New scaling in elastic data



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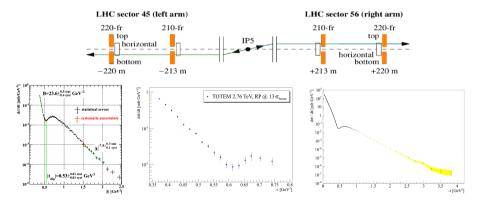
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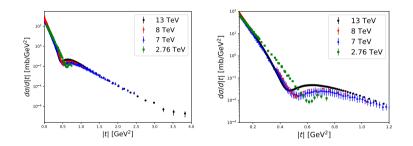
- The starting point: elastic $d\sigma/dt$ measurements from TOTEM
- Definition of a new scaling variable
- Interpretation in the *b*-parameter space
- Phys. Lett. B 830 (2022) 137141, ArXiv 2204.08328

TOTEM elastic $pp \ d\sigma/dt$ cross section measurements

- Elastic *pp* $d\sigma/dt$ measurements: tag both intact protons in TOTEM Roman Pots 2.76, 7, 8 and 13 TeV
- Very precise measurements at 2.76, 7, 8 and 13 TeV: Eur. Phys. J. C 80 (2020) no.2, 91; EPL 95 (2011) no. 41004; Nucl. Phys. B 899 (2015) 527; Eur. Phys. J. C79 (2019) no.10, 861

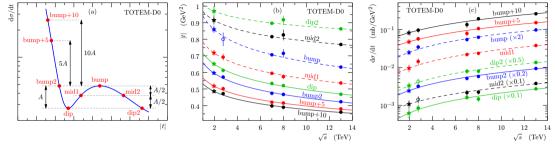


The starting point: elastic $d\sigma/dt$ measurements from TOTEM



- The TOTEM collaboration measured elastic $pp \ d\sigma/dt$ differential cross sections as a function of t for different center-of-mass energies of 2.76, 7, 8, and 13 TeV
- We see that the data points vary together especially in the dip/bump region

The starting point: elastic $d\sigma/dt$ measurements from TOTEM



- Extrapolation of TOTEM data at 2.76, 7, 8, 13 TeV down to the Tevatron energy of 1.96 TeV to compare between pp and $p\bar{p}$ elastic energies
- 8 reference points were identified on TOTEM elastic $pp \ d\sigma/dt$ and their variation of the t and $d\sigma/dt$ as a function of \sqrt{s}
- Striking observation: the same \sqrt{s} dependence for all characteristic points, namely $|t| = a \log(\sqrt{s}) + b$ and $d\sigma/dt = c\sqrt{s} + d$ in the dip/bump region
- It means that all points vary together and are not independent, curves are "parallel"

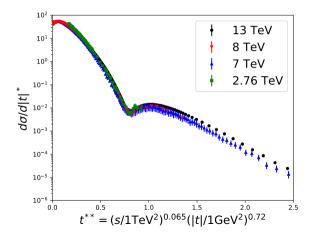
- Find a new variable, which we call t^{**} for which s^{-α}dσ/dt as a function of t^{**} does no longer depend on √s, where α is a constant to be fitted to data
- ${\, \bullet \,}$ We use the Quality Factor method to fit α

$$\mathrm{QF} = \left[\Sigma_i rac{(v_{i+1} - v_i)^2 imes \Delta v_{i+1} imes \Delta v_i}{(u_{i+1} - u_i)^2 + \epsilon^2}
ight] \, .$$

where the Δv_i are the uncertainties on v_i and ϵ is a small constant to regularize divergences when $u_{i+1} = u_i$.

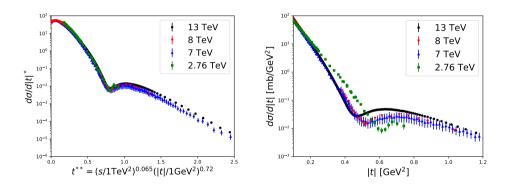
- The u_i and v_i are respectively: $\ln(t^{**})$ and $\ln[(s^{-\alpha}d\sigma/dt]]$
- The QF method is well adapted when there is no analytic expression for $\ln[(s^{-\alpha}d\sigma/dt]]$: fit α so that there is a continuous description of data

The first approach of a new scaling in data



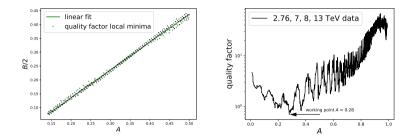
- We introduce the variable
 - $t^* = (s/|t|)^A \times |t|$ which is inspired by geometric scaling in terms of saturation models which is natural since we look for a new scaling in data
- $t^{**} = t^*/s^B$, A and B being parameters to be fitted to data
- We definitely observe that dσ/dt* shows scaling as a function of t**

The first approach of a new scaling in data: zoom in the dip/bump region



- Left: $d\sigma/dt^*$ as a function of t^{**} in the bump/dip region: scaling
- Right: $d\sigma/dt$ as a function of |t| in the bump/dip region: no scaling

A and B parameters are not independent



• We noticed that we have a full valley of parameters that are possible for A and B $(t^{**} = (s/|t|)^A \times t/s^B)$

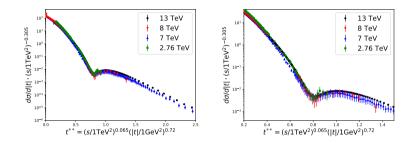
- A and B correlation: B = A = 0.065 obtained by fitting the B value that leads to a minimum of QF for a given A value (QF are found to be similar)
- It means that we can have a 1-parameter fit only

