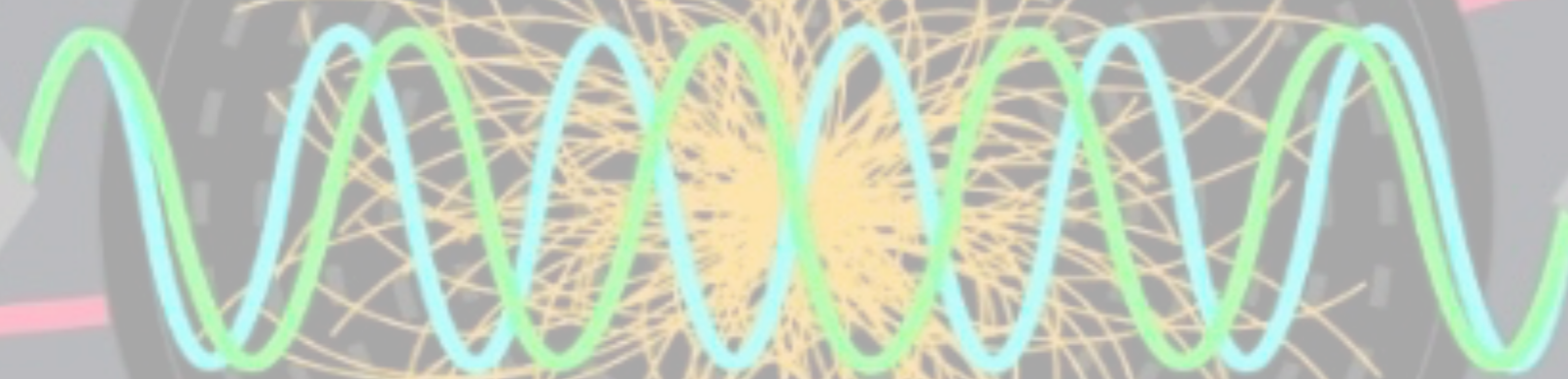


XXV DAE - HEP

$D_q^* \rightarrow D_q \gamma$: Probing the inner structure of D-meson

(Based on arXiv: 2301:XXX)

In collaboration with Prof. Namit Mahajan



Anshika Bansal

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13/12/2022

SYMPOSIUM 2022

Outline

- ❖ Motivation
- ❖ Introduction
- ❖ Light Cone Sum Rules in a Nutshell
- ❖ $D_q^* D_q \gamma$ coupling in LCSR
- ❖ Results
- ❖ Summary and discussion

Motivation

- Distribution amplitudes (DAs) are very crucial universal non-perturbative input for theoretical computations.
- DAs for heavy meson case are modelled using the heavy quark expansion. No precise form is known so far.
[Grozin and Neubert, PRD 55 (1997) 272-290]
- Exclusive decay of B-meson indicates that the first inverse moment of these DAs is a very important parameter.

$$\lambda_M^{-1}(\mu) = \int_0^\infty \frac{d\omega}{\omega} \phi_+^M(\omega, \mu)$$

- For B-meson case, this parameter is constrained experimentally using $B^- \rightarrow \ell^- \nu_\ell \gamma$ decays (provides only the lower limit).
- Using $B^- \rightarrow \ell^- \nu_\ell \gamma$ and QCD sum rules, the value ranges between (0.45 ± 0.15) GeV. However, non-leptonic decays demands $\lambda_B \approx 0.2$ GeV.
[Lee and Neubert, PRD 72, 094028(2005), Baneke et. al, Eur.Phys.J.C 71 (2011) 1818]
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Our Objective

To estimate the inverse moment of the D-meson distribution amplitude using the experimental input on the branching ratio of radiative D_q^* decays.

$D_q^* \rightarrow D_q \gamma$ Decays: An Introduction

- The amplitude for $D_q^* \rightarrow D_q \gamma$ ($q = u, d, s$) is:

Coupling

$$\mathcal{M}(D_q^* \rightarrow D_q(p)\gamma(k)) = e g_{D_q} \epsilon_{\mu\nu\rho\sigma} k^\rho \epsilon_\gamma^\sigma v^\nu \epsilon_{D_q^*}^\mu$$

$\frac{e g_{D_q}}{2}$ is the transition magnetic moment.

- The decay width:

$$\Gamma(D_q^*(p') \rightarrow D_q(p)\gamma(k)) = \frac{\alpha_{em}}{3} |g_{D_q}|^2 |\vec{k}|^3$$

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Experimental Data

Channel	Branching Ratio	Decay widths	g_{D_q}
$D^{*+} \rightarrow D^+ \gamma$ ($q = d$)	$(1.6 \pm 0.4) \%$	$(83.4 \pm 1.8) \text{ KeV}$	0.47
$D^{*0} \rightarrow D^0 \gamma$ ($q = u$)	$(35.3 \pm 0.9) \%$	$< 2.1 \text{ MeV}$	< 10.98
$D_s^{*+} \rightarrow D_s^+ \gamma$ ($q = s$)	$(93.5 \pm 0.7) \%$	$< 1.9 \text{ MeV}$	< 16.27

[PDG]

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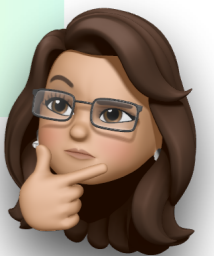
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Answer is **YES!!** We will do this using the method of **Light Cone Sum Rules**.

Light Cone Sum Rules in a Nutshell

Basic Idea

To calculate the hadronic objects of interest using the analytic properties of the correlation function involved.

