## Gamma-ray production in neutral-current interactions

#### **Artur Ankowski**

University of Wrocław, INFN & Dept. of Physics, "Sapienza" Università di Roma

In collaboration with O. Benhar, T. Mori, R. Yamaguchi, and M. Sakuda

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#### **Outline**

#### **1** Introduction

- Neutral current vs. charged current
- How are the NC events detected?
- Could nuclear deexcitation provide an additional signature?

#### 2 Case of NC neutrino-oxygen scattering

- Nuclear structure and deexcitation scheme
- Estimate of the cross section

#### **3** Summary

## Introduction

## Neutral current vs. charged current

Structure of the NC and CC is the same,

$$\langle N'(p') | \hat{J}^{\mu} | N(p) \rangle =$$

$$\langle N'(p') | \gamma^{\mu} (F_1 + F_2) - \frac{(p+p')^{\mu}}{2M} F_2$$

$$+ \gamma^{\mu} \gamma_5 F_A + \gamma_5 \frac{q^{\mu}}{M} F_P | N(p) \rangle$$

Both *n*'s and *p*'s may undergo NCE scattering

#### Structure of the NC and CC is the same, but

$$F_i^{z,(p)} = -\frac{1}{2} \left( F_i^{(n)} + F_i^s \right) + \left( \frac{1}{2} - 2\sin^2 \theta_W \right) F_i^{(p)},$$

$$F_A^{z,(p)} = \frac{1}{2} \left( F_A + F_A^s \right),$$

$$F_i^{z,(n)} = -\frac{1}{2} \left( F_i^{(p)} + F_i^s \right) + \left( \frac{1}{2} - 2\sin^2 \theta_W \right) F_i^{(n)},$$

$$F_A^{z,(n)} = -\frac{1}{2} \left( F_A - F_A^s \right),$$

W.M. Alberico, S.M. Bilenky, C.Maieron, Phys. Rep. 358, 227 (2002)

#### Structure of the NC and CC is the same, but

$$\begin{split} F_i^{z,(p)} &= -\frac{1}{2} \left( F_i^{(n)} + \overline{F}_i^s \right) + \left( \frac{1}{2} - 2 \sin^2 \theta_W \right) F_i^{(p)}, \\ F_A^{z,(p)} &= \frac{1}{2} \left( F_A + F_A^s \right), \\ F_i^{z,(n)} &= -\frac{1}{2} \left( F_i^{(p)} + \overline{F}_i^s \right) + \left( \frac{1}{2} - 2 \sin^2 \theta_W \right) F_i^{(n)}, \\ F_A^{z,(n)} &= -\frac{1}{2} \left( F_A - F_A^s \right), \end{split}$$

 $F_i^S$  are vanishing, see e.g.

D. Androić *et al.*, PRL 104, 012001 (2010); Z. Ahmed *et al.*, arXiv:1107.0913

#### Structure of the NC and CC is the same, but

$$F_i^{z,(p)} = -\frac{1}{2}F_i^{(n)} + \left(\frac{1}{2} - 2\sin^2\theta_W\right)F_i^{(p)},$$

$$F_A^{z,(p)} = \frac{1}{2}\left(F_A + F_A^s\right),$$

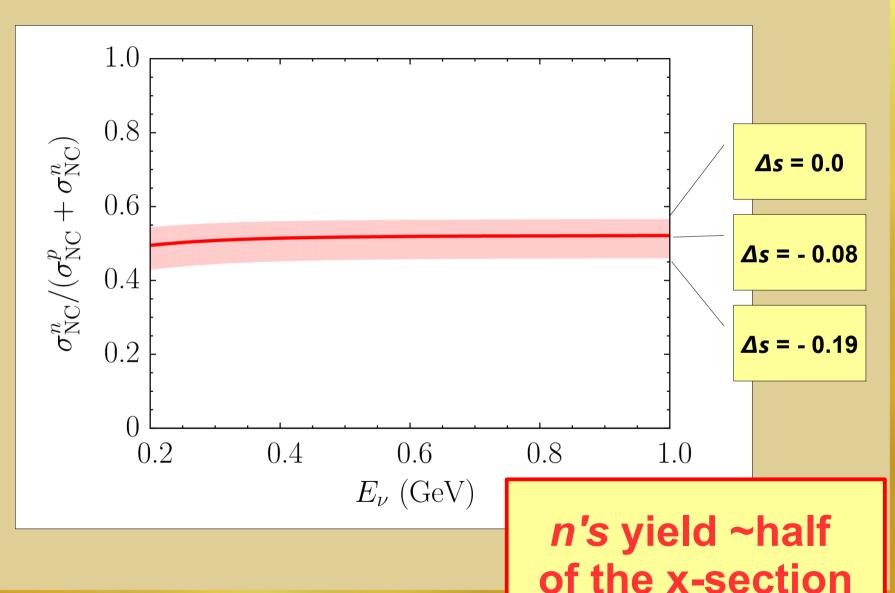
$$F_i^{z,(n)} = -\frac{1}{2}F_i^{(p)} + \left(\frac{1}{2} - 2\sin^2\theta_W\right)F_i^{(n)},$$

