

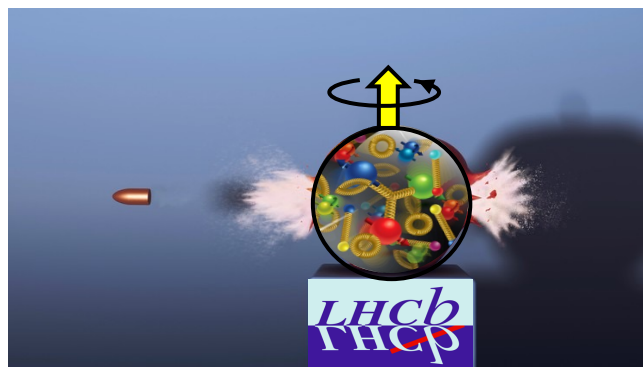


University
of Ferrara



Fixed target experiments at LHC - strong2020 workshop

The LHCspin project



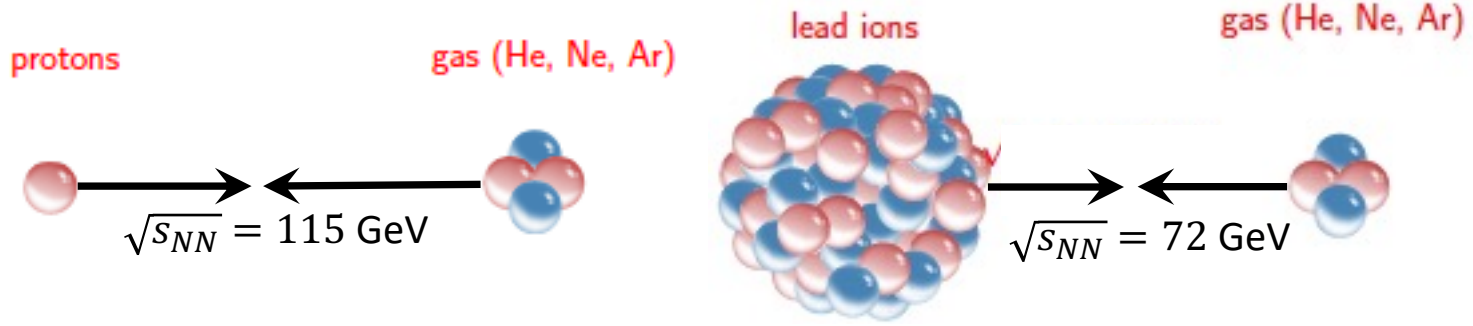
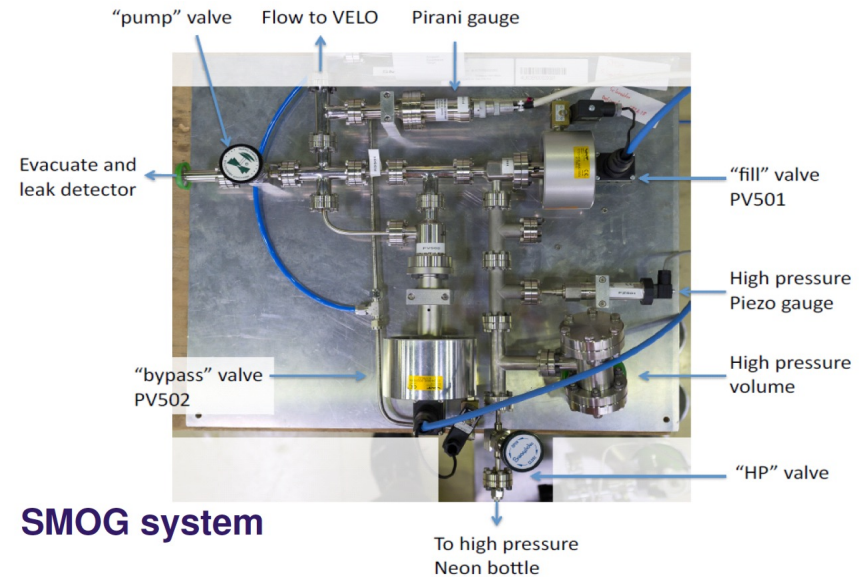
V. Carassiti¹
G. Ciullo¹
P. Di Nezza²
P. Lenisa¹
L. L. Pappalardo¹
M. Santimaria²
E. Steffens³

¹ University of Ferrara and INFN, ² INFN - [Laboratori Nazionali di Frascati](#), ³ University of Erlangen

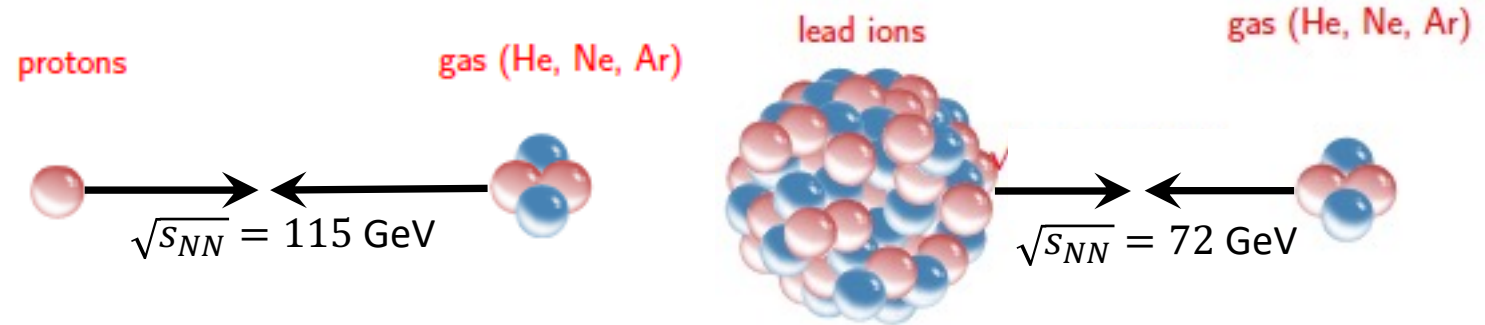
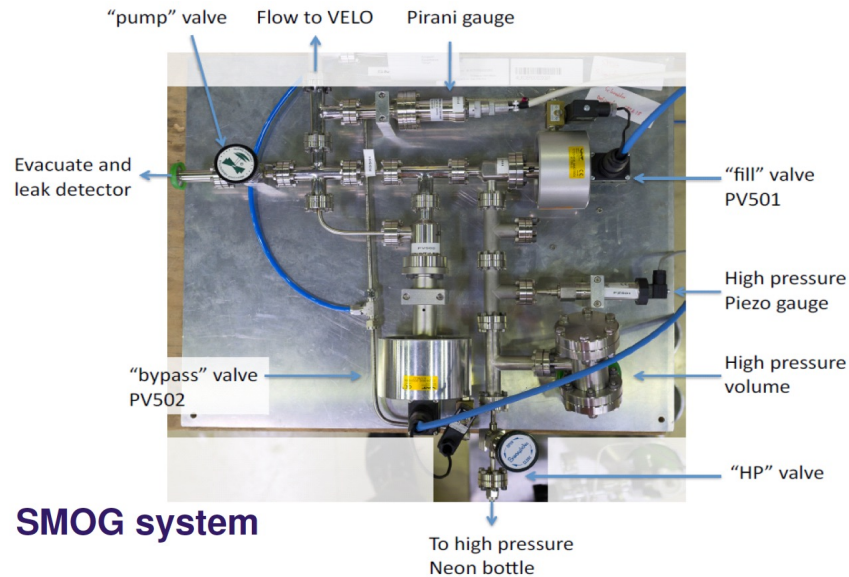
In collaboration with:

R.Engels (fz-juelich), J.Depner (Erlangen), K.Grigoryev (fz-juelich),
S. Mariani (INFN-FI), A.Nass (fz-juelich), F.Rathmann (fz-juelich),
D.Reggiani (PSI-Zurich), M. Statera A.Vasilyev (Gatchina),

Fixed-Target collisions at LHCb



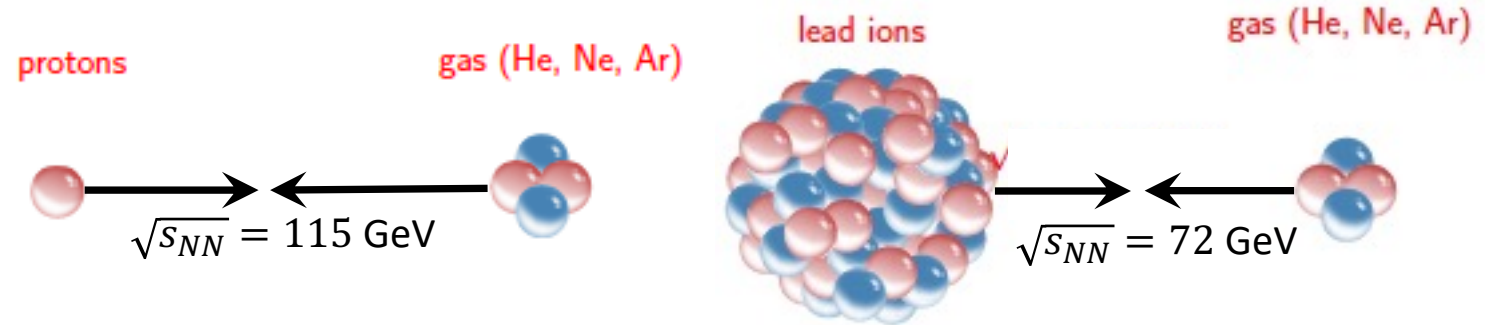
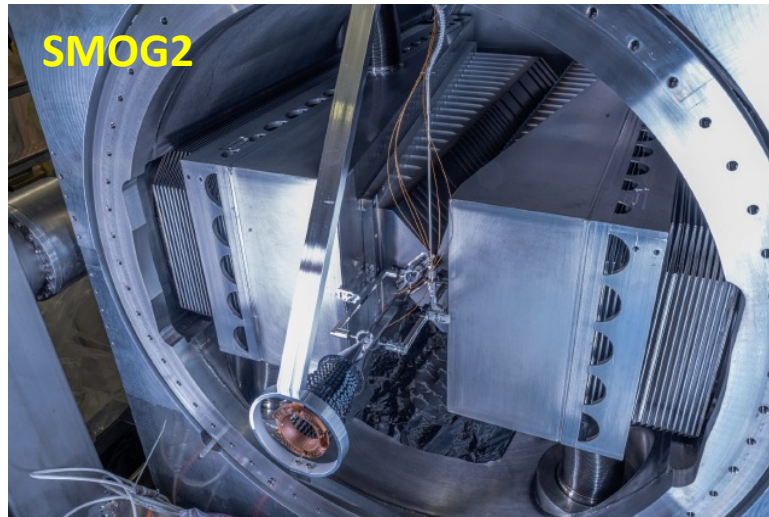
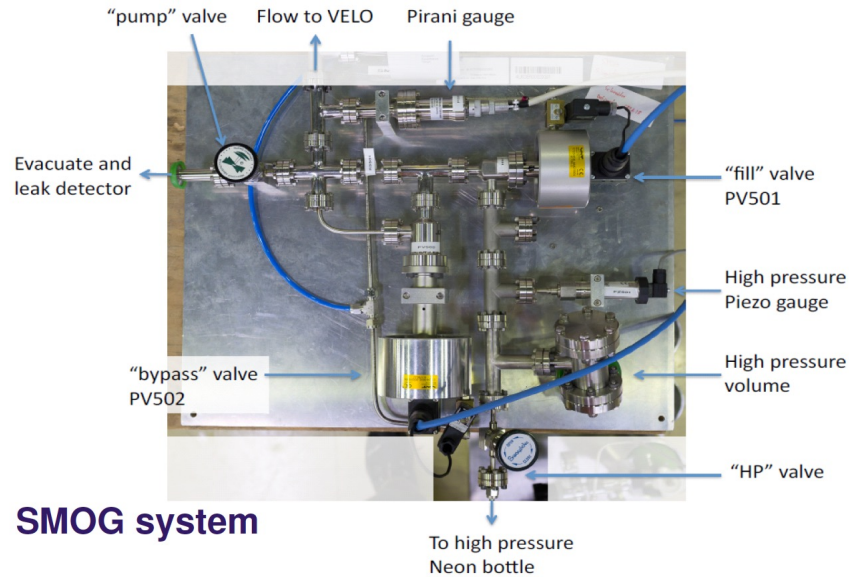
Fixed-Target collisions at LHCb



Already lots of interesting analyses with SMOG:

- ✓ Charm production in pHe and pAr
- ✓ **Charm production in pNe and PbNe (→ Emilie, Frederic)**
- ✓ **Prompt and detached antiproton production in pHe (→ Saverio)**
- ✓ Λ_c polarization in pNe
- ✓ Strangeness enhancement in PbNe vs pNe
- ✓ Cold Nuclear Matter effects in light-hadrons production in p-gas
- ✓ ...

Fixed-Target collisions at LHCb



Already lots of interesting analyses with SMOG:

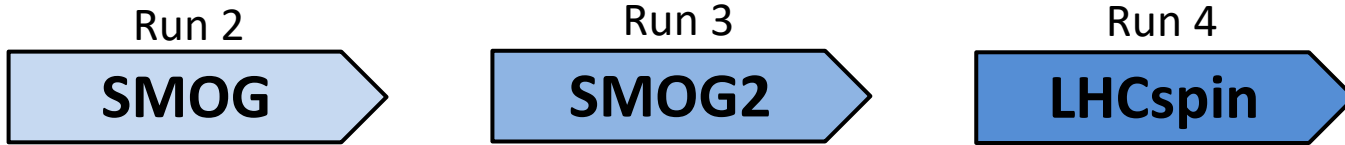
- ✓ Charm production in pHe and pAr
- ✓ **Charm production in pNe and PbNe (→ Emilie, Frederic)**
- ✓ **Prompt and detached antiproton production in pHe (→ Saverio)**
- ✓ Λ_c polarization in pNe
- ✓ Strangeness enhancement in PbNe vs pNe
- ✓ Cold Nuclear Matter effects in light-hadrons production in p-gas
- ✓ ...

...and many more to come with SMOG2! (→ Edoardo)

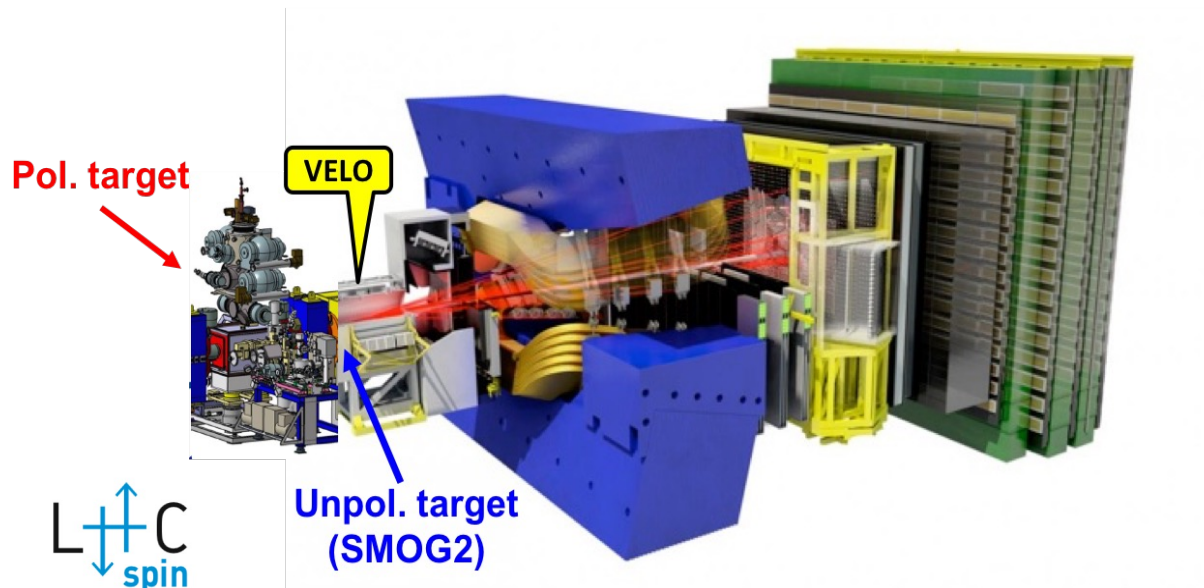
The LHCspin project



The SMOG2 realization sets the basis for the development of a future **polarized gas target for LHCb**



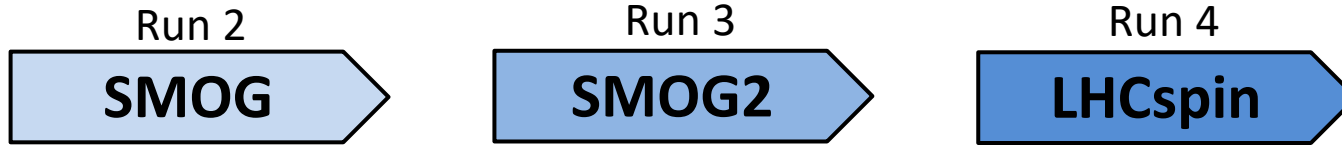
The **LHCspin** project aims to bring spin physics at the LHC through the implementation of a new-generation **polarized gaseous fixed target** in the **LHCb** spectrometer.



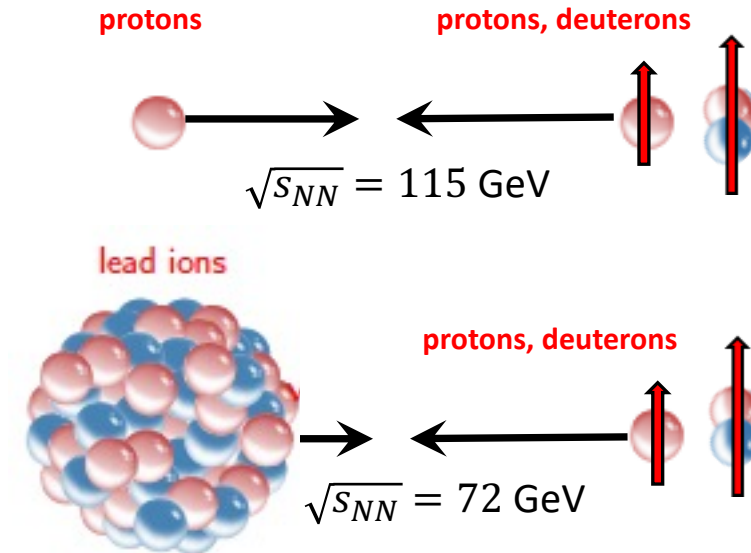
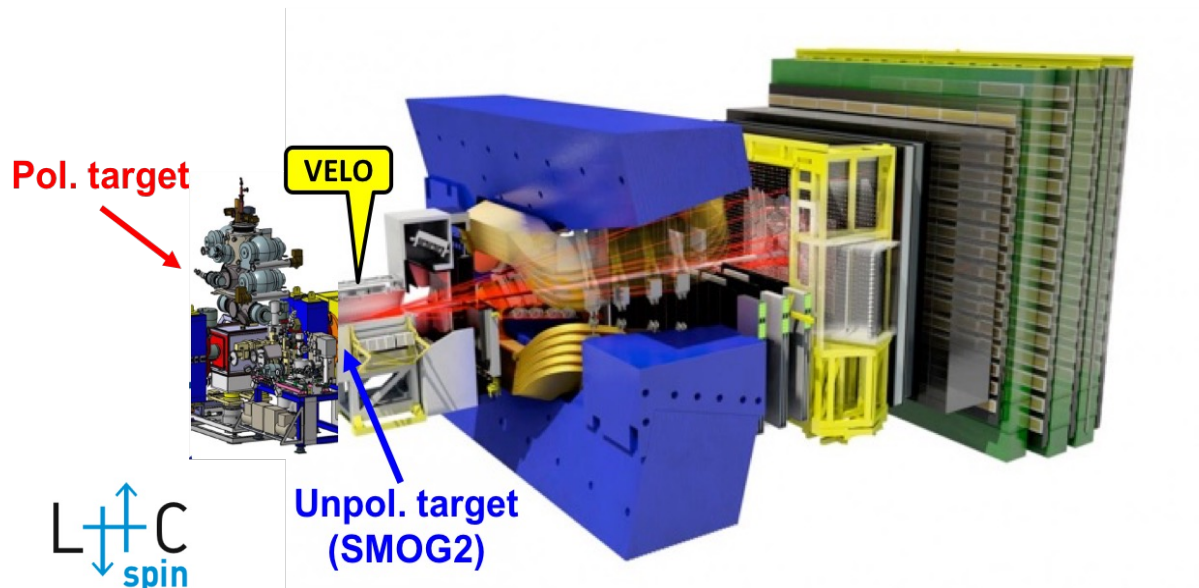
The LHCspin project



The SMOG2 realization sets the basis for the development of a future **polarized gas target for LHCb**



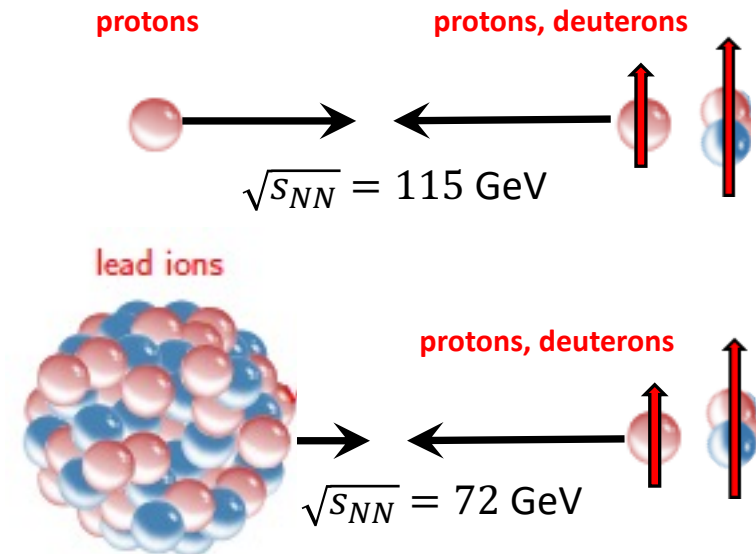
The **LHCspin** project aims to bring spin physics at the LHC through the implementation of a new-generation **polarized gaseous fixed target** in the **LHCb** spectrometer.



The LHCspin project



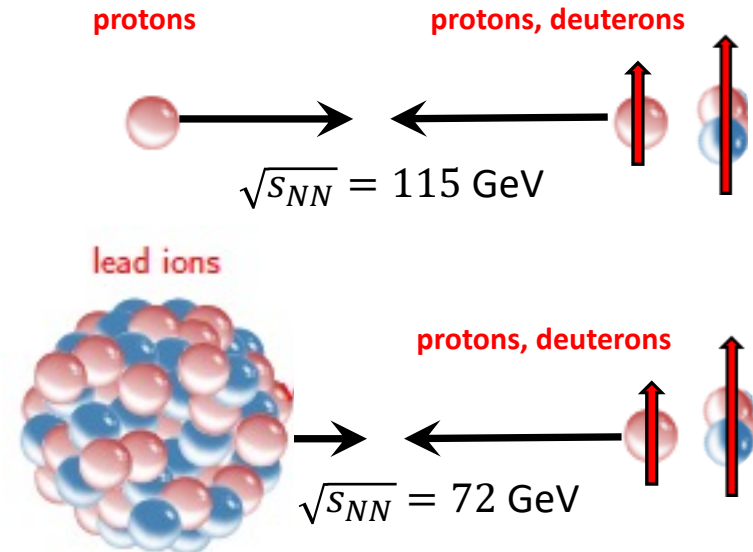
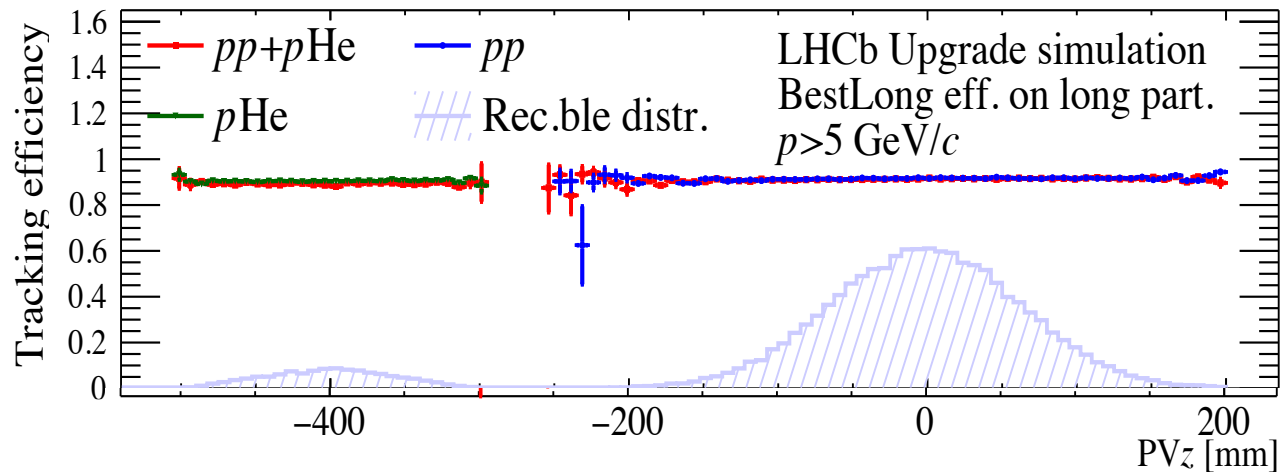
- ✓ polarized gas target technology well established (HERMES @ DESY, ANKE @ COSY with high performance)
- ✓ Target experts from HERMES and COSY involved in first person in the design of the apparatus
- ✓ marginal impact on LHC beam lifetime and LHCb mainstream physics program and performances
- ✓ can run in parallel with collider mode (interaction regions well displaced)
- ✓ can benefit from both protons and heavy-ion beams
- ✓ allows also injection of unpolarized gases ($H_2, D_2, He, N_2, O_2, Ne, Ar, \dots$)
- ✓ broad physics program (next slides)



The LHCspin project



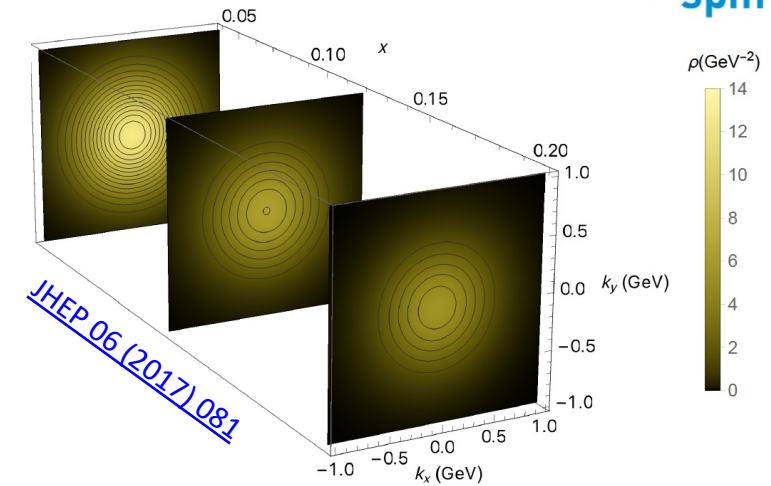
- ✓ polarized gas target technology well established (HERMES @ DESY, ANKE @ COSY with high performance)
- ✓ Target experts from HERMES and COSY involved in first person in the design of the apparatus
- ✓ marginal impact on LHC beam lifetime and LHCb mainstream physics program and performances
- ✓ can run in parallel with collider mode (interaction regions well displaced)
- ✓ can benefit from both protons and heavy-ion beams
- ✓ allows also injection of unpolarized gases ($H_2, D_2, He, N_2, O_2, Ne, Ar, \dots$)
- ✓ broad physics program (next slides)



The physics goals of LHCspin



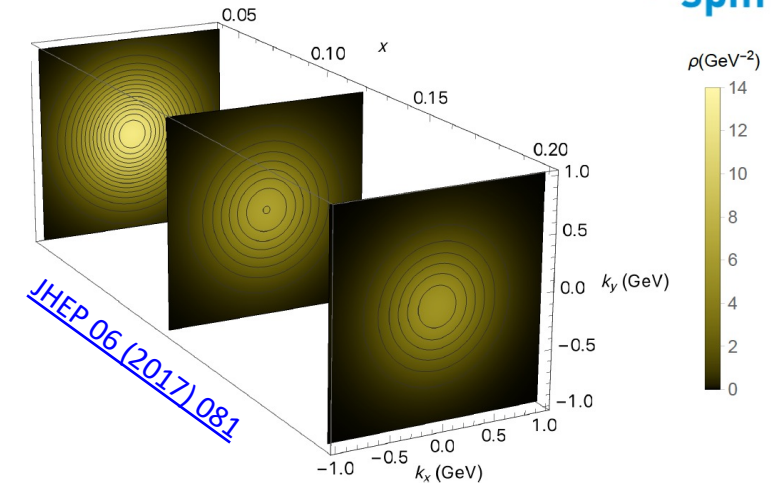
- **Multi-dimensional nucleon structure in a poorly explored kinematic domain**
- Measure experimental observables sensitive to both **quarks and gluons TMDs**
- Make use of new probes (charmed and beauty mesons)
- Complement present and future SIDIS results
- Test non-trivial process dependence of quarks and (especially) gluons TMDs
- Extend our understanding of the strong force in the non-perturbative regime



The physics goals of LHCspin



- **Multi-dimensional nucleon structure in a poorly explored kinematic domain**
- Measure experimental observables sensitive to both **quarks and gluons TMDs**
- Make use of new probes (charmed and beauty mesons)
- Complement present and future SIDIS results
- Test non-trivial process dependence of quarks and (especially) gluons TMDs
- Extend our understanding of the strong force in the non-perturbative regime



quark pol.

