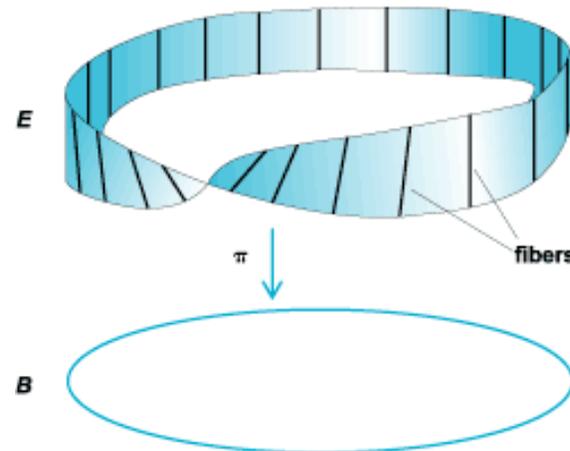
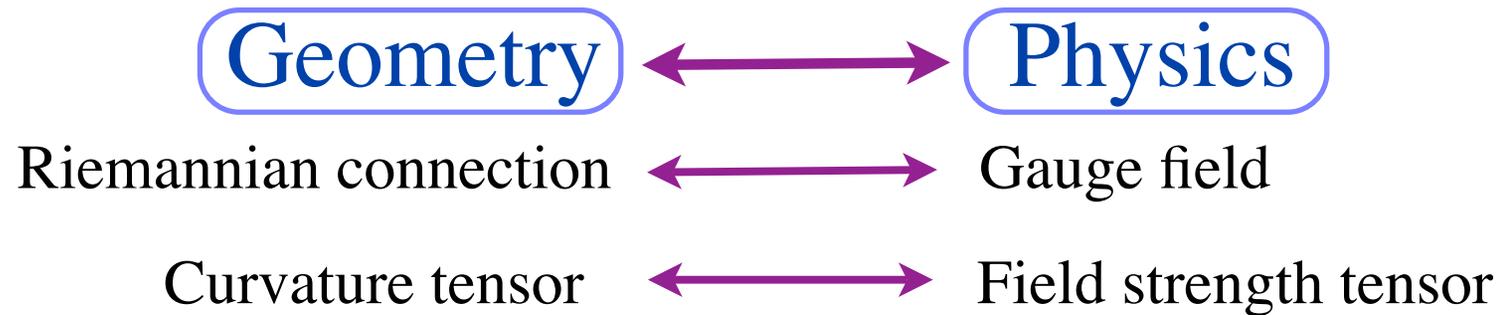


Water



Earth

Atoms and Geometry, 2010 A.D.



Möbius strip, the simplest nontrivial example of a fiber bundle

Gauge theories “live” in a fiber bundle space that possesses non-trivial topology (knots, links, twists,...)

AdS/CFT Correspondence:

Conformal gauge theory is dual to gravity

The chain of ideas:

1. String model pre-dates QCD as a description of hadron spectrum and hadron scattering amplitudes (Regge poles, Veneziano amplitude, Hagedorn spectrum,...)
2. Quantum theory of strings cannot be made consistent in (3+1) dimensions; the minimum number of dimensions is 5.
Also: massless spin 2 state - graviton?
3. It has long been expected that QCD in the strong coupling and large N is described by string theory (at large N , planar diagrams dominate; they describe worldsheets of strings)

4. How to make this string description consistent on the quantum level?

Add a 5th dimension to QCD and let gravity live there

BUT: what is the metric?

$$ds^2 = R^2 w^2(z) (dx_{3+1}^2 + dz^2)$$

what is the form of $w(z)$?

5. Consider instead a conformal theory, such as maximally super-symmetric N=4 Yang-Mills; then the requirement of conformal invariance fixes the metric of the 5th dimension uniquely - it is an Anti- de Sitter space AdS_5 :

