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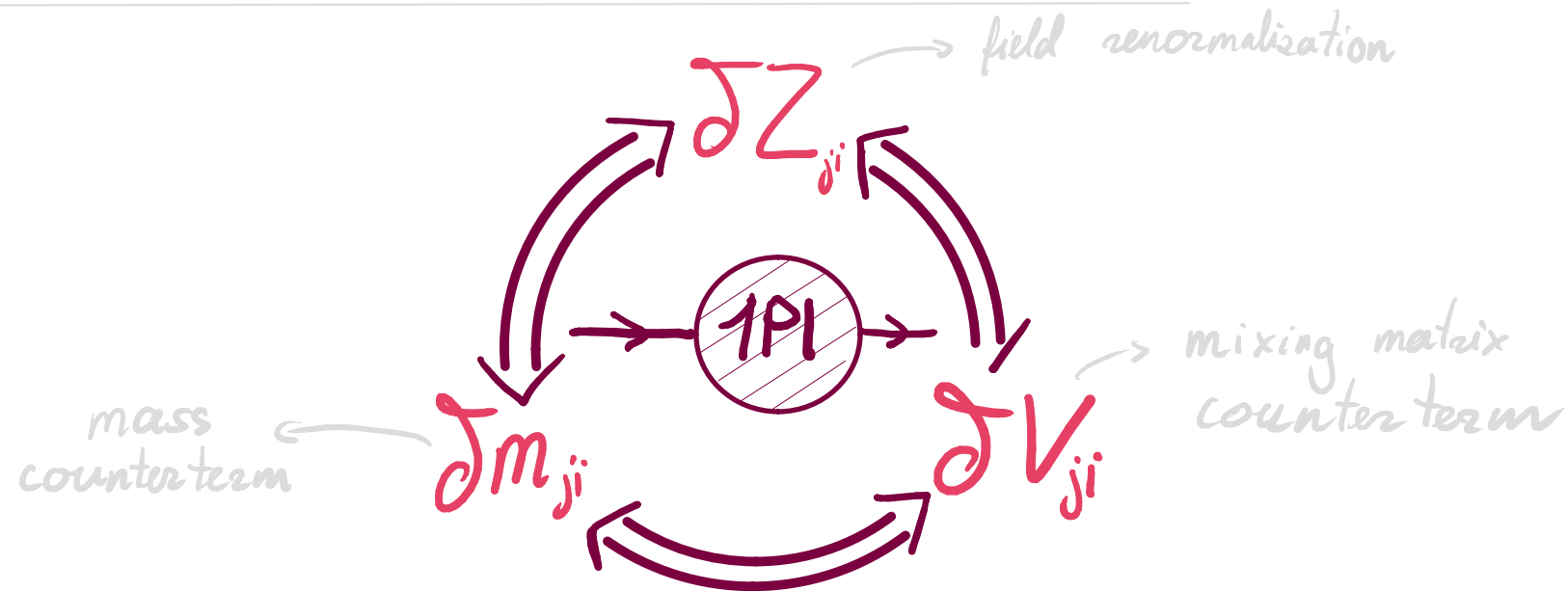
On-Shell Renormalisation of Scalar Sectors

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Contents

- Inspiration : fermions, Nielsen identities... arXiv:2107.03361
- Mixing Renormalization
- Scalars @ 1-loop + problems
- Summary / Outlook

Fermions I



Renormalization and mixing:

$$\psi_i^o \rightarrow Z_{jk} \psi_k$$

flavour indices

$$m_{ji}^o \rightarrow m_i + \delta m_{ji}$$

Fermions II

On-Shell conditions:

spinor $\not{p} u_i = m_i u_i$

$$\Sigma_{ji}^R(\not{p}) u_i = 0$$

self-energy

Includes tadpoles!

$$\xrightarrow{1\text{-loop}} \left[(m_i^2 - m_j^2) \delta Z_{ji} - m_j \delta m_{ji} - \delta m_{ji}^+ m_i \right] u_i = -(\not{p} + m_j) \Sigma_{ji}(\not{p}) u_i$$

$(m_i^2 - m_j^2)$?

- δZ^A and δm are degenerate!
- δZ^H is as usual

$$(\not{p} + m_j) \Sigma_{ji}(\not{p}) u_i \Big|_{uv} \overset{\text{Hermitian}}{\sim} \frac{1}{E_{uv}} \times \{ m_i, m_j, 2m_i m_j, m_i^2 + m_j^2 \}$$

↳ associated with δm !

from 1-loop PV functions

NO $m_i^2 - m_j^2$!

$$\hookrightarrow \dots \Big|_{uv}^A \sim \frac{1}{E_{uv}} \times \{ m_i^2 - m_j^2 \}$$

Fermions III

Nielsen Identity:

$$\partial_\xi \Sigma_{ji}(\phi) = \left[\not{1} \Sigma + \Sigma \not{\Lambda} \right]_{ji}$$

→ All orders

→ Analogous for Scalars

→ $\Lambda^{(-)} = 0$ @ tree-level

→ $\Lambda^{(-)} = \Lambda^{(-)}(\phi)$

⇓ 1-loop

$$\partial_\xi \Sigma_{ji}(\phi) = \left[\not{1}(\phi - m_i) + (\phi - m_j) \not{\Lambda} \right]_{ji}$$

⇓ OS

$$(\not{\phi} + m_j) \partial_\xi \Sigma_{ji}(\phi) u_i = (m_i^2 - m_j^2) \not{\Lambda} u_i$$

Fermions IV

$$\delta Z_{ji}^A U_i \equiv -(\not{p} + m_j) \Sigma_{ji} U_i + \text{h.c.} \quad \left(m_i^2 - m_j^2 \right)$$

$$\delta m = \delta m(\Sigma, \delta Z^A)$$

1-loop!