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Matching the two gives estimates for the hadronic objects

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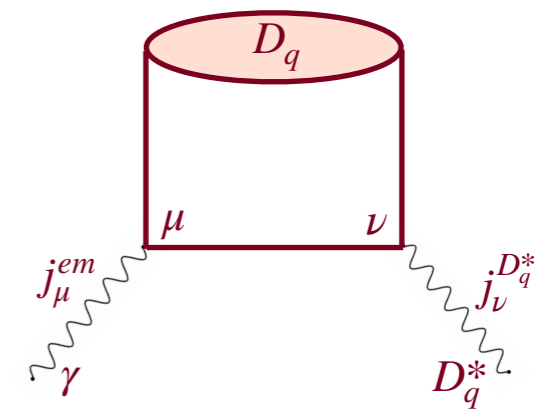
Borel Transformation

(To suppress the effect of continuum and higher resonances to reduce the uncertainty due to duality approximation)

Coupling in LCSR

- The correlation function involved:

$$T_{\mu\nu} = -ie \int d^4x e^{ik \cdot x} \langle D_q(p) | T \left\{ j_\mu^{em}(x) j_\nu^{D_q^*}(0) \right\} | 0 \rangle$$

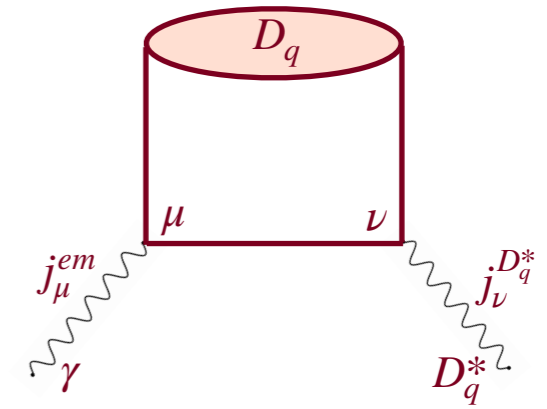


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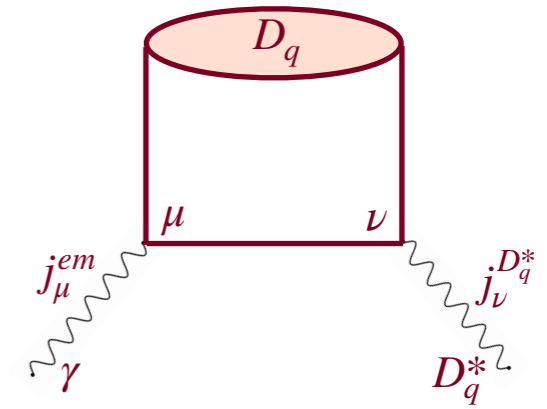
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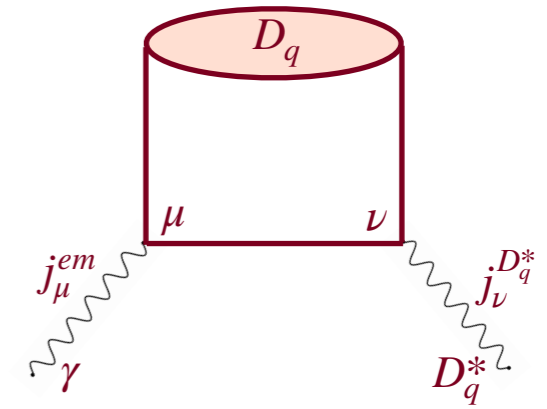
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- Bi-quark operator between vacuum & D-state written in terms of D-meson DAs as:

$$\langle D(p) | \bar{c}_\alpha(0) [0, x] q_\beta(x) | 0 \rangle = \frac{if_D m_D}{4} \int_0^\infty d\omega e^{i\omega v \cdot x} \left[(1 + v^\mu \gamma_\mu) \left\{ \phi_+^D(\omega) - \frac{\phi_+^D(\omega) - \phi_-^D(\omega)}{2v \cdot x} x_\mu \gamma^\mu \right\} \gamma_5 \right]_{\beta\alpha}$$

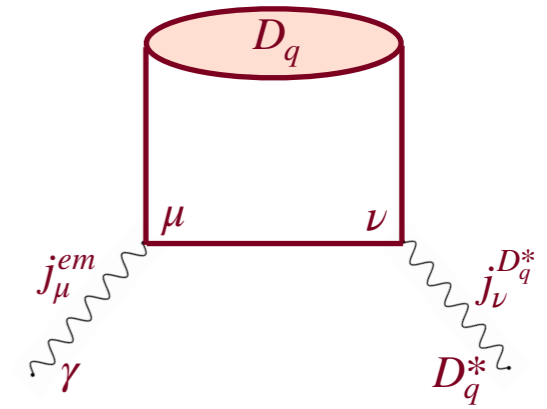
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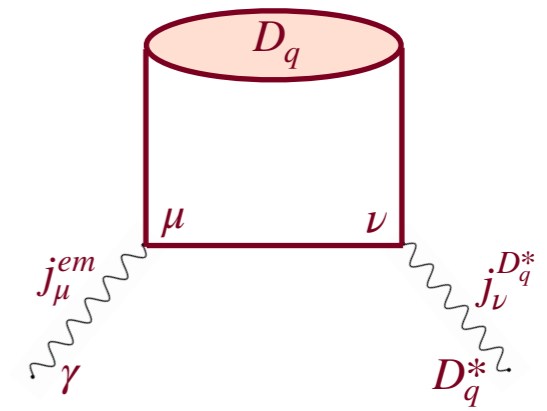
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Momentum of the D-meson carried by the light quark