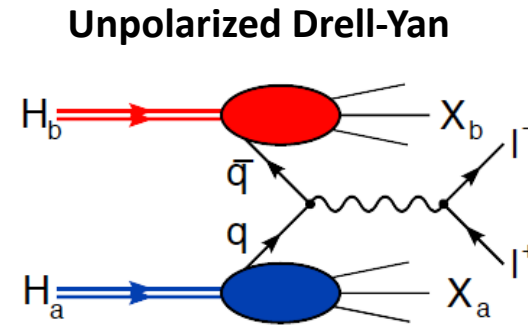
	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

nucleon pol.

- **Significant experimental progress in the last 15 years!**
- **main results from SIDIS** (HERMES, COMPASS, JLAB, → EIC)
- **Drell-Yan** in h-h collisions offers a complementary approach (COMPASS, RHIC)
- Several extractions already available from global analyses
- Now entering the precision era

Quark TMDs

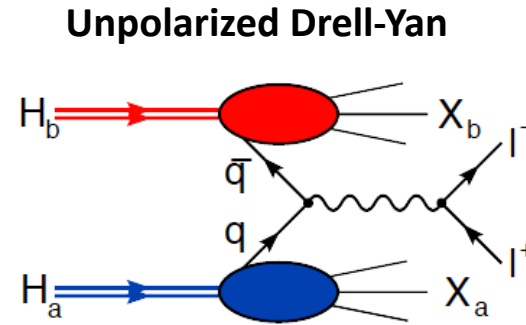
		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



- Theoretically cleanest hard h-h scattering process
- LHCb has excellent μ -ID & reconstruction for $\mu^+\mu^-$
- **dominant:** $\bar{q}(x_{beam}) + q(x_{target}) \rightarrow \mu^+\mu^-$
- suppressed: $q(x_{beam}) + \bar{q}(x_{target}) \rightarrow \mu^+\mu^-$
- beam sea quarks probed at small x
- target valence quarks probed at large x

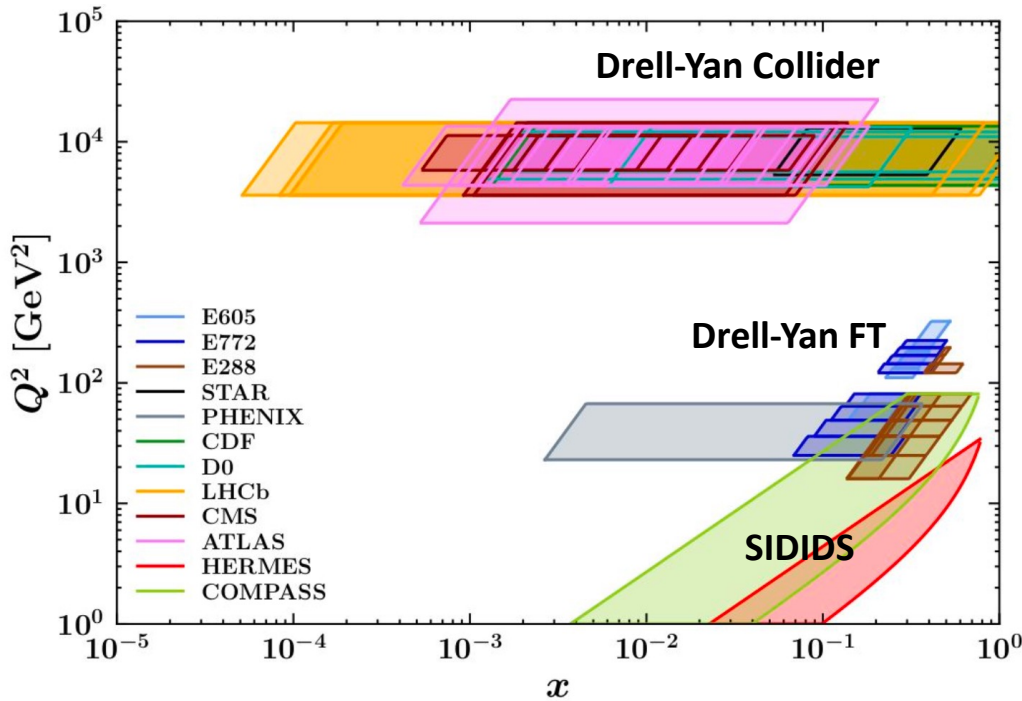
Quark TMDs

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



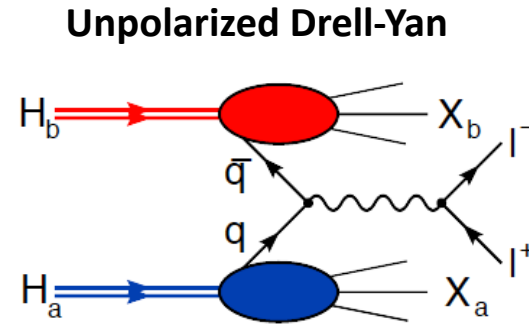
- Theoretically cleanest hard h-h scattering process
- LHCb has excellent μ -ID & reconstruction for $\mu^+\mu^-$
- **dominant:** $\bar{q}(x_{beam}) + q(x_{target}) \rightarrow \mu^+\mu^-$
- suppressed: $q(x_{beam}) + \bar{q}(x_{target}) \rightarrow \mu^+\mu^-$
- beam sea quarks probed at small x
- target valence quarks probed at large x

From Andrea's talk



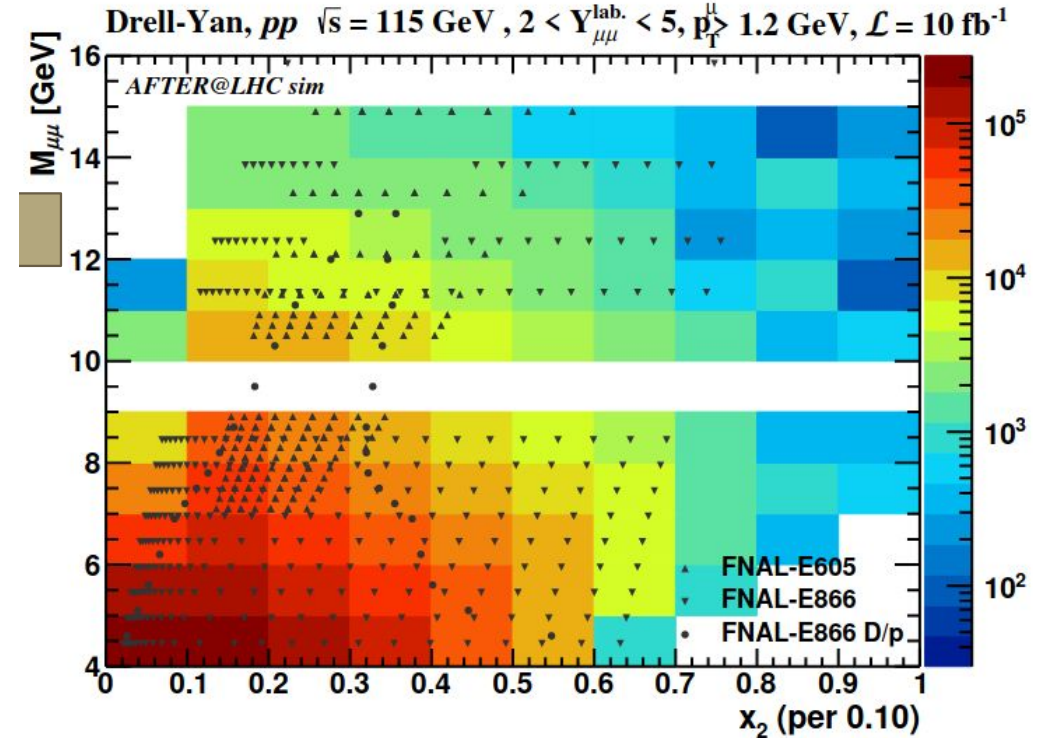
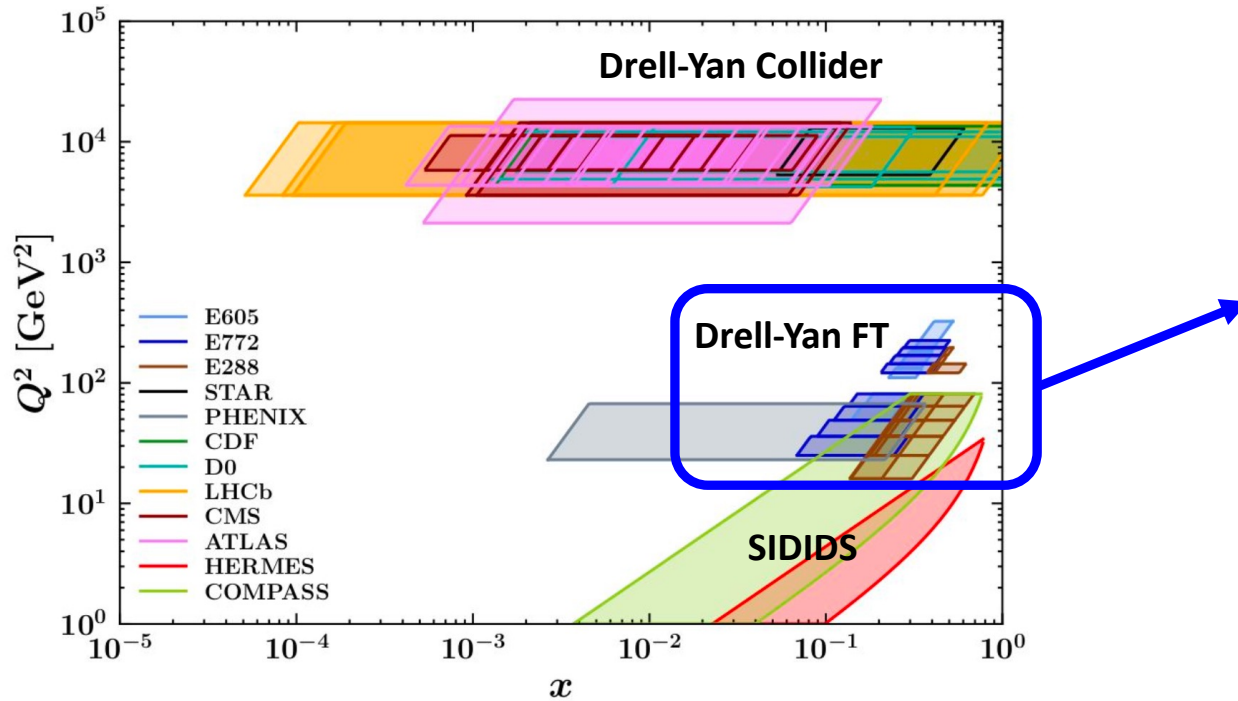
Quark TMDs

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



- Theoretically cleanest hard h-h scattering process
- LHCb has excellent μ -ID & reconstruction for $\mu^+\mu^-$
- **dominant:** $\bar{q}(x_{beam}) + q(x_{target}) \rightarrow \mu^+\mu^-$
- suppressed: $q(x_{beam}) + \bar{q}(x_{target}) \rightarrow \mu^+\mu^-$
- beam sea quarks probed at small x
- target valence quarks probed at large x

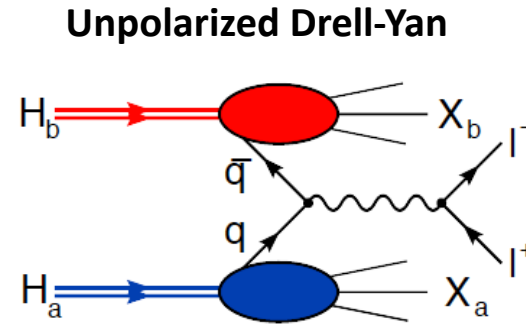
From Andrea's talk



[arXiv:1807.00603]

Quark TMDs

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

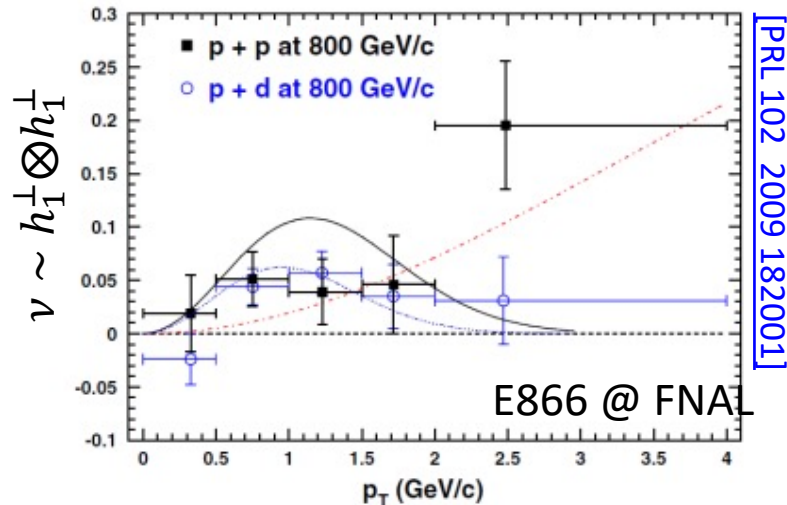


- Theoretically cleanest hard h-h scattering process
- LHCb has excellent μ -ID & reconstruction for $\mu^+\mu^-$
- **dominant:** $\bar{q}(x_{beam}) + q(x_{target}) \rightarrow \mu^+\mu^-$
- suppressed: $q(x_{beam}) + \bar{q}(x_{target}) \rightarrow \mu^+\mu^-$
- beam sea quarks probed at small x
- target valence quarks probed at large x

Sensitive to unpol. and BM TMDs for $q_T \ll M_U$

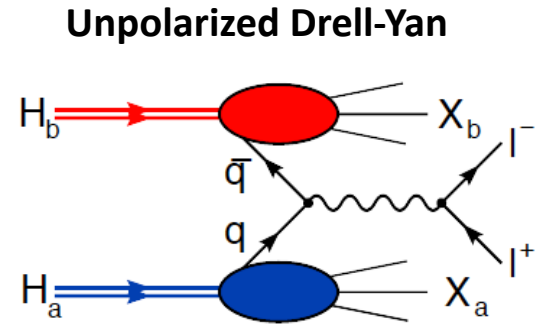
$$d\sigma_{UU}^{DY} \propto f_1^{\bar{q}} \otimes f_1^q + \cos 2\phi h_1^{\perp, \bar{q}} \otimes h_1^{\perp, q}$$

violation of Lam-Tung relation



Quark TMDs

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

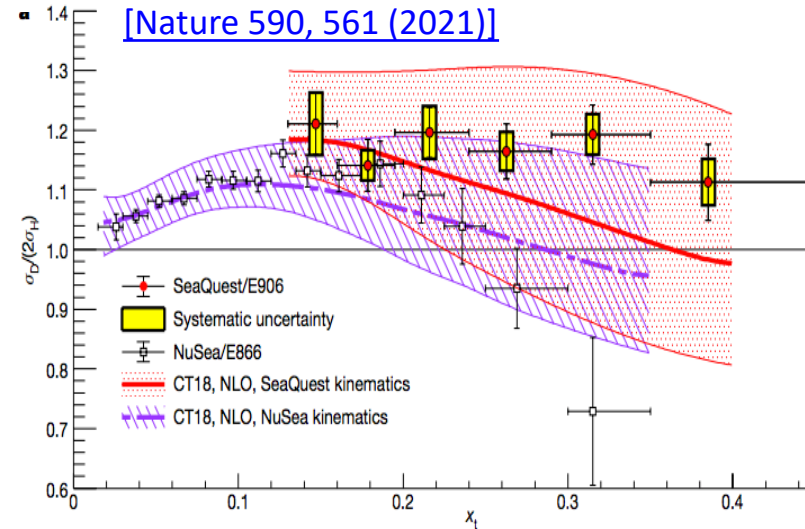
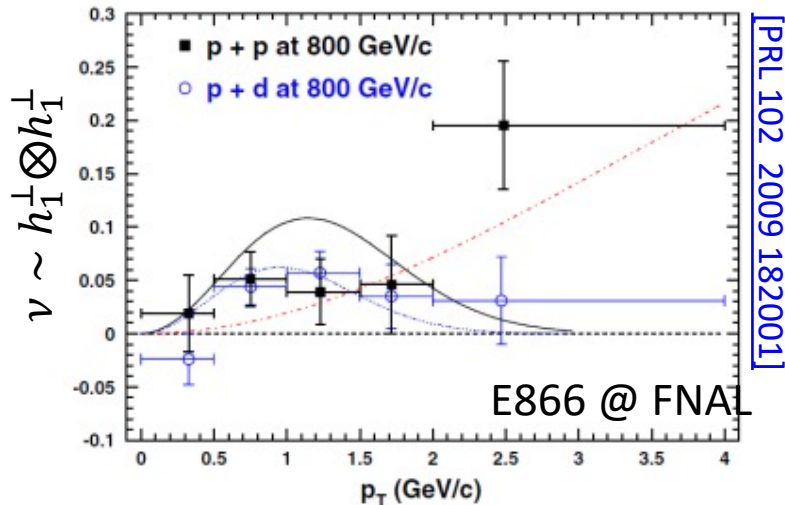


- Theoretically cleanest hard h-h scattering process
- LHCb has excellent μ -ID & reconstruction for $\mu^+\mu^-$
- **dominant:** $\bar{q}(x_{beam}) + q(x_{target}) \rightarrow \mu^+\mu^-$
- **suppressed:** $q(x_{beam}) + \bar{q}(x_{target}) \rightarrow \mu^+\mu^-$
- beam sea quarks probed at small x
- target valence quarks probed at large x

Sensitive to unpol. and BM TMDs for $q_T \ll M_U$

$$d\sigma_{UU}^{DY} \propto f_1^{\bar{q}} \otimes f_1^q + \cos 2\phi h_1^{\perp, \bar{q}} \otimes h_1^{\perp, q}$$

violation of Lam-Tung relation

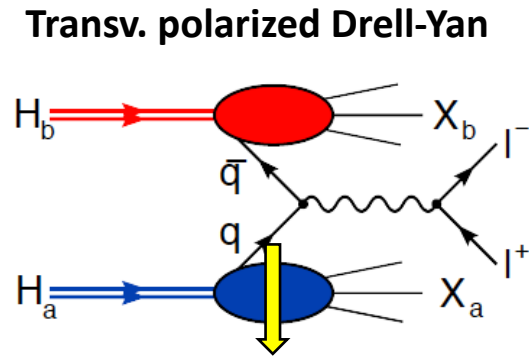


- Lattice QCD: $\bar{s}(x) \neq s(x)$ [\[arXiv:1809.04975\]](https://arxiv.org/abs/1809.04975)
- **proton sea more complex than originally thought!**

- H & D targets allow to study the **antiquark content of the nucleon**
- SeaQuest (E906): $\bar{d}(x) > \bar{u}(x) \Rightarrow$ **sea is not flavour symmetric!**

Quark TMDs

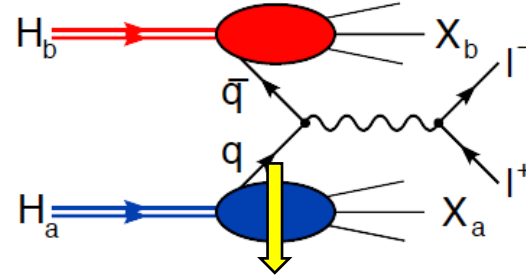
		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



Quark TMDs

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

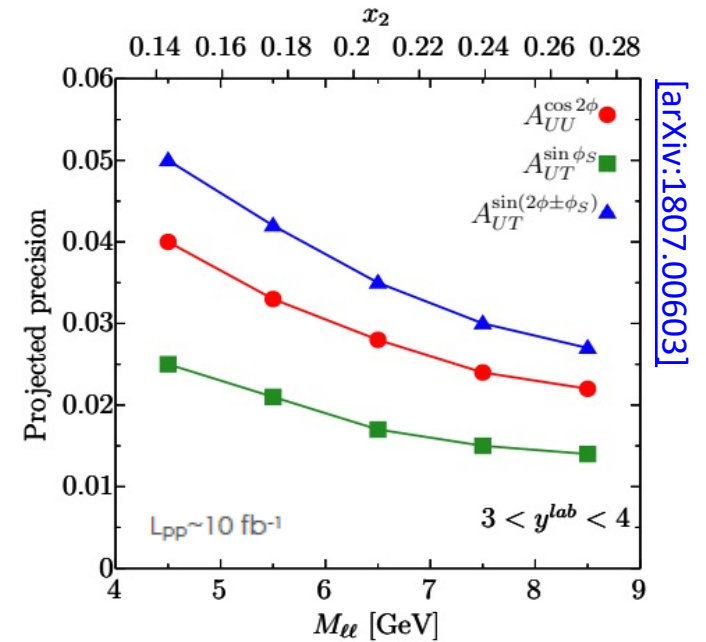
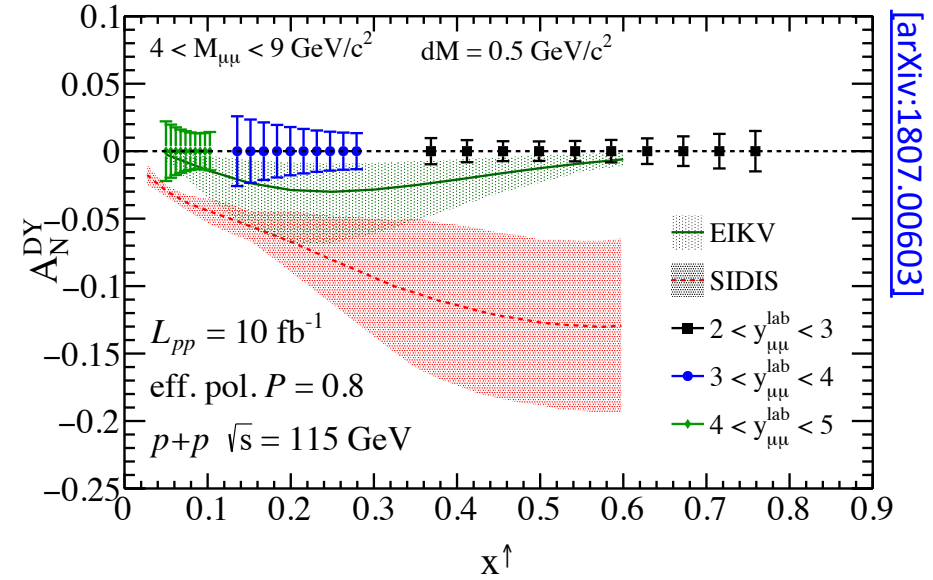
Transv. polarized Drell-Yan



- Sensitive to quark TMDs through TSSAs

$$A_N^{DY} = \frac{1}{P} \frac{\sigma_{DY}^\uparrow - \sigma_{DY}^\downarrow}{\sigma_{DY}^\uparrow + \sigma_{DY}^\downarrow} \Rightarrow A_{UT}^{\sin\phi_S} \sim \frac{f_1^q \otimes f_{1T}^{\perp q}}{f_1^q \otimes f_1^q}, \quad A_{UT}^{\sin(2\phi - \phi_S)} \sim \frac{h_1^{\perp q} \otimes h_1^q}{f_1^q \otimes f_1^q}, \dots$$

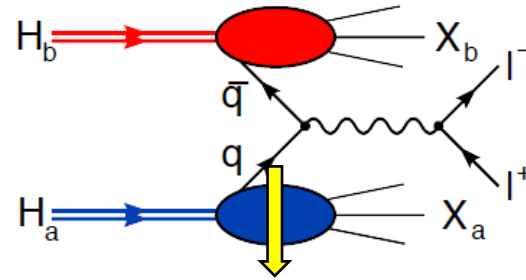
(ϕ : azimuthal orientation of lepton pair in dilepton CM)



Quark TMDs

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

Transv. polarized Drell-Yan



- Sensitive to quark TMDs through TSSAs

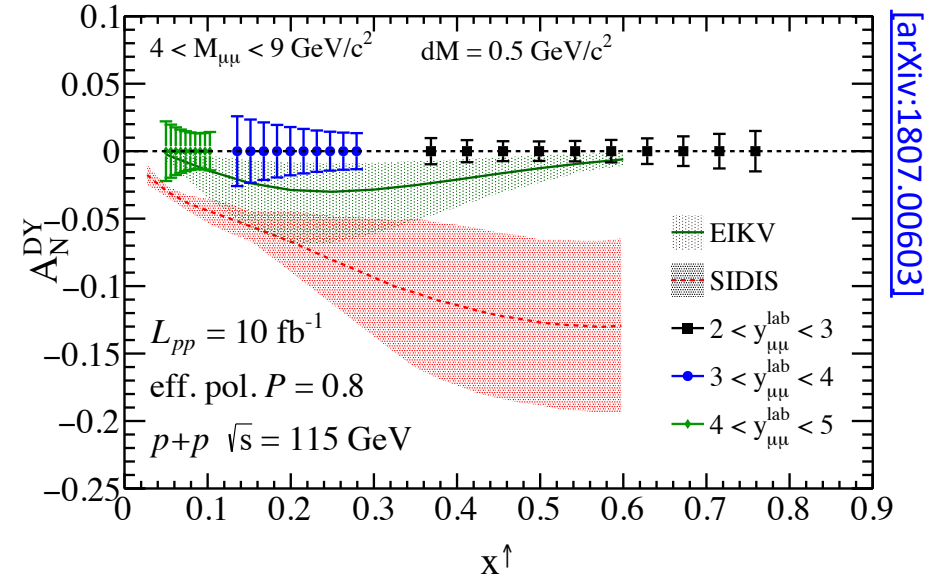
$$A_N^{DY} = \frac{1}{P} \frac{\sigma_{DY}^\uparrow - \sigma_{DY}^\downarrow}{\sigma_{DY}^\uparrow + \sigma_{DY}^\downarrow} \Rightarrow A_{UT}^{\sin\phi_S} \sim \frac{f_1^q \otimes f_{1T}^{\perp q}}{f_1^q \otimes f_1^q}, \quad A_{UT}^{\sin(2\phi - \phi_S)} \sim \frac{h_1^{\perp q} \otimes h_1^q}{f_1^q \otimes f_1^q}, \dots$$

(ϕ : azimuthal orientation of lepton pair in dilepton CM)

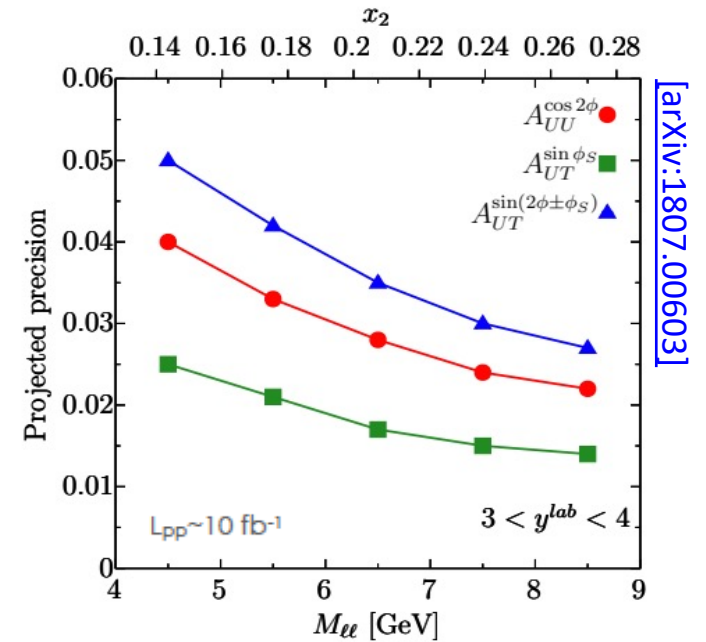
- Extraction of qTMDs does not require knowledge of FF
- Verify sign change of Sivers function wrt SIDIS

$$f_{1T}^\perp|_{DY} = -f_{1T}^\perp|_{SIDIS}$$

- Test flavour sensitivity using both H and D targets



[arXiv:1807.00603]



[arXiv:1807.00603]

Gluon TMDs

		gluon pol.		
		U	Circularly	Linearly
nucleon pol.	U	f_1^g		$h_1^{\perp g}$
	L		g_{1L}^g	$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

Theory framework well consolidated ...**but experimental access still extremely limited!**

Gluon TMDs

		gluon pol.		
		U	Circularly	Linearly
nucleon pol.	U	f_1^g		$h_1^{\perp g}$
	L		g_{1L}^g	$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

Theory framework well consolidated ...but experimental access still extremely limited!

Similar naming/notation of quark TMDs, but there are important differences!

- the **linearity gTMD** (h_1^g) is completely unrelated to the quark transversity (h_1^q), and has no collinear counterpart
- **different naïve-time-reversal properties**

	T-even	T-odd
q	h_1^q	$h_1^{\perp q}$
g	$h_1^{\perp g}$	h_1^g

Gluon TMDs

		gluon pol.		
		U	Circularly	Linearly
nucleon pol.	U	f_1^g		$h_1^{\perp g}$
	L		g_{1L}^g	$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

Theory framework well consolidated ...but experimental access still extremely limited!

