



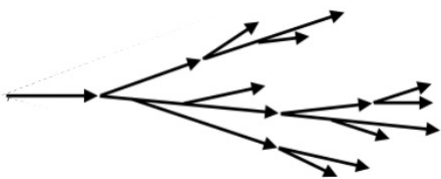
Direct observation of the QCD dead-cone effect

Nima Zardoshti (CERN)
On behalf of the ALICE Collaboration



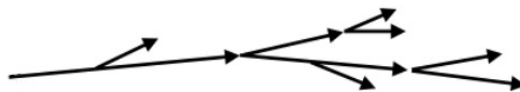
Gluon-initiated shower

Broader shower profile
Higher number of emissions



Quark-initiated shower

narrower shower profile
Fewer emissions in the shower



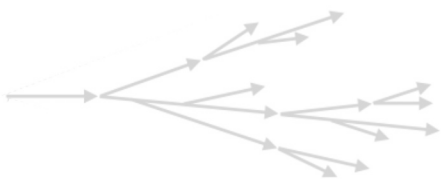
Casimir Colour factors

Different emission properties due to the different amount of colour charge carried by quarks and gluons

$$\frac{C_A}{C_F} = \frac{9}{4}$$

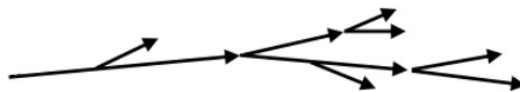
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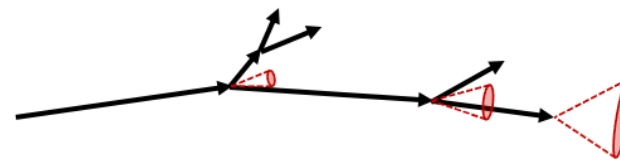
Quark-initiated shower

narrower shower profile
Fewer emissions in the shower



Heavy-quark-initiated shower

Suppression of small angle emissions
Harder fragmentation



Casimir Colour factors

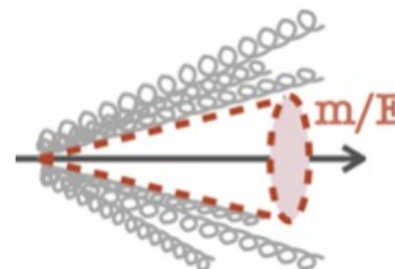
Different emission properties due to the different amount of colour charge carried by quarks and gluons

$$\frac{C_A}{C_F} = \frac{9}{4}$$

The dead-cone effect

A suppression of emissions in a cone of size m/E around the direction of the emitter

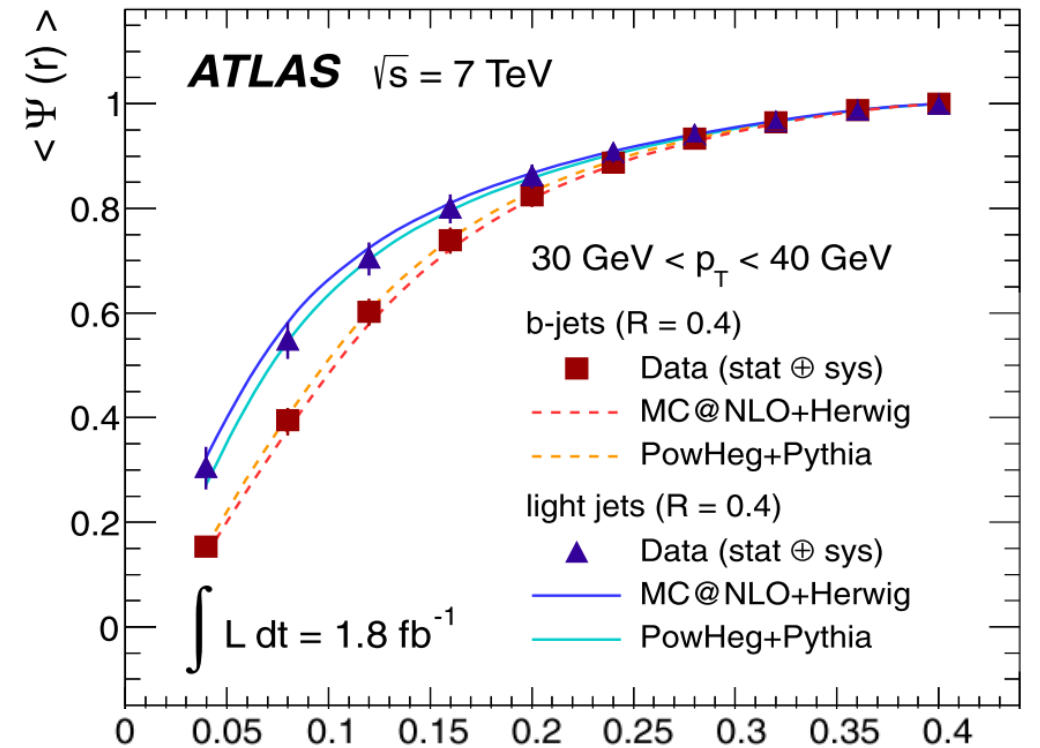
Sizeable effect for low energy heavy quarks



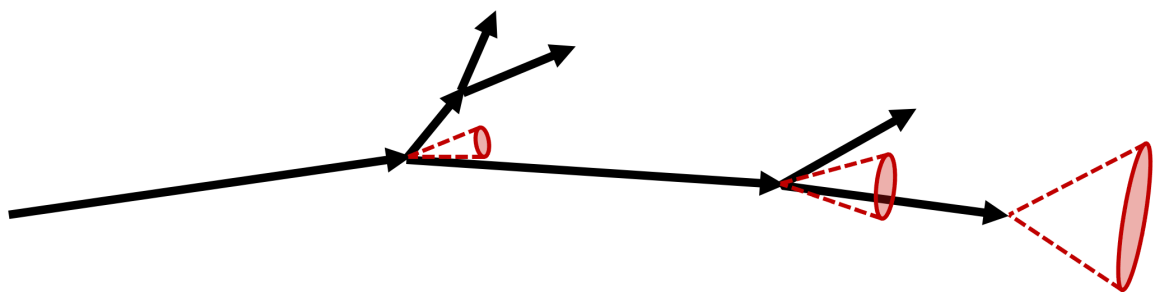
p_T density around the initial scattered b-quark direction is depleted compared to the density around light quarks and gluons

Consequence of the suppression of emissions in the dead cone of the b-quark

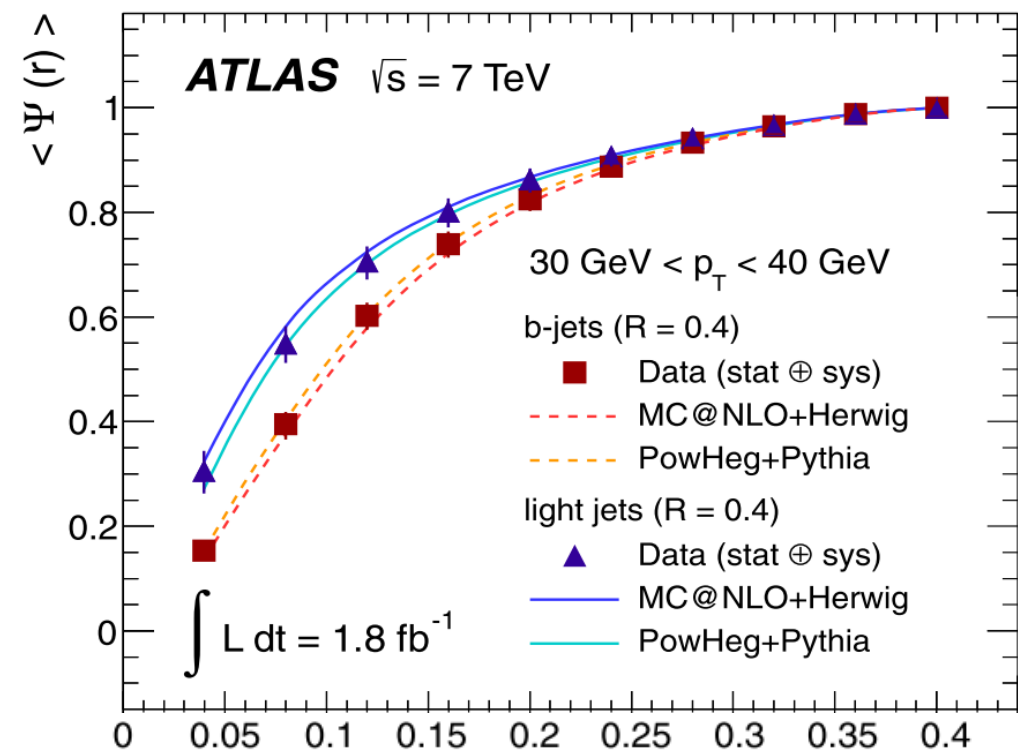
$$\Psi(r) = \frac{p_T(0, r)}{p_T(0, R)}; \quad r \leq R$$



Requirements for a direct observation of the dead cone



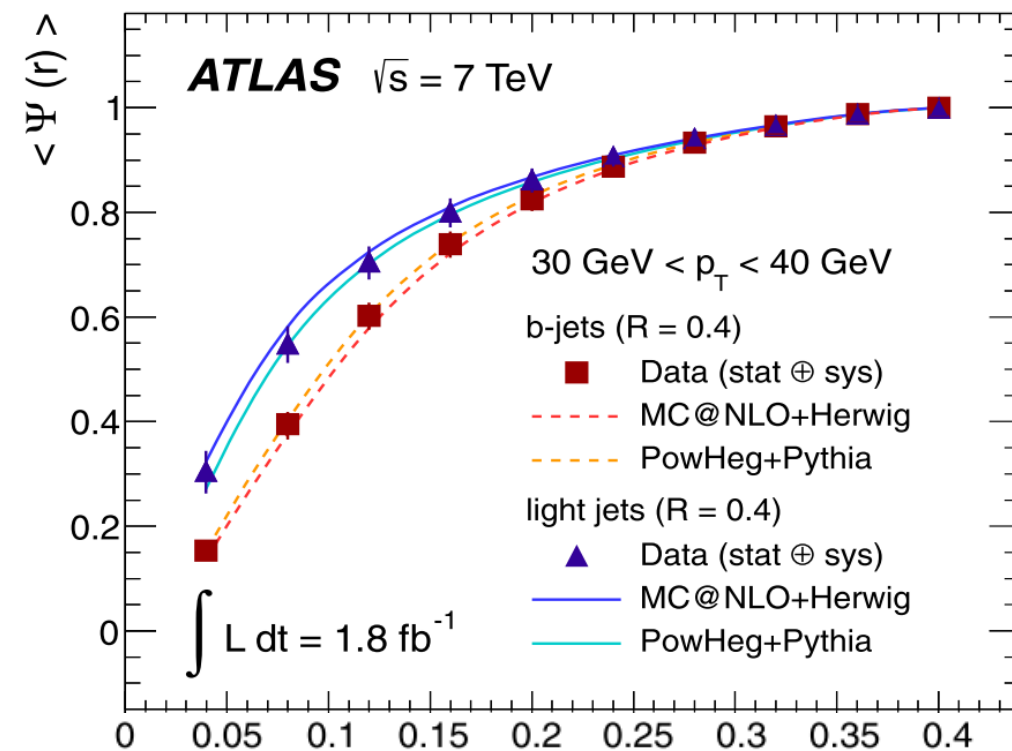
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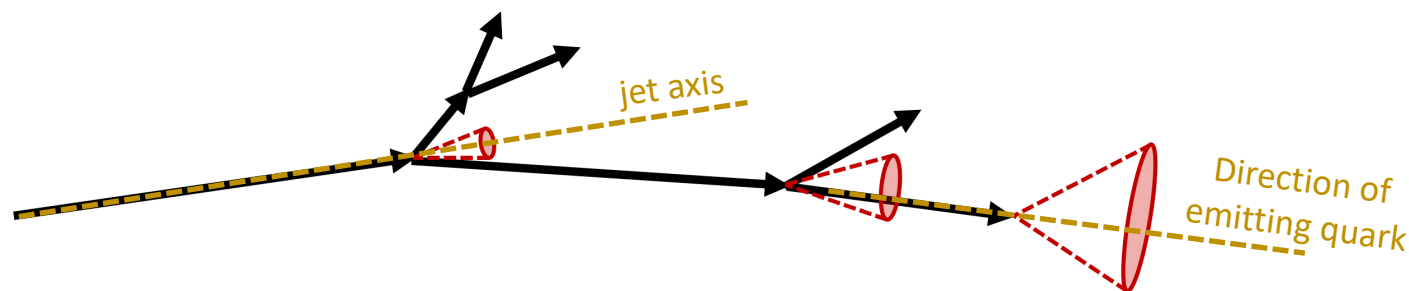
Requirements for a direct observation of the dead cone

The dead-cone angle appears in emissions at the partonic level - need to reconstruct the dynamically evolving direction of the heavy quark

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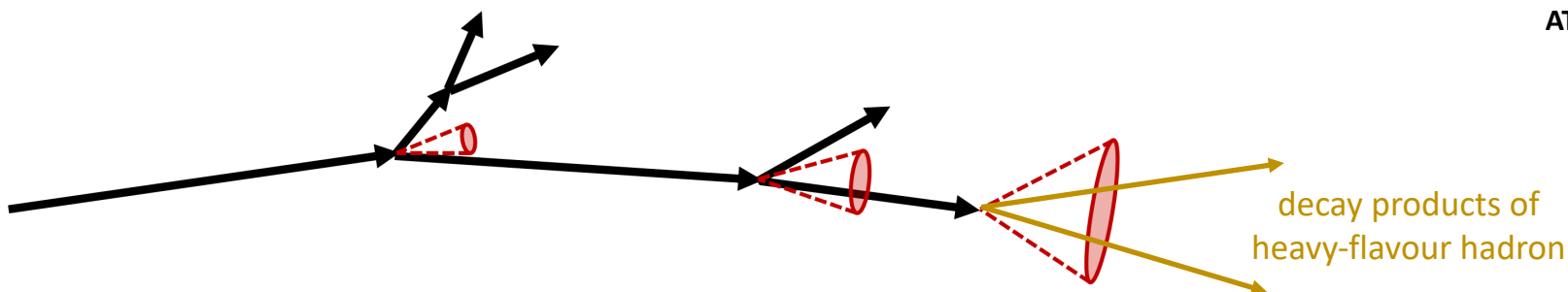
ATLAS Collaboration, Eur. Phys. J. C (2013) 73:2676



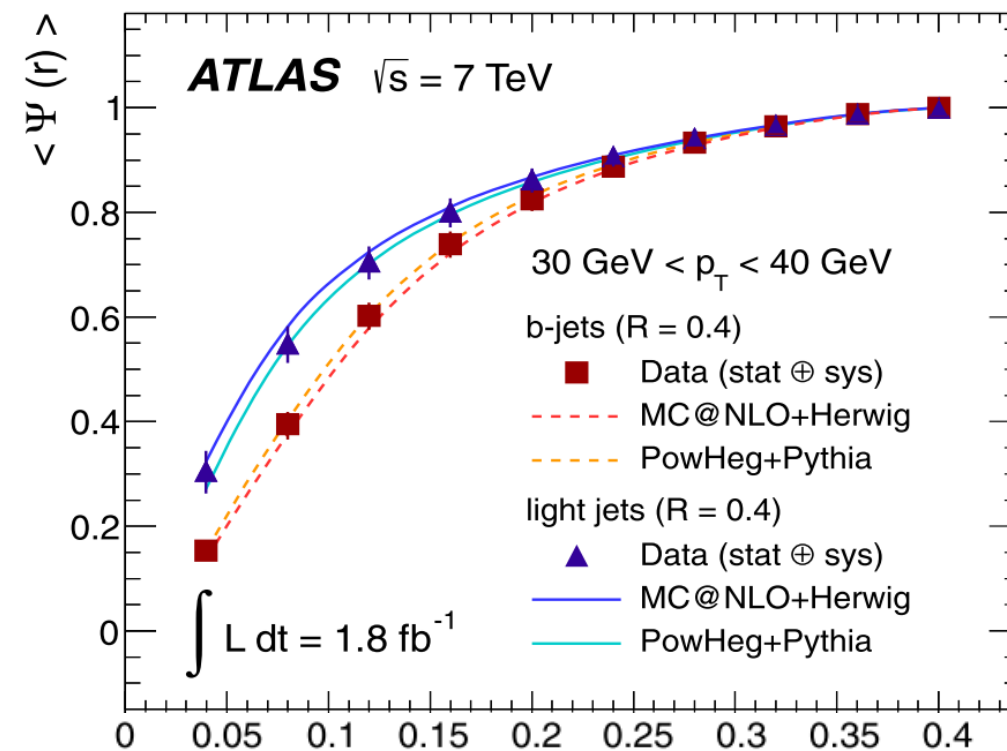
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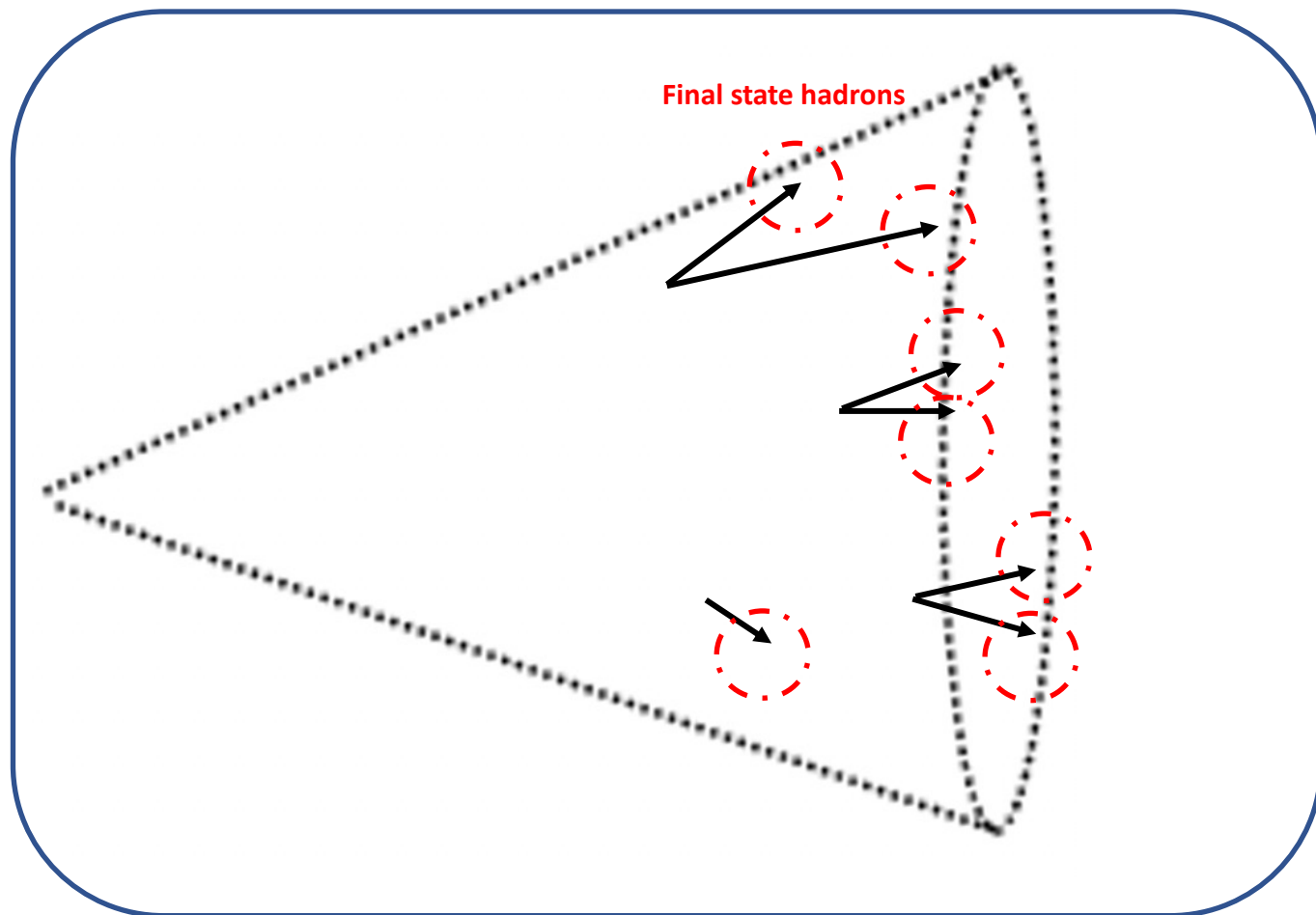
Sources such as the decay products of the heavy-flavour hadron can populate the dead-cone region



