



Wright  
Laboratory



Brookhaven<sup>™</sup>  
National Laboratory

# Jets and jet structure In heavy ion collisions

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Summer student lectures @ Prague  
26-28 June 2022

[raghavke.me](http://raghavke.me)

# The basics

**What** are 'jets' and 'substructure'?

**Why** do we study JSS?

**How** does one experimentally measure JSS?

## Physics with jets

**Fundamental** Quantum ChromoDynamics

**Studying transport properties** of the Quark-Gluon Plasma

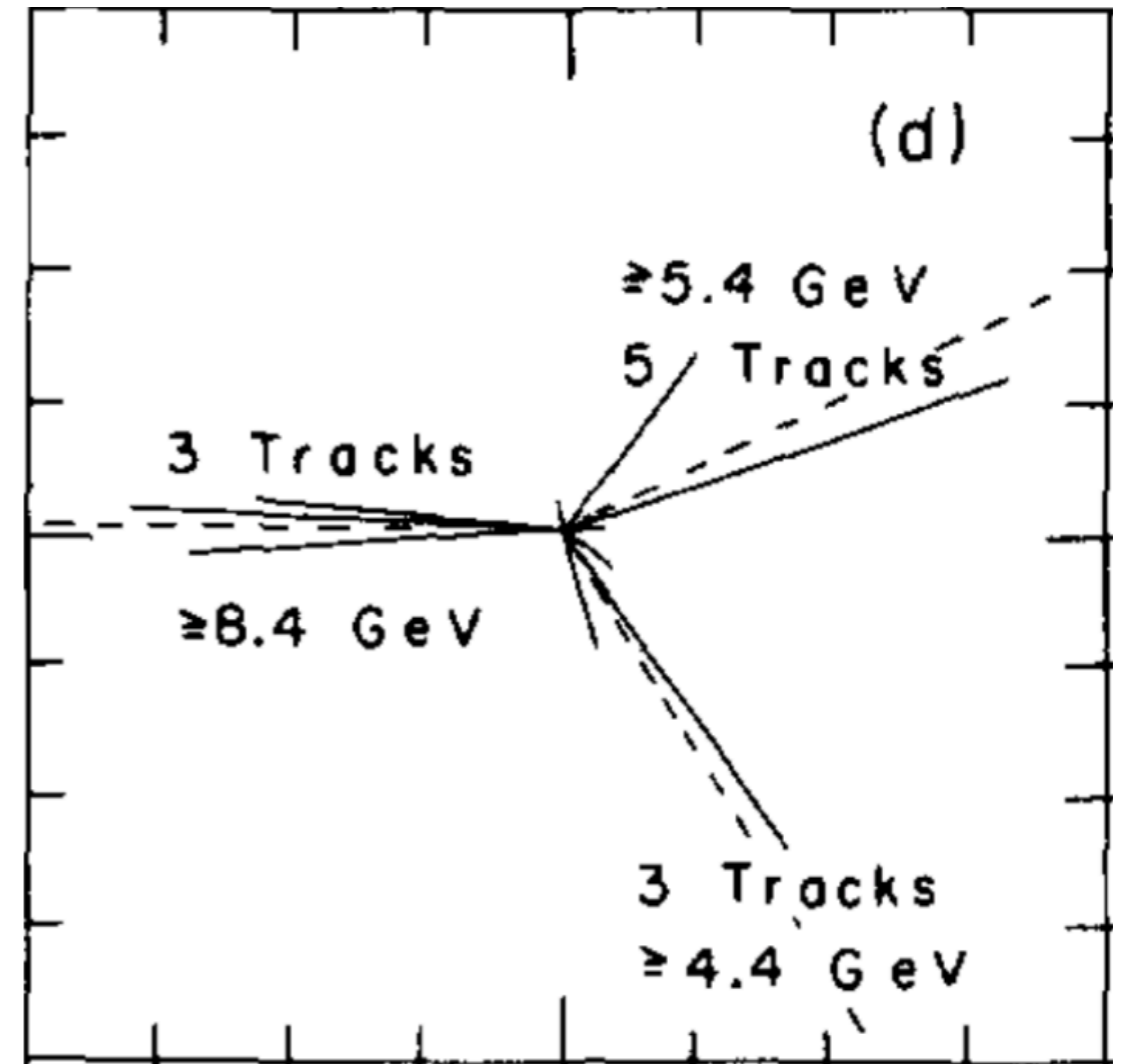
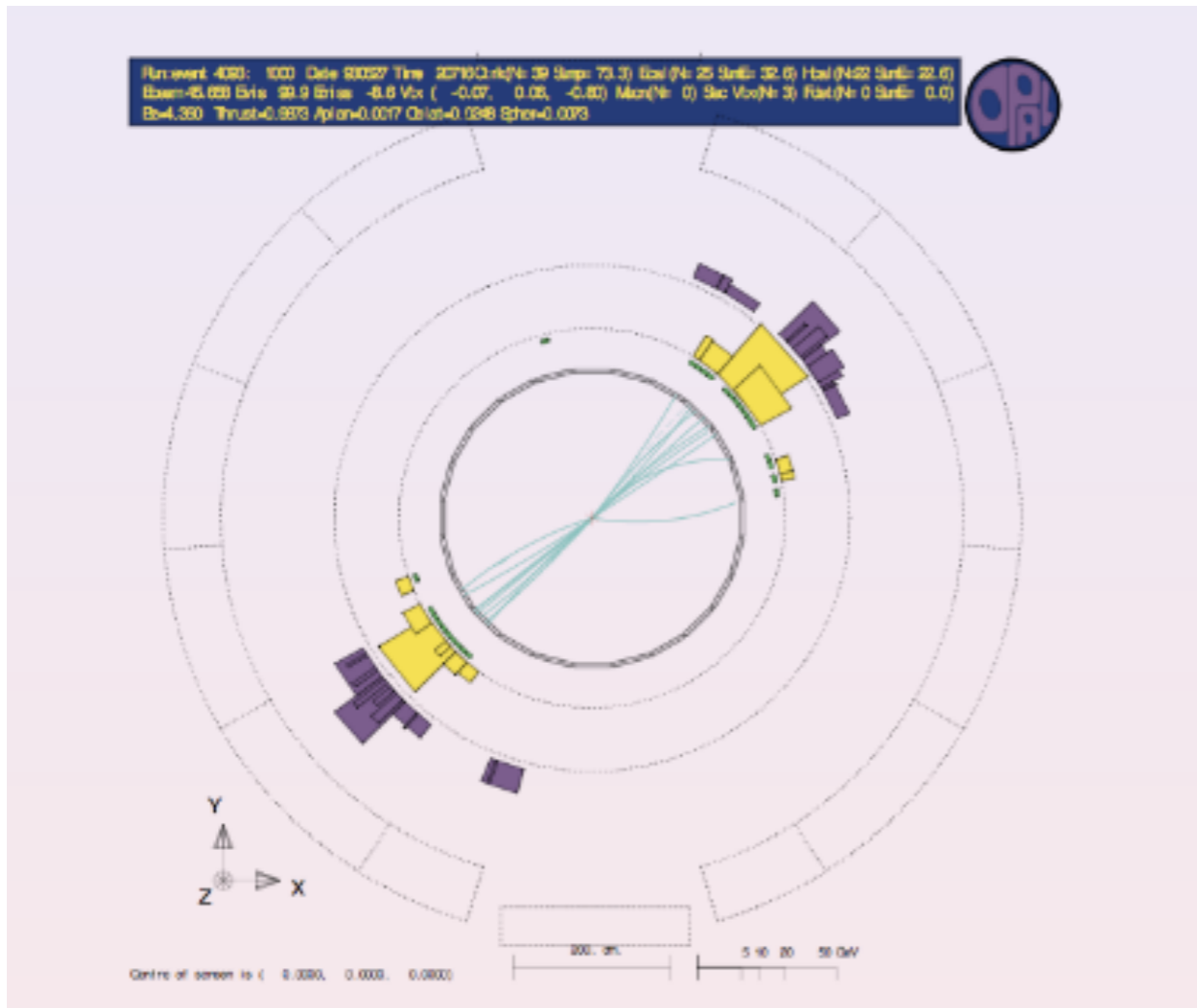
# What are Jets?



# These jets

OPAL Phys. Lett, B 265 462-474 (1991)

TASSO Phys. Lett, V 86B number 2 (1979)



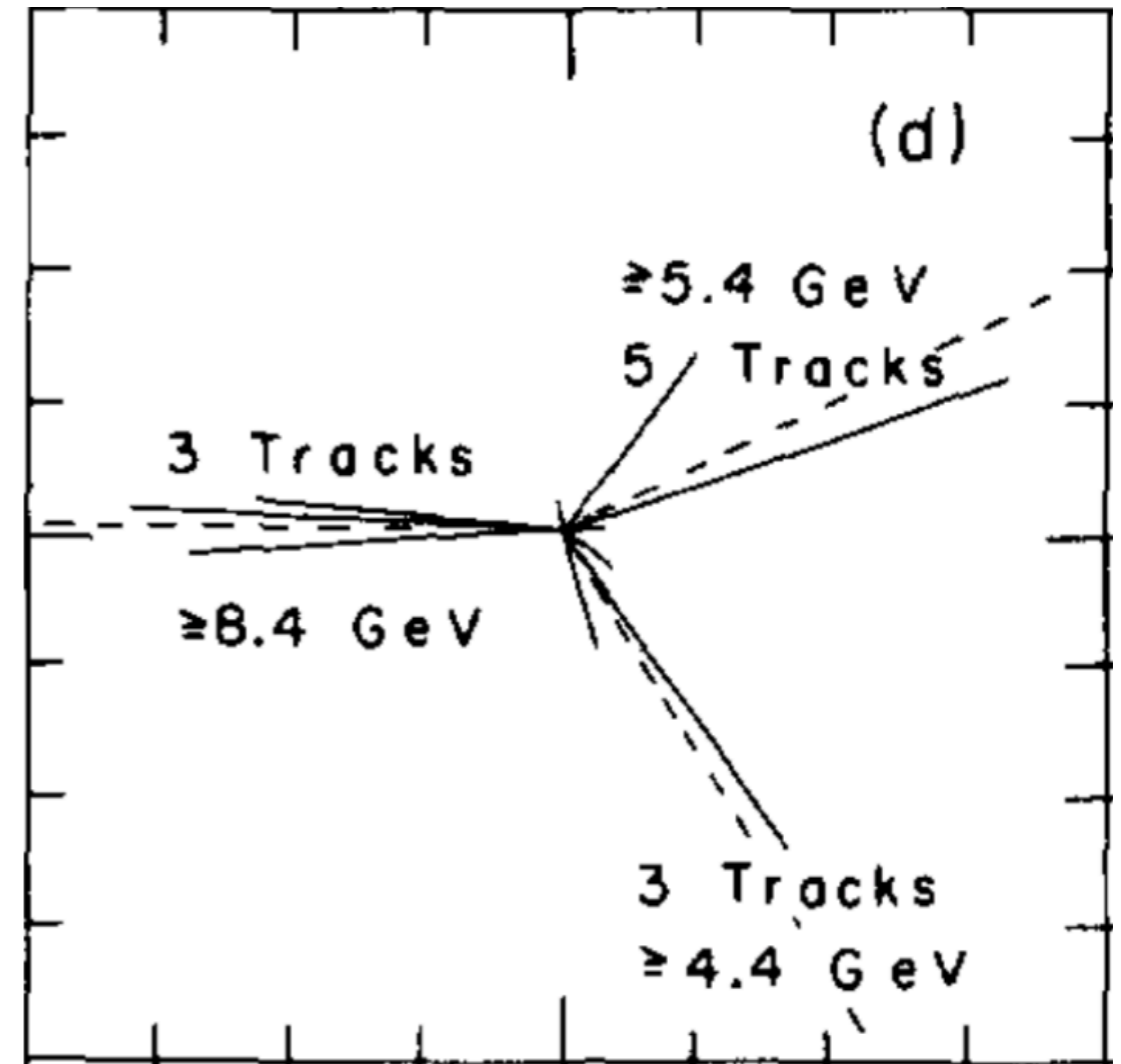
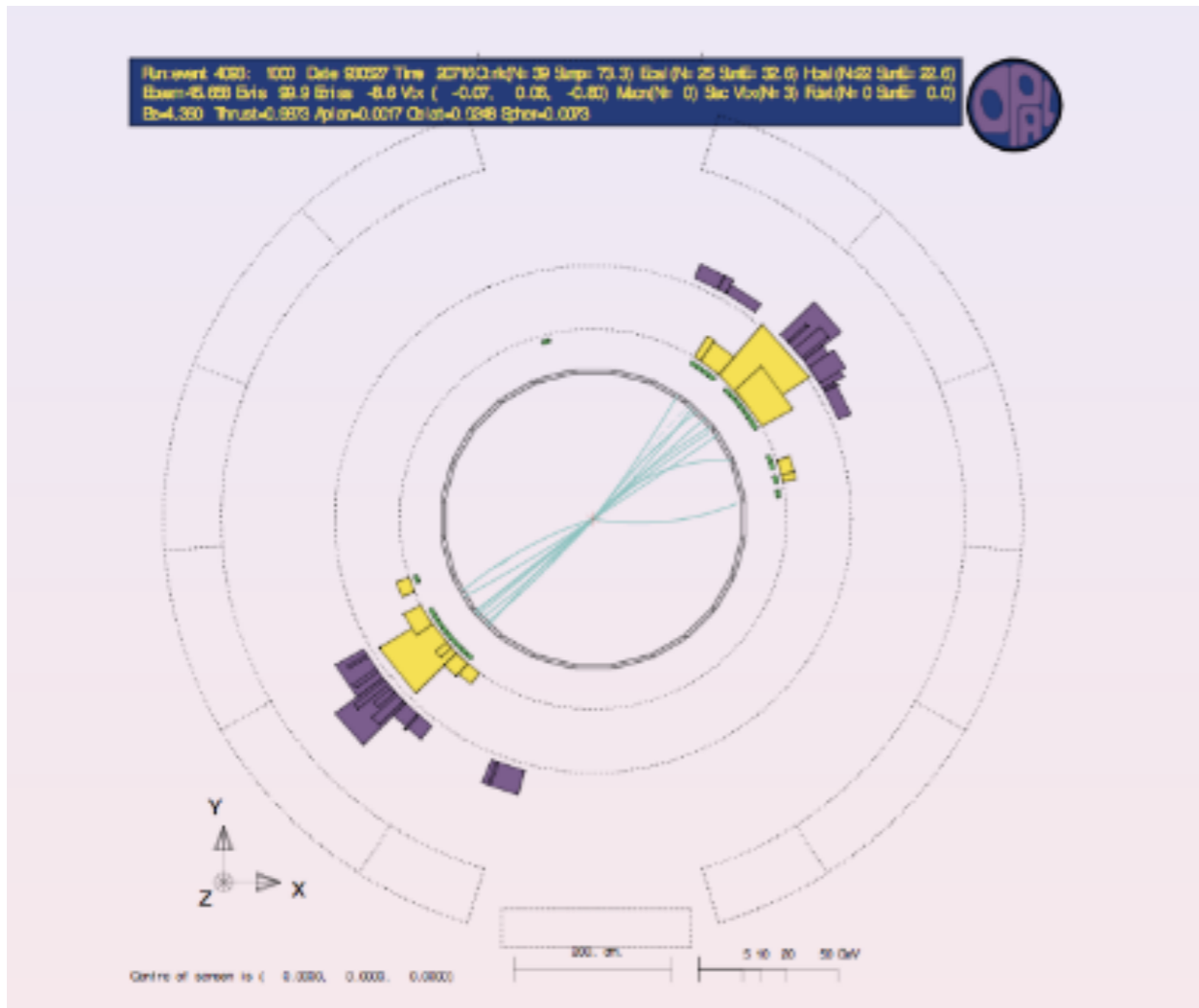
**QUIZ!**

where are the jets?

# These jets

OPAL Phys. Lett, B 265 462-474 (1991)

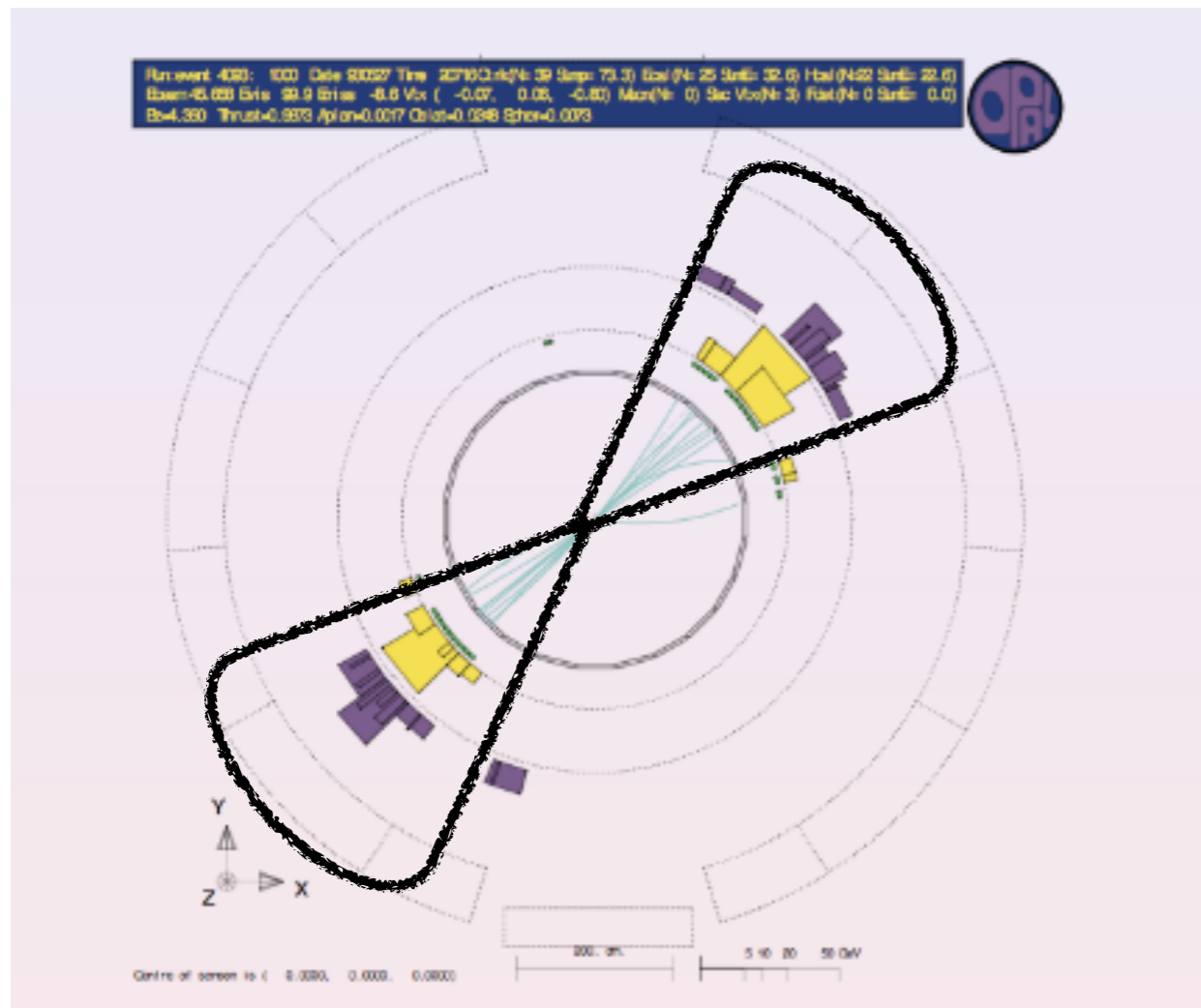
TASSO Phys. Lett, V 86B number 2 (1979)



- Jets - collection of objects detected in an experiment

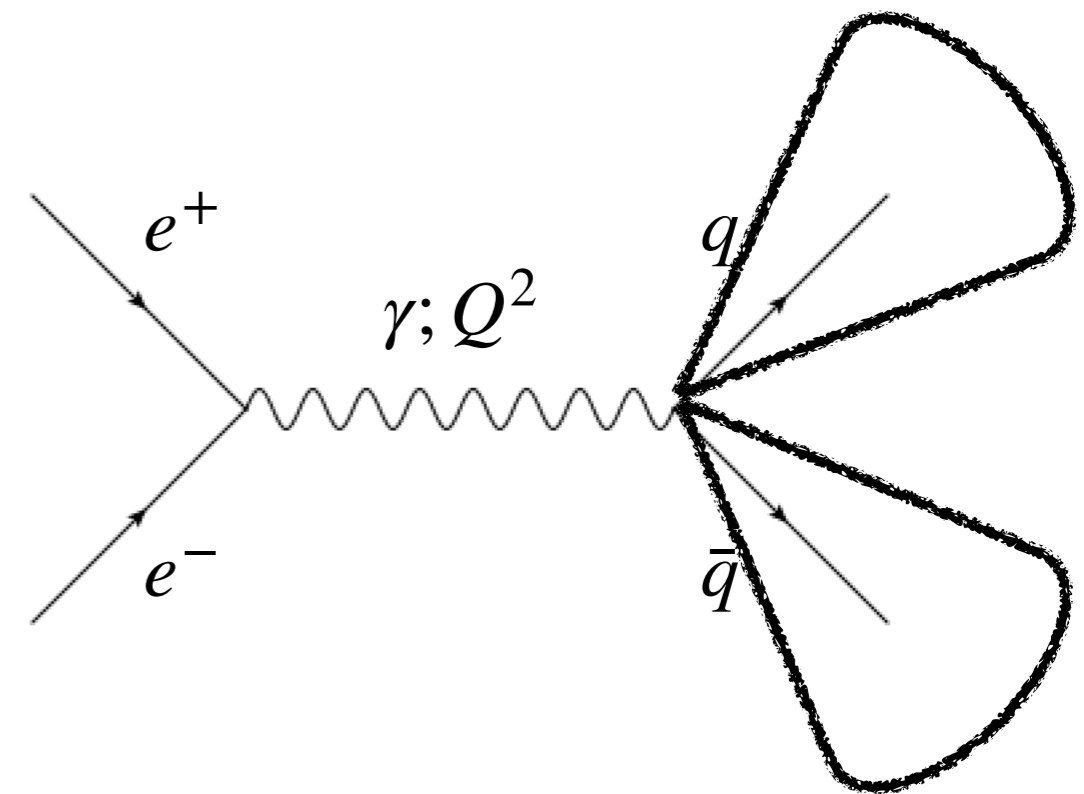
**In order to define our jets,  
we need to know where they come from**

# Where did these jets come from?



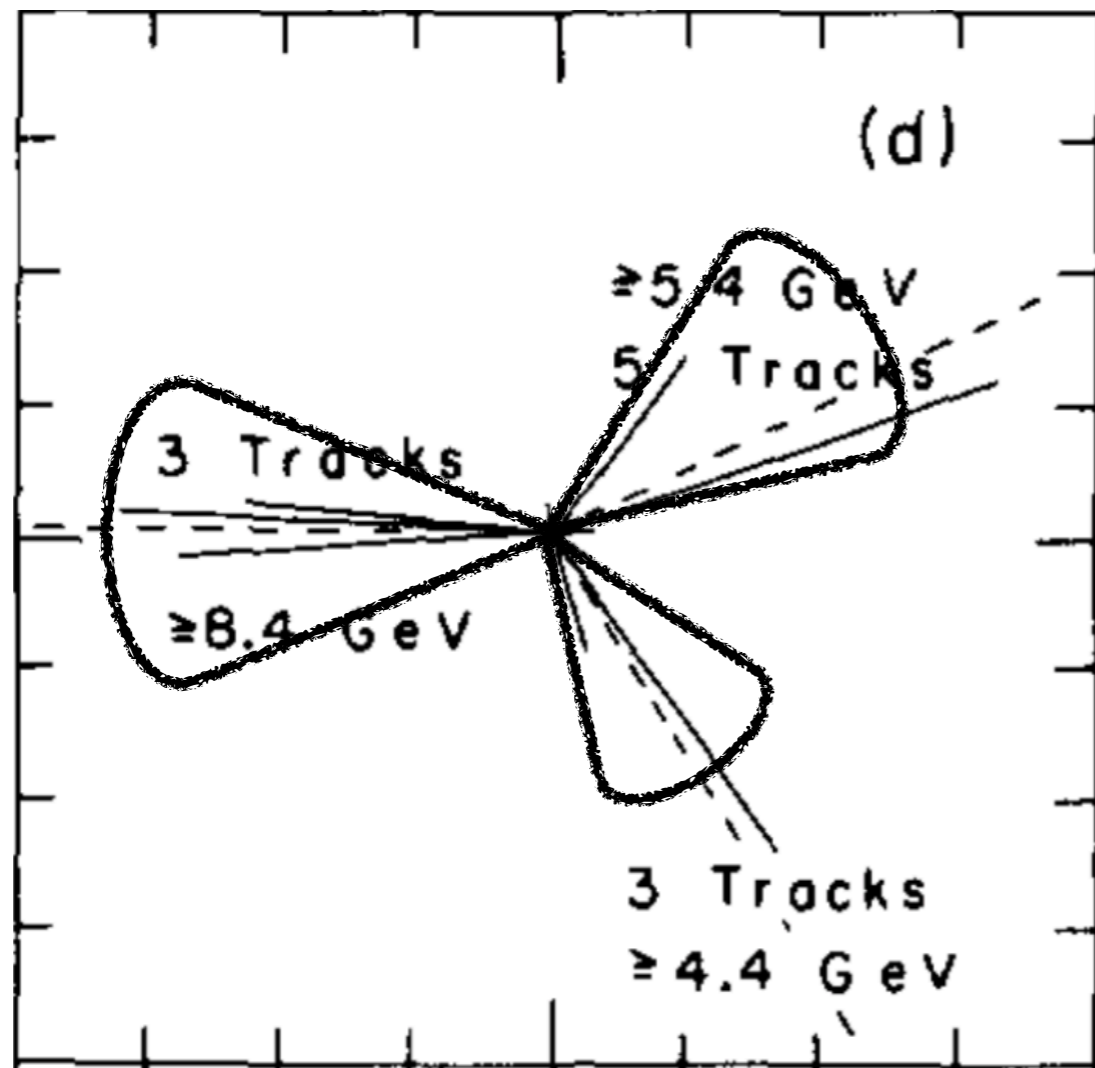
$$\sqrt{s} \approx 91 \text{ GeV}$$

$$e^+ + e^- \rightarrow q + \bar{q}$$



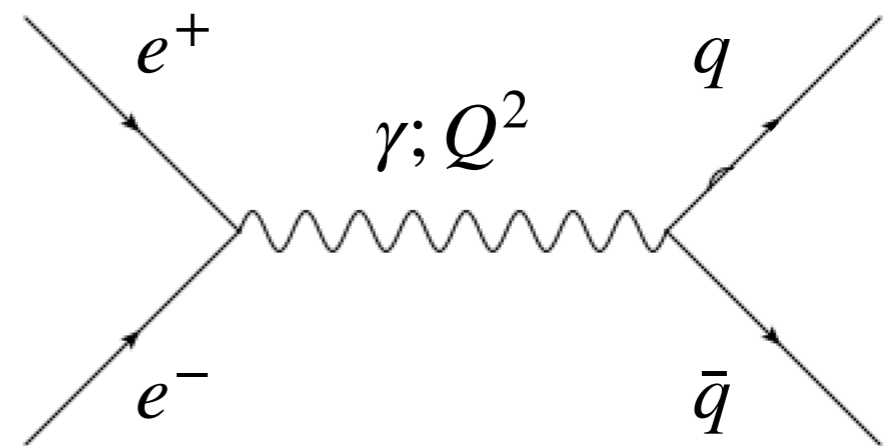
- Electron positron scattering results in a virtual photon - very clean and simple event to study

# Where did these jets come from?



$$\sqrt{s} \approx 20 \text{ GeV}$$

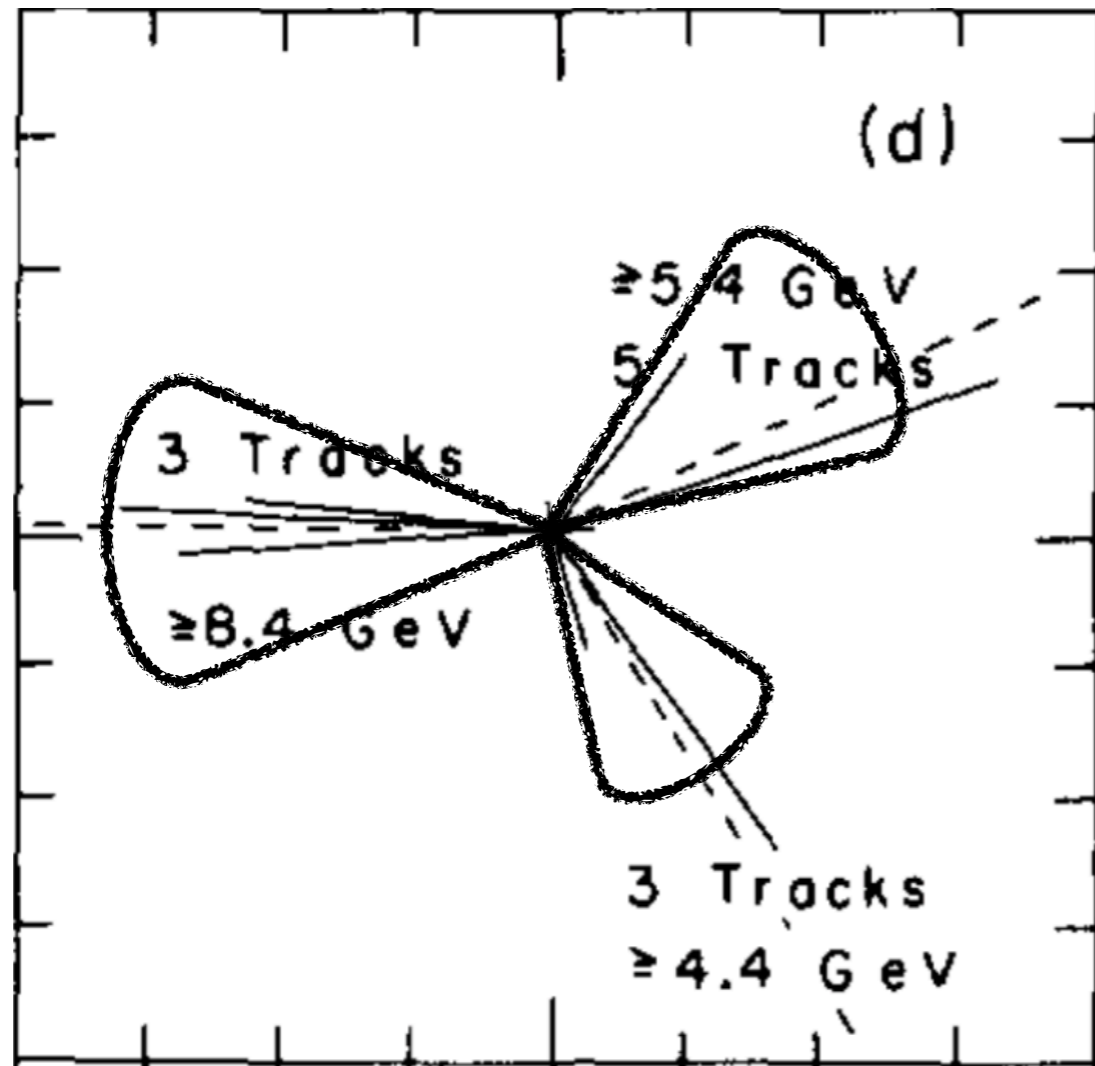
$$e^+ + e^- \rightarrow q + \bar{q}$$



**QUIZ!**

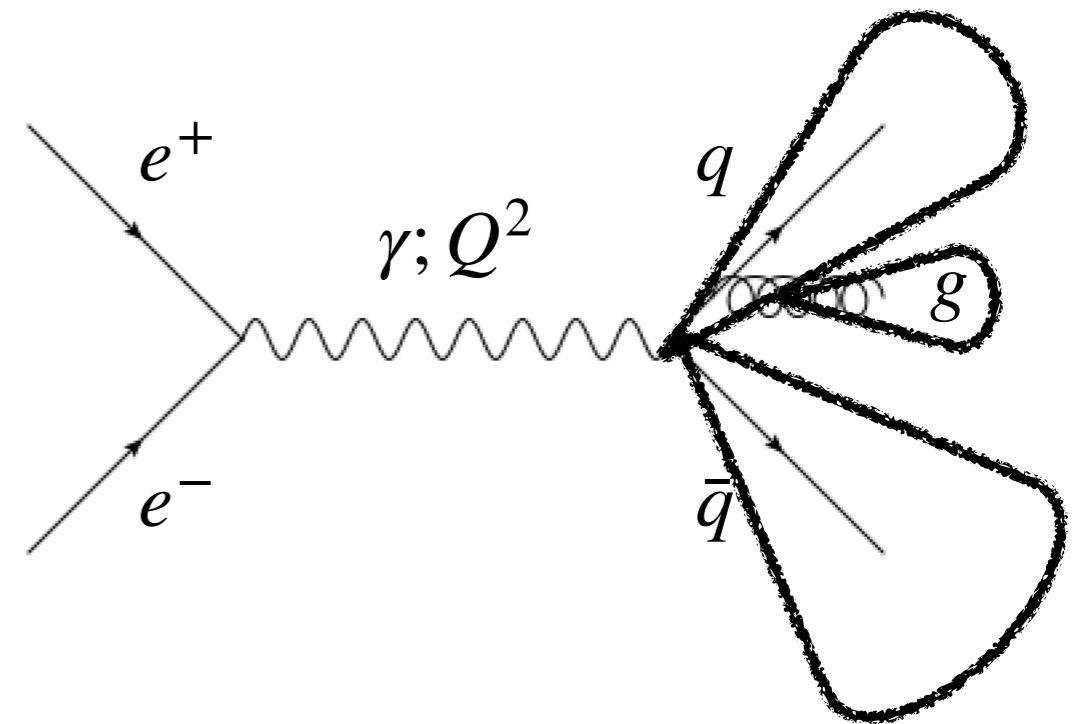
whats the third jet here?

# Where did these jets come from?



$$\sqrt{s} \approx 20 \text{ GeV}$$

$$e^+ + e^- \rightarrow q + \bar{q} + g$$




















- Discovery of the gluon! (Which one is it on the left?)

**Jets are parton proxies!**



# Standard Model of Elementary Particles

		three generations of matter (fermions)			interactions / force carriers (bosons)	
		I	II	III		
QUARKS	mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
	charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
	spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
						
		<b>up</b>	<b>charm</b>	<b>top</b>	<b>gluon</b>	<b>higgs</b>
LEPTONS	mass	$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	
	charge	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0	
	spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
						
		<b>down</b>	<b>strange</b>	<b>bottom</b>	<b>photon</b>	
LEPTONS	mass	$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 91.19 \text{ GeV}/c^2$	
	charge	-1	-1	-1	0	
	spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
						
		<b>electron</b>	<b>muon</b>	<b>tau</b>	<b>Z boson</b>	
LEPTONS	mass	$< 1.0 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$\approx 80.39 \text{ GeV}/c^2$	
	charge	0	0	0	$\pm 1$	
	spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
						
		<b>electron neutrino</b>	<b>muon neutrino</b>	<b>tau neutrino</b>	<b>W boson</b>	


















QUARKS

LEPTONS


















GAUGE BOSONS  
VECTOR BOSONS

SCALAR BOSONS

# Jet producing partons

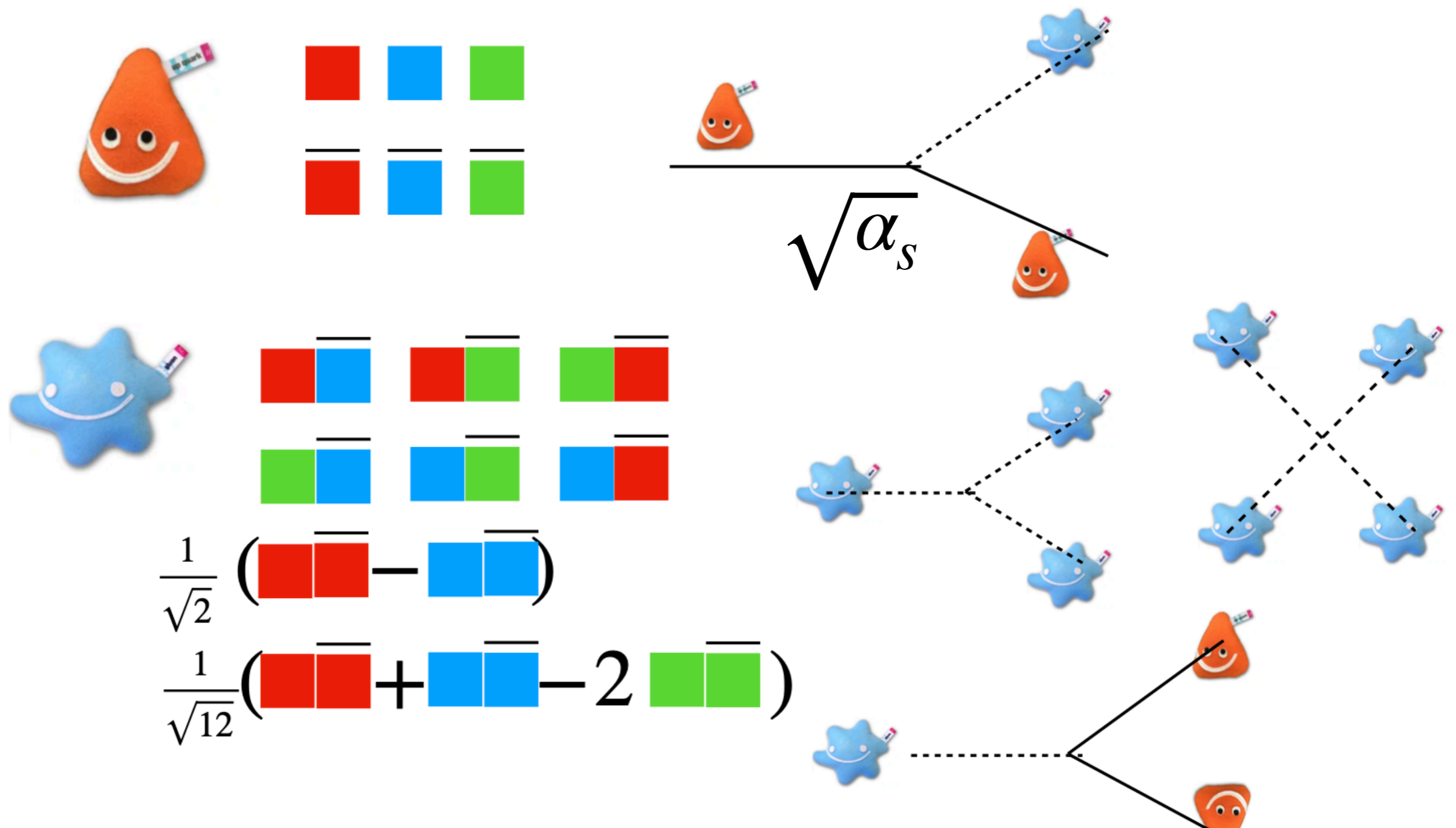
	<p>mass <math>\approx 2.2 \text{ MeV}/c^2</math></p> <p>charge <math>\frac{2}{3}</math></p> <p>spin <math>\frac{1}{2}</math></p>  <p><b>up</b></p>	<p>mass <math>\approx 1.28 \text{ GeV}/c^2</math></p> <p>charge <math>\frac{2}{3}</math></p> <p>spin <math>\frac{1}{2}</math></p>  <p><b>charm</b></p>	<p>mass <math>\approx 173.1 \text{ GeV}/c^2</math></p> <p>charge <math>\frac{2}{3}</math></p> <p>spin <math>\frac{1}{2}</math></p>  <p><b>top</b></p>	<p>0</p> <p>0</p> <p>1</p>  <p><b>gluon</b></p>	<p>mass <math>\approx 124.97 \text{ GeV}/c^2</math></p> <p>0</p> <p>0</p> <p>0</p>  <p><b>H</b></p> <p><b>higgs</b></p>
<b>QUARKS</b>	<p>mass <math>\approx 4.7 \text{ MeV}/c^2</math></p> <p>charge <math>-\frac{1}{3}</math></p> <p>spin <math>\frac{1}{2}</math></p>  <p><b>down</b></p>	<p>mass <math>\approx 96 \text{ MeV}/c^2</math></p> <p>charge <math>-\frac{1}{3}</math></p> <p>spin <math>\frac{1}{2}</math></p>  <p><b>strange</b></p>	<p>mass <math>\approx 4.18 \text{ GeV}/c^2</math></p> <p>charge <math>-\frac{1}{3}</math></p> <p>spin <math>\frac{1}{2}</math></p>  <p><b>bottom</b></p>	<p>0</p> <p>0</p> <p>1</p>  <p><b><math>\gamma</math></b></p> <p><b>photon</b></p>	<b>SCALAR BOSONS</b>
	<p>mass <math>\approx 0.511 \text{ MeV}/c^2</math></p> <p>-1</p> <p><math>\frac{1}{2}</math></p>  <p><b>e</b></p> <p><b>electron</b></p>	<p>mass <math>\approx 105.66 \text{ MeV}/c^2</math></p> <p>-1</p> <p><math>\frac{1}{2}</math></p>  <p><b><math>\mu</math></b></p> <p><b>muon</b></p>	<p>mass <math>\approx 1.7768 \text{ GeV}/c^2</math></p> <p>-1</p> <p><math>\frac{1}{2}</math></p>  <p><b><math>\tau</math></b></p> <p><b>tau</b></p>	<p>mass <math>\approx 91.19 \text{ GeV}/c^2</math></p> <p>0</p> <p>1</p>  <p><b>Z</b></p> <p><b>Z boson</b></p>	<b>GAUGE BOSONS</b>
<b>LEPTONS</b>	<p>mass <math>&lt; 1.0 \text{ eV}/c^2</math></p> <p>0</p> <p><math>\frac{1}{2}</math></p>  <p><b><math>\nu_e</math></b></p> <p><b>electron neutrino</b></p>	<p>mass <math>&lt; 0.17 \text{ MeV}/c^2</math></p> <p>0</p> <p><math>\frac{1}{2}</math></p>  <p><b><math>\nu_\mu</math></b></p> <p><b>muon neutrino</b></p>	<p>mass <math>&lt; 18.2 \text{ MeV}/c^2</math></p> <p>0</p> <p><math>\frac{1}{2}</math></p>  <p><b><math>\nu_\tau</math></b></p> <p><b>tau neutrino</b></p>	<p>mass <math>\approx 80.39 \text{ GeV}/c^2</math></p> <p><math>\pm 1</math></p> <p>1</p>  <p><b>W</b></p> <p><b>W boson</b></p>	<b>VECTOR BOSONS</b>

# Today's partons!

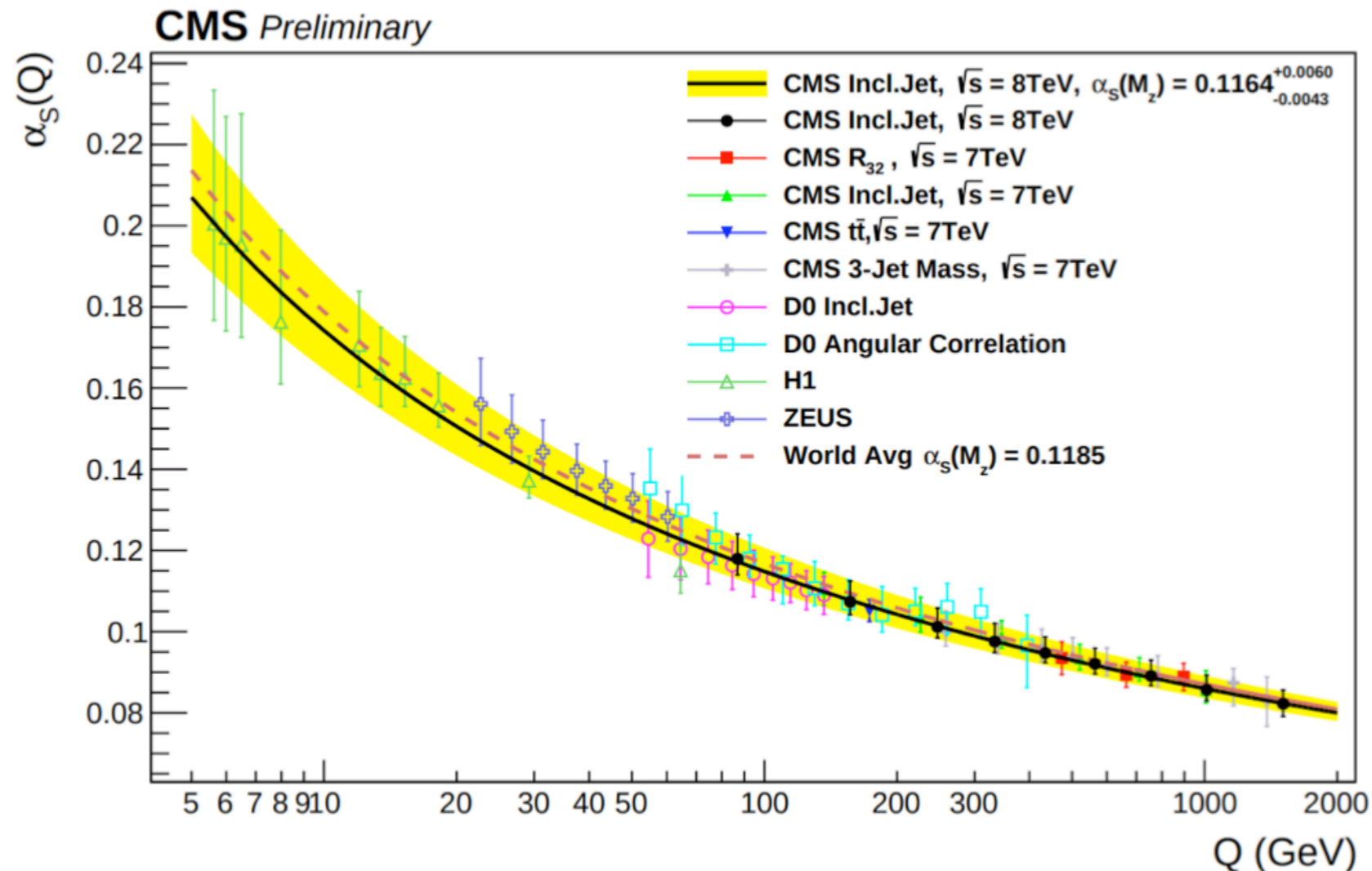
	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
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<b>QUARKS</b>	 <b>up</b>	 <b>charm</b>	 <b>top</b>	 <b>gluon</b>	 <b>higgs</b>
	$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	 <b>down</b>	 <b>strange</b>	 <b>bottom</b>	 <b>photon</b>	
<b>LEPTONS</b>	$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 91.19 \text{ GeV}/c^2$	
	-1	-1	-1	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	 <b>electron</b>	 <b>muon</b>	 <b>tau</b>	 <b>Z boson</b>	
	$< 1.0 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$\approx 80.39 \text{ GeV}/c^2$	
	0	0	0	$\pm 1$	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	 <b>electron neutrino</b>	 <b>muon neutrino</b>	 <b>tau neutrino</b>	 <b>W boson</b>	
				<b>GAUGE BOSONS</b>	<b>SCALAR BOSONS</b>
				<b>VECTOR BOSONS</b>	