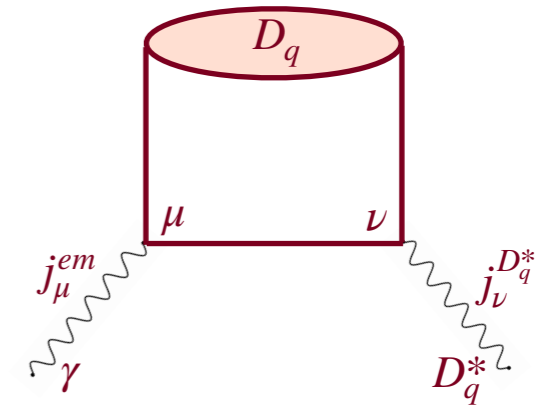
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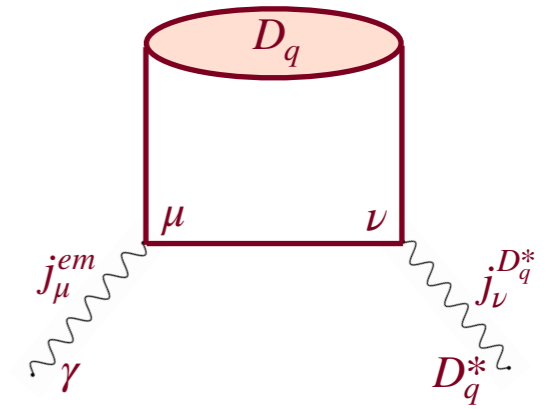
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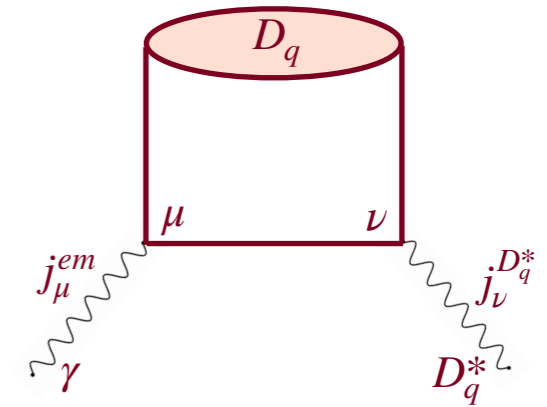
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Exponential Model

[Grozin and Neubert, PRD 55 (1997) 272-290]

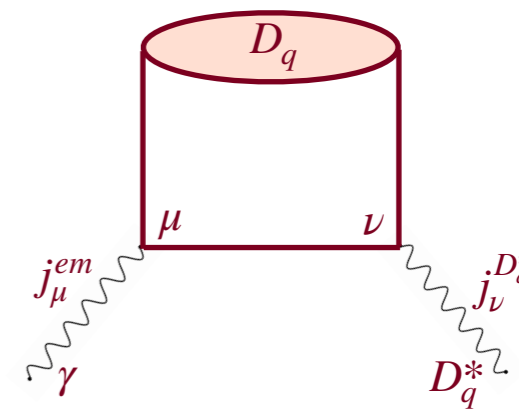
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$\lambda_D = \omega_0$ in the exponential model \implies our objective is to find out ω_0

Dispersion Relation

- Using unitarity:

$$T_{\mu\nu}^{had}(p, k) = ie \langle D_q(p) | T \left\{ j_{\mu}^{em}(x) j_{\nu}^{D_q^*}(0) \right\} | 0 \rangle \sim \langle D_q(p) | j_{\mu}^{em} | D_q^*(p+k) \rangle \langle D_q^*(p+k) | j_{\nu}^{D_q^*}(0) | 0 \rangle + \sum_n \langle D_q(p) | j_{\mu}^{em} | n \rangle \langle n | j_{\nu}^{D_q^*}(0) | 0 \rangle$$

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Contribution from
higher resonance
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- Finally performing the Borel Transformation:

$$\mathcal{B}_{M^2} \left(\frac{1}{(m^2 - q^2)^k} \right) = \frac{1}{(k-1)!} \frac{\exp(-m^2/M^2)}{M^{2(k-1)}}$$

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$$G_{D_q^* D_q}(-k^2) = \frac{1}{f_{D_q^*} m_{D_q^*}} \int_0^{s_0} ds e^{\frac{(m_{D_q^*}^2 - s)}{M^2}} \frac{1}{\pi} \text{Im} (T^{QCD}(s, Q^2))$$

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To be fixed by demanding that the sum rule should be saturated by the lowest state and the contributions coming from the heavier and continuum states are suppressed.

Results

- The analytic results for g_{D_q} is consistent with the heavy quark and chiral symmetry prediction, according to which $g_{D_q} \sim \frac{Q_c}{m_c}$.

[Amundson et. Al, PLB 296 (1992) 415-419]