invariance w.r.t. $x \rightarrow \lambda x$ fixes $w(z) = \frac{1}{z}$

supersymmetry - S_5 , so the theory lives in $AdS_5 \times S_5$

J.Maldacena

$N = 4$ SU(N) Yang-Mills
theory

=

String theory on
AdS₅ x S⁵

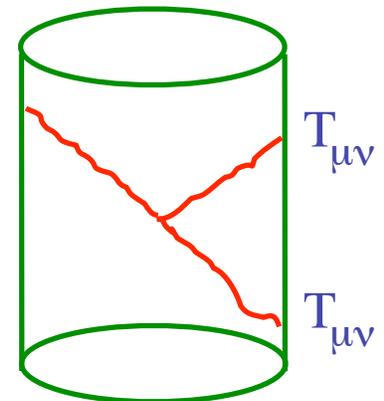
Radius of curvature

$$R_{S^5} = R_{AdS_5} = (g_{YM}^2 N)^{1/4} l_s$$

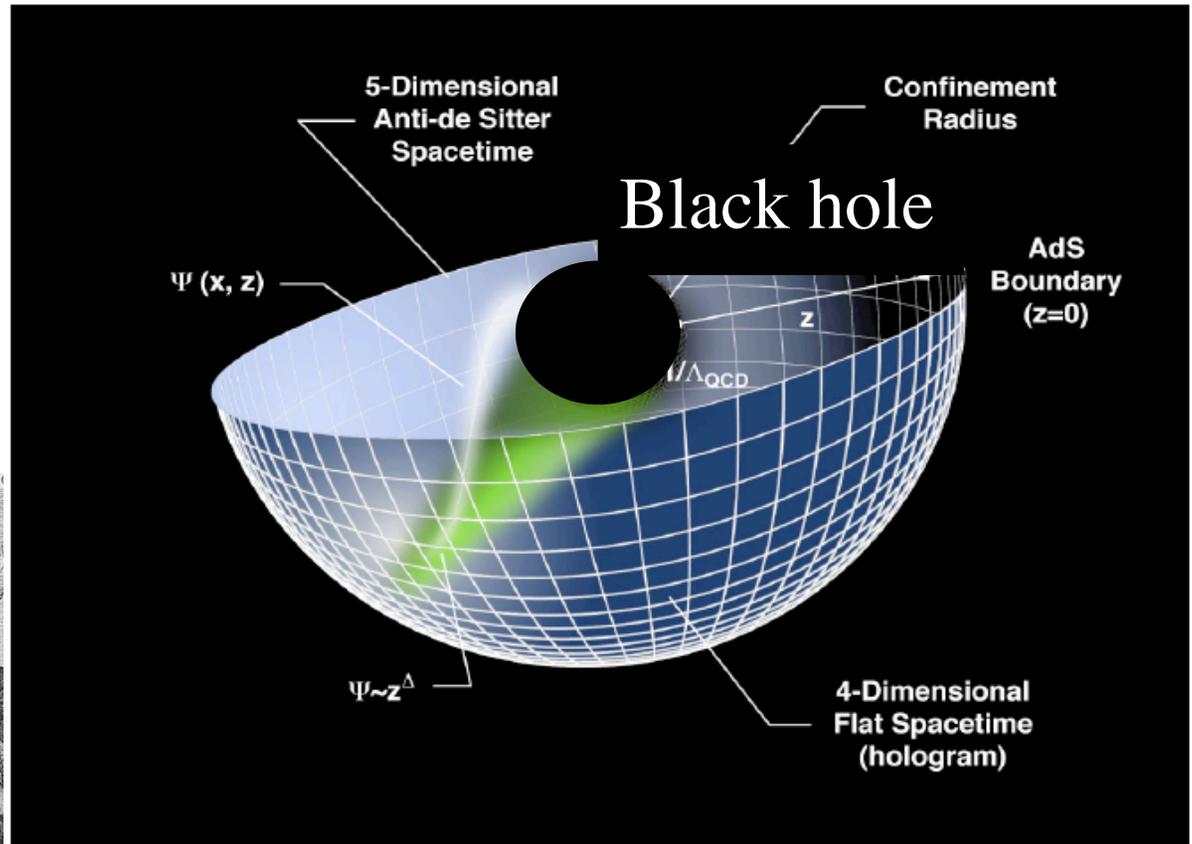
Duality:

$g^2 N$ small \rightarrow R small \rightarrow perturbation theory
 $g^2 N$ large \rightarrow R large \rightarrow classical gravity

$T_{\mu\nu}$



The metaphor of the cave, 2011 A.D.



“The prisoners would take the shadows to be real things and the echoes to be real sounds, not just reflections of reality, since they are all they had ever seen or heard.”

