

Black Holes in teleparallel gravity

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CENTER OF
APPLIED SPACE TECHNOLOGY
AND MICROGRAVITY



1. Teleparallelism

- Teleparallel Geometry
- Symmetry
- Teleparallel Gravity

2. Black Holes in $f(T, B, \phi)$ teleparallel gravity

- Born-Infeld $f(T)$ -gravity
- Teleparallel perturbations of GR
- Scalar-Torsion gravity

3. Conclusion and Outlook

Geometric fields

Geometric fields

Tetrad components $\theta^a{}_\mu(x)$

An independent flat and metric compatible connection $\Gamma^\sigma{}_{\mu\nu}(x)$

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- technically: basis 1-forms $\theta^a = \theta^a{}_\mu dx^\mu$
- practically: 16 field components $\theta^a{}_\mu$ with inverse $e_a{}^\mu$ $\theta^a{}_\mu e_a{}^\nu = \delta^\nu{}_\mu$, $\theta^a{}_\nu e_b{}^\nu = \delta^a{}_b$
- the metric is a derived object $g_{\mu\nu} = \eta_{ab} \theta^a{}_\mu \theta^b{}_\nu$, $\eta_{ab} = \text{diag}(-, +, +, +)$

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- tetrad basis representation $\Gamma^\mu{}_{\nu\rho} = e_a{}^\mu (\partial_\rho \theta^a{}_\nu + \omega^a{}_{b\rho} \theta^b{}_\nu)$
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Curvature (rotation)

Non-Metricity (changing lengths)

Torsion (non closing)

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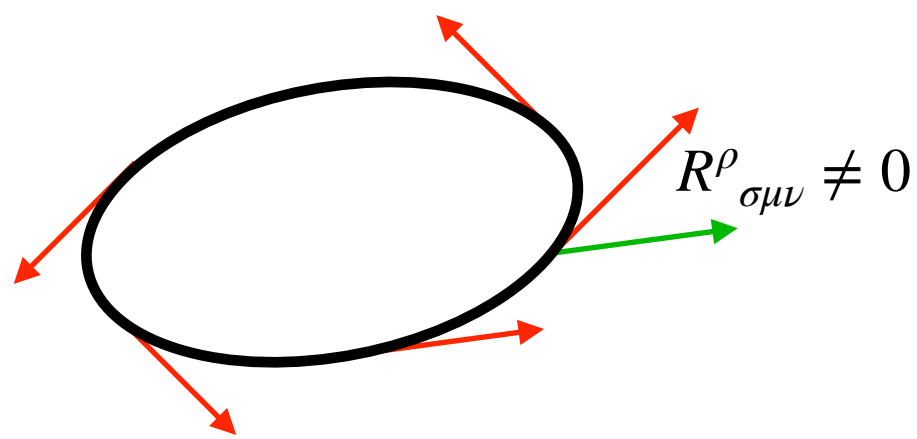
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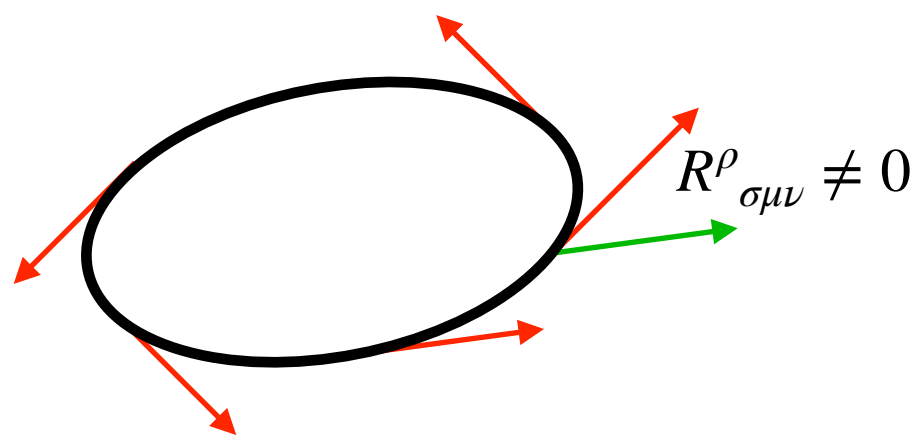
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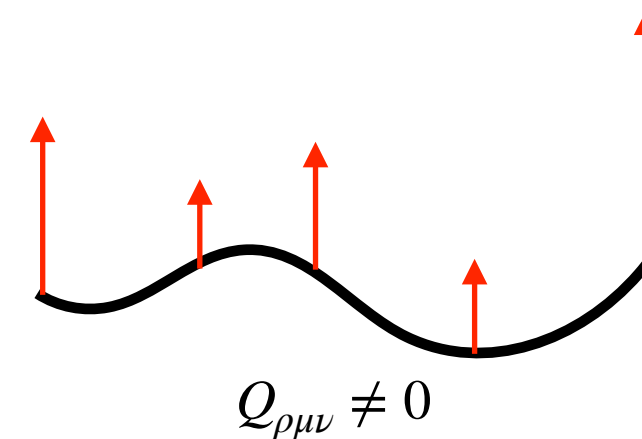
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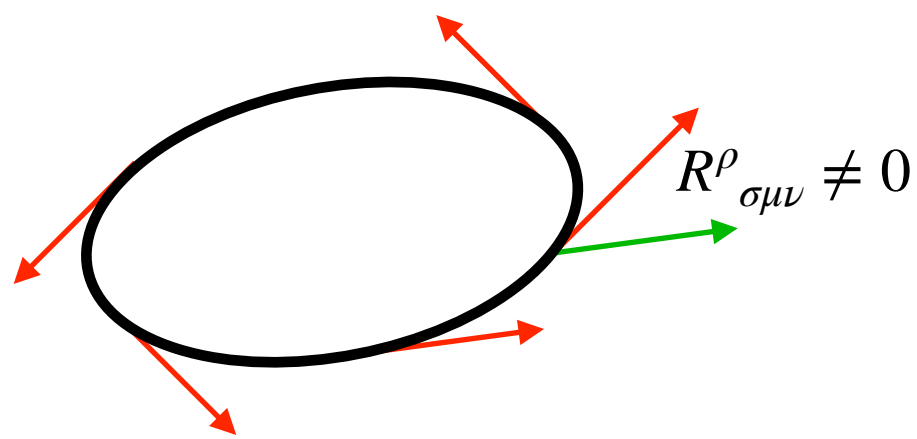
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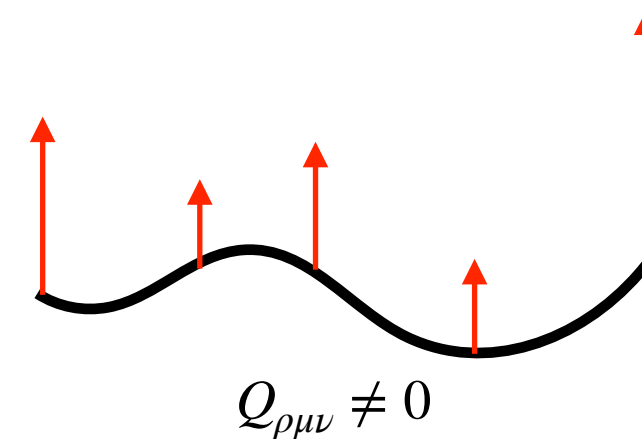
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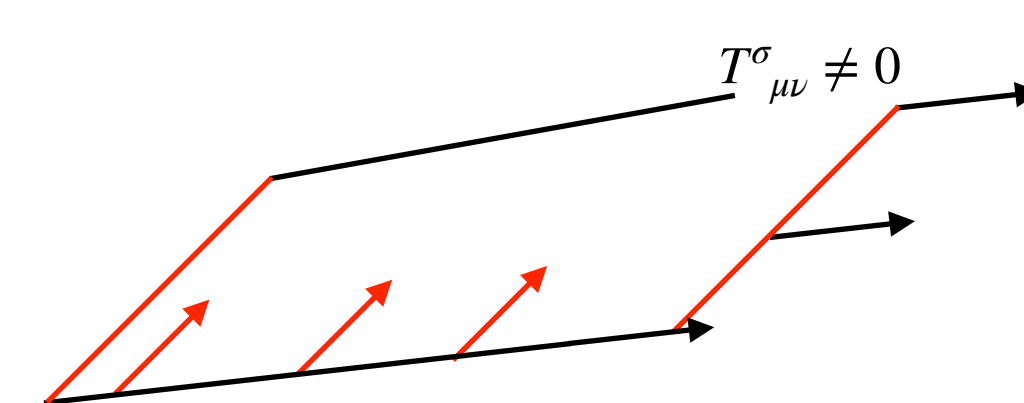
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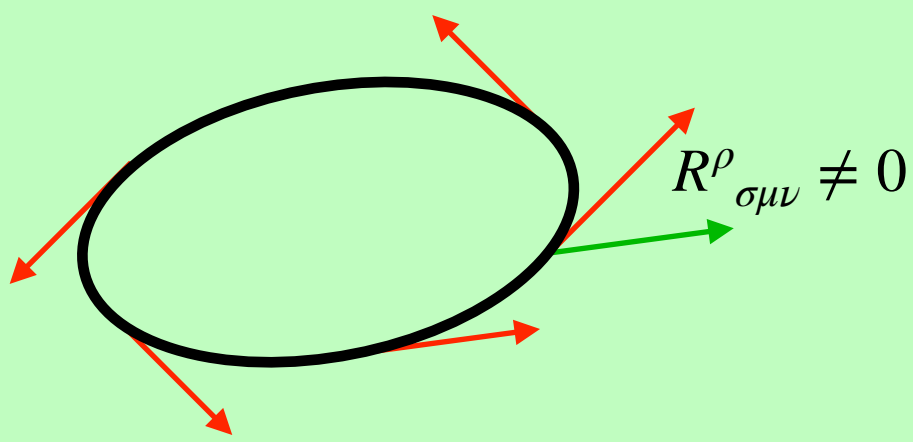
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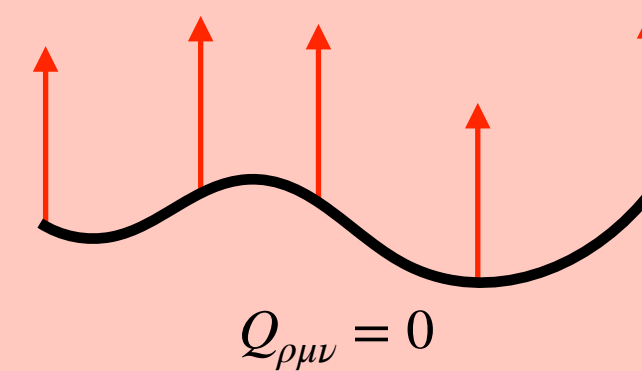
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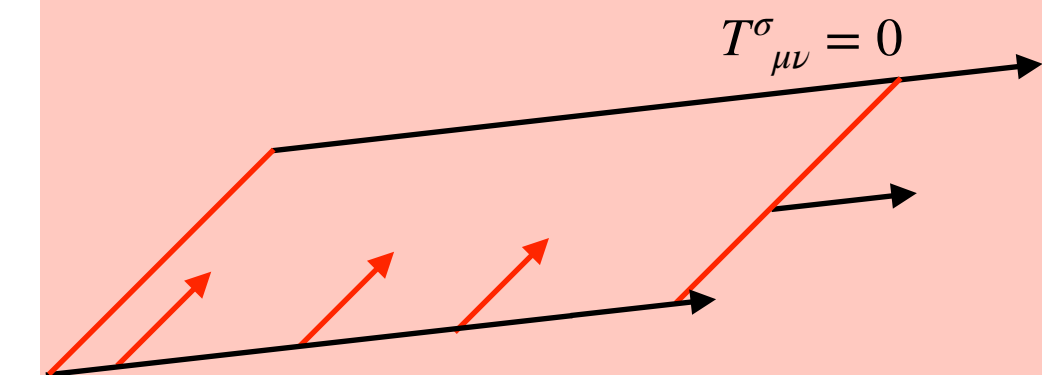
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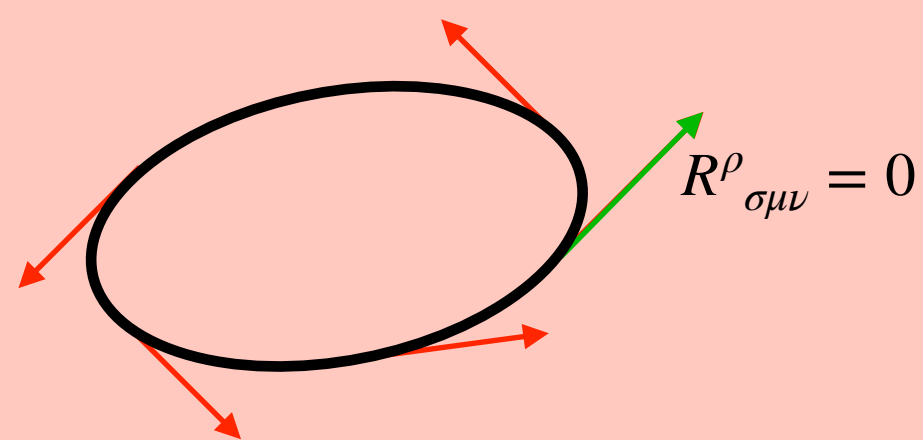
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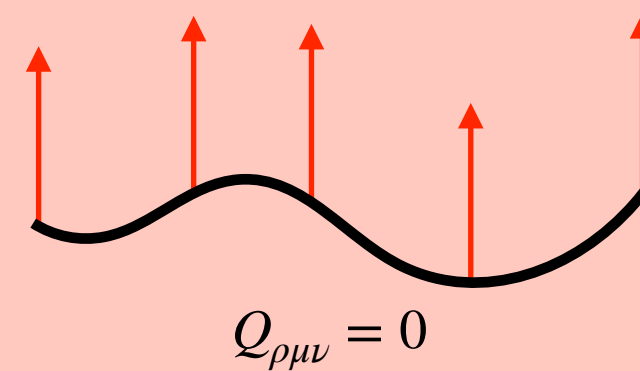
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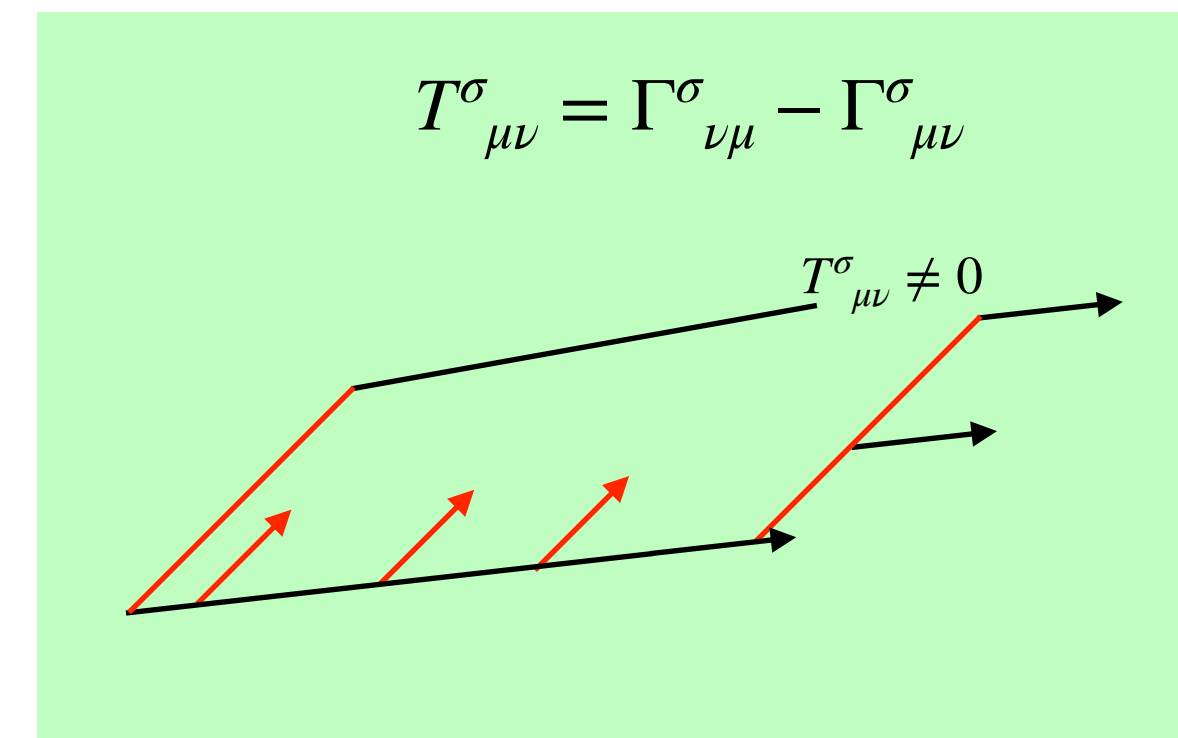
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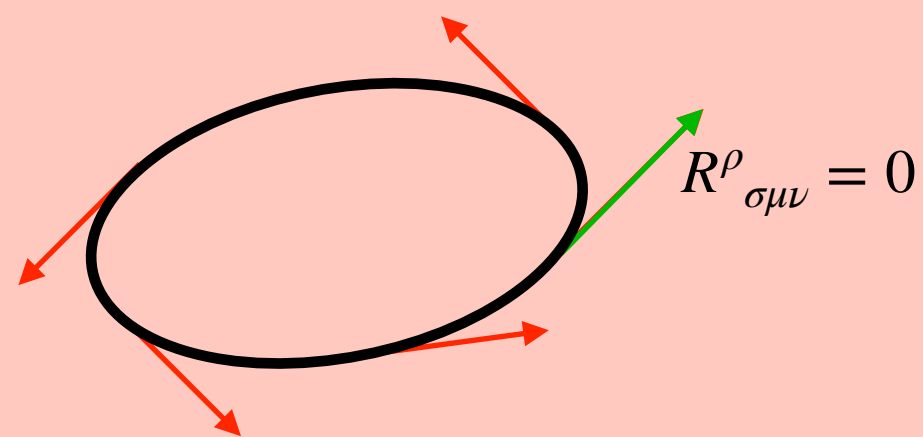
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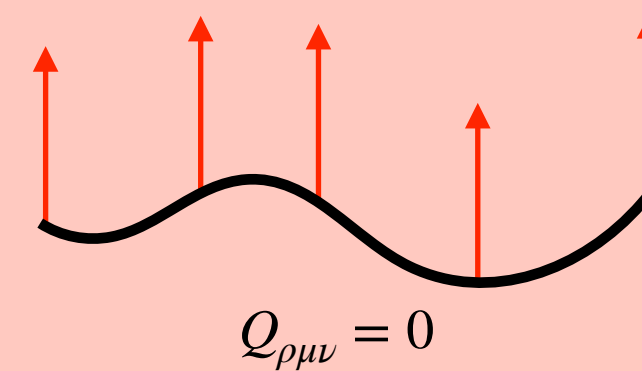
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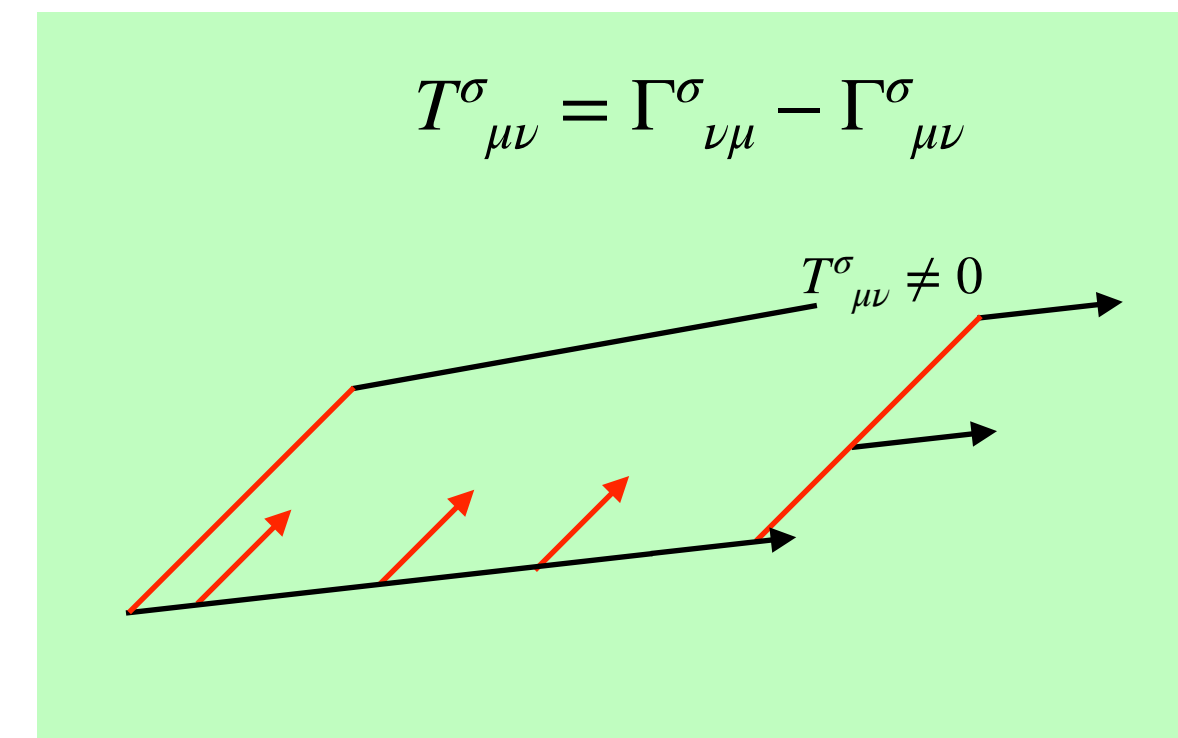
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Without loss of generality one can always work with the pair $(\theta^a{}_\mu, \delta^a{}_b)$, called **Weitzenböck** gauge.