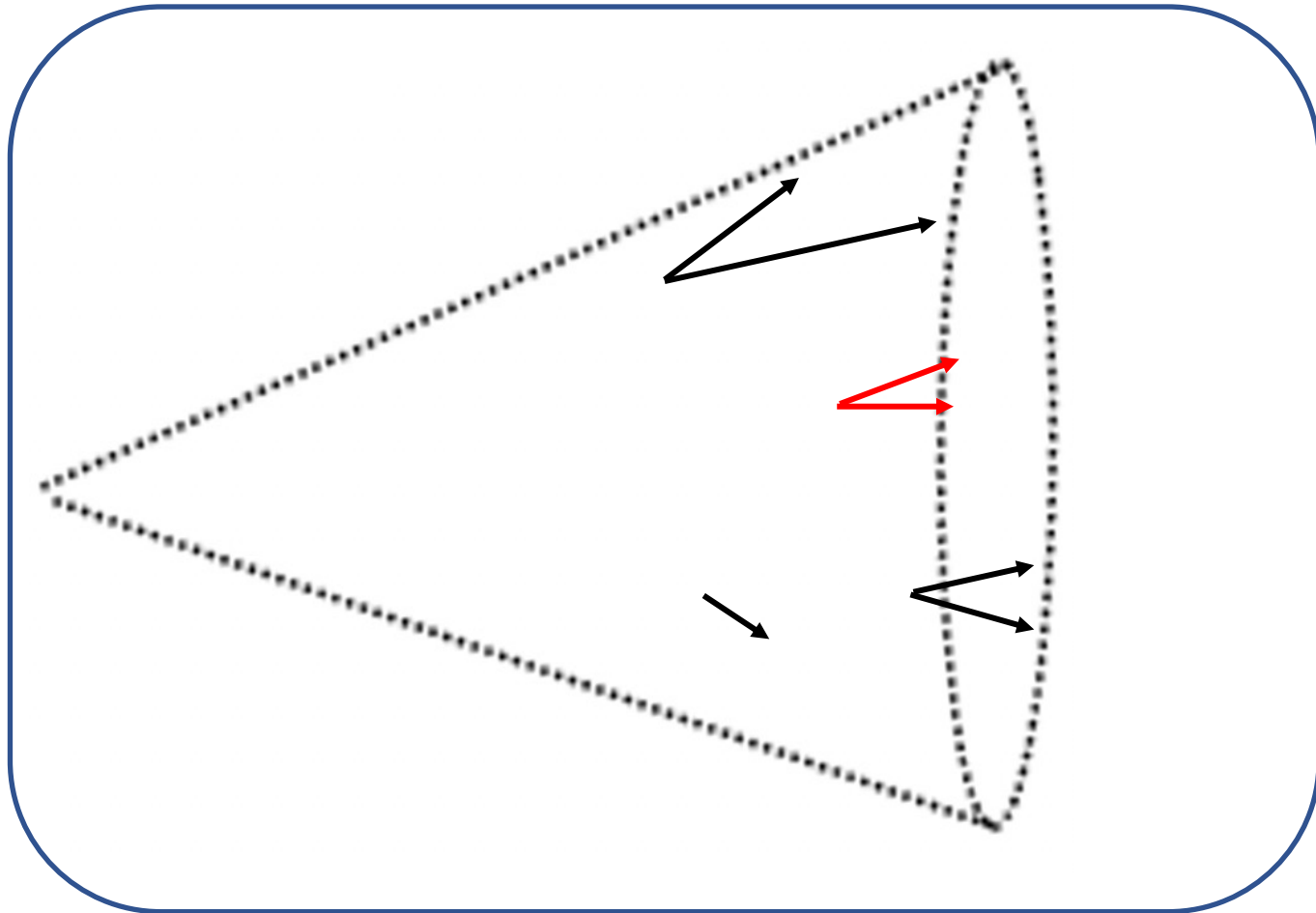
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ATLAS Collaboration, Eur. Phys. J. C (2013) 73:2676

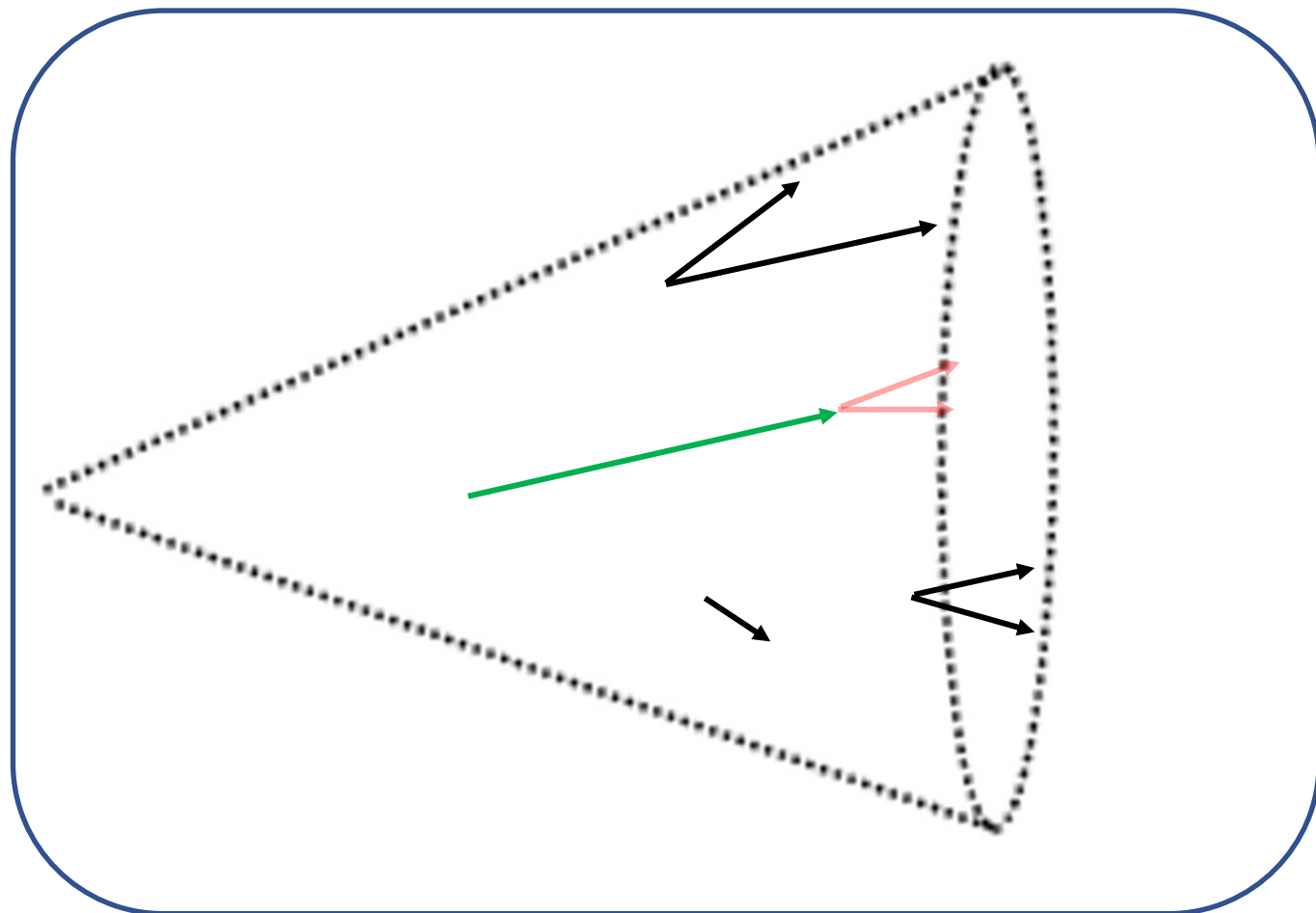


Jet clustering algorithms select the final state particles belonging to a given shower



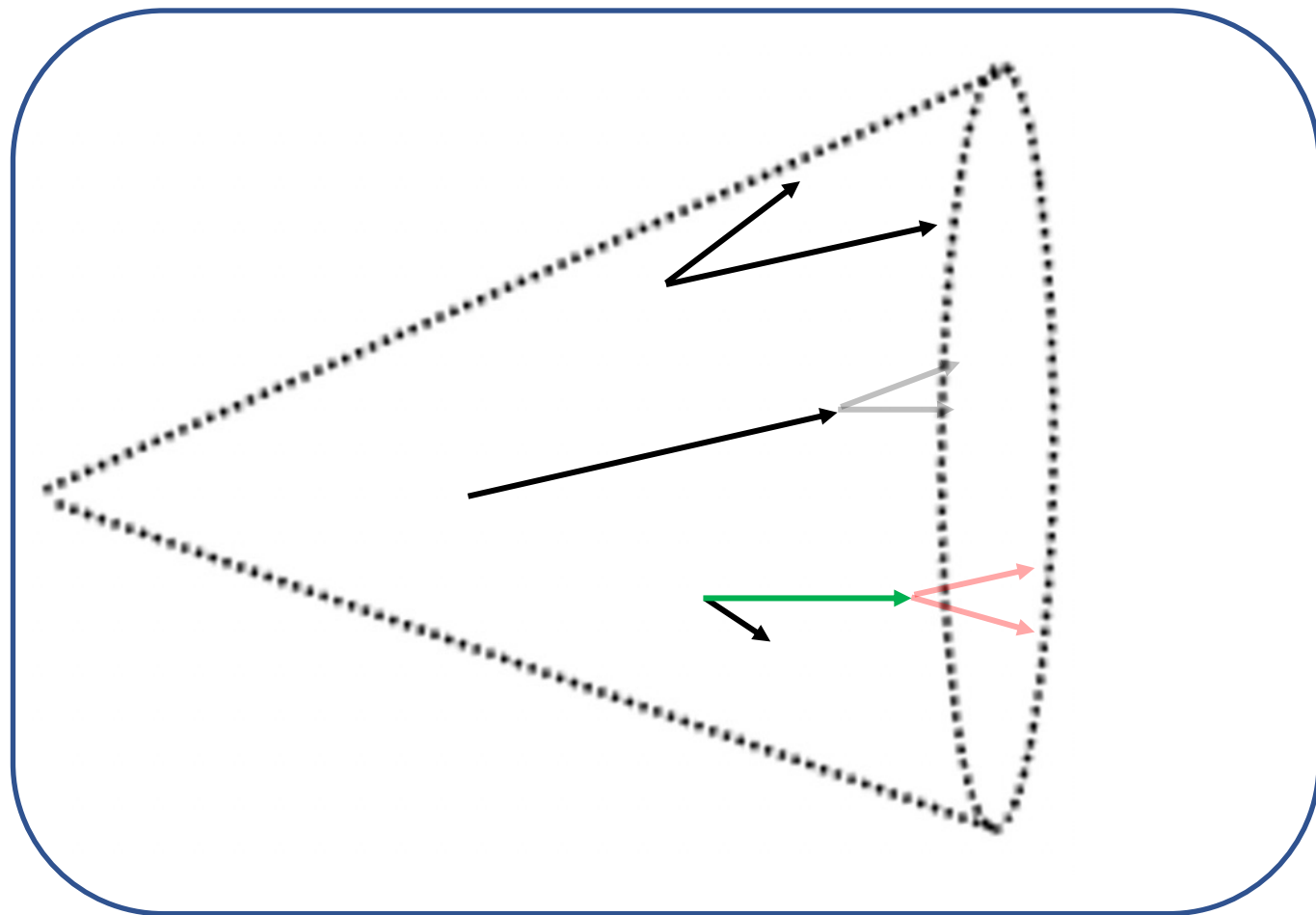
Jet clustering algorithms select the final state particles belonging to a given shower

Take advantage of the angular ordering of QCD emissions to reconstruct the shower using the C/A algorithm