# Quantum ChromoDynamics (QCD)

## Interactions with color



# Unique features of QCD



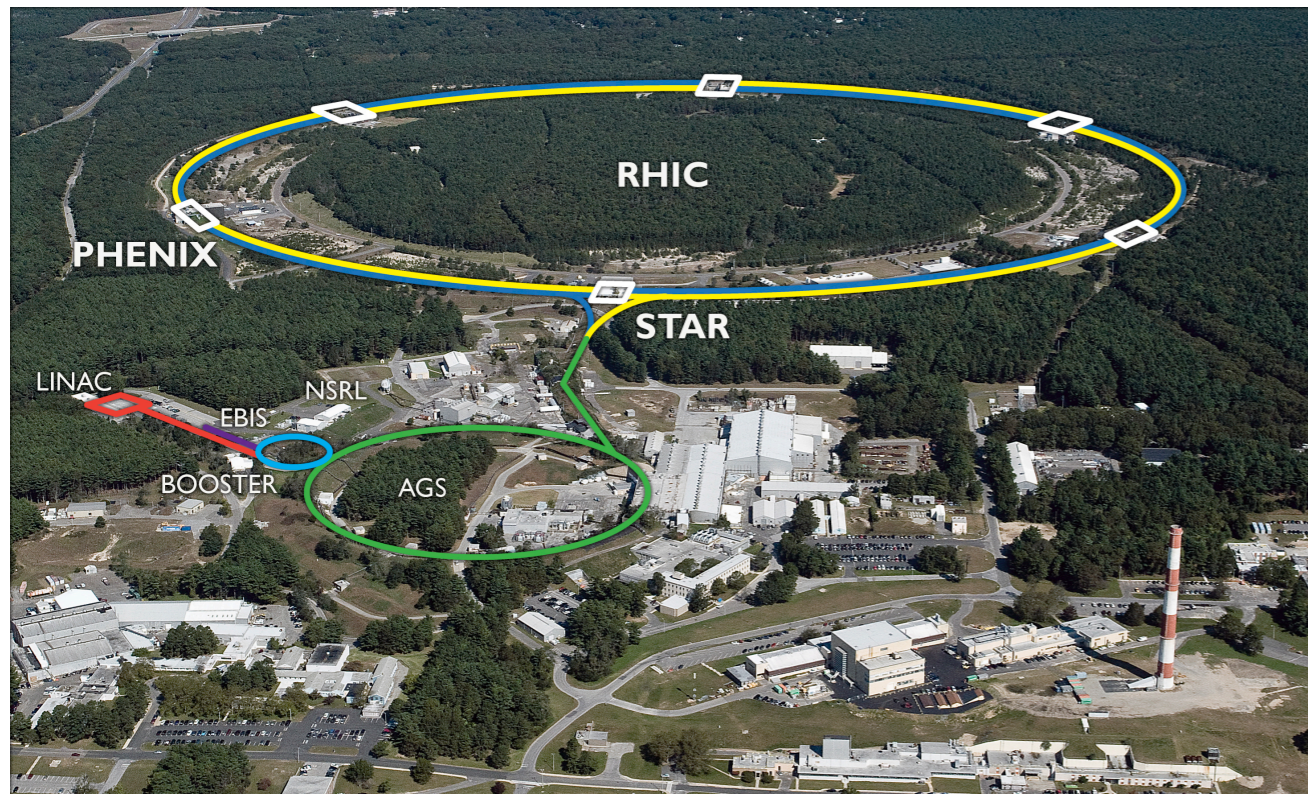
Small energy scales - large value of coupling constant leads to non-perturbative corrections and **confinement/hadronization**



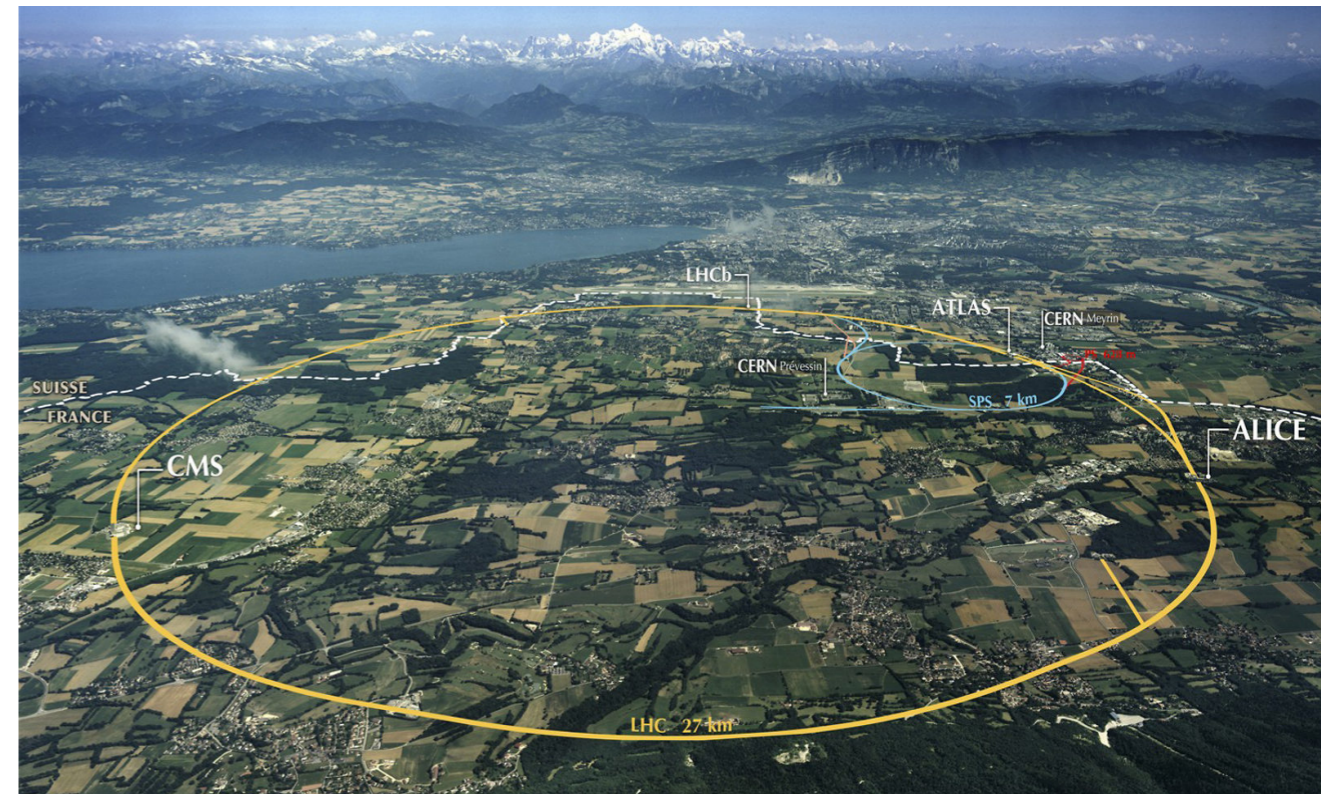
2004

Large energy scales - small values of coupling leads to perturbative calculations - **asymptotic freedom**

# Current accelerators

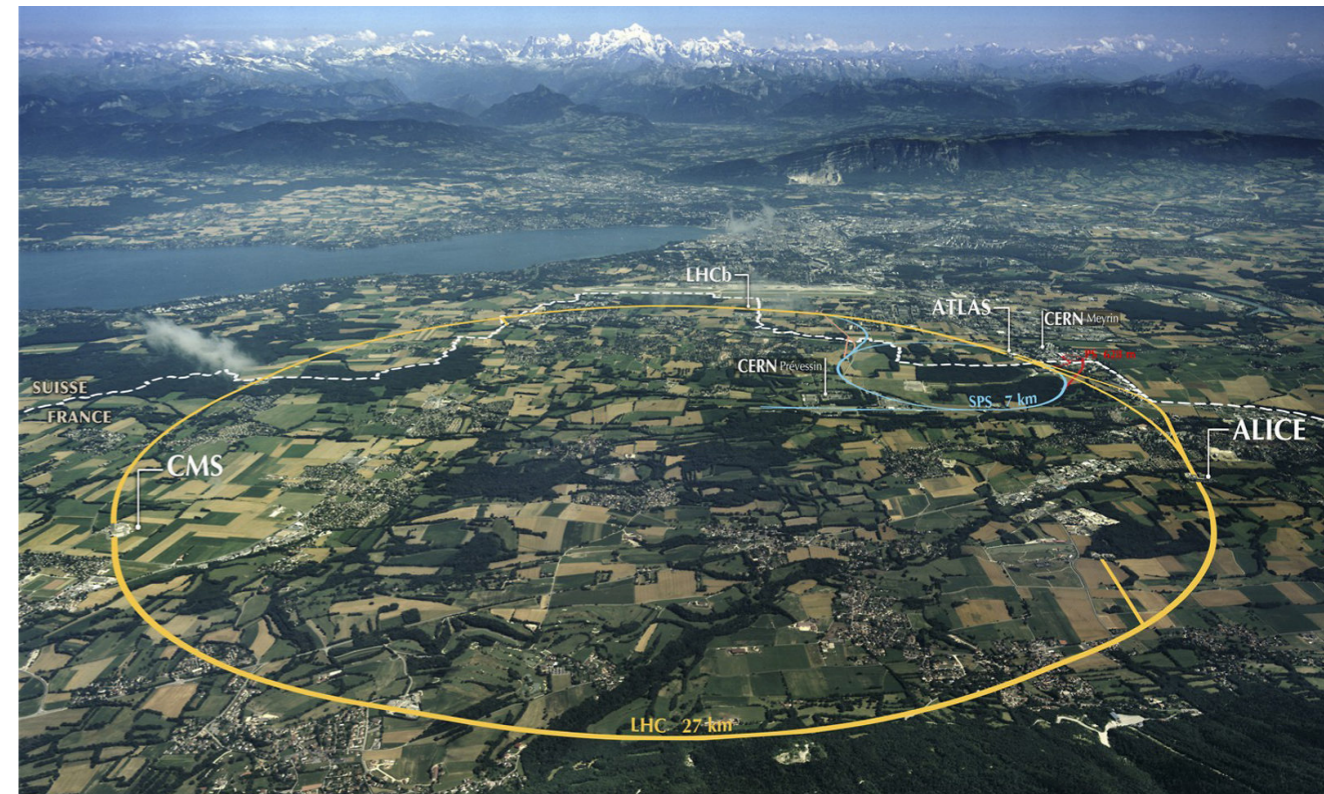
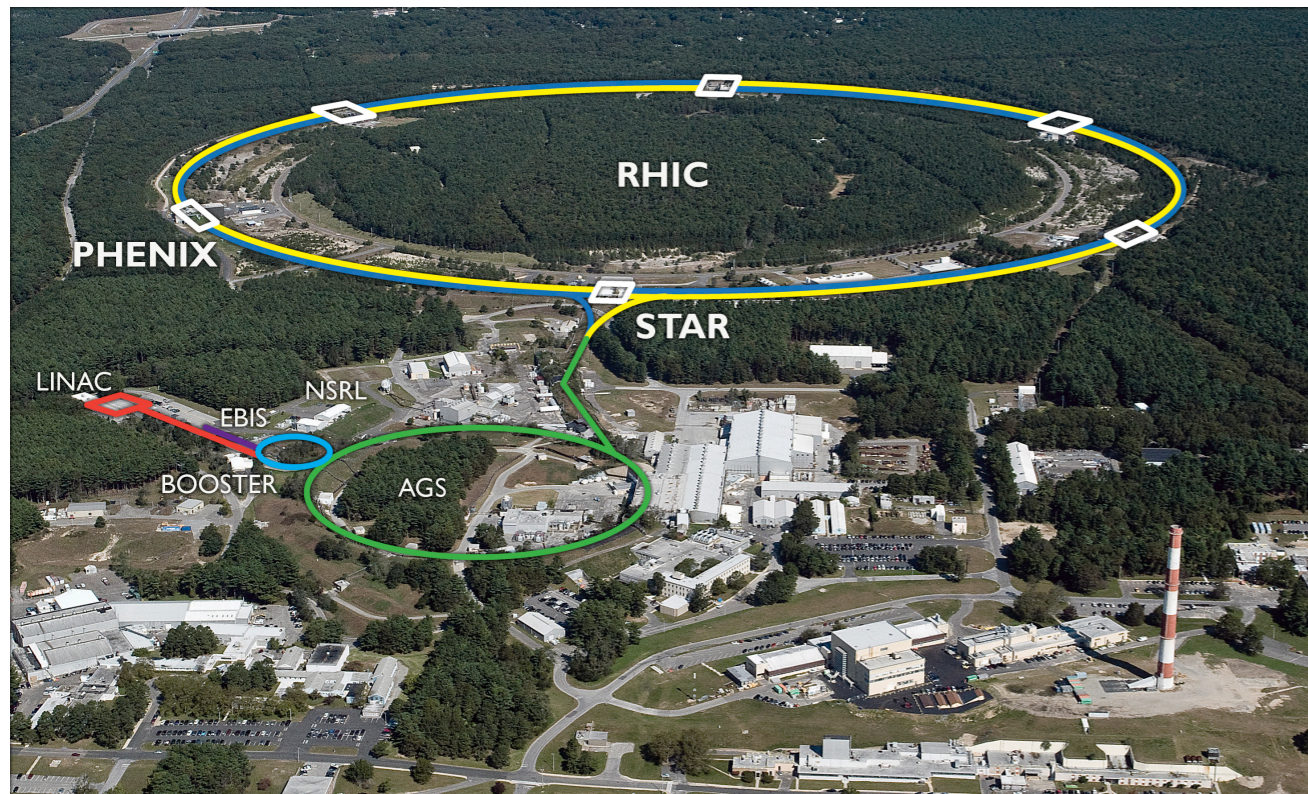


- Relativistic Heavy Ion Collider
- $\sqrt{s_{NN}} = 7.7 - 200 \text{ GeV}$
- Particle Species - p+p, p+Au, d+Au, Au+Au, Cu+Au, Ur+Ur, etc...



- Large Hadron Collider
- $\sqrt{s_{NN}} = 2.76, 5.02 \text{ TeV}$
- Particle Species - p+p, p+Pb, Pb+Pb, Ar+Ar

# Current accelerators



- Relativistic Heavy Ion Collider

- $\sqrt{s_{NN}} = 7.7 - 200 \text{ GeV}$

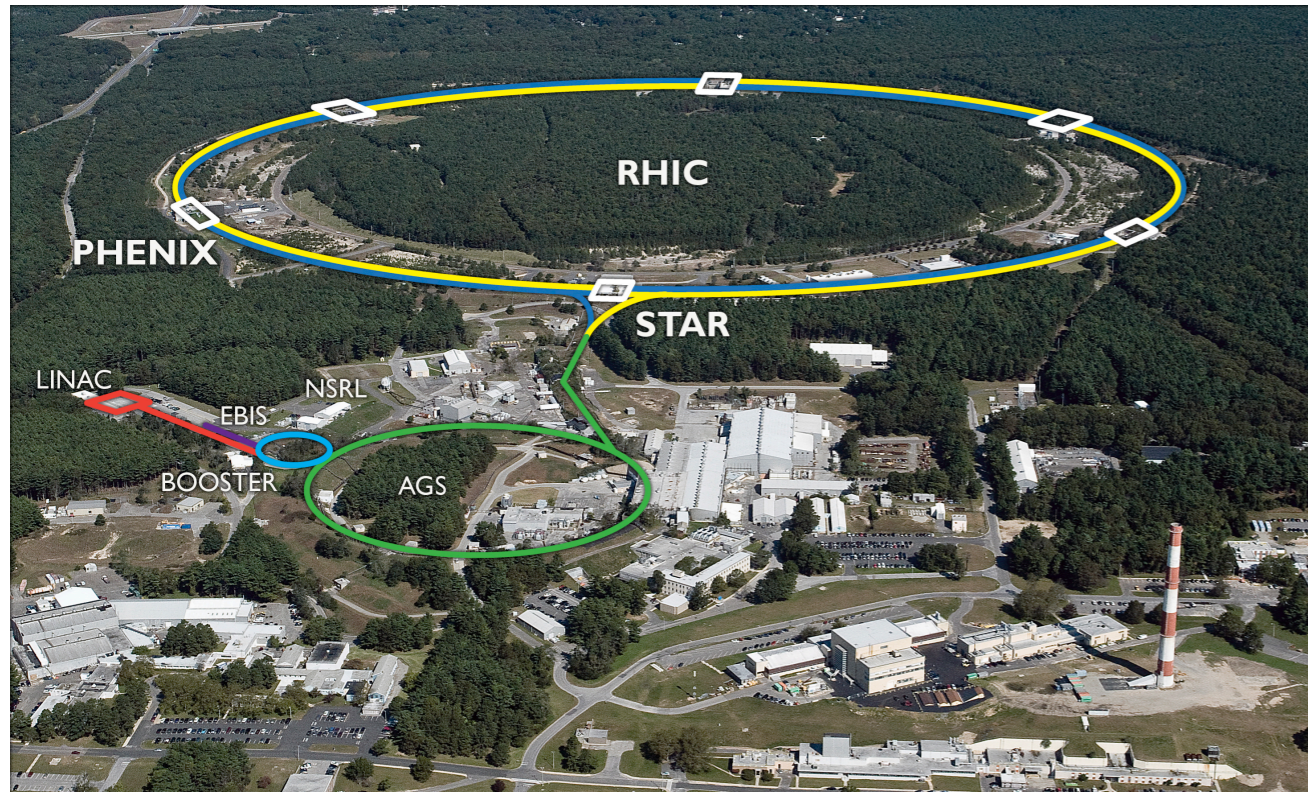
- Particle Species - p+p, p+Au, d+Au, Au+Au, Cu+Au, Ur+Ur, etc...

- Large Hadron Collider

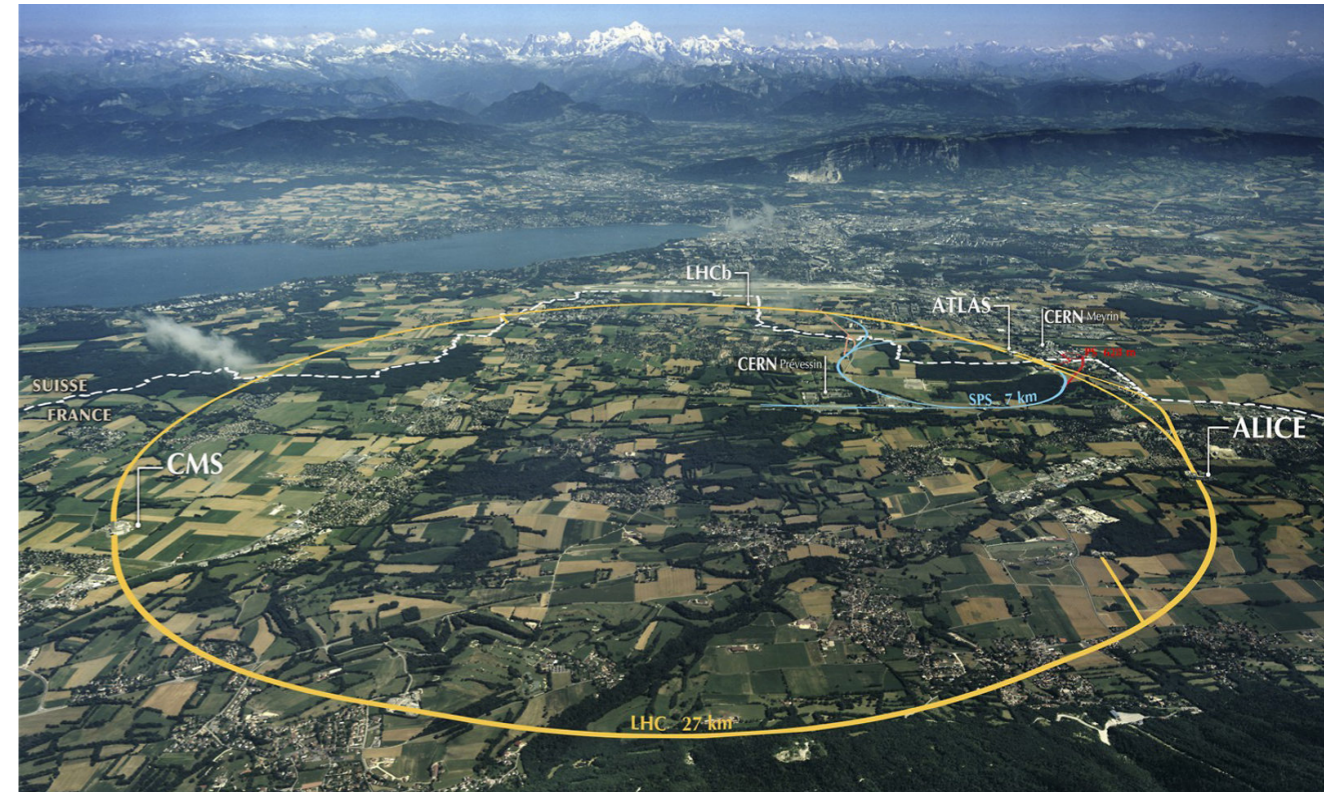
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# Current accelerators



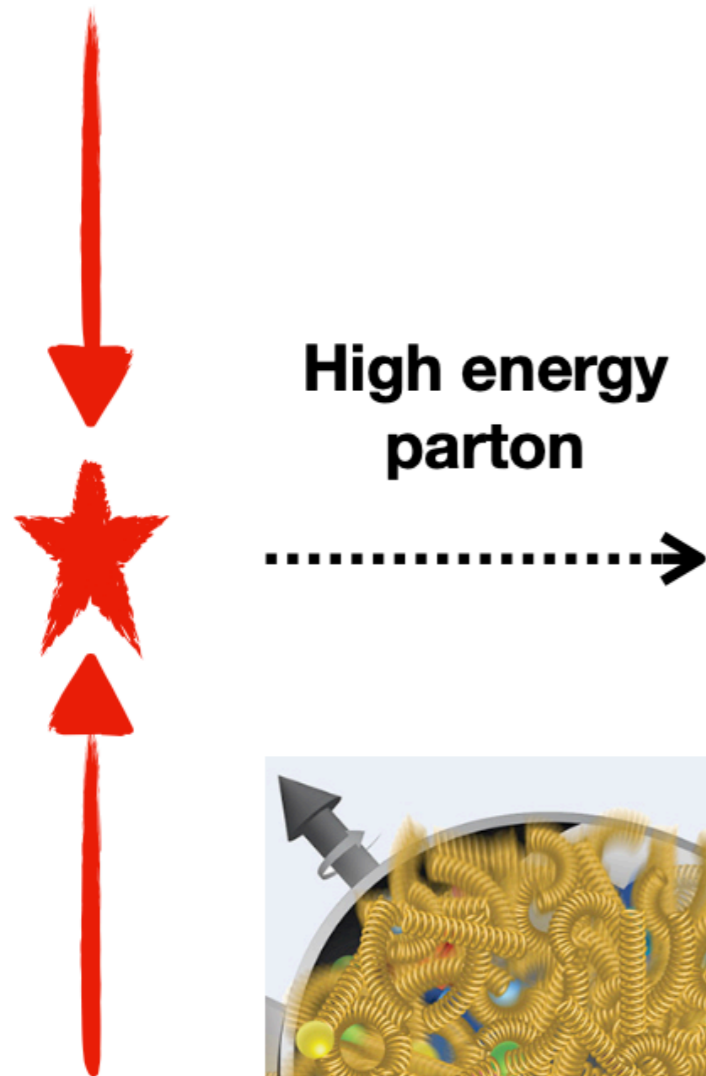
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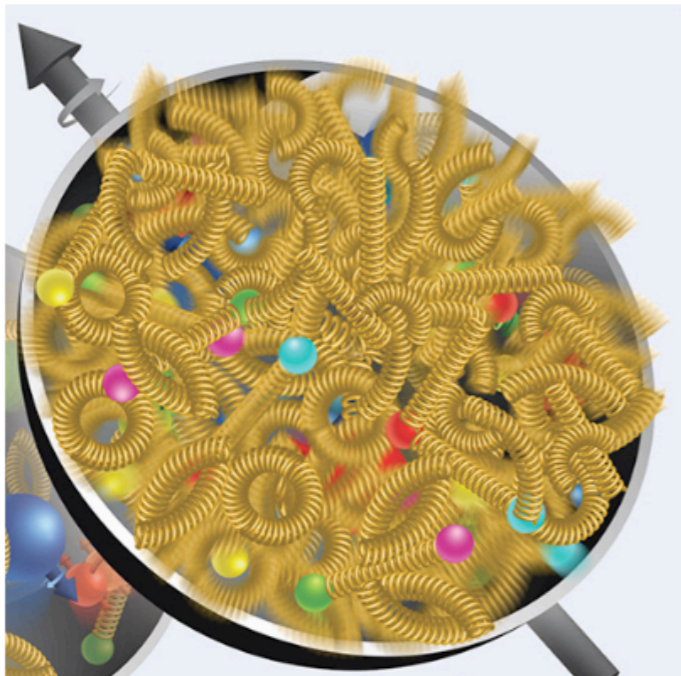
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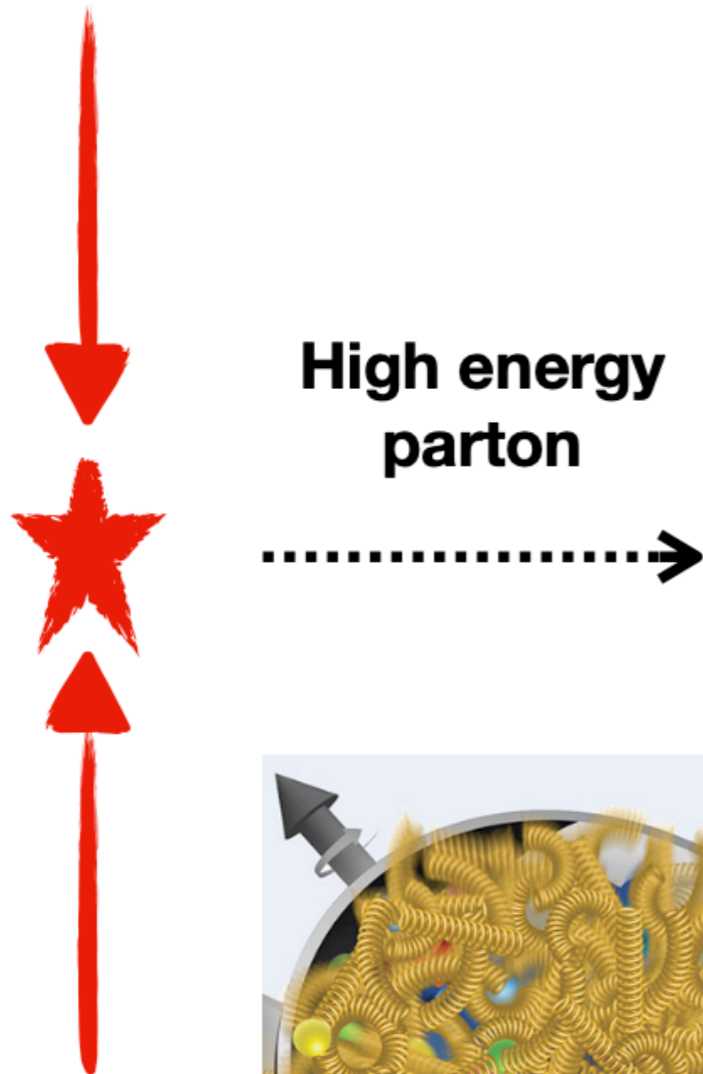
# Observing quarks/gluons



High energy (hard) scattering of valence quarks, sea quarks or gluons

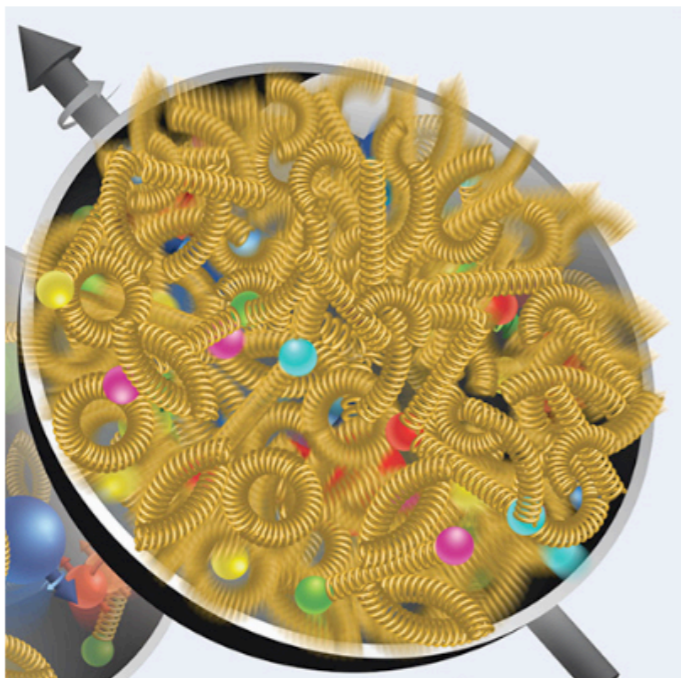


# Observing quarks/gluons



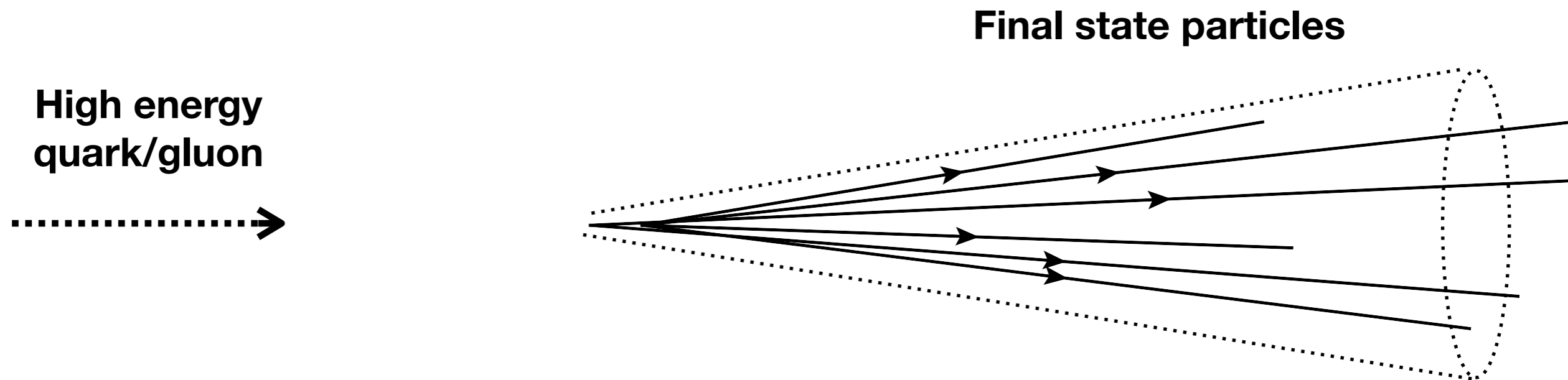
High energy (hard) scattering of valence quarks, sea quarks or gluons

What do we **see** in the experiment?



You already know the answer :)

# Observing quarks/gluons in nature - Jets



Collimated collection of hadrons resulting from the '**metamorphosis**' of partons due to fragmentation and hadronization

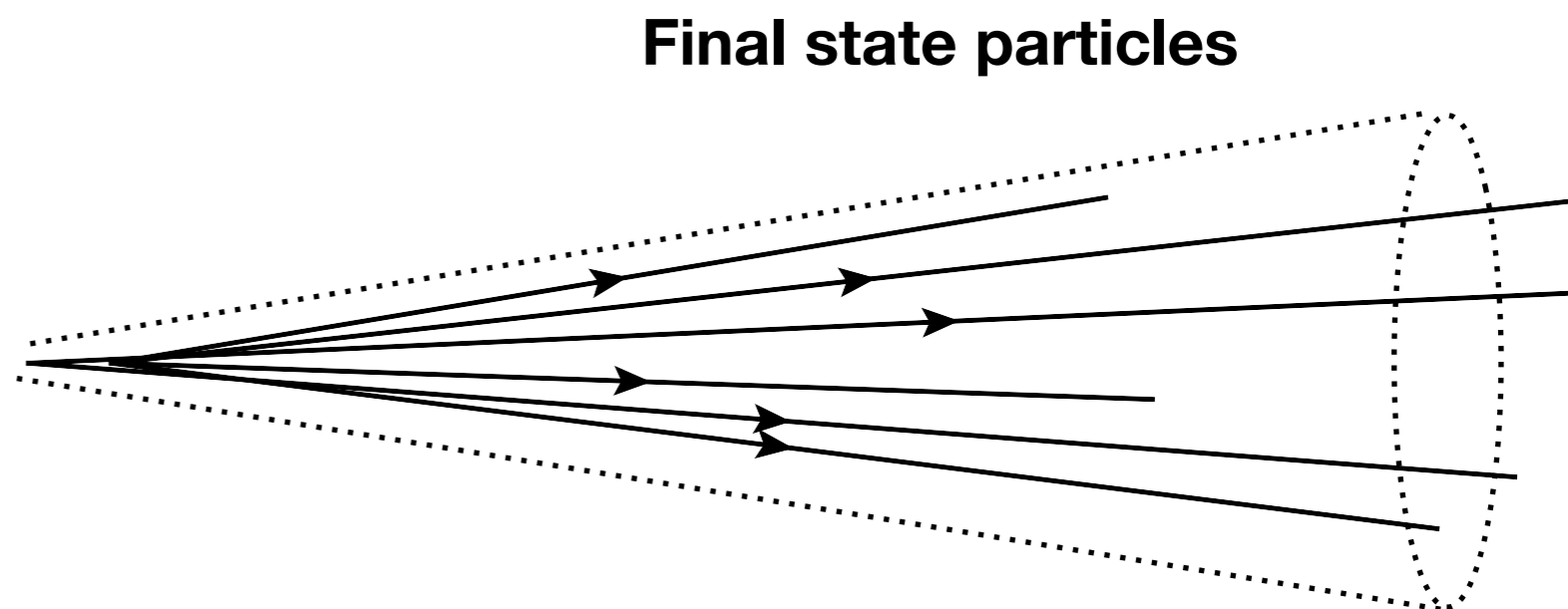
*Gaillard et. al, Nucl. Phys. B111 (1976) 253-271*



# Metamorphosis

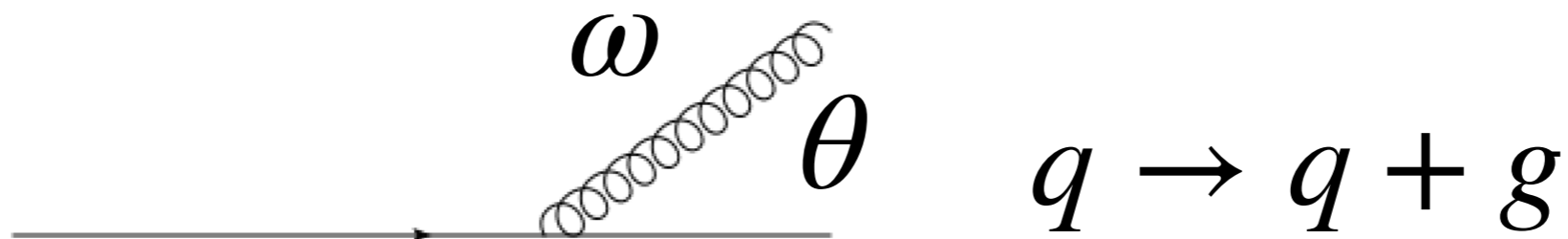
- How do we go from

High energy  
quark/gluon  
----->



- Simplified version - there are two steps here
- Parton shower followed by hadronization
- What is a parton shower?

- Lets start with the basics of QCD - splitting functions

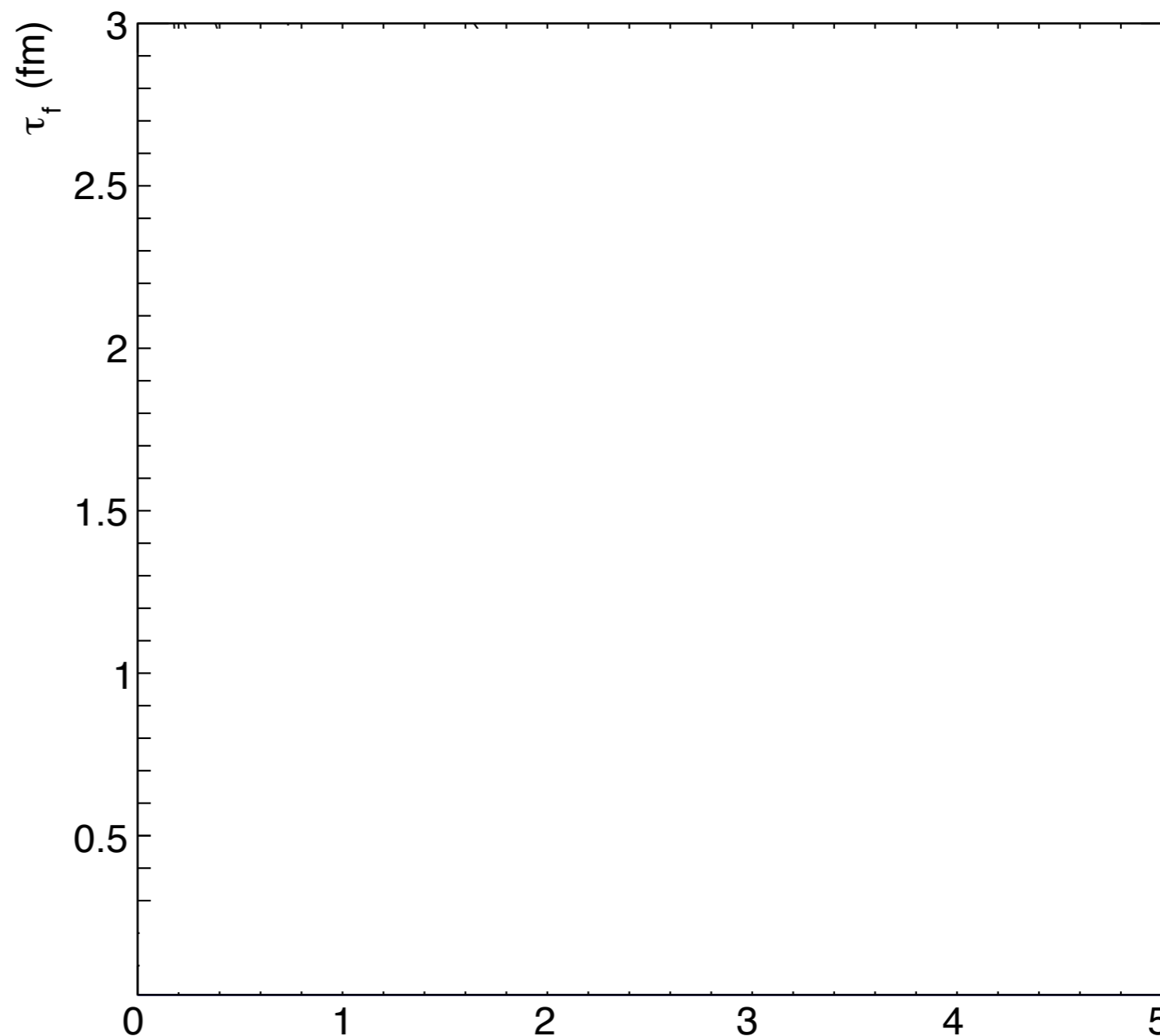


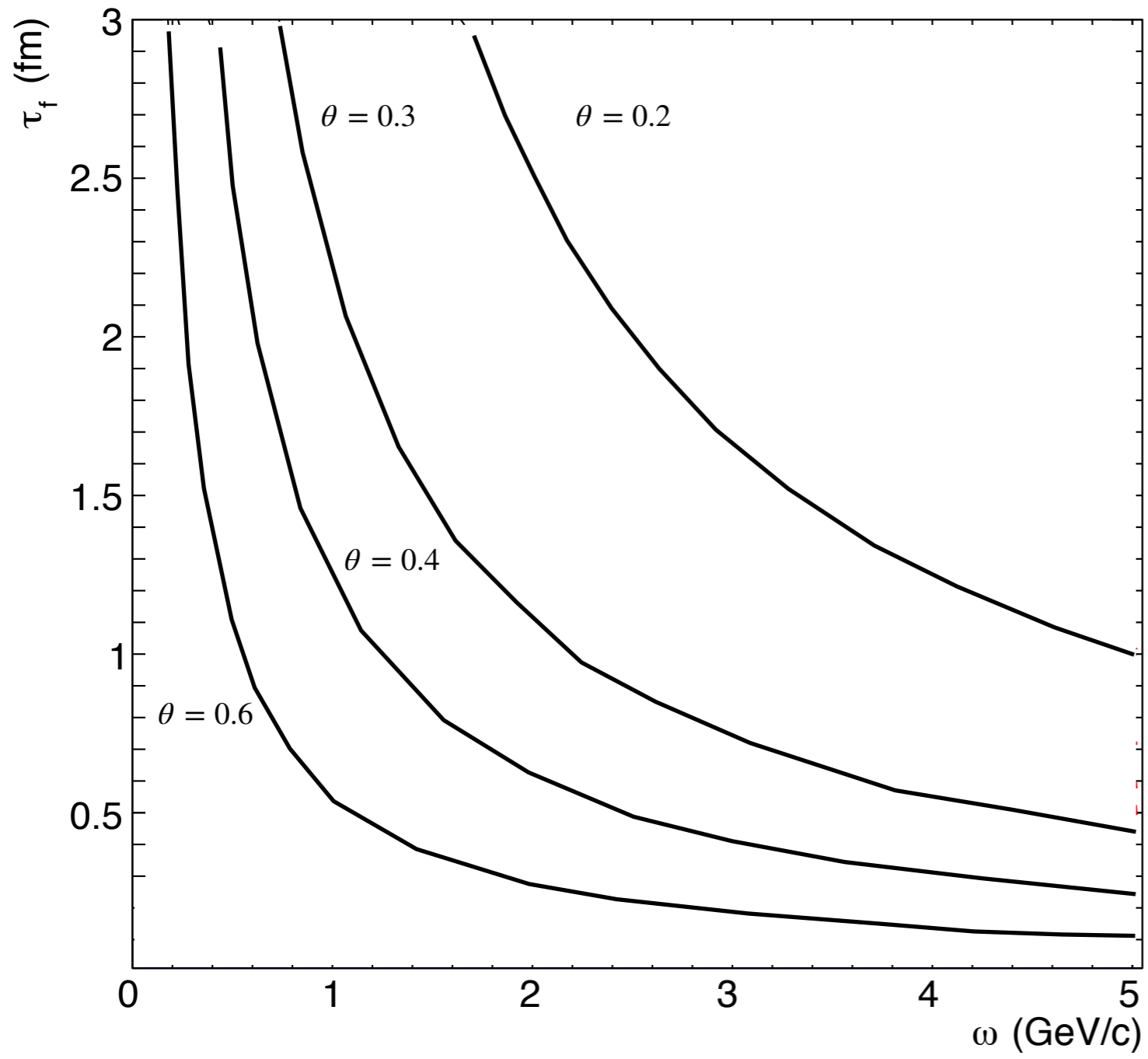
### Splitting probabilities

$$dP \approx \frac{d\theta}{\theta} \frac{d\omega}{\omega}$$

### formation times

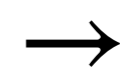
$$\tau_f^{vac} \cong \frac{\omega}{k_T^2} = \frac{1}{\theta^2 \omega}$$





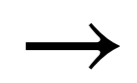
$$\tau_f^{vac} \cong \frac{\omega}{k_T^2} = \frac{1}{\theta^2 \omega}$$

Narrow emissions



Later times

Wider emissions



earlier times

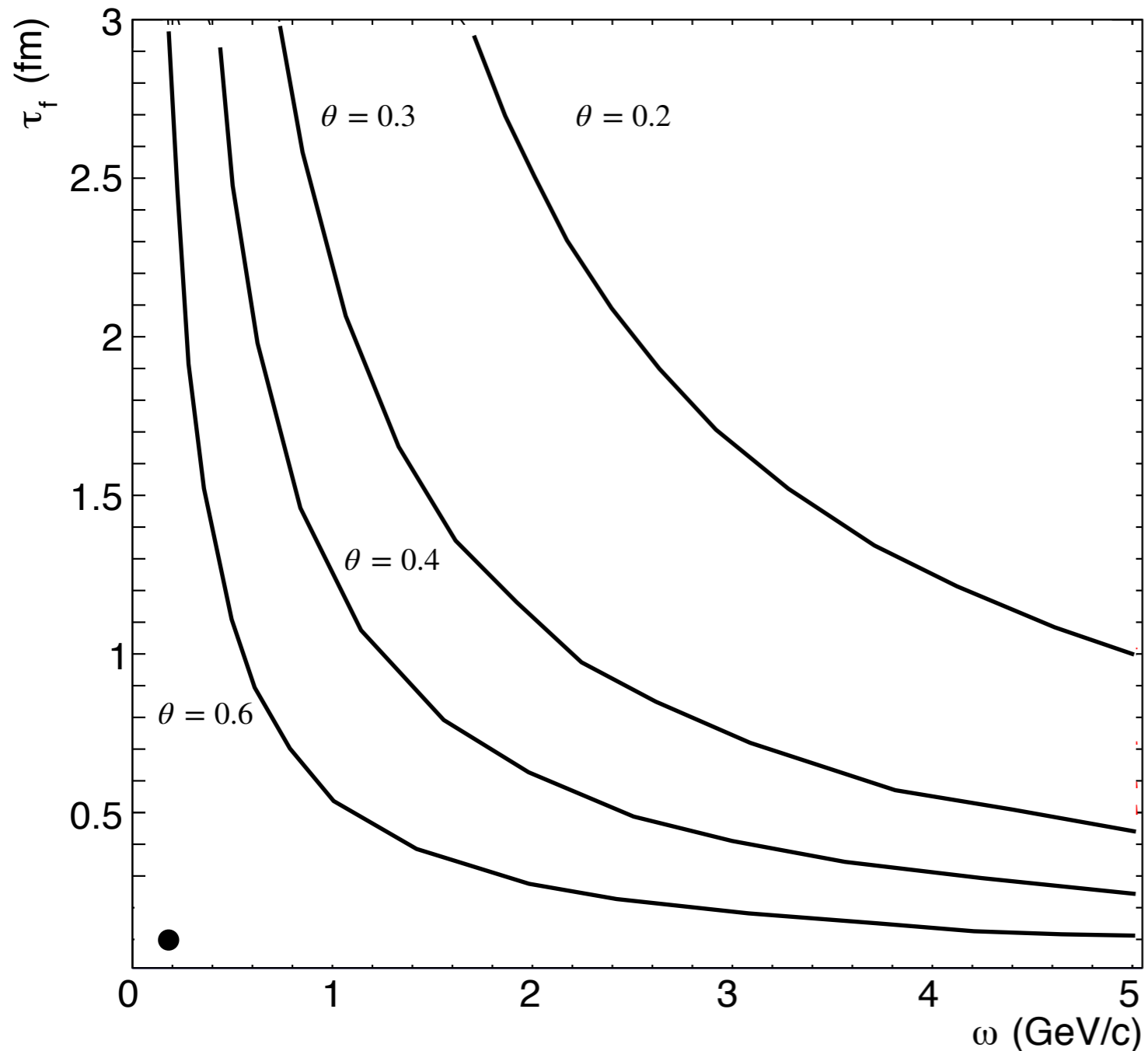
0(0.00)  $\xrightarrow{(131.08, 102.77, 1942.15, 21)}$   
13.1 GeV

**Gloun jet**

**Particle Gun in PYTHIA 8**

**Parton Shower**

*Final, 100 e+trajac*



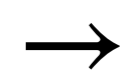
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Narrow emissions

