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(cosmological) tests of acceleration: why and what

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testing acceleration...



Outline

- quantifying the acceleration
- does inflation look like dark energy?
- beyond w
 - formalism
 - examples
 - what can we learn
- outlook + summary

measuring dark things (in cosmology)

Einstein: $G_{\mu\nu} = 8\pi G T_{\mu\nu}$

(determined by the metric) → geomet

stuff
(what is it?)

assuming FLRW:

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho$$

$$\dot{\rho} = -3\frac{\dot{a}}{a}(\rho + p)$$

our favourite theory

cosmologists observe the
geometry of space time

depends on the **total**
energy momentum tensor

That is what we measure!

something
else

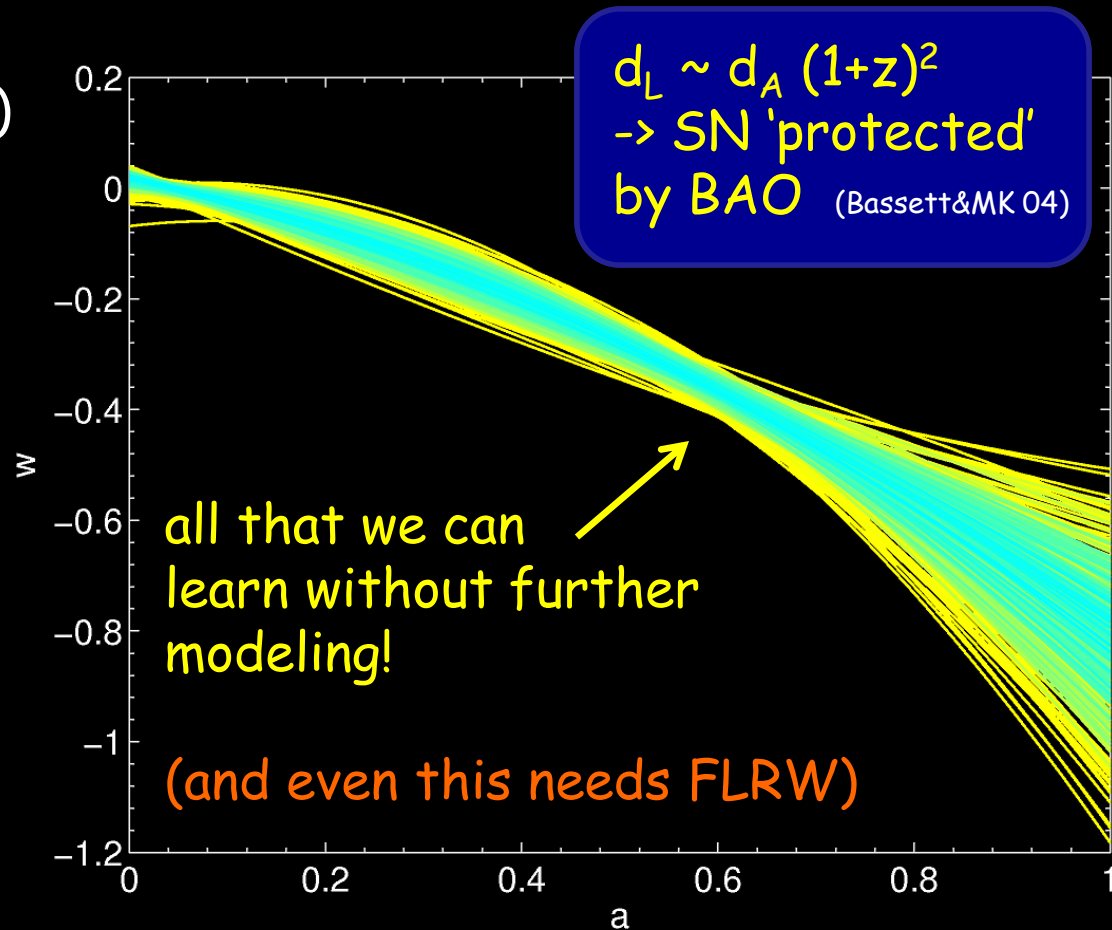


constraints on the total w

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho \quad \dot{\rho} = -3\frac{\dot{a}}{a}(\rho + p) \quad \rightarrow \text{rewrite } p = w\rho$$

- quadratic expansion of $w(a)$
- fit to Union SNe, BAO and CMB peak location
→ just distances, no perturbations
- best: $\chi^2 = 309.8$
- Λ CDM: $\chi^2 = 311.9$
- **w const.: $\chi^2 = 391.3$**

definitely not $w=0$!



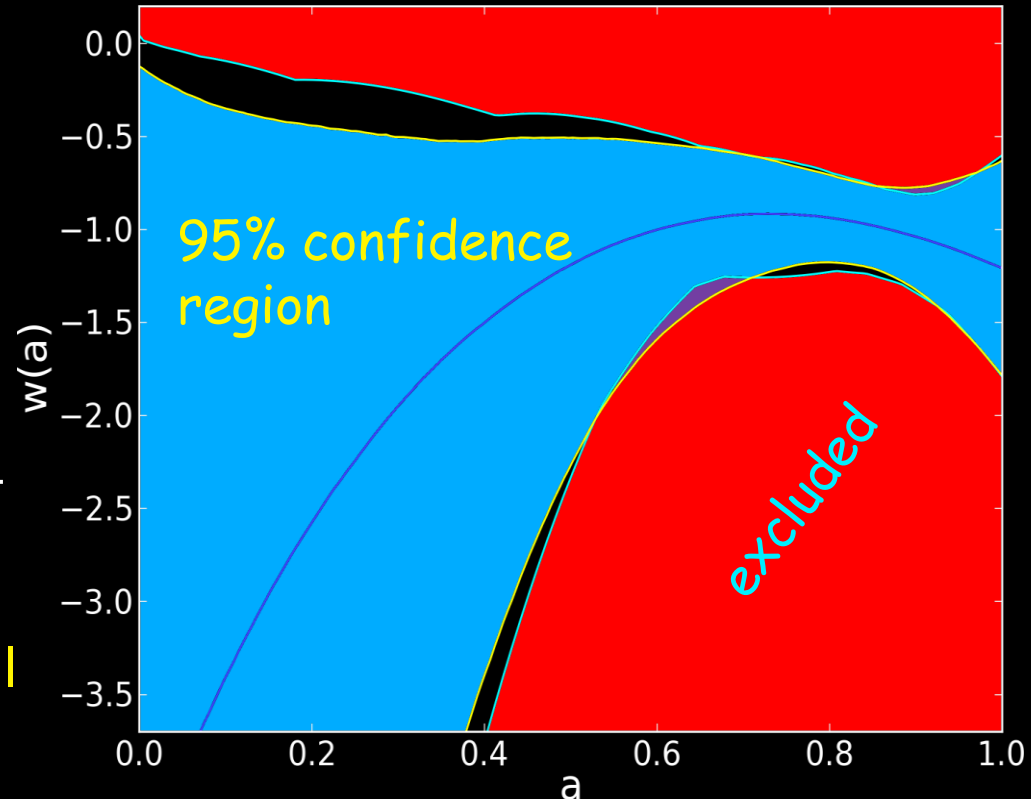
dark energy $w(a)$

total w : weighted combination of w_{DE} and $w_{DM}(=0)$

→ what is what?

