

Leptoquarks (LQ)

- Coloured BSM particles, couple to a lepton and a quark.
- Can be scalar or vector.

Motivations 1

- Radiative Majorana mass of neutrinos.
 - 4.2σ experimental excess on $(g - 2)_\mu$.
- $$\Delta a_\mu^{Ex-SM} = (2.51 \pm 0.59) \times 10^{-9}$$

Motivation 2

Contributes to b -decay anomalies:

- NC anomaly: $b \rightarrow s\mu^+\mu^-$
- CC anomaly: $b \rightarrow c\tau^-\bar{\nu}$

Model 1: $\tilde{R}_2 + S_1$

$$\tilde{R}_2(3, 2, 1/6) = (\tilde{R}_2^{+2/3}, \tilde{R}_2^{-1/3})^T, S_1(\bar{3}, 1, 1/3)$$

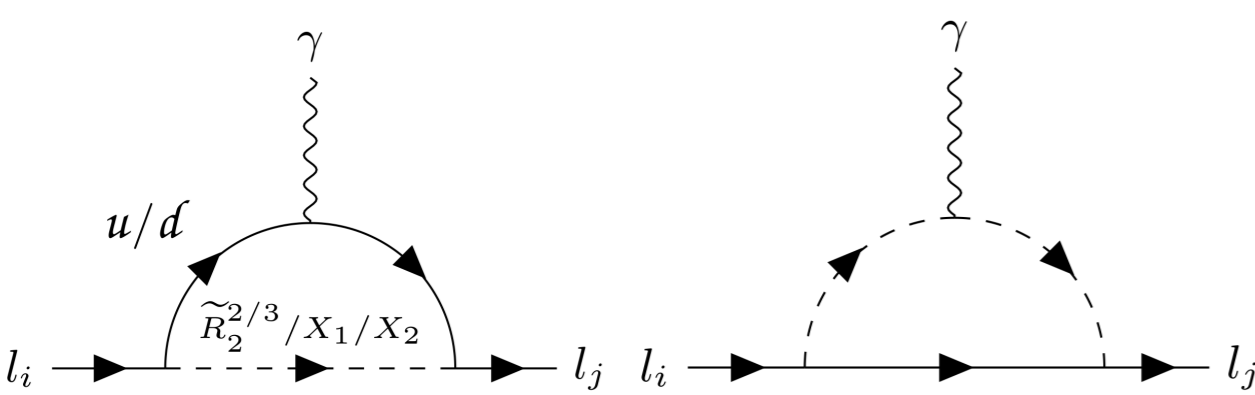
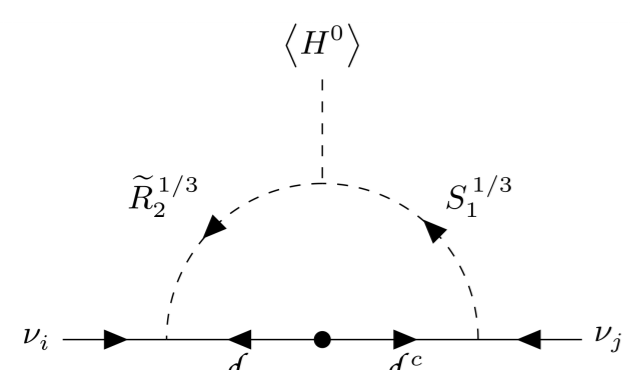
$$\mathcal{L} \supseteq \mathcal{Y}_1^L \bar{Q}_L^c S_1 (i\sigma_2) L_L + \mathcal{Y}_1^R \bar{u}_R^c S_1 l_R + \mathcal{Y}_2 \bar{d}_R \tilde{R}_2^T (i\sigma_2) L_L + \kappa H^\dagger \tilde{R}_2 S_1 + h.c. [1]$$

After EWSB, $\tilde{R}_2^{-1/3}$ and $S_1^{-1/3}$ mix \rightarrow
 $\tan 2\theta_{LQ} = \frac{-\sqrt{2}\kappa v_h}{m_S^2 - m_R^2}$