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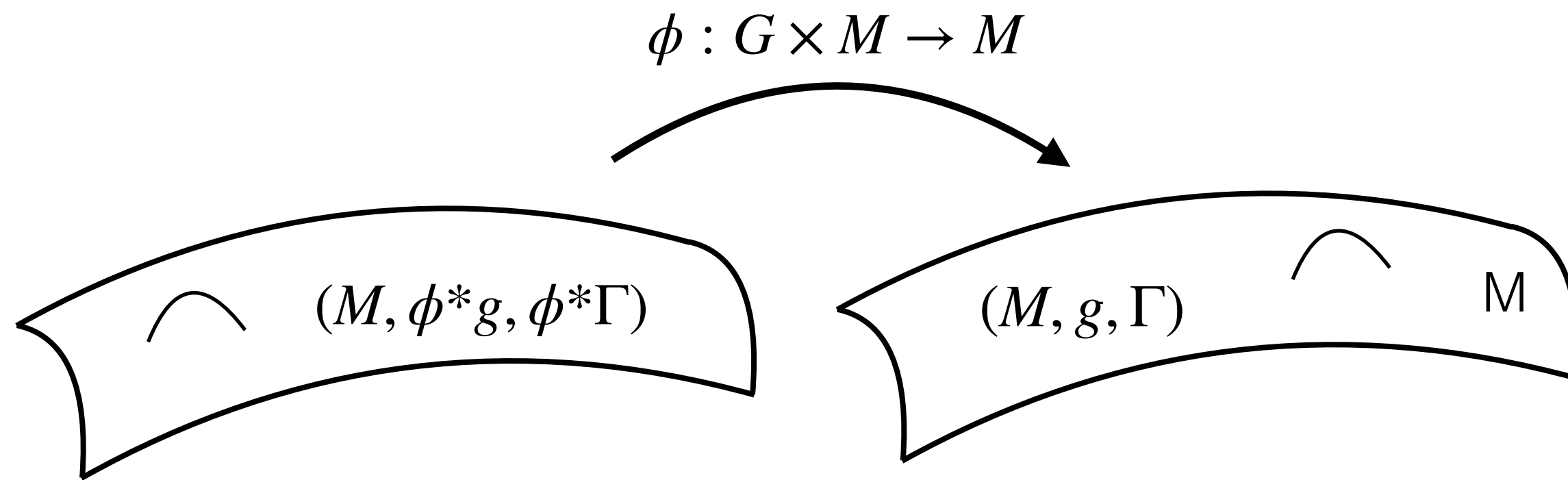
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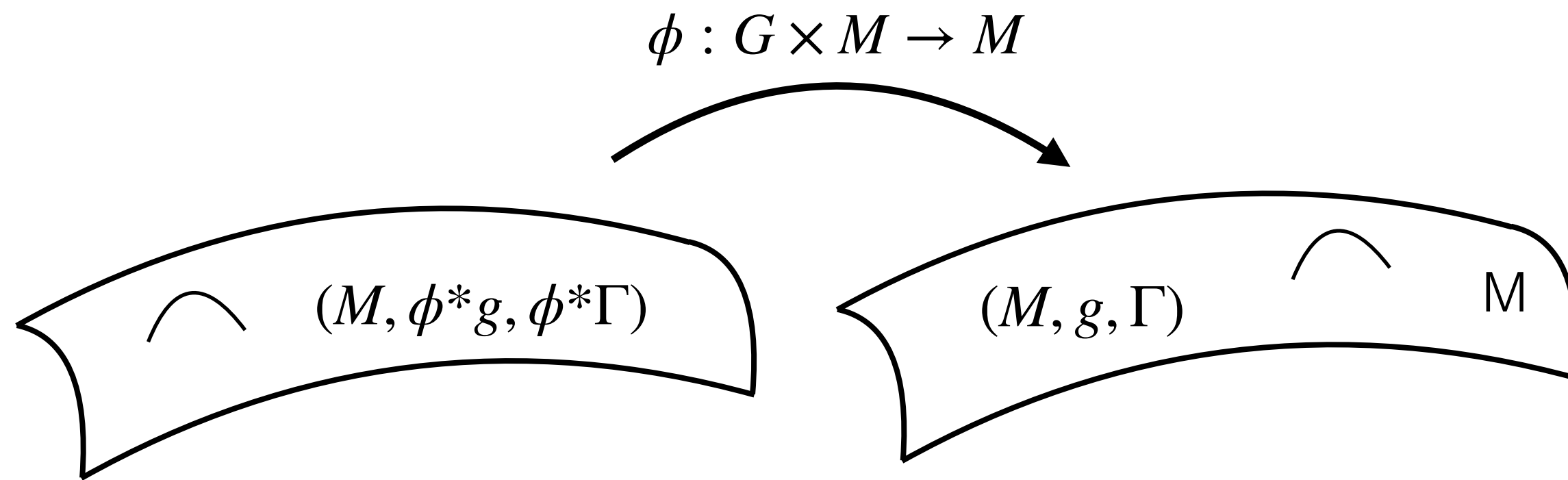
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Symmetries on a manifold M