→ need specific model for DE!

- canonical scalar field model
[$\leftrightarrow c_s^2=1, \sigma=0$]
- WMAP-7yr +
SN-Ia compilation
- regularised transition of $w=-1$
- cubic expansion of $w(a)$
- cosmological constant fits well
- $|1+w| < 0.2$ at $a \sim 0.8$ @ 2σ



To make this figure, we needed to fix the perturbations -
but what can/should be fixed?

very early dark energy?

The cosmological constant ($w=-1$, no perturbations) fits the data very well. Why look further?

Because we don't like it!

Inflation is usually modeled as a period of accelerated expansion, just like dark energy

- Is this unavoidable?
- Was inflation due to a cosmological constant?
- What would we have observed during inflation?

causal sources after COBE?

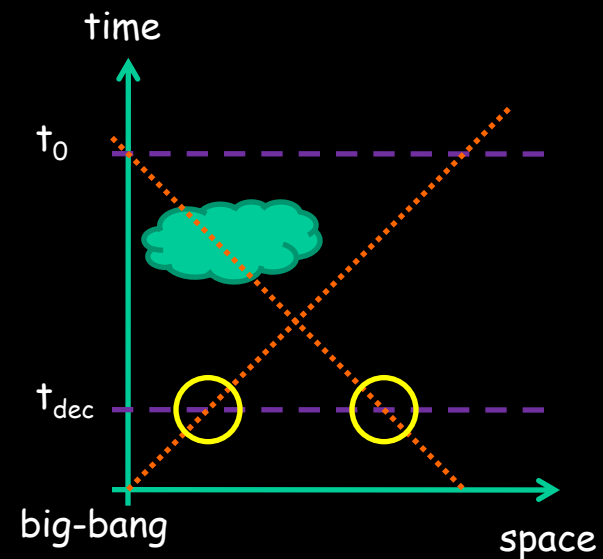
COBE observed fluctuations correlated on scales much larger than the horizon at last scattering!

-> Horizon problem

-> is this not proof of "acausal" physics?

NO!

Can create them at late times with time-dependent potentials (ISW).



(a) causality constraints

(Scodeller, MK & Durrer, 2009)

outgoing spherical shells of energy with velocity v

accepted points in MCMC chain
(sorry for ugly figure)

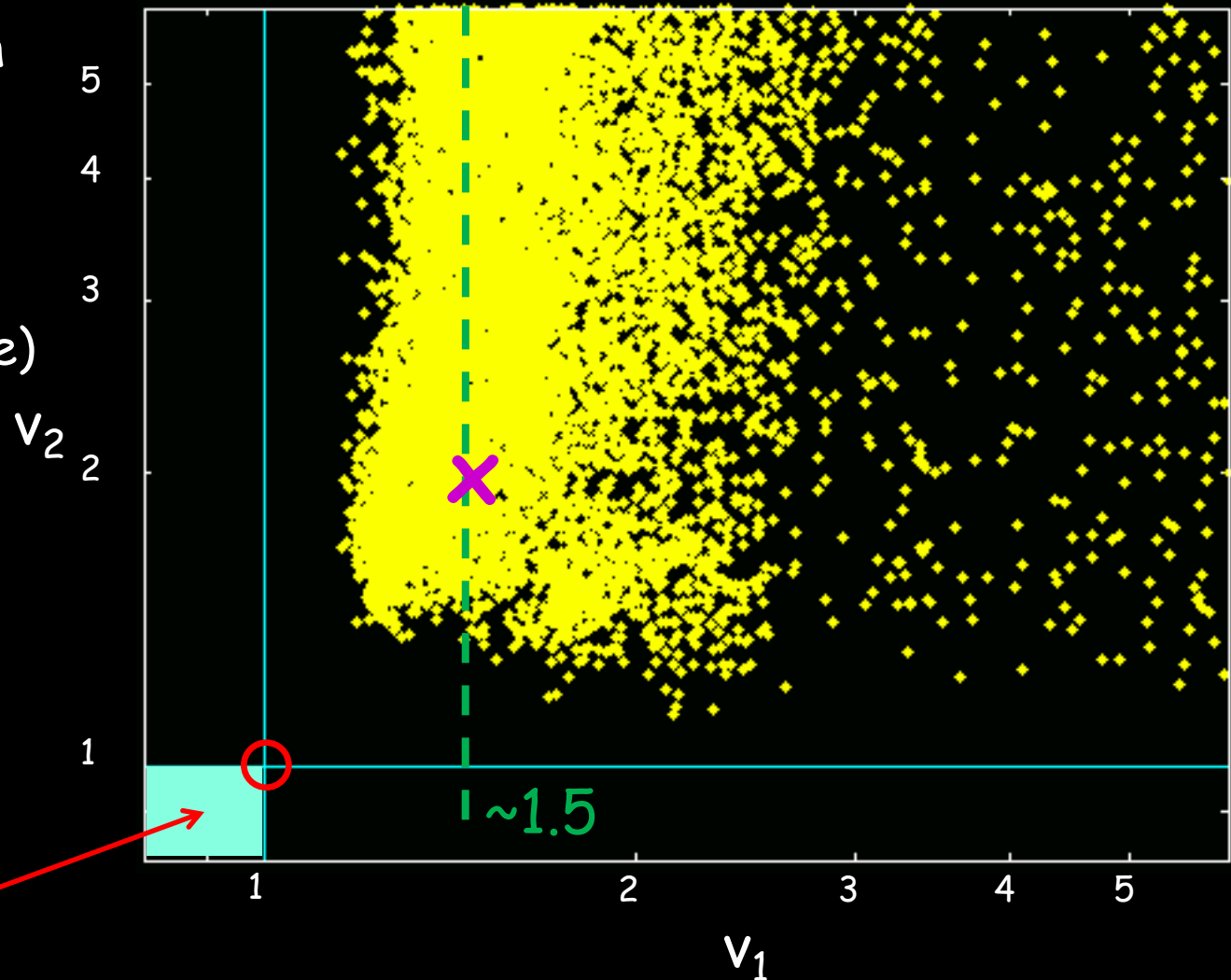
good fit requires:

$$v_1 > 1.2$$

$$v_2 > 1.3$$

other parameters roughly as always

causal region



TE cross-polarisation

Polarisation induced at last scattering and reionisation
[Spergel & Zaldarriaga, 1997] -- TE shows a dip around $l \sim 100$:

adiabatic density mode $\sim \cos(kc_s t_{dec})$
velocity mode: derivative $\sim \sin(kc_s t_{dec})$

