

Sterile neutrino searches with a Deuterated Liquid Scintillator

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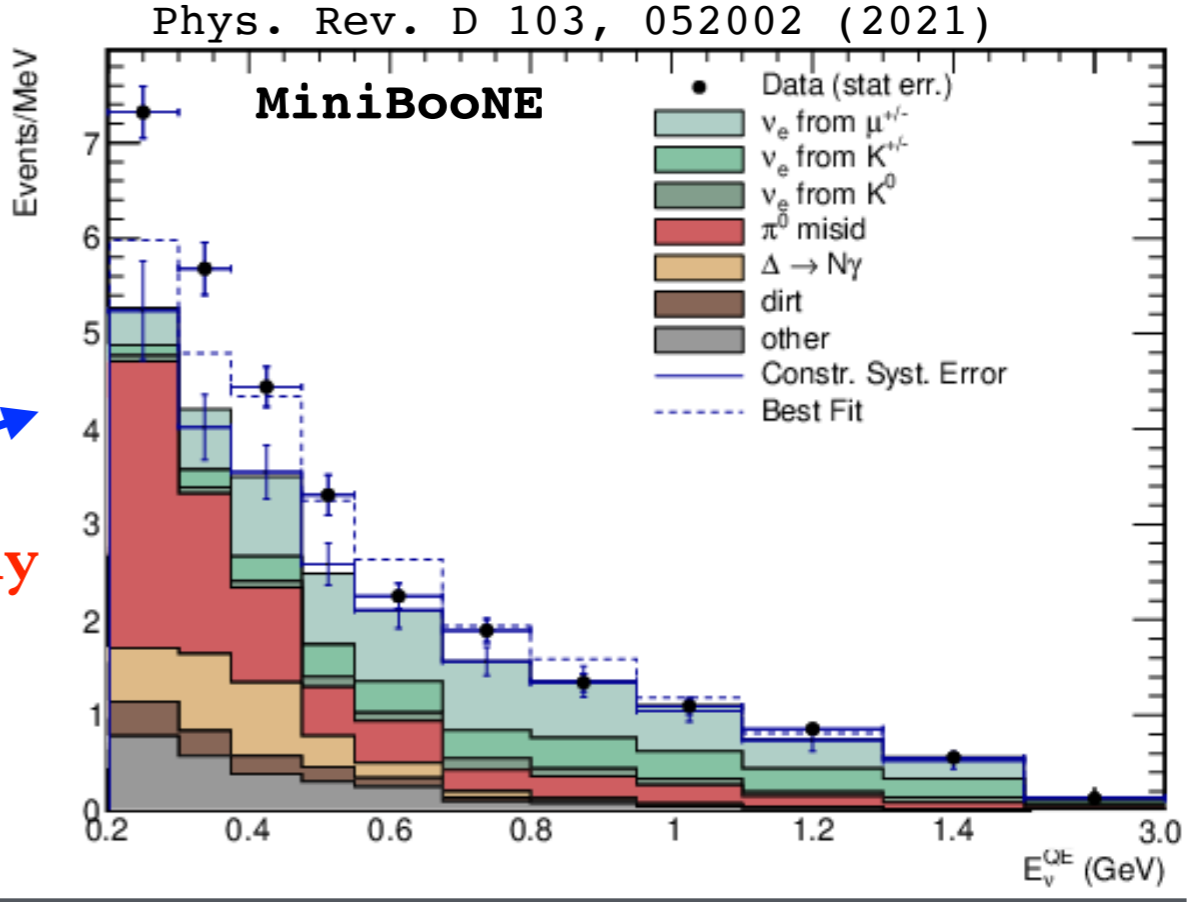
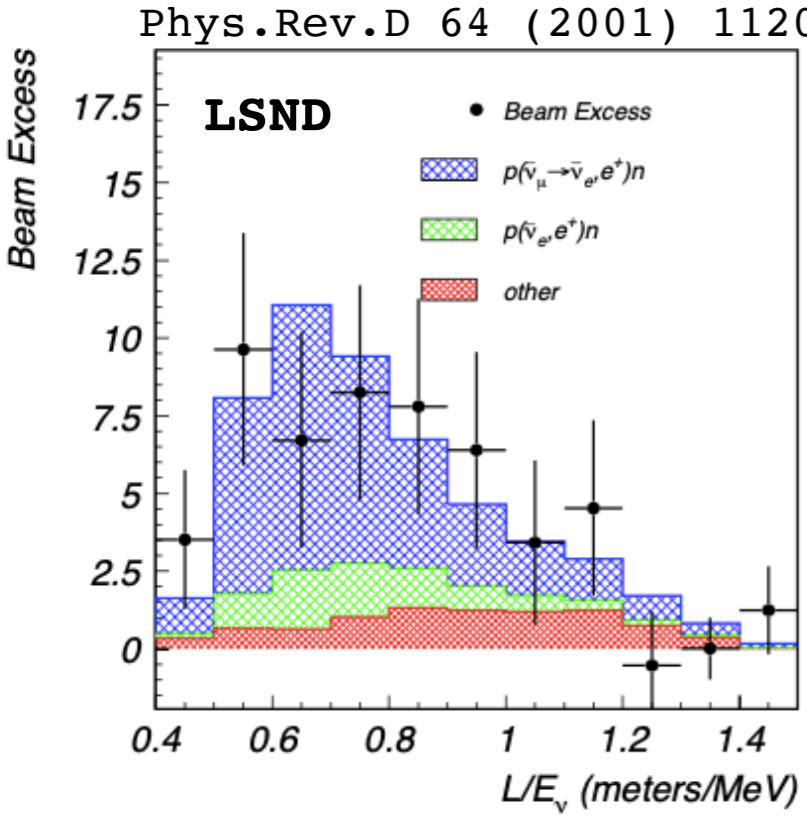
Ongoing work in collaboration with D. Bhatia, D. Indumathi,
V. Datar and M.V.N. Murthy

DAE BRNS HEP SYMPOSIUM 2022, IISER Mohali

- * **Short baselines anomalies and sterile neutrinos**
- * **Two flavor neutrino oscillation probabilities**
- * **Deuterated liquid scintillator**
- * **Sterile neutrino searches**
- * **Challenges with background events**