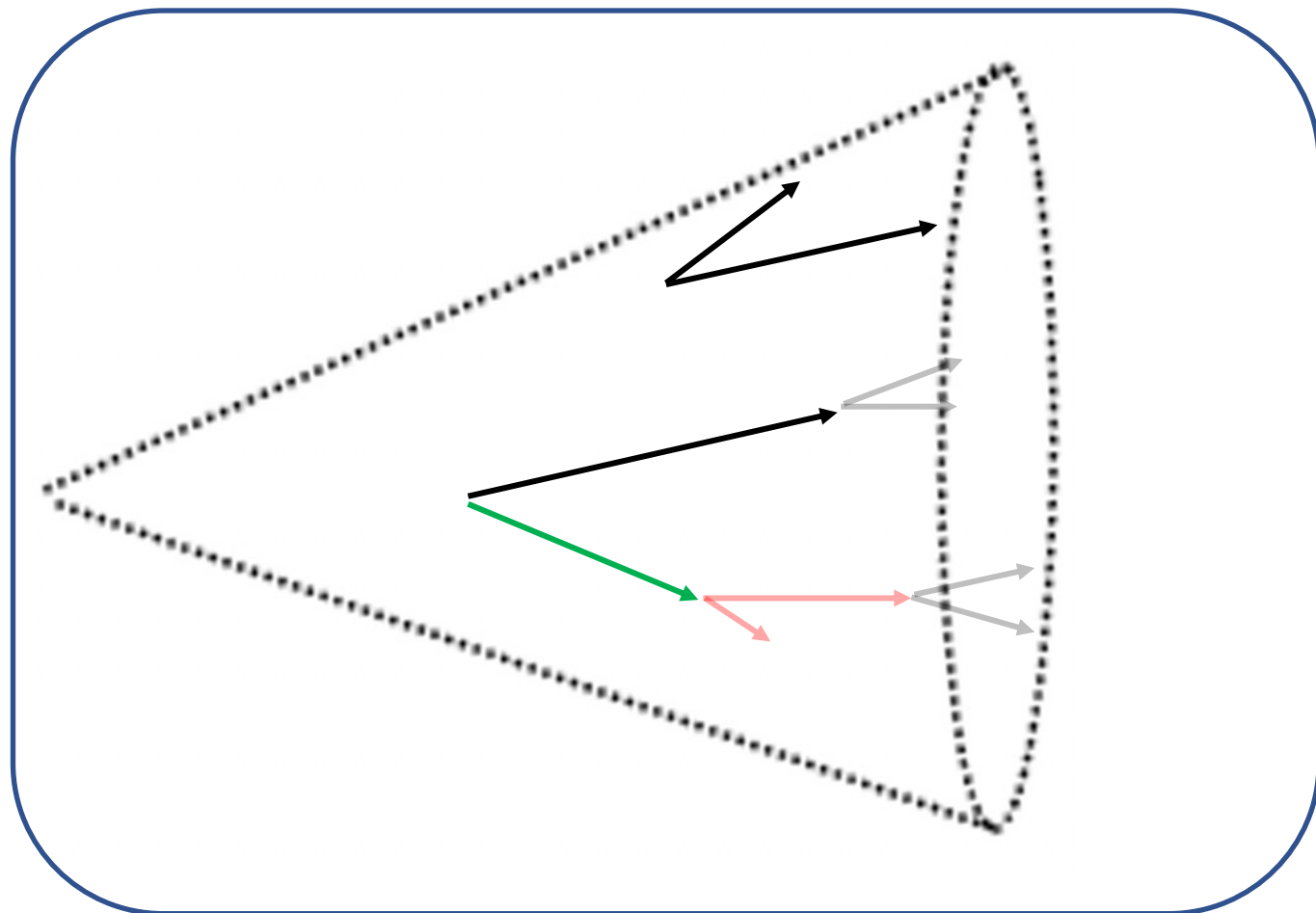
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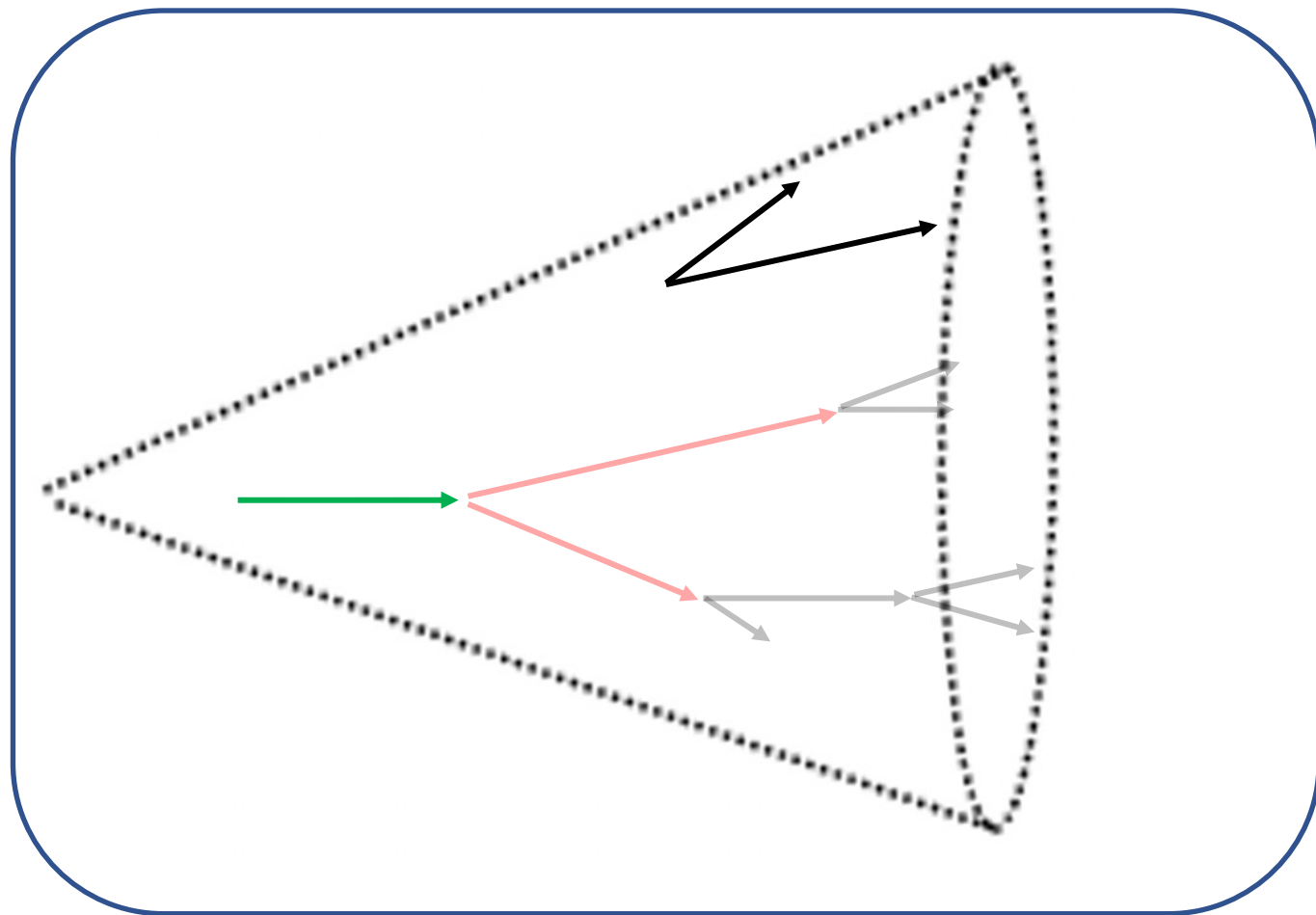
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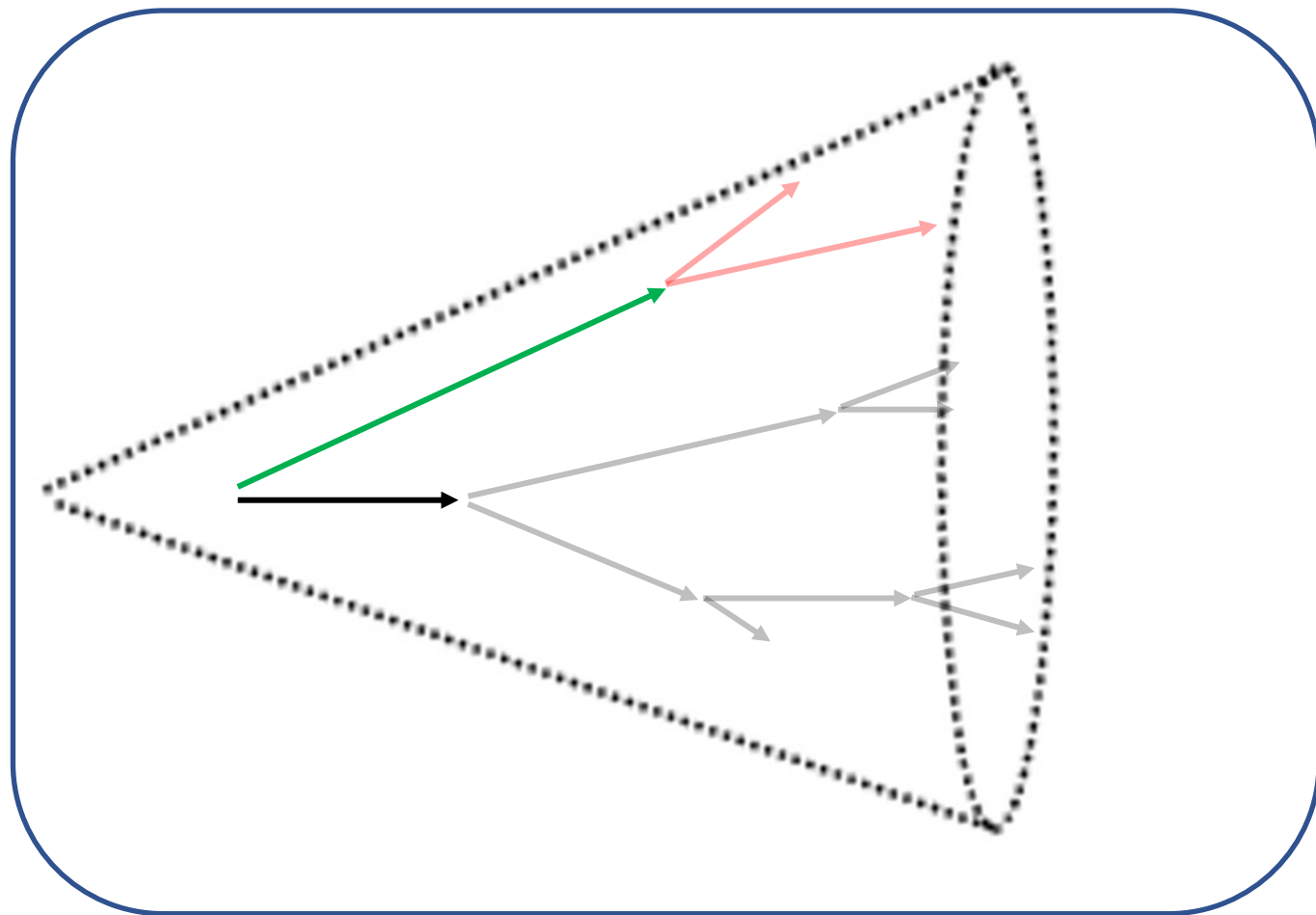
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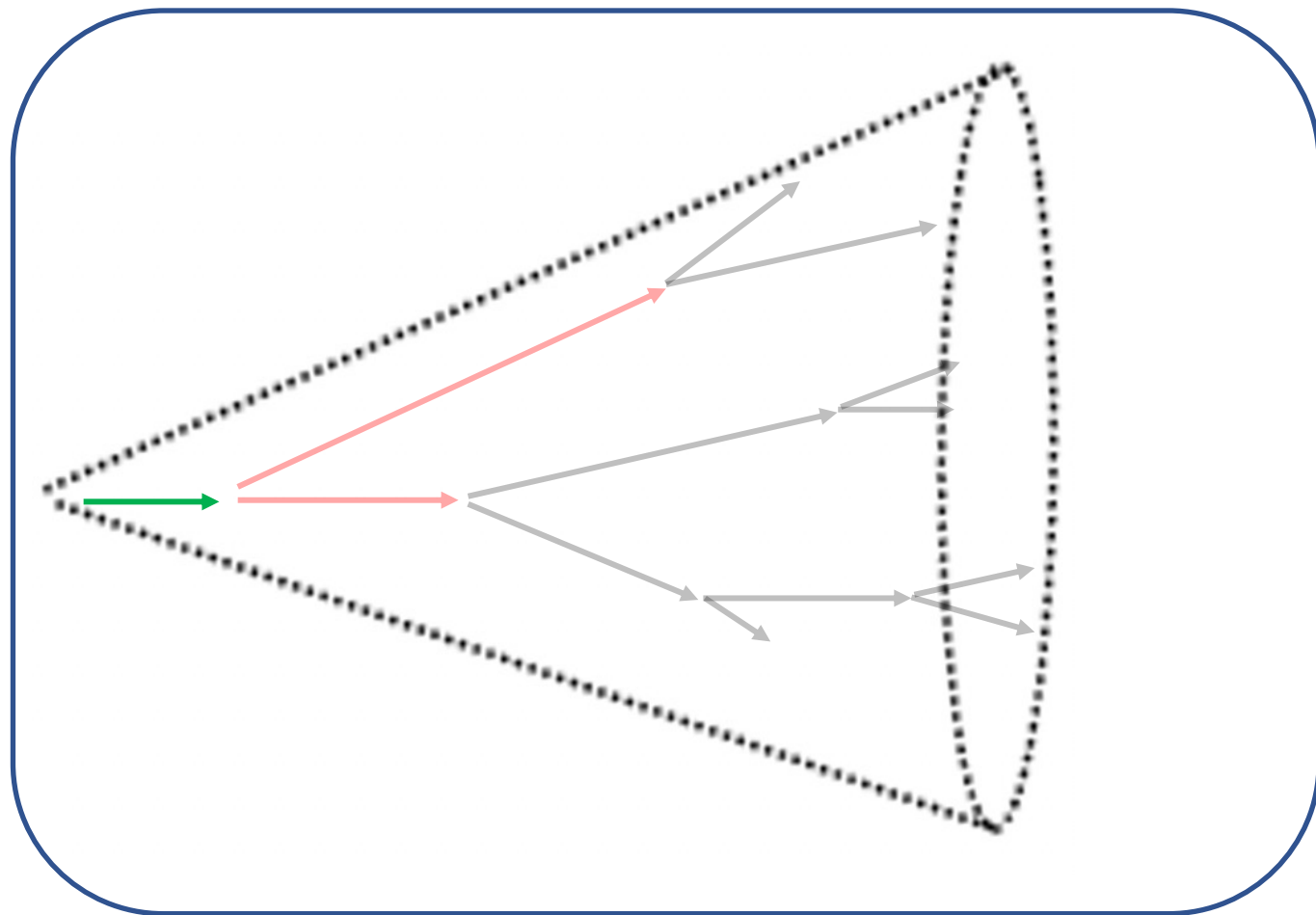
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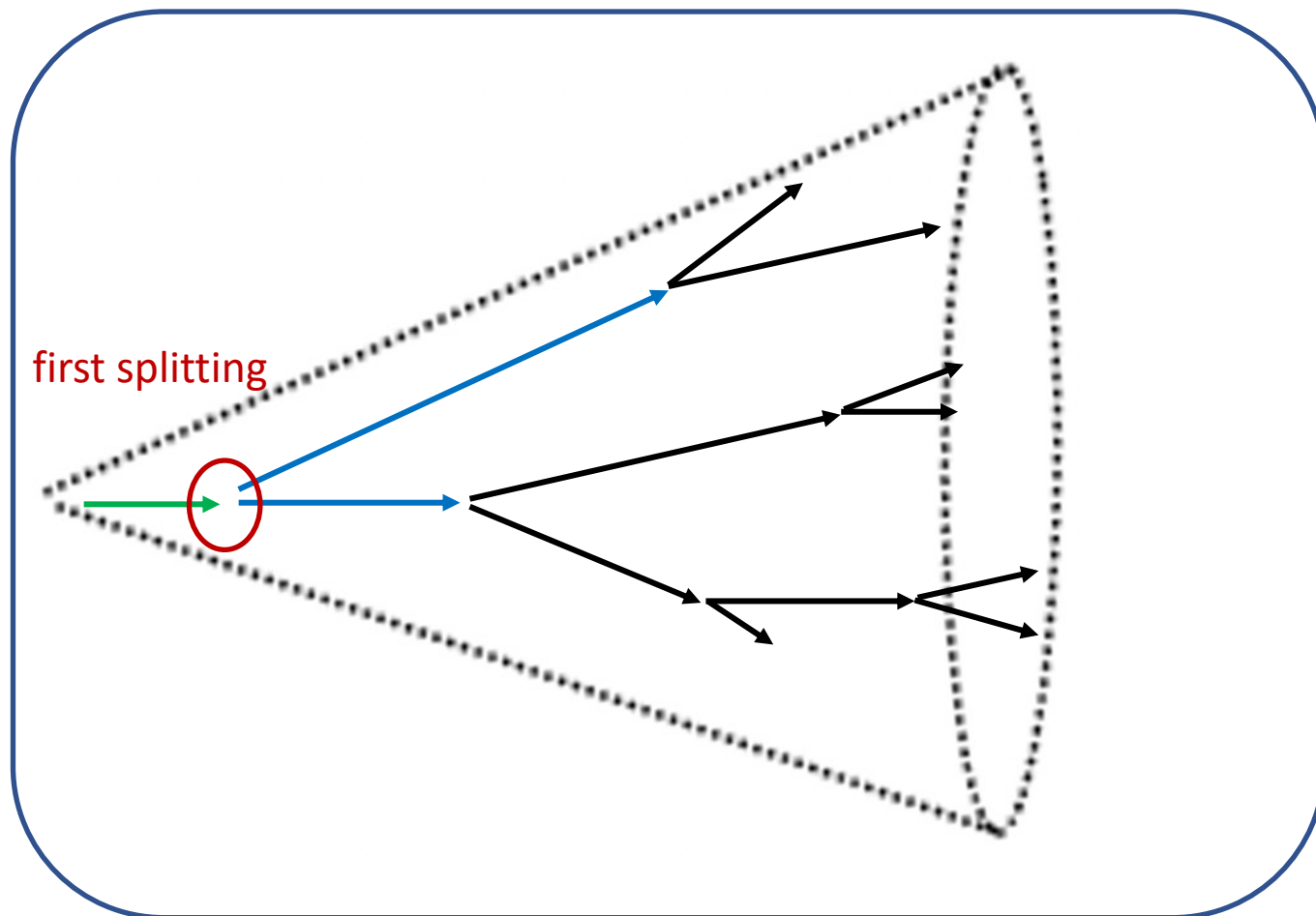
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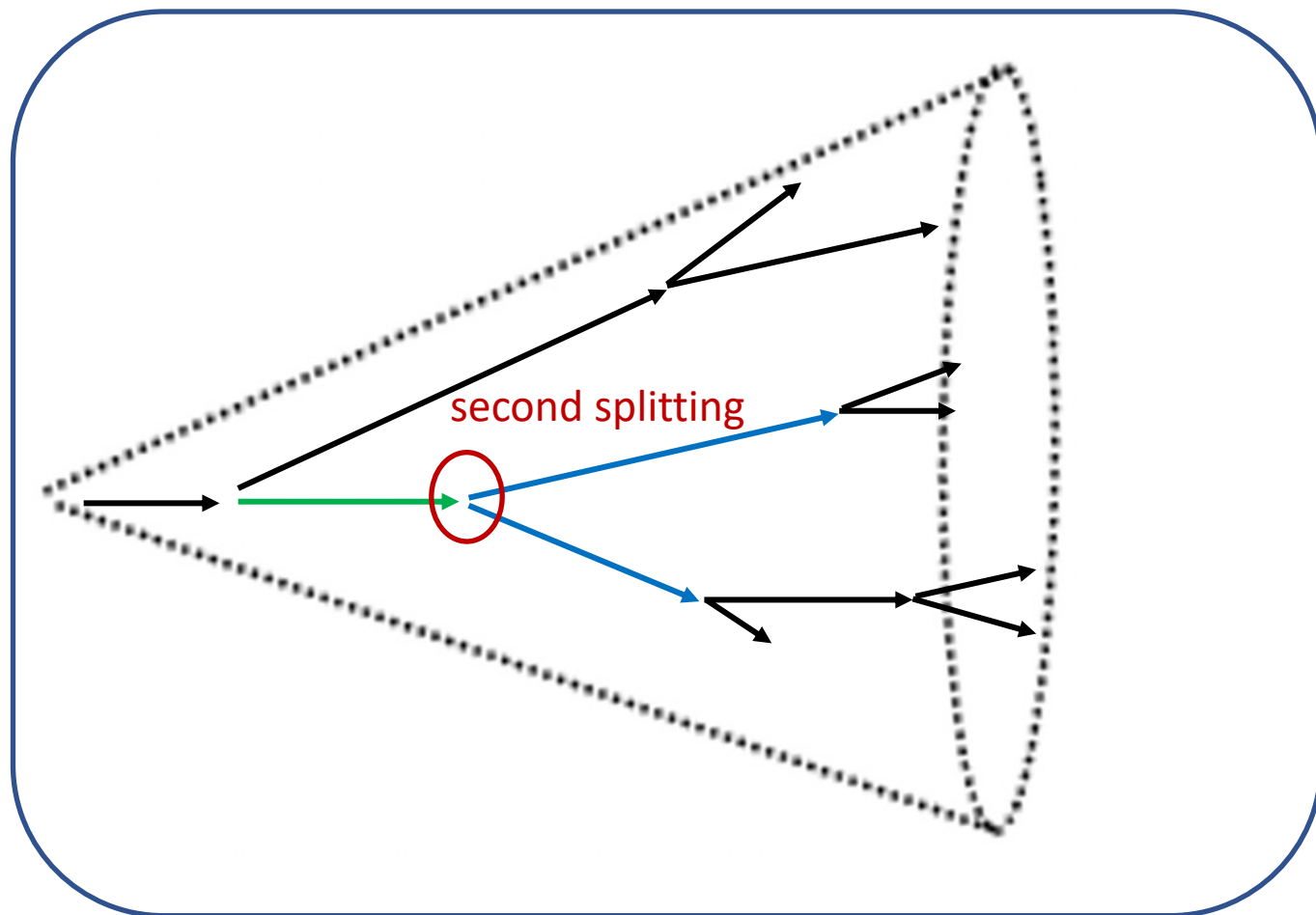
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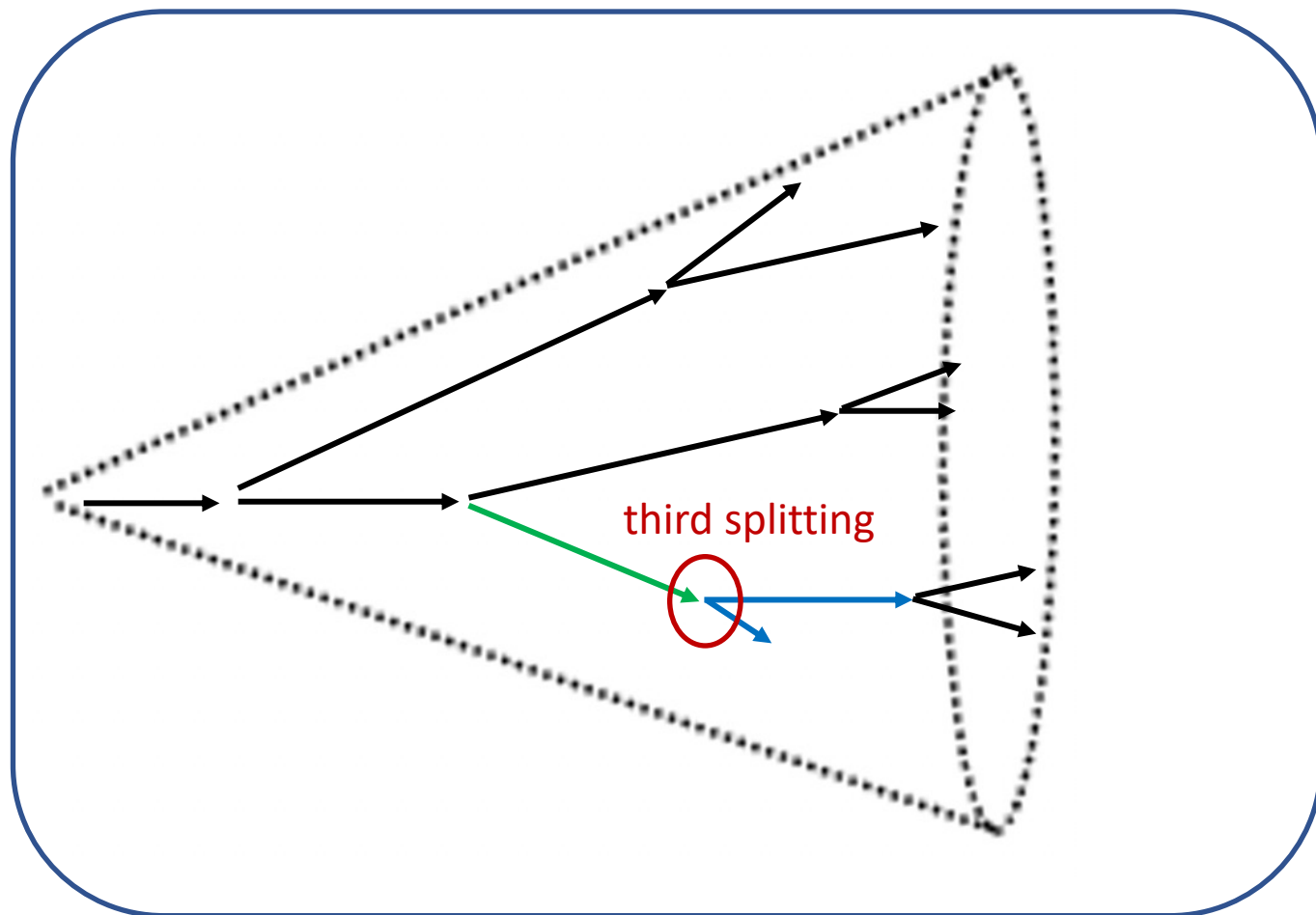
Unwind the clustering history and follow a given branch of emissions through the shower