$$F_A^{z,(n)} = -\frac{1}{2}\left(F_A - F_A^s\right),$$

$$F_A^s = \frac{\Delta s}{(1 + \frac{Q^2}{M_A^2})^2}$$

$$\Delta s = -0.08 \pm 0.26$$
, MiniBooNE  $\Delta s = -0.12 \pm 0.09$ , BNL E734

## Neutron contribution to the NC <sup>16</sup>O cross section



• The total NC x-section of a symmetric nucleus is **largely independent of**  $\Delta s$ :

When  $\Delta s = 0.0$  is used instead of  $\Delta s = -0.08$ , the cross section changes by **less than 0.3%** for 0.2 < E < 5 GeV.

• In the following  $\Delta s = 0.0$ .

#### **Detection**

- Neutrons (~50% of NC events) do not emit
   Cherenkov light.
- In water Cherenkov detectors, the threshold momentum for observation of proton is 1.07 GeV/c.
- Hence, an additional signature for NC event might be very useful.

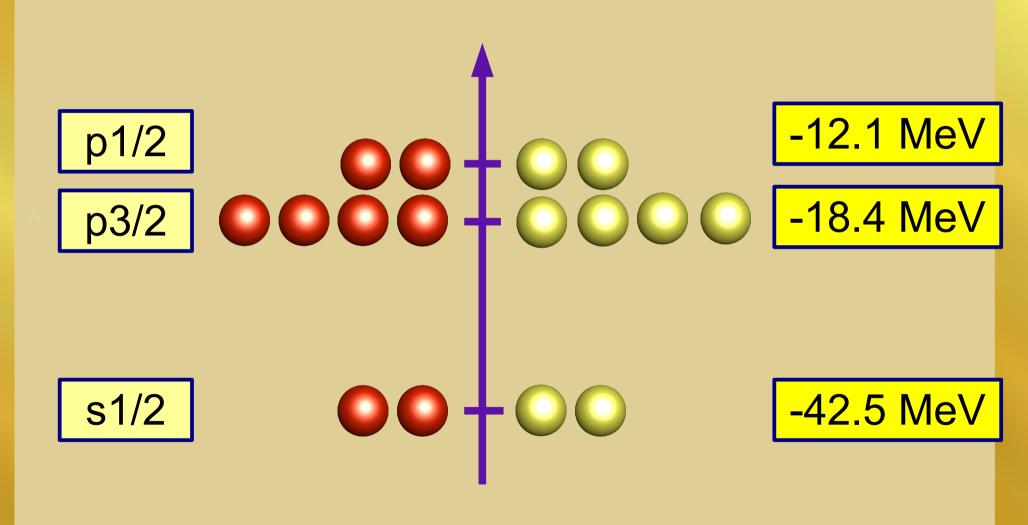
# Nuclear deexcitation as a prompt signal of NC event

 Nucleon knockout often leaves residual nucleus in an excited state.

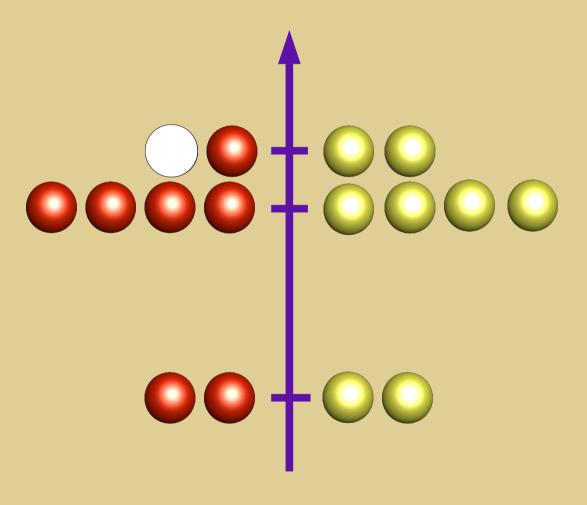
- Deexcitation may yield photons above detection threshold.
- Such photons could provide useful signal, especially for water Cherenkov detectors.