Mass eigenstates

- $\tilde{R}_2^{+2/3}$: pure doublet
- $X_1^{-1/3}$: mostly singlet
- $X_2^{-1/3}$: mostly doublet

$$M_\nu \simeq \frac{3 \sin 2\theta_{LQ}}{32\pi^2} \ln \left(\frac{M_1^2}{M_2^2} \right) \times [\mathcal{Y}_1^L m_d \mathcal{Y}_2^T + \mathcal{Y}_2 m_d (\mathcal{Y}_1^L)^T]$$



$(g - 2)_e$
 $\mu \rightarrow e\gamma$
other CLFV bounds

Benchmark: 3 sets of \mathcal{Y} , $M_{LQ} \sim 1.5$ TeV, $|\theta_{LQ}| \sim 0.618$ rad

Pair production at the LHC/FCC

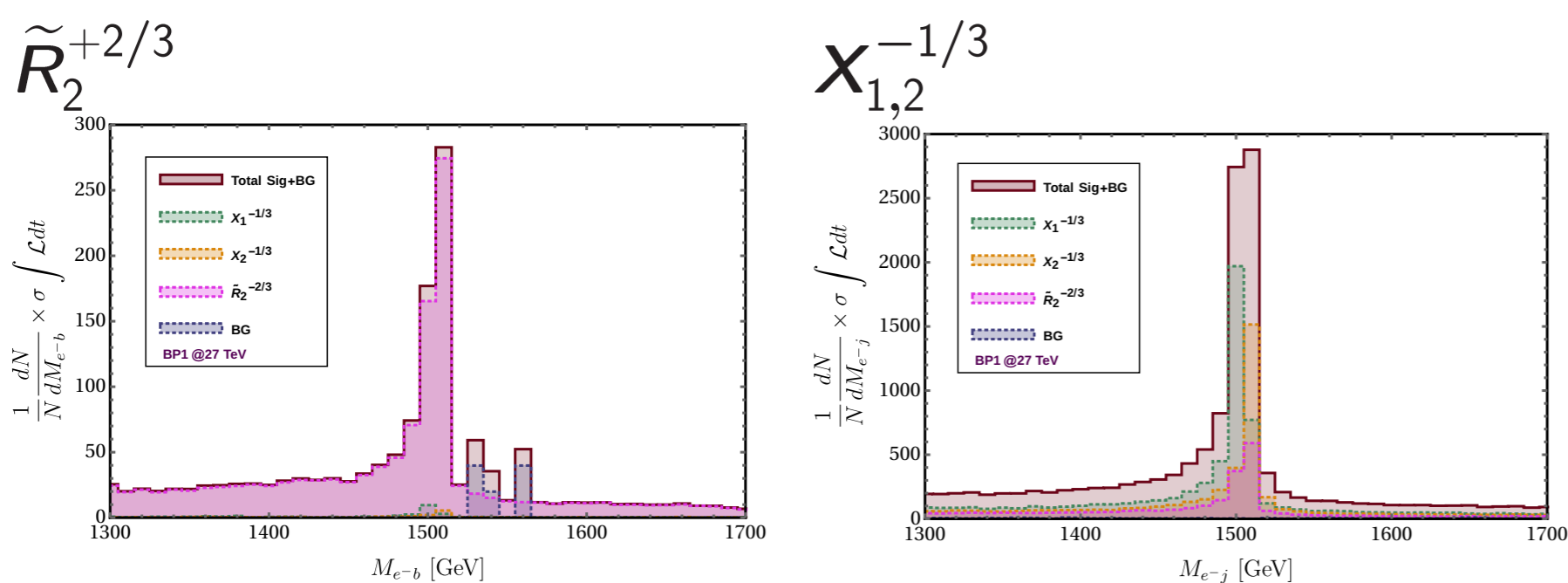
Model signatures

- 2 light- or c -jet + OSD
- 5σ at 14 TeV, 130 fb^{-1}