Jet clustering algorithms select the final state particles belonging to a given shower

Take advantage of the angular ordering of QCD emissions to reconstruct the shower using the C/A algorithm

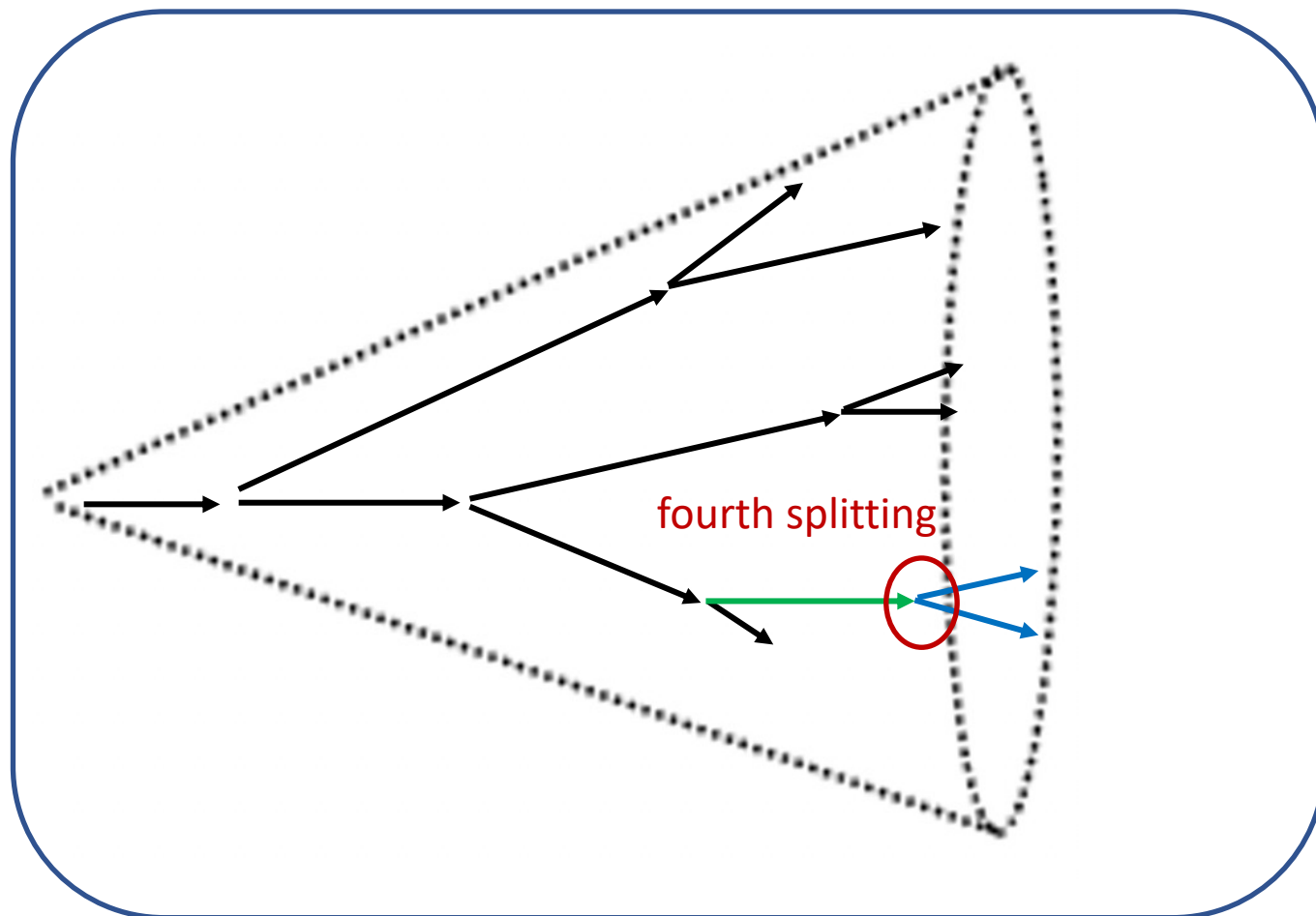
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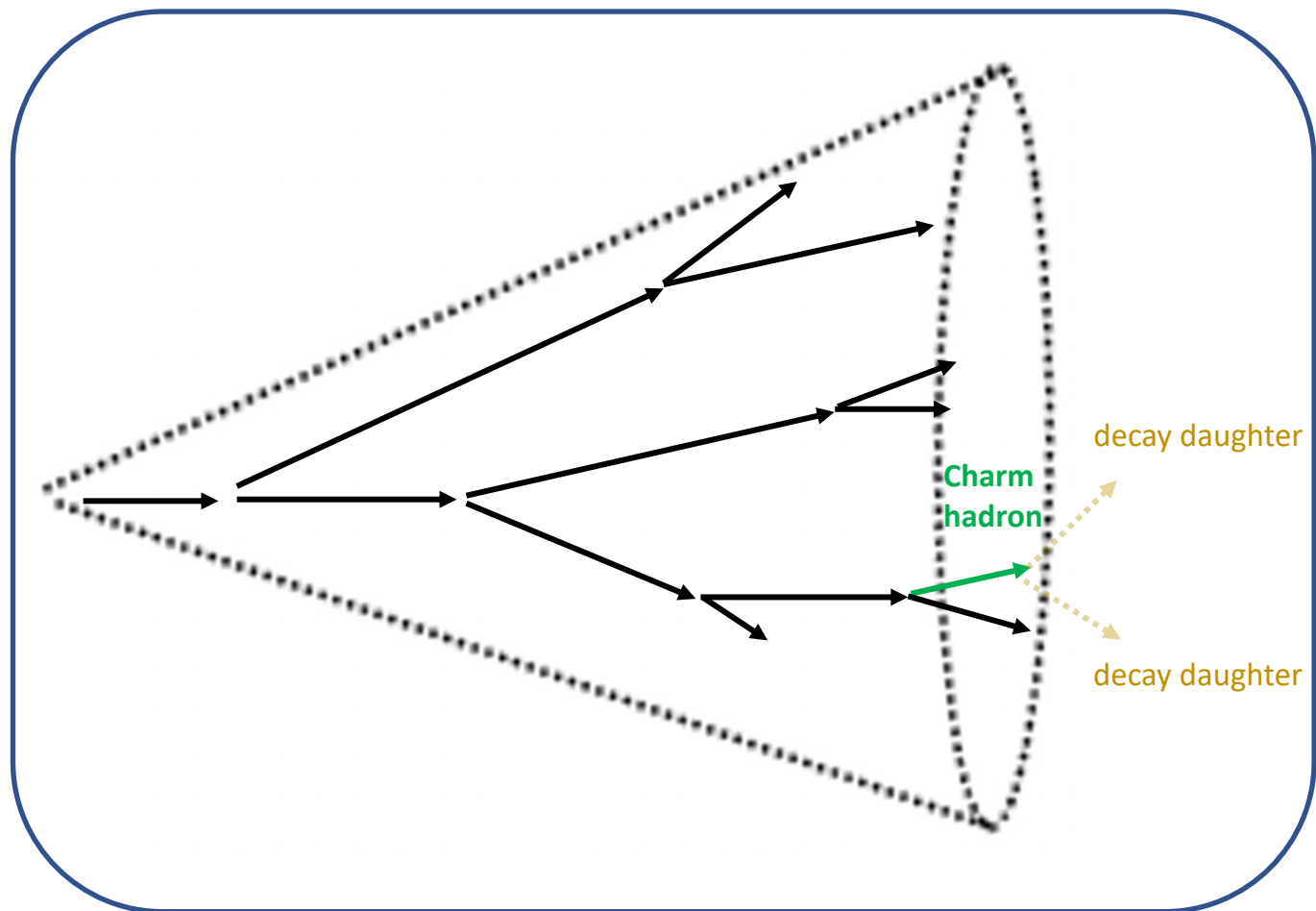


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Unwind the clustering history and follow a given branch of emissions through the shower

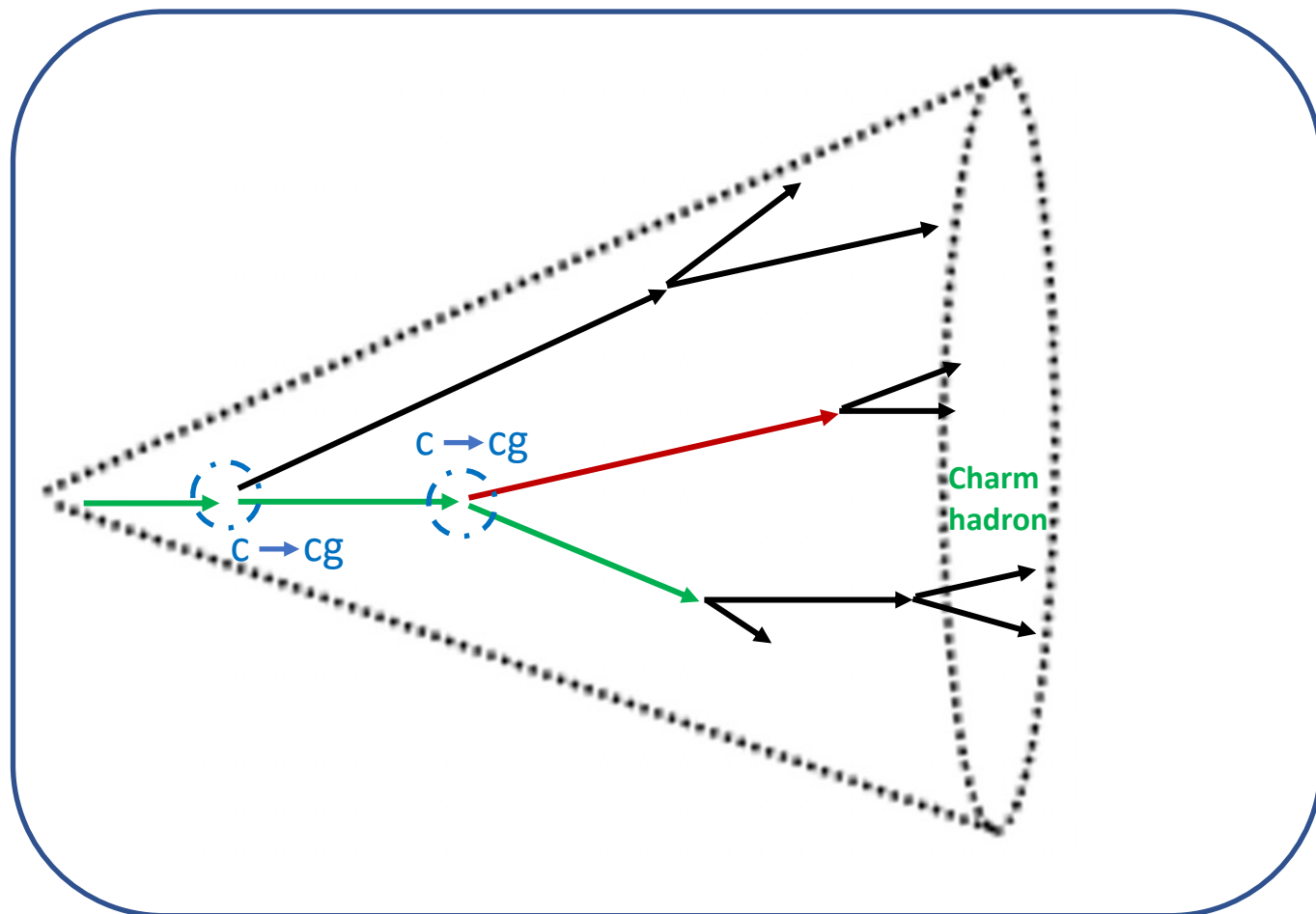
The parton flavour is not necessarily conserved along the followed branch



Heavy-flavour quarks retain their flavour throughout the shower evolution

Only $\text{Quark}_{\text{HF}} \rightarrow \text{Quark}_{\text{HF}} + \text{Gluon emissions}$

Due to their large mass HF quarks production is suppressed during hadronisation

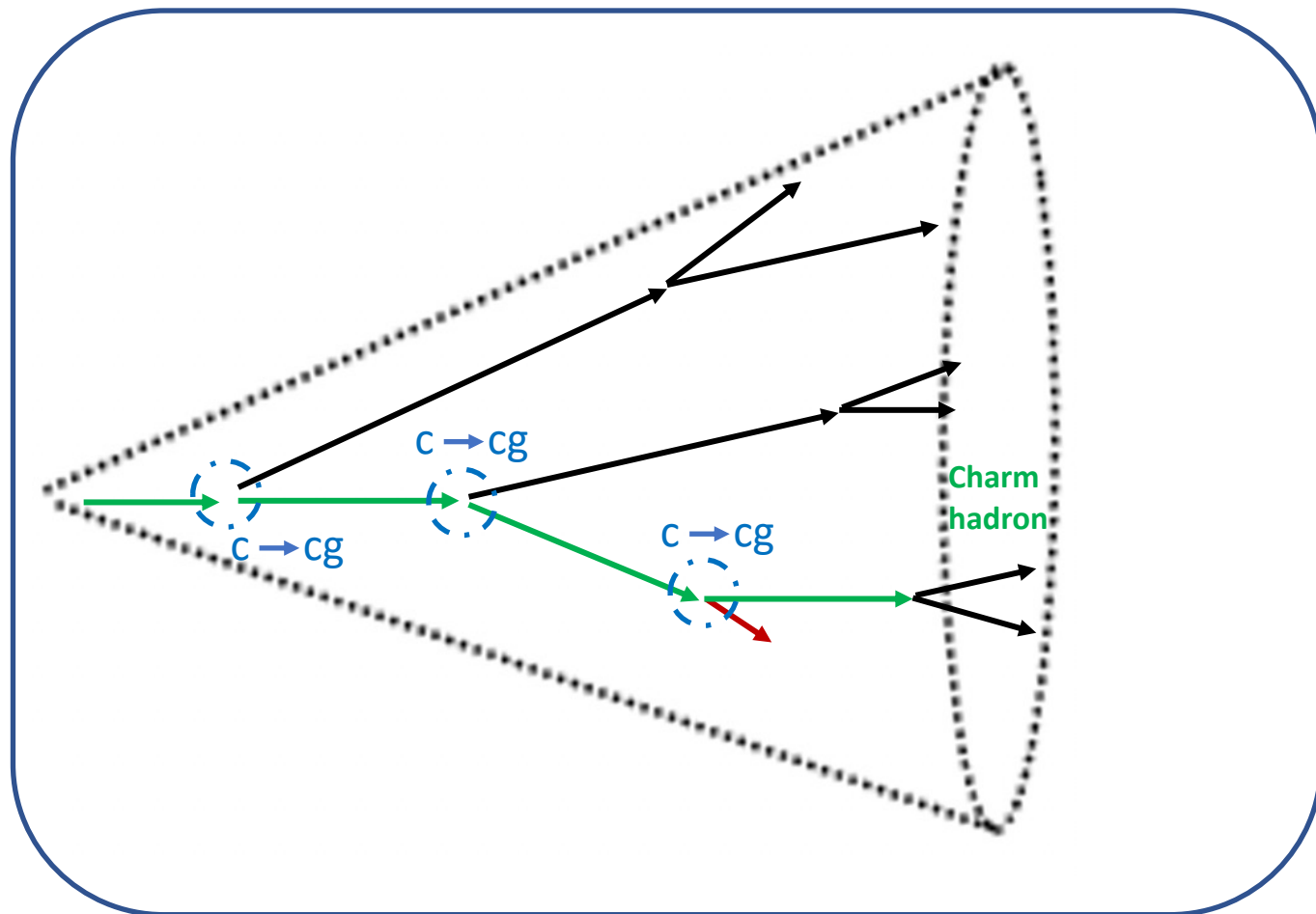


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The shower of a jet tagged via a reconstructed heavy-flavour hadron would contain a traceable branch of $\text{Quark}_{\text{HF}} \rightarrow \text{Quark}_{\text{HF}} + \text{Gluon emissions}$

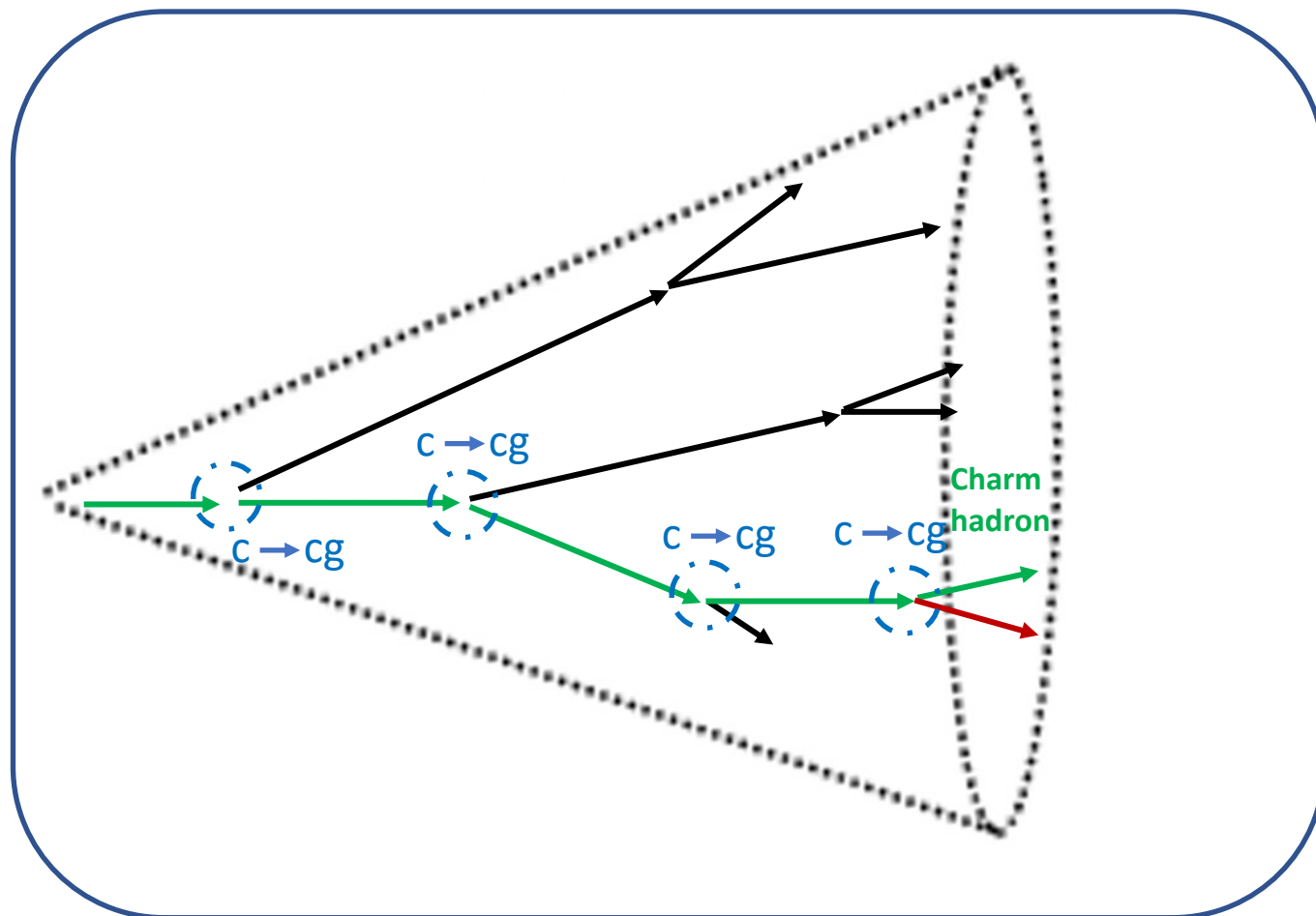


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L. Cunqueiro, M. Ploskon, Phys. Rev. D 99, 074027 (2019)