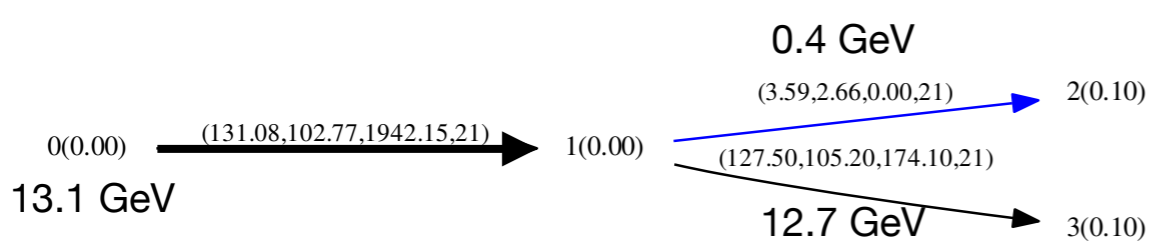


Later times

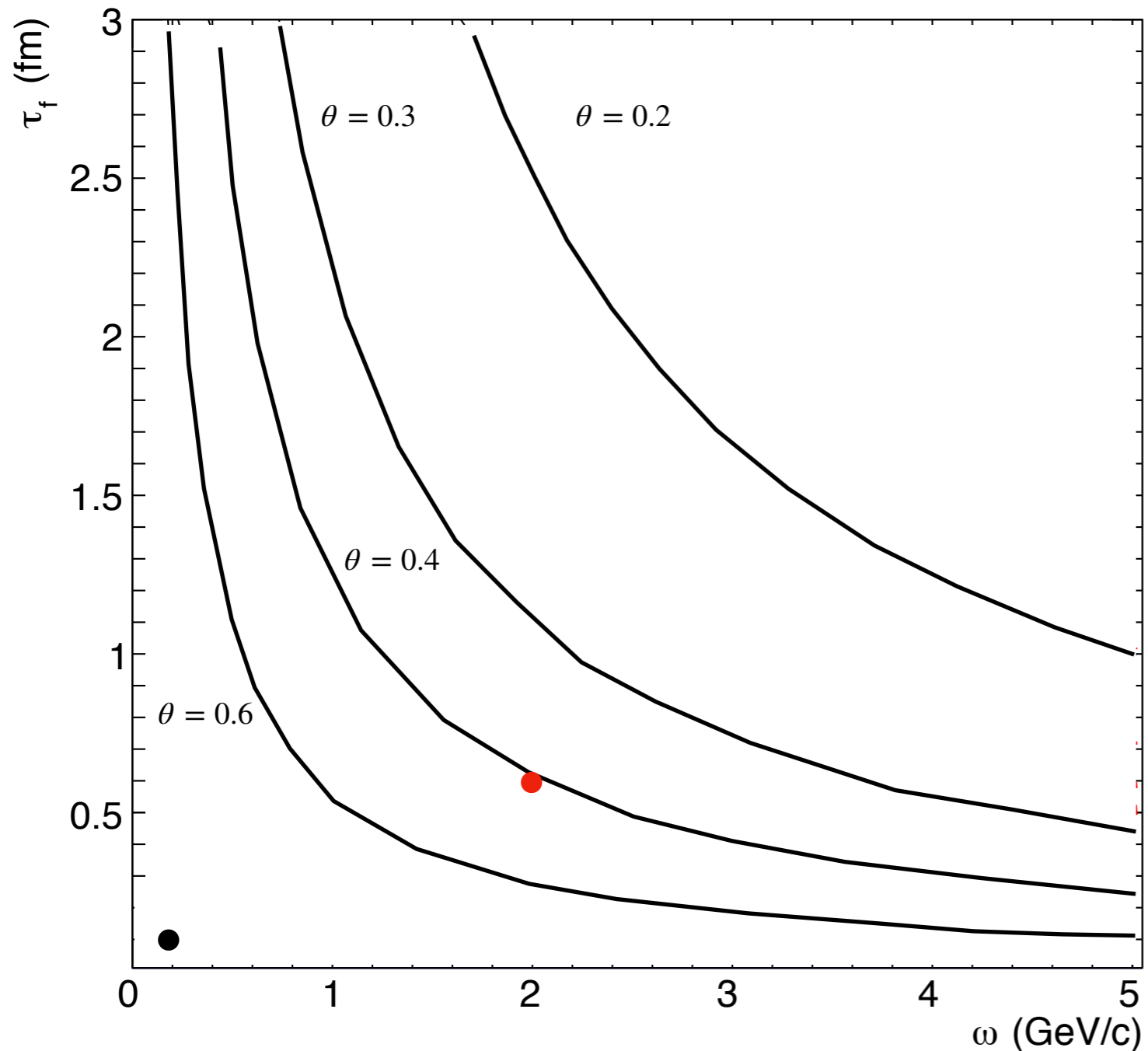
Wider emissions



earlier times



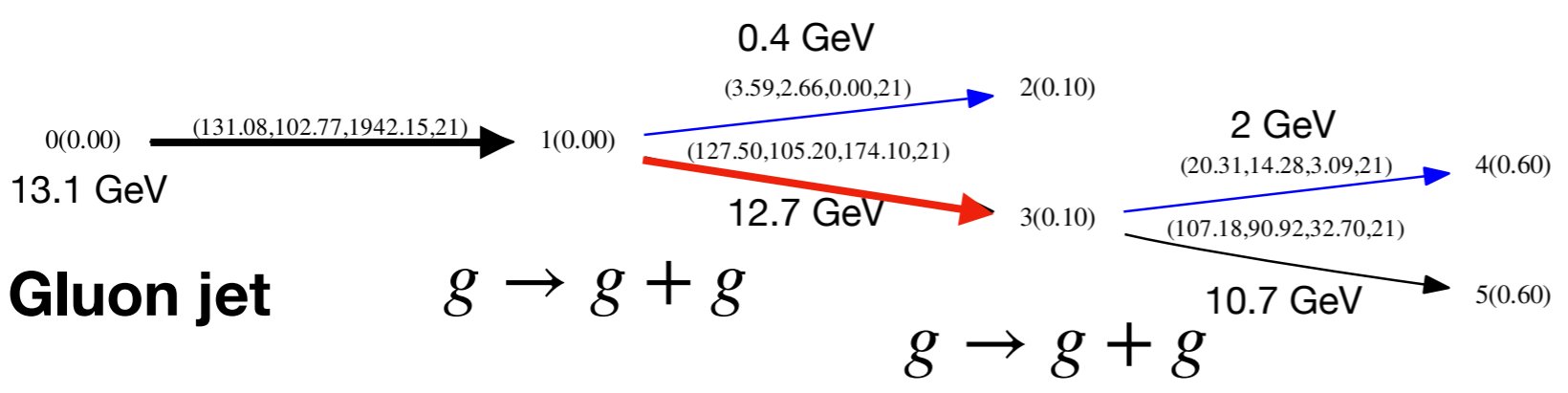
**Gluon jet**       $g \rightarrow g + g$



$$\tau_f^{vac} \cong \frac{\omega}{k_T^2} = \frac{1}{\theta^2 \omega}$$

Narrow emissions  
 →  
 Later times

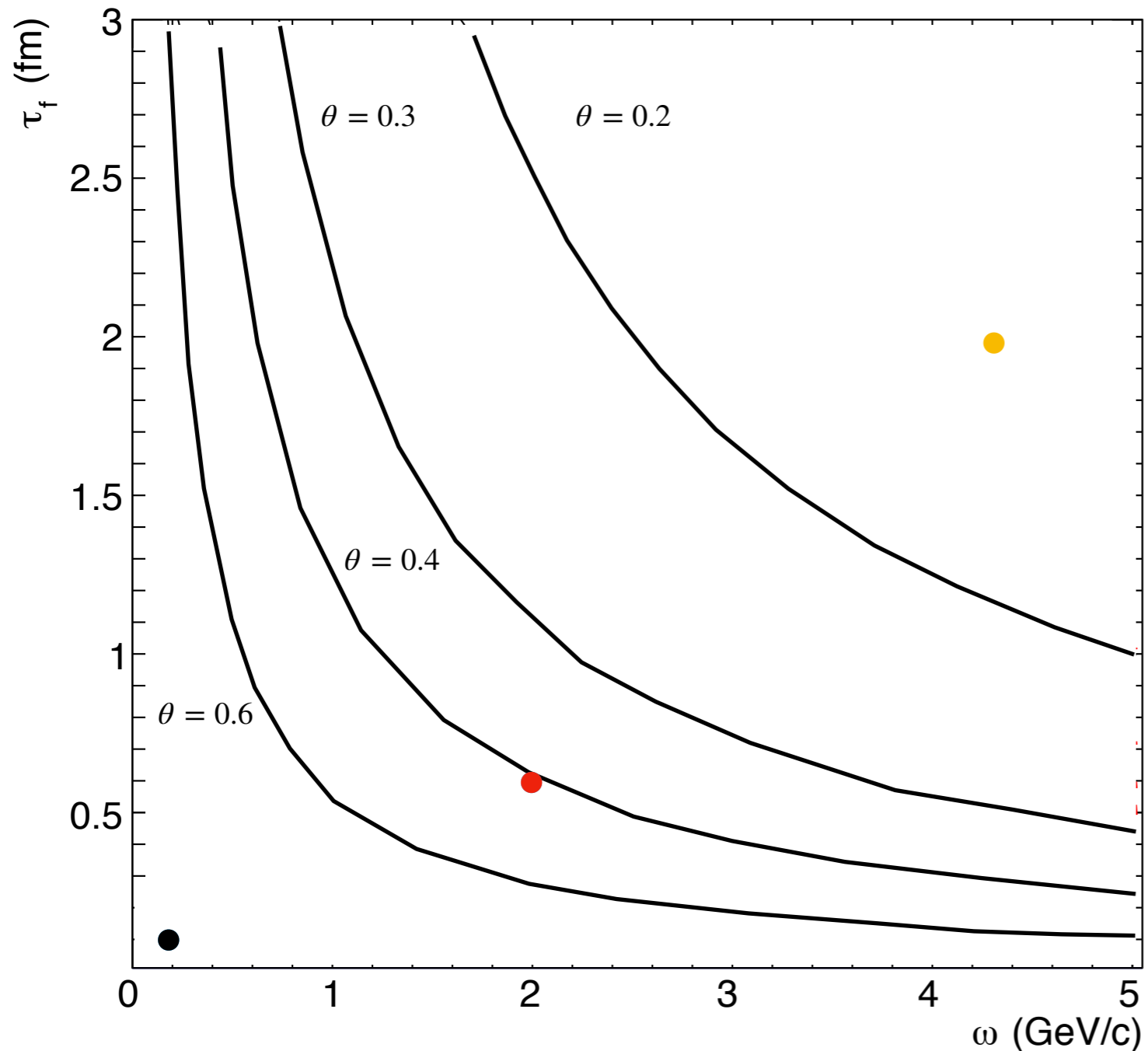
Wider emissions  
 →  
 earlier times



**Particle Gun in PYTHIA 8**

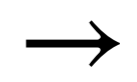
**Parton Shower**





$$\tau_f^{vac} \cong \frac{\omega}{k_T^2} = \frac{1}{\theta^2 \omega}$$

Narrow emissions

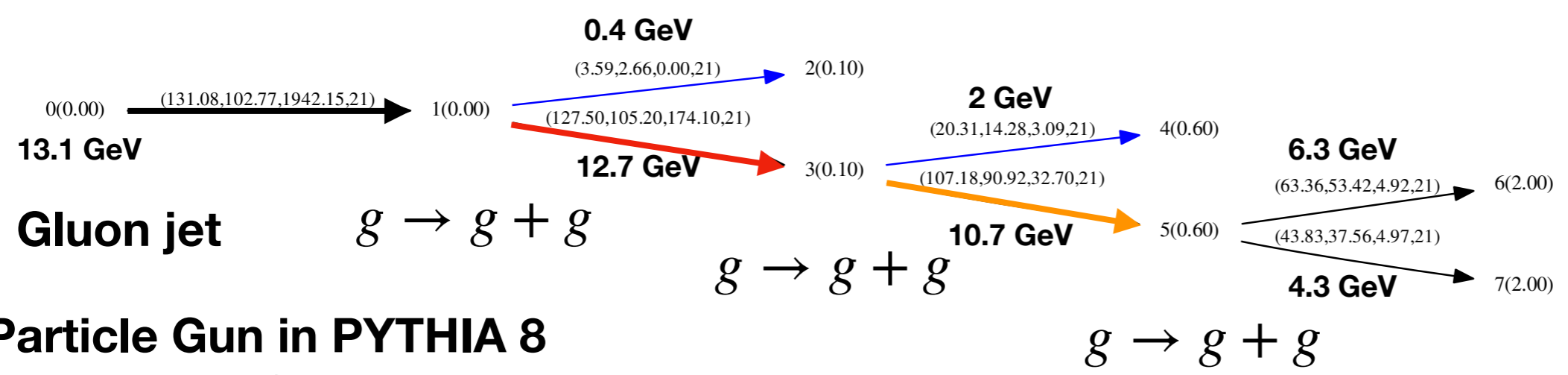


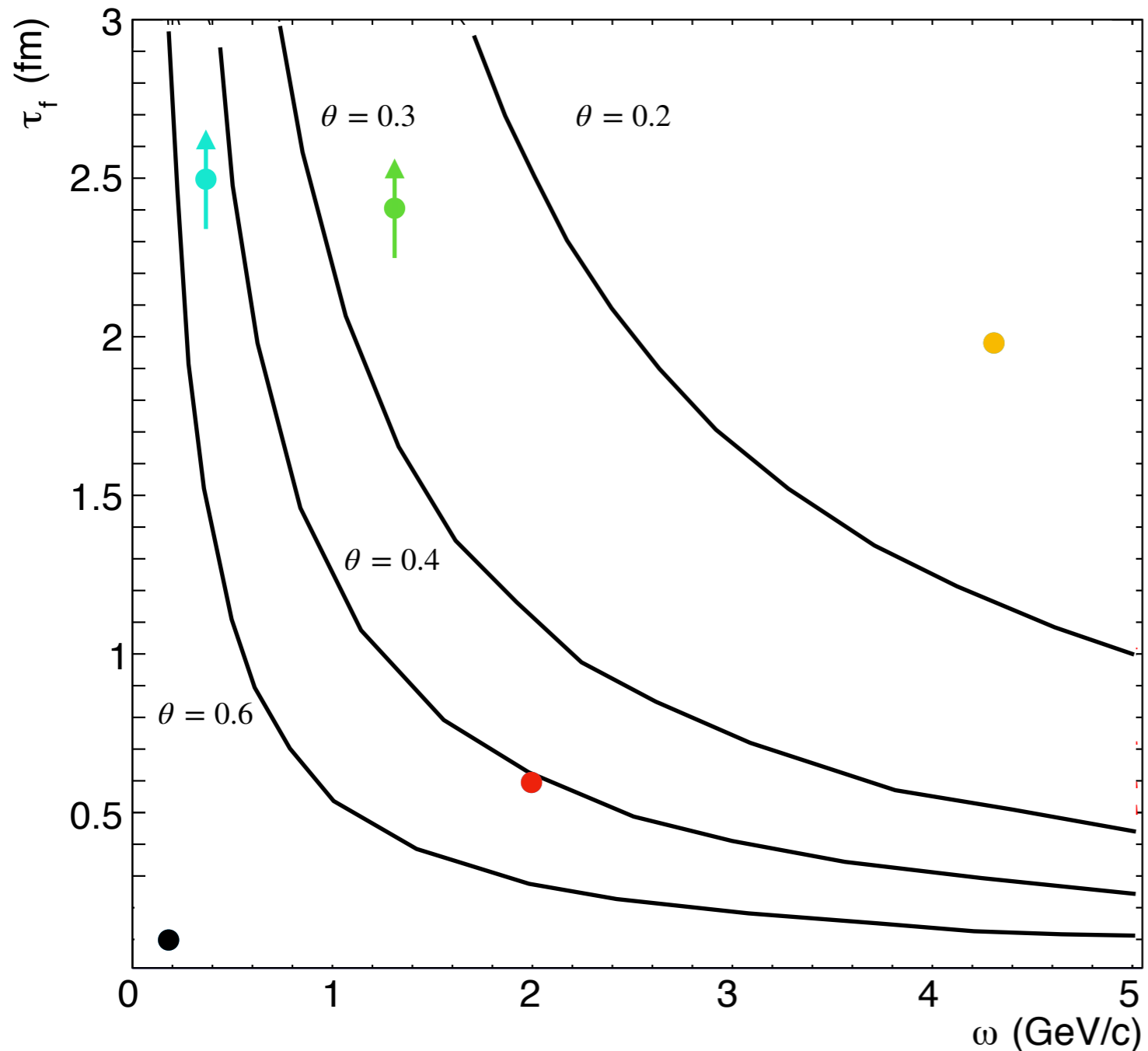
Later times

Wider emissions



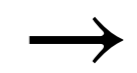
earlier times





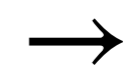
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Narrow emissions

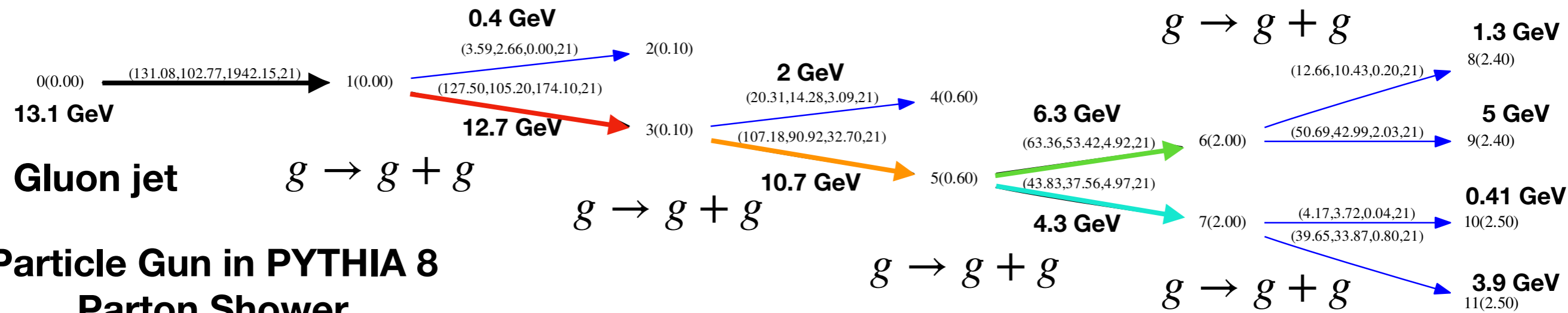


Later times

Wider emissions



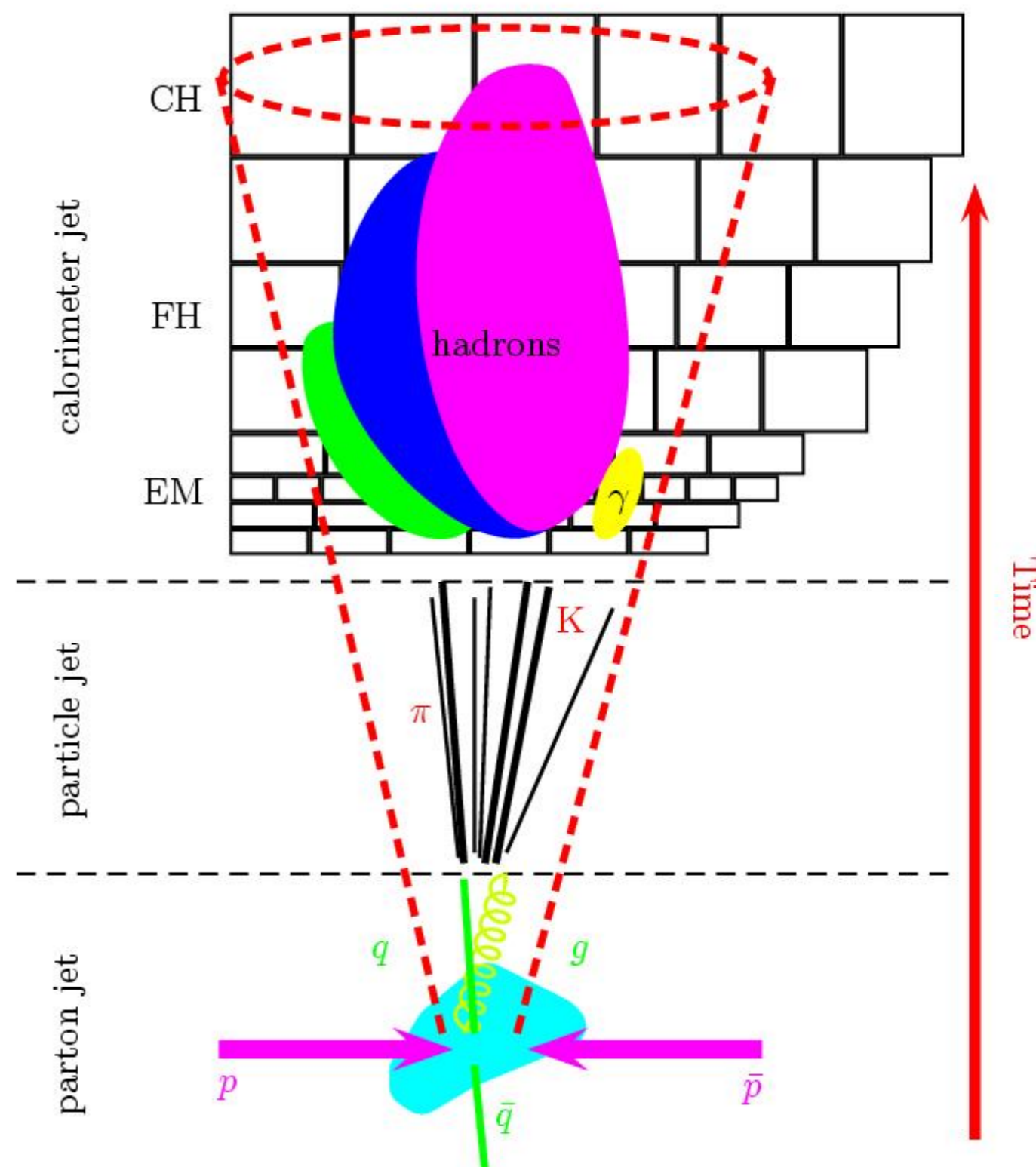
earlier times



# Jets in experiment

- Jets are defined by a clustering algorithm, **characteristic angular scale (Jet R)** and their **momentum (p<sub>T</sub>)**
- Jets can also have an **inherent angular scale (θ)** and a **momentum fraction (z)** due to vacuum splitting

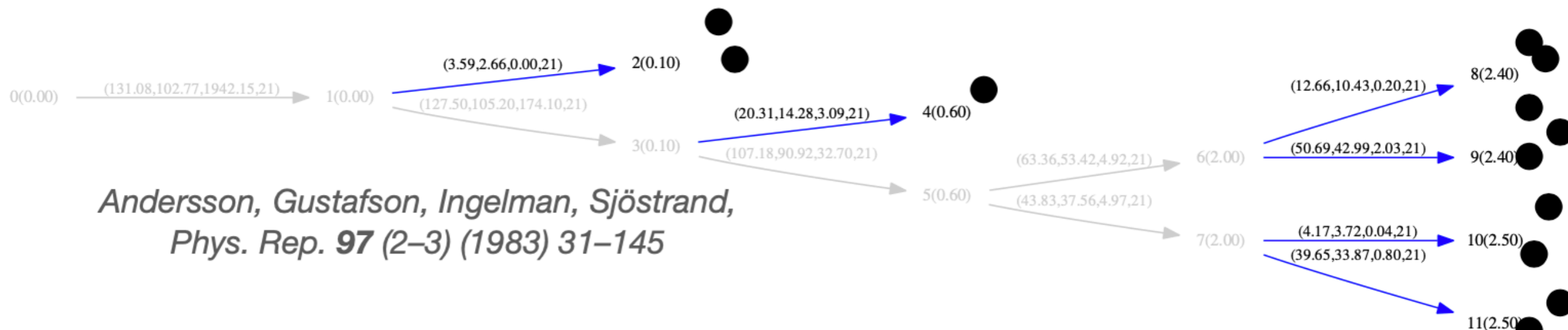
$$dP \approx \frac{dz}{z} \frac{d\theta}{\theta}$$



# What do we want to measure?

- We want to translate an *intrinsic* (and unmeasurable) parton shower to **experimentally accessible** observable(s)
- We start with the hadronized remnants of the parton shower

Lund string model decided that the 6 partons converted into 12 final-state particles



Andersson, Gustafson, Ingelman, Sjöstrand,  
Phys. Rep. **97** (2–3) (1983) 31–145

# What do we want to measure?

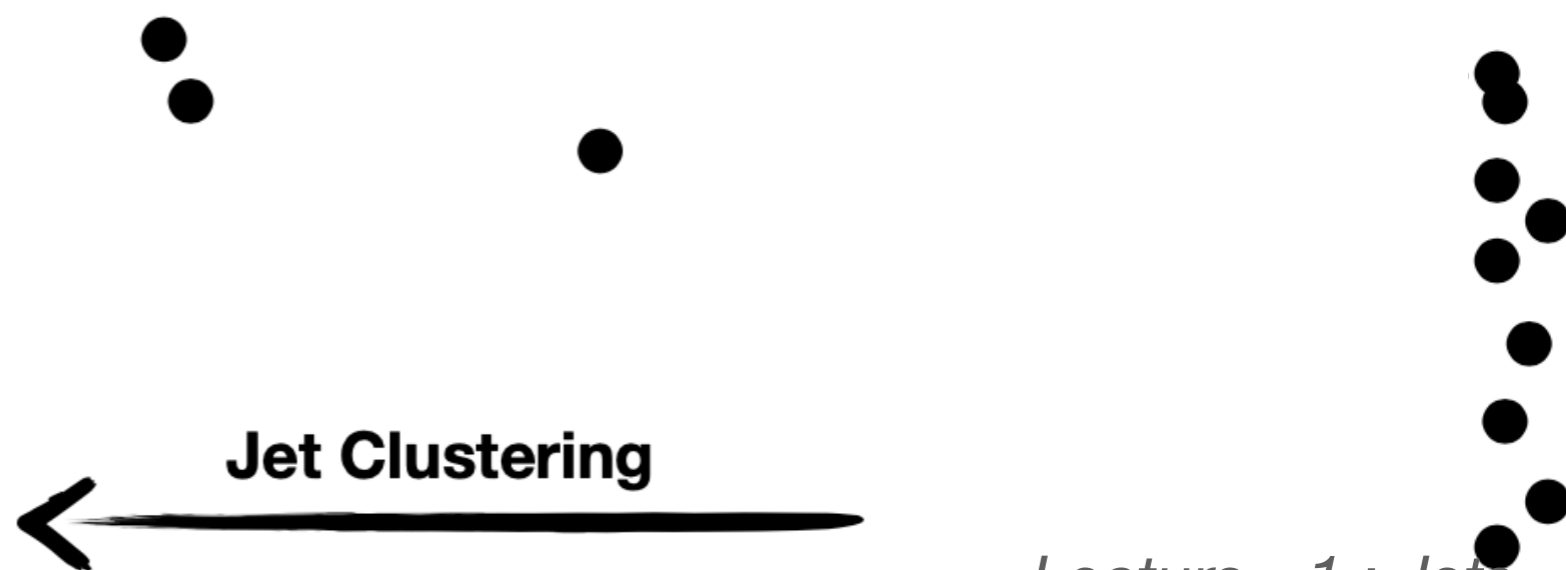
- We want to translate an *intrinsic* (and unmeasurable) parton shower to **experimentally accessible** observable(s)
- Run jet clustering algorithms starting from the final state hadrons based on momentum weighted distance criterion

$$\min(p_{ti}^{2p}, p_{tj}^{2p}) \Delta R_{ij}^2 / R^2$$

*FastJet 3.3.4*

*Cacciari, Salam, Soyez,  
Eur. Phys. J. C 72 (2012) 1896*

*Cacciari, Salam,  
Phys. Lett. B 641 (2006) 57*



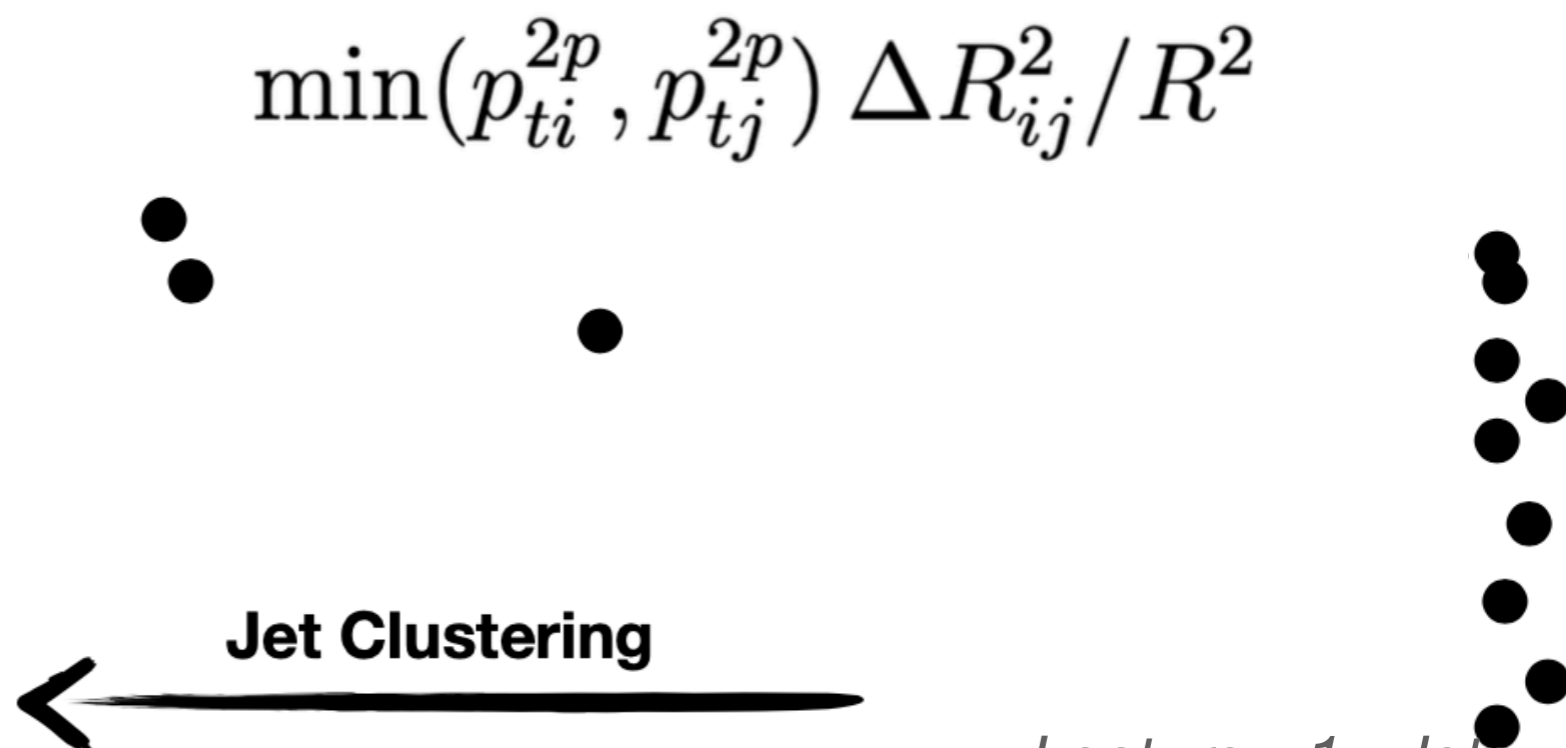
# What do we want to measure?

- We want to translate an *intrinsic* (and unmeasurable) parton shower to **experimentally accessible** observable(s)
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$$p = 1 \quad k_T$$

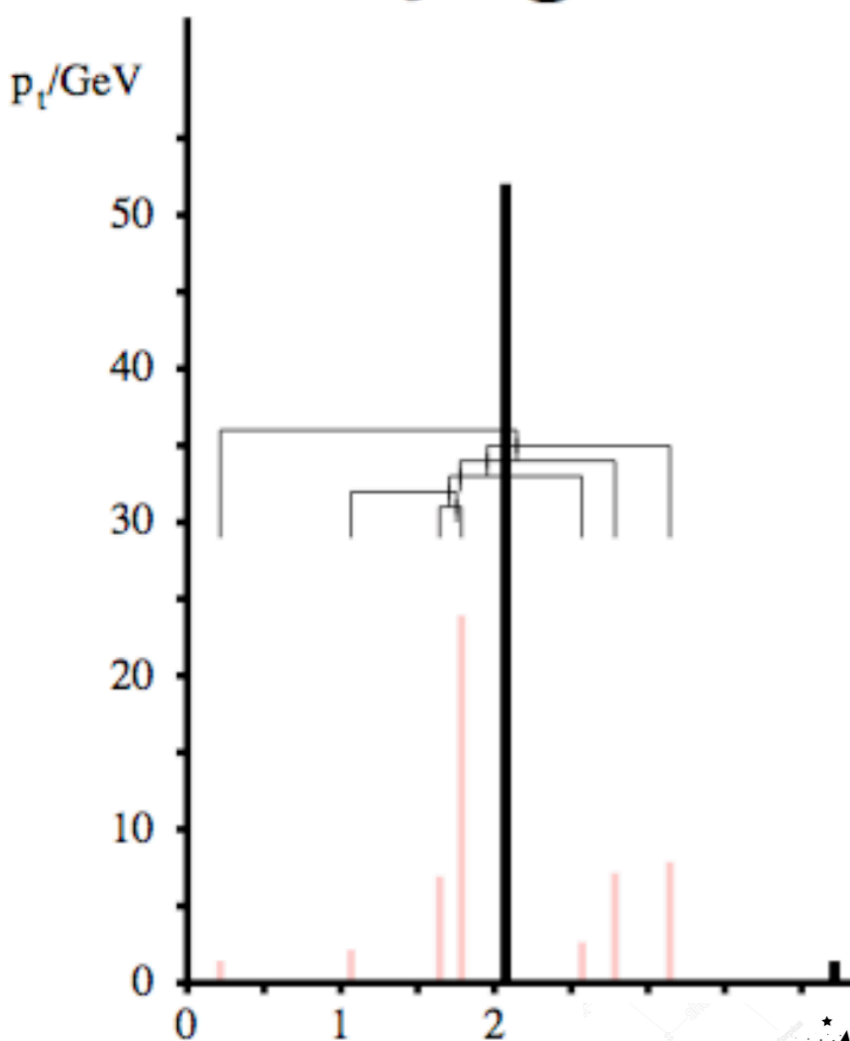
$$p = 0 \quad C/A$$

$$p = -1 \quad \text{anti-}k_T$$

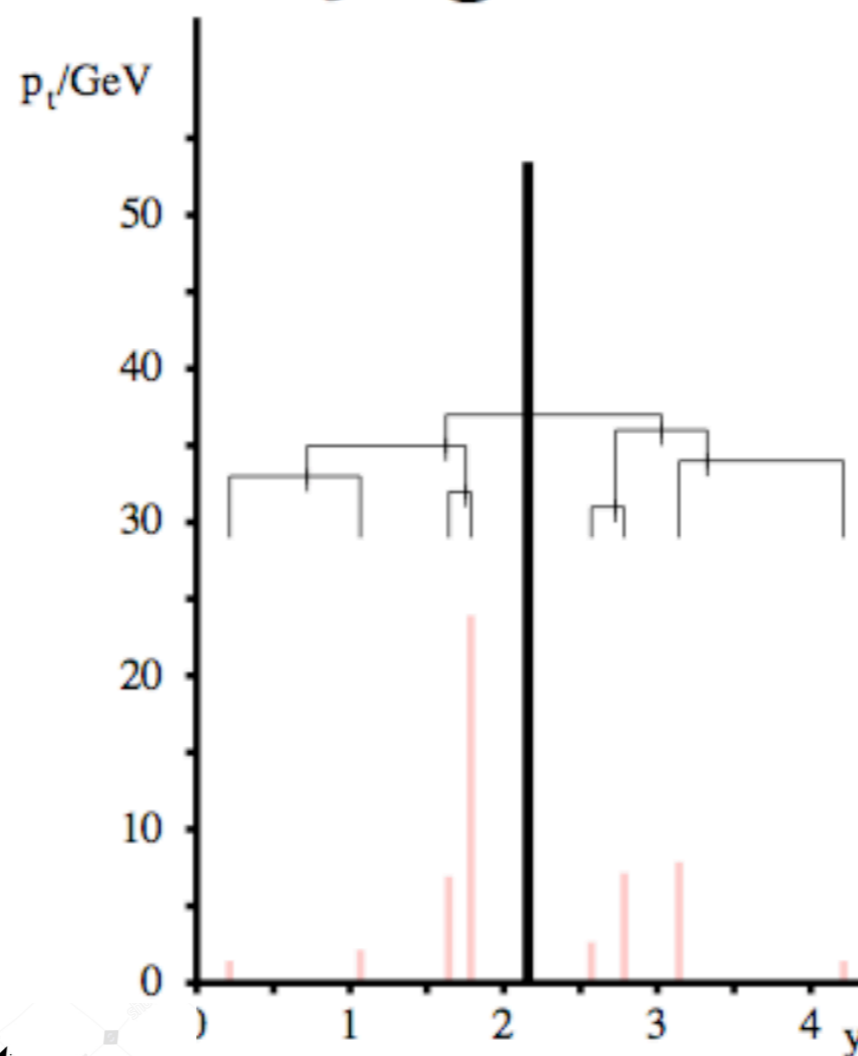


# Clustering algorithms

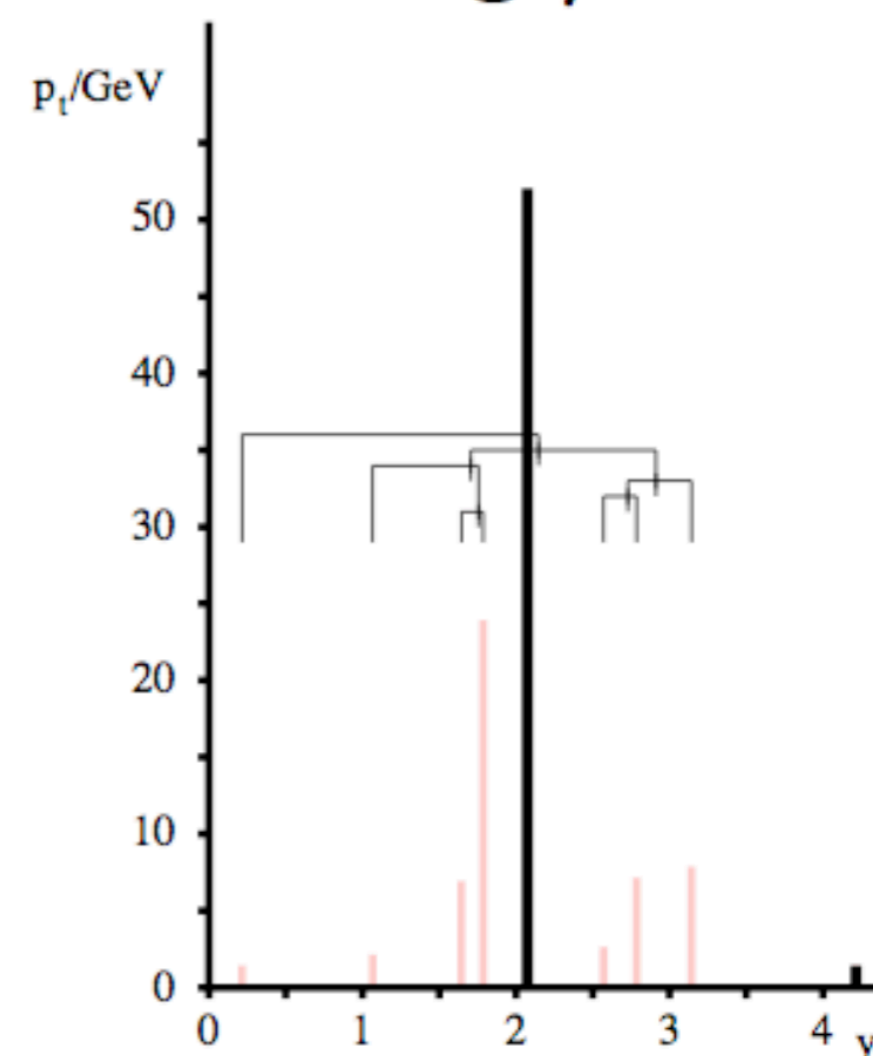
anti- $k_t$  algorithm



$k_t$  algorithm



Cambridge/Aachen

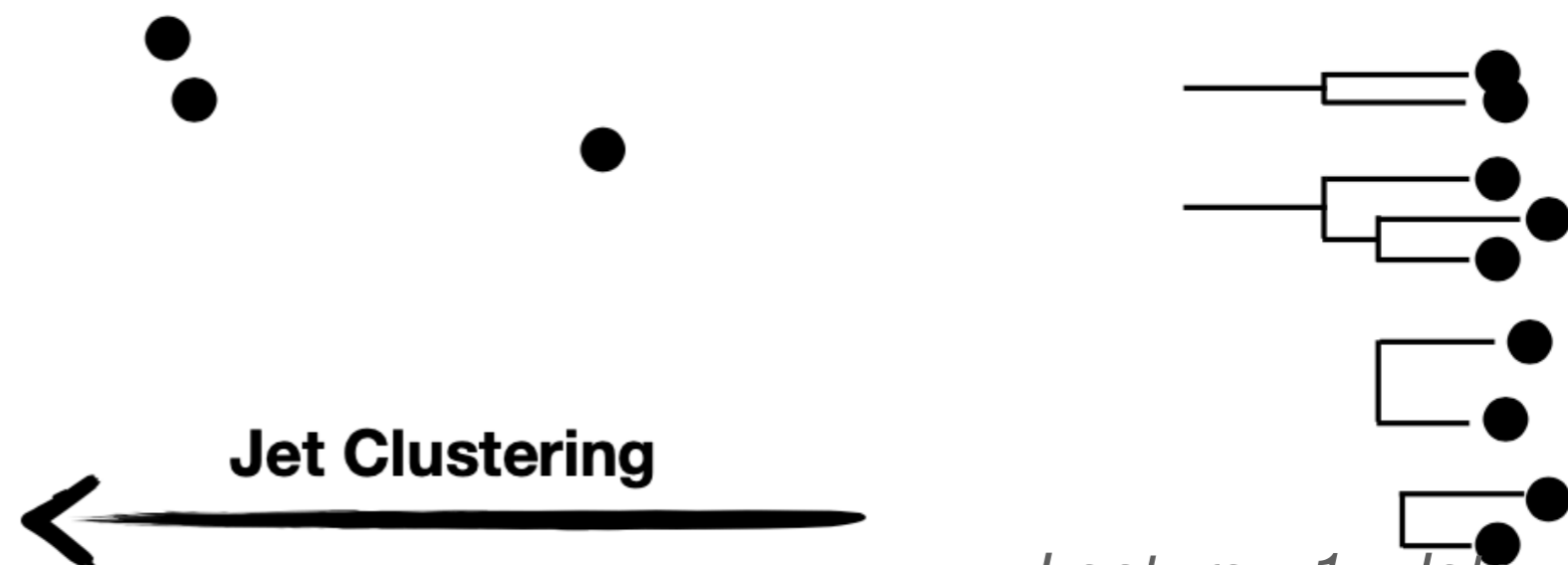


Which of the clustering trees is 'realistic'?

# What do we want to measure?

- We want to translate an *intrinsic* (and unmeasurable) parton shower to **experimentally accessible** observable(s)
- Run jet clustering algorithms starting from the final state hadrons based on momentum weighted distance criterion

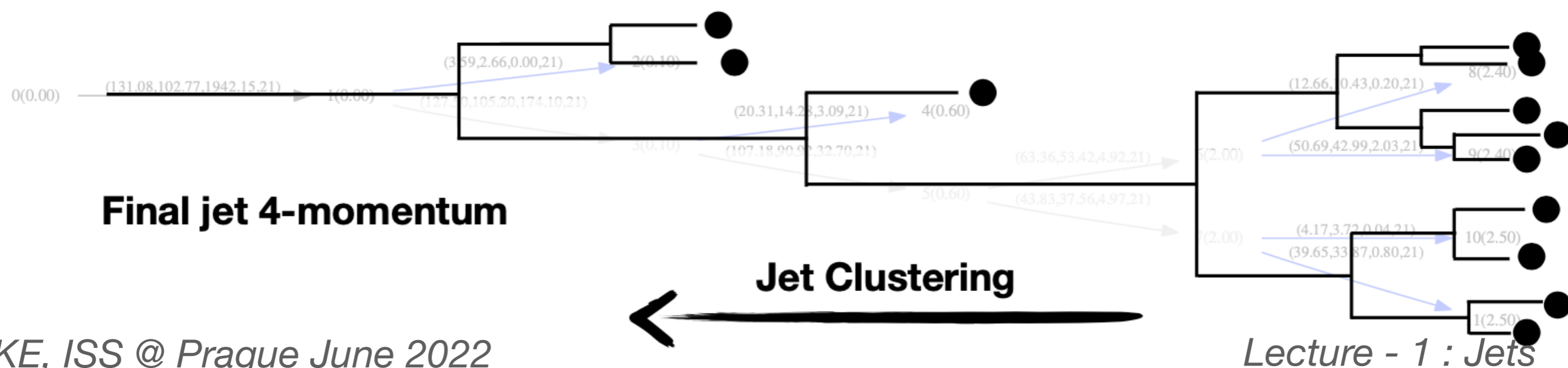
$$\min(p_{ti}^{2p}, p_{tj}^{2p}) \Delta R_{ij}^2 / R^2$$





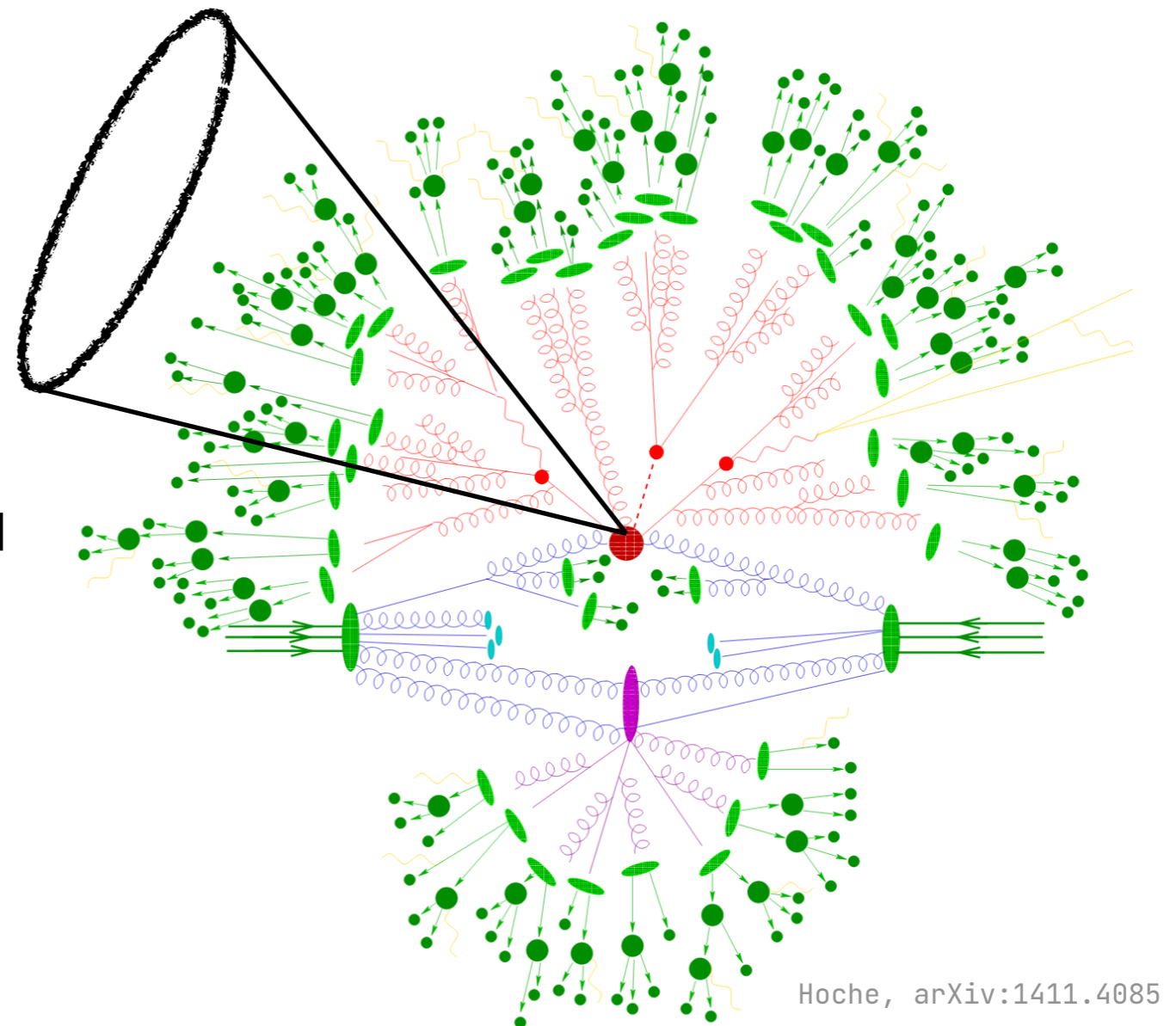
# What do we want to measure?