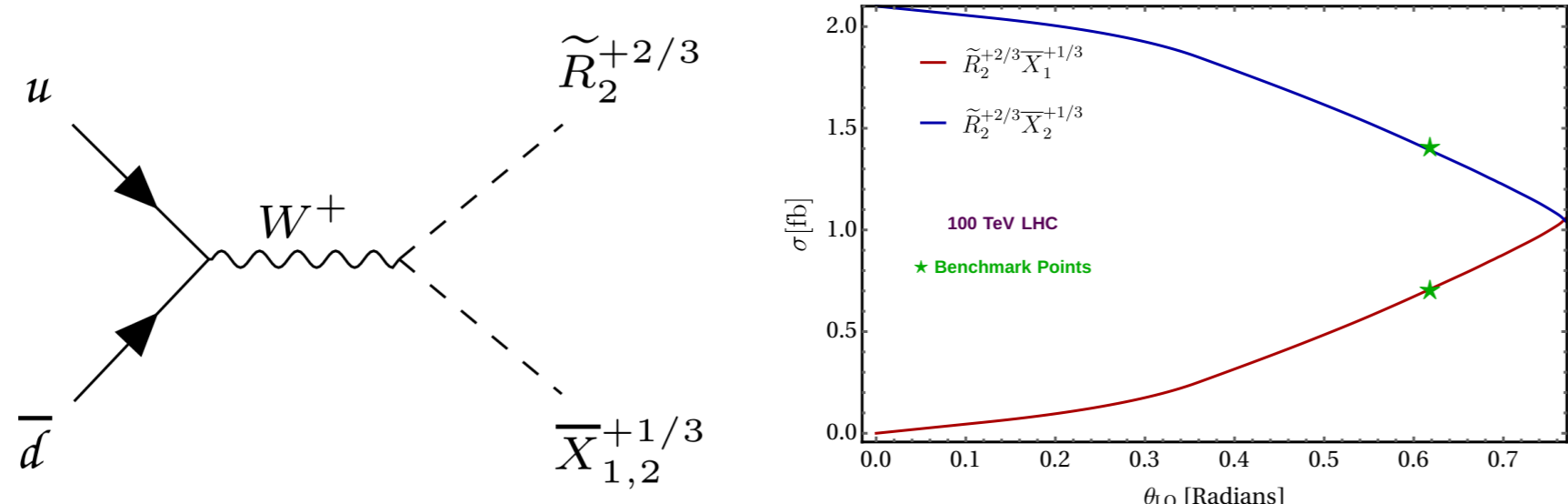
Distinguishing pure vs mixed LQ states

- $\tilde{R}_2^{+2/3} \rightarrow d\ell^+$ via \mathcal{Y}_2 only, $X_{1,2}^{-1/3} \rightarrow u\ell^-, d\nu$ via $\mathcal{Y}_1^{L,R}, \mathcal{Y}_2$.
- Four complementary final states are analyzed
- discerns between $\tilde{R}_2^{+2/3}$ and $X_{1,2}^{-1/3}$.
- Very difficult to segregate $X_1^{-1/3}$ and $X_2^{-1/3}$.

Invariant mass reconstruction



$pp \rightarrow W^+ \rightarrow \tilde{R}_2^{+2/3} X_{1,2}^{-1/3}$
More doublet % \rightarrow higher $\sigma_{Prod} \rightarrow$ probe of θ_{LQ}