Short-baseline anomalies

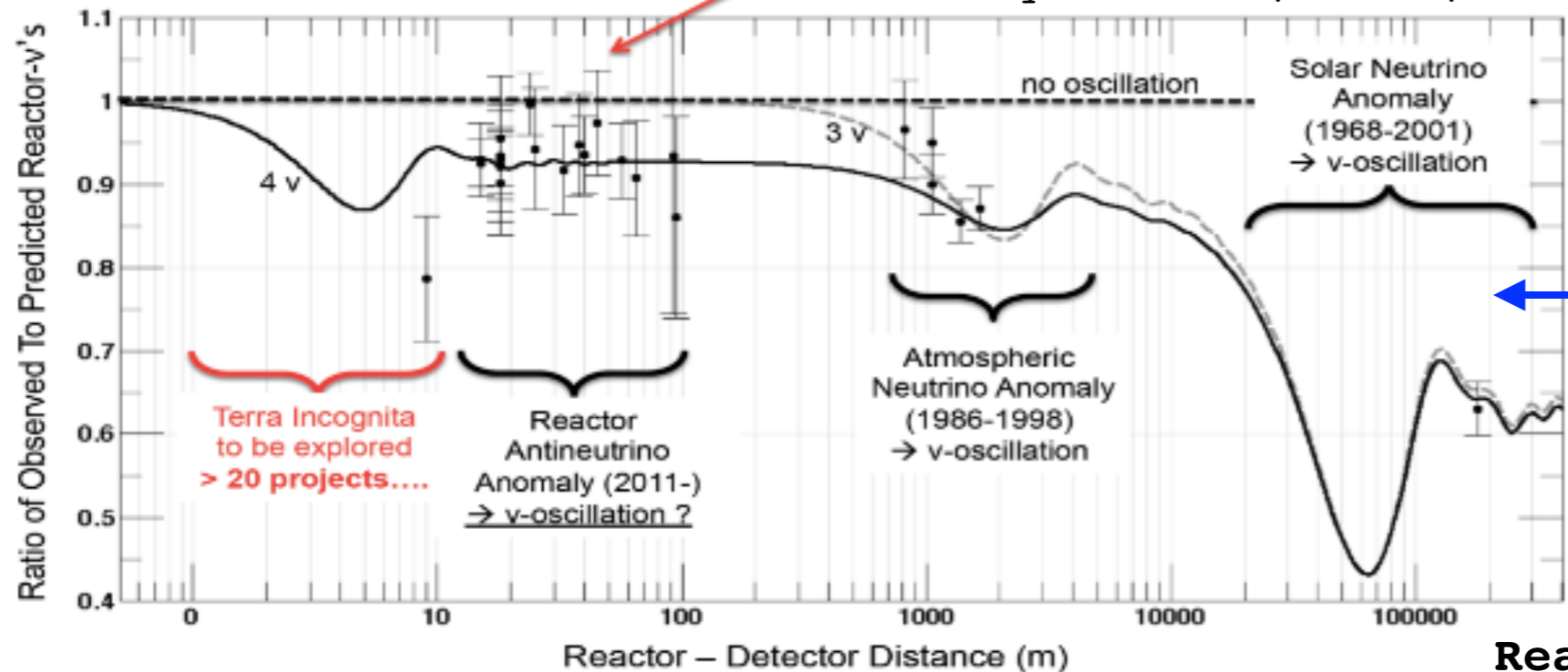
A bunch of unaccounted flavor transitions when neutrino travel short distances !



Appearance anomaly
Surplus of $\bar{\nu}_e$

Observed/predicted averaged event ratio: $R=0.927 \pm 0.023$ (3.0σ)

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Disappearance anomaly
Deficit of $\bar{\nu}_e$

Reactor antineutrino anomaly

- * For the baselines and neutrino energies of the short-baseline experiments, the data cannot be explained by oscillations due to the solar or atmospheric mass-squared differences.
- * The measurement of the invisible decay width of the Z boson limits the number of light SM active neutrinos to $N_\nu = 2.9840 \pm 0.0082$.
- * The simplest way to explain the SBL anomalies is to assume the existence of new light neutrinos which are similar to active neutrinos in masses, mixings except that they do not undergo weak interactions.
- * Conventionally, these new neutrinos have been called “Sterile”. Sterile neutrinos are SM singlets.
- * SBL anomalies can be explained if Active and Sterile neutrinos are mixed. In the simplest case of the 3+1 model, a sterile neutrino of mass $\sim 1\text{eV}$ is needed along with the 3 active neutrinos.

Probabilities in two flavors

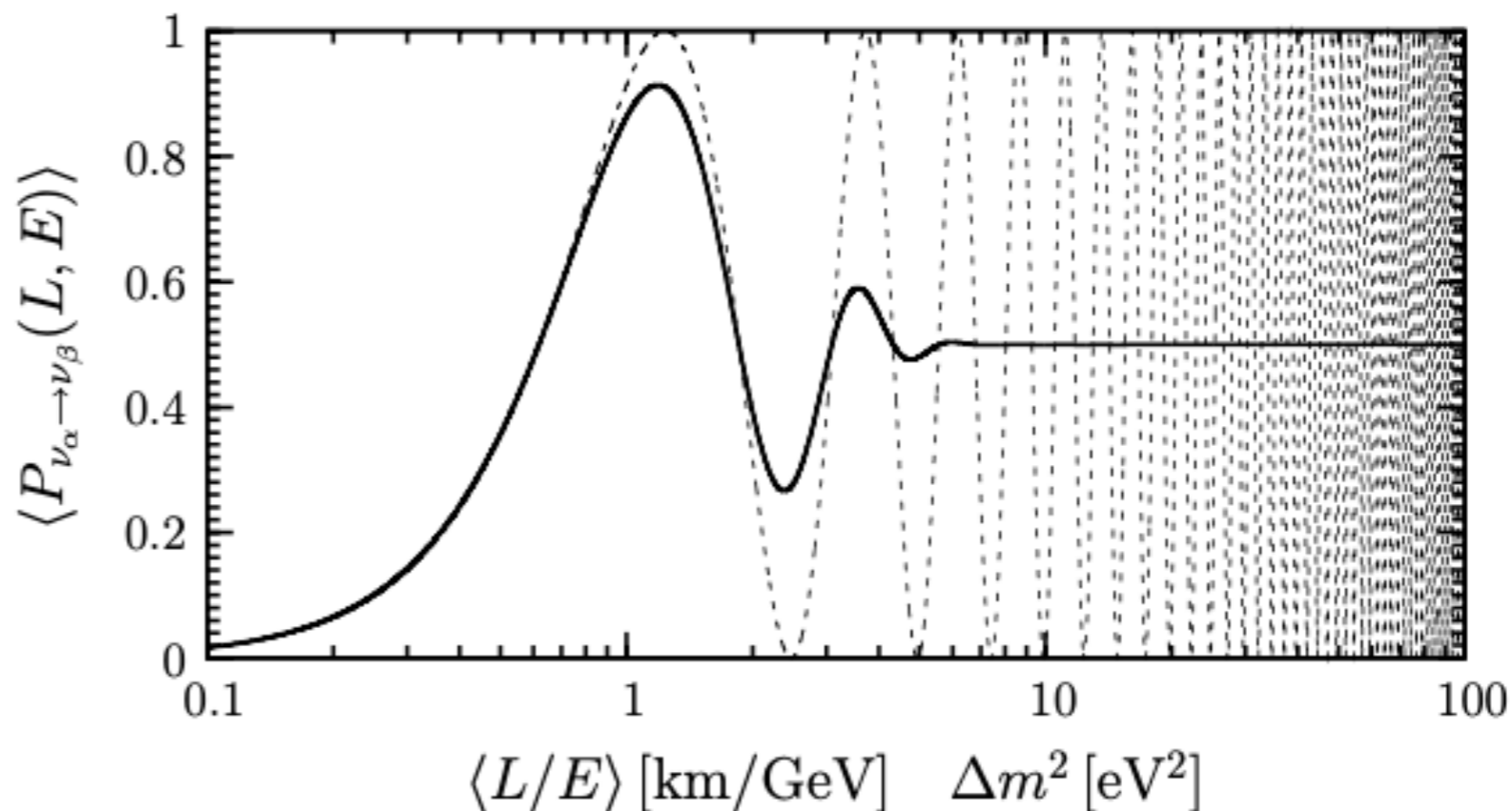
$$P(\nu_\alpha \rightarrow \nu_\beta) = |\langle \nu_\beta | \nu_\alpha(\vec{x}, t) \rangle|^2$$

