Spontaneous Matter Genesis

Matthew McCullough

In collaboration with John March-Russell To appear on arXiv soon....

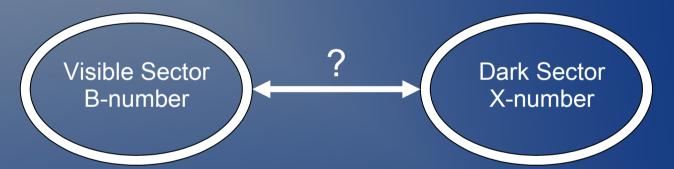
PONT 2011
Palais des Papes, Avignon

Thursday, April 21st

Asymmetric Dark Matter

$$\Omega_B \sim \Omega_{DM}$$

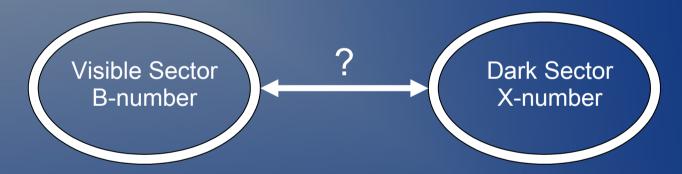
- A clue or a coincidence?
- If ADM then options include:



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Sequestered Sectors



Coupled Sectors



Populate individually?

Coupled Sectors



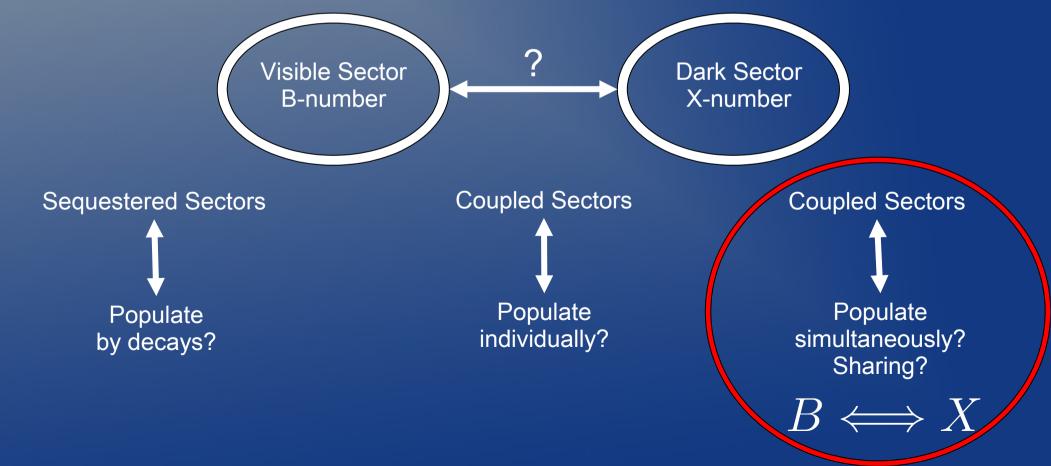
Populate simultaneously? Sharing?

$$B \iff X$$

Asymmetric Dark Matter

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Sharing operator breaks:

$$U(1)_B \times U(1)_X \Rightarrow U(1)_{\{B,X\}}$$

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- Follow the 3 "Palais des Papes" Commandments:
 - Though shall violate baryon number
 - Though shall respect neither C nor CP
 - Though shall not find oneself in thermal equilibrium

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- Generalized Sakharov:
 - Need violation of $U(1)_{\{B,X\}}$
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- Many recent examples: Shelton & Zurek, Davoudiasl et al, Haba & Matsumoto, Buckley & Randall...

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- In a CPT-invariant theory!

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 - More recently: Carroll and Shu, 2005.

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 - More recently: Carroll and Shu, 2005.
- Breaking CPT means breaking Lorentz
 - Bad idea now
 - Could happen in early Universe...

