



Automated Evaluation of Feynman integrals using GKZ hypergeometric systems

Based on: [arxiv.2211.01285](https://arxiv.org/abs/2211.01285)

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Need for precision calculations

- SM is not sufficient to explain the entire spectrum of observable physical phenomena - DM, dark energy, gravity ...
- Ongoing searches for BSM physics: higher luminosities and energy frontiers!
- Technological and financial challenges
- Extract more insights from whatever data we already have
- Fixed order precision calculations - involve multiloop multiscale Feynman diagrams. Complementary to the EFT approach¹.
- Calculate the corresponding integrals - usually very complicated.

¹Jonas Klappert. "Precision calculations for Higgs physics in the standard model and beyond". PhD thesis. RWTH Aachen U., 2020. DOI: [10.18154/RWTH-2020-05782](https://doi.org/10.18154/RWTH-2020-05782).

A proof-of-concept implementation demonstrating the utility of recent mathematical developments in evaluating such Feynman integrals in the form of a *Mathematica* package *FeynGKZ*².

²B. Ananthanarayan et al. “FeynGKZ: a Mathematica package for solving Feynman integrals using GKZ hypergeometric systems”. In: (Nov. 2022). arXiv: 2211.01285 [*hep-th*].

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Feynman integrals

The momentum representation

- Typically involve tensor structures in numerator - do tensor reduction
- Calculate the scalar integrals
- Momentum representation:

$$I_{\Gamma}(\nu, D) = \int \prod_{r=1}^l \frac{d^D k_r}{i\pi^{\frac{D}{2}}} \frac{1}{\prod_{j=1}^n (-q_j^2 + m_j^2)^{\nu_j}} \quad (1)$$

l : number of loops

D : the space-time dimension

$\nu = (\nu_1, \dots, \nu_n)$: propagator powers

k_r -s and q_j -s are the loop-momenta and internal-momenta for the Feynman graph Γ .

The Lee-Pomeransky representation

- An alternate form³:

$$\begin{aligned} I_{\Gamma}(\nu, D) &= \frac{\Gamma(\frac{D}{2})}{\Gamma(\frac{D}{2} - \omega)} \left(\prod_{i=1}^n \int_{\alpha_i=0}^{\infty} \frac{d\alpha_i \alpha_i^{\nu_i-1}}{\Gamma(\nu_i)} \right) G(\alpha)^{-\frac{d}{2}} \\ &= \frac{\Gamma(\frac{D}{2})}{\Gamma(\frac{D}{2} - \omega)\Gamma(\nu)} \int_{\mathbb{R}_+^n} d\alpha \alpha^{\nu-1} G(\alpha)^{-\frac{d}{2}} \end{aligned} \tag{2}$$

- Lee-Pomeransky polynomial: $G(\alpha) = U(\alpha) + F(\alpha)$.

³Roman N. Lee and Andrei A. Pomeransky. "Critical points and number of master integrals". In: *JHEP* 11 (2013), p. 165. DOI: [10.1007/JHEP11\(2013\)165](https://doi.org/10.1007/JHEP11(2013)165). arXiv: [1308.6676 \[hep-ph\]](https://arxiv.org/abs/1308.6676).

The Lee-Pomeransky representation (contd.)

- **Generalized** G-polynomial:

$$G_z(\alpha) = \sum_{a_j \in A} z_j \alpha^{a_j} = \sum_{j=1}^N z_j \prod_{i=1}^n \alpha_i^{a_{ij}} \quad (3)$$

$z_j \rightarrow$ **generic/indeterminate**

- **Generalized** Feynman integral:

$$I_{G_z}(\nu, \nu_0) = \Gamma(\nu_0) \int_{\mathbb{R}_+^n} d\alpha \alpha^{\nu-1} G_z(\alpha)^{-\nu_0} \quad (4)$$

where, $\nu_0 = \frac{D}{2}$

The associated GKZ system and its solutions

The associated GKZ system

$I_{G_z}(\nu, \nu_0)$ satisfies a **holonomic** system of PDEs called a **GKZ hypergeometric system**⁴.