TE: $\sin(2kc_s t_{dec})$

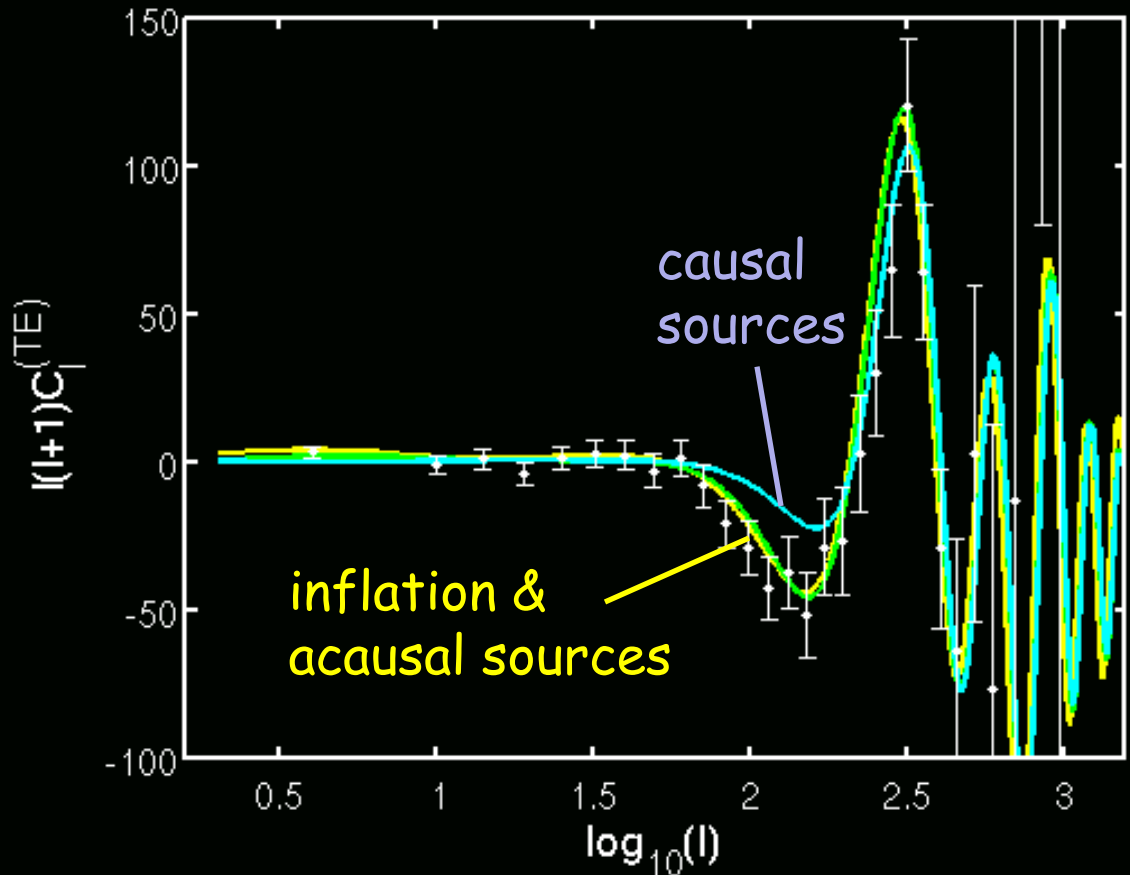
peak: $kt_{dec} \approx 0.66$

horizon: $kt_{dec} \sim 1/v$

$\rightarrow v \sim 1.5$

possibilities:

- inflation
- acausal physics
- huge reionisation fine-tuning (?)



w during inflation

(Ilic, MK, Liddle & Frieman, 2010)

- Scalar field inflaton:

$$+w = \frac{2\dot{\phi}^2}{3H^2} - \frac{2}{3H^2} \frac{\ddot{\phi}}{\dot{\phi}} \frac{\dot{\phi}}{H} \frac{1}{H^2} \frac{d}{dt} \left(\frac{\dot{\phi}}{H} \right)$$

and $r = T/S \sim 24(1+w)$

- Link to dw/da :

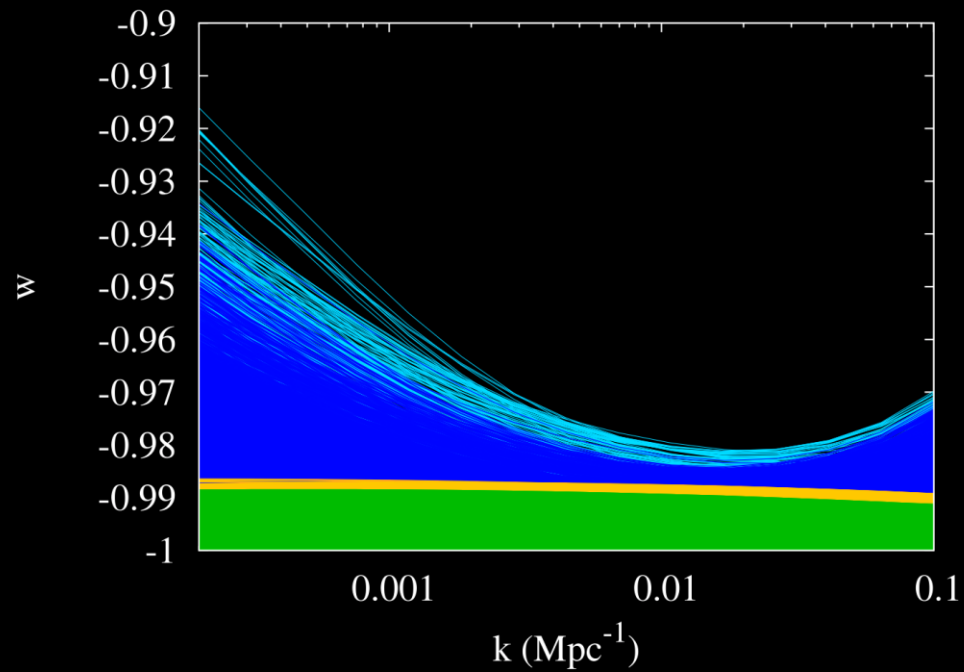
$$\frac{dw}{da} = \frac{2\dot{\phi}^2}{3H^2} \frac{d}{da} \left(\frac{\dot{\phi}}{H} \right) - \frac{2}{3H^2} \frac{d}{da} \left(\frac{\ddot{\phi}}{\dot{\phi}} \frac{\dot{\phi}}{H} \right)$$