- ϕ is an action of a Lie Group G as diffeomorphism from M to M
- ϕ is a symmetry of $(M, g, \Gamma) \Leftrightarrow g = \phi^*g, \Gamma = \phi^*\Gamma$
- infinitesimally ϕ is encoded in a vector field $X = X^\sigma \partial_\sigma$

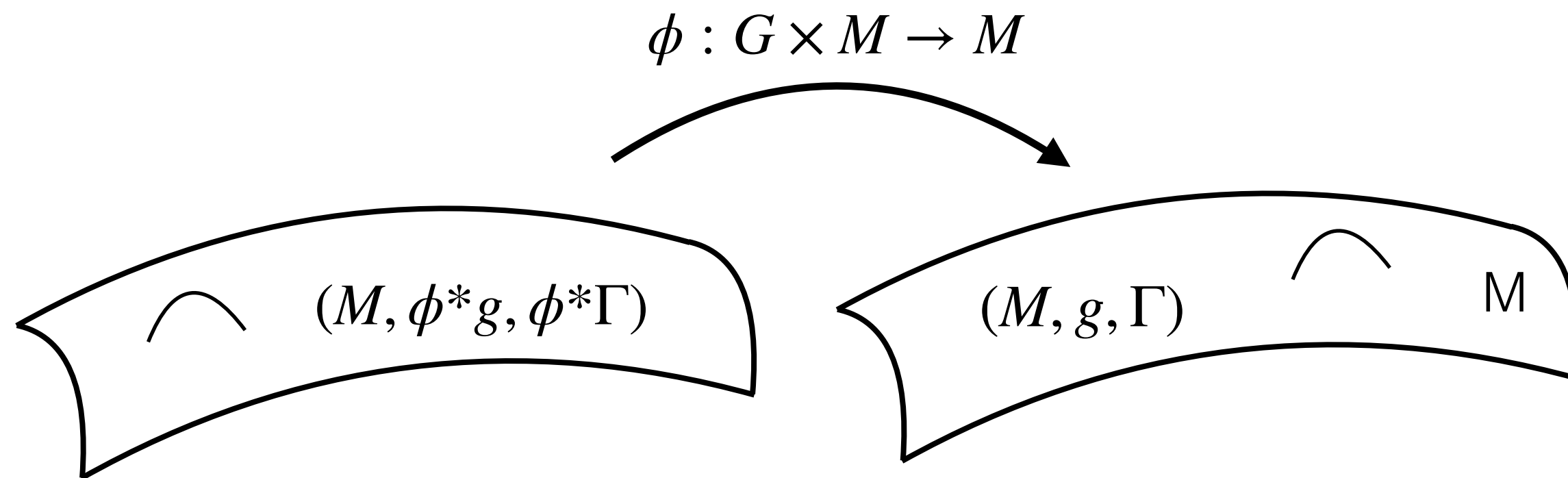
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$$\theta^a{}_\mu(x), \Lambda^a{}_b(x) \Rightarrow \omega^a{}_{b\mu} \text{ and } \Gamma^\rho{}_{\mu\nu}, \text{ Torsion } T^\sigma{}_{\mu\nu} = \Gamma^\sigma{}_{\nu\mu} - \Gamma^\sigma{}_{\mu\nu}$$

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Symmetries on a manifold M



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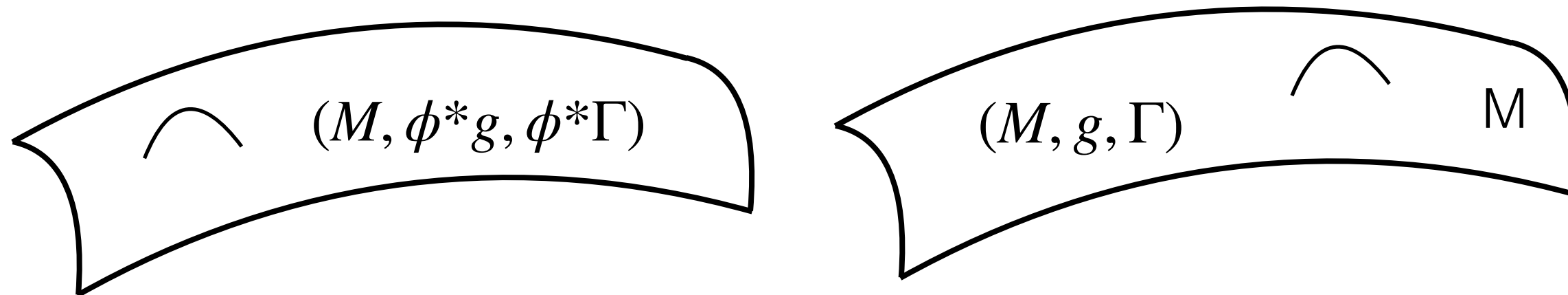
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In teleparallel variables