Similar naming/notation of quark TMDs, but there are important differences!

- the **linearity gTMD** (h_1^g) is completely unrelated to the quark transversity (h_1^q), and has no collinear counterpart
- **different naïve-time-reversal properties**

	T-even	T-odd
q	h_1^q	$h_1^{\perp q}$
g	$h_1^{\perp g}$	h_1^g

- Also the gTMD phenomenology is enriched by the **process dependence** originating by ISI/FSI encoded in the **gauge links**.
- **The gluon correlator depends on 2 path-dependent gauge links**, resulting in a more complex process dependence



- Depending on their combinations, **there are 2 independent versions of each gTMD** that can be probed in different processes and can have different magnitude and width and different x and k_T dependencies!

Gluon TMDs

		gluon pol.		
		U	Circularly	Linearly
nucleon pol.	U	f_1^g		$h_1^{\perp g}$
	L		g_{1L}^g	$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

Theory framework well consolidated ...but experimental access still extremely limited!

Similar naming/notation of quark TMDs, but there are important differences!

- the **linearity gTMD** (h_1^g) is completely unrelated to the quark transversity (h_1^q), and has no collinear counterpart
- different naïve-time-reversal properties**

	T-even	T-odd
q	h_1^q	$h_1^{\perp q}$
g	$h_1^{\perp g}$	h_1^g

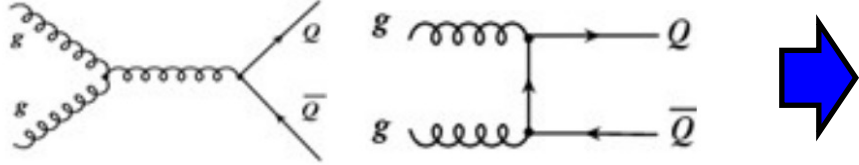
- Also the gTMD phenomenology is enriched by the **process dependence** originating by ISI/FSI encoded in the **gauge links**.
- The gluon correlator depends on 2 path-dependent gauge links**, resulting in a more complex process dependence



- Depending on their combinations, **there are 2 independent versions of each gTMD** that can be probed in different processes and can have different magnitude and width and different x and k_T dependencies!
- E.g. there are 2 types of f_1^g and $h_1^{\perp g}$: $[+ +] = [- -]$ **Weizsacker-Williams (WW)** ; $[+ -] = [- +]$ **DiPole (DP)**
- 2 indep. GSF: $f_{1T}^{\perp g[+,+]}$ **“f-type”** \rightarrow antisymm. colour structure ; $f_{1T}^{\perp g[+,-]}$ **“d-type”** \rightarrow symm. colour structure

Probing the g TMDs

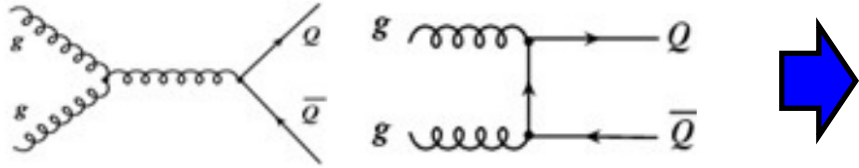
In high-energy hadron collisions, heavy quarks are dominantly produced through gg fusion:



The most efficient way to access the gluon dynamics inside the proton at LHC is to **measure heavy-quark observables**

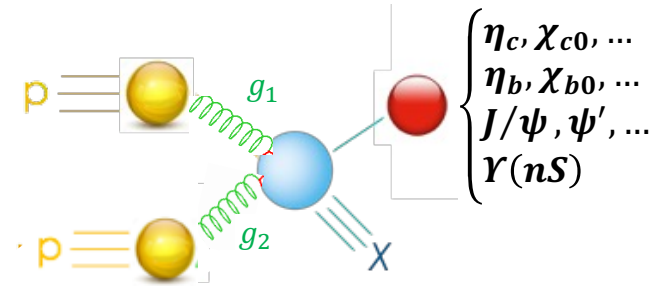
Probing the gTMDs

In high-energy hadron collisions, heavy quarks are dominantly produced through gg fusion:



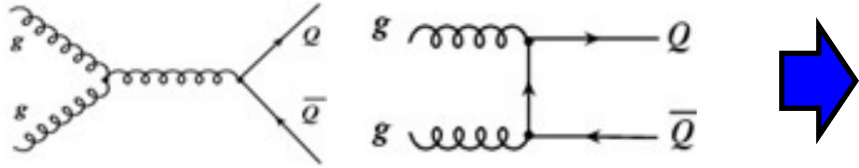
The most efficient way to access the gluon dynamics inside the proton at LHC is to **measure heavy-quark observables**

- **Inclusive quarkonia production in (un)polarized pp interaction** ($pp^{(\uparrow)} \rightarrow [Q\bar{Q}]X$) turns out to be an ideal observable to access gTMDs (assuming TMD factorization)



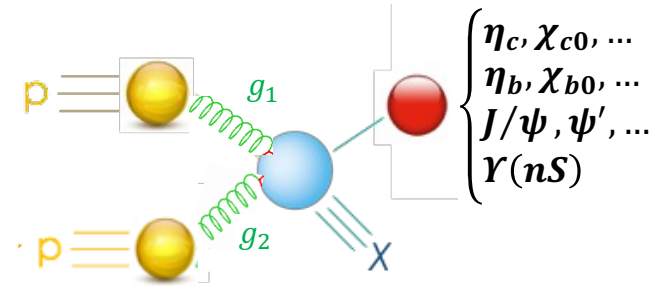
Probing the gTMDs

In high-energy hadron collisions, heavy quarks are dominantly produced through gg fusion:



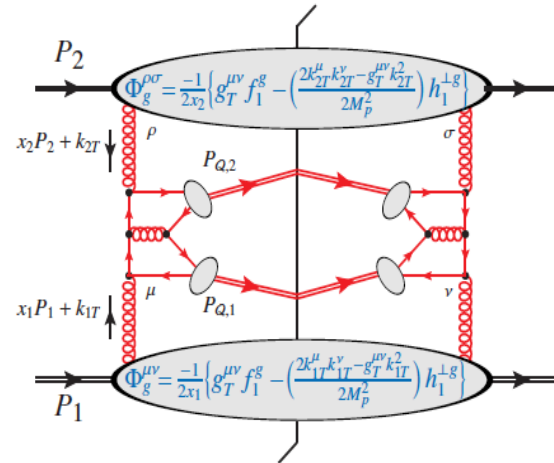
The most efficient way to access the gluon dynamics inside the proton at LHC is to **measure heavy-quark observables**

- **Inclusive quarkonia production in (un)polarized pp interaction** ($pp^{(\uparrow)} \rightarrow [Q\bar{Q}]X$) turns out to be an ideal observable to access gTMDs (assuming TMD factorization)



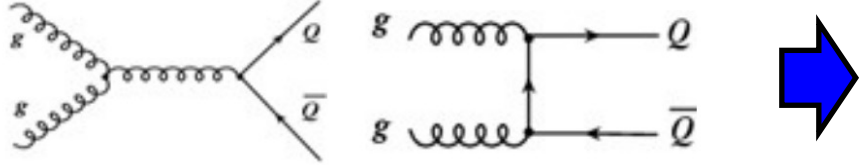
- TMD factorization requires $q_T(Q) \ll M_Q$. Can look at **associate quarkonia production**, where only the relative q_T needs to be small:

E.g.: $pp^{(\uparrow)} \rightarrow J/\psi + J/\psi + X$ (\rightarrow Alice)



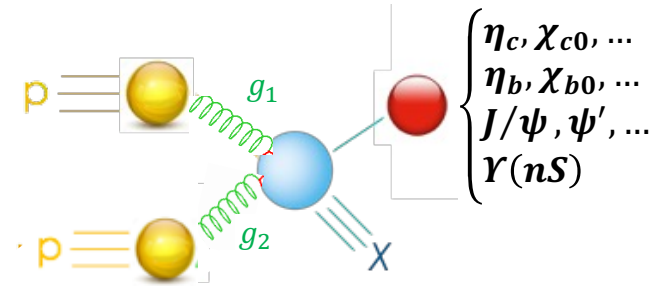
Probing the g TMDs

In high-energy hadron collisions, heavy quarks are dominantly produced through gg fusion:



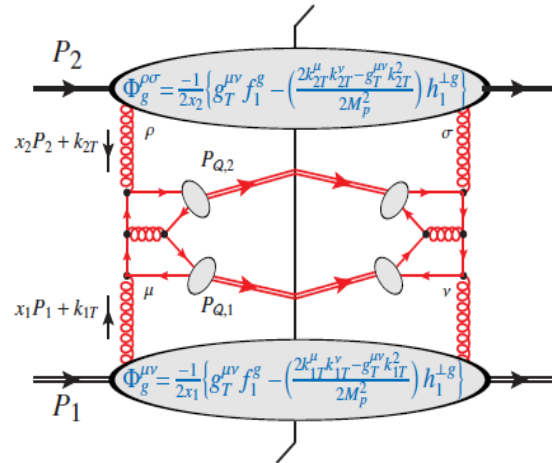
The most efficient way to access the gluon dynamics inside the proton at LHC is to **measure heavy-quark observables**

- **Inclusive quarkonia production in (un)polarized pp interaction** ($pp^{(\uparrow)} \rightarrow [Q\bar{Q}]X$) turns out to be an ideal observable to access g TMDs (assuming TMD factorization)

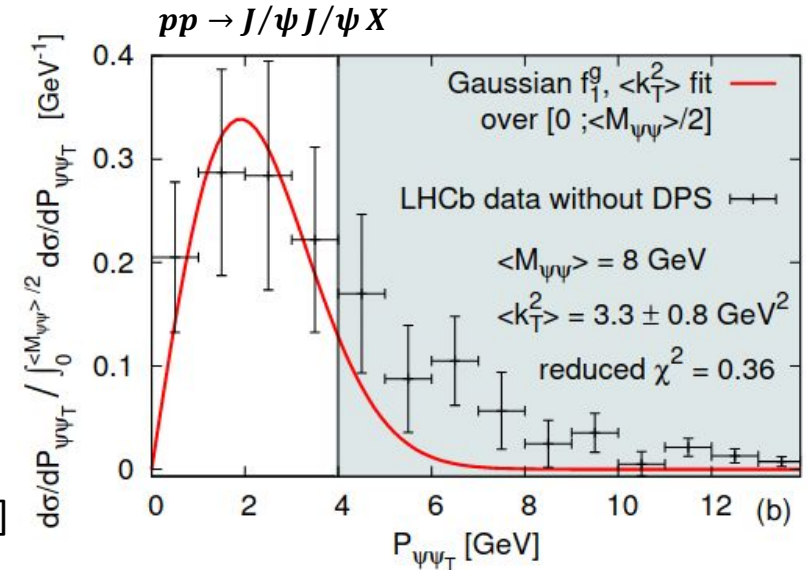


- TMD factorization requires $q_T(Q) \ll M_Q$. Can look at **associate quarkonia production**, where only the relative q_T needs to be small:

E.g.: $pp^{(\uparrow)} \rightarrow J/\psi + J/\psi + X$ (\rightarrow Alice)

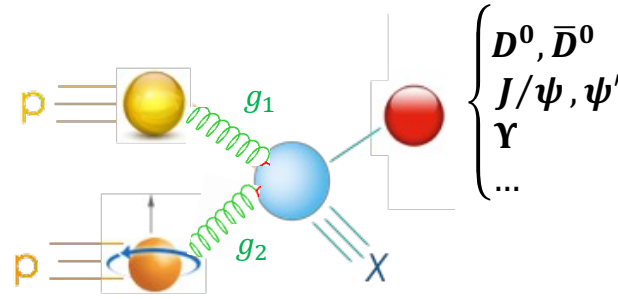


First extraction of f_1^g from LHCb di- J/ψ production data at 13 TeV [[Lansberg et. al.](#)]



Probing the gluon Sivers funct.

$$\Gamma_T^{\mu\nu}(x, \mathbf{p}_T) = \frac{x}{2} \left\{ g_T^{\mu\nu} \frac{\epsilon_T^{\rho\sigma} p_{T\rho} S_{T\sigma}}{M_p} f_{1T}^{\perp g}(x, \mathbf{p}_T^2) + \dots \right\}$$

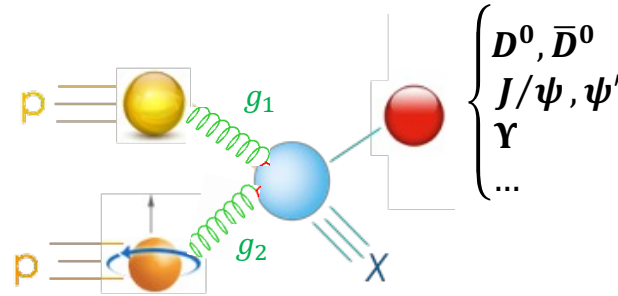


- Sheds light on spin-orbit correlations of unpol. gluons inside a transv. pol. proton
- sensitive to color exchange among IS and FS and to gluon OAM
- expected to be quite small (quasi-saturation of Burkardt sum rule by $f_{1T}^{\perp q}$ and QCD predictions in large- N_c limit)

		gluon pol.		
		U	Circularly	Linearly
nucleon pol.	U	f_1^g		$h_1^{\perp g}$
	L		g_{1L}^g	$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

Probing the gluon Sivers funct.

$$\Gamma_T^{\mu\nu}(x, \mathbf{p}_T) = \frac{x}{2} \left\{ g_T^{\mu\nu} \frac{\epsilon_T^{\rho\sigma} p_{T\rho} S_{T\sigma}}{M_p} f_{1T}^{\perp g}(x, \mathbf{p}_T^2) + \dots \right\}$$



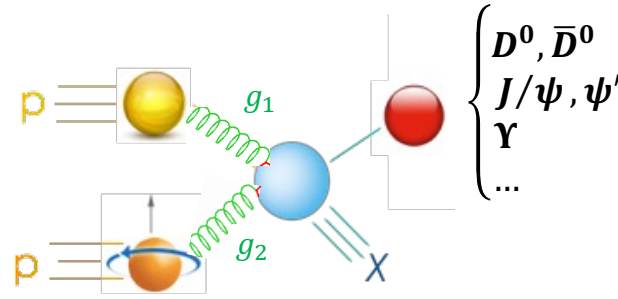
		gluon pol.	
		U	Linearly
nucleon pol.	U	f_1^g	$h_1^{\perp g}$
	L		$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	$h_{1T}^g, h_{1T}^{\perp g}$
		g_{1L}^g	

- Sheds light on spin-orbit correlations of unpol. gluons inside a transv. pol. proton
- sensitive to color exchange among IS and FS and to gluon OAM
- expected to be quite small (quasi-saturation of Burkardt sum rule by $f_{1T}^{\perp g}$ and QCD predictions in large- N_c limit)
- can be accessed through the measurement of the TSSAs in **inclusive heavy meson production**

$$A_N = \frac{1}{P} \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow} \propto [f_{1T}^{\perp g}(x_a, k_{\perp a}) \otimes f_g(x_b, k_{\perp b}) \otimes d\sigma_{gg \rightarrow QQg}] \sin \phi_S + \dots$$

Probing the gluon Sivers funct.

$$\Gamma_T^{\mu\nu}(x, \mathbf{p}_T) = \frac{x}{2} \left\{ g_T^{\mu\nu} \frac{\epsilon_T^{\rho\sigma} p_{T\rho} S_{T\sigma}}{M_p} f_{1T}^{\perp g}(x, \mathbf{p}_T^2) + \dots \right\}$$

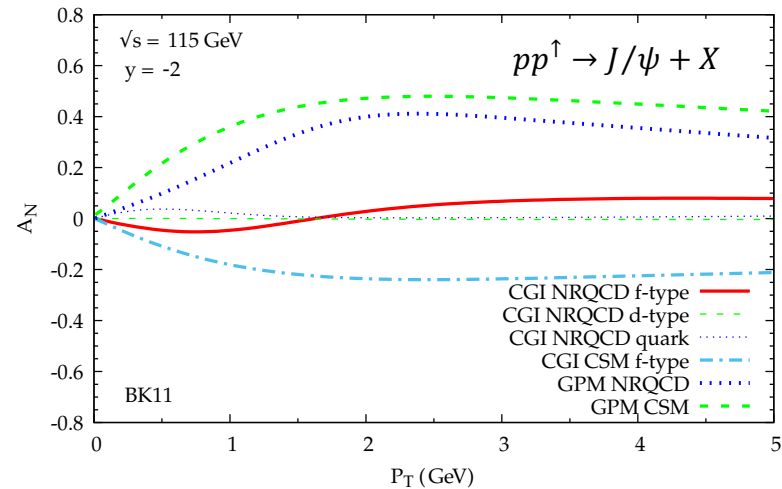
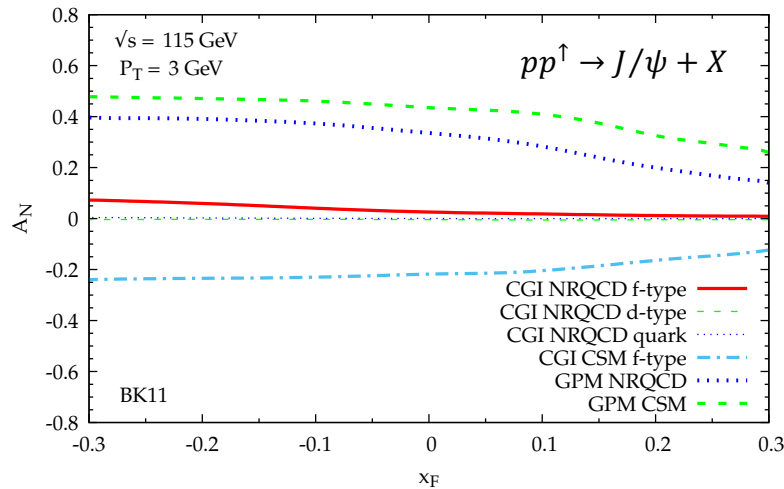


		gluon pol.		
		U	Circularly	Linearly
nucleon pol.	U	f_1^g		$h_1^{\perp g}$
	L		g_{1L}^g	$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

- Sheds light on spin-orbit correlations of unpol. gluons inside a transv. pol. proton
- sensitive to color exchange among IS and FS and to gluon OAM
- expected to be quite small (quasi-saturation of Burkardt sum rule by $f_{1T}^{\perp q}$ and QCD predictions in large- N_c limit)
- can be accessed through the measurement of the TSSAs in **inclusive heavy meson production**

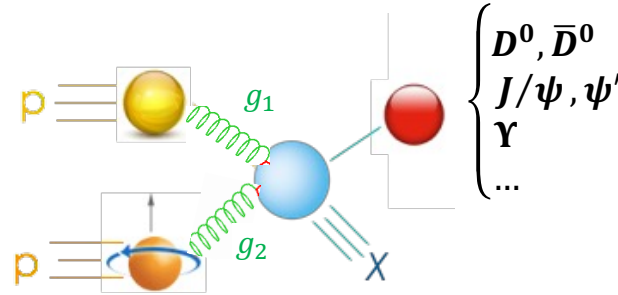
$$A_N = \frac{1}{P} \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow} \propto [f_{1T}^{\perp g}(x_a, k_{\perp a}) \otimes f_g(x_b, k_{\perp b}) \otimes d\sigma_{gg \rightarrow QQg}] \sin \phi_S + \dots$$

Predictions for pol. FT meas. at LHC (LHCspin-like) [\[Phys. Rev. D 102, 094011 \(2020\)\]](#)



Probing the gluon Sivers funct.

$$\Gamma_T^{\mu\nu}(x, \mathbf{p}_T) = \frac{x}{2} \left\{ g_T^{\mu\nu} \frac{\epsilon_T^{\rho\sigma} p_{T\rho} S_{T\sigma}}{M_p} f_{1T}^{\perp g}(x, \mathbf{p}_T^2) + \dots \right\}$$

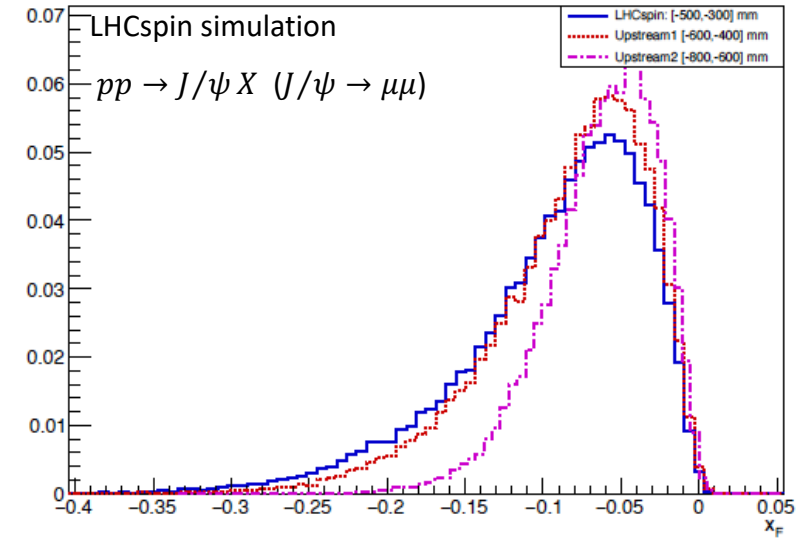
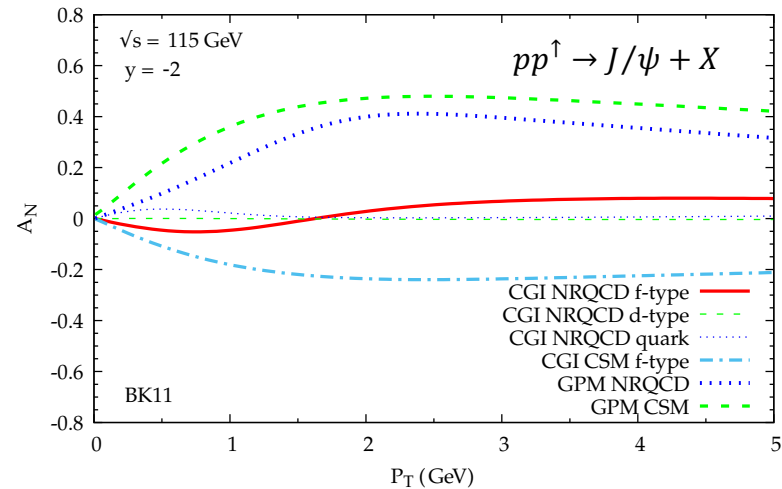
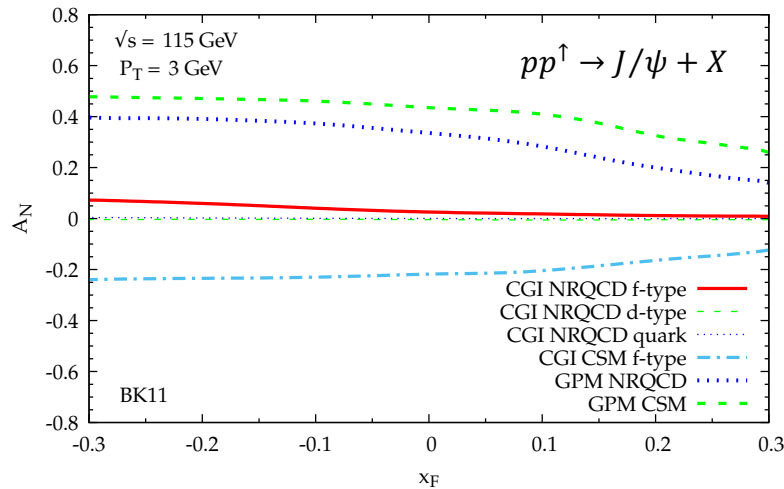


		gluon pol.		
		U	Circularly	Linearly
nucleon pol.	U	f_1^g		$h_1^{\perp g}$
	L		g_{1L}^g	$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

- Sheds light on spin-orbit correlations of unpol. gluons inside a transv. pol. proton
- sensitive to color exchange among IS and FS and to gluon OAM
- expected to be quite small (quasi-saturation of Burkardt sum rule by $f_{1T}^{\perp q}$ and QCD predictions in large- N_c limit)
- can be accessed through the measurement of the TSSAs in **inclusive heavy meson production**

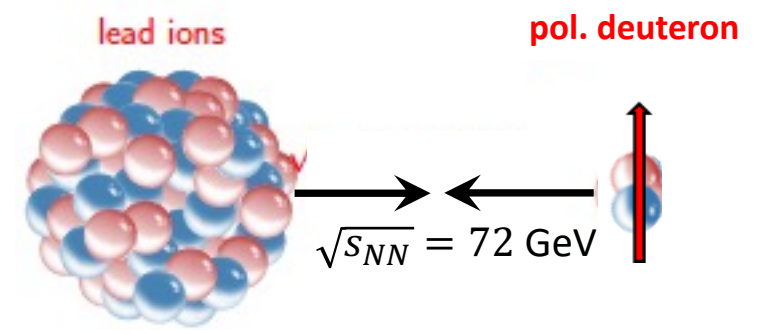
$$A_N = \frac{1}{P} \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow} \propto [f_{1T}^{\perp g}(x_a, k_{\perp a}) \otimes f_g(x_b, k_{\perp b}) \otimes d\sigma_{gg \rightarrow QQg}] \sin \phi_S + \dots$$

Predictions for pol. FT meas. at LHC (LHCspin-like) [\[Phys. Rev. D 102, 094011 \(2020\)\]](#)



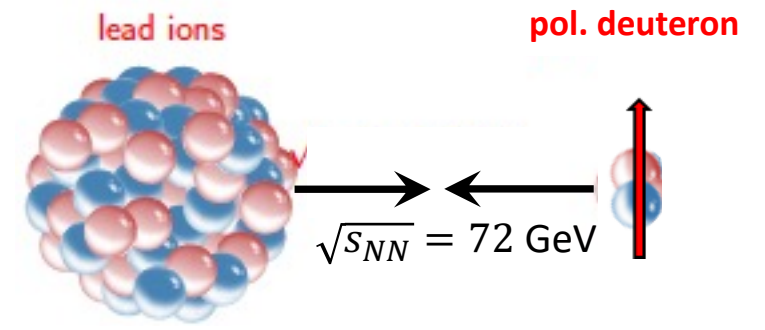
Merging spin physics with heavy-ion physics

- probe collective phenomena in heavy-light systems through **ultra-relativistic collisions of heavy nuclei with trasv. pol. deuterons**
- polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the **elliptic flow** relative to the polarization axis (**ellipticity**).

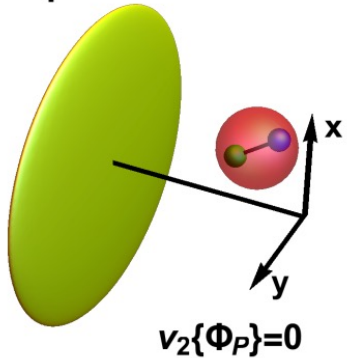


Merging spin physics with heavy-ion physics

- probe collective phenomena in heavy-light systems through **ultra-relativistic collisions of heavy nuclei with trasv. pol. deuterons**
- polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the **elliptic flow** relative to the polarization axis (**ellipticity**).



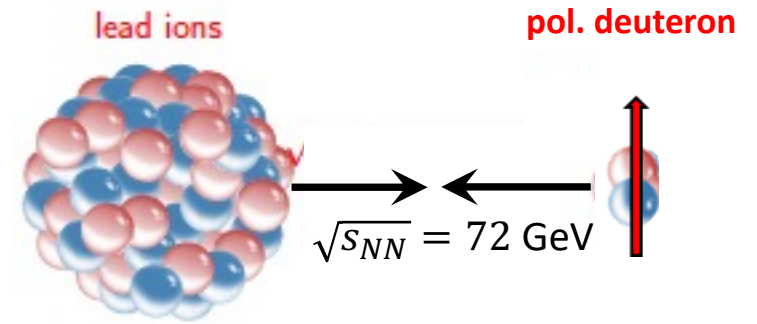
unpolarized d+A



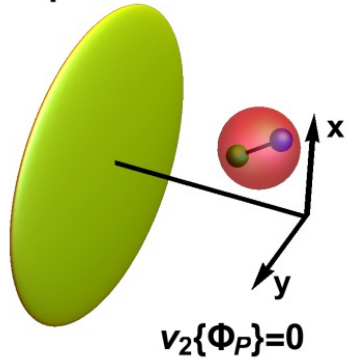
Unpol. deuterons: the fireball is azimuthally symmetric and $v_2 \approx 0$.

Merging spin physics with heavy-ion physics

- probe collective phenomena in heavy-light systems through **ultra-relativistic collisions of heavy nuclei with trasv. pol. deuterons**
- polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the **elliptic flow** relative to the polarization axis (**ellipticity**).