- We learned from data that $d\sigma/dt^*$ scales as a function of $t^{**} = s^{0.065} \times |t|^{1-A}$ where $t^* = (s/|t|)^A \times |t|$
- We have

$$\frac{d\sigma}{dt^*} = \frac{d\sigma}{dt}\frac{dt}{dt^*} = \frac{d\sigma}{dt} \times s^{A\frac{A-1.065}{1-A}} \times f(t^{**}) = (s)^{-\alpha}\frac{d\sigma}{dt}f(t^{**})$$

- Since $d\sigma/dt^*$ scales, it does not depend on *s*, which means the *s* dependence on $d\sigma/dt$ is imposed by scaling
- $s^{-\alpha}d\sigma/dt$ should not depend on s or scales by definition with $\alpha = \frac{-A(A-1.065)}{1-A}$

Scaling in TOTEM elastic data



- We use the QF method to fit the A parameter (using all data from TOTEM at 2.76, 7.8 and 13 TeV): A = 0.28
- $s^{-0.305} d\sigma/dt$ scales as a function of t^{**} (0.305 = $\frac{-A(A-1.065)}{1-A}$ and $t^{**} = s^{0.065} \times |t|^{1-A}$)
- All data shown or zoomed into the dip and bump region where scaling is supposed to be valid (low |t| corresponds to the QED Coulomb region and high |t| the perturbative QCD domain)

Going to the *b*-parameter space

• Relation between the profile function Γ and the amplitude A:

$$\operatorname{Re}(\Gamma(s,b)) = \frac{1}{4\pi i s} \int_0^\infty dq \, q \, J_0(qb) \, A(s,t=-q^2)$$

and

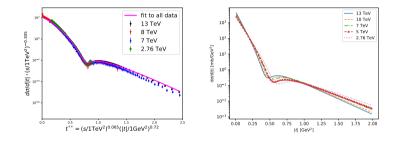
$$\frac{\mathrm{d}\sigma}{\mathrm{d}|t|} = \frac{1}{16\pi s^2} |A(s,t)|^2 = |\mathcal{A}(s,t)|^2$$

• We fit the amplitude to TOTEM data using the formula

$$\begin{array}{lll} \mathcal{A}(s,t) &=& i \big(\mathcal{A}_1(s,t) + \mathcal{A}_2(s,t) \big) e^{i t} \\ \mathcal{A}_1(s,t) &=& \mathsf{N}_1(s) e^{-B_1(s)|t|} \\ \mathcal{A}_2(s,t) &=& \mathsf{N}_2(s) e^{-B_2(s)|t|} e^{i \phi} \end{array}$$

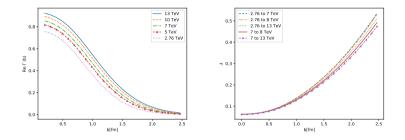
where $N_1(s) = N_1^0(s/1 \text{ TeV}^2)^{\alpha/2}$, $N_2(s) = N_2^0(s/1 \text{ TeV}^2)^{\alpha/2}$, $B_1(s) = B_1^0(s/1 \text{ TeV}^2)^{\gamma/2}$ and $B_2(s) = B_2^0(s/1 \text{ TeV}^2)^{\gamma/2}$

Fit to TOTEM elastic data



- $\alpha = 0.305$ and $\gamma/2 \equiv 0.065/(1 A) = 0.065/0.72 \approx 0.09$ are fixed by scaling
- six free parameters in the fit to $\mathcal{A}(s,t)$: N_1^0 , N_2^0 , B_1^0 , B_2^0 , ϕ , and θ , and predictions of $d\sigma/dt$ for different s from the fits
- Fit quality: $\chi^2/dof = 1.08$ for $0.2 < t^{**} < 1.5$ in the dip-bump region (476 data points) that avoid very low |t| (Coulomb QED region) and high |t|, perturbative QCD domain $(\chi^2/dof = 8.7$ for the full domain in |t|, 599 data points)

The profile function



Use top formula to compute the profile function (see previous slide)
We define λ as

• $\lambda = (\alpha - \gamma)/2 + \text{term}$ vanishing when $b \to 0$, $\lambda = 0.06$ which means that scaling predicts a universal behavior of λ at small b

Possible interpretation of a new scaling and conclusion

- Values of λ at small b (0.06) are compatible with expectations from a dense object, such as a black disc
- Values of λ reach higher values around 0.3 for b = 1 fm, which is reminiscent of the power-law exponent in the small-x limit of QCD, described by the perturbative BFKL evolution equation at next-to-leading logarithmic accuracy
- Scaling together with the value of λ at low b, could be interpreted as having a large density of gluons inside colorless gluonic compounds (responsible for diffraction) that reach the black disc limit at small b. At higher b, the density of gluons is smaller and in principle describable by BFKL dynamics
- In this sense, we can interpret our results as the presence of a black disk in the proton at high energy with a high density of gluons

- TOTEM measured very precisely the elastic $d\sigma/dt$ at high energy at the LHC (2.76, 7, 8, and 13 TeV)
- It seems that the points on the $d\sigma/dt$ do not vary independently from each other, but rather in a correlated way
- We analyzed the behavior of the differential cross section of proton-proton elastic scattering as a function of t and s at LHC energies and found that $d\sigma/dt$ at $\sqrt{s} = 2.76$, 7, 8, and 13 TeV exhibits scaling. The data fall onto a universal curve after mapping them with $d\sigma/dt \rightarrow d\sigma/dt(s)^{-0.305}$ and $|t| \rightarrow (s)^{0.065} |t|^{0.72}$
- This might be related to the presence of a black disk in the proton at high energy