- We want to translate an *intrinsic* (and unmeasurable) parton shower to **experimentally accessible** observable(s)
- Run jet clustering algorithms starting from the final state hadrons based on momentum weighted distance criterion

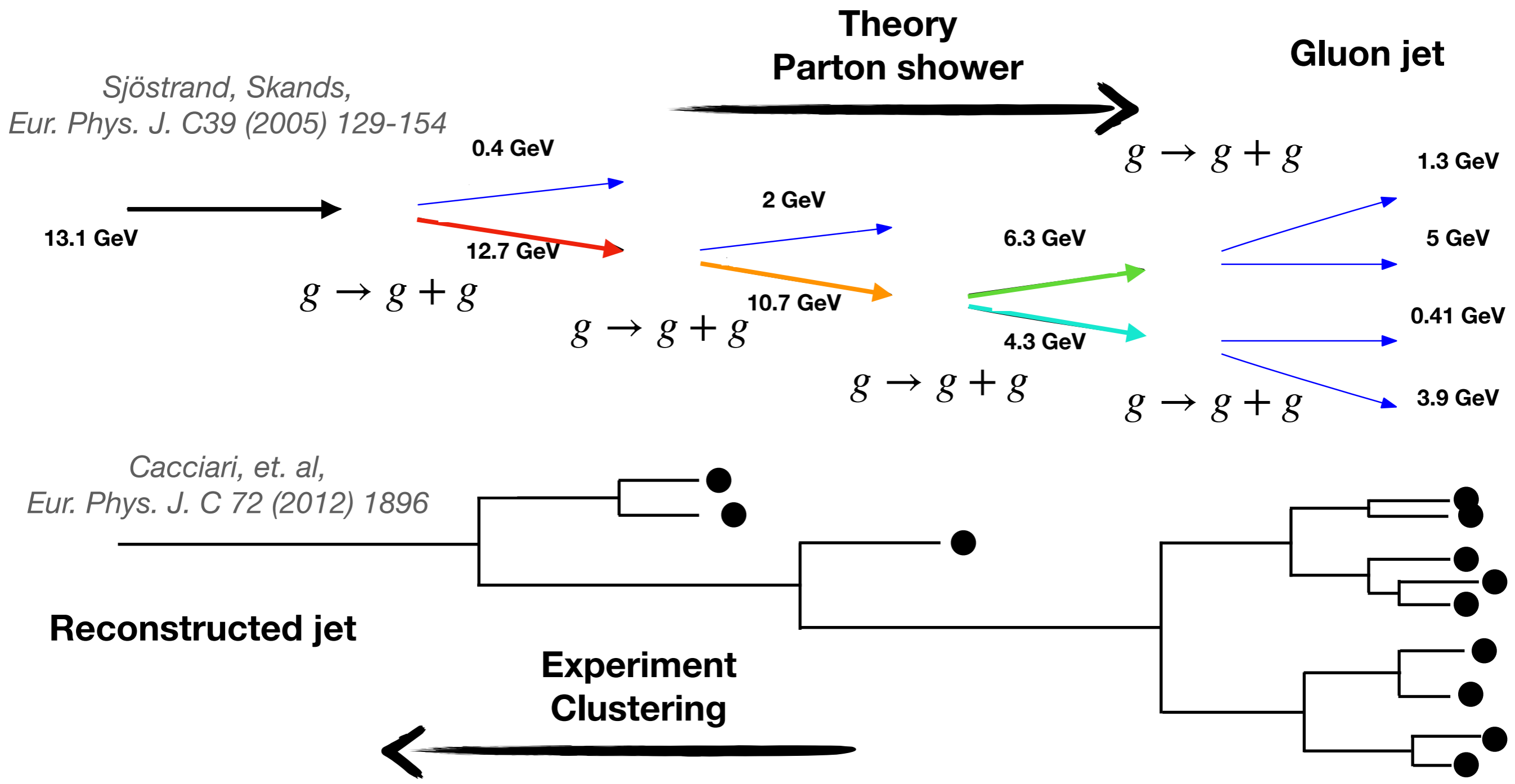


# Recap - 1

- Jet — Collection of particles grouped by an iterative algorithm
- With an appropriate jet radius and constituent momentum threshold, you hope to recover the originating parton
- Jets work as proxy to hard scattered  $q/g$

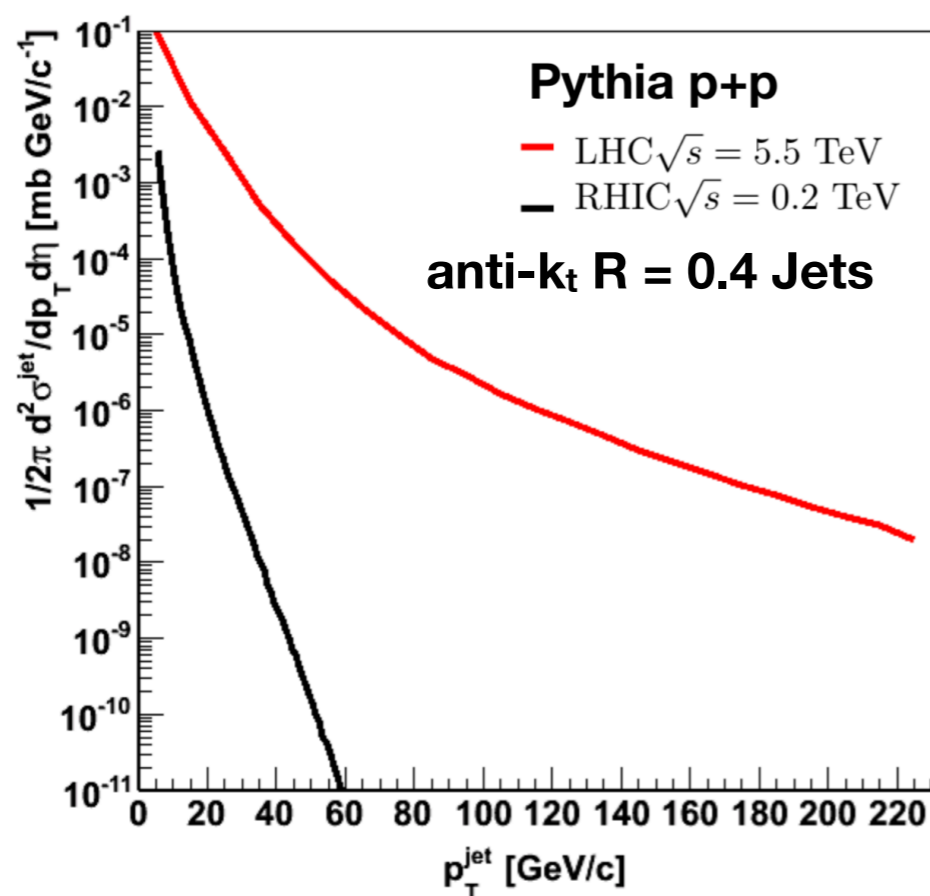


# Jet correspondence



*Cacciari, et. al, Eur. Phys. J. C 72 (2012) 1896*

# Ingredients in calculating jet spectra



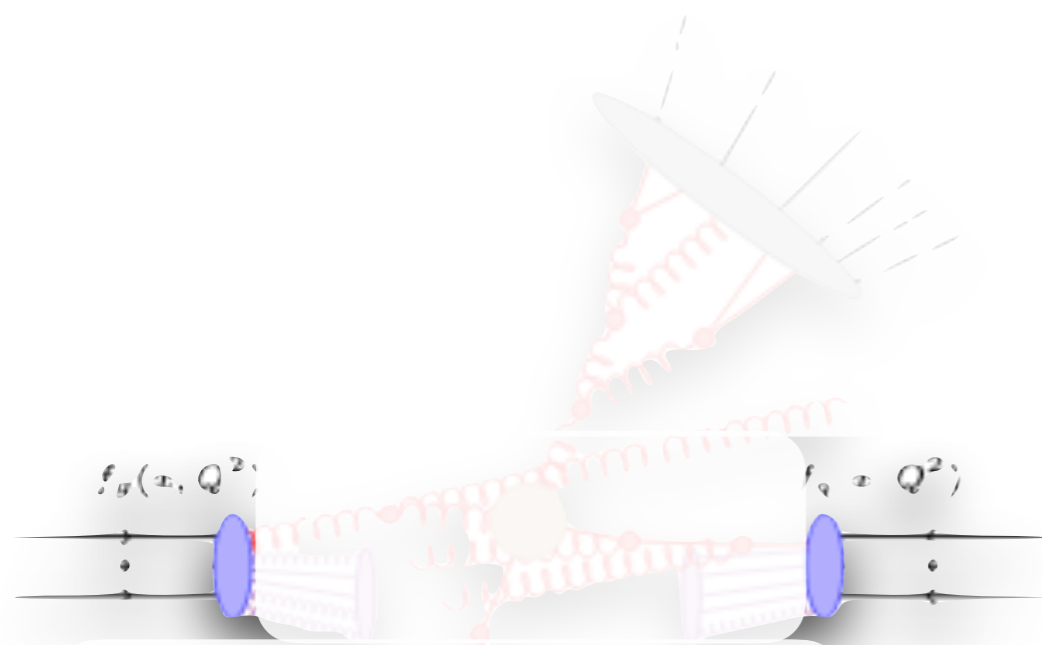
- Calculable via QCD factorization
- Steeper cross-section at RHIC compared to the LHC

Jet production in high energy collisions of hadrons can be described in terms of following ingredients:

- *Initial state of hadrons*
- *Hard collision of partons*
- *Parton Shower*
- *Underlying Event (UE)*
- *Hadronization*

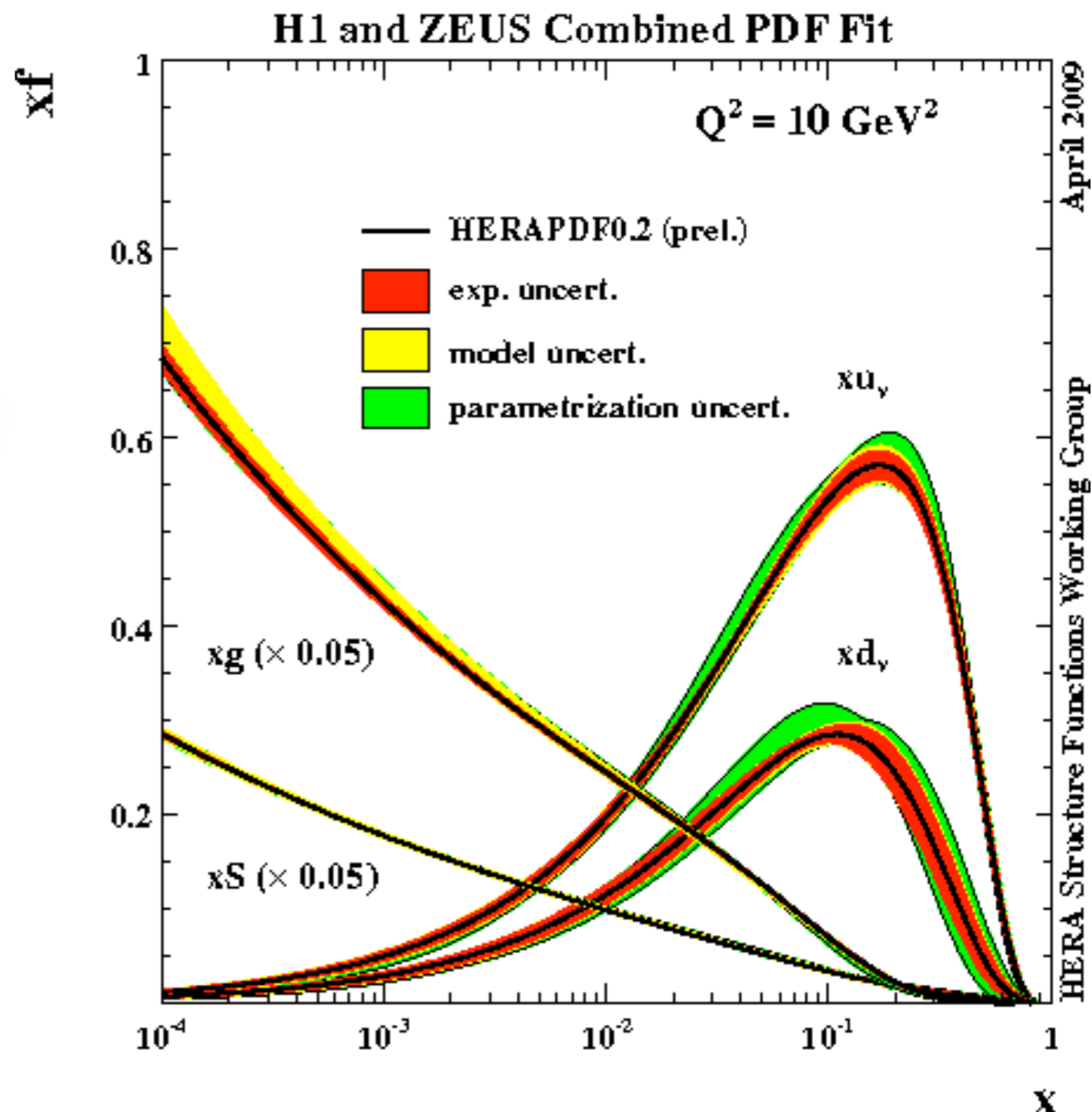
$$d\sigma_{pp \rightarrow \text{jet}+X}(Q^2) = \sum_{a,b} \int \underbrace{f_a(x_1, Q^2) f_b(x_2, Q^2)}_{\text{proton structure}} \underbrace{d\hat{\sigma}_{a+b \rightarrow \text{jet}+X}(x_1, x_2, Q^2)}_{\text{hard process+PS+Had.}} dx_1 dx_2$$

$$d\sigma_{pp \rightarrow \text{jet}+X}(Q^2) = \sum_{a,b} \int \underbrace{f_a(x_1, Q^2) f_b(x_2, Q^2)}_{\text{proton structure}} \underbrace{d\hat{\sigma}_{a+b \rightarrow \text{jet}+X}(x_1, x_2, Q^2)}_{\text{hard process+PS+Had.}} dx_1 dx_2$$



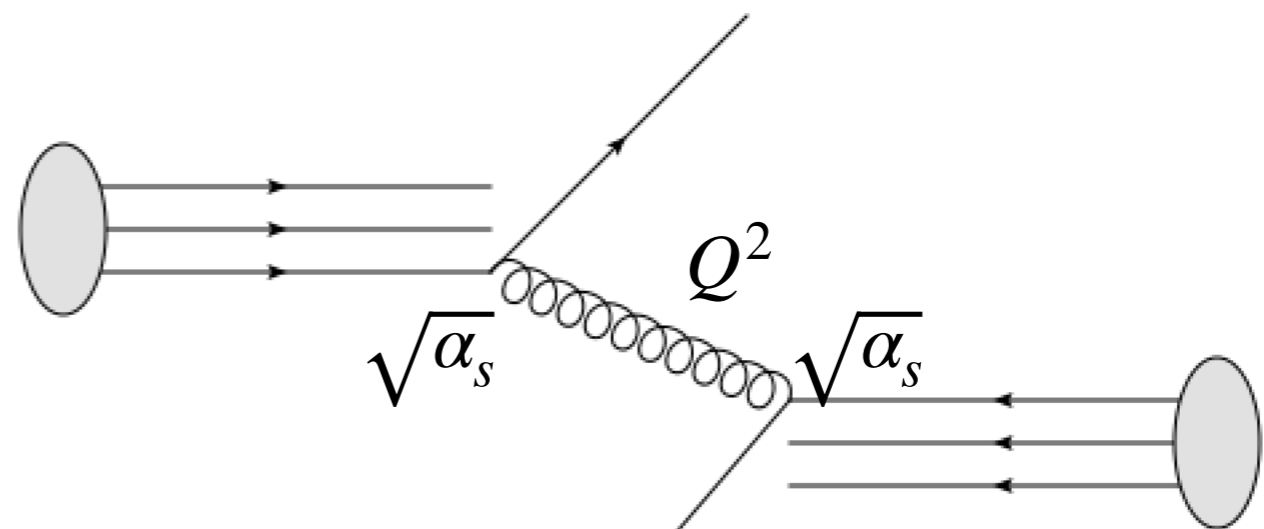
# Parton Distribution Functions

$$f_a(x, Q^2)$$



$$d\sigma_{pp \rightarrow \text{jet}+X}(Q^2) = \sum_{a,b} \int \underbrace{f_a(x_1, Q^2) f_b(x_2, Q^2)}_{\text{proton structure}} \underbrace{d\hat{\sigma}_{a+b \rightarrow \text{jet}+X}(x_1, x_2, Q^2)}_{\text{hard process} + \text{PS} + \text{Had.}} dx_1 dx_2$$

# Hard process - matrix element calculation



$$M \propto \frac{\alpha_s}{Q^2}$$

$$d\sigma_{pp \rightarrow \text{jet}+X}(Q^2) = \sum_{a,b} \int \underbrace{f_a(x_1, Q^2) f_b(x_2, Q^2)}_{\text{proton structure}} \underbrace{d\hat{\sigma}_{a+b \rightarrow \text{jet}+X}(x_1, x_2, Q^2)}_{\text{hard process} + \text{PS} + \text{Had.}} dx_1 dx_2$$

# Parton Shower

$$\frac{d}{d \ln Q^2} \begin{pmatrix} q \\ g \end{pmatrix} = \begin{pmatrix} P_{q \leftarrow q} & P_{q \leftarrow g} \\ P_{g \leftarrow q} & P_{g \leftarrow g} \end{pmatrix} \otimes \begin{pmatrix} q \\ g \end{pmatrix}$$

Splitting functions are:

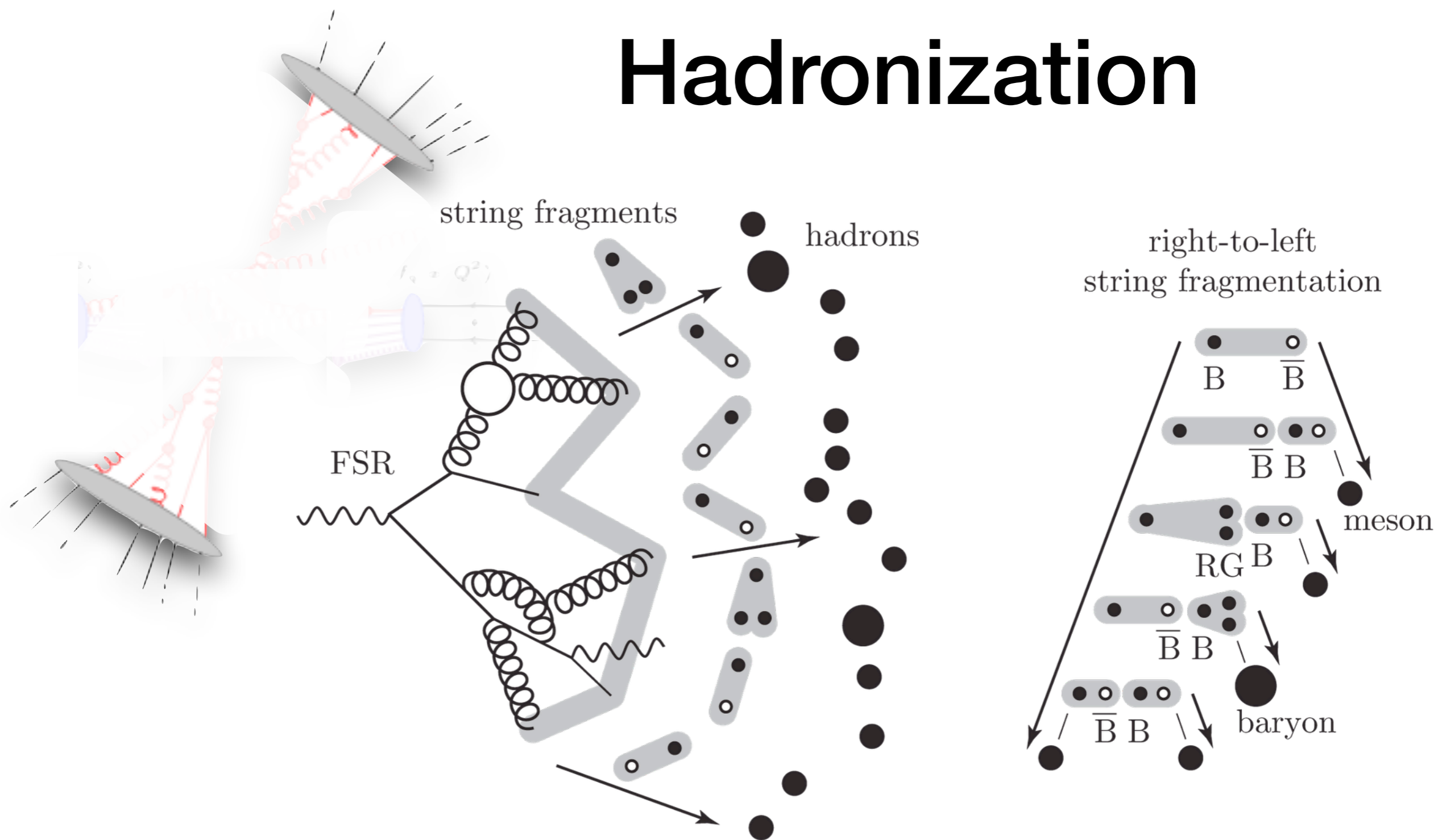
$$P_{qg}(z) = T_R [z^2 + (1-z)^2],$$

$$P_{gg}(z) = 2C_A \left[ \frac{z}{(1-z)_+} + \frac{1-z}{z} + z(1-z) \right] + \delta(1-z) \frac{(11C_A - 4n_f T_R)}{6}.$$

$$P_{gq}(z) = C_F \left[ \frac{1 + (1-z)^2}{z} \right]$$

$$d\sigma_{pp \rightarrow \text{jet}+X}(Q^2) = \sum_{a,b} \int \underbrace{f_a(x_1, Q^2) f_b(x_2, Q^2)}_{\text{proton structure}} \underbrace{d\hat{\sigma}_{a+b \rightarrow \text{jet}+X}(x_1, x_2, Q^2)}_{\text{hard process} + \text{PS} + \text{Had.}} dx_1 dx_2$$

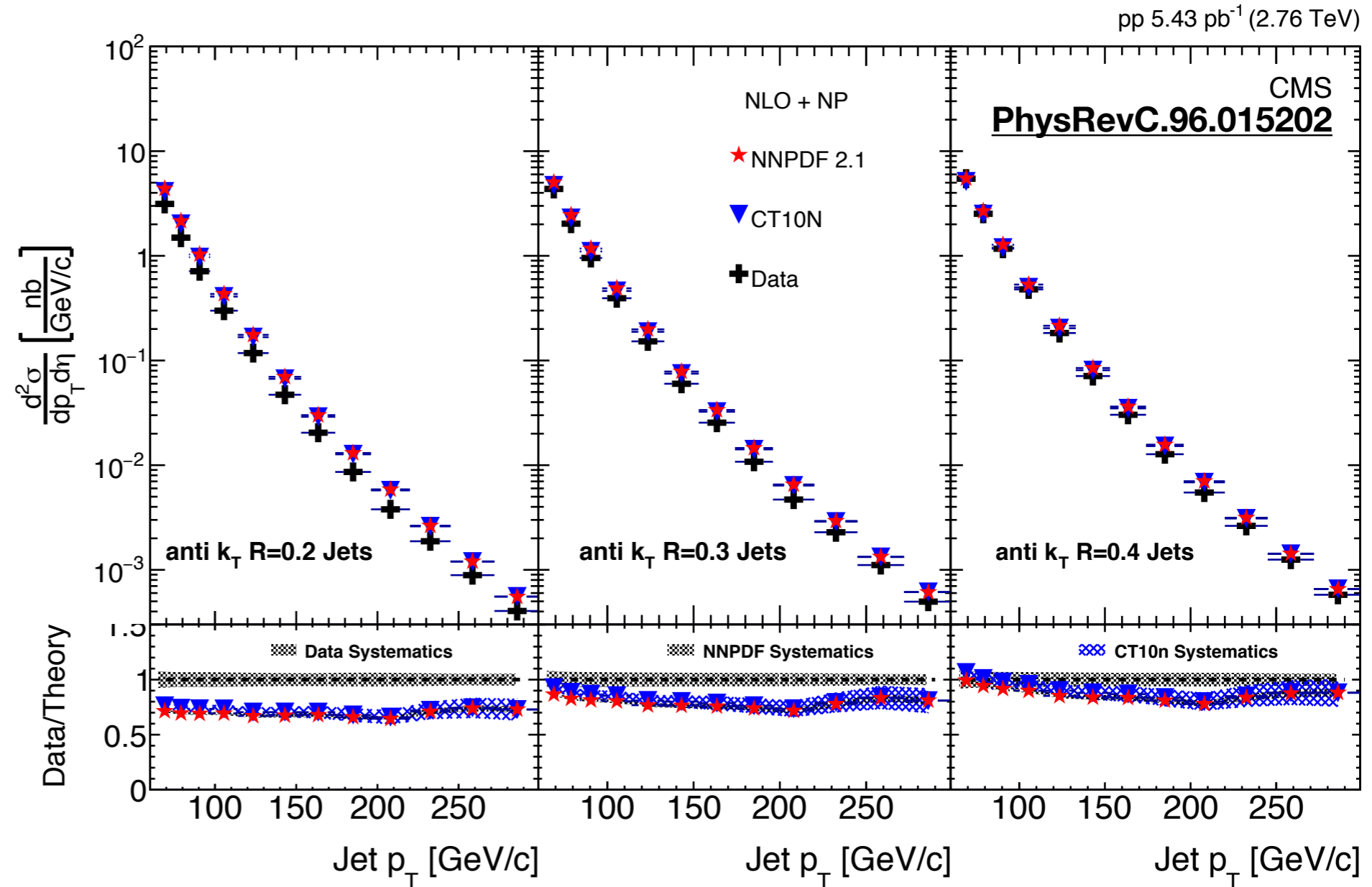
# Hadronization



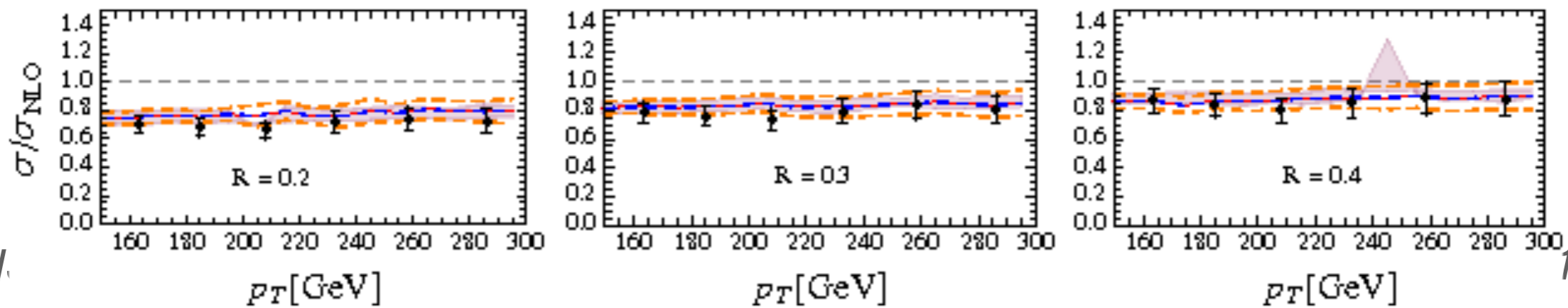


# pp jet Spectra @ LHC

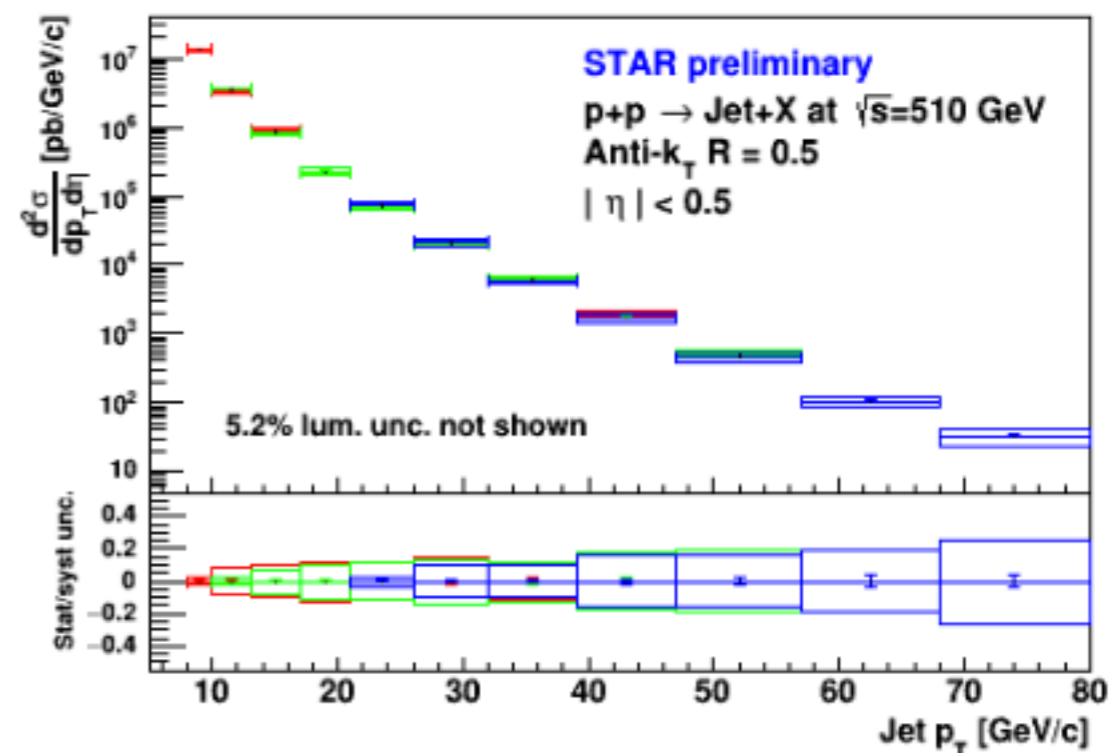
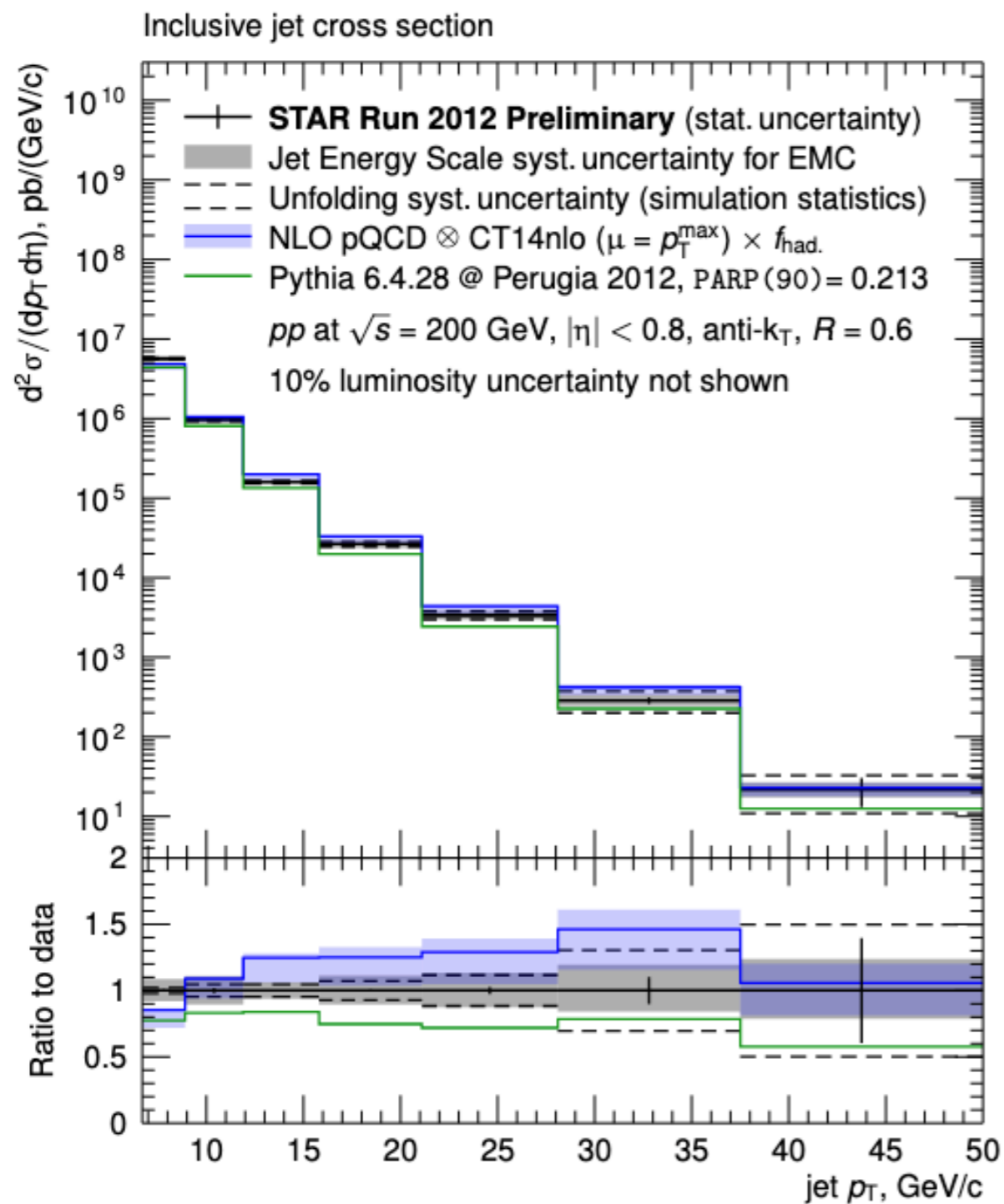
- Unfolded Jet Cross section overestimated by NLO+NP
- Resummations in In r vital
- At NNLO these corrections are quite significant



Xiaohui Liu et.al PRL 119 212001



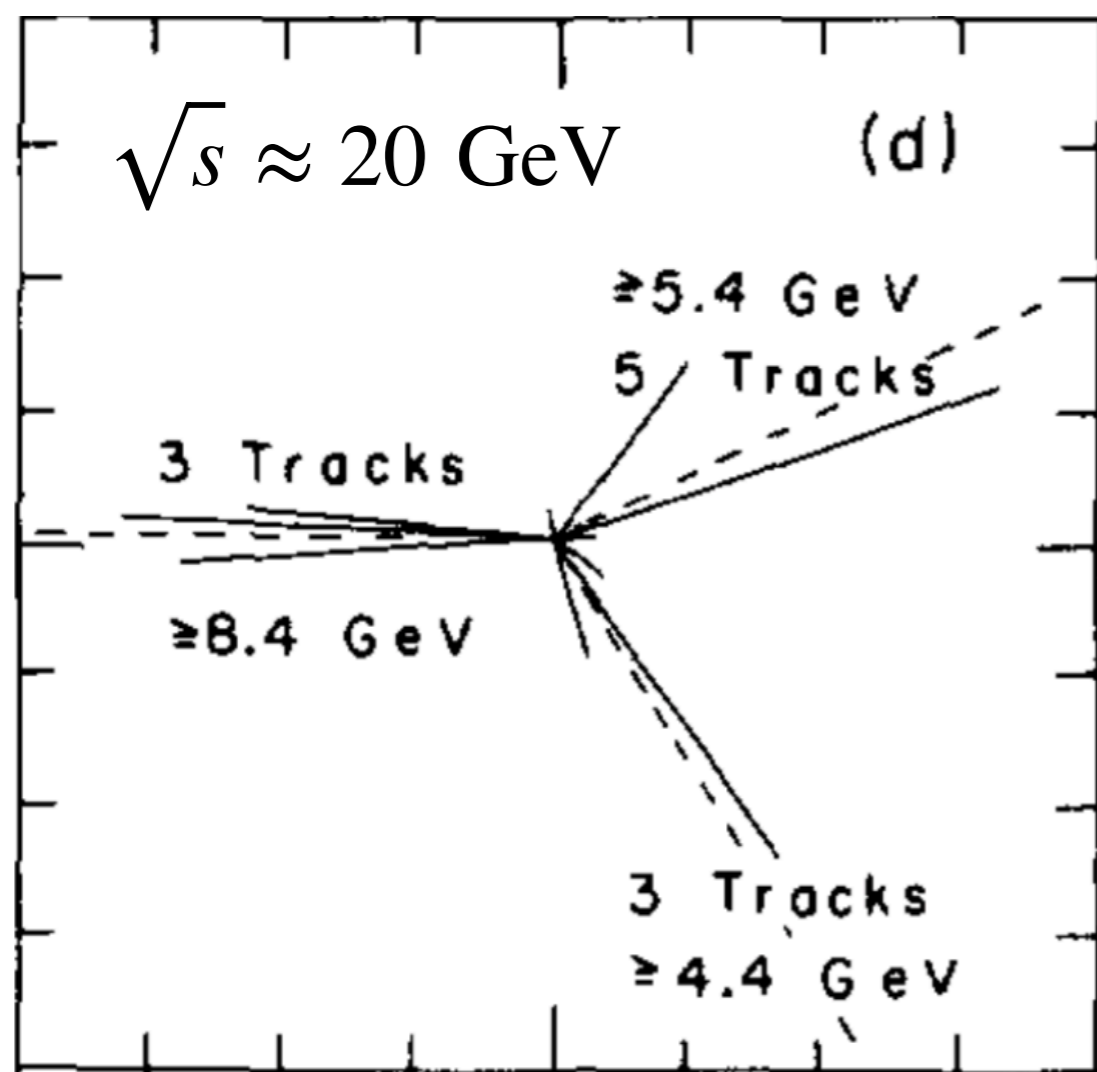
# pp jet spectra @ RHIC



- Steeper spectra also consistent with NLO but highlighting regions tension
- Measurements at varying  $\sqrt{s}$  will provide insights into energy dependence of various MC parameters

# Observing quarks/gluons in nature - Jets

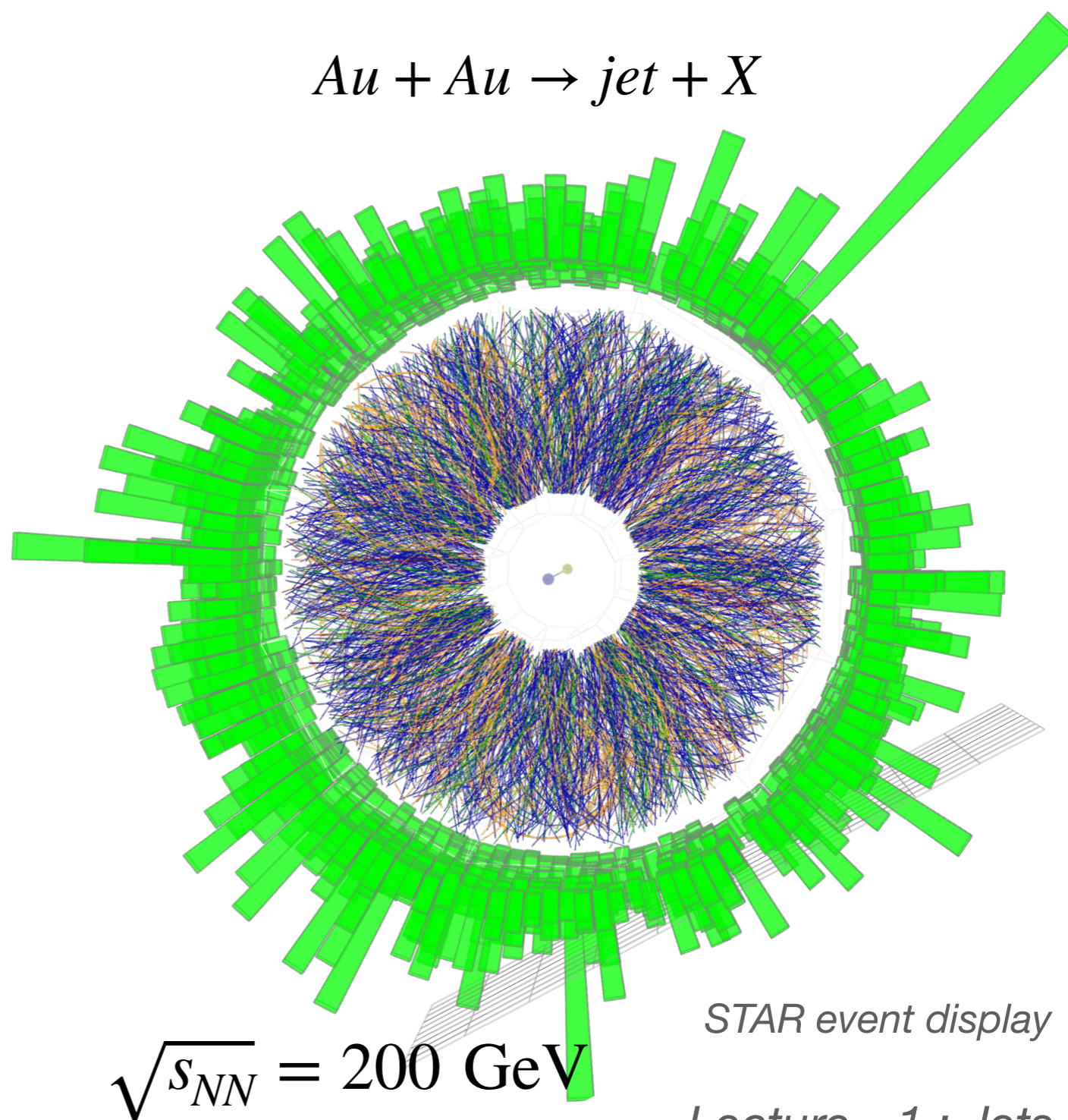
$$e^+ + e^- \rightarrow q + \bar{q} + g$$



TASSO Phys. Lett, V 86B number 2 (1979)

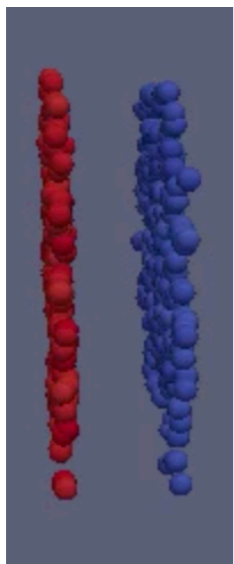
RKE, ISS @ Prague June 2022

$$Au + Au \rightarrow \text{jet} + X$$

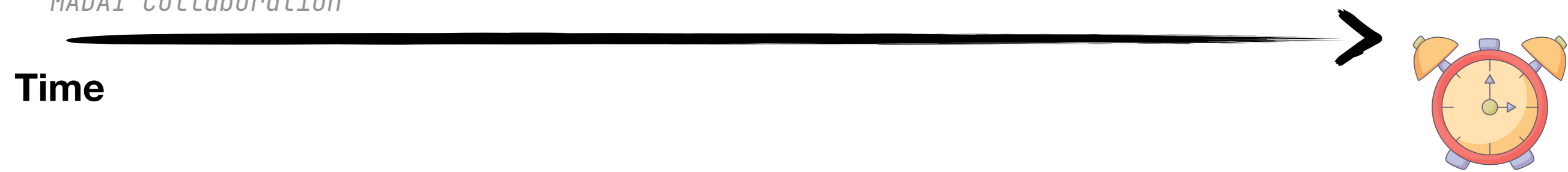


STAR event display

Lecture - 1 : Jets



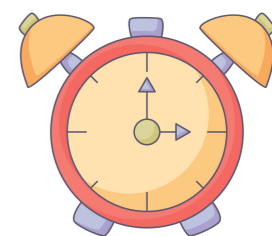
MADAI Collaboration





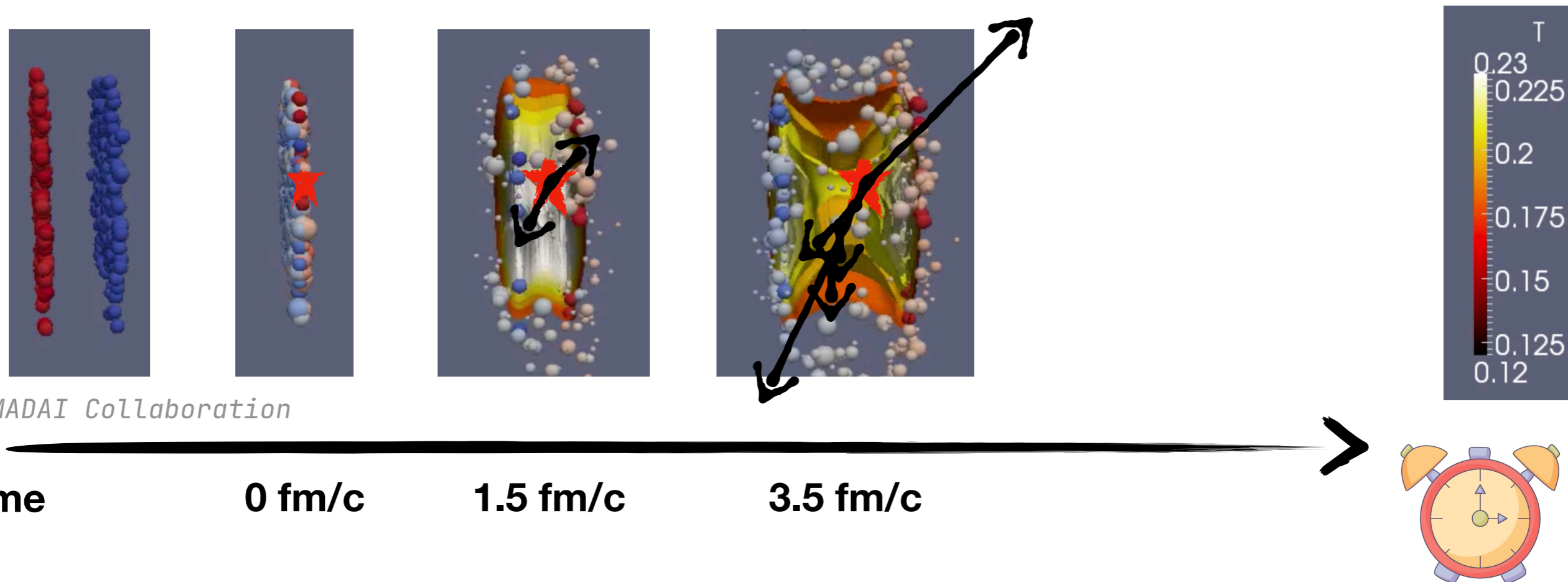
MADAI Collaboration

Time 0 fm/c 1.5 fm/c ➔



**0 fm/c**      Jets are created nearly instantaneously

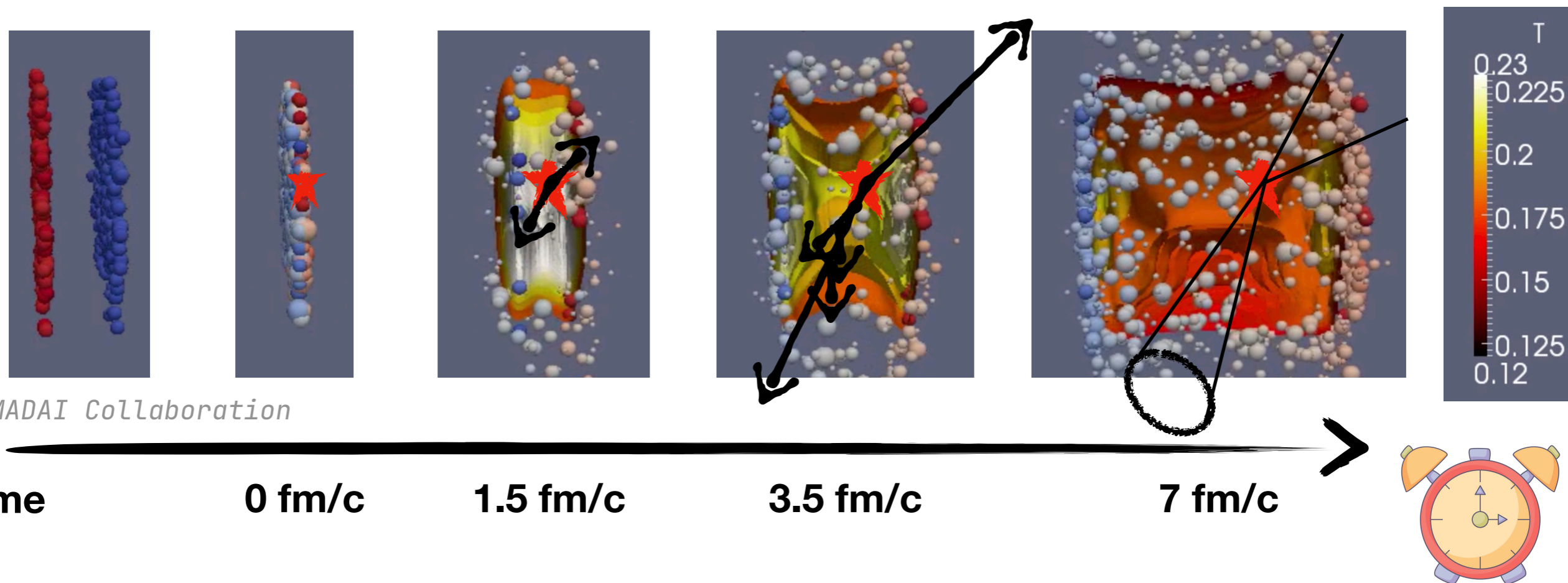
**0.4 fm/c**      QGP is formed and starts to expand/cool



**0 - 5 fm/c**

Concurrent evolution of jet shower and QGP

# Jets and the QGP



**$> 10 \text{ fm/c}$**

Resulting modification to jet observables offers clue to medium properties

# Jets in Heavy Ions - What to expect?

Energy Loss of Energetic Partons in Quark-Gluon Plasma:  
Possible Extinction of High  $p_T$  Jets in Hadron-Hadron Collisions.