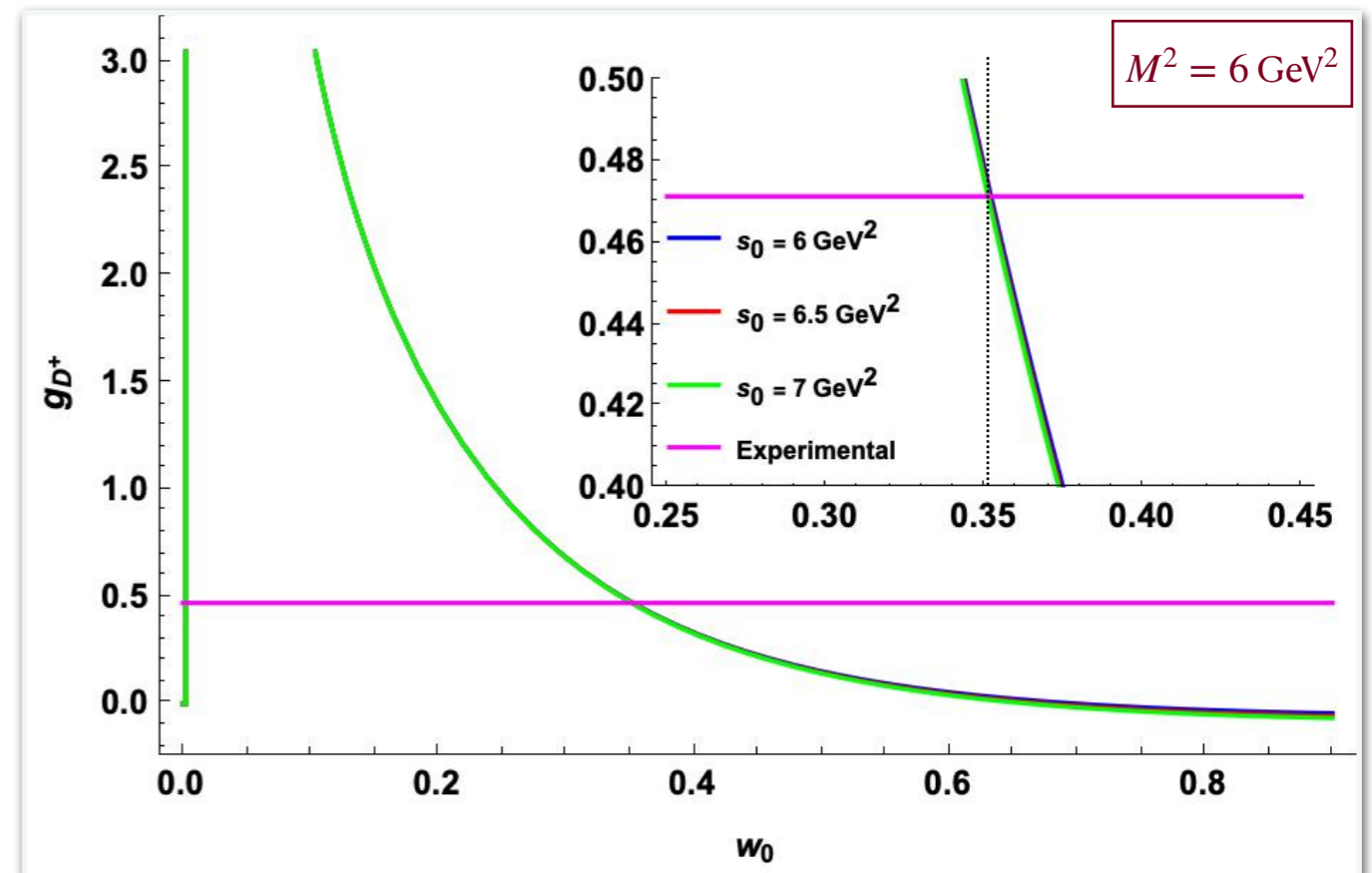
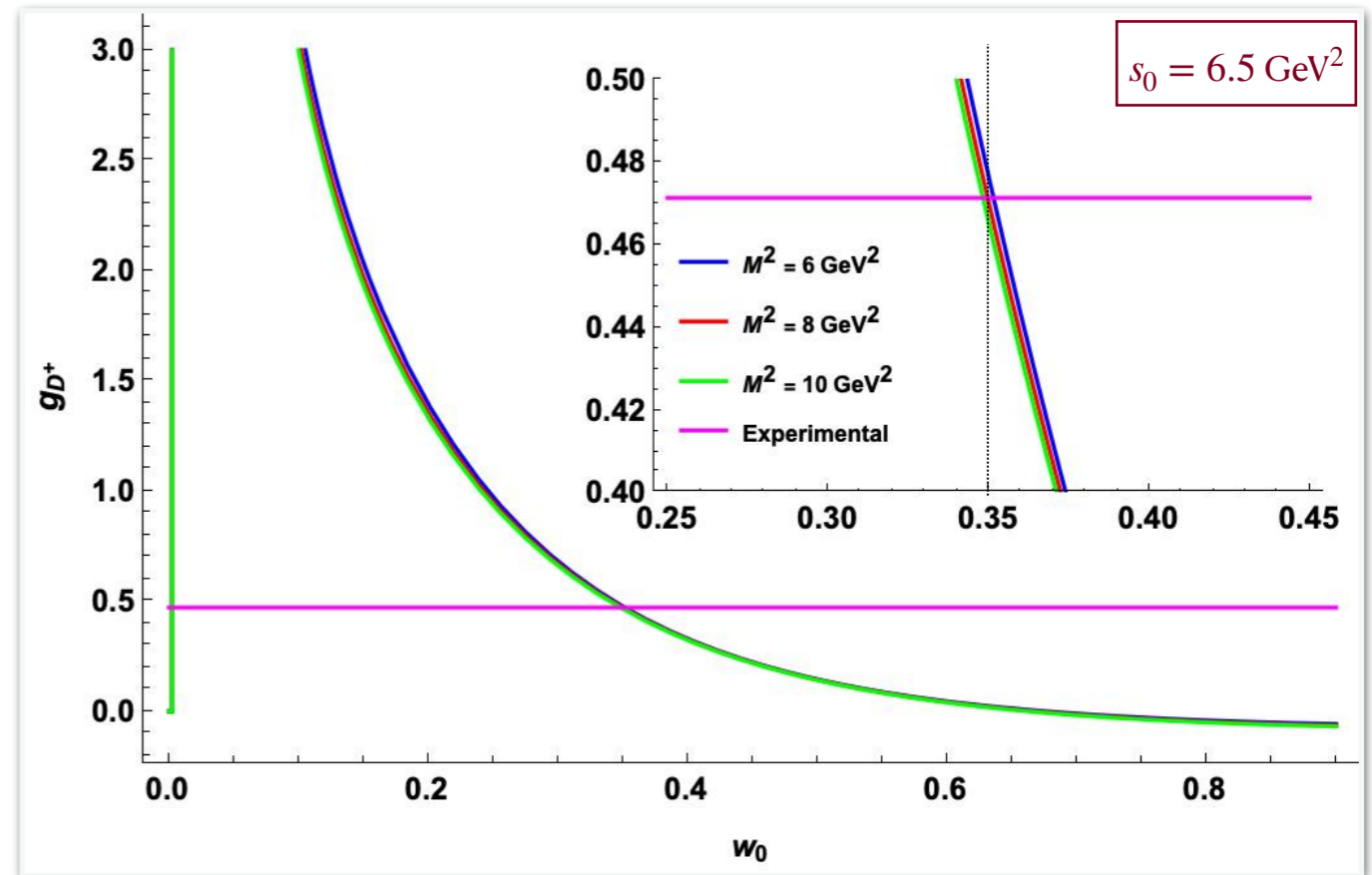
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