Geometric fields

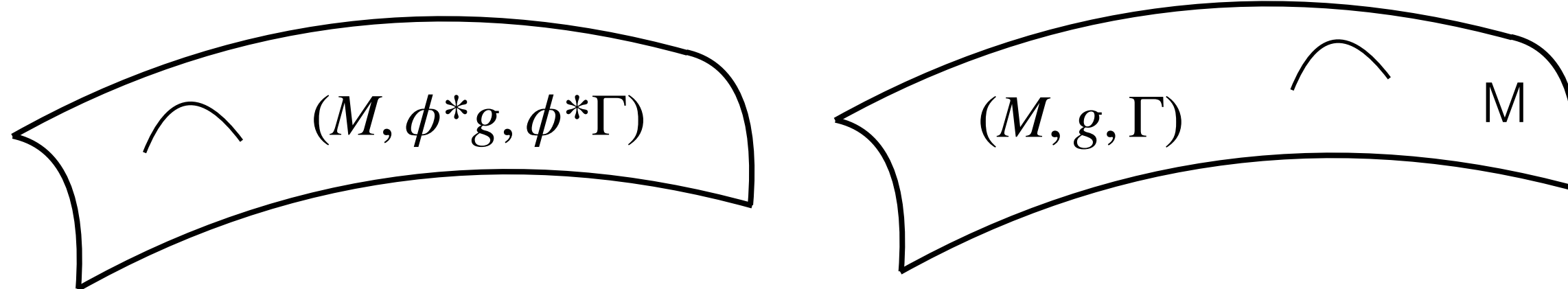
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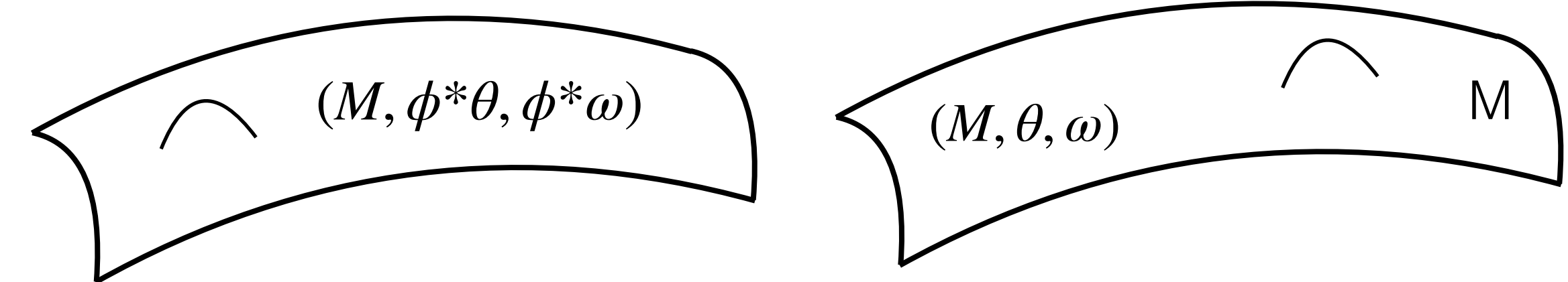
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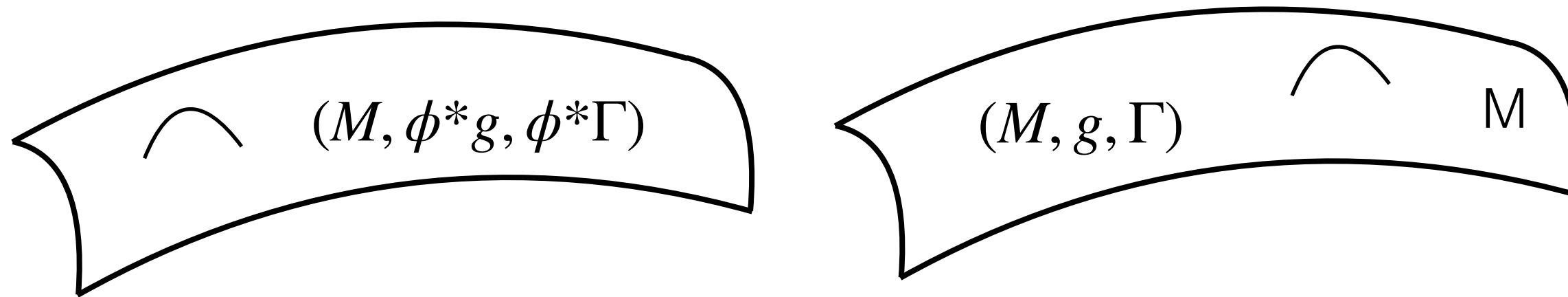
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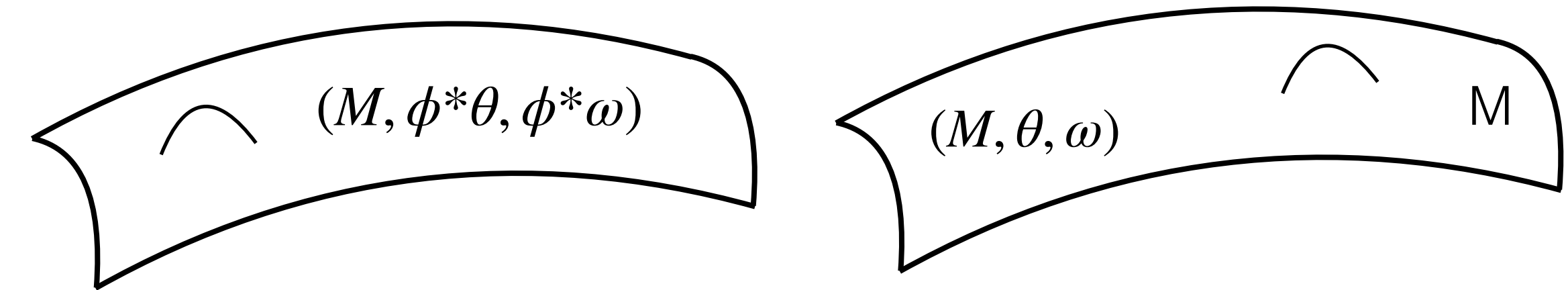
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- $g = \phi^*g, \Gamma = \phi^*\Gamma$ independent of Lorentz frame
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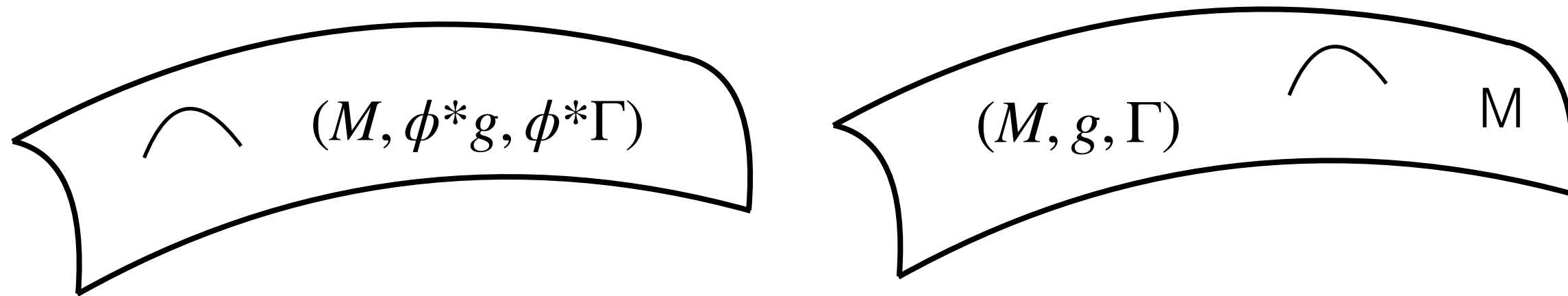
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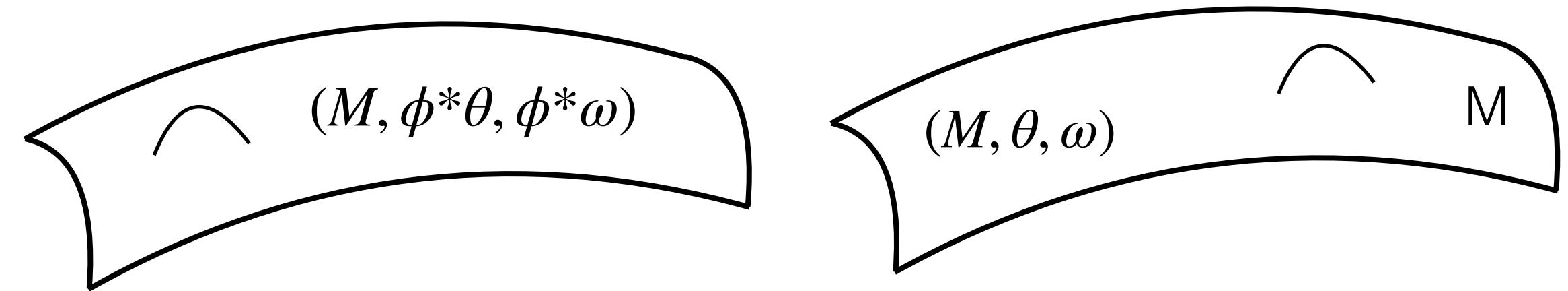
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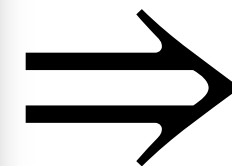


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- $G = SO(3), \mathfrak{g} = \mathfrak{so}(3)$
- $X_z = \partial_\phi,$
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The general spherically symmetric tetrad, $C_i = C_i(t, r)$

$$\theta^a{}_\mu = \begin{pmatrix} C_1 & C_2 & 0 & 0 \\ C_3 \sin \vartheta \cos \varphi & C_4 \sin \vartheta \cos \varphi & C_5 \cos \vartheta \cos \varphi - C_6 \sin \vartheta & -\sin \vartheta (C_5 \sin \varphi + C_6 \cos \vartheta \cos \varphi) \\ C_3 \sin \vartheta \sin \varphi & C_4 \sin \vartheta \sin \varphi & C_5 \cos \vartheta \sin \varphi + C_6 \cos \vartheta & \sin \vartheta (C_5 \cos \varphi - C_6 \cos \vartheta \sin \varphi) \\ C_3 \cos \vartheta & C_4 \cos \vartheta & -C_5 \sin \vartheta & C_6 \sin^2 \vartheta \end{pmatrix}$$

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The standard spherically symmetric metric

$$g = (C_1^2 - C_3^2)dt^2 - (C_4^2 - C_2^2)dr^2 - (C_5^2 + C_6^2)(d\vartheta^2 + \sin^2 \vartheta d\varphi^2) - (C_3 C_4 - C_1 C_2)dt dr$$

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Teleparallel theories of gravity

$$S[\theta] = \int d^4x |\theta| f(T^\sigma{}_{\mu\nu}, \partial T^\sigma{}_{\mu\nu}, \dots)$$

Geometric fields

$$\theta^a{}_\mu(x), \Lambda^a{}_b(x) \Rightarrow \omega^a{}_{b\mu} \text{ and } \Gamma^\rho{}_{\mu\nu}, \text{ Torsion } T^\sigma{}_{\mu\nu} = \Gamma^\sigma{}_{\nu\mu} - \Gamma^\sigma{}_{\mu\nu}$$

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The torsion scalars

$$\mathbf{v}_\mu = T^\rho{}_{\rho\mu}, \quad \mathbf{a}_\mu = \frac{1}{6} \epsilon_{\mu\nu\rho\sigma} T^{\nu\rho\sigma}, \quad \mathbf{t}_{\mu\nu\rho} = T_{(\mu\nu)\rho} + \frac{1}{3} (g_{\rho(\mu} \mathbf{v}_{\nu)} - g_{\nu\mu} \mathbf{v}_\rho)$$

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Most prominent theories in the literature

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$$S[\theta] = \int d^4x |\theta| T = \int d^4x |\theta| (-\dot{R} + B)$$

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- f(T,B,ϕ,X) gravity

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f(T,B,Φ,X)-Gravity and antisymmetric field eqs in spherical symmetry

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- $C_3 = 0 = C_6$ coordinate choices $C_2 = 0, C_5 = \xi r$ ($\xi = \pm 1$)

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$$E_{[\vartheta\varphi]} = 0 \Leftrightarrow C_1 C_6 (f'_T + f'_B) = 0$$

- $C_1 = 0 = C_5$, coordinate choices $C_4 = 0$, $C_6 = \chi r$ ($\chi = \pm 1$)

$$\theta^a{}_{\pm\mu} = \begin{pmatrix} 0 & iC_2 & 0 & 0 \\ iC_3 \sin \vartheta \cos \varphi & 0 & -r\chi \sin \varphi & -r\chi \sin \vartheta \cos \vartheta \cos \varphi \\ iC_3 \sin \vartheta \sin \varphi & 0 & r\chi \cos \varphi & r\chi \sin \vartheta \cos \vartheta \sin \varphi \\ iC_3 \cos \vartheta & 0 & 0 & r\chi \sin^2 \vartheta \end{pmatrix}$$

Geometric fields

$$\theta^a{}_\mu(x), \Lambda^a{}_b(x) \Rightarrow \omega^a{}_{b\mu} \text{ and } \Gamma^\rho{}_{\mu\nu}, \text{ Torsion } T^\sigma{}_{\mu\nu} = \Gamma^\sigma{}_{\nu\mu} - \Gamma^\sigma{}_{\mu\nu}$$

Weitzenböck gauge

$$(\theta^a{}_\mu, \delta^a_b) \Rightarrow \omega^a{}_{b\mu} = 0, \Gamma^\mu{}_{\nu\rho} = e_a{}^\mu \partial_\rho \theta^a{}_\nu,$$

Teleparallel Killing equations:

$$(\mathcal{L}_X \theta)^a{}_\mu = -\lambda^a{}_b \theta^b{}_\mu, \quad (\mathcal{L}_X \omega)^a{}_{b\mu} = \partial_\mu \lambda^a{}_b = 0$$

The torsion scalars

$$\mathbf{v}_\mu = T^\rho{}_{\rho\mu}, \quad \mathbf{a}_\mu = \frac{1}{6} \epsilon_{\mu\nu\rho\sigma} T^{\nu\rho\sigma}, \quad \mathbf{t}_{\mu\nu\rho} = T_{(\mu\nu)\rho} + \frac{1}{3} (g_{\rho(\mu} \mathbf{v}_{\nu)} - g_{\nu\mu} \mathbf{v}_\rho)$$

$$T_{\text{vec}} = \mathbf{v}_\mu \mathbf{v}^\mu, \quad T_{\text{ax}} = \mathbf{a}_\mu \mathbf{a}^\mu, \quad T_{\text{ten}} = \mathbf{t}_{\mu\nu\rho} \mathbf{t}^{\mu\nu\rho}$$

and the torsion scalar

$$\mathbb{T} = -\frac{2}{3} T_{\text{vec}} + \frac{3}{2} T_{\text{ax}} + \frac{2}{3} T_{\text{ten}} = -\dot{R} + \underbrace{\dot{\nabla}_\mu (2T^\sigma{}_\mu)}_B$$

f(T,B,Φ,X)-Gravity and antisymmetric field eqs in spherical symmetry

$$S[\theta] = \int d^4x |\theta| f(T, B, \Phi, X), \quad X = \frac{1}{2} \partial_\mu \Phi \partial^\mu \Phi, \quad T = \text{TEGR}$$

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$$g = A^2 dt^2 - B^2 dr^2 - r^2 (d\vartheta^2 + \sin^2 \vartheta d\varphi^2)$$