The successes

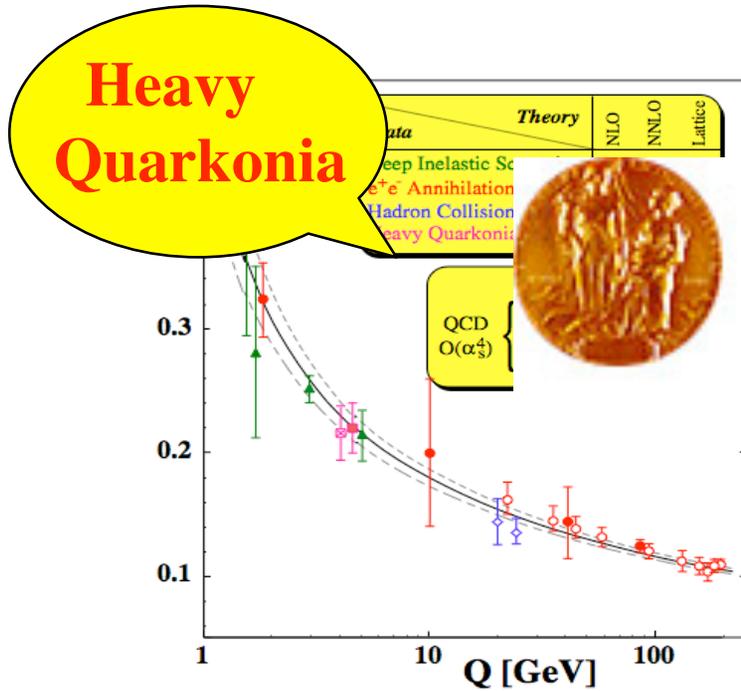
QCD = quarks + geometry

$$\mathcal{L} = -\frac{1}{4}G_{\mu\nu}^a G_{\mu\nu}^a + \sum_f \bar{q}_f^a (i\gamma_\mu D_\mu - m_f) q_f^a;$$

$$D_\mu = \partial_\mu - igA_\mu^a t^a$$

Elegant, consistent, and correct theory

Asymptotic Freedom: “atoms” revealed



At short distances,
the strong force becomes weak
(**anti-screening**) -
one can access the “asymptotically
free” regime in hard processes

and in super-dense matter
(inter-particle distances $\sim 1/T$)

$$\alpha_s(Q) \simeq \frac{4\pi}{b \ln(Q^2/\Lambda^2)}$$

number
of colors

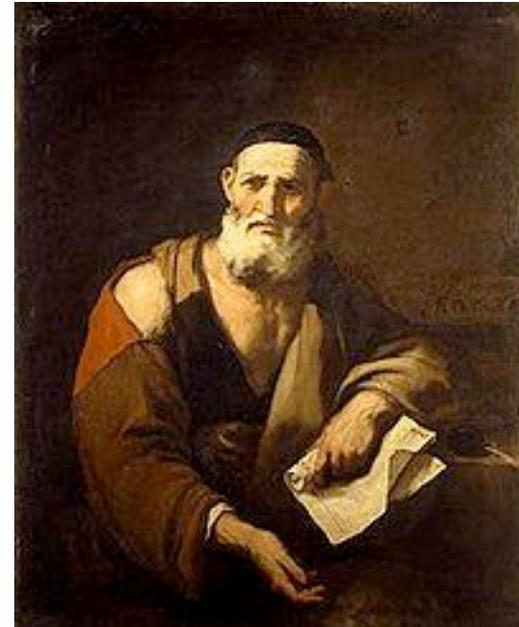
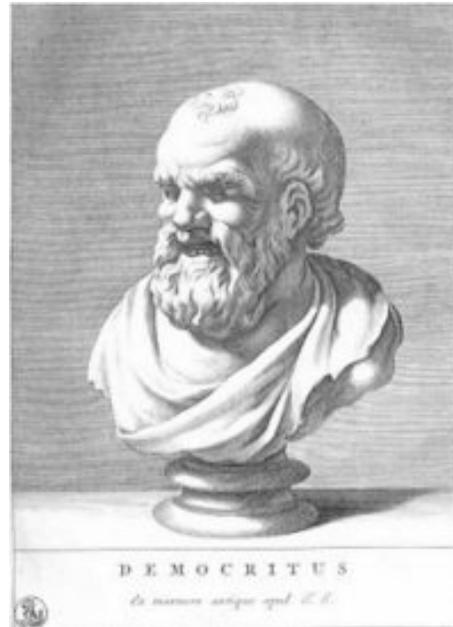
number
of flavors

$$b = (11N_c - 2N_f)/3$$

But: Strong confining interaction at large distances

Where do we stand? (V B.C.)

- ☐ “Atoms”
- ☐ Geometry
- ☐ Void



Where do we stand? (2010 A.D.)

- “Atoms” identified: quarks and leptons
- Geometry (gauge field dynamics)
- Void (the structure of the vacuum)

The next step:
from particles (“atoms”) to fields (geometry)

QCD: understanding the dynamics of gauge fields (geometry)

Problem

Measurement

| | | |
|---|---|--|
| <input type="checkbox"/> Weak/vacuum fields | ↔ | <input type="checkbox"/> Jets, parton fragmentation |
| <input type="checkbox"/> Strong static fields | ↔ | <input type="checkbox"/> Small x distributions in nuclei |
| <input type="checkbox"/> Real-time dynamics | ↔ | <input type="checkbox"/> EM probes, jets, quarkonia |
| <input type="checkbox"/> Gauge fields with boundary conditions/ event horizons | ↔ | <input type="checkbox"/> Bulk behavior, quarkonia , soft photons and dileptons, ... |
| <input type="checkbox"/> Low-energy effective Theory of Everything: hydrodynamics | ↔ | <input type="checkbox"/> Transport properties: shear and bulk viscosities, vorticity |
| <input type="checkbox"/> Topology of gauge fields | ↔ | <input type="checkbox"/> Local parity violation, spin |