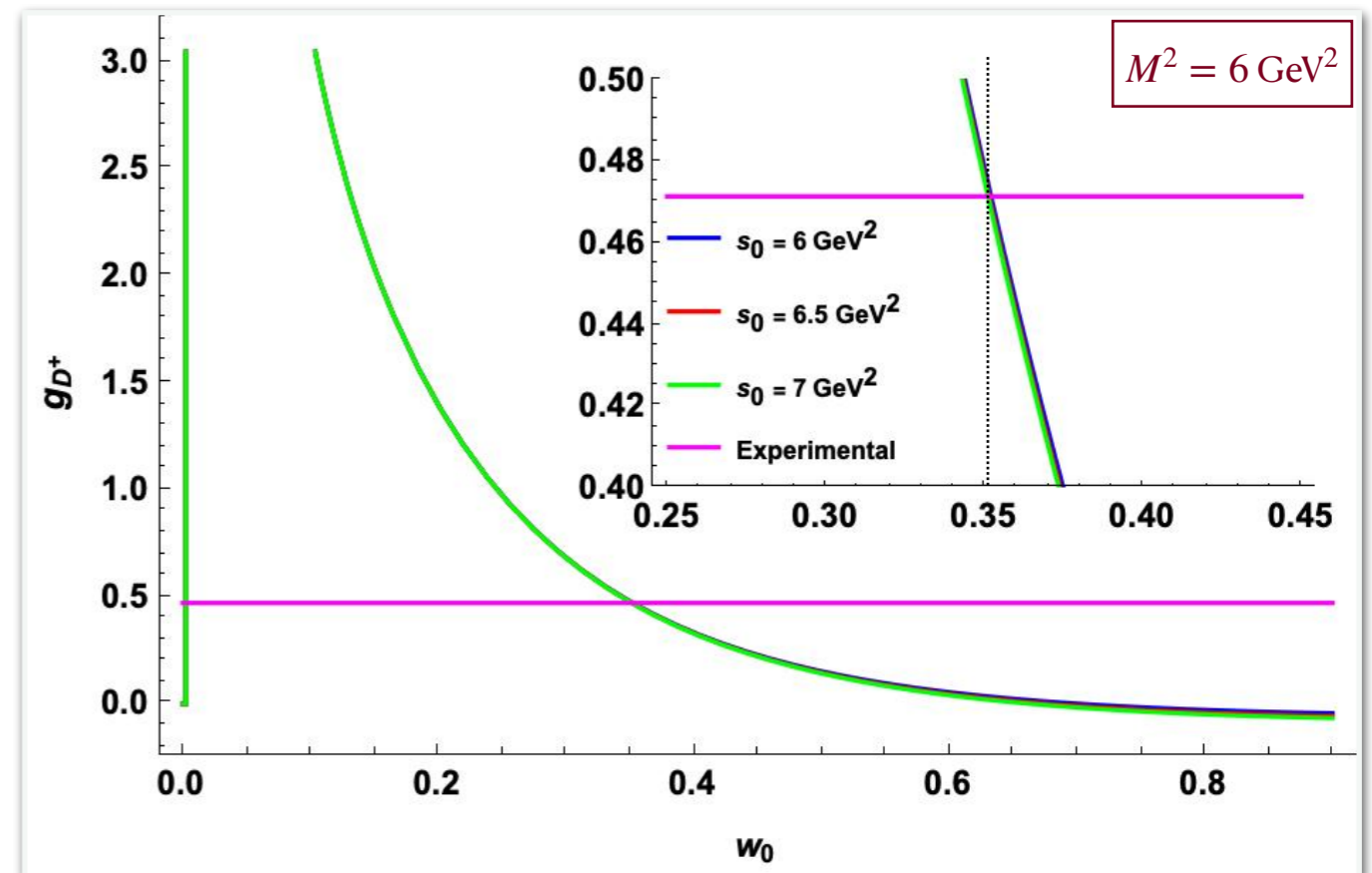
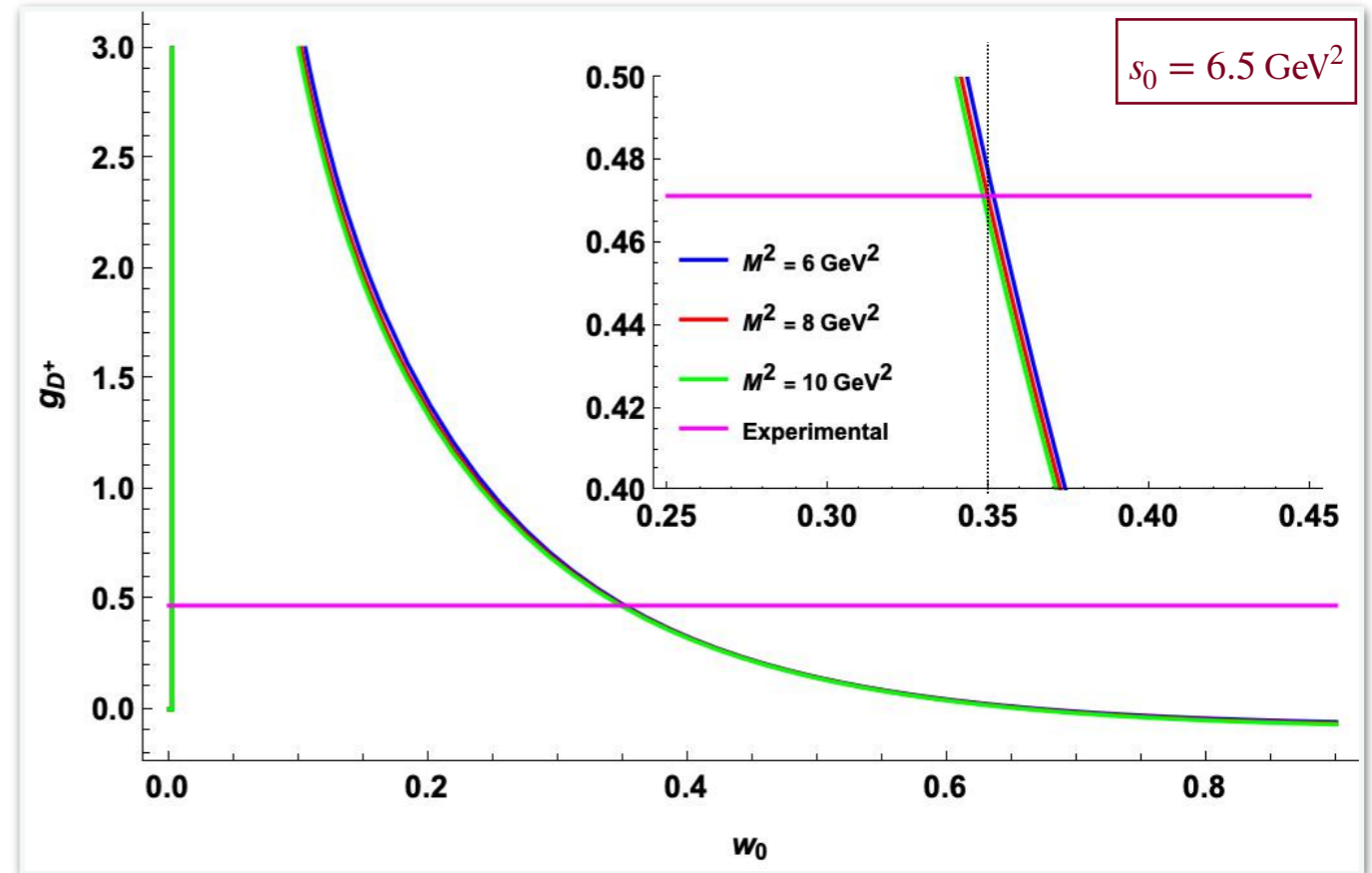
