

Cosmological perturbations in the UV-protected inflation

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Based on: Cristiano Germani, Gerasimos Rigopoulos, & Y. W.
in preparation

PONT Avignon, 18th April 2011

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Cosmological perturbations in gravitationally enhanced friction models of inflation

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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A scalar field ϕ is a good candidate of an Inflaton as

$$\rho = \frac{1}{2}\dot{\phi}^2 + V, p = \frac{1}{2}\dot{\phi}^2 - V$$

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

Slow Roll Inflation

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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$$\rho = \frac{1}{2}\dot{\phi}^2 + V, \quad p = \frac{1}{2}\dot{\phi}^2 - V$$

By geometrical identity (Raychaudhuri eq.)

$$\ddot{a} \propto -(\rho + 3p) \propto -(\dot{\phi}^2 - V)$$

Slow Roll Inflation

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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$$\begin{aligned} \dot{\phi}^2 &\ll V, \text{ Inflation happens ("slow roll")} \\ \dot{\phi}^2 &\sim V, \text{ Inflation ends} \end{aligned}$$

Slow Roll Inflation

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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Q: Do we know any scalar field?

Higgs Boson!

The failure of Higgs boson in Inflation:

$$V(\Phi) \simeq \lambda\Phi^4$$

$$H = \dot{a}/a \simeq \text{const} \text{ (almost de Sitter)}$$

$$\ddot{\Phi} \ll 3H\dot{\Phi} \text{ (slow roll)}$$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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$$\Phi \gg M_p$$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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$$\Phi \gg M_p$$

$R \gg M_p^2$ for the Standard Model value of $\lambda \sim 10^{-1}$

Inflation happens during the Quantum Gravity regime!!

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

How to lower the curvature during Inflation?

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

How to lower the curvature during Inflation?

Repeat:

$$R \sim H^2 \propto \frac{V(\Phi)}{M_p^2} \propto \frac{\Phi^4}{M_p^2}$$

$$\dot{\Phi} \simeq -M_p \Phi$$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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$$\dot{\Phi} \simeq -\frac{M_p \Phi}{\mu}$$

$$\epsilon = -\frac{\dot{H}}{H^2} \propto \frac{\dot{\Phi}^2}{H^2 M_p^2} \sim \frac{M_p^2}{\Phi^2 \mu^2} \ll 1 \Rightarrow \Phi \gg \frac{M_p}{\mu}$$



$$\frac{M_p^2}{\mu^4} \ll R \ll M_p^2$$

Solution: $3\mu H\dot{\Phi} \simeq -V'$ **Increase the friction!**

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

New Higgs Inflation

[Germani & Kehagias 2010]

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$$\frac{M_p^2}{\mu^4} \ll R \ll M_p^2$$

Quantum Gravity regime is avoided during Inflation!

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Gravitationally Enhanced Friction (GEF)

To form eventually Galaxies, the friction should only be efficient at high energies:

$$3\mu H = f(H), \quad \frac{df}{dH} \geq 0$$

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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If no new d.o.f. are added, the scalar e.o.m. can only be

$$\mu \left(\ddot{\Phi} + 3H\dot{\Phi} \right) = -V' \rightarrow t_{\text{eff}} \simeq \frac{t}{\sqrt{\mu}} \text{ as } \frac{\dot{\mu}}{\mu H} \ll 1$$

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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If $H \gg M$ the scalar field clock is **frozen** w.r.t. the observer and friction is enhanced.

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

Gravitationally Enhanced Friction: Realization

In order to realize covariantly the enhanced friction, we promote the rescaling to all coords.

$$\partial_\mu \rightarrow \sqrt{\mu} \partial_\mu, \quad \mu = 1 + \frac{H^2}{M^2}$$

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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The enhanced friction is covariantly realized by shifting the kinetic action

$$g^{\mu\nu} \partial_\mu \Phi \partial_\nu \Phi \rightarrow \left(g^{\mu\nu} - \frac{G^{\mu\nu}}{M^2} \right) \partial_\mu \Phi \partial_\nu \Phi$$

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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Thanks to Bianchi identities, this action has the properties:

- It is shift (Galilean) invariant $\Phi \rightarrow \Phi + c + c_\mu x^\mu$

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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- Propagates only spin 0 (scalar) and spin 2 (graviton) particles (no higher derivatives, Lapse and Shift are still Lagrange multipliers)

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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- It is shift (Galilean) invariant $\Phi \rightarrow \Phi + c + c_\mu x^\mu$
- Propagates only spin 0 (scalar) and spin 2 (graviton) particles (no higher derivatives, Lapse and Shift are still Lagrange multipliers)
- Makes harder for the scalar field to roll down the potential!

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

UV Protected Inflation

[Germani & Kehagias 2011]

In large field scenarios

$\Phi \gg \Lambda_{cut-off} \leftarrow$ not only M_p !

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

UV Protected Inflation

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The (unknown) UV completed theory may spoil the effective Inflaton potential by higher powers of $\Phi^2/\Lambda_{cut-off}^2$

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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In small field scenarios ($\Phi \ll \Lambda_{cut-off}$) there is no problem!