Quarkonium and asymptotic freedom (Coulomb potential)

Spectral representation in the t-channel: $V(R) = \sum_m \sigma(m^2) \frac{\exp(-mR)}{R}$

Diagrammatic equation: $\text{Disc}_t \left(\text{Diagram with loop and s-channel line} \right) = \text{Diagram with loop and t-channel line} \propto \sigma \left(\text{Diagram with vertex} \right)$

If physical particles can be produced (positive spectral density),
then unitarity implies screening

$$\left\{ \frac{d \alpha_s^{-1}(R)}{d \ln R} \right\}^{\text{phys}} \propto \frac{1}{3} N + \frac{2}{3} n_f$$

Gluons Quarks
(transverse)

Coulomb potential in QCD - II

Missing non-Abelian effect: instantaneous Coulomb exchange dressed by (zero modes of) transverse gluons

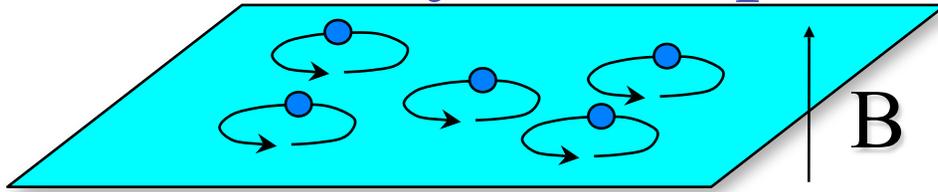
$$\sum_n [0 + \perp \rightarrow 0]^n = \left| \begin{array}{c} 0 \quad 0 \\ \text{---} \bullet \text{---} \\ \text{---} \perp \end{array} \right| + \left| \begin{array}{c} 0 \quad 0 \quad 0 \\ \text{---} \bullet \text{---} \bullet \text{---} \\ \text{---} \perp \quad \perp \end{array} \right| + \dots$$

Negative sign
(the shift of the ground level
due to perturbations - **unstable vacuum!**):

$$\delta E \equiv E - E_0 = \sum_n \frac{|\langle 0 | \delta V | n \rangle|^2}{E_0 - E_n} < 0$$

Anti-screening $\left\{ \frac{d \alpha_s^{-1}(R)}{d \ln R} \right\}^{\text{stat}} \propto -4N$

Asymptotic freedom and the instability of the perturbative vacuum

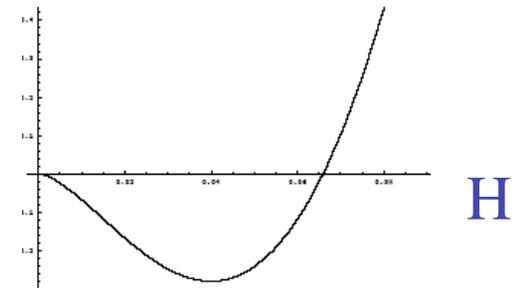


The effective potential: sum over 2D Landau levels

$$V_{\text{pert}}(H) = \frac{gH}{4\pi^2} \int dp_z \sum_{n=0}^{\infty} \sum_{s_z=\pm 1} \sqrt{2gH(n+1/2-s_z) + p_z^2}.$$

Paramagnetic response of the vacuum: V

$$\text{Re } V_{\text{pert}}(H) = \frac{1}{2} H^2 + (gH)^2 \frac{b}{32\pi^2} \left(\ln \frac{gH}{\mu^2} - \frac{1}{2} \right)$$



1. The lowest level $n=0$ of radius $\sim (gH)^{-1/2}$ is **unstable!**

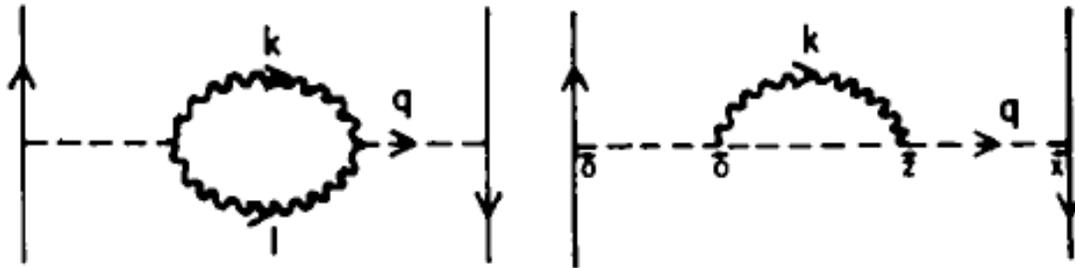
2. Strong fields \longleftrightarrow Short distances

Instability of perturbative
QCD vacuum;
What is the true ground state?