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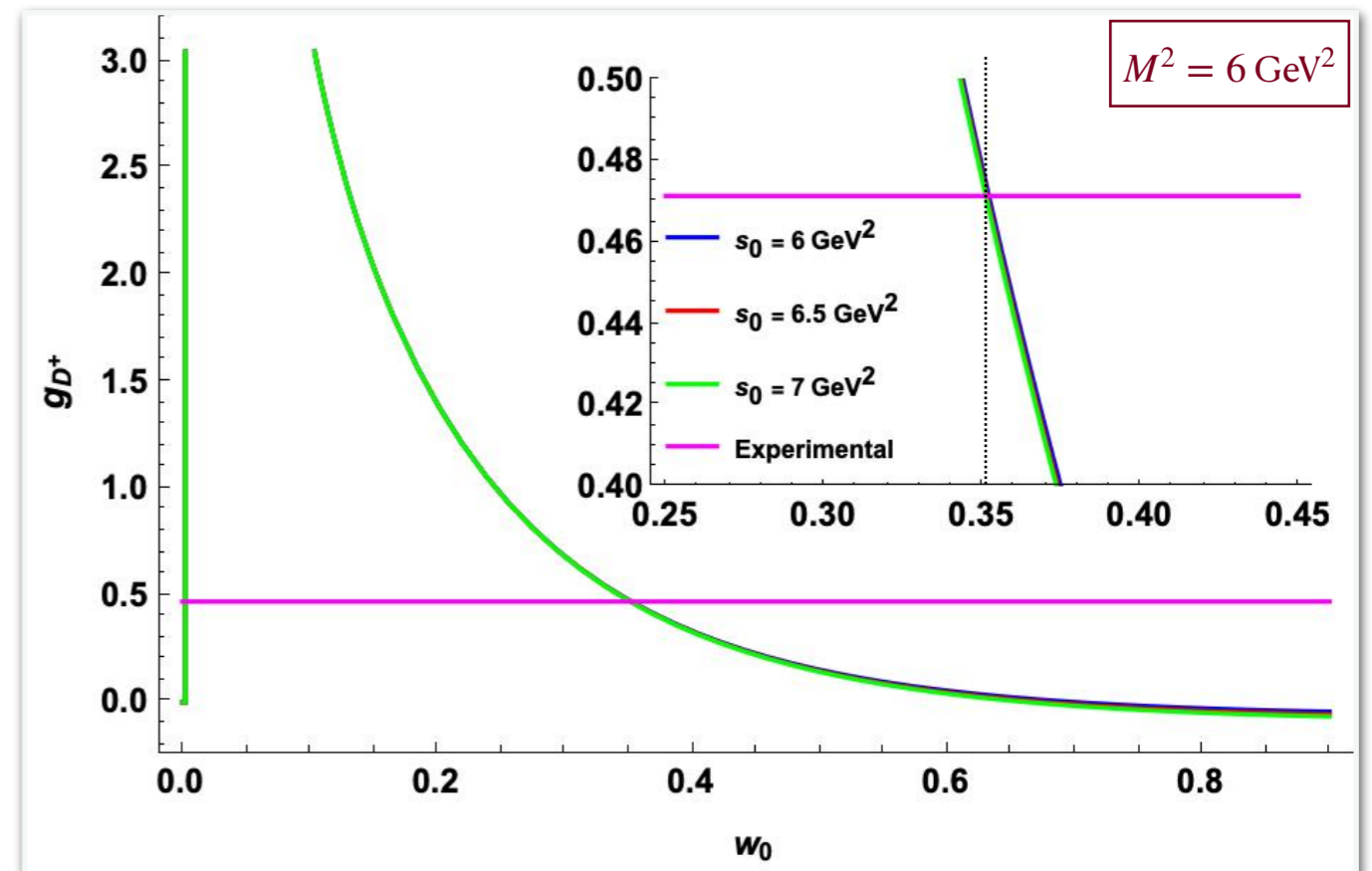
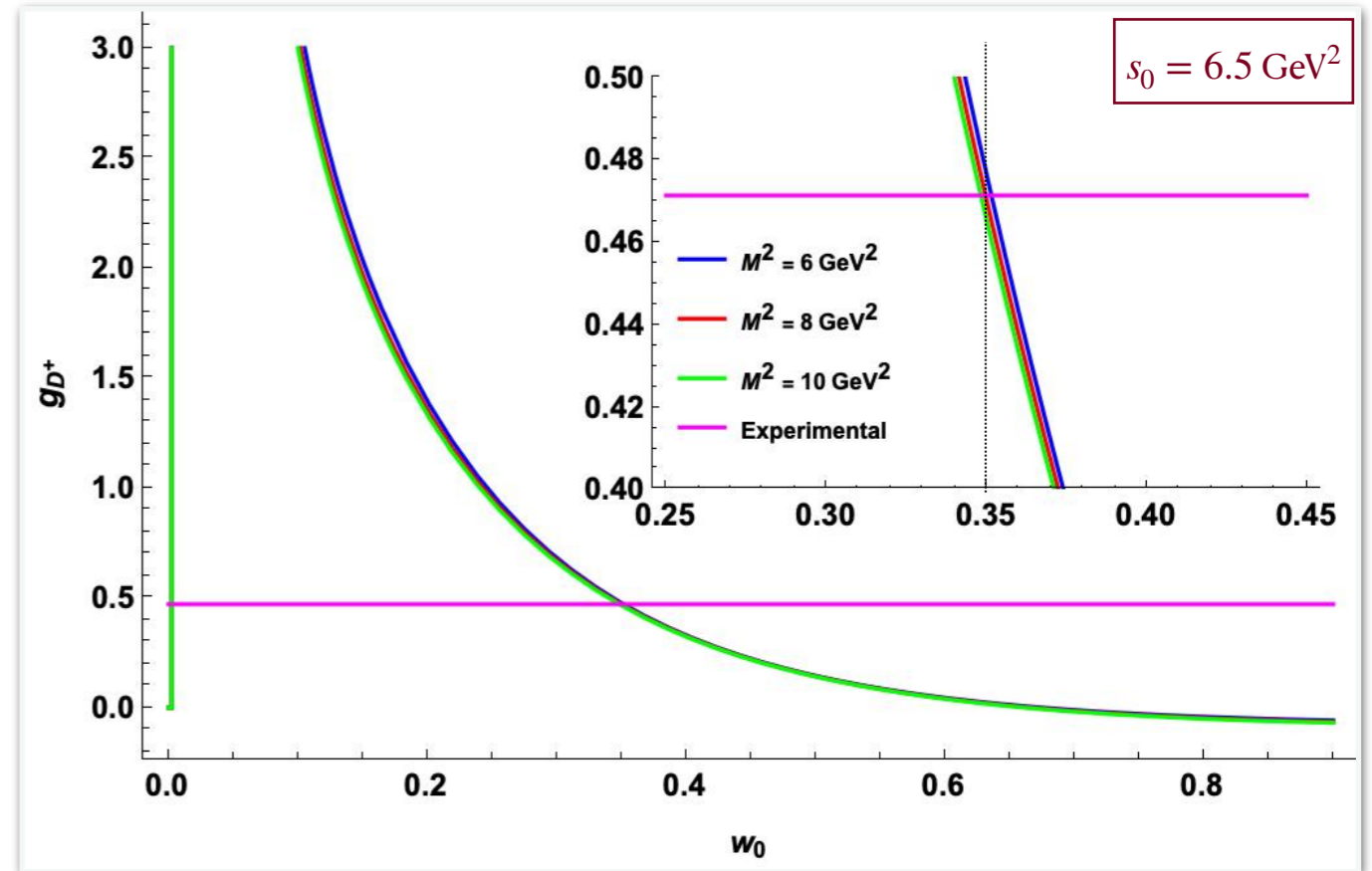