# NC neutrino-oxygen scattering

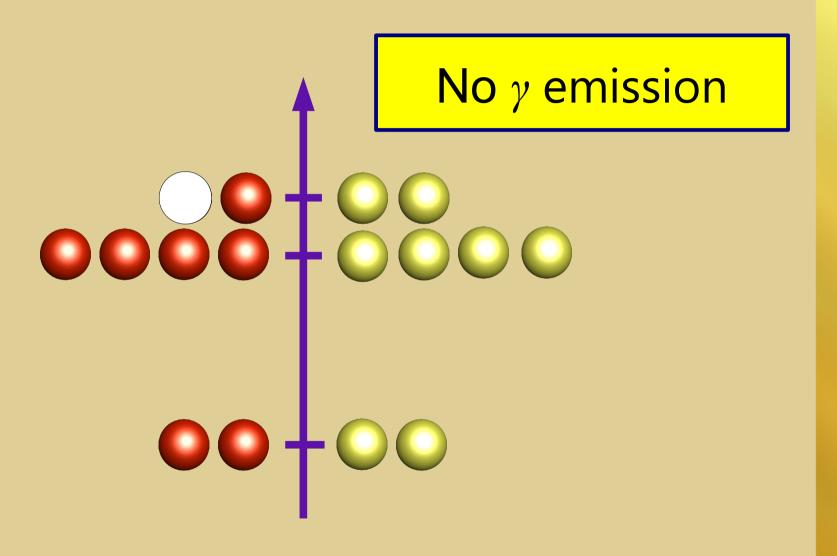
## Structure of the oxygen nucleus



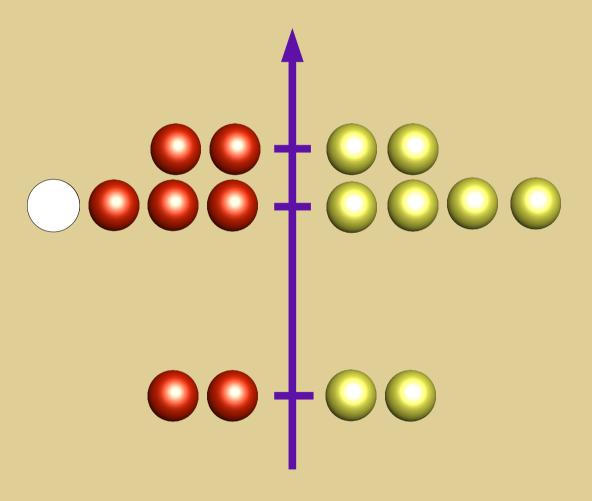
### p1/2 knockout



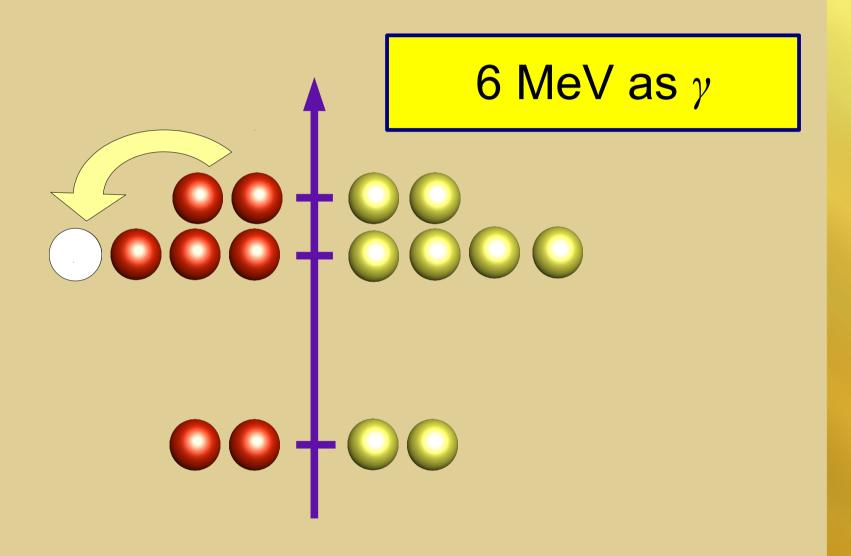
### p1/2 knockout



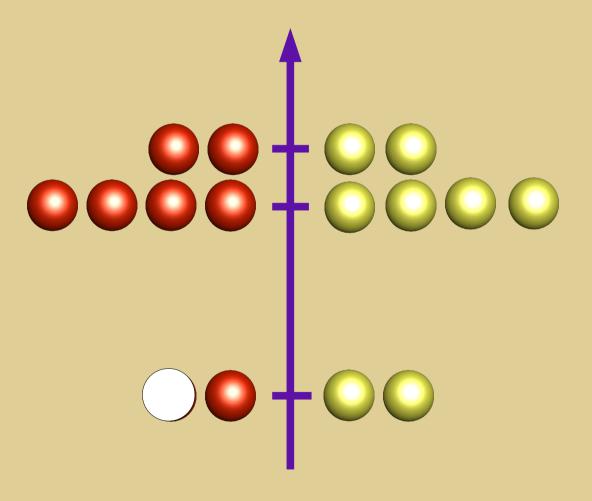