Lorentz Violation

Consider scalar field in expanding Universe

$$\ddot{\phi} + 3H\dot{\phi} + \frac{dV}{d\phi} = 0$$

- Post-inflation ϕ displaced, but spatially homogeneous.
- If $\overline{m_{\phi} \ll H}$ then critically damped;

$$\dot{\phi} \approx \frac{-1}{3H} \frac{dV}{d\phi} \approx \frac{-1}{5g_{\star}^{1/2}(T)} \frac{M_P m_{\phi}^2 \phi}{T^2}$$

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• Lorentz violation: $\partial_{\mu}\phi \equiv \{\dot{\phi}, {f 0}\}$

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Equilibrium Thermodynamics (ET)

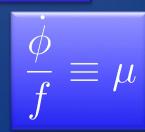
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Scalar coupled to current:

$$\mathcal{L} \supset \frac{\partial_{\mu}\phi}{f} J^{\mu}_{\{B,X\}} \Rightarrow \frac{\dot{\phi}}{f} (n_{+} - n_{-})$$

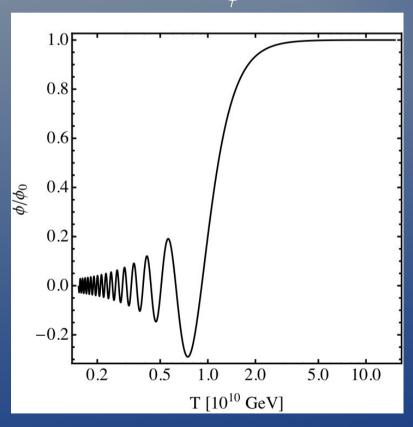
• Generates effective chemical potential:



- Sakharov in spontaneous genesis:
 - C-violation; rolling ϕ derivatively coupled to $J^{\mu}_{\{B,X\}}$
 - Thermal equilibrium
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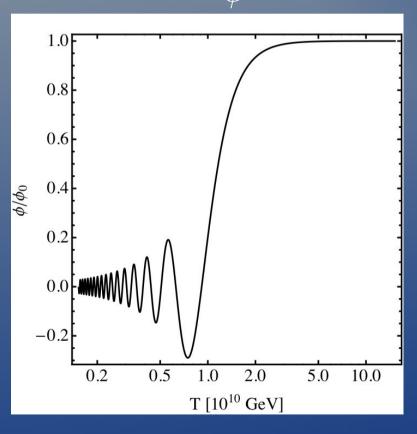
- Sakharov in spontaneous genesis:
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 - Need violation of $U(1)_{\{B,X\}}$
- Assume extra interactions which violate $\,U(1)_X\,$
- Interactions freeze out at $\overline{T_X}$, whenever $\Gamma_X \lesssim H$
- Sharing means generation of $U(1)_{\{B,X\}}$ asymmetry
- Sharing freezes out at T_S

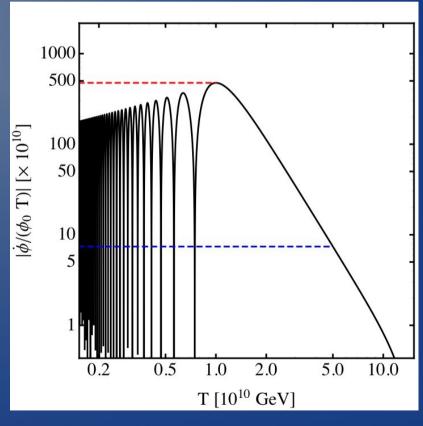
- Want X-violation to freeze out before ϕ starts oscillating
- E.g. for $m_\phi=1~{
 m TeV}$, field evolves as:



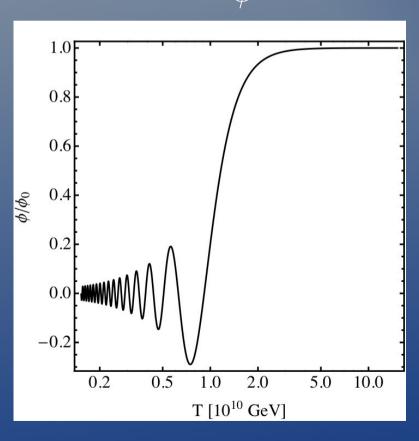
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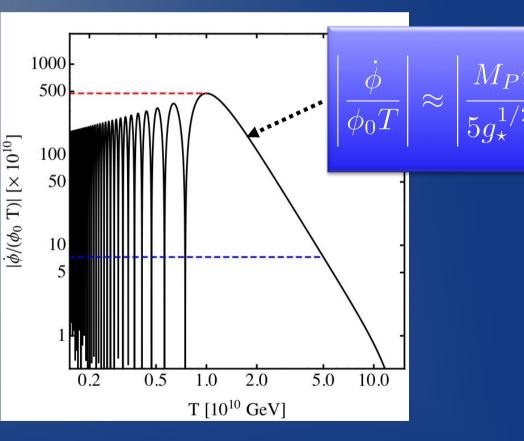
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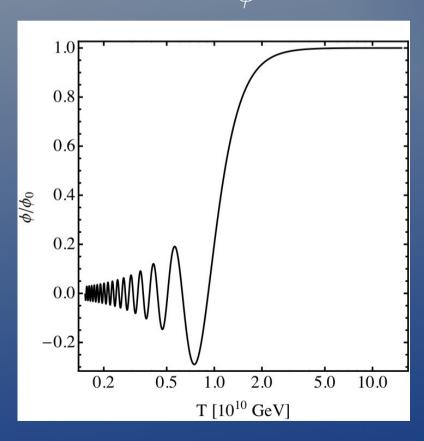