Geometric fields

$$\theta^a{}_\mu(x), \Lambda^a{}_b(x) \Rightarrow \omega^a{}_{b\mu} \text{ and } \Gamma^\rho{}_{\mu\nu}, \text{ Torsion } T^\sigma{}_{\mu\nu} = \Gamma^\sigma{}_{\nu\mu} - \Gamma^\sigma{}_{\mu\nu}$$

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- setting $f'_T = f'_B = 0$ yields TEGR or $f(R)$

- $C_3 = 0 = C_6$ coordinate choices $C_2 = 0, C_5 = \xi r$ ($\xi = \pm 1$)

$$\theta^a{}_{\pm\mu} = \begin{pmatrix} C_1 & 0 & 0 & 0 \\ 0 & C_4 \sin \vartheta \cos \varphi & r\xi \cos \vartheta \cos \varphi & -r\xi \sin \vartheta \sin \varphi \\ 0 & C_4 \sin \vartheta \sin \varphi & r\xi \cos \vartheta \sin \varphi & r\xi \sin \vartheta \cos \varphi \\ 0 & C_4 \cos \vartheta & -r\xi \sin \vartheta & 0 \end{pmatrix}$$

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1. Teleparallelism
 - Teleparallel Geometry
 - Symmetry
 - Teleparallel Gravity

2. Black Holes in $f(T, B, \phi)$ teleparallel gravity
 - Born-Infeld $f(T)$ -gravity
 - Teleparallel perturbations of GR
 - Scalar-Torsion gravity

3. Conclusion and Outlook

Geometric fields

$$\theta^a{}_\mu(x), \Lambda^a{}_b(x) \Rightarrow \omega^a{}_{b\mu} \text{ and } \Gamma^\rho{}_{\mu\nu}, \text{ Torsion } T^\sigma{}_{\mu\nu} = \Gamma^\sigma{}_{\nu\mu} - \Gamma^\sigma{}_{\mu\nu}$$

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$$S[\theta] = \int d^4x |\theta| f(T, B, \phi)$$

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$$S[\theta] = \int d^4x |\theta| f(T, B, \phi)$$

The theory

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$$S[\theta] = \int d^4x |\theta| f(T, B, \phi)$$

The theory

$$f = \hat{\lambda} \left(\sqrt{1 + \frac{T}{\hat{\lambda}}} - 1 \right)$$

Geometric fields

$\theta^a{}_\mu(x), \Lambda^a{}_b(x) \Rightarrow \omega^a{}_{b\mu}$ and $\Gamma^\rho{}_{\mu\nu}$, Torsion $T^\sigma{}_{\mu\nu} = \Gamma^\sigma{}_{\nu\mu} - \Gamma^\sigma{}_{\mu\nu}$

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- Spherically symmetric tetrads solving the antisymmetric field equations

Real

$$\theta^a_{1\mu} = \begin{pmatrix} \mathcal{A}(r) & 0 & 0 & 0 \\ 0 & \mathcal{B}(r) \sin \vartheta \cos \varphi & \xi r \cos \vartheta \cos \varphi & -r \xi \sin \vartheta \sin \varphi \\ 0 & \mathcal{B}(r) \sin \vartheta \sin \varphi & \xi r \cos \vartheta \sin \varphi & \xi r \sin \vartheta \cos \varphi \\ 0 & \mathcal{B}(r) \cos \vartheta & -r \xi \sin \vartheta & 0 \end{pmatrix} \quad \xi = \pm 1$$

Complex

$$\theta^a_{2\mu} = \begin{pmatrix} 0 & i\mathcal{B}(r) & 0 & 0 \\ i\mathcal{A}(r) \sin \vartheta \cos \varphi & 0 & -\chi r \sin \varphi & -r \chi \sin \vartheta \cos \vartheta \cos \varphi \\ i\mathcal{A}(r) \sin \vartheta \sin \varphi & 0 & \chi r \cos \varphi & -r \chi \sin \vartheta \cos \vartheta \sin \varphi \\ i\mathcal{A}(r) \cos \vartheta & 0 & 0 & \chi r \sin^2 \vartheta \end{pmatrix} \quad \chi = \pm 1$$

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With metric

$$ds^2 = - \mathcal{A}^2 dt^2 + \mathcal{B}^2 dr^2 + r^2 d\Omega^2$$

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$$\mathcal{A}(r)^2 = 1 - \frac{2M}{r} - \frac{2M}{r\lambda} \mathcal{T}, \quad \mathcal{B}(r)^2 = \frac{r^4 \lambda^4}{16M^4 \mathcal{A}(r)^2} \left(1 + \frac{\lambda^2 r^2}{4M^2} \right)^{-2}$$

$$\mathcal{T} = \tan^{-1} \left(\frac{\lambda r}{2M} \right), \quad \lambda = M\sqrt{\hat{\lambda}}$$

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- Schwarzschild limit for $\lambda \rightarrow \infty$

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- Schwarzschild limit for $\lambda \rightarrow \infty$
- weak field expansion PPN parameters

$$\gamma = -1, \quad \beta - 1 = \frac{8}{(2\lambda - \pi)^2} \\ \Rightarrow \lambda \gtrsim 140$$

Geometric fields

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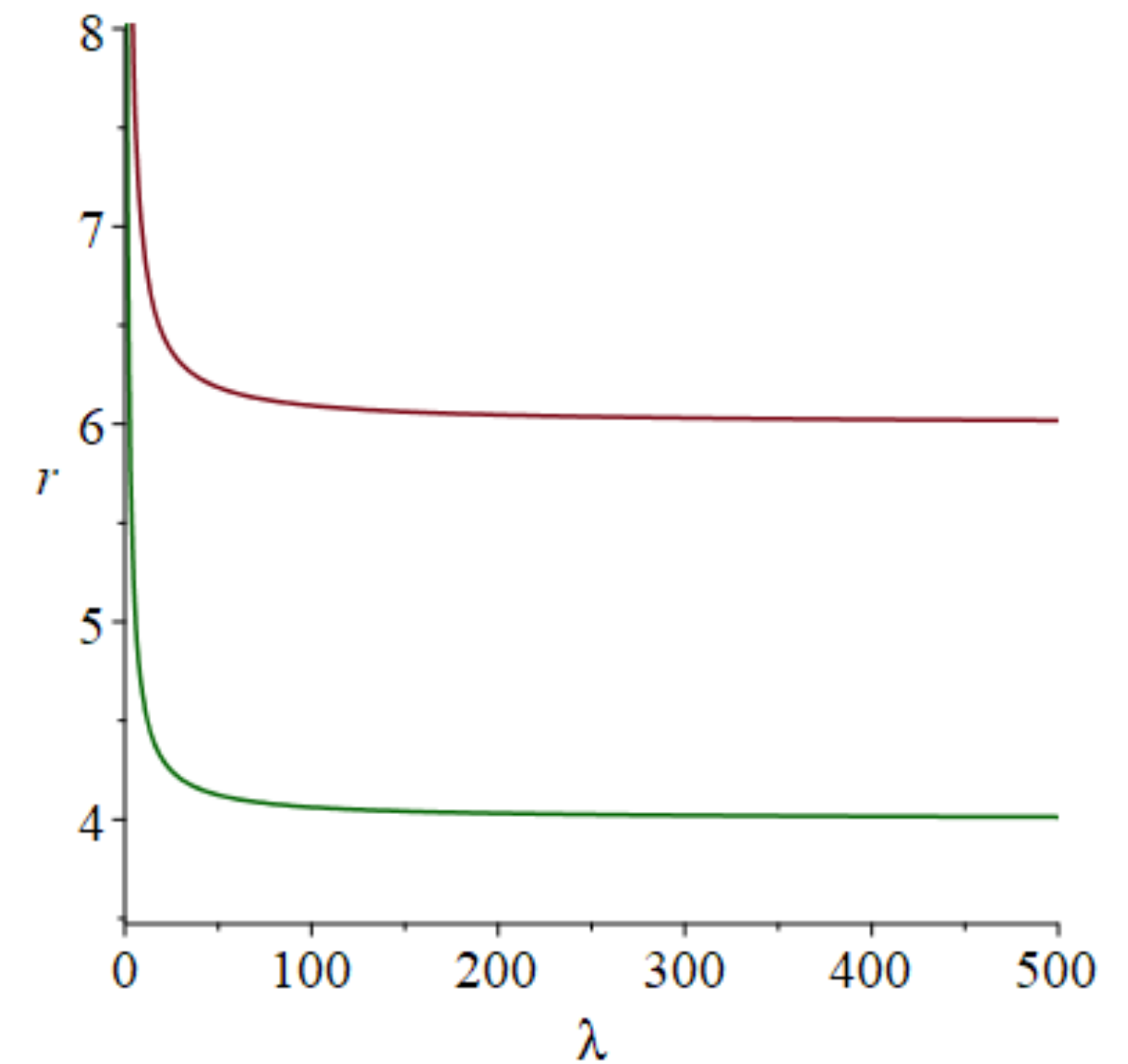
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$$\Rightarrow \lambda \gtrsim 140$$

- marginally stable and marginally bound orbits



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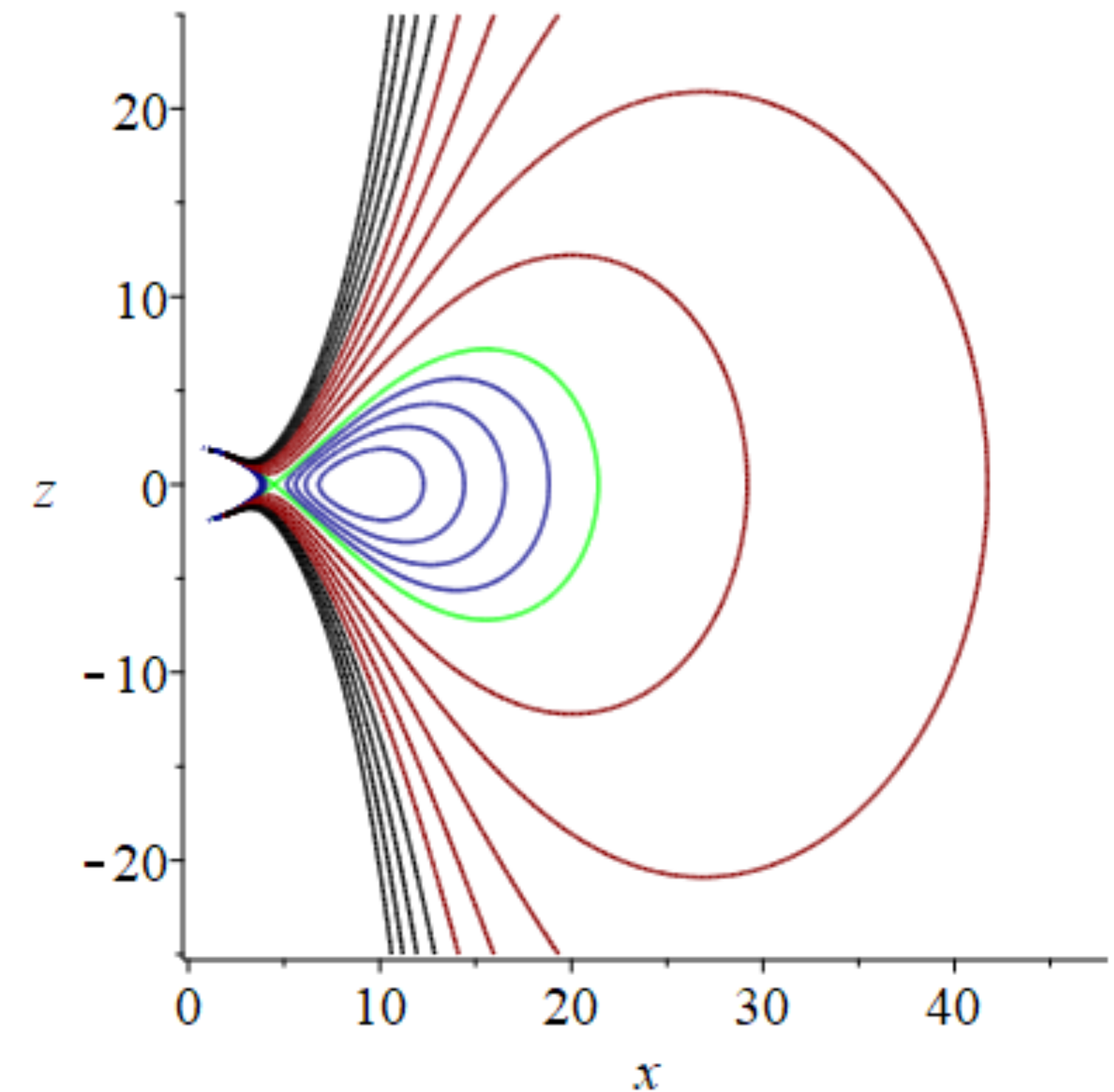
- Schwarzschild limit for $\lambda \rightarrow \infty$
- weak field expansion PPN parameters

$$\gamma = -1, \quad \beta - 1 = \frac{8}{(2\lambda - \pi)^2}$$

$$\Rightarrow \lambda \gtrsim 140$$

- marginally stable and marginally bound orbits
- Equipotential surfaces of Thick Disc models

$$\lambda = 140$$



Geometric fields

$$\theta^a{}_\mu(x), \Lambda^a{}_b(x) \Rightarrow \omega^a{}_{b\mu} \text{ and } \Gamma^\rho{}_{\mu\nu}, \text{ Torsion } T^\sigma{}_{\mu\nu} = \Gamma^\sigma{}_{\nu\mu} - \Gamma^\sigma{}_{\mu\nu}$$