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

- Suppose some global (chiral) symmetry is broken at energies $f > \text{TeV}$ (like in the QCD axion case)

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

- Suppose some global (chiral) symmetry is broken at energies $f > \text{TeV}$ (like in the QCD axion case)
- a Pseudo Nambu-Goldstone Boson Φ is produced with a potential **at one loop (chiral anomaly)**

$$V(\Phi) = \Lambda^4 \cos^2\left(\frac{\Phi}{f}\right) \simeq \Lambda^4 \left(1 - \frac{\Phi^2}{f^2}\right)$$

which is protected by the restoration of global shift symmetry $\Phi \rightarrow \Phi + c$ as $\Lambda \rightarrow 0$ up to M_p^{-1}

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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- If $\Lambda \sim 10^{16}$ GeV (GUT scale), Inflation is produced with

$$n_s - 1 \propto \epsilon \simeq -\frac{M_p^2}{8\pi f^2}$$

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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- so $n_s - 1 \simeq -0.04 \rightarrow f > M_p$
 \Rightarrow **the model may not be trusted!**

Resolution

Once again we can increase the friction so that

$$\epsilon \rightarrow \frac{\epsilon_{old}}{\mu} \Rightarrow n_s - 1 \sim -\frac{M_p^2}{8\pi f \mu}$$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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$$\epsilon \rightarrow \frac{\epsilon_{old}}{\mu} \Rightarrow n_s - 1 \sim -\frac{M_p^2}{8\pi f \mu}$$

Then for large enough μ , $f \ll M_p$

The model is Natural!

(i.e. no UV modifications of the potential)

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Resolution

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Resolution

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Resolution

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This is due to the fact that the new coupling $G^{\alpha\beta} \partial_\alpha \Phi \partial_\beta \Phi$ is the unique that

- *Does not* introduce new degrees of freedom
- *Is invariant* under the global unbroken symmetry
 $\Phi \rightarrow \Phi + c$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

What about the “eta” problem?

QG does not like global symmetries so that

$$V \rightarrow V \times \left(1 + c \frac{\phi^2}{M_p^2} + \dots\right)$$

Then

$$\eta_{qg} = \frac{V''}{3V} M_p^2 \simeq \eta + c \frac{2}{3} \geq 1, \text{ if } c \sim \mathcal{O}(1)!$$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Again the increment of friction solves the problem as

$$\eta \rightarrow \frac{\eta}{\mu} \Rightarrow c_{\text{eff}} = \frac{c}{\mu} \ll 1$$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Cosmological perturbations in GEF

ADM form

$$ds^2 = -N^2 dt^2 + h_{ij}(dx^i + N^i dt)^2$$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

**Cosmological
perturbations**

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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- Vary wrt the constraints N, N^i , substitute back into the action and canonically normalize ζ and γ_{ij}
- $N = \frac{\Gamma}{H} \dot{\zeta}$, $N^i = -\frac{\Gamma}{H} \partial_i \zeta + \frac{\Sigma}{H^2} \partial_i \partial^{-2} \dot{\zeta}$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Cosmological perturbations in GEF

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$$ds^2 = -N^2 dt^2 + h_{ij}(dx^i + N^i dt)^2$$

- Use the gauge $\delta\Phi = 0$
- then: $h_{ij} = a^2 \left[(1 + 2 \underbrace{\zeta}_{\text{scalar perturbation}}) \delta_{ij} + \underbrace{\gamma_{ij}}_{\text{tensor perturbation}} \right]$
- Vary wrt the constraints N, N^i , substitute back into the action and canonically normalize ζ and γ_{ij}
- $N = \frac{\Gamma}{H} \dot{\zeta}$, $N^i = -\frac{\Gamma}{H} \partial_i \zeta + \frac{\Sigma}{H^2} \partial_i \partial^{-2} \dot{\zeta}$
- $\Gamma(\dot{\Phi}, H, M) \simeq 1 + \epsilon$, $\Sigma(\dot{\Phi}, H, M) \simeq \mathcal{O}(\epsilon)$ for $H \gg M$

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Matching with the best fit WMAP 7-years

$$\Phi_0 \sim 10^{-2}M_p, H \sim 10^{-4}M_p, M \sim 10^{-8}M_p \text{ (Higgs)}$$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Since the new non-linear interaction and $c_s^2 < 1$, any
Non-Gaussianity?

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Red tilt in UV Protected Inflation

■ $n_s - 1 \simeq -6\epsilon + \frac{2}{3} \frac{M_p^2 V''}{V} \frac{M^2}{H^2} < 0 \Rightarrow$ Red spectrum!

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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$$\frac{\Lambda^2}{M_p^2} \sim 10^{-4} \frac{\Phi_0}{f}, \quad \frac{M^2}{H^2} \sim 10 \frac{f^2}{M_p^2} \text{ (pNGB)}$$

- Consistent with the theoretical hierarchy of scales:

$$M_p^4 \gg \Lambda^4 \gg M^2 M_p^2, \quad \frac{M M_p}{\Lambda^2} \ll \frac{f}{M_p} \ll 1$$

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Red tilt in UV Protected Inflation

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- Additional data to fix?

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Preliminary: Non-Gaussianity in GEF

$$L_3 = \epsilon^2 [aC_1 M_p^2 \zeta (\partial_i \zeta)^2 + aC_2 M_p \dot{\zeta} (\partial_i \zeta)^2 + a^3 C_3 M_p^2 \zeta \dot{\zeta}^2 + a^3 C_4 M_p \dot{\zeta}^3 + aC_5 \dot{\zeta}^2 \partial_i^2 \zeta] + \mathcal{O}(\epsilon^3)$$

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Preliminary: Non-Gaussianity in GEF

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Preliminary: Non-Gaussianity in GEF

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Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

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Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

Conclusions

Slow Roll Inflation

The failure of
Higgs boson

Lowering the
curvature during
Inflation

Gravitationally
Enhanced Friction

UV Protected
Inflation

Cosmological
perturbations

Conclusions

Conclusions

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Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions

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- Consistent with WMAP 7-years result.
- Non-Gaussianities from Inflation models with GEF are small.

Slow Roll Inflation

The failure of Higgs boson

Lowering the curvature during Inflation

Gravitationally Enhanced Friction

UV Protected Inflation

Cosmological perturbations

Conclusions