Assumption 1: All energy eigenstates are produced with the same 3-momentum i.e. $p_1 = p_2 = p$

Assumption 2: Neutrinos are relativistic i.e.

$$(1) \quad t = T = L$$

$$(2) \quad E_i = p + m_i^2/2p$$



$$P(\nu_e \rightarrow \nu_e) = 1 - \sin^2 2\theta \sin^2 \frac{\Delta m^2 L}{4E_\nu}$$

Disappearance Channel

$$P(\nu_e \rightarrow \nu_x) = \sin^2 2\theta \sin^2 \frac{\Delta m^2 L}{4E_\nu}$$

Appearance Channel

Type of experiment	L	E	Δm^2 sensitivity
Reactor SBL	~ 10 m	~ 1 MeV	~ 0.1 eV ²
Accelerator SBL (Pion DIF)	~ 1 km	$\gtrsim 1$ GeV	$\gtrsim 1$ eV ²
Accelerator SBL (Muon DAR)	~ 10 m	~ 10 MeV	~ 1 eV ²
Accelerator SBL (Beam Dump)	~ 1 km	$\sim 10^2$ GeV	$\sim 10^2$ eV ²
Reactor LBL	~ 1 km	~ 1 MeV	$\sim 10^{-3}$ eV ²
Accelerator LBL	$\sim 10^3$ km	$\gtrsim 1$ GeV	$\gtrsim 10^{-3}$ eV ²
ATM	20 – 10^4 km	0.5 – 10^2 GeV	$\sim 10^{-4}$ eV ²
Reactor VLB	$\sim 10^2$ km	~ 1 MeV	$\sim 10^{-5}$ eV ²
Accelerator VLB	$\sim 10^4$ km	$\gtrsim 1$ GeV	$\gtrsim 10^{-4}$ eV ²
SOL	$\sim 10^{11}$ km	0.2 – 15 MeV	$\sim 10^{-12}$ eV ²

As an example: [Kudankulam Nuclear Power Plant](#)

- * Tirunelveli, Tamil Nadu
- * Pressurized light-water reactors
- * Total 6 units commissioned - 2 operational and 4 more to be built by 2027
- * Power of each unit: 3 GWth

$$\begin{aligned}\phi(E_\nu) = & f_{235U} \exp(0.870 - 0.160E_\nu - 0.091E_\nu^2) \\ & + f_{239Pu} \exp(0.896 - 0.239E_\nu - 0.0981E_\nu^2) \\ & + f_{238U} \exp(0.976 - 0.162E_\nu - 0.0790E_\nu^2) \\ & + f_{241Pu} \exp(0.793 - 0.080E_\nu - 0.1085E_\nu^2)\end{aligned}$$

Phys. Rev. D 39 (1989) 3378

For a PWR type reactor:

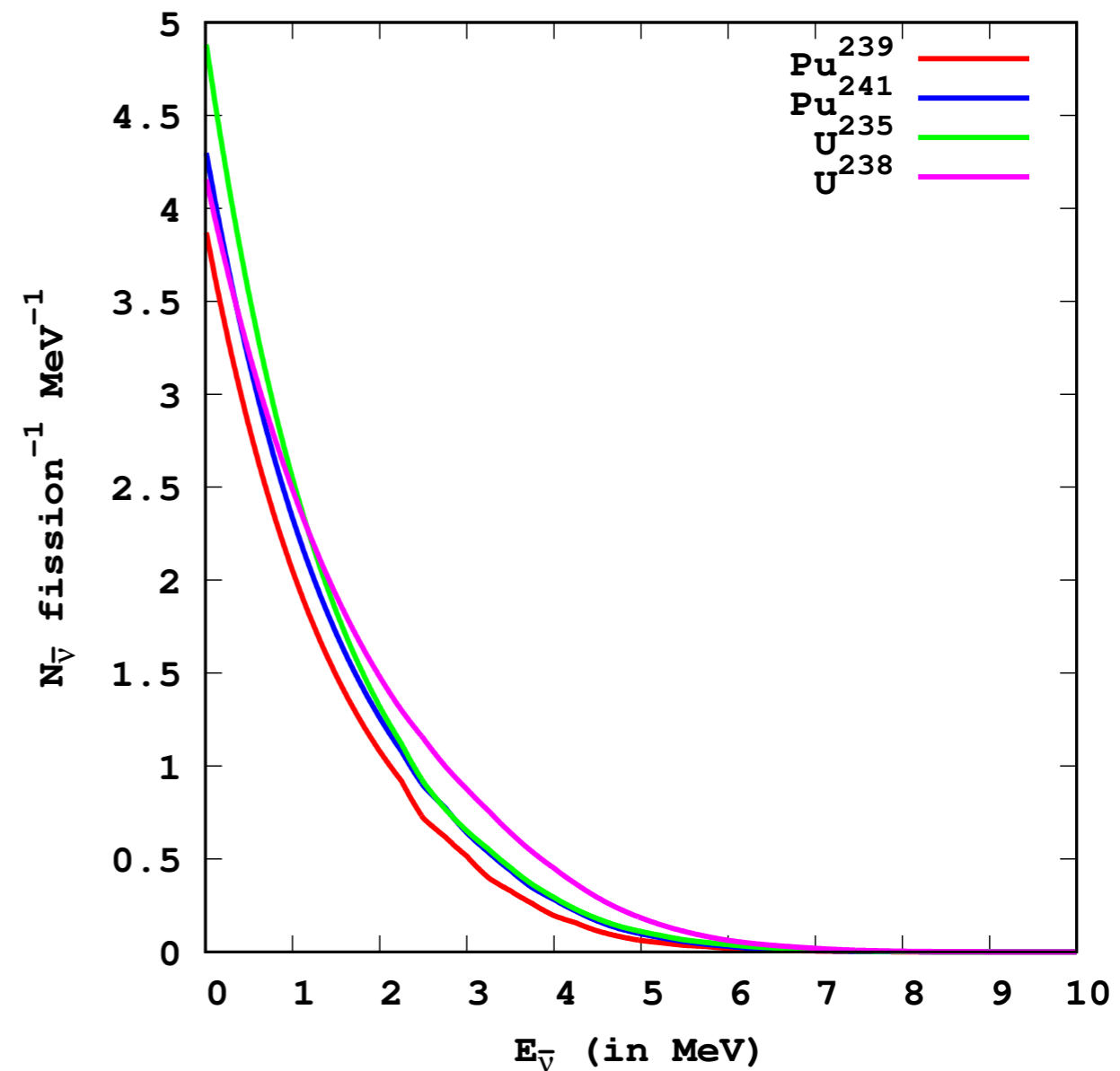
$$f_{235U} = 0.58, f_{239Pu} = 0.30, f_{238U} = 0.07, f_{241Pu} = 0.05$$

(changes with time)