$$\frac{dw}{da} = \frac{2}{3} \left(\frac{\ddot{\phi}}{\dot{\phi}} \frac{\dot{\phi}}{H} \right) - \frac{4}{3} \frac{\dot{\phi}}{H} \frac{d}{da} \left(\frac{\dot{\phi}}{H} \right)$$

$n_s \neq 1 \Rightarrow \varepsilon \neq 0$ or $\eta \neq 0$
 $\Rightarrow w \neq -1$ and/or w not constant

WMAP 5yr constraints on w :

- $(1+w) < 0.02$
- No deviation from $w=-1$ necessary (but in the middle of long slow-roll period, not clear if representative of dark energy)



- $w \sim -1$ appears natural during observable period of inflation
- but it was not an (even effective) cosmological constant!

measuring dark things (in cosmology)

Einstein eq. (possibly effective):

$$G_{\mu\nu} - 8\pi G T_{\mu\nu}^{(bright)} = 8\pi G T_{\mu\nu}^{(dark)}$$

directly measured

given by metric:

- $H(z)$
- $\Phi(z, k), \Psi(z, k)$

- inferred from lhs
- obeys conservation laws
- can be characterised by
 - $p = w(z) \rho$
 - $\delta p = c_s^2(z, k) \delta \rho, \pi(z, k)$

linear perturbation equations

metric: $ds^2 = -(1 + 2\psi)dt^2 + a(t)^2(1 - 2\phi)dx^2$

conservation equations (in principle for full dark sector)

$$\delta'_i = 3(1 + w_i)\phi' - \frac{V_i}{Ha^2} - \frac{3}{a} \left(\frac{\delta p_i}{\rho_i} - w_i \delta_i \right) \quad \delta p = c_s^2 \delta \rho + 3Ha(c_s^2 - c_a^2)\rho \frac{V}{k^2}$$

$$V'_i = -(1 - 3w_i)\frac{V_i}{a} + \frac{k^2}{Ha} \left(\frac{\delta p_i}{\rho_i} + (1 + w_i)(\psi - \sigma_i) \right)$$

(vars: $\delta = \delta\rho/\rho$, $V \sim$ divergence of velocity field, δp , σ anisotropic stress)

Einstein equations (common, may be modified if not GR)

$$k^2 \phi = -4\pi Ga^2 \sum_i \rho_i \left(\delta_i + 3Ha \frac{V_i}{k^2} \right)$$

$$k^2 (\phi - \psi) = 12\pi Ga^2 \sum_i (1 + w_i) \rho_i \sigma_i$$

(Bardeen 1980)

simplified observations

- **Curvature** from **radial & transverse BAO**
- **$w(z)$** from **SN-Ia, BAO** directly (and contained in most other probes)
- In addition 5 quantities, e.g. **ϕ , ψ , bias, δ_m , V_m**
- Need **3 probes** (since 2 cons eq for DM)
- e.g. 3 power spectra: **lensing, galaxy, velocity**
- **Lensing** probes **$\phi + \psi$**
- **Velocity** probes **ψ** (z-space distortions?)
- And **galaxy $P(k)$** then gives bias
- **what do we learn if we do this?**

some model predictions

$$k^2\phi = -4\pi G a^2 Q \rho_m \Delta_m \quad \psi = (1 + \eta)\phi$$

scalar field: $S = \int d^4x \sqrt{-g} \left(\frac{1}{2} \partial_\mu \phi \partial^\mu \phi + V(\phi) \right)$

One degree of freedom: $V(\phi) \leftrightarrow w(z)$
 therefore other variables fixed: $c_s^2 = 1, \sigma = 0$
 $\rightarrow \eta = 0, Q(k \gg H_0) = 1, Q(k \sim H_0) \sim 1.1$

(naïve) DGP: compute in 5D, project result to 4D

Lue, Starkmann 04
 Koyama, Maartens 06
 Hu, Sawicki 07

$$\eta = \frac{2}{3\beta - 1} \quad Q = 1 - \frac{1}{3\beta}$$

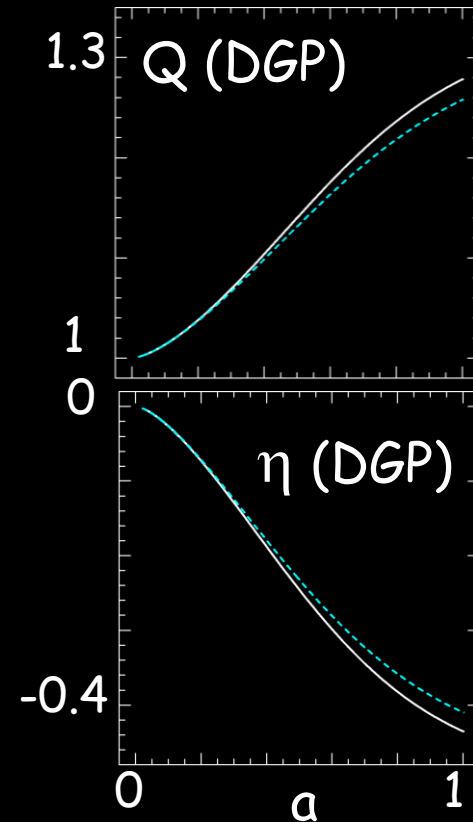
implies large
DE perturb.