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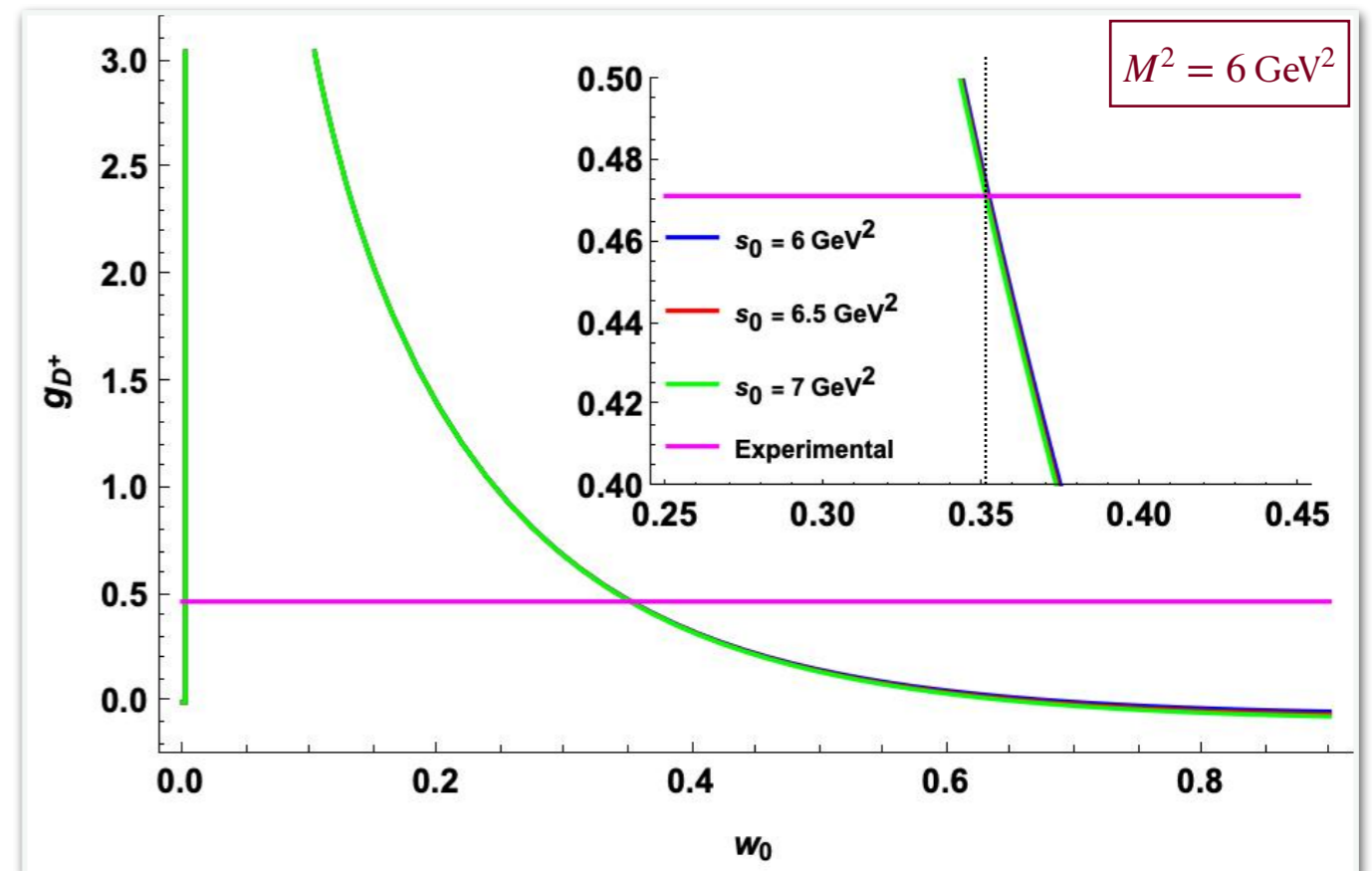