Quarkonia: the potential problems...

$$D_{\mu\nu}^{ab} = i \delta^{ab} \frac{1}{\bar{k}^2} \quad \mu = \nu = 0 \quad \text{Coulomb gauge}$$

$$= i \delta^{ab} \left(\delta_{ij} - \frac{k_i k_j}{\bar{k}^2} \right) \frac{1}{k^2 + i\epsilon} \quad \mu = i, \nu = j.$$



Screening

Anti-Screening

THE STATIC POTENTIAL IN QUANTUM CHROMODYNAMICS

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Received 18 May 1977

Quarkonia: the potential problems...

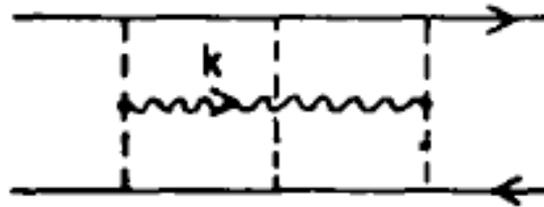
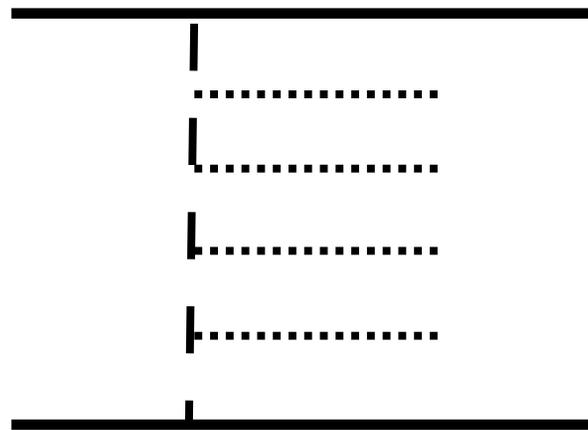


Fig. 4. A three loop graph with a singular static limit.

Higher Fock states; classified in NRQCD

Bodwin, Braaten, Lepage; ...



energy of transverse gluons -
constant per unit of length
(planar, large N limit);
... a string?

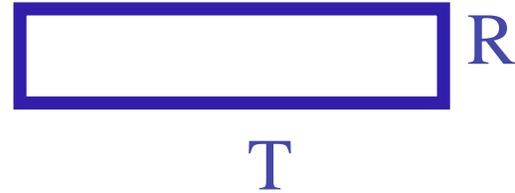
... can be seen as parts of the QCD string?

e.g., the “gluon chain” model

Greensite and Thorn

Scale invariance and confinement

Consider a rectangular Wilson loop:



$$W(C) = \exp \left(ig \int_C A_\mu dx^\mu \right)$$

It is related to the potential $V(R)$ acting between the charges Q and \bar{Q} :

$$W(C) \rightarrow \exp(-TV(R))$$

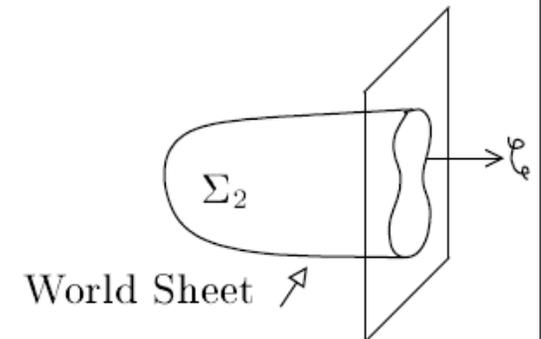
Scale transformation: $T \rightarrow \lambda T$; $R \rightarrow \lambda R$;

the only solution: Coulomb potential

$$V(R) \sim \frac{1}{R}$$

Insights from holography

$$V(L) = - \lim_{T \rightarrow \infty} \frac{1}{T} \ln \langle W(C_{L \times T}) \rangle .$$



Of course, no confinement in N=4 SYM: scale invariance dictates that the expectation value of W must scale with T/L - thus Coulomb potential

$$V(L) \sim 1/L$$

Nevertheless, an interesting lesson: solving for the minimal surface in super-gravity and evaluating its area, one finds

$$V(L) = - \frac{4\pi^2 \sqrt{\lambda}}{\Gamma^4(1/4) L} .$$

$$N \rightarrow \infty, \quad \lambda = g_{YM}^2 N .$$

Why square root?