### p3/2 knockout



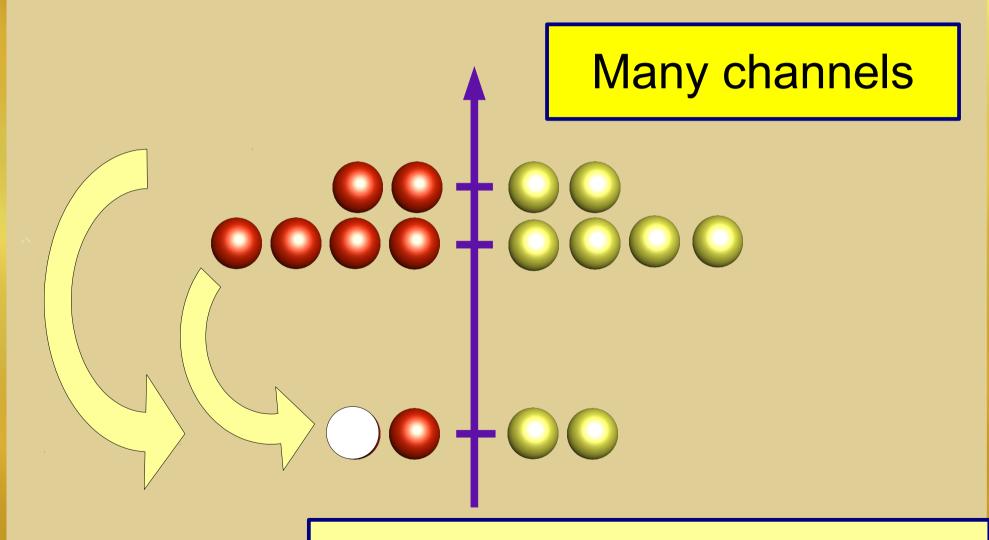
## p3/2 knockout



#### s1/2 knockout

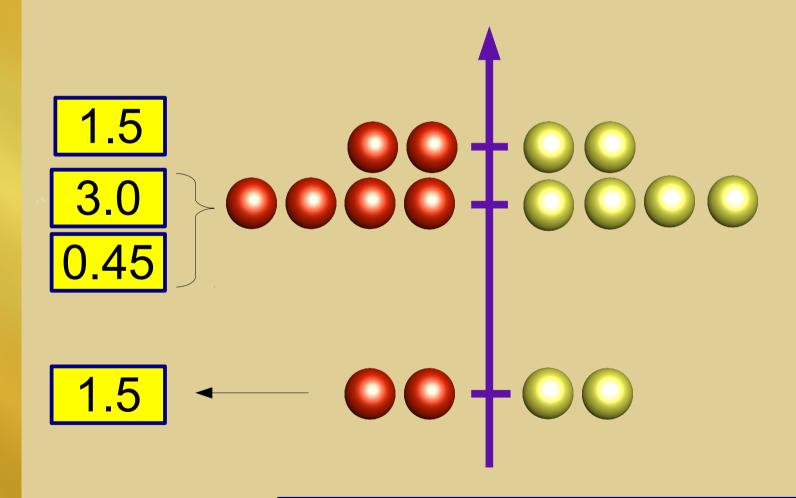


#### s1/2 knockout



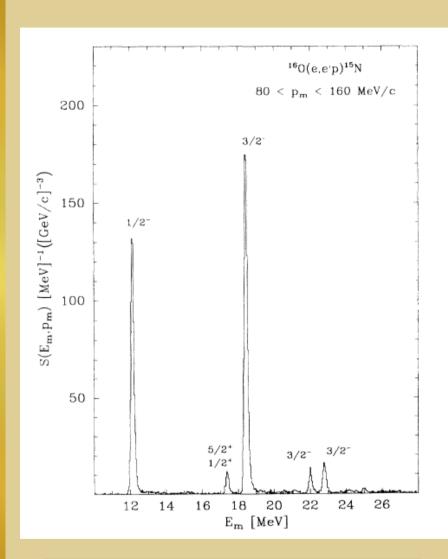
Branching ratio for  $E_{\gamma} > 4$  MeV is ~12%, Ejiri, PRC 48, 1442 (1993)