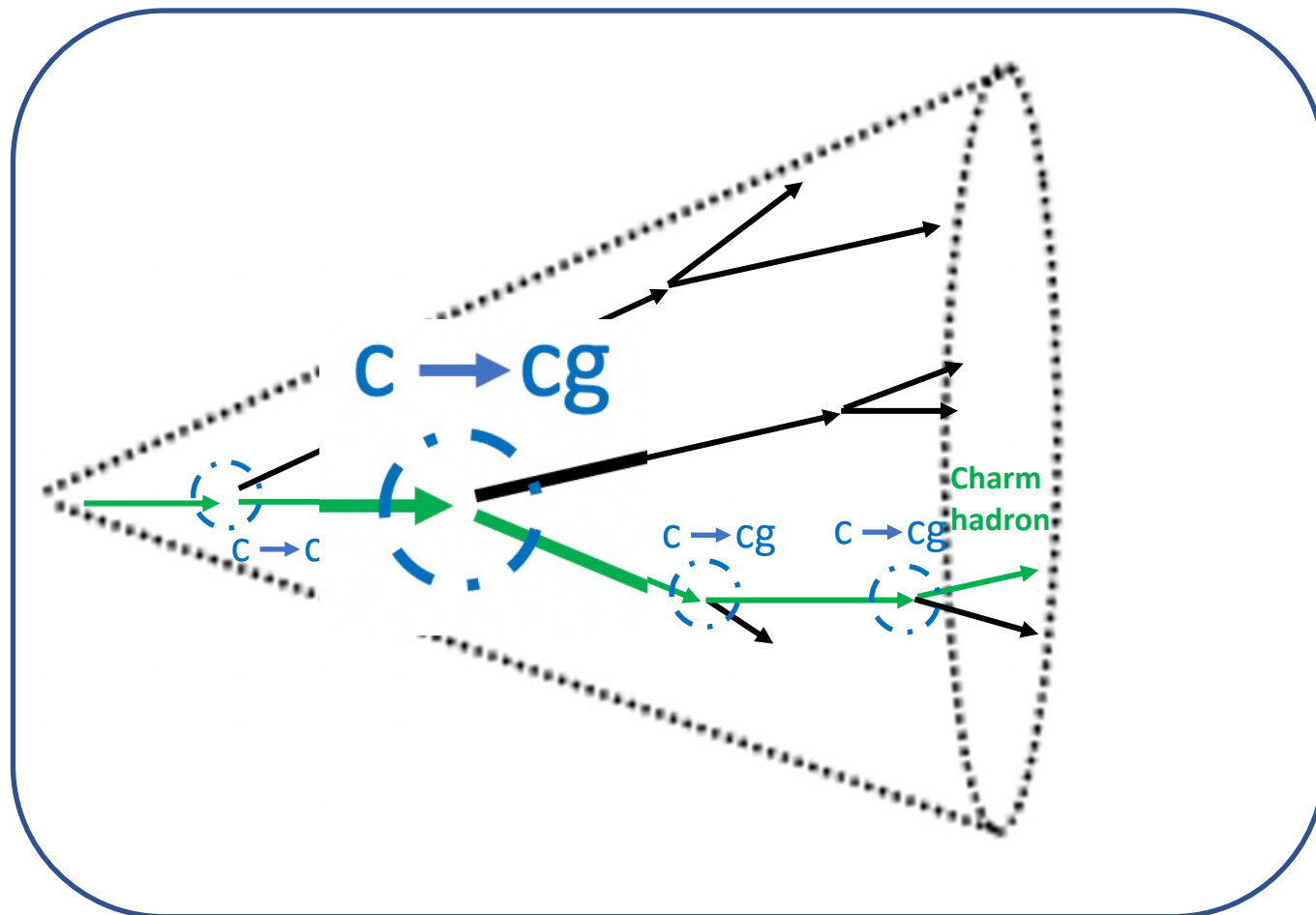
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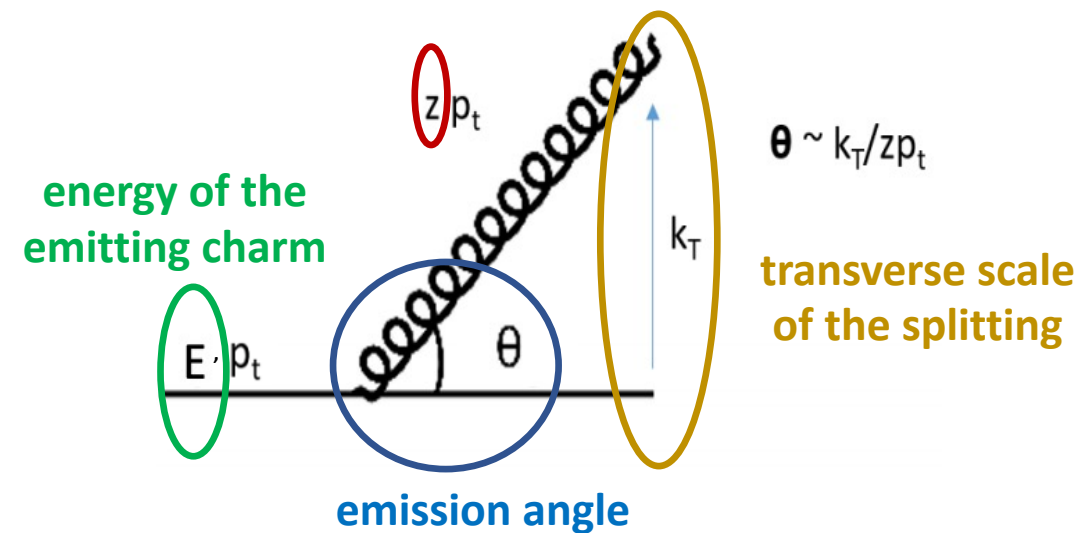
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Can follow a quark sample of emissions



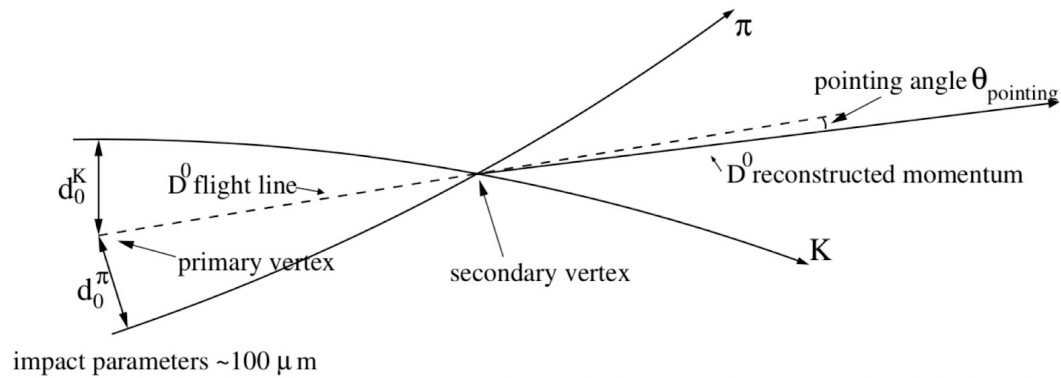
fraction of momentum carried by the emitted gluon



The angle of each gluon emission directly probes the dead-cone effect

The energy of the charm quark at each emission point directly sets the size of the dead-cone region

The transverse scale of each splitting will be used to suppress non-perturbative effects

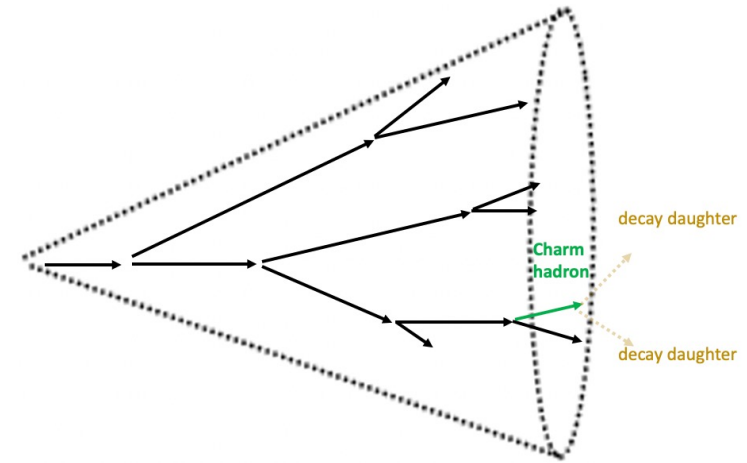


Reconstruct the D⁰ meson

D⁰ mesons are reconstructed through the D⁰ → K⁻ π⁺ (and charge conjugate) decay channel

Selections on decay topology and PID are used to identify D⁰-meson candidates

$$2 < p_T^{D^0} < 36 \text{ GeV}/c$$



Jet Finding

The analysis is performed using Run 2 pp collisions at

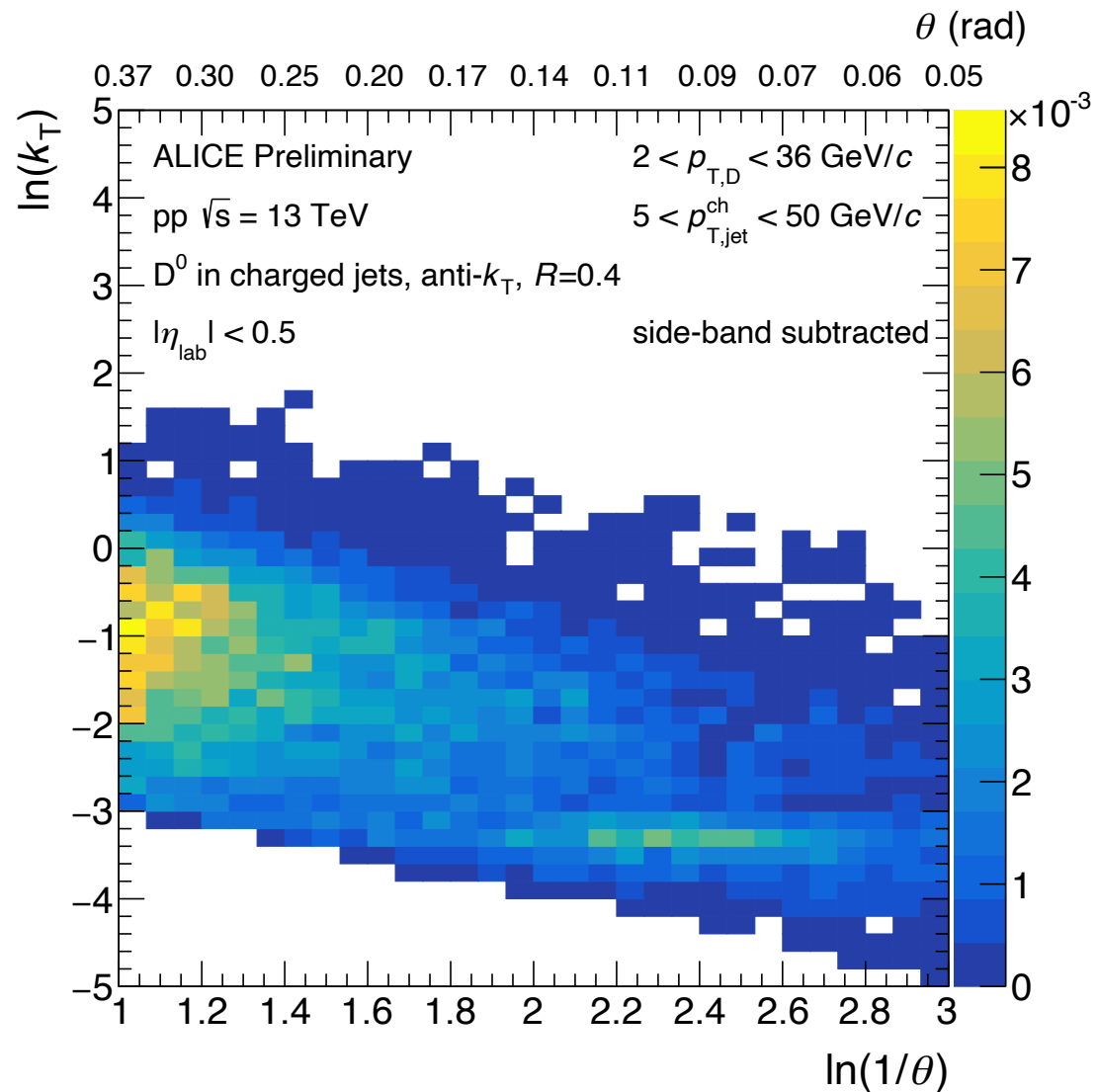
$$\sqrt{s} = 13 \text{ TeV} \quad \mathcal{L}_{\text{int}} = 25 \text{ nb}^{-1}$$

The D⁰ decay daughters are replaced by the D⁰ prior to jet finding

Jet finding is performed using the anti- k_T algorithm with R = 0.4

$$5 < p_T^{\text{jet}} < 50 \text{ GeV}/c$$

Jets with a D⁰-meson candidate amongst their constituents are selected

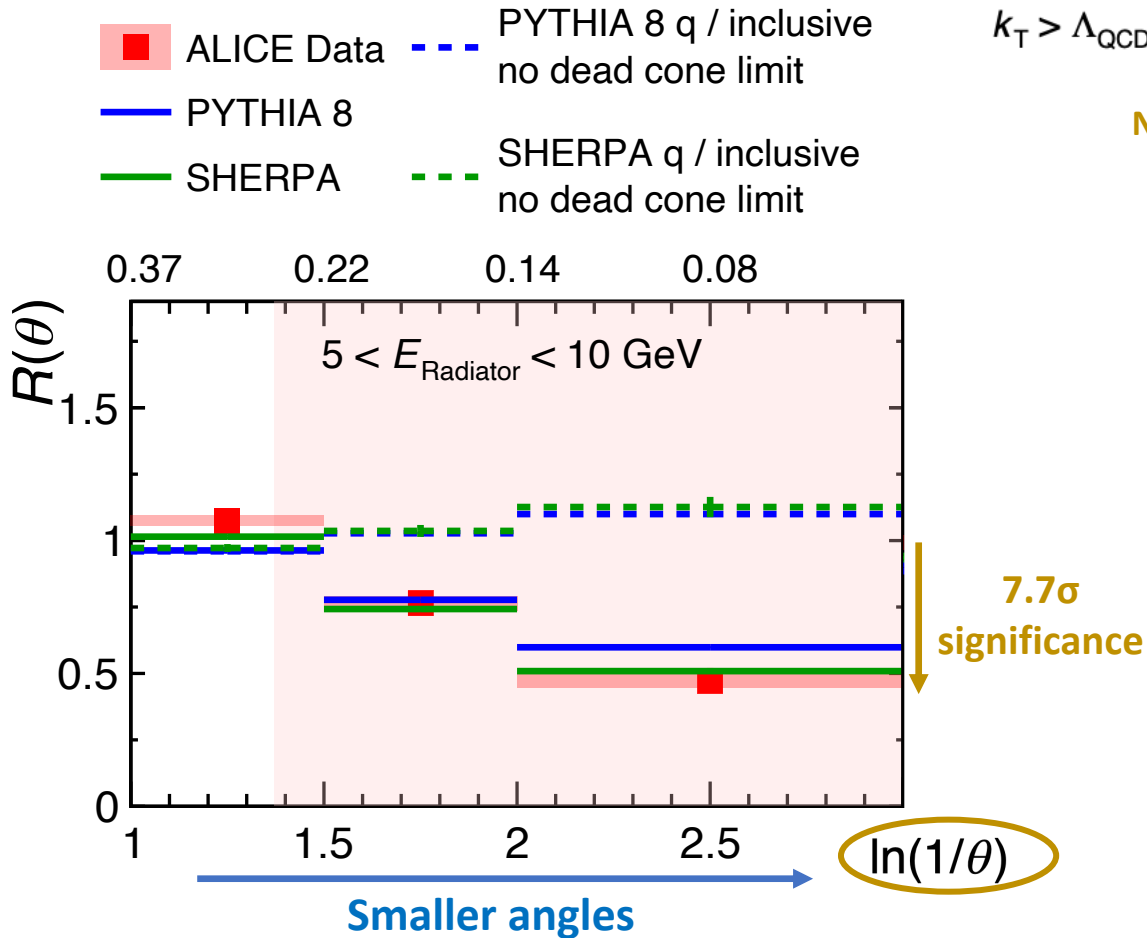


Primary Lund plane of c->cg emissions

Measurement of a Lund plane of a particular splitting flavour in QCD

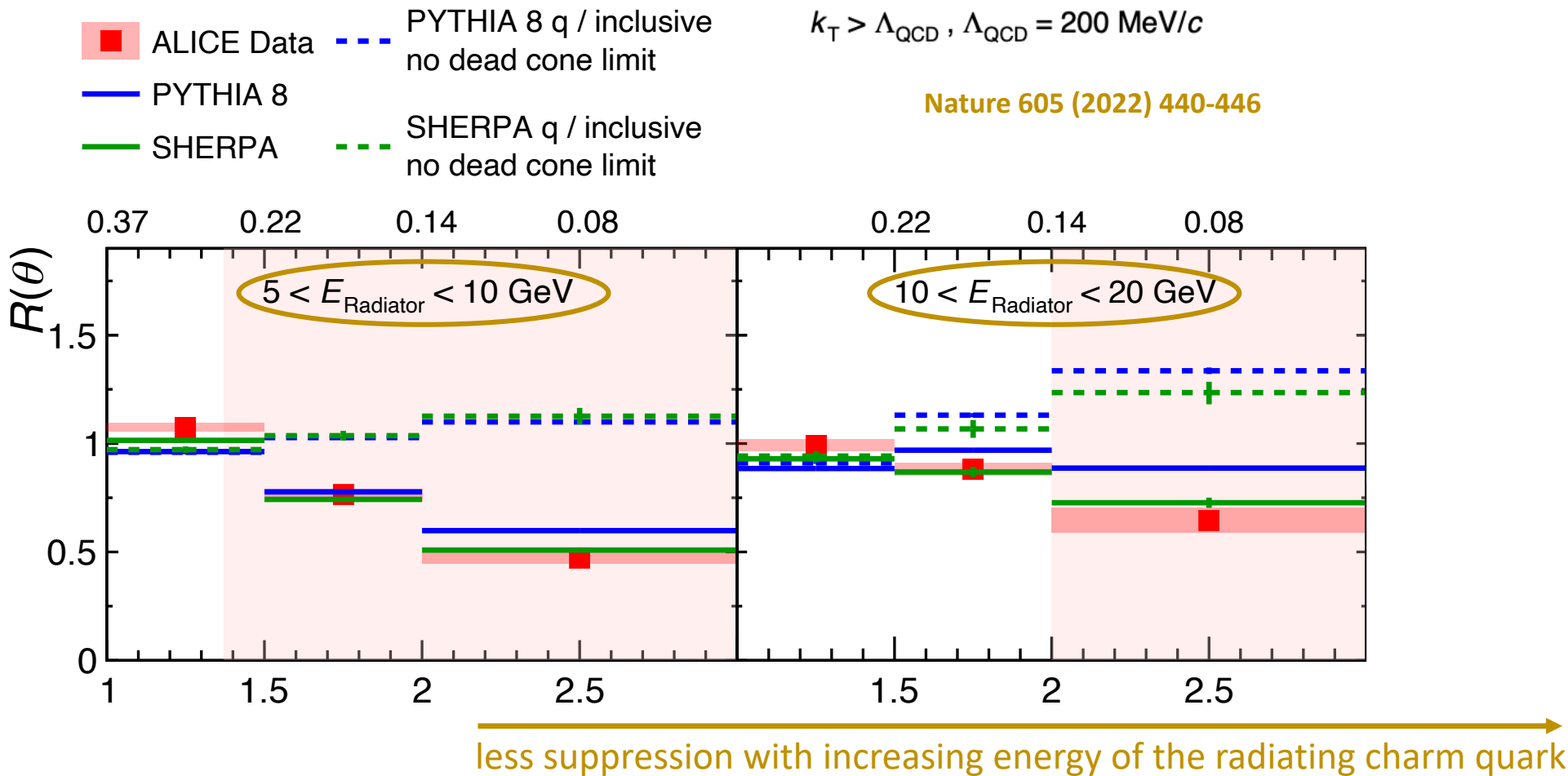
$$R(\theta) = \frac{1}{N^{\text{D}^0 \text{ jets}}} \frac{dn^{\text{D}^0 \text{ jets}}}{d\ln(1/\theta)} \bigg/ \frac{1}{N^{\text{inclusive jets}}} \frac{dn^{\text{inclusive jets}}}{d\ln(1/\theta)} \bigg|_{k_T, E_{\text{Radiator}}}$$