“not your grandfather’s Coulomb potential”,
a weaker one! screening?

Insights from holography

$$\alpha(\lambda) = \begin{cases} \frac{\lambda}{4\pi} + \dots, & \text{for } \lambda \rightarrow 0; \\ \frac{4\pi^2 \sqrt{2\lambda}}{\Gamma^4(1/4)} + O(1), & \text{for } \lambda \rightarrow \infty. \end{cases} \quad V(L) = -\frac{\alpha(L)}{L}$$

Indeed, the square root of the coupling (but not the pre-factor) can be reproduced on the Field Theory side by resummation of planar diagrams:

Erickson, Semenoff

Szabo, Zarembo,

hep-th/9911088

$$-\ln \left\{ 1 + \text{diagram}_1 + \text{diagram}_2 + \dots \right\} = \frac{\lambda T}{4\pi L} - \frac{\lambda^2 T}{(2\pi)^3 L} \ln \frac{T}{L} + \dots$$

$$\sum_{\text{ladders}} \begin{array}{c} S \quad T \\ \text{diagram} \\ 0 \quad 0 \end{array} = 1 + \sum_{\text{ladders}} \begin{array}{c} S \quad T \\ \text{diagram} \\ 0 \quad 0 \end{array}$$

$$\alpha(\lambda) = \frac{\sqrt{\lambda}}{\pi} - 1 + O(1/\sqrt{\lambda}),$$

Infinite number of higher Fock states, but their effect may be resummed?

J/Ψ as a probe of QGP

Volume 178, number 4

PHYSICS LETTERS B

9 October 1986

J/ψ SUPPRESSION BY QUARK-GLUON PLASMA FORMATION ☆

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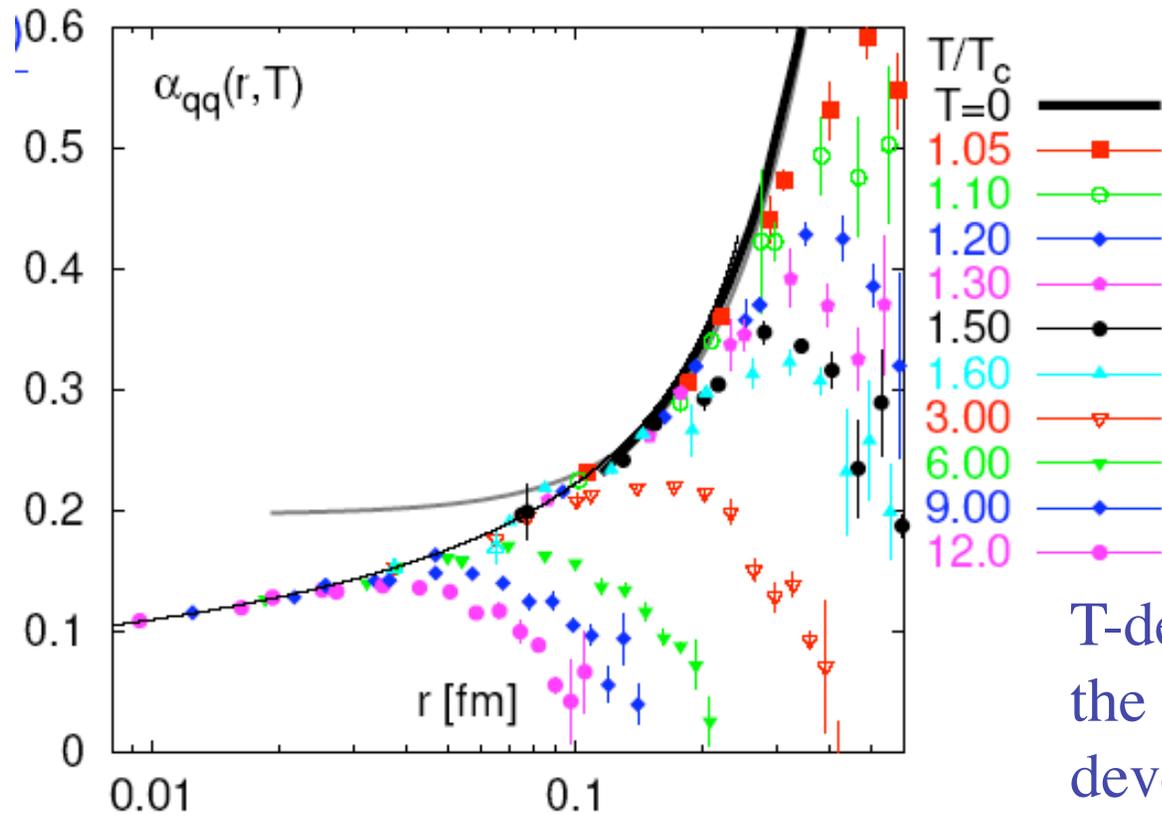
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Received 17 July 1986