Ideals

Let $P = \mathbb{F}[x_1, \dots, x_n]$ be some polynomial ring in x_1, \dots, x_n over \mathbb{F} . $\mathcal{I} \subset P$ is said to be an ideal if

- $0 \in \mathcal{I}$
- $f + g \in \mathcal{I} \quad \forall f, g \in \mathcal{I}$
- $f \cdot g \in \mathcal{I} \quad \forall f \in P, g \in \mathcal{I}$

Thus, $\langle S \rangle = \sum_i f_i g_i$; $f \in P, g \in S$ is the ideal spanned by $S \subset P$.

⁴I.M. Gelfand, Mikhail M. Kapranov, and Andrey V. Zelevinsky. “Discriminants, Resultants, and Multidimensional Determinants”. In: 1994; I.M. Gelfand, Mikhail M. Kapranov, and Andrei Zelevinsky. “Generalized Euler integrals and A-hypergeometric functions”. In: *Advances in Mathematics* 84 (1990), pp. 255–271.

The associated GKZ system (contd.)

- We describe the GKZ system as follows:

$$H_{\mathcal{A}}(\underline{\nu}) = I_{\mathcal{A}} \cup \langle \mathcal{A} \cdot \theta + \underline{\nu} \rangle \quad (5)$$

$$\begin{aligned} \mathcal{A} &= \{a_{ij}; i \in \{1, \dots, n+1\}, j \in \{1, \dots, N\} \mid a_{ij} = 1; i = 1\} \\ \underline{\nu} &= (\nu_0, \nu_1, \dots, \nu_n)^T \end{aligned} \quad (6)$$

- $\mathcal{A} \rightarrow (n+1) \times N$ matrix; $n+1 \leq N$
- $\theta = (\theta_1, \dots, \theta_N)^T$; $\theta_i = z_i \partial_i \rightarrow$ Euler operators
- **Assume:** $(1, \dots, 1)$ lies in \mathbb{Q} -row span of \mathcal{A}

The associated GKZ system (contd.)

- $H_{\mathcal{A}}(\underline{\nu})I_{G_z}(\nu, \nu_0) = 0$
- $I_{G_z}(\nu, \nu_0) \rightarrow$ *GKZ hypergeometric function!*⁵

⁵René Pascal Klausen. “Hypergeometric Series Representations of Feynman Integrals by GKZ Hypergeometric Systems”. In: (2020). eprint: [1910.08651](#); Leonardo de la Cruz. “Feynman integrals as A-hypergeometric functions”. In: (2019). eprint: [1907.00507](#).

- **Algebraically:** the **SST algorithm**⁶ → the **Gröbner deformation** method
- **Geometrically:** the **triangulation** method
- Both are equivalent!
- In this talk, focus on the **geometric** picture.

⁶Mutsumi Saito, Bernd Sturmfels, and Nobuki Takayama. *Gröbner deformations of hypergeometric differential equations*. Vol. 6. Springer Science & Business Media, 2013.

Solving the GKZ system (contd.)

- We saw:

$$\mathcal{A} = \begin{pmatrix} 1 \\ A \end{pmatrix} = \begin{pmatrix} 1 & 1 & \dots & 1 \\ a_1 & a_2 & \dots & a_N \end{pmatrix} \in \mathbb{Z}_{\geq 0}^{(n+1) \times N} \quad (7)$$

- A defines an assembly of N points (a point configuration) in \mathbb{Z}^n

$$\text{Conv}(A) := \left\{ \sum_{j=1}^N k_j a_j \mid k \in \mathbb{R}_{\geq 0}^N, \sum_{j=1}^N k_j = 1 \right\} \quad (8)$$

- **Newton polytope** of $G_Z(\alpha)$:

$$\Delta_{G_Z} := \text{Conv}(A) \quad (9)$$

Solving the GKZ system (contd.)

- Triangulate Δ_{G_2} !
- Triangulation structure: $T = \{\sigma_1, \dots, \sigma_r\}$.
- $\sigma_i \subset \{1, \dots, N\}$ is some index set.