Compare the angular distribution of charm-quark emissions to those of light quarks and gluons



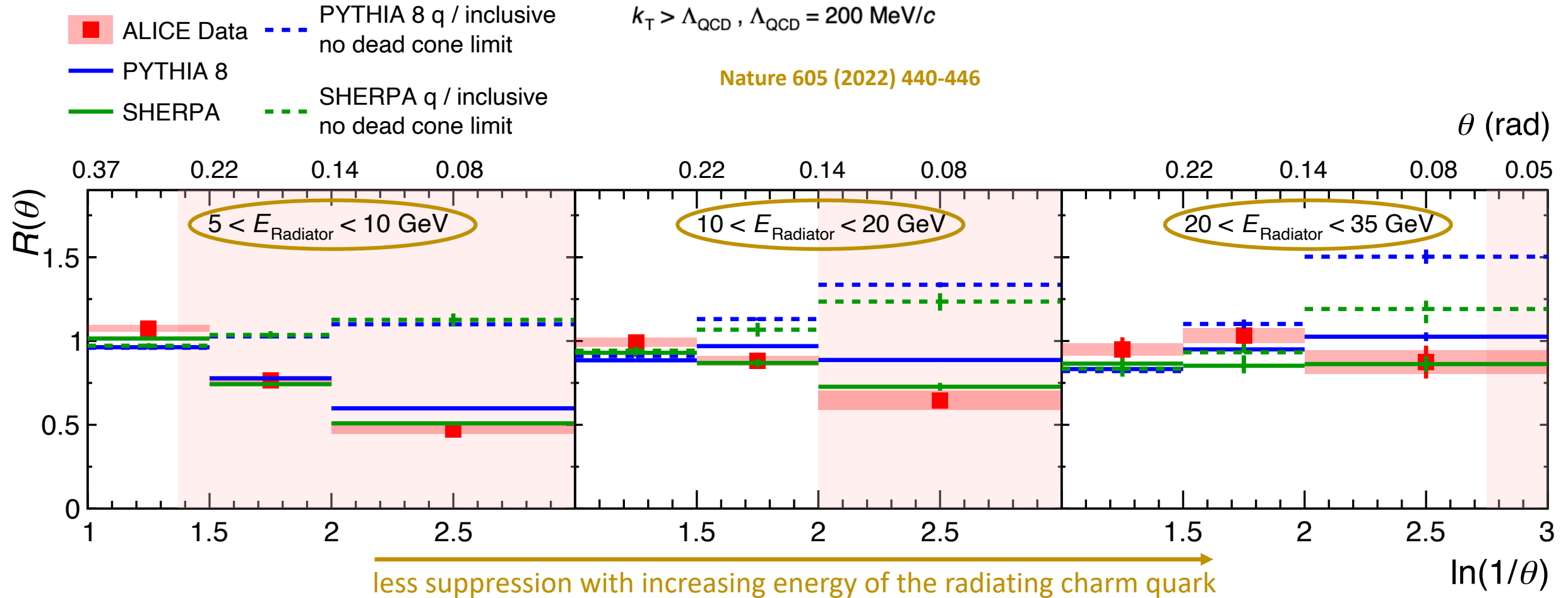
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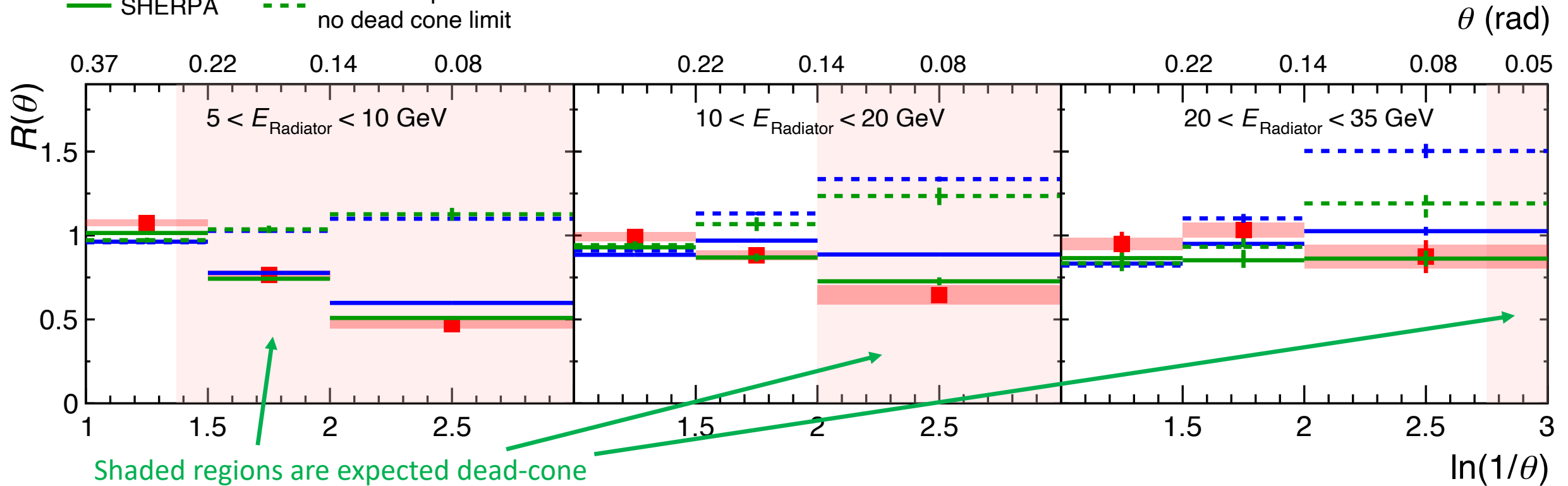
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Compare the angular distribution of charm-quark emissions to those of light quarks and gluons

- ALICE Data
- PYTHIA 8
- SHERPA
- - - PYTHIA 8 q / inclusive no dead cone limit
- - - SHERPA q / inclusive no dead cone limit

$k_T > \Lambda_{\text{QCD}}, \Lambda_{\text{QCD}} = 200 \text{ MeV}/c$

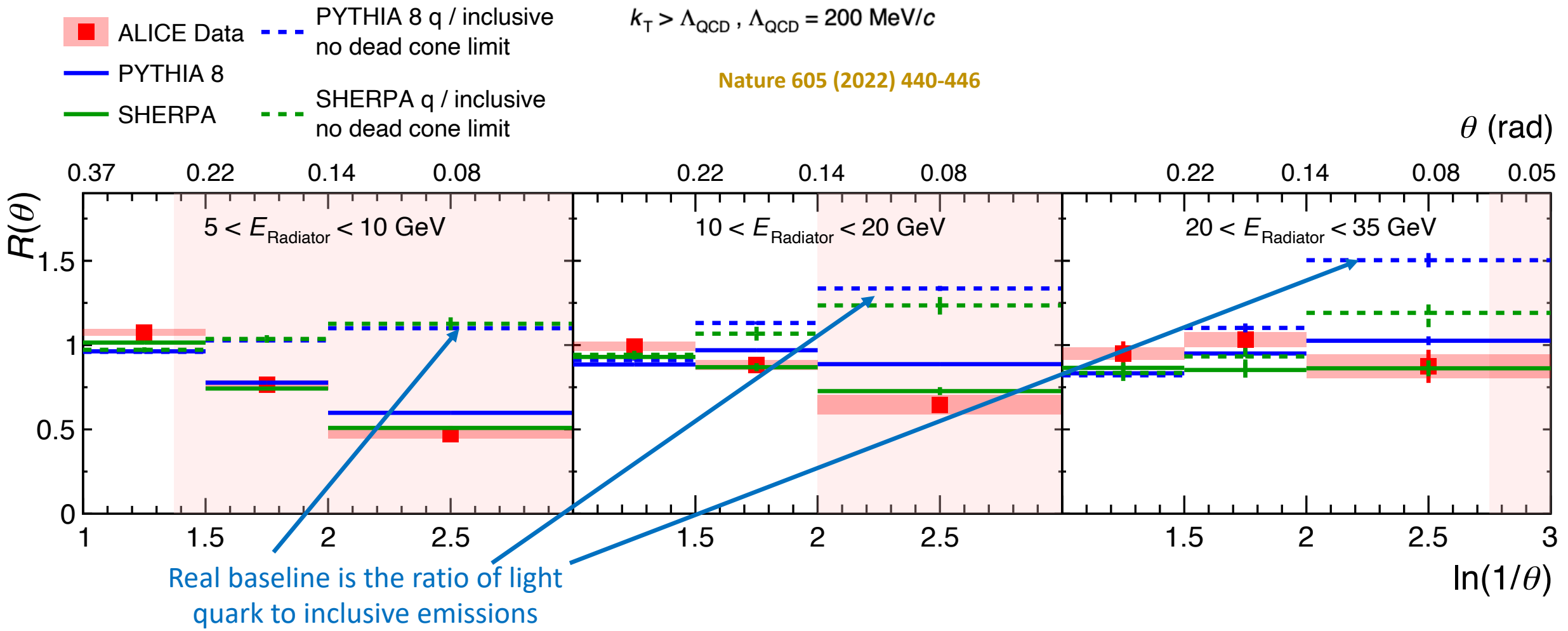
Nature 605 (2022) 440-446



Shaded regions are expected dead-cone region for a charm mass of 1.275 GeV/c

$$R(\theta) = \frac{1}{N^{\text{D}^0 \text{ jets}}} \frac{dn^{\text{D}^0 \text{ jets}}}{d \ln(1/\theta)} \bigg/ \frac{1}{N^{\text{inclusive jets}}} \frac{dn^{\text{inclusive jets}}}{d \ln(1/\theta)} \bigg|_{k_T, E_{\text{Radiator}}}$$

Compare the angular distribution of charm-quark emissions to those of light quarks and gluons

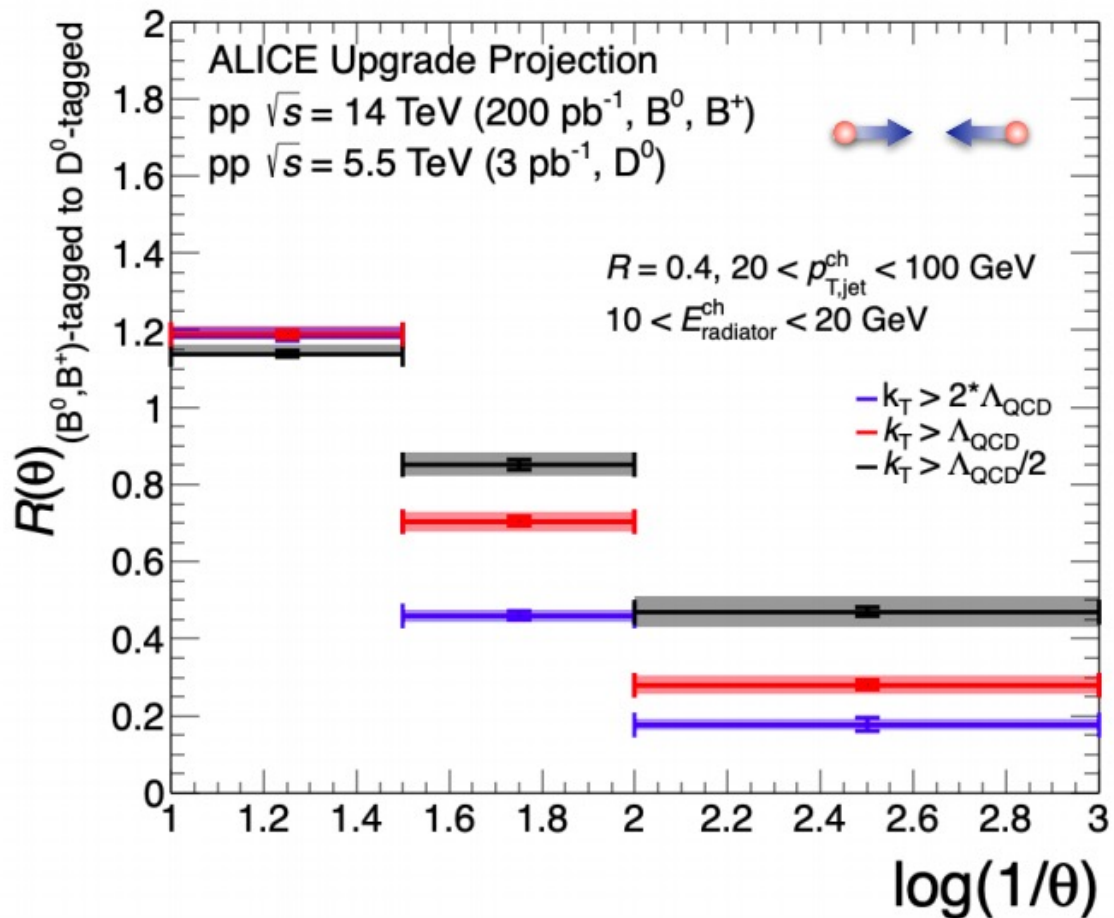


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Compare the angular distribution of charm-quark emissions to those of light quarks and gluons

Dead cone of B⁺ - tagged jets

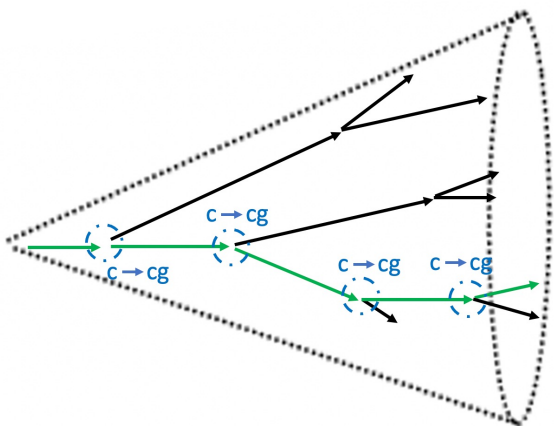
Dead cone of D⁰ - tagged jets



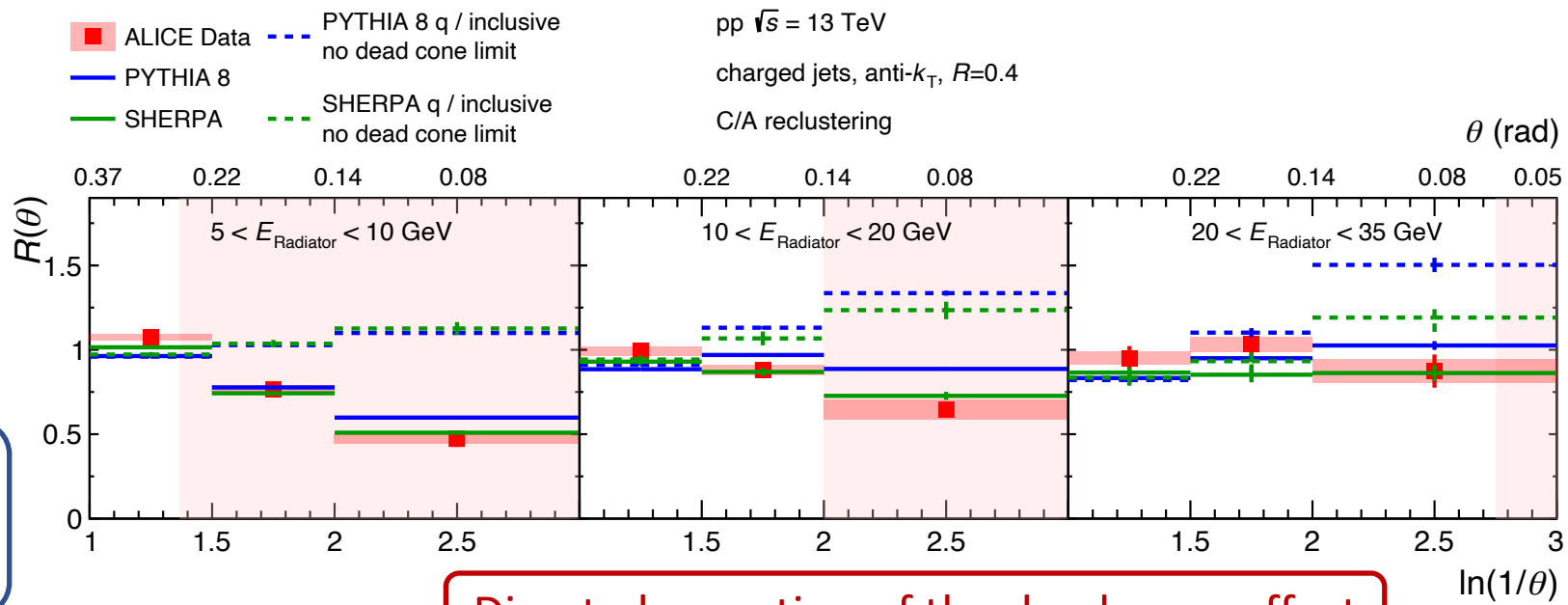
Low p_T charm and inclusive sensitive to both mass and Casimir colour effects (Run 2)