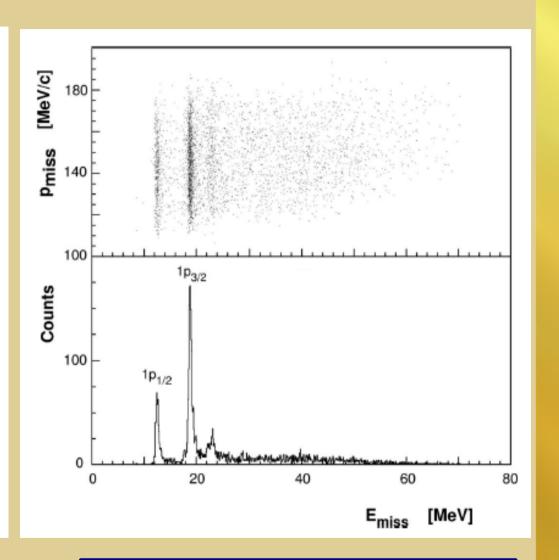
# Spectroscopic strengths from the (e,e'p) Saclay data



M. Bernheim et al., NP A375, 381 (1982)

#### Finite widths





M. Leuschner *et al.*, PRC 49, 955 (1994)

K.G. Fissum *et al.*, PRC 70, 034606 (2004)

## Cross section for $\gamma$ emission accompanying NC event

It is a product of

- NC cross section for the knockout from each shell
- branching ratios for deexcitation by γ emission

$$B(p_{\frac{1}{2}}) = 0\%, \quad B(p_{\frac{3}{2}}) = 100\%, \quad B(s_{\frac{1}{2}}) \approx 15\% \text{ for } E_{\gamma} > 6 \text{ MeV}$$

- F. Ajzenberg-Selove, NP A523, 1 (1991);
- H. Ejiri, PRC 48, 1442 (1993);
- K. Kobayashi et al., arXiv:nucl-ex/0604006

## Cross section for $\gamma$ emission accompanying NC event

- The cross sections for each shell and for the whole nucleus are calculated using the spectral function obtained by Omar Benhar.
- Accuracy of the spectral function, accounting for nucleon-nucleon correlations, has been shown against (e,e'p) data.

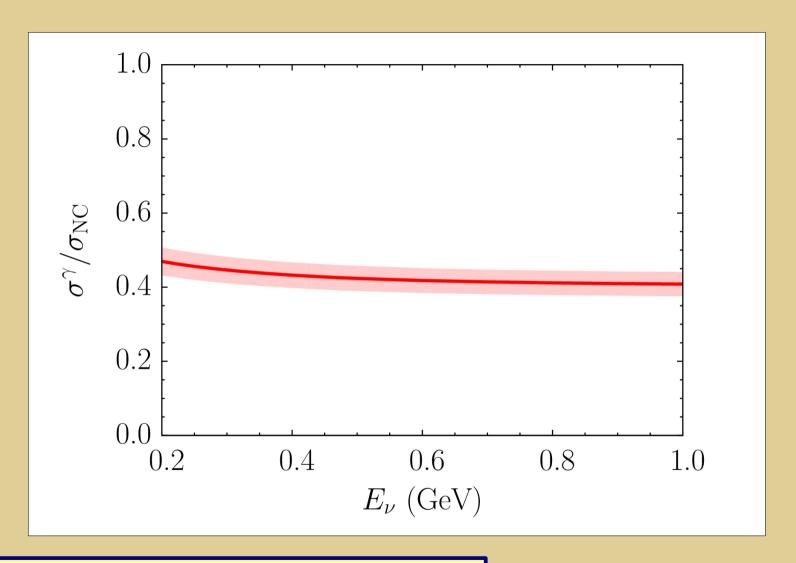
# Cross section for $\gamma$ emission accompanying NC event