Solving the GKZ system (contd.)

Can always obtain a [regular triangulation](#)!⁷

⁷Israel M. Gelfand, Mikhail M. Kapranov, and Andrei Zelevinsky. "HYPERGEOMETRIC FUNCTIONS, TORIC VARIETIES AND NEWTON POLYHEDRA". In: 1991.

Solving the GKZ system (contd.)

Can always obtain a **unimodular regular triangulation** ($\text{vol}_0(\sigma_i) = 1$)!⁸

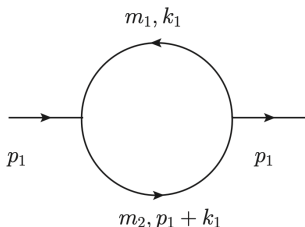
⁸W. Bruns and J. Gubeladze. *Polytopes, Rings, and K-Theory*. ISBN: 9780387763569; Finn F. Knudsen. "Construction of nice polyhedral subdivisions". In: 1973.

Solving the GKZ system (contd.)

- Regular triangulations can be used to construct a basis for the **finite-dimensional** solution space of $H_{\mathcal{A}}(\underline{\nu})$
- Each element: Γ -series
- Whole solution: linear combination of the Γ -series elements
- Unimodularity: **one $\sigma_j \rightarrow$ one Γ -series**
- Might as well use just the unimodular regular triangulations to construct a basis!

Demonstration

Bubble diagram with two unequal masses



The corresponding integral in momentum-representation:

$$I_{\Gamma}(\nu_1, \nu_2, D; p_1^2) = \int \frac{d^D k_1}{i\pi^{\frac{D}{2}}} \frac{1}{(-k_1^2 + m_1^2)^{\nu_1} (-(p_1 + k_1)^2 + m_2^2)^{\nu_2}} \quad (10)$$

with two unequal masses m_1 and m_2 , and external momentum p_1 .

Bubble diagram with two unequal masses (contd.)

After successfully loading the package and installing its dependencies, specify the integral in its momentum representation as:

```
In[3] := MomentumRep = {{k1, m1, a1}, {p1 + k1, m2, a2}};  
LoopMomenta = {k1};  
InvariantList = {p12 → -s};  
Dim = 4 - 2ε;  
Prefactor = 1;
```

Bubble diagram with two unequal masses (contd.)

Now derive the \mathcal{A} -matrix:

```
In[4] := FindAMatrixOut = FindAMatrix[{MomentumRep, LoopMomenta,  
InvariantList, Dim, Prefactor}, UseMB → False];
```

```
Prints ⇒ The Symanzik polynomials →  $U = x_1 + x_2$   
          ,  $F = m_1^2 x_1^2 + s x_1 x_2 + m_1^2 x_1 x_2 + m_2^2 x_1 x_2 + m_2^2 x_2^2$ 
```

```
The Lee-Pomeransky polynomial →  $G =$   
           $x_1 + m_1^2 x_1^2 + x_2 + s x_1 x_2 + m_1^2 x_1 x_2 + m_2^2 x_1 x_2 + m_2^2 x_2^2$ 
```

```
The associated  $\mathcal{A}$ -matrix →  $\begin{pmatrix} 1 & 1 & 1 & 1 & 1 \\ 2 & 1 & 1 & 0 & 0 \\ 0 & 1 & 0 & 2 & 1 \end{pmatrix}$ , which has  $\text{codim} = 2$ .
```

```
Normalized Volume of the associated Newton Polytope → 3
```

```
Time Taken 1.50005 seconds
```

Bubble diagram with two unequal masses (contd.)