Weitzenböck gauge

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f(T,B,ϕ) gravity

$$S[\theta] = \int d^4x |\theta| f(T, B, \phi)$$

The theory

$$f = \hat{\lambda} \left(\sqrt{1 + \frac{T}{\hat{\lambda}}} - 1 \right)$$

- Spherically symmetric tetrads solving the antisymmetric field equations

Complex

$$\theta^a{}_{2\mu} = \begin{pmatrix} 0 & i\mathcal{B}(r) & 0 & 0 \\ i\mathcal{A}(r)\sin\vartheta\cos\varphi & 0 & -\chi r\sin\varphi & -r\chi\sin\vartheta\cos\vartheta\cos\varphi \\ i\mathcal{A}(r)\sin\vartheta\sin\varphi & 0 & \chi r\cos\varphi & -r\chi\sin\vartheta\cos\vartheta\sin\varphi \\ i\mathcal{A}(r)\cos\vartheta & 0 & 0 & \chi r\sin^2\vartheta \end{pmatrix} \chi = \pm 1$$

With metric

$$ds^2 = -\mathcal{A}^2 dt^2 + \mathcal{B}^2 dr^2 + r^2 d\Omega^2$$

Non-perturbative solution

$$\mathcal{A}(r)^2 = 1 - \frac{2M}{r} - \frac{2M}{r\lambda} \mathcal{T}, \mathcal{B}(r)^2 = \frac{r^4 \lambda^4}{16M^4 \mathcal{A}(r)^2} \left(1 + \frac{\lambda^2 r^2}{4M^2} \right)^{-2}$$

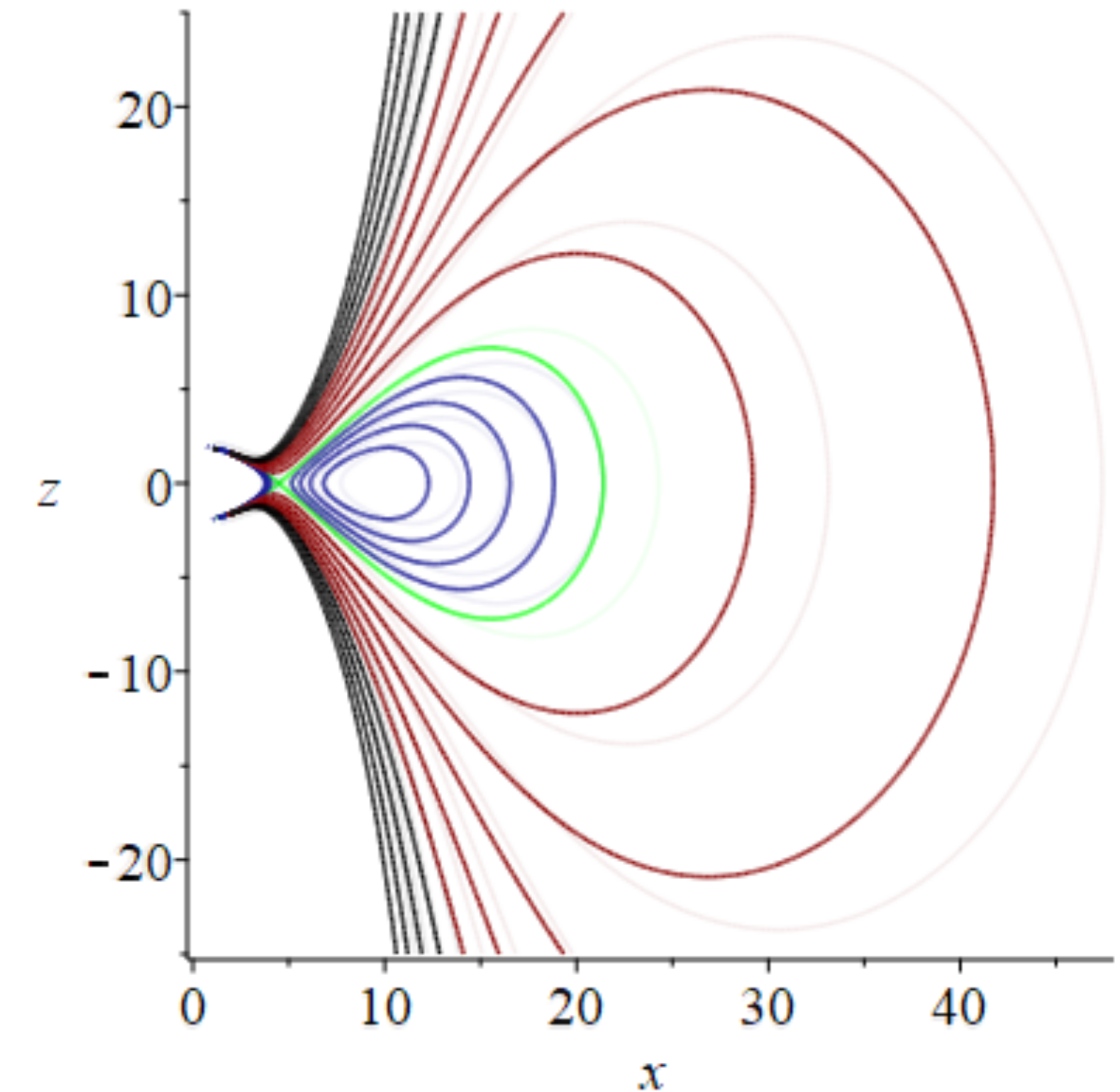
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- Equipotential surfaces of Thick Disc models

$$\lambda = 140 \text{ vs } \lambda = 10$$



Geometric fields

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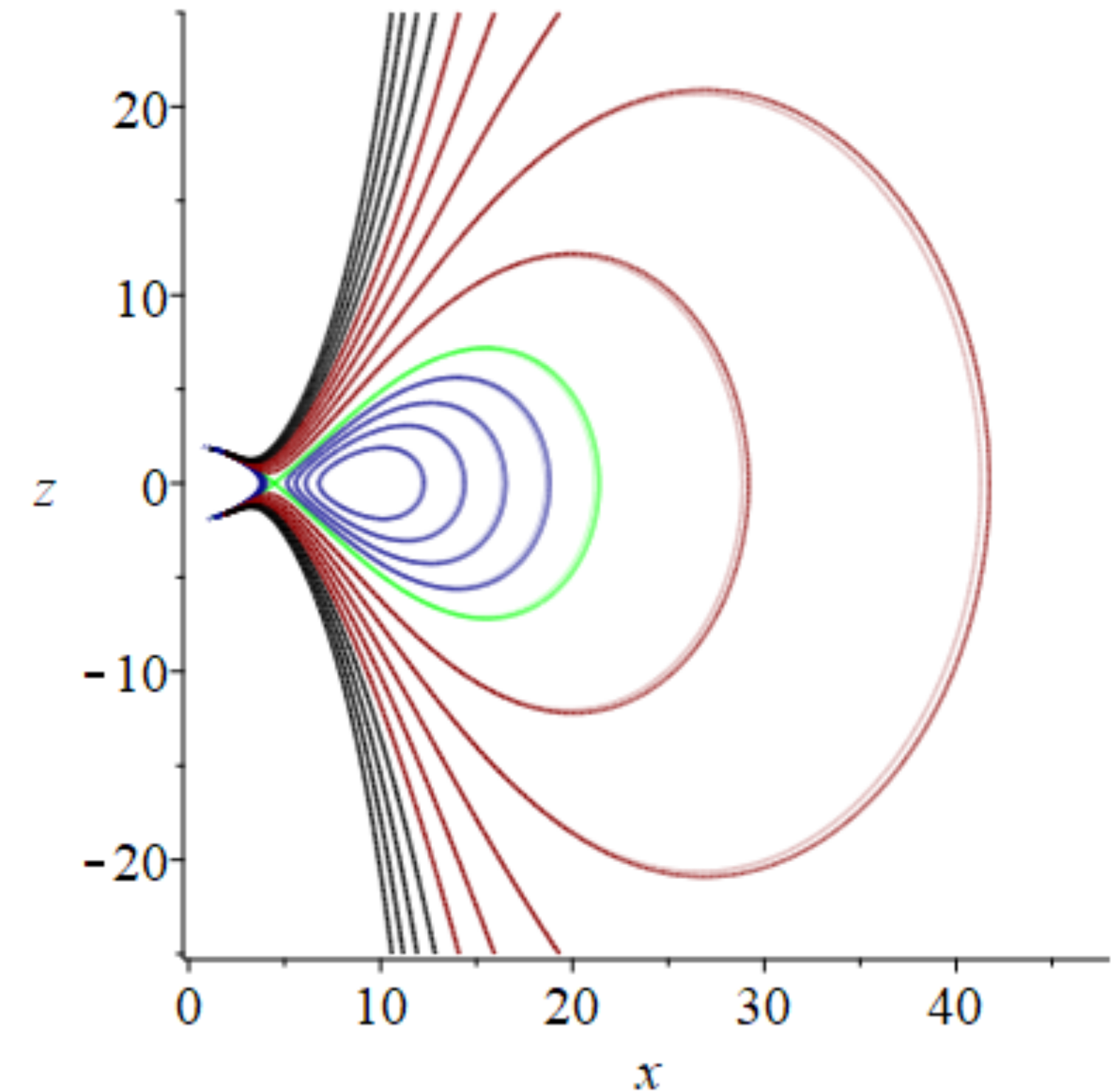
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Geometric fields

