

# NNLL resummation of groomed additive observables

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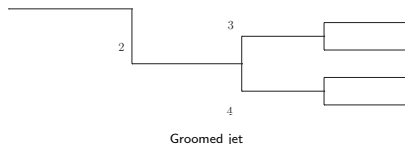
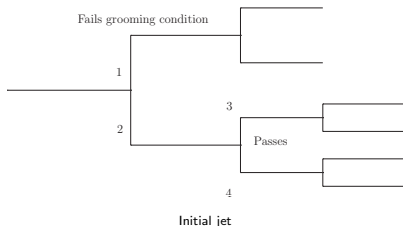
- 1 Introduction and motivation
- 2 Recap of NLL resummation
- 3 NNLL resummation in the small  $z_{\text{cut}}$  limit
- 4 Including finite  $z_{\text{cut}}$  effects at NLL
- 5 Results
- 6 Summary and future work

### **Groomed jet shape observables are a good choice for direct comparison between perturbative QCD and measurements:**

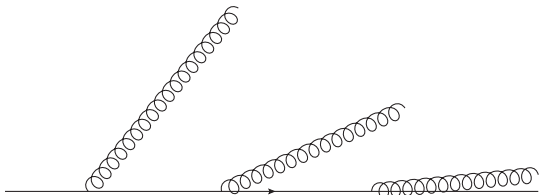
- less susceptible pileup and underlying event
- hadronisation corrections are smaller
- Resummations are easier due to lack of non global logs

- pQCD calculations have been compared to measurements for several angularities and the jet mass (CMS, ATLAS)
- Calculations compared for the groomed jet mass are either NLL + NLO (Marzani, Schunk, Soyez) or NNLL in the small  $z_{\text{cut}}$  limit (Frye, Larkoski, Schwartz, Yan).
- Other groomed angularities are calculated to NLL accuracy with NLO matching (Caletti, Fedkevych, Marzani, Reichelt, Schumann, Soyez).
- Aim is to get NNLL at small  $z_{\text{cut}}$  with finite  $z_{\text{cut}}$  at NLL for any additive observable and demonstrate NLO matching for a selection of observables.
- This talk:  $e^+e^- \rightarrow q\bar{q}$  split into two hemispheres, both of which are groomed with mMDT Dasgupta, Fregoso, Marzani, Salam (Soft Drop Larkoski, Marzani, Soyez, Thaler with  $\beta = 0$ ).

- 1 Start with a hemisphere (jet) clustered with the Cambridge Aachen (C/A) algorithm.
- 2 Undo the last clustering in the sequence to obtain two branches,  $i$  and  $j$ .
- 3 If the softer branch does not satisfy  $\frac{\min(E_j, E_i)}{E_i + E_j} > z_{\text{cut}}$  then it is discarded and the groomer returns to step 2.
- 4 If  $\frac{\min(E_i, E_j)}{E_i + E_j} > z_{\text{cut}}$  is satisfied, the groomer stops and the groomed hemisphere contains all of the particles in both  $i$  and  $j$ .



Can consider strongly ordered primary emissions inclusive of branchings



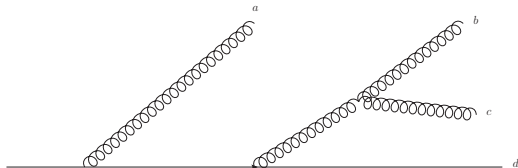
$$\Sigma^{\text{NLL}'}(v; z_{\text{cut}}) = \frac{1}{\sigma} \int_0^v \frac{d\sigma}{dv'} dv' = \left( 1 + \frac{C_F \alpha_s}{\pi} C_1 \right) \exp[-R^{\text{NLL}'}(v, z_{\text{cut}})]$$

$$R^{\text{NLL}'}(v, z_{\text{cut}}) = \int_{z_{\text{cut}}}^1 \int^{z^2 Q^2} \frac{C_F \alpha_s^{\text{CMW}}(k_t^2)}{\pi} \left( \frac{2}{z} + \delta(1-z) \gamma_0^{\text{h.c.}} \right) \frac{dk_t^2}{k_t^2} dz \Theta(V_{\text{s.c.}}(z, k_t) - v)$$

$$\gamma_0^{\text{h.c.}} = -\frac{3}{2}$$

The resummation can be structured as an inclusive piece and a clustering correction:

$$\Sigma^{\text{NNLL}}(v; z_{\text{cut}}) = \Sigma_{\text{inc.}}(v; z_{\text{cut}}) + \Sigma_{\text{clust.}}(v; z_{\text{cut}})$$