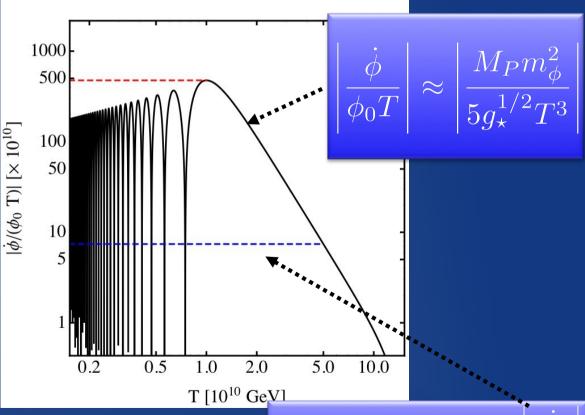
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$$\left| |n_{+} - n_{-}| \propto \left| \frac{\mu}{T} \right| = \left| \frac{\phi}{Tf} \right| \right|$$

Assume a SUSY model, with:

$$\mathcal{L} \supset -\frac{1}{2}m_{\phi}^{2}\phi^{2} + \frac{\partial_{\mu}\phi}{f}J_{X}^{\mu} + \int d^{2}\theta \ M_{X}\overline{X}X$$

$$+ \int d^{2}\theta \ W_{MSSM} + \mathcal{L}_{soft}$$

$$+ \int d^{2}\theta \ \frac{1}{M_{S}^{2}}X^{2}U^{c}D^{c}D^{c}$$

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Dark sector:

 $U(1)_X$ symmetry Chemical potential from scalar Dirac mass for dark matter

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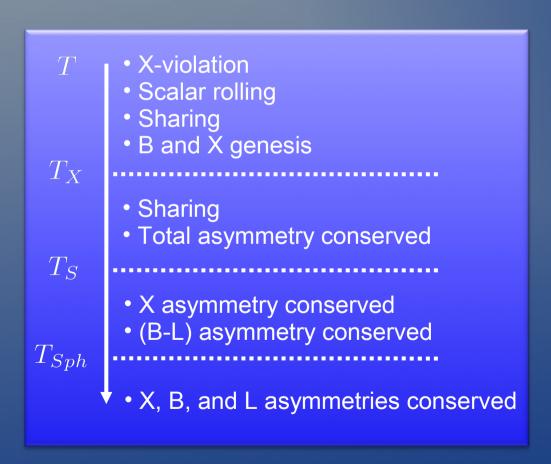
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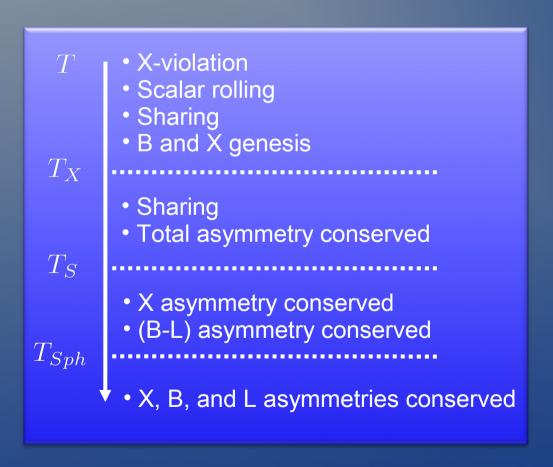
Sharing operator:

Breaks $U(1)_X \times U(1)_B \to U(1)_{X-2B}$ If $M_S \gtrsim 1~{
m TeV}$ Freezes out at $T_S \gtrsim 70~{
m GeV}$

Schematically:



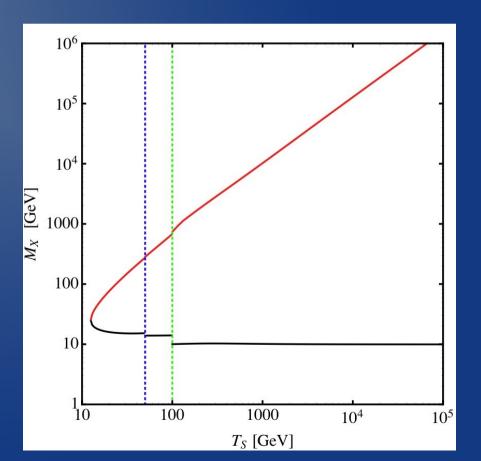
Schematically:



• Require correct abundances:

$$\frac{X(T_S)}{B(T_{Sph})} = \left(\frac{X}{B-L}\right)_{T_S} \left(\frac{B-L}{B}\right)_{T_{Sph}} = 4.69 \frac{m_B}{m_X}$$

ullet Dark matter mass just depends on T_S

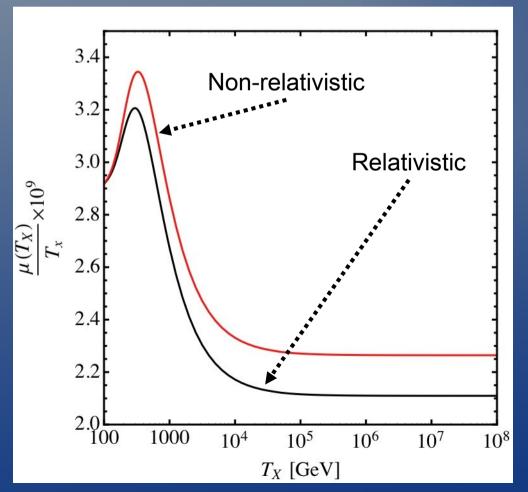


• Can track back to find required chemical potential at T_X

$$X(T_X, \mu_X(T_X)) = \left(\frac{X}{X + 2(B - L)}\right)_{T_X} \left(\frac{X + 2(B - L)}{B - L}\right)_{T_S} \left(\frac{B - L}{B}\right)_{T_{Sph}} B_{T_{Sph}}$$

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$$\mu_X(T_X) \approx 2.2 \times 10^{-9} T_X$$

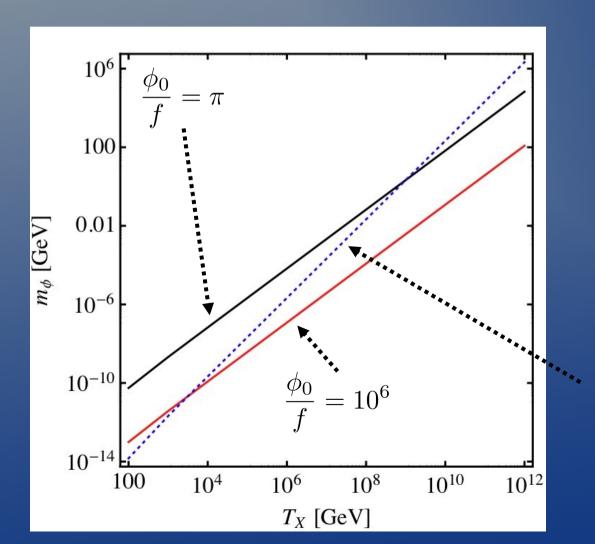
$$\Rightarrow$$

$$m_{\phi} \approx \left(g_{\star}^{1/2} \frac{T_X^3}{M_P} \frac{f}{\phi_0}\right)^{1/2} \times 10^{-4}$$

• But damped roll requires:

$$m_{\phi} < 1.66 g_{\star}^{1/2} \frac{T_X^2}{M_P}$$

• Can track back to find required chemical potential at T_X ... and use this to find scalar mass.