If high energy heavy ion collisions lead to the formation of a hot quark-gluon plasma, then colour screening prevents $c\bar{c}$ binding in the deconfined interior of the interaction region. To study this effect, the temperature dependence of the screening radius, as obtained from lattice QCD, is compared with the J/ψ radius calculated in charmonium models. The feasibility to detect this effect clearly in the dilepton mass spectrum is examined. It is concluded that J/ψ suppression in nuclear collisions should provide an unambiguous signature of quark-gluon plasma formation.

Screening in QGP



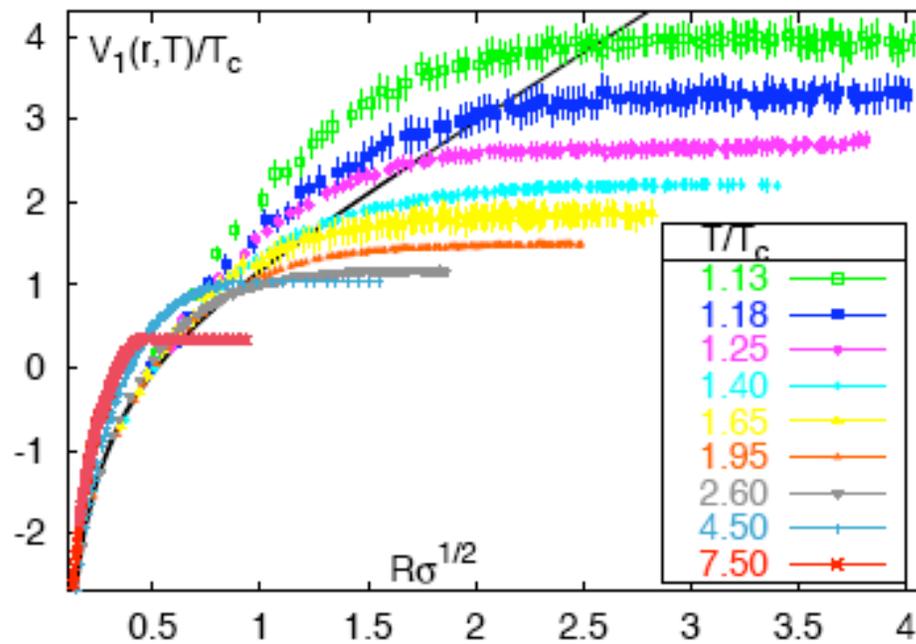
Strong force is
screened by
the presence of
thermal gluons
and quarks

T-dependence of
the running coupling
develops in
the non-perturbative region
at $T < 3 T_c$; $\Delta E/T > 1$