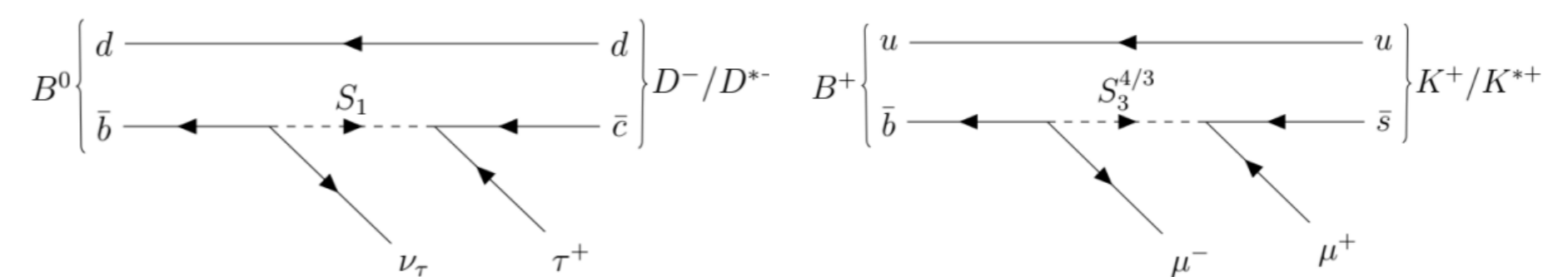
Model 2: S_1 and S_3 separately

$$S_1(\bar{3}, 1, 1/3)$$

$$\mathcal{L}_{S_1} \supseteq \bar{Q}^c i\tau_2 Y_{S_1}^{\alpha} L^\alpha S_1 + \bar{u}_R^c i Z_{S_1}^{\alpha} \ell_R^\alpha S_1 + h.c.$$

$$S_3(\bar{3}, 3, 1/3) = (S_3^{4/3}, S_3^{1/3}, S_3^{-2/3})^T$$

$$\mathcal{L}_{S_3} \supseteq \bar{Q}^c i Y_{S_3}^{\alpha} i\tau_2 \tau \cdot S_3 L^\alpha + h.c. [2]$$

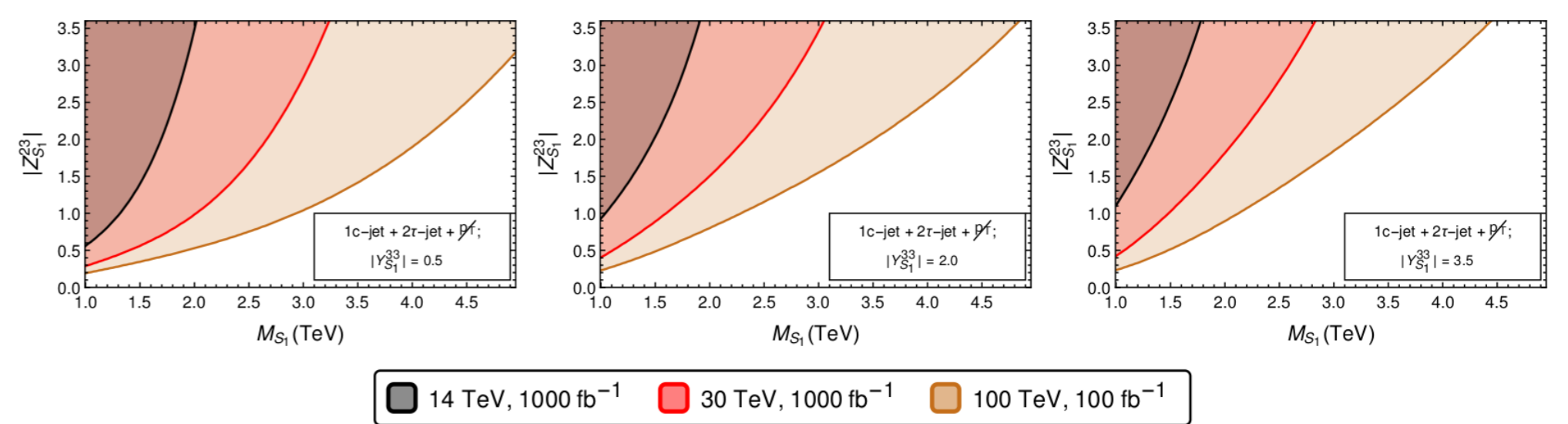


$$Y_{S_1}^{33}, Z_{S_1}^{23}: b \rightarrow c\tau^-\bar{\nu} \quad Y_{S_3}^{22}, Y_{S_3}^{32}: b \rightarrow s\mu^+\mu^-$$