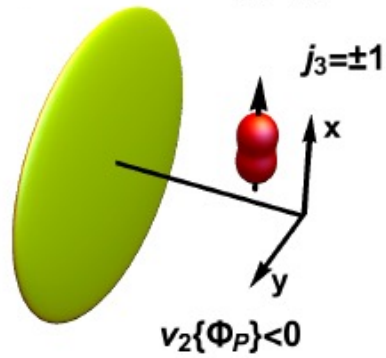


unpolarized d+A



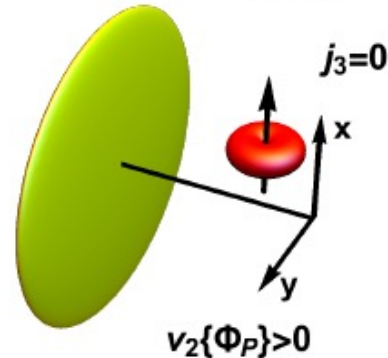
Unpol. deuterons: the fireball is azimuthally symmetric and $v_2 \approx 0$.

d↑+A



$j_3 = \pm 1 \rightarrow$ prolate fireball stretched along the pol. axis, corresponds to $v_2 < 0$

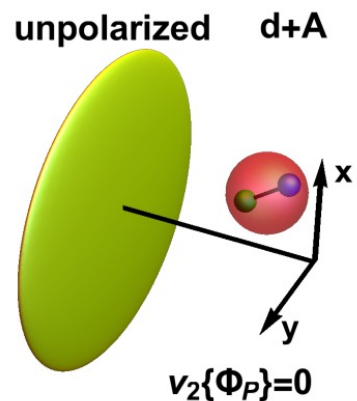
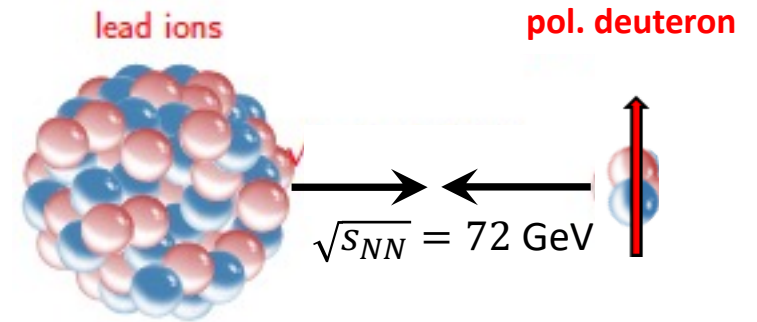
d↑+A



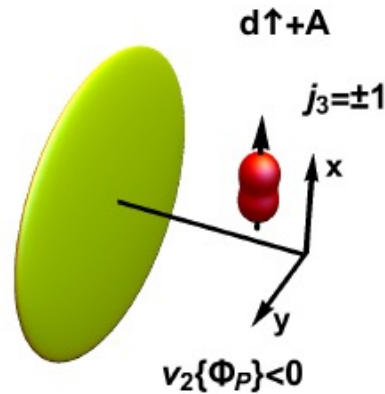
$j_3 = 0 \rightarrow$ oblate fireball corresponds to $v_2 > 0$

Merging spin physics with heavy-ion physics

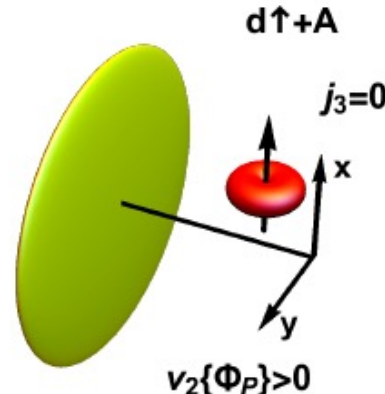
- probe collective phenomena in heavy-light systems through **ultra-relativistic collisions of heavy nuclei with trasv. pol. deuterons**
- polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the **elliptic flow** relative to the polarization axis (**ellipticity**).



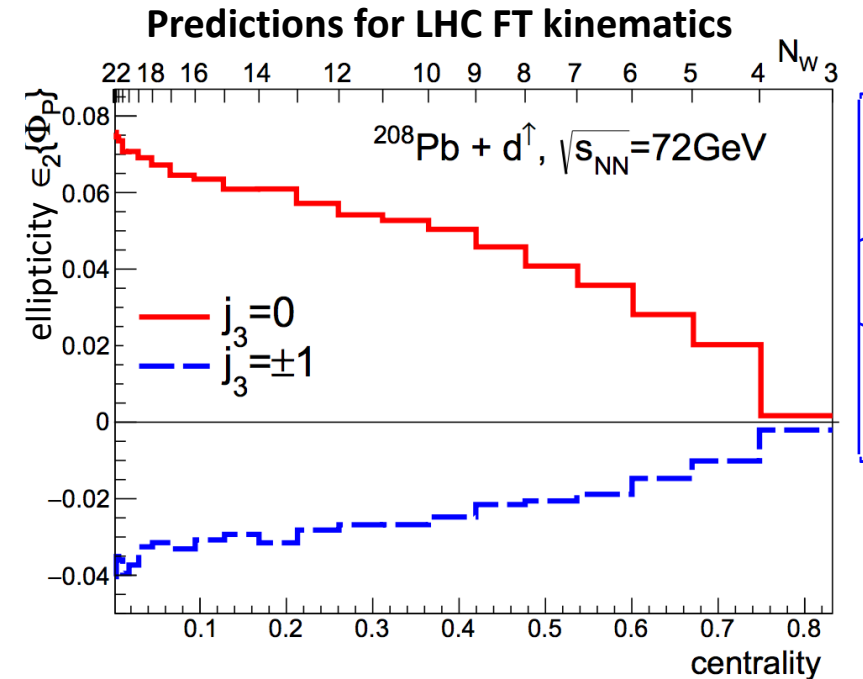
Unpol. deuterons: the fireball is azimuthally symmetric and $v_2 \approx 0$.



$j_3 = \pm 1 \rightarrow$ prolate fireball stretched along the pol. axis, corresponds to $v_2 < 0$



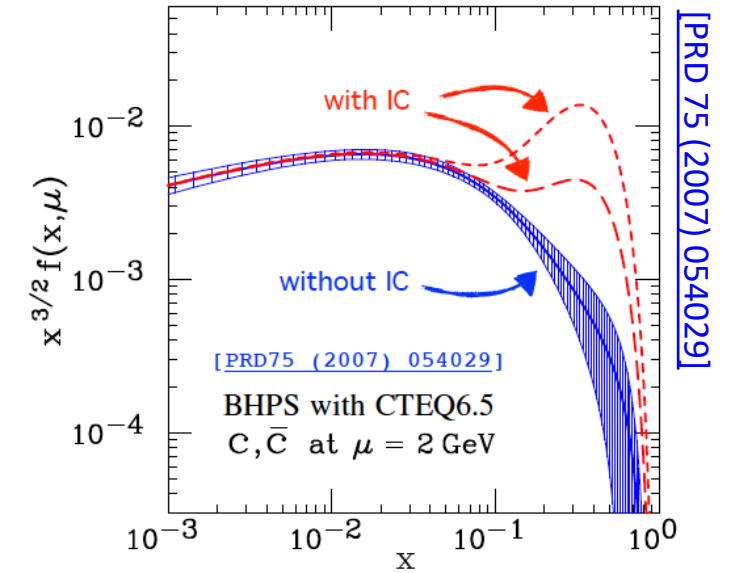
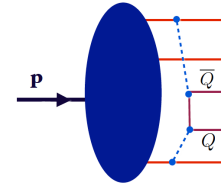
$j_3 = 0 \rightarrow$ oblate fireball corresponds to $v_2 > 0$



[PRC 101 (2020) 024901]

More physics reach with unpolarized FT reactions

- **Intrinsic heavy-quark** [S.J. Brodsky et al., Adv.High Energy Phys. 2015 (2015) 231547]
 - 5-quark Fock state of the proton may contribute at high x !
 - **charm PDFs** at large x could be larger than obtained from conventional fits



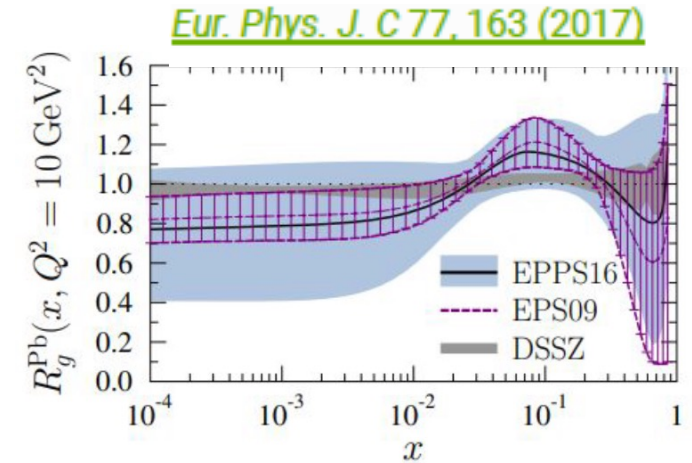
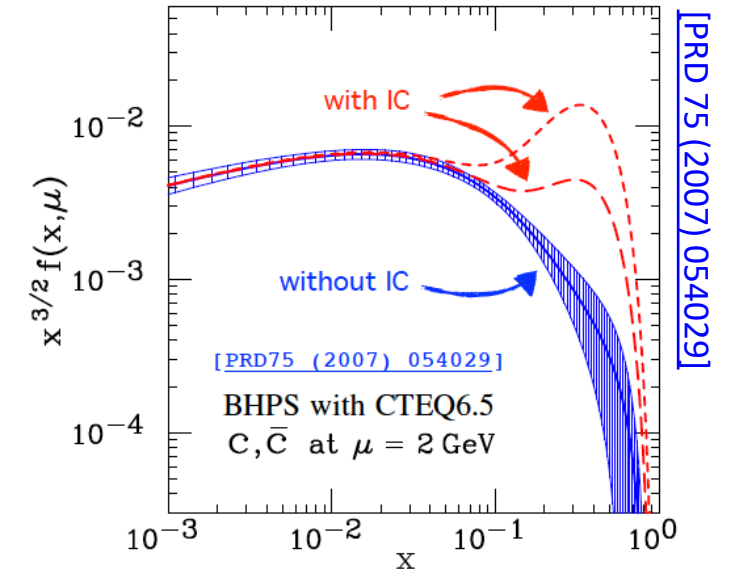
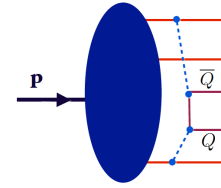
[PRD 75 (2007) 054029]

[PRD75 (2007) 054029]

BHPS with CTEQ6.5
C, C-bar at $\mu = 2$ GeV

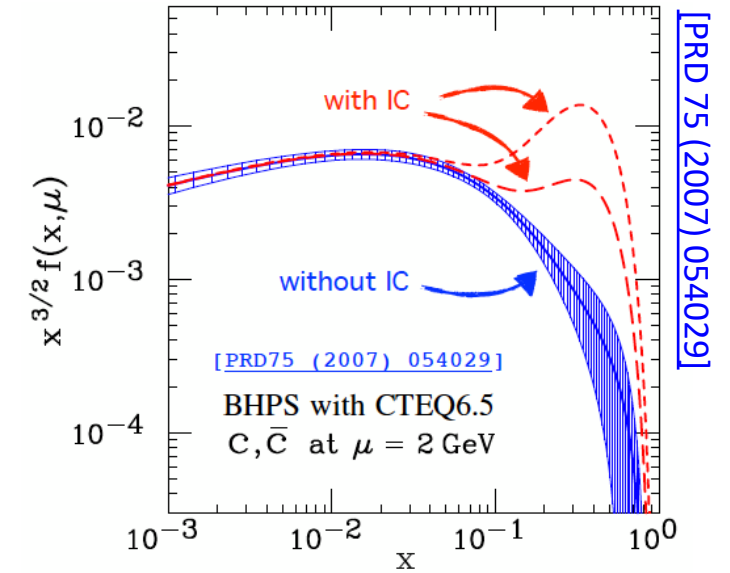
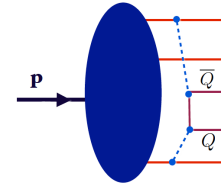
More physics reach with unpolarized FT reactions

- **Intrinsic heavy-quark** [S.J. Brodsky et al., Adv.High Energy Phys. 2015 (2015) 231547]
 - 5-quark Fock state of the proton may contribute at high x !
 - **charm PDFs** at large x could be larger than obtained from conventional fits
- **pA collisions** (using unpolarized gas: He, N, Ne, Ar, Kr, Xe)
 - constraints on nPDFs (e.g. on poorly understood **gluon antishadowing at high x**)
 - studies of parton energy-loss and absorption phenomena in the cold medium
 - reactions of interest for cosmic-ray physics and DM searches



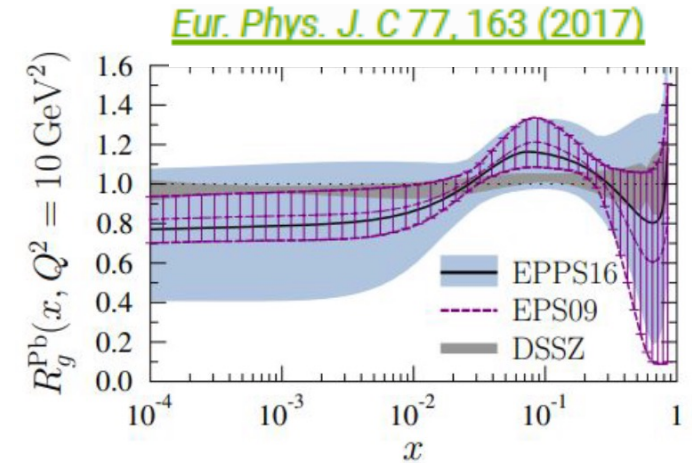
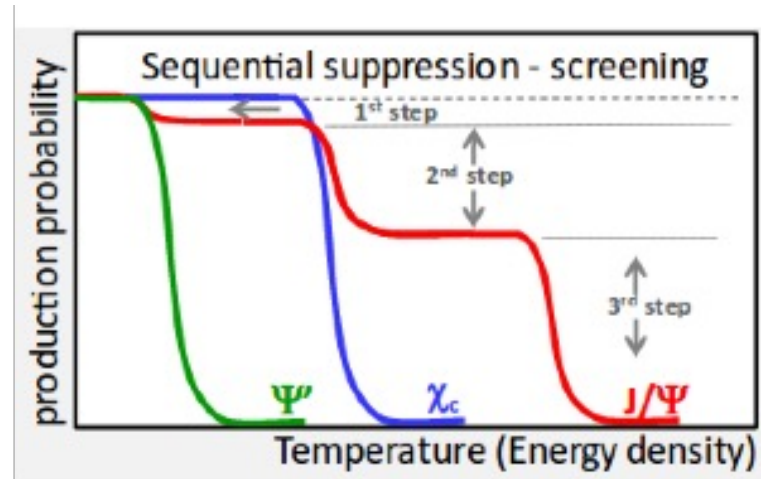
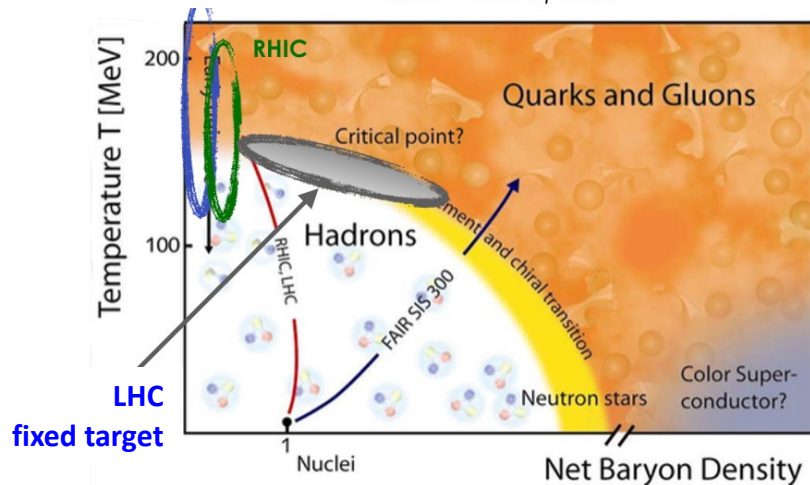
More physics reach with unpolarized FT reactions

- **Intrinsic heavy-quark** [S.J. Brodsky et al., Adv.High Energy Phys. 2015 (2015) 231547]
 - 5-quark Fock state of the proton may contribute at high x !
 - **charm PDFs** at large x could be larger than obtained from conventional fits
- **pA collisions** (using unpolarized gas: He, N, Ne, Ar, Kr, Xe)
 - constraints on nPDFs (e.g. on poorly understood **gluon antishadowing at high x**)
 - studies of parton energy-loss and absorption phenomena in the cold medium
 - reactions of interest for cosmic-ray physics and DM searches
- **PbA collisions at $\sqrt{s_{NN}} \approx 72$ GeV** (using unpolarized gas: He, N, Ne, Ar, Kr, Xe)
 - Study of **QGP formation** (search for predicted **sequential quarkonium suppression**)



LHC @ 5.02 TeV

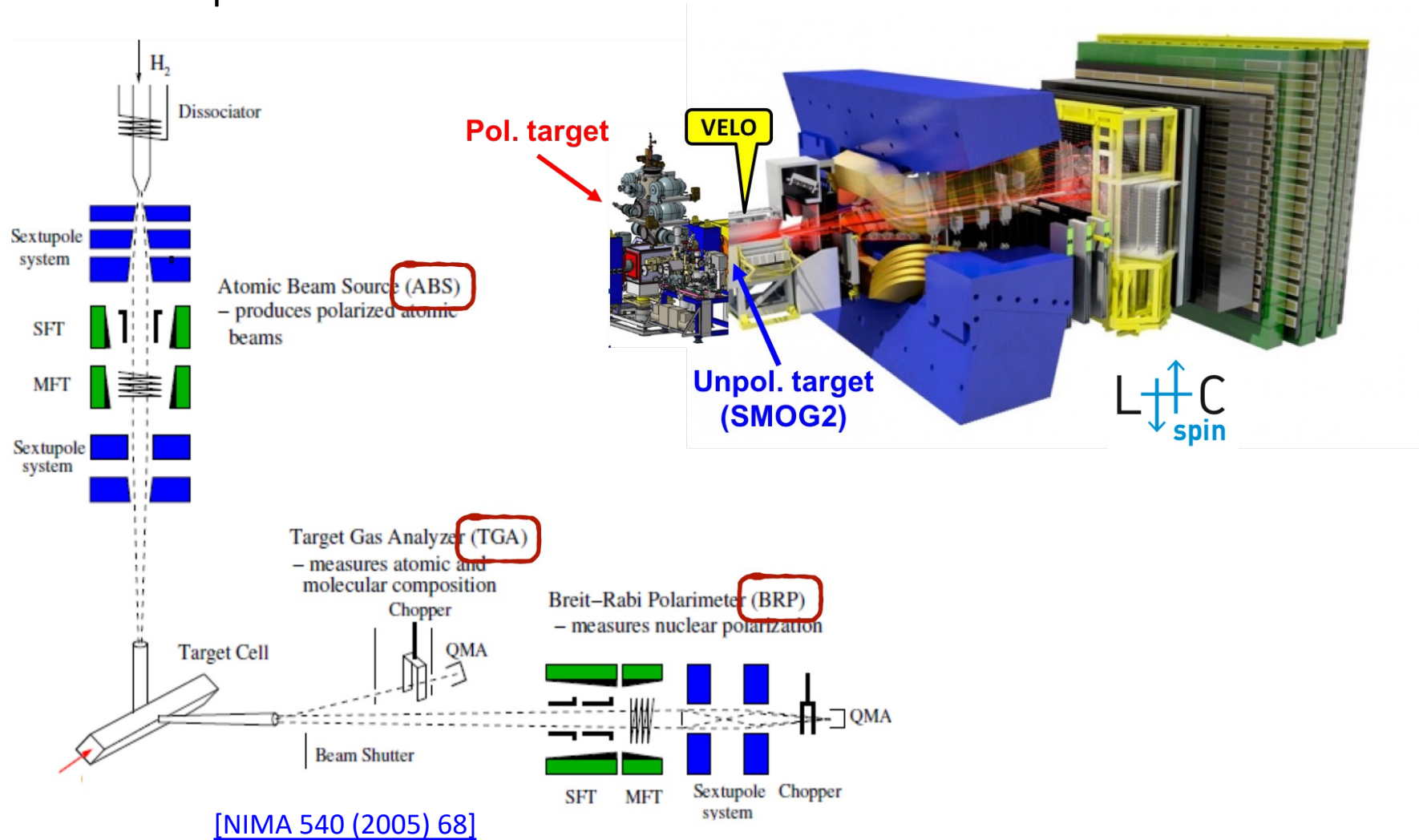
QCD Phase-Space



$c\bar{c}$ states: $J/\psi, \chi_c, \psi', \dots$
 Different binding energies, different dissociation temperatures \rightarrow **medium thermometer**

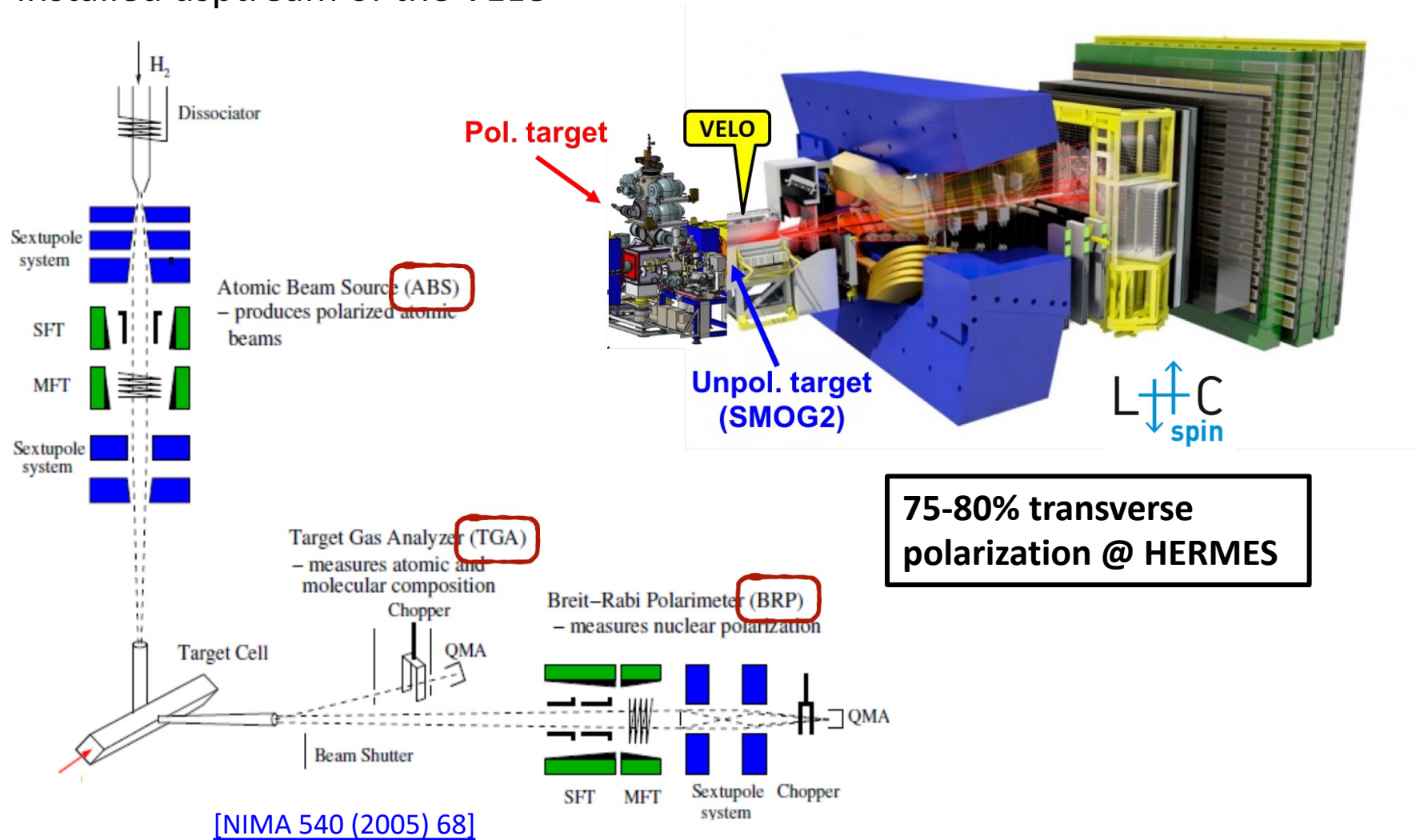
The LHCspin apparatus

The LHCspin apparatus consists of a **new-generation HERMES-like polarized gaseous fixed target** to be installed upstream of the VELO



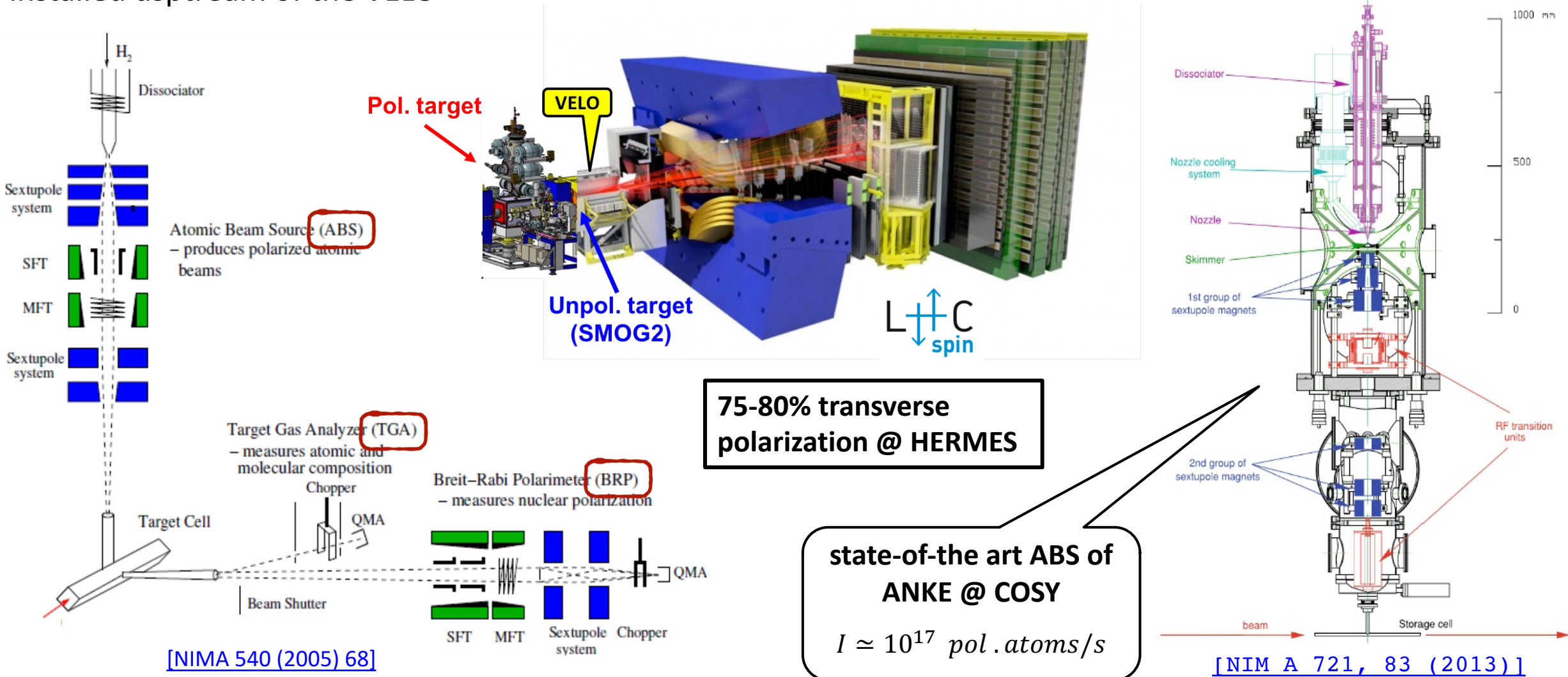
The LHCspin apparatus

The LHCspin apparatus consists of a **new-generation HERMES-like polarized gaseous fixed target** to be installed upstream of the VELO