Huber-Mueller Fluxes

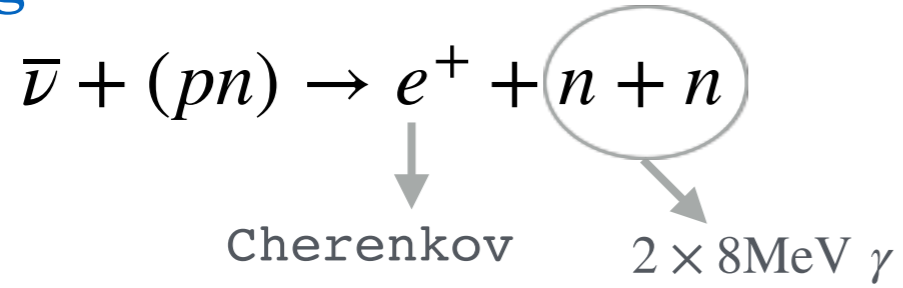
Phys. Rev. C 83 (2011) 054615

Phys. Rev. C 84 (2011) 024617

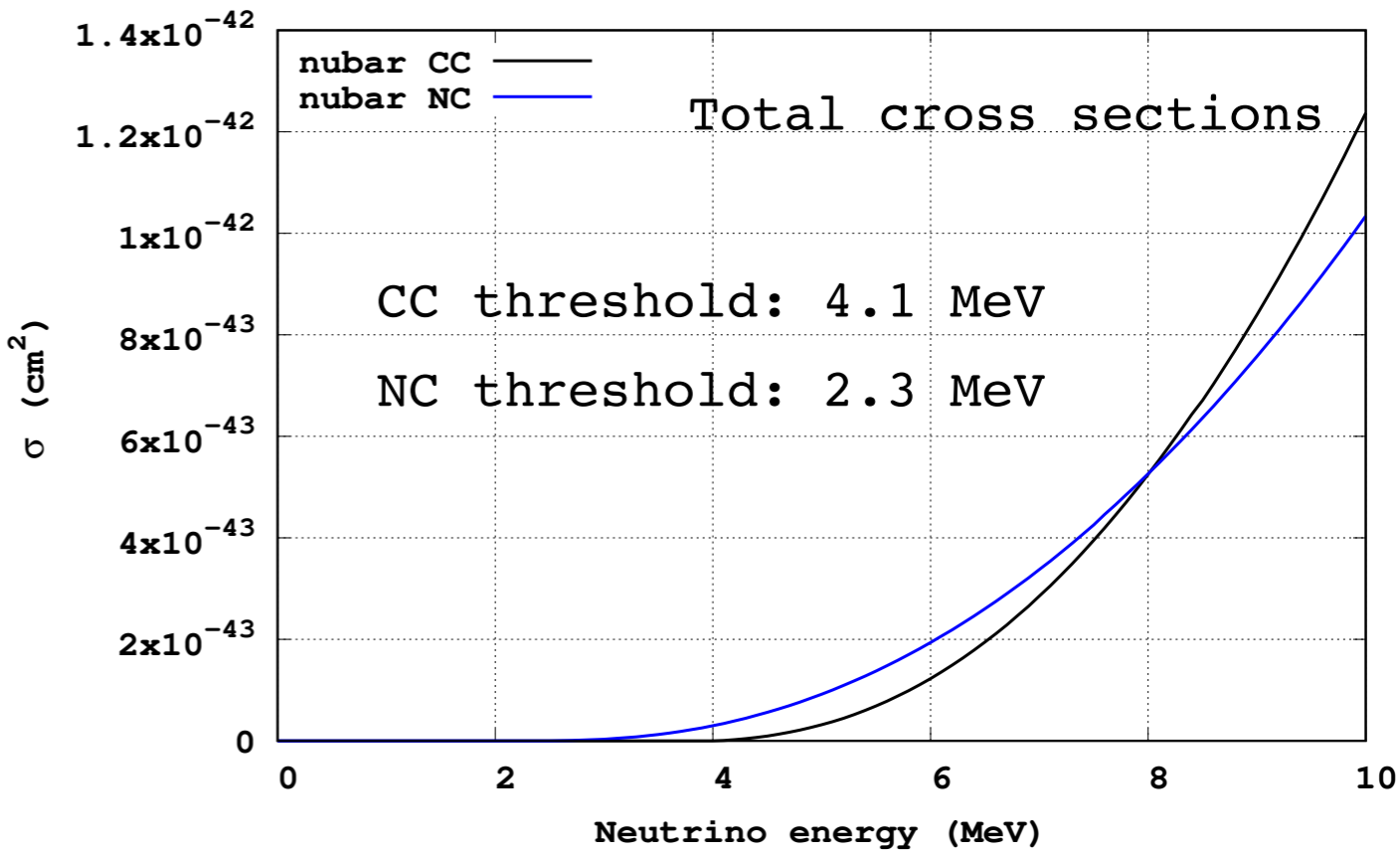
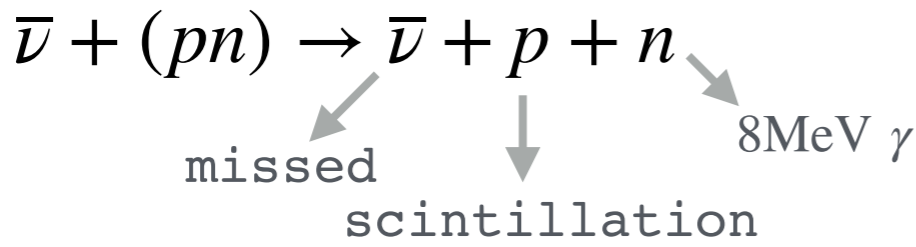


- * We consider a liquid scintillator which is 100% deuterated such as C_nD_{2n}
- * A deuterated liquid scintillator (DLS) will be sensitive to all neutrino flavors. Therefore, it can see both CC and NC events.
- * SNO experiments could conclusively establish neutrino oscillations by observing both CC and NC events. However, it could only do a counting experiment.
- * As India is one of largest producer of heavy water, availability of deuterium is not a problem.
- * DLS can detect neutrinos mainly via. dissociation of the deuteron.
- * Observing NC events in addition to CC events cannot provide crucial information regarding non-unitarity of the 3×3 PMNS matrix - especially in new physics scenarios with sterile neutrinos.
- * DLS may be able to provide spectral information of the NC events in addition to total rates.
- * However, it comes with its own set of challenges - especially with respect to background suppression and observation of scintillation signal at very low energies.

