

Lattice simulations of inflation and reheating



Kaloian Lozanov UIUC, Cosmology and HEPheno



The equation of state after inflation

KL and M. Amin (2017, 18, 19)

 $w = \frac{\text{pressure}}{\text{energy density}}$

$$S = \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \Big]$$

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assumption: self-couplings dominate over others

at sufficiently late times:

$$w = \begin{cases} 0 & \text{if } n = 1\\ 1/3 & \text{if } n > 1 \end{cases}$$

(even without couplings to other fields!)



Inflaton dynamics



oscillatory phase

Inflaton (homogeneous) dynamics ϕ $V(\phi)$ inflation $\phi \times a^{3/(n+1)}$ 0.0 inflation ends: 50 50 oscillatory phase x 100 $\bar{\phi}(t) \propto a(t)^{-3/(n+1)}$

Inflaton (homogeneous) dynamics ϕ $V(\phi)$ inflation $w \equiv \frac{p}{\rho} = \frac{\phi^2/2 - (\nabla\phi)^2/6 - V(\phi)}{\dot{\phi}^2/2 + (\nabla\phi)^2/2 + V(\phi)}$ inflation ends: oscillatory phase $\bar{\phi}(t) \propto a(t)^{-3/(n+1)}$





- parametric resonance of $\delta \phi(t, \mathbf{x})$
- $\bar{\phi}$ fragments

KL and M. Amin (2017) KL and M. Amin (2018)

Equation of state



(e-folds after inflation) ΔN

$$\Delta N \equiv \int_{a_{\rm end}}^{a} d\ln a$$

Equation of state

- $M \ll m_{\rm pl}$ (efficient resonance)



(e-folds after inflation) ΔN

$$\Delta N \equiv \int_{a_{\rm end}}^{a} d\ln a$$

Non-perturbative decay (parametric self-resonance)



 $m^2 \equiv V'(\bar{\phi}_{\rm osc})/\bar{\phi}_{\rm osc}$

Non-perturbative decay (parametric self-resonance)

n = 1.5 $V(\bar{\phi}) \propto |\bar{\phi}|^{2n}$ $\frac{\Re(\mu_k)}{m}$ 2.0inflaton vev amplitude = const 0.161.5 $\phi_{\rm osc}/M$ 1.0 $\bar{\phi} \sim M$ 0.080.5 $\delta \phi_k \propto \exp(\pm \mu_k t)$ 0.00 0.00.51.01.52.00.0 $\kappa = k/am$ inflaton field wavenumber

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inflaton field wavenumber

comoving wavenumber

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n=3Power spectrum: 2.0 $\Re(\mu_k)$ inflaton vev amplitude $\left< \delta \phi(x)^2 \right> \equiv \int \mathcal{P}_{\delta \phi} d \ln k$ 2 mconst 0.036 10^{-1} 1.5 10^{-3} ${\cal P}_{\delta\phi}/ar{\phi}_{
m osc}^2$ $\phi_{\rm osc}/M$ 10^{-5} 1.0 10^{-7} 0.018 10^{-9} 0.5 10^{-11} 0.000 10^{-13} 0.00.51.01.52.0 10^{-3} 10^{-2} 0.0 10^{-1} k/m_0 $\kappa = k/am$ inflaton field wavenumber comoving wavenumber

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inflaton field wavenumber

Non-perturbative decay (parametric self-resonance)



 $V(\bar{\phi}) \propto |\bar{\phi}|^{2n}$ $\bar{\phi} \sim M \ll m_{\rm pl}$

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inflaton field wavenumber

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Towards radiation domination n > 1

Non-perturbative decay (parametric self-resonance)



Towards radiation domination n > 1

Non-perturbative decay (parametric self-resonance)



Towards radiation domination n > 1 $M \ll m_{\rm pl}$ $M \sim m_{\rm pl}$

- slow production of $\delta \phi(t, \mathbf{x})$
- $\bar{\phi}$ fragments gradually

$$\Delta N_{\rm fr} \approx \frac{n+1}{3} \ln \left(10^3 \frac{M}{m_{\rm pl}} \right)$$

at sufficiently late times:

 ϕ virialized + turbulent $\rightarrow w = \frac{1}{3}$



• $\bar{\phi}$ fragments quickly

Spectral index: $n_{\rm s}$

Tensor-to-scalar ratio: r

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$$50 \le N_* \le 60$$

Planck Collaboration (2015)