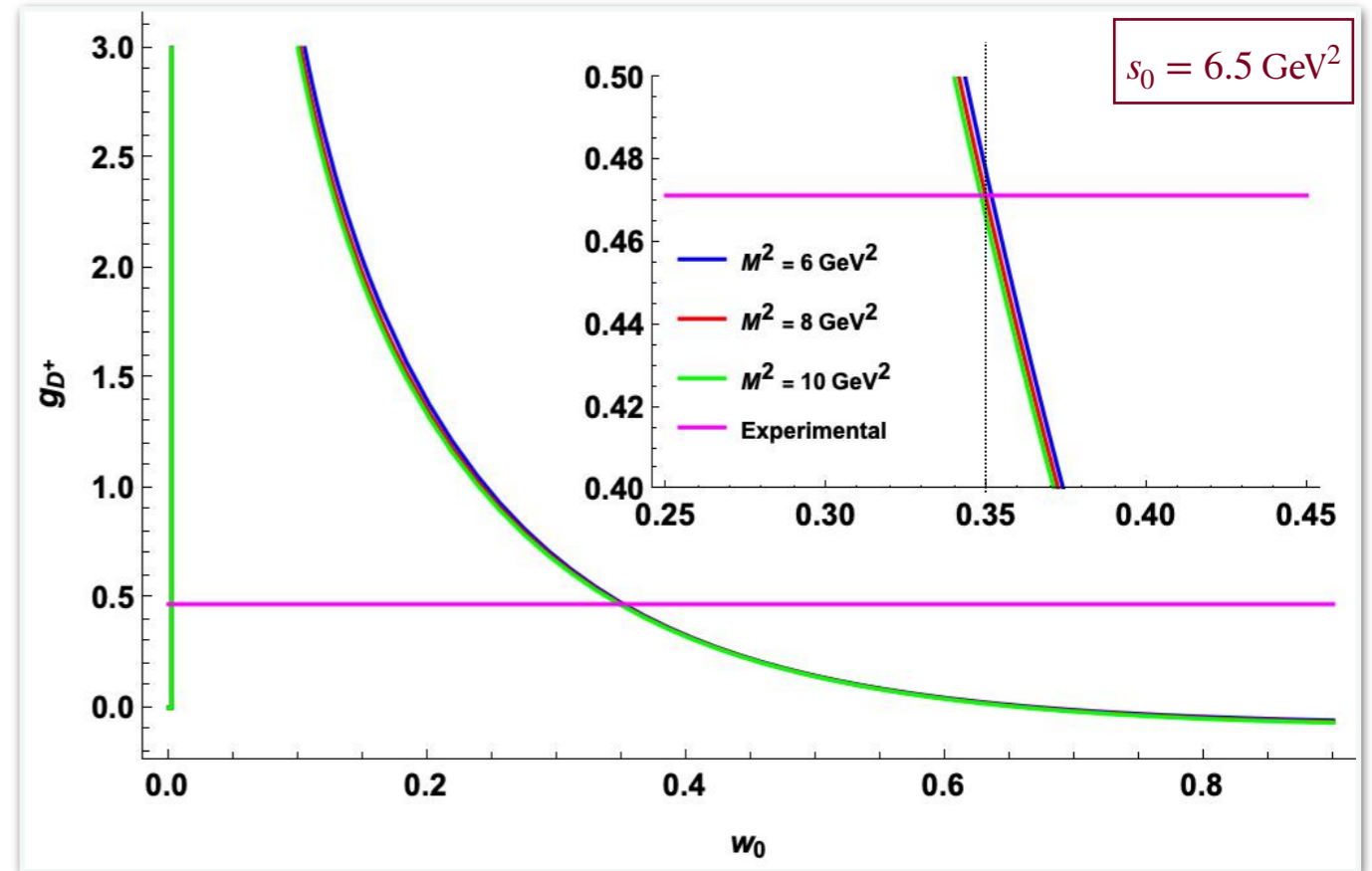
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Thank you !!!