Compute the unimodular regular triangulations⁹:

```
In[5]:=  Triangulations = FindTriangulations[FindAMatrixOut];

Prints =>  Finding all regular triangulations ...
           Found 5 Regular Triangulations, out of which 3 are Unimodular
           The 3 Unimodular Regular Triangulations →
           1 ::  {{1,2,3},{2,3,4},{3,4,5}}
           2 ::  {{1,2,3},{2,4,5},{2,3,5}}
           3 ::  {{2,4,5},{1,3,5},{1,2,5}}
           Time Taken 0.126965 seconds
```

⁹Jörg Rambau. "TOPCOM: Triangulations of Point Configurations and Oriented Matroids". In: *Proceedings of the International Congress of Mathematical Software*. 2002. URL: <http://www.zib.de/PaperWeb/abstracts/ZR-02-17>.

Bubble diagram with two unequal masses (contd.)

Calculate the Γ -series:

```
In[7]:= SeriesSolution = SeriesRepresentation[Triangulations,2];
```

```
Prints => Unimodular Triangulation -> 2
```

```
Number of summation variables -> 2
```

```
Non-generic limit -> {z1 -> m1^2, z2 -> s + m1^2 + m2^2, z3 -> 1, z4 -> m2^2, z5 -> 1}
```

```
The series solution is the sum of following 3 terms.
```

```
Term 1 ::
```

$$\left(\left((-1)^{-n_1-n_2} \Gamma[-2+\epsilon+a_1-n_1-n_2] \Gamma[4-2\epsilon-a_1-a_2+n_2] \right. \right. \\ \left. \Gamma[a_2+2n_1+n_2] (m_1^2)^{2-\epsilon-a_1} \left(\frac{m_1^2 m_2^2}{(s+m_1^2+m_2^2)^2} \right)^{n_1} \left(\frac{m_1^2}{s+m_1^2+m_2^2} \right)^{n_2} \right. \\ \left. (s+m_1^2+m_2^2)^{-a_2} \right) / \left(\Gamma[a_1] \Gamma[4-2\epsilon-a_1-a_2] \Gamma[a_2] \right. \\ \left. \Gamma[1+n_1] \Gamma[1+n_2] \right)$$

```
Term 2 ::
```

$$\left(\left((-1)^{-n_1-n_2} \Gamma[-2+\epsilon+a_2-n_1-n_2] \Gamma[4-2\epsilon-a_1-a_2+n_2] \right. \right. \\ \left. \Gamma[a_1+2n_1+n_2] (m_2^2)^{2-\epsilon-a_2} \left(\frac{m_1^2 m_2^2}{(s+m_1^2+m_2^2)^2} \right)^{n_1} \left(\frac{m_2^2}{s+m_1^2+m_2^2} \right)^{n_2} \right. \\ \left. (s+m_1^2+m_2^2)^{-a_1} \right) / \left(\Gamma[a_1] \Gamma[4-2\epsilon-a_1-a_2] \Gamma[a_2] \right. \\ \left. \Gamma[1+n_1] \Gamma[1+n_2] \right)$$

```
Term 3 ::
```

$$\left(\left((-1)^{-n_1-n_2} \Gamma[2-\epsilon-a_2+n_1-n_2] \Gamma[2-\epsilon-a_1-n_1+n_2] \right. \right. \\ \left. \Gamma[-2+\epsilon+a_1+a_2+n_1+n_2] \left(\frac{m_1^2}{s+m_1^2+m_2^2} \right)^{n_1} \left(\frac{m_2^2}{s+m_1^2+m_2^2} \right)^{n_2} \right. \\ \left. (s+m_1^2+m_2^2)^{2-\epsilon-a_1-a_2} \right) / \left(\Gamma[a_1] \Gamma[4-2\epsilon-a_1-a_2] \right. \\ \left. \Gamma[a_2] \Gamma[1+n_1] \Gamma[1+n_2] \right)$$

```
Time Taken 0.066558 seconds
```


Bubble diagram with two unequal masses (contd.)