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$$N_* = 66.89 - \frac{1}{12} \ln g_{\rm th} + \frac{1}{4} \ln \frac{V_*^4}{m_{\rm pl}^4 \rho_{\rm end}} - \ln \frac{k_*}{a_0 H_0} + \frac{3\bar{w}_{\rm int} - 1}{4} \Delta N_{\rm rad}$$

$$\Delta N_{\rm rad} \equiv \int_{a_{\rm end}}^{a_{\rm rad}} d\ln a$$

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reheating

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$$\Delta N_{\rm fr} = \begin{cases} 1 & \text{if} & M \ll m_{\rm pl} \\ \frac{n+1}{3} \ln \left(10\frac{\kappa}{\Delta\kappa}\frac{M}{m_{\rm pl}}\right) & \text{if} & M \sim m_{\rm pl} \end{cases}$$

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$$w_{\rm int}(\Delta N) = \begin{cases} \frac{n-1}{n+1} & \text{if} & 0 < \Delta N < \Delta N_{\rm rad} \\ \frac{1}{3} & \text{if} & \Delta N > \Delta N_{\rm rad} \end{cases}$$

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Kallosh and Linde (2013) Carrasco, Kallosh and Linde (2015)









Planck Collaboration (2015)



Planck Collaboration (2015)

$$V(\phi) \propto \left| 1 - e^{-\phi/M} \right|^{2n}$$
 α -attractors (E-models)











Planck Collaboration (2015)



$$V(\phi) \propto \left[1 + \left|\frac{\phi}{M}\right|^{2n}\right]^{\frac{q}{2n}} - 1$$
 Monodromy



Silverstein and Westphal (2008) McAllister, et al. (2014)

$$V(\phi) \propto \left[1 + \left|\frac{\phi}{M}\right|^{2n}\right]^{\frac{1}{4n}} - 1$$
 Monodromy
 $q = 1/2$



Silverstein and Westphal (2008) McAllister, et al. (2014)











Equation of state

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Non-perturbative decay (parametric <u>self</u>-resonance)



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n = 1



• $\bar{\phi}$ forms oscillons (stable)

KL and Amin (2017, 18, 19)

See also Amin, Easther, Finkel, Flauger, Hertzberg (2011)

matter-like eos: w = 0

couplings to other fields?

See also Antusch, Figueroa, Marschall, Torrenti (2020) Antusch, Marschall, Torrenti (2022)

 $M \sim m_{\rm pl}$

- $\delta \phi(t, \mathbf{x})$ production shut off
- $\bar{\phi}_{\rm osc}(t) = {\rm pressureless} \ {\rm dust}$

• stochastic GWs from fragmentation

KL and M. Amin (2019)

• stochastic GWs from fragmentation

KL and M. Amin (2019)

Oscillons



stochastic GWs from fragmentation ulletTransients

KL and M. Amin (2019)

 10^{-9} 10^{-10} 10^{-11} 10^{-12} 10^{-13} 10^{-14}



lacksquare



lacksquare



ullet





• stochastic GWs from fragmentation

KL and M. Amin (2019)

• (early) dark energy (n=2)

Agrawal, Cyr-Racine, Pinner, Randall (2019) Weiner, Adshead, Giblin (2020) M. Amin, KL and T. Smith (in progress)

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- matter-antimatter asymmetry...

KL and M. Amin (2014)

Oscillons & baryogenesis

KL and M. Amin, PRD 90, 083528 (2014)

+ U(1)dynamics at the end of inflation

inflaton/anti-inflaton asymmetry

complex

inflaton ϕ

decay

very different dynamics from homogeneous case!

matter-antimatter asymmetry $\eta \approx 6 \times 10^{-10}$





KL and M. Amin, JCAP 1606 032 (2016)

$$S = \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - (D_{\mu}\phi)^{\dagger} D^{\mu}\phi - V(|\phi|) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \Big]$$

See also Figueroa, Garcia-Bellido, Torrenti, (2015, 16)

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inflaton amplitude

wavenumber

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Non-Abelian code in collaboration with M. Amin, A. Caravano, KL, E. Komatsu... (in progress)



<u>Abelian code: KL, M. Amin (2020)</u> **GFiRe: Gauge Field integrator for Reheating** <u>NEW LATTICE CODE!!!</u> JCAP 058 04 (2020)

See also Figueroa, Florio, Torrenti, Valkenburg (2020, 2021)

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{\alpha \phi}{4f} F_{\mu\nu} \tilde{F}^{\mu\nu} \right]$$

Anber and Sorbo (2010), Barnaby and Peloso (2011)

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Axion inflaton:

Anber and Sorbo (2010), Barnaby and Peloso (2011)

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 $\xi \gtrsim 1$

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 $\xi \gtrsim 1$
 $\xi \gg 1$
Anber and Sorbo (2010), Barnaby and Peloso (2011)

$$S = \int d^{4}x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{\alpha \phi}{4f} F_{\mu\nu} \tilde{F}^{\mu\nu} \Big]$$
Axion inflaton: $\phi(\mathbf{x}, t) = \overline{\phi}(t) + \delta \phi(\mathbf{x}, t)$
 $\xi = \alpha \dot{\overline{\phi}}/(2fH)$
 $\xi \gtrsim 1$
 $\xi \gg 1$
Weak δA_{μ} Backreaction

Anber and Sorbo (2010), Barnaby and Peloso (2011)

$$\begin{split} S &= \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{\alpha \phi}{4f} F_{\mu\nu} \tilde{F}^{\mu\nu} \Big] \\ & \text{Axion inflaton: } \phi(\mathbf{x}, t) = \overline{\phi}(t) + \delta \phi(\mathbf{x}, t) \\ & \xi &= \alpha \overline{\phi} / (2fH) \\ & \xi &\gtrsim 1 \\ & \xi &\gg 1 \\ & \text{Weak } \delta A_\mu \text{ Backreaction} \\ \end{split}$$

Anber and Sorbo (2010), Barnaby and Peloso (2011)

$$S = \int d^{4}x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{\alpha \phi}{4f} F_{\mu\nu} \tilde{F}^{\mu\nu} \Big]$$
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$$\xi = \alpha \overline{\phi} / (2fH)$$

$$\xi \gtrsim 1$$

$$\xi \gg 1$$
Weak δA_{ν} Backreaction
$$\xi \gg 1$$
Strong δA_{ν} Backreaction

Perturbation theory:

Meerburg and Pajer (2013) Linde, Mooij and Pajer (2013) Garcia-Bellido, Peloso and Unal (2016) Domcke, Muia, Pieroni, and Witkowski (2017) Domcke, Guidetti, Welling and Westphal (2020)

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 $\xi = \alpha \overline{\phi}/(2fH), \quad \zeta = -H \delta \phi / \overline{\phi}$
 $\xi \gtrsim 1$
Weak δA_{μ} Backreaction
 $\xi \gg 1$
Strong δA_{μ} Backreaction

Perturbation theory:

Meerburg and Pajer (2013) Linde, Mooij and Pajer (2013) Garcia-Bellido, Peloso and Unal (2016) Domcke, Muia, Pieroni, and Witkowski (2017) Domcke, Guidetti, Welling and Westphal (2020)

Anber and Sorbo (2010), Barnaby and Peloso (2011) <u>Lattice theory: Angelo Caravano</u>, Komatsu, KL, Weller (2021a,b, 2022)

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Gauge fields, inflation & reheating Anber and Sorbo (2010), Barnaby and Peloso (2011) Angelo Caravano, Komatsu, KL, Weller (2021a,b, 2022) $S = \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{\alpha \phi}{4f} F_{\mu\nu} \tilde{F}^{\mu\nu} \right]$ Axion inflaton: $\phi(\mathbf{x}, t) = \overline{\phi}(t) + \delta \phi(\mathbf{x}, t)$ $\xi = \alpha \overline{\phi} / (2fH), \quad \zeta = -H\delta\phi/\overline{\phi}$ $\xi \gtrsim 1$ $\xi \gg 1$ Weak δA_{μ} Backreaction Strong δA_{μ} Backreaction 10^{-6} 10 $N_{e} = 0.0$ = 1.4710-3 - 10^{-7} $N_{e} = 5.99$ = 0.0 10^{-5} = 1.48 10^{-8} N_e = 3.90 10^{-7} $\mathcal{P}_{\zeta}^{a} = \overline{\mathcal{P}_{\mathrm{vac}}} + \overline{\mathcal{P}}_{\mathrm{vac}}^{2} f(\xi) e^{4\pi\xi}$ $N_{e} = 4.52$ 10^{-9} $N_{o} = 6.52$ 10^{1} 10² 10² $\mathcal{P}_{ m vac} = H^4/(2\pi \dot{\overline{\phi}})^2$ k/mk/m

 \mathcal{P}_{i}

Anber and Sorbo (2010), Barnaby and Peloso (2011) Angelo Caravano, Komatsu, KL, Weller (2021a,b, 2022)

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 $\sigma^2 = \langle \zeta^2 \rangle$