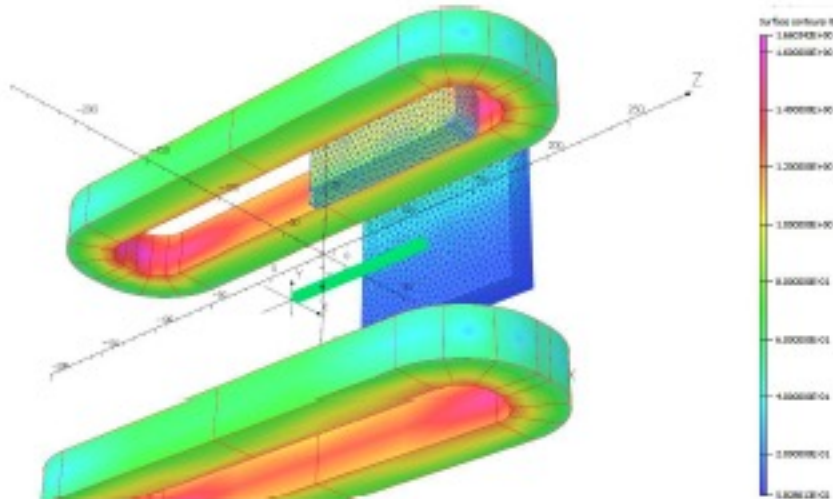
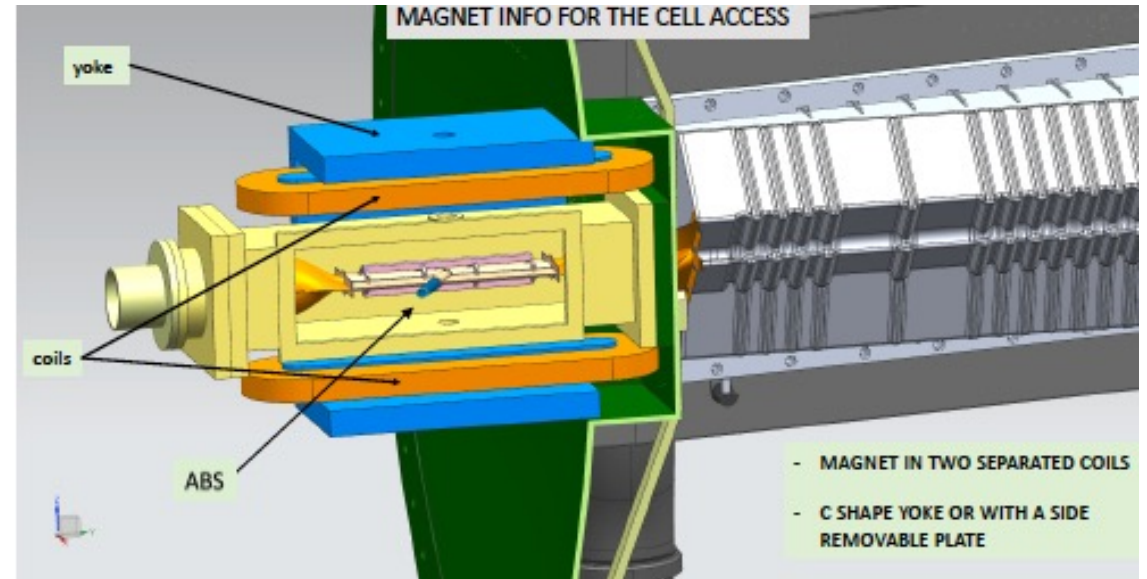
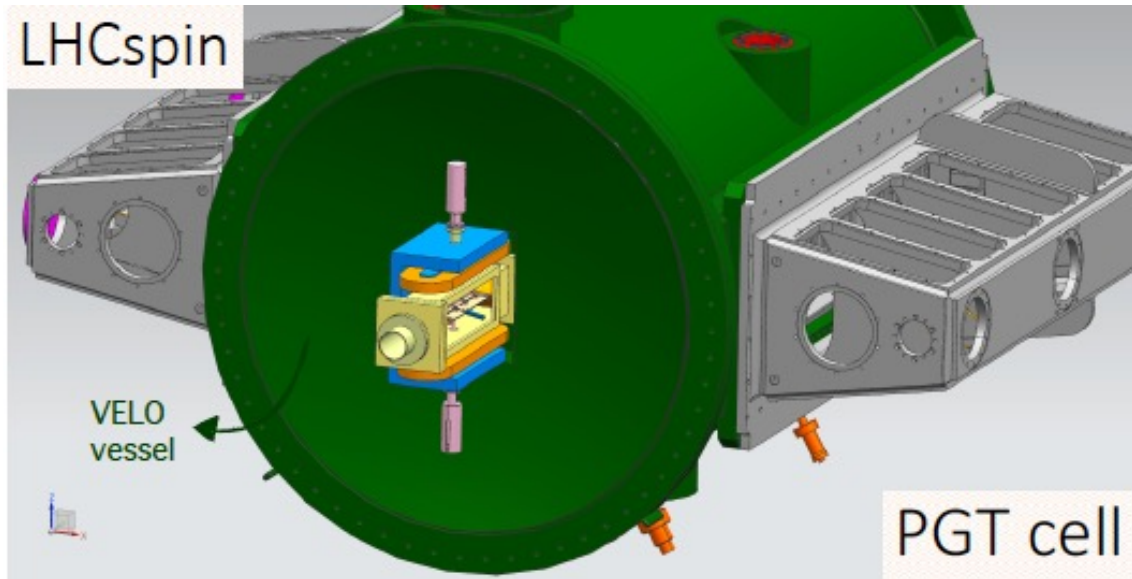


The LHCspin apparatus

The LHCspin apparatus consists of a **new-generation HERMES-like polarized gaseous fixed target** to be installed upstream of the VELO

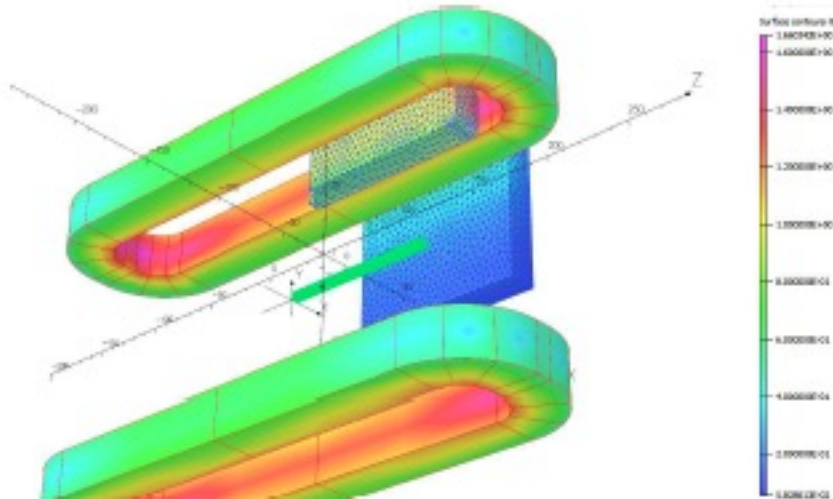
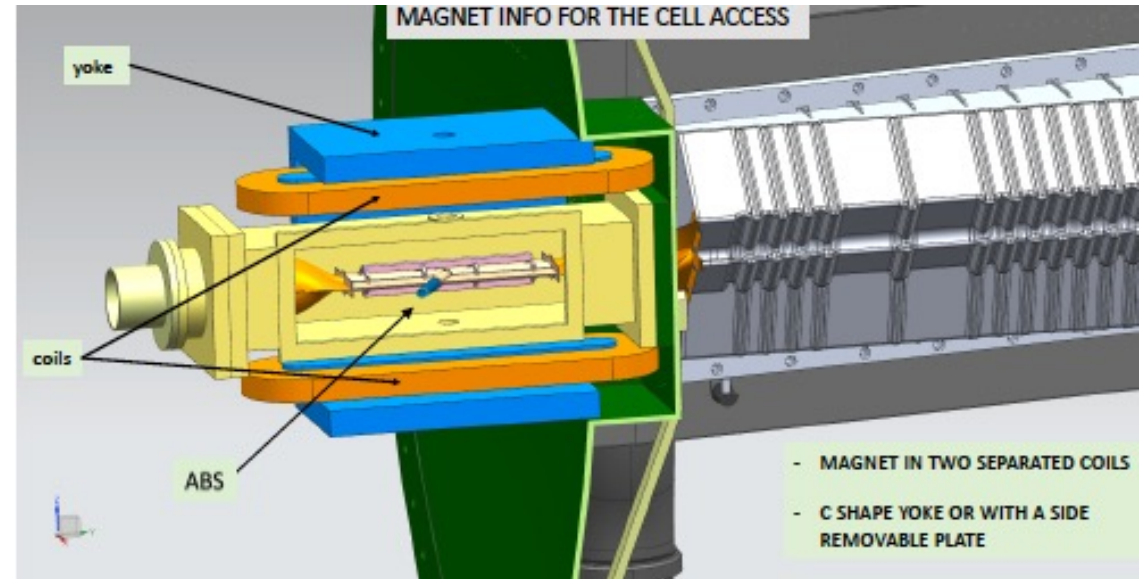
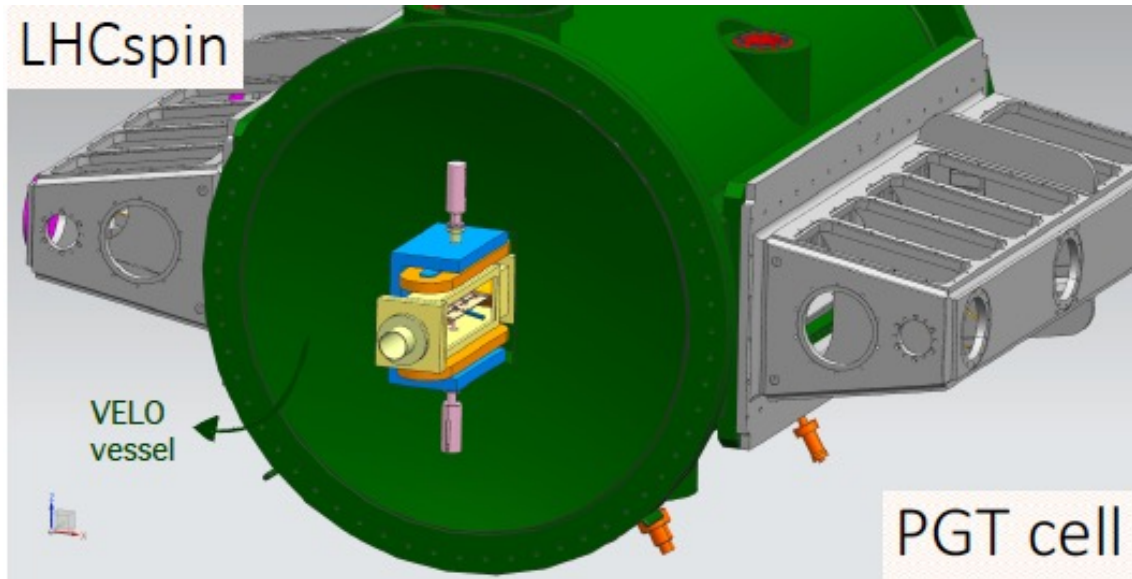


The LHCspin apparatus [\[PoS \(SPIN2018\)\]](#)



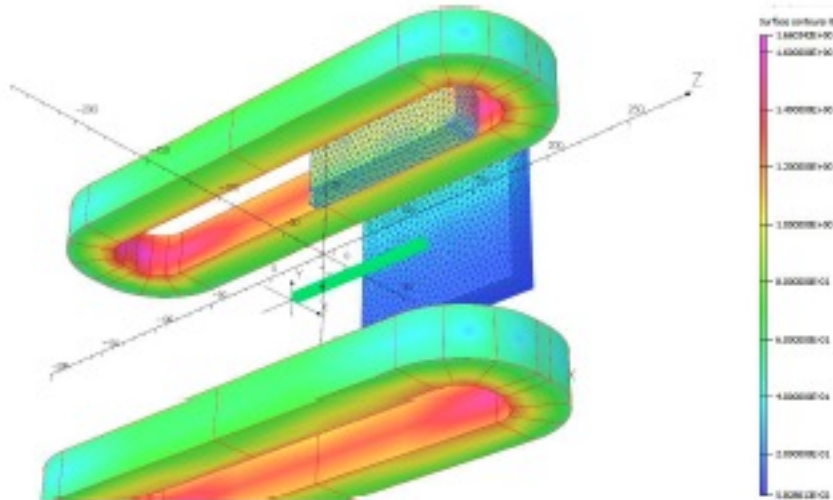
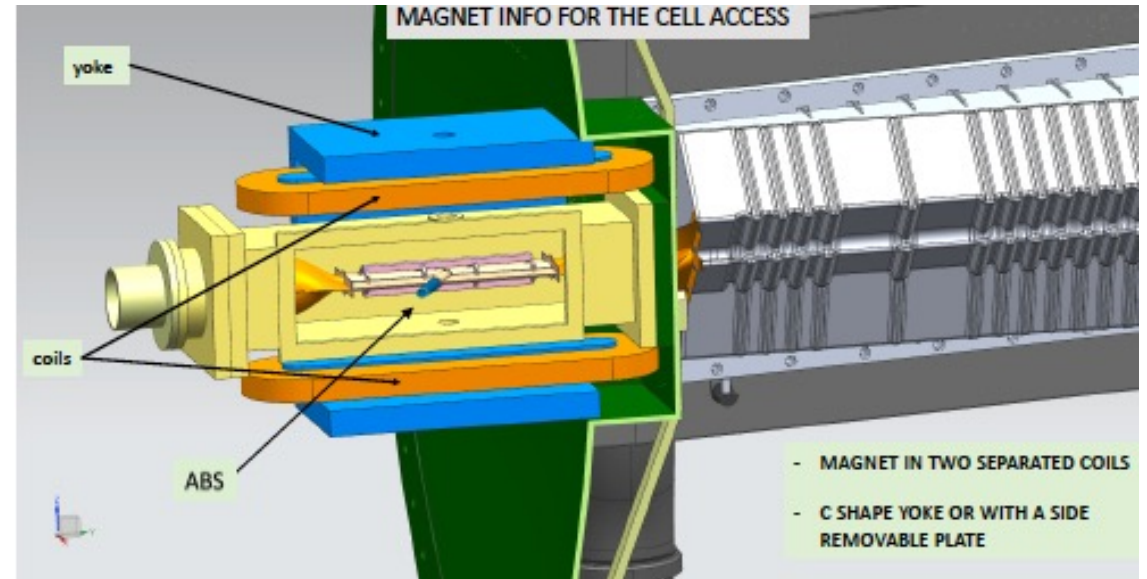
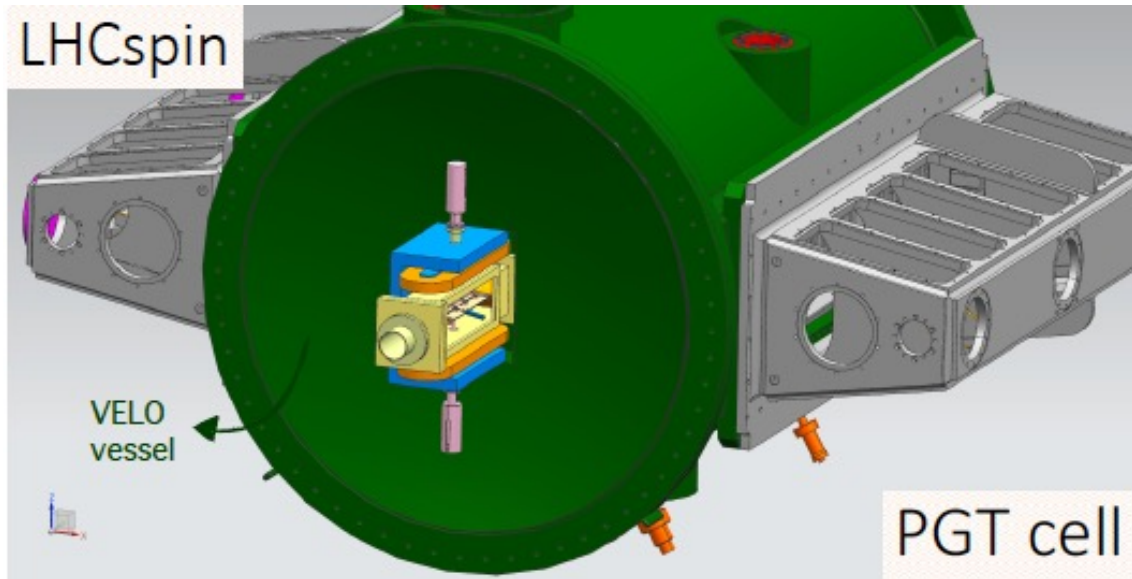
- Compact superconductive dipole magnet for static transverse field to maintain polarization inside the cell and avoid beam-induced depolarization
- Required $B = 300 \text{ mT}$ with $\Delta B/B \sim 10\%$

The LHCspin apparatus [\[PoS \(SPIN2018\)\]](#)



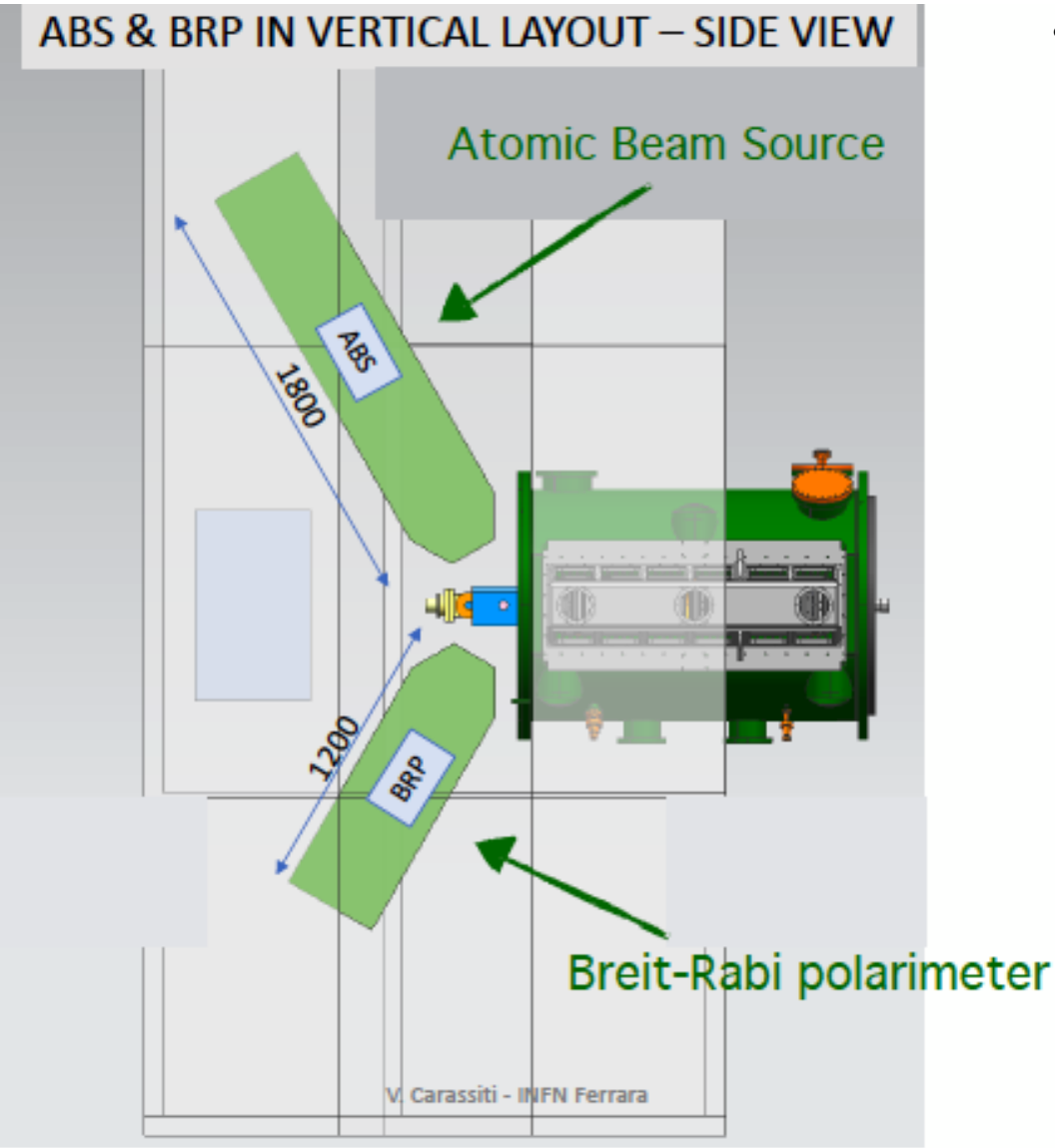
- Compact superconductive dipole magnet for static transverse field to maintain polarization inside the cell and avoid beam-induced depolarization
- Required $B = 300 \text{ mT}$ with $\Delta B/B \sim 10\%$
- **Need to modify main flange of VELO vessel (inward)**
- **No need for additional detectors**

The LHCspin apparatus [PoS (SPIN2018)]



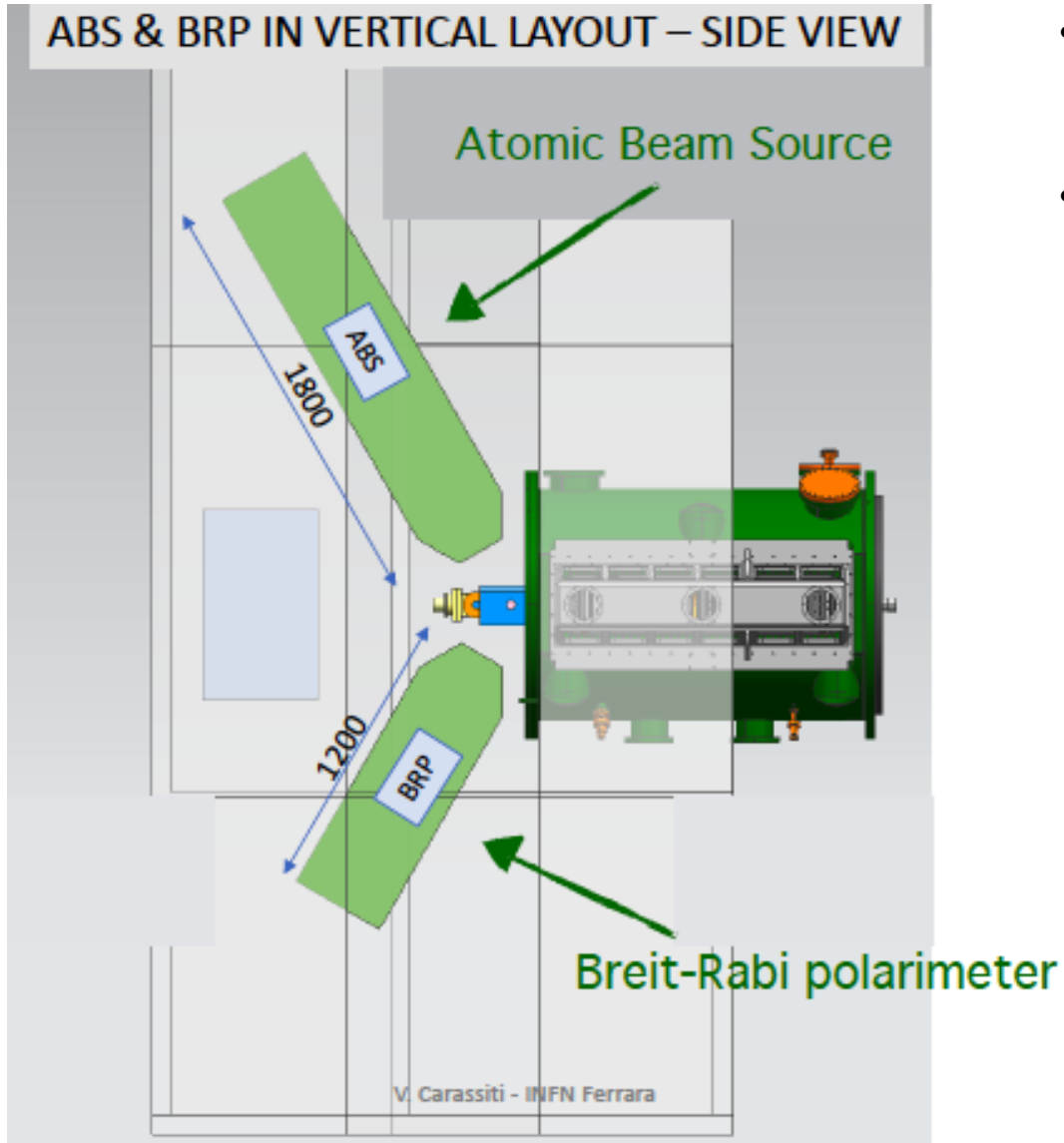
- Compact superconductive dipole magnet for static transverse field to maintain polarization inside the cell and avoid beam-induced depolarization
- Required $B = 300 \text{ mT}$ with $\Delta B/B \sim 10\%$
- **Need to modify main flange of VELO vessel (inward)**
- **No need for additional detectors**
- **Possibility to switch from dipole magnet to solenoid to realize a Longitudinal polarized target in Run5**

The LHCspin apparatus [\[PoS \(SPIN2018\)\]](#)



- Need to develop a new-generation compact ABS and diagnostic system to fit into the limited available space in the VELO alcove

The LHCspin apparatus [\[PoS \(SPIN2018\)\]](#)



- Need to develop a new-generation compact ABS and diagnostic system to fit into the limited available space in the VELO alcove
- Coating studies for the cell walls are ongoing. Crucial for target polarization (back-up slides)

The jet target hypothesis

Alternative solution with **jet target** also under evaluation:

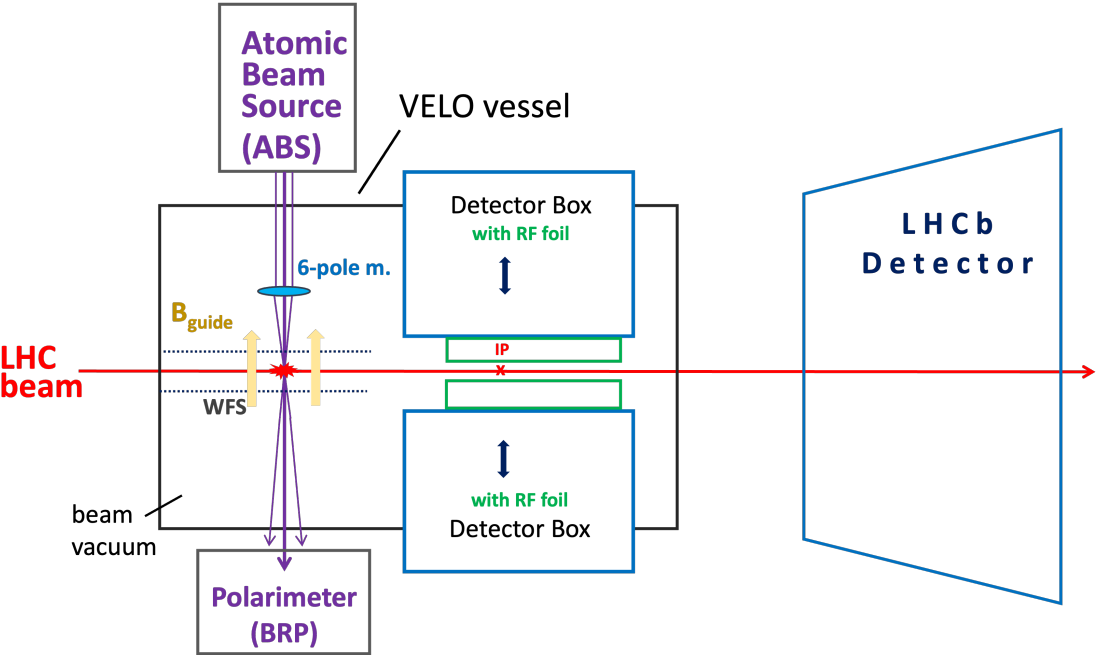
- lower density ($\sim 10^{12}$ atoms/cm²) → about a factor of 40 smaller
- higher polarization (up to 90%)
- lower systematics in P measurement

The jet target hypothesis

Alternative solution with **jet target** also under evaluation:

- lower density ($\sim 10^{12}$ atoms/cm²) \rightarrow about a factor of 40 smaller
- higher polarization (up to 90%)
- lower systematics in P measurement

Pure Jet Target

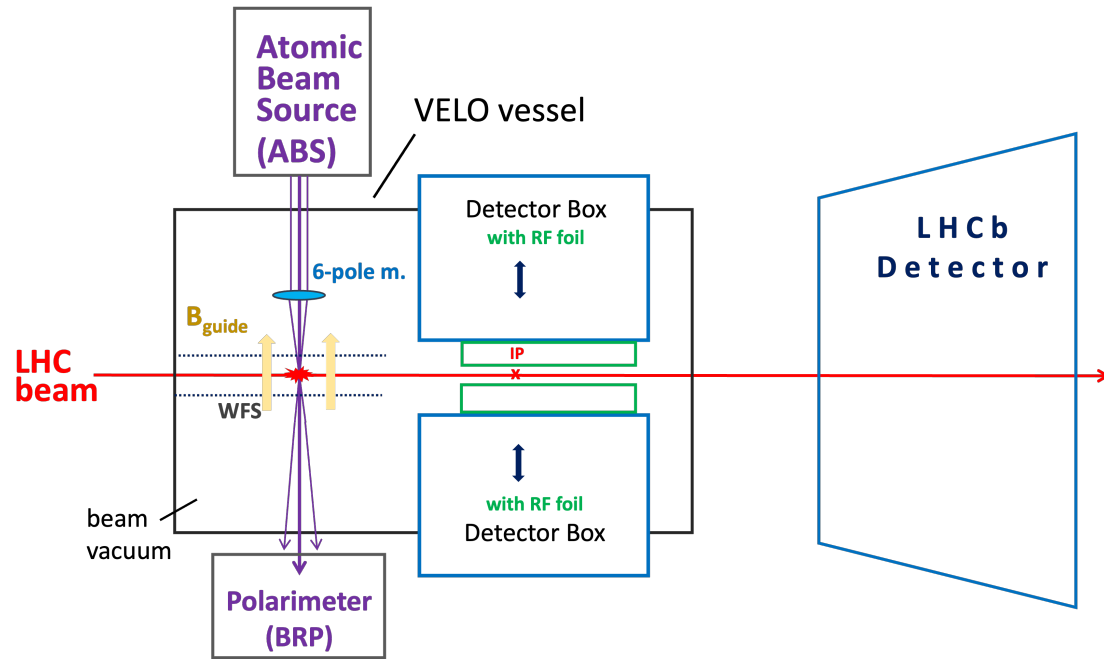


The jet target hypothesis

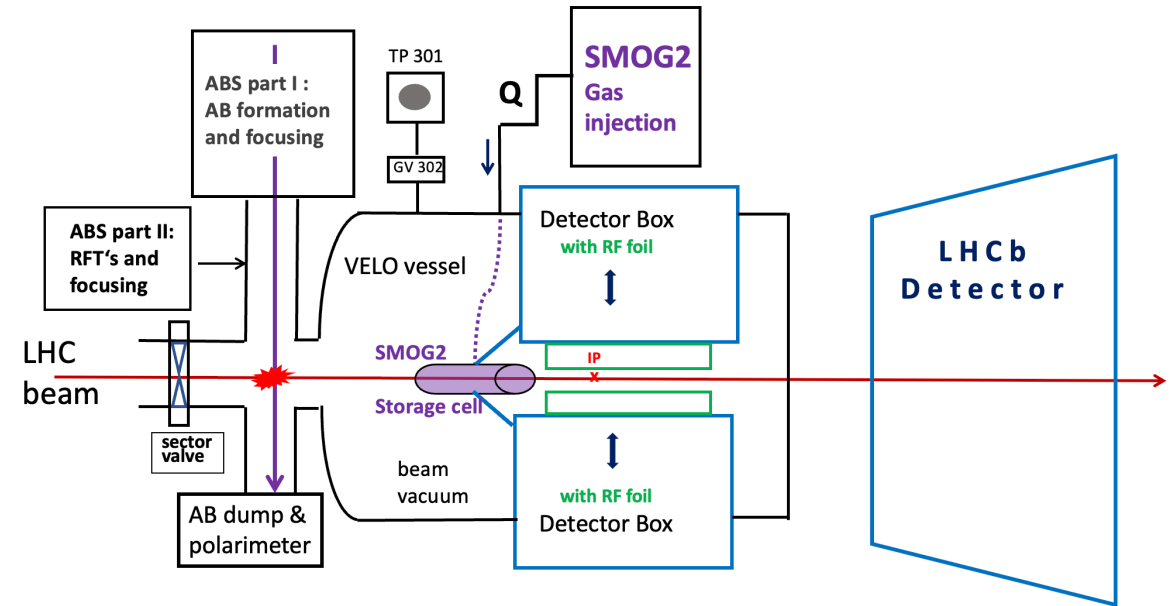
Alternative solution with **jet target** also under evaluation:

- lower density ($\sim 10^{12}$ atoms/cm²) \rightarrow about a factor of 40 smaller
- higher polarization (up to 90%)
- lower systematics in P measurement

Pure Jet Target



Jet Target + SMOG2



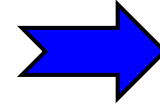
Expected performance

Target

- $I_0 = 6.5 \cdot 10^{16} \text{ s}^{-1}$ (HERMES)
- $C_{\text{tot}} = 17.4 \text{ l/s}$ (20 cm cell)
- $\theta = 3.7 \cdot 10^{13} \text{ atoms/cm}^2$

Beam

- $1.2 \cdot 10^{11} \text{ p/bunch}$ (RUN3)
- 2808 bunches
- $I_{\text{beam}} = 3.8 \cdot 10^{18} \text{ p/s}$



$$\mathcal{L}_{pH}(300 \text{ K}) \approx 1.4 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$$

$$L_{pH}(\text{Run 4}) \approx 5 \text{ fb}^{-1}$$

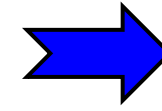
Expected performance

Target

- $I_0 = 6.5 \cdot 10^{16} s^{-1}$ (HERMES)
- $C_{tot} = 17.4$ l/s (20 cm cell)
- $\theta = 3.7 \cdot 10^{13}$ atoms/cm²

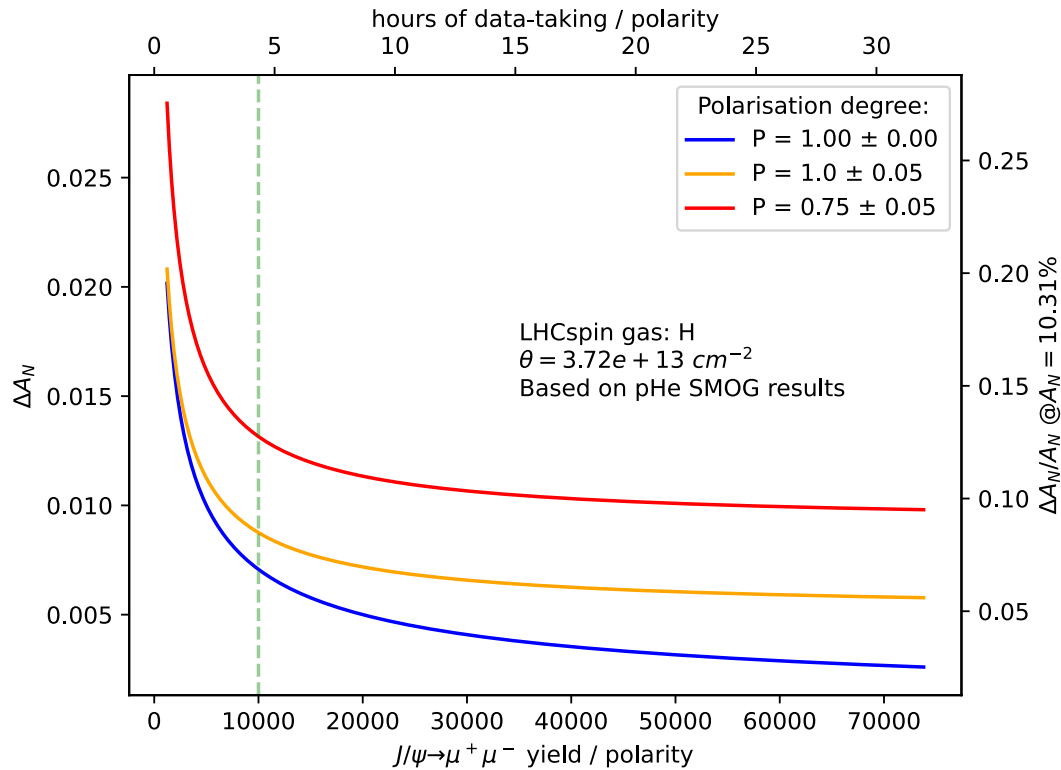
Beam

- $1.2 \cdot 10^{11}$ p/bunch (RUN3)
- 2808 bunches
- $I_{beam} = 3.8 \cdot 10^{18}$ p/s



$$\mathcal{L}_{pH}(300\text{ K}) \approx 1.4 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$$

$$L_{pH}(\text{Run 4}) \approx 5 \text{ fb}^{-1}$$



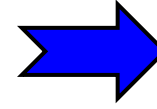
Expected performance

Target

- $I_0 = 6.5 \cdot 10^{16} \text{ s}^{-1}$ (HERMES)
- $C_{\text{tot}} = 17.4 \text{ l/s}$ (20 cm cell)
- $\theta = 3.7 \cdot 10^{13} \text{ atoms/cm}^2$

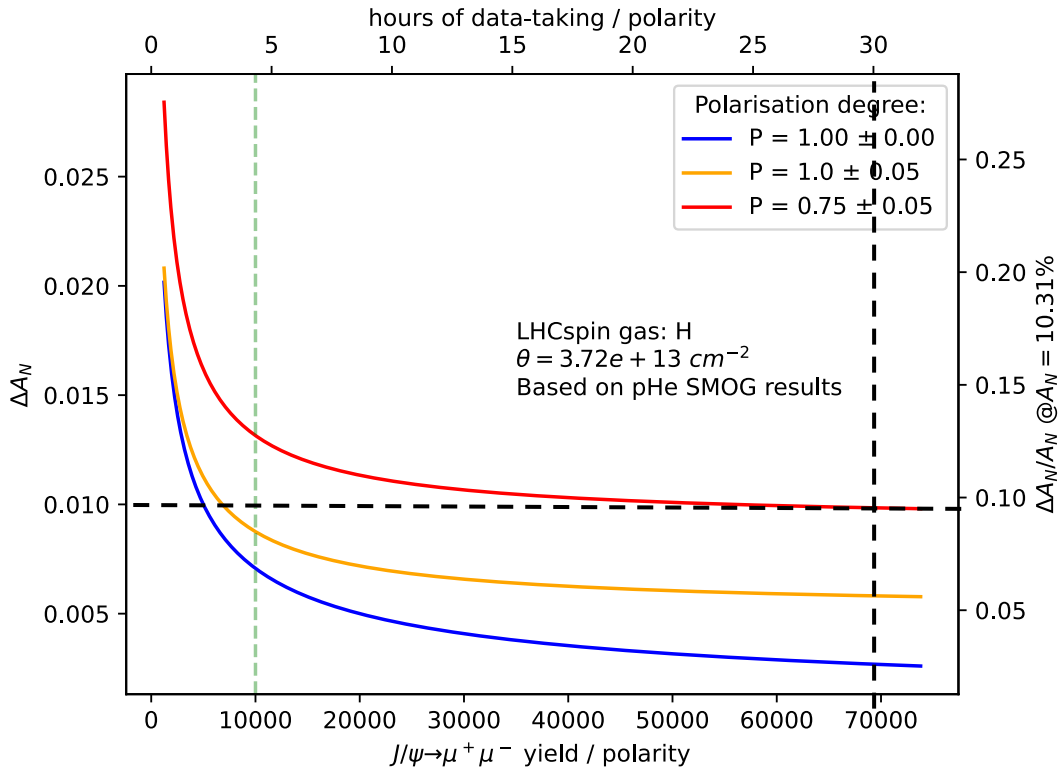
Beam

- $1.2 \cdot 10^{11} \text{ p/bunch}$ (RUN3)
- 2808 bunches
- $I_{\text{beam}} = 3.8 \cdot 10^{18} \text{ p/s}$



$$\mathcal{L}_{pH}(300 \text{ K}) \approx 1.4 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$$

$$L_{pH}(\text{Run 4}) \approx 5 \text{ fb}^{-1}$$



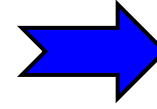
Expected performance

Target

- $I_0 = 6.5 \cdot 10^{16} \text{ s}^{-1}$ (HERMES)
- $C_{\text{tot}} = 17.4 \text{ l/s}$ (20 cm cell)
- $\theta = 3.7 \cdot 10^{13} \text{ atoms/cm}^2$

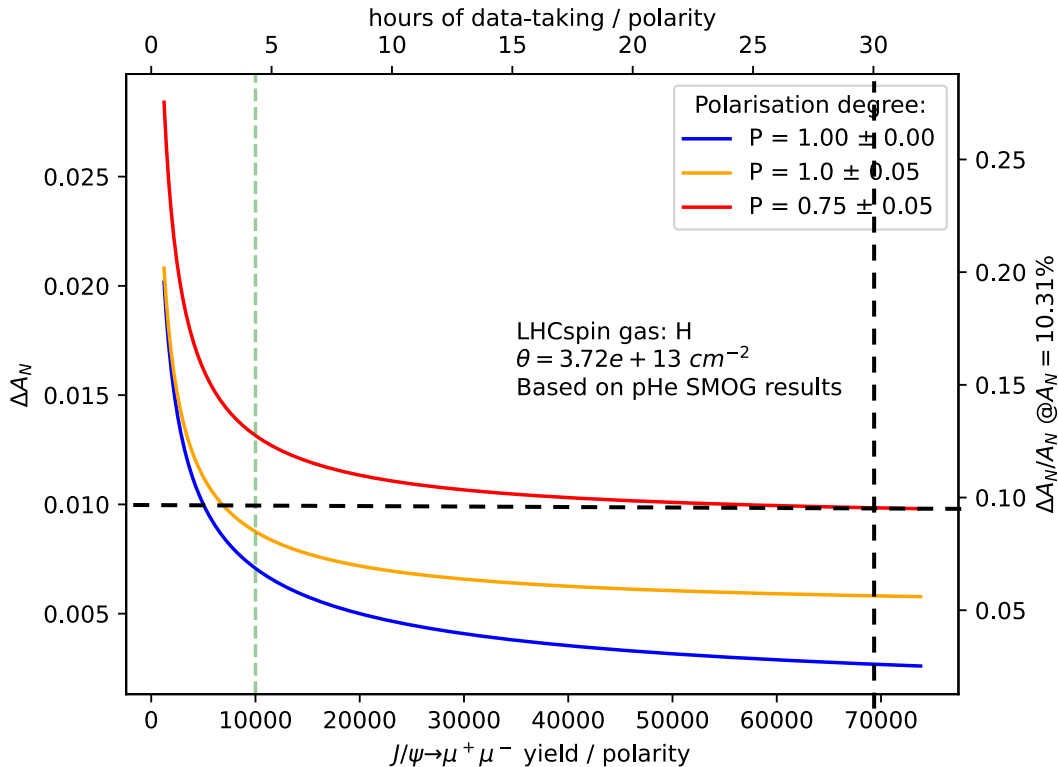
Beam

- $1.2 \cdot 10^{11} \text{ p/bunch}$ (RUN3)
- 2808 bunches
- $I_{\text{beam}} = 3.8 \cdot 10^{18} \text{ p/s}$



$$\mathcal{L}_{pH}(300 \text{ K}) \approx 1.4 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$$

$$L_{pH}(\text{Run 4}) \approx 5 \text{ fb}^{-1}$$



Expected yields for Run4 (Run4+Run5):

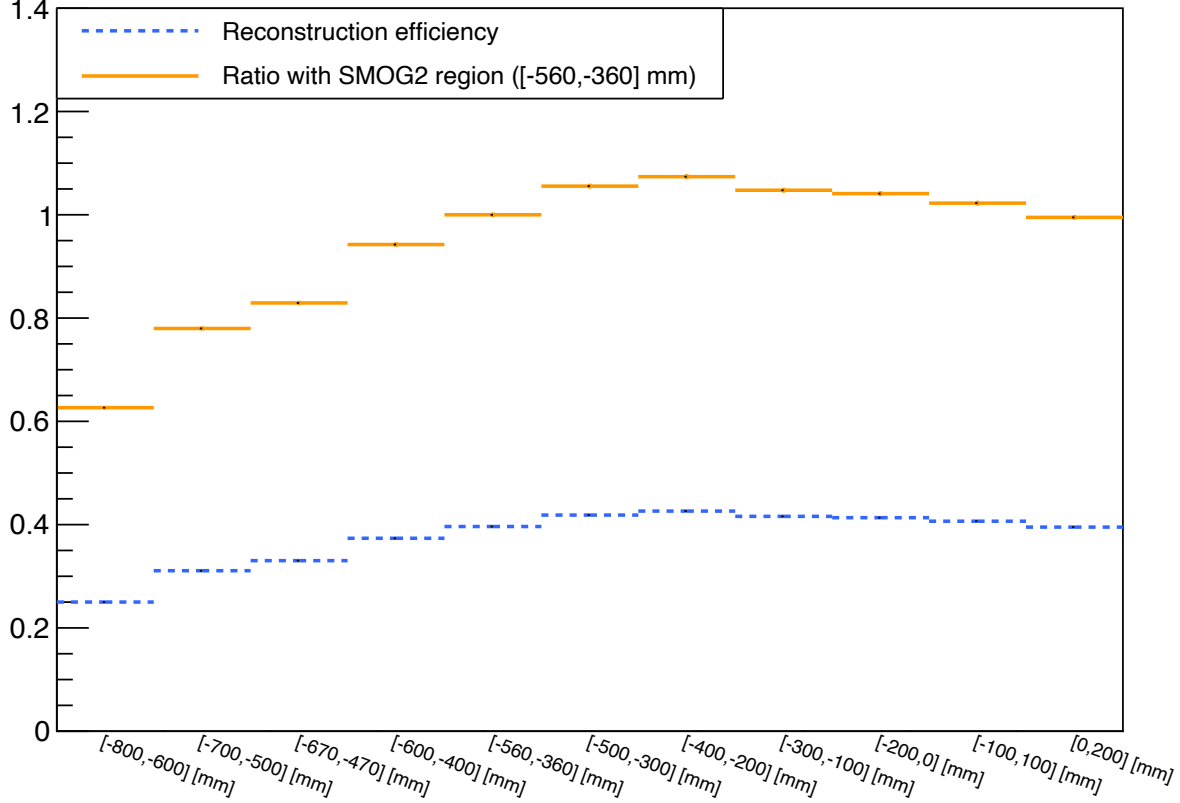
Channel	Events / week	Total events
$J/\psi \rightarrow \mu^+ \mu^-$	194k (434k)	23M (75M)
$\psi(2S) \rightarrow \mu^+ \mu^-$	3.5k (7.7k)	414k (1.3M)
$D^0 \rightarrow K^- \pi^+$	976k (2.2M)	117M (380M)
$J/\psi J/\psi \rightarrow \mu^+ \mu^- \mu^+ \mu^-$	77 (170)	930 (3000)
Drell Yan ($5 < M_{\mu\mu} < 9 \text{ GeV}$)	110 (250)	13k (43k)
$\Upsilon \rightarrow \mu^+ \mu^-$	83 (187)	10k (32k)
$\Lambda_c^+ \rightarrow p K^- \pi^+$	19k (43k)	2.3M (7.5M)

assumptions:

- 120 weeks/RUN
- 84h/week
- $Stat(\text{Run5}) \sim \sqrt{5} Stat(\text{Run4})$

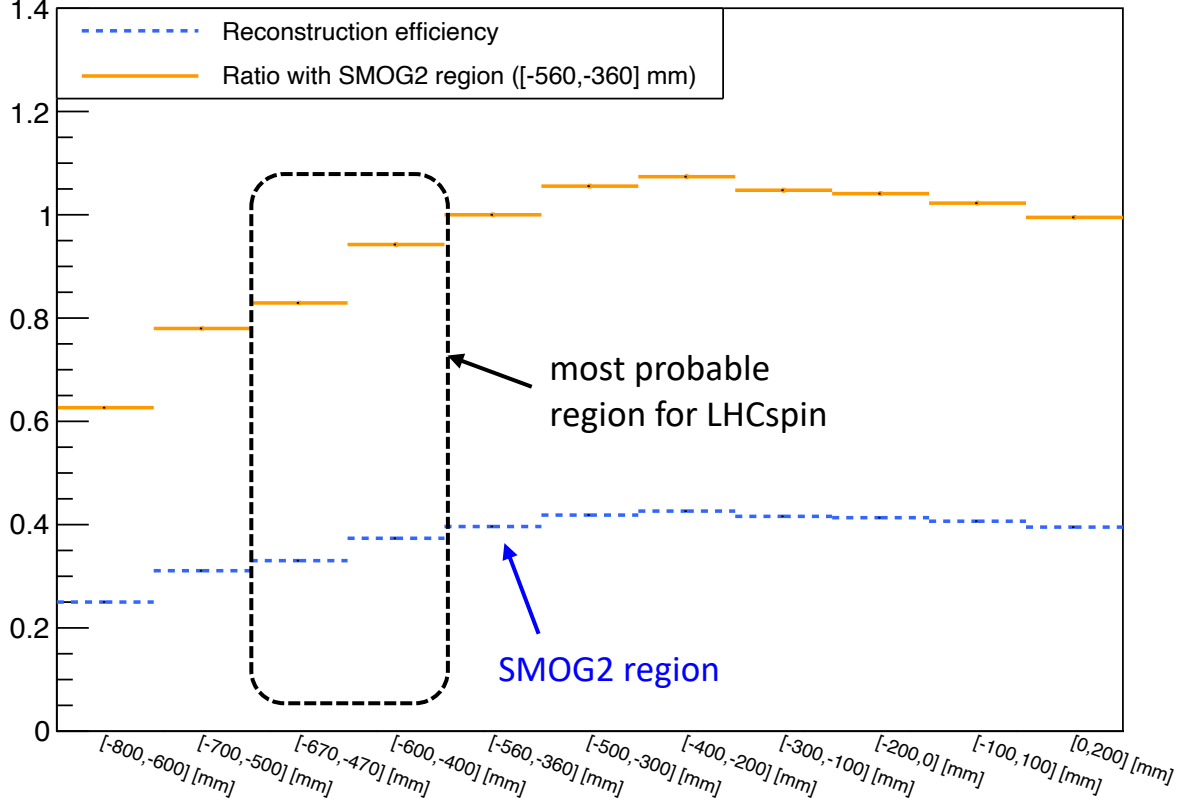
Reconstruction efficiencies

$J/\Psi \rightarrow \mu^+\mu^- \in_{rec}(PV)$ vs cell position



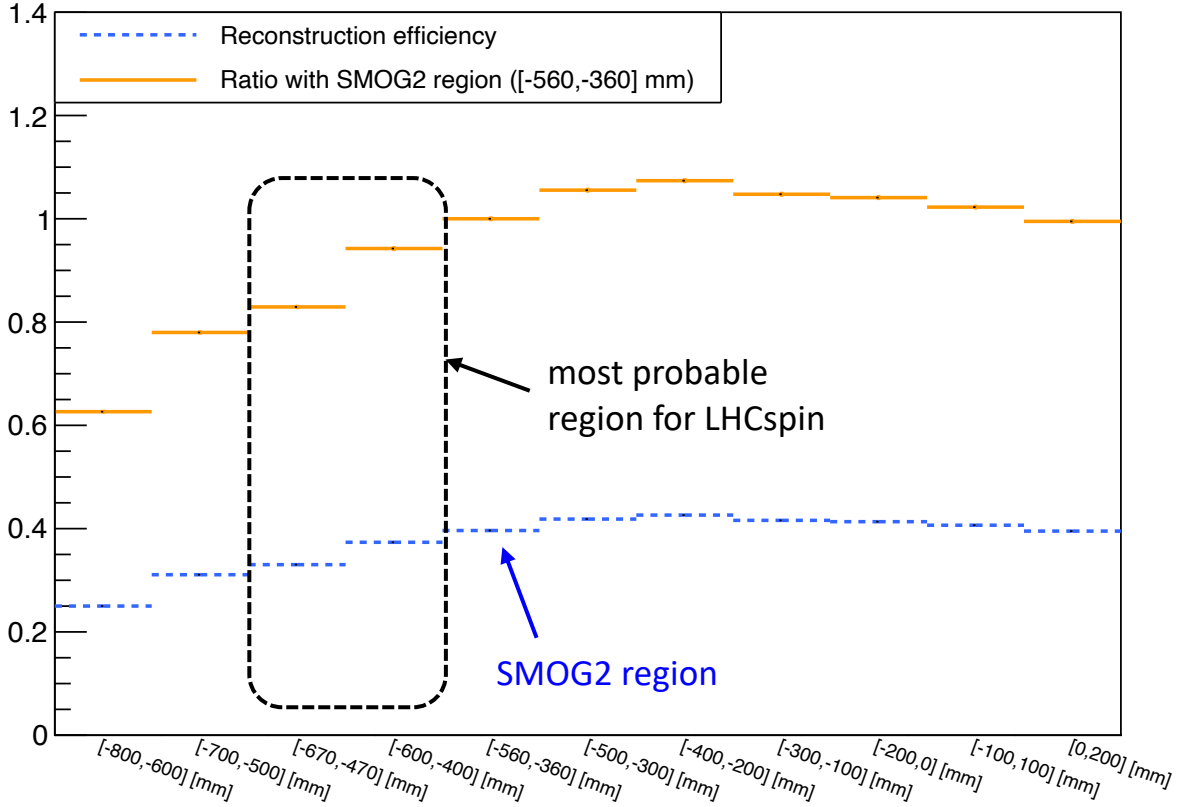
Reconstruction efficiencies

$J/\Psi \rightarrow \mu^+\mu^- \in_{rec}(PV)$ vs cell position

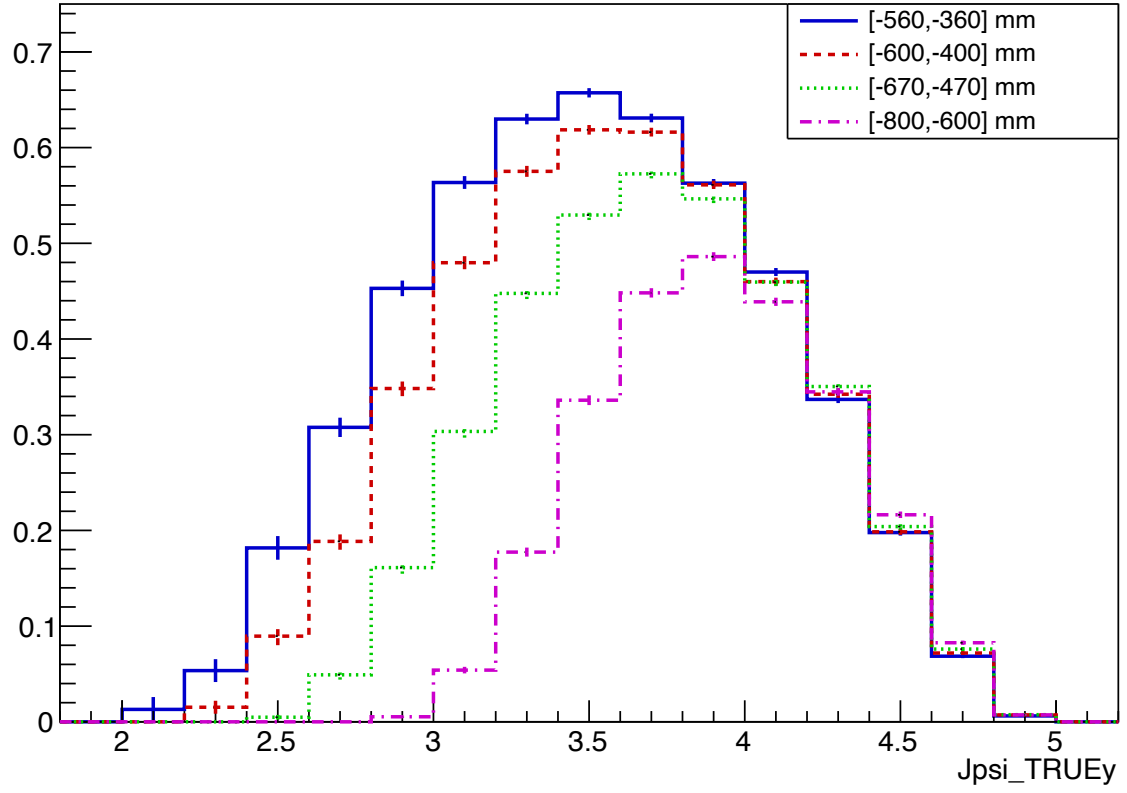


Reconstruction efficiencies

$J/\Psi \rightarrow \mu^+\mu^- \in_{rec}(PV)$ vs cell position

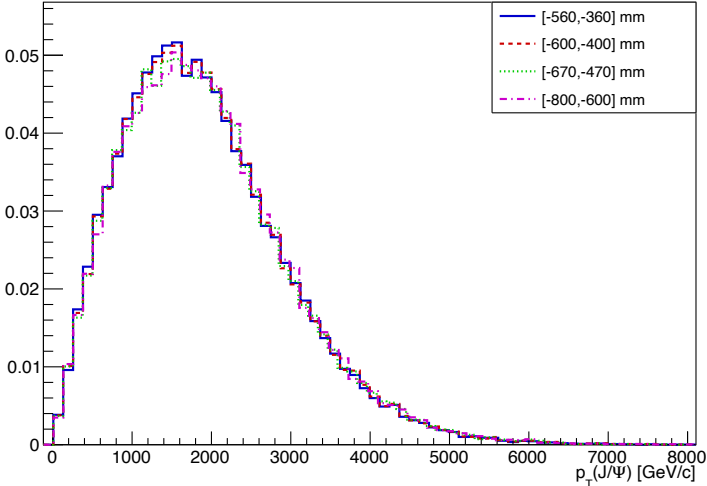


$J/\Psi \rightarrow \mu^+\mu^-$ PV X track reconstruction efficiency

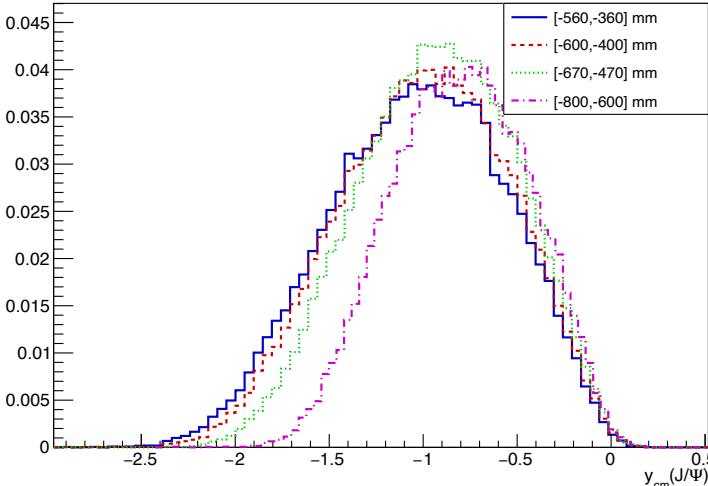


Kinematic coverage

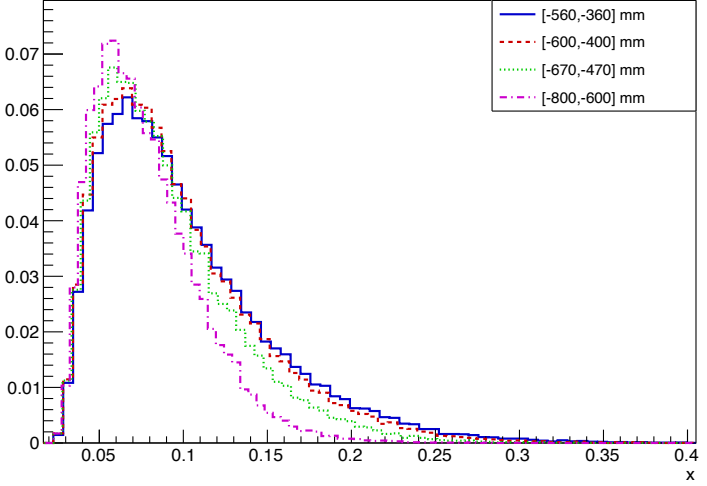
$J/\Psi \rightarrow \mu^+\mu^-$



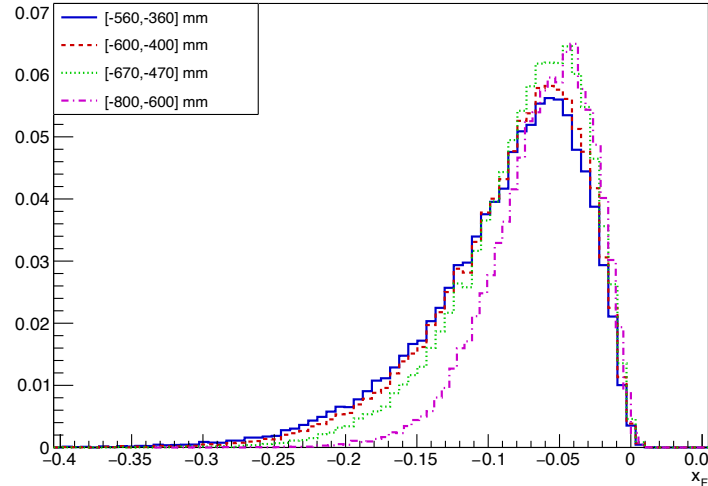
$J/\Psi \rightarrow \mu^+\mu^-$



$J/\Psi \rightarrow \mu^+\mu^-$

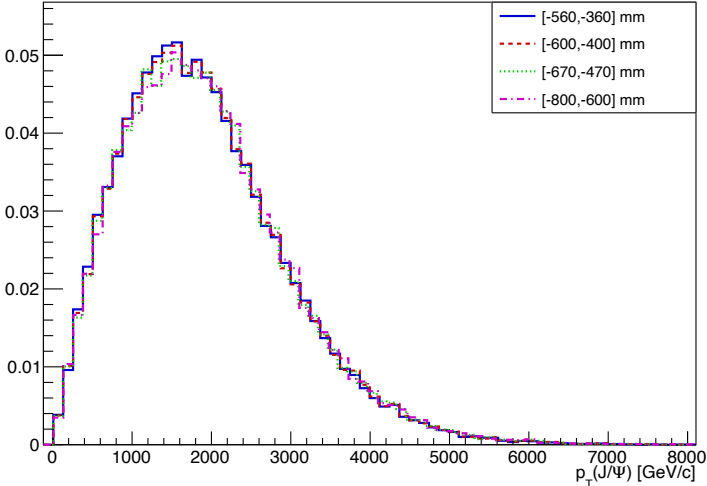


$J/\Psi \rightarrow \mu^+\mu^-$

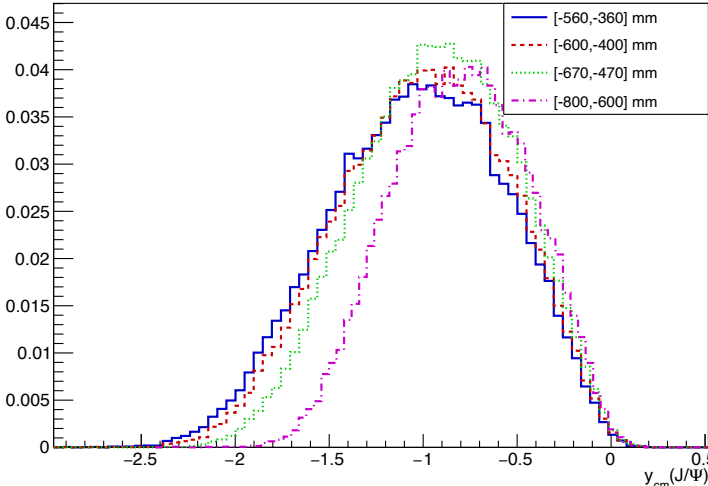


Kinematic coverage

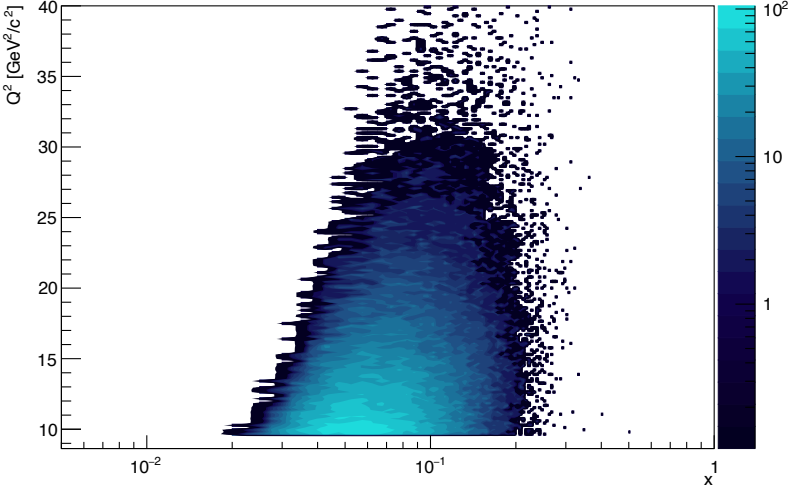
$J/\Psi \rightarrow \mu^+\mu^-$



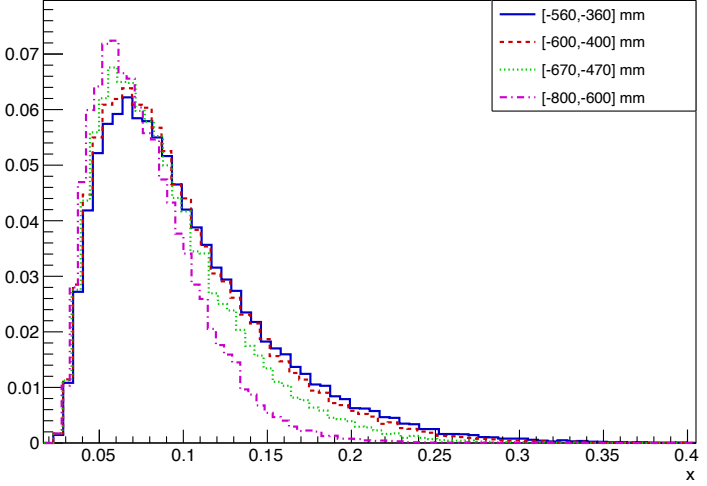
$J/\Psi \rightarrow \mu^+\mu^-$



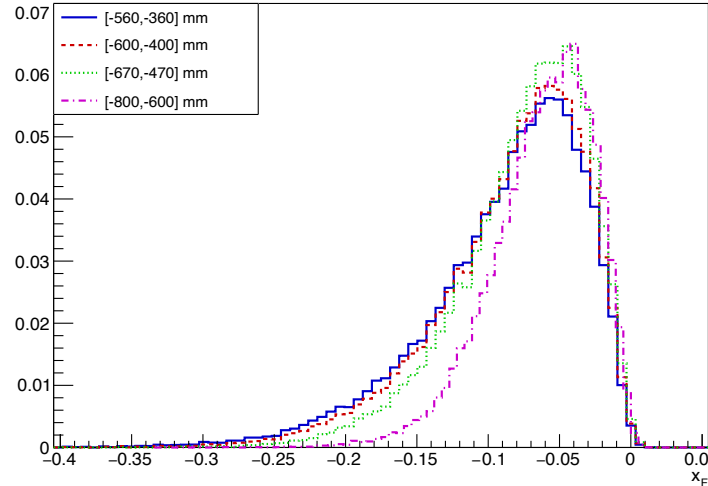
$J/\Psi \rightarrow \mu^+\mu^- [-670,-470] \text{ mm}$



$J/\Psi \rightarrow \mu^+\mu^-$

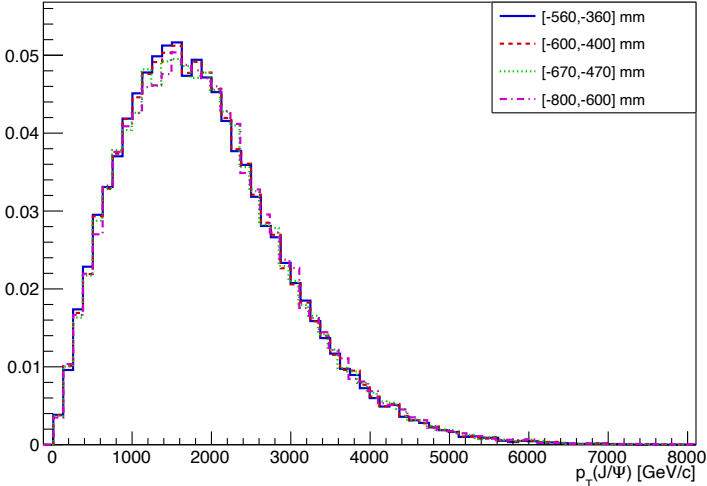


$J/\Psi \rightarrow \mu^+\mu^-$

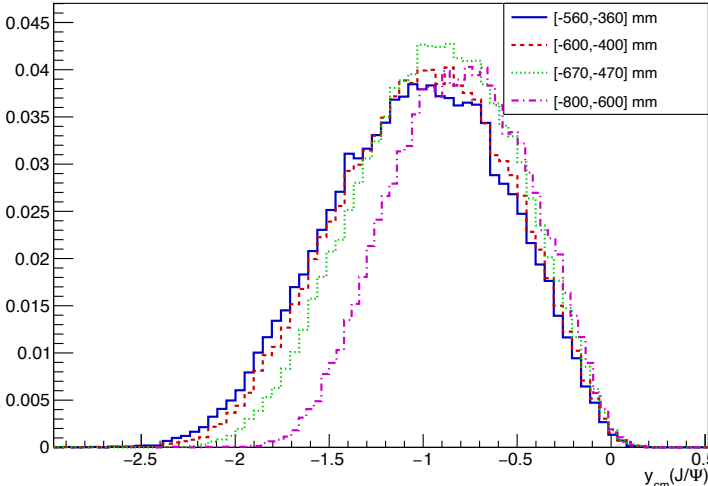


Kinematic coverage

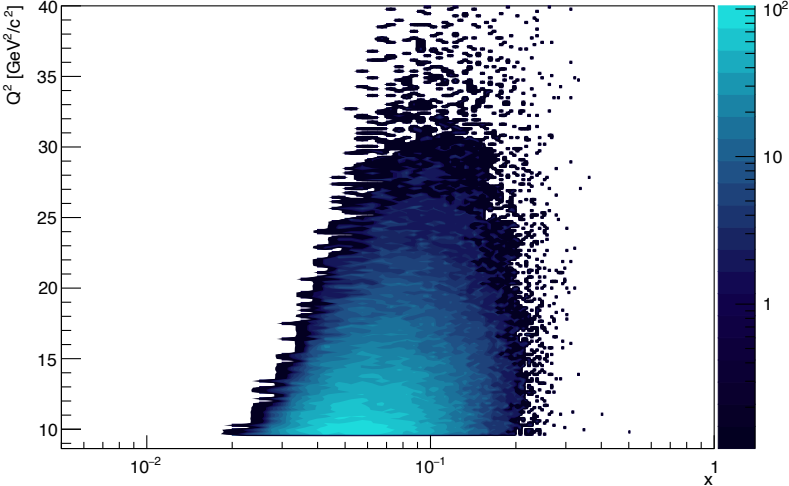
$J/\Psi \rightarrow \mu^+\mu^-$



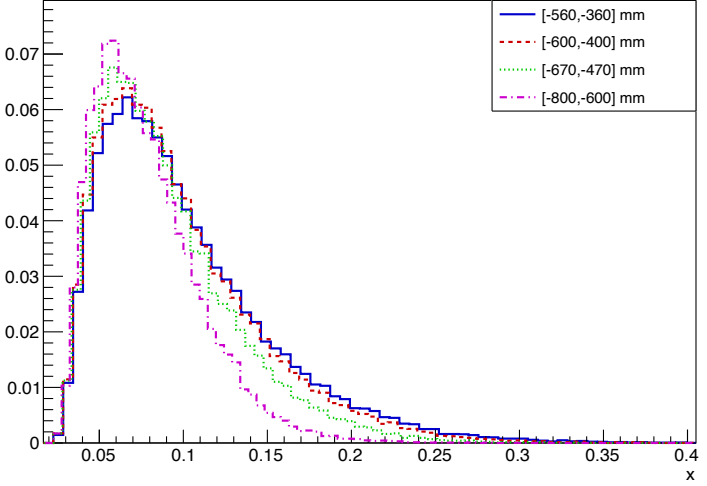
$J/\Psi \rightarrow \mu^+\mu^-$



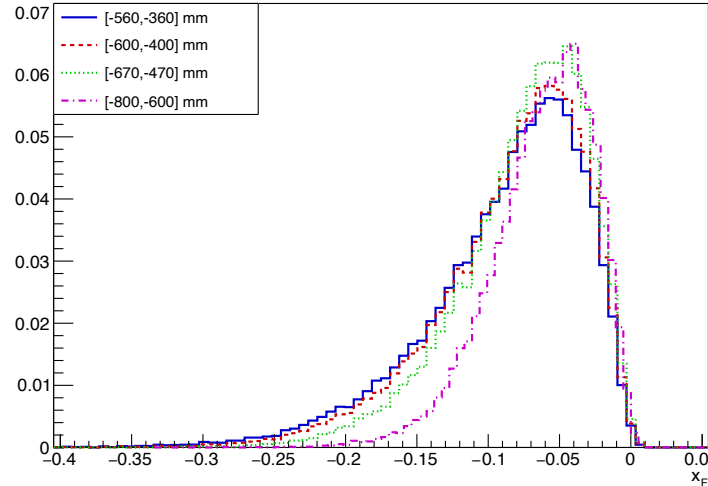
$J/\Psi \rightarrow \mu^+\mu^- [-670,-470] \text{ mm}$



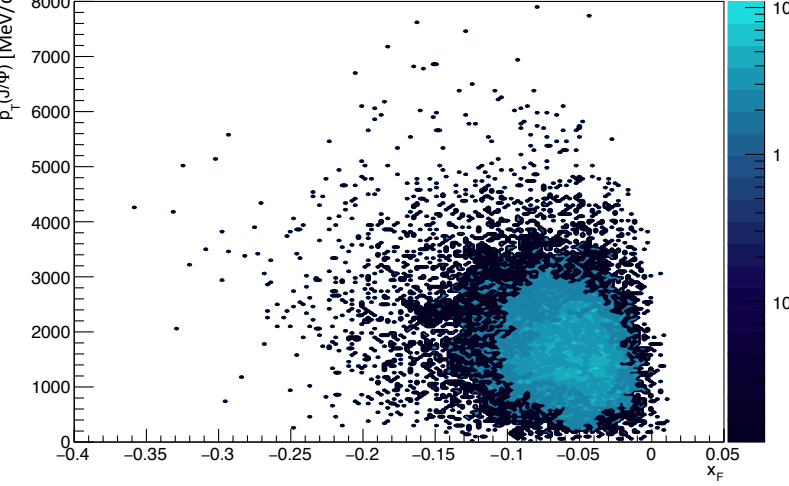
$J/\Psi \rightarrow \mu^+\mu^-$



$J/\Psi \rightarrow \mu^+\mu^-$



$J/\Psi \rightarrow \mu^+\mu^- [-670,-470] \text{ mm}$



A preliminary analysis tool for pseudo-data

A pseudo-data set based on a transversely Pol. H target has been generated to study the interplay between statistical and systematic (due to the measurement of the polarization) uncertainties.

A preliminary analysis tool for pseudo-data

A pseudo-data set based on a transversely Pol. H target has been generated to study the interplay between statistical and systematic (due to the measurement of the polarization) uncertainties.

Similar approach used at HERMES (Appendix C of [[JHEP, 12:010, 2020](#)]):

- Use official LHCb MC data for inclusive production of $J/\psi \rightarrow \mu^+ \mu^-$ in fixed-target configuration (PYTHIA8 + EPOS)
- Assign to each simulated event a target polarization state (\uparrow or \downarrow) using a random extraction modulated with a model for the cross section (in this way we introduce a spin-dependence in the simulation)
- The model assumes a dominant $\sin \phi$ modulation (e.g. sensitive to the gluon Sivers) plus a suppressed $\sin 2\phi$ modulation (to account e.g. for possible higher-twist contributions). Both terms depend mildly on the kinematics (x, p_T):

$$\rho = \frac{1}{2} \left[1 + \left(a_1 + a_2 \frac{x - \bar{x}}{x_{max}} + a_3 \frac{p_T - \bar{p}_T}{p_{T \ max}} \right) \sin \phi + \left(b_1 + b_2 \frac{x - \bar{x}}{x_{max}} + b_3 \frac{p_T - \bar{p}_T}{p_{T \ max}} \right) \sin 2\phi \right]$$

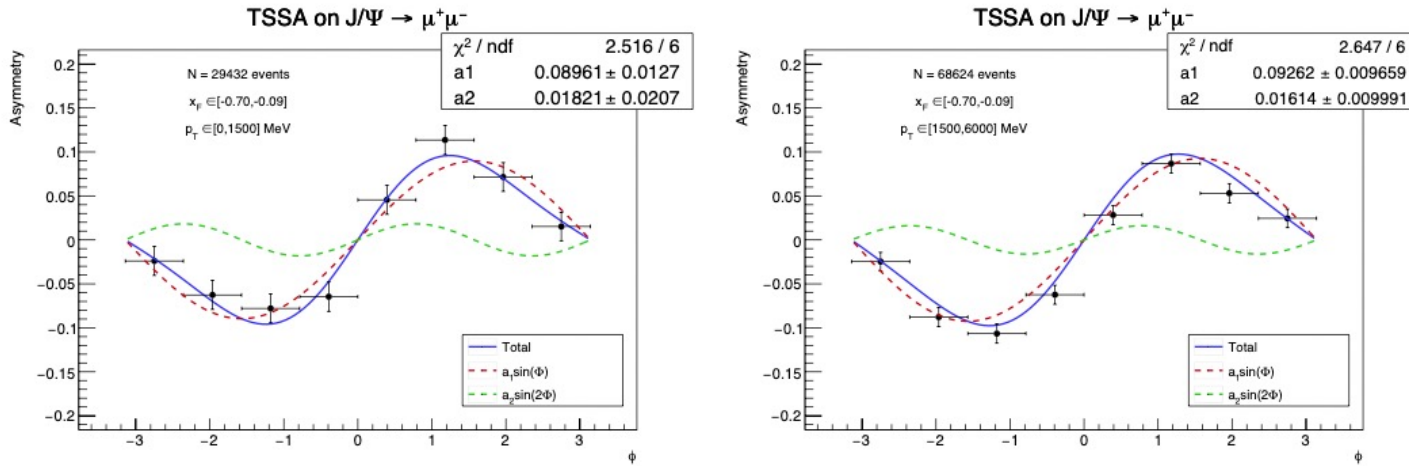
- Using these pseudo-data the TSSA is computed in the usual way:

$$A_N = \frac{1}{P} \frac{N^\uparrow - N^\downarrow}{N^\uparrow + N^\downarrow}$$

and the uncertainties on $N^{\uparrow(\downarrow)}$ (Poisson) and P (systematic) propagated accordingly.

A preliminary analysis tool for pseudo-data

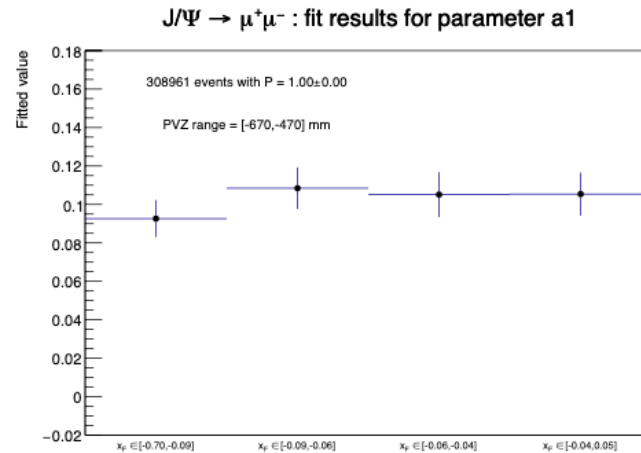
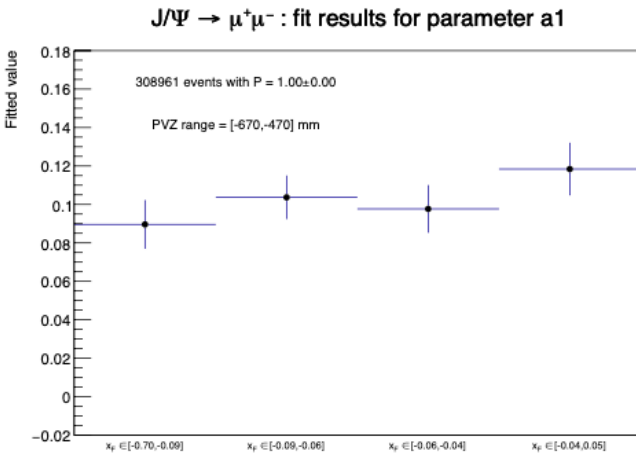
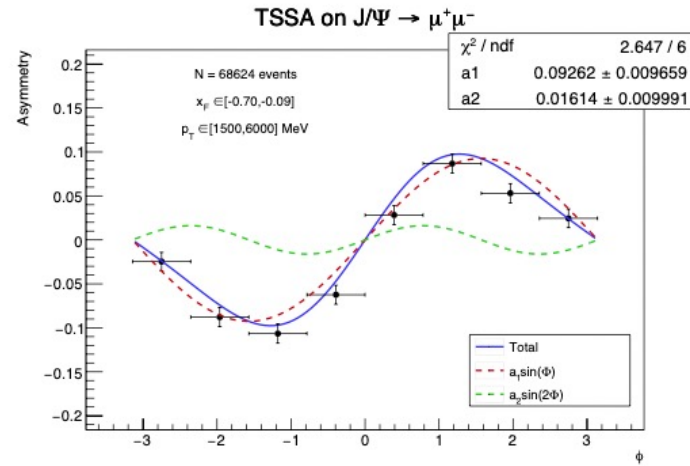
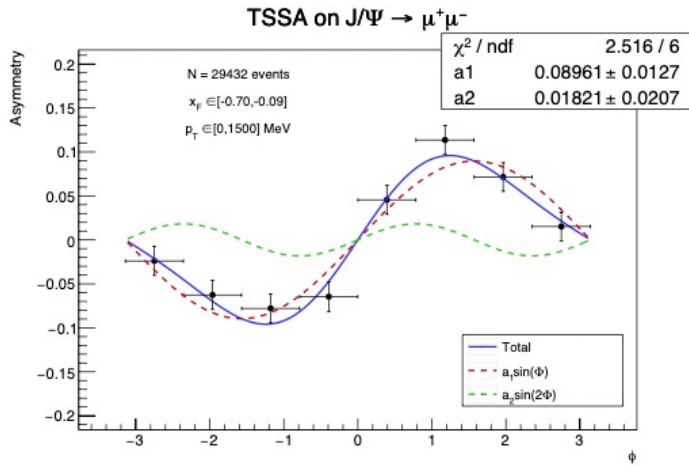
- The data points are binned of x_F and p_T (2D binning), represented vs. ϕ and fitted with $f = a_1 \sin \phi + a_2 \sin 2\phi$ where the free parameters a_1 and a_2 represent the amplitude of the corresponding azimuthal modulation



- The extracted parameters a_1 and a_2 are consistent with those used to generate the model (no bias is observed)
- With the available MC statistics (corresponding to 2 weeks of data-taking) there is no sensitivity for the $\sin 2\phi$ term

A preliminary analysis tool for pseudo-data

- The data points are binned of x_F and p_T (2D binning), represented vs. ϕ and fitted with $f = a_1 \sin \phi + a_2 \sin 2\phi$ where the free parameters a_1 and a_2 represent the amplitude of the corresponding azimuthal modulation



- The extracted parameters a_1 and a_2 are consistent with those used to generate the model (no bias is observed)
- With the available MC statistics (corresponding to 2 weeks of data-taking) there is no sensitivity for the $\sin 2\phi$ term
- The amplitudes a_1 are the reported vs. x_F in bins of p_T (and vice-versa)
- A mild kinematic dependence is observed consistent with the model

Statistical vs Systematics uncertainties

- The analysis tool described above allows to study the interplay between statistical uncertainties and systematic uncertainties (due to the measurement of the polarization) under different data-taking scenarios

p_T (MeV)	x_F	a_1 ($\Delta P = 0\%$)	a_1 ($\Delta P = 5\%$)	a_1 ($\Delta P = 20\%$)	a_1 ($\Delta P = 50\%$)
[0,1500]	[-0.70,-0.09]	0.090 ± 0.013	0.089 ± 0.013	0.087 ± 0.014	0.087 ± 0.022
[0,1500]	[-0.09,-0.06]	0.104 ± 0.011	0.104 ± 0.012	0.103 ± 0.016	0.100 ± 0.027
[0,1500]	[-0.06,-0.04]	0.098 ± 0.012	0.098 ± 0.013	0.097 ± 0.016	0.094 ± 0.027
[0,1500]	[-0.04,0.05]	0.118 ± 0.014	0.117 ± 0.014	0.114 ± 0.017	0.113 ± 0.030
[1500,6000]	[-0.70,-0.09]	0.093 ± 0.010	0.092 ± 0.010	0.090 ± 0.013	0.089 ± 0.023
[1500,6000]	[-0.09,-0.06]	0.108 ± 0.011	0.108 ± 0.011	0.108 ± 0.015	0.107 ± 0.027
[1500,6000]	[-0.06,-0.04]	0.105 ± 0.012	0.105 ± 0.012	0.104 ± 0.015	0.103 ± 0.026
[1500,6000]	[-0.04,0.05]	0.105 ± 0.011	0.105 ± 0.012	0.102 ± 0.015	0.102 ± 0.026

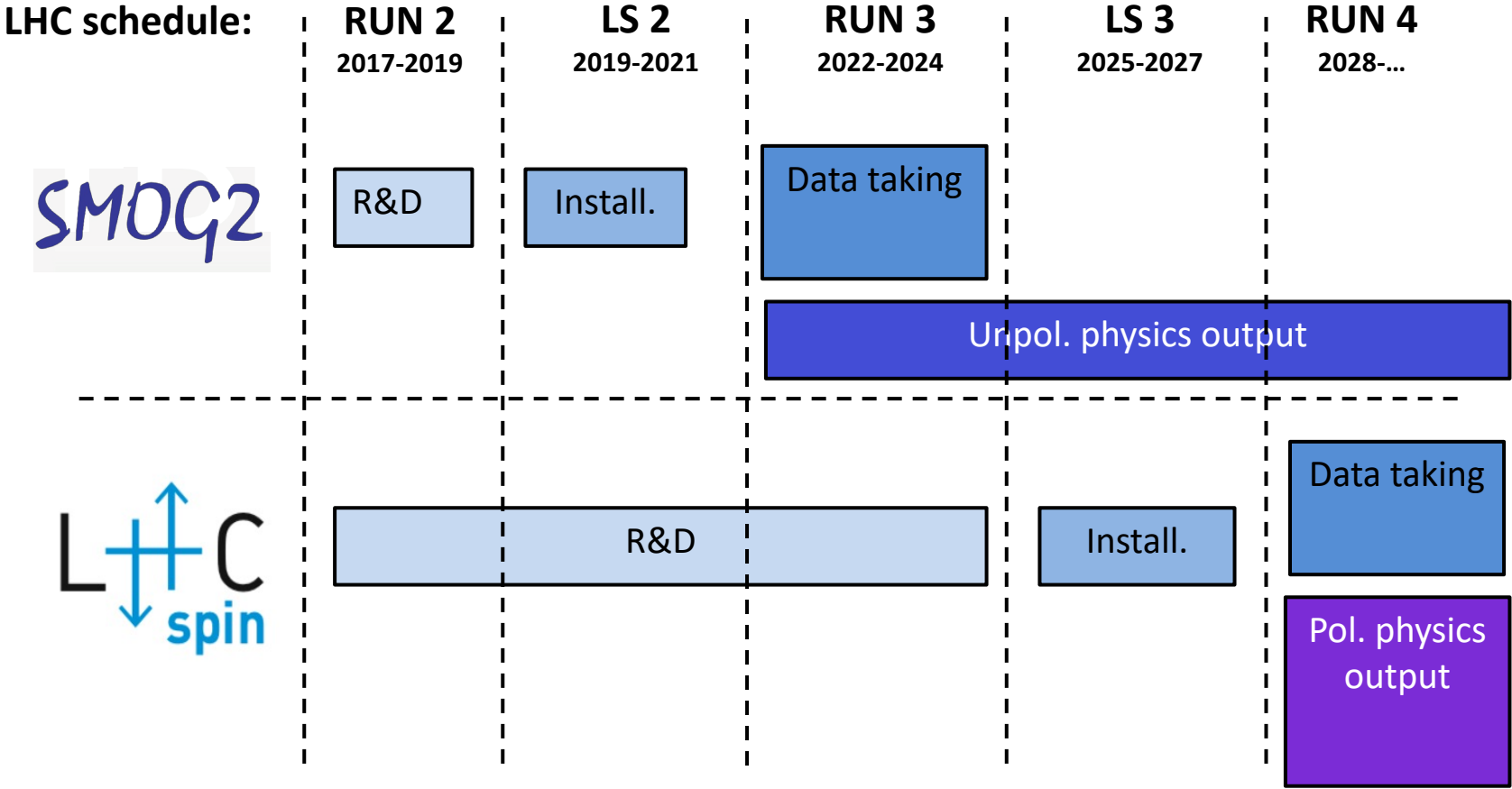
Statistical vs Systematics uncertainties

- The analysis tool described above allows to study the interplay between statistical uncertainties and systematic uncertainties (due to the measurement of the polarization) under different data-taking scenarios

p_T (MeV)	x_F	a_1 ($\Delta P = 0\%$)	a_1 ($\Delta P = 5\%$)	a_1 ($\Delta P = 20\%$)	a_1 ($\Delta P = 50\%$)
[0,1500]	[-0.70,-0.09]	0.090 ± 0.013	0.089 ± 0.013	0.087 ± 0.014	0.087 ± 0.022
[0,1500]	[-0.09,-0.06]	0.104 ± 0.011	0.104 ± 0.012	0.103 ± 0.016	0.100 ± 0.027
[0,1500]	[-0.06,-0.04]	0.098 ± 0.012	0.098 ± 0.013	0.097 ± 0.016	0.094 ± 0.027
[0,1500]	[-0.04,0.05]	0.118 ± 0.014	0.117 ± 0.014	0.114 ± 0.017	0.113 ± 0.030
[1500,6000]	[-0.70,-0.09]	0.093 ± 0.010	0.092 ± 0.010	0.090 ± 0.013	0.089 ± 0.023
[1500,6000]	[-0.09,-0.06]	0.108 ± 0.011	0.108 ± 0.011	0.108 ± 0.015	0.107 ± 0.027
[1500,6000]	[-0.06,-0.04]	0.105 ± 0.012	0.105 ± 0.012	0.104 ± 0.015	0.103 ± 0.026
[1500,6000]	[-0.04,0.05]	0.105 ± 0.011	0.105 ± 0.012	0.102 ± 0.015	0.102 ± 0.026

- A 5% systematic uncertainty on P has no impact on the total uncertainty on a_1
- For $\Delta P = 20\%$ the systematic uncertainty amounts to 30-40% of the statistical uncertainty
- For $\Delta P = 50\%$ the systematic uncertainty approximately equals the statistical uncertainty
- We expect $\Delta P \approx 10 - 15\%$ for the storage cell hypothesis (and close to 0 for the jet target hypothesis)**

The time schedule of the project



Conclusions

- The FT program at LHCb is active since Run 2, now greatly enriched with SMOG2
- LHCspin is the natural evolution: a polarized fixed target at LHCb will bring spin-physics for the first time at the LHC and will open the way to a broad and ambitious physics program
- Novel approaches and reactions will be exploited for studies of the 3D nucleon structure
- First insights into the yet unknown gluon TMDs (such as the GSF) will be possible thanks to the excellent capabilities of LHCb in reconstructing quarkonia states and heavy mesons.
- Cutting-edge unpolarized physics will also be at reach (cold nuclear matter effects, intrinsic charm, QGP studies, etc.)
- The R&D calls for a new generation of polarized gas targets. A very challenging but worth the effort!