Check for an expression in terms of known hypergeometric functions¹⁰:

```
In[8]:= GetClosedForm[SeriesSolution];

Prints => Closed form found with Olsson!
Term 1 ::

$$\frac{1}{\Gamma[a_1]} \Gamma[-2 + \epsilon + a_1]$$


$$H3\left[a_2, 4 - 2\epsilon - a_1 - a_2, 3 - \epsilon - a_1, \frac{m_1^2 m_2^2}{(s + m_1^2 + m_2^2)^2}, \frac{m_1^2}{s + m_1^2 + m_2^2}\right]$$


$$m_1^4 (m_1^2)^{-\epsilon - a_1} (s + m_1^2 + m_2^2)^{-a_2}$$

Term 2 ::

$$\frac{1}{\Gamma[a_2]} \Gamma[-2 + \epsilon + a_2]$$


$$H3\left[a_1, 4 - 2\epsilon - a_1 - a_2, 3 - \epsilon - a_2, \frac{m_1^2 m_2^2}{(s + m_1^2 + m_2^2)^2}, \frac{m_2^2}{s + m_1^2 + m_2^2}\right]$$


$$m_2^4 (m_2^2)^{-\epsilon - a_2} (s + m_1^2 + m_2^2)^{-a_1}$$

Term 3 ::

$$\left( \left( G1\left[-2 + \epsilon + a_1 + a_2, 2 - \epsilon - a_1, 2 - \epsilon - a_2, -\frac{m_2^2}{s + m_1^2 + m_2^2}\right], -\frac{m_1^2}{s + m_1^2 + m_2^2} \right) \Gamma[2 - \epsilon - a_1] \Gamma[2 - \epsilon - a_2] \right.$$


$$\left. \Gamma[-2 + \epsilon + a_1 + a_2] (s + m_1^2 + m_2^2)^{2 - \epsilon - a_1 - a_2} \right) / (\Gamma[a_1]$$


$$\Gamma[4 - 2\epsilon - a_1 - a_2] \Gamma[a_2])$$

Time Taken 0.05827 seconds
```

¹⁰B. Ananthanarayan et al. "Olsson.wl : a *Mathematica* package for the computation of linear transformations of multivariable hypergeometric functions". In: (Dec. 2021). arXiv: 2201.01189 [cs.MS].

Bubble diagram with two unequal masses (contd.)

Evaluate the sum of the Γ -series terms numerically:

```
In[9] := SumLim = 30;  
ParameterSub = { $\epsilon \rightarrow 0.001$ ,  $a_1 \rightarrow 1$ ,  $a_2 \rightarrow 1$ ,  $s \rightarrow 10$ ,  $m_1 \rightarrow 0.4$ ,  $m_2 \rightarrow 0.3$ };  
NumericalSum[SeriesSolution, ParameterSub, SumLim];
```

```
Prints  $\Rightarrow$  Numerical result = 997.382  
Time Taken 0.222572 seconds
```

Summary and Future Works

- Very rich mathematical structures appear in context of computing multiloop multiscale Feynman integrals
- In-depth analysis of such structures might furnish insights for developing novel computational frameworks and algorithms for evaluating these integrals.
- Laurent expansion of hypergeometric functions in the dim-reg parameter ϵ - critical bottleneck in this approach. Known automated implementations: *HypExp*, *XSummer*, also private in-house implementations.

- Tackle the issue of ϵ -expansion of hypergeometric functions: a new algorithm has recently been proposed in [arxiv.2208.01000](https://arxiv.org/abs/2208.01000)¹¹.
- Convert the standing proof-of-concept implementation into a performance-driven one. Current performance bottlenecks stem from an excess of dependence on *Mathematica*, particularly in the numerical summation step.
- Explore the scope of *FeynGKZ* in evaluating stringy canonical forms, inspired by [arxiv.2005.07395](https://arxiv.org/abs/2005.07395)¹². To be taken up soon.

¹¹Souvik Bera. “ ϵ -Expansion of Multivariable Hypergeometric Functions Appearing in Feynman Integral Calculus”. In: (Aug. 2022). arXiv: 2208.01000 [math-ph].




¹²Song He et al. “Stringy canonical forms and binary geometries from associahedra, cyclohedra and generalized permutohedra”. In: *JHEP* 10 (2020), p. 054. DOI: 10.1007/JHEP10(2020)054. arXiv: 2005.07395 [hep-th].






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