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Strong δA_{μ} Backreaction



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 $\xi \ge 1$ non-Gaussianity

 $\xi \gg 1$ Strong δA_{μ} Backreaction

Weak δA_{μ} Backreaction



Anber and Sorbo (2010), Barnaby and Peloso (2011) Angelo Caravano, Komatsu, KL, Weller (2021a,b, 2022)

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$\xi \gtrsim 1$ non-Gaussianity

Weak δA_{μ} Backreaction



 $\xi \gg 1$

Strong δA_{μ} Backreaction



 C/σ

Anber and Sorbo (2010), Barnaby and Peloso (2011) Angelo Caravano, Komatsu, KL, Weller (2021a,b, 2022)

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Weak δA_{μ} Backreaction



Strong δA_{μ} Backreaction

 $\xi \gg 1$ Gaussianity!



 C/σ

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 $\xi \gtrsim 1$ Weak δA_{μ} Backreaction

Strong δA_{μ} Backreaction

$$\sigma^2 = \langle \zeta^2 \rangle$$

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 $\xi \gtrsim 1$

Weak δA_{μ} Backreaction

$$\kappa_3 = \frac{\langle \zeta^3 \rangle}{\sigma^3}$$

Strong δA_{μ} Backreaction

$$\sigma^2 = \langle \zeta^2 \rangle$$

Anber and Sorbo (2010), Barnaby and Peloso (2011) Angelo Caravano, Komatsu, KL, Weller (2021a,b, 2022)

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 $\xi \gtrsim 1$

Weak δA_{μ} Backreaction

$$\kappa_3 = \frac{\langle \zeta^3 \rangle}{\sigma^3}$$
$$\kappa_4 = \frac{\langle \zeta^4 \rangle - 3\sigma^4}{\sigma^4}$$

Strong δA_{μ} Backreaction

$$\sigma^2 = \langle \zeta^2 \rangle$$

Anber and Sorbo (2010), Barnaby and Peloso (2011) Angelo Caravano, Komatsu, KL, Weller (2021a,b, 2022)

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Weak δA_{μ} Backreaction

Strong δA_{μ} Backreaction

 $\xi \gg 1$

$$\kappa_{3} = \frac{\langle \zeta \rangle}{\sigma^{3}}$$

$$\kappa_{4} = \frac{\langle \zeta^{4} \rangle - 3\sigma^{4}}{\sigma^{4}}$$

$$\kappa_{5} = \frac{\langle \zeta^{5} \rangle - 10 \langle \zeta^{3} \rangle \sigma^{2}}{\sigma^{5}}$$

$$\sigma^{2} = \langle \zeta^{2} \rangle$$

123

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 $\kappa_3 = \frac{\langle \zeta^3 \rangle}{-3}$

 $\sigma^2 = \langle \zeta^2$

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Weak δA_{μ} Backreaction

Strong δA_{μ} Backreaction

 $\xi \gg 1$



 N_e

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Strong δA_{μ} Backreaction

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 $\sigma^2 = \langle \zeta^2$

 $\kappa_4 = \frac{\langle \zeta^4 \rangle - 3\sigma^4}{4}$

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ho}$



 N_e

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Anber and Sorbo (2010), Barnaby and Peloso (2011) Angelo Caravano, Komatsu, KL, Weller (2021a,b, 2022)

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Weak δA_{μ} Backreaction

Strong δA_{μ} Backreaction

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 N_e



 N_e

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$$\xi \gtrsim 1$$

Weak δA_{μ} Backreaction

Strong δA_{μ} Backreaction



$$\phi_k'' + 2\mathcal{H}\phi_k' + k^2\phi_k = a^2\frac{\alpha}{4f}\int d^3q \,F_{\mu\nu q}\tilde{F}_{k-q}^{\mu\nu}$$





Anber and Sorbo (2010), Barnaby and Peloso (2011) Angelo Caravano, Komatsu, KL, Weller (2021a,b, 2022)

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Weak δA_{μ} Backreaction

Strong δA_{μ} Backreaction

 $\xi \gg 1$



$$\begin{split} \phi_k'' + 2\mathcal{H}\phi_k' + k^2\phi_k &= a^2\frac{\alpha}{4f}\int d^3q \,F_{\mu\nu q}\tilde{F}_{k-q}^{\mu\nu} \\ q/\mathcal{H} &\lesssim \xi \end{split}$$



 N_e





Conclusions

Reheating at the end of inflation:

• very rich dynamics

Future plans:

- more realistic models (e.g. gauge fields)
- observational signatures
 -expansion history, relics, GWs