↳ Explicit computation @ 1-loop for SM and dRDM fermions

- Can be shown that the same "logic" is valid to ALL ORDERS in perturbation theory!

Mixing Renormalization I

- What happens to δV_{ji} ? (mixing matrix counterterm)

Standard approach: $\delta V_{ji} \sim [\delta Z^A V - V \delta Z^A]_{ji}^{uv} \neq 0$

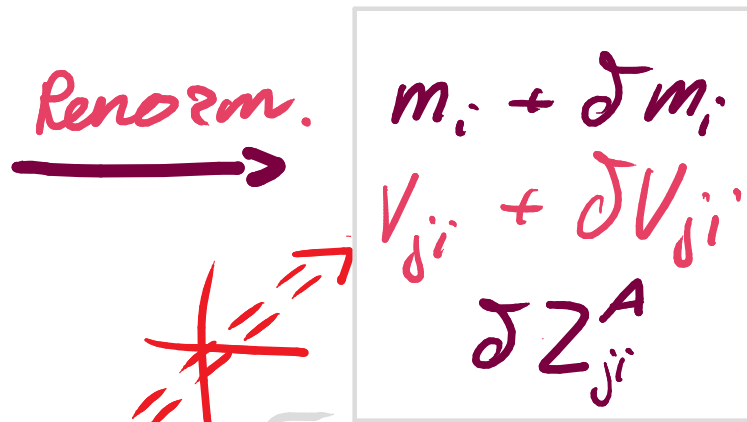
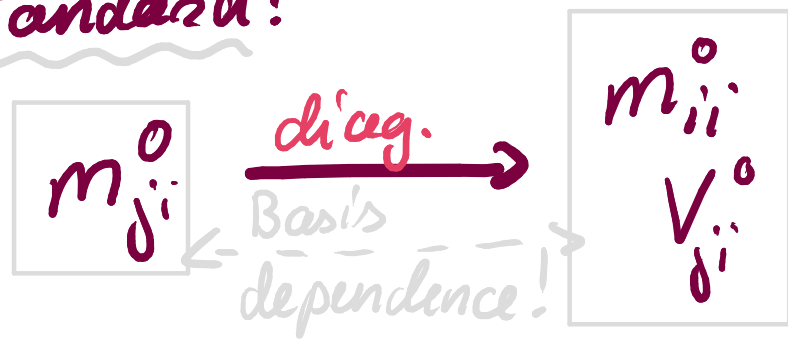
But now $[\delta Z_{ji}^A]^{uv} = 0$ and $[\delta m_{ji}]^{uv} \neq 0$

$$\left[\delta V_{ji} \right]^{uv} = 0$$

- UV divergences stay in the mass term
- Is it needed to renormalize mixing matrices?

Mixing Renormalization II

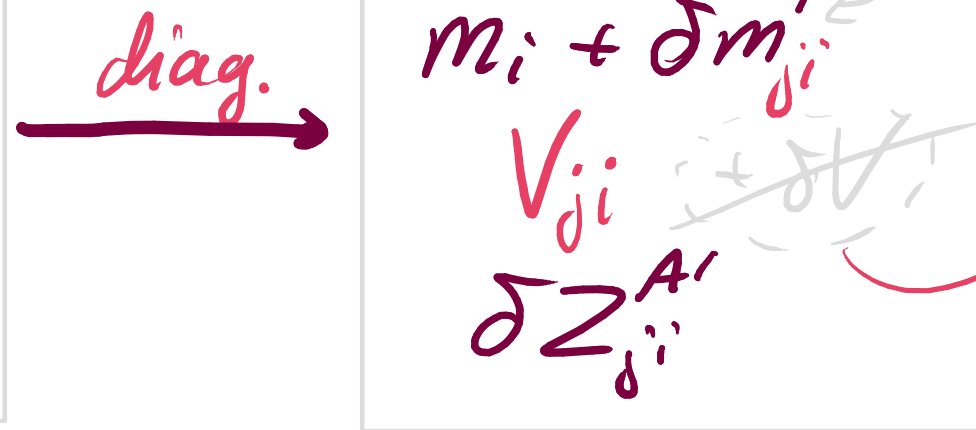
Standard:



Finite rotation
can give:

$$\cancel{V_{ji}} + \delta V_{ji}'$$

Proposed:



NO
counterterm!

- Renorm. and finite rotations commute only if:
 $\delta V_{ji} = 0$ and δm_{δ^i} !