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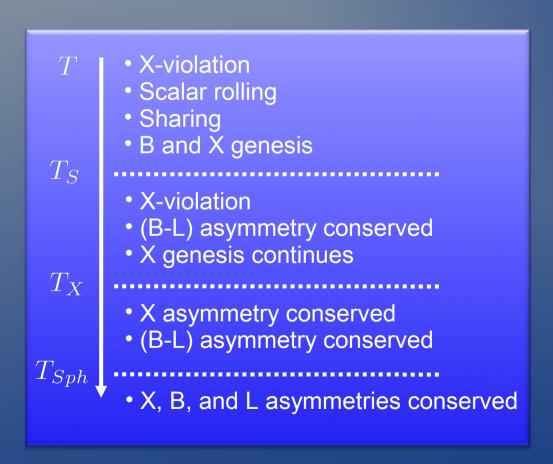
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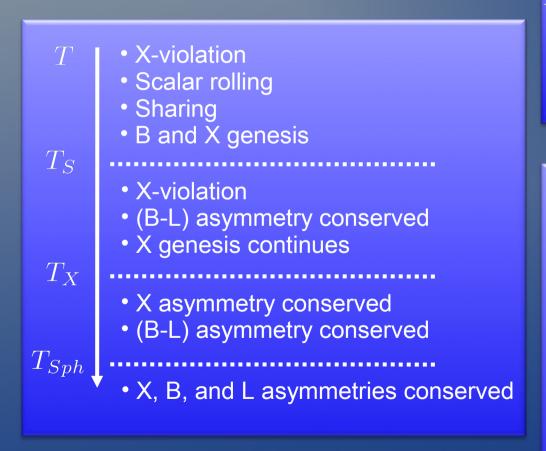
Genesis After Sharing

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$$X(T_S, \mu_X(T_S)) = \left(\frac{X}{B-L}\right)_{T_S} \left(\frac{B-L}{B}\right)_{T_{Sph}} B_{T_{Sph}}$$

• Which gives us:

$$\mu_X(T_S) \approx 2.2 \times 10^{-9} T_S$$

Potential sourced by rolling scalar:

$$\mu_X(T) \propto \frac{1}{T^2}$$

• Thus:

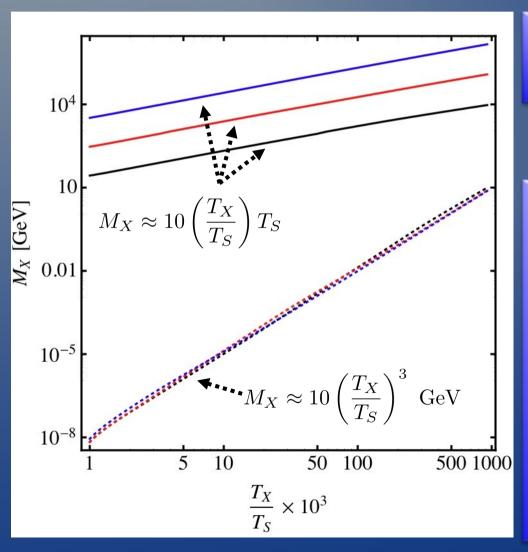
$$\mu_X(T_X) = \frac{T_S^2}{T_X^2} \mu_X(T_S)$$

• So we can find:

$$\rho_{DM} = M_X X(T_X, \mu_X(T_X)) s(T_{now})$$

Genesis After Sharing

Broadens asymmetric dark matter mass range



• Blue: $T_S = 1 \text{ TeV}$

• Red: $T_S = 10 \text{ TeV}$

• Black: $T_S = 100 \text{ TeV}$

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Dark Sector Models

- Require violation of $U(1)_{\{B,X\}}$ at high T
- In visible sector constrained
- In dark sector?
 - Dark sphalerons:
 - » Strongly first order not required
 - » Only one family as CP-violation not required

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 - Explicit violation
 - » Spontaneously broken dark GUT?
- Require efficient annihilation of symmetric component
 - Additional light bosons?

Spontaneous Matter Genesis Conclusions

- Spontaneous genesis in the dark sector is appealing:
 - Damped scalars very natural
 - All in thermodynamic equilibrium
 - No need for additional CP-violation
 - X-violating operators much less constrained now than B-violation

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 - Damped scalars very natural
 - All in thermodynamic equilibrium
 - No need for additional CP-violation
 - X-violating operators much less constrained now than B-violation
- Connecting dark to visible sector appealing
 - Spontaneous genesis allows for broad range of dark matter masses
 - Predicts: Light scalar, direct detection...