Our estimate is that

• ~7-MeV photons (from the s1/2 knockout) follow **≤3.5%** NC events

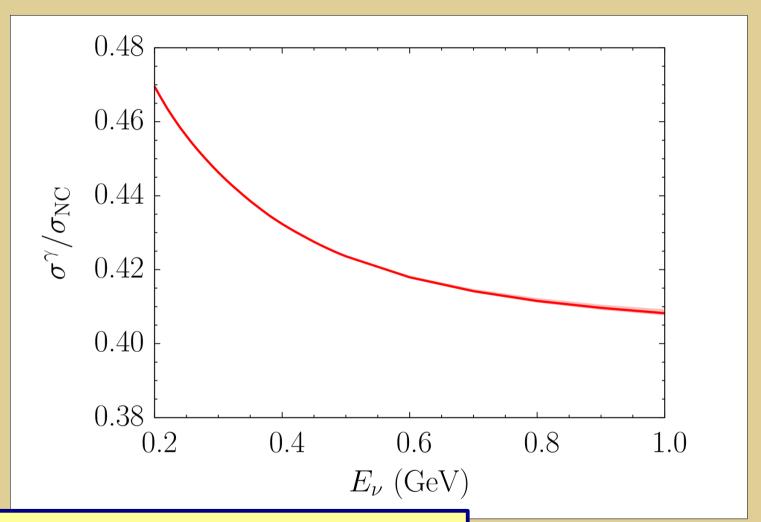
• 6-MeV photons (from the p3/2 knockout) follow **~40%** NC events

#### 6-MeV γ production vs. NC



Ratio weakly dependent on E

#### 6-MeV γ production vs. NC



Ratio largely independent of  $M_A$ . Band for  $M_A$  1.03 – 1.39 GeV

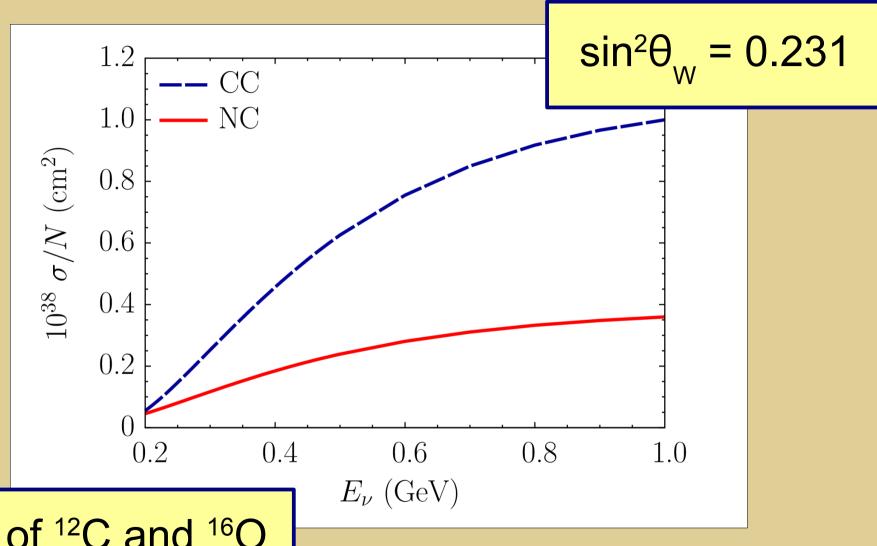
#### What features of oxygen are essential?

- 1 Large contribution to the cross section from a narrow level deexciting into photons
- **2** Deexcitation with  $E\gamma > 5$  MeV providing clear signal
- **3** Available spectral function

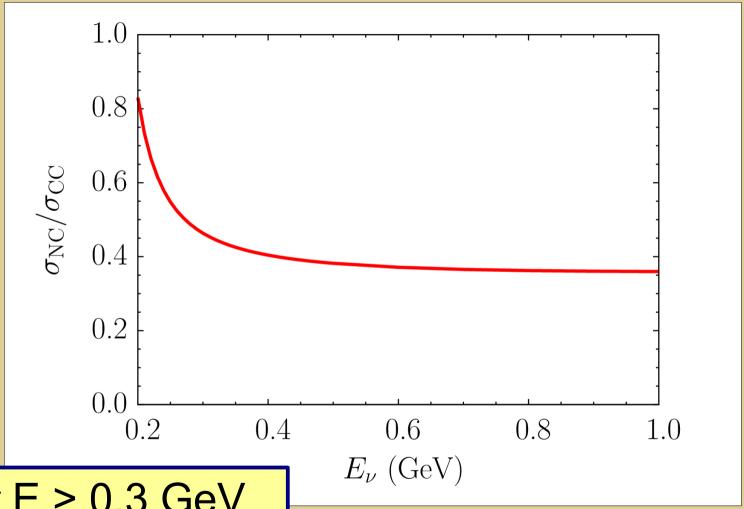
#### **Summary**

- **1** Deexcitation into photons of  $E_{\gamma} > 5$  MeV following the NC interaction may provide a useful signature for water Cherenkov detectors
- **2** The ratio  $\sigma(\gamma)/\sigma(NC)$  is largely independent of the axial mass value
- 3 In the important for T2K region  $E \lesssim 1$  GeV, the ratio is ~40% for the 6-MeV  $\gamma$ 's from p3/2 knockout. The  $E\gamma > 6$  MeV photons from s1/2 hole are just 2-3%.