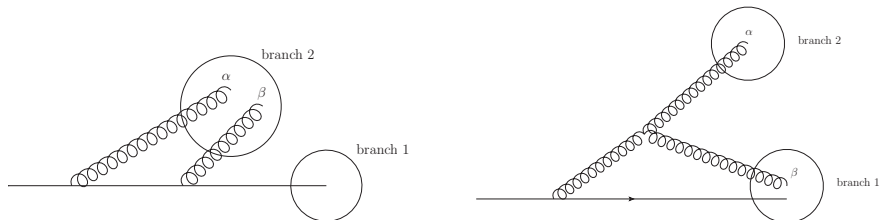
- hard collinear part proceeds the same as for ungroomed observables, e.g ARES (Banfi et al.)
- Can distil what we need into a modification of  $C_1$  and evaluating the Sudakov factor up to NNLL accuracy.

$$\Sigma_{\text{inc.}}(v, z_{\text{cut}}) = \frac{1}{\sigma} \int_0^v \frac{d\sigma}{dv'} dv' = \left( 1 + \frac{C_F \alpha_s}{\pi} C^{\text{r.c.}}(v) \right) \exp[-R^{\text{NNLL}}(v, z_{\text{cut}})]$$

$$R^{\text{NNLL}}(v, z_{\text{cut}}) = \int_{z_{\text{cut}}}^1 \int^{z^2 Q^2} \frac{C_F \alpha_s(k_t^2)}{\pi} \left( \frac{2}{z} \left( 1 + \frac{\alpha_s(k_t^2) K^{\text{CMW}}}{2\pi} \right) + \delta(1-z) \left( \gamma_0^{\text{h.c.}} + \frac{\alpha_s(k_t^2)}{2\pi} \gamma_1^{\text{h.c.}} \right) \right) \frac{dk_t^2}{k_t^2} dz \Theta(V_{\text{s.c.}}(z, k_t) - v)$$

- soft part is essentially the same as at NLL as soft logs are logs of  $z_{\text{cut}}$
- Multiple emission effects turn out to be N<sup>3</sup>LL





$$z_\alpha < z_{\text{cut}}, \quad z_\beta < z_{\text{cut}}, \quad z_{\text{cut}} < z_\alpha + z_\beta$$

$$\Sigma^{\text{clust}}(v, z_{\text{cut}}) = \mathcal{F}_{\text{clust.}}(v) \exp[-R^{\text{NLL}}(v, z_{\text{cut}})]$$

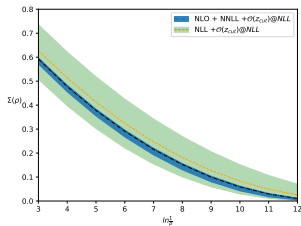
$$\mathcal{F}_{\text{clust.}}(v) = \left(\frac{\alpha_s}{2\pi}\right)^2 \times$$

$$\left( C_F^2 \frac{4\pi}{3} \text{Cl}_2\left(\frac{\pi}{3}\right) - C_F C_A 1.161 - C_F T_R n_f 1.754 \right) \frac{\ln\left(v^{\frac{2}{a+b}}\right)}{1 + \beta_0 \alpha_s(Q^2) \ln\left(v^{\frac{2}{a+b}}\right)}$$

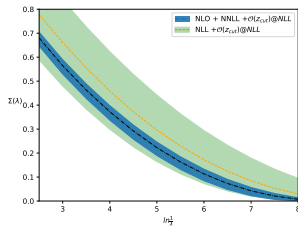
- NLL terms suppressed by powers of  $z_{\text{cut}}$  could be of a similar size to NNLL terms.
- Can match NNLL small  $z_{\text{cut}}$  calculation to finite  $z_{\text{cut}}$  NLL calculation (Dasgupta, Fregoso, Marzani, Salam)

$$\Sigma(v, z_{\text{cut}}) = \left( 1 + C^{\text{r.c.}}(v) + \mathcal{F}_{\text{clust.}}, 1 \right) \exp \left( \begin{array}{cc} -R_q^{\text{NNLL}} - R_q z_{\text{cut}} - R_{q \rightarrow g} & R_{g \rightarrow q} \\ R_{q \rightarrow g} & -R_g^{\text{NLL}} - R_g z_{\text{cut}} - R_{g \rightarrow q} \end{array} \right) \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad (1)$$

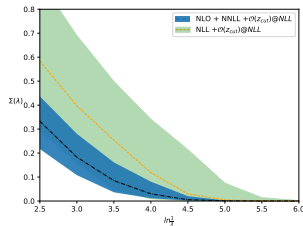
- Method applied to three observables: heavy hemisphere mass, Width ( $\lambda_1^1$ ) and Les Houches angularity ( $\lambda_{0.5}^1$ )
- NLO matching using multiplicative scheme
- Uncertainty from varying resummation and renormalisation scales to introduce variation beyond the stated accuracy.