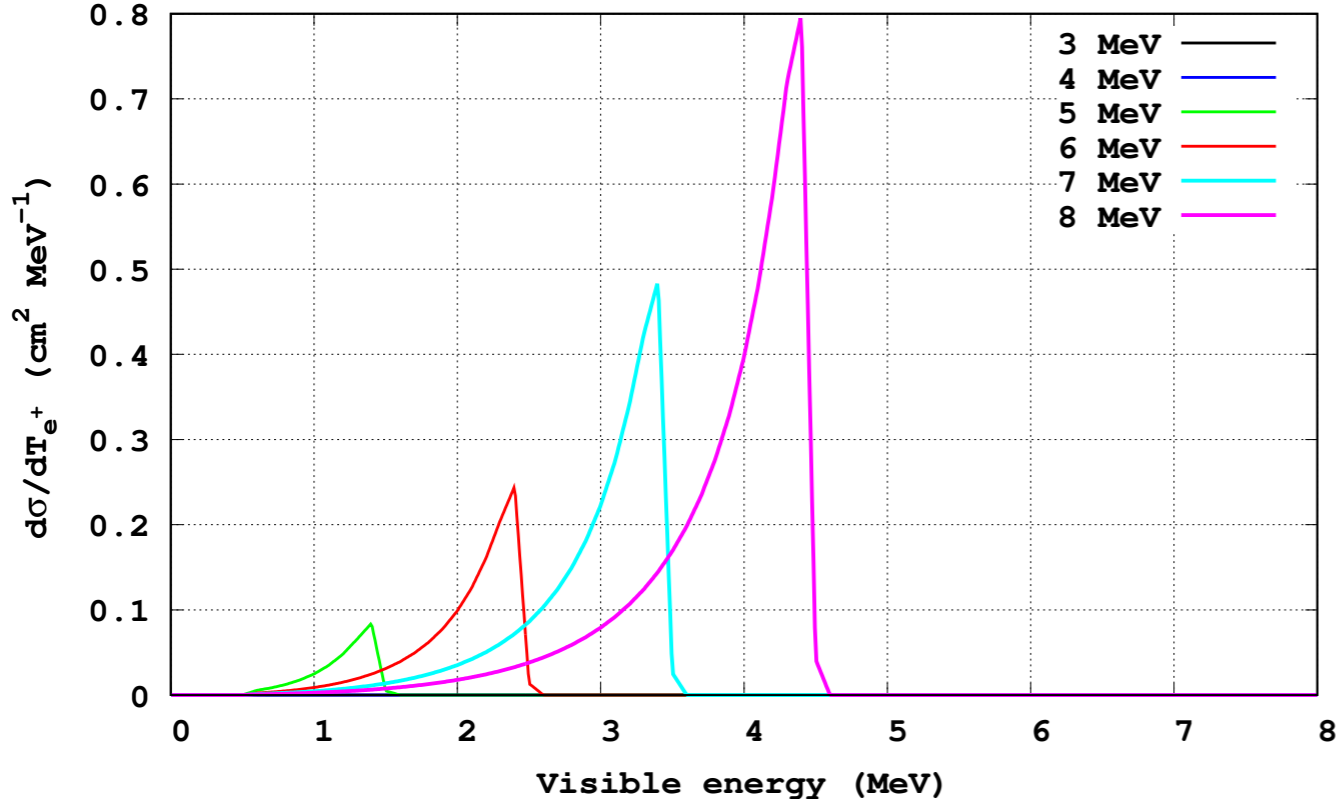
Charged Current:



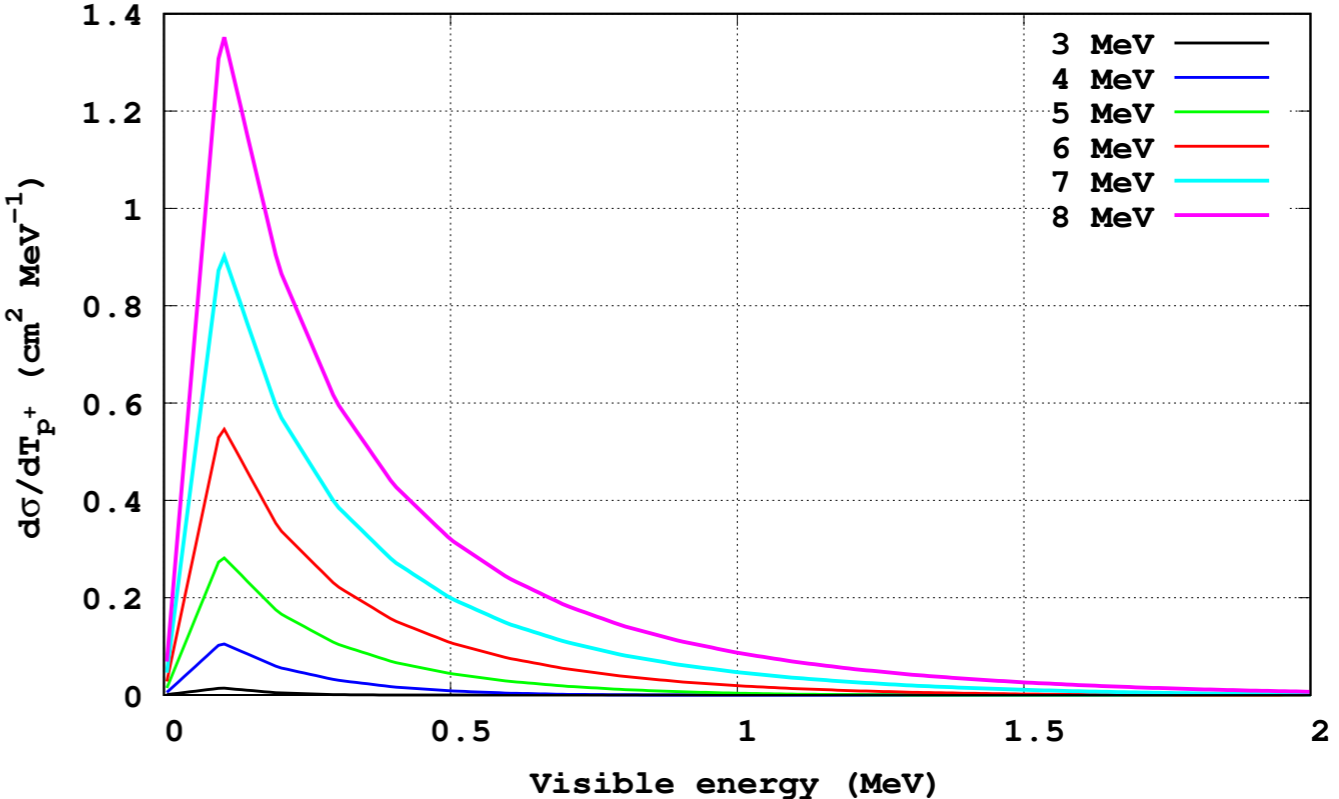
Neutral Current:



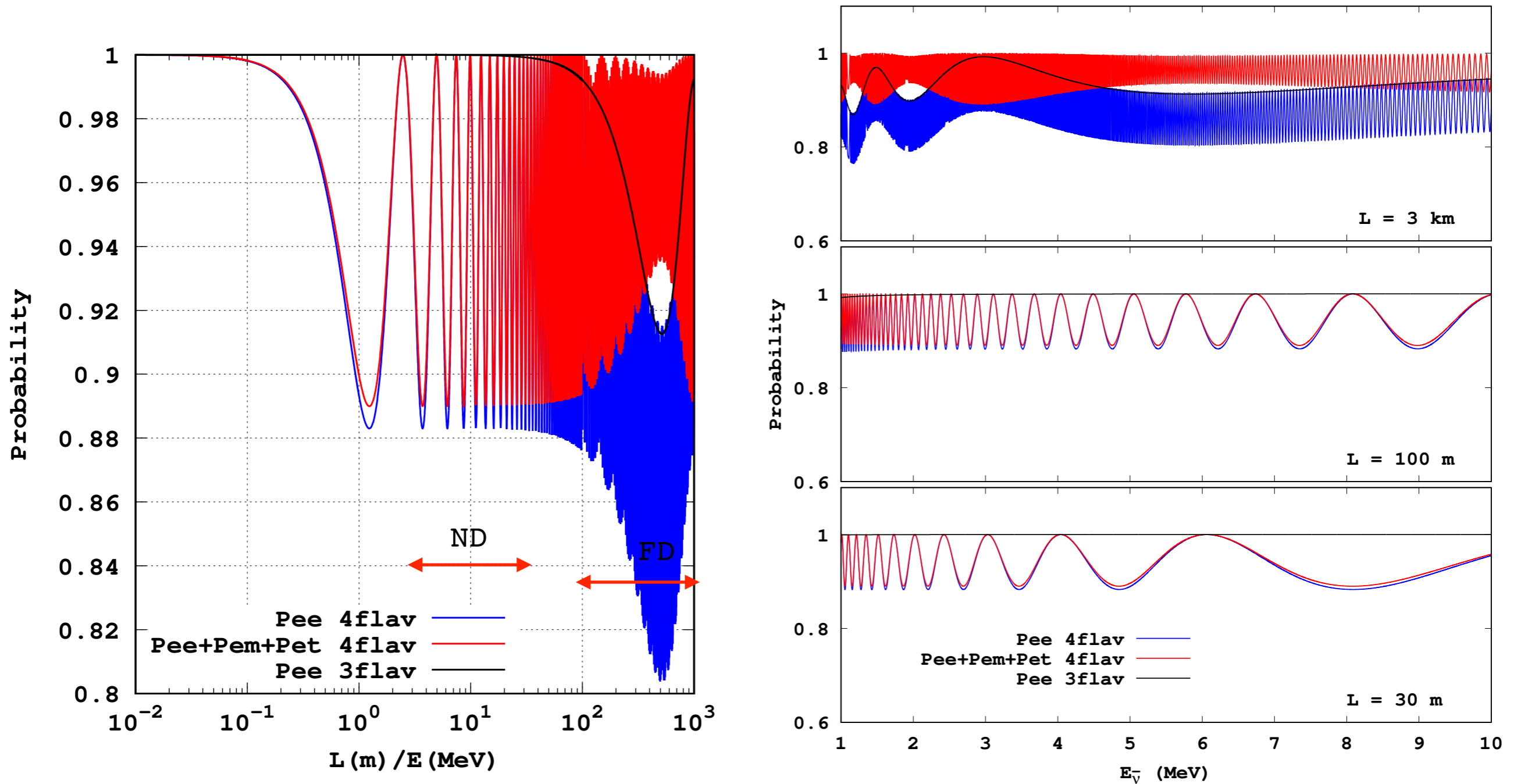
Differential xsec: CC



Differential xsec: NC



Probability plots



At near detector distances: $P_{ee} \approx P_{ee} + P_{e\mu} + P_{e\tau} \approx 1 - P_{es} \implies$ 2-neutrino oscillation: θ_{14} and Δm_{41}^2

Inensitive to Earth matter, mass ordering, octant or CP violation

Full 4-flavor effects only at far detector: sensitivity to additional oscillation parameters