Scalar-Tensor:

Boisseau, Esposito-Farese, Polarski, Starobinski 2000,
 Acquaviva, Baccigalupi, Perrotta 04

$$\mathcal{L} = F(\phi)R - \partial_\mu \phi \partial^\mu \phi - 2V(\phi) + 16\pi G^* \mathcal{L}_{\text{matter}}$$

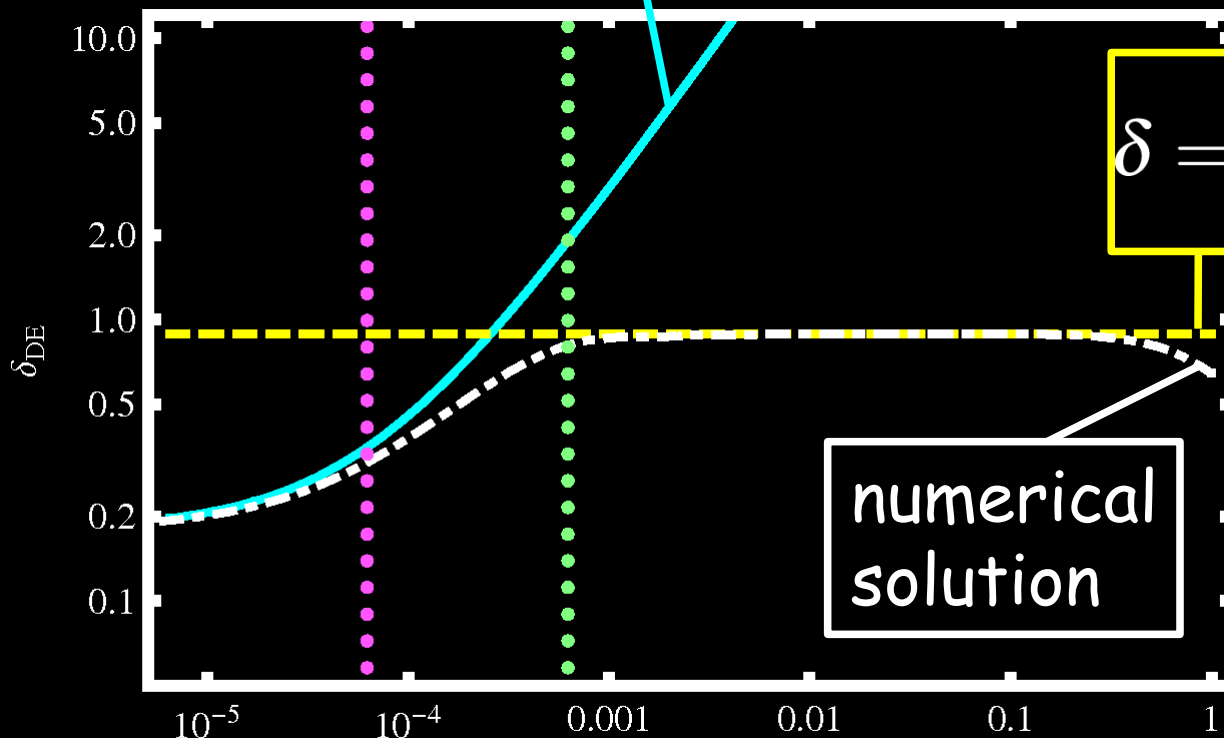
$$\eta = \frac{F'^2}{F + F'^2} \quad Q = \frac{G^* 2(F + F'^2)}{F G_0 2F + 3F'^2}$$



behaviour of scalar field δ

model $\{w, c_s, \sigma=0\}$; matter dom.: $\Phi = \text{constant}$, $\delta_m \sim a$

$$\delta = \delta_0(1+w) \left(\frac{a}{1-3w} + \frac{3H_0^2 \Omega_m}{k^2} \right) \rightarrow \delta(w=-0.8) \leq 1/20 \delta(w=0) \text{ on subhorizon scales}$$

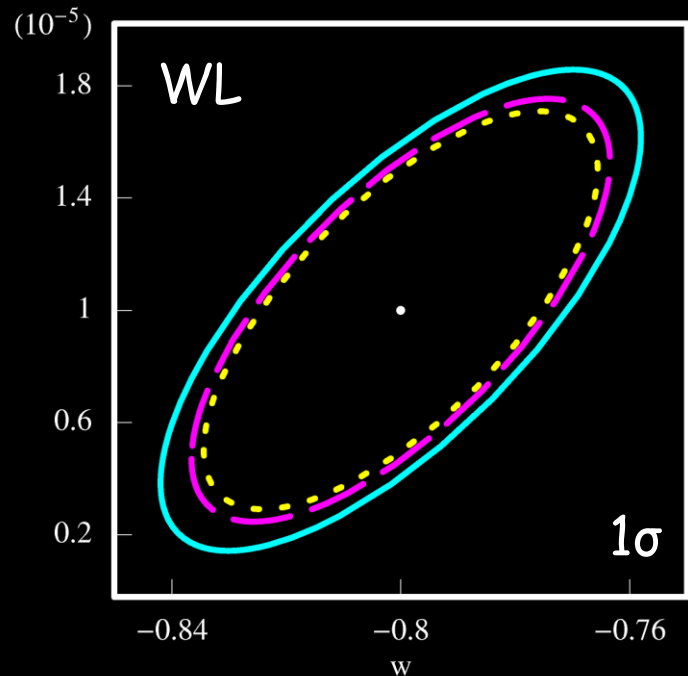
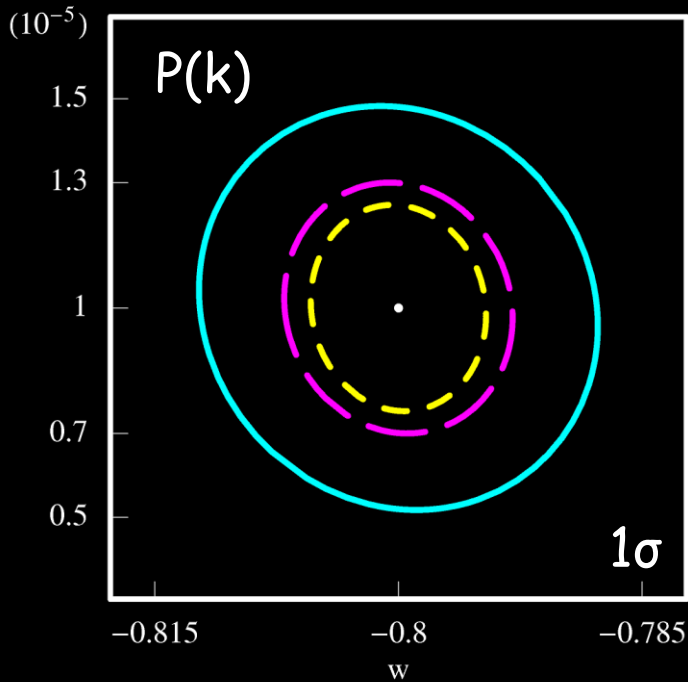


$$\delta = \delta_0 \frac{3}{2} (1+w) \frac{H_0^2 \Omega_m}{c_s^2 k^2}$$

- $w = -0.8$
- $c_s = 0.1$
- $k = 200 H_0$

can we see the DE sound horizon?