Mass

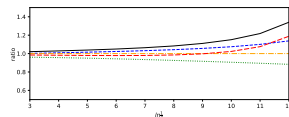
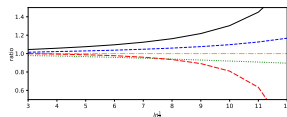
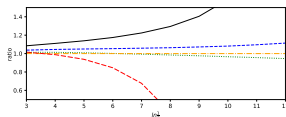
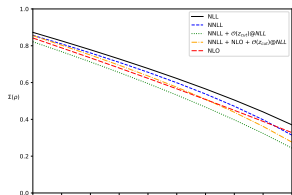
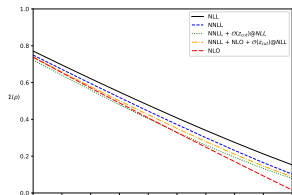
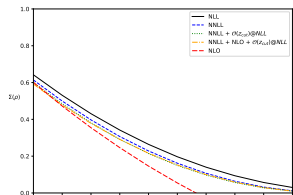


Width



Les Houches Angularity

## Heavy hemisphere mass

 $z_{\text{cut}} = 0.1$  $z_{\text{cut}} = 0.2$  $z_{\text{cut}} = 0.3$

- General method for NNLL resummation of additive observables computed on jets groomed with mMDT
- First results with small  $z_{\text{cut}}$  at NNLL +  $\mathcal{O}(z_{\text{cut}})$  at NLL + NLO matching
- future work: extension to gluon jets
- Other extensions: non-additive observables, Soft drop with  $\beta \neq 0$

# Backup Slides

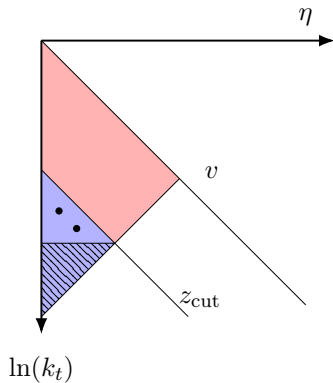
$$V_{\text{s.c.}}(k_1, \dots, k_n) = \sum_i^n V_{\text{s.c.}}(k_i)$$

Can ignore the contribution of all emissions softer than  $z_{\text{cut}}$  to the observable

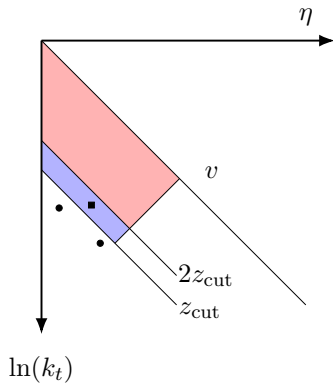
$$\rightarrow \Sigma^{\text{NLL}}(v, z_{\text{cut}}) = \frac{\exp[R(v, z_{\text{cut}}) - \gamma_E R'(v, z_{\text{cut}})]}{\Gamma[1 + R'(v, z_{\text{cut}})]} = \exp[-R(v, z_{\text{cut}})] + \text{N}^3\text{LL} \quad (2)$$



# Clustering Lund Planes



Independent emission



Correlated emission

$$\Sigma(v) = \Sigma_{\text{NNLL}}(v) \left[ 1 + \left( \Sigma^{(1)}(v) - \Sigma_{\text{NNLL}}^{(1)}(v) \right) + \left( \Sigma^{(2)}(v) - \Sigma_{\text{NNLL}}^{(2)}(v) \right) - \Sigma_{\text{NNLL}}^{(1)}(v) \left( \Sigma^{(1)}(v) - \Sigma_{\text{NNLL}}^{(1)}(v) \right) \right].$$

$\Sigma^{(n)}$  is the fixed order result up to  $\alpha_s^n$ .

$$R(v, z_{\text{cut}}) \rightarrow R^{\text{NLL}}(v, z_{\text{cut}}) + R^{\text{remainder}}(xv, z_{\text{cut}})$$

where

$$R^{\text{remainder}}(v, z_{\text{cut}}) = R^{\text{NNLL}}(v, z_{\text{cut}}) - R^{\text{NLL}}(v, z_{\text{cut}})$$

and

$$\alpha_s(Q) \rightarrow \alpha_s(xQ) + \alpha_s^2(Q)\beta_0 \ln(x)$$