Events plots

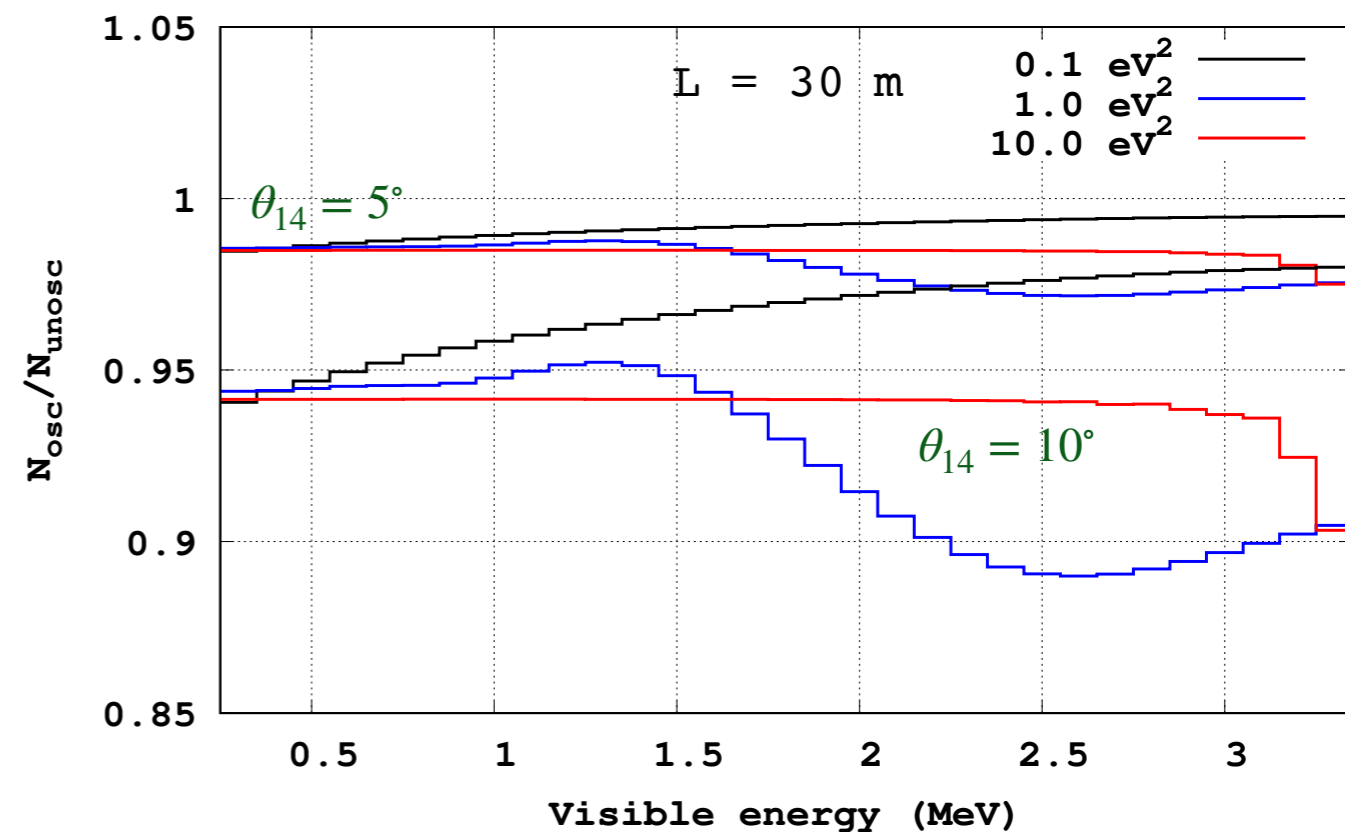
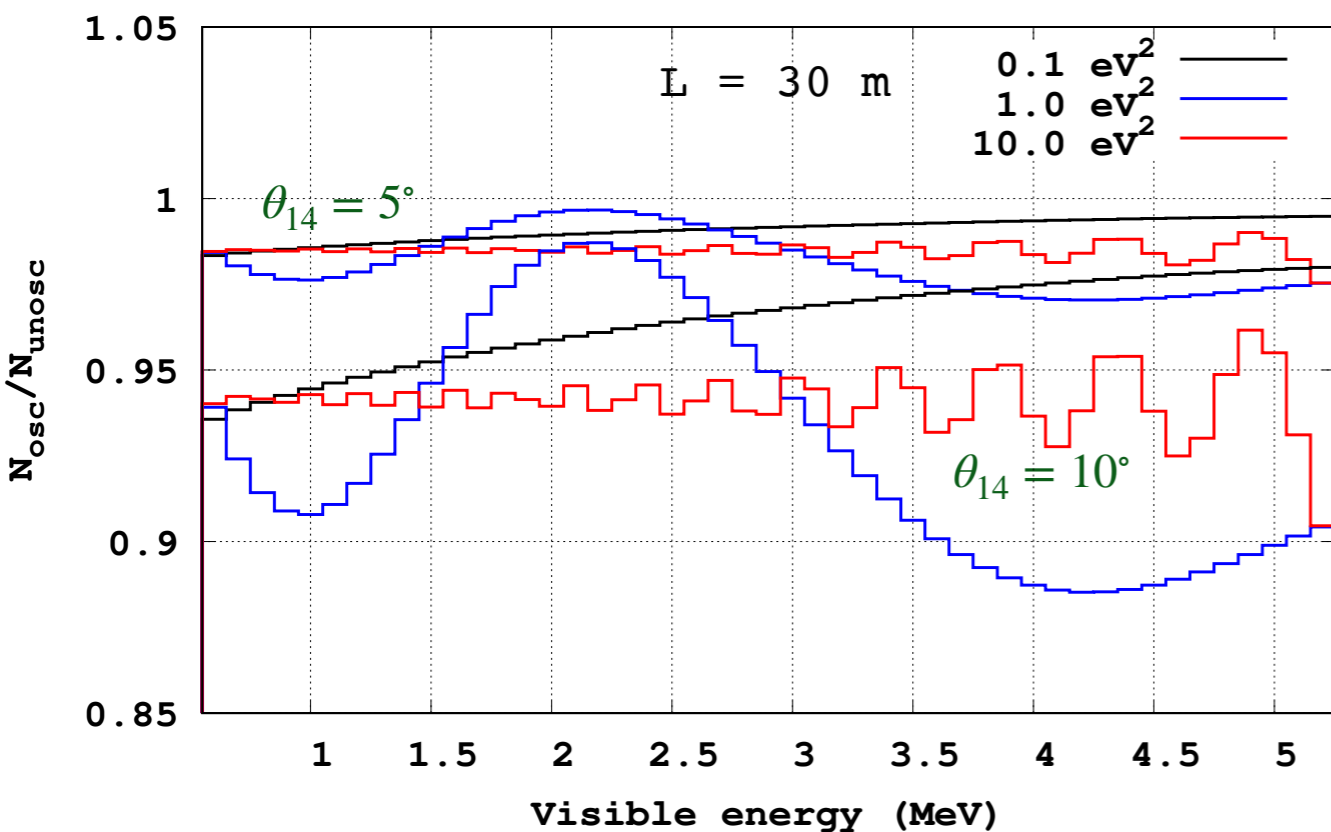
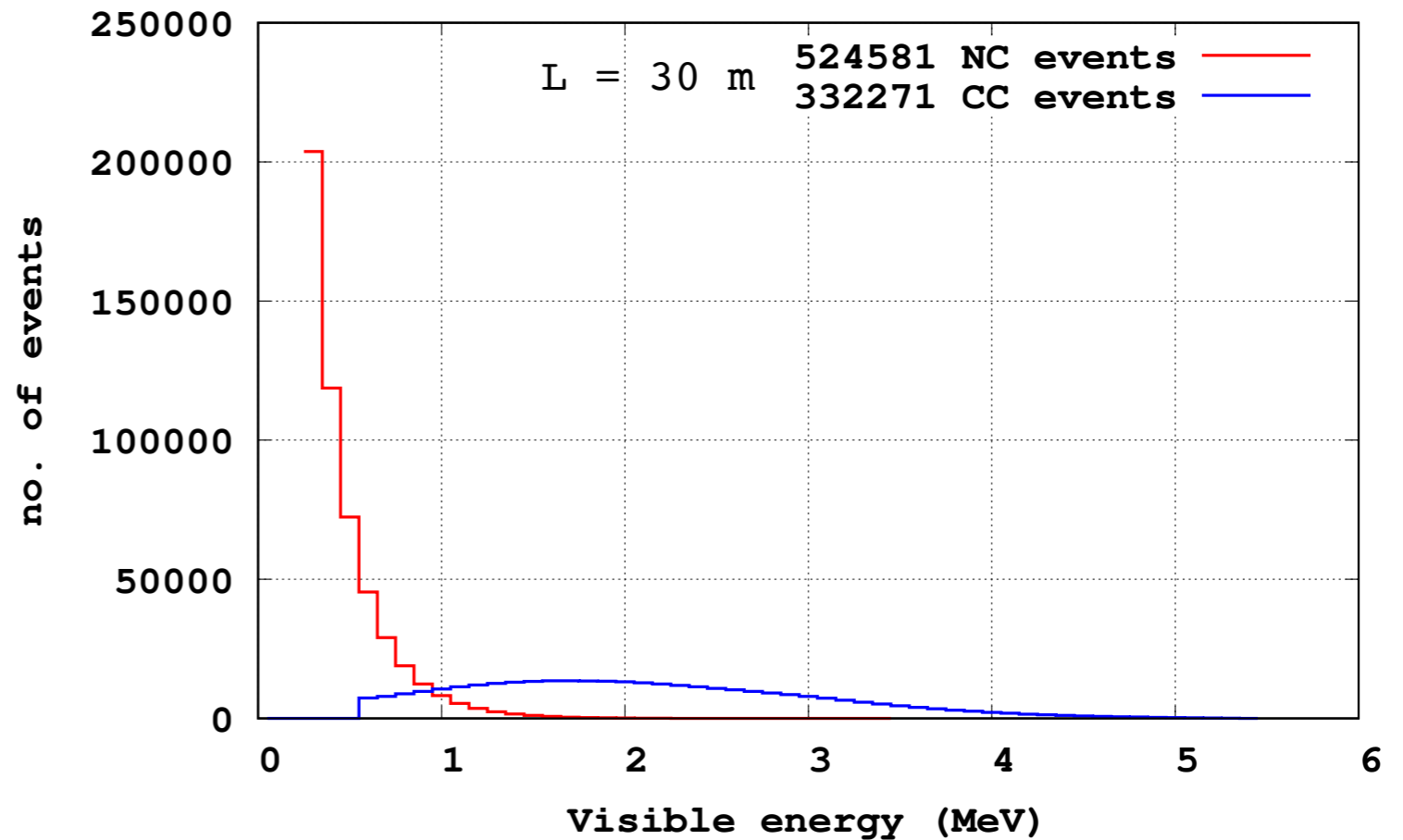
Visible Energy:

For CC: e^+

For NC: p

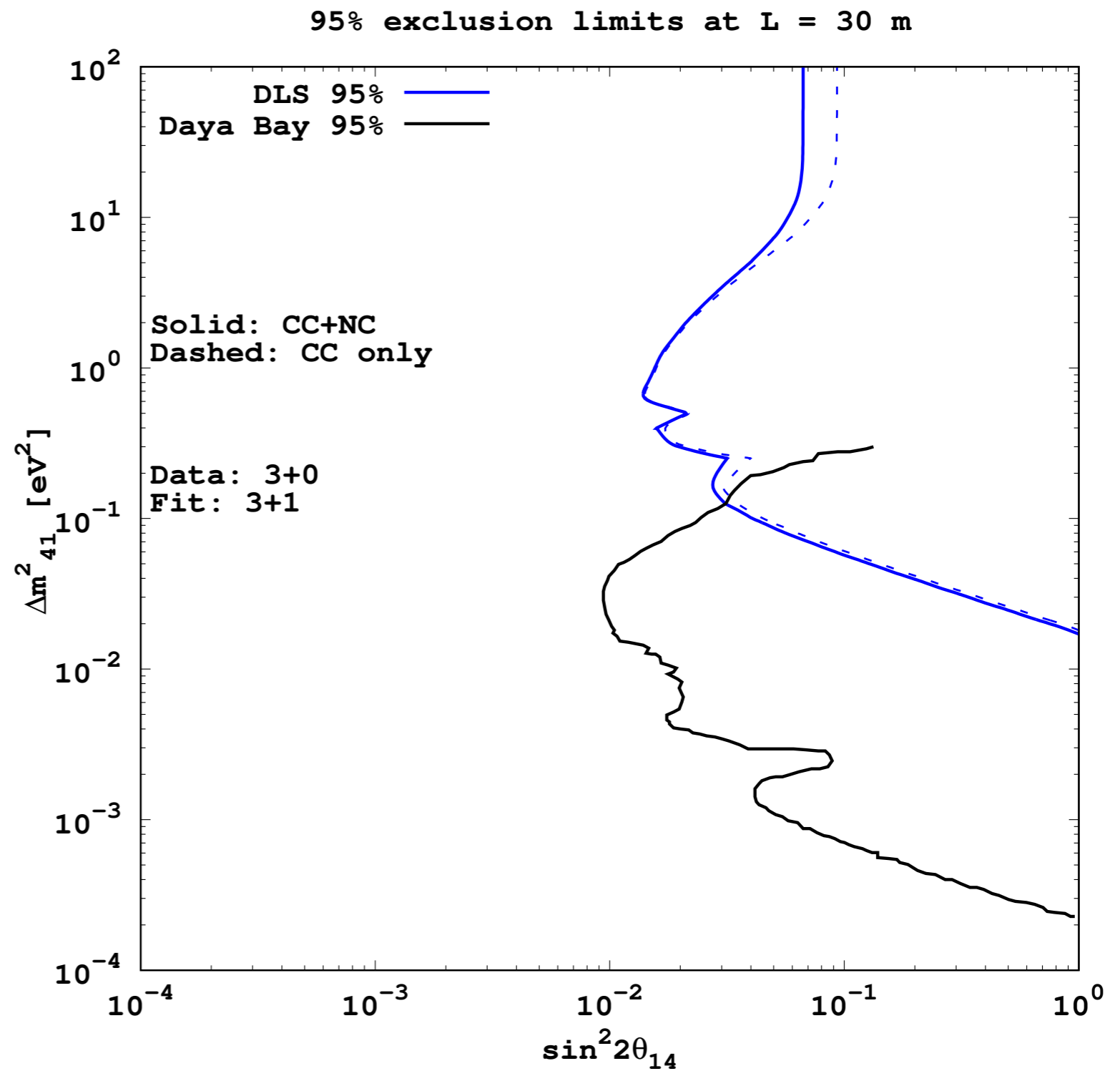
NC events pile up
between 0.0 - 3.5 MeV

Δm^2 and $\sin^2 2\theta$ - driven features in
both NC and CC spectra



Results comparable with Daya Bay
(with much less stats)

Complementarity in sensitivity



5% normalization & 5% tilt error

- * Below 200 KeV, the scintillation from radioactive C inundate the detector
- * Other radioactive impurities of heavier elements can also produce photons that can dissociate the deuterium - problem for NC events
- * Main problem in case of neutron tagging is cosmic muon induced backgrounds.
- * Possible to mitigate via:
 - a) muon veto detector
 - b) timing information of the muon events via. plastic scintillators
 - c) going underground
 - d) statistical subtraction using reactor downtime
- * Millisecond window between prompt event neutron capture

Several worldwide efforts for sterile neutrino searches

Here we explore the possibility of a DLS detector

DLS can in-principle see both CC and NC events

Sterile oscillation exclusion sensitivity in electron channel comparable to present global constraints

Main challenges in the analysis comes from cosmic muon induced background - more work going on!

Supernova neutrinos, Geoneutrinos, Solar neutrinos, low energy atmospheric neutrinos, neutrino tomography

THANK YOU FOR YOUR ATTENTION!

$$\begin{aligned}P_{\nu_e \rightarrow \nu_e} &= 1 - 4(1 - |U_{e4}|^2)|U_{e4}|^2 \sin^2(1.27\Delta m_{41}^2 L/E), \\P_{\nu_\mu \rightarrow \nu_\mu} &= 1 - 4(1 - |U_{\mu4}|^2)|U_{\mu4}|^2 \sin^2(1.27\Delta m_{41}^2 L/E), \\P_{\nu_\mu \rightarrow \nu_e} &= 4|U_{\mu4}|^2|U_{e4}|^2 \sin^2(1.27\Delta m_{41}^2 L/E),\end{aligned}$$

$$\begin{aligned}\sin^2 2\theta_{ee} &= 4(1 - |U_{e4}|^2)|U_{e4}|^2, \\ \sin^2 2\theta_{\mu\mu} &= 4(1 - |U_{\mu4}|^2)|U_{\mu4}|^2, \\ \sin^2 2\theta_{\mu e} &= 4|U_{\mu4}|^2|U_{e4}|^2,\end{aligned}$$

$$\begin{aligned}\sin^2 2\theta_{ee} &= \sin^2 2\theta_{14} &= 4(1 - |U_{e4}|^2)|U_{e4}|^2 \\ \sin^2 2\theta_{\mu\mu} &= 4 \cos^2 \theta_{14} \sin^2 \theta_{24} (1 - \cos^2 \theta_{14} \sin^2 \theta_{24}) &= 4(1 - |U_{\mu4}|^2)|U_{\mu4}|^2 \\ \sin^2 2\theta_{\tau\tau} &= 4 \cos^2 \theta_{14} \cos^2 \theta_{24} \sin^2 \theta_{34} (1 - \cos^2 \theta_{14} \cos^2 \theta_{24} \sin^2 \theta_{34}) &= 4(1 - |U_{\tau4}|^2)|U_{\tau4}|^2 \\ \sin^2 2\theta_{\mu e} &= \sin^2 2\theta_{14} \sin^2 \theta_{24} &= 4|U_{\mu4}|^2|U_{e4}|^2 \\ \sin^2 2\theta_{e\tau} &= \sin^2 2\theta_{14} \cos^2 \theta_{24} \sin^2 \theta_{34} &= 4|U_{e4}|^2|U_{\tau4}|^2 \\ \sin^2 2\theta_{\mu\tau} &= \sin^2 2\theta_{24} \cos^4 \theta_{14} \sin^2 \theta_{34} &= 4|U_{\mu4}|^2|U_{\tau4}|^2\end{aligned}$$