two large surveys to $z_{\max} = 2, 3, 4$
 fiducial model has $w = -0.8$
 → only if $c_s < 0.01$ can we measure it!
 (for $w = -0.9$ we need $c_s < 0.001$)



$P(k) + WL$			
c_s^2	σ_{w_0}	$\sigma_{c_s^2}/c_s^2$	σ_W/W
10^{-5}	0.00639	0.15	0.11
10^{-4}	0.00581	0.41	0.36
10^{-3}	0.00547	0.87	1.02
10^{-2}	0.00531	2.48	2.39
10^{-1}	0.00528	14.79	13.14
1	0.00524	22.05	21.29

(Sapone & MK 2009; Sapone, MK & Amendola 2010)

the importance of η / σ

(Saltas & MK 2011, cf talk yesterday afternoon)

scalar-tensor theories: $\eta = \frac{F'^2}{F + F'^2}$

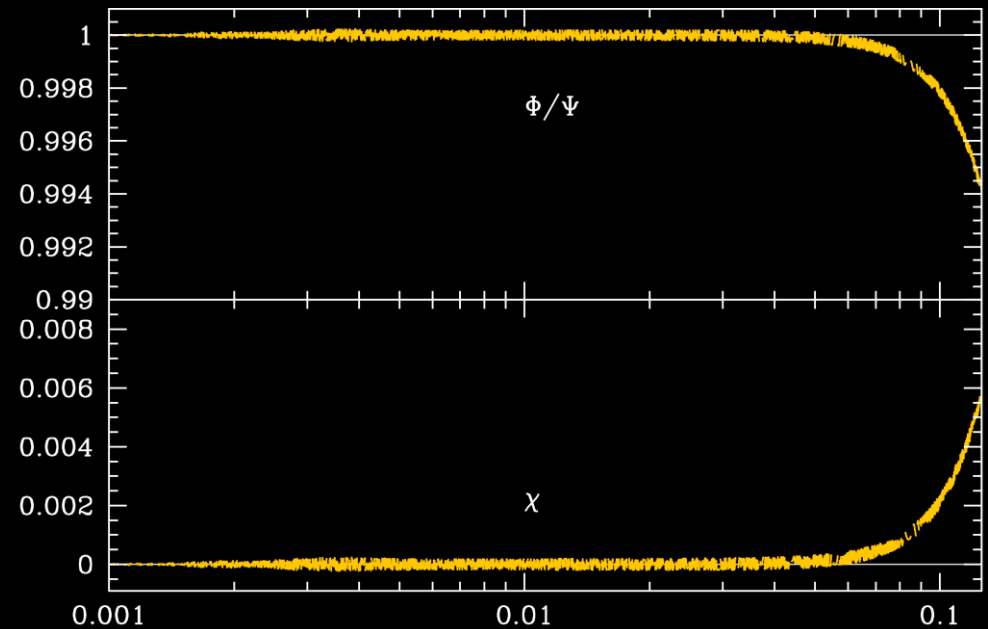
$f(R)$ theories: $\eta = 0 \leftrightarrow f''(R) = 0$; $R+f(G)$: $\eta = 0 \leftrightarrow f''(G) = 0$

$f(R,G)$: in de Sitter background requires mass of effective scalar to diverge \rightarrow instabilities, dS cannot be reached dynamically

also in DGP $\eta \neq 0$!

canonical scalar field: $\eta = 0$
 \rightarrow standard 'GR' model: at late times only very small anisotropic stress from relativistic particles

$\rightarrow \eta$ can rule out whole classes of models!



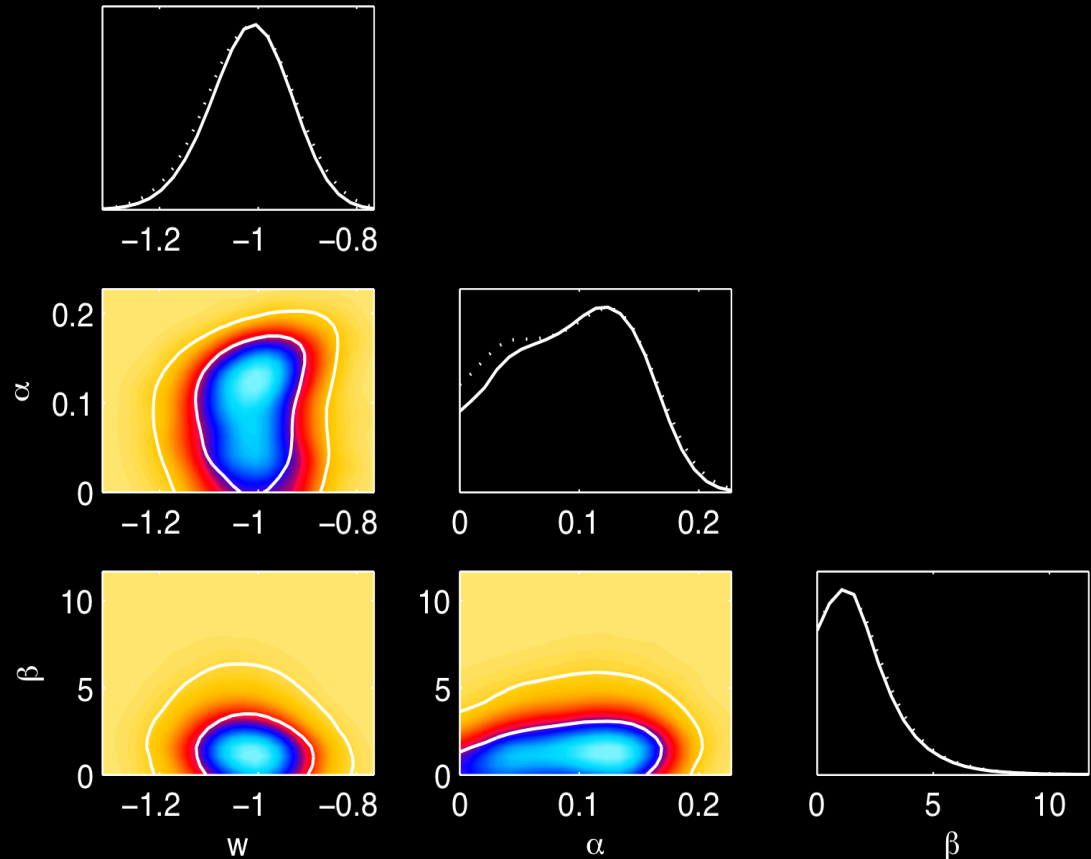
(Pogosian & Silvestri 2008)

current constraints on $\sigma \sim (\phi - \psi)$

(from Lukas Hollenstein, private communication)