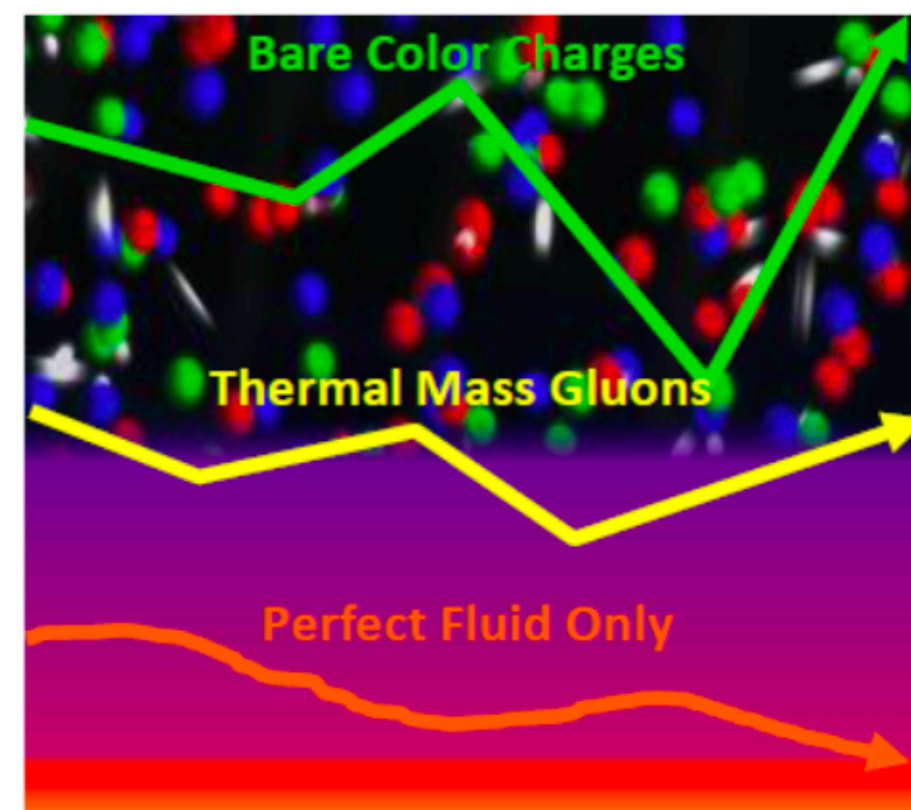
J.D.Bjorken, [FERMILAB-Pub-82/59-THY](#)

Image credit - 2015 LRP

hadron-hadron collisions with high associated multiplicity and with transverse energy  $dE_T/dy$  in excess of 10 GeV per unit rapidity, it is possible that quark-gluon plasma is produced in the collision. If so, a

produced secondary high- $p_T$  quark or gluon might lose tens of GeV of its initial transverse momentum while plowing through quark-gluon plasma produced in its local environment. High energy hadron jet experiments

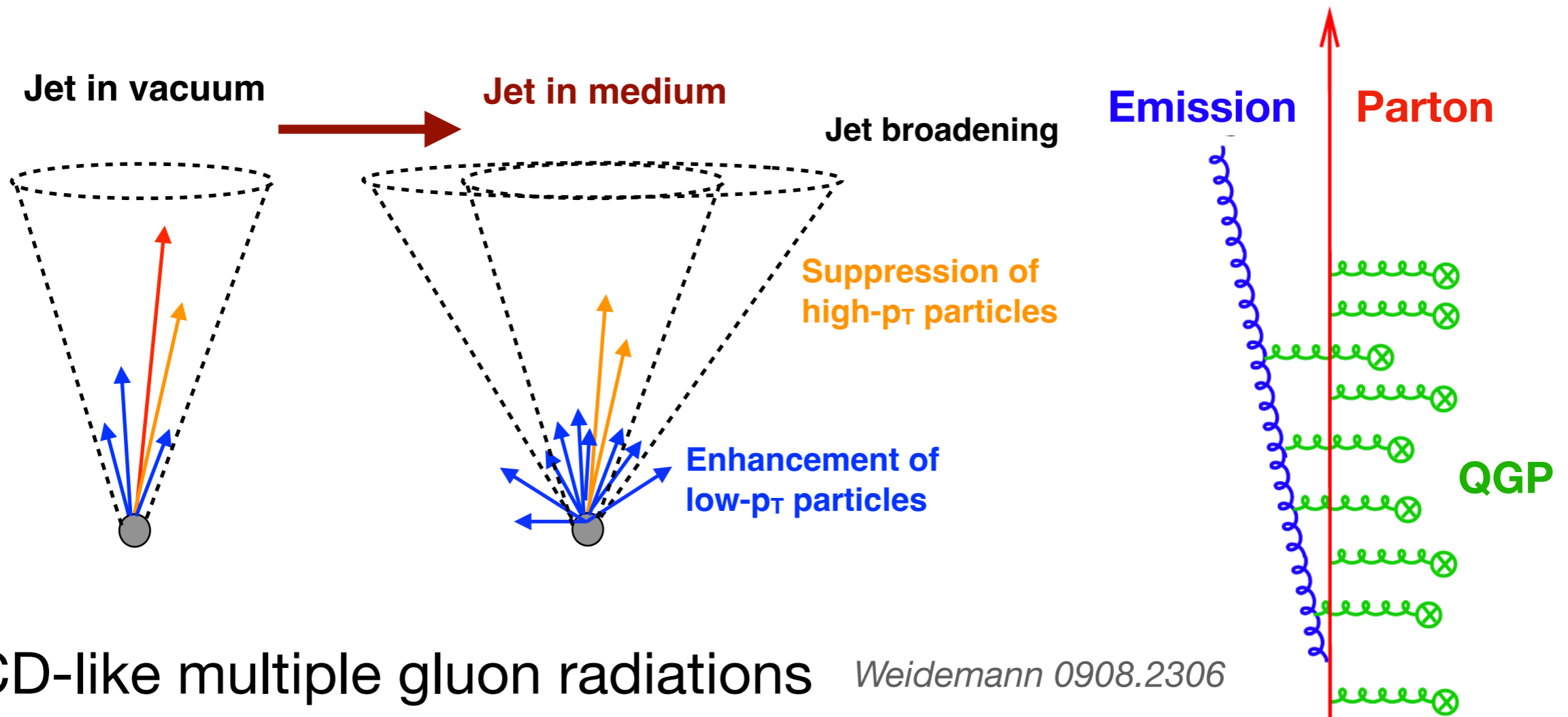
should be analysed as function of associated multiplicity to search for this effect. An interesting signature may be events in which the hard collision occurs near the edge of the overlap region, with one jet escaping without absorption and the other fully absorbed.



- High  $p_T$  particle/jet yield suppressed and energy loss
- Modification of the hard scattered partons due to scatterings in the medium - could affect both jet axis and distributions of jet constituents



# Mechanisms of energy loss

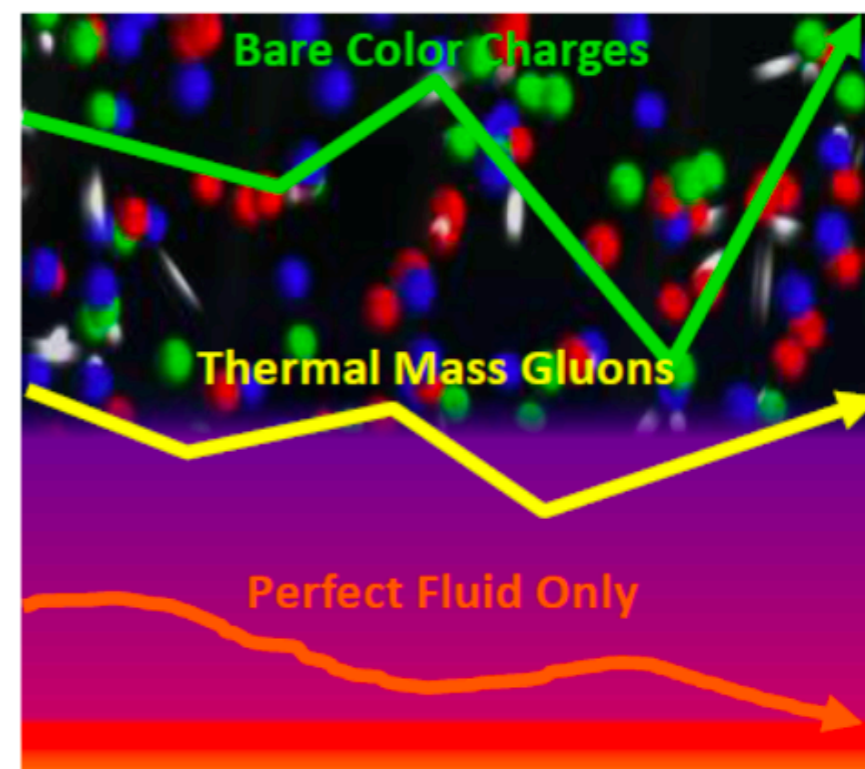
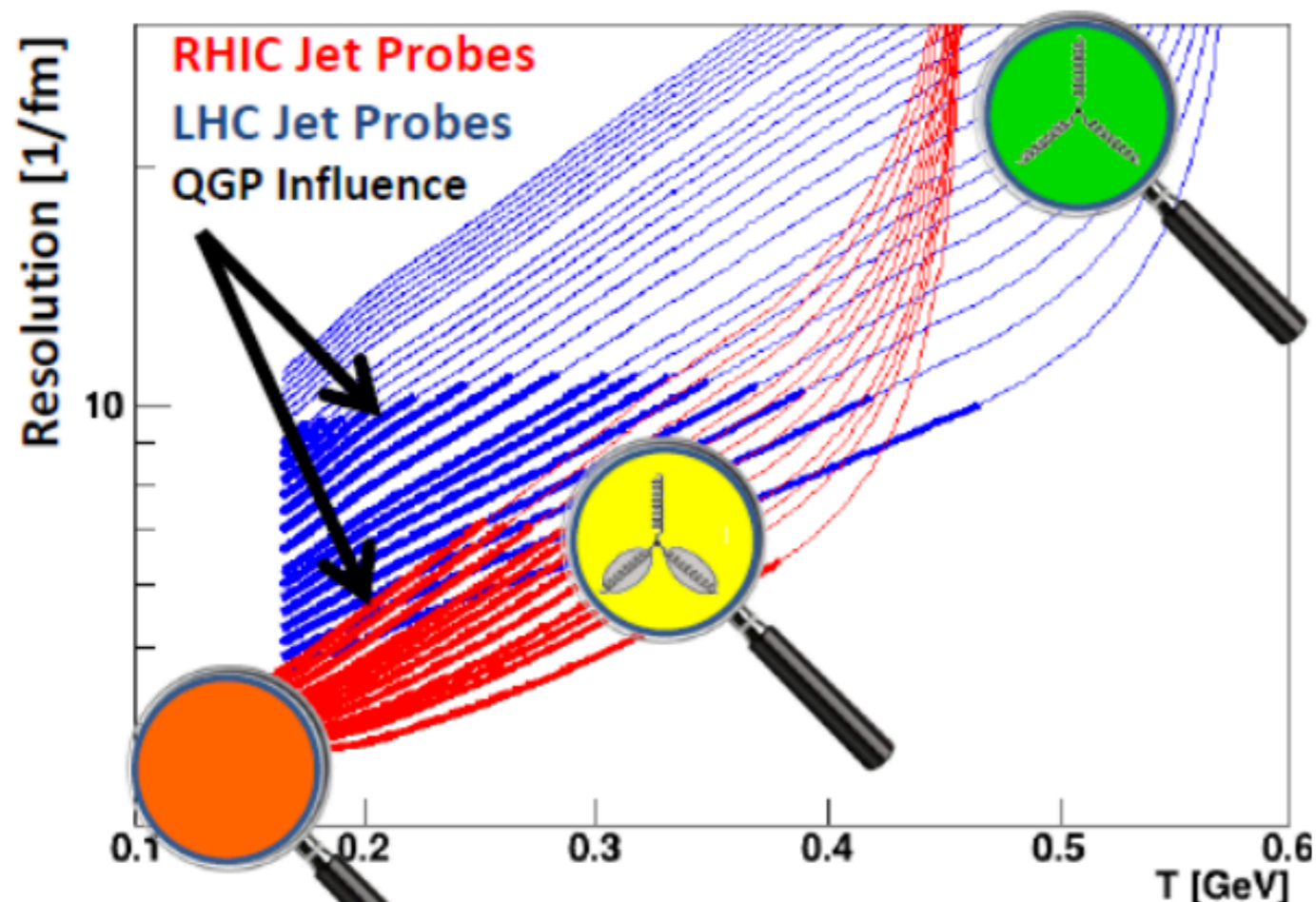


- pQCD-like multiple gluon radiations *Weidemann 0908.2306*
- medium induced scatterings (inelastic), *Zapp et al. JHEP 03 (2013) 080*
- AdS/CFT energy loss, *Liu et al. Phys. Rev. Lett. 97 (2006) 182301*
- color coherence/decoherence, *Solano et al. JHEP 10 (2014) 019*
- modified partonic splitting functions etc...

*Qin, Wang,  
Int. J. Mod. Phys. E 24 (2015) 11, 1530014*

# Jets in Heavy Ions - probe of parton energy loss

The 2015  
LONG RANGE PLAN  
for NUCLEAR SCIENCE



Microscopic properties of the QGP Medium  
- structure at varying scales

This is inherently a two step process that is not mutually exclusive

**Understand jet energy loss → Extract medium properties**

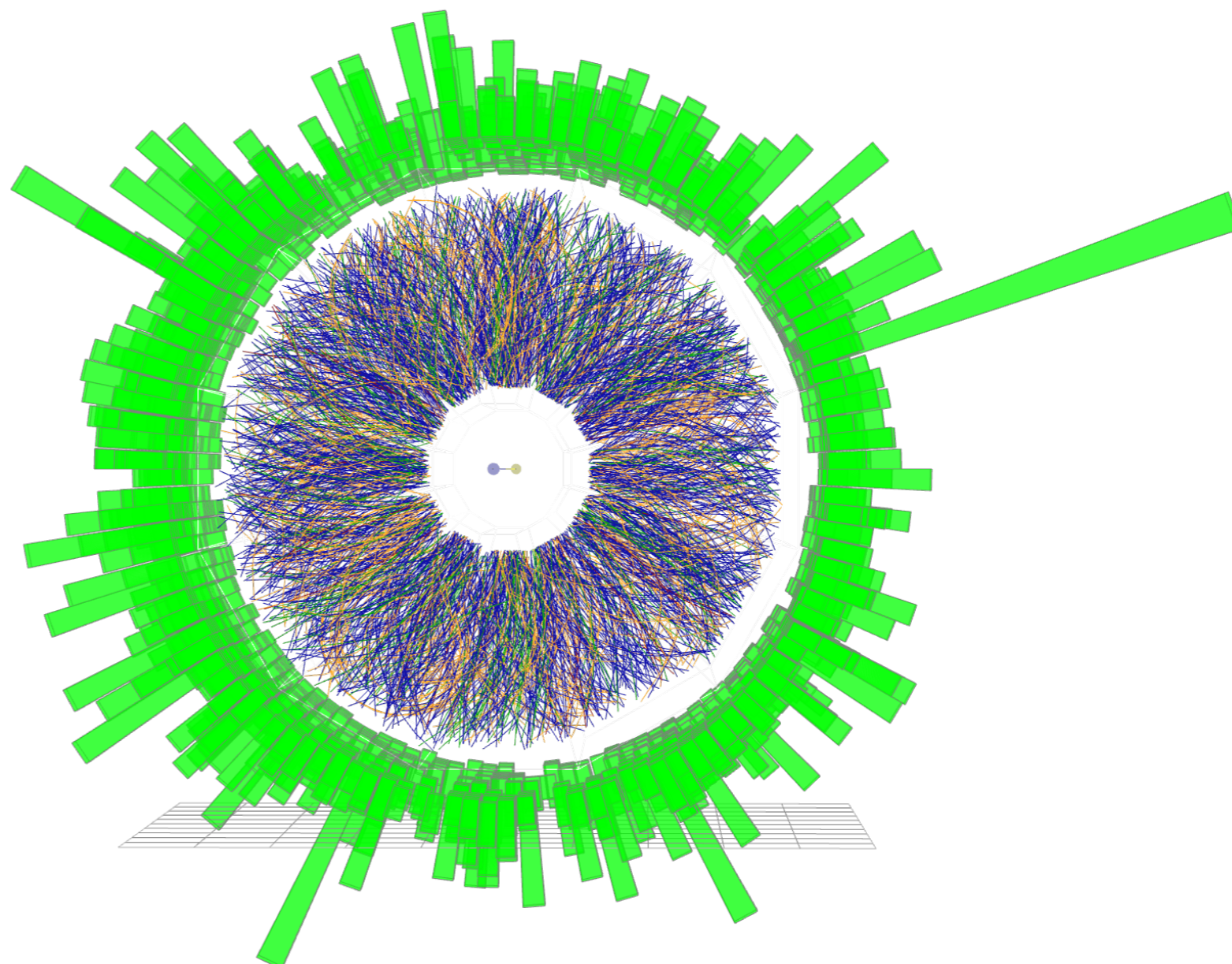
# What makes jet measurements in heavy-ion collisions hard?



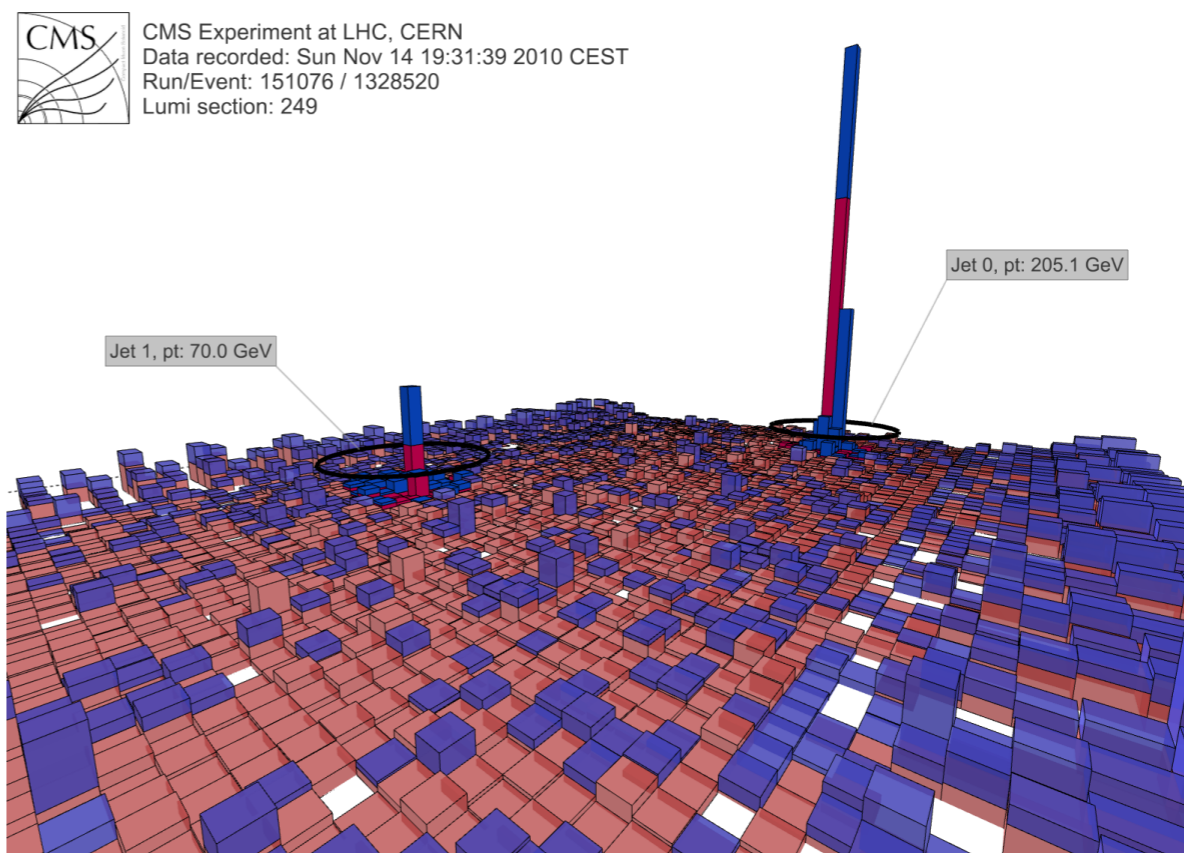
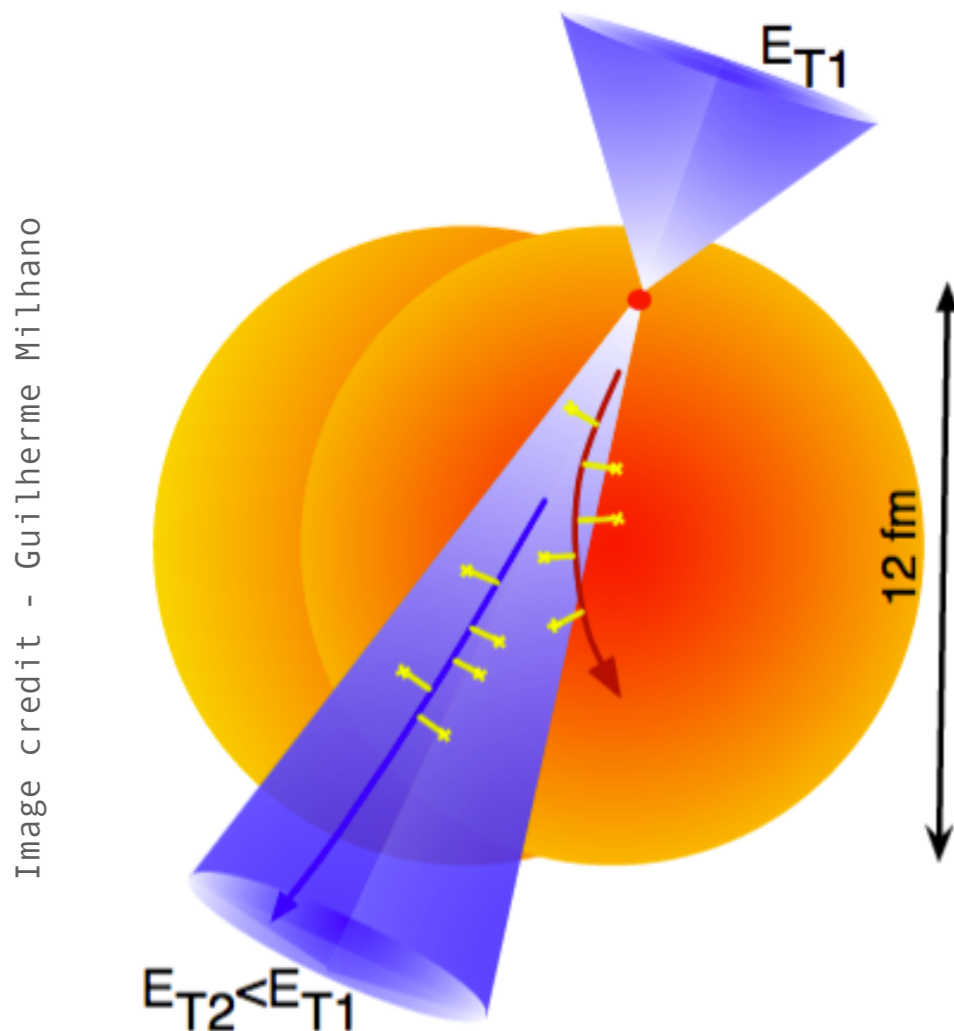
# What makes jet measurements in heavy-ion collisions hard?



Where are the jets here?

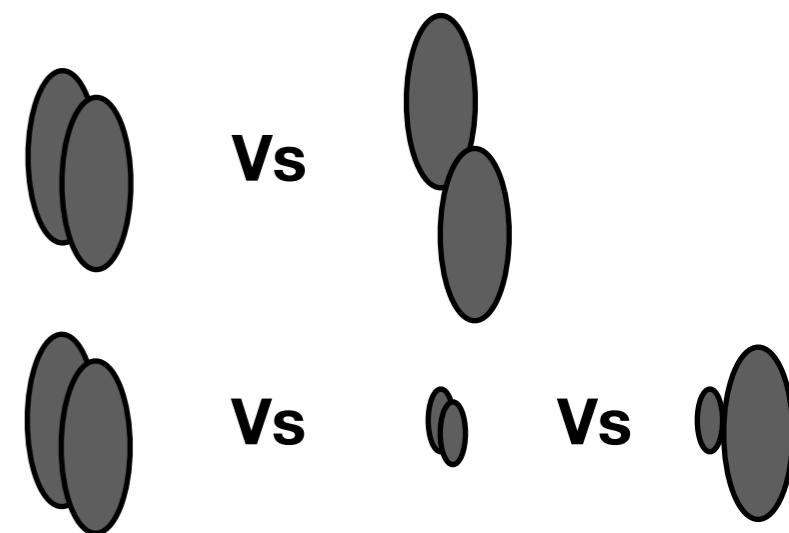


# Experimentally measuring Jet Quenching



In comparisons of jet observables to a reference

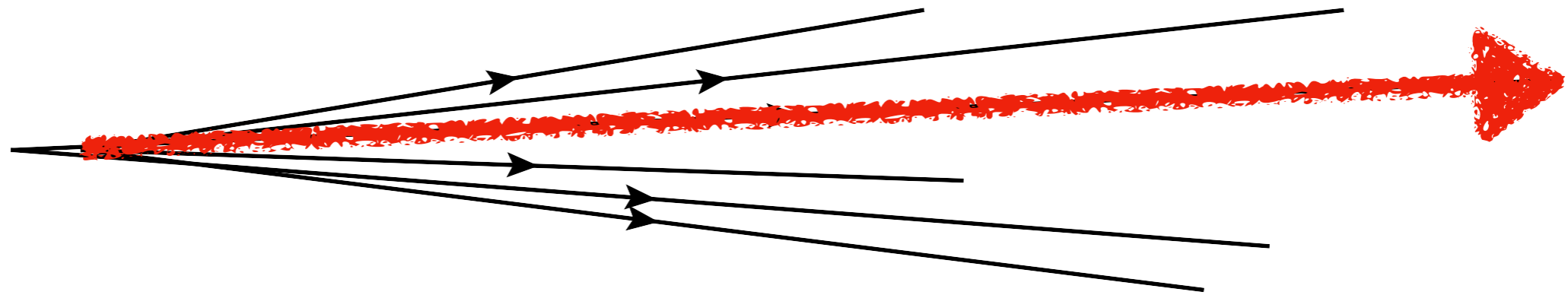
- Central vs Peripheral collisions
- A+A vs p+p or p+A (with the assumption that medium induced modifications are weaker or non-existent in the latter systems )



# Recap - 2

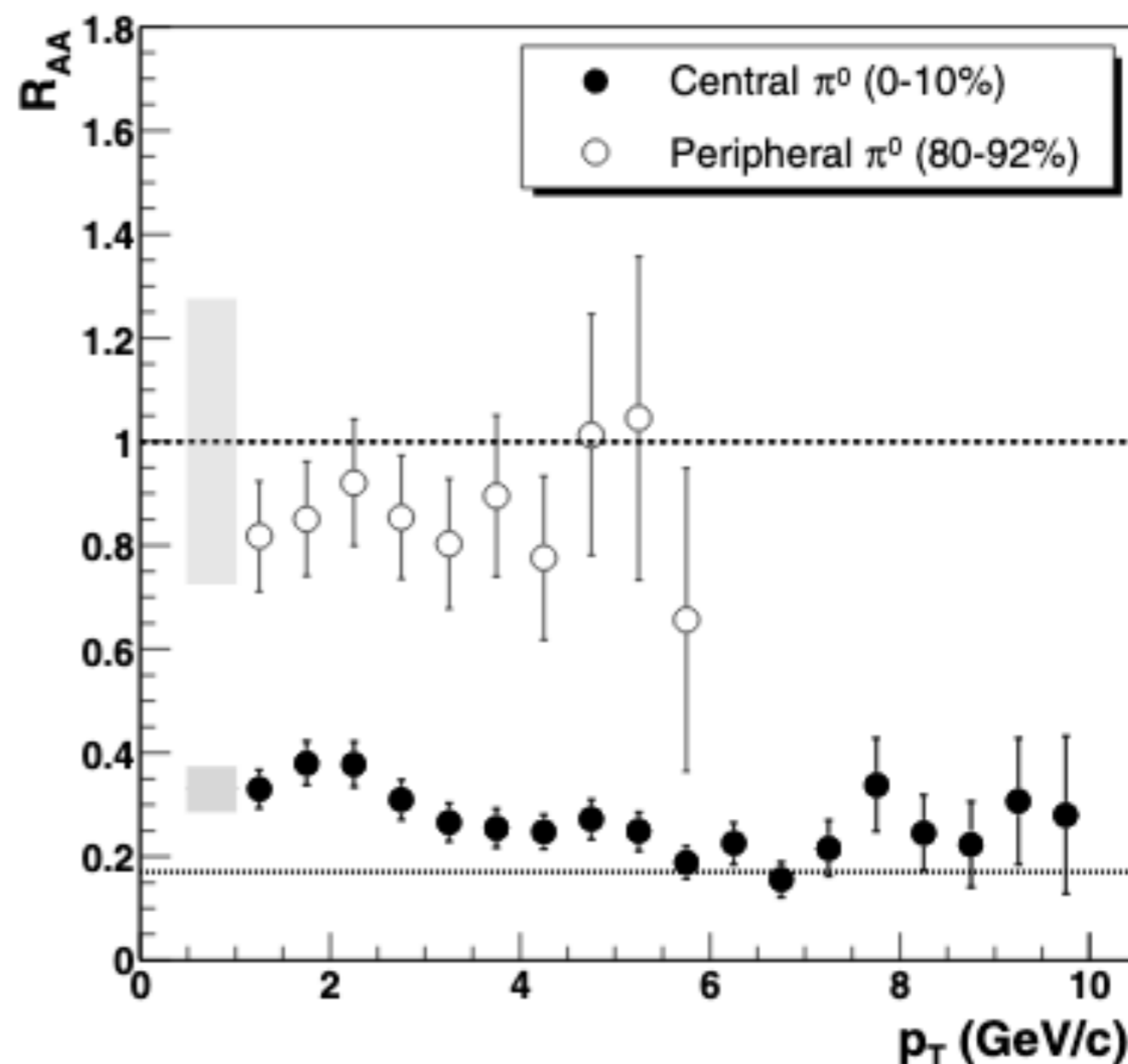
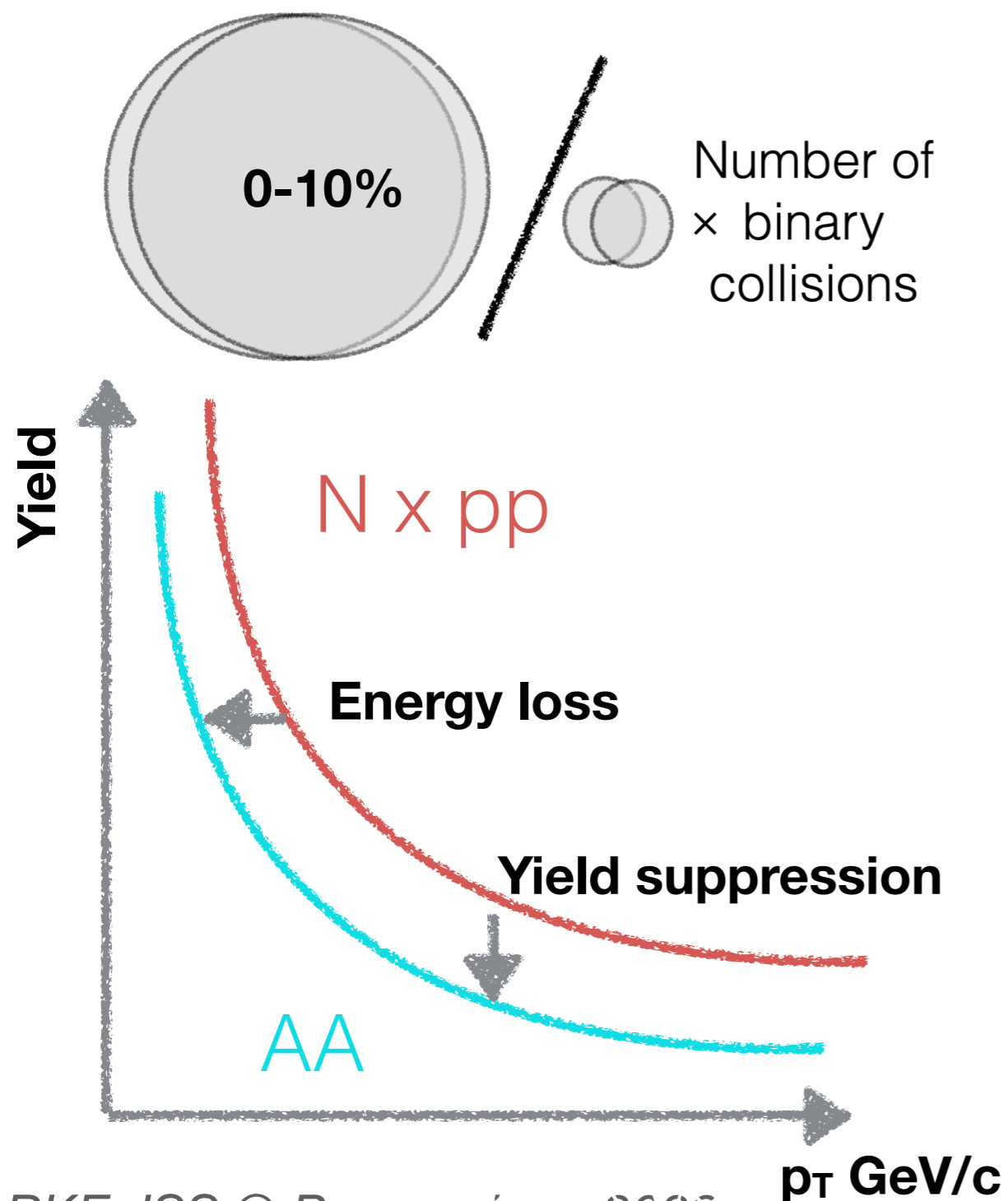
- Jets are calculable via QCD factorization and thus heralded as calibrated probes
- Produced early in the collisions and traverse the QGP lifetime
- Expect parton jets to lose energy in the medium due to potential scatterings and radiations
- Need to reconstruct and compare jet properties to reverse engineer medium impact
- Extracting jets from background is an issue

# 0th generation Jet Measurements



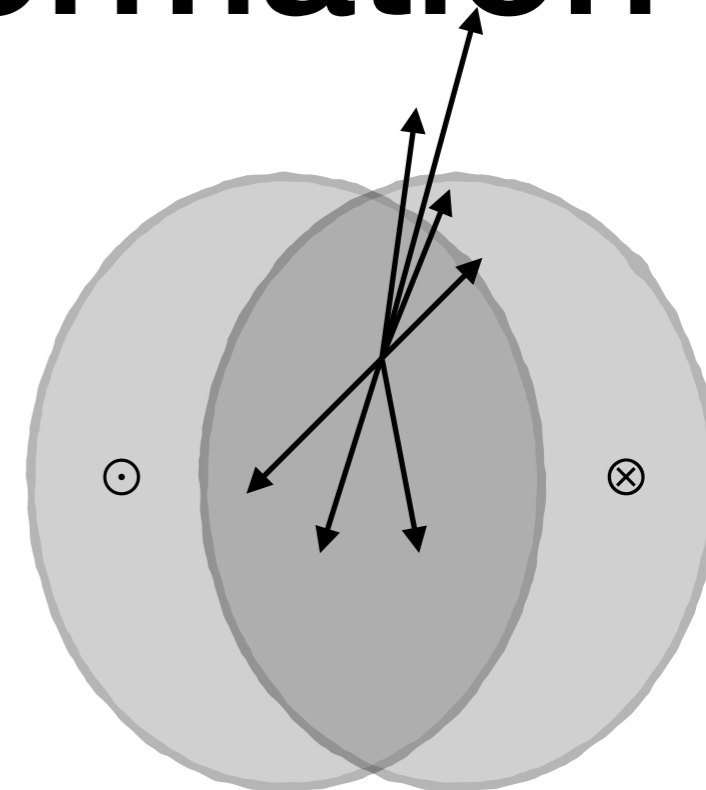
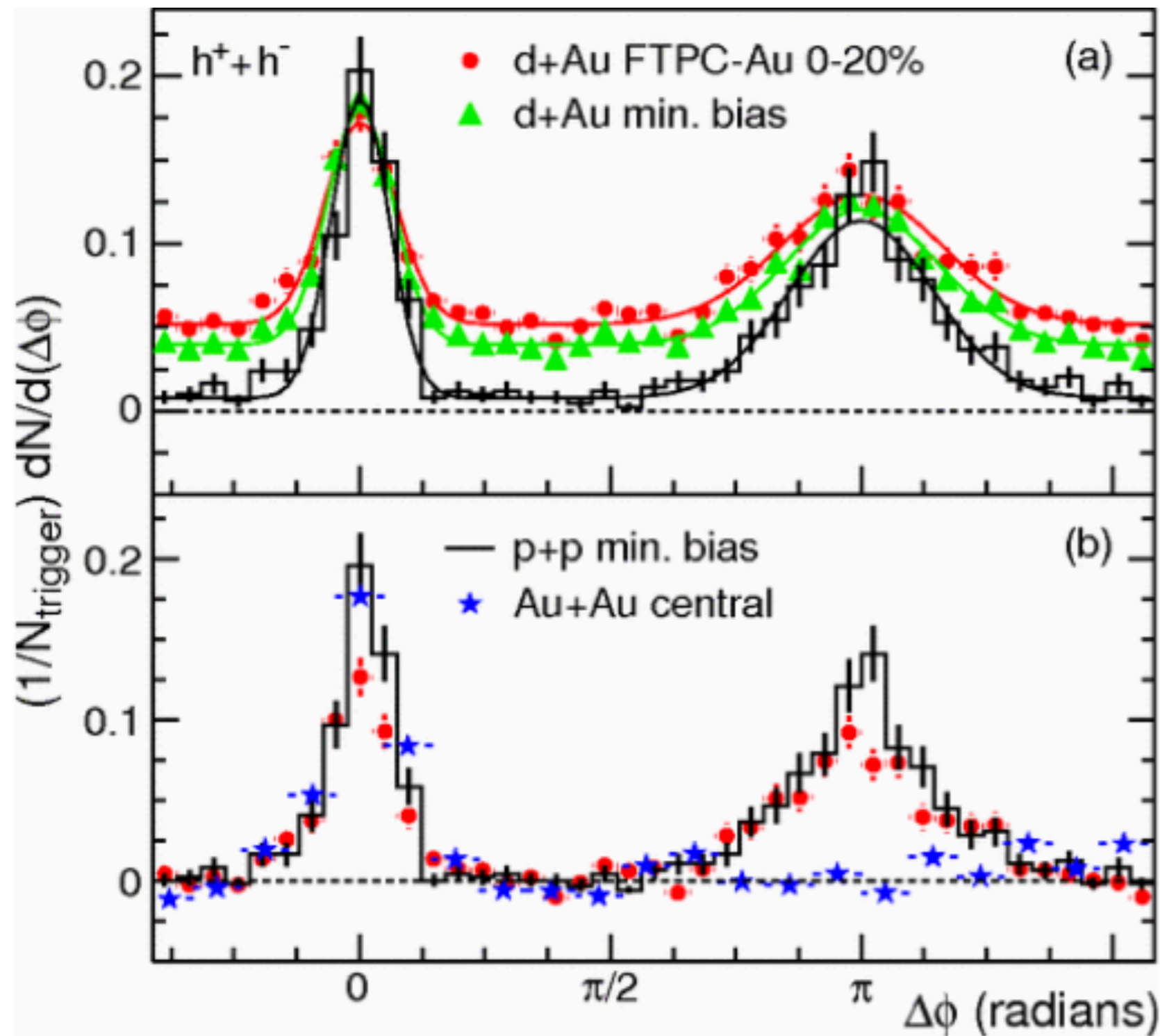
**High momentum hadron as  
proxy for jet**

# Signature of energy loss and evidence for QGP formation





# Signature of energy loss and evidence for QGP formation



- Disappearance of the away-side jet particles!

# 0th gen - what did we learn?

- High momentum hadron production is **significantly suppressed** in central heavy ion collisions
- **Disappearance of back-to-back production**, i.e. away-side yield in 2-particle correlations
- Signature of quenching - but not learning much about jets or parton energy loss - need to reconstruct jets!

*ATLAS Heavy Ion Publications*

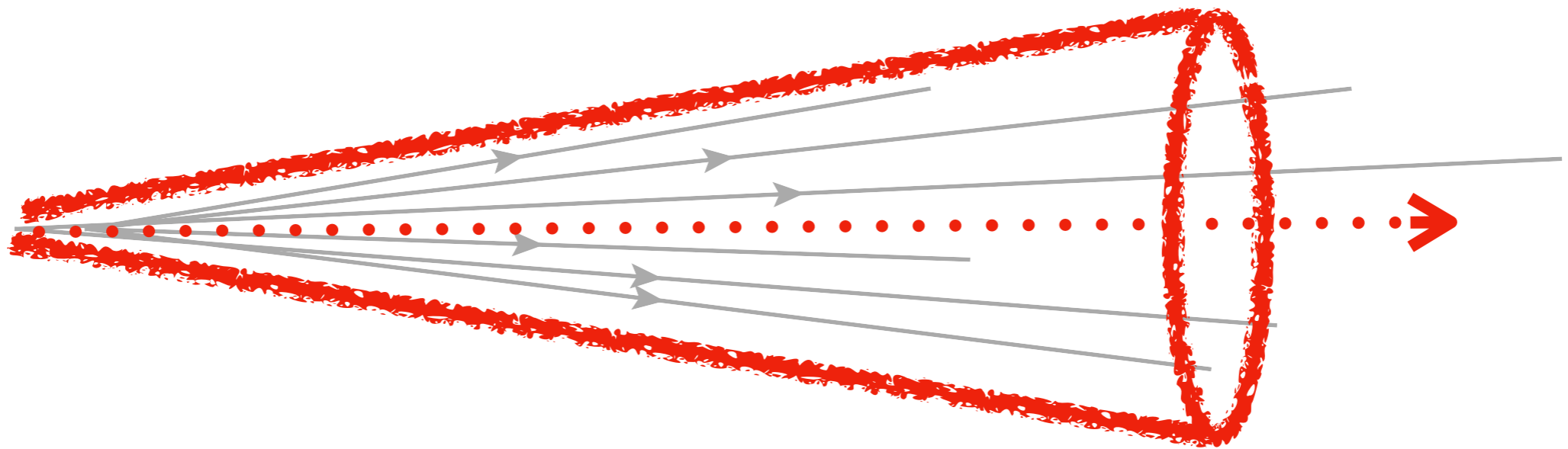
*ALICE Heavy Ion Publications*

*CMS Heavy Ion Publications*

*STAR Publications*

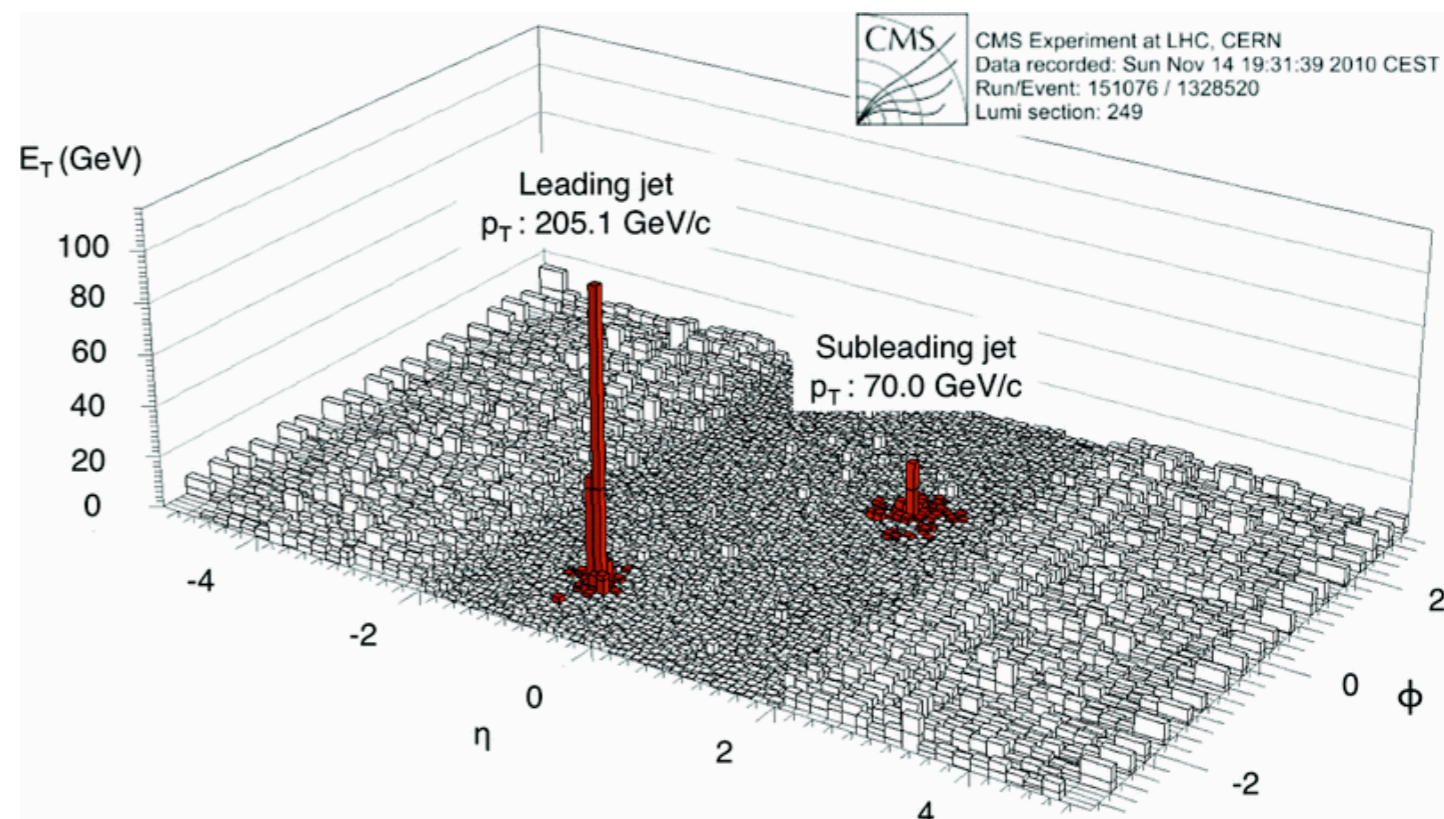
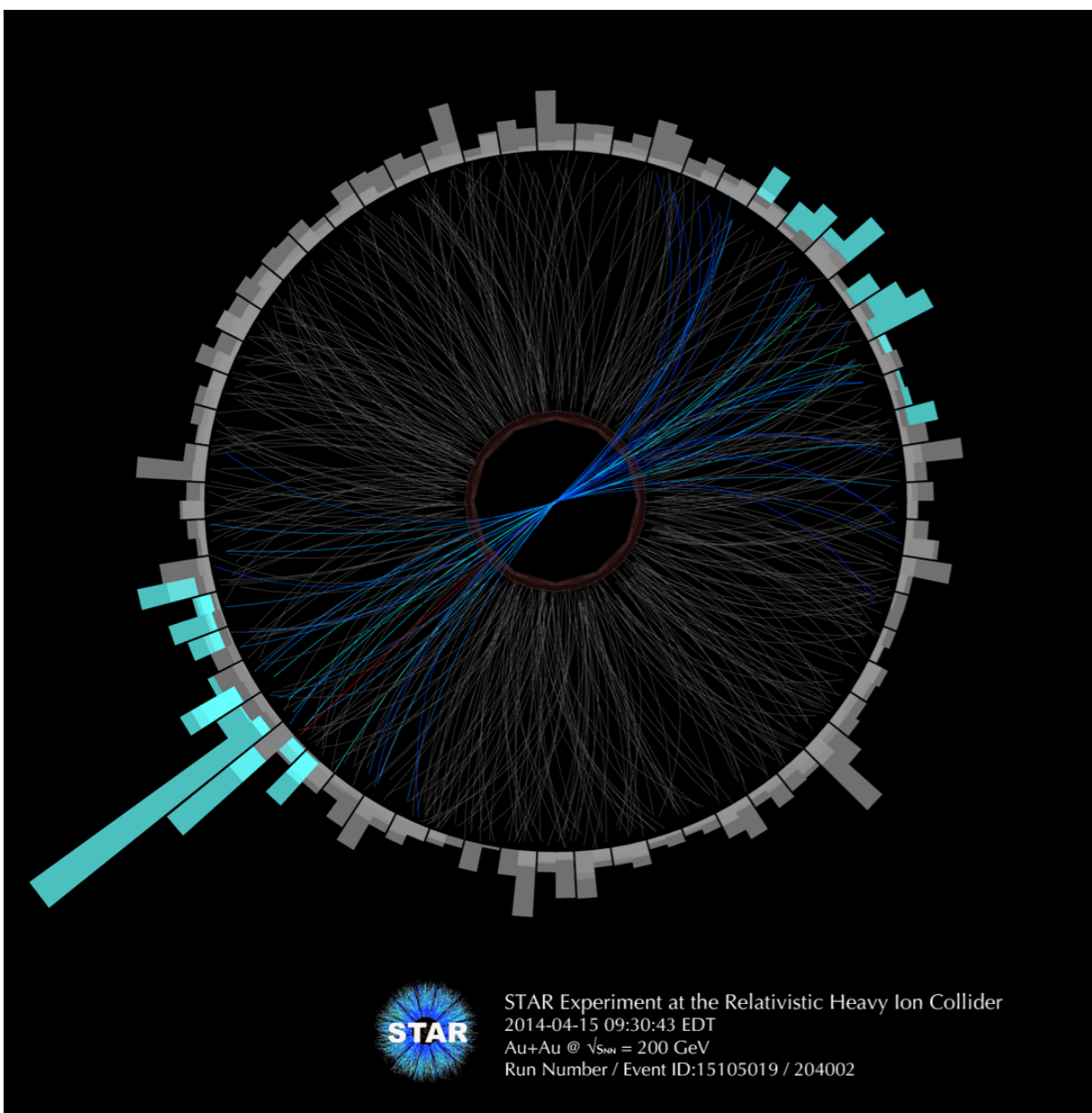
*PHENIX Publications*