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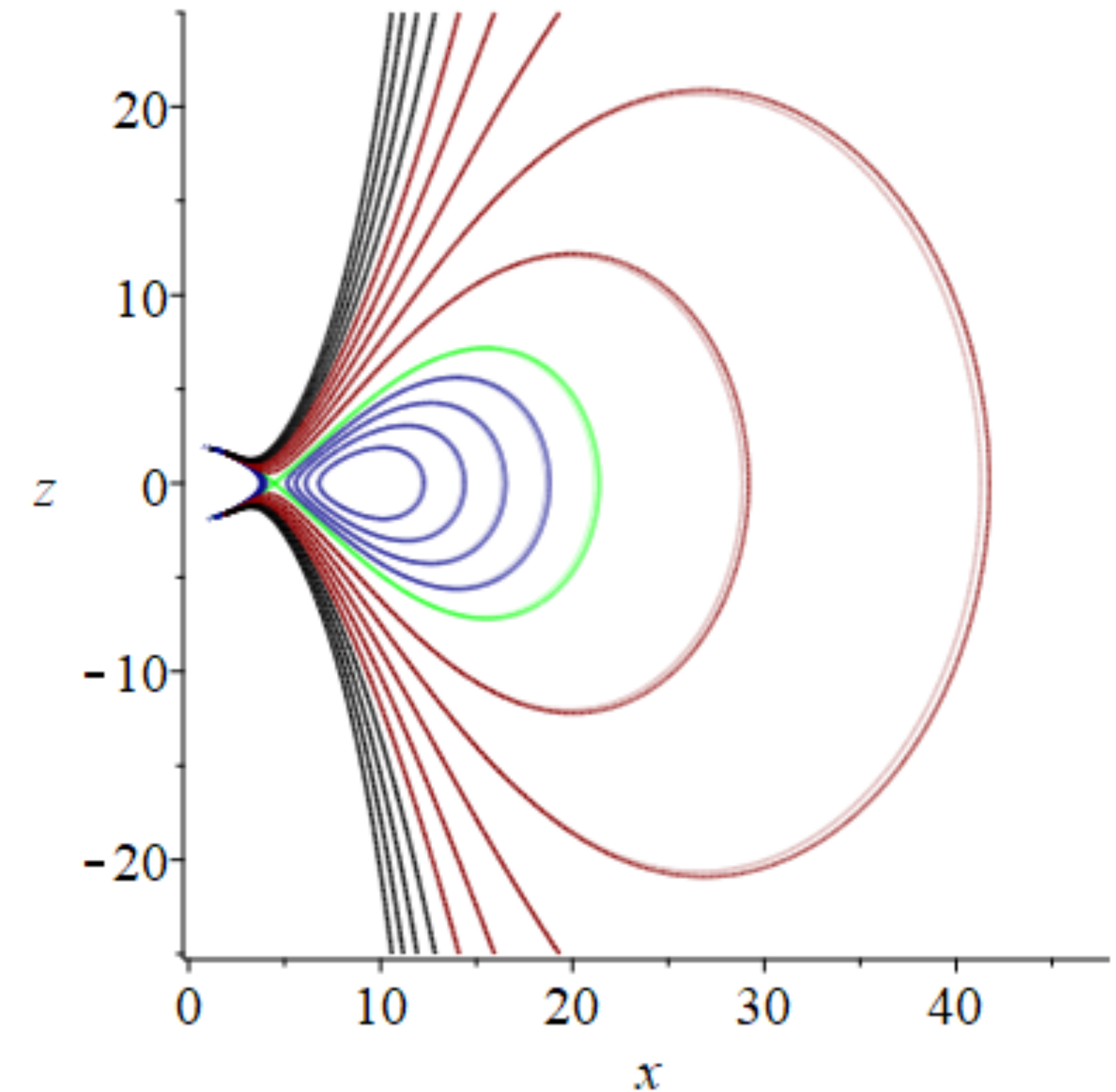
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$$\lambda = 140 \text{ vs } \lambda \rightarrow \infty$$

torus with cusp, tori, bound structures, inner surfaces



Geometric fields

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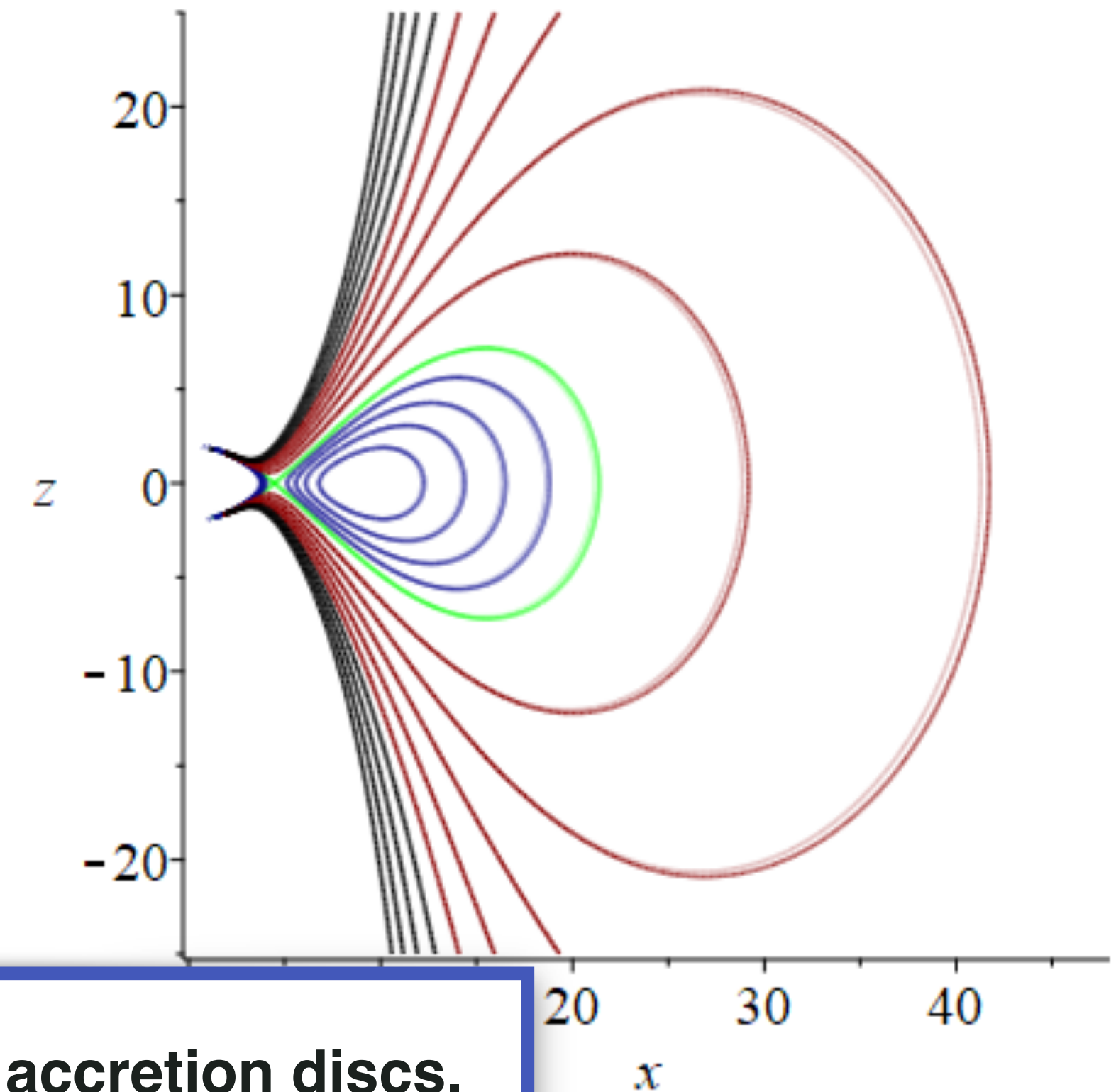
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$$\Rightarrow \lambda \gtrsim 140$$

- marginally stable and marginally bound orbits

- Equipotential surfaces of Thick Disc models

$$\lambda = 140 \text{ vs } \lambda \rightarrow \infty$$



In spherical symmetry: no strong effect of λ on accretion discs.

Geometric fields

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$$S[\theta] = \int d^4x |\theta| f(T, B, \phi)$$

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Complex

$$\theta^a_{2\mu} = \begin{pmatrix} 0 & i\mathcal{B}(r) & 0 & 0 \\ i\mathcal{A}(r) \sin \vartheta \cos \varphi & 0 & -\chi r \sin \vartheta & -r \chi \sin \vartheta \cos \vartheta \cos \varphi \\ i\mathcal{A}(r) \sin \vartheta \sin \varphi & 0 & \chi r \cos \vartheta & -r \chi \sin \vartheta \cos \vartheta \sin \varphi \\ i\mathcal{A}(r) \cos \vartheta & 0 & 0 & \chi r \sin^2 \vartheta \end{pmatrix} \quad \chi = \pm 1$$

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Perturbative solutions

Geometric fields

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Geometric fields

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- exists for the complex and real tetrad

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- complex: independent of χ
- real: dependent on ξ

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Deflection of light

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Real

$$\theta^a_{1\mu} = \begin{pmatrix} \mathcal{A}(r) & 0 & 0 & 0 \\ 0 & \mathcal{B}(r) \sin \vartheta \cos \varphi & \xi r \cos \vartheta \cos \varphi & -r \xi \sin \vartheta \sin \varphi \\ 0 & \mathcal{B}(r) \sin \vartheta \sin \varphi & \xi r \cos \vartheta \sin \varphi & \xi r \sin \vartheta \cos \varphi \\ 0 & \mathcal{B}(r) \cos \vartheta & -r \xi \sin \vartheta & 0 \end{pmatrix} \quad \xi = \pm 1$$

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$$\theta^a_{2\mu} = \begin{pmatrix} 0 & i\mathcal{B}(r) & 0 & 0 \\ i\mathcal{A}(r) \sin \vartheta \cos \varphi & 0 & -\chi r \sin \varphi & -r \chi \sin \vartheta \cos \vartheta \cos \varphi \\ i\mathcal{A}(r) \sin \vartheta \sin \varphi & 0 & \chi r \cos \varphi & -r \chi \sin \vartheta \cos \vartheta \sin \varphi \\ i\mathcal{A}(r) \cos \vartheta & 0 & 0 & \chi r \sin^2 \vartheta \end{pmatrix} \quad \chi = \pm 1$$

With metric

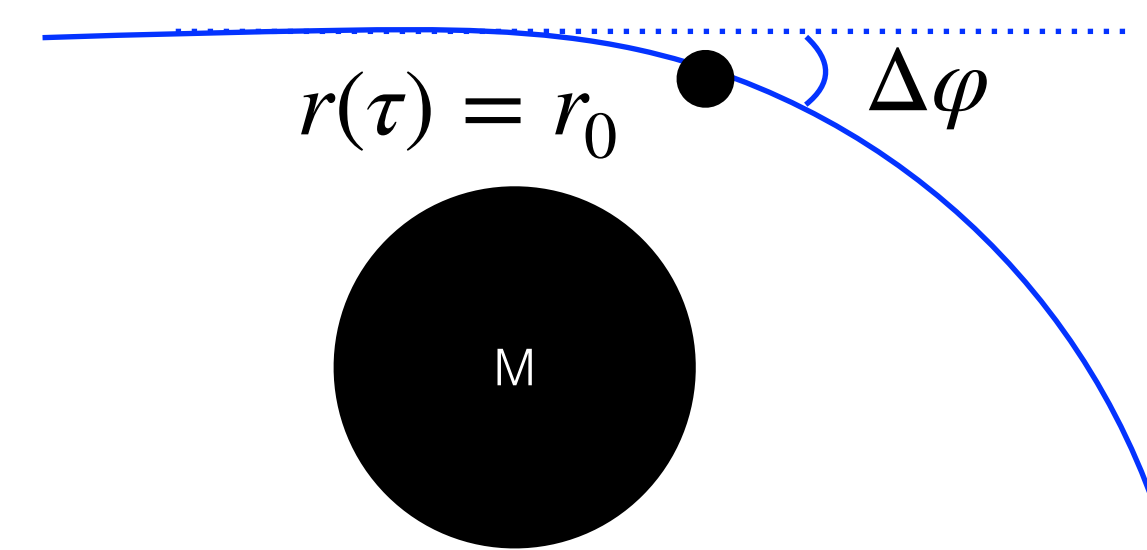
$$ds^2 = -\mathcal{A}^2 dt^2 + \mathcal{B}^2 dr^2 + r^2 d\Omega^2$$

Perturbative solutions

$$\mathcal{A}^2(r) = 1 - \frac{2M}{r} + \epsilon a(r), \mathcal{B}^2(r) = \left(1 - \frac{2M}{r}\right)^{-1} + \epsilon b(r)$$