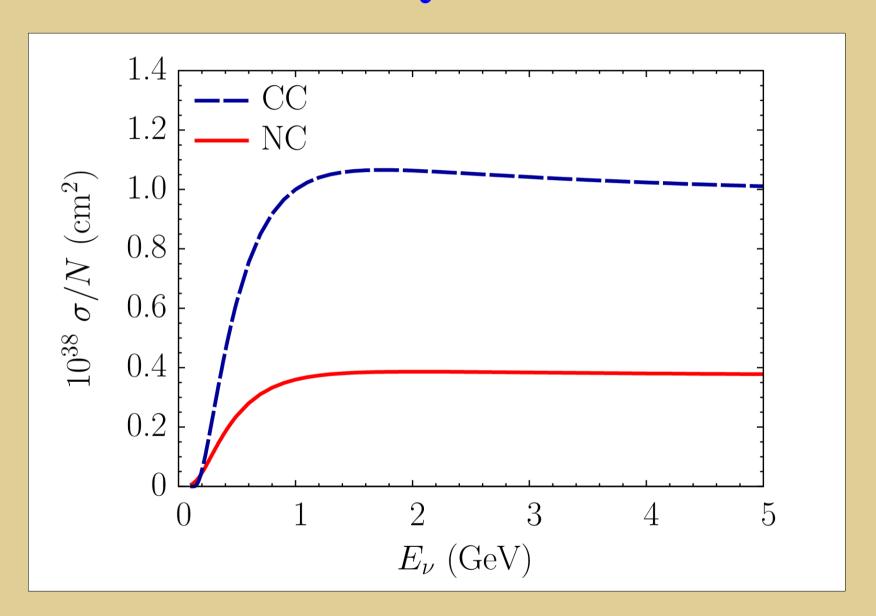
Back-up slides

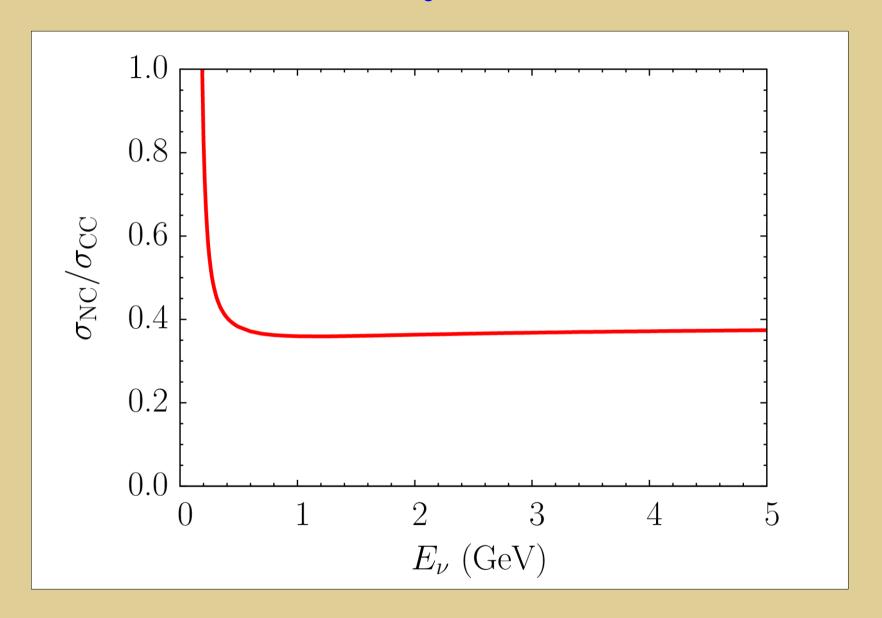


σ/N of <sup>12</sup>C and <sup>16</sup>O agree to 2.5%

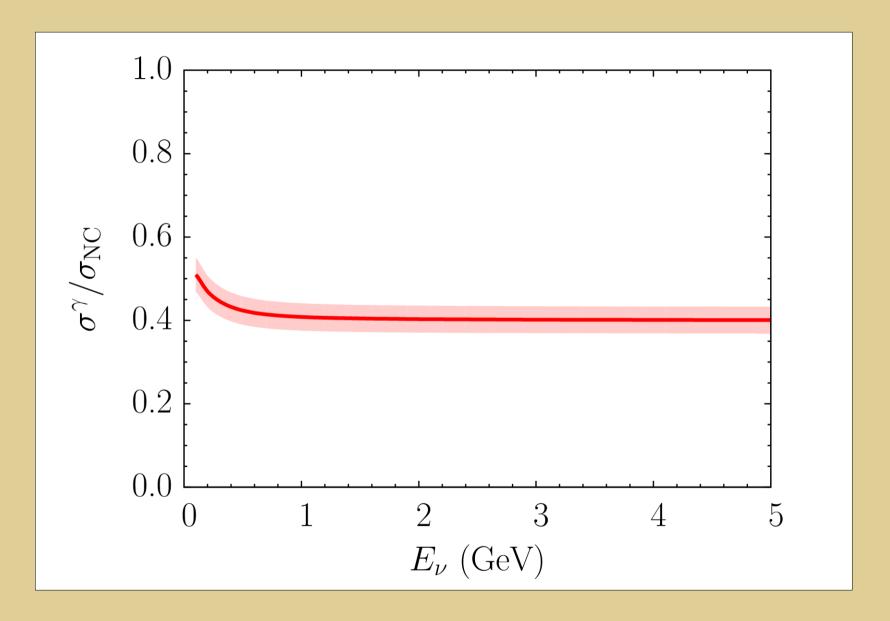


For E > 0.3 GeV  $\sigma_{NC} \approx 0.37 \sigma_{CC}$ 