# 1st generation Jet Measurements



**Fully reconstructed jets as proxies  
for hard scattered partons**

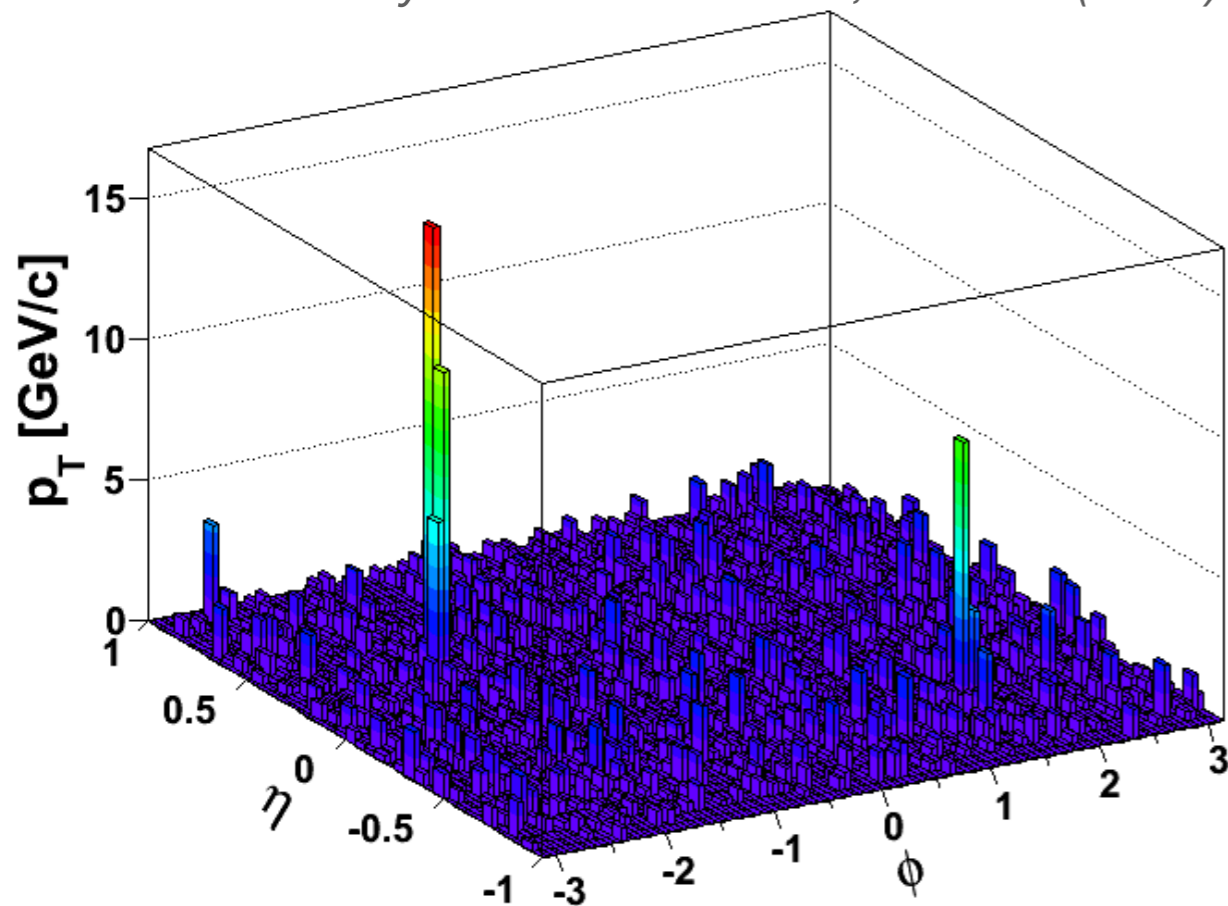
# Visually observing jet quenching



- Dijet-Balancing is a useful tool to 'see' jet quenching

# Selecting a di-jet sample at STAR

*STAR Phys. Rev. Lett. 119 6, 062301 (2017)*

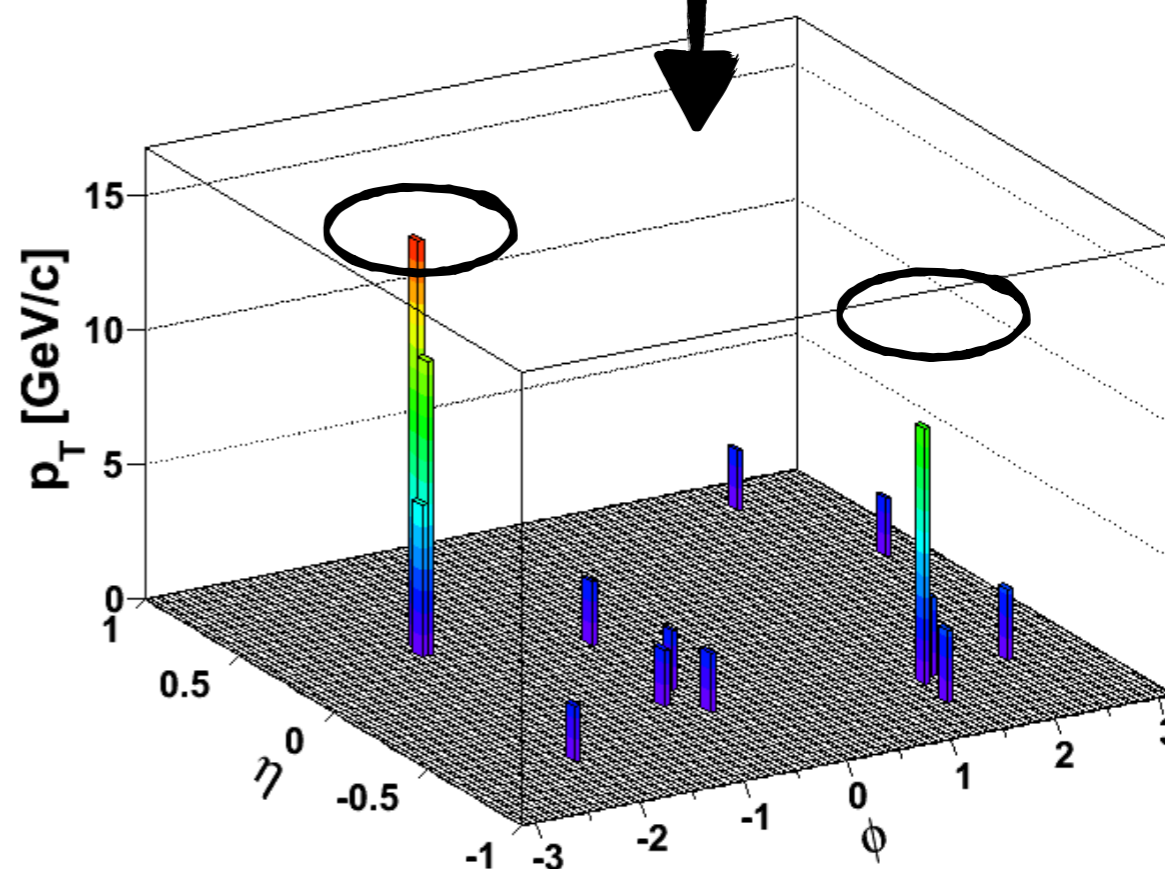


- Significant bias in jet fragmentation but lets utilize the bias to our advantage!

- Resulting di-jets free of combinatorial jets and background

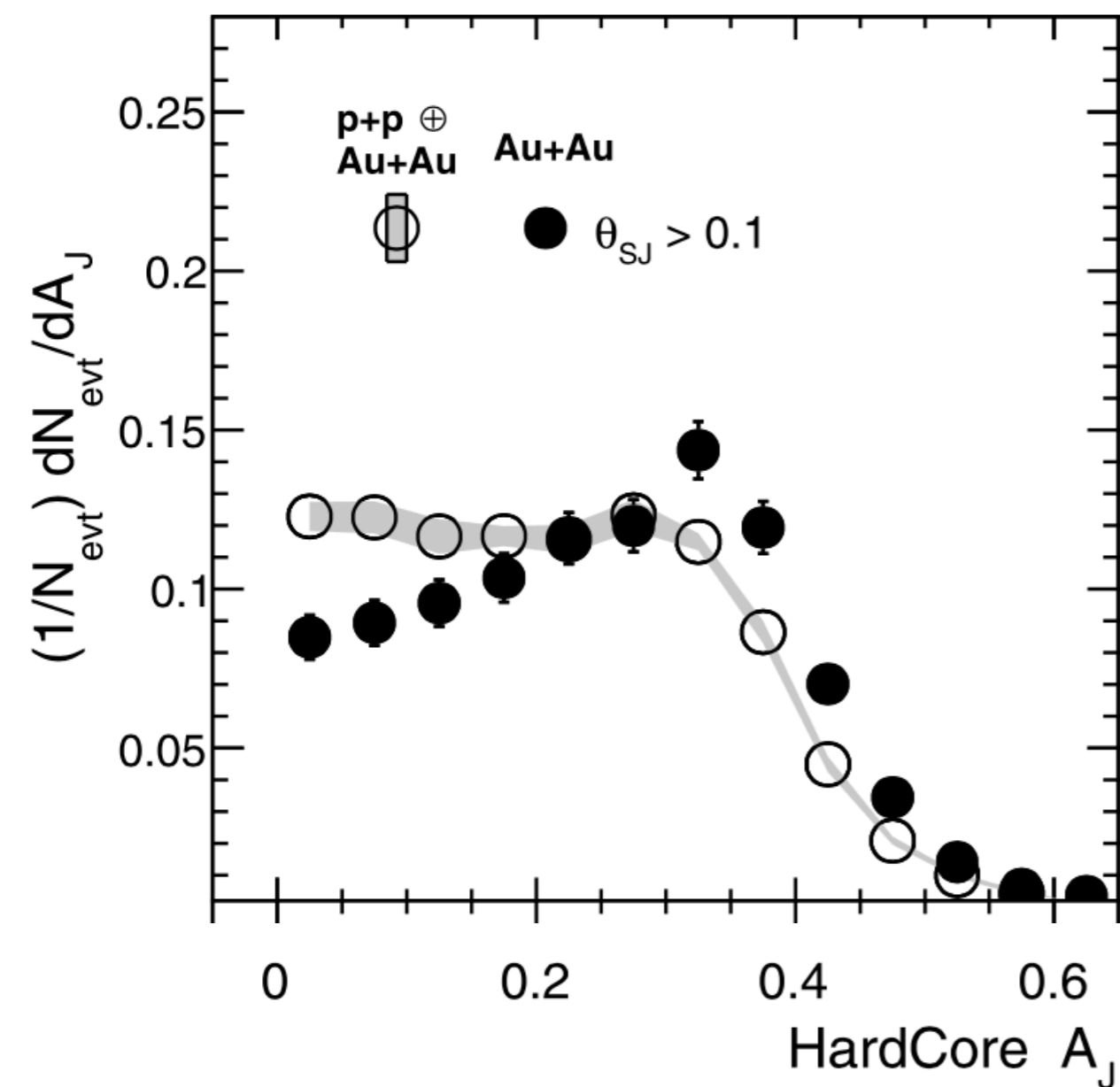
## HardCore Selection

$$p_T > 2 \text{ GeV/c}$$

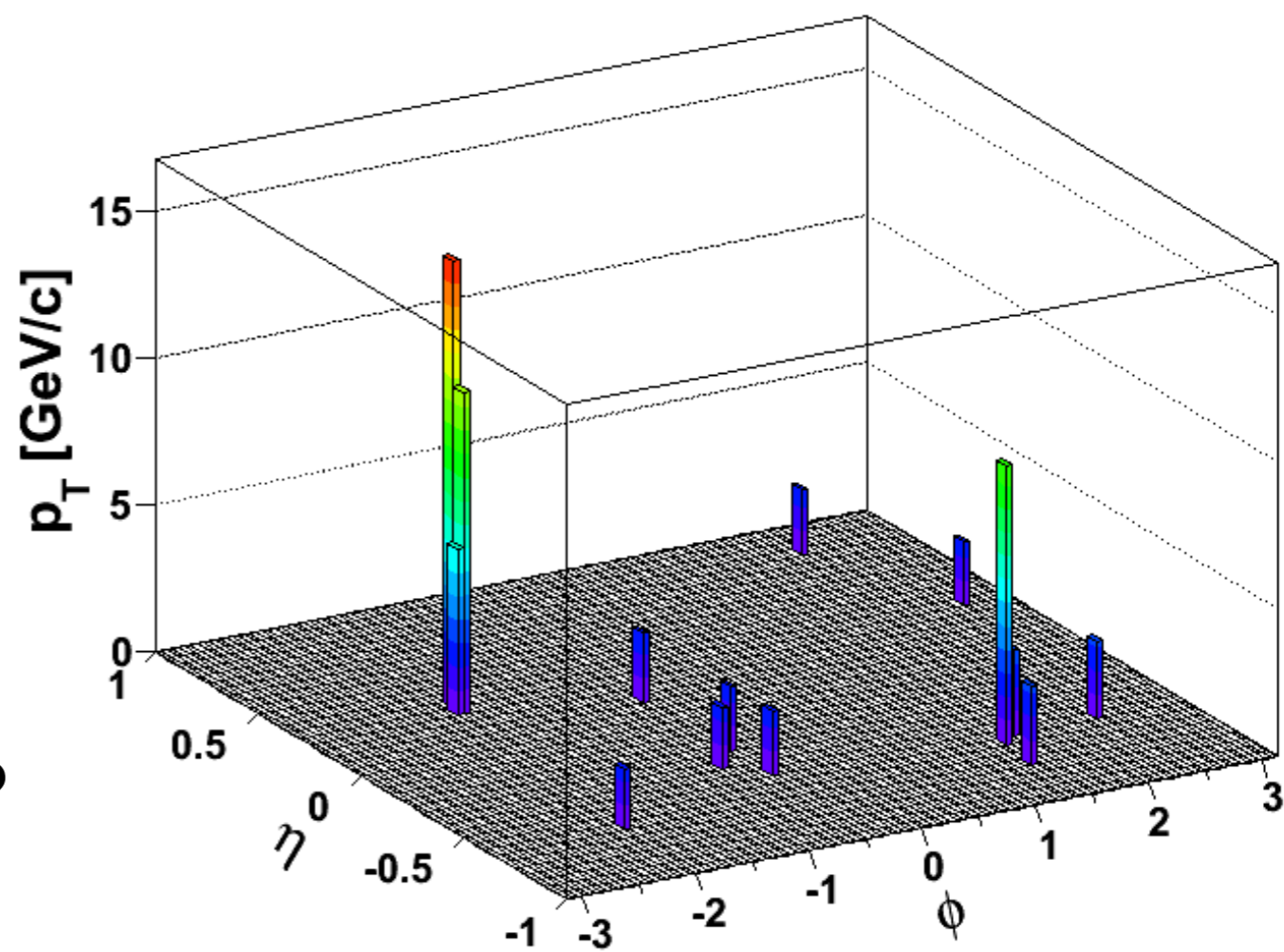


$$A_J = \frac{p_{T,\text{jet}}^{\text{Trigger}} - p_{T,\text{jet}}^{\text{Recoil}}}{p_{T,\text{jet}}^{\text{Trigger}} + p_{T,\text{jet}}^{\text{Recoil}}}$$

- HardCore Dijets in Au+Au are imbalanced compared to p+p ⊕ Au+Au



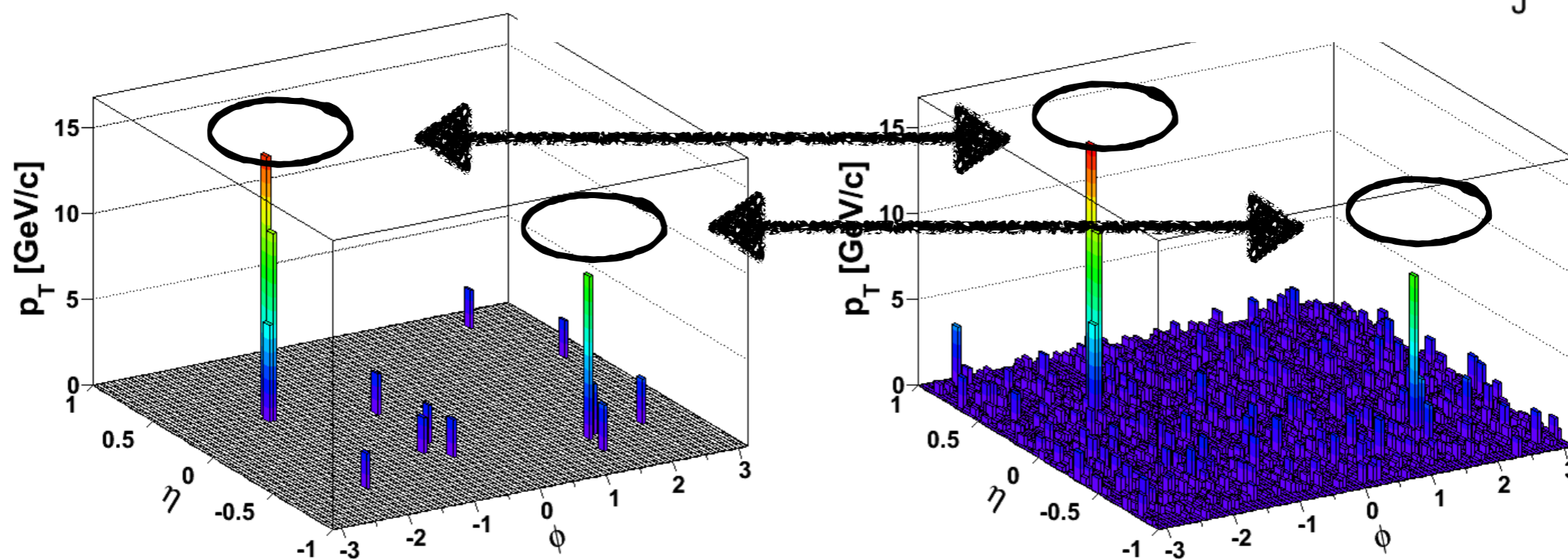
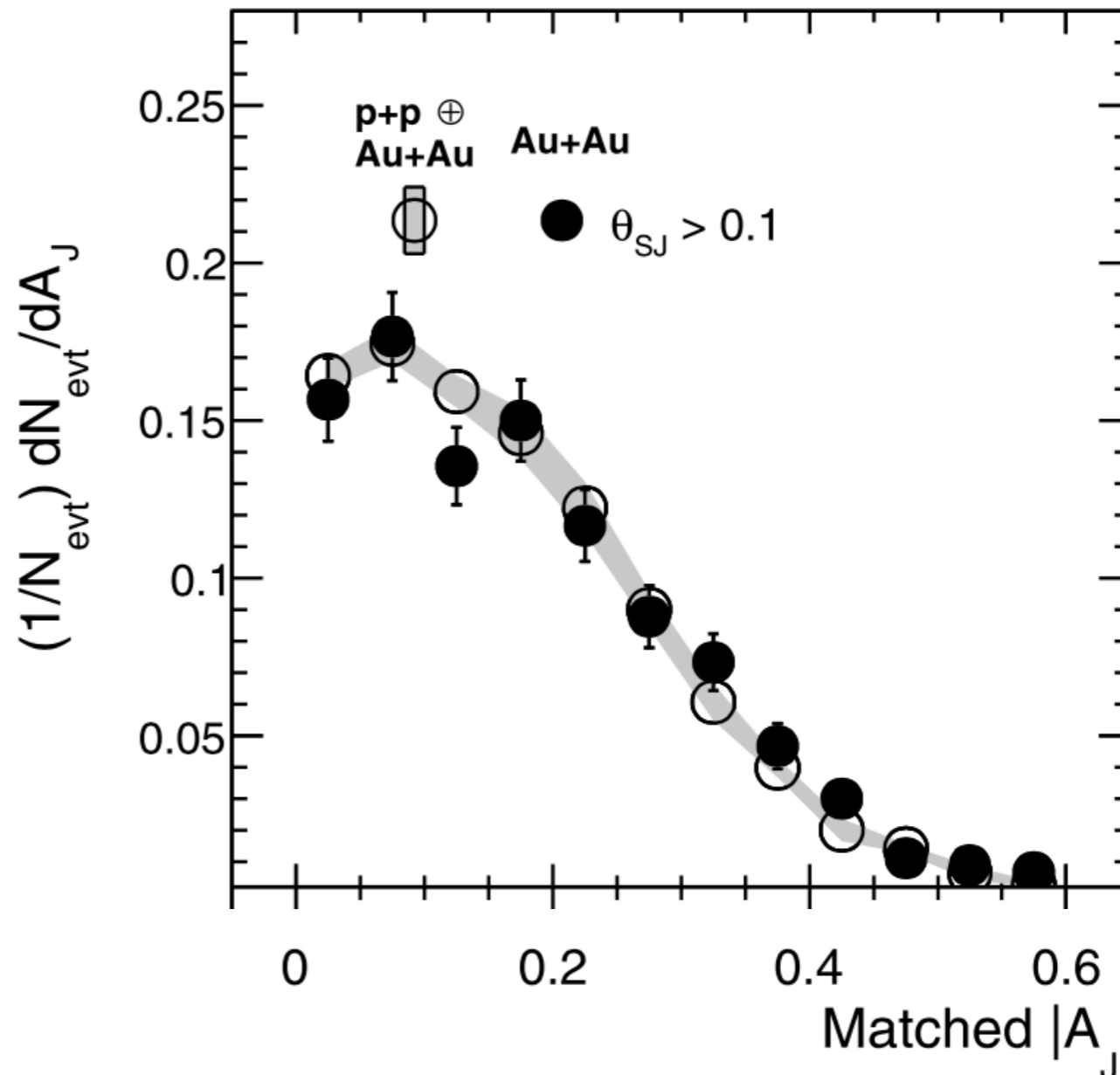
- High  $p_T$  particles ( $p_T > 2$  GeV) lose energy in the medium



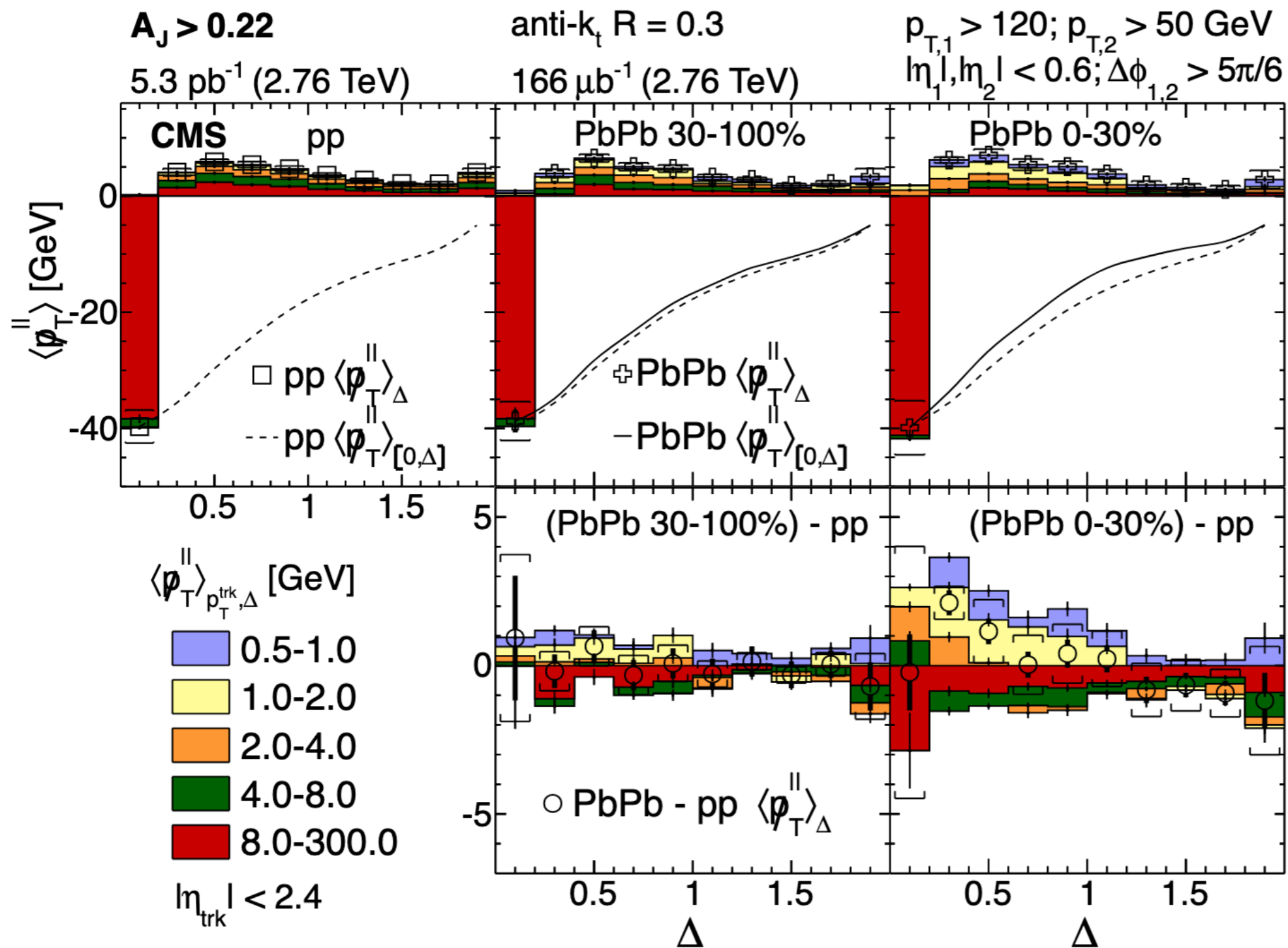
- Where is the quenched energy?

$$A_J = \frac{p_{T,\text{jet}}^{\text{Trigger}} - p_{T,\text{jet}}^{\text{Recoil}}}{p_{T,\text{jet}}^{\text{Trigger}} + p_{T,\text{jet}}^{\text{Recoil}}}$$

- Matched jets recover the lost energy  
( $0.2 < p_T < 2$  GeV)  
within the jet cone  
( $R = 0.4$ )!



# At the LHC?



**Balancing only occurs much further after from the jet axis**  
 $\Delta R \approx 2$  the entire hemisphere!

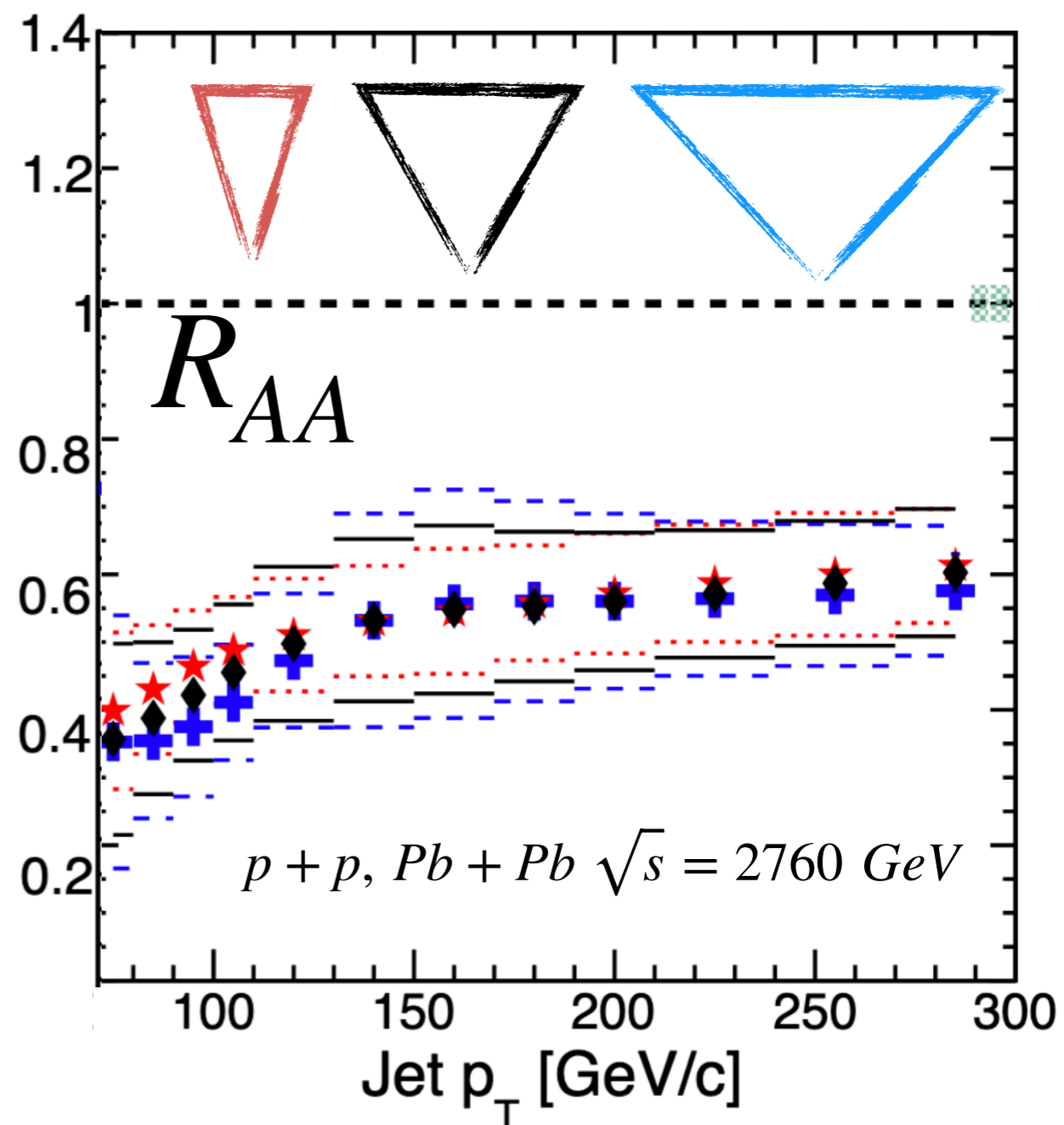
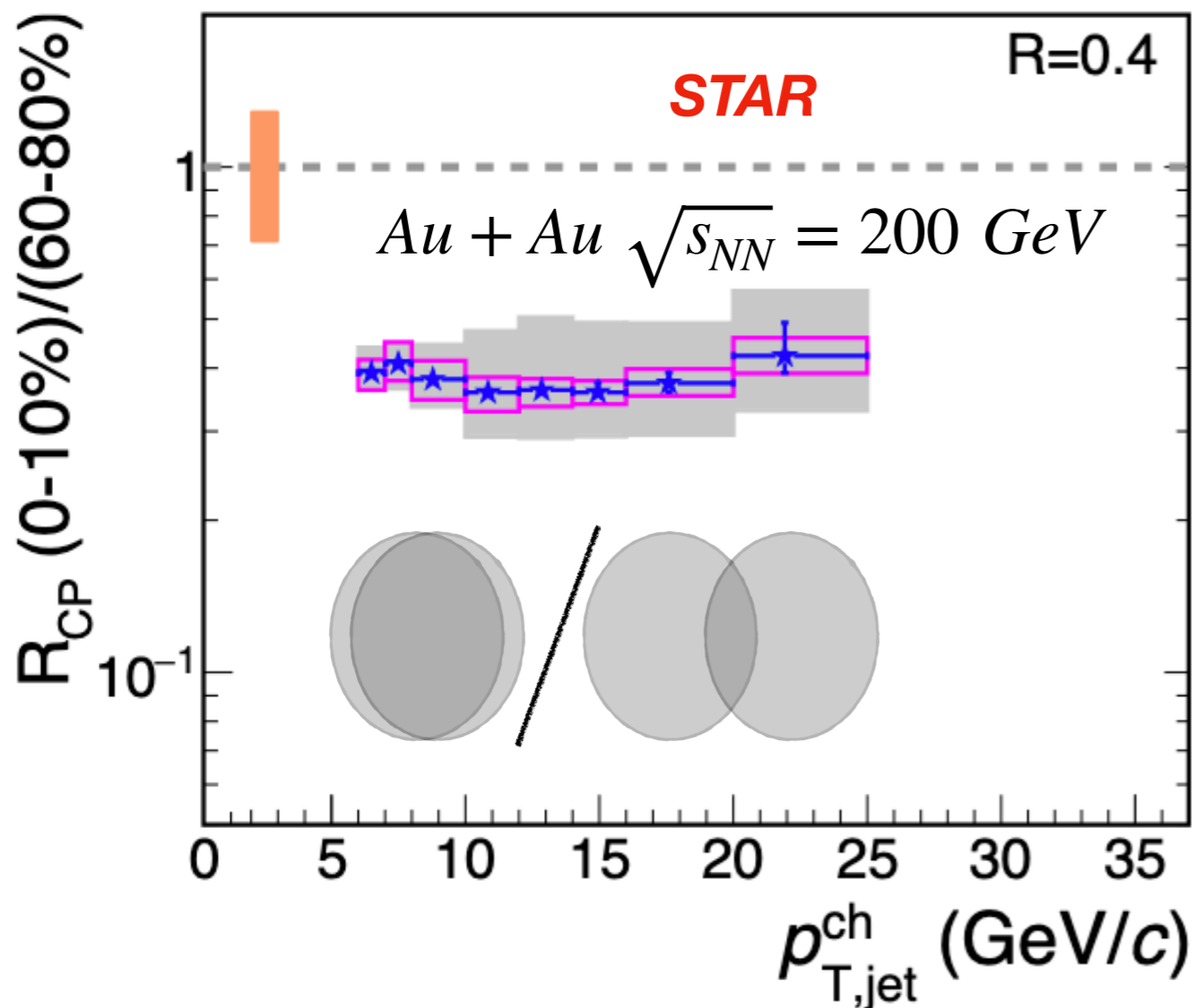
**Enhancement of particles in heavy ions around the jet cone**



# Inclusive jets @ RHIC and LHC

CMS Phys. Rev. C 96 (2017) 015202

STAR Phys. Rev. C 102 5, 054913 (2020)

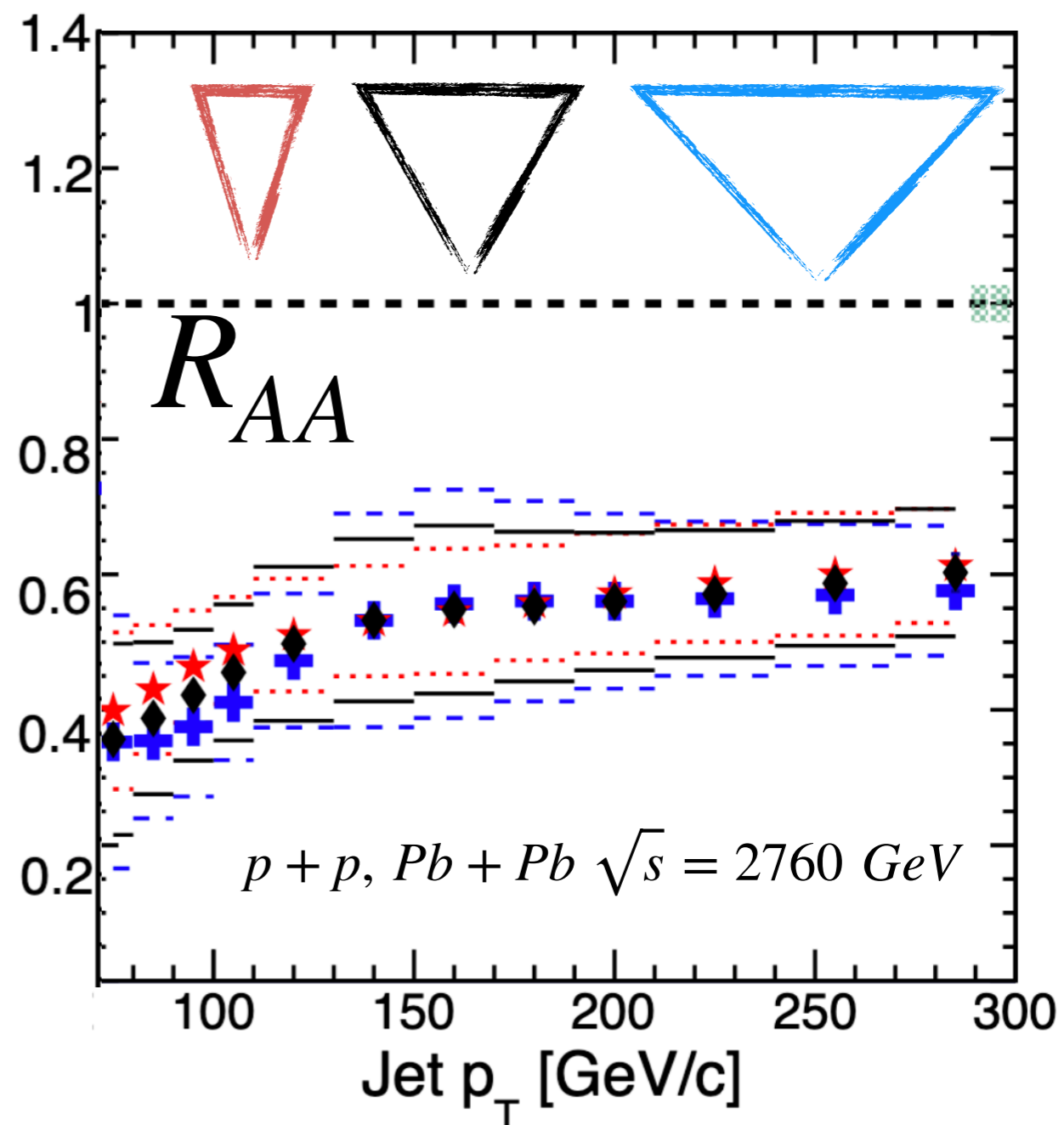
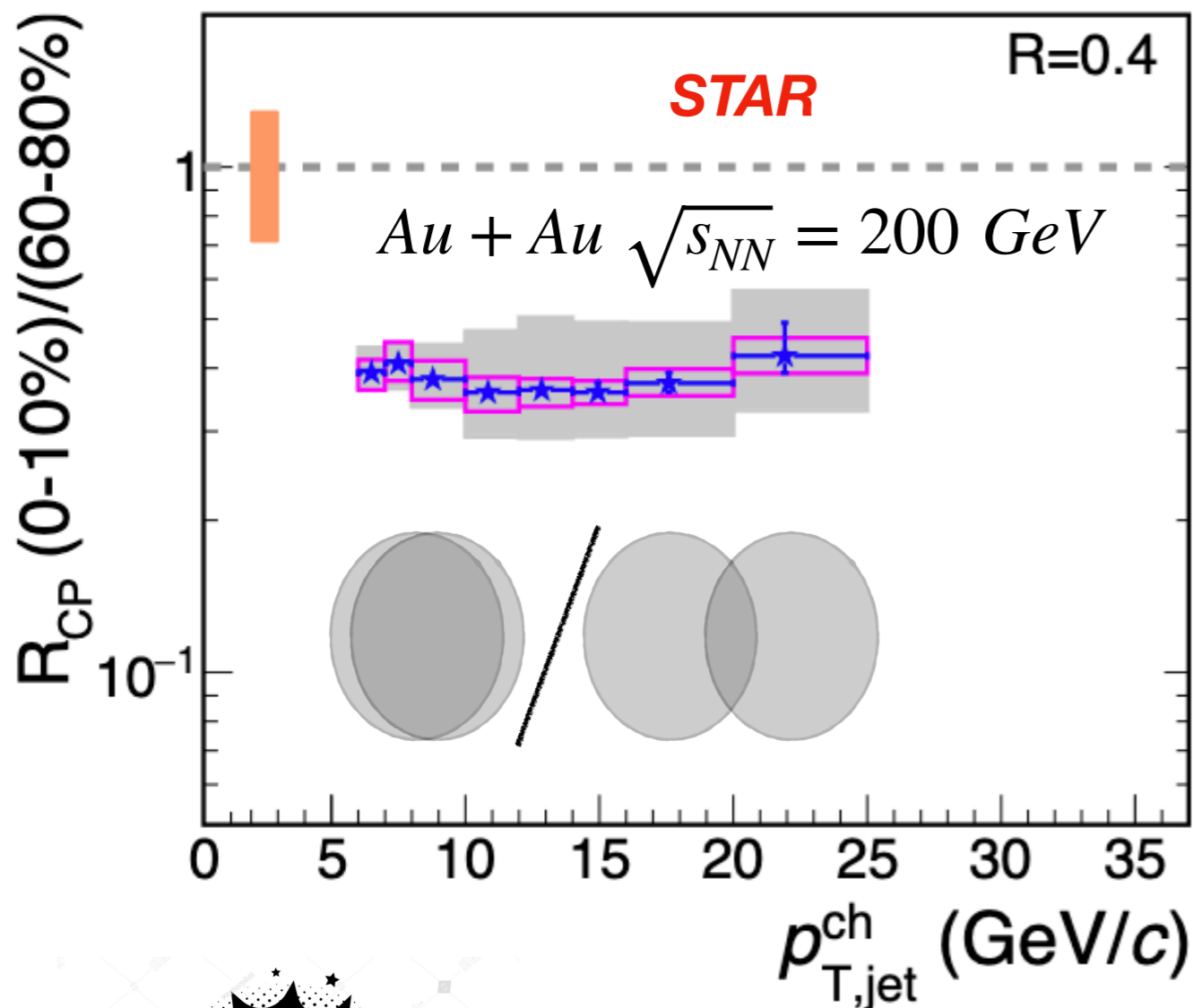


- Similar levels of suppression  $R_{AA}$  ( $R_{CP}$ )  $\sim 0.4-0.6$

# Inclusive jets @ RHIC and LHC

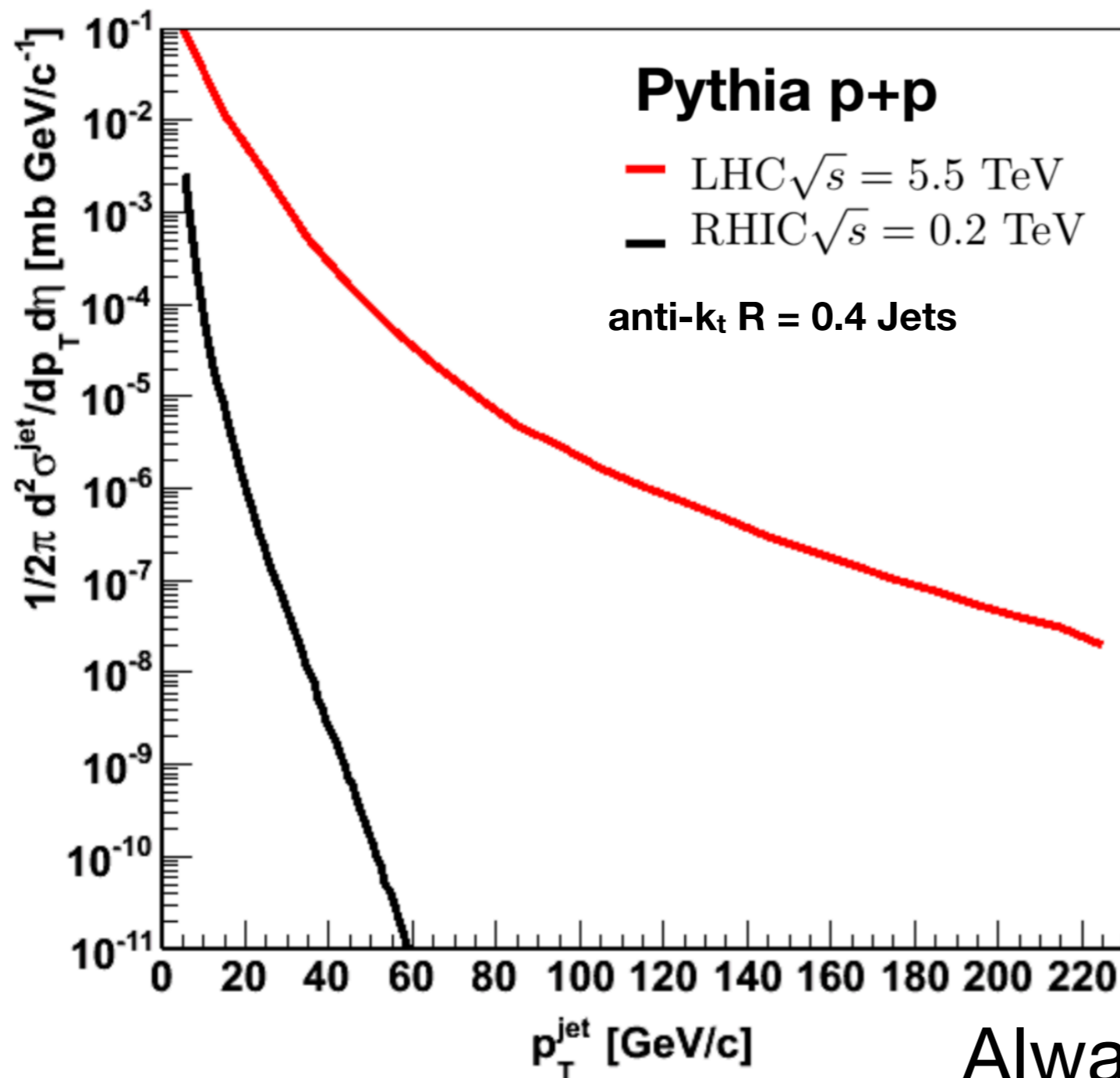
CMS Phys. Rev. C 96 (2017) 015202

STAR Phys. Rev. C 102 5, 054913 (2020)



Is the energy loss the same?