Beauty and charm-quark emissions isolate mass effects

High p_T charm and inclusive comparisons isolate Casimir colour factors

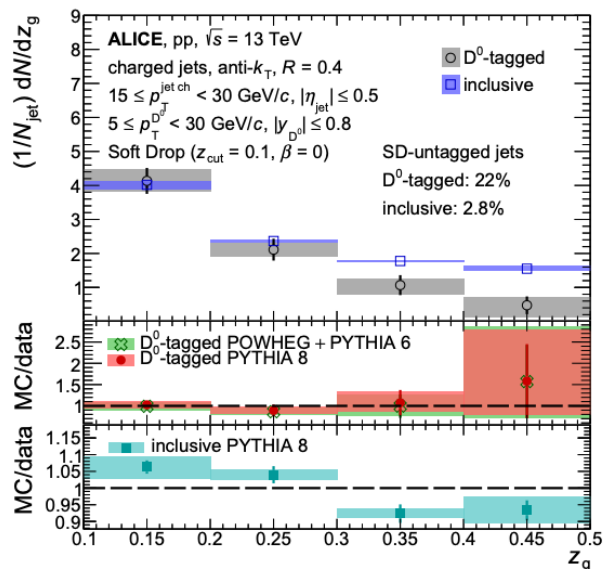


New experimental technique to reconstruct emissions of heavy-flavour quarks in the parton shower



Direct observation of the dead-cone effect

See G. M. Innocenti's talk on Tuesday for more details



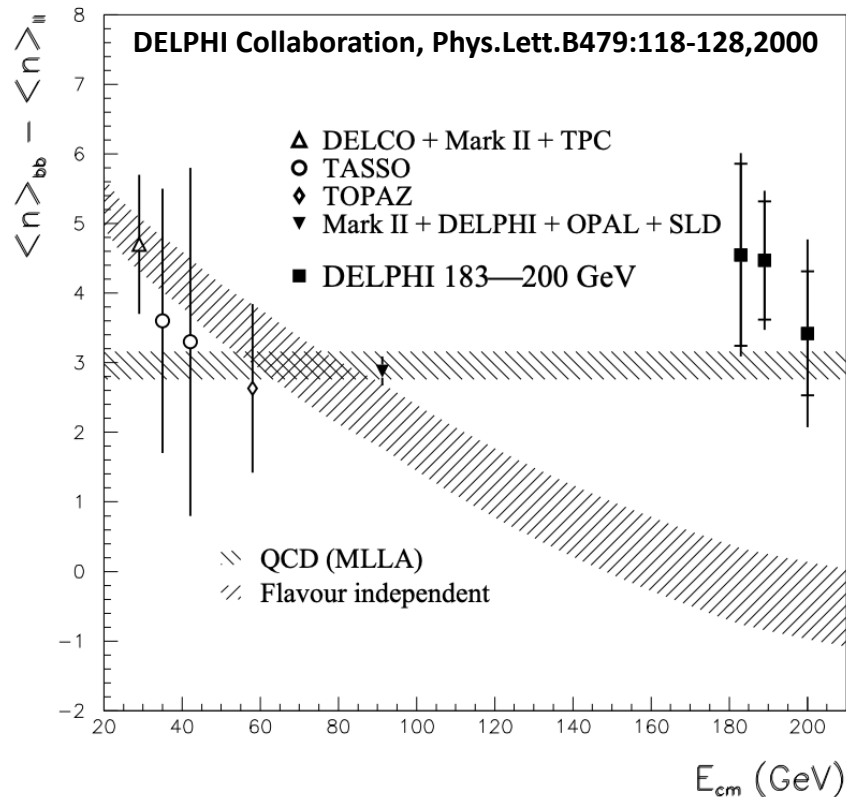
Constraining the role of parton mass and Casimir colour factors in the parton shower



ALICE

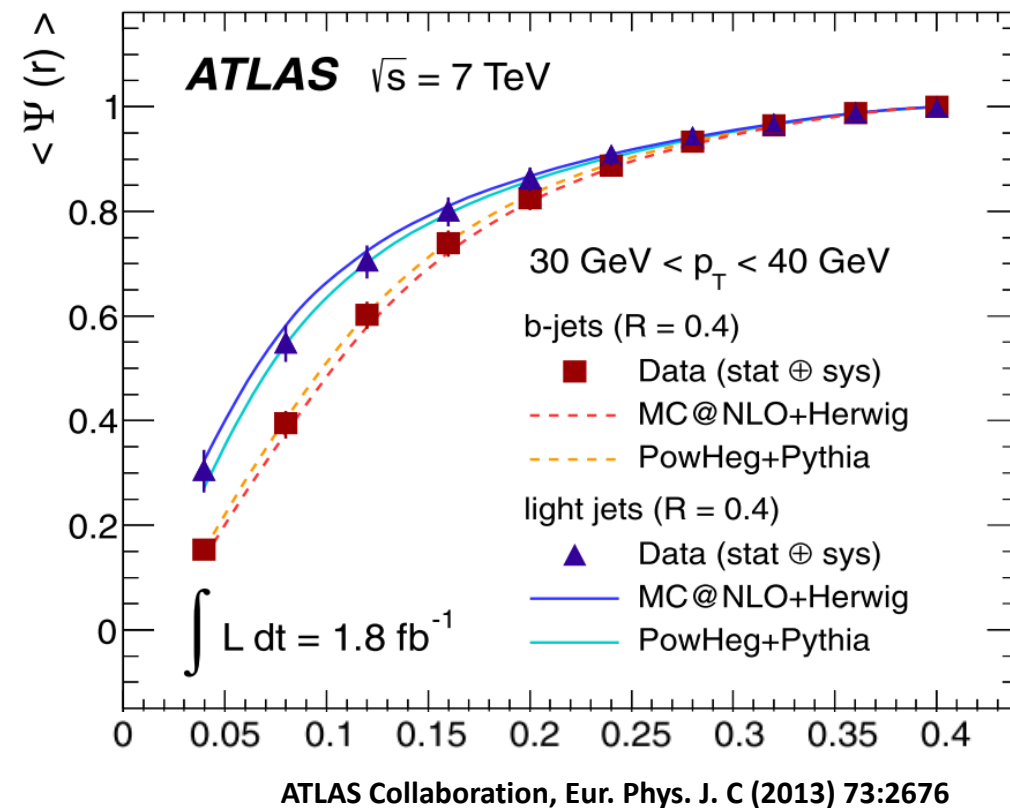
Backup

36



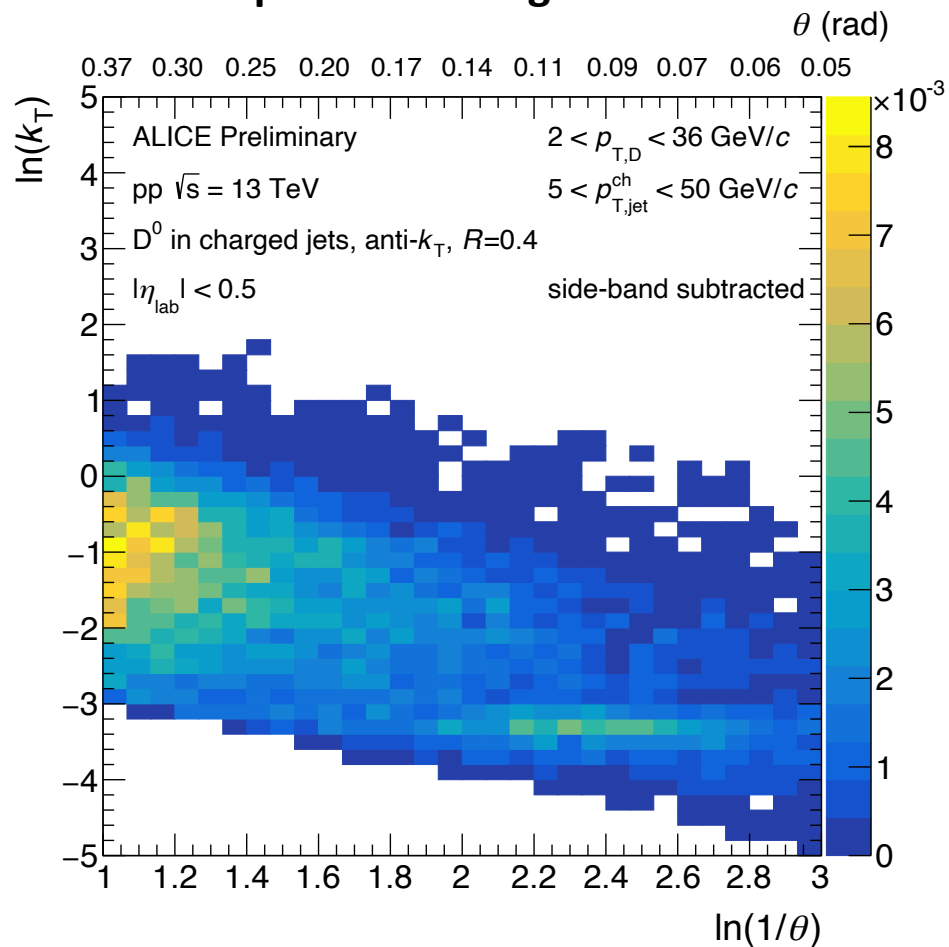
Difference in average multiplicity between events containing a b-quark jet and those with a light-quark jet compatible with flavour dependent QCD calculations

$$\Psi(r) = \frac{p_T(0, r)}{p_T(0, R)}; \quad r \leq R$$

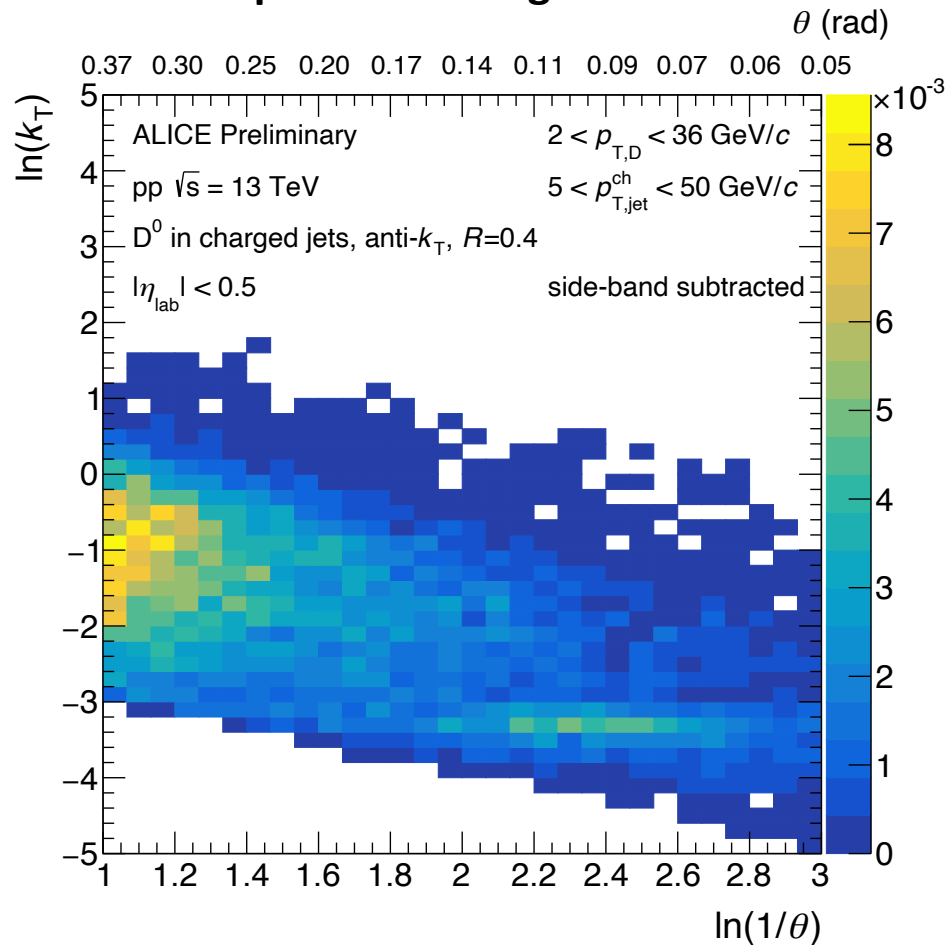


p_T density around the initial scattered b-quark direction is depleted compared to the density around light quarks and gluons

Lund plane of $c \rightarrow cg$ emissions

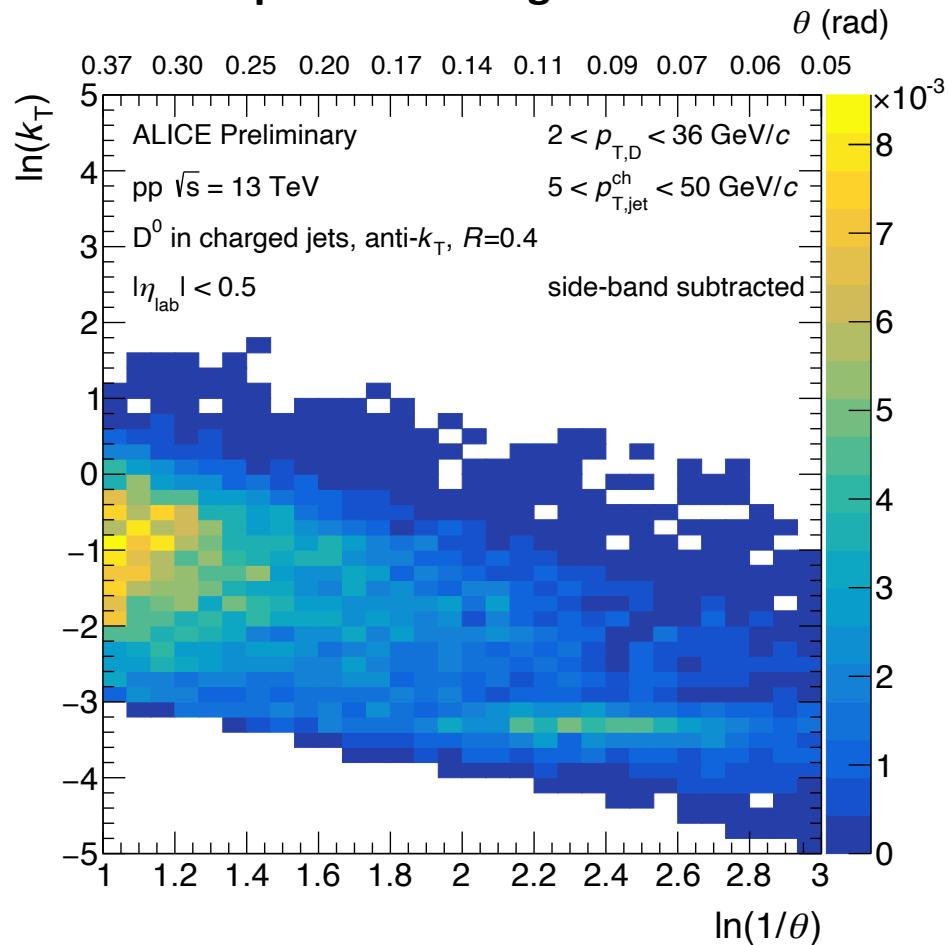


Lund plane of $c \rightarrow cg$ emissions



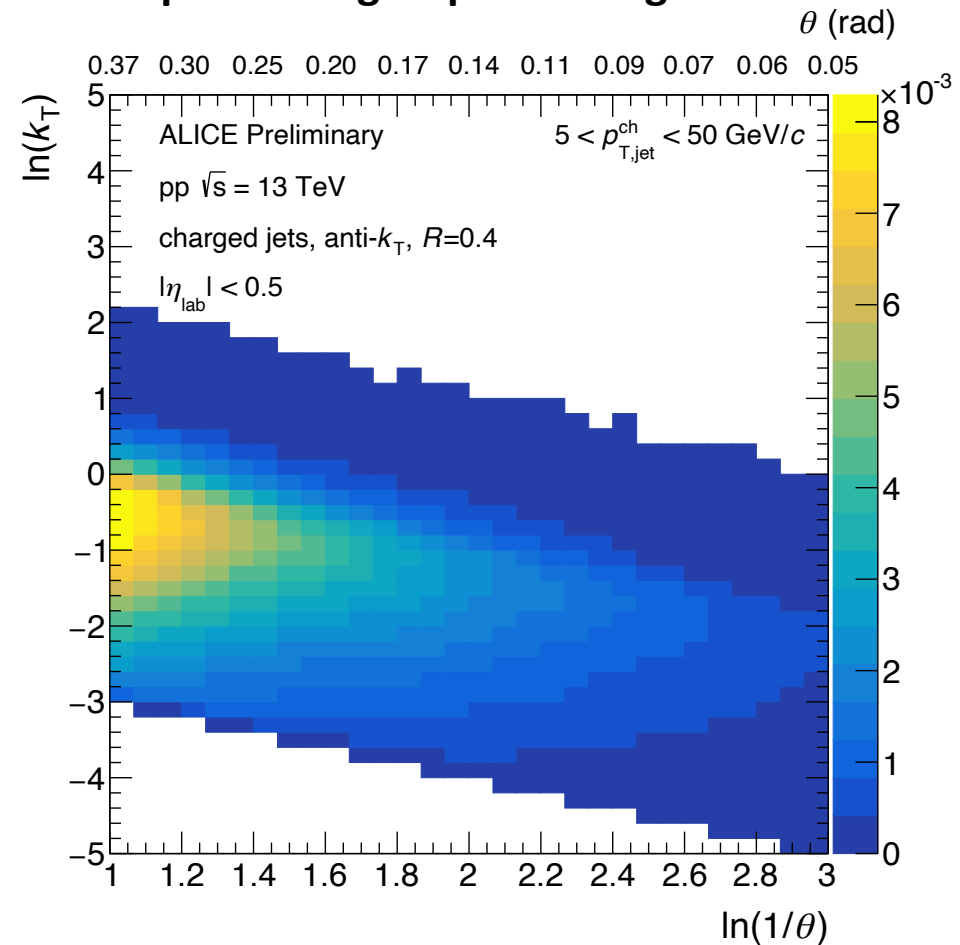
Following the branch containing the D^0 coincides with following the hardest branch in 99% of cases

Lund plane of c -> cg emissions



Following the branch containing the D^0 coincides with following the hardest branch in 99% of cases

Lund plane of light quark and gluon emissions



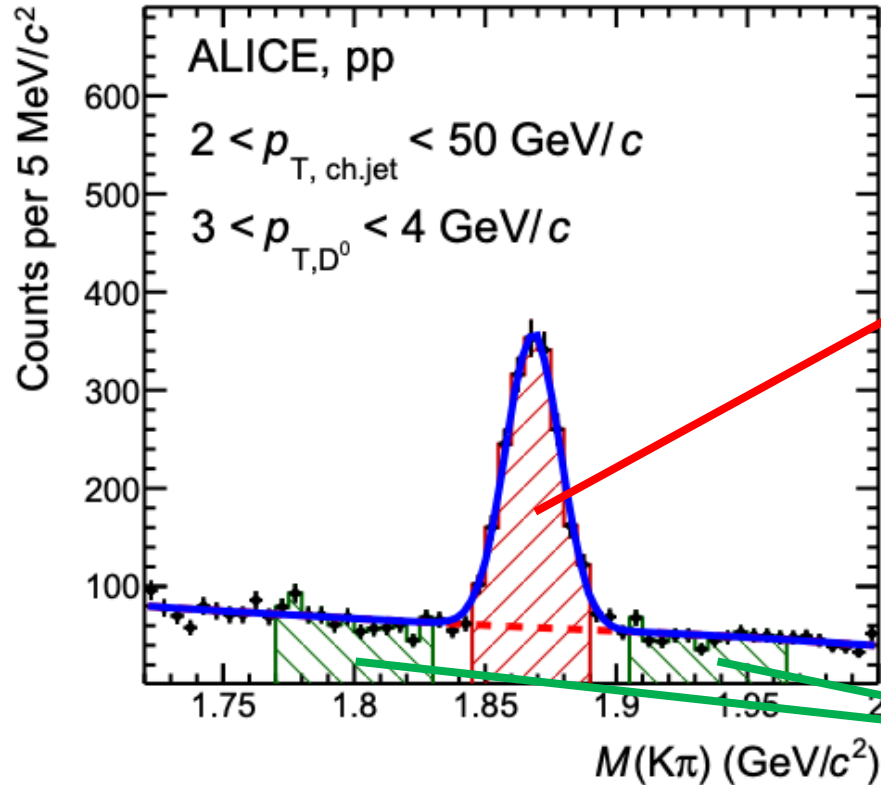
The hardest branch is followed at each declustering step