Inspired by modified gravity models: $\sigma \sim \alpha \Delta_m + \beta \psi$

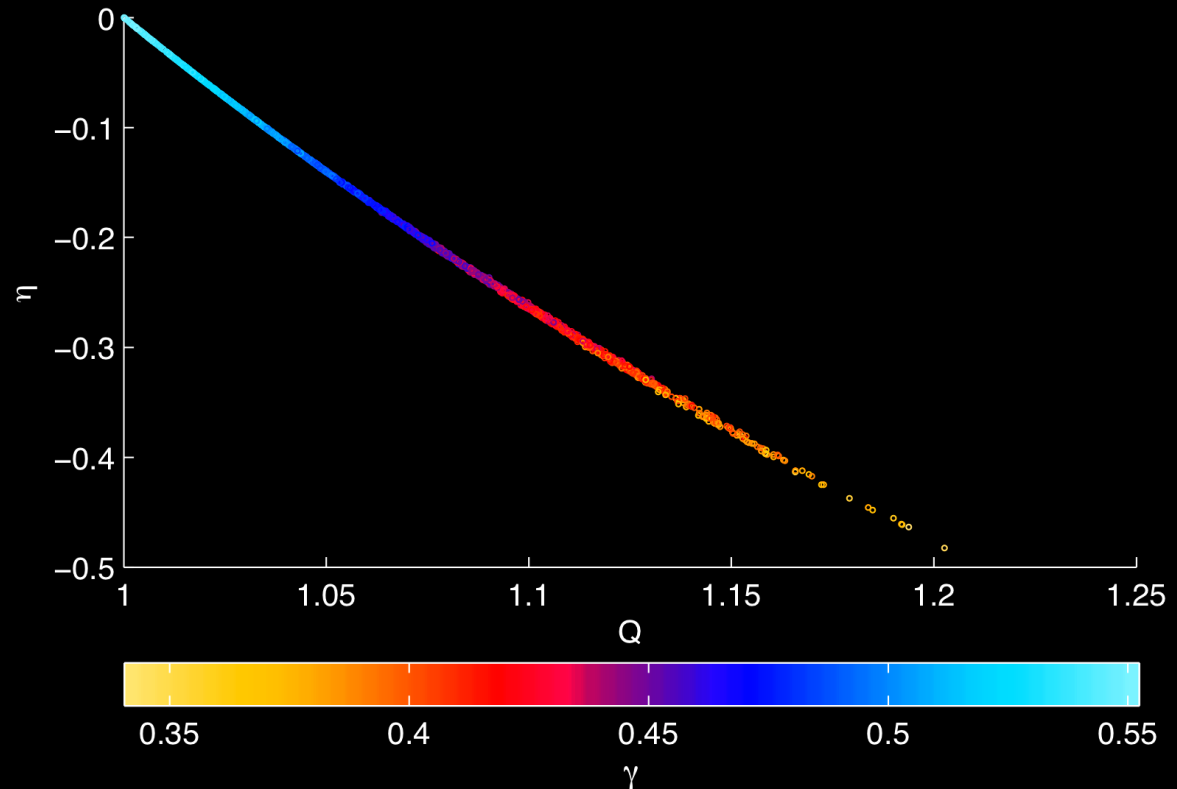
- WMAP-7yr & SN-Ia data compilation
- w, α, β constant
- $c_s = 1$ ($\rightarrow Q \approx 1$ for $\sigma = 0$)
- w consistent with -1
- α, β consistent with 0
- no signs of anything strange going on



expressed as Q and eta

(again Lukas Hollenstein)

- Projection on Q and η (on small scales)
- Need to vary also c_s^2 (δp) to access more of parameter space
- **really** need **2** extra parameters!
- again consistent with standard model



→ current constraints are weak, $O(1)$ in Q and η !
→ no deviation from standard cosmology ('GR')

current state of constraints

methodology:

'just' stick a model into a likelihood for as much data as you believe

model:

binned or parametrised variables, e.g. $\{w, Q, \eta\}$

data:

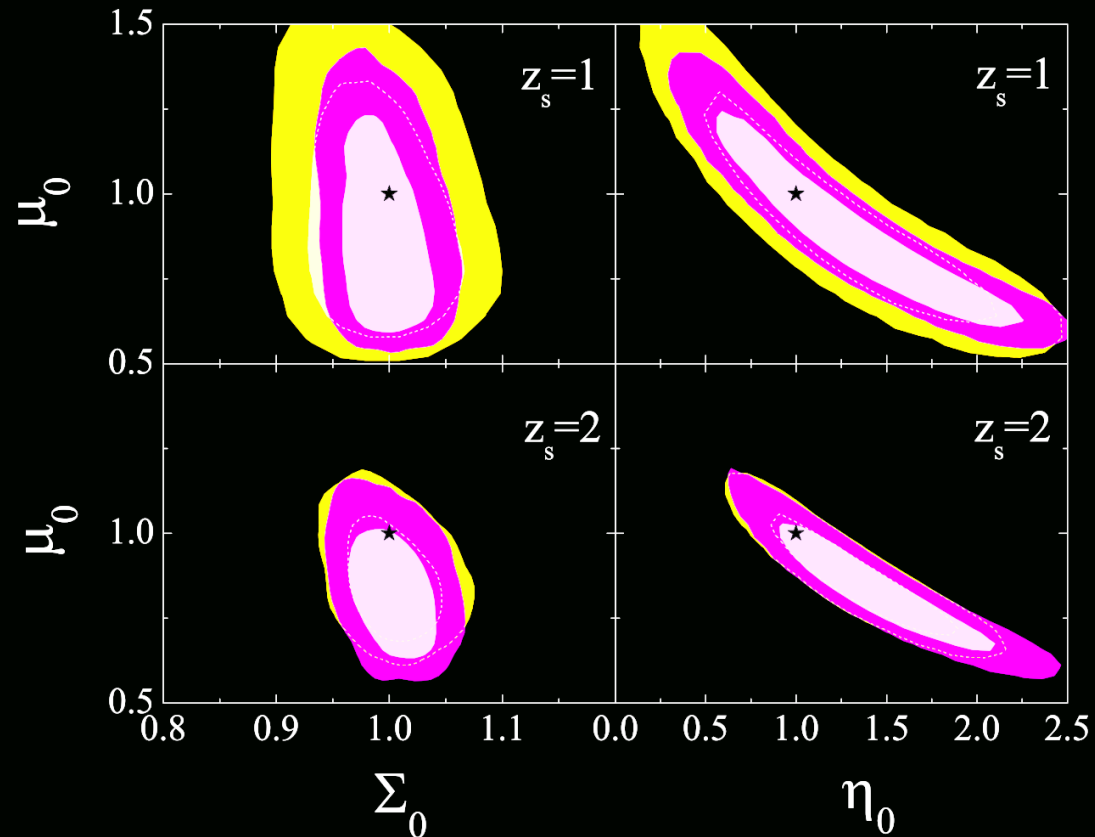
SN-Ia & BAO: constrain w

CMB: other params + **ISW**

WL: beware systematics

result:

weak constraints, no deviations from LCDM

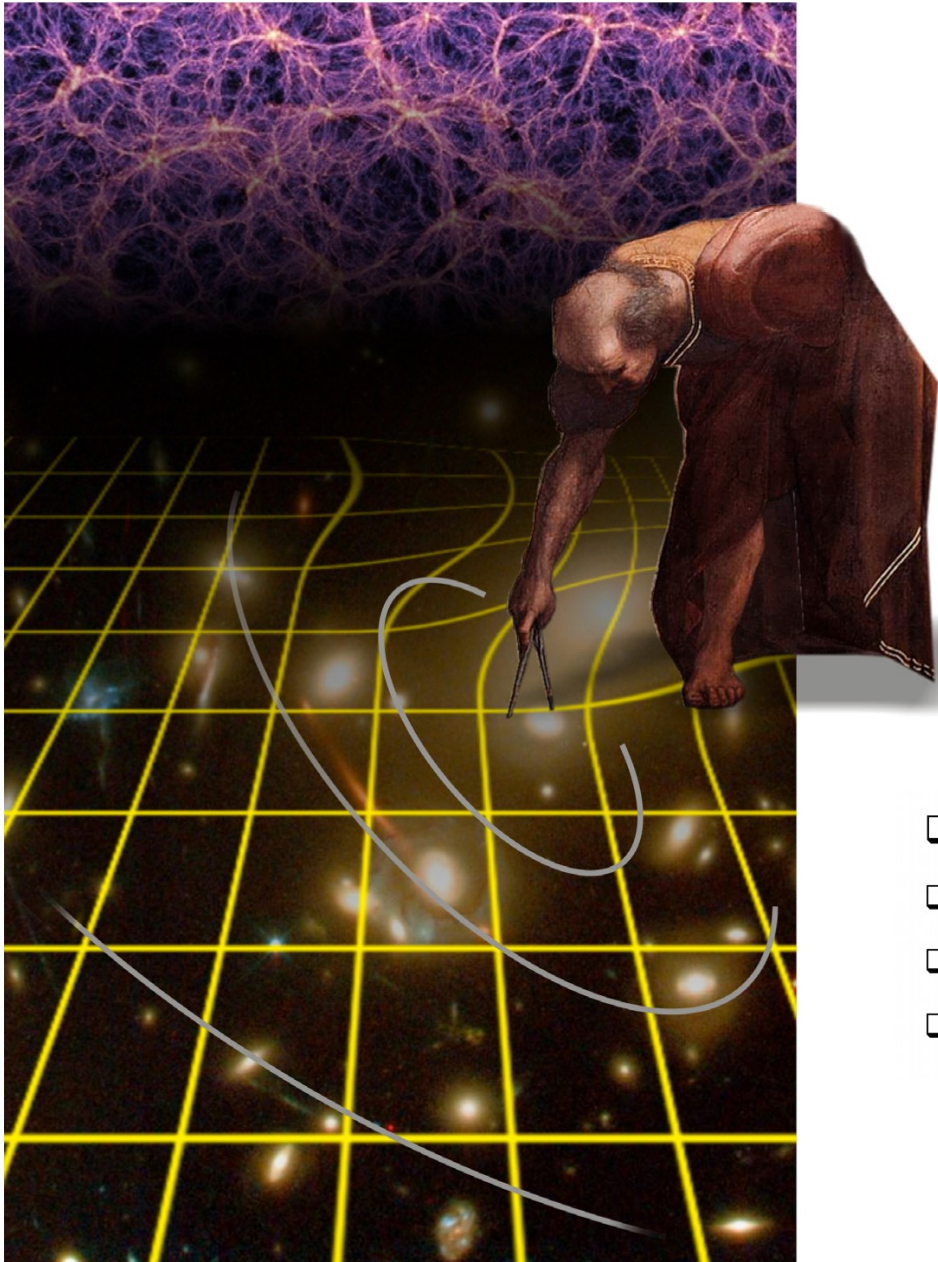


(Zhao et al 2010)

WL+CMB+ISW, model w/ transition

μ : modified Poisson eqn. in $\psi \sim Q$

Σ : lensing ($\phi+\psi$), $\eta \sim \phi/\psi$



Euclid

Mapping the Geometry of the Dark Universe

- ❑ the nature of the Dark Energy
- ❑ the nature of the Dark Matter
- ❑ the initial conditions (Inflation Physics)
- ❑ modifications to Gravity



The Euclid Mission

Primary probes:

all-sky Vis+NIR imaging and spectroscopic survey

- Weak Lensing
- Galaxy Clustering, BAO

Additional Probes: cluster counts, redshift space distortions, integrated Sachs-Wolfe effect

huge legacy data set!

other science: strong lensing, galaxy evolution, star formation, supernovae, extrasolar planets (even Earth-sized planets in habitable zone!)



conclusions

- ✓ if metric is close to FLRW, then acceleration is detected at very high significance
- ✓ behaviour is compatible with cosmological constant
- ✓ even when taking into account perturbations
- ✓ but same would have been true during inflation

- ✓ we need to improve measurements of perturbations as they are a good model discriminator
- ✓ example: anisotropic stress and modifications of GR
- ✓ first goal should be: Kill ~~Bill~~ Lambda

- ✓ need to combine probes, e.g. lensing and velocities
- ✓ Euclid would be a great mission for this purpose