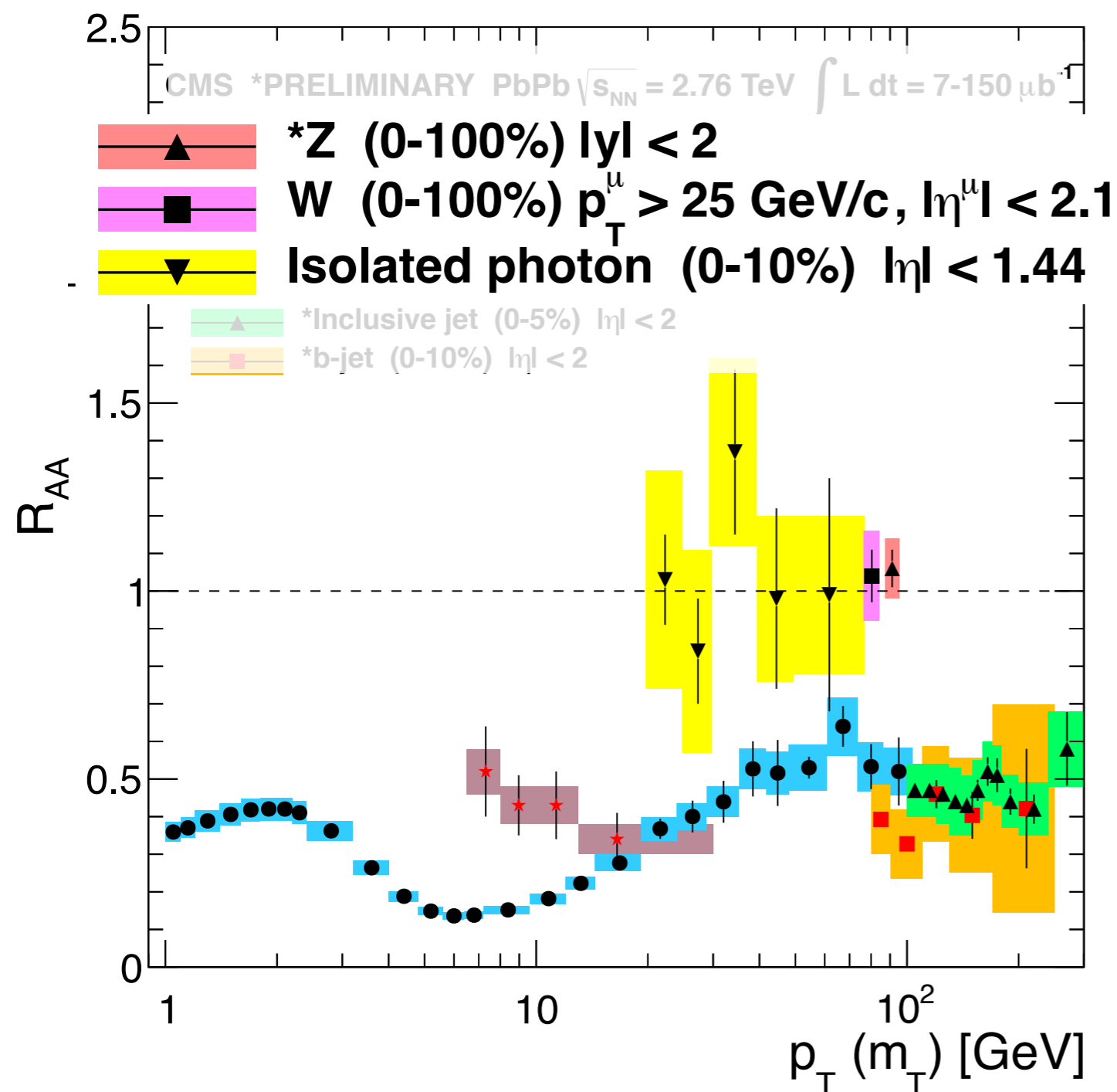
# Lets hold on a second



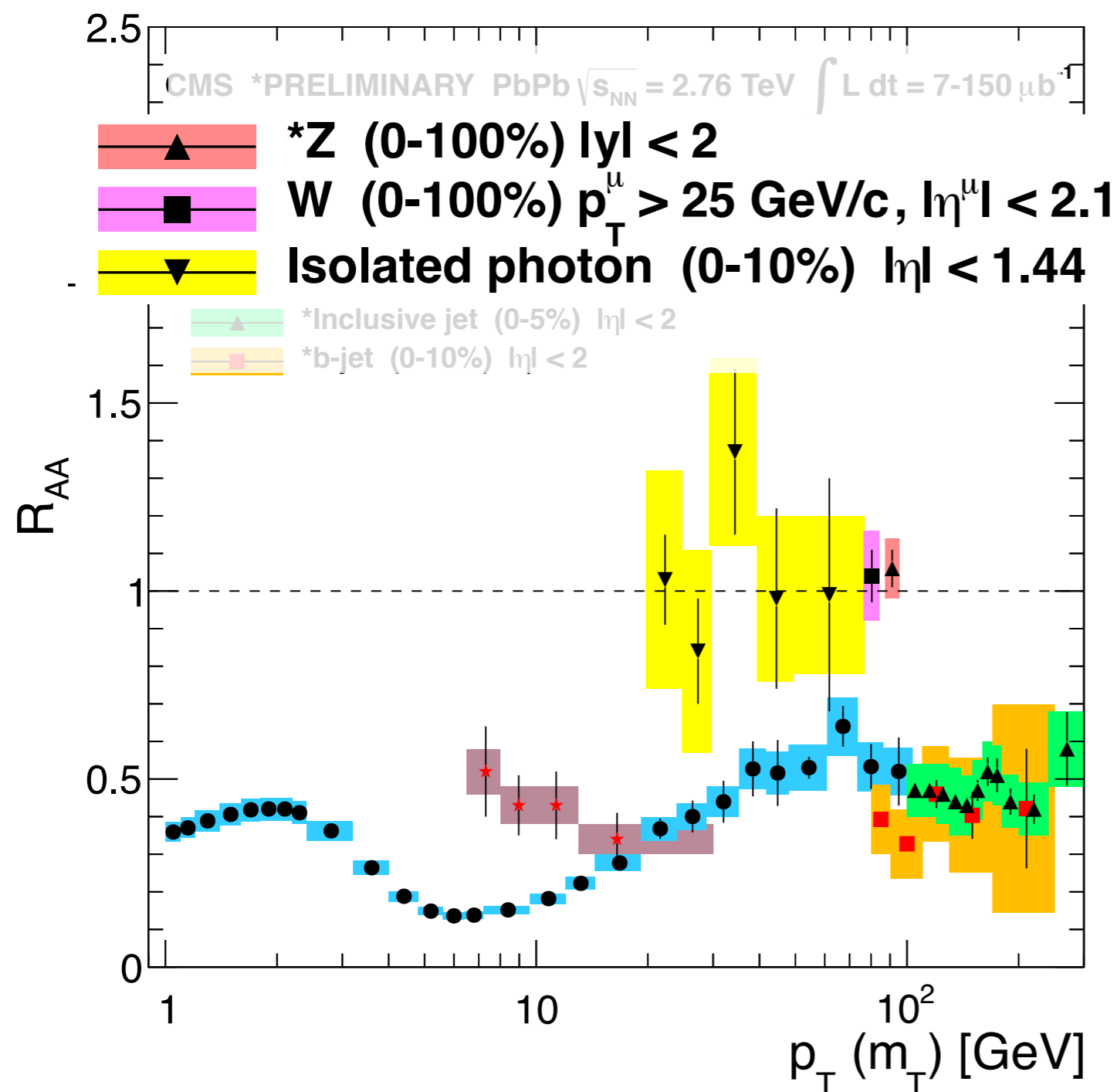
- Energy loss effects shifts the spectrum left
- Similar levels of RAA for a steeper spectrum implies what?
- At fixed jet  $p_T$  - do jets at RHIC or LHC lose more fractional energy?

Always keep the underlying spectrum in mind!

# Lets hold on another second, why do we trust RAA?

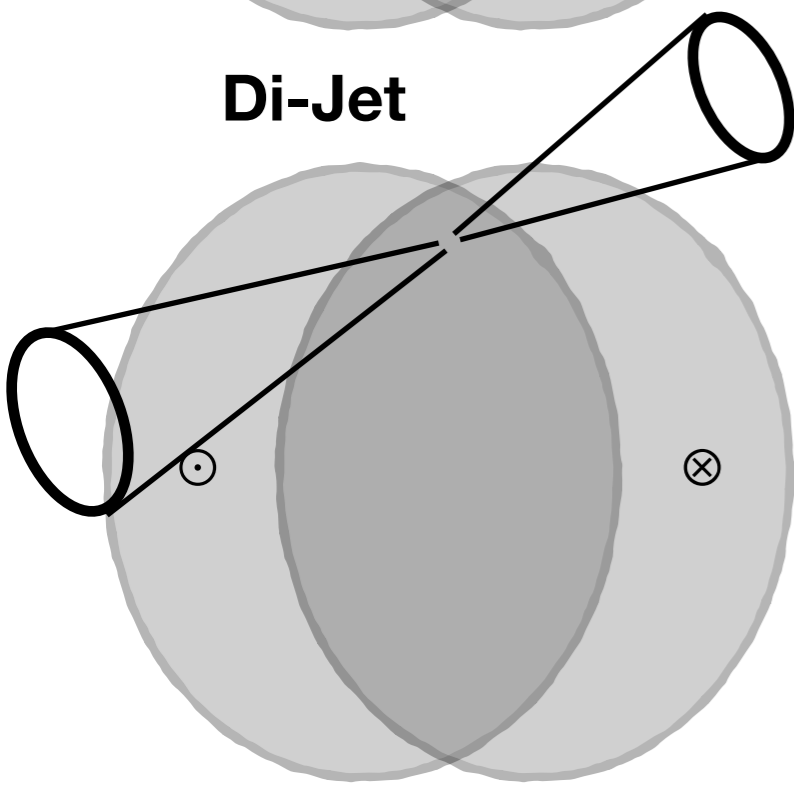
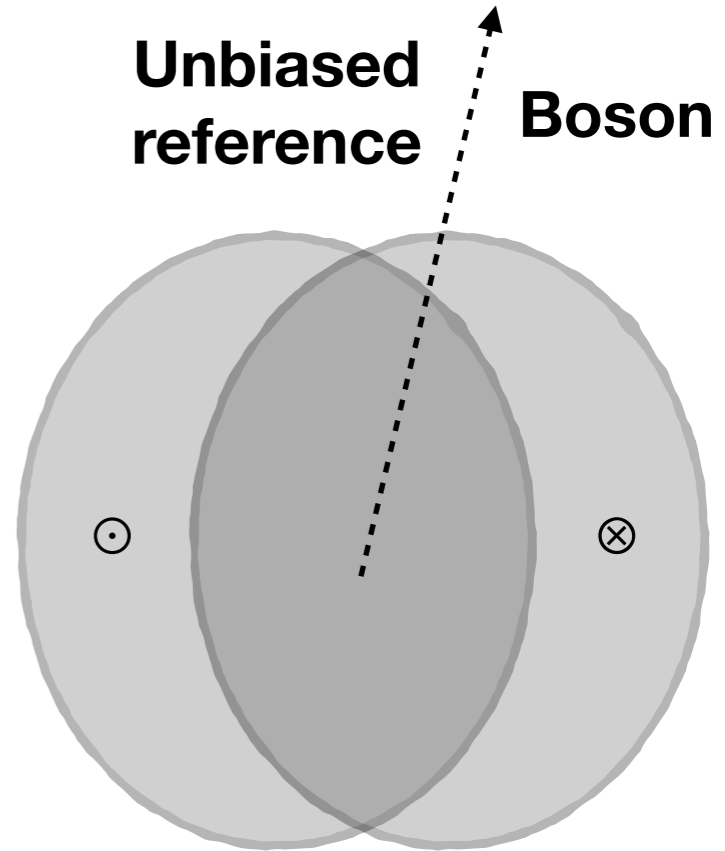
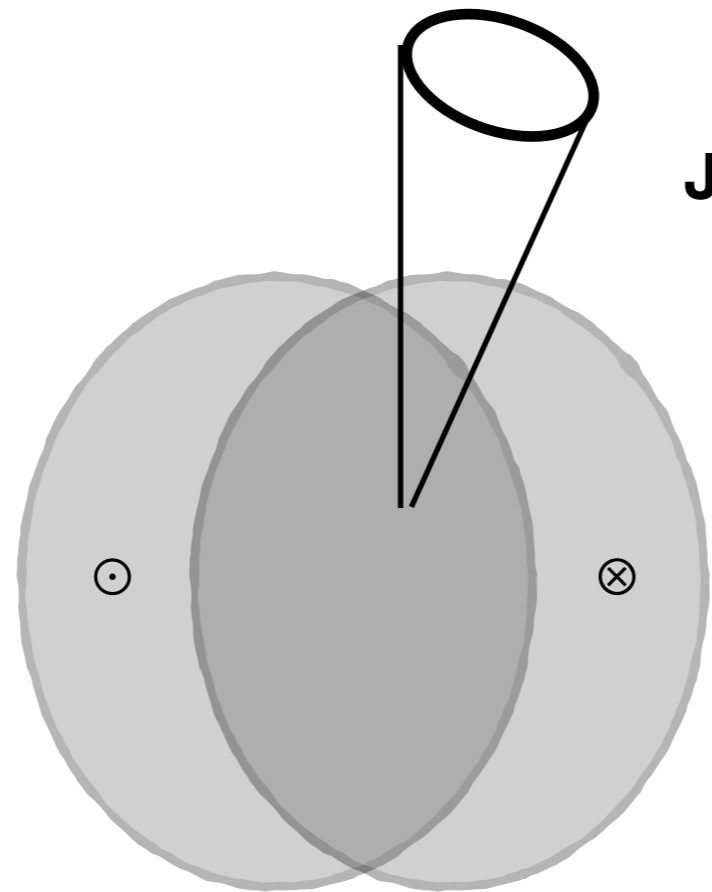
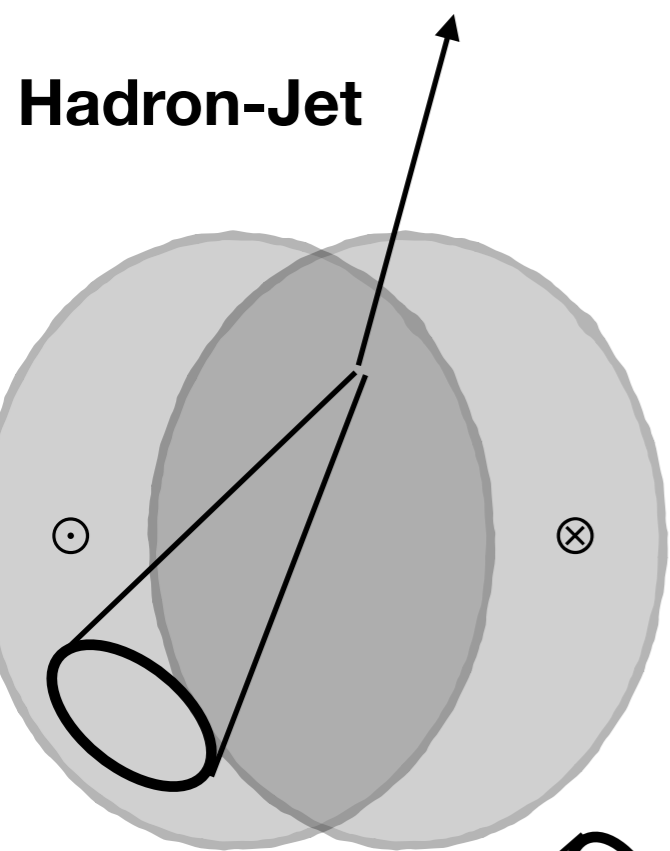


# Lets hold on another second, why do we trust RAA?

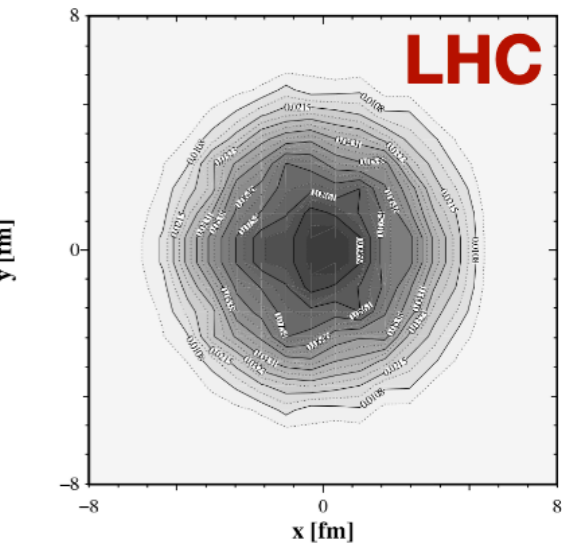
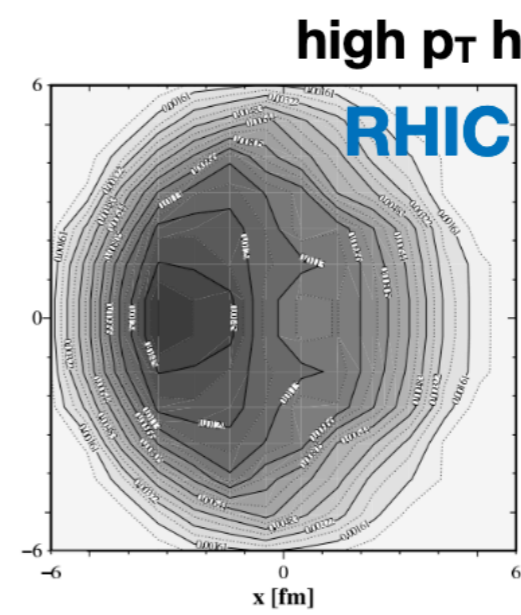


- Scaling by the number of binary collisions reproduces the yield of electro-weak bosons  $\gamma, W^\pm, Z^0$

# Not all events are created equal



- Each of these event selections come with their own biases
- Important to understand their effect on an observable



Renk, PRC 85 064908 (2012), 87 024905 (2013)

# 1st gen - what did we learn?

- Colored probes are opaque whereas QGP appearing transparent to EW probes ( $\gamma$ , Z, W)
- $R_{AA}$  Nuclear modification factor (comparing yield in AA w.r.t binary collisions scaled pp) for  $\gamma/Z \sim 1$ , hadrons  $\sim 0.2$  and **Jet  $R_{AA} \sim 0.5$**  (even at high  $p_T$ ! With mild momentum dependence)
- **Large momentum asymmetry** in Di-jet,  $\gamma/Z$ +Jet pairs
- Large spread of quenched energy - **Broadening effect**

*ATLAS Heavy Ion Publications*

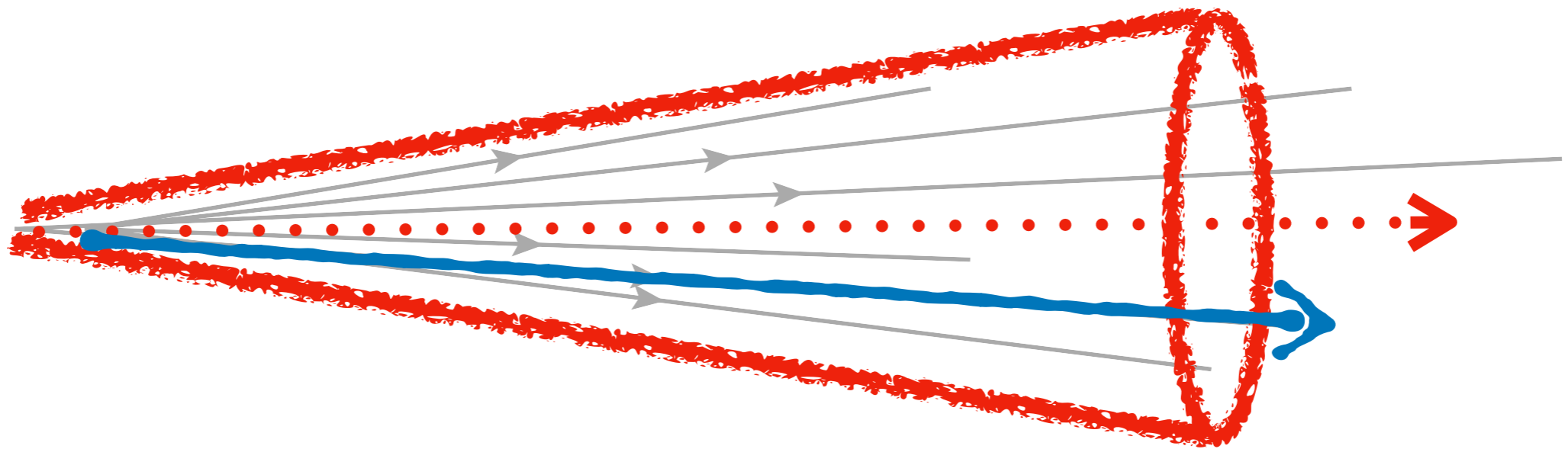
*ALICE Heavy Ion Publications*

*CMS Heavy Ion Publications*

*STAR Publications*

*PHENIX Publications*

# 2nd generation Jet Measurements

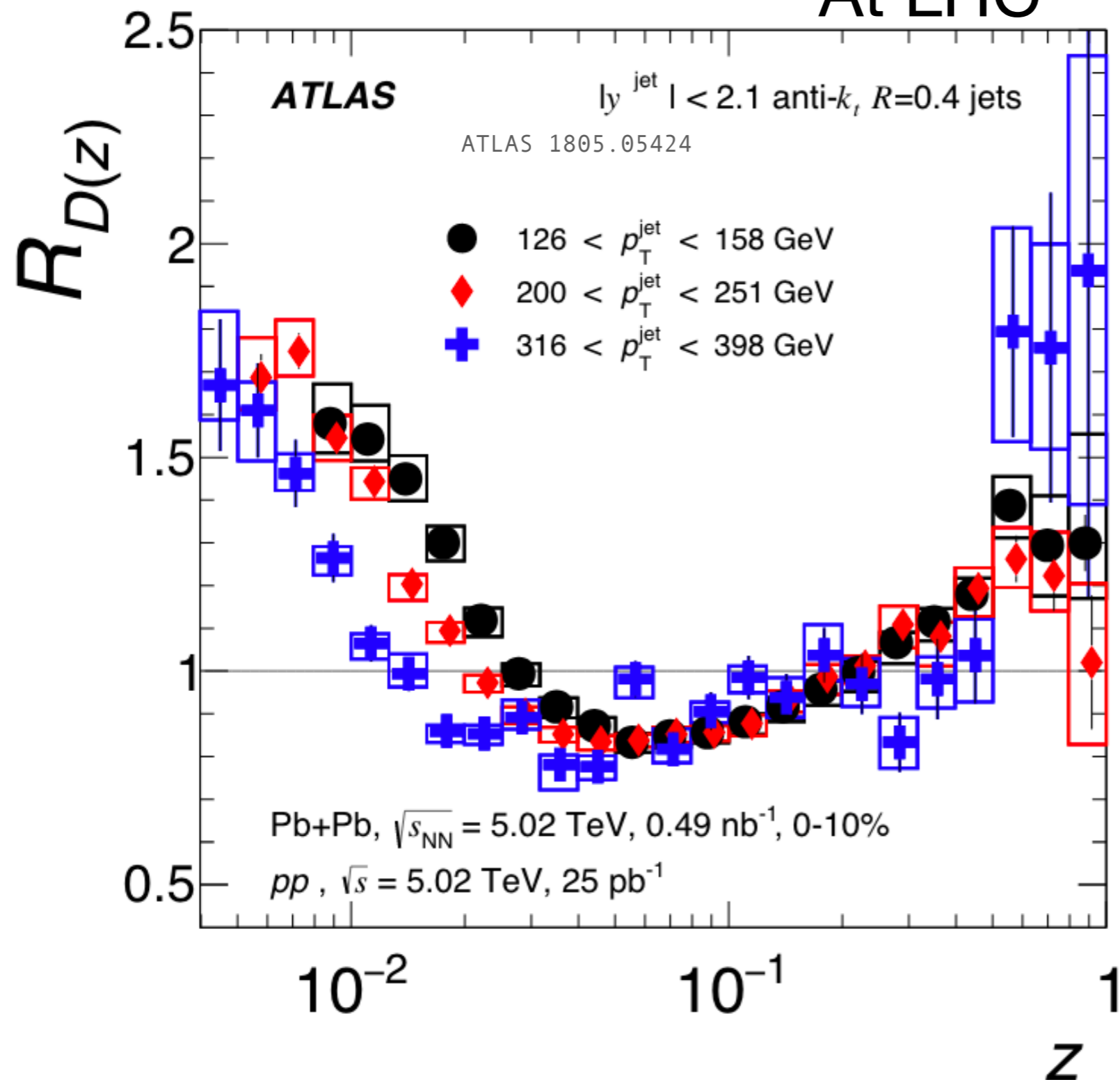


**Jet structure looks at particle production within the jet**



# Intra-jet particle production

At LHC



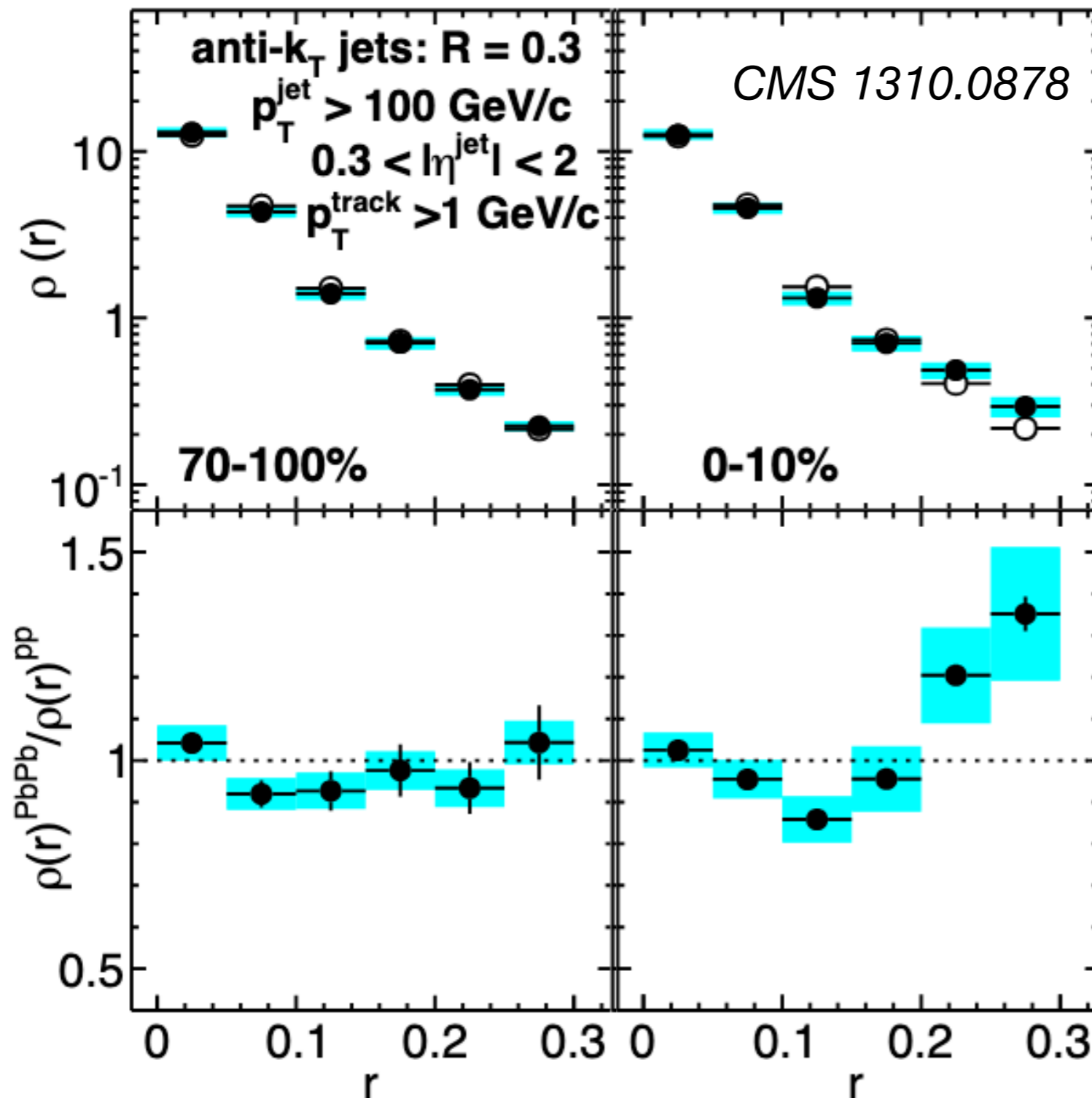
- Enhancement of low  $z$  hadrons at similar  $p_T$  (3.5 GeV) - points to a medium scale!
- Interesting observation of possible enhancement at high  $z$

$$z = \frac{p_{T, \text{trk}}}{p_{T, \text{jet}}}$$

# Intra-jet particle production

At LHC

CMS,  $\sqrt{s_{NN}} = 2.76$  TeV

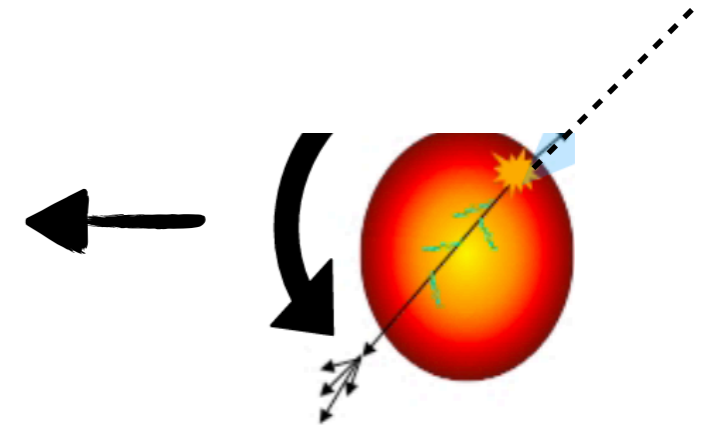
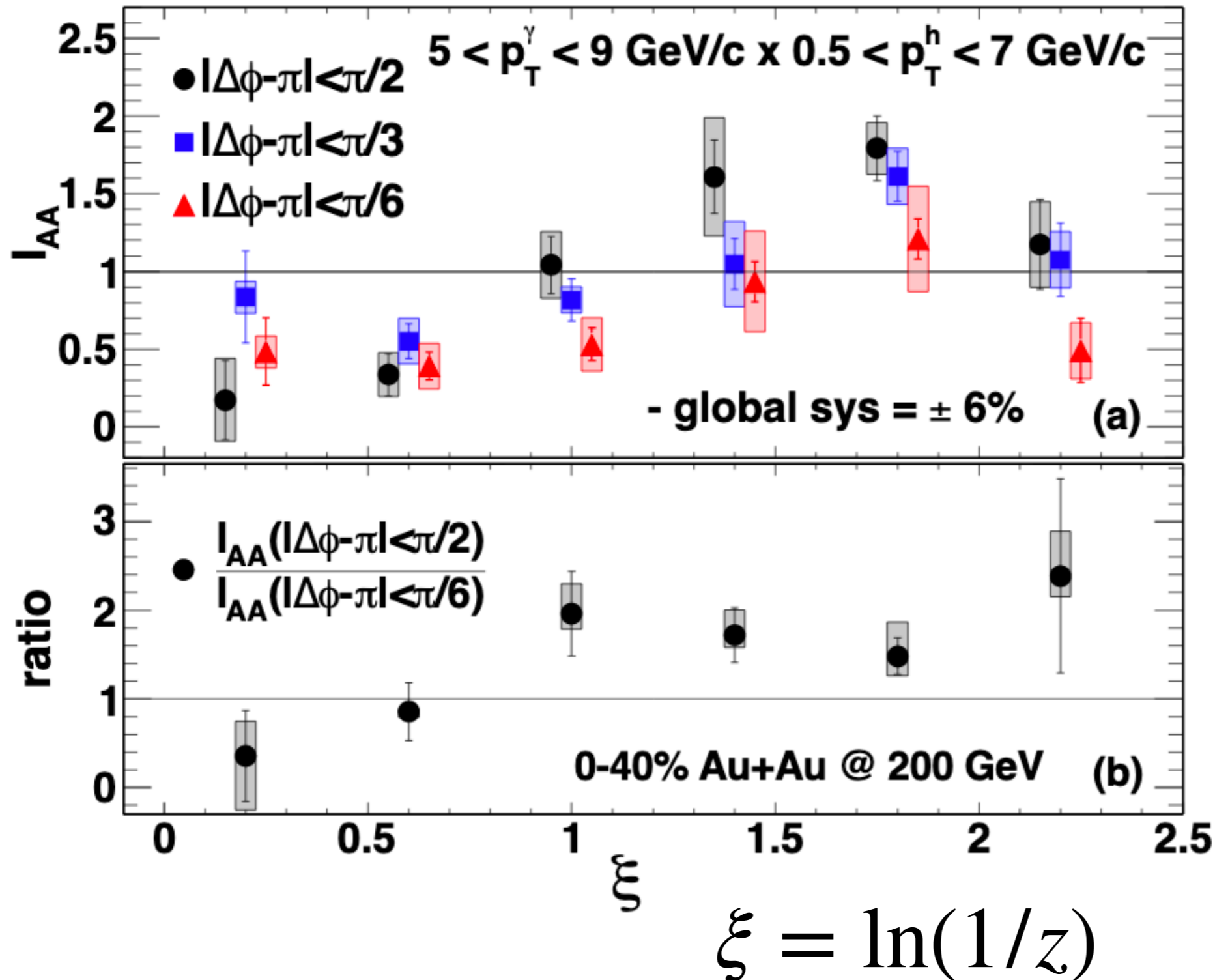


- Enhancement of hadrons around the edges of the jet cone (and extending beyond)
- Interesting observation of possible narrowing at the jet core!
- Convolution of effects including varying parton flavor and kinematic selection

# Jet/ $\gamma$ -Hadron correlations

At RHIC

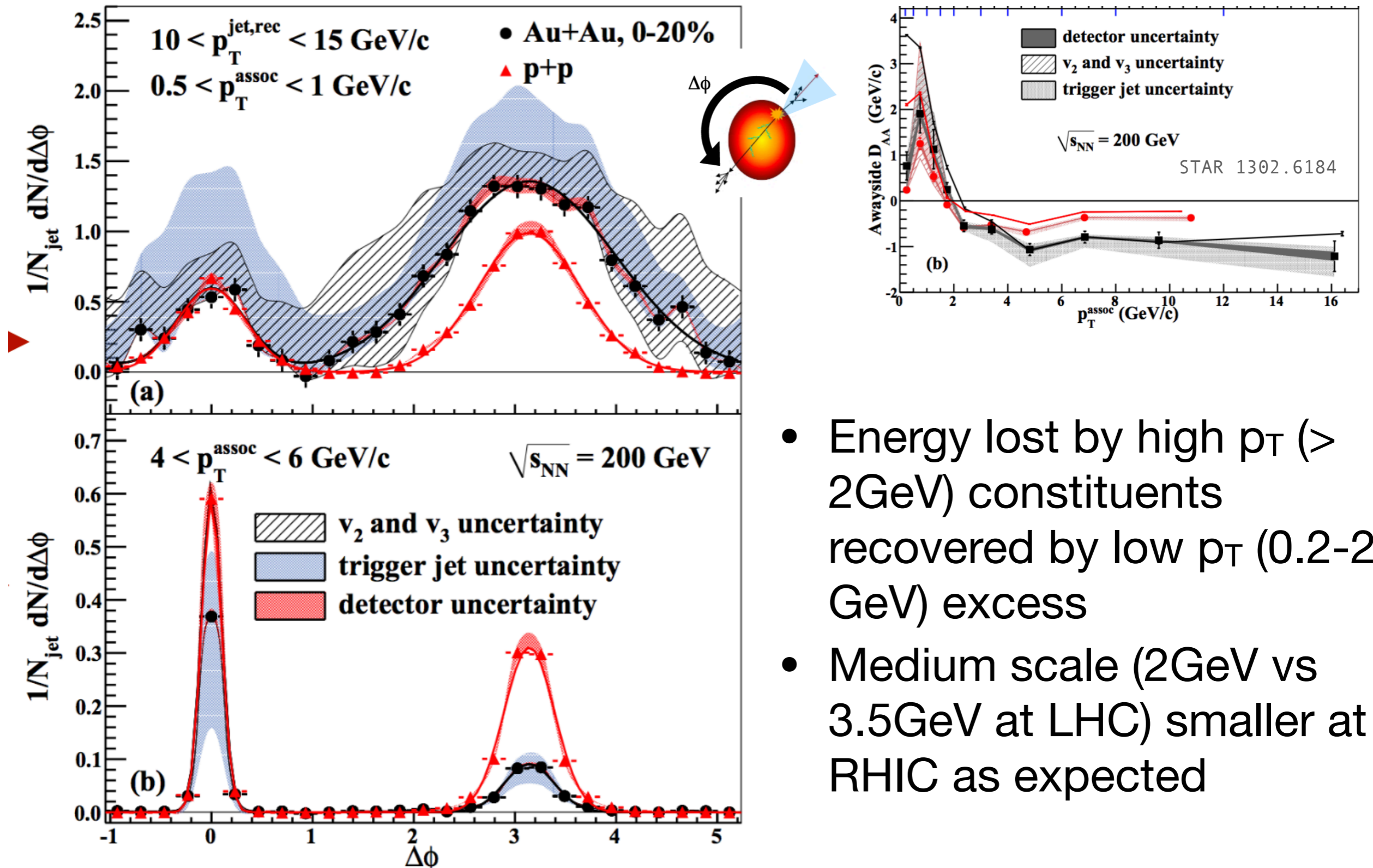
PHENIX 1212.3323



- Suppression of low  $\xi$  (high  $p_T$  tracks) recoiling off photons
- Clear and consistent signal of enhancement and broadening of recoil jet's low  $p_T$  constituents.

# Jet/ $\gamma$ -Hadron correlations

At RHIC



- Energy lost by high  $p_T$  ( $> 2\text{GeV}$ ) constituents recovered by low  $p_T$  (0.2-2 GeV) excess
- Medium scale (2GeV vs 3.5GeV at LHC) smaller at RHIC as expected

# 2nd gen - what did we learn?

- Observation of **fragmentation modification at low z and around the jet** - Highlights need for and use of calibrated probes with good reference
- Modification appears to occur at a fixed energy scale
- Large spread of quenched energy - **Broadening effect**
- But we still don't know the 'How' of energy loss

*ATLAS Heavy Ion Publications*

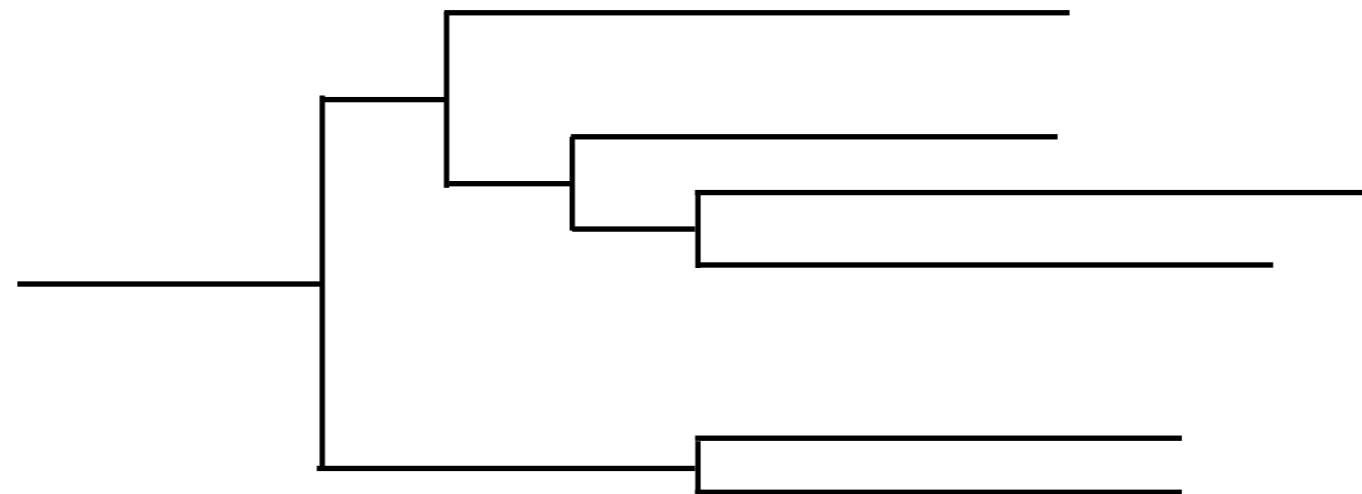
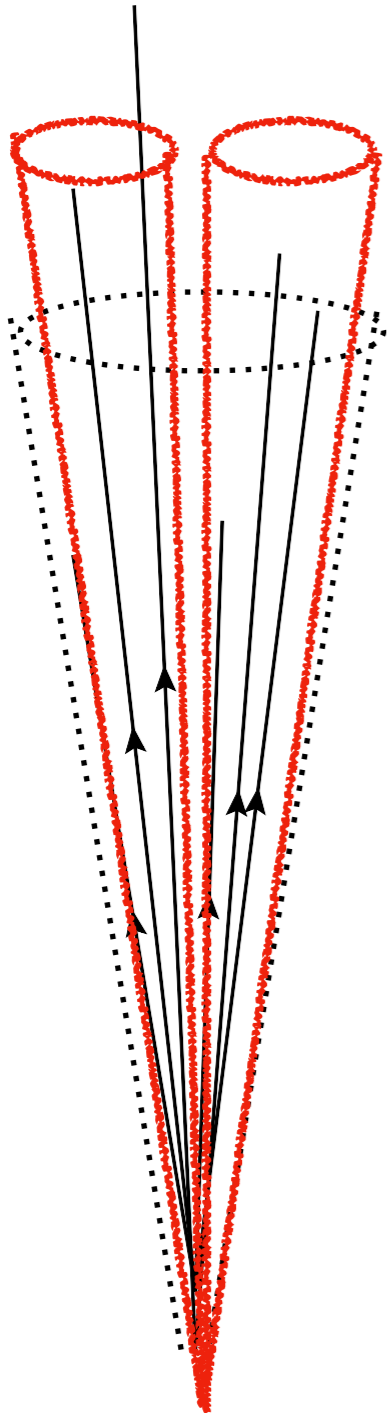
*ALICE Heavy Ion Publications*

*CMS Heavy Ion Publications*

*STAR Publications*

*PHENIX Publications*

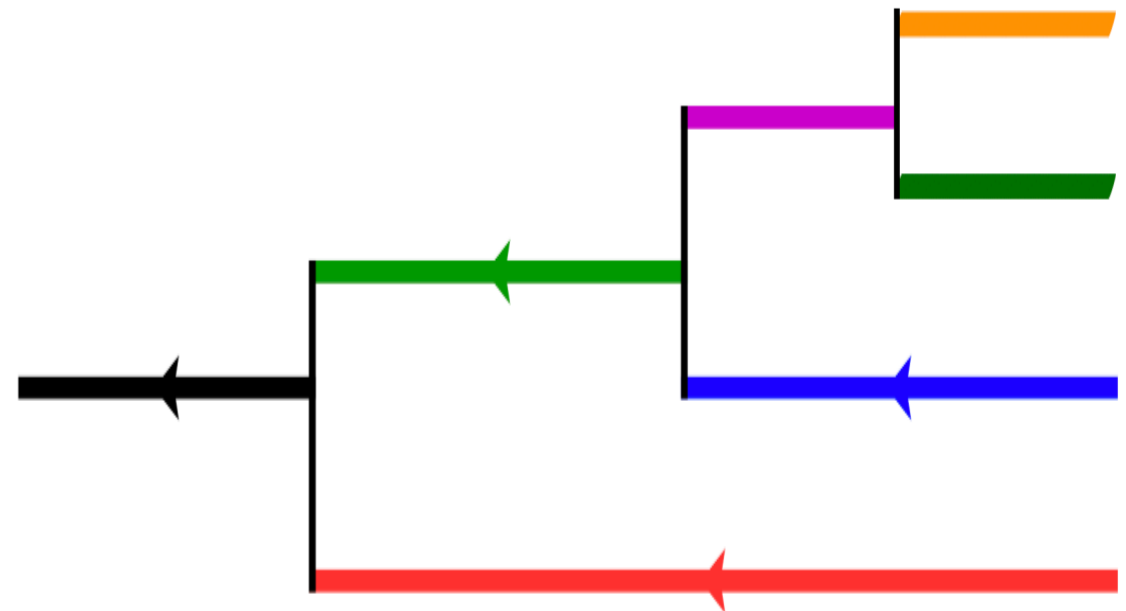
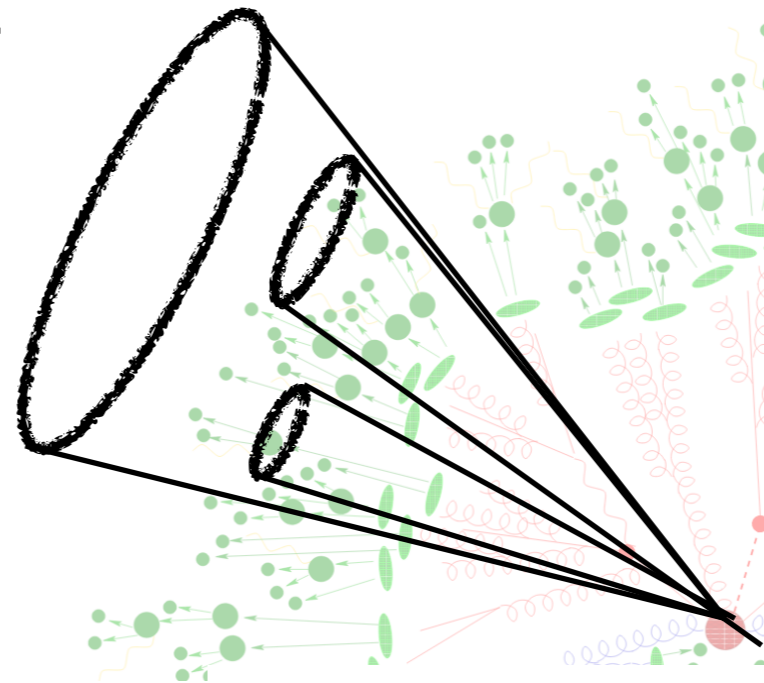
# 3rd generation Jet Measurements



**Exploit clustering  
information to look inside  
jets and study the evolution**

# What is Jet SubStructure

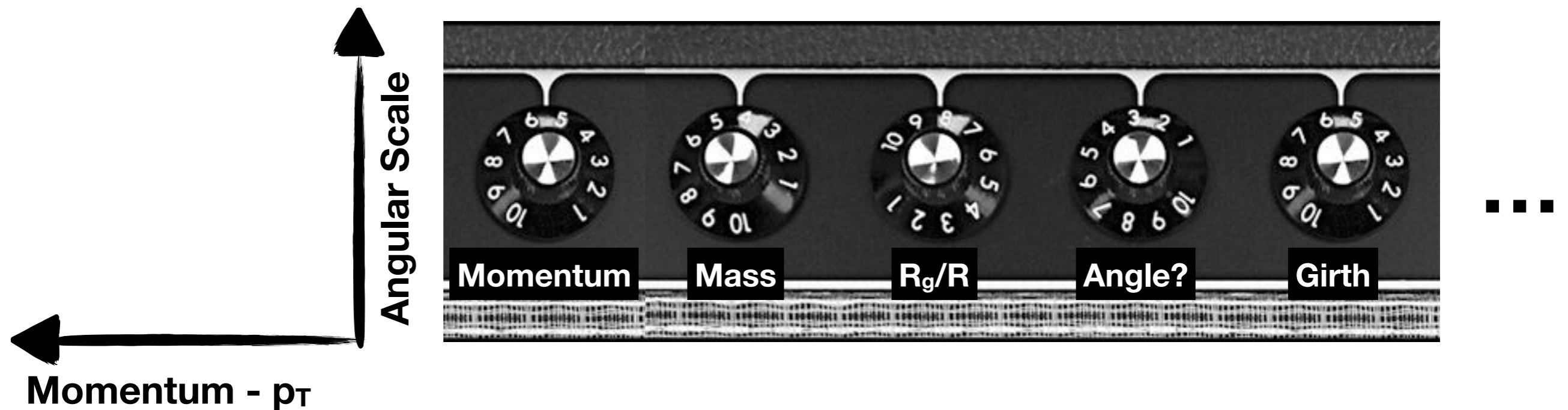
- Structures within jets - built via correlations or algorithms
- SoftDrop - Utilizing the clustering information to build a tree
- JSS observables such as  $z$ ,  $\theta$  are of course dependent on the algorithm and jet definitions etc...