Can we apply this
to
Scalar Sectors?

Scalars I

Similar to fermions @ 1-loop

$$\Pi_{ji}^R(p^2) = \Pi_{ji}(p^2) + \delta Z_{ji}^+(p^2 - m_i^2) + (p^2 - m_j^2) \delta Z_{ji} - \delta m_{ji}^2$$

↳ Self-energy

$$\partial_\xi \Pi_{ji}(p^2) = (\Lambda \Pi)_{ji} + (\Pi \Lambda^+)_{ji} \xrightarrow{1\text{-loop}} \partial_\xi \Pi_{ji}(p^2) = \Lambda_{ji}(p^2 - m_i^2) + (p^2 - m_j^2) \Lambda_{ji}^+$$

$$\xrightarrow{OS} \partial_\xi \Pi_{ji}(m_i^2) = (m_i^2 - m_j^2) \Lambda_{ji}^+$$

$$\Rightarrow \delta Z_{ji}^A = -\frac{1}{2} (\Pi_{ji}(m_i^2) + \Pi_{ji}(m_j^2)) \Big|_{m_i^2 - m_j^2}$$

$$\delta m = \delta m(\Pi, \delta Z^A)$$

→ Extend to all orders, but...

↳ Explicit computation @ 1-loop in dHDM

Scalars II

(preliminary)

- Physical states MIX with Goldstones and longitudinal modes

- N_1 is for the reduced functional

Reduced functional $\tilde{\Gamma} = \Gamma^R - \int dx \mathcal{L}^{GF}$ ← Gauge-fixing (Faddeev-Popov)

← Always renormalized („classical“)

1-loop $\Pi_{HG}^R = \Pi_{HG} + \delta Z_{HG}^T (p^2 - m_H^2) + p^2 \delta Z_{HG} - \delta m_{HG}^2$

“Higgs” “Goldstone”

← from the “Higgs potential”

propagator is still $\frac{i}{p^2 - m_G^2}$

- Renormalize @ $p^2 = m_H^2$ and $p^2 = 0$

- $m_H^2 \delta Z_{HG}^A - \delta m_{HG}^2 = \frac{1}{2} (\Pi_{HG}(0) + \Pi_{HG}(m_H^2))$

- No $m_H^2 - m_G^2$!

Scalars III

(preliminary)

$$\partial_\xi \Pi_{HG}^R = \Lambda_{Hk}^{\tilde{}} \tilde{\Pi}_{kG}^+ + \tilde{\Pi}_{Hk}^+ \Lambda_{kG}^+ + \partial_\xi \Pi^{GF} @ \text{tree-level}$$

not diagonal!

1-loop

k^n also includes longitudinal modes!

$$\Rightarrow p^\nu \tilde{\Pi}'_{zG}$$

(e.g. Z_L^ν)

$$\partial_\xi \Pi_{HG}^R = \Lambda_{HG} p^2 + p^2 \Lambda'_{H2} \tilde{\Pi}'_{zG} + (p^2 - m_H^2) \Lambda_{HG}^+ + p^2 \tilde{\Pi}_{H2}^+ \Lambda_{zG}^+ \rightarrow 0$$

OS

$$m_H^2 \partial_\xi \delta Z_{HG}^A - \partial_\xi \delta \mathcal{M}_{HG}^2 = -\frac{m_H^2}{2} \left(-\Lambda_{HG}^+(0) + \Lambda_{HG}(m_H^2) + \Lambda'_{H2}(m_H^2) \tilde{\Pi}'_{zG} \right)$$

not very distinct...

+ terms from longitudinal modes

Scalars IV

(preliminary)

the real **PROBLEM!**

Higher orders depend on all previous orders!

exclude tree-level

(order)

$$\partial_{\xi} \left[(m_i^2 - m_j^2) \delta Z_{ji}^{(n)} - \delta m_{ji}^{(n)} \right] = -\partial_{\xi} \left[\Gamma(m_i^2) + \Gamma(m_i^2) \delta Z \right]_{ji}^{(n)}$$

physical indices

short-hand notation; includes c.t.'s up to order n

• includes non-physical parts

• $\delta Z_{ss'}$ only for scalars,

but $\partial_{\xi} \Gamma$ also includes

longitudinal modes!

↳ STI? Other nice correlations?

Can $(m_i^2 - m_j^2)$ and $\partial_{\xi} \delta m_{ji}^2 = 0$ be preserved to all orders?

Summary / Outlook

- 1-loop physical states OK
 - ▷ δZ_{ji}^A as coeff. of $m_i^2 - m_j^2$; also UV finite
 - ▷ Solve for δm_{ji}
 - ▷ $\delta V_{ji} = 0$, $\partial_\xi \delta m_{ji} = 0$
 - ▷ Universal, process-independent...
- 1-loop non-physical: no $m_i^2 - m_j^2 \Rightarrow$ How to select δZ^A ?
 - ▷ δm^2 and δZ^H ensure OS and $\partial_\xi \delta m^2 \neq 0$?

- ! Not obvious n^{th} order extension
 - ▷ depends on all previous orders + longitudinal modes

OUTLOOK: ALL ORDER EXTENSION