- exists for the complex and real tetrad
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Deflection of light



$$\Delta\varphi_{\text{real}} = \frac{4M}{r_0} + \epsilon \left(\frac{(\xi - 1)(M(4(44 - 9\pi)\alpha + 8(29 - 6\pi)\beta + 6(34 - 7\pi)\gamma) + \pi r_0(6\alpha + 8\beta + 7\gamma))}{2\xi r_0^3} \right)$$

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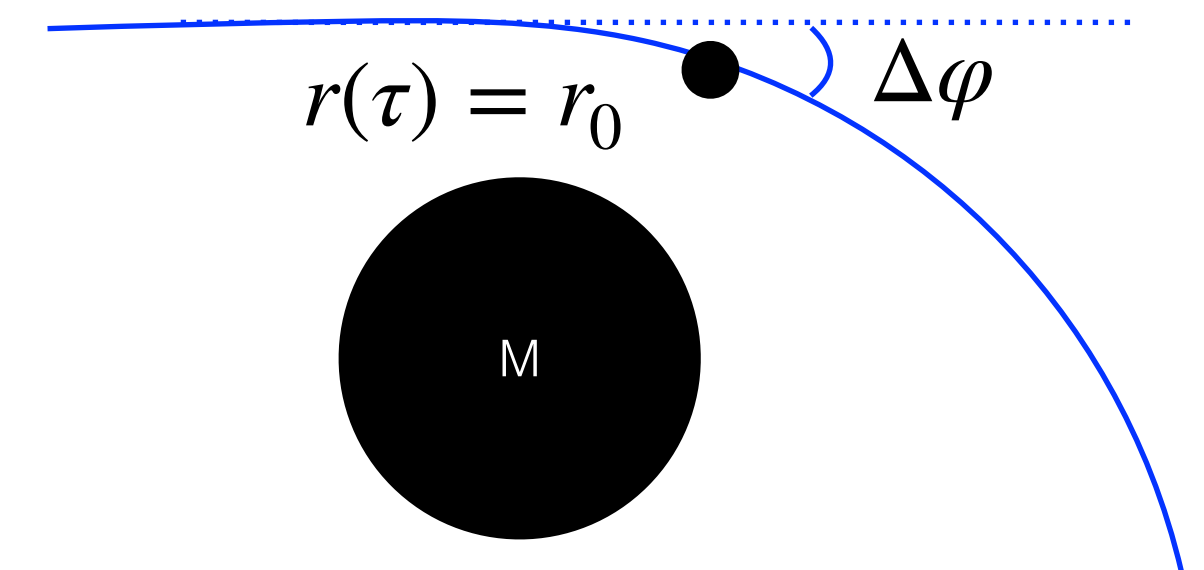
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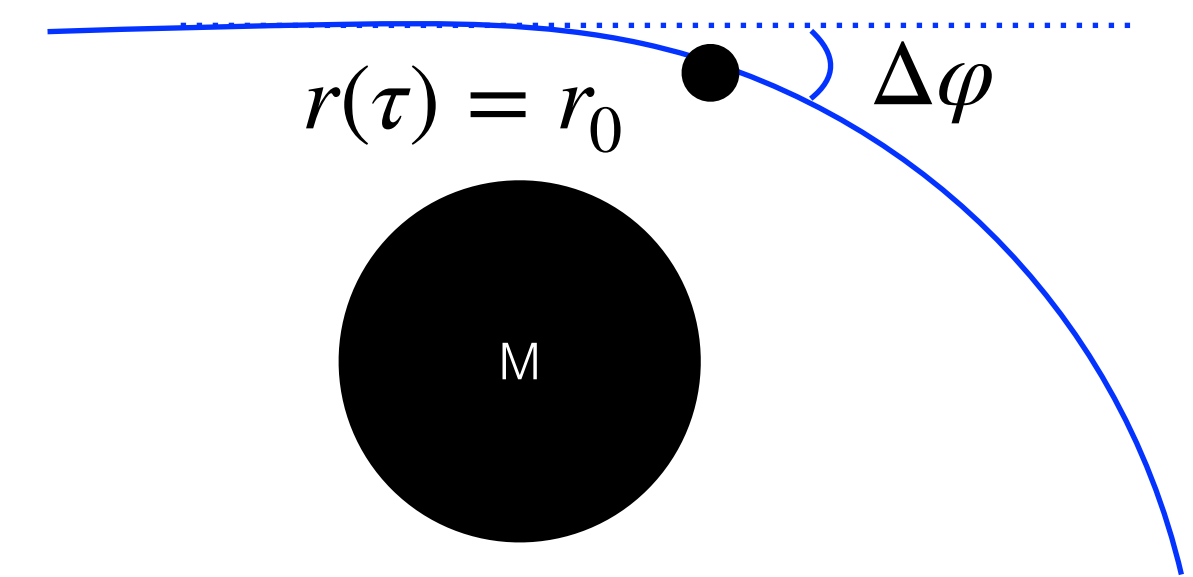
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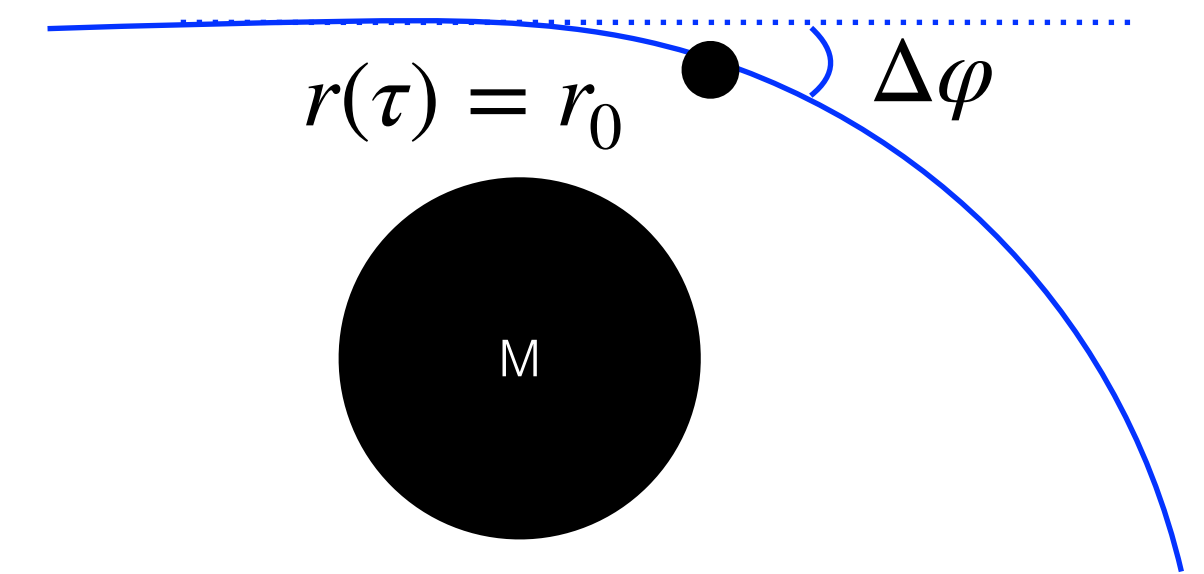
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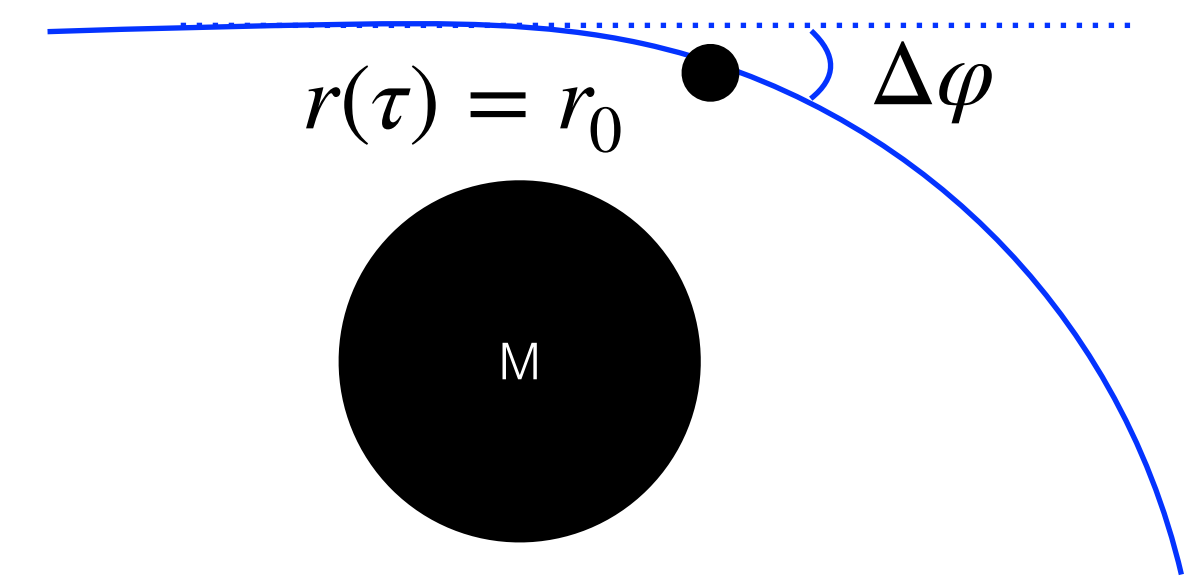
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- real tetrad with $\xi = 1$ least constraint



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or $\frac{(m+1)}{m-1} \frac{1}{\beta} \left(\alpha \dot{R} + \kappa^2 (\psi V - 4V) \right) \leq 0$

1. Teleparallelism

- Teleparallel Geometry
- Symmetry
- Teleparallel Gravity

2. Black Holes in $f(T, B, \phi)$ teleparallel gravity

- Born-Infeld $f(T)$ -gravity
- Teleparallel perturbations of GR
- Scalar-Torsion gravity

3. Conclusion and Outlook

Conclusion and outlook

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Tetrad components $\theta^a{}_\mu(x)$, Lorentz transformations $\Lambda^a{}_b(x)$ generating $\omega^a{}_{b\mu}$ and $\Gamma^\rho{}_{\mu\nu}$, Torsion $T^\sigma{}_{\mu\nu} = \Gamma^\sigma{}_{\nu\mu} - \Gamma^\sigma{}_{\mu\nu}$

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Tetrad components $\theta^a{}_\mu(x)$, Lorentz transformations $\Lambda^a{}_b(x)$ generating $\omega^a{}_{b\mu}$ and $\Gamma^\rho{}_{\mu\nu}$, Torsion $T^\sigma{}_{\mu\nu} = \Gamma^\sigma{}_{\nu\mu} - \Gamma^\sigma{}_{\mu\nu}$

Lorentz transformations as gauge transformations

Affine connection, metric and torsion (in coordinate basis), are invariant, possible choice $(\theta^a{}_\mu, \delta^a_b)$ **Weitzenböck** gauge.

Teleparallel theories of gravity $S[\theta] = \int d^4x |\theta| f(T^\sigma{}_{\mu\nu}, \partial T^\sigma{}_{\mu\nu}, \dots)$

$$f = T, \quad f = f(T, B, \phi), \quad f = \hat{\lambda} \left(\sqrt{1 + \frac{T}{\hat{\lambda}}} - 1 \right), \quad f = f(T, B) = T + \frac{\epsilon}{2} (\alpha T^2 + \beta B^2 + \gamma BT), \quad f(T, B, \phi) = -A(\phi)T + 2\beta X + C(\phi)B - 2\kappa^2 V(\phi)$$

Black Holes

- Born-Infeld gravity: non-perturbative solutions exist, thick accretion disc has been constructed, constraint on $\hat{\lambda}$ from weak field limit

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Black Holes to be continued

- axially symmetric scalarized black holes are under investigation, no teleparallel generalization of Kerr has been found
- impact on accretion discs in axial symmetry?

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Thank you for your attention

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