# Key Idea

## Use jet-substructure as a selection tool

Identify jet observable(s) sensitive to the parton shower kinematics



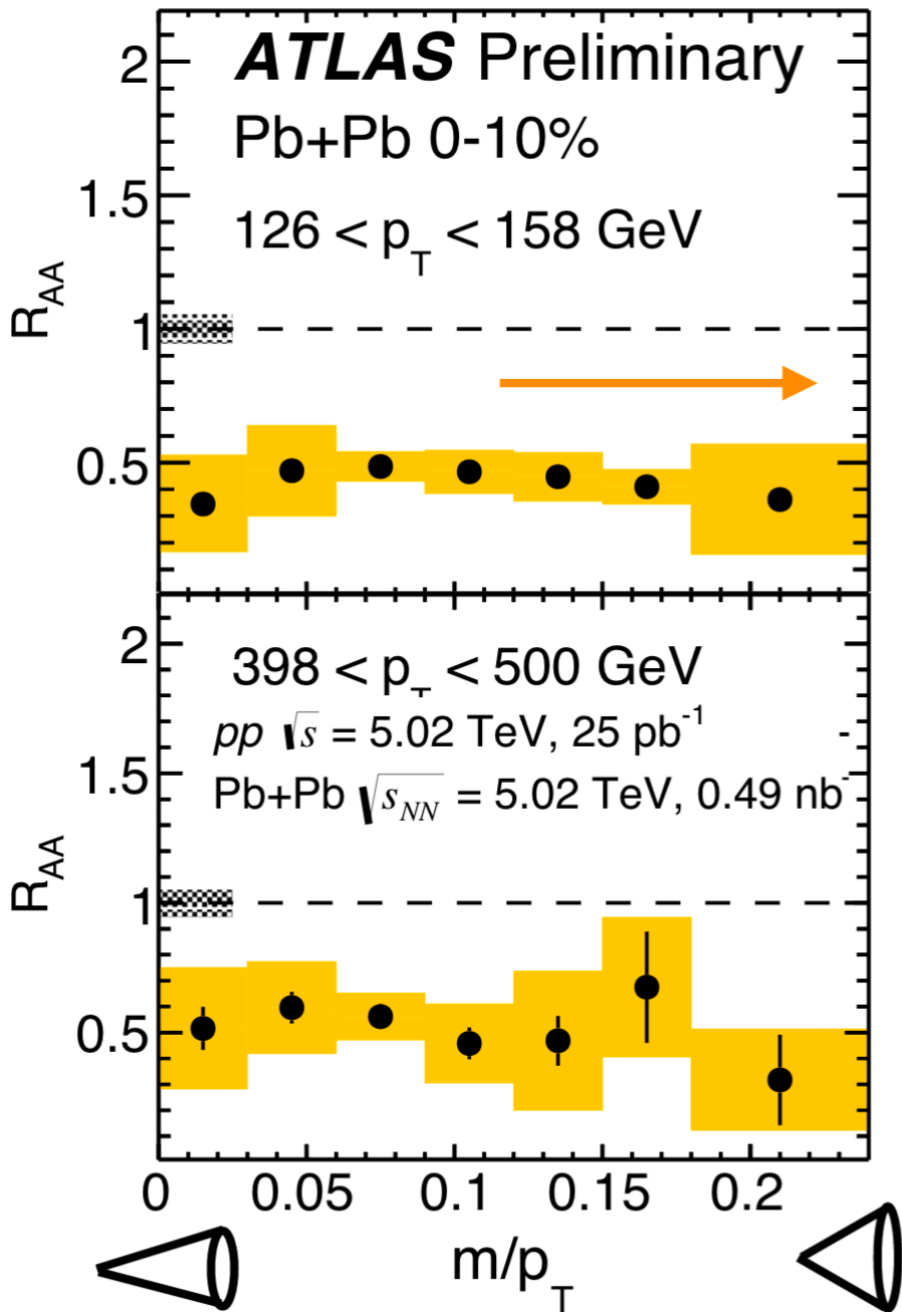
**Partonic energy loss via a differential study in momentum scale and angular scale**



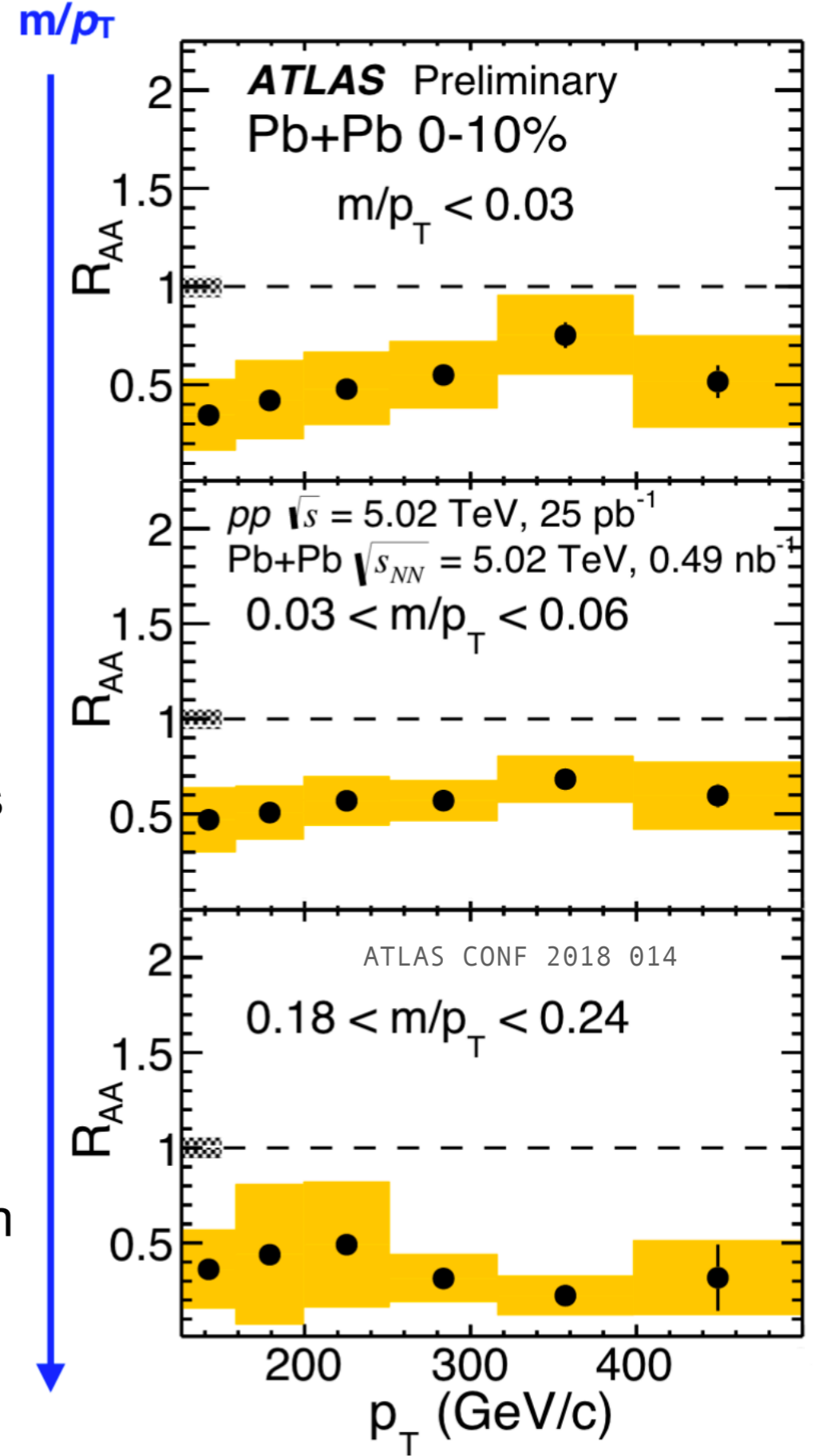
# R<sub>AA</sub> for various Mass selections



Figures taken from Laura Haver's talk at WWND 2019



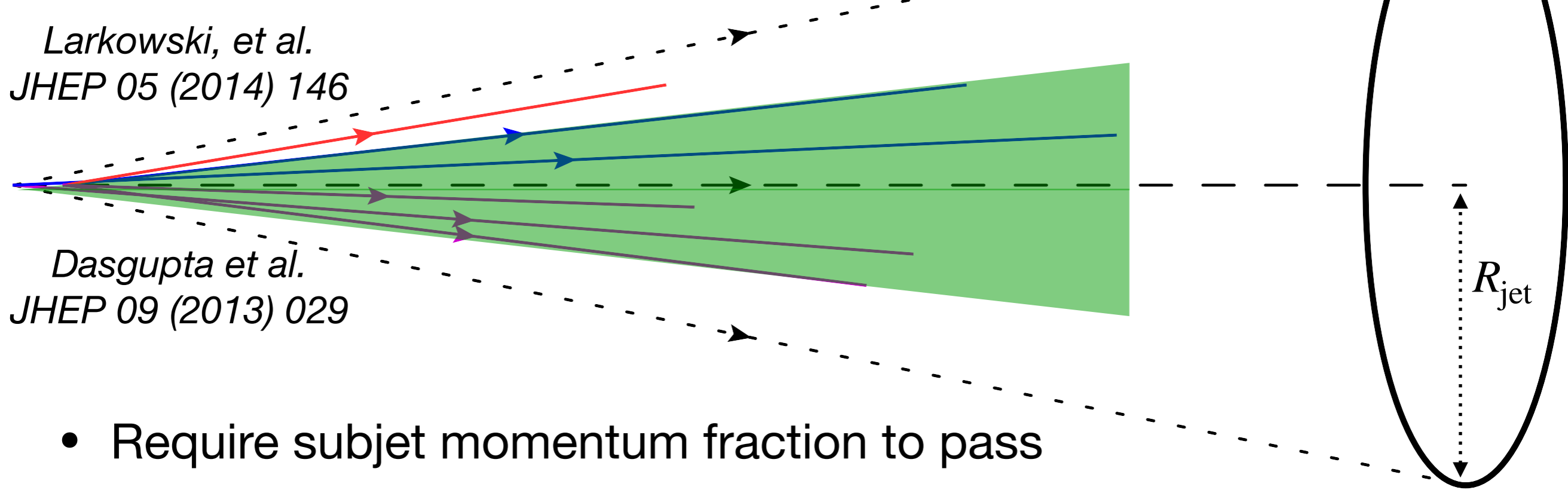
- Observe no significant effect of varying m/p<sub>T</sub>
- Jet p<sub>T</sub> possibly too high?
- m/p<sub>T</sub> bins include jets of varying mass/ smeared resolution scales?
- Mass cancellation effects from quenching vs medium response possible?



# SoftDrop grooming

*Larkowski, et al.*  
*JHEP 05 (2014) 146*

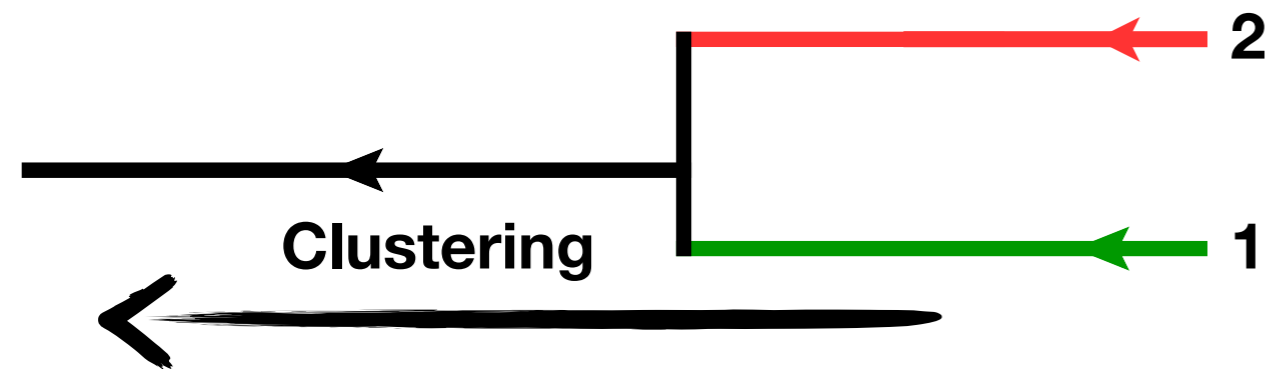
*Dasgupta et al.*  
*JHEP 09 (2013) 029*



- Require subjet momentum fraction to pass

$$z_g = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}} > z_{\text{cut}} (R_g / R_{\text{jet}})^\beta$$

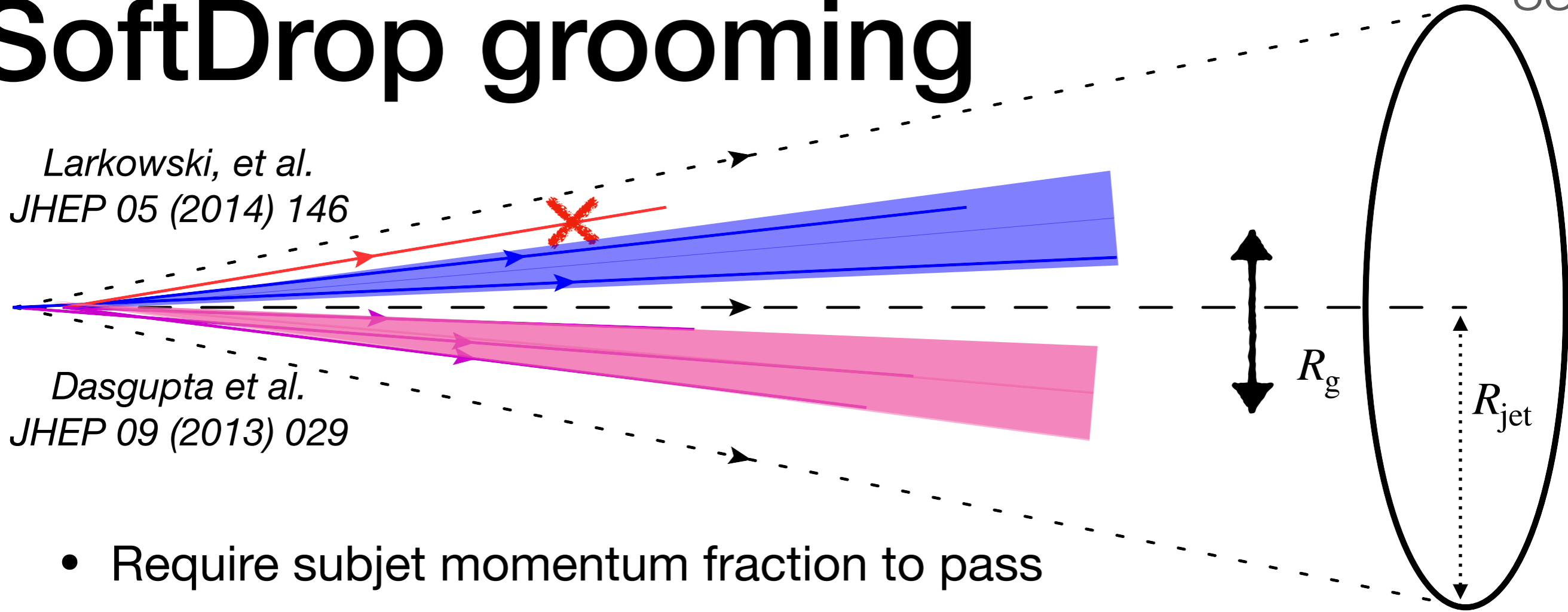
$z_{\text{cut}} = 0.1$   
 $\beta = 0$



# SoftDrop grooming

Larkowski, et al.  
JHEP 05 (2014) 146

Dasgupta et al.  
JHEP 09 (2013) 029

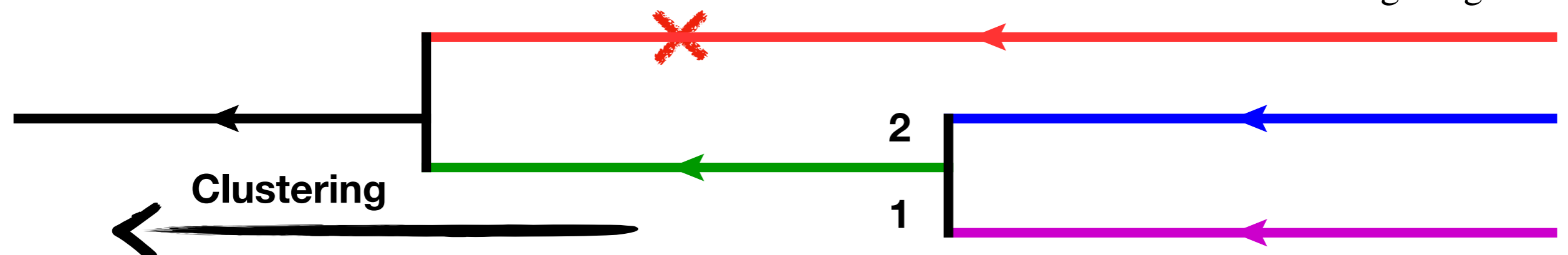


- Require subjet momentum fraction to pass

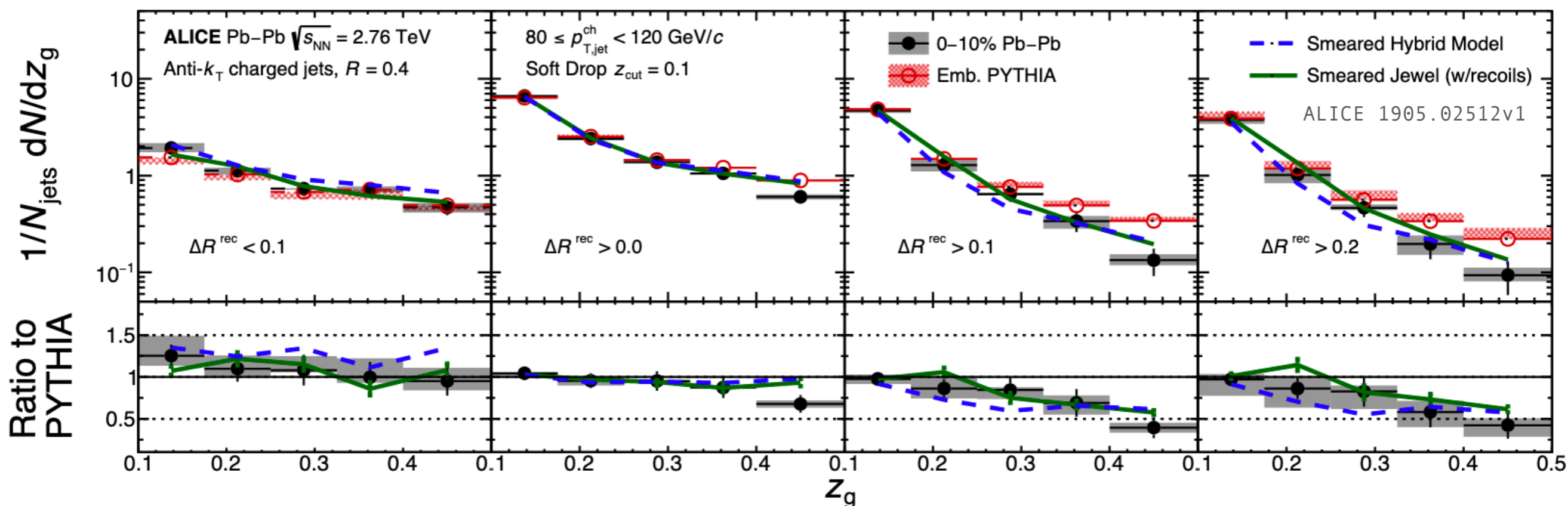
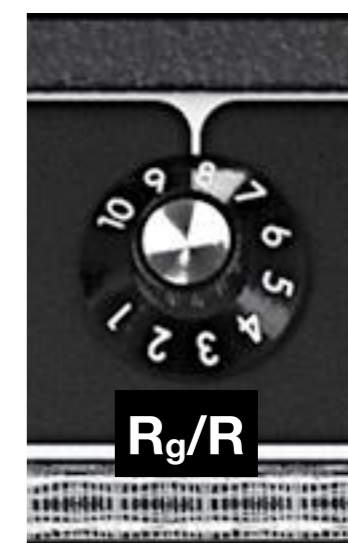
$$z_g = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}} > z_{cut} (R_g / R_{jet})^\beta$$

$z_{cut} = 0.1$   
 $\beta = 0$

- With the two surviving branches (first hard split) - we define observables that characterize jet substructure  $z_g, R_g$



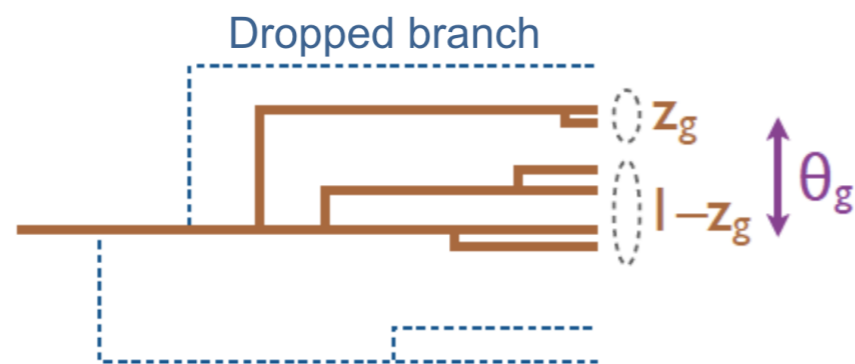
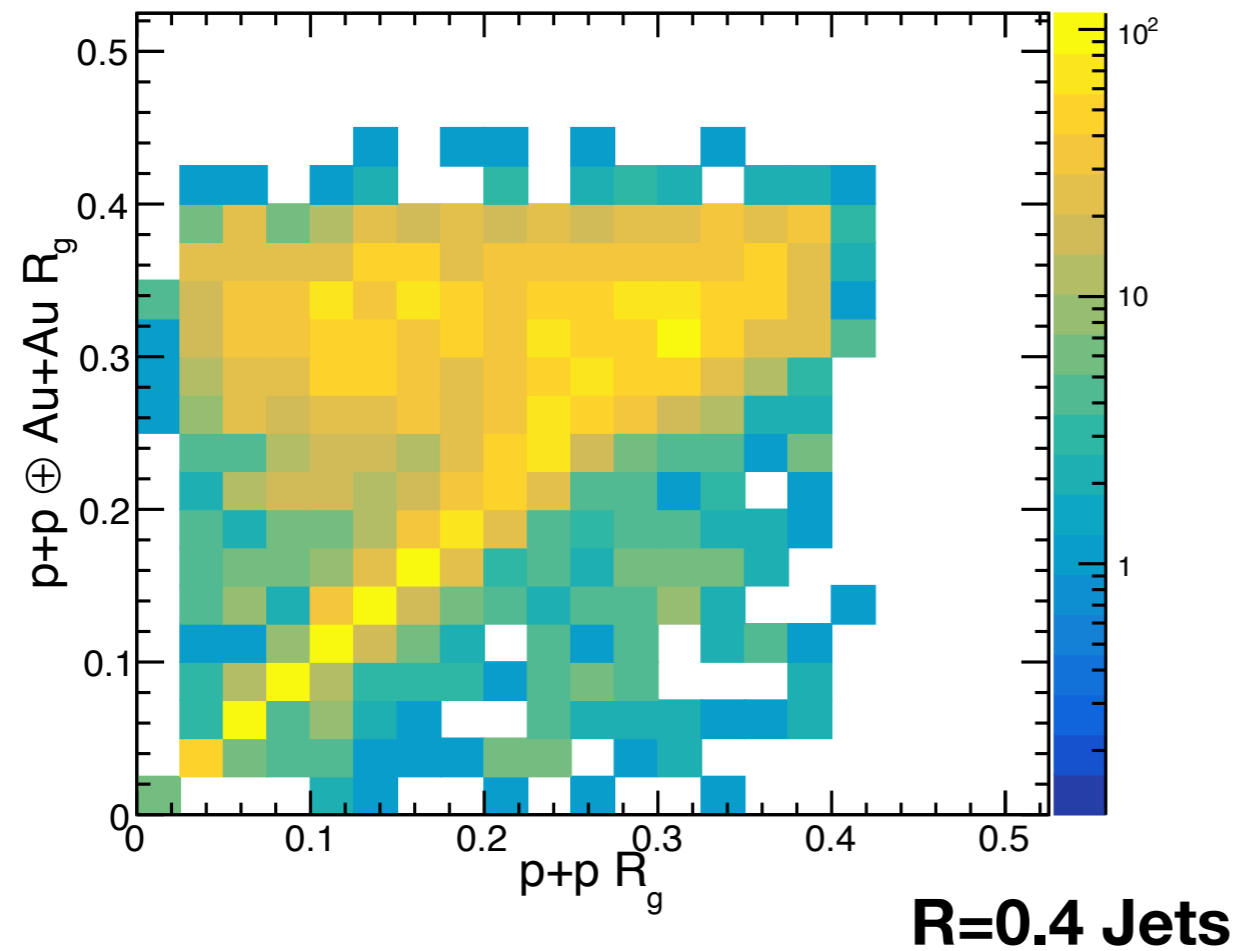
# Softdrop splitting at varying opening angles



- Suppression of wide angle  $z_g$  splits and enhancement of narrow angle  $z_g$
- MC models are generally able to reproduce the trend but further systematic studies are needed to discriminate these models

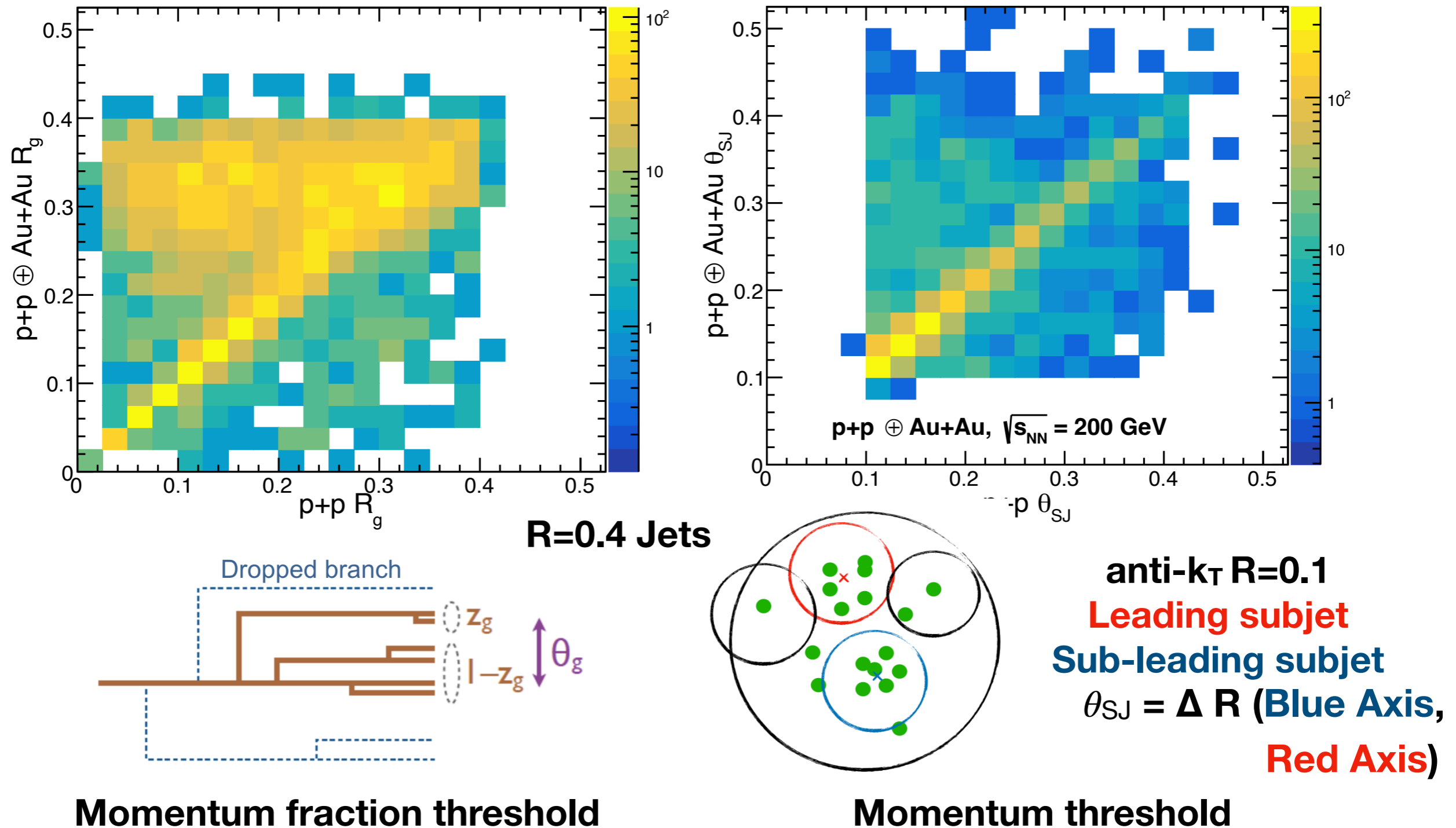
# Choosing a robust observable

*STAR Phys. Rev. C 105, 044906 (2022)*

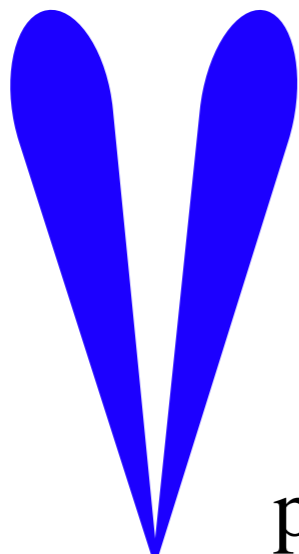


# Choosing a robust observable

STAR Phys. Rev. C 105, 044906 (2022)



$0.1 < \theta < 0.2$     $0.2 < \theta < 0.3$

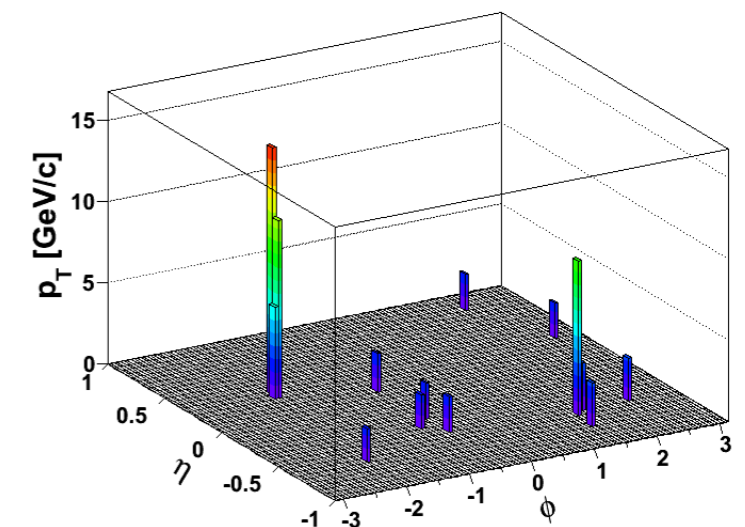
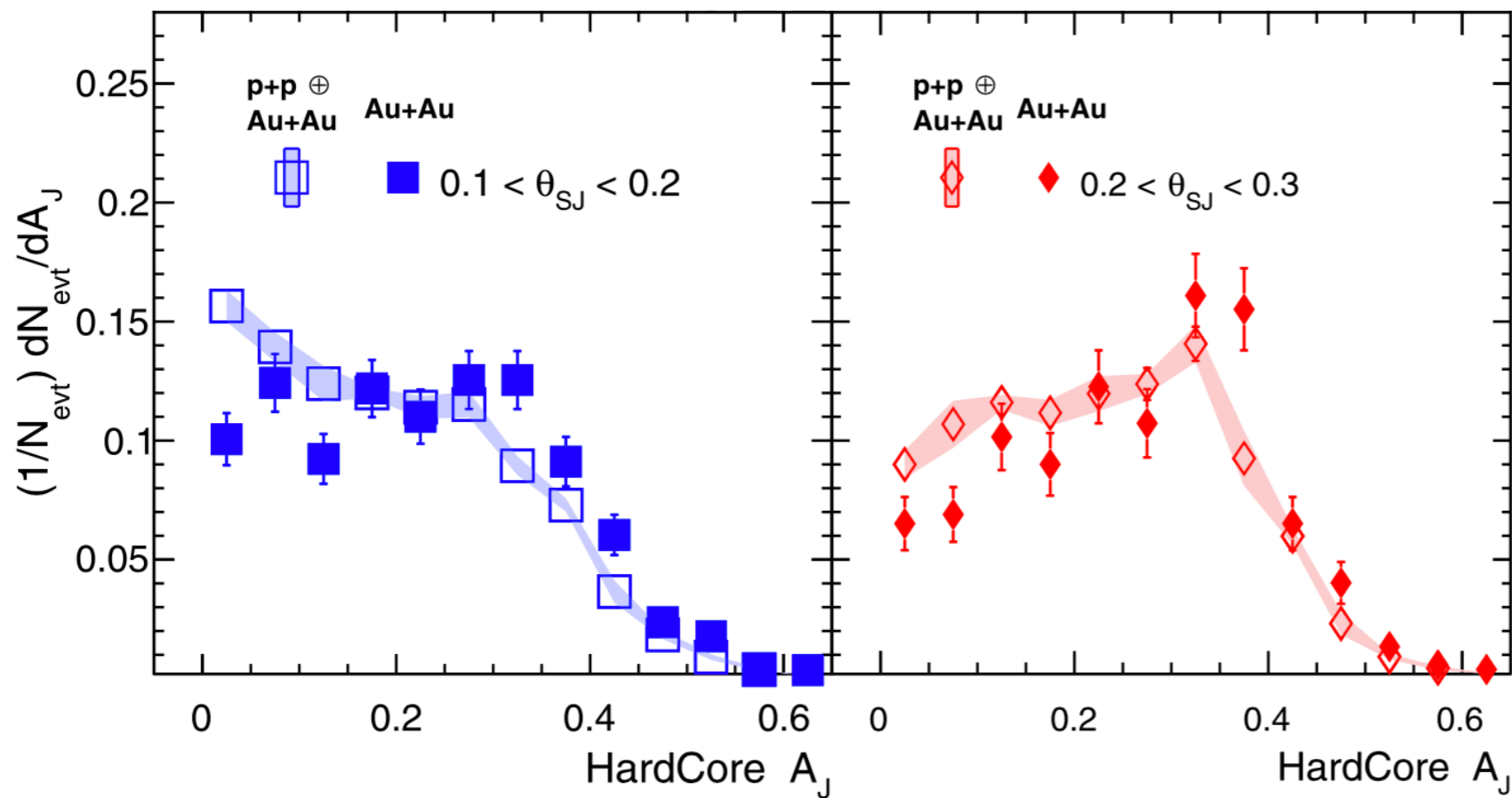


$$A_J = \frac{p_{T,jet}^{Trigger} - p_{T,jet}^{Recoil}}{p_{T,jet}^{Trigger} + p_{T,jet}^{Recoil}}$$

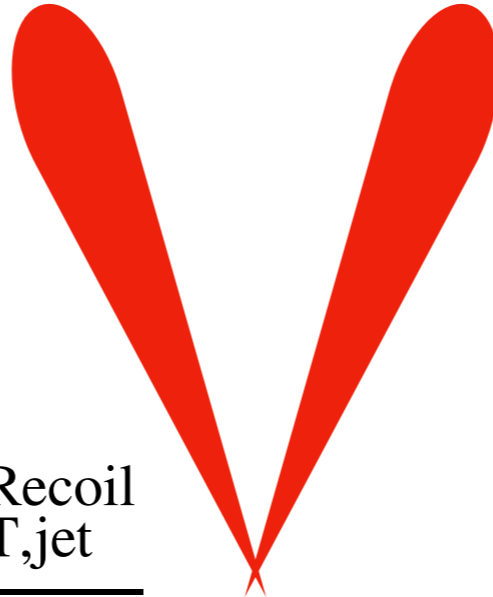
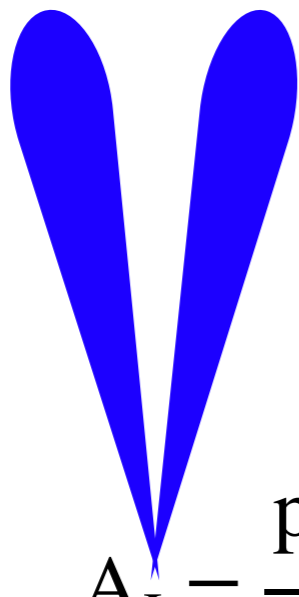
STAR 2109.09793



- Imbalance is independent on opening angle



$0.1 < \theta < 0.2$     $0.2 < \theta < 0.3$

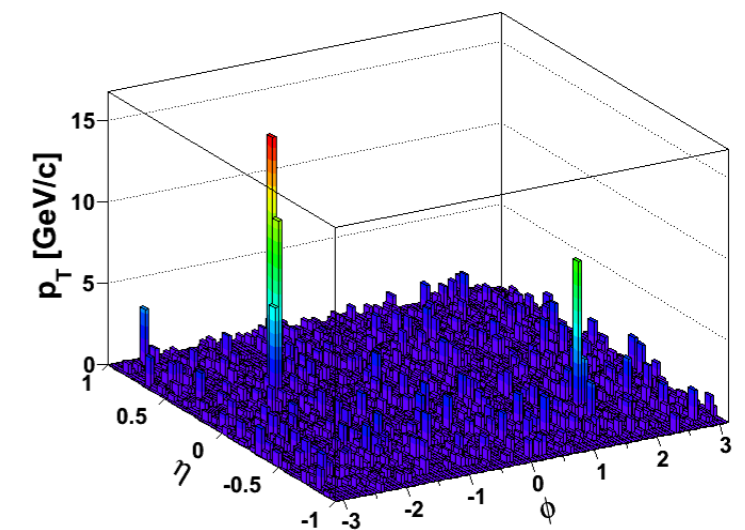
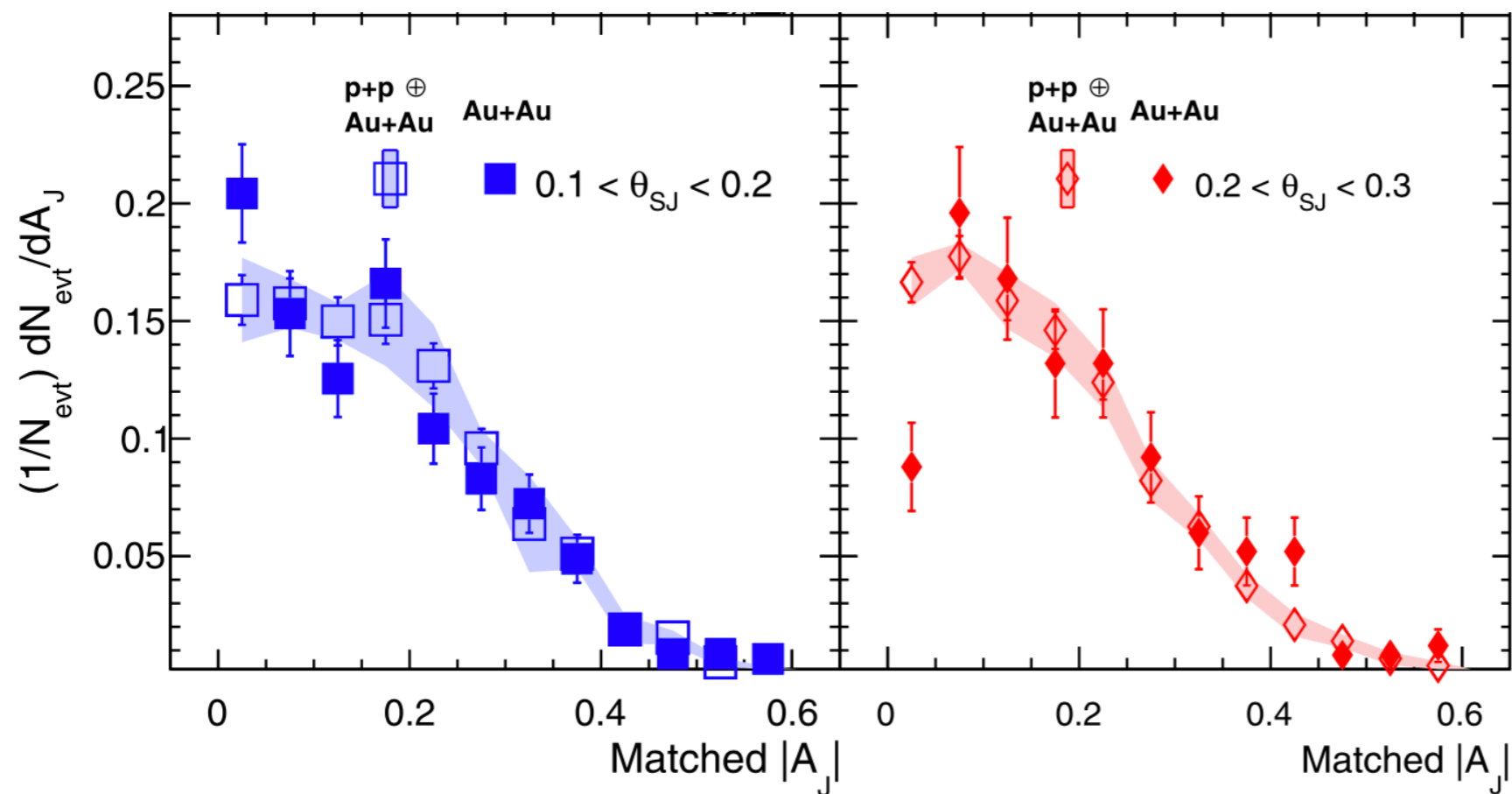


$$A_J = \frac{p_{T,\text{jet}}^{\text{Trigger}} - p_{T,\text{jet}}^{\text{Recoil}}}{p_{T,\text{jet}}^{\text{Trigger}} + p_{T,\text{jet}}^{\text{Recoil}}}$$

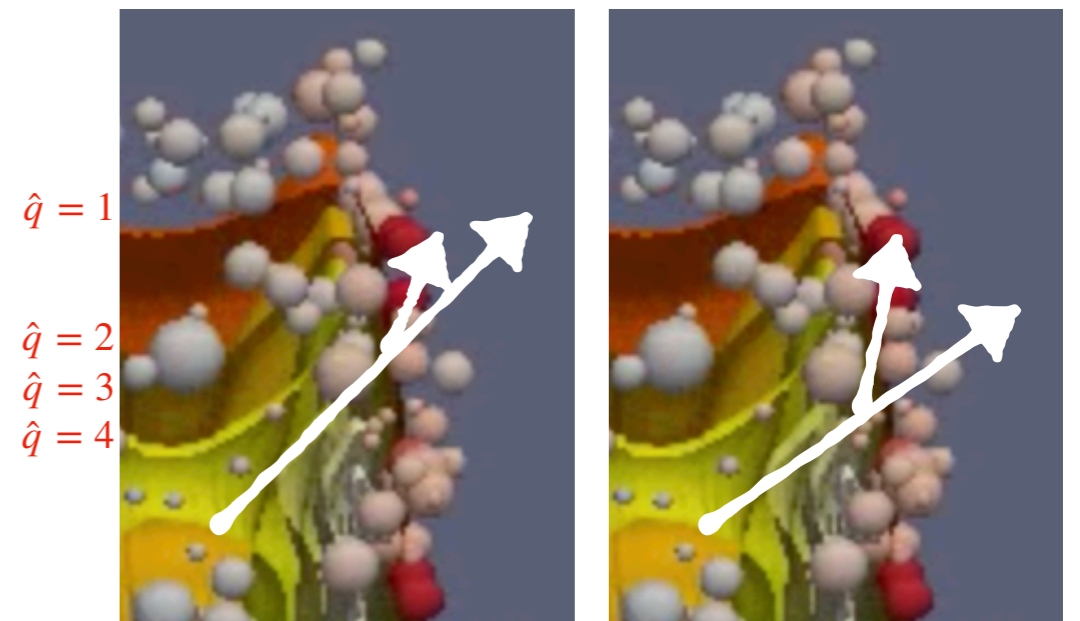
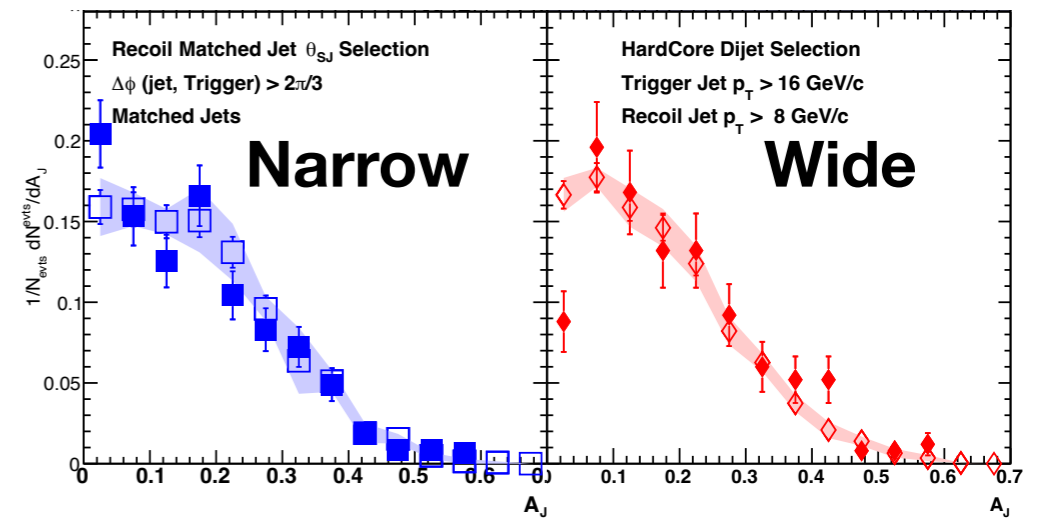
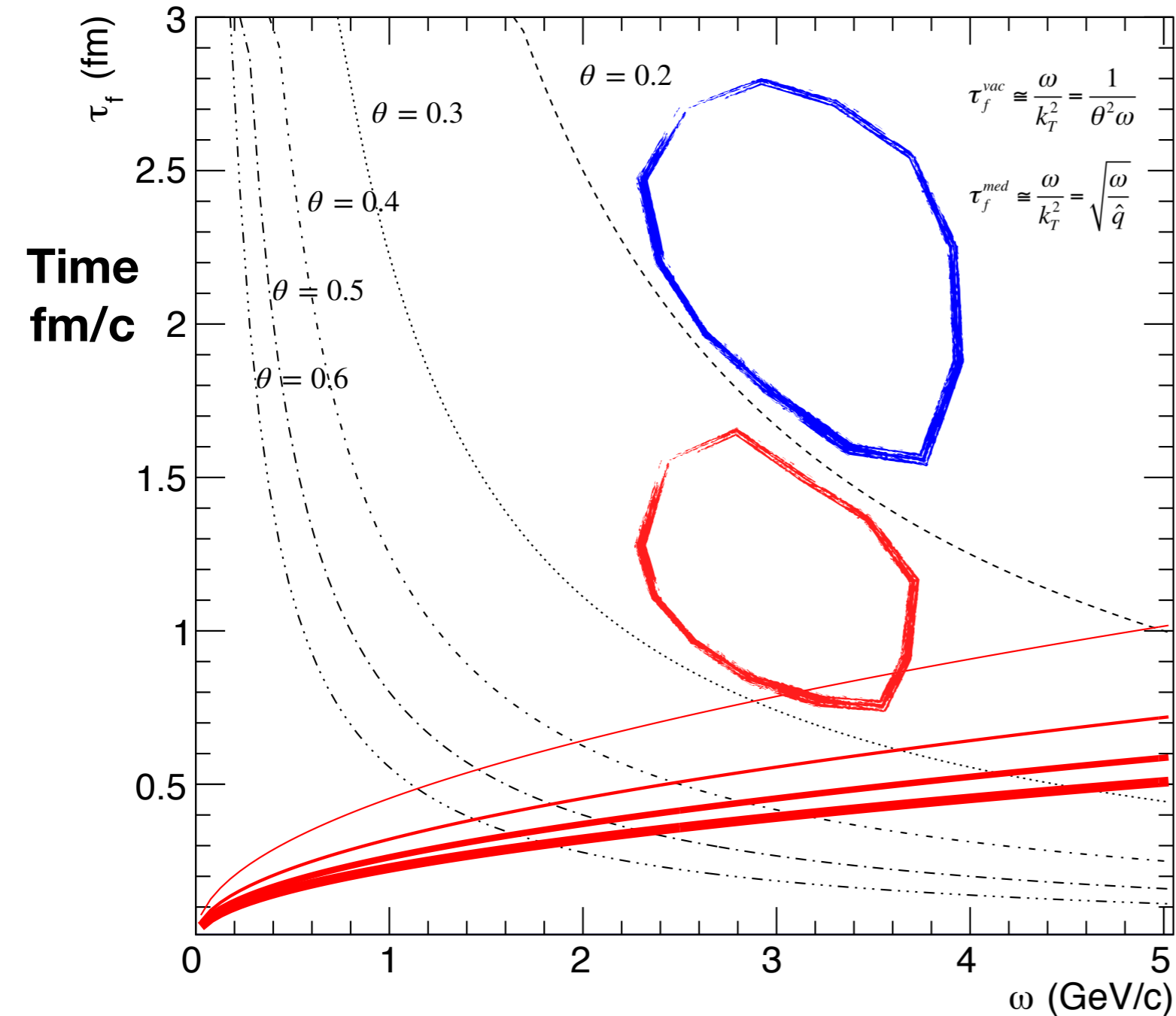
STAR 2109.09793



- Balancing is also independent on opening angle







- Energy loss for these dijets is an experimental observation of soft radiation from a single color charge!
- Potential upper limit on the coherence length  $\lambda_{\perp} \sim \frac{1}{\hat{q}t_f} \leq 0.1$

# 3rd gen jet measurements ongoing!

- Jet substructure enables a **systematic exploration of parton-QGP interactions**
- **Tagging jets** of particular angular scales and studying their **calibrated energy loss** can point us towards quantitative measurements of the QGP's microscopic properties
- Differential measurements further constrain theoretical scenarios and **probe medium at varying resolution scales**
- Consistent picture of energy loss at RHIC for specially selected dijets via **soft gluon emission from a single color charge**

# Backup