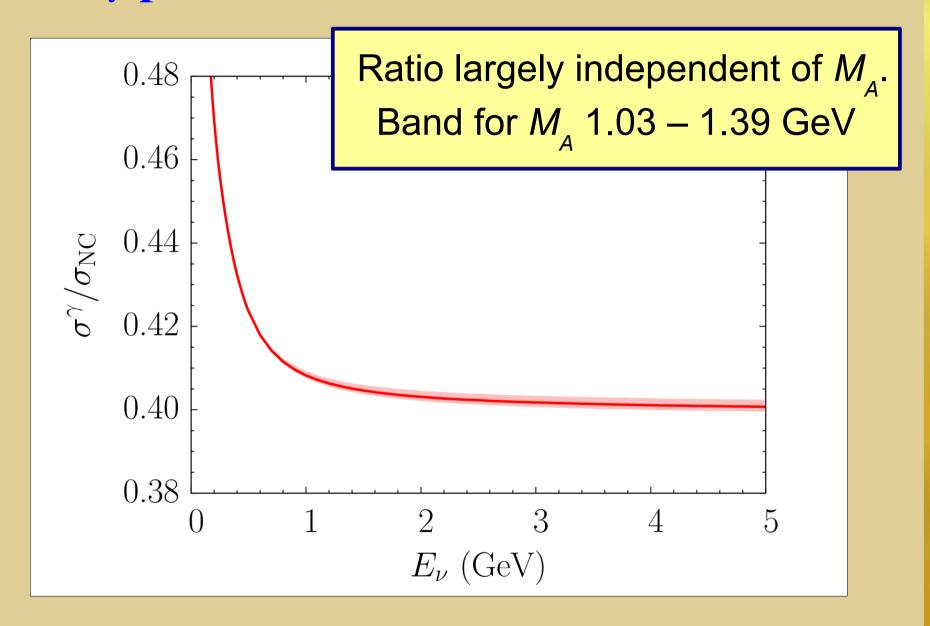




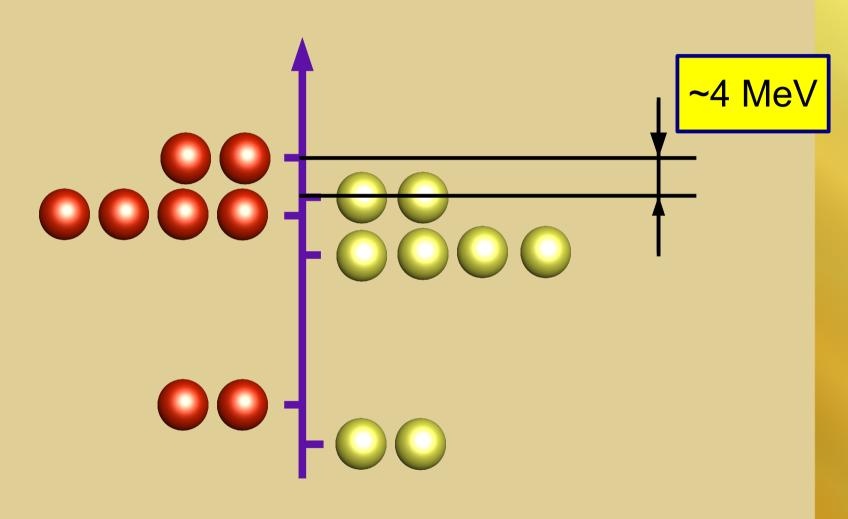
### 6-MeV y production x-section vs. NC x-section



#### 6-MeV y production x-section vs. NC x-section



## Structure of the oxygen nucleus



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