O.Kaczmarek, F. Karsch, P.Petreczky,
F. Zantow, hep-lat/0309121

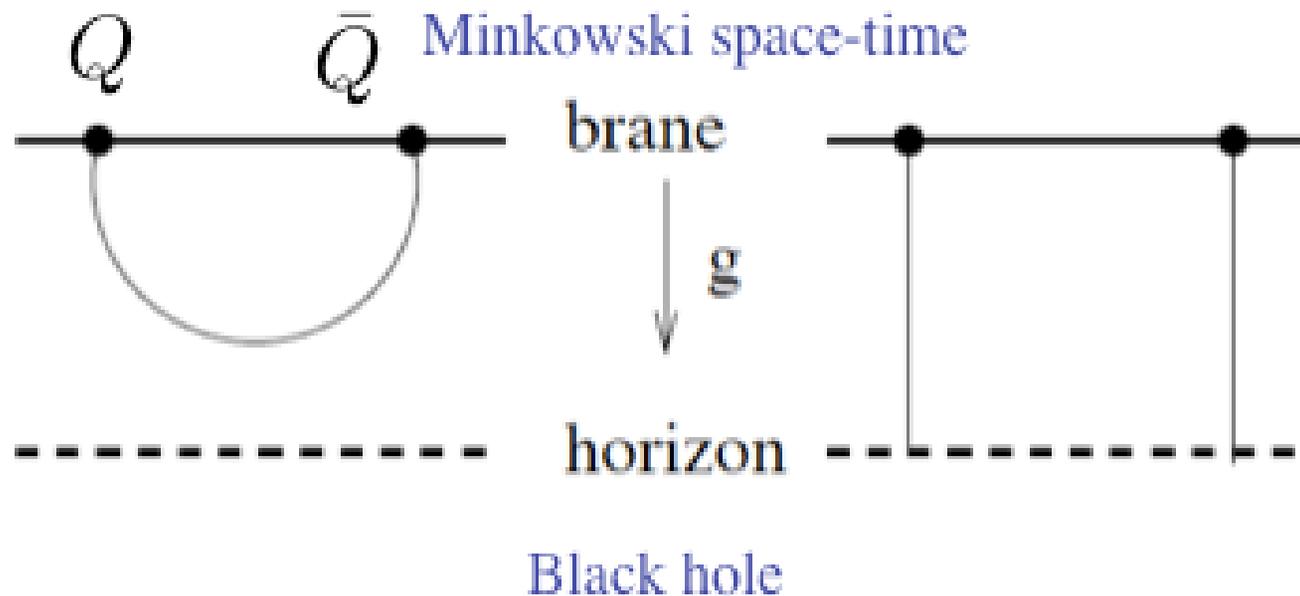
- “cold” plasma

Heavy quark internal energy above T_c



O.Kaczmarek, F. Karsch, P.Petreczky,
F. Zantow, hep-lat/0309121

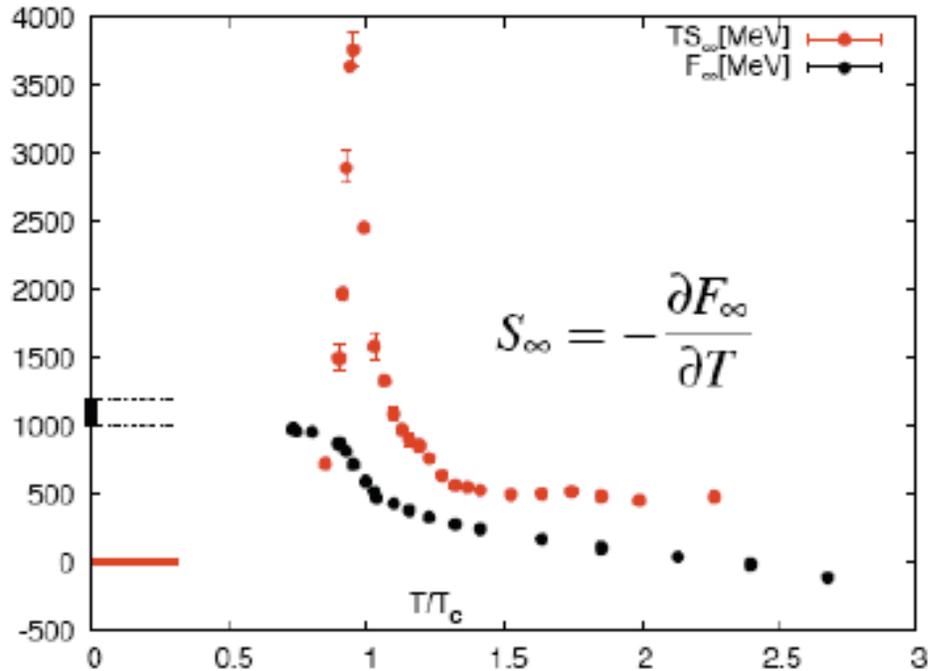
Color screening and black holes



What is the effect
on the entropy?

Shuryak and Zahed,
hep-th/0308073

Entropy: coupling to real states =>
 non-instantaneous, non-local effective potential

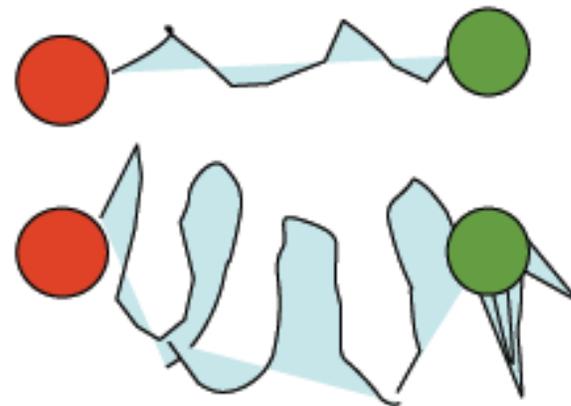


String contribution to
 the black hole entropy:
 a hint for understanding
 better quarkonia
 in the QCD plasma?

$$F = U - TS$$

$$\delta F = \left(\frac{\partial U}{\partial l} - T \frac{\partial S}{\partial l} \right) \delta l = 0$$

Kosterlitz-Thouless phase transition:
 Frautschi, Polyakov



Summary

QCD is a fundamental theory responsible for 95% of the mass of the visible Universe

Heavy quarkonium has been, and remains to be, an important testing ground for new approaches to QCD, such as gauge/gravity

New data (LHC, RHIC, ...) appear and drive the development of theory