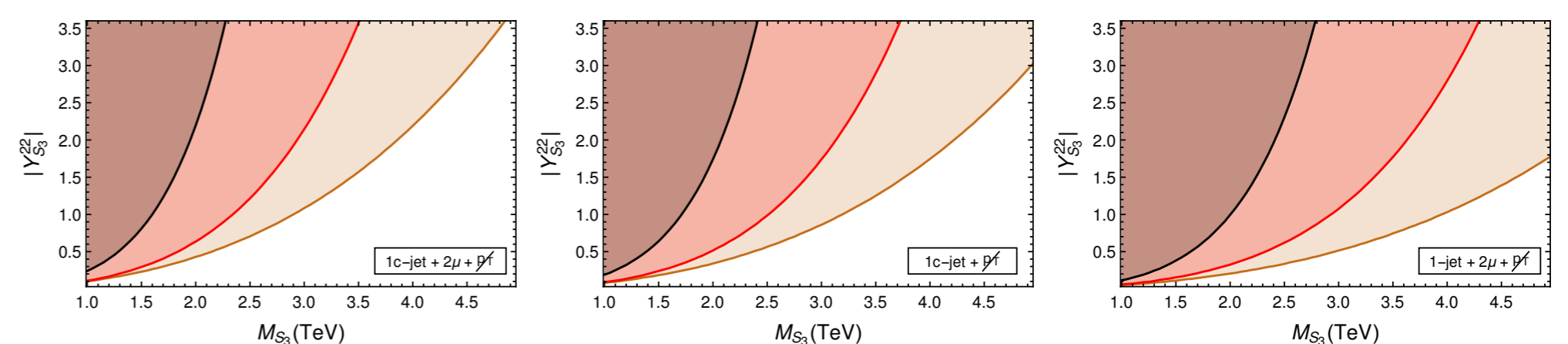
Single production at the LHC/FCC

$q - g \rightarrow LQ \ell/\nu$: direct probes of Y, Z .

Reach for S_1

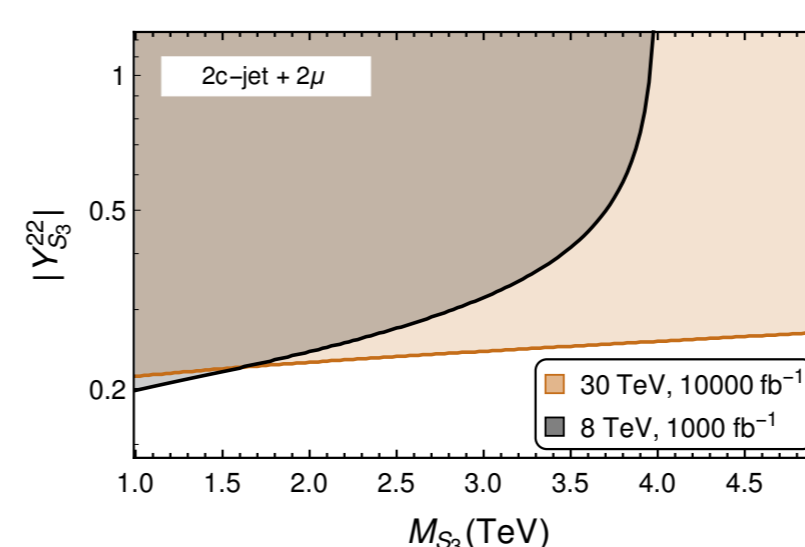


Reach for S_3



Pair production at a $\mu^+\mu^-$ collider

Reach for S_3 at muon collider



Advantages

- No ISR QCD background
- Cleaner signal
- Better reach and precision

Acknowledgements

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Bibliography

- [1] S. Parashar, A. Karan, Avnish, P. Bandyopadhyay, and K. Ghosh, *Phys.Rev.D* 106 (2022) 9, 095040.
- [2] P. Bandyopadhyay, A. Karan, R. Mandal, and S. Parashar, *Eur. Phys. J. C* 82, 916 (2022).