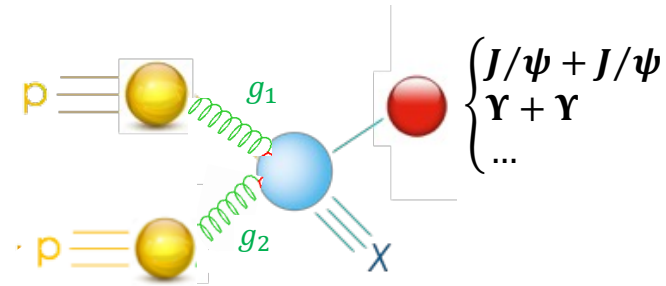
If approved, LHCspin will make LHCb the first experiment simultaneously running in collider and fixed-target mode with polarized targets, opening a whole new range of explorations.

Backup

Probing the g TMDs

$$\frac{d\sigma}{dM_{QQ}dY_{QQ}d^2P_{QQ\tau}d\Omega} = \frac{\sqrt{M_{QQ}^2 - 4M_Q^2}}{(2\pi)^2 8s M_{QQ}^2} \left\{ F_1(M_{QQ}, \theta_{CS}) \mathcal{C} [f_1^g f_1^g] (x_{1,2}, P_{QQ\tau}) \right.$$

$$+ F_2(M_{QQ}, \theta_{CS}) \mathcal{C} [w_2 h_1^{\perp g} h_1^{\perp g}] (x_{1,2}, P_{QQ\tau})$$



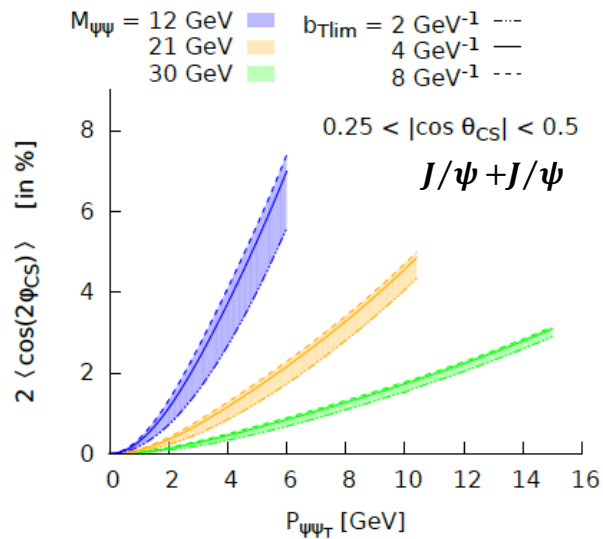
gluon pol.

	U	Circularly	Linearly
U	f_1^g		$h_1^{\perp g}$
L		g_{1L}^g	$h_{1L}^{\perp g}$
T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

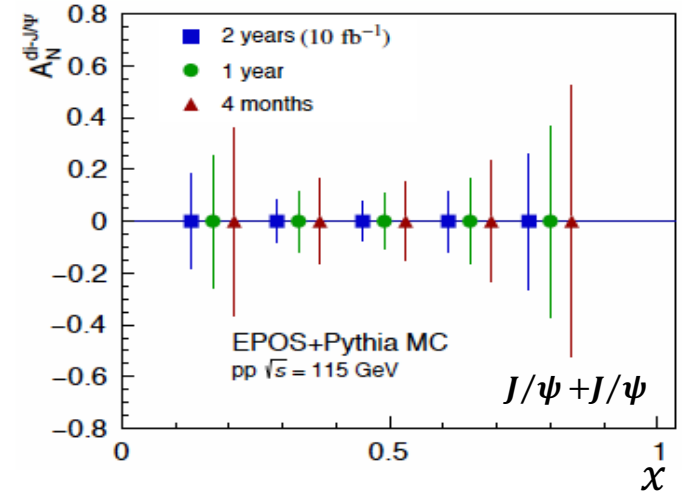
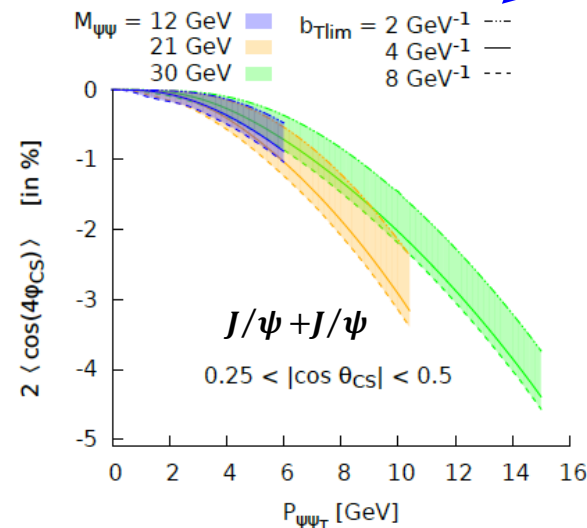
nucleon pol.

$$+ \left(F_3(M_{QQ}, \theta_{CS}) \mathcal{C} [w_3 f_1^g h_1^{\perp g}] (x_{1,2}, P_{QQ\tau}) + F'_3(M_{QQ}, \theta_{CS}) \mathcal{C} [w'_3 h_1^{\perp g} f_1^g] (x_{1,2}, P_{QQ\tau}) \right) \cos 2\phi_{CS} + F_4(M_{QQ}, \theta_{CS}) \mathcal{C} [w_4 h_1^{\perp g} h_1^{\perp g}] (x_{1,2}, P_{QQ\tau}) \cos 4\phi_{CS}$$

Predictions based on CSM + TMD evolution for $x_1 \sim x_2 \sim 10^{-3}$ at forward rapidity [EPJ C 80, 87 (2020)]



Azimuthal amplitudes
~5%!



A synergic attack to g TMDs

[D. Boer: Few-body Systems 58, 32 (2017)]

	DIS	DY	SIDIS	$pA \rightarrow \gamma \text{jet } X$	$ep \rightarrow e' Q \bar{Q} X$ $ep \rightarrow e' j_1 j_2 X$	$pp \rightarrow \eta_{c,b} X$ $pp \rightarrow H X$	$pp \rightarrow J/\psi \gamma X$ $pp \rightarrow \Upsilon \gamma X$
$f_1^g^{[+,+]}$ (WW)	×	×	×	×	✓	✓	✓
$f_1^g^{[+,-]}$ (DP)	✓	✓	✓	✓	×	×	×

- Can be measured at the EIC
- Can be measured at RHIC & LHC (including LHCb+SMOG2/LHCspin)
- Can be measured at RHIC and LHCb+LHCspin

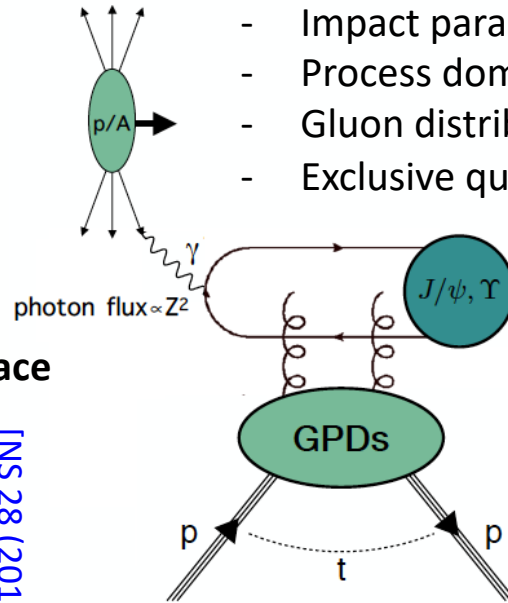
	$pp \rightarrow \gamma \gamma X$	$pA \rightarrow \gamma^* \text{jet } X$	$ep \rightarrow e' Q \bar{Q} X$ $ep \rightarrow e' j_1 j_2 X$	$pp \rightarrow \eta_{c,b} X$ $pp \rightarrow H X$	$pp \rightarrow J/\psi \gamma X$ $pp \rightarrow \Upsilon \gamma X$
$h_1^{\perp g [+,+]}$ (WW)	✓	×	✓	✓	✓
$h_1^{\perp g [+,-]}$ (DP)	×	✓	×	×	×

	DY	SIDIS	$p^\dagger A \rightarrow h X$	$p^\dagger A \rightarrow \gamma^{(*)} \text{jet } X$	$p^\dagger p \rightarrow \gamma \gamma X$ $p^\dagger p \rightarrow J/\psi \gamma X$ $p^\dagger p \rightarrow J/\psi J/\psi X$	$ep^\dagger \rightarrow e' Q \bar{Q} X$ $ep^\dagger \rightarrow e' j_1 j_2 X$
$f_{1T}^{\perp g [+,+]}$ (WW)	×	×	×	×	✓	✓
$f_{1T}^{\perp g [+,-]}$ (DP)	✓	✓	✓	✓	×	×

UPC and gGPDs

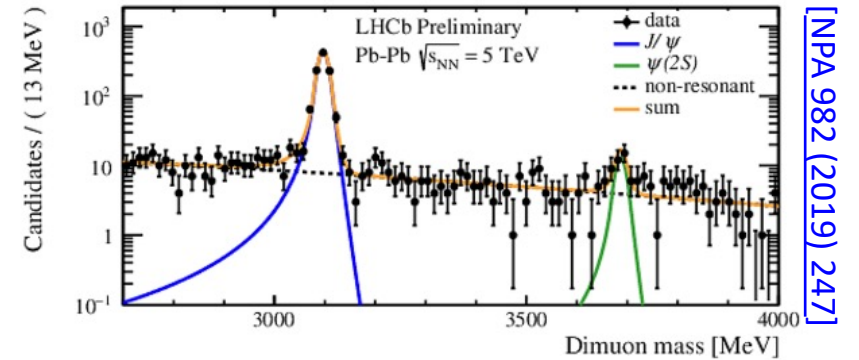
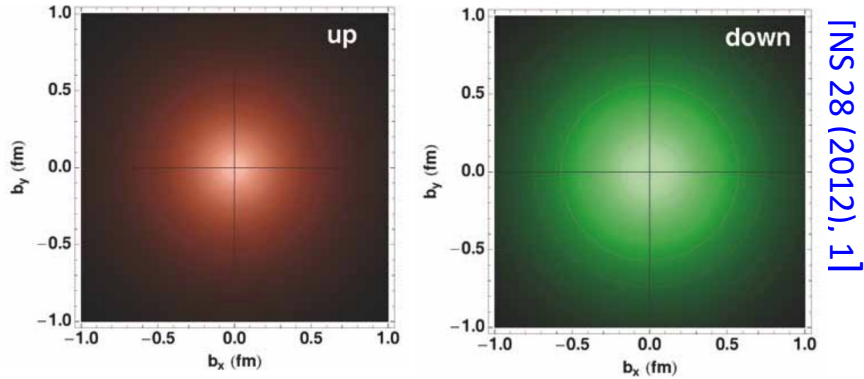
GPD	U	L	T
U	H		\mathcal{E}_T
L		\tilde{H}	$\tilde{\mathcal{E}}_T$
T	E	\tilde{E}	H_T, \tilde{H}_T

Can be accessed at LHC in **Ultra-Peripheral collisions (UPC)**

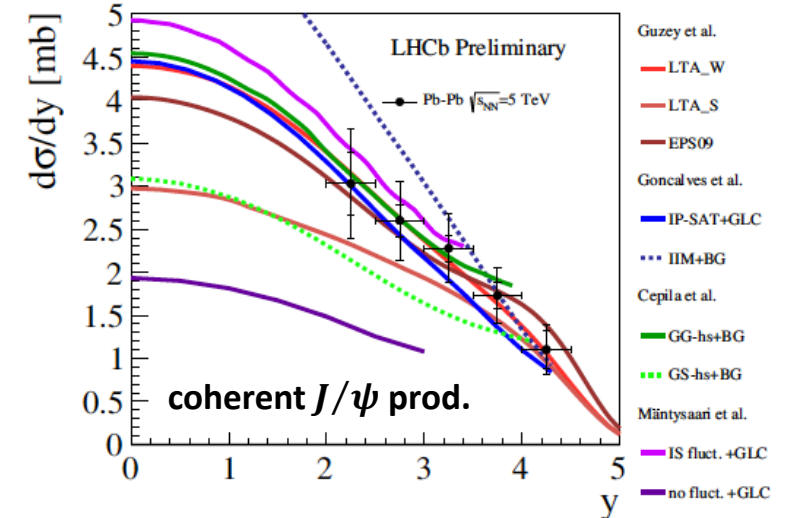


- Impact parameter larger than sum of radii
- Process dominated by EM interaction
- Gluon distributions probed by pomeron exchange
- Exclusive quarkonia prod. sensitive to gluon GPDs [\[PRD 85 \(2012\), 051502\]](#)

3D maps of parton densities in coordinate space



First results from LHCb in PbPb UPC



LHCspin could allow to access the GPD E^g (a key ingredient of the Ji sum rule)

$$J^g = \frac{1}{2} \int_0^1 dx \left(H^g(x, \xi, 0) + E^g(x, \xi, 0) \right)$$

Cell coating R&D

The inner coating of the storage cell is a crucial aspect of the R&D. It is needed to:

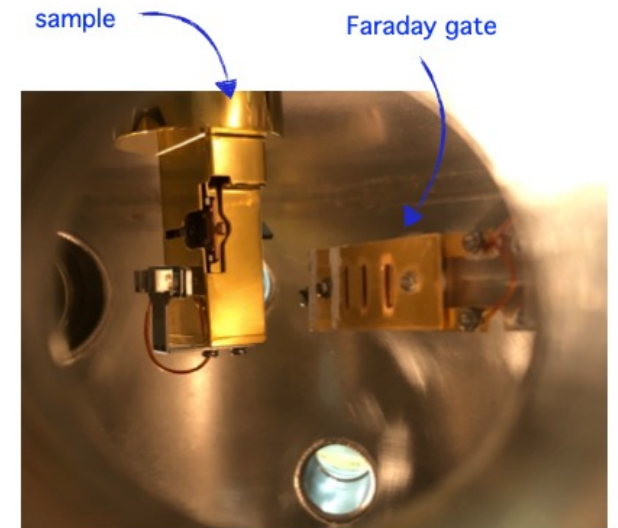
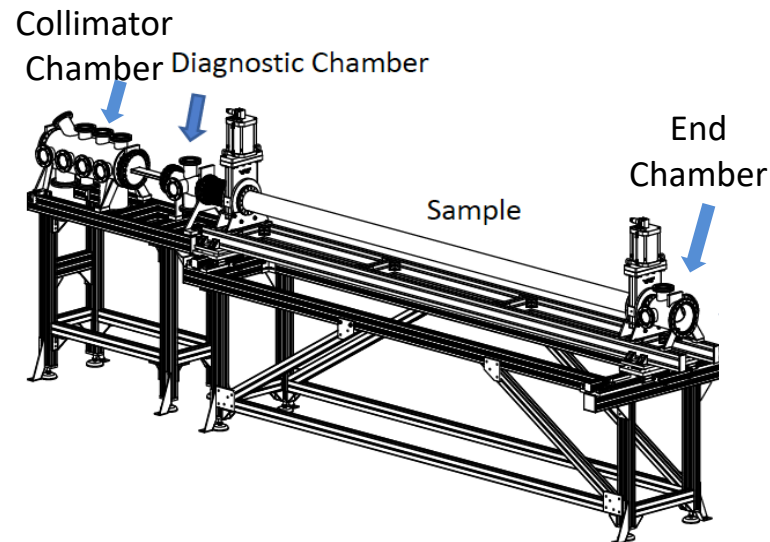
- ✓ **minimize e-cloud** related beam instabilities → ensure low Secondary Electron Yield (SEY)
- ✓ **minimize H depolarization** due to wall collisions → can be monitored through measurement of H recombination

Teflon and Drifilm (HERMES) not compatible with LHC requirements. **Amorphous Carbon** (a-C) is allowed but may induce depolarization. Possible solution (a-la HERMES): generate a **thin layer of ice on top of a-C coating**

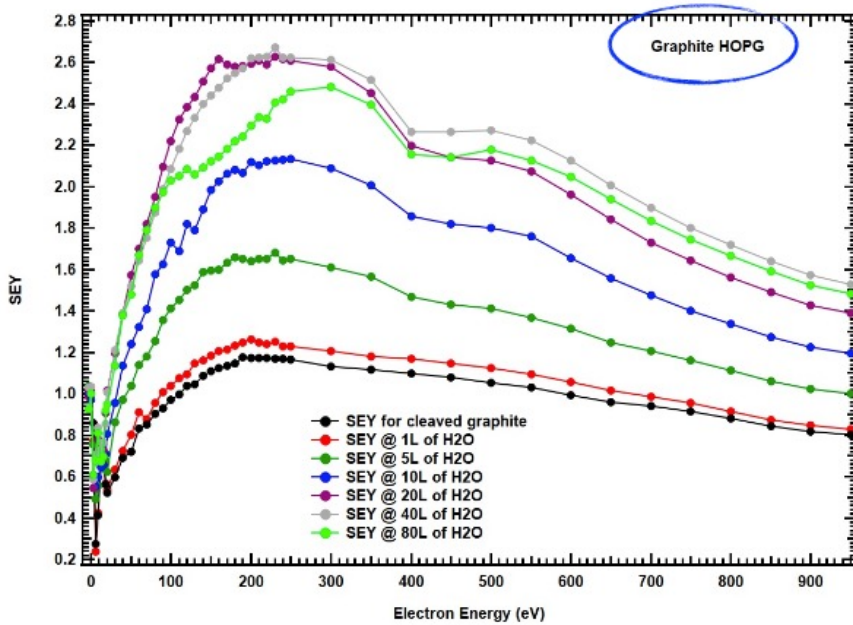
- reduces the depolarization
- ensures a renewable surface
- requires to cool down the cell to $\sim 100\text{ K}$
- could cause a larger SEY
- Need to be investigated → **dedicated R&D**

The **ARYA project at INFN-LNF**:

- existing surface-coating laboratory has been equipped with UHV ultrapure water dosing system
- SEY measurement through electron gun on target vs H_2O dose



Cell coating R&D

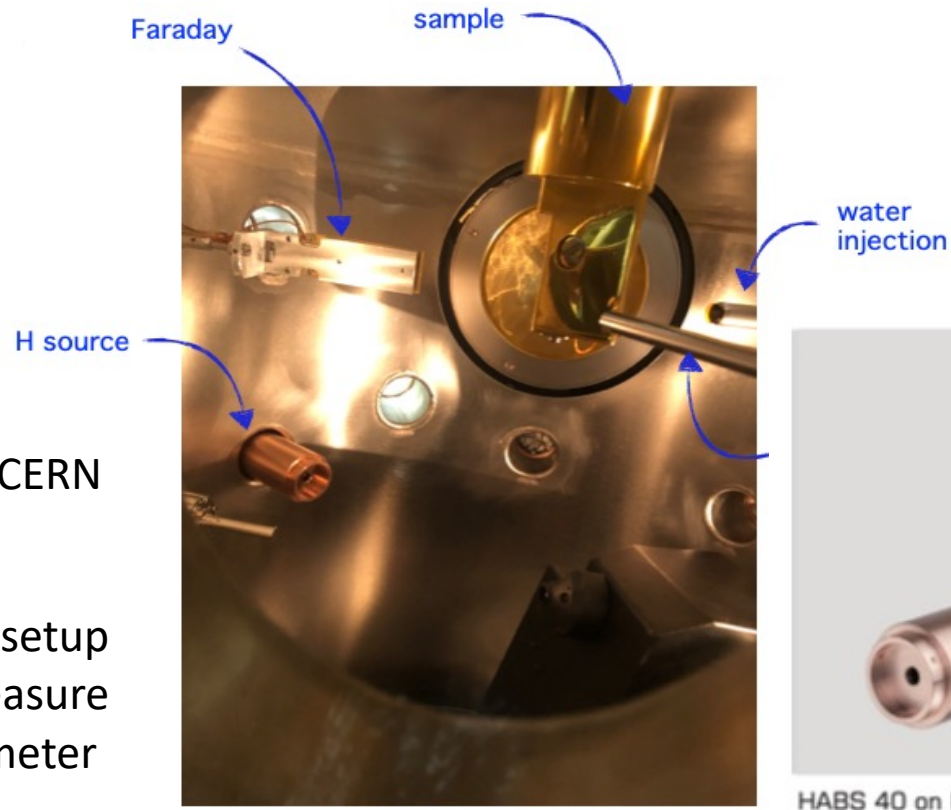


- SEY measurements on graphite vs incident electron energy with 1, 4, 10, 20, 40, 80 monolayers of H_2O at 90K
- Max SEY = 2.6: impact on LHC is under evaluation

Next steps:

- Measurement on the actual a-C sample from CERN
- **Measurement of H recombination**

H recombination studies: with the same setup inject H by means of an atomic source and measure H/H_2 fraction vs H_2O dose with a mass spectrometer



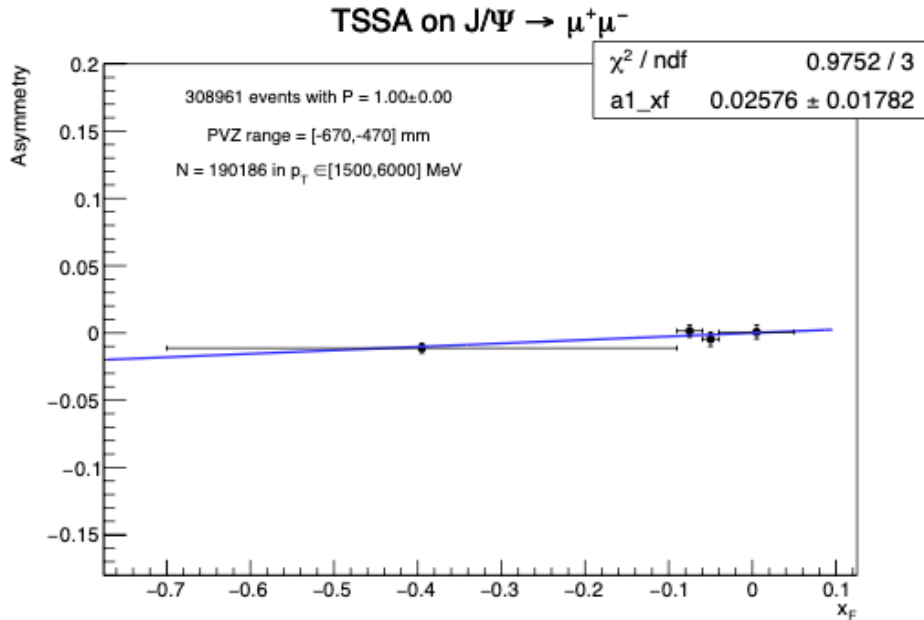
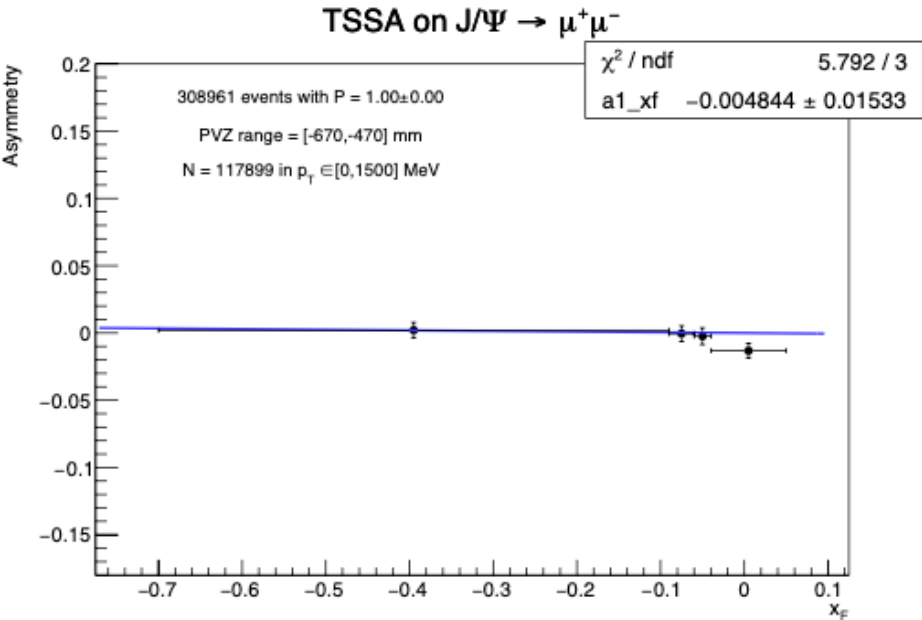
Independent depolarization studies on a-C ongoing in Juelich (dedicated laboratory)

Main reactions or interest (...an incomplete wishlist)

- $pp \rightarrow \mu^+ \mu^- + X$ ($pp \rightarrow e^+ e^- + X$)
 - $pd \rightarrow \mu^+ \mu^- + X$ ($pd \rightarrow e^+ e^- + X$)
 - $pp^\uparrow \rightarrow \mu^+ \mu^- + X$ ($pp^\uparrow \rightarrow e^+ e^- + X$)
 - $pd^\uparrow \rightarrow \mu^+ \mu^- + X$ ($pd^\uparrow \rightarrow e^+ e^- + X$)
 - $pp^{(\uparrow)} \rightarrow \eta_c + X$ ($pp^{(\uparrow)} \rightarrow \chi_{c,b} + X$)
 - $pp^{(\uparrow)} \rightarrow J/\psi + X$
 - $pp^{(\uparrow)} \rightarrow \Upsilon + X$
 - $pp^{(\uparrow)} \rightarrow J/\psi + J/\psi + X$
 - $pp^{(\uparrow)} \rightarrow J/\psi + \gamma + X$
 - $pp^{(\uparrow)} \rightarrow \Upsilon + \gamma + X$
 - pA, PbA ($A = He, Ne, Ar, Kr, \dots$)
- unpolarized TMDs of valence and sea quarks and momentum distrib. of sea quarks
- TMDs of valence and sea quarks
- Pol and unpol gluon PDFs
- Nuclear matter effects, QGP, etc

A preliminary analysis tool for pseudo-data

- A simpler approach is also developed which consists in evaluating the TSSA directly from the yields in each 2D bin



- A linear fit is also shown to quantify the mild kinematic dependence of the asymmetry