σ_{fit}

μ_{fit}

Gaussian width

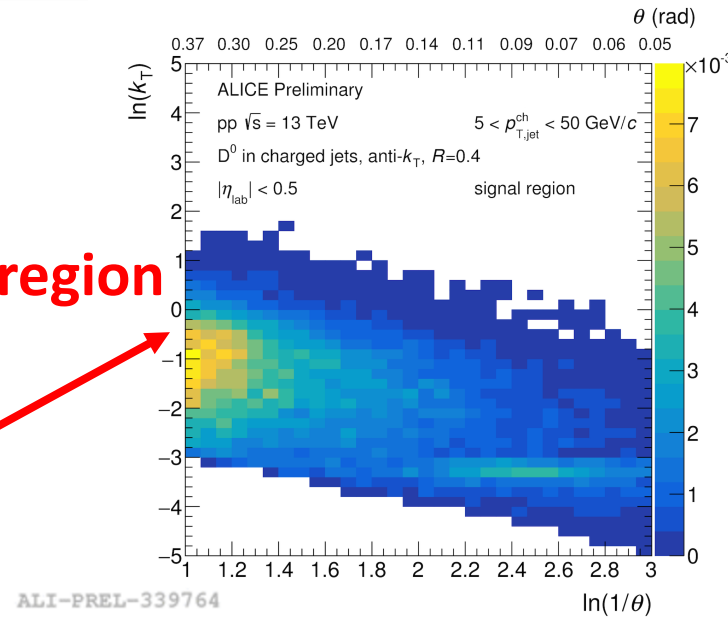
Gaussian mean



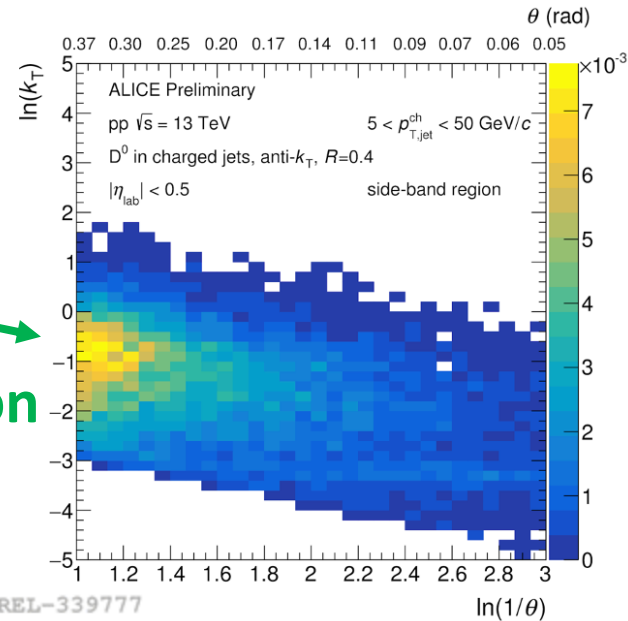
signal region

sideband region

Performed in intervals of p_{T,D^0}



ALI-PREL-339764



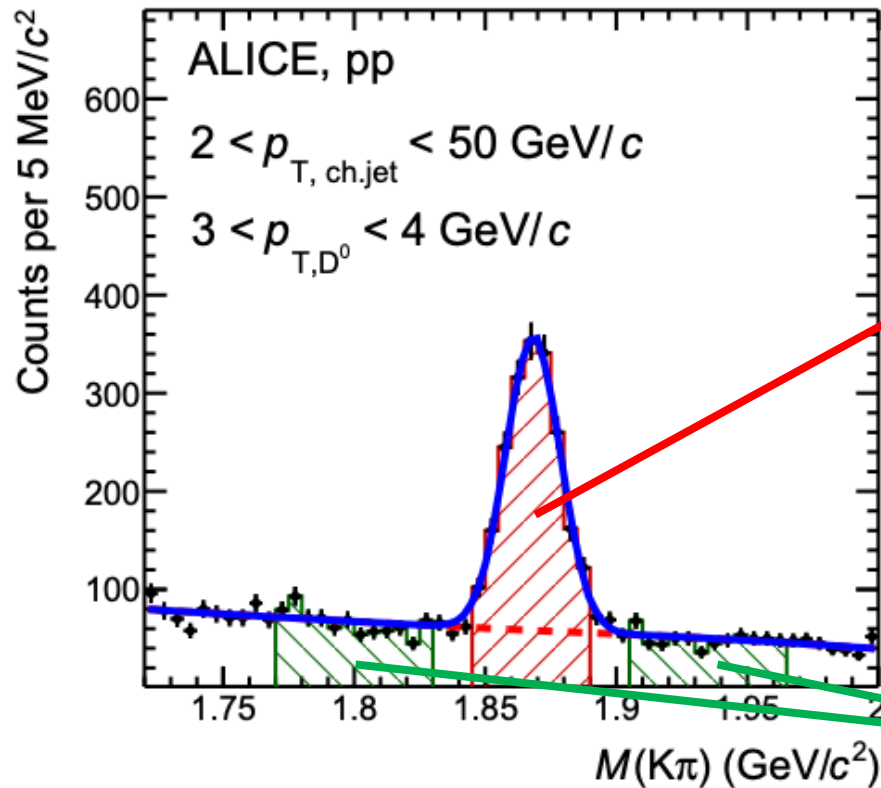
ALI-PREL-339777

$$|M - \mu_{\text{fit}}| < 2\sigma_{\text{fit}}$$

Contains almost all of the signal and some background candidates

$$4\sigma_{\text{fit}} < |M - \mu_{\text{fit}}| < 9\sigma_{\text{fit}}$$

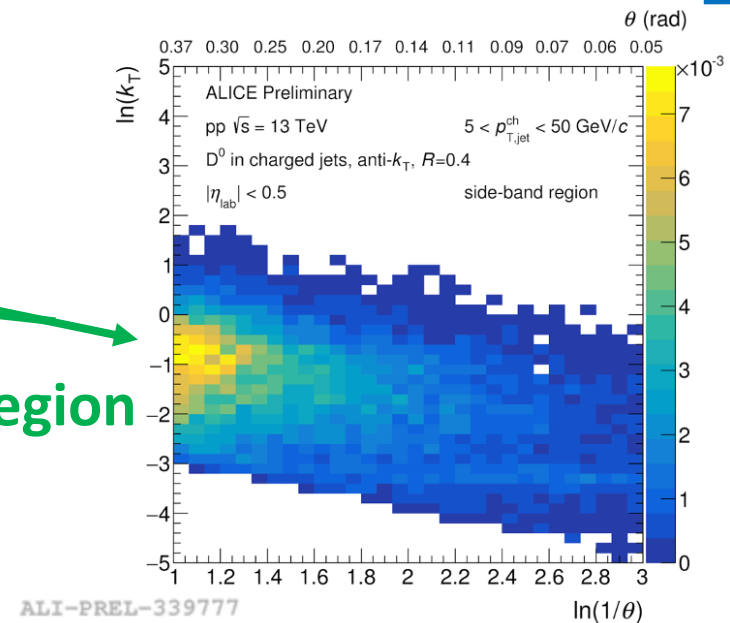
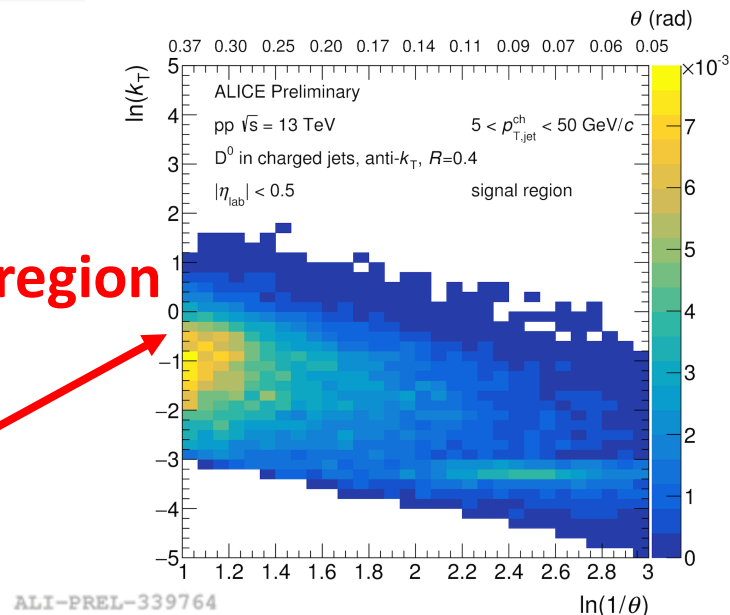
Entirely composed of background candidates



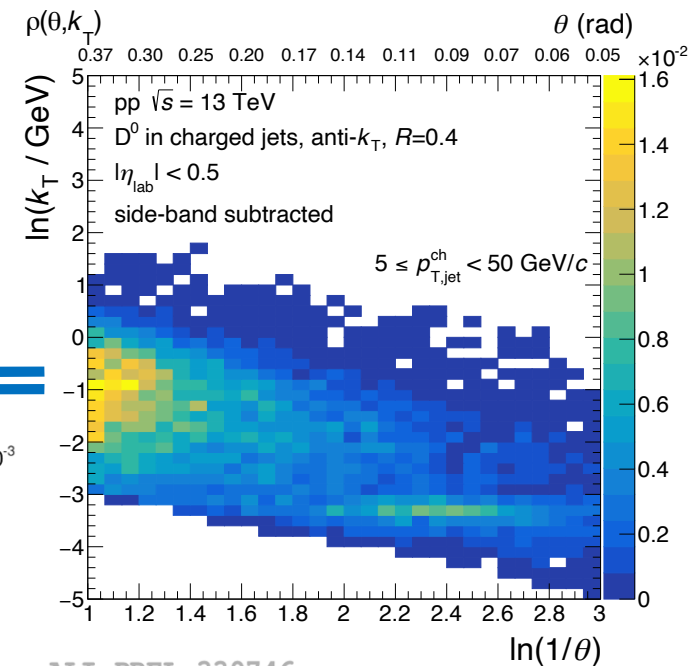
Performed in intervals of p_{T,D^0}

signal region

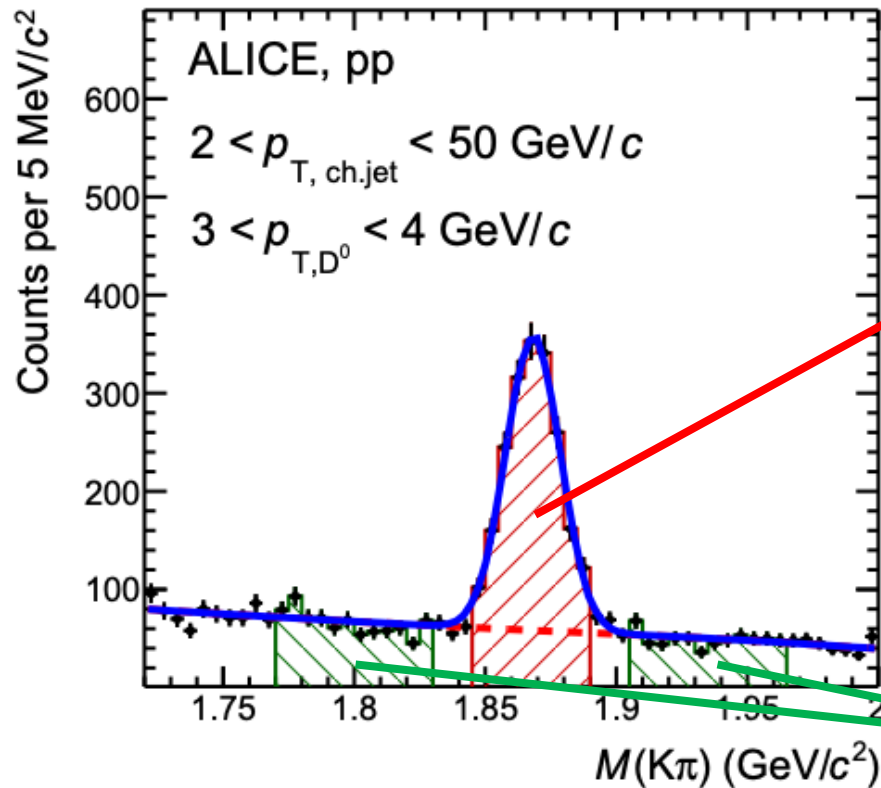
sideband region



subtracted



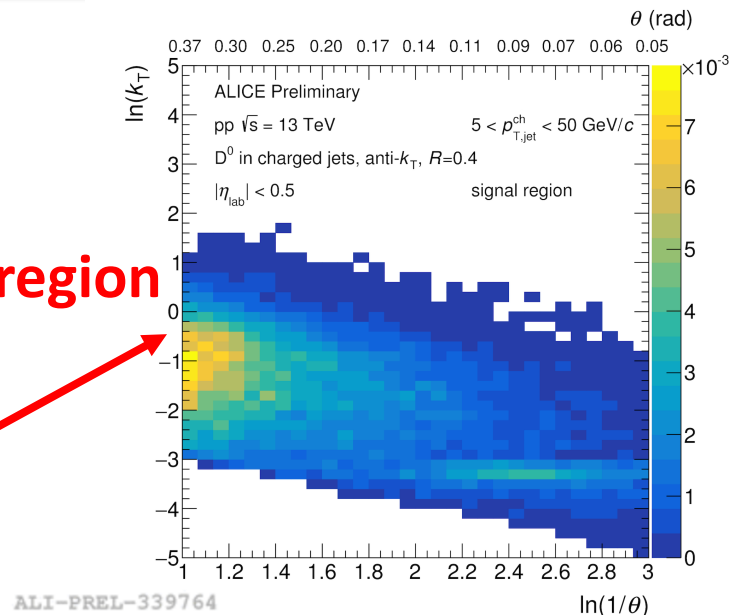
Purely signal D^0 -tagged jet distribution extracted



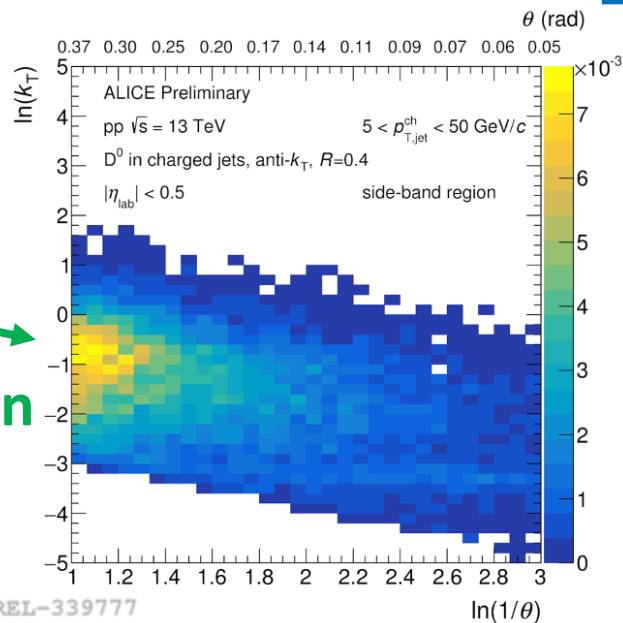
Performed in intervals of p_{T,D^0}

signal region

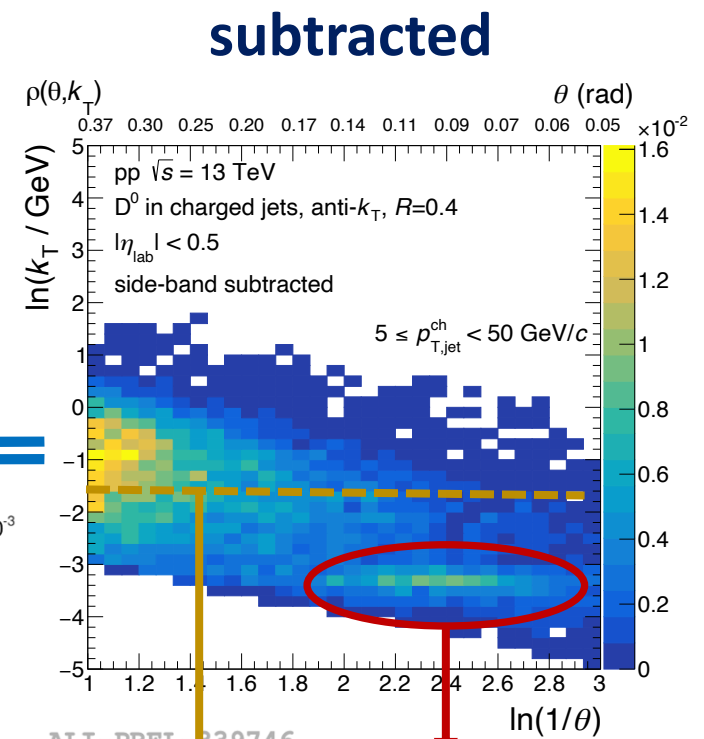
sideband region



ALI-PREL-339764



ALI-PREL-339777



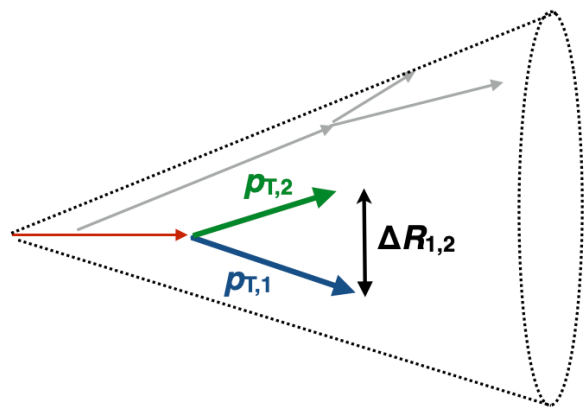
ALI-PREL-339746

Soft π from D^* decays

Selection in k_T of emissions suppresses non-perturbative effects

See G. M. Innocenti's talk on Tuesday for more details

$$dP_{i \rightarrow jk} = \frac{d\theta}{\theta} dz P_{i \rightarrow jk}(z)$$



Soft Drop grooming condition

$$z = \frac{p_{T,2}}{p_{T,1} + p_{T,2}} > z_{\text{cut}} \left(\frac{\Delta R_{1,2}}{R} \right)^\beta$$

$$z = 0.1, \beta = 0$$

A. J. Larkoski et al. , JHEP 1405 (2014) 146

Converges onto the QCD splitting function for the first splitting that passes Soft Drop

Emissions from charm-quarks have a steeper splitting probability than light quarks and gluons

Fewer symmetric splittings

