

A new lattice for the beta-beam decay ring to enlarge the stability limit

A. CHANCÉ J. PAYET¹

CEA Saclay IRFU/SACM

4th August 2011

¹presenter

NuFact'11 Beta-Beam Decay ring status

Why a new lattice?



- Collective effects (space charge, beam loading, wake fields...) must be investigated.
- The head tail effects (excitation of the tail of the beam by the head due to tranversal fields) are one of the big issues for the stability in the decay ring.
- Two solutions have been under study to mitigate this effect:
 - a new lattice with a smaller transition gamma;
 - making an amplitude detuning with octupoles, which sometimes enables to push the intensity limit.
 For the decay ring, we will see the impact of this solution on the optics.

How to decrease the transition gamma

The momentum compaction α_P is defined by:

$$\alpha_P = \frac{1}{\gamma_T^2} = \frac{1}{L} \oint \frac{D_x(s) \, ds}{\rho}$$

L is the total length, D_X the horizontal dispersion and ρ the curvature radius. α_P can be increased by:

- increasing the average dispersion in the dipoles.
 - \Rightarrow Larger apertures.
 - \Rightarrow Few degrees of freedom.
 - \Rightarrow Difficult to significantly change the value of α_P .
- increasing the integration length.
 - ⇒ "Wiggler" scheme: the sign of the curvature radius alternates with the dispersion.
 - \Rightarrow The integrate $\oint \frac{D_x(s) ds}{\rho}$ increases.

We have chosen the second solution.

- © Simpler arcs: FODO lattices.
 - © More compact arcs: enlarged duty factor (39%).
 - ③ More dipoles needed (208 against 172 before).



- © Simpler arcs: FODO lattices.
 - © More compact arcs: enlarged duty factor (39%).
 - ③ More dipoles needed (208 against 172 before).



- © Simpler arcs: FODO lattices.
 - © More compact arcs: enlarged duty factor (39%).
 - ③ More dipoles needed (208 against 172 before).



- ③ Simpler arcs: FODO lattices.
 - © More compact arcs: enlarged duty factor (39 %).
 - ③ More dipoles needed (208 against 172 before).



Optical functions in the decay ring





- Straight section length: 2706 m.
- Total length: 2200 π m.
- Dipoles of 7.46 T (for ⁶He²⁺) with an angle of $\pi/84$ rad.

Optical functions in the injection chicane



Optical functions at the center of the chicane:

NuFact'11 Beta-Beam Decay ring status

Optical functions in the arcs



- Phase advance of $\pi/2$ in both planes per FODO period.
- ⇒ Cancels some non linearities brought by the sextupoles.

Dynamic aperture



rms beam sizes: $\sigma_x = 1.83 \text{ mm}$ $\sigma_y = 0.76 \text{ mm}.$

Advantages of the new lattice:

Very simple arcs.

- Two sextupole families to cancel the chromaticity.
- Large dynamic aperture in the momentum range of the beam.

Tracking with PTC_MADX and $\delta=0$



- Step of 1 for n
- Beam sizes: $\sigma_x = 1.83$ mm and $\sigma_y = 0.76$ mm.
- The beam stays elliptic up to 10 σ_X .
- Small non linearities.

Summary of the new lattice

- The transition gamma γ_T for the new lattice is 18.7 (to compare with the old value of 27).
- The value of the average betatron functions are : $< \beta_X >= 125$ m (against 134 m before) $< \beta_Y >= 160$ m (against 175 m before).
- The fractional part of the tune is kept equal to the one of the reference lattice. We have then:

$$Q_x = 22.228$$
 and $Q_y = 19.16$

- The arcs are more compact and much simpler.
- The dynamic aperture is still large (more than 20σ in both planes) and the tracking shows that the beam keeps an elliptic shape for the first 10 sigmas.



- BUT the aimed intensities are still over this limit (more than twice).
- A studied solution is to add octupoles in the lattice to make an amplitude detuning. The expected amplitude detuning are:

$$\frac{\partial Q_x}{\partial \varepsilon_x} = 425 \text{ m}^{-1}, \ \frac{\partial Q_x}{\partial \varepsilon_y} = -878 \text{ m}^{-1}, \ \frac{\partial Q_y}{\partial \varepsilon_y} = 1155 \text{ m}^{-1}$$

- How to make the amplitude detuning?
 - Where to put the octupoles?
 - Which strength?
 - Which impact on the dynamic aperture?

Location of the octupoles

(e)

The intensity limit is mostly sensitive to the derivative $\frac{\partial Q_y}{\partial \varepsilon_y}$. The two other derivatives are less relevant (possibility to optimize their value with optical reasons). We will use two octupole families. Two locations were studied:

- The arcs. The phase advance in the FODO lattices in the arcs is $\pi/2$ rad. The octupoles can be compensated and their contribution to the dynamic aperture should be small. The integrated octupole strengthes are 360 T/m² and 1000 T/m².
- The long straight section. There is plenty of place to insert them there. Larger betatron functions imply weaker octupoles. The integrated octupole strengthes are 32 T/m² and 63 T/m².

Dynamic aperture with octupoles



Dynamic aperture with another $\partial Q_x / \partial \varepsilon_x$



Summary of the amplitude detuning

- We have studied the impact of an octupole detuning in the decay ring.
- Although the phase advances are not optimal in the long straight section, it is more interesting to put the octupoles in this section. The needed octupoles are much weaker for this solution with a similar (or even larger) dynamic aperture.
- If only the derivative $\partial Q_y / \partial \varepsilon_y$ is significant to reduce the collective effects, it is possible to optimize the octupole strenghtes to enlarge significantly the dynamic aperture.
- Amplitude detuning is affordable if only the dynamic aperture is considered.
- BUT The gain is very small on the intensity limit (see Hansen's presentation).

Thank you for your attention!

3 E Tune variations (